#### Non-Perturbative Yang-Mills Beyond One-Loop Order

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#### Who? Seth Grable

From? Department of Physics University of Colorado Boulder

When? August 2024



We want to calulate YM, QCD, and non-abelian theories!



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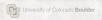
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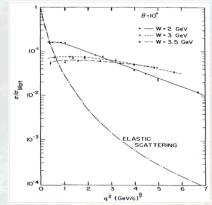
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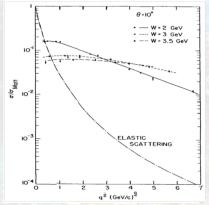






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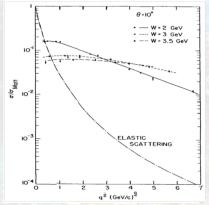
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Associated physics: Confinement, Bjorken scaling, QCD foundations.

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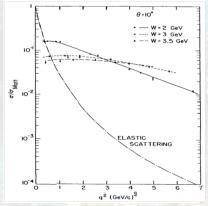


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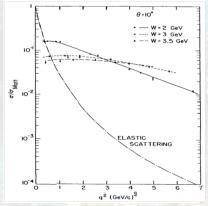
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Perturbation theory at high energies



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What works:

- Perturbation theory at high energies
- Lattice calculations at zero chemical potential



#### YM is highly non-linear, and strongly interacting



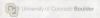
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YM is highly non-linear, and strongly interacting What doesn't work:



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  What doesn't work:
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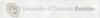


We need non-perturbative analytic methods for non-abelian gauge theories!



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 We have them for scalar QFT, Ads/CFT, and condensed matter (SYK).



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Background-field theory calculations

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- Background-field theory calculations
- Hubbard-Stratonovich transformation
- Laplace's method.



Main Idea:



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$$Z = \int \mathcal{D}A e^{-\frac{1}{4g_0^2} \int_{x} (\mathcal{F}^{a}_{\mu\nu})^2}.$$
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background field + fluctuations.

$$A^{a}_{\mu}(x) = B^{a}_{\mu}(x) + a^{a}_{\mu}(x)$$
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This allows for an expansion around  $B^a_{\mu}$  in terms of  $a^a_{\mu}$ .



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 $S_0$  gives the quadratic contributions in  $a^a_\mu$ ,  $S_I$  gives the cubic and quartic terms.



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- $S_0$  gives the quadratic contributions in  $a^a_{\mu}$ ,  $S_I$  gives the cubic and quartic terms.
- **The general strategy: reduce the complexity of**  $S_1$ .
  - $S_0$ , quadratic contributions, are calculable (Gaussian).



How can we simplify  $S_l$  (cubic and quartic terms)?



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$$S_{I} = \int_{x} g_{0}(D_{\mu}a_{\nu}^{a})f^{abc}a_{\mu}^{b}a_{\nu}^{c} + \frac{g_{0}^{2}}{4}(f^{abc}a_{\mu}^{b}a_{\nu}^{c})^{2}.$$
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- However, I discovered that IR divergence perfectly cancels for  $n_f = 12$ .

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- **Elements of Confinement for QCD with Twelve** Massless Quarks: arXiv:2310.12203





#### **Stable background field configurations** $\overline{B}$



 $n_f = 12$  Results

#### Stable background field configurations $\overline{B}$

$$\bar{B}_{\sigma_8} \approx 4.27 \Lambda_{\overline{MS}}^2$$

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(5)

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Romatschke and I further found a vanishing Polyakov loop expectation value under a critical temperature  $T_c \approx .81 \Lambda_{\overline{MS}}$ , and a non-zero expectation value for  $T > T_c$ .

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- Romatschke and I further found a vanishing Polyakov loop expectation value under a critical temperature  $T_c \approx .81\Lambda_{\overline{MS}}$ , and a non-zero expectation value for  $T > T_c$ .
  - We see this as a validation and proof of concept for applying the BGFM to non-abelian theories.



#### $n_f = 12$ is a special case.



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- General strategy: reduce *S*<sub>1</sub> complexity.
- Technical method: Hubbard-Stratonovich transformation



Simple Example:



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 $HS \rightarrow$ 

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x can be integrated out: It's Gaussian!



# Gap Equation from Laplace's Method

Left with some integral that looks like:



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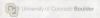
### Gap Equation from Laplace's Method

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  - Energy gap in the electronic spectrum of a superconductor.



#### What to do next?



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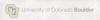
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**Calculate Pure Yang-Mills in** d = 3 + 1 at T = 0



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- What to do next?
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- The effective action is gauge invariant.



Zeta-function regularization, and running coupling renormalization allows for an exact continuum limit calculation.



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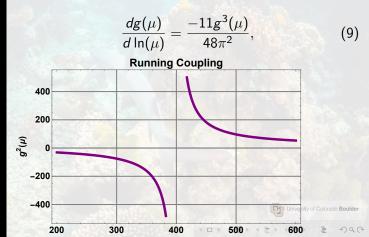
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$$\frac{dg(\mu)}{d\ln(\mu)} = \frac{-11g^3(\mu)}{48\pi^2},$$
(9)



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Independent gap equations are given for the background and auxiliary fields.



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 Functional dependence on Δ: -ln Z(Δ, B)

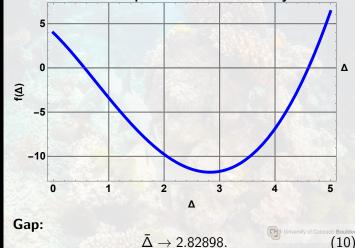


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**Functional dependence on**  $\Delta$ :  $-\ln Z(\Delta, B)$ 

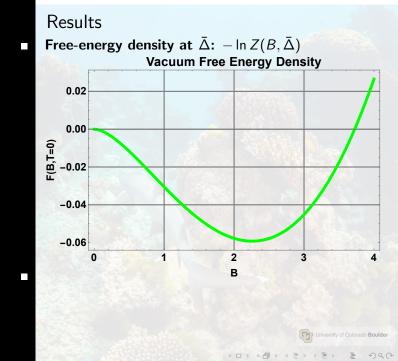
Functional Dependence of the Auxillary field  $\Delta$ 

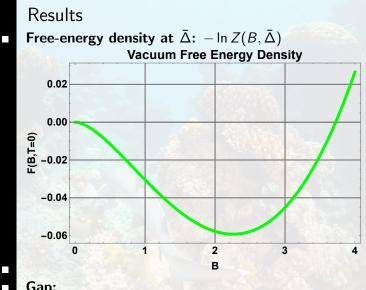




# Results **Free-energy density at** $\bar{\Delta}$ : $-\ln Z(B, \bar{\Delta})$







Gap:

$$\bar{B} \rightarrow \frac{2.25885\Lambda_{\rm YM}^2}{\lambda'}$$
,  $\bar{B} \rightarrow 0$ , University of Colorado Boundary



#### Vacuum Pressure:



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#### Vacuum Pressure:

$$\frac{\ln Z}{\beta V} = 6 \times 0.0592377 \Lambda_{\rm YM}^4.$$
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Lattice calculations give:  $m_{Glue}^a \approx 3 \sim 4\Lambda_{QCD}$ 

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Non-Perturbative Yang-Mills Beyond One-Loop Order: arXiv:2407.13042



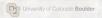
#### Calculating the propagators



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Calculating the propagatorsConfinement/Polyakov loop



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Calculating the propagators
 Confinement/Polyakov loop
 Finite temp YM



Calculating the propagators
 Confinement/Polyakov loop
 Finite temp YM
 Full QCD



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- Calculating the propagators
- Confinement/Polyakov loop
- Finite temp YM
- Full QCD
- Finite temperature and chemical potential QCD

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# THANK YOU FOR LISTENING!THANK YOU UNIVERSITY OF ADELAIDE!



# THANK YOU FOR LISTENING! THANK YOU UNIVERSITY OF ADELAIDE! Non-Perturbative Yang-Mills Beyond One-Loop Order: arXiv:2407.13042



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## Fun!



Background fields:



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't Hooft, 1973 Computation of the quantum effects due to a four-dimensional pseudoparticle

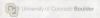
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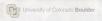
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### Glueballs:



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#### Glueballs:

Berg, 1980 Plaquette-plaquette correlations in the su (2) lattice gauge theory



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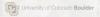
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$$\mathcal{F}^{a}_{\mu\nu} = F^{a}_{\mu\nu} + D^{ac}_{\mu} \ a^{c}_{\nu}(x) - D^{ac}_{\nu} \ a^{c}_{\mu}(x) + f^{abc} a^{b}_{\mu}(x) a^{c}_{\nu}(x)$$
(14)



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where:

$$D_{\mu}^{ac} = \partial_{\mu}\delta^{ac} + f^{abc}B_{\mu}^{b}(x) \tag{15}$$

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Square  $\mathcal{F}^{a}_{\mu\nu}$  to get:

$$Z = \int B \int \mathcal{D}a \mathcal{D}\bar{c} \mathcal{D}c e^{-\int_{x} \frac{1}{4g_{0}^{2}} (F_{\mu\nu}^{a})^{2} - S_{0} - S_{I}}$$
(16)

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#### With:

$$S_{0} = \int_{x} \frac{1}{2} a_{\mu}^{a} \Big[ - (D^{2})^{ac} \delta_{\mu\nu} + 2F_{\mu\nu}^{c} f^{abc} \Big] a_{\nu}^{c} \qquad (17)$$

$$S_{I} = \int_{X} g_{0}(D_{\mu}a_{\nu}^{a})f^{abc}a_{\mu}^{b}a_{\nu}^{c} + \frac{g_{0}^{2}}{4}(f^{abc}a_{\mu}^{b}a_{\nu}^{c})^{2}.$$
 (18)



## Gauge Fixing and Local Gauge Invariance

Gauge Fixing condition:



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### Gauge Fixing and Local Gauge Invariance

Gauge Fixing condition:

 $D_{\mu}^{ac}a_{\mu}^{c}(x)=0$ 



### Gauge Fixing and Local Gauge Invariance

Gauge Fixing condition:

$$D_{\mu}^{ac}a_{\mu}^{c}(x)=0$$

After gauge fixing  $S_0$  is invariant to:

$$\begin{aligned} a^{a}_{\mu}(x) &\to a^{a}_{\mu}(x) - f^{abc}\beta^{b}(x)a^{c}_{\mu}(x) \\ B^{a}_{\mu}(x) &\to B^{a}_{\mu}(x) + B^{a}_{\mu}(x)D_{\mu}\beta^{a}(x) \\ c^{a}(x) &\to c^{a}(x) - f^{abc}\beta^{b}(x)c^{c}(x). \end{aligned}$$
(19)

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### Gauge Fixing and Local Gauge Invariance

Gauge Fixing condition:

$$D^{ac}_{\mu}a^{c}_{\mu}(x)=0$$

After gauge fixing  $S_0$  is invariant to:

$$\begin{aligned} a^{a}_{\mu}(x) &\to a^{a}_{\mu}(x) - f^{abc}\beta^{b}(x)a^{c}_{\mu}(x) \\ B^{a}_{\mu}(x) &\to B^{a}_{\mu}(x) + B^{a}_{\mu}(x)D_{\mu}\beta^{a}(x) \\ c^{a}(x) &\to c^{a}(x) - f^{abc}\beta^{b}(x)c^{c}(x). \end{aligned}$$
(19)

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The effective action remains invariant to this transformation after applying a HS transformation.



#### The Eigenspectrum of $S_0$

$$D_{\mu}^{ac} = \partial_{\mu}\delta^{ac} + \frac{i}{2}A^{ac}F_{\mu\nu}x_{\nu}$$
(20)

where A is the hermitian matrix consisting of the sum generators in the adjoint representation.



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#### The Eigenspectrum of $S_0$

$$D_{\mu}^{ac} = \partial_{\mu}\delta^{ac} + \frac{i}{2}A^{ac}F_{\mu\nu}x_{\nu}$$
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where A is the hermitian matrix consisting of the sum generators in the adjoint representation. Then:

$$-(D^{2})^{ac} = -(\partial_{0}^{2})\delta^{ac} - (\partial_{1}^{2})\delta^{ac} + i(AB)^{ac}(x_{1}\partial_{0} - x_{0}\partial_{1}) + \frac{1}{4}(A^{2}B^{2})^{ac}(x_{0}^{2} + x_{1}^{2}) - (\partial_{2}^{2})\delta^{ac} - (\partial_{3}^{2})\delta^{ac} + i(AB)^{ac}(x_{3}\partial_{2} - x_{2}\partial_{3}) + \frac{1}{4}(A^{2}B^{2})^{ac}(x_{2}^{2} + x_{3}^{2}).$$
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In the diagonal basis of A these zero eigenvalues



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In the diagonal basis of A these zero eigenvalues Let A' be a six-by-six diagonal matrix containing the non-zero eigenvalues of A



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- In the diagonal basis of A these zero eigenvalues
   Let A' be a six-by-six diagonal matrix containing the non-zero eigenvalues of A
  - Now use:

$$c_{\mu} = \left[\frac{\partial_{\mu} \mathbb{1}}{(B\mathbf{A}')^{\frac{1}{2}}} + \frac{1}{2}(B\mathbf{A}')^{\frac{1}{2}}\hat{x}_{\mu}\right]$$

$$c_{\mu}^{\dagger} = \left[-\frac{\partial_{\mu} \mathbb{1}}{(B\mathbf{A}')^{\frac{1}{2}}} + \frac{1}{2}(B\mathbf{A}')^{\frac{1}{2}}\hat{x}_{\mu}\right].$$
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(22)

$$- (D^{2})' = B\mathbf{A}' \Big[ (c_{0}^{\dagger} + ic_{1}^{\dagger}) (c_{0} - ic_{1}) + 1 \Big] + B\mathbf{A}' \Big[ (c_{2}^{\dagger} + ic_{3}^{\dagger}) (c_{2} - ic_{3}) + 1 \Big],$$
(23)

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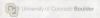


For zero eigenvalues of A contributions of  $D^2$  vanish in dim-reg or zeta-reg.



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- The remaining contributions are quantum harmonic oscillators



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- For zero eigenvalues of A contributions of  $D^2$  vanish in dim-reg or zeta-reg.
- The remaining contributions are quantum harmonic oscillators
  - Eigenspectrum of  $S_0$ :

 $(\Lambda_a^+)_{m,n} = (2n+1)B\lambda_a + (2m+1)B\lambda_a + 2B\lambda_a$  $(\Lambda_a^-)_{m,n} = (2n+1)B\lambda_a + (2m+1)B\lambda_a - 2B\lambda_a,$  (24)

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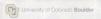


### Evaluating the Path Integral

HS transformation is applied to Z of the form

$$\int_{-\infty}^{\infty} \mathcal{D}\sigma \int_{0}^{\infty} \mathcal{D}\xi \mathbf{Re}[e^{i\int_{x}\xi_{\mu\nu}^{a}(\sigma_{\mu\nu}^{a}-f^{abc}a_{\mu}^{b}a_{\nu}^{c})}] = 1, \quad (25)$$

under the ansatz that  $\xi^a_{\mu\nu}$  is diagonal in Lorentz space.



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### Evaluating the Path Integral

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under the ansatz that  $\xi^{\rm a}_{\mu\nu}$  is diagonal in Lorentz space.

Some elbow grease gives:

$$Z = \int \mathcal{D}a\mathcal{D}c\mathcal{D}\bar{c} \operatorname{\mathbf{Re}}\left[\exp\left[-\int_{x}\frac{1}{4g_{0}^{2}}(\bar{F}_{\mu\nu}^{ab})^{2} + \frac{(B\Delta_{\mu\nu}^{a})^{2}}{g_{0}^{2}} + a_{\mu}^{a}\left[-(D^{2})^{ac}\delta_{\mu\nu} + 2A^{ac}F_{\mu\nu} + \Delta BA^{ac}\delta_{\mu\nu}\right]a_{\nu}^{c} + \bar{c}^{a}\left[(-(D_{\mu})^{2})^{ac}\right]c^{c}\right]_{\bar{B},\bar{\Delta}}.$$

(26)

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### Heat Kernels

$$\ln \det \theta = -\left[\frac{d}{ds}\frac{1}{\Gamma[s]}\int_0^\infty d\tau \tau^{s-1} K_\theta\right]_{s=0},\qquad(27)$$

where

$$K_{\theta} = \mathbf{Tr}_{\mu\nu}^{ab} \sum_{n,m} e^{-\tau\theta/\mu^2}$$
(28)



### Heat Kernels

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where

$$K_{\theta} = \operatorname{Tr}_{\mu\nu}^{ab} \sum_{n,m} e^{-\tau\theta/\mu^2}$$
(28)

$$\mathcal{K}_{\text{Ghost}} = \sum_{l} \beta V \frac{(B\lambda^{l})^{2}}{16\pi^{2}} \left[ \frac{1}{\sinh^{2}(\frac{B\lambda^{l}\tau}{\mu^{2}})} \right]$$
  
&  
$$\mathcal{K}_{\text{Glue}}^{R_{0}} = \sum_{l} \beta V \frac{(B\lambda^{l})^{2}}{4\pi^{2}} \left( e^{-\frac{\tau B\lambda^{l} \Delta}{\mu^{2}}} \right) \left[ 2 + \frac{1}{\sinh^{2}(\frac{B\lambda^{l}\tau}{\mu^{2}})} \right].$$
  
(29)



### Indet $\theta$ 's **Ghosts**:

$$\operatorname{\mathsf{n}\,det} \theta_{\operatorname{\mathsf{Ghost}}} = -\sum_{l} \beta V \frac{(B\lambda^{l})^{2}}{48\pi^{2}} \left[ \ln\left(\frac{B\lambda^{l}}{\mu^{2}}\right) + \ln\left(\frac{2e}{A^{12}}\right) \right]$$
(30)



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### Indet $\theta$ 's **Ghosts**:

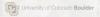
$$\ln \det \theta_{\text{Ghost}} = -\sum_{I} \beta V \frac{(B\lambda^{I})^{2}}{48\pi^{2}} \left[ \ln \left( \frac{B\lambda^{I}}{\mu^{2}} \right) + \ln \left( \frac{2e}{A^{12}} \right) \right]$$
(30)

Glues:

 $-\frac{1}{2}\ln\det[\theta_{\mathbf{Glue}}^{R_0}] = \sum_{l}\frac{\beta V(B\lambda^l)^2}{8\pi^2} \times$  $\frac{d}{ds} \left| \frac{1}{\Gamma(s)} \left( 2 \int_0^\infty d\tau \tau^{s-1} e^{-\tau \frac{B\lambda/\Delta}{\mu^2}} \right) \right|^{\frac{1}{2}}$  $+\int_{0}^{\infty} d\tau \sum_{n=0}^{2} \frac{\tau^{n-1} (-\frac{B\lambda'\Delta}{\mu^{2}})^{n}}{n!} \frac{1}{\sinh^{2} \left(\frac{B\lambda'\tau}{\mu^{2}}\right)}$  $+\int_{0}^{\infty} d\tau \sum_{n=3}^{\infty} \frac{\tau^{n-1} \left(-\frac{B\lambda^{l}\Delta}{\mu^{2}}\right)^{n}}{n!} \frac{1}{\sinh^{2} \left(\frac{B\lambda^{l}\tau}{\mu^{2}}\right)^{n}} \int_{s=0}^{\infty} \int_{s=0}^{\infty} d\tau d\tau d\tau$ (21)



### Evaluating the Path Integral **The equation**:



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## Evaluating the Path Integral **The equation:**

$$\frac{\ln Z}{\beta V} = -\frac{(B\lambda')^2}{g_0^2} - \frac{(B\lambda')^2}{48\pi^2} \left[ 11 \ln \left(\frac{B\lambda'}{\mu^2}\right) + C + f(\Delta) + g(\Delta, B, \mu) \right] \Big|_{\bar{B}\bar{\Delta}}$$
(32)

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## Evaluating the Path Integral **The equation:**

$$\frac{\ln Z}{\beta V} = -\frac{(B\lambda^{\prime})^2}{g_0^2} - \frac{(B\lambda^{\prime})^2}{48\pi^2} \left[ 11 \ln \left(\frac{B\lambda^{\prime}}{\mu^2}\right) + C + f(\Delta) + g(\Delta, B, \mu) \right] \Big|_{\bar{B}\bar{\Delta}}$$
(32)

where

$$C = \ln\left(\frac{e}{2A^{12}}\right)$$

$$g(\Delta, B, \mu) = 3\Delta^{2} \ln\left(\frac{B\lambda^{l}}{\mu^{2}}\right) + 12\ln(\Delta)$$

$$f(\Delta) = -\Delta\left(\ln\left(64\pi^{6}\right) - 3\Delta\ln(2)\right)$$

$$+ 12\Delta\ln\left(\Gamma\left[\frac{\Delta+2}{2}\right]\right) - 24\zeta^{(1,0)}\left(-\frac{1}{2}\frac{\Delta+2}{2}\right)$$

$$= -\frac{\Delta}{2}\left(\frac{1}{2}\frac{\Delta}{2}\right)$$

$$= -\frac{2}{2}\left(\frac{1}{2}\frac{\Delta}{2}\right)$$

$$= -\frac{2}{2}\left(\frac{1}{2}\frac{\Delta}{2}\right)$$



### Renormalization

Use:

$$\frac{1}{g^2(\mu)} = \frac{1}{g_0^2} + \frac{1}{48\pi^2}g(\Delta, B, \mu)$$
(34)

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#### Renormalization

### Use: $\frac{1}{g^2(\mu)} = \frac{1}{g_0^2} + \frac{1}{48\pi^2}g(\Delta, B, \mu) \quad (34)$ Gives: $\frac{\ln Z}{\beta V} = -\sum_{a} \frac{(B\lambda')^2}{48\pi^2} \left[ 11\ln\left(\frac{B\lambda'}{\Lambda_{\rm YM}^2}\right) + C + f(\Delta) \right]_{\bar{B},\bar{\Delta}}.$

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(35)

#### Renormalization

# Use: $\frac{1}{g^{2}(\mu)} = \frac{1}{g_{0}^{2}} + \frac{1}{48\pi^{2}}g(\Delta, B, \mu) \quad (34)$ Gives: $\frac{\ln Z}{\beta V} = -\sum_{a} \frac{(B\lambda')^{2}}{48\pi^{2}} \left[ 11 \ln \left(\frac{B\lambda'}{\Lambda_{YM}^{2}}\right) + C + f(\Delta) \right]_{\bar{B},\bar{\Delta}} \quad (35)$

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#### Solve the gap equations and get results!

