

QCHSC 2024

The XVth Quark Confinement and the Hadron Spectrum Conference

ALICE explores strangeness and nucleosynthesis in hadronic collisions

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University and INFN Bologna

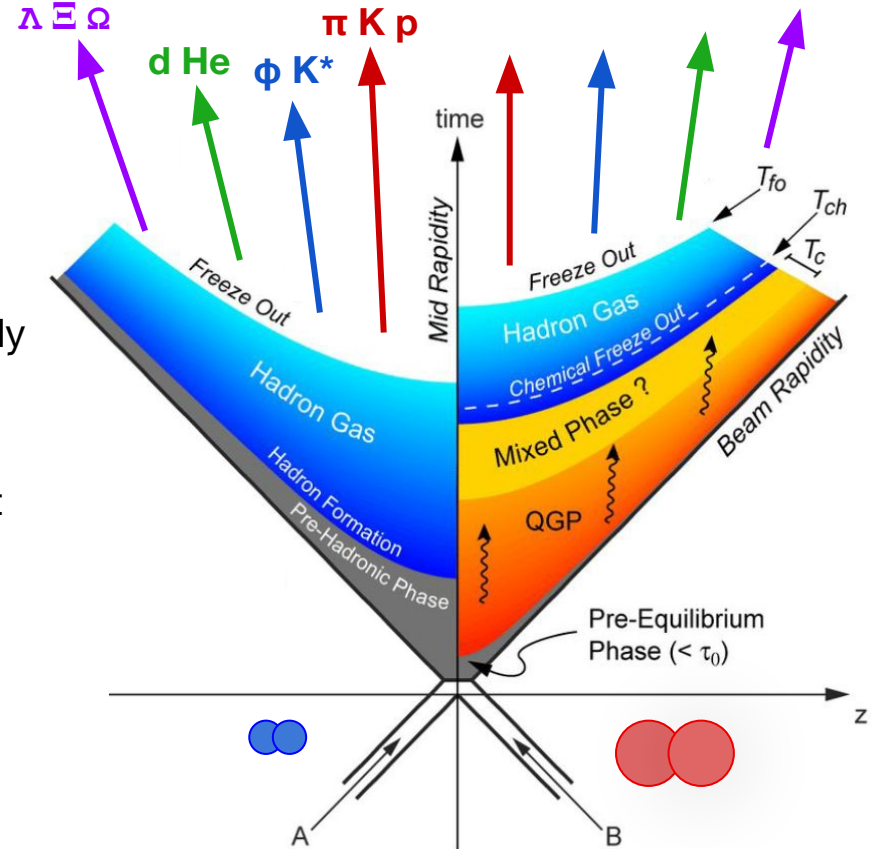


Light flavour production in hadronic collisions

Light flavor (u, d, s) quarks **thermally produced** in the QGP → soft probes are key to study the properties of the medium

Large interest also in **small systems** (pp, pA) not only as a reference to AA

- different **system sizes** are sensitive to different production mechanisms
- striking similarities observed **between small and large systems**

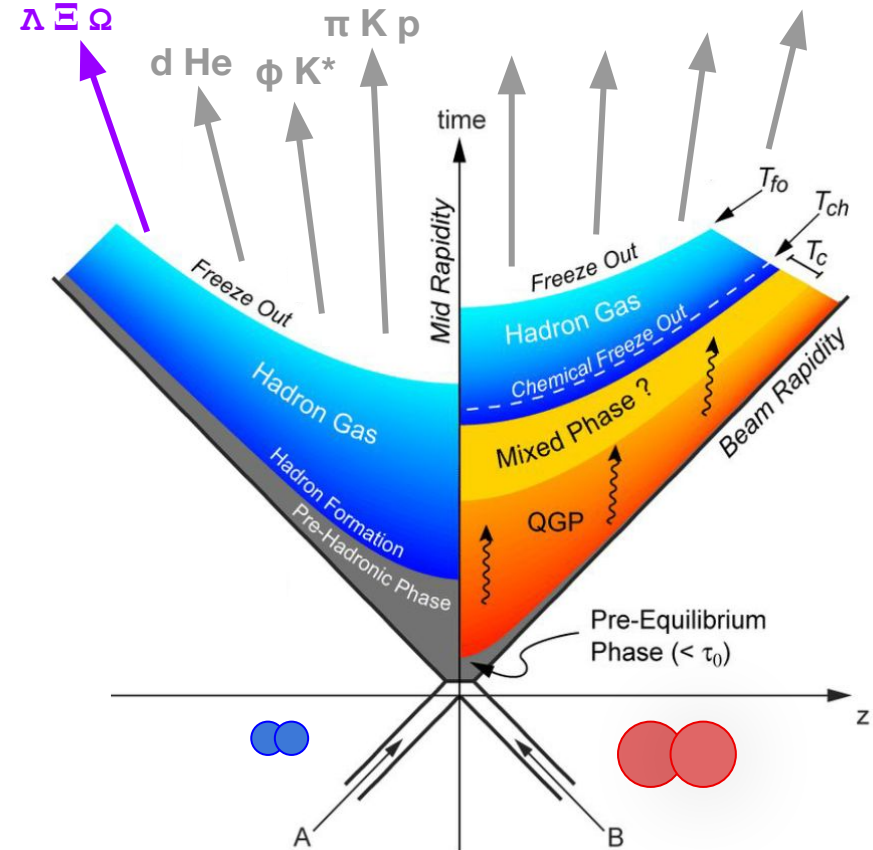


Light flavour production in hadronic collisions

Strangeness production:

Strangeness enhancement in AA is historically considered a signature of QGP formation, with Minimum-Bias pp as reference

Striking **similarities** observed **between small and large systems** as a function of multiplicity



Light flavour production in hadronic collisions

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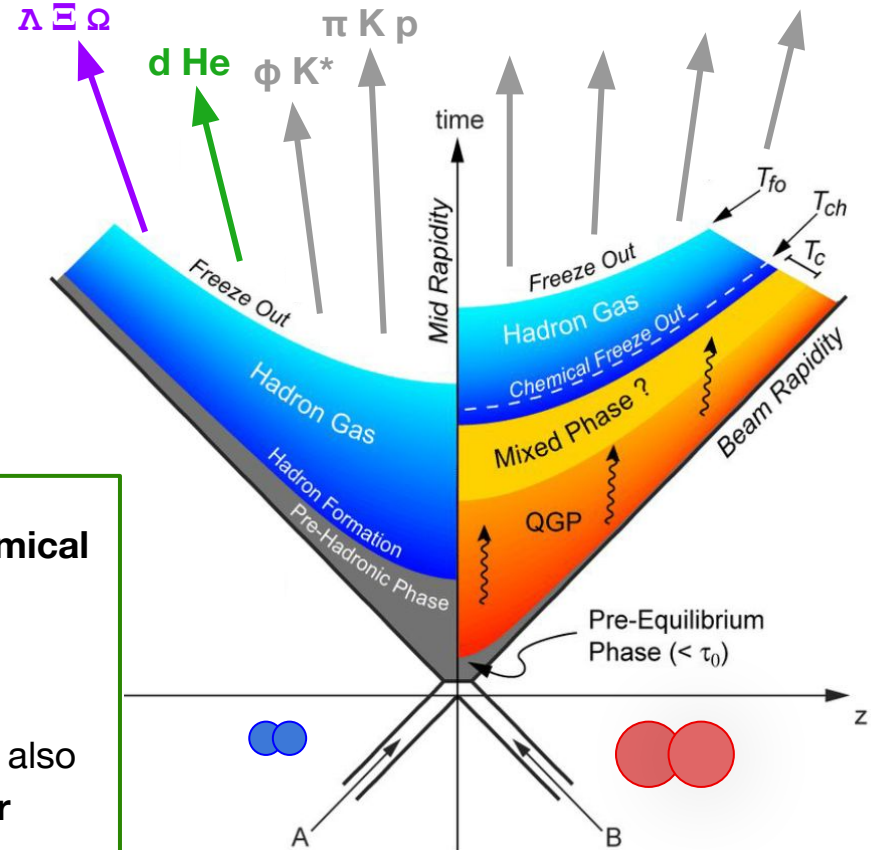
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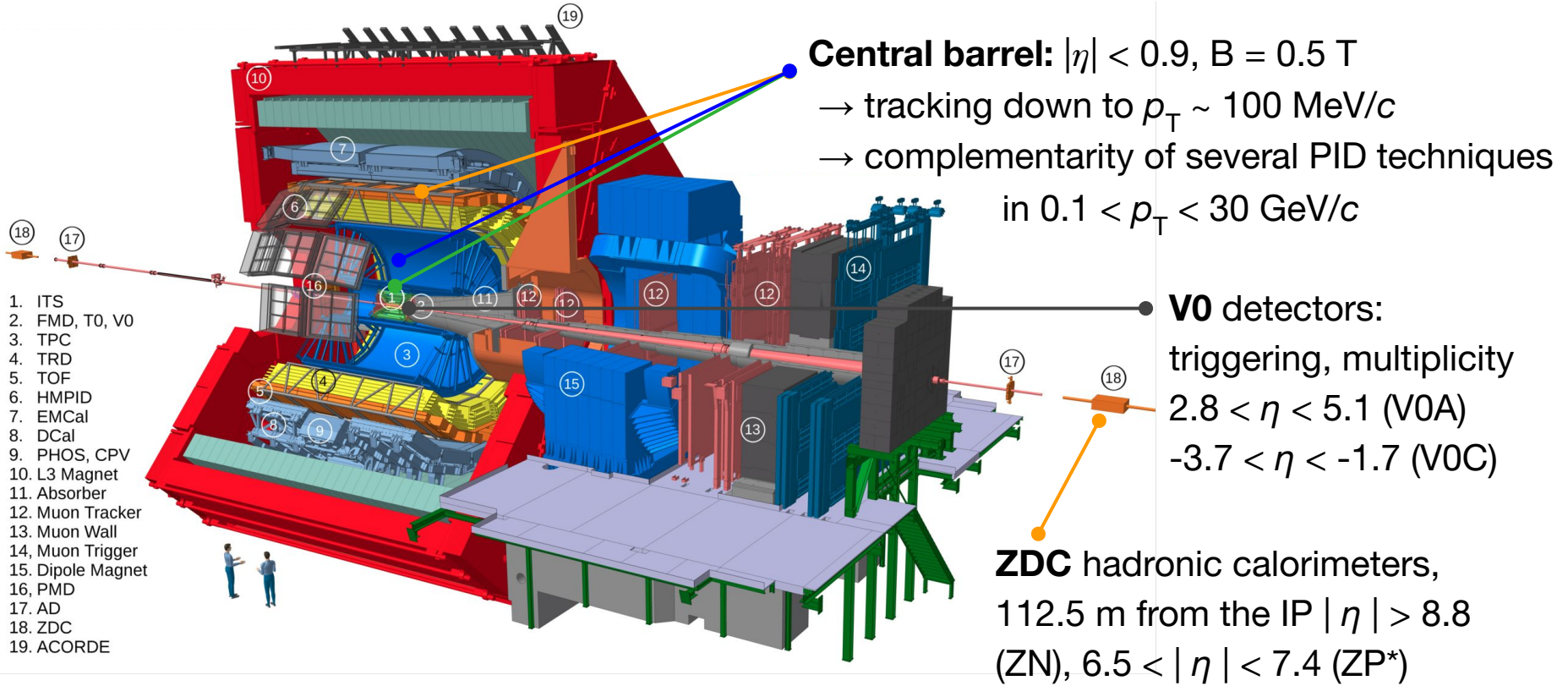
(Anti)nuclei production:

Constrain the **properties** of the system in AA at the **chemical freeze-out** (T, μ_B) and understand **nucleosynthesis mechanisms** in high-energy collisions

The production of antinuclei measured at accelerators is also crucial input in **astrophysical searches for Dark Matter**



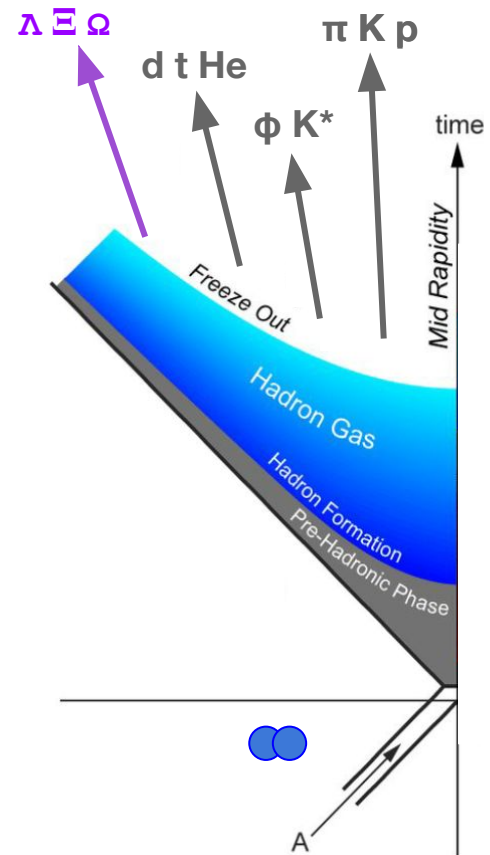
The ALICE detector in Run 1 and 2



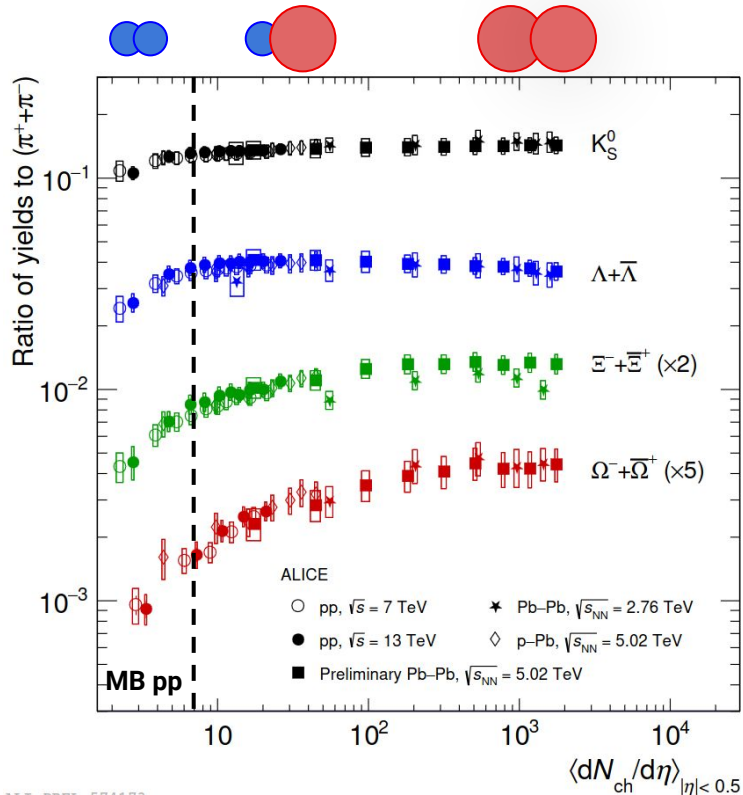
1. ITS
2. FMD, TO, V0
3. TPC
4. TRD
5. TOF
6. HMPID
7. EMCal
8. DCal
9. PHOS, CPV
10. L3 Magnet
11. Absorber
12. Muon Tracker
13. Muon Wall
14. Muon Trigger
15. Dipole Magnet
16. PMD
17. AD
18. ZDC
19. ACORDE

*considering LHC beam optics ZP acceptance for protons is $7.0 < |\eta| < 8.7$

Strangeness production in pp collisions at the LHC



Strangeness production across collision systems



Continuous evolution of strange hadron yield ratios to pions **with the charged-particle multiplicity** observed at the LHC, smoothly connecting different systems and energies

Strangeness production increases with particle multiplicity, saturating for central Pb-Pb

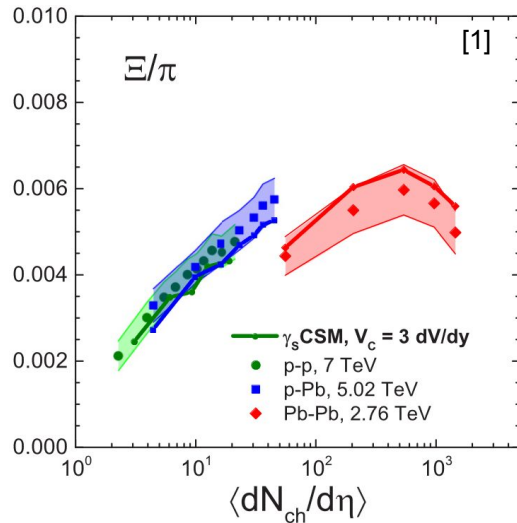
Strange content **hierarchy**:

$$|S_{\Omega^\pm}| > |S_{\Xi^\pm}| > |S_\Lambda| \approx |S_{K_S^0}|$$

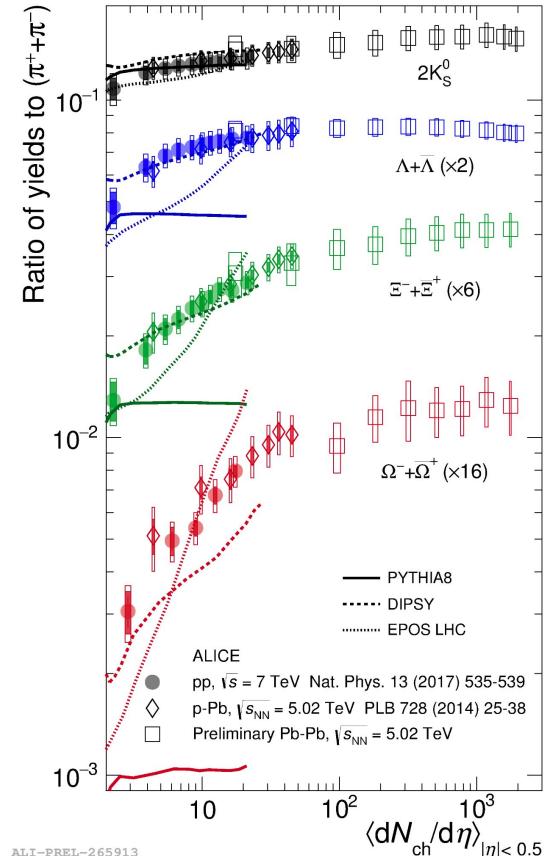
Modelling strange hadron production

Several approaches attempt to describe strangeness hadronization in small systems:

- **canonical suppression + running of Υ_S** in statistical hadronization models (CSM) [1]



- **string/rope hadronization** models including colour reconnection (CR) effects (PYTHIA/DIPSY) [2]
- **two-component (core-corona) models (EPOS)** [3]



[1] [Phys. Rev. C 100, 054906 \(2019\)](#) [2] [Phys. Rev. D 100, 074023 \(2019\)](#) [3] [Phys. Rev. C 101, 024912 \(2020\)](#)

Multi-differential analysis by ALICE

Several ALICE new results from **multi-differential analyses** to explore strangeness production in small collision systems

- Study strangeness production as a function of **new observables**, sensitive to different collision properties:
 - **Effective energy** → **This talk!**
 - Transverse **spherocity** → [JHEP 05 \(2024\) 184](#)
- Explore strangeness production **in- and out-of-jets** with different techniques
→ [JHEP 07 \(2023\) 136](#), [arxiv:2405.14511](#) , [arxiv:2405.19855](#)
- Investigate **multiple strange hadron production in high multiplicity events** → **This talk!**

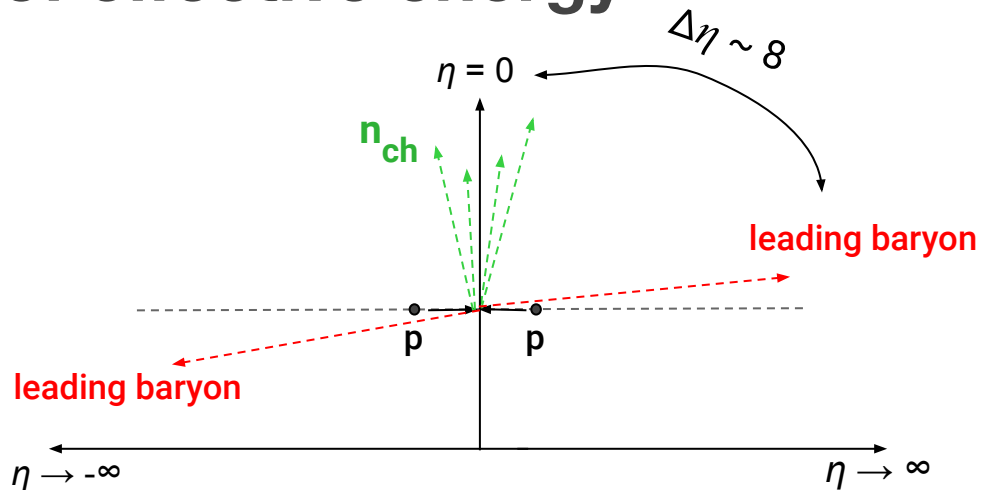
The concept of effective energy

Two-dimensional analysis as a function of:

- **Charged-multiplicity at midrapidity**
proxy for local effects,
e.g. jet production
- **Leading energy**
proxy for global effects, e.g. the
initial effective energy in the collision

$$E_{\text{eff}} = \sqrt{s} - E_{\text{leading}} \rightarrow \text{First studied at the CERN ISR in the 80's [1,2]}$$

Independent proxies given the large η separation



[1] [M. Basile et al., Phys. Lett. B 95 \(1980\) 311](#)

[2] [A. Akindinov et al., EPJ C 50 \(2007\) 341–352](#)

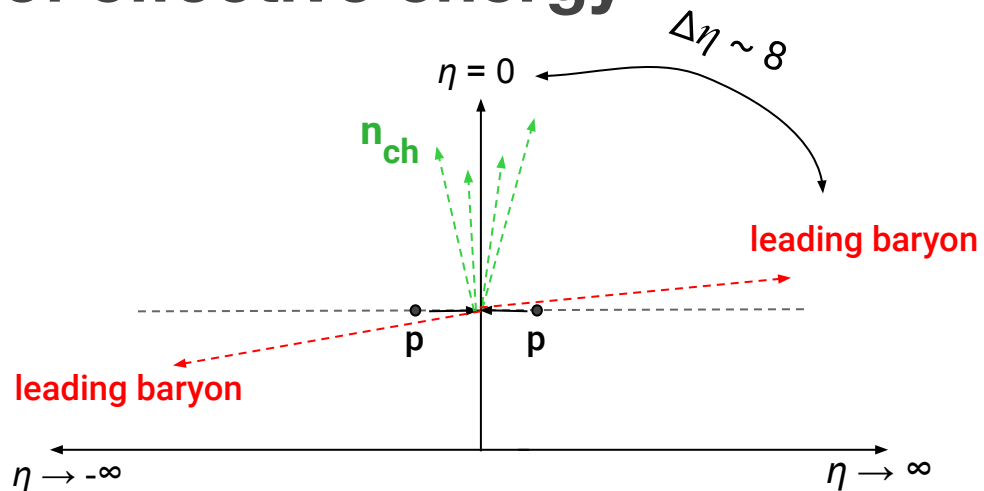
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Independent proxies given the large η separation



What is the interplay of global and local effects for strange hadron production?

[1] [M. Basile et al., Phys. Lett. B 95 \(1980\) 311](#)

[2] [A. Akindinov et al., EPJ C 50 \(2007\) 341–352](#)

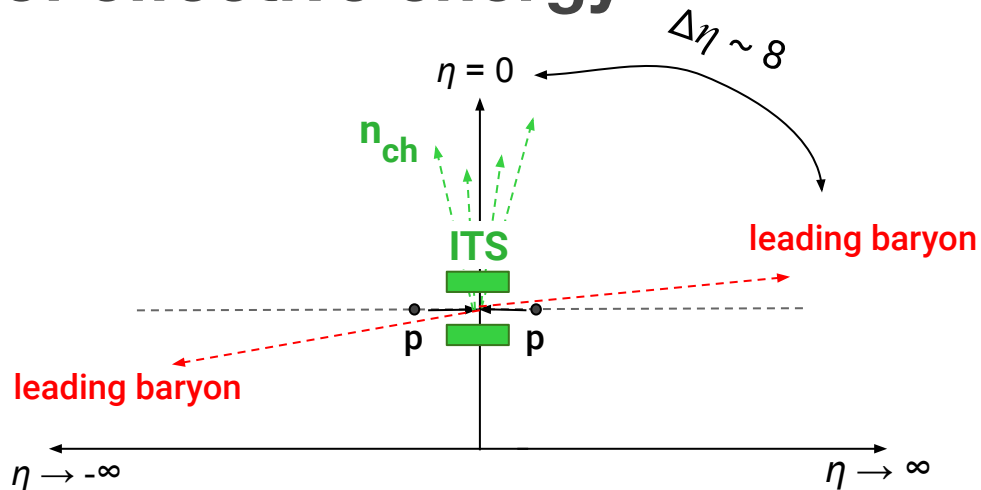
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Independent proxies given the large η separation



ALICE can measure:

- midrapidity multiplicity (**SPD**)

The concept of effective energy

Two-dimensional analysis as a function of:

- **Charged-multiplicity at midrapidity**

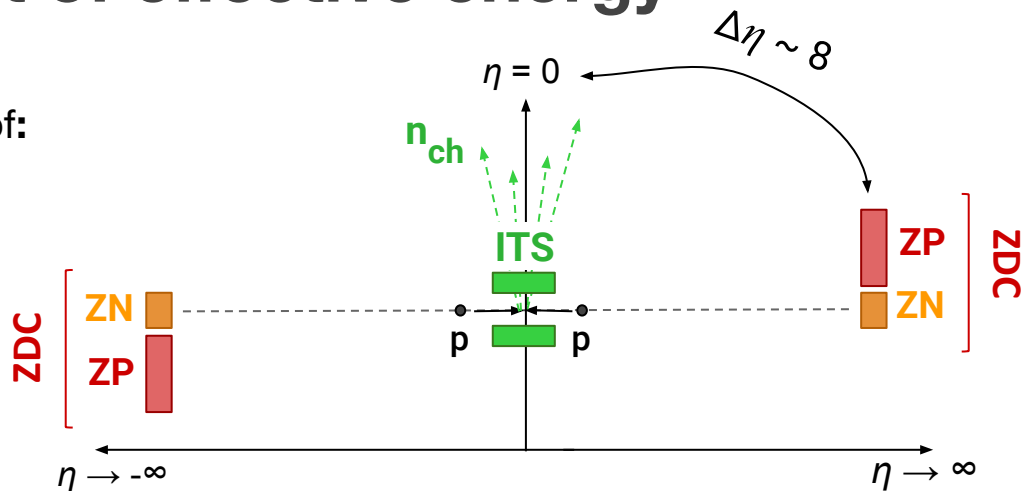
proxy for local effects,
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- **Leading energy**

proxy for global effects, e.g. the
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$$E_{\text{eff}} = \sqrt{s} - E_{\text{leading}} \approx \sqrt{s} - E_{\text{ZDC}}$$

Independent proxies given the large η separation



ALICE can measure:

- midrapidity multiplicity (**SPD**)
- leading energy (**ZDC**)

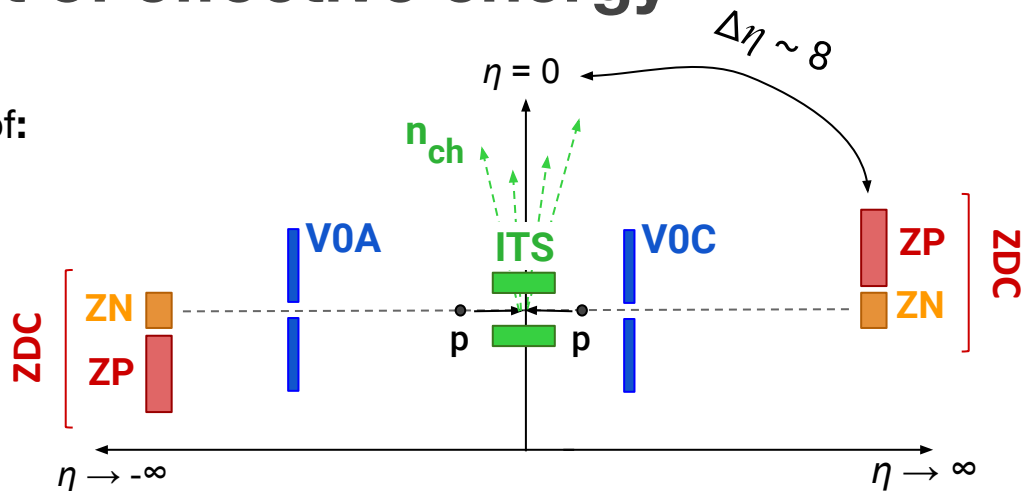
The concept of effective energy

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Independent proxies given the large η separation



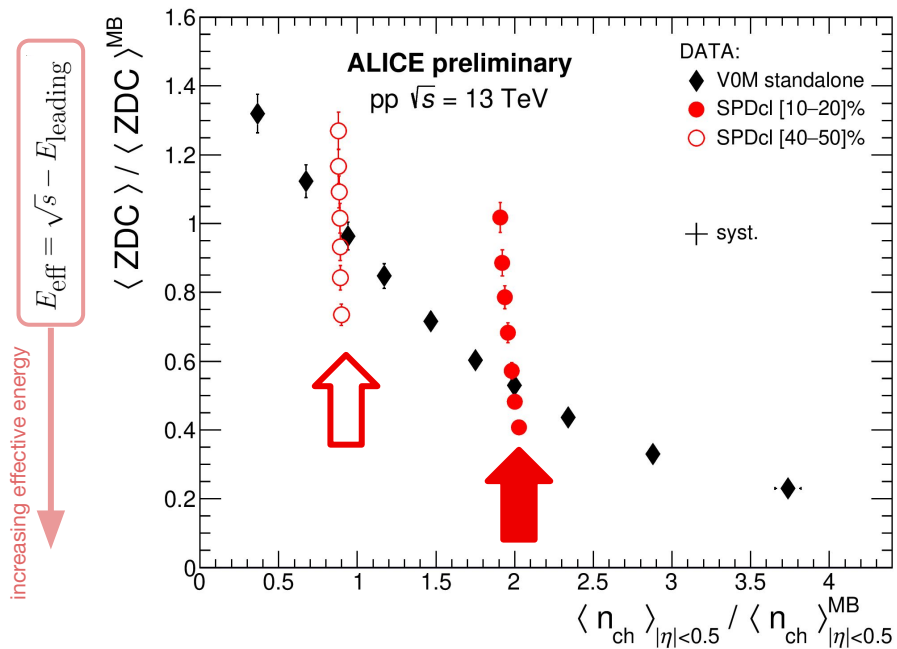
ALICE can measure:

- midrapidity multiplicity (**SPD**)
- leading energy (**ZDC**)
- forward multiplicity (**VOM** = V0A+V0C)

Event classes with a differential approach

The leading energy decreases with increasing particle multiplicity produced at midrapidity [1]

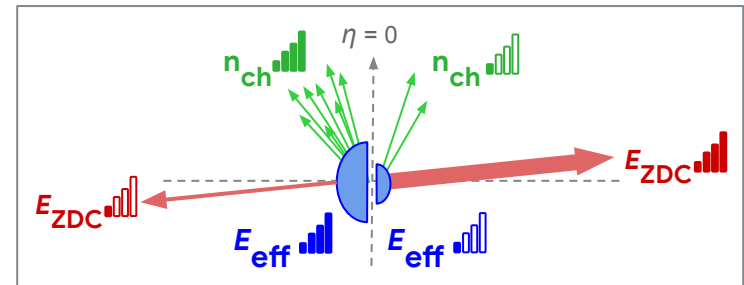
New multi-differential event classes: **similar midrapidity multiplicity** and **different leading energies**



◆ Standalone

● High local multiplicity (midrapidity)

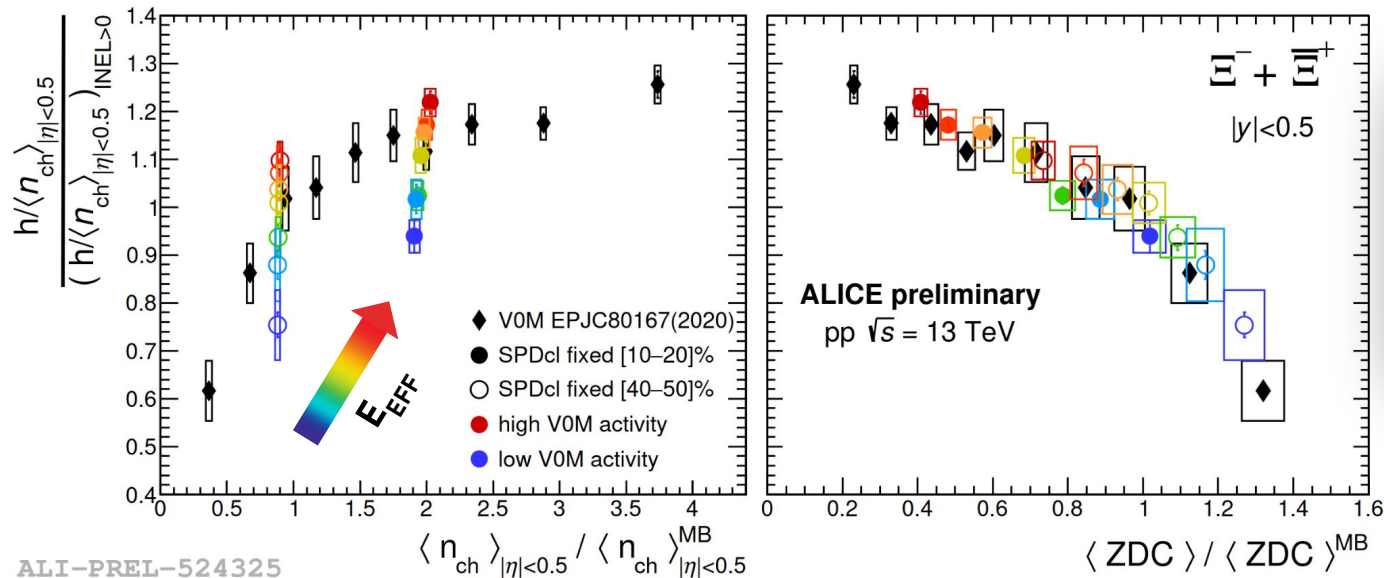
○ Low local multiplicity (midrapidity)



Strangeness production at fixed multiplicity

In events with the same particle multiplicity produced:

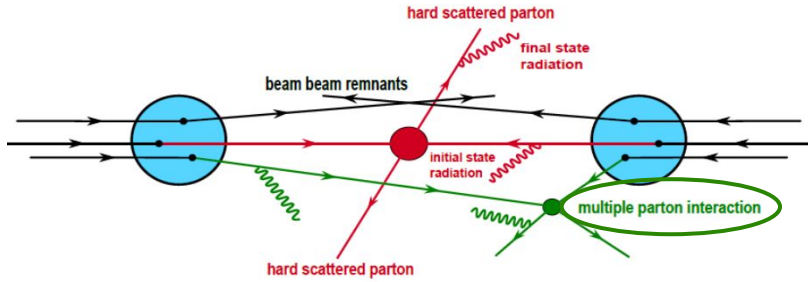
- **increase** in Ξ production per charged particle is observed for **decreasing forward energy** (ZDC)
- universal scaling trends with ZDC energy, **compatible within uncertainties**



Ξ enhancement
observed at fixed
multiplicity,
correlated with the
effective energy

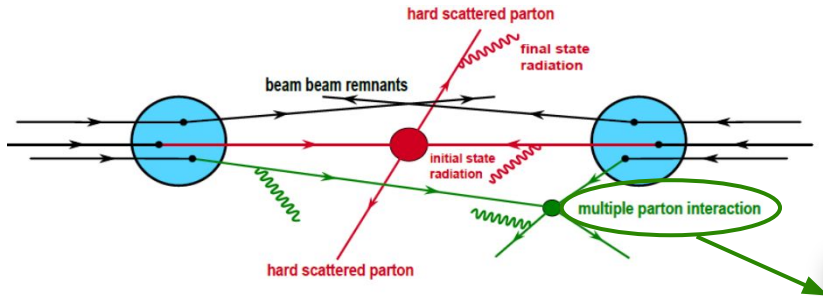
ALI-PREL-524325

Connection to the MPIs in PYTHIA



In PYTHIA, the number of **Multiple Parton Interactions** strongly **influence** the **string hadronization processes** responsible for strange hadron production

Connection to the MPIs in PYTHIA



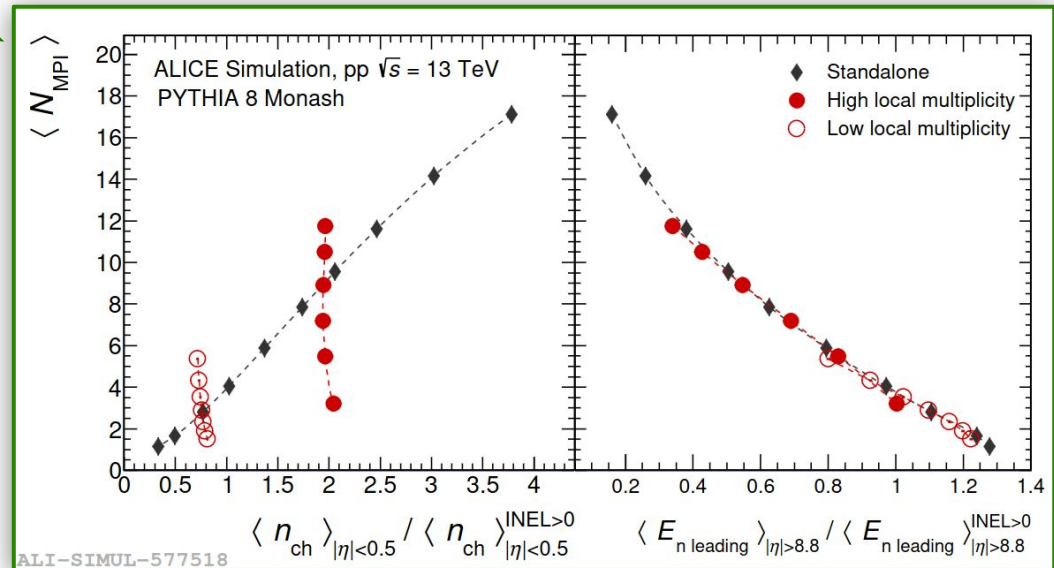
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NEW!

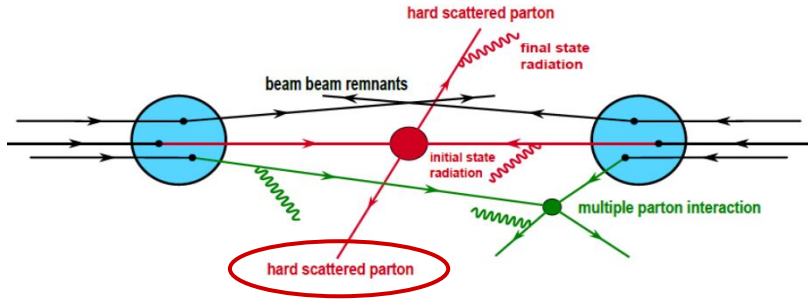
MPIs increase at fixed local multiplicity with decreasing leading energies

Universal dependence with the leading energy, i.e. common for all selections

Leading energy \rightarrow a powerful observable to probe the dependence of particle production on the number of MPIs

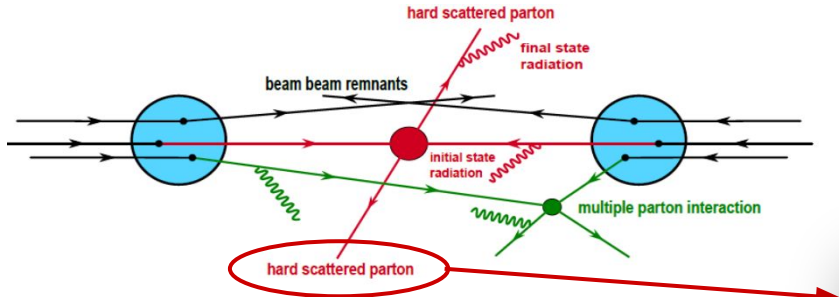


Connection to hard-scattering processes in PYTHIA



The presence of jets at midrapidity is studied in PYTHIA considering the $\langle p_T^\pi \rangle_{|y|<0.5}$, proxy for the p_T of the **hard parton scattering process**

Connection to hard-scattering processes in PYTHIA

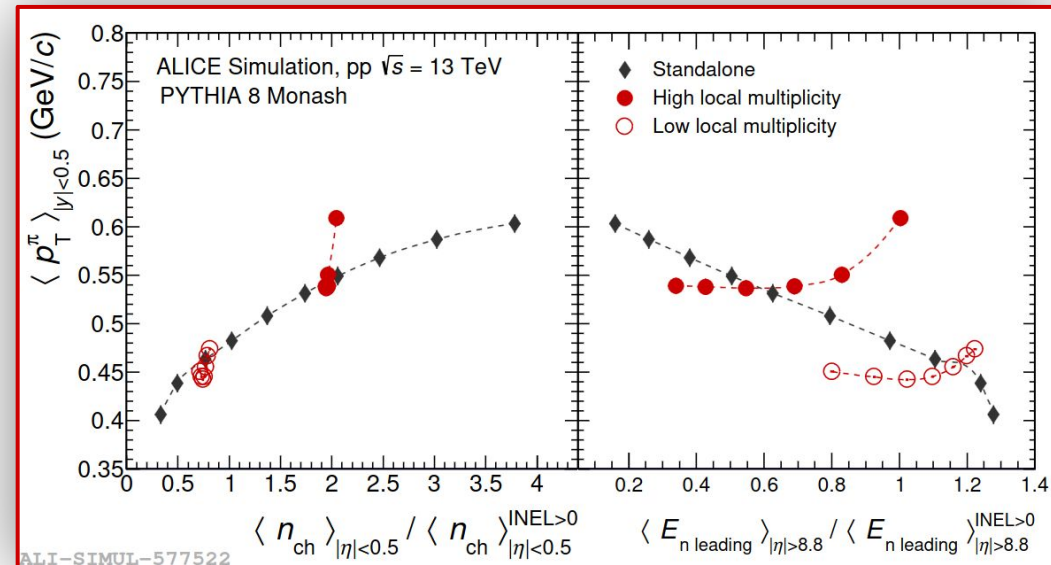


Very mild dependence of $\langle p_T^\pi \rangle_{|y|<0.5}$ with the leading energy

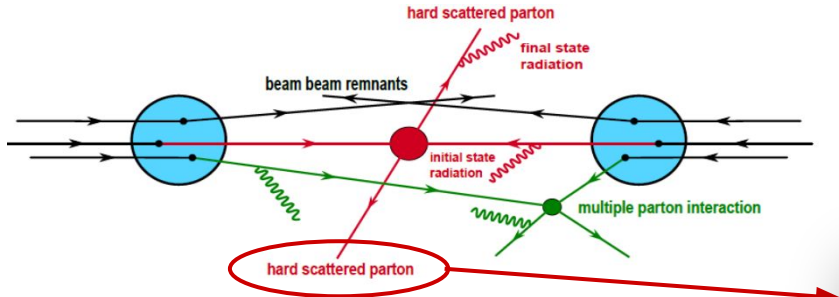
Local phenomena, such as jets at midrapidity, are correlated with local observables, such as the charged-multiplicity

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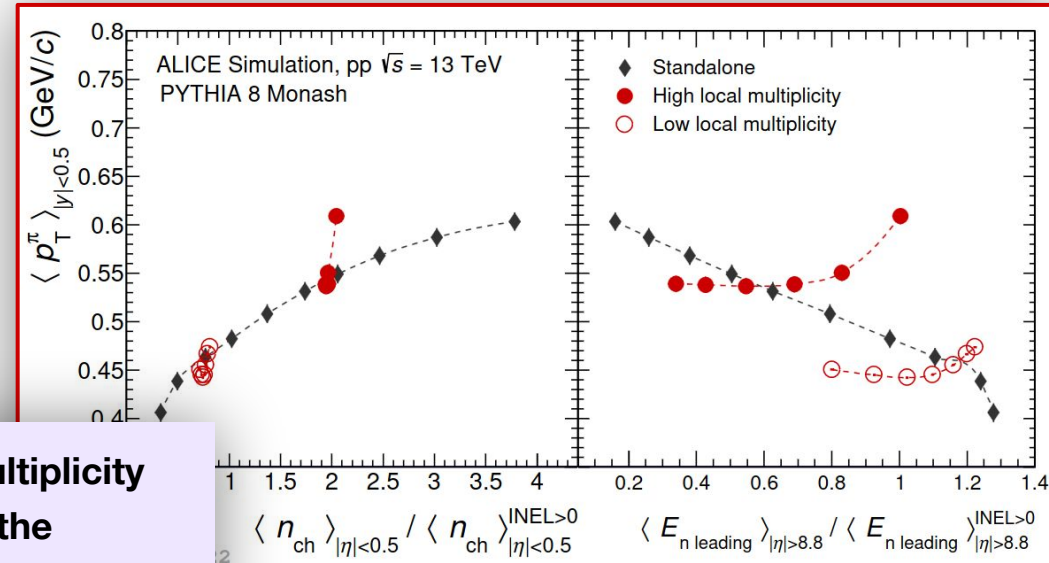


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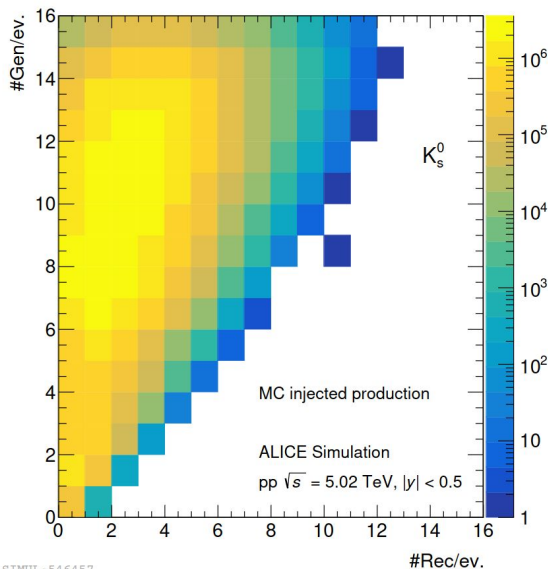


The observed Ξ enhancement at fixed multiplicity is connected to global properties of the event, e.g. MPI, the effective energy

Multiple strange hadron production

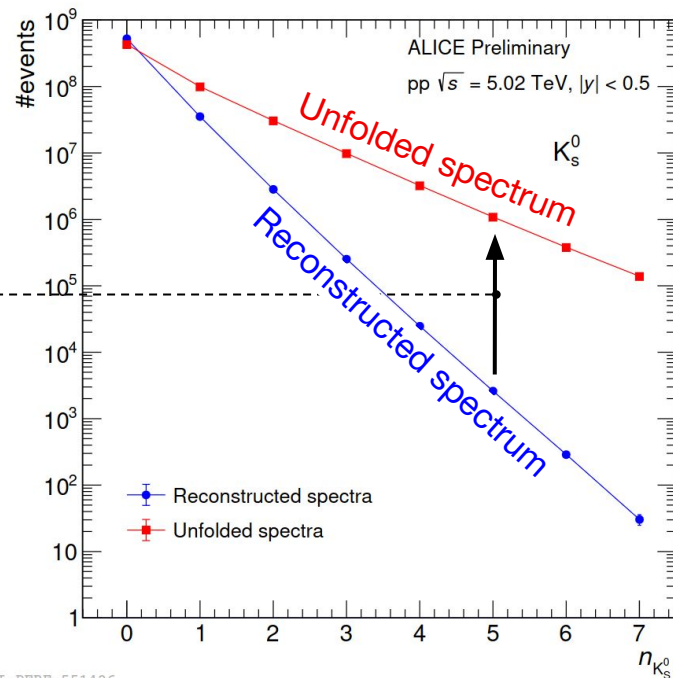
Can strangeness enhancement in small systems arise from **multiple strange hadron production in high multiplicity events?**

ALICE measured the **multiplicity distribution** for the production of (multi-)**strange hadrons** in pp collisions



ALI-SIMUL-546457

The detector response is taken into account via **Bayesian unfolding**



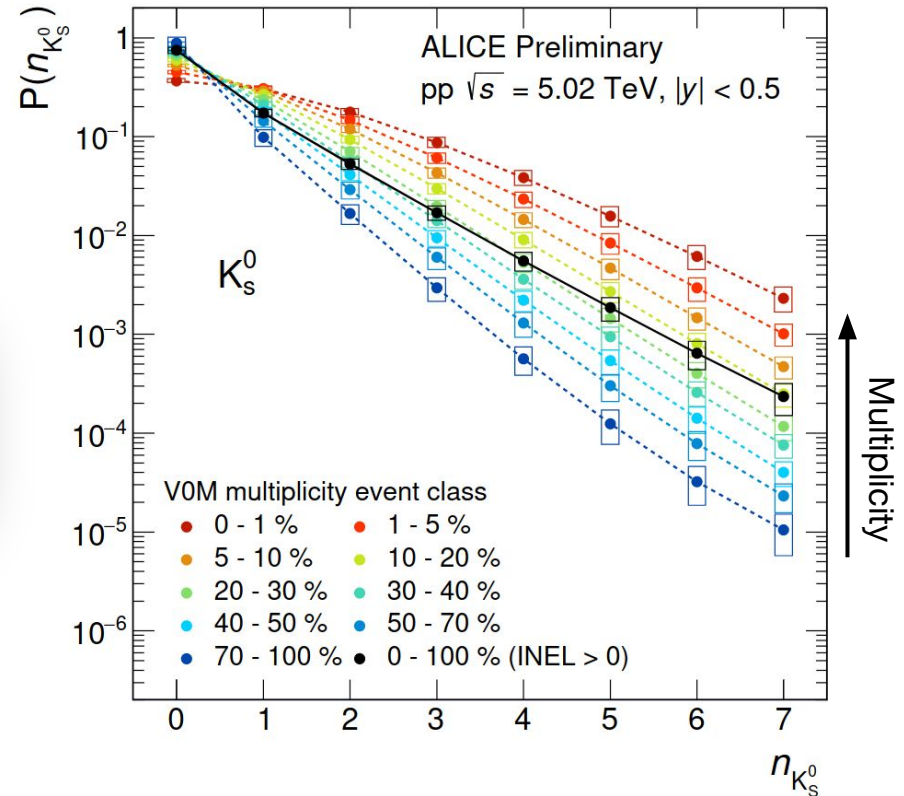
LT-PRR-551426

Strange particle multiplicity distribution

Probability to produce n strange particles of a given species per event $P(n_s)$
→ spanning across large ranges of strange hadron count
→ measured as a function of charged-multiplicity

The **probability** to produce $n \geq 1$ strange hadron per event **increases** with the charged-particle **multiplicity**

Same feature observed for all other (multi-)strange hadrons (Λ , Ξ , Ω)



Multiple strange hadron production yields

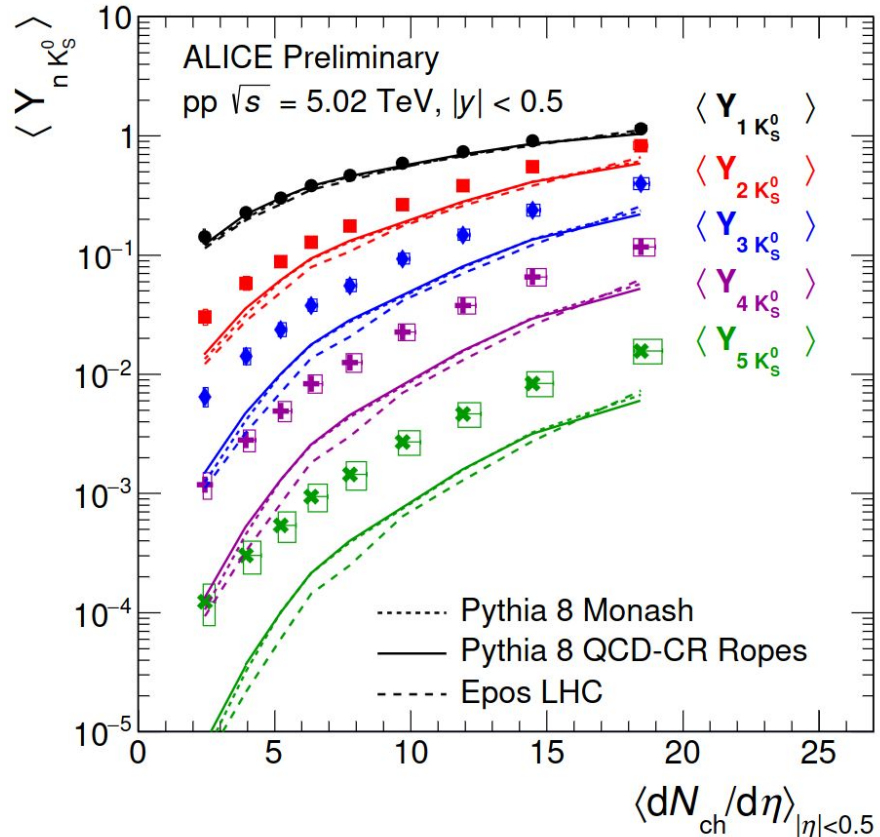
Average production yield of 1, 2, 3, ... particles/event is calculated as:

$$\langle Y_{k-part} \rangle = \sum_{n=k}^{\infty} \frac{n!}{k!(n-k)!} P(n)$$

The n-particle yields increase **more than linearly** with the charged-particle multiplicity

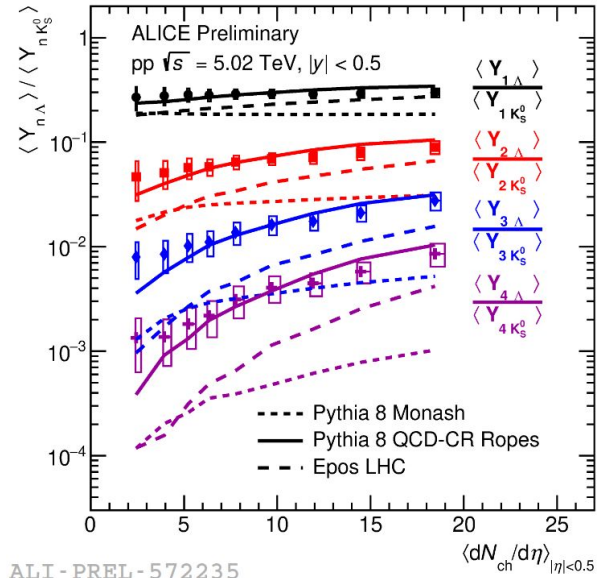
Model comparison:

the increase with multiplicity is qualitatively reproduced by PYTHIA Monash, QCD-CR Ropes and EPOS LHC, but all underestimate the n-particle yields for $n > 1$



Yield ratios with $\Delta S = 0$

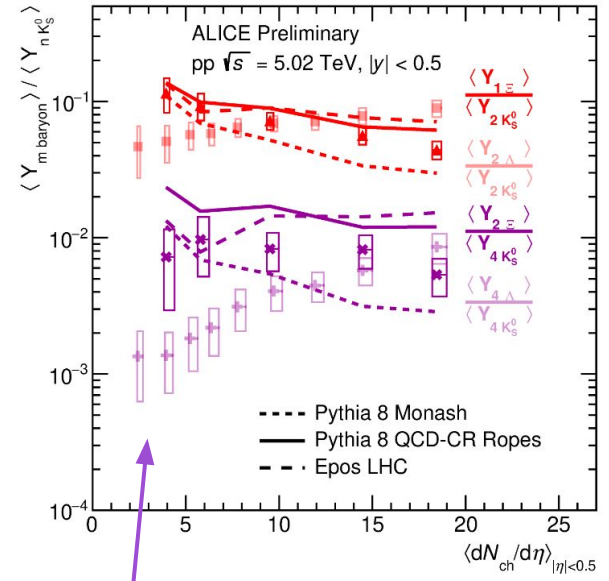
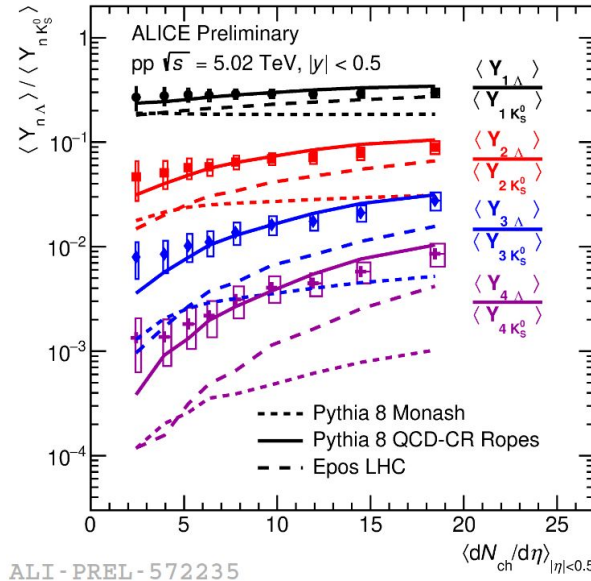
$\langle Y_{n,\Lambda} \rangle / \langle Y_{n,K_S^0} \rangle$
 $\Delta S = 0 \rightarrow$ **increase vs multiplicity** for multiple hadron production



ALI-PREL-572235

Yield ratios with $\Delta S = 0$

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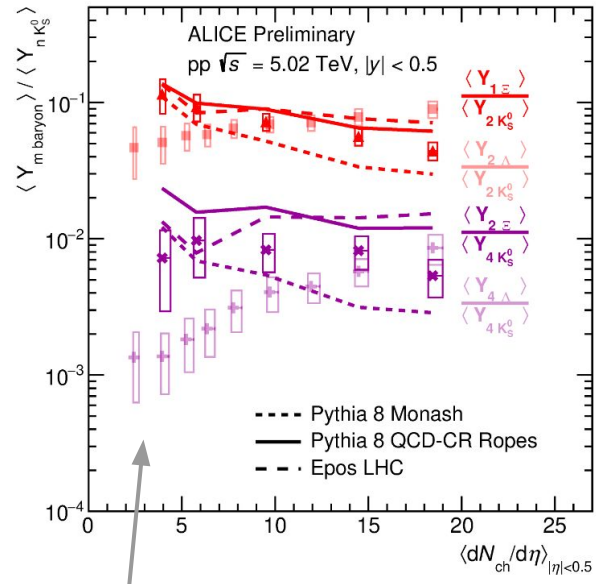
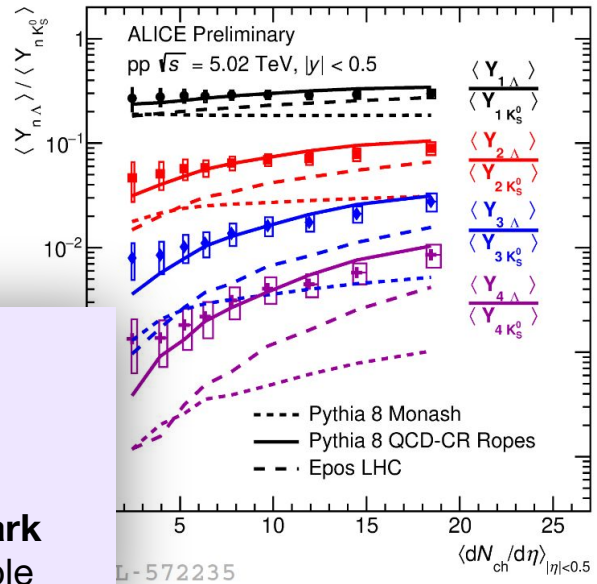
$\langle Y_{m \text{ baryon}} \rangle / \langle Y_{n K_S^0} \rangle$
 $\Delta S = 0 \rightarrow$ test hadron production at fixed number of strange quarks and different light quark content
Decreasing trend with multiplicity

Yield ratios with $\Delta S = 0$

$\langle Y_{n,\Lambda} \rangle / \langle Y_{n,K_S^0} \rangle$
 $\Delta S = 0 \rightarrow$ **increase vs multiplicity** for multiple hadron production

Yield ratios with $\Delta S = 0$ show more than “simple” **strange baryon enhancement**

The **probability of light (u,d) quark (re-)combination** also plays a role

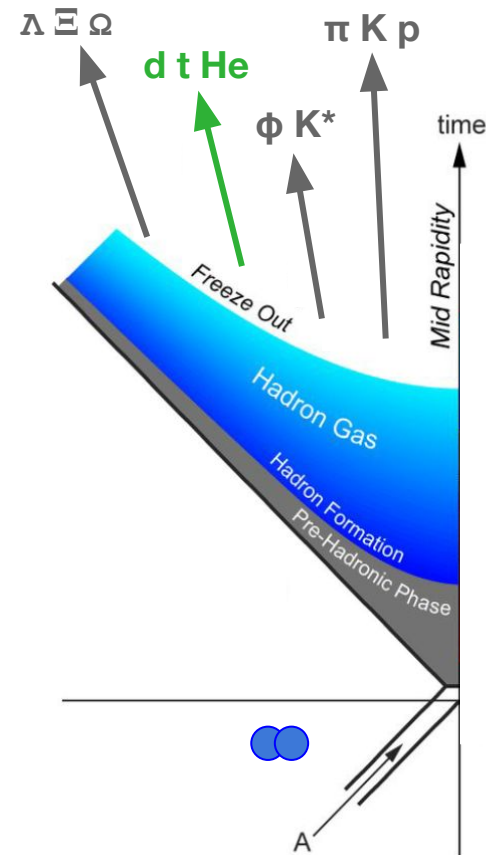


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 $\Delta S = 0 \rightarrow$ test hadron production at fixed number of strange quarks and different light quark content

Decreasing trend with multiplicity

Model comparison:
 All trends are rather well reproduced by Pythia 8 QCD-CR Ropes, which includes color reconnection effects

(Anti)nuclei measurements with ALICE and relevance for astrophysics



LHC: an (anti)nucleus factory

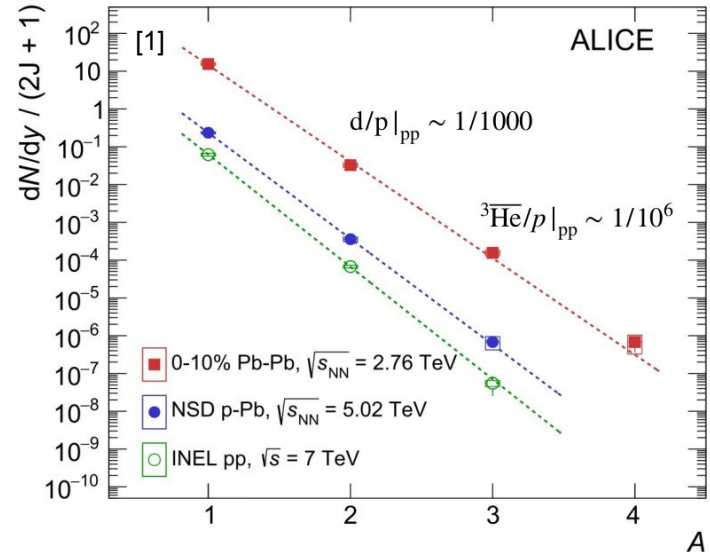
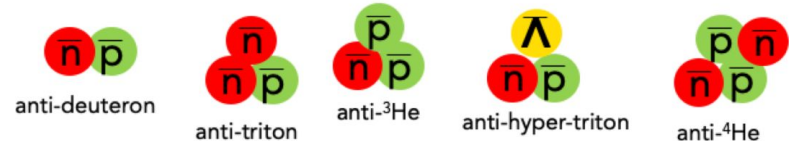
Years of light (anti)nuclei measurements at the LHC → improved understanding of their production mechanisms, but still **not fully understood**

Ideal conditions at LHC energies:
antimatter / matter ~ 1

Rarely produced in high-energy collisions
→ require large integrated luminosity

Two main classes of models to describe light (anti)nuclei production

- **statistical hadronisation** model (SHM)
- **coalescence** model



[1] [Acharya et al., Phys. Lett. B 800, 135043 \(2020\)](#)

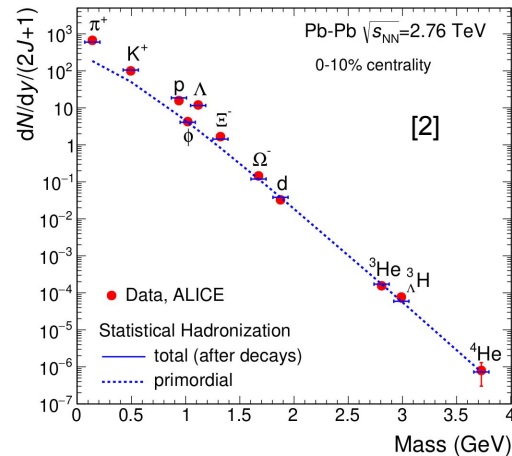
Modelling the production of (anti)nuclei

Statistical models (SHM)

Hadron production from a system in **thermal and hadrochemical equilibrium** [1,2]

$$dN/dy \propto V \exp\left(-\frac{m}{T_{ch}}\right)$$

Extension to pp in the canonical ensemble → **CSM**



Coalescence models

Nucleons close in the phase space at the freeze-out can form a nucleus via **coalescence** [1,2]

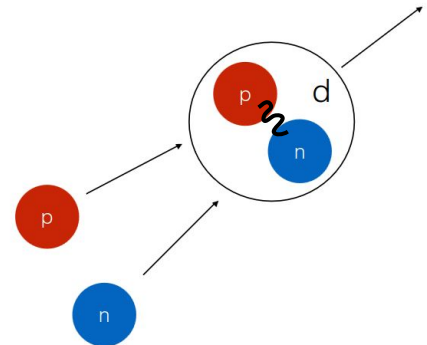
Coalescence parameter B_A :

- is related to the coalescence probability
- depends on the source size and on the nucleus wave function

$$B_A = \frac{E_A \frac{d^3 N_A}{d^3 p_A}}{\left(E_p \frac{d^3 N_p}{d^3 p_p}\right)^A}$$

where:

- A is the mass number of the nucleus
- $p_p = p_A / A$



[1] [V. Vovchenko et al., PLB 785 \(2018\) 171-17](#)
 [2] [Andronic, A. et al., Nature 561, 321–330 \(2018\)](#)

[1] [J. I. Kapusta, Phys.Rev. C21, 1301 \(1980\)](#)
 [2] [Scheibl, Heinz, Phys.Rev.C59:1585-1602 \(1999\)](#)

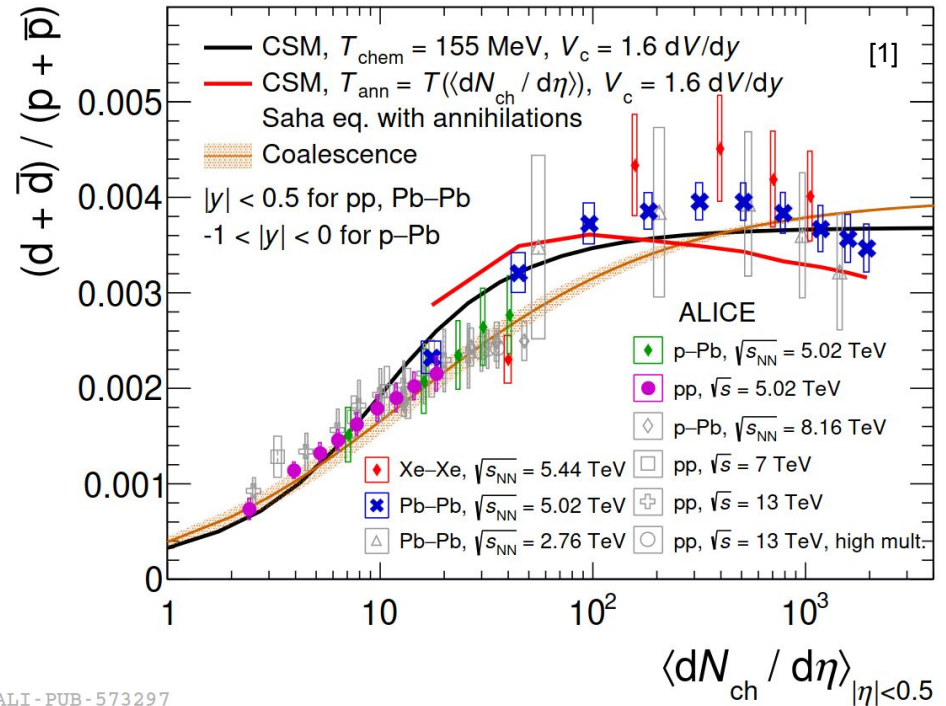
Yield ratios: d/p

d/p ratio evolves smoothly with multiplicity

→ dependence on the system size

Both models are able to describe the data:

- CSM: canonical suppression
- Coalescence model: interplay between source size and nuclear size



ALI - PUB - 573297

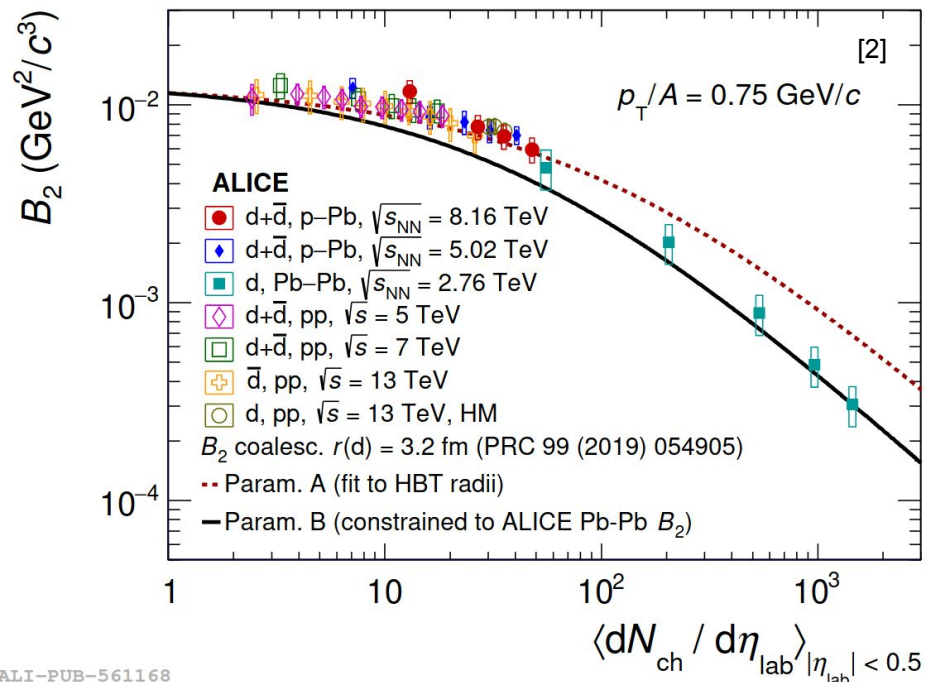
[1] [arxiv:2405.19826](https://arxiv.org/abs/2405.19826)

Coalescence probability B_A

$$B_A = \frac{E_A \frac{d^3 N_A}{d^3 p_A}}{\left(E_P \frac{d^3 N_P}{d^3 p_P} \right)^A}$$

For high multiplicity Pb–Pb significant B_2 drop

→ can be explained within the coalescence model due to the **increase in the source size**



ALI-PUB-561168

[1] [arxiv:2405.19826](https://arxiv.org/abs/2405.19826)

[2] [Phys. Lett. B 846 \(2023\) 137795](https://arxiv.org/abs/2307.13779)

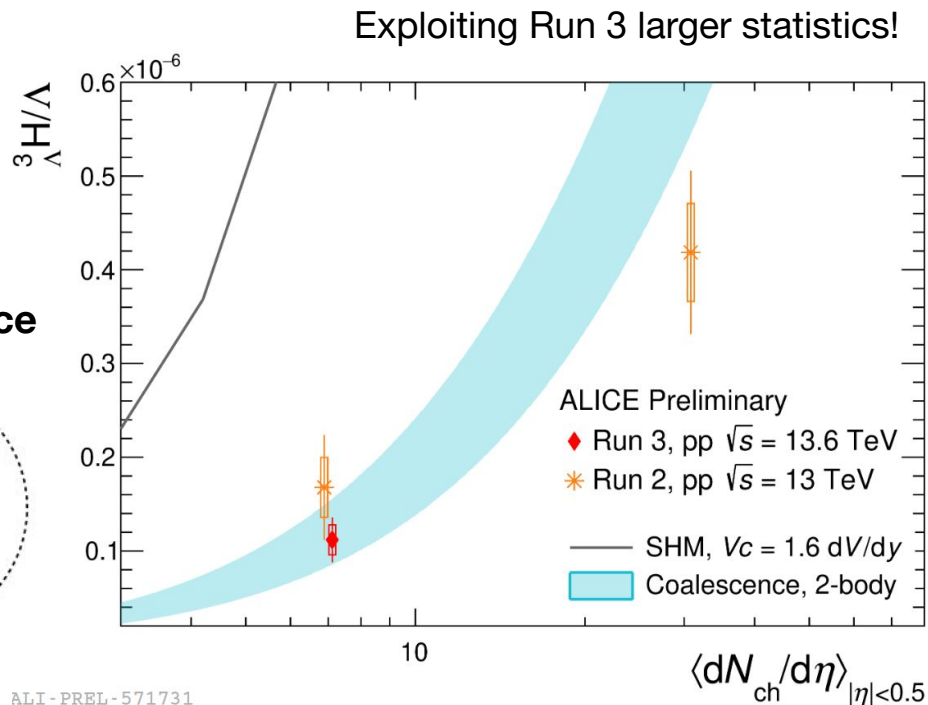
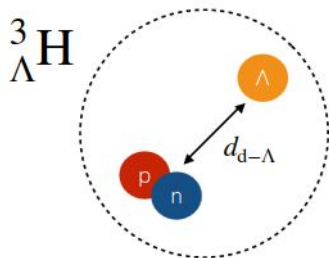
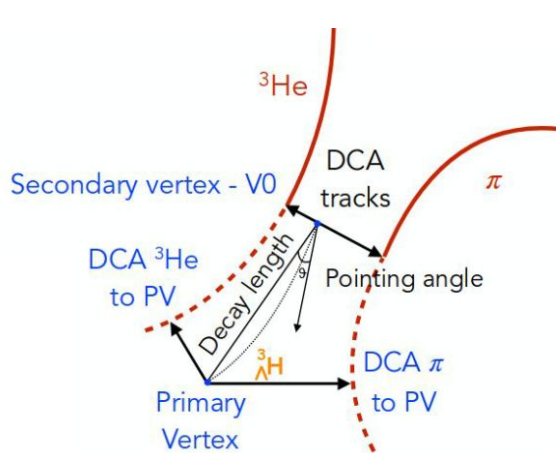
Hypertriton production

Hypertriton has a larger radius wrt deuteron:

$$r({}^3_{\Lambda}\text{H}) = 10.79 \text{ fm}^{(1)}, r(d) = 1.96 \text{ fm}$$

SHM and coalescence predictions significantly diverge at low multiplicity

→ ${}^3_{\Lambda}\text{H}/\Lambda$ in pp collisions favours coalescence



(1) [F. Hildenbrand and H.-W. Hammer, Phys. Rev. C 100, 034002](#) 33

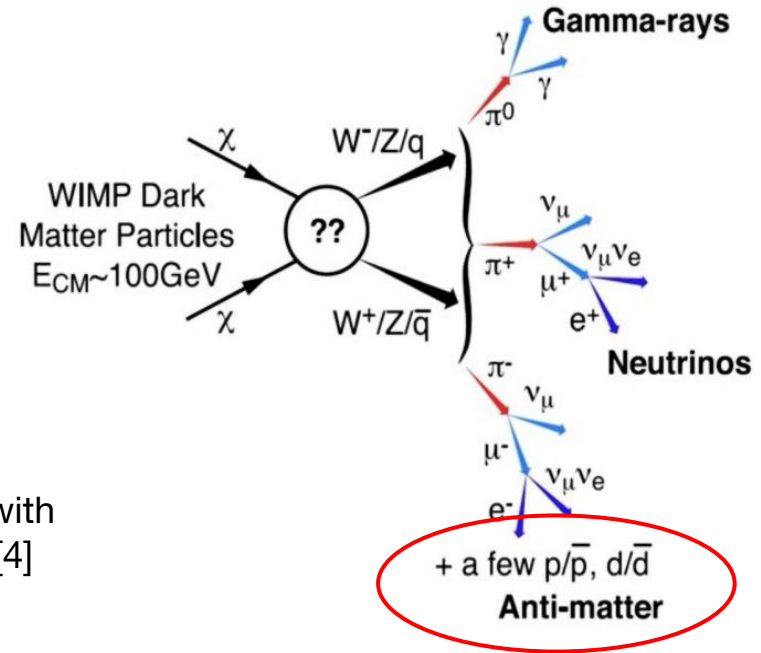
Light antinuclei as probes for Dark Matter

Cosmic ray antideuteron and antihelium nuclei have been suggested as possible **probes for Dark Matter (DM) WIMPs** (mass \sim few GeV – few TeV) [1-3]

- Anti-p and anti-n can be produced by WIMP **annihilation into SM channels**
- anti-deuterons and anti- ^3He are produced via **coalescence** of anti-nucleons

Subject for indirect DM searches with space-based experiments such as **AMS-02** (ongoing)

Since ~ 10 years AMS has collected > 220 billions CRs with few candidate events compatible with antihelium mass [4]



[1] [M. Korsmeier, et al., Phys. Rev. D 97, 103011 \(2018\)](#)

[3] [E. Carlson, et al., Phys. Rev. D 89, 076005 \(2014\)](#)

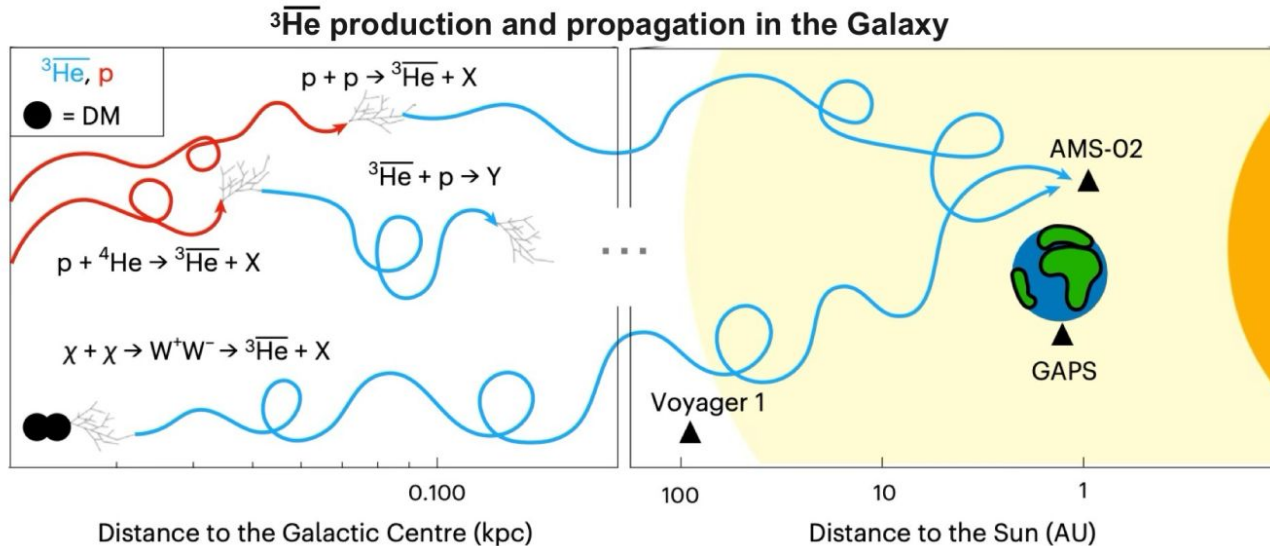
[2] [M. Cirelli, et al., JHEP 08, 009 \(2014\)](#)

[4] [Ting, S., CERN Colloquium \(2016\)](#)

(Anti)nuclei in cosmic rays

The largest fractions of primary CR are protons and helium, **no antinuclei as primary CR**

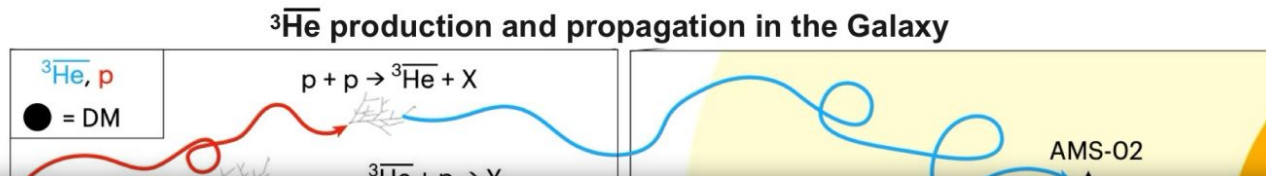
Secondary anti-p, anti-d, anti- ^3He produced by **interaction of primary CR** with the InterStellar Medium ISM (pp, p-He, ...) constitute a **background** for the DM signal



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From ALICE:

- **Antinuclei formation mechanism** (coalescence) → constraint with LHC data
- **Measured antinucleus inelastic cross section** to account for absorption by ISM

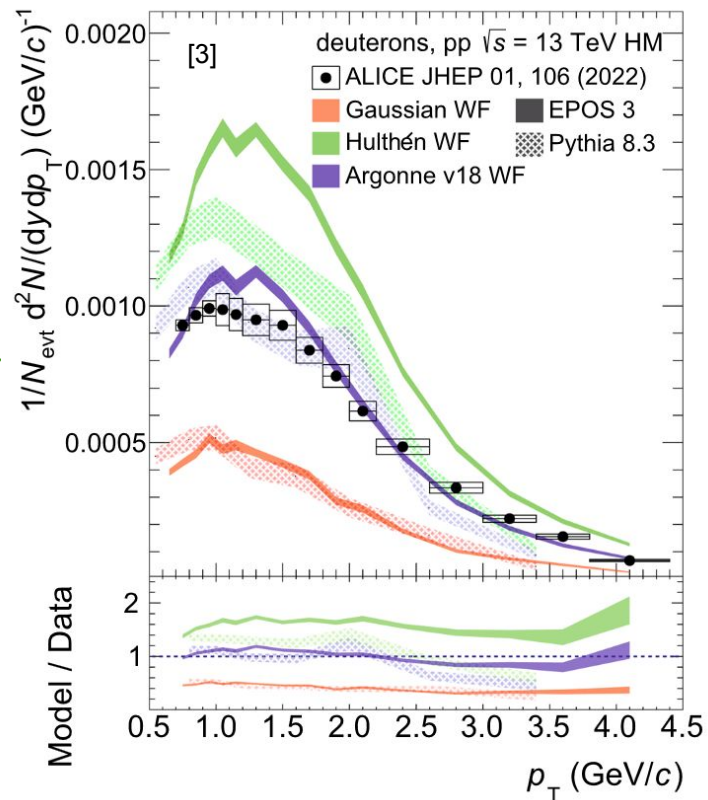
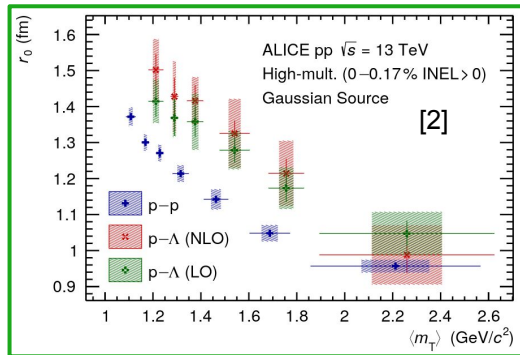
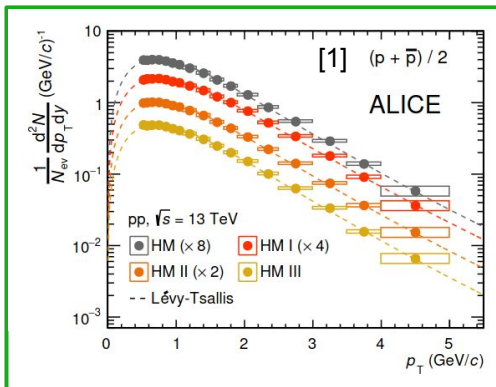
ALICE experimental input to coalescence modelling

A recent study shows that ALICE measurements of:

- proton production **yields**
- proton **source radius**

in high-multiplicity pp collisions, allow for **prediction** of the deuteron spectrum via **coalescence** with **no free parameters**

Wigner-formalism based coalescence afterburner



[1] [JHEP 01 \(2022\) 106](#)

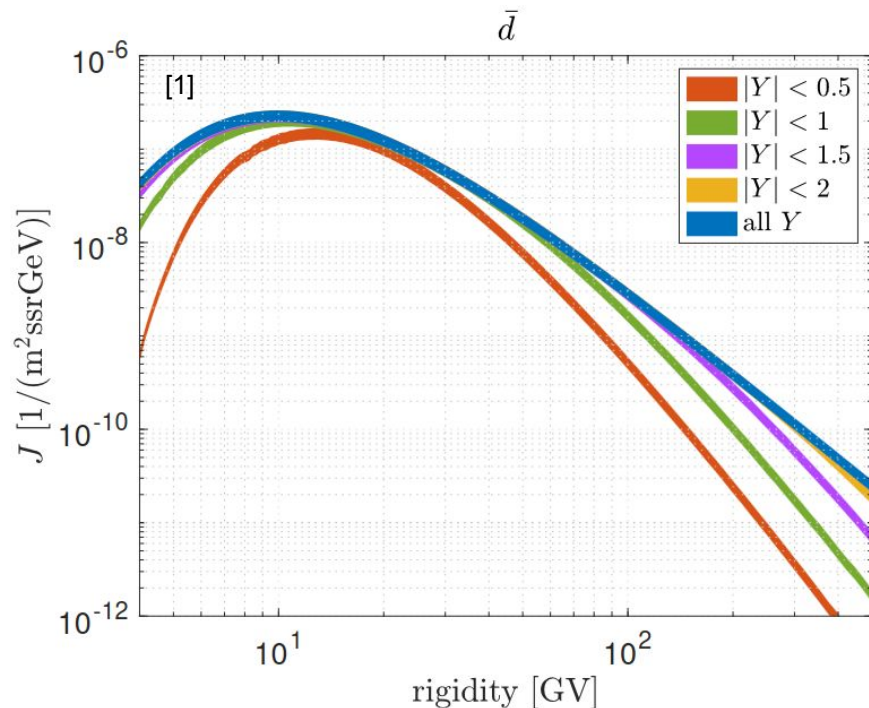
[2] [Phys. Lett. B 811 \(2020\) 135849](#)

[3] [Mahlein, M., Eur. Phys. J. C 83, 804 \(2023\)](#) 37

Rapidity dependence of coalescence

Model predictions based on ALICE measurements are used as input to calculate **antideuteron flux from cosmic rays** [1]

But typically ALICE measurements cover **midrapidity** ($|y| < 0.5$) while astrophysical models extrapolate to the **forward** region



[1] [K. Blum, Phys. Rev. C, 109, L031904](#)

Rapidity dependence of coalescence

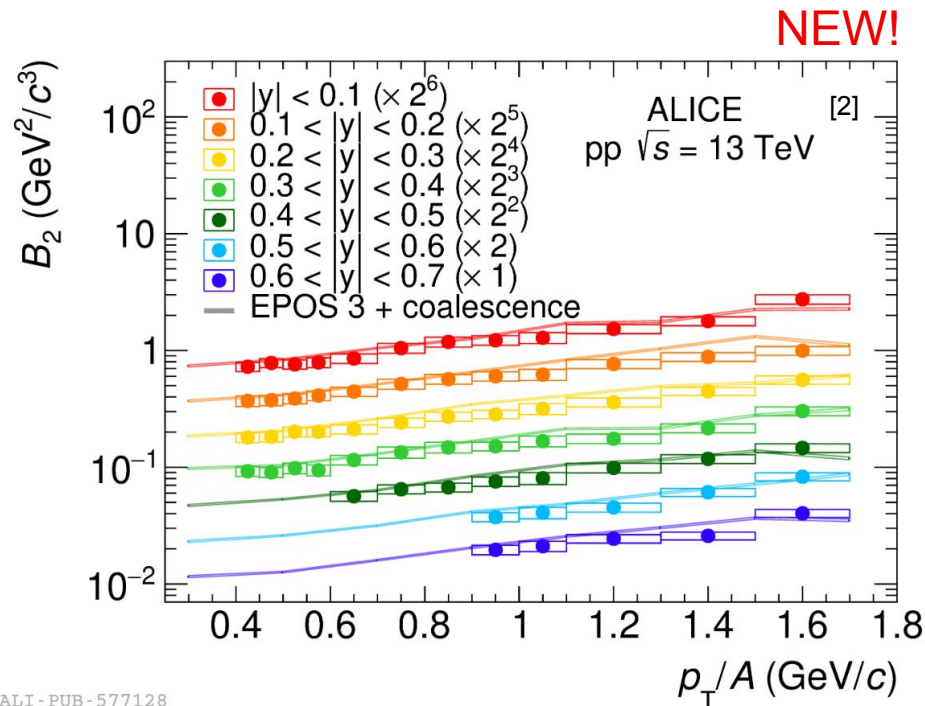
$$B_A = \frac{E_A \frac{d^3N_A}{d^3p_A}}{\left(E_p \frac{d^3N_p}{d^3p_p}\right)^A}$$

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The current acceptance of the ALICE detector allows to extend the measurement of antideuteron **up to $|\Delta y| < 0.7$** [2]

→ no strong dependence on rapidity is observed



ALI-PUB-577128

[1] [K. Blum, Phys. Rev. C. 109, L031904](#)

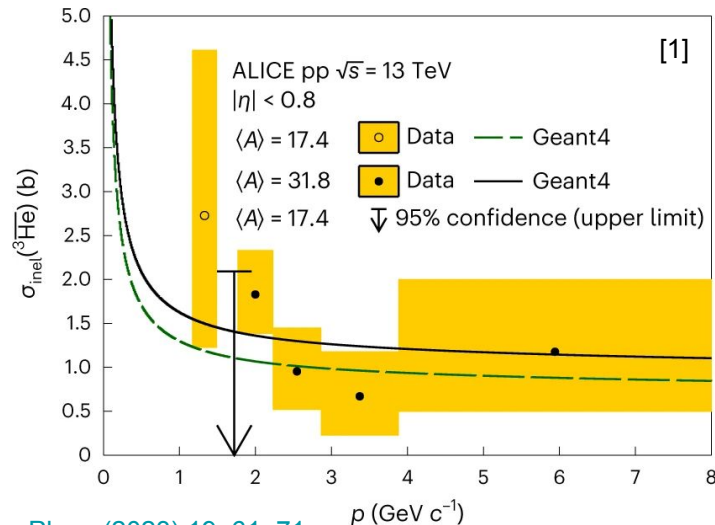
[2] [arxiv:2407.10527](#)

Measurement of the inelastic ${}^3\overline{\text{He}}$ cross section

ALICE measured the **inelastic cross section** for antinuclei using the LHC as antimatter factory and the experimental apparatus as a target → **significant impact on ${}^3\overline{\text{He}}$ propagation in space**

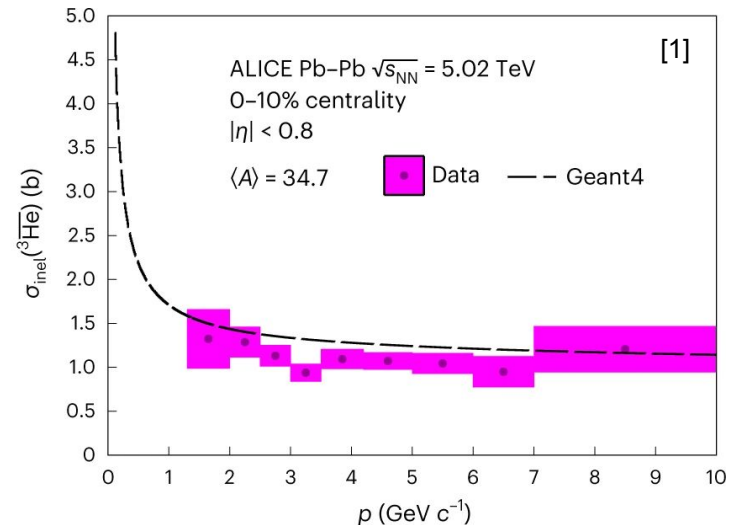
Antimatter-to-matter ratio:

measurement of reconstructed ${}^3\overline{\text{He}}/{}^3\text{He}$ ratio and compare to MC expectations



TOF/TPC-matching ratio:

measurement of reconstructed ${}^3\overline{\text{He}}(\text{TOF})/{}^3\overline{\text{He}}(\text{TPC})$ ratio and compare to MC expectations



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In the last years ALICE tried to push our understanding of strangeness and nuclei production mechanisms in high-energy hadronic collisions

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- First measurement of **multiple strange hadron production** and yields ratio with $\Delta S = 0$

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Important insights from nuclei production and applications to Astrophysics:

- LHC can be used as **antimatter factory** to study the production of light (anti)nuclei
- Yield ratios of (anti)nuclei and hypernuclei to **constrain formation mechanisms** with applications to indirect Dark Matter searches
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More to come with LHC Run 3 increased statistics!