

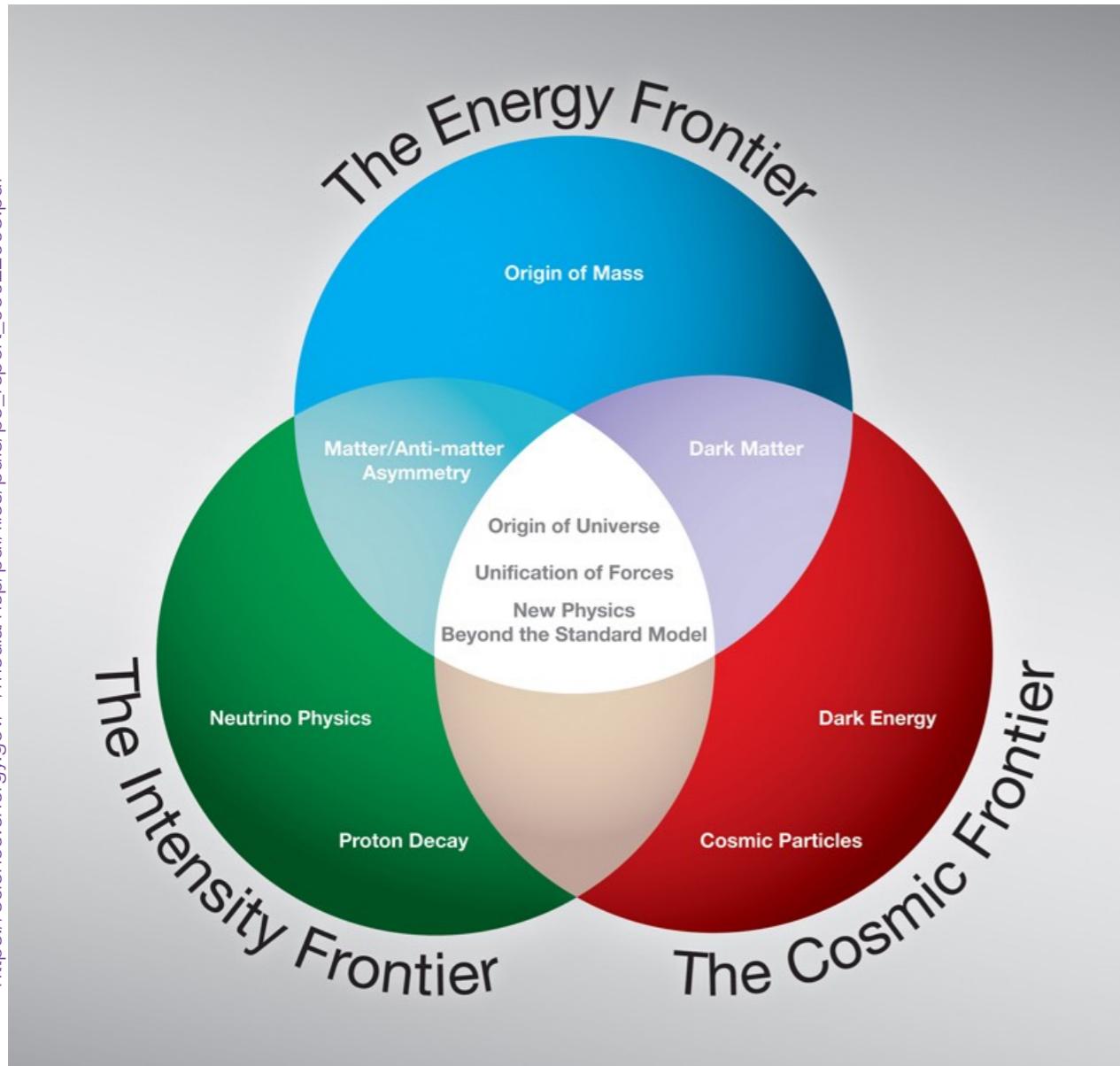


A review of Axion physics

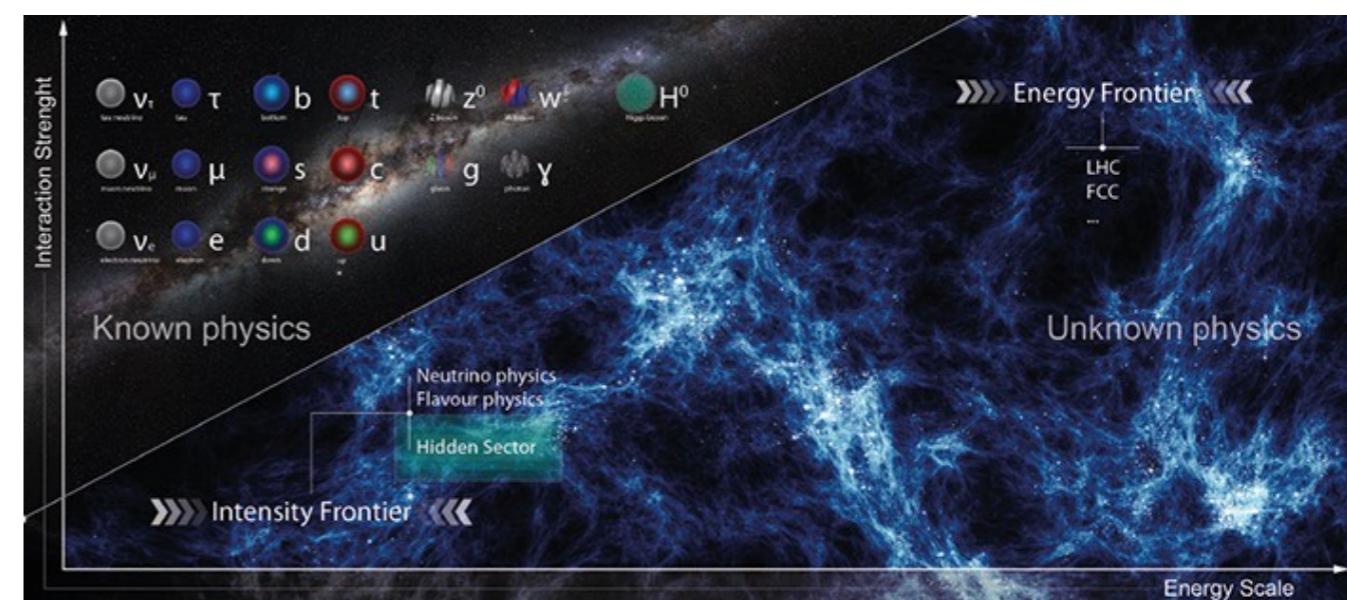
Seokhoon Yun

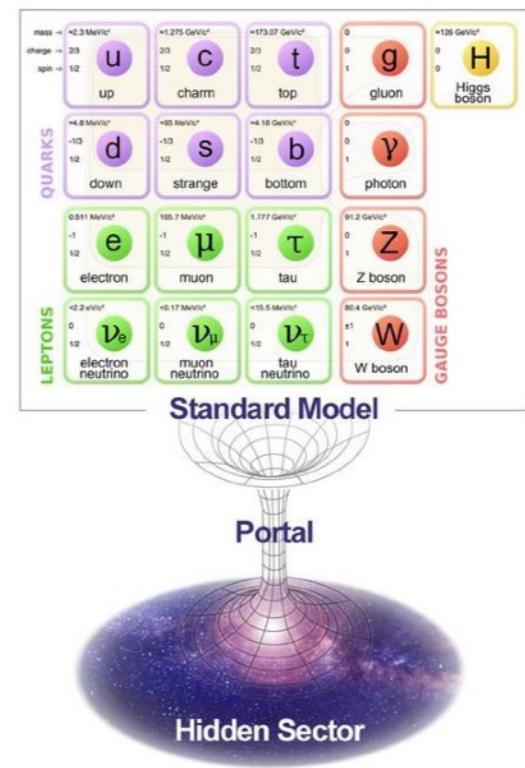


Frontiers of particle physics



- **Energy frontier**
 - high-energy high-collision rate accelerators (e.g., LHC)
- **Intensity frontier**
 - medium or low-energy ultra-high-collision rate
- **Cosmic frontier**
 - universe & astroparticle





Portal physics

01

Higgs portal

spin-0 scalar

02

Axion portal

spin-0 pseudoscalar

03

Vector portal

spin-1 vector

04

Neutrino portal

spin-1/2 fermion

$$|S|^2 H^\dagger H$$

$$(a/f_a) F_{\mu\nu} \tilde{F}^{\mu\nu}$$

$$F_{\mu\nu} X^{\mu\nu}$$

$$\bar{L} H N$$

$$\left(\partial_\mu a / 2f_a\right) \bar{\psi} \gamma^\mu \gamma^5 \psi$$

Outline

01 Strong CP problem: the motivation of axions

02 QCD axion (KSVZ & DFSZ) | Axion-like particles

03 Effective field theory framework on axions

e.g., anomalous coupling to the electromagnetic field ($g_{a\gamma}$)

04 Current statue of searches

- *Experimental*
- *Astrophysical*
- *Cosmological*

Low energy QCD

- The QCD Lagrangian around the confinement scale

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{4}G^{a\mu\nu}G^a_{\mu\nu} + i\bar{q}\gamma^\mu D_\mu q - \frac{g_s^2}{32\pi^2}\theta_{\text{QCD}}G^{a\mu\nu}\tilde{G}^a_{\mu\nu} - (\bar{q}_L M_q q_R + \text{h.c.})$$

$$q_{L,R} = P_{L,R} (u, d, s)^T, \quad M_q = \text{diag} [m_u, m_d, m_s]$$

- CP violation in the strong interactions

$$\bar{\theta} = \theta_{\text{QCD}} + \arg \text{Det } M_q$$

Naturally $\mathcal{O}(1)$?

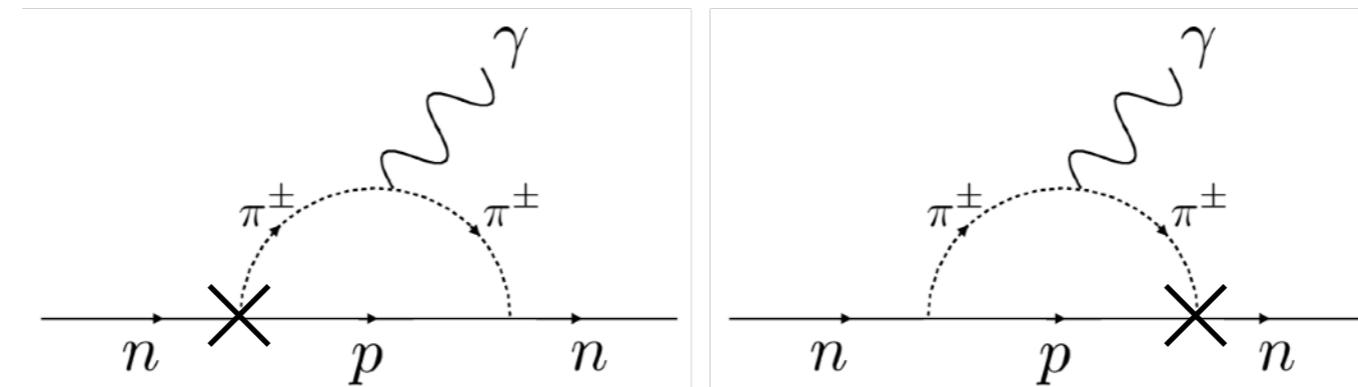
Strong CP problem @ quantum level

- P & T violating pion-nucleon coupling for $\bar{\theta} \ll 1$

$$\mathcal{L}_{\bar{\theta}\pi N} = -\bar{\theta} \frac{c_+ \mu}{f_\pi} \pi^a N \sigma^a N, \quad \text{where } \mu = \frac{m_u m_d}{m_u + m_d} \text{ & } c_+ \approx 1.7$$

baryon mass splitting

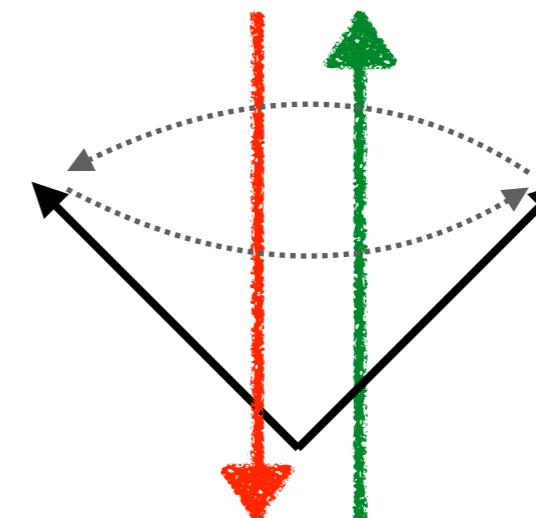
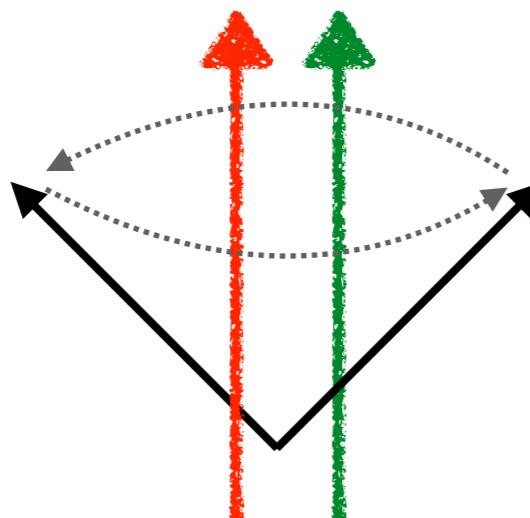
- Electric dipole moment of the neutron



$$d_n F_{\mu\nu} \bar{n} \gamma^{\mu\nu} i \gamma^5 n \quad \text{with} \quad d_n \sim 3 \times 10^{-16} \bar{\theta} e \text{ cm}$$

Experimental bound on d_n

- Conceptual way to measure d_n via a (Larmor) precession in \vec{E} & \vec{B}

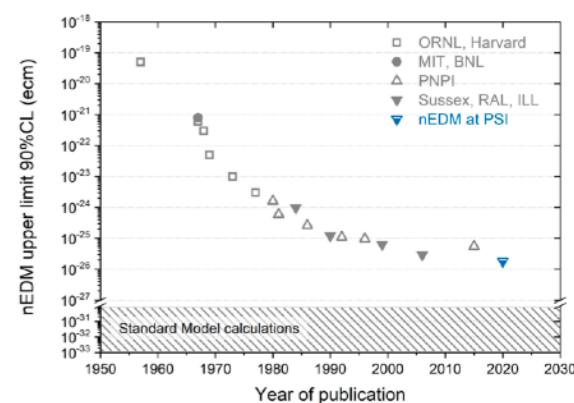


$$\nu = 2\mu B + 2dE$$

$$\nu = 2\mu B - 2dE$$

- Bound on electric dipole moment of the neutron

$$d_n \sim 3 \times 10^{-16} \bar{\theta} e \text{ cm} \leq 10^{-26} e \text{ cm} \quad \text{☞} \quad \bar{\theta} < 10^{-10}$$



Axion

- Solution for the strong CP problem
- Global Peccei-Quinn symmetry anomalous to the SM - **PQ fermions**

$$\Psi \rightarrow e^{i\alpha\gamma^5/2}\Psi \quad \rightleftharpoons \quad \frac{g_s^2}{32\pi^2} \alpha G_{\mu\nu} \tilde{G}^{\mu\nu}$$

- Spontaneous breaking of the PQ symmetry - **PQ scalar field**

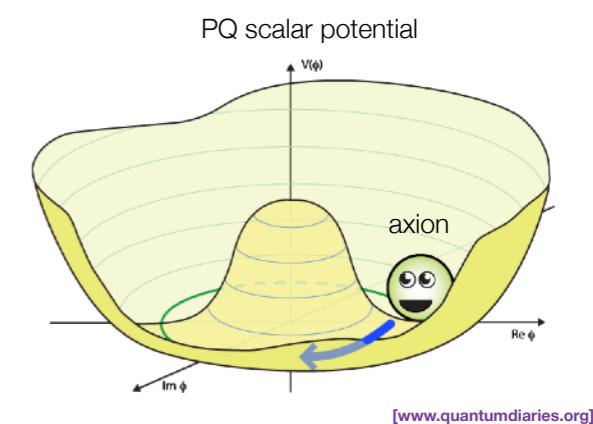
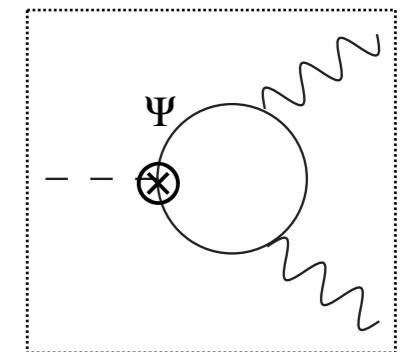
$$V_\varphi = \lambda_\varphi \left(|\varphi|^2 - f_a^2/2 \right) \quad \rightleftharpoons \quad \varphi \rightarrow \left(f_a/\sqrt{2} \right) e^{ia/f_a}$$

- NG boson a , the so-called ‘axion’

$$\mathcal{L}_{\text{PQ}} \supset -y_\Psi \varphi \bar{\Psi}_L \Psi_R + \text{h.c.} \quad \rightleftharpoons \quad \bar{\theta} \rightarrow \bar{\theta} + a/f_a$$

- With the minimum of $V_{\bar{\theta}}$ at $\bar{\theta} + a/f_a = 0$,

“The strong CP problem is solved dynamically”



Axion models & properties



[J.E. Kim, 1979] [M.A. Shifman, A.I. Vainshtein and V.I. Zakharov, 1980]

Fermion extension of SM
e.g., heavy PQ colored fermion

PQ scalar

☞ axion (axial degree)



[M. Dine, W. Fischler and M. Srednicki, 1980] [A.R. Zhitnitsky, 1980]

Scalar extension of SM
e.g., multi Higgs doublets

- Goldstone boson nature ☞ axion mass and coupling $\propto f_a^{-1}$

$$\textbf{Axion mass} \simeq 5.7 \text{ meV} \left(\frac{10^9 \text{ GeV}}{f_a} \right)$$

$$\textbf{Coupling to SM} \propto \frac{v_{\text{EW}}}{f_a} \sim \frac{100 \text{ GeV}}{f_a}$$

- Axion mass independent of f_a (but still light), if there is an explicit breaking of PQ symmetry

“Axion-like particles”

Effective axion couplings

- In the non-linearly realized PQ-symmetry basis
i.e., invariant under the shift $a \rightarrow a + \text{constant}$

$$\frac{\partial_\mu a}{2f_a} \left(\sum_\psi c_\psi \bar{\psi} \gamma^\mu \psi + \sum_\phi c_\phi \phi^\dagger i \overleftrightarrow{\partial}_\mu \phi \right) + \frac{a}{f_a} \sum_V c_V \frac{g_V^2}{32\pi^2} F_V^{\mu\nu} \tilde{F}_{V\mu\nu}$$

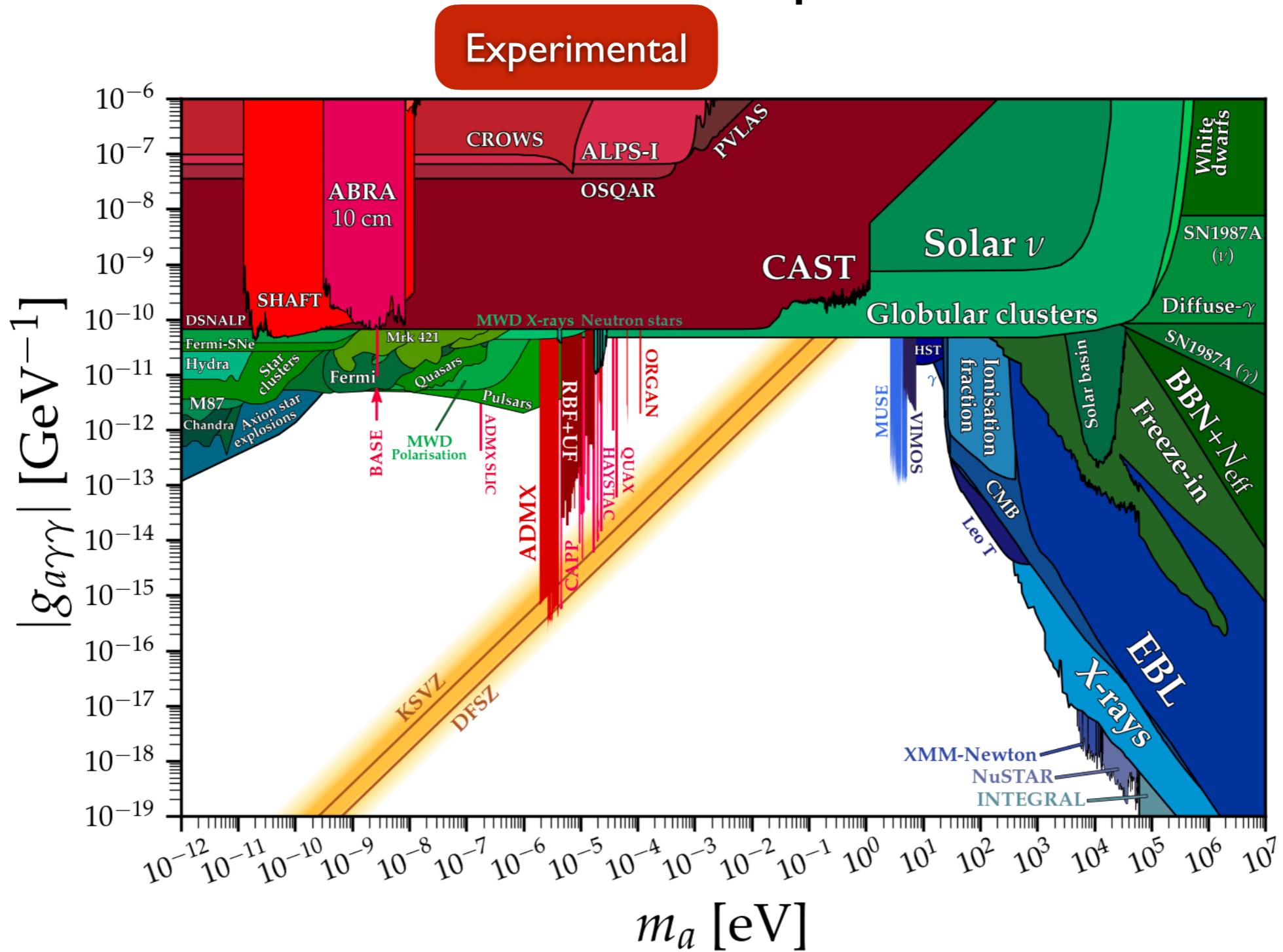
$$\equiv J_{\text{PQ}}^\mu$$

$$\boxed{\frac{g_{a\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} = - g_{a\gamma} a \vec{E} \cdot \vec{B}}$$

- Can be converted into the linearly-realized basis through the axionic rotation,
which relies on the conservation of PQ-current (i.e., $\partial_\mu J_{\text{PQ}}^\mu$)

☞ Effective operators with non-vanishing PQ-charge contribute to axion physics

Current status on $\frac{g_{a\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu}$

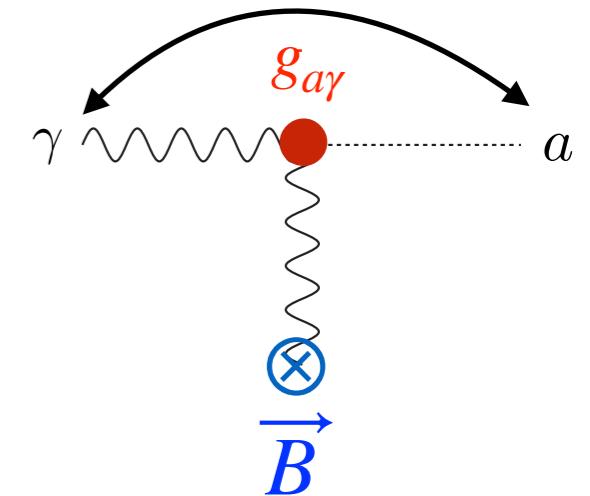


Experimental searches

- Axion-photon conversion in a magnetic background
- Light-Shining-Through-Walls
- Photon polarizations (Dichroism & Birefringence)

Axion-Photon oscillation

“Primakoff” process



- Equation of motion in the presence of a background (transverse) magnetic field \vec{B}

$$\mathcal{L}_{a \& \gamma} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{1}{2}\partial_\mu a \partial^\mu a - \frac{1}{2}m_a^2 a^2 - \frac{g_{a\gamma\gamma}}{4}a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

————— = $g_{a\gamma} \vec{E} \cdot \vec{B} a$

$$\partial_\mu F^{\mu\nu} = g_{a\gamma} \tilde{F}^{\mu\nu} \partial_\mu a$$

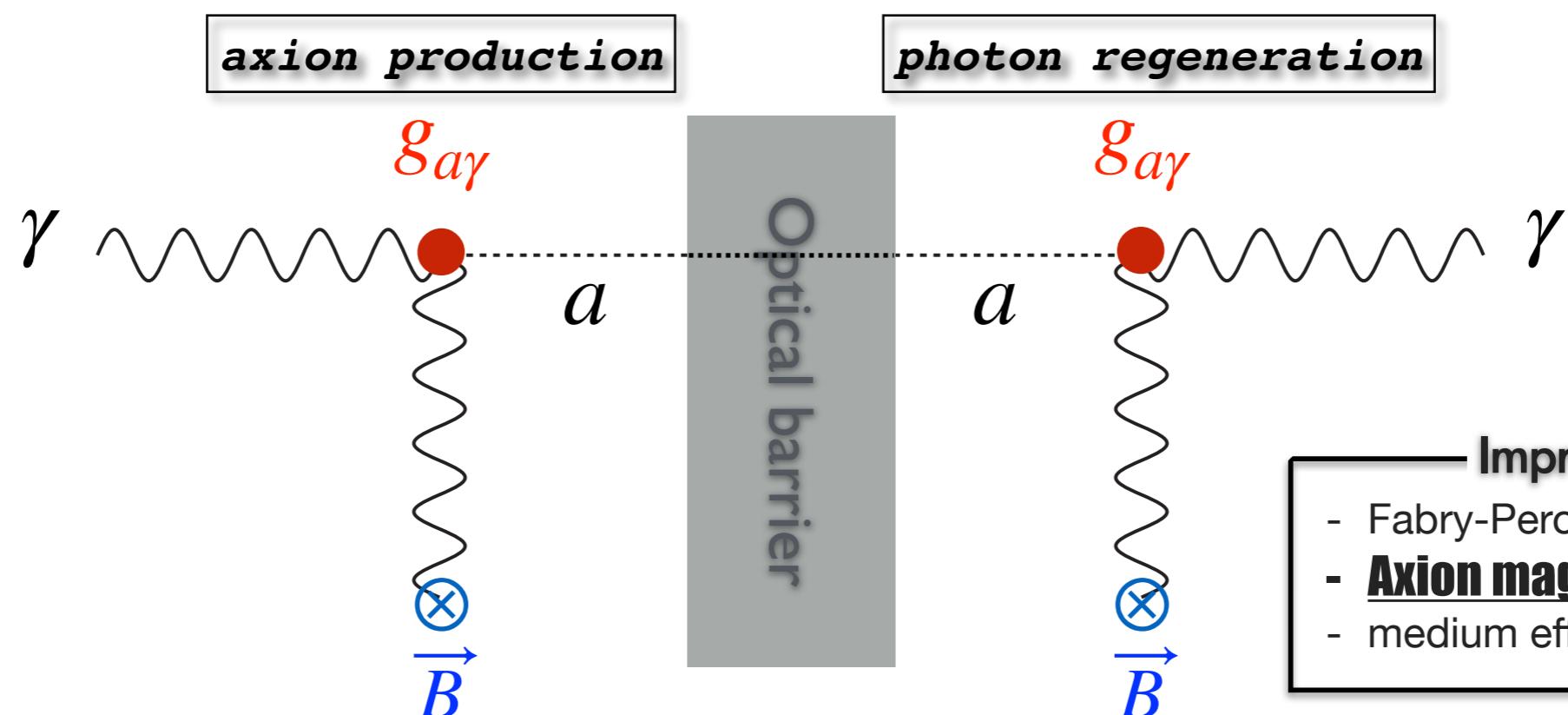
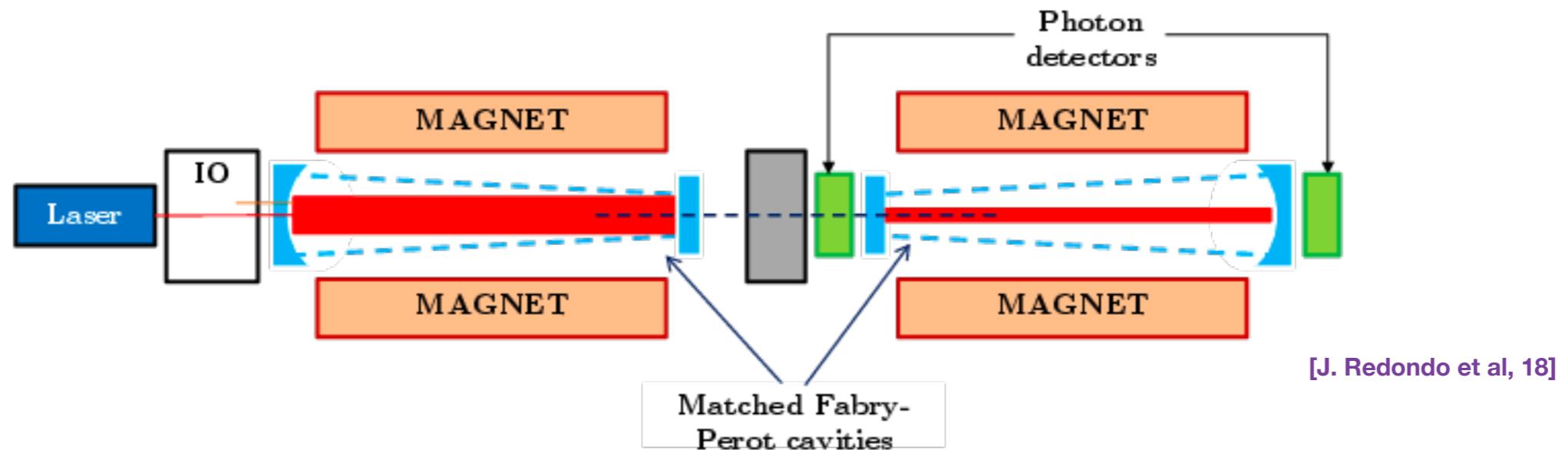
$$\partial_\mu \partial^\mu a = -m_a^2 a - \frac{1}{4}g_{a\gamma} F_{\mu\nu} \tilde{F}^{\mu\nu}$$



$$\partial_\mu \partial^\mu \vec{A} = g_{a\gamma} \vec{B} \partial_t a$$

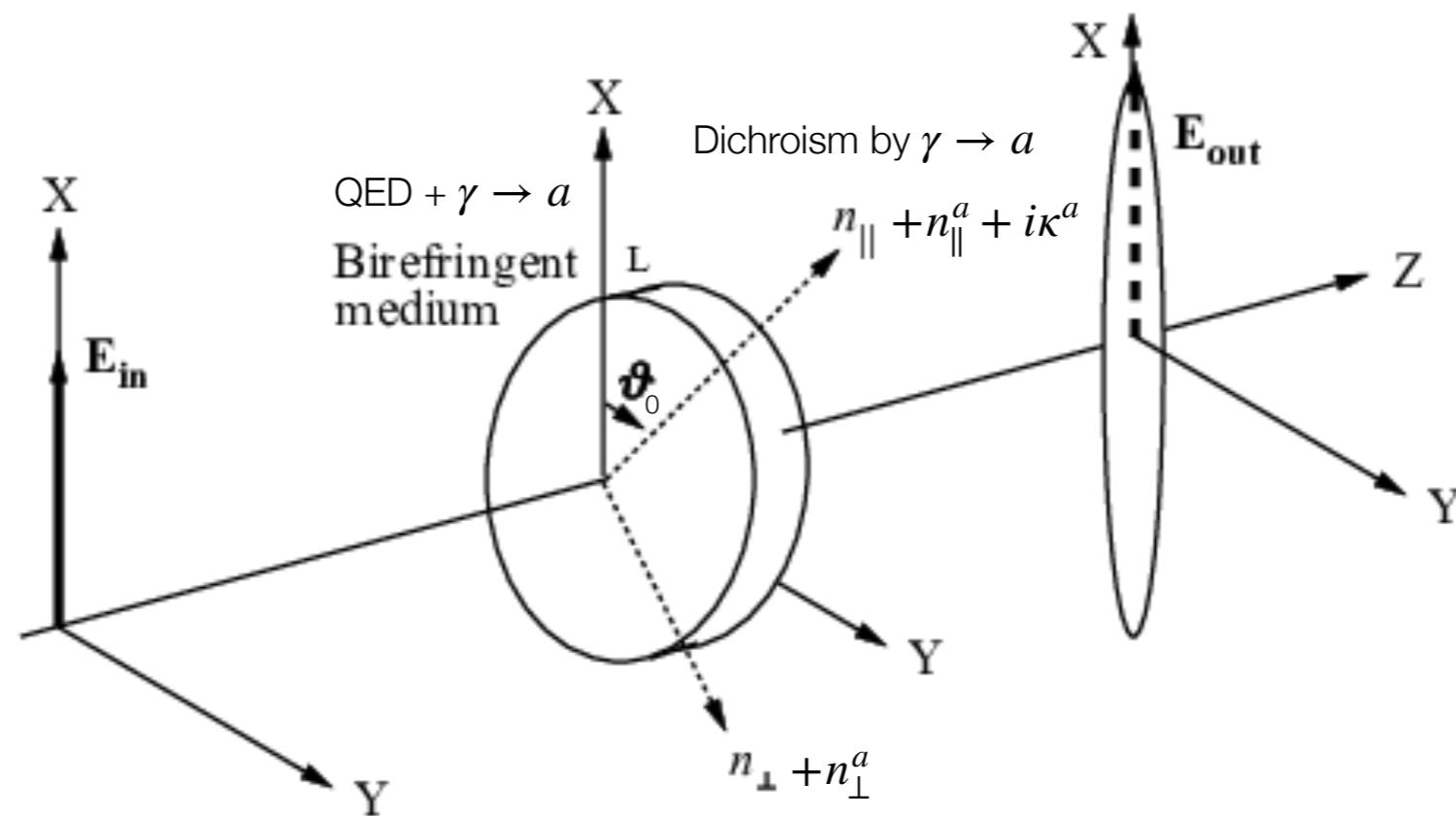
$$\partial_\mu \partial^\mu a = -m_a^2 a - g_{a\gamma} \vec{B} \cdot \partial_t \vec{A}$$

Light-Shining-Through-Walls



Photon polarization

[PVRAS, arXiv:2005.12913]

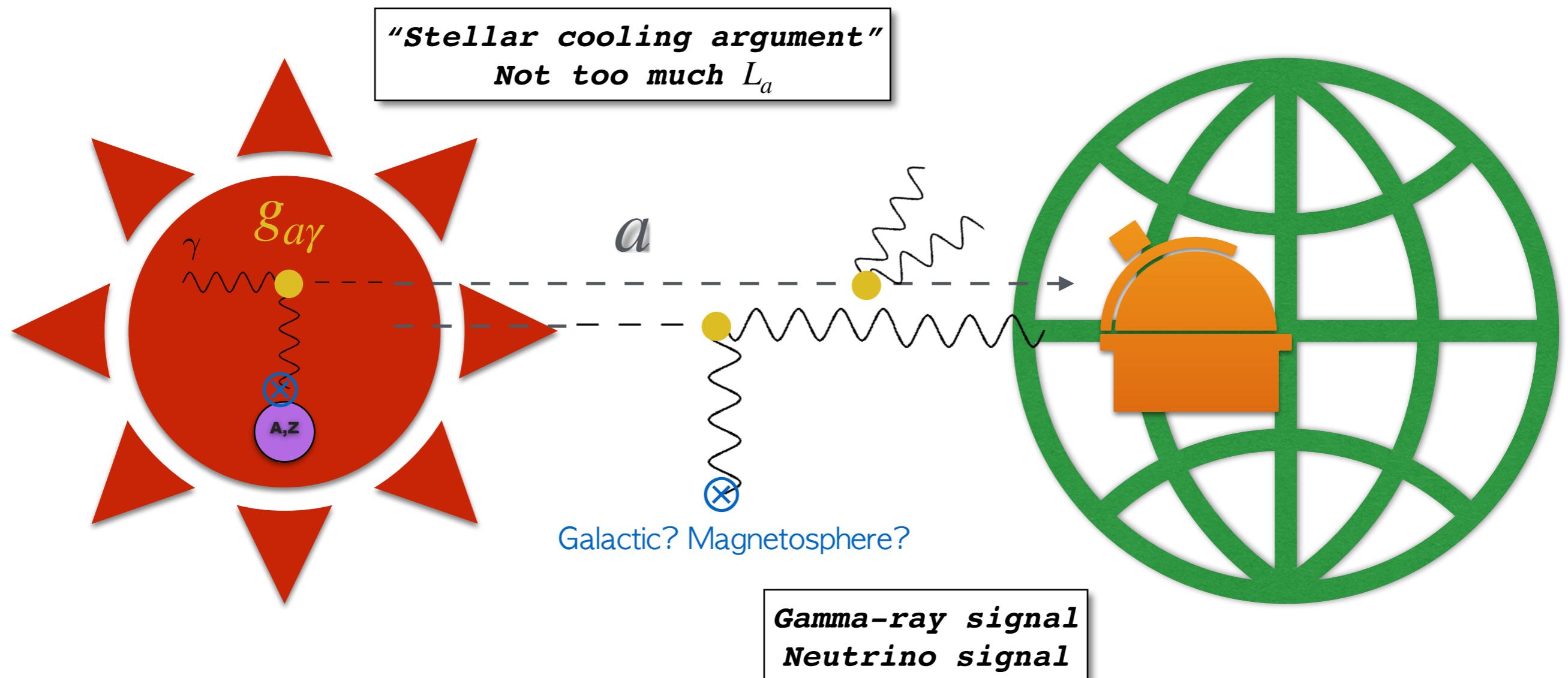


1. Initial linearly polarized light to the \hat{x} -direction & propagating to the \hat{z} -direction
2. Passing through a magnetic field background
 - **Birefringence: ellipticity of the γ -polarization (phase retardation by QED & $\gamma \rightarrow a$)**
 - **Dichroism: rotation of the γ -polarization (depletion by $\gamma \rightarrow a$)**

Astrophysical searches

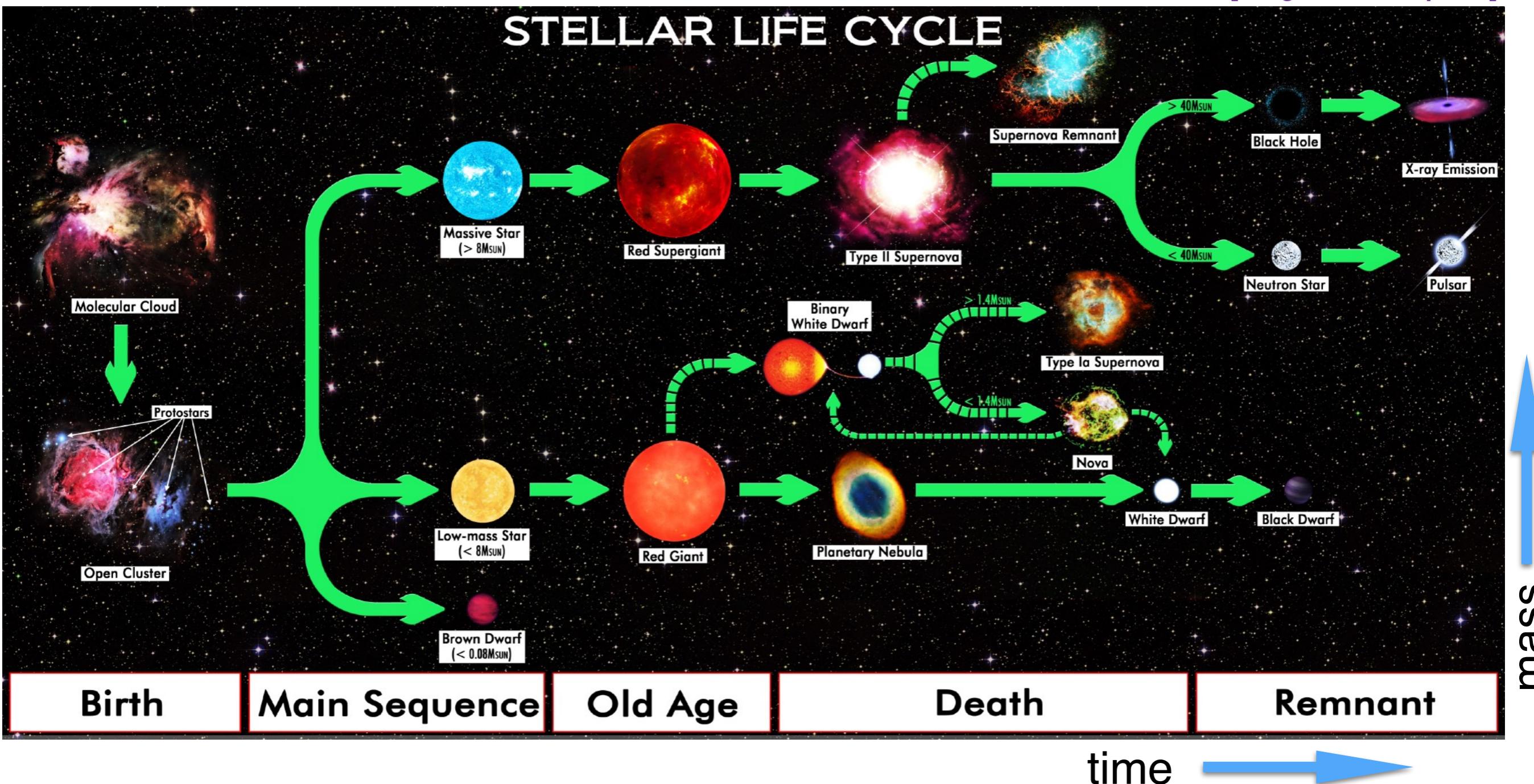
- Axion production: Primakoff process
- Implications:
 - 1) Stellar cooling argument
 - 2) γ -ray signal from axion-photon conversion in a magnetic field
 - 3) γ -ray signal from axion decay into 2γ

Astrophysical searching scheme



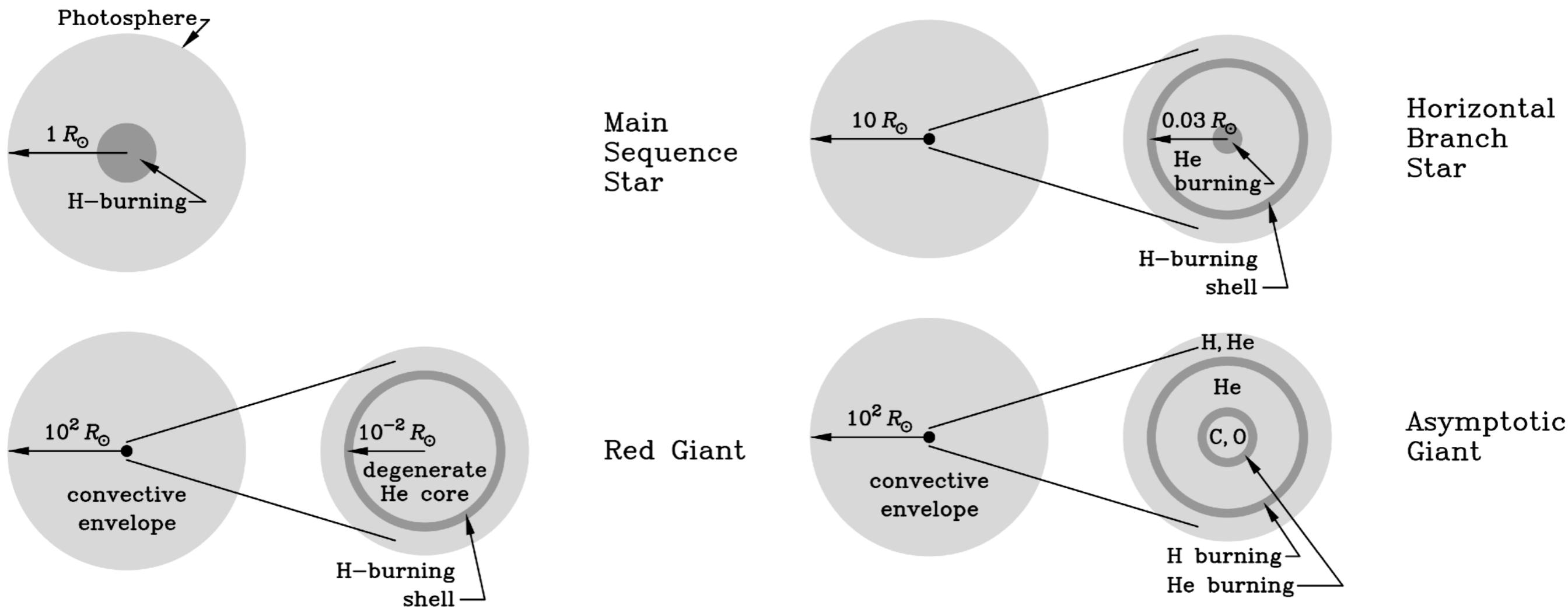
Stellar evolution

[image from Wikipedia]



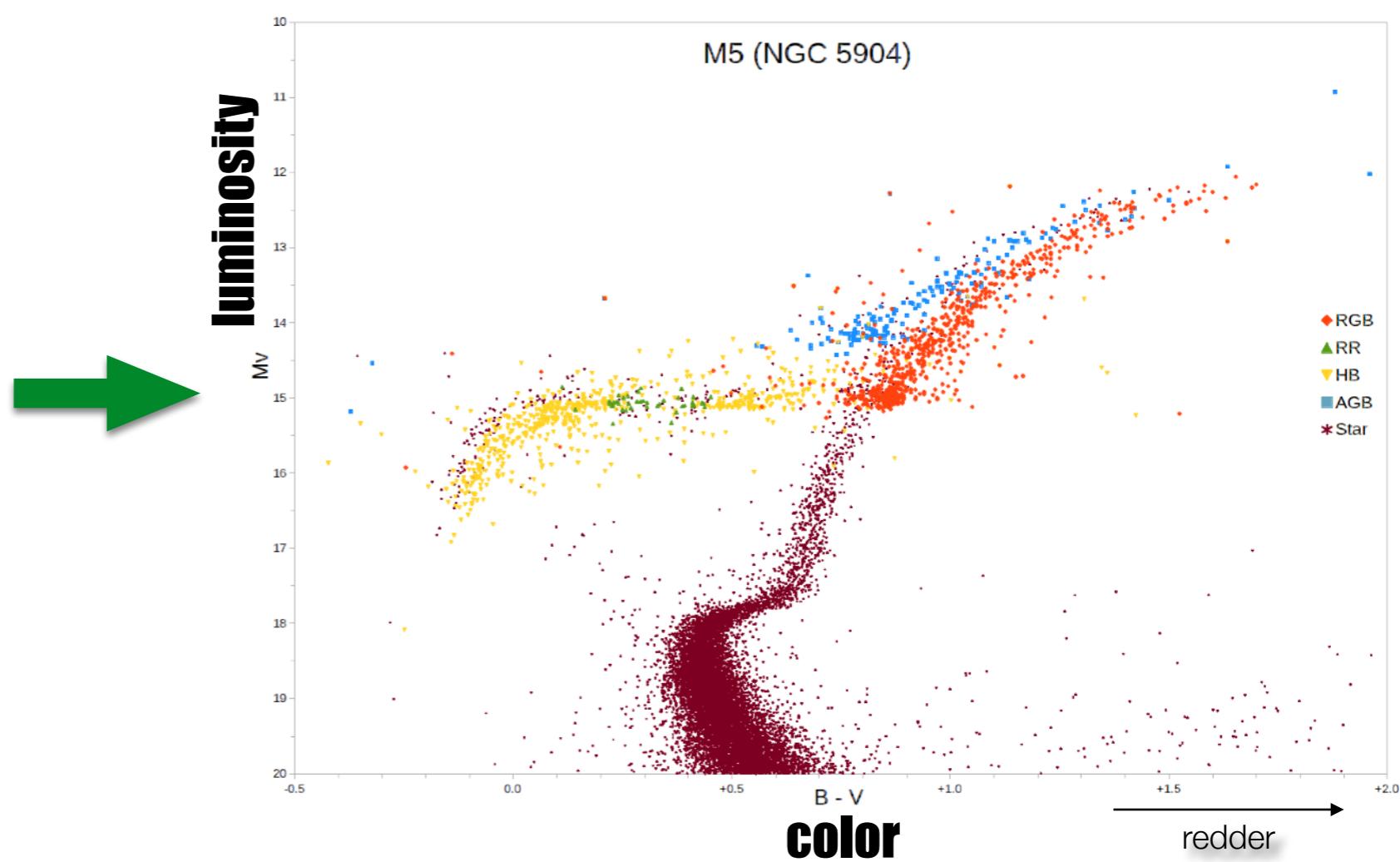
Stellar evolution for low ρ_c

[Raffelt, “Stars as laboratories for Fundamental Physics”]

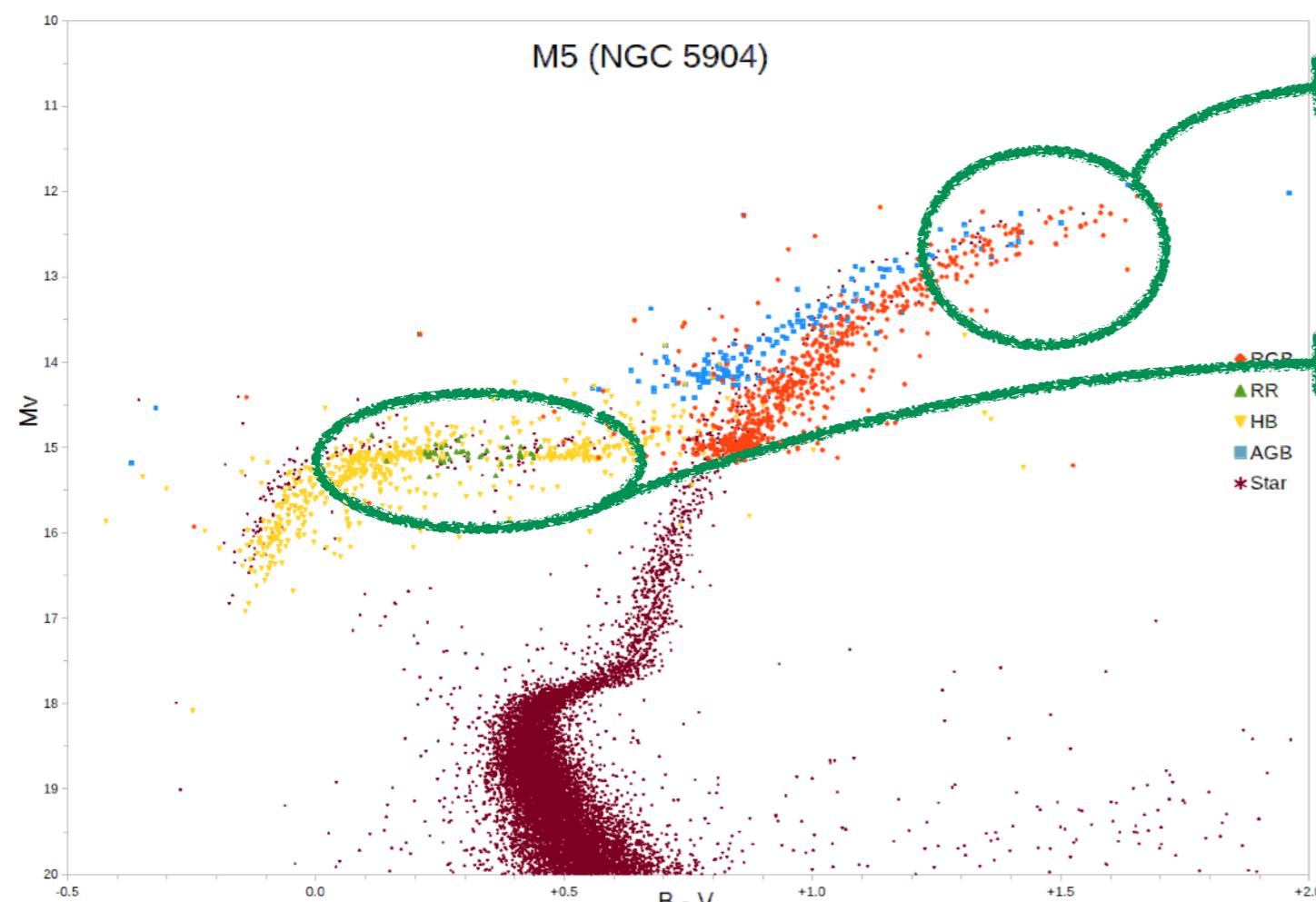


Globular cluster

- Very clean laboratories for studies of stellar evolution
- Coeval and almost equal chemical compositions
- Only different initial mass



Stellar cooling argument on GC



- Delay of helium ignition in red-giant
 - brightness of the tip of a red-giant branch
 - core mass increase at helium ignition
- Lifetime of stars on horizontal-branch
 - duration of helium burning
 - the number ratio of stars on HB vs. RGB
- Dominantly neutrino emission channel

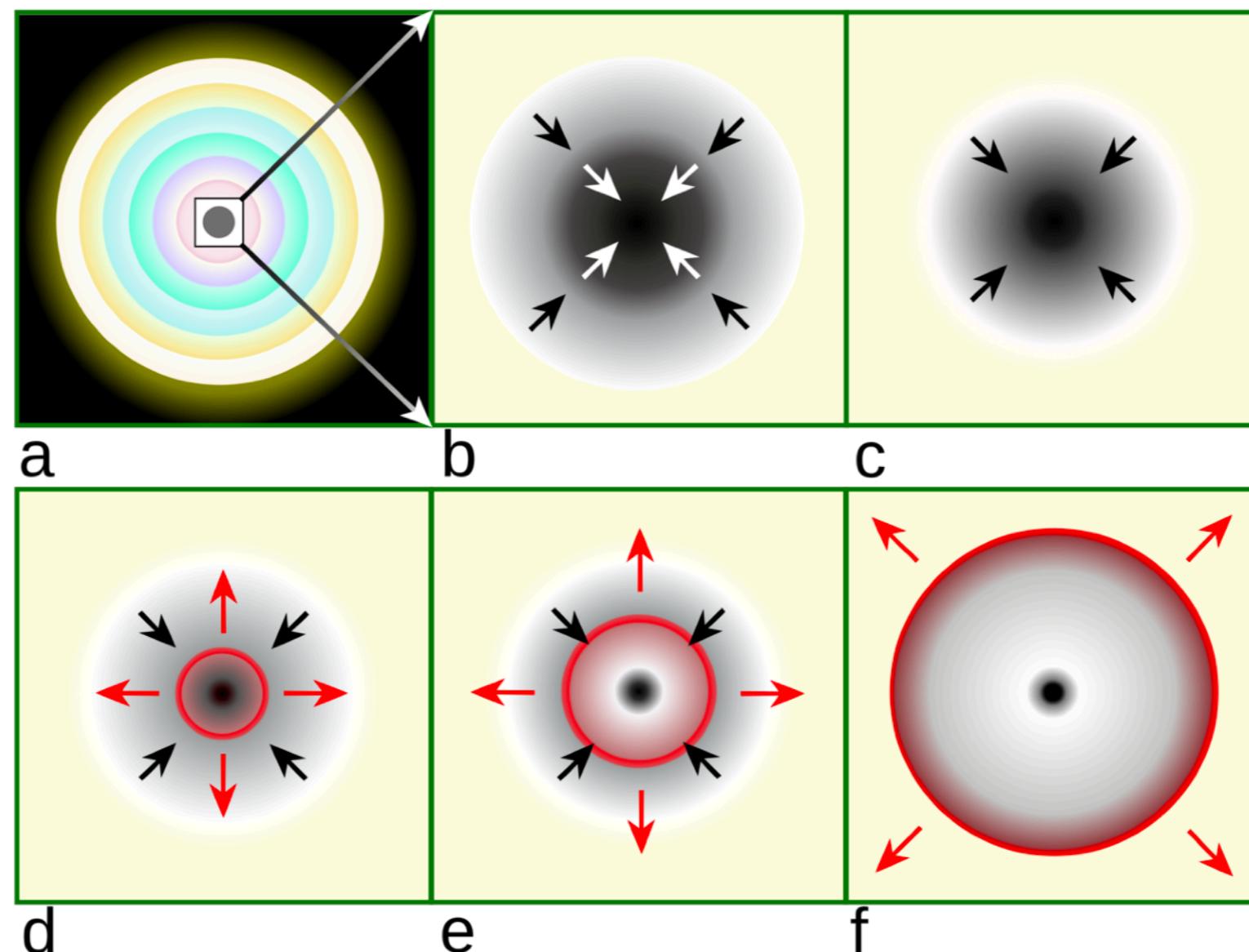
$$\epsilon_x \lesssim 10 \text{ erg g}^{-1} \text{ s}^{-1}$$

"Raffelt's criterion"

ν VS x

Core-collapse supernovae

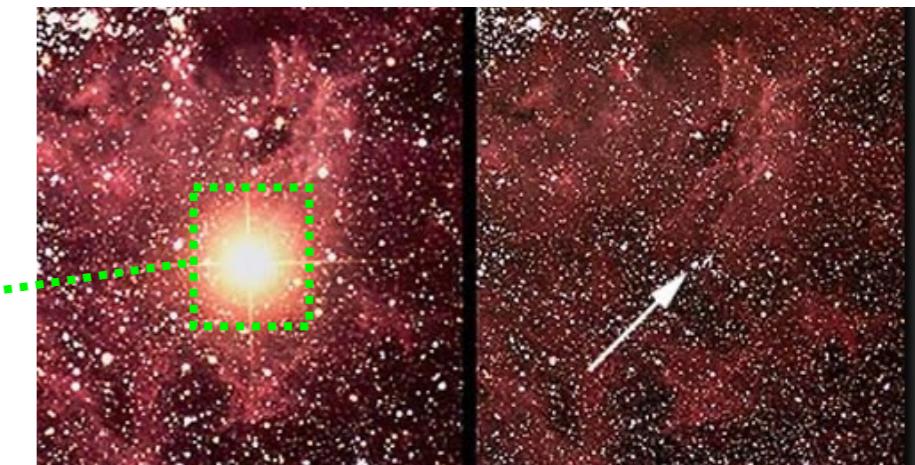
[image from Wikipedia]



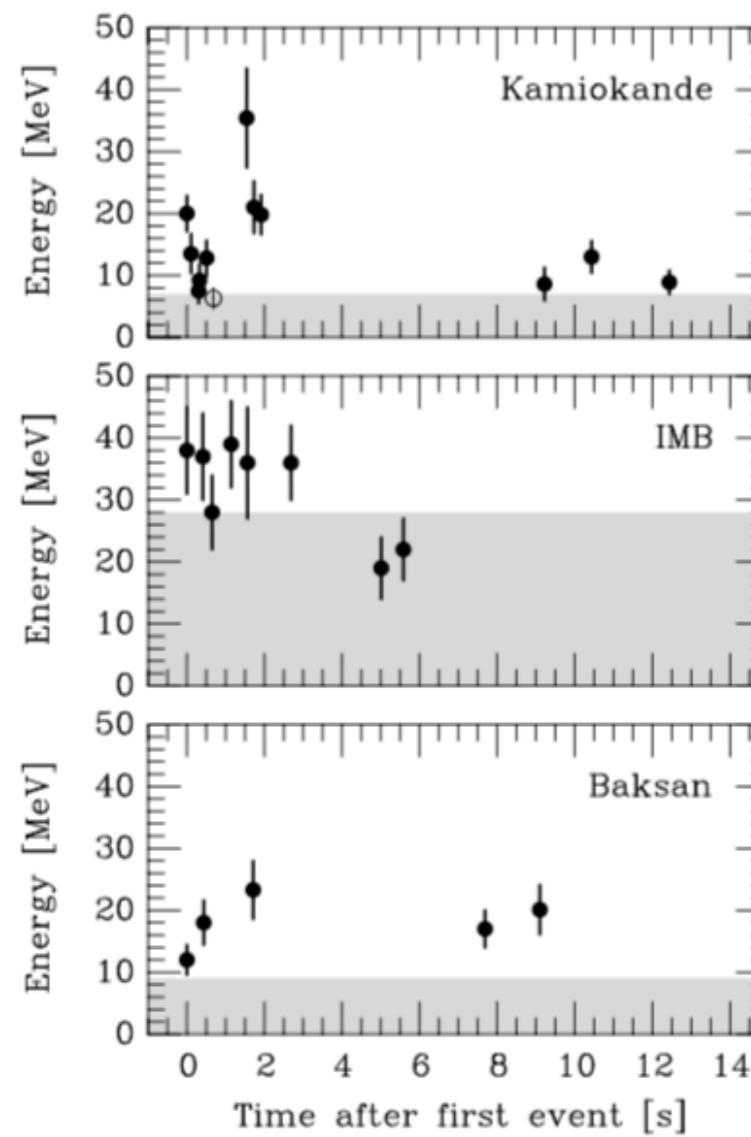
SN1987A

Remnant?

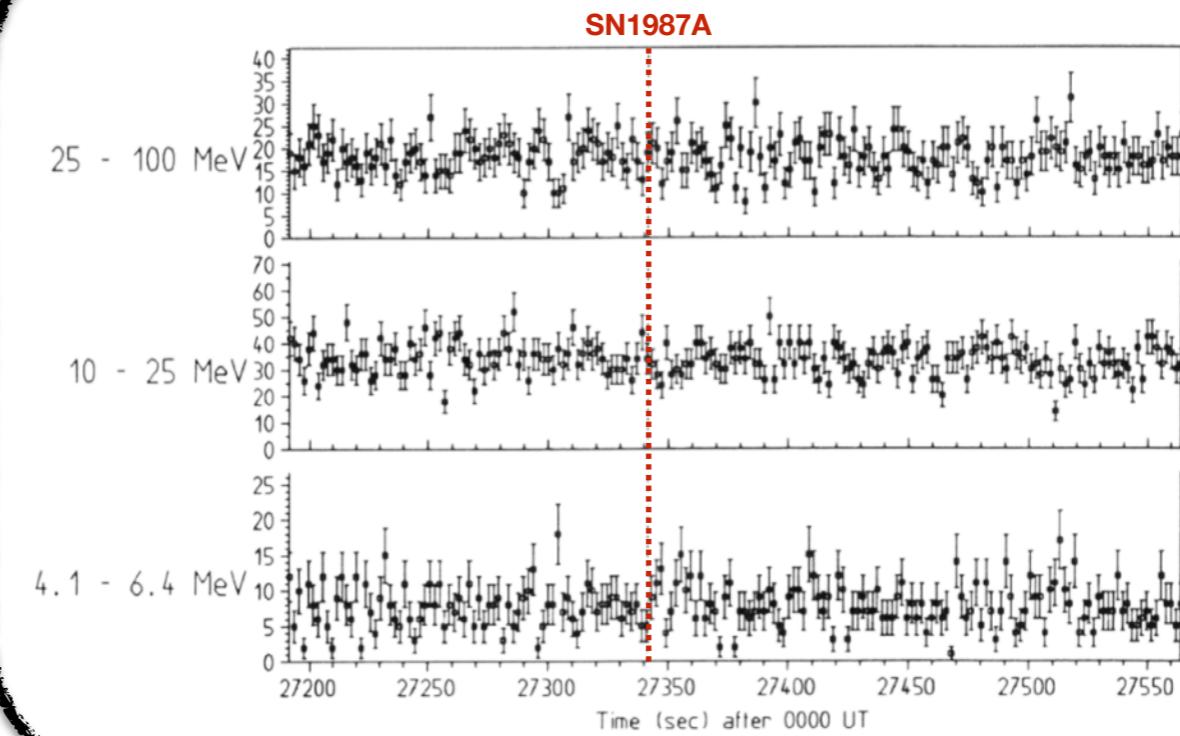
[Chandra, 17]



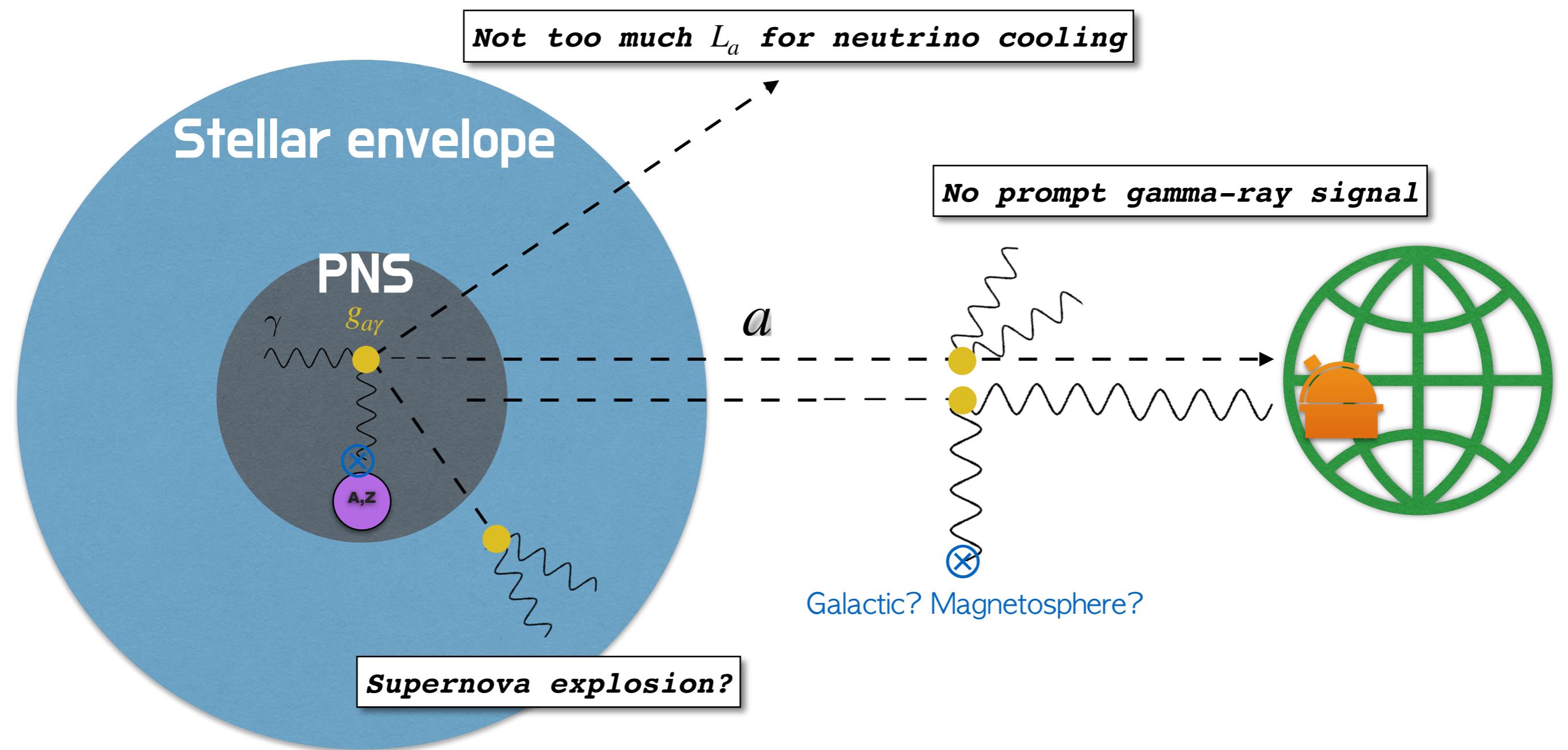
Observed neutrino signal



Lack of prompt gamma-ray signal



Searches on supernova axions

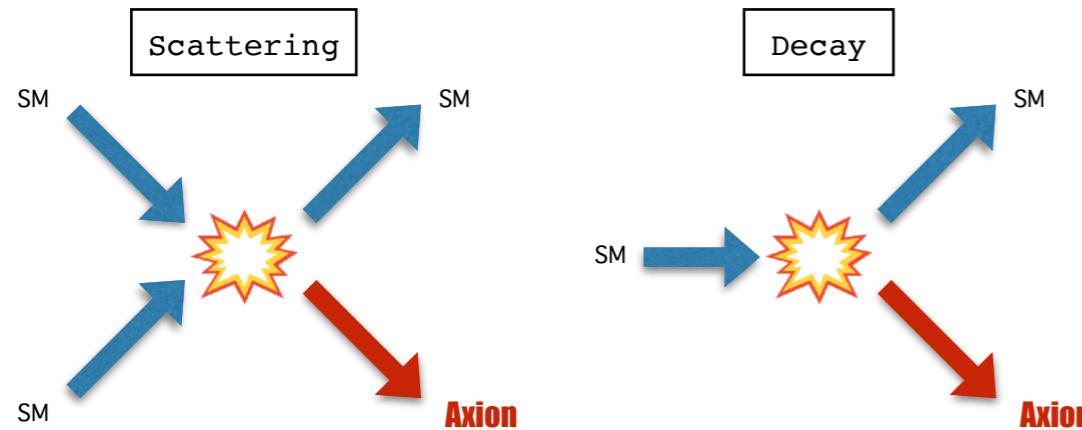


Cosmological searches

- Cosmic axions
 - particles (dark radiation) ➡ CMB spectrum, BBN, etc
 - waves (dark matter) ➡ Microwave cavities, atomic clock, etc

Cosmological axion relic

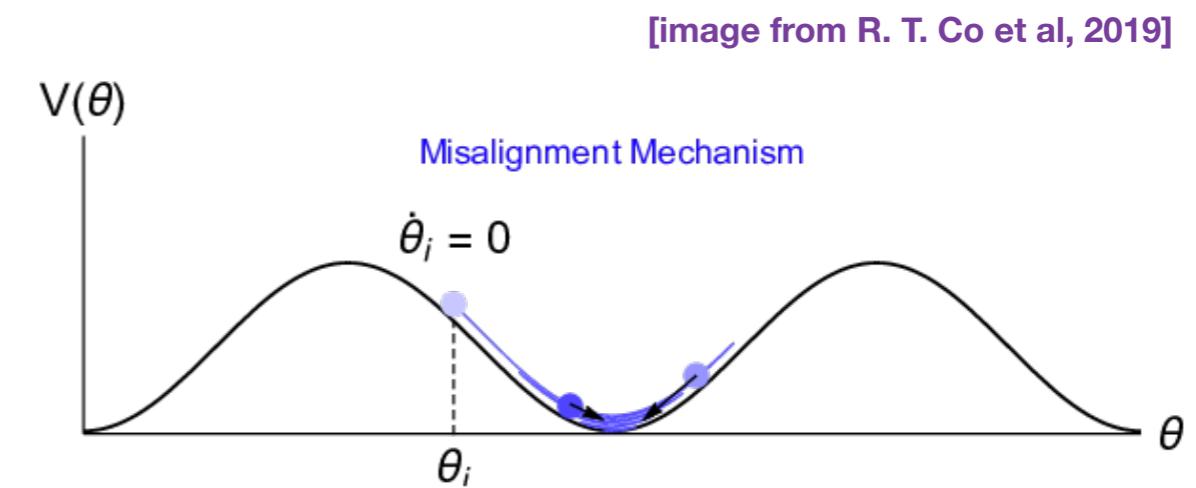
Cosmic axion background



[SY et al, 21][SY et al, 22], etc

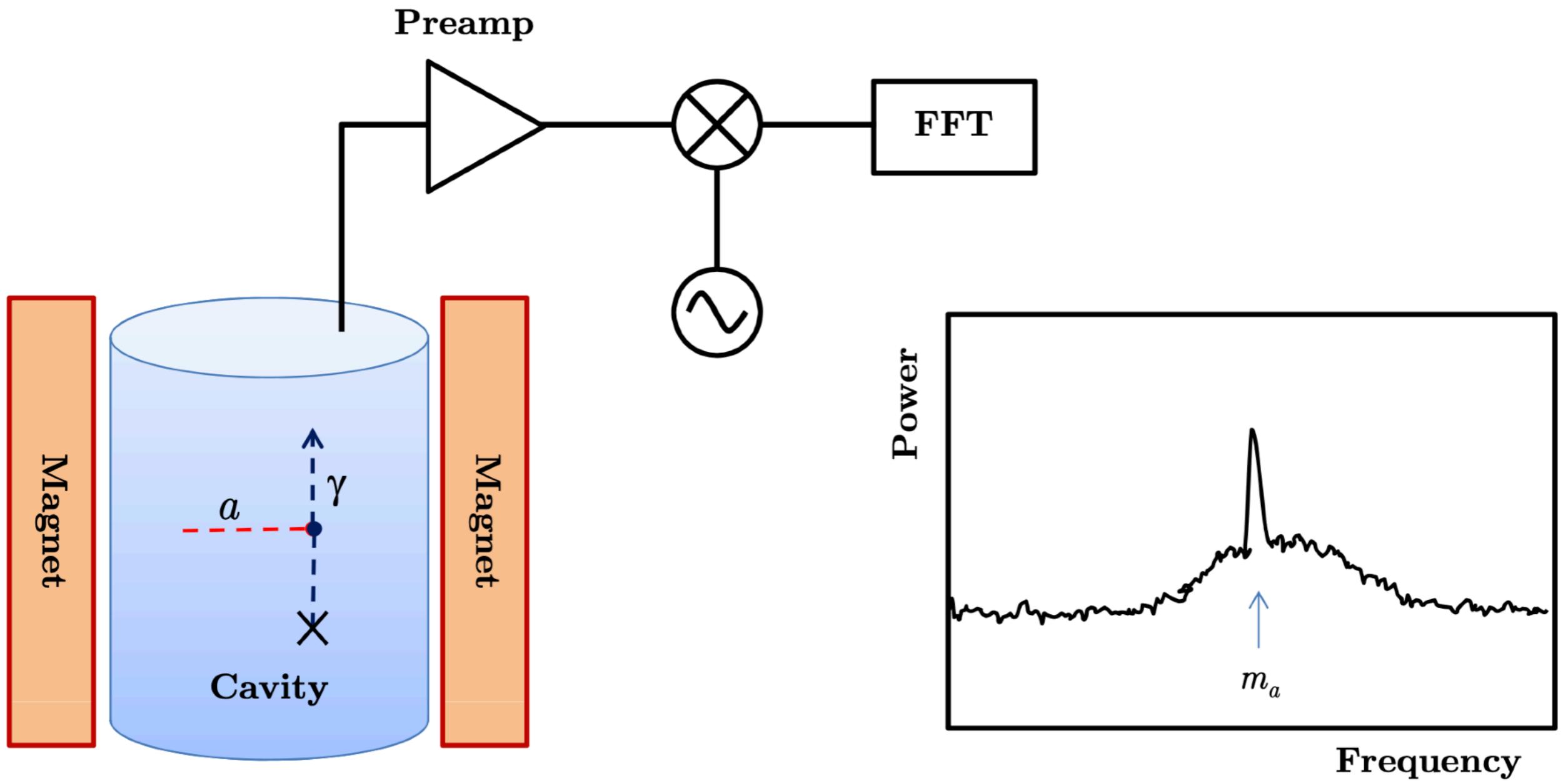
- Thermally or non-thermally
- Additional radiation degree
- CMB spectrum, BBN, etc

Wave dark matter



- Ultra light dark matter
- (coherent) Wave feature
- Microwave cavity, atomic clock, etc

Microwave cavity

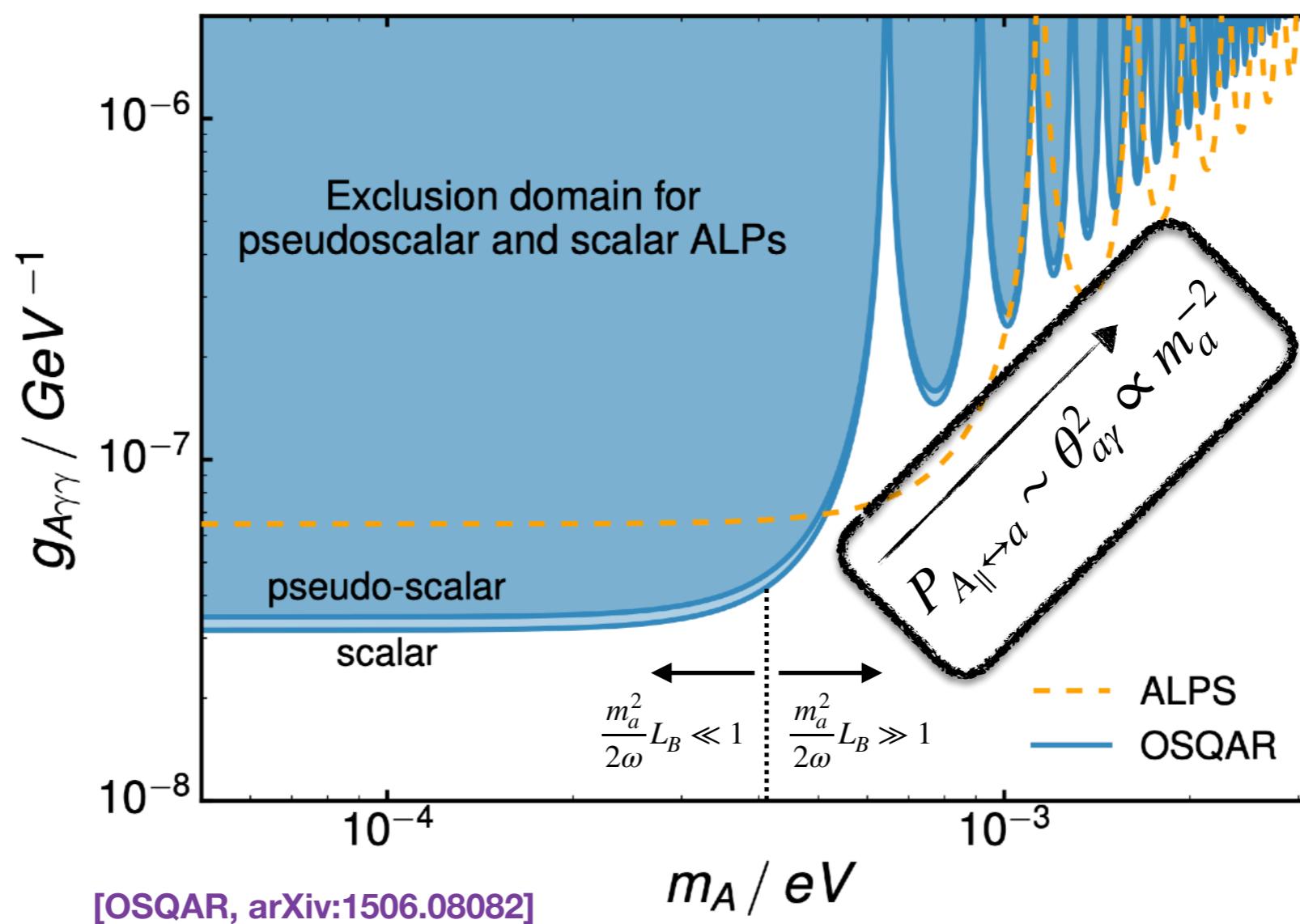


Axion magnetic resonance

- Enhancement of axion-photon conversion in a spatially/temporally *rotating* magnetic field
- Applications to LSTW & Helioscope

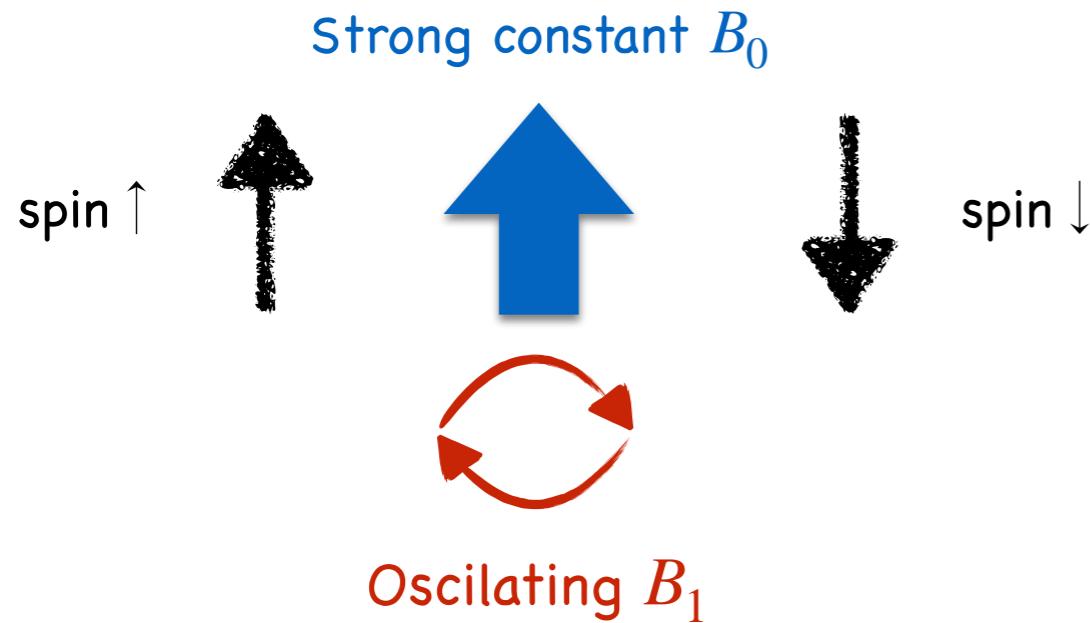
Light-Shining-Through-Walls

- In the conventional setups, a background magnetic field is constant



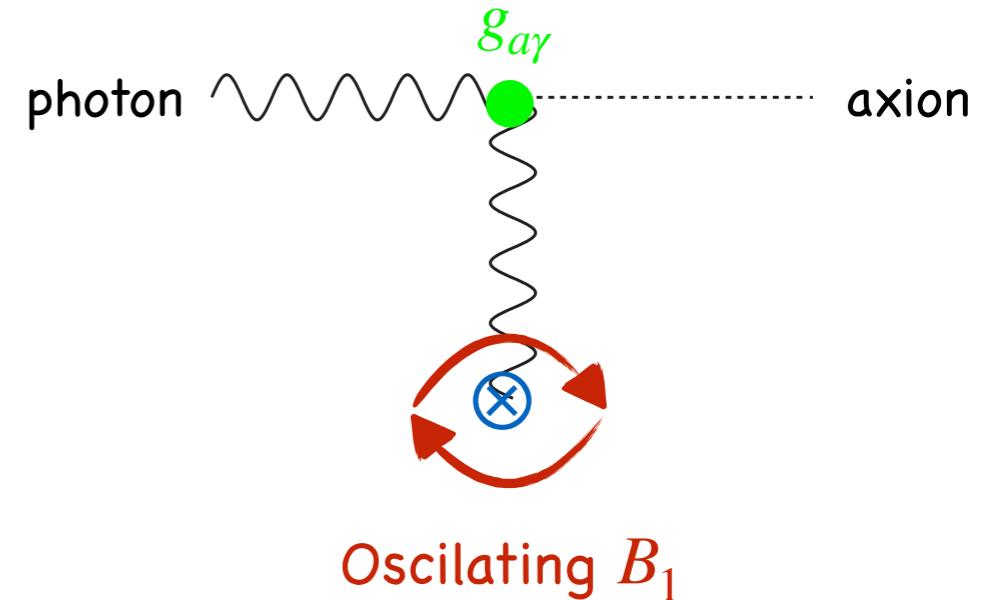
Axion magnetic resonance

Nuclear magnetic resonance



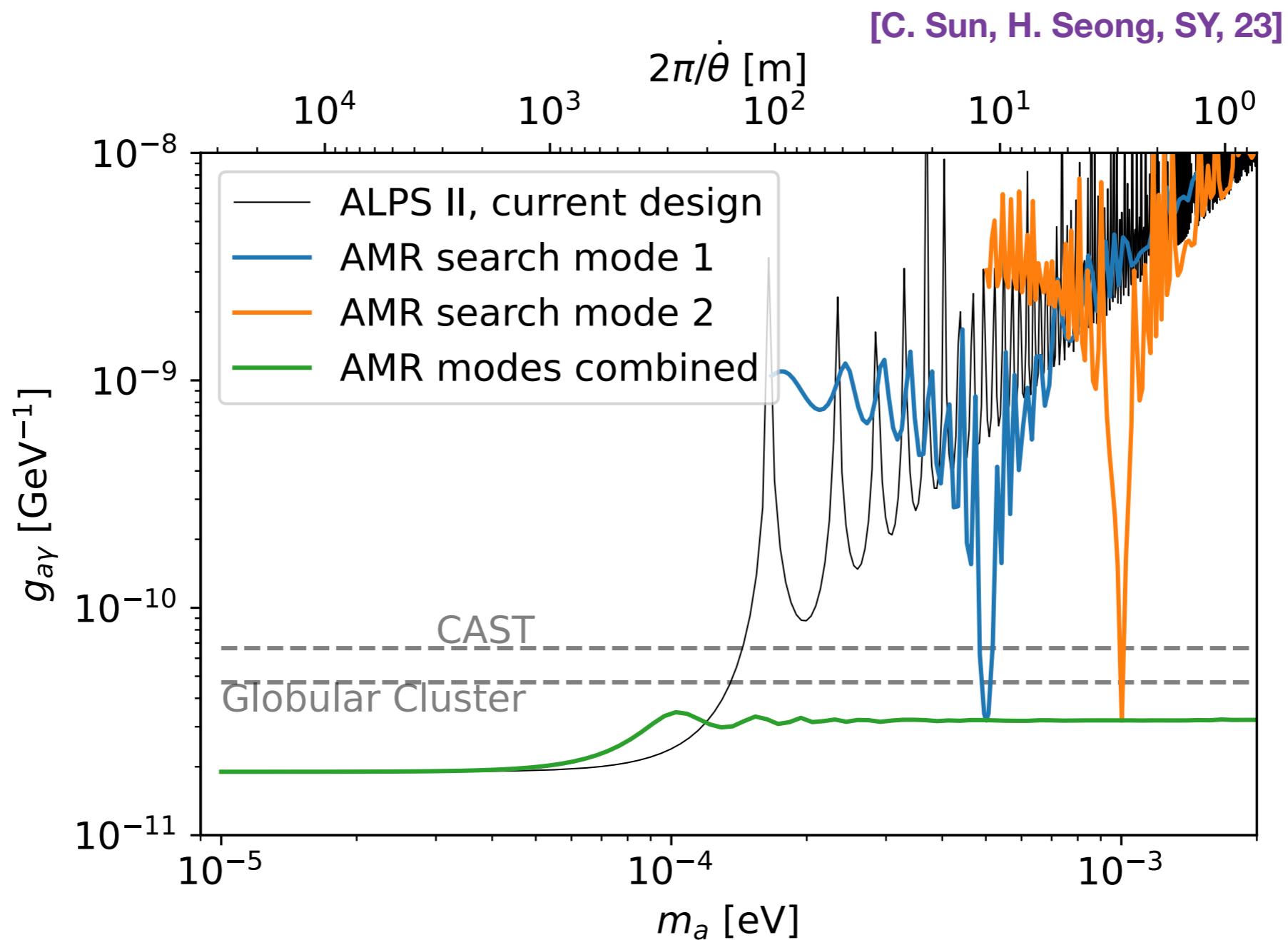
- Two states: spin $\uparrow \downarrow$
- Larmor precession frequency $\omega_0 = \mu B_0$
- Transition in oscillating B_1 with $\dot{\theta} = d\hat{B}_1/dt$
- Rabi frequency $\sqrt{(\omega_0 - \dot{\theta})^2 + (\mu B_1)^2}$

Axion magnetic resonance



- Two state: photon & axion
- Axion momentum transfer $\omega_a = m_a^2/2E$
- Transition in oscillating B_1 with $\dot{\theta} = d\hat{B}_1/dt$
- Rabi frequency $\sqrt{(\omega_a - \dot{\theta})^2 + (g_{a\gamma} B/2)^2}$

Axion magnetic resonance



Thank you!

Back up

Effective Chiral Lagrangian

- Quark-antiquark condensation as the chiral breaking

$$\begin{pmatrix} \pi^0 + \frac{1}{\sqrt{3}}\eta & \sqrt{2}\pi^+ & \sqrt{2}K^+ \\ \sqrt{2}\pi^- & -\pi^0 + \frac{1}{\sqrt{3}}\eta & \sqrt{2}K^0 \\ \sqrt{2}K^- & \sqrt{2}\bar{K}^0 & -\frac{2}{\sqrt{3}}\eta' \end{pmatrix} + \frac{2}{\sqrt{6}}\eta'$$

$\langle \bar{q}_L q_R \rangle = \Lambda_{\text{QCD}}^3 \Sigma$ where $\Sigma = e^{i\Phi}$ with Φ the meson-field matrix (octet+singlet)

- Effective theory of mesons "Chiral perturbation theory"

$$M_q \rightarrow LM_q R^\dagger, \quad \bar{\theta} \rightarrow \bar{\theta} + \arg \text{Det} [LR^\dagger], \quad \Sigma \rightarrow R \Sigma L^\dagger$$

- Meson potential at the leading order - based on the symmetry argument

$$-V = \left(\Lambda^4 e^{i\bar{\theta}} \text{Det} \Sigma + \text{h.c.} \right) + \left(\frac{m_\pi^2 f_\pi^2}{2(m_u + m_d)} \text{tr} [M_q e^{i\Phi}] + \text{h.c.} \right)$$

$$U(1)_A \Rightarrow \eta'$$

The vacuum in the meson space

- Meson vacuum in the presence of $\bar{\theta} \neq 0$, i.e., $\left. \frac{\partial V}{\partial \Phi} \right|_{\bar{\theta}} = 0$

$$\langle \Phi \rangle = \text{diag} [\phi_u, \phi_d, \phi_s] \Rightarrow -V = \Lambda^4 \cos \left(\sum \phi_q + \bar{\theta} \right) + \frac{m_\pi^2 f_\pi^2}{m_u + m_d} \sum m_q \cos \phi_q$$

- In the limit of $m_{u,d} \ll m_s \ll \Lambda_{\text{QCD}}$,

$$\text{i)} \quad \bar{\theta} + \sum \phi_q \approx 0, \quad \text{ii)} \quad m_u \sin \phi_u = m_d \sin \phi_d = m_s \sin \phi_s$$

$$\boxed{\sin \phi_u \simeq -\frac{m_d \sin \bar{\theta}}{\sqrt{m_u^2 + m_d^2 + 2m_u m_d \cos \bar{\theta}}}, \quad \sin \phi_d \simeq -\frac{m_u \sin \bar{\theta}}{\sqrt{m_u^2 + m_d^2 + 2m_u m_d \cos \bar{\theta}}}, \quad \sin \phi_s \simeq -\frac{m_u m_d \sin \bar{\theta}}{m_s \sqrt{m_u^2 + m_d^2 + 2m_u m_d \cos \bar{\theta}}}}$$

if $\bar{\theta} \ll 1$



$$\langle \Phi \rangle = -\frac{M_q^{-1}}{\text{tr } M_q^{-1}} \bar{\theta}$$

The potential of $\bar{\theta}$ (leading to axion potential)

$$-V = \Lambda^4 \cos \left(\sum \phi_q + \bar{\theta} \right) + \frac{m_\pi^2 f_\pi^2}{m_u + m_d} \sum m_q \cos \phi_q$$

$$\sin \phi_u \simeq -\frac{m_d \sin \bar{\theta}}{\sqrt{m_u^2 + m_d^2 + 2m_u m_d \cos \bar{\theta}}}, \quad \sin \phi_d \simeq -\frac{m_u \sin \bar{\theta}}{\sqrt{m_u^2 + m_d^2 + 2m_u m_d \cos \bar{\theta}}}, \quad \sin \phi_s \simeq -\frac{m_u m_d \sin \bar{\theta}}{m_s \sqrt{m_u^2 + m_d^2 + 2m_u m_d \cos \bar{\theta}}}$$



$$V_{\bar{\theta}} = -f_\pi^2 m_\pi^2 \sqrt{\frac{m_u^2 + m_d^2 + 2m_u m_d \cos \bar{\theta}}{(m_u + m_d)^2}}$$

Stellar object

“ Star

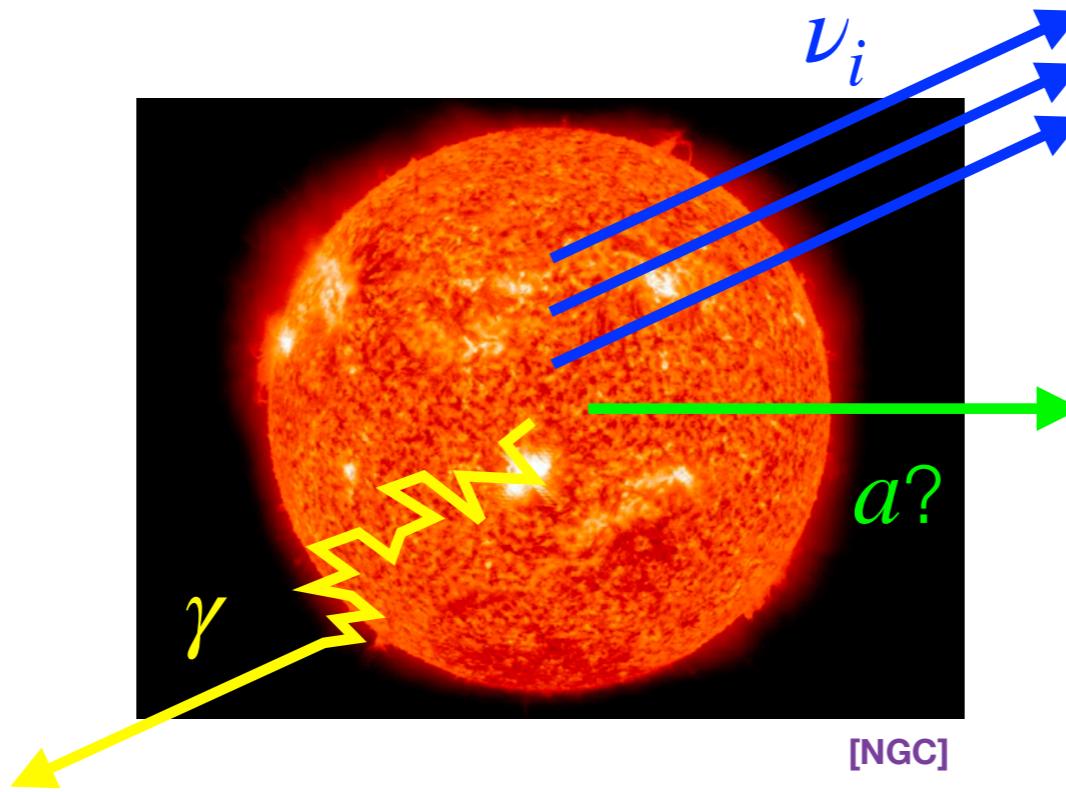
From Wikipedia, the free encyclopedia

This article is about the astronomical object. For other uses, see [Star \(disambiguation\)](#).

A **star** is an [astronomical object](#) consisting of a luminous [spheroid](#) of [plasma](#) held together by its own [gravity](#). ”

- can be used as laboratories to probe fundamental interactions.
 - observed photon or neutrino signal
 - e.g. *γ-ray spectral modulation which can be explained by axion-photon oscillations*
 - Stellar evolution
 - extra energy leakage could change the evolutionary path which fits well to the standard scenario*

Stellar argument on Sun



- Standard solar model calibrated to match the present-day solar radius & luminosity

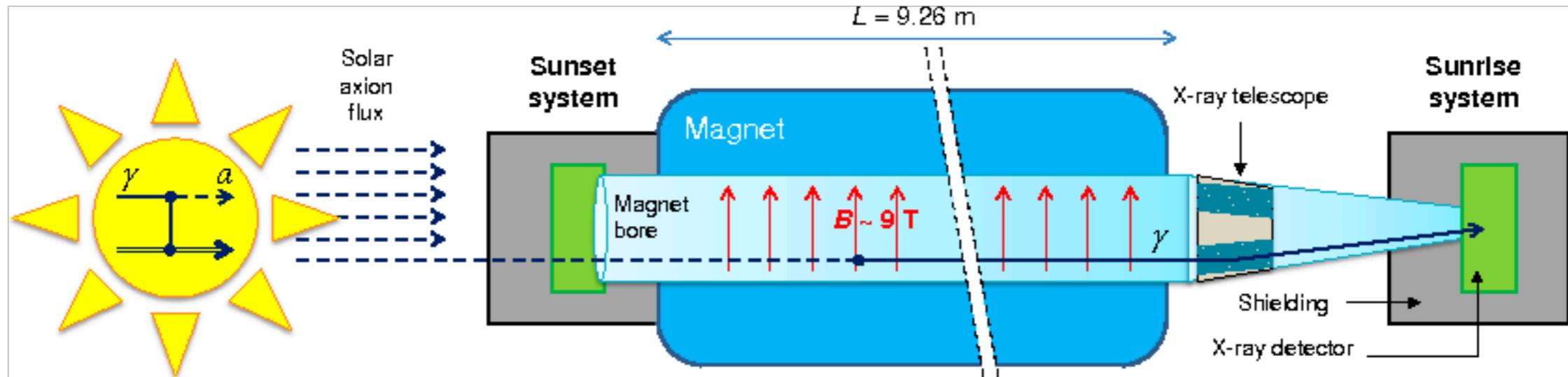
$$R_{\odot} = 6.9598 \times 10^{10} \text{ cm} \quad L_{\odot} = 3.8418 \times 10^{33} \text{ erg s}^{-1}$$

- Observables:
 - 1) Helioseismology (sound speed, surface helium)
 - 2) Solar neutrino fluxes $\Phi(^8\text{B})$ and $\Phi(^7\text{Be})$

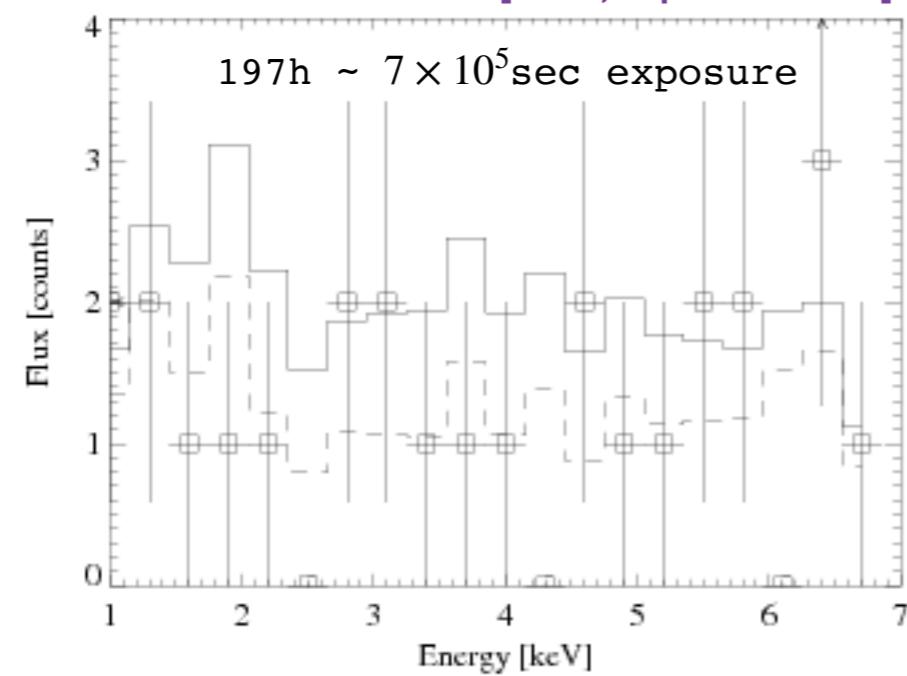
Axion cooling leading to observable deviations

Axion helioscope

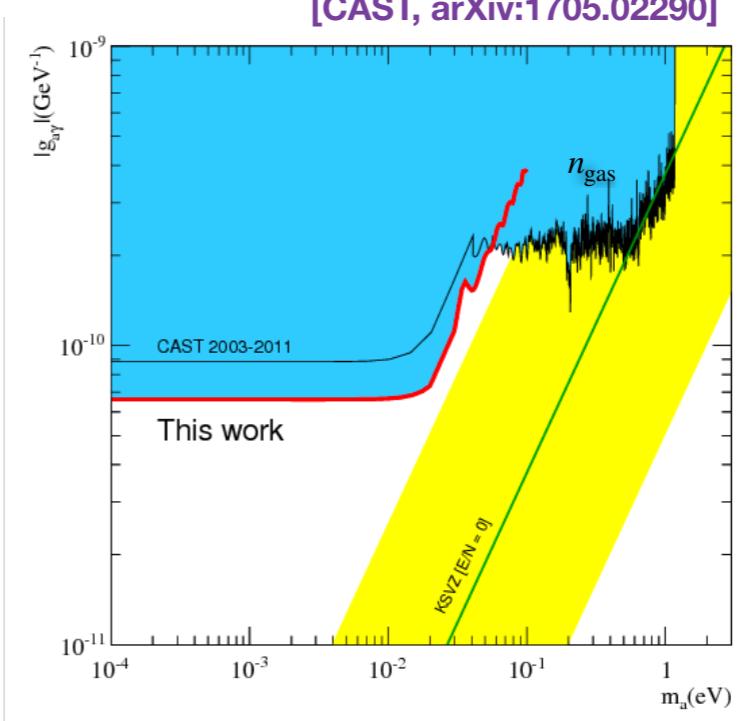
[CAST, arXiv:1705.02290]



[CAST, hep-ex/0702006]



[CAST, arXiv:1705.02290]



Stellar cooling argument on HB

triple- α reaction

- A helium-burning core and a hydrogen-burning shell after the tip of red-giants

$$\rho \approx 10^4 \text{ g cm}^{-3} \quad T \approx 10^8 \text{ K}$$

- The number ratio R of HB vs. RGB in Globular clusters agreed with standard predictions to within 10%
- Lifetime of HB t_{He} (i.e., duration of helium burning phase) agrees within 10%

$$t_{\text{He}} \text{ reduced by } \frac{L_{3\alpha}}{L_a + L_{3\alpha}}$$



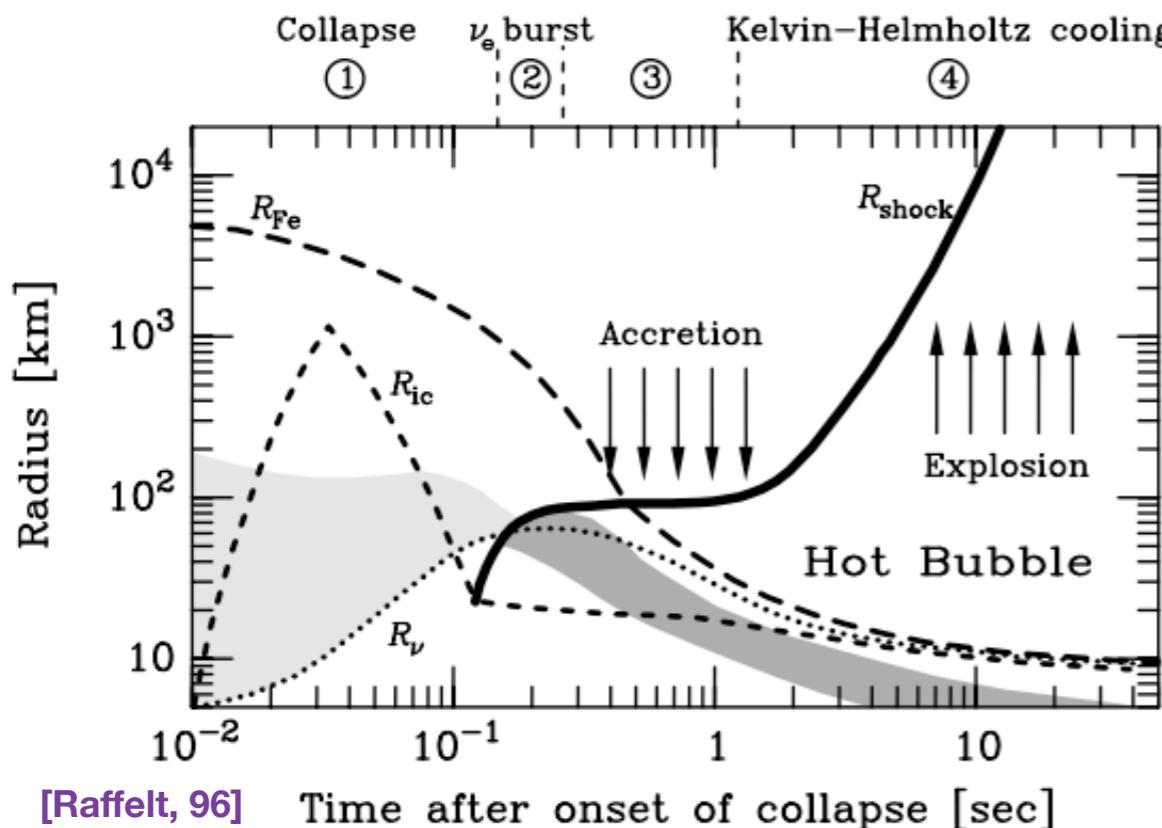
$$L_a \lesssim 10 \text{ erg g}^{-1} \text{ s}^{-1}$$

[Raffelt]

$$L_{3\alpha} \sim 20L_{\odot} \approx 8 \times 10^{34} \text{ erg s}^{-1}$$

$$M_c^{\text{HB}} \sim 0.5 M_{\odot}$$

Stellar cooling argument on SN



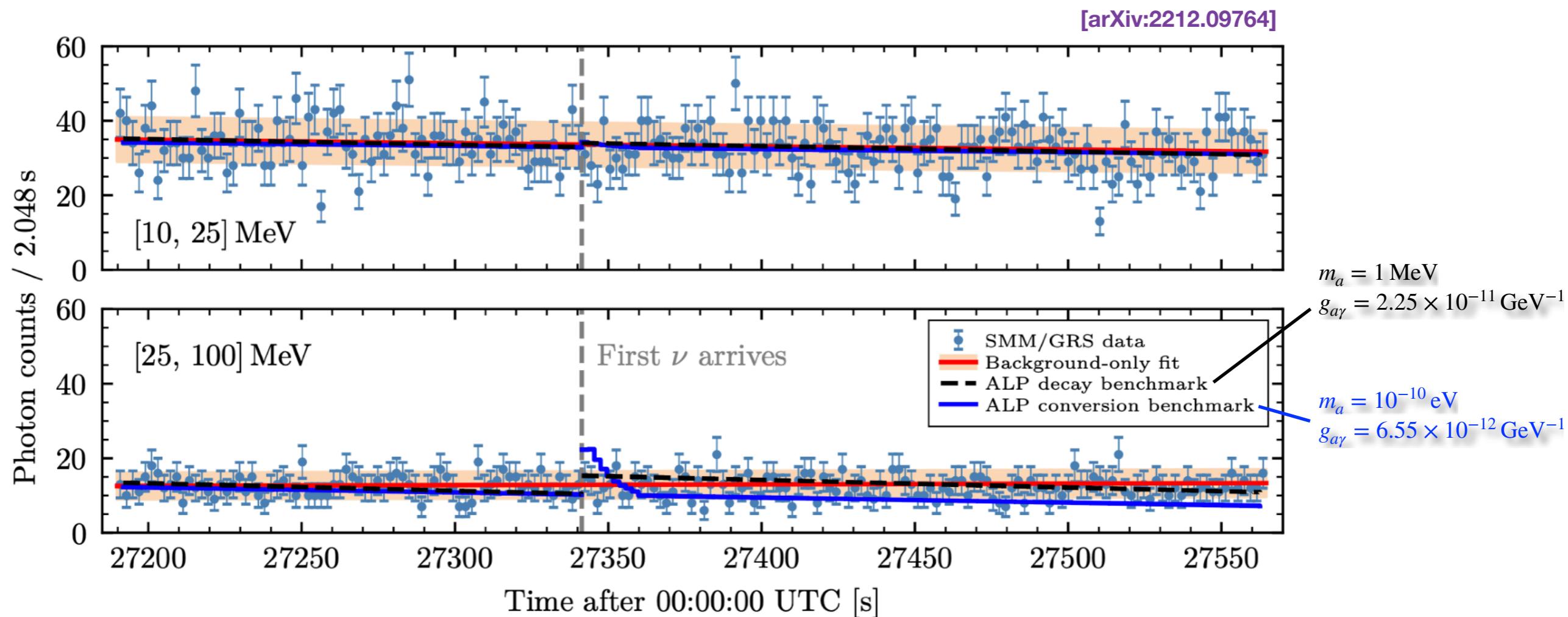
- The “**delayed explosion scenario**” ($D\nu M$)
 - Thermal relaxation of the proto-neutron star as the source of neutrino emission for $> 5\text{ sec}$
 - the energy deposition of the neutrino flow causes the explosion ($\mathcal{O}(1)\%$ of total neutrino bulk energy)



$$L_x \lesssim 3 \times 10^{52} \text{ erg s}^{-1}$$

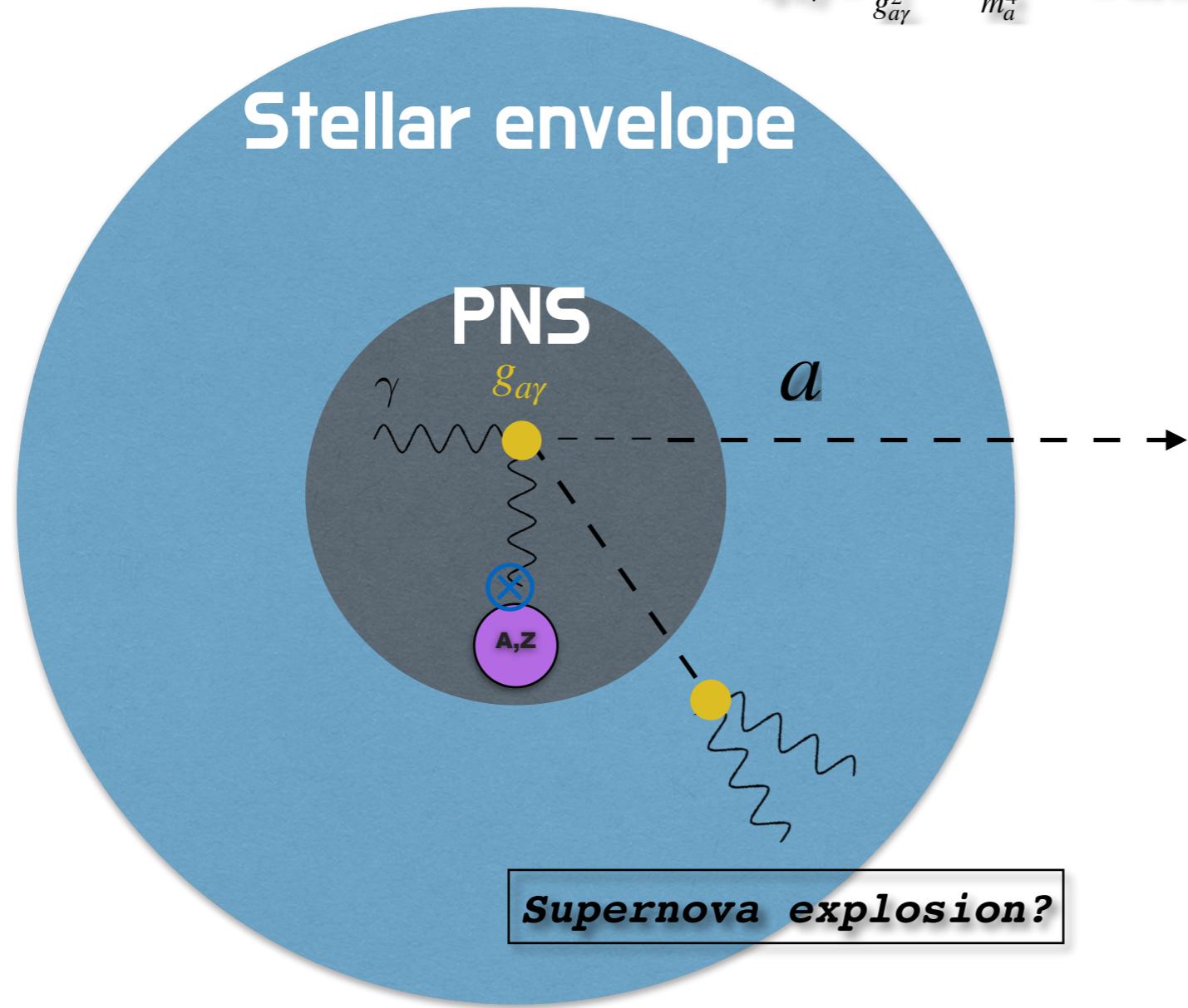
“**Raffelt’s criterion**”

SN1987A gamma-ray burst bound



Low energy supernova?

$$\lambda_{a \rightarrow 2\gamma} = \frac{64\pi}{g_{a\gamma}^2} \frac{\sqrt{\omega^2 - m_a^2}}{m_a^4} = 4.0 \times 10^{13} \text{ cm} \left(\frac{g_{a\gamma}}{10^{-9} \text{ GeV}^{-1}} \right)^{-2} \left(\frac{\omega}{100 \text{ MeV}^{-1}} \right) \left(\frac{m_a}{10 \text{ MeV}} \right)^{-4}$$

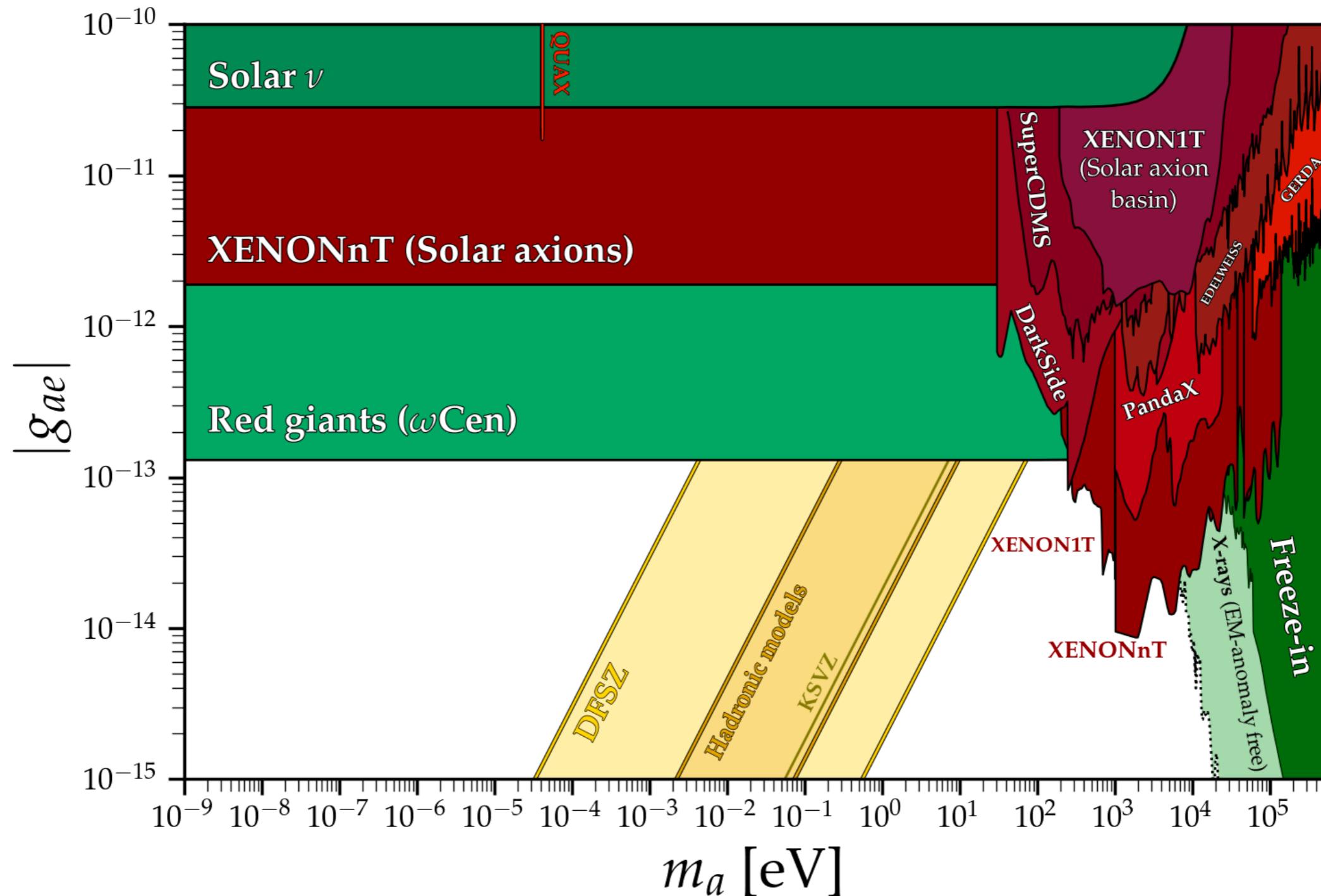


- Observables:
 - 1) Light curves
 - 2) Spectroscopic ejecta velocities

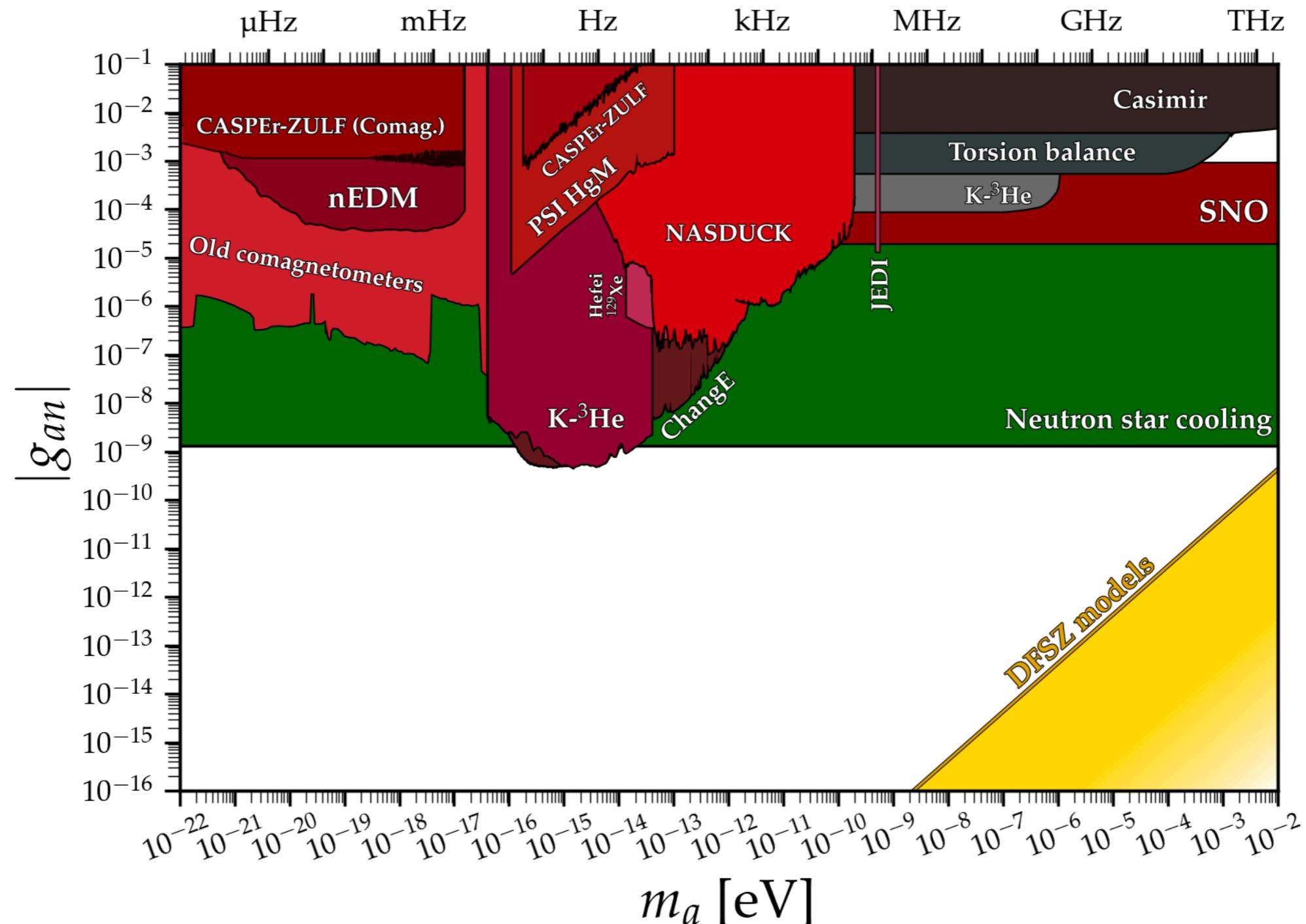
[arXiv:2201.09890]

$E_{\text{exp}} < 0.1 \text{ Bethe} = 10^{50} \text{ erg}$

Current status on $i g_{aee} a \bar{e} \gamma^5 e$



Current status on $i g_{ann} a \bar{n} \gamma^5 n$



Current status on $i g_{ann} a \bar{n} \gamma^5 n$

