

# Back to the phase space: Thermal Axions

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DI PADOVA

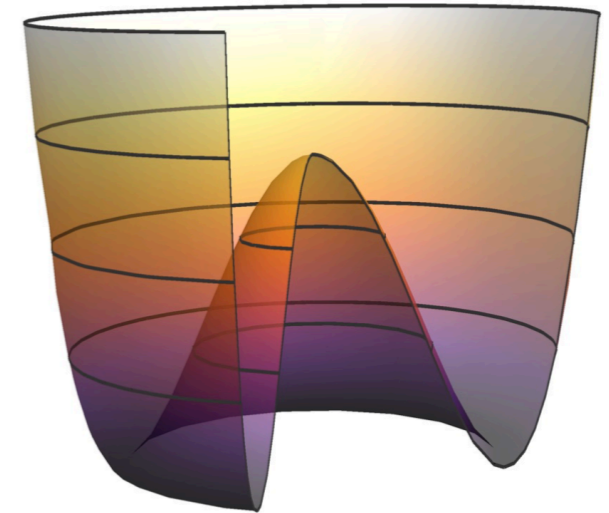


Cosmology, Astrophysics, Theory and Collider Higgs 2024 (CATCH22+2)  
Dublin Institute for Advanced Studies (DIAS), 2 May 2024

# The QCD Axion

## New global $U(1)_{PQ}$ symmetry

- spontaneously broken at the scale  $f_a$  (with  $f_a \gg$  weak scale)
- anomalous under strong interactions



$$\frac{a}{f_a} \frac{\alpha_s}{8\pi} G^{\mu\nu} \tilde{G}_{\mu\nu}$$

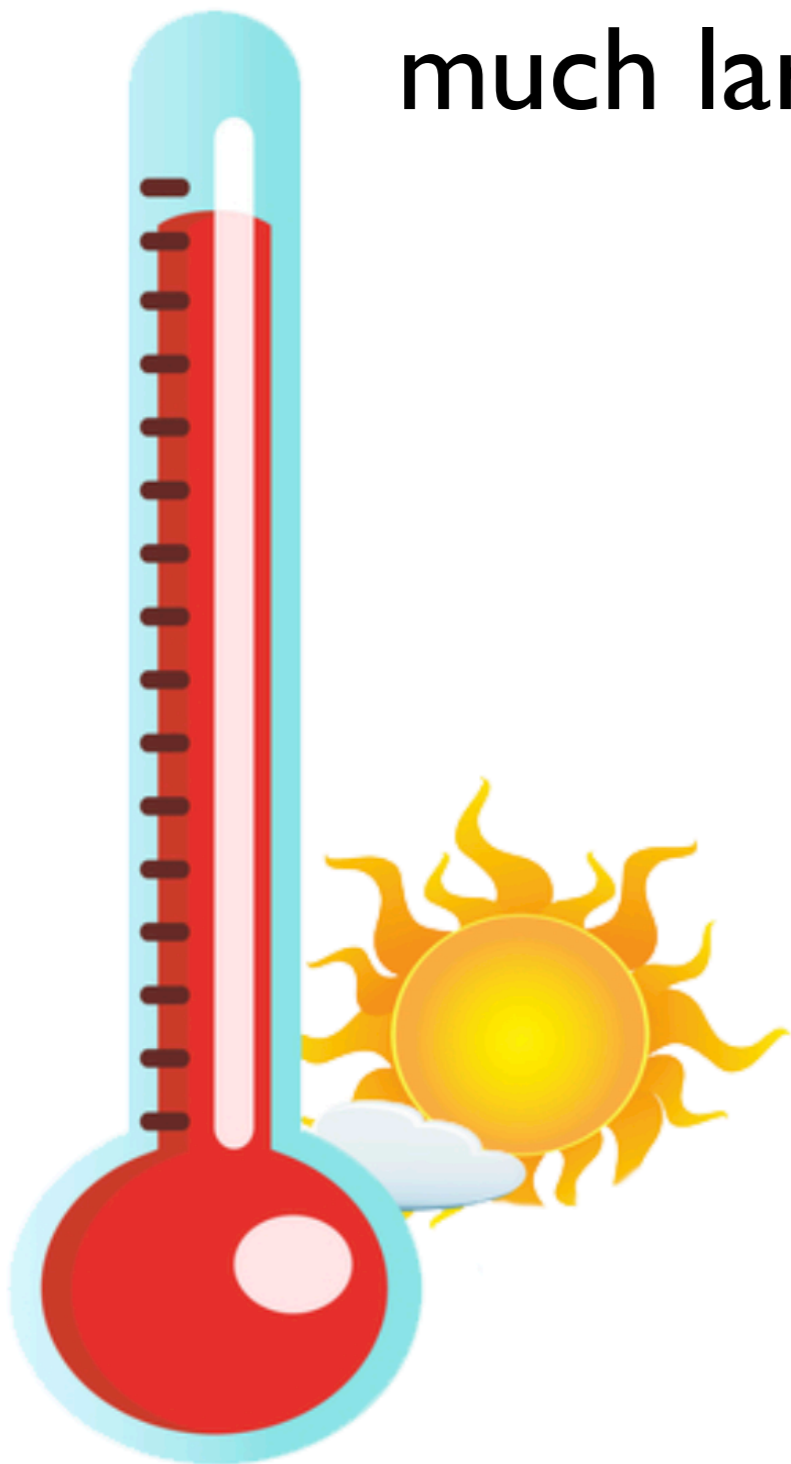
$$m_a \simeq 5.7 \left( \frac{10^{12} \text{ GeV}}{f_a} \right) \mu\text{eV}$$

Results in this talk mostly about the QCD axion

Easy to extend to ALPs (especially when the mass is negligible)

# Hot Axions

Axions produced with kinetic energy much larger than their mass (i.e. “hot”)



Additional radiation at:

- BBN ( $m_a \lesssim \text{MeV}$ )
- CMB formation ( $m_a \lesssim 0.3 \text{ eV}$ )

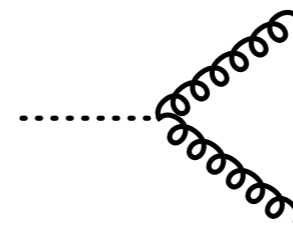
$$\rho_{\text{rad}} = \left[ 1 + \frac{7}{8} \left( \frac{T_\nu}{T_\gamma} \right)^4 N_{\text{eff}} \right] \rho_\gamma$$

$$\Delta N_{\text{eff}} = \frac{8}{7} \left( \frac{11}{4} \right)^{4/3} \frac{\rho_a}{\rho_\gamma}$$

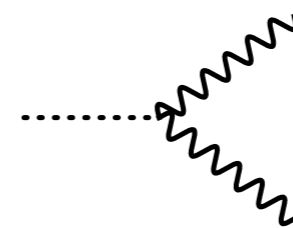
# Thermal Production

Scatterings and/or decays involving  
primordial thermal bath particles  
(axion energy  $\gg m_a$ , i.e. “hot”)

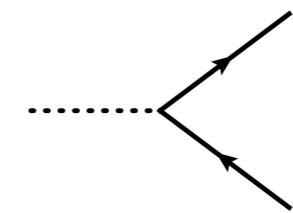
Unavoidable  
Production Source!



$$\frac{\alpha_s}{8\pi} \frac{a}{f_a} G^{\mu\nu} \tilde{G}_{\mu\nu}$$



$$c_{\gamma\gamma} \frac{\alpha_{\text{em}}}{8\pi} \frac{a}{f_a} F^{\mu\nu} \tilde{F}_{\mu\nu}$$



$$c_\psi \frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \gamma^5 \psi$$



# Thermal Production

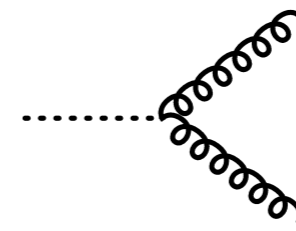
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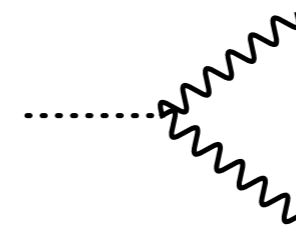
## Computing $\Delta N_{\text{eff}}$ - I

### Instantaneous decoupling

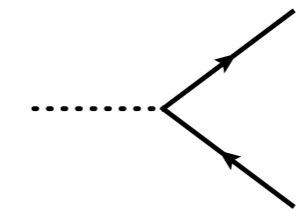
- Assume they thermalize at early times
- Estimate the decoupling temperature,  $\Gamma(T_D) = H(T_D)$ , and the resulting  $\Delta N_{\text{eff}}$



$$\frac{\alpha_s}{8\pi} \frac{a}{f_a} G^{\mu\nu} \tilde{G}_{\mu\nu}$$



$$c_{\gamma\gamma} \frac{\alpha_{\text{em}}}{8\pi} \frac{a}{f_a} F^{\mu\nu} \tilde{F}_{\mu\nu}$$



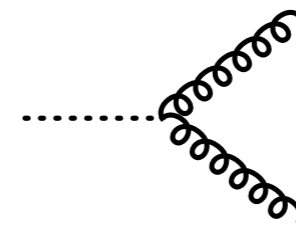
$$c_{\psi} \frac{\partial_{\mu} a}{f_a} \bar{\psi} \gamma^{\mu} \gamma^5 \psi$$

$$\Delta N_{\text{eff}} \simeq 0.027 \left( \frac{106.75}{g_{*s}(T_D)} \right)^{4/3}$$

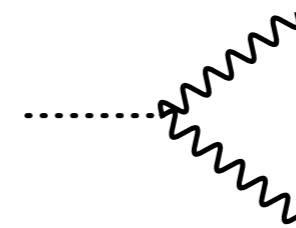
# Thermal Production

Scatterings and/or decays involving primordial thermal bath particles (axion energy  $\gg m_a$ , i.e. “hot”)

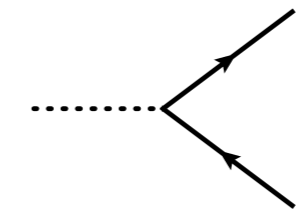
Unavoidable  
Production Source!



$$\frac{\alpha_s}{8\pi} \frac{a}{f_a} G^{\mu\nu} \tilde{G}_{\mu\nu}$$



$$c_{\gamma\gamma} \frac{\alpha_{\text{em}}}{8\pi} \frac{a}{f_a} F^{\mu\nu} \tilde{F}_{\mu\nu}$$



$$c_{\psi} \frac{\partial_{\mu} a}{f_a} \bar{\psi} \gamma^{\mu} \gamma^5 \psi$$

## Computing $\Delta N_{\text{eff}}$ - II

### Boltzmann equation for $n_a$

- Track the number density of axions
- Convert the asymptotic result to  $\Delta N_{\text{eff}}$  via the equilibrium distribution

$$\frac{dn_a}{dt} + 3Hn_a = \sum_{\alpha} \gamma_{\alpha}$$

$$\Delta N_{\text{eff}} \simeq 74.85 Y_a^{4/3}$$

# Thermal Production



Equilibrium thermodynamics for the conversion to energy  
Spectral distortions neglected  
Maxwell-Boltzmann statistics (i.e., no quantum effects)  
Static thermal bath (i.e., no energy exchanged)

## Computing $\Delta N_{\text{eff}}$ - II

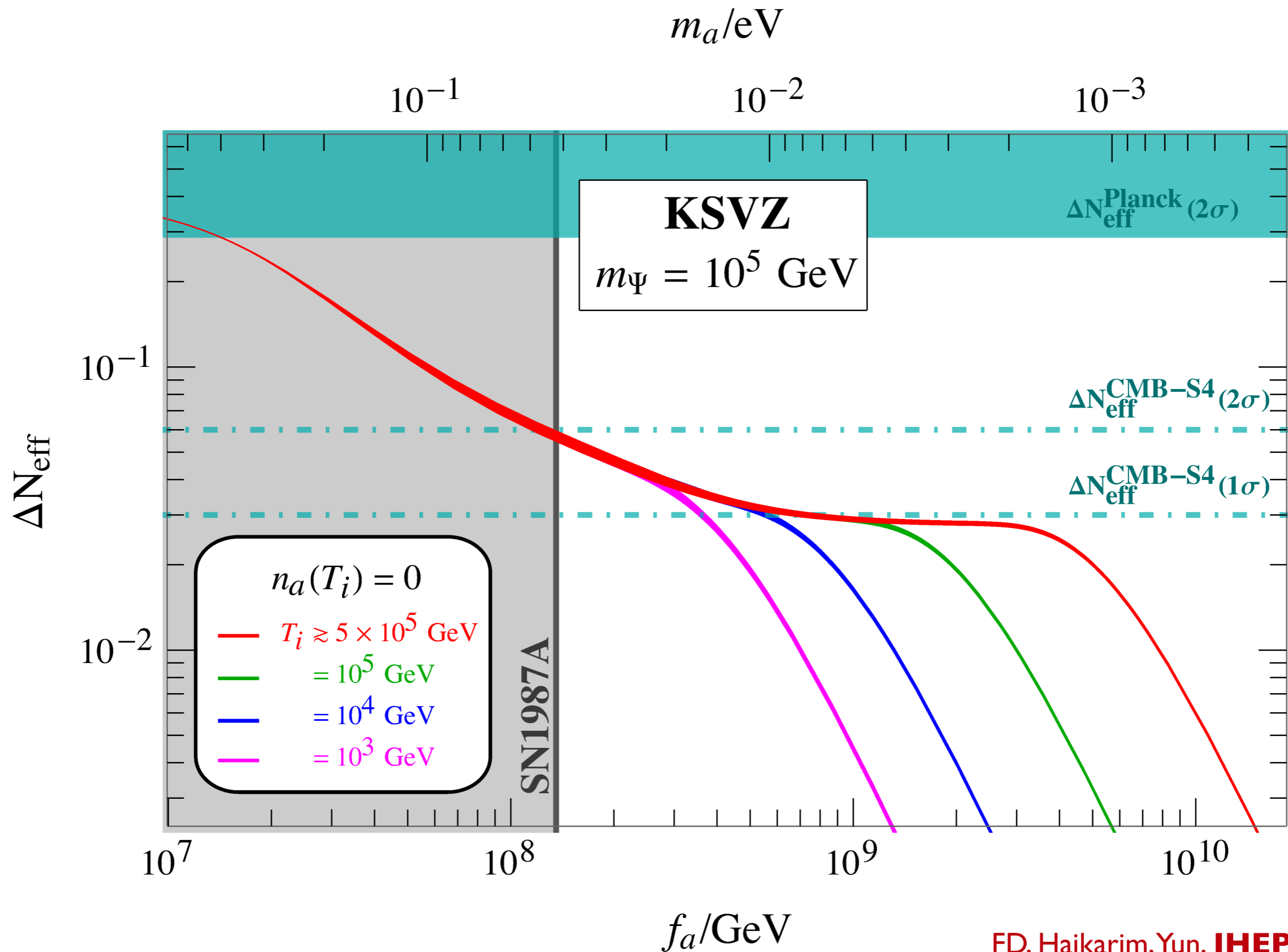
### **Boltzmann equation for $n_a$**

- Track the number density of axions
- Convert the asymptotic result to  $\Delta N_{\text{eff}}$  via the equilibrium distribution

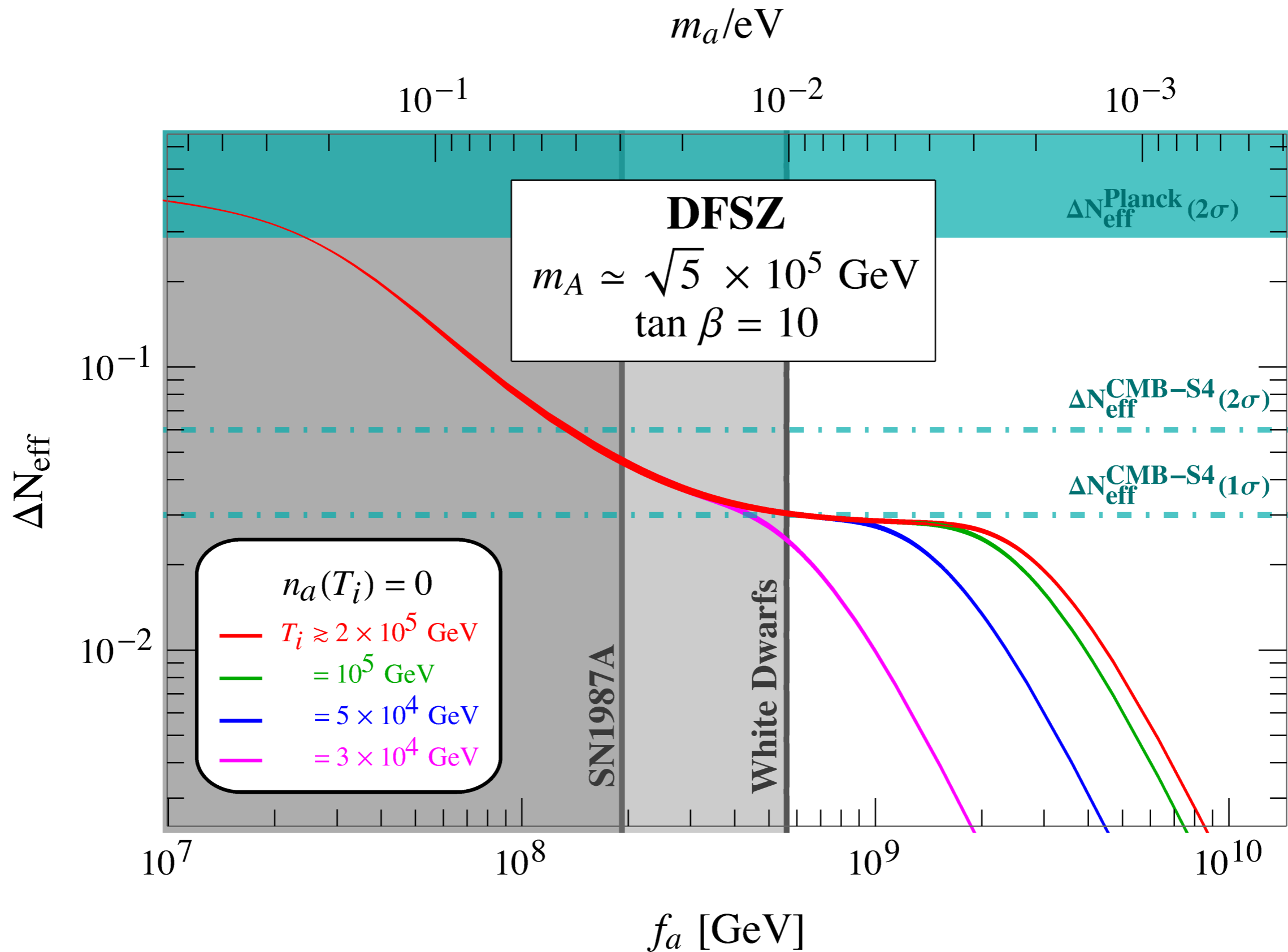
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# KSVZ Axion



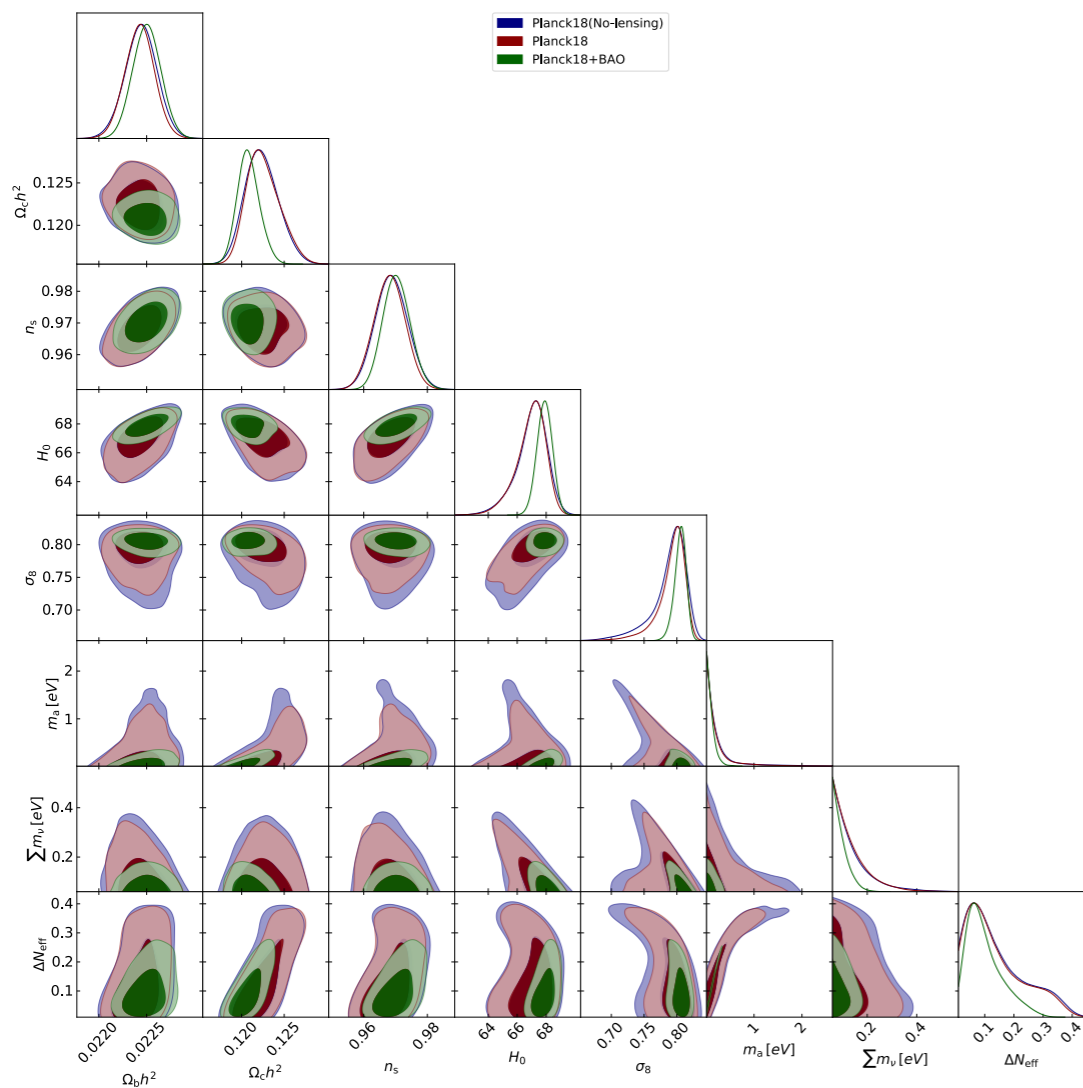
# DFSZ Axion





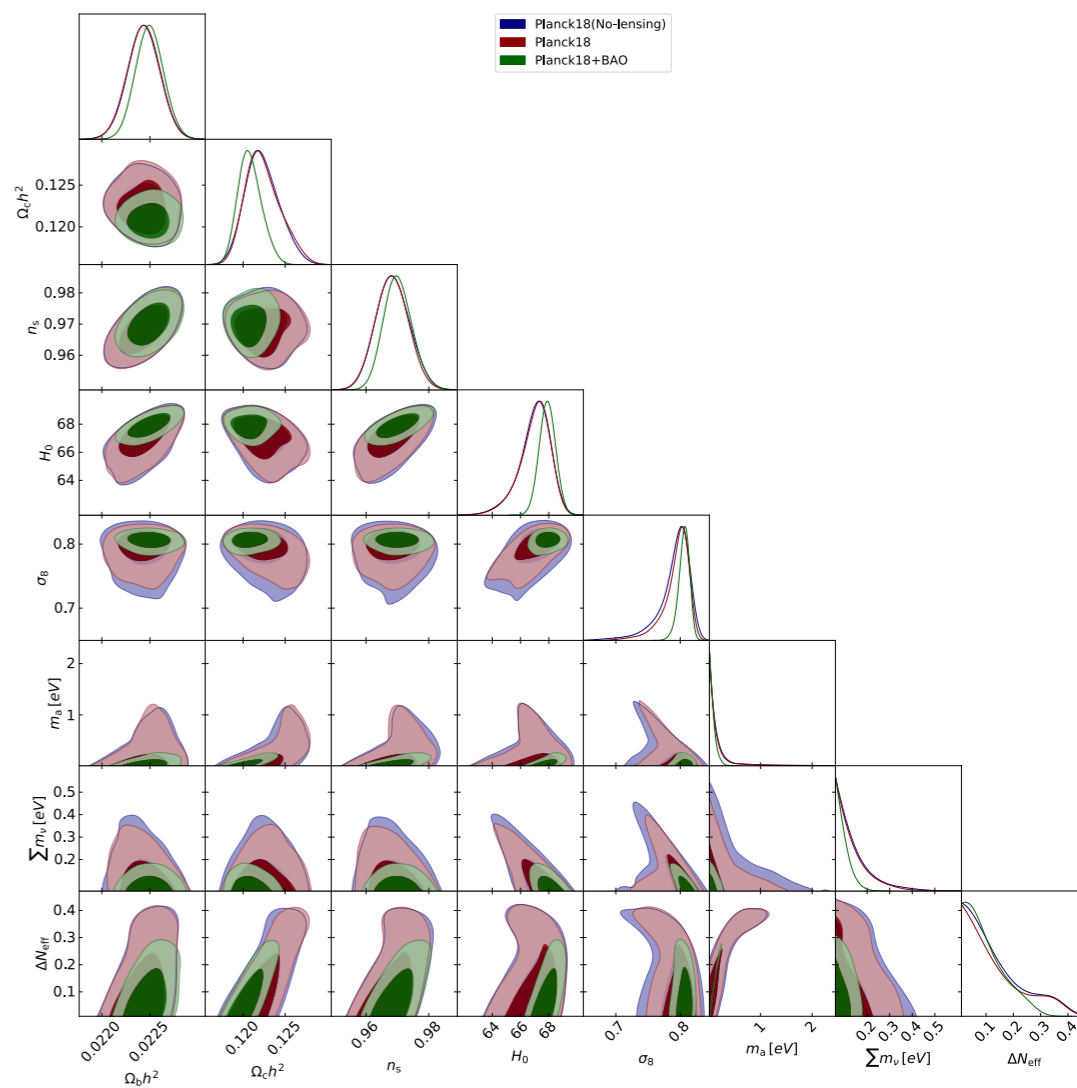
# Axion Mass Bound

## KSVZ



$$m_a \leq 0.282(0.420) \text{ eV}$$

## DFSZ



$$m_a \leq 0.209(0.293) \text{ eV}$$

# Back to the Phase-Space

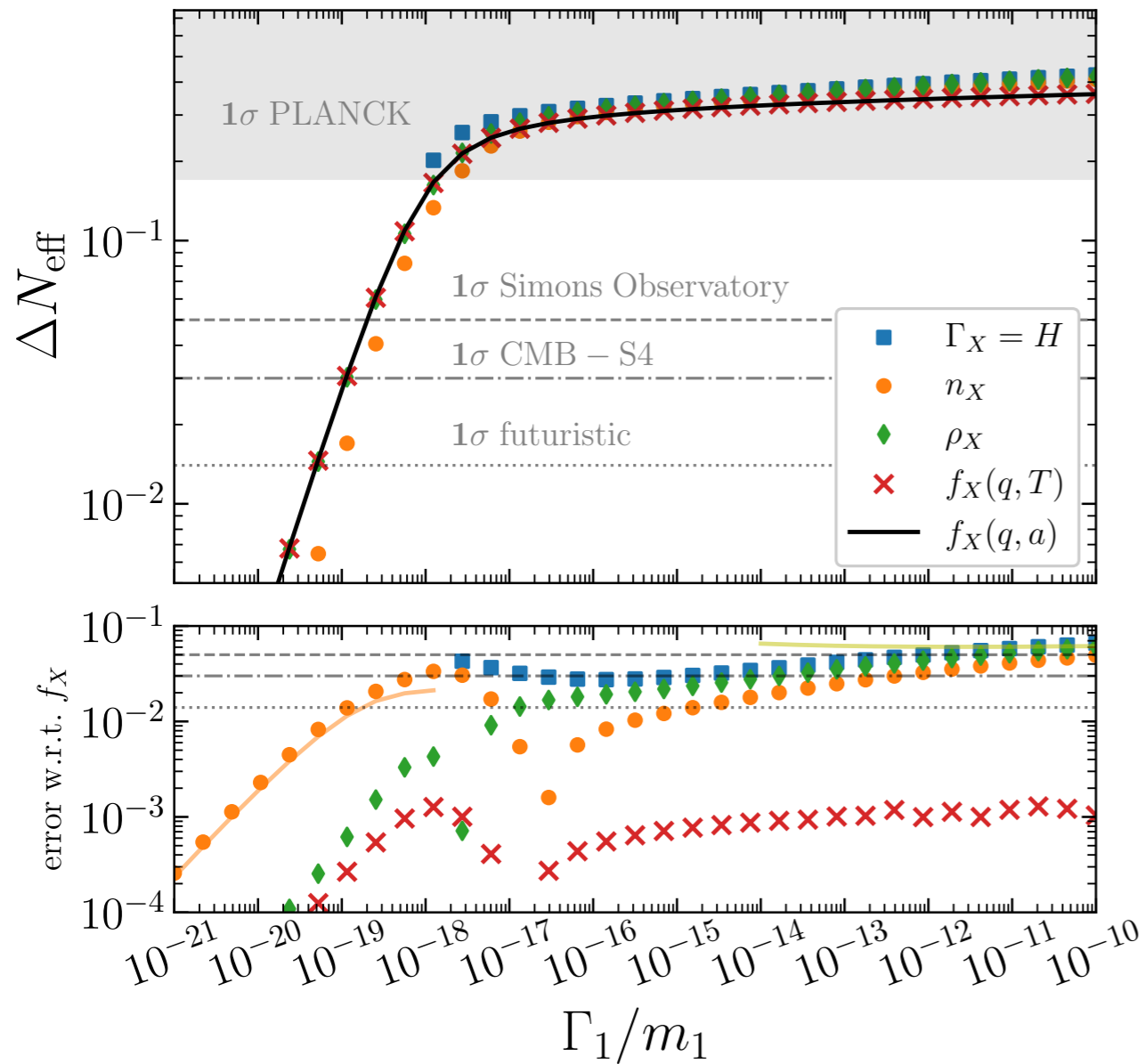
Model-independent analysis:  $\mathcal{B}_1 \dots \mathcal{B}_n \rightarrow \mathcal{B}_{n+1} \dots \mathcal{B}_m X$   
generic production of a light  $X$

$$\frac{df_X(k, t)}{dt} = \left( 1 - \frac{f_X(k, t)}{f_X^{\text{eq}}(k, t)} \right) \mathcal{C}_{n \rightarrow m X}(k, t)$$

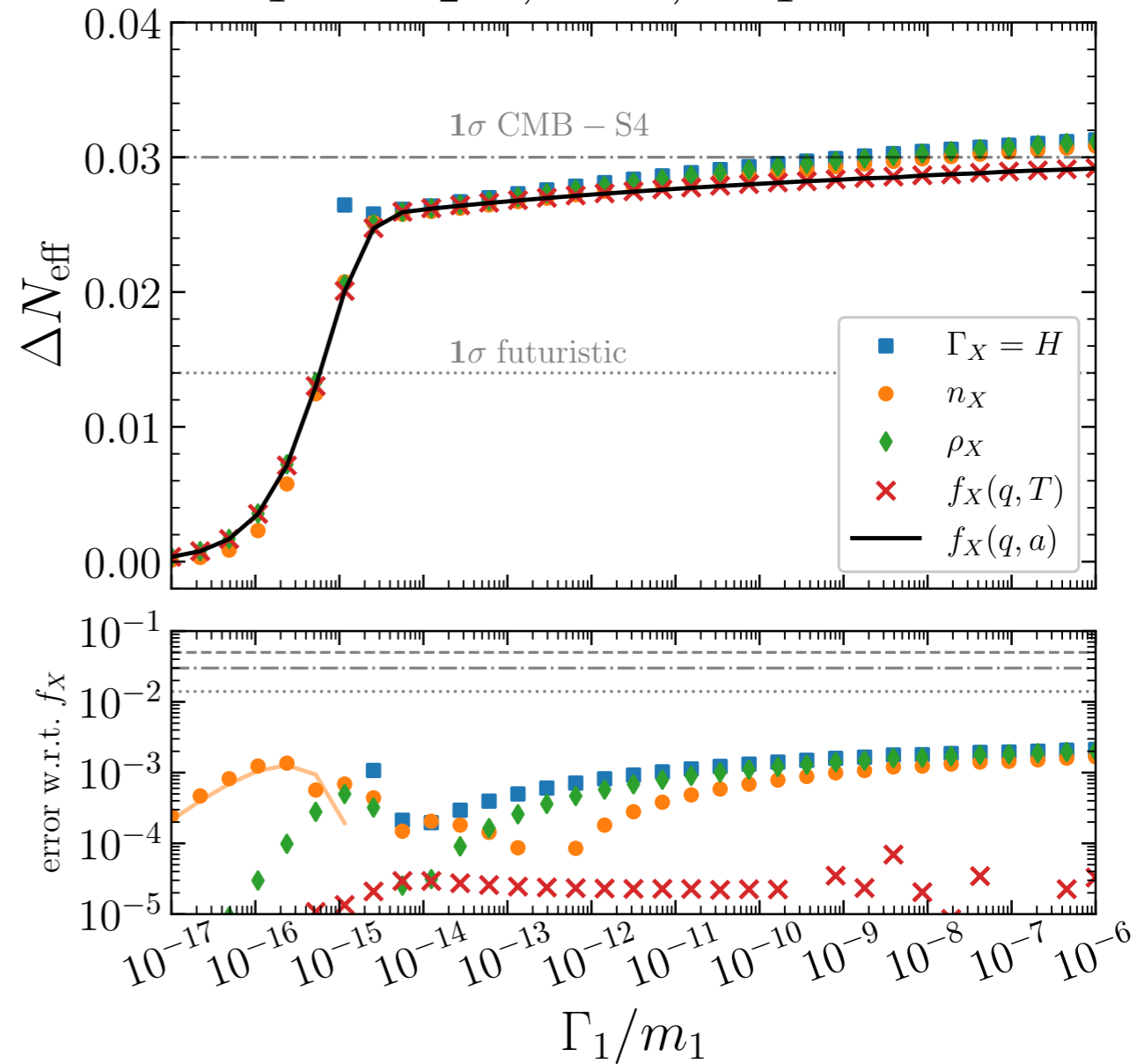
1. Keep track of phase-space and compute the energy density
2. Quantum statistical effects take into account
3. Energy exchanged with the thermal bath accounted for

# Error in predicting $\Delta N_{\text{eff}}$

$\mathcal{B}_1 \rightarrow \mathcal{B}_2 X$ , MB,  $m_1 = 1 \text{ GeV}$



$\mathcal{B}_1 \rightarrow \mathcal{B}_2 X$ , MB,  $m_1 = 1 \text{ TeV}$



# Axion-Fermion Interactions

$$\mathcal{L}_{\text{int}} = \frac{\partial_{\mu} a}{2f_a} \sum_{\psi} c_{\psi} \bar{\psi} \gamma^{\mu} \gamma_5 \psi$$

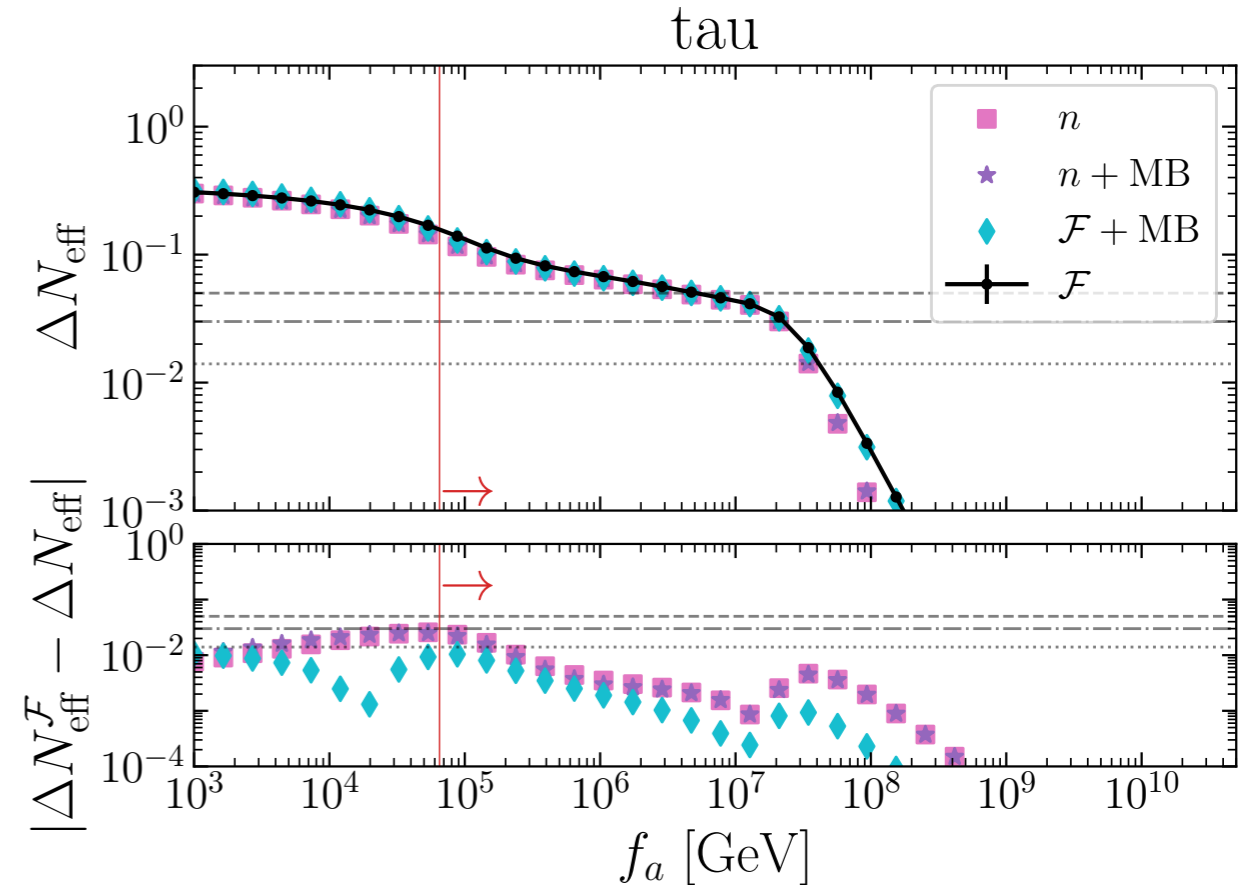
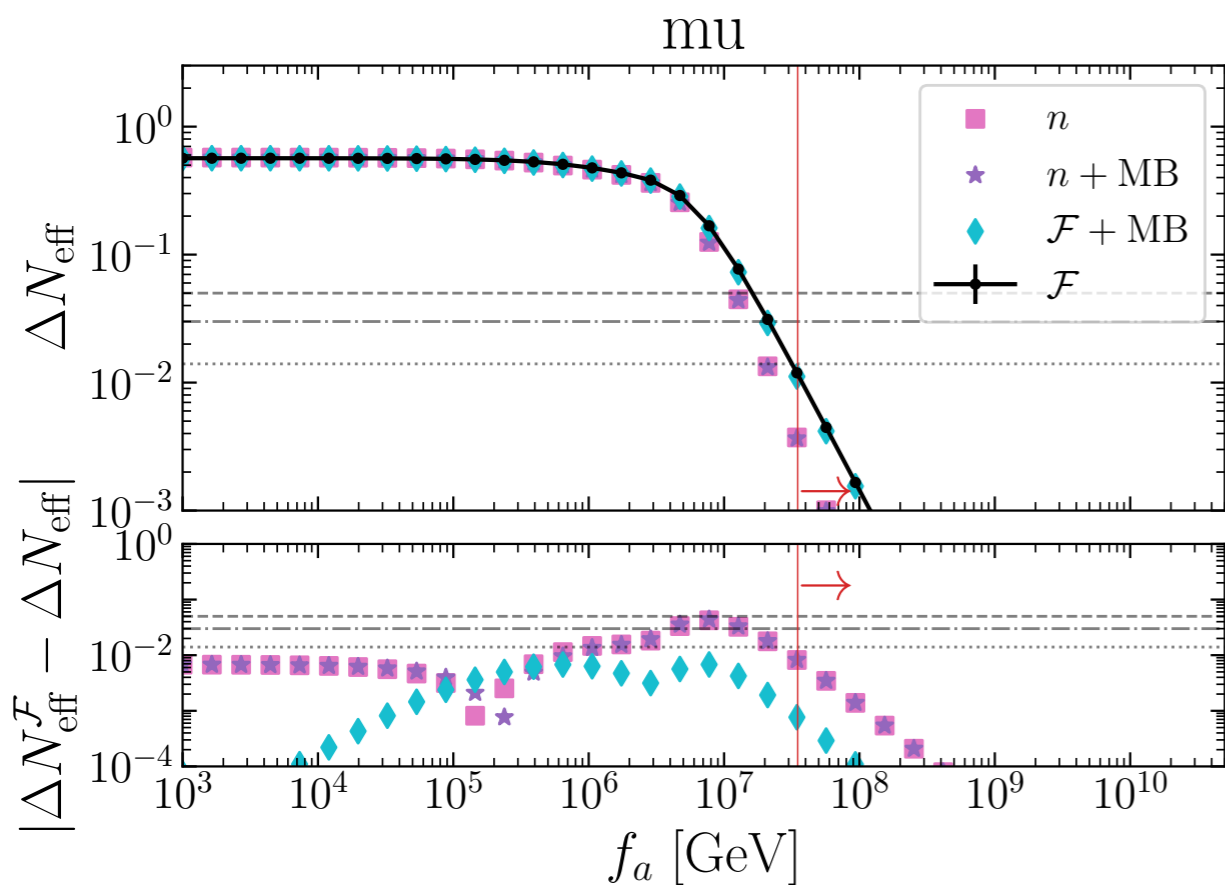
Recent studies performed  
by tracking the axion  
number density

Baumann et al, Phys.Rev.Lett. 117 (2016)  
Ferreira, Notari, Phys.Rev.Lett. 120 (2018)  
FD et al, **JCAP 11 (2018)**  
Arias-Aragón et al., **JCAP 11 (2020)**  
Arias-Aragón et al., **JCAP 03 (2021)**  
Green et al., JCAP 02 (2022)

Will it change if we go back to the phase space?

# Axion-Fermion Interactions

PRELIMINARY

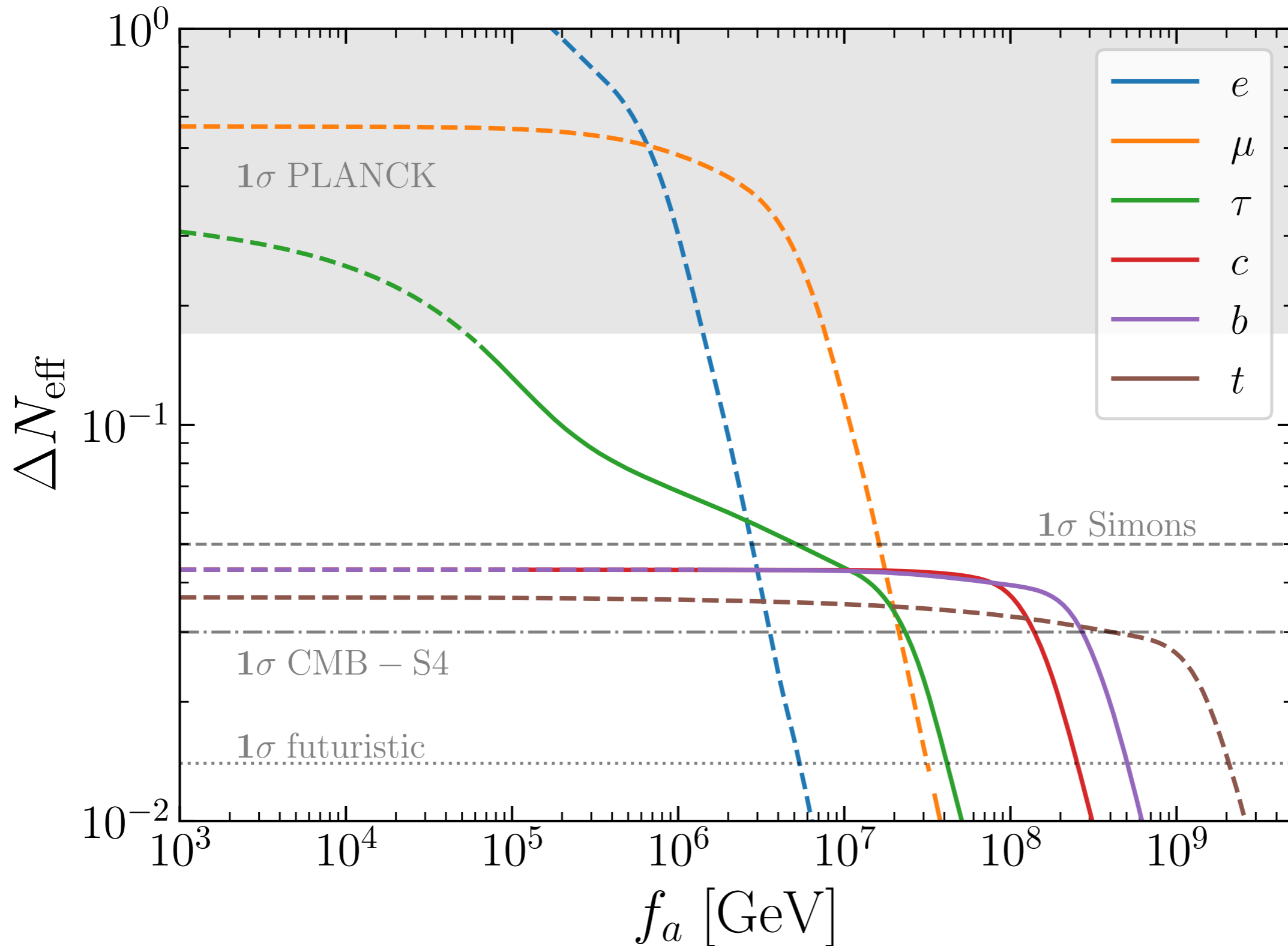


Difference detectable by future CMB-S4 surveys!

- MUON: effect maximum in regions in tension with stellar bounds
- TAU: effect maximum in allowed regions



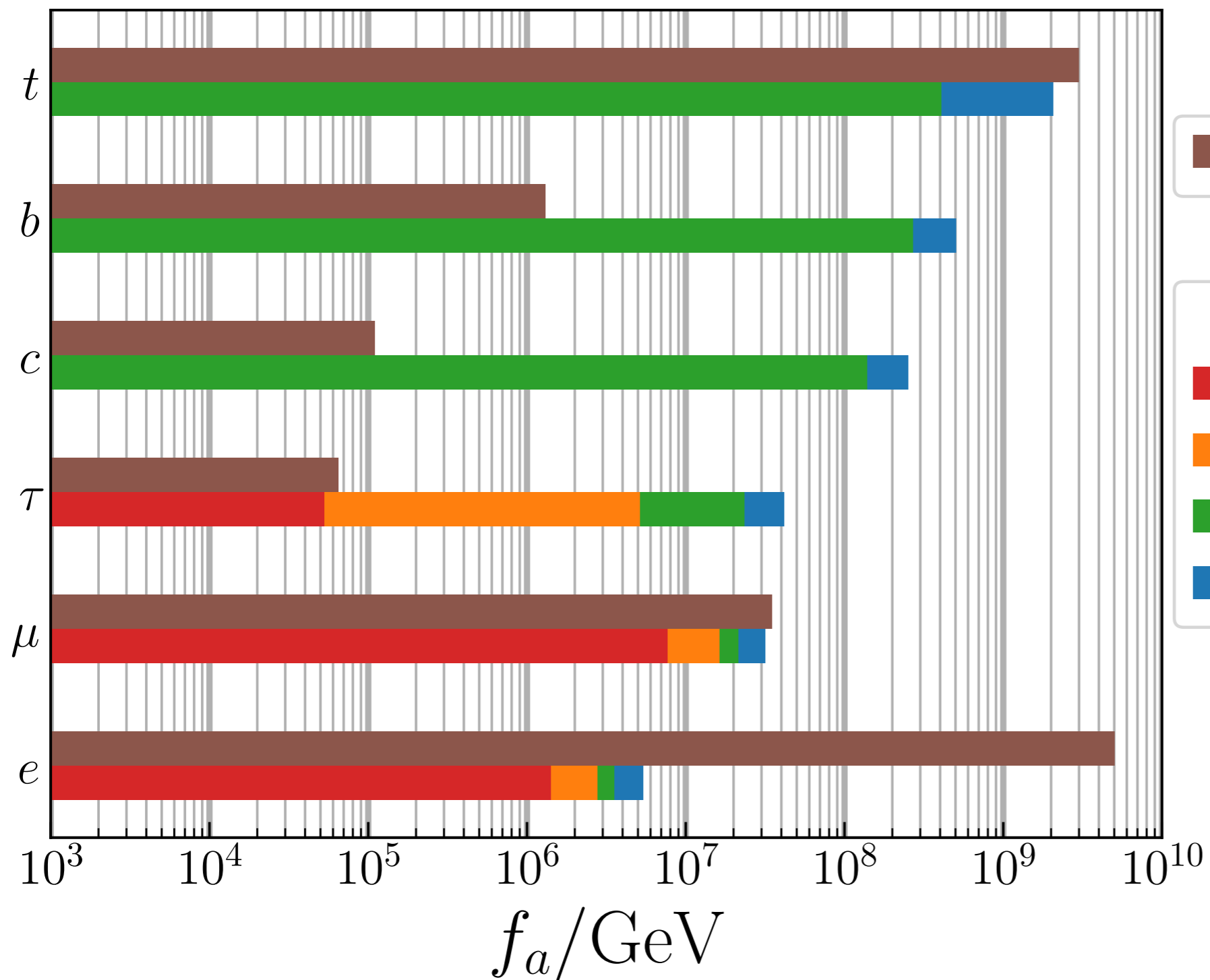
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PRELIMINARY

# Axion-Fermion Interactions

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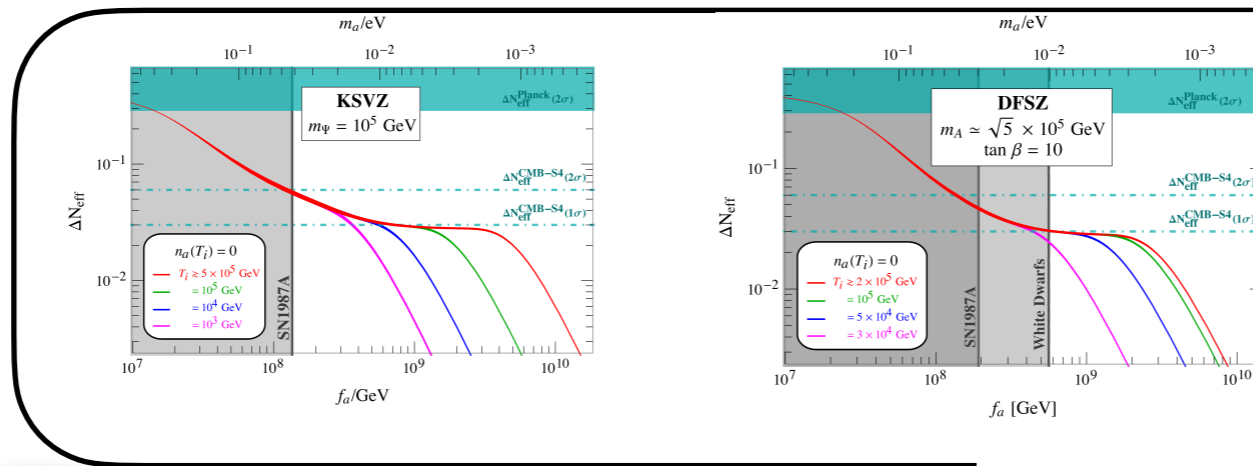


astrophysics

$\Delta N_{\text{eff}}$

- Planck
- Simons
- CMB-S4
- futuristic

# The Way Back to the Phase Space



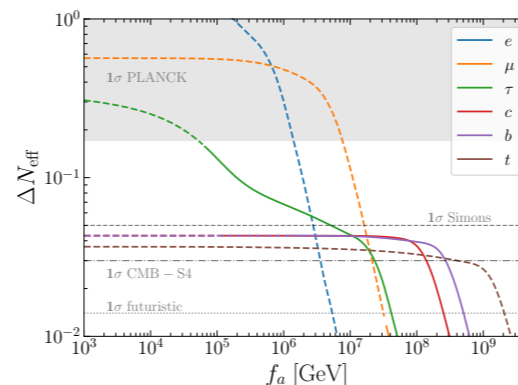
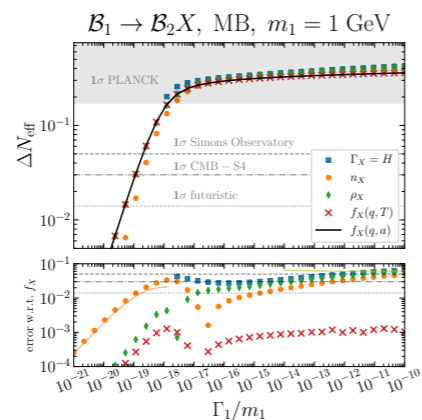
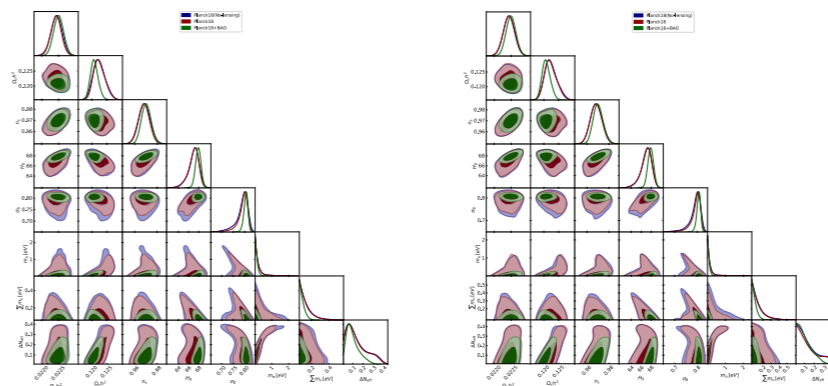
$\Delta N_{\text{eff}}$  tracking the number density

FD, Hajkarim, Yun, **JHEP 10 (2021)**

FD, Hajkarim, Yun, **Phys.Rev.Lett. 128 (2022)**

FD, Di Valentino, Giarè, Hajkarim, Melchiorri, Mena, Renzi, Yun, **JCAP 09 (2022)**

Axion cosmological mass bound

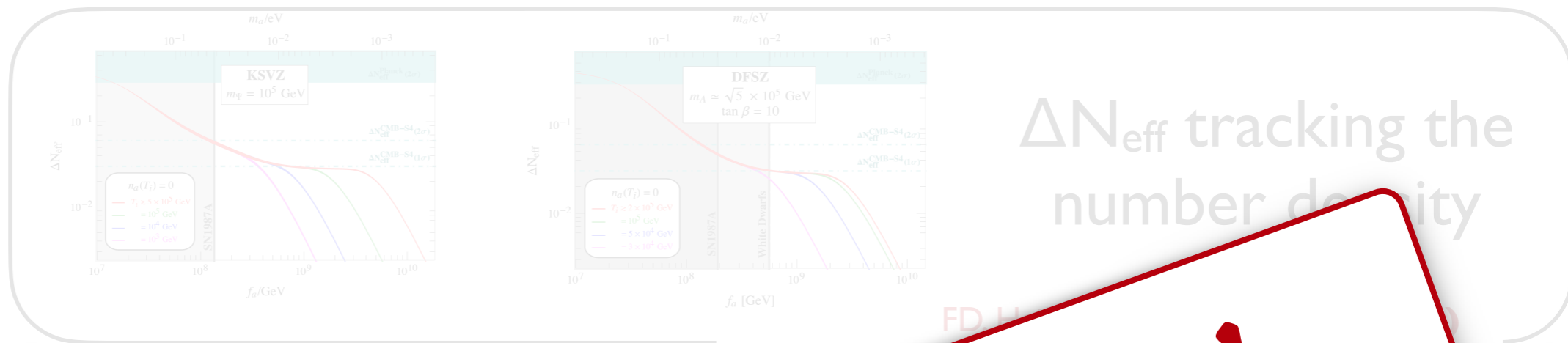


Importance of a phase space analysis

FD, Hajkarim, Lenoci, **JCAP 03 (2024)**

FD, Lenoci, **in preparation**

# The Way Back to the Phase Space

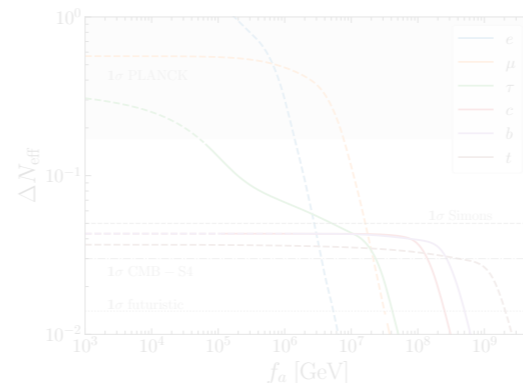
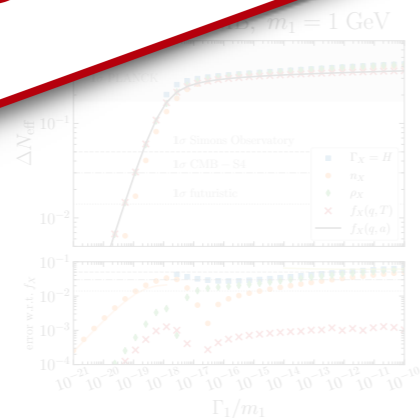


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Axion cosmological mass bounds

**THANK YOU!**



Importance of a phase space analysis

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