Exclusion bounds for neutral gauge bosons

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based on [2402.14786] with Z. Trócsányi



CATCH22+2 Dublin, May 1-5, 2024

Outline

• Theory considerations

• Light Z' bosons

• Heavy Z' bosons

Conclusions

Theory considerations for neutral gauge bosons

Motivation for extra neutral gauge bosons (Z')

- Z' appears after breaking a U(1) gauge group (or higher)
- Fifth fundamental interaction?
- Breaking a larger gauge group with a scalar VeV \rightarrow the unbroken subgroup has U(1)-s (e.g.: GUT, SUSY, string)
- Z' can connect to a secluded sector in the SM
- A discovery would have a lot of consequences: extended scalar $[\chi]$ (make Z' massive) and extended fermion sectors $[3 \times \nu_R]$ (cancel gauge anomalies)

Minimal extension of the SM

- SM gauge group + $U(1)_z$
- Covariant derivative is modified:

$$D_{\mu}^{U(1)} = -i \left(y \, \mathbf{z} \right) \begin{pmatrix} g_{y} & -g_{z} \, \boldsymbol{\eta} \\ 0 & g_{z} \end{pmatrix} \begin{pmatrix} B_{\mu} \\ B_{\mu}' \end{pmatrix}$$

z charges are defined

at $\eta(\mu_0) = 0$

- $\eta \propto \text{kinetic mixing } (F'_{\mu\nu}F^{\mu\nu})$
- Rotate to mass eigenstates:

$$\begin{pmatrix} B_{\mu} \\ W_{\mu}^{3} \\ B_{\mu}' \end{pmatrix} = \begin{pmatrix} c_{W} & -s_{W} & 0 \\ s_{W} & c_{W} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{Z} & -s_{Z} \\ 0 & s_{Z} & c_{Z} \end{pmatrix} \begin{pmatrix} A_{\mu} \\ Z_{\mu} \\ Z_{\mu}' \end{pmatrix}$$

z charge remarks

- 2 free z charges (due to anomaly cancellation + Yukawa masses +normalization)
- z_N ~ neutrino mass generation mechanism: tree level Majorana mass term is allowed if $z_\chi + 2z_N = 0$ (~ $\chi \, \bar{\nu}_R \nu_R^c$)
- Z_{ϕ} appears only in the combination: $Z_{\phi} \frac{\eta}{2} \rightarrow$ use this one

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- In the branching ratios of Z' the combination appears:

$$\mathbf{Z} = \frac{z_{\phi} - \eta/2}{z_N}$$
 e.g.: $\text{Br}(Z' \to \text{inv.}) = \frac{(3+n_N)}{(3+n_N)+2(1-2c_W^2 \mathbf{Z})^2}$ for $M_{Z'} < m_{\mu}$

• B-L model: $z_{\phi} = 0 \rightarrow \mathbf{Z} = \mathbf{0}$

Other free parameters:

- $M_{Z'}$ (or rather $\xi = M_{Z'}/M_Z$ to treat diff. mass scales)
- Either the mixing angle s_z or the new gauge coupling g_z :

$$-s_{z}c_{z}\frac{(1-\xi^{2})}{\rho} = \frac{2}{\sqrt{g_{Y}^{2}+g_{L}^{2}}} g_{z}\left(z_{\phi} - \frac{\eta}{2}\right)$$

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This is BSM only!

- From global fits one has: $\rho = 1.00038 \pm 0.00020$
- The tree level model prediction is:

$$\rho = \frac{M_W^2}{M_Z^2 c_W^2} = 1 + (\xi^2 - 1) s_z^2 \qquad \begin{bmatrix} 2305.1193 \\ 2306.01836 \end{bmatrix}$$



Constraints

• Indirect: EW. Precision Observables $\rightarrow \rho$ -parameter

$$|s_z| < 4.5 \cdot 10^{-3} \text{ or } |z_N g_z| < \frac{1.7 \cdot 10^{-3}}{|Z|} @ 95\% \text{ C.L.}$$

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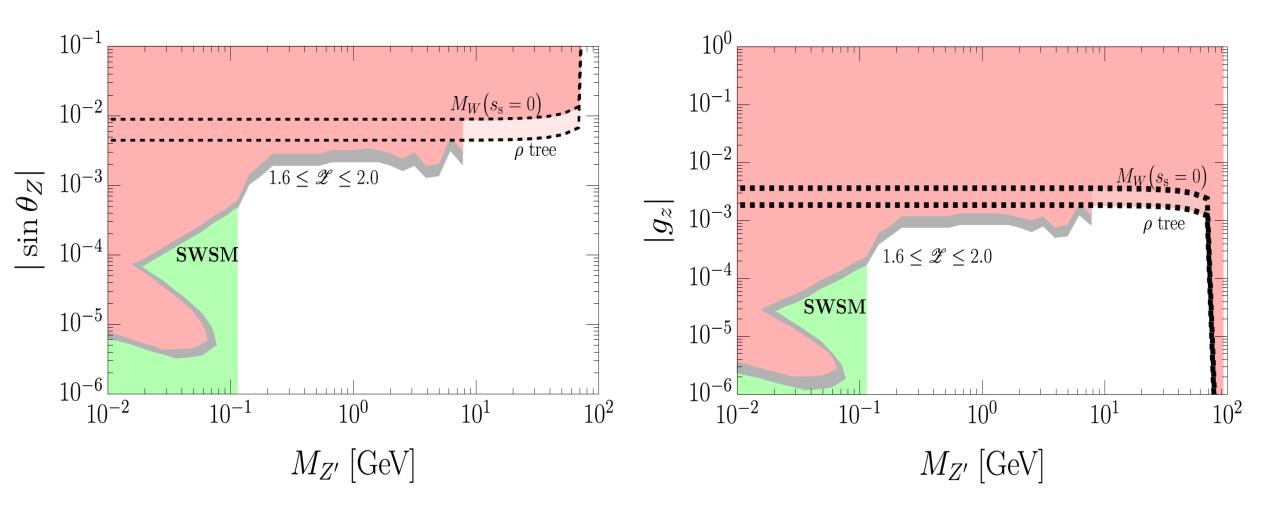
- **Direct:** searches for $Z' \to \text{inv.}$: **BaBar** [1702.03327], **NA64** [1906.00176]
- Direct: searches for $Z' \rightarrow e^+e^-$: FASER [2308.05587]

Far from a complete list

• These experiments search for dark photons (A'):

$$\mathcal{L} \supset \frac{1}{2} m_{A'}^2 A'^2 + (e\epsilon) A'_{\mu} \sum_{f} Q_f \bar{f} \gamma^{\mu} f \to \epsilon(M_{Z'}, |s_z|, \mathbf{Z}) \text{ or } \epsilon(M_{Z'}, |z_N g_z|, \mathbf{Z})$$
[1801.04847]

SWSM: $z_{\phi} = 1$ and $z_N = 1/2$



Uncertainty due to sterile neutrinos + running of η : thickness of gray line

Heavy Z' bosons

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$$|s_z| < 2.5 \cdot 10^{-3} \left[\frac{1 \text{ TeV}}{M_{Z'}} \right] \text{ or } |z_N g_z| < \frac{0.11}{|\mathbf{Z}|} \left[\frac{M_{Z'}}{1 \text{ TeV}} \right] @ 95\% \text{ C.L.}$$

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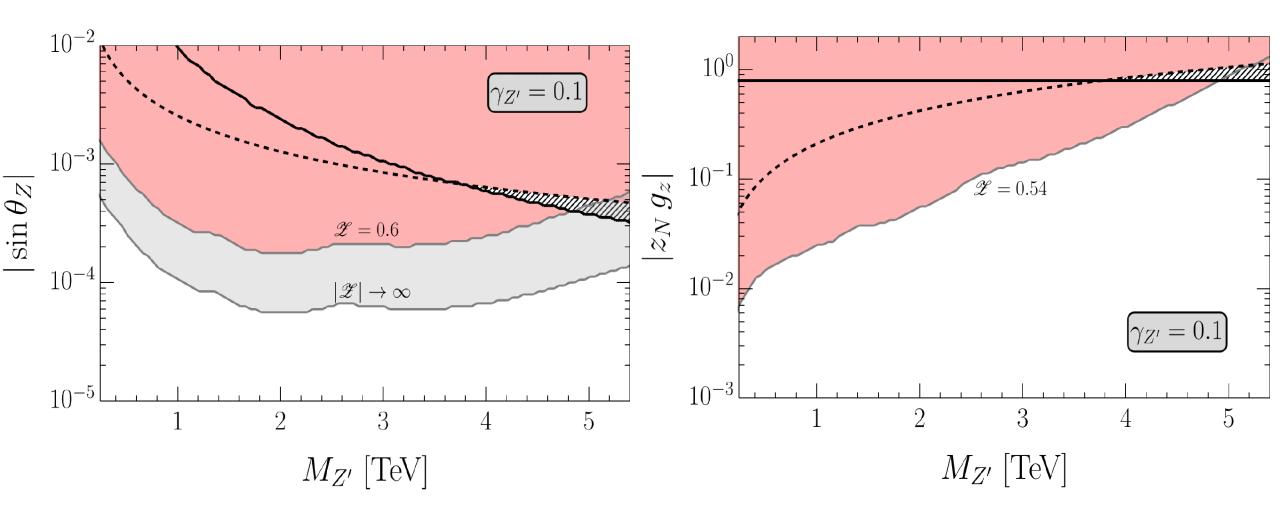
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• **Direct:** searches in the LHC (**ATLAS** [1903.06248] and **CMS** [2103.02708]) in dilepton final states:

$$\sigma = \frac{4\pi^2}{3s} \frac{\Gamma_{Z'}}{M_{Z'}} \operatorname{Br}(Z' \to \ell^+ \ell^-) \sum_{q} \operatorname{Br}(Z' \to \overline{q}q) w_q(s, M_{Z'})$$

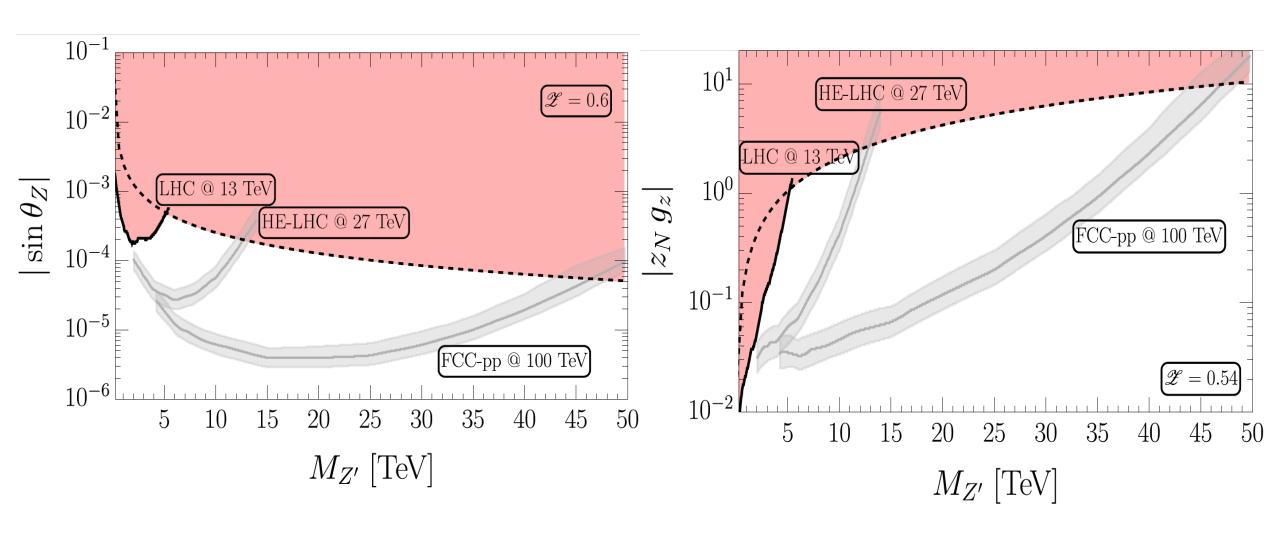
• experiments use assumptions for $\gamma_{Z'} = \Gamma_{Z'}/M_{Z'}$

Sample exclusion bounds



There is a **Z** value corresponding to a *loosest* bound!

Projections for future colliders



...using detector simulations for the HE-LHC and FCC-hh

Interesting process

• For very large $M_{Z'}$ the decay

$$Z' \rightarrow Z + W^+ + W^-$$

might dominate over the leptonic decay of the Z'!

• The ratio of the branching fractions (also the cross sections' in the NWA):

$$\frac{\sigma(pp \to Z' + X \to ZWW + X)}{\sigma(pp \to Z' + X \to \ell\ell + X)} = \frac{\text{Br}(Z' \to Z + W^+ + W^-)}{\text{Br}(Z' \to \ell^+ + \ell^-)} = 0.4 \left(\frac{Z^2}{2 - 6Z + 5Z^2}\right) \left[\frac{M_{Z'}}{10 \text{ TeV}}\right]^2$$

Potentially relevant for FCC-hh

Conclusions

- Useful parametrization: different U(1) extensions can be investigated on the same footing
- ρ can be used to quickly assess the constraints from EWPO
- $(M_{Z'}, s_z)$ bounds from ρ are model independent!
- Direct searches provide more stringent constraints
- In the $p + p \rightarrow Z' + X \rightarrow \ell^+\ell^- + X$ searches there is a least severe bound: model independent constraints



z charge assignment

field	$SU(3)_{c}$	$SU(2)_{\rm L}$	y	z
Q_{L}	3	2	$\frac{1}{6}$	$z_q = \frac{1}{3}(z_\phi - z_N)$
U_{R}	3	1	$\frac{2}{3}$	$z_u = \frac{1}{3}(4z_\phi - z_N)$
$D_{ m R}$	3	1	$-\frac{1}{3}$	$z_d = -\frac{1}{3}(2z_\phi + z_N)$
$\ell_{ m L}$	1	2	$-\frac{1}{2}$	$z_{\ell} = z_N - z_{\phi}$
$N_{ m R}$	1	1	0	z_N
e_{R}	1	1	-1	$z_e = z_N - 2z_\phi$
ϕ	1	2	$\frac{1}{2}$	z_{ϕ}
χ	1	1	0	$z_{\chi} = -1$

- Cancel anomalies
 - +Yukawa mass terms
 - = Fix all, but two z charges
- χ : new singlet scalar
- *N*: right handed (sterile) neutrinos

• Choose z_N and z_ϕ to be free

Branching ratios

• Light Z' with $M_{Z'} < m_{\mu}$:

Br(Z'
$$\rightarrow$$
 inv.) = $\frac{(3 + n_N)}{(3 + n_N) + 2(1 - 2 c_W^2 \mathbf{Z})^2}$

• Heavy **Z**':

$$Br(Z' \to \ell^+ \ell^-) = \frac{2 - 6 \mathbf{Z} + 5 \mathbf{Z}^2}{16 - 32 \mathbf{Z} + \mathbf{Z}^2 (41 + C_{w,s} \xi^2)}$$

The param $C_{w,s} \simeq 1.4 \cdot 10^{-4}$ is from three-body decays

V-A couplings

• Most of the Z' phenomenology depends on:

$$\mathcal{L}_{NC}^{(Z')} = -\frac{e}{2s_w c_w} Z'_{\mu} \sum_{f} \bar{f} \gamma^{\mu} (v_{Z',f} - a_{Z',f} \gamma^5) f$$

- $v_{Z',f}$ and $a_{Z',f}$ are pretty simple for $\xi \gg 1$ and $\xi \ll 1$
- For instance: $a_{Z',f}$ is negligible for $\xi \ll 1$
- but $a_{Z',f} = \pm \frac{1}{2} s_z \xi^2$ for $\xi \gg 1$
- The vector cps. depend on $(s_z, \xi, \mathbf{Z}) \leftrightarrow (z_N g_z, \mathbf{Z})$

Matching to dark photons

- Relate similar processes (which is searched for) [1801.04847]
- For NA64 and BaBar one has

$$e\epsilon = \frac{|v_{Z',\ell}|}{2s_W c_W} \sqrt{\text{Br}(Z' \to \text{inv.})}$$

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• For the FASER, the A' is also required to decay in the detector \rightarrow solve the matching equations

$$\Gamma(\pi^0 \to A' + \gamma) \operatorname{Br}(A' \to e^+ e^-) = \Gamma(\pi^0 \to Z' + \gamma) \operatorname{Br}(Z' \to e^+ e^-)$$

$$m_{A'}\Gamma_{A'}=M_{Z'}\Gamma_{Z'}$$

Matching to dark photons

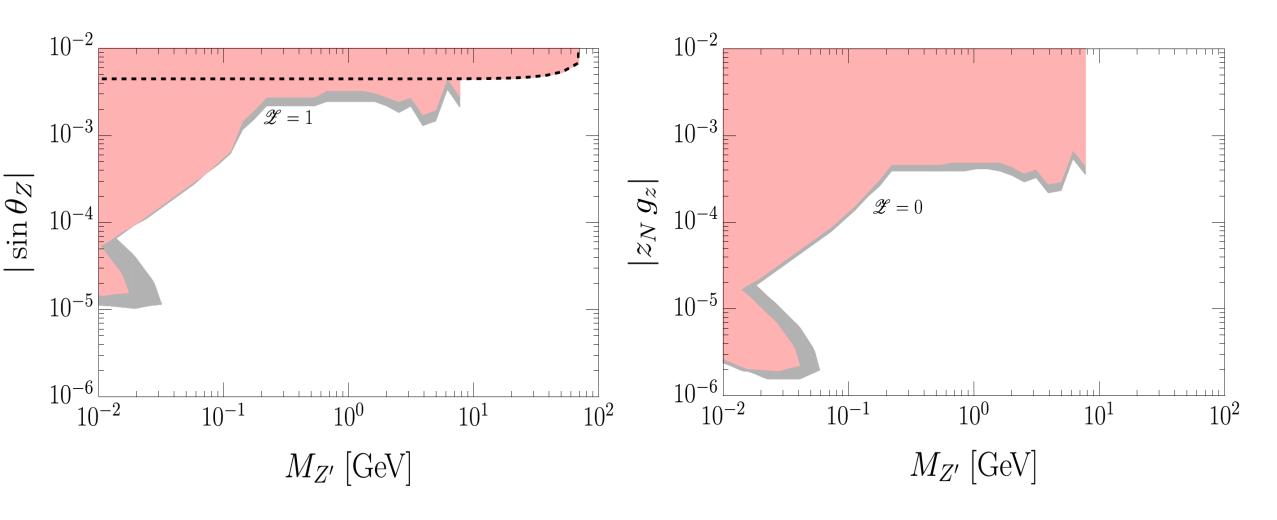
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• For the FASER, the A' is also required to decay in the detector \rightarrow solve the matching equations \rightarrow for $M_{Z'} \ll m_{\pi}^0$:

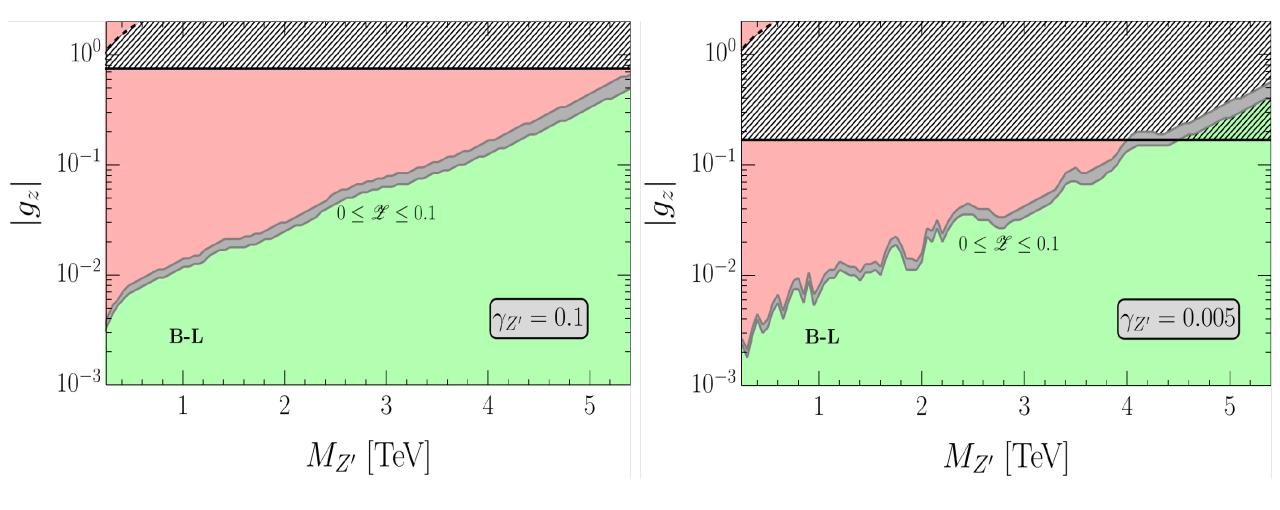
$$e\epsilon = \frac{|v_{Z',\ell}|}{2 s_W c_W} \sqrt{\text{Br}(Z' \to e^+ e^-)} \text{ and } M_{Z'} = \text{Br}(Z' \to e^+ e^-) M_{A'}$$

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