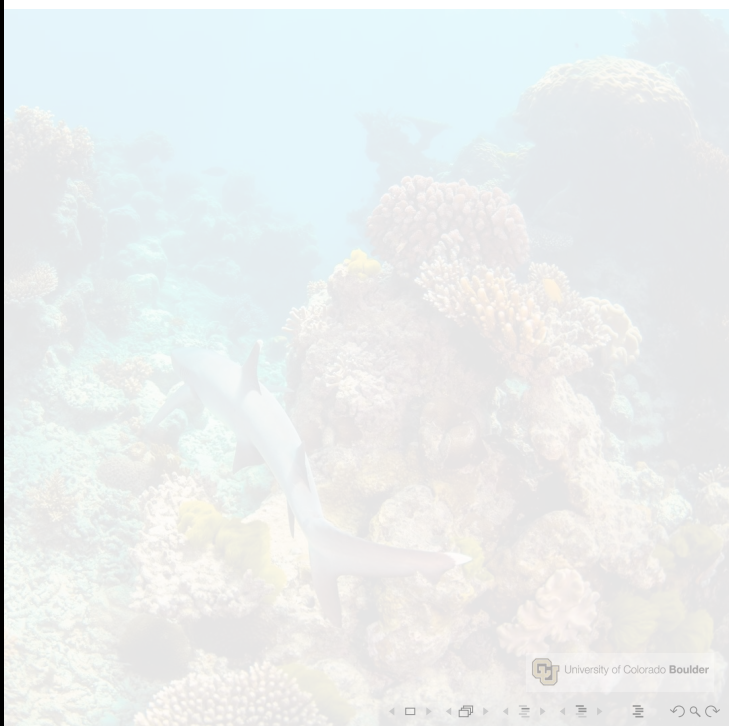


Non-Perturbative Yang-Mills Beyond One-Loop Order

Who? **Seth Grable**

From? Department of Physics
University of Colorado Boulder

When? August 2024



YM calculations: The Inspiration

- We want to calculate YM, QCD, and non-abelian theories!
- Potential cool calculations:

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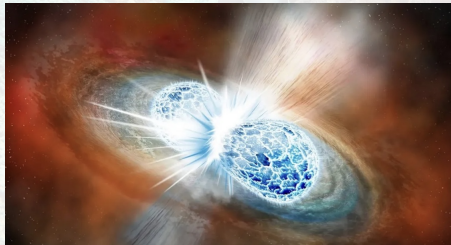
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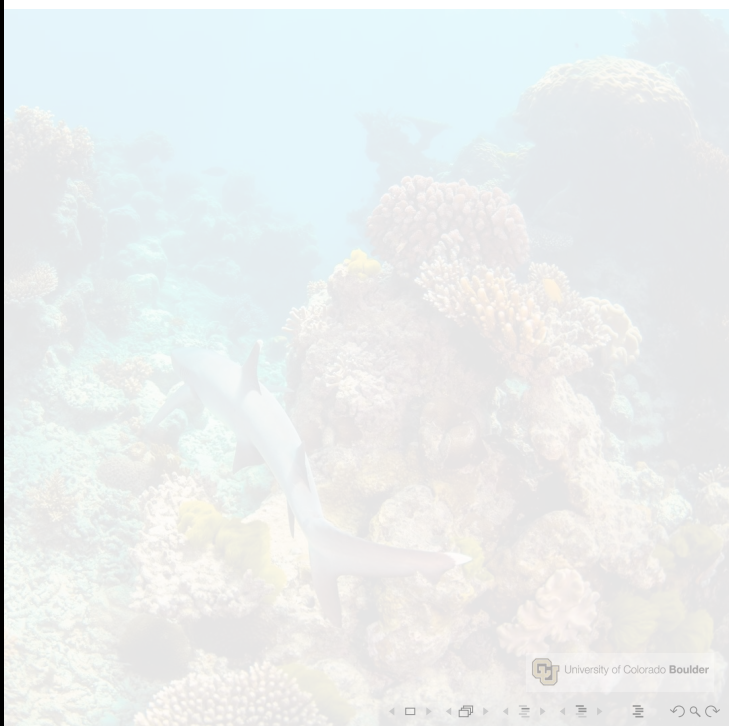
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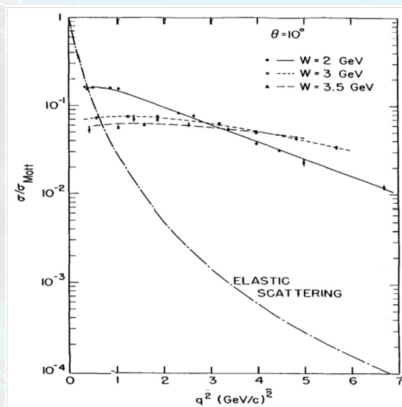
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- Next gen of gravitational wave detectors in the coming decade: Einstein Telescope, NEMO, Cosmic Explorer, LISA





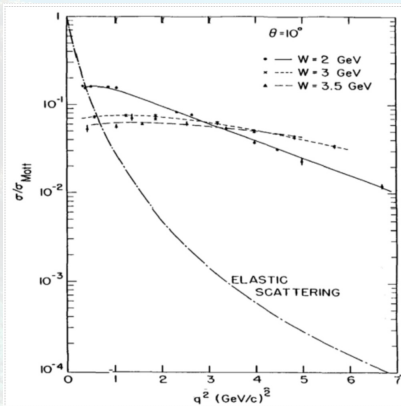
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Award-winning feature: Asymptotic freedom



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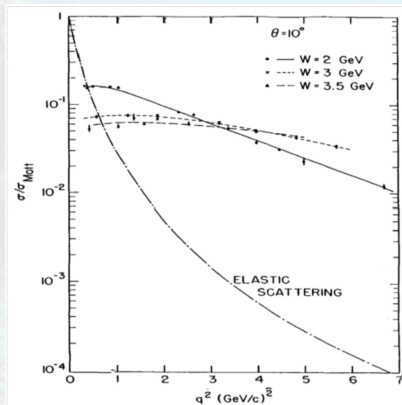
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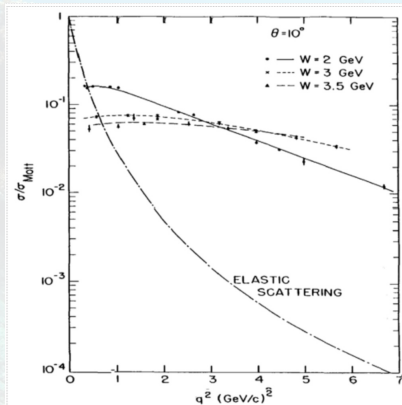


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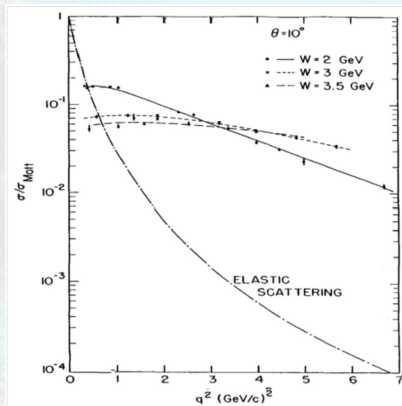
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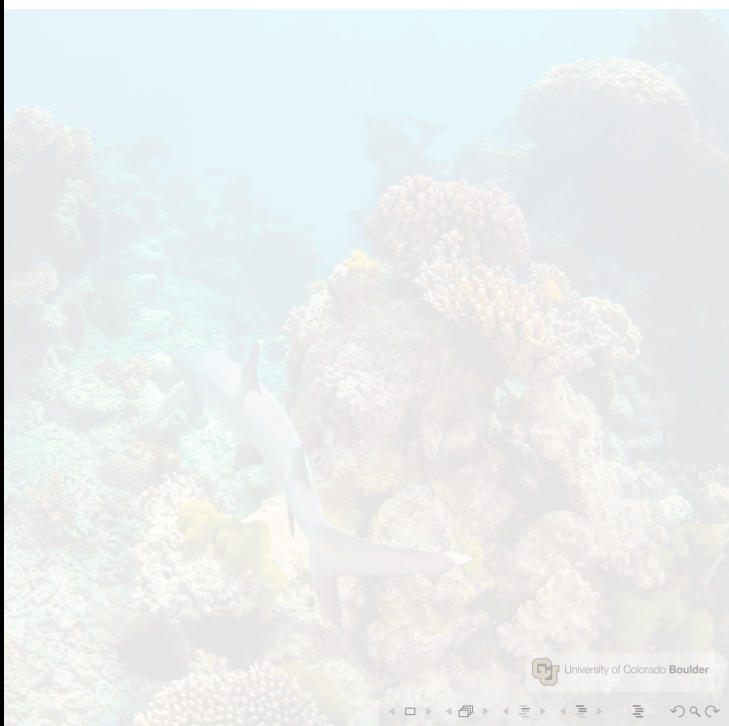


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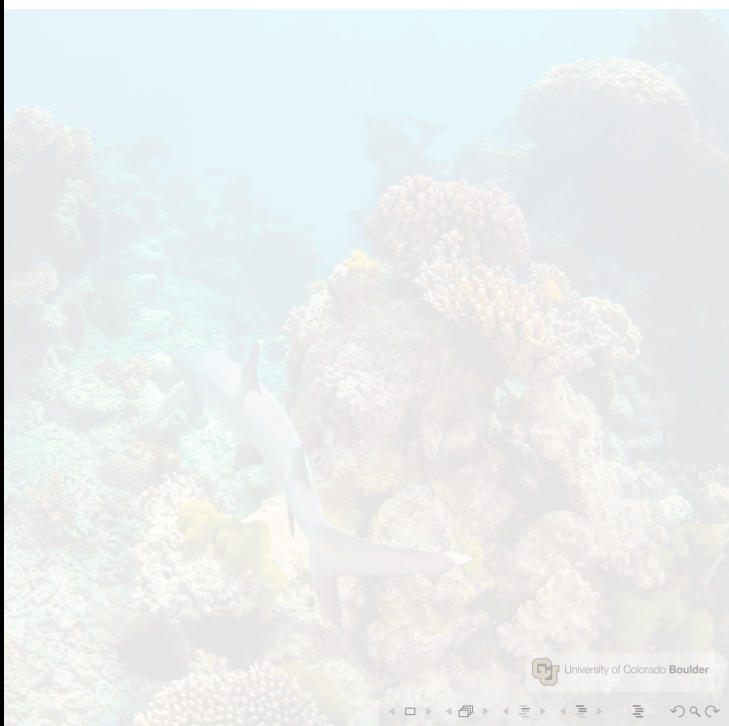
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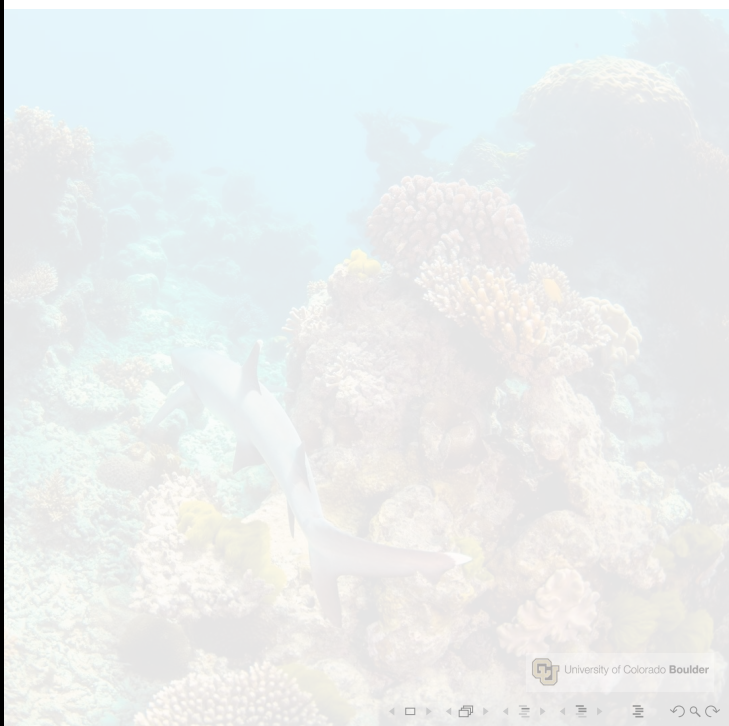
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 - **Laplace's method.**



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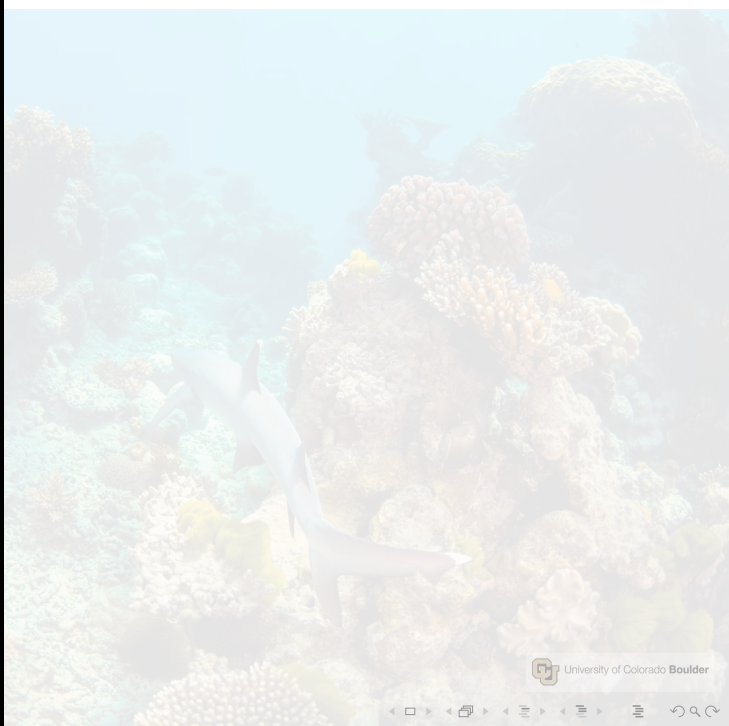
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$$A_\mu^a(x) = B_\mu^a(x) + a_\mu^a(x) \quad (2)$$

- This allows for an expansion around B_μ^a in terms of a_μ^a .



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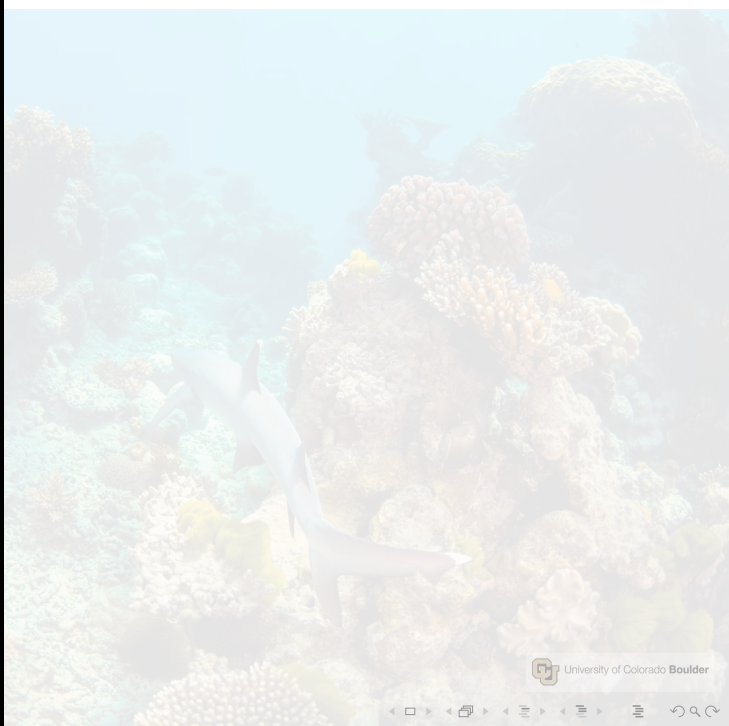
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- The general strategy: reduce the complexity of S_I .
- S_0 , quadratic contributions, are calculable (Gaussian).



The Background Field Method

- How can we simplify S_I (cubic and quartic terms)?

$$S_I = \int_x g_0 (D_\mu a_\nu^a) f^{abc} a_\mu^b a_\nu^c + \frac{g_0^2}{4} (f^{abc} a_\mu^b a_\nu^c)^2. \quad (4)$$

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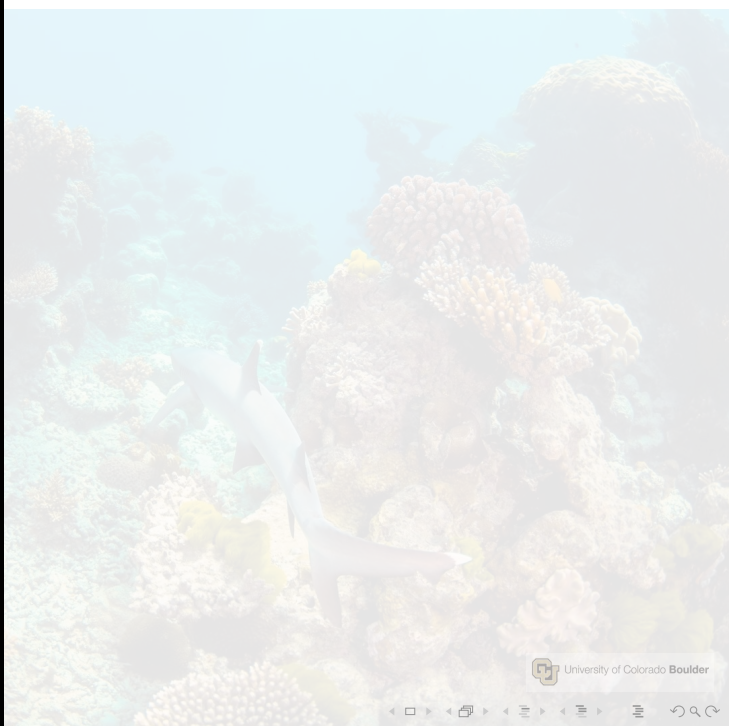
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- Just drop S_I ?
- This gives the one-loop β -function, but also IR divergences generic n_f .
- However, I discovered that IR divergence perfectly cancels for $n_f = 12$.



$n_f = 12$ Results

- Stable background field configurations \bar{B}

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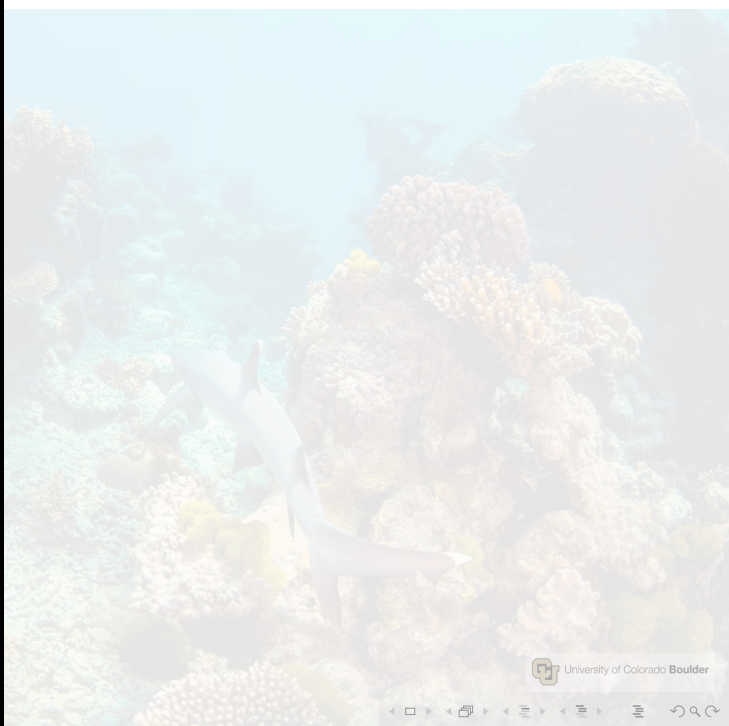
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$$\bar{B}_{\sigma_8} \approx 4.27 \Lambda_{\overline{MS}}^2 \quad (5)$$

- Romatschke and I further found a vanishing Polyakov loop expectation value under a critical temperature $T_c \approx .81 \Lambda_{\overline{MS}}$, and a non-zero expectation value for $T > T_c$.



Generic n_f

- $n_f = 12$ is a special case.

Generic n_f

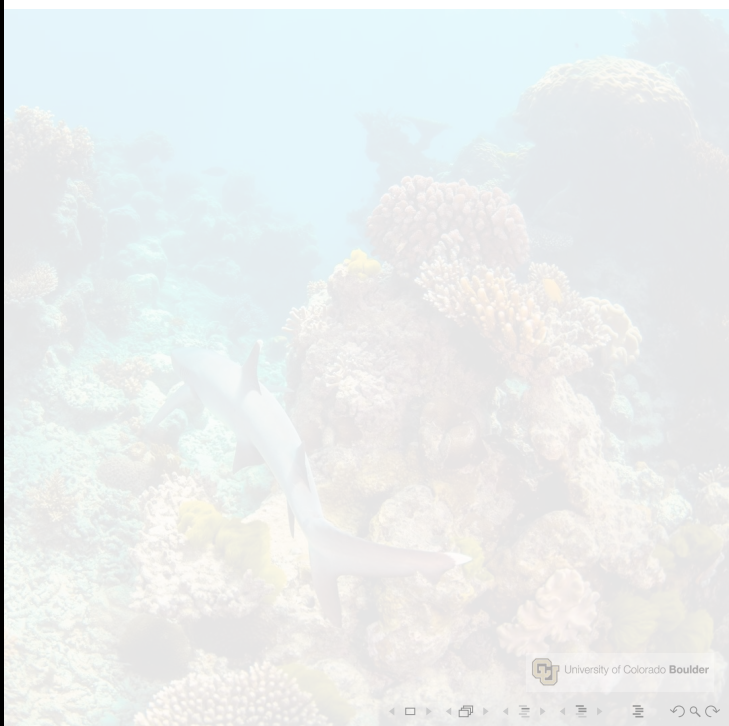
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Hubbard-Stratonovich transformation

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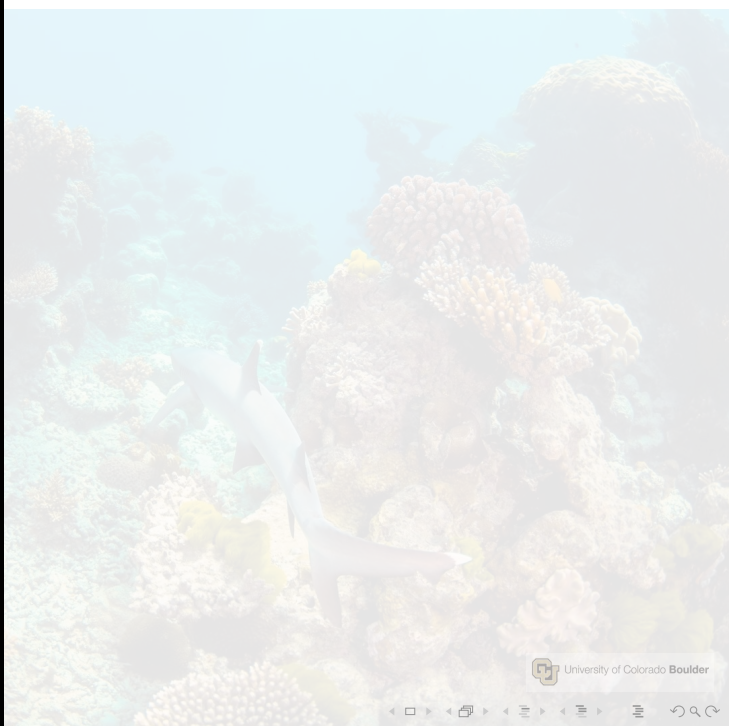
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- $i\Delta$ acts like an additional "effective mass" coming from the x^4 term



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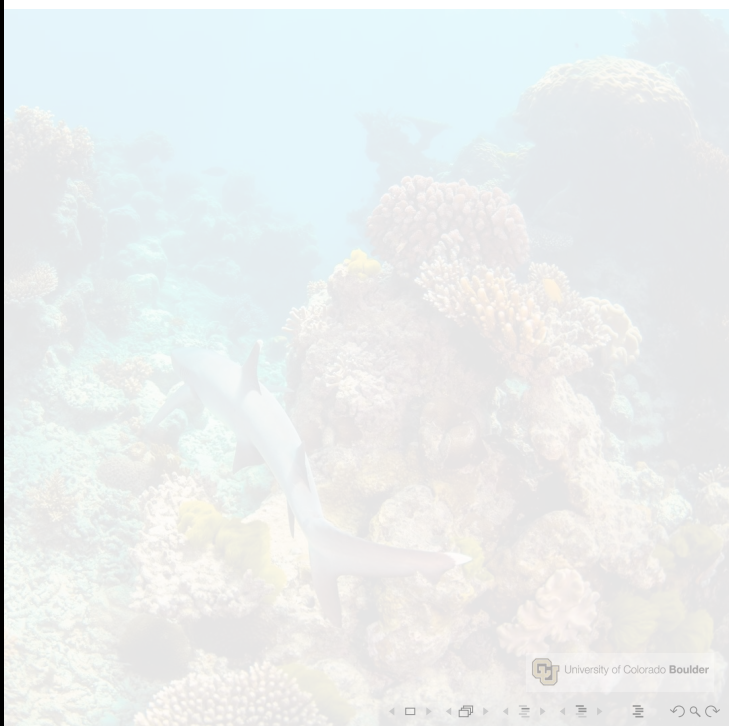
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- “Gapped Hamiltonian”



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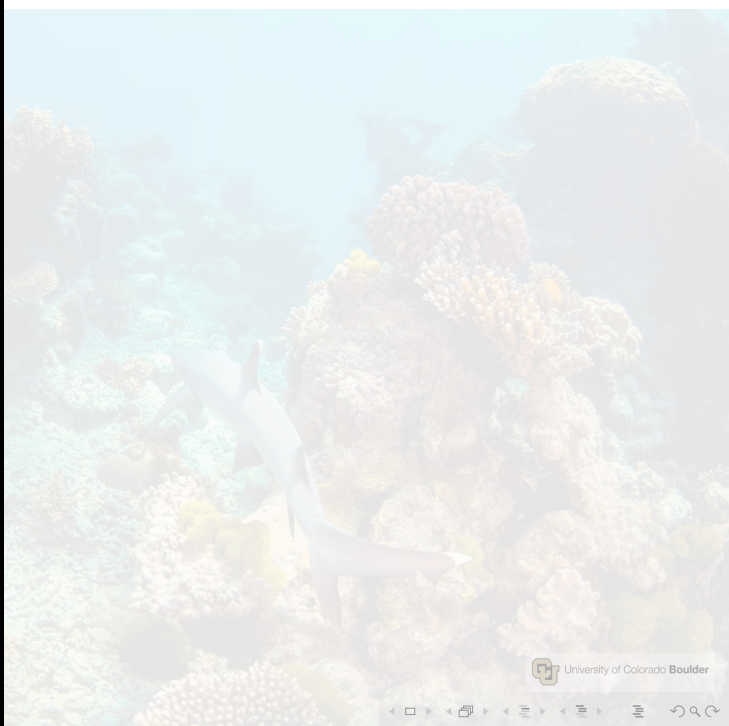
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- The effective action is gauge invariant.



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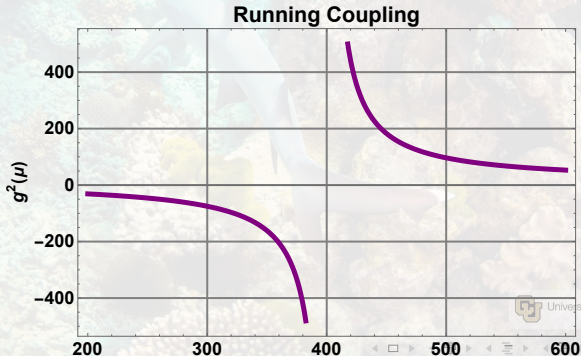
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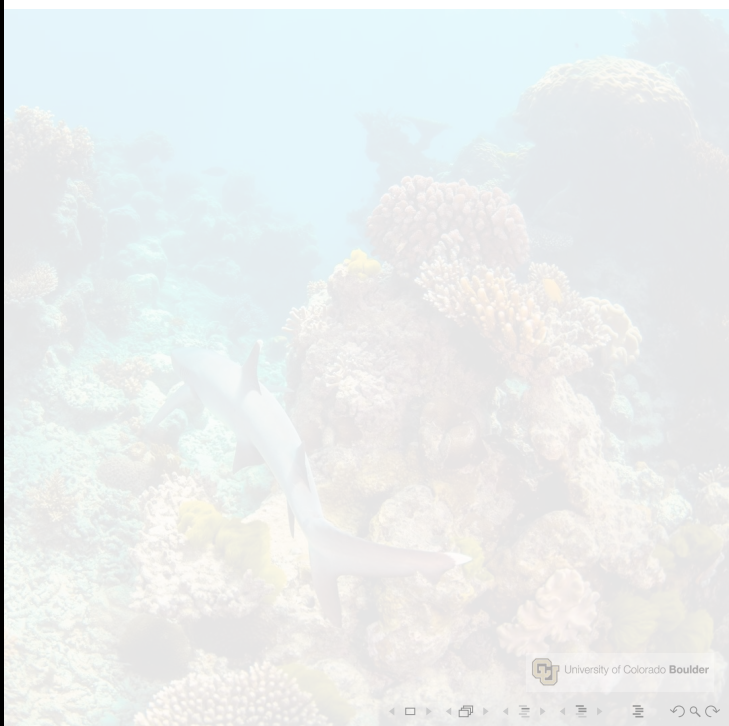
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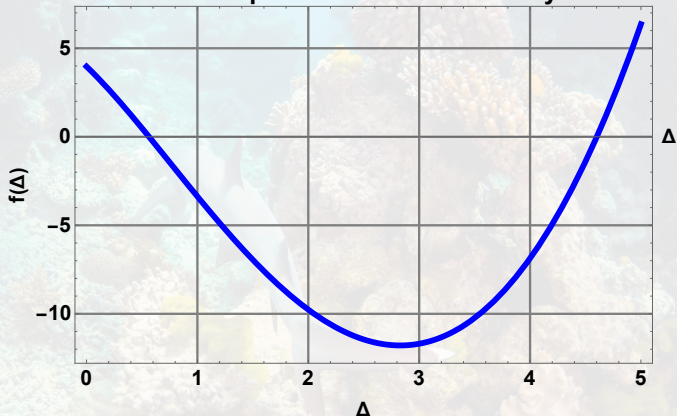
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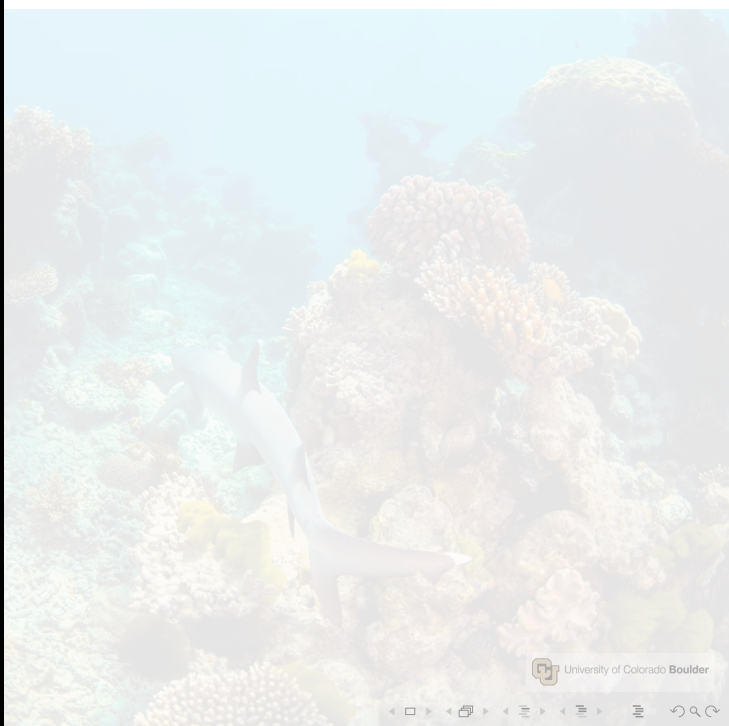
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Functional Dependence of the Auxillary field Δ



Gap:

$$\bar{\Delta} \rightarrow 2.82898.$$



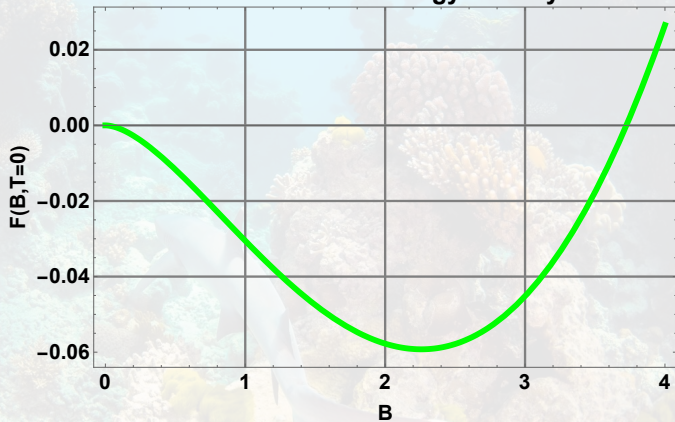
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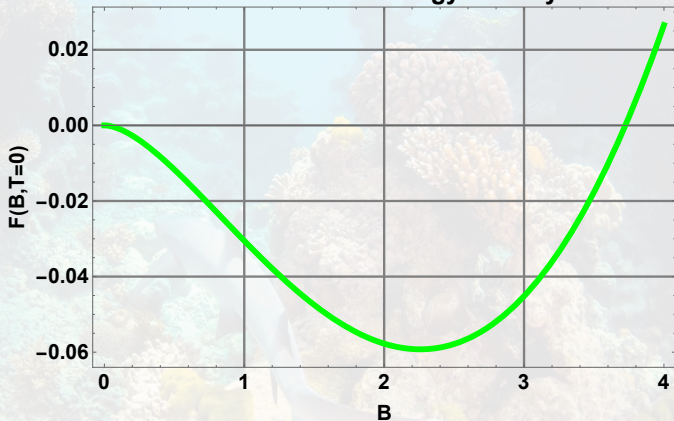
Vacuum Free Energy Density



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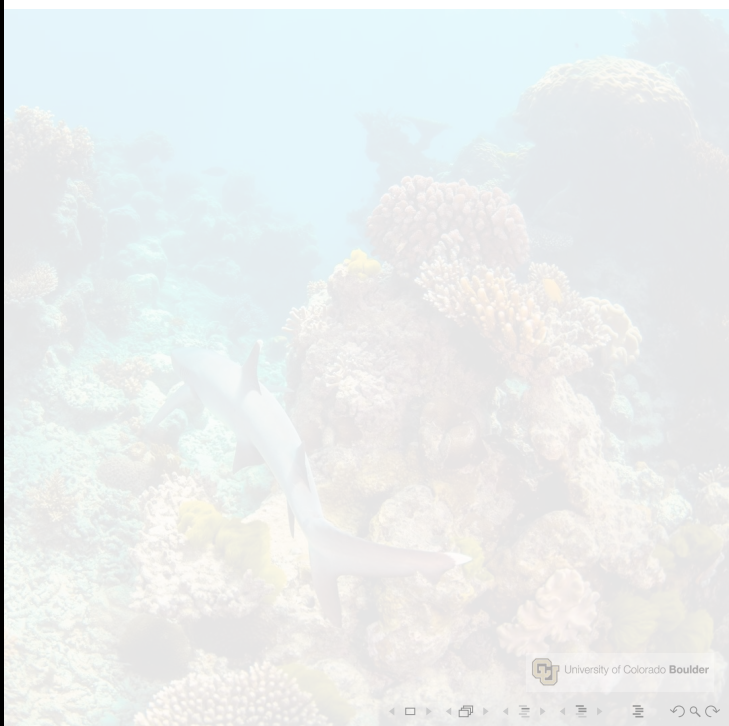
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Vacuum Free Energy Density



Gap:

$$\bar{B} \rightarrow \frac{2.25885\Lambda_{\text{YM}}^2}{\lambda'}, \quad \bar{B} \rightarrow 0, \quad (11)$$



Results

- **Vacuum Pressure:**

Results

■ Vacuum Pressure:

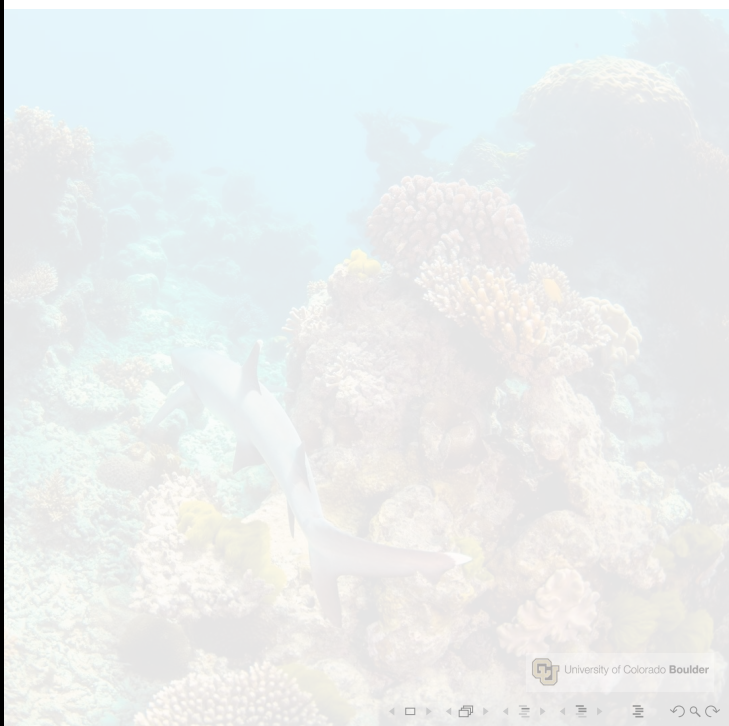
$$\frac{\ln Z}{\beta V} = 6 \times 0.0592377 \Lambda_{\text{YM}}^4. \quad (12)$$

Results

- Vacuum Pressure:

$$\frac{\ln Z}{\beta V} = 6 \times 0.0592377 \Lambda_{\text{YM}}^4. \quad (12)$$

- The scalar glueball mass is read off as:



Future Work

- **Calculating the propagators**

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Future Work

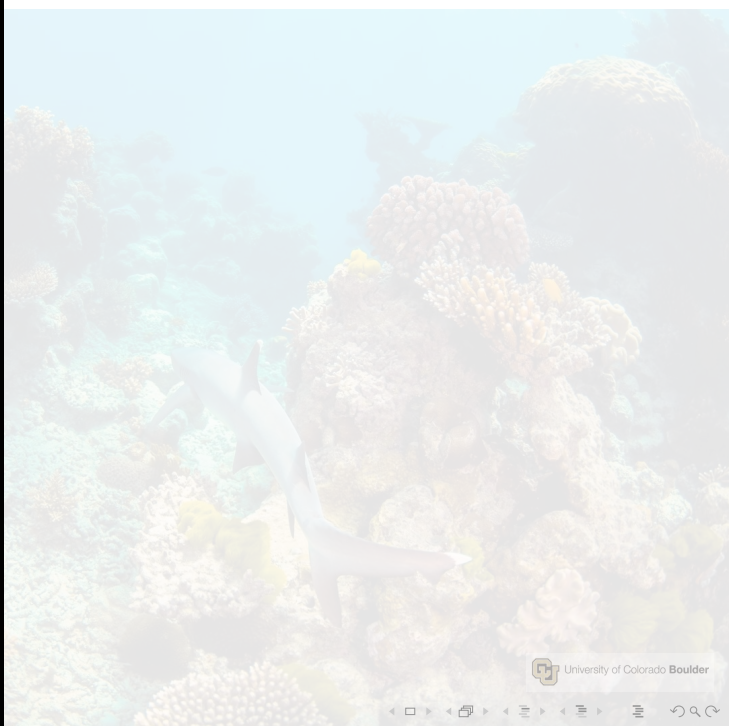
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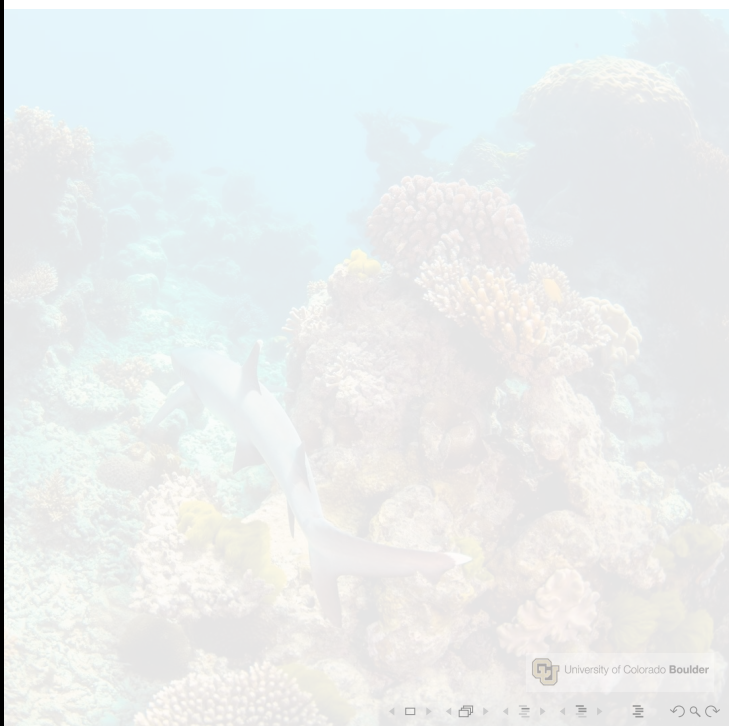
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Fun!



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Bibliography

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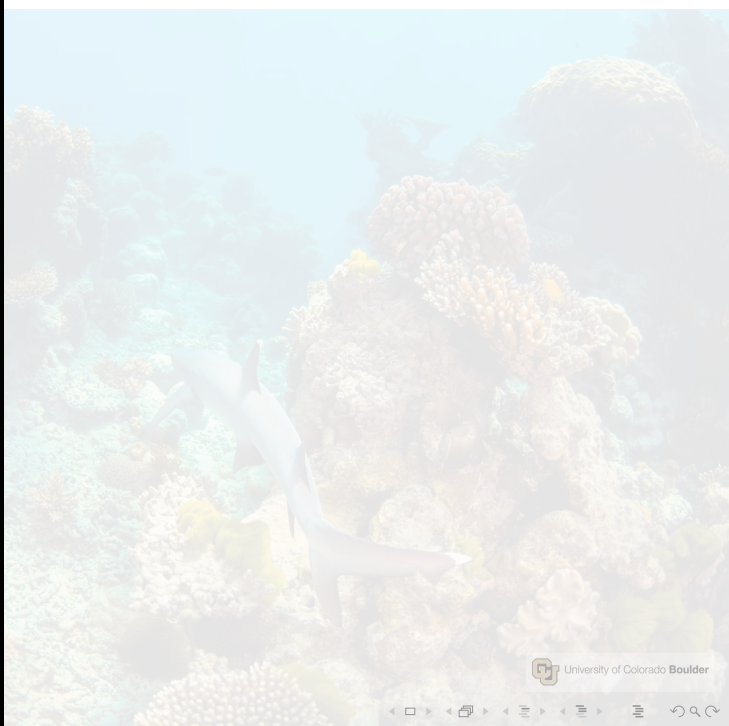
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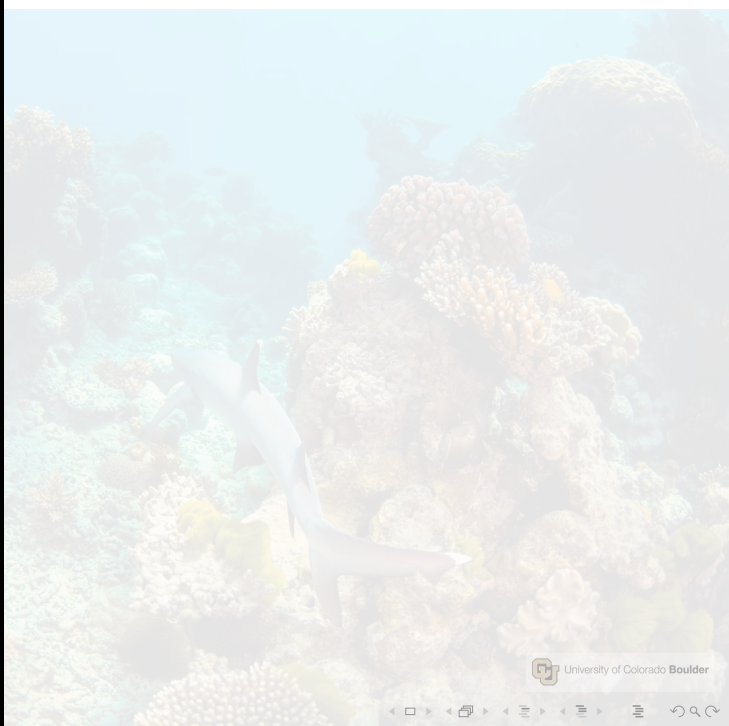
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$$S_0 = \int_x \frac{1}{2} a_{\mu}^a \left[- (D^2)^{ac} \delta_{\mu\nu} + 2 F_{\mu\nu}^c f^{abc} \right] a_{\nu}^c \quad (17)$$

$$S_I = \int_x g_0 (D_{\mu} a_{\nu}^a) f^{abc} a_{\mu}^b a_{\nu}^c + \frac{g_0^2}{4} (f^{abc} a_{\mu}^b a_{\nu}^c)^2. \quad (18)$$



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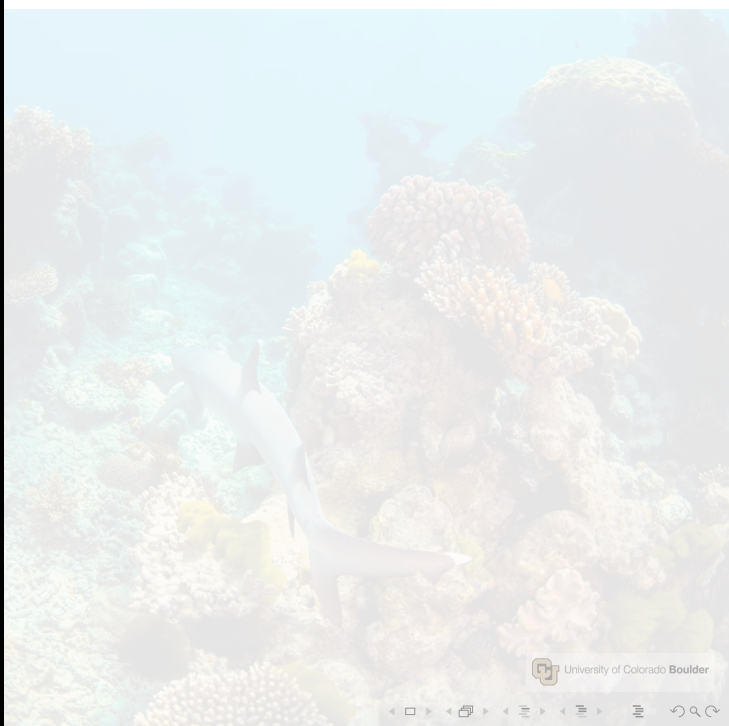
Gauge Fixing and Local Gauge Invariance

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- **After gauge fixing S_0 is invariant to:**

$$\begin{aligned} a_\mu^a(x) &\rightarrow a_\mu^a(x) - f^{abc} \beta^b(x) a_\mu^c(x) \\ B_\mu^a(x) &\rightarrow B_\mu^a(x) + B_\mu^a(x) D_\mu \beta^a(x) \\ c^a(x) &\rightarrow c^a(x) - f^{abc} \beta^b(x) c^c(x). \end{aligned} \quad (19)$$



The Eigenspectrum of S_0

$$D_{\mu}^{ac} = \partial_{\mu} \delta^{ac} + \frac{i}{2} A^{ac} F_{\mu\nu} x_{\nu} \quad (20)$$

where **A** is the hermitian matrix consisting of the sum generators in the adjoint representation.

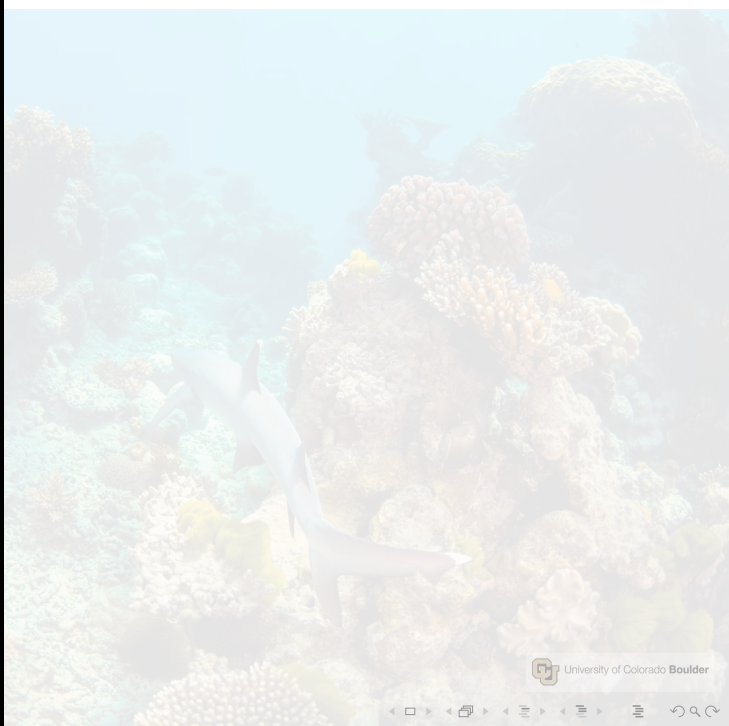
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Then:

$$\begin{aligned} -(D^2)^{ac} = & -(\partial_0^2) \delta^{ac} - (\partial_1^2) \delta^{ac} + i(AB)^{ac} (x_1 \partial_0 - x_0 \partial_1) \\ & + \frac{1}{4} (A^2 B^2)^{ac} (x_0^2 + x_1^2) \\ & - (\partial_2^2) \delta^{ac} - (\partial_3^2) \delta^{ac} + i(AB)^{ac} (x_3 \partial_2 - x_2 \partial_3) \\ & + \frac{1}{4} (A^2 B^2)^{ac} (x_2^2 + x_3^2). \end{aligned} \quad (21)$$



The Eigenspectrum of S_0 Continued..

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The Eigenspectrum of S_0 Continued..

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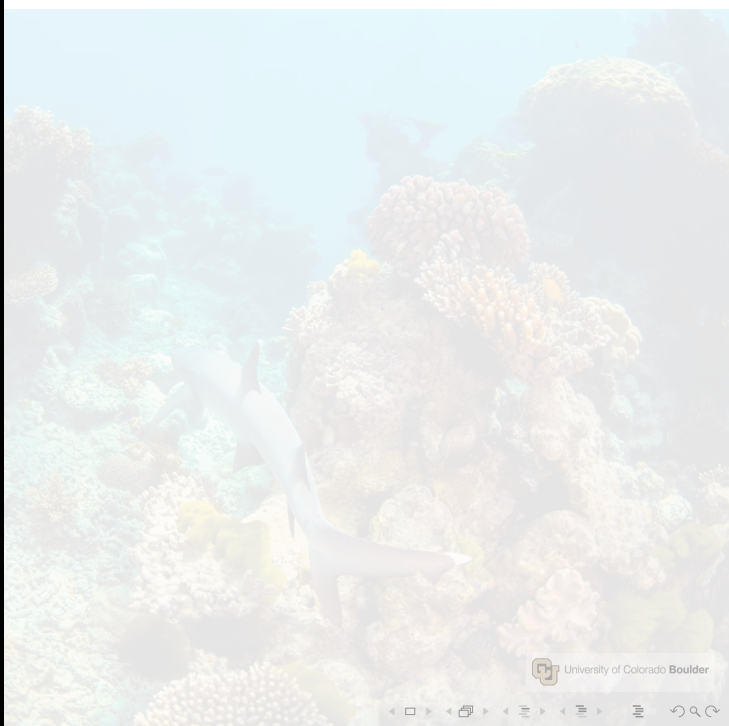
$$c_\mu = \left[\frac{\partial_\mu \mathbb{1}}{(\mathbf{BA}')^{\frac{1}{2}}} + \frac{1}{2}(\mathbf{BA}')^{\frac{1}{2}} \hat{x}_\mu \right] \quad (22)$$
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$$-(D^2)' = \mathbf{BA}' \left[(c_0^\dagger + ic_1^\dagger)(c_0 - ic_1) + 1 \right] + \mathbf{BA}' \left[(c_2^\dagger + ic_3^\dagger)(c_2 - ic_3) + 1 \right], \quad (23)$$



The Eigenspectrum of S_0 Continued..

- For zero eigenvalues of A contributions of D^2 vanish in dim-reg or zeta-reg.

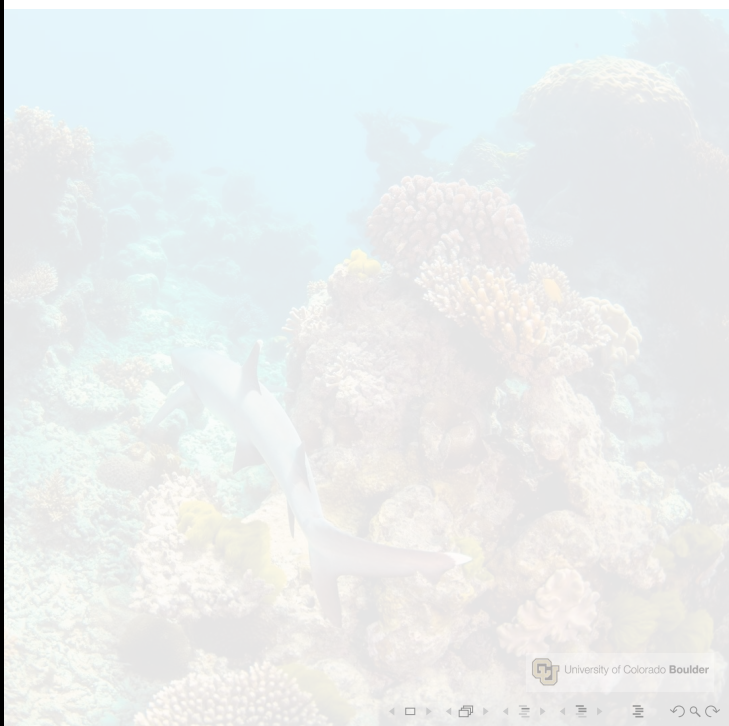
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$$\begin{aligned}(\Lambda_a^+)_{m,n} &= (2n + 1)B\lambda_a + (2m + 1)B\lambda_a + 2B\lambda_a \\(\Lambda_a^-)_{m,n} &= (2n + 1)B\lambda_a + (2m + 1)B\lambda_a - 2B\lambda_a,\end{aligned}\quad (24)$$



Evaluating the Path Integral

- HS transformation is applied to Z of the form

$$\int_{-\infty}^{\infty} \mathcal{D}\sigma \int_0^{\infty} \mathcal{D}\xi \operatorname{Re}[e^{i \int_x \xi_{\mu\nu}^a (\sigma_{\mu\nu}^a - f^{abc} a_{\mu}^b a_{\nu}^c)}] = 1, \quad (25)$$

under the ansatz that $\xi_{\mu\nu}^a$ is diagonal in Lorentz space.

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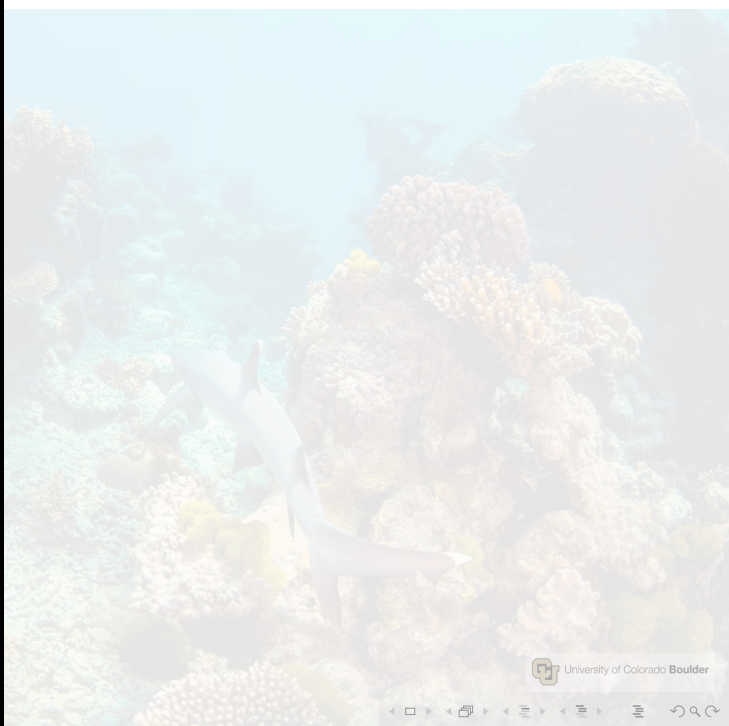
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- Some elbow grease gives:

$$\begin{aligned} Z = \int \mathcal{D}a \mathcal{D}c \mathcal{D}\bar{c} \operatorname{Re} \left[\exp \left[- \int_x \frac{1}{4g_0^2} (\bar{F}_{\mu\nu}^{ab})^2 + \frac{(B\Delta_{\mu\nu}^a)^2}{g_0^2} \right. \right. \\ \left. \left. + a_{\mu}^a \left[-(D^2)^{ac} \delta_{\mu\nu} + 2A^{ac} F_{\mu\nu} \right. \right. \right. \\ \left. \left. \left. + \Delta B A^{ac} \delta_{\mu\nu} \right] a_{\nu}^c + \bar{c}^a \left[-(D_{\mu})^2 \right]^{ac} c^c \right] \right]_{\bar{B}, \bar{\Delta}}. \end{aligned} \quad (26)$$



Heat Kernels

$$\ln \det \theta = - \left[\frac{d}{ds} \frac{1}{\Gamma[s]} \int_0^\infty d\tau \tau^{s-1} K_\theta \right]_{s=0}, \quad (27)$$

where

$$K_\theta = \mathbf{Tr}_{\mu\nu}^{ab} \sum_{n,m} e^{-\tau\theta/\mu^2} \quad (28)$$

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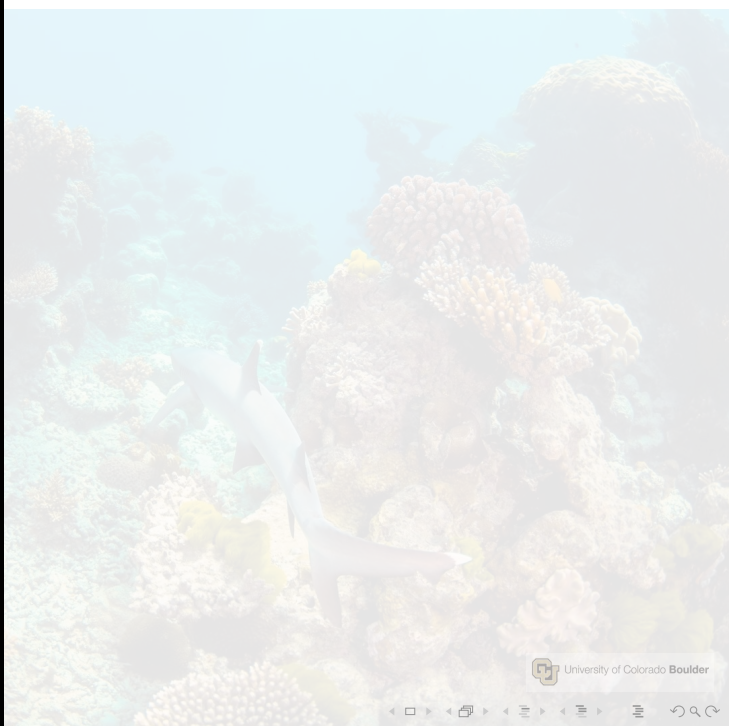
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$$K_{\text{Ghost}} = \sum_l \beta V \frac{(B\lambda')^2}{16\pi^2} \left[\frac{1}{\sinh^2\left(\frac{B\lambda'\tau}{\mu^2}\right)} \right]$$

&

$$K_{\text{Glue}}^{R_0} = \sum_l \beta V \frac{(B\lambda')^2}{4\pi^2} \left(e^{-\frac{\tau B\lambda'\Delta}{\mu^2}} \right) \left[2 + \frac{1}{\sinh^2\left(\frac{B\lambda'\tau}{\mu^2}\right)} \right]. \quad (29)$$



Indet θ 's

■ **Ghosts:**

$$\ln \det \theta_{\text{Ghost}} = - \sum_l \beta V \frac{(B\lambda')^2}{48\pi^2} \left[\ln \left(\frac{B\lambda'}{\mu^2} \right) + \ln \left(\frac{2e}{A^{12}} \right) \right] \quad (30)$$

Indet θ 's

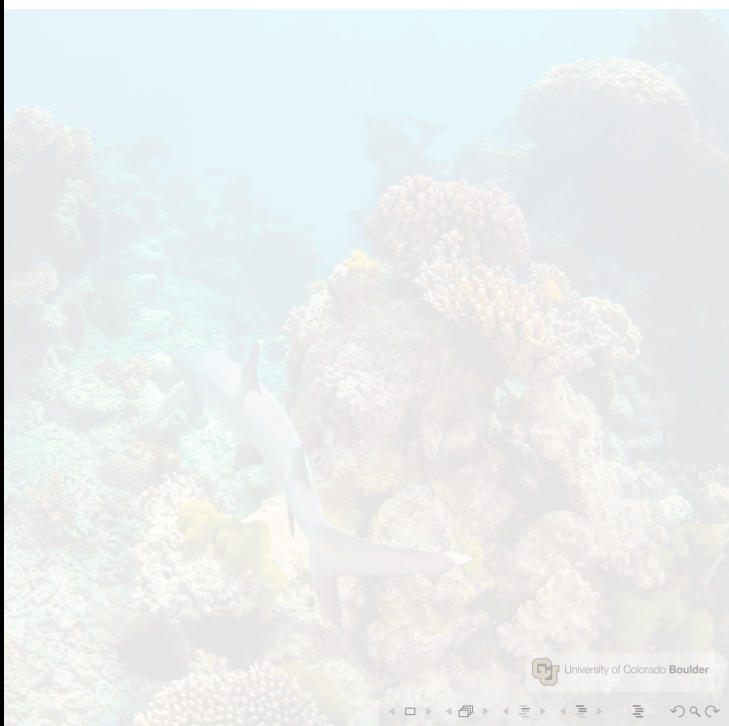
Ghosts:

$$\ln \det \theta_{\text{Ghost}} = - \sum_I \beta V \frac{(B\lambda')^2}{48\pi^2} \left[\ln \left(\frac{B\lambda'}{\mu^2} \right) + \ln \left(\frac{2e}{A^{12}} \right) \right] \quad (30)$$

Glues:

$$- \frac{1}{2} \ln \det[\theta_{\text{Glue}}^{R_0}] = \sum_I \frac{\beta V (B\lambda')^2}{8\pi^2} \times$$

$$\frac{d}{ds} \left[\frac{1}{\Gamma(s)} \left(2 \int_0^\infty d\tau \tau^{s-1} e^{-\tau \frac{B\lambda'\Delta}{\mu^2}} \right. \right. \\ \left. \left. + \int_0^\infty d\tau \sum_{n=0}^2 \frac{\tau^{n-1} \left(-\frac{B\lambda'\Delta}{\mu^2}\right)^n}{n!} \frac{1}{\sinh^2 \left(\frac{B\lambda'\tau}{\mu^2} \right)} \right. \right. \\ \left. \left. + \int_0^\infty d\tau \sum_{n=3}^\infty \frac{\tau^{n-1} \left(-\frac{B\lambda'\Delta}{\mu^2}\right)^n}{n!} \frac{1}{\sinh^2 \left(\frac{B\lambda'\tau}{\mu^2} \right)} \right) \right]_{s=0}$$



Evaluating the Path Integral

■ The equation:

$$\frac{\ln Z}{\beta V} = -\frac{(B\lambda')^2}{g_0^2} - \frac{(B\lambda')^2}{48\pi^2} \left[11 \ln \left(\frac{B\lambda'}{\mu^2} \right) + C + f(\Delta) + g(\Delta, B, \mu) \right] \Bigg|_{\bar{B}\bar{\Delta}} \quad (32)$$

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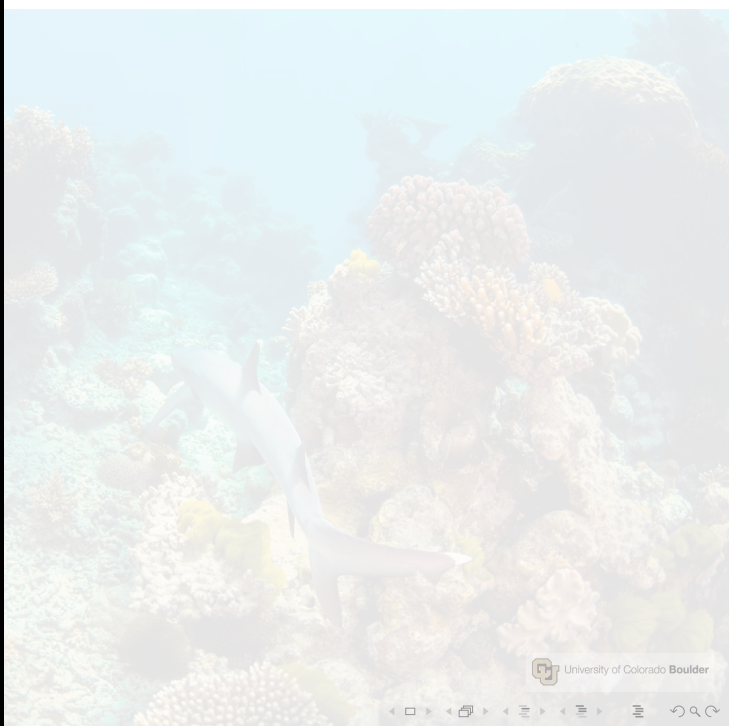
where

$$C = \ln \left(\frac{e}{2A^{12}} \right)$$

$$g(\Delta, B, \mu) = 3\Delta^2 \ln \left(\frac{B\lambda'}{\mu^2} \right) + 12 \ln(\Delta)$$

$$f(\Delta) = -\Delta \left(\ln(64\pi^6) - 3\Delta \ln(2) \right)$$

$$+ 12\Delta \ln \left(\Gamma \left[\frac{\Delta+2}{2} \right] \right) - 24\zeta^{(1,0)} \left(-1, \frac{\Delta+2}{2} \right)$$



Renormalization

■ Use:

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■ Solve the gap equations and get results!

