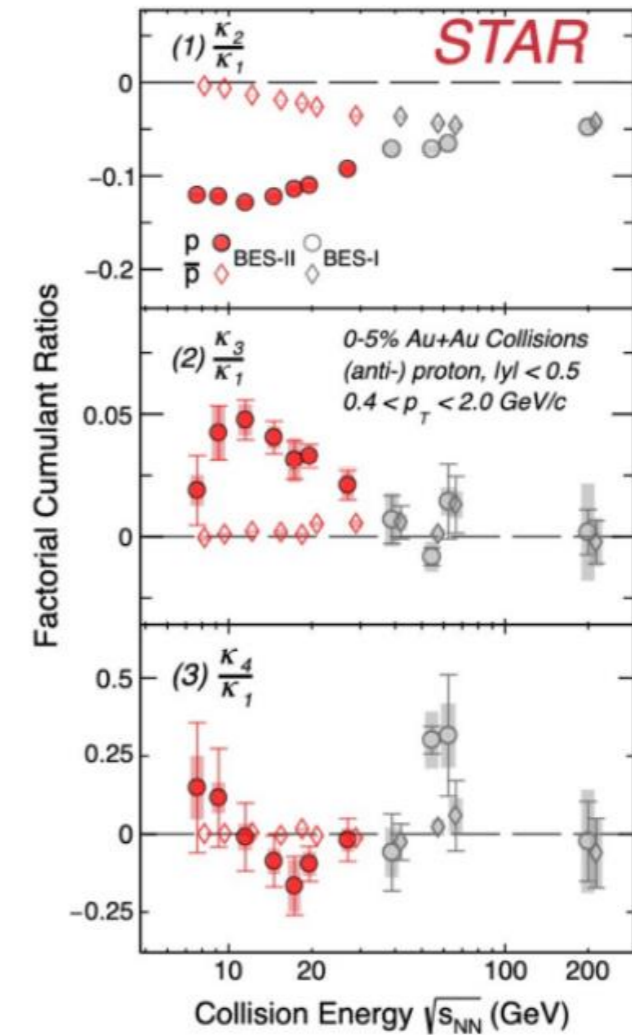
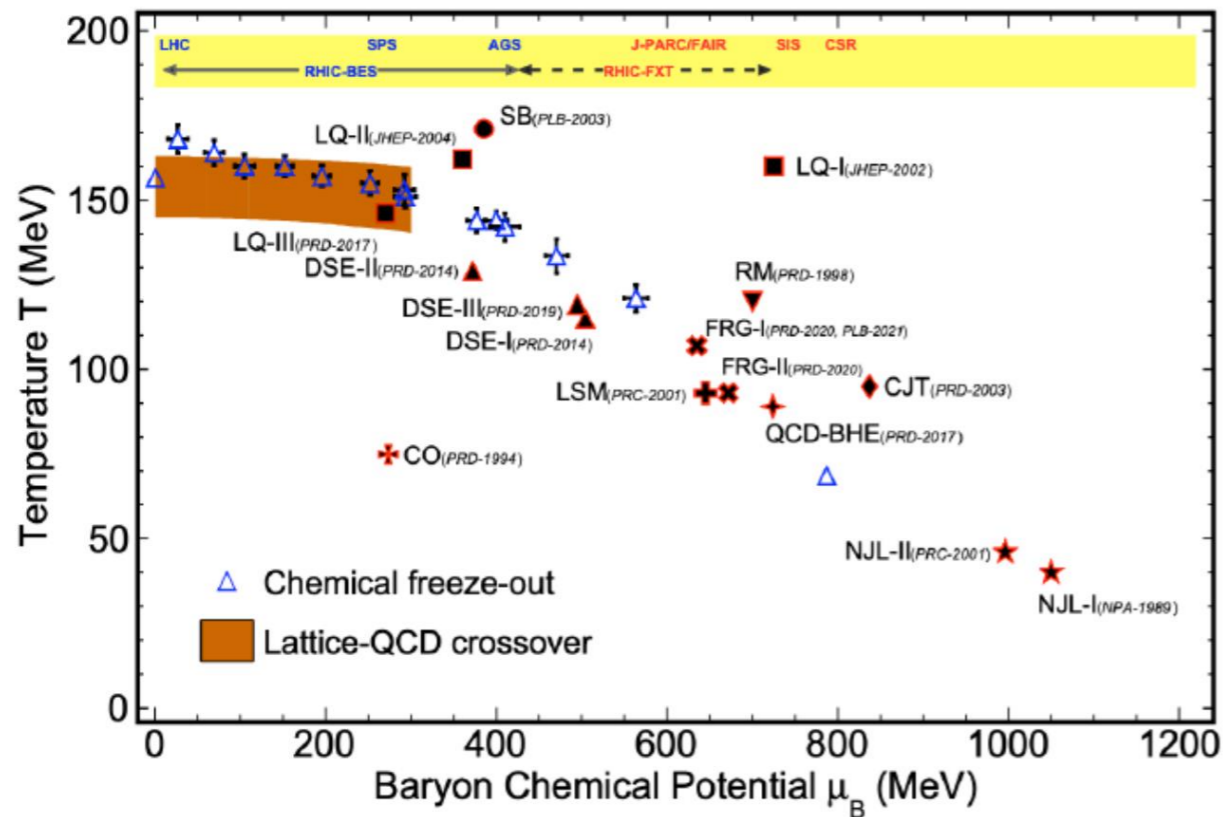
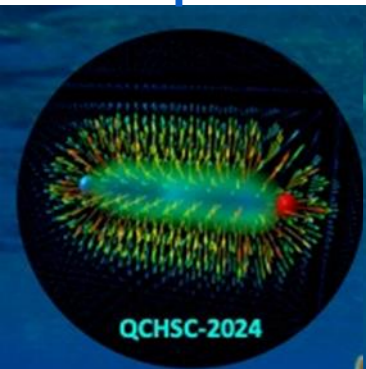


# Fluctuations of conserved charges in heavy ion collisions and lattice QCD

Rene Bellwied (University of Houston)



Special thanks to my UH theory colleagues: C. Ratti, V. Vovchenko, J. Jahan



**XVth Quark Confinement and the Hadron Spectrum Conference**  
 Cairns Convention Centre, Cairns, Queensland, Australia  
 19-24 August 2024 (inclusive)



# The QCD Phase Diagram

What we know

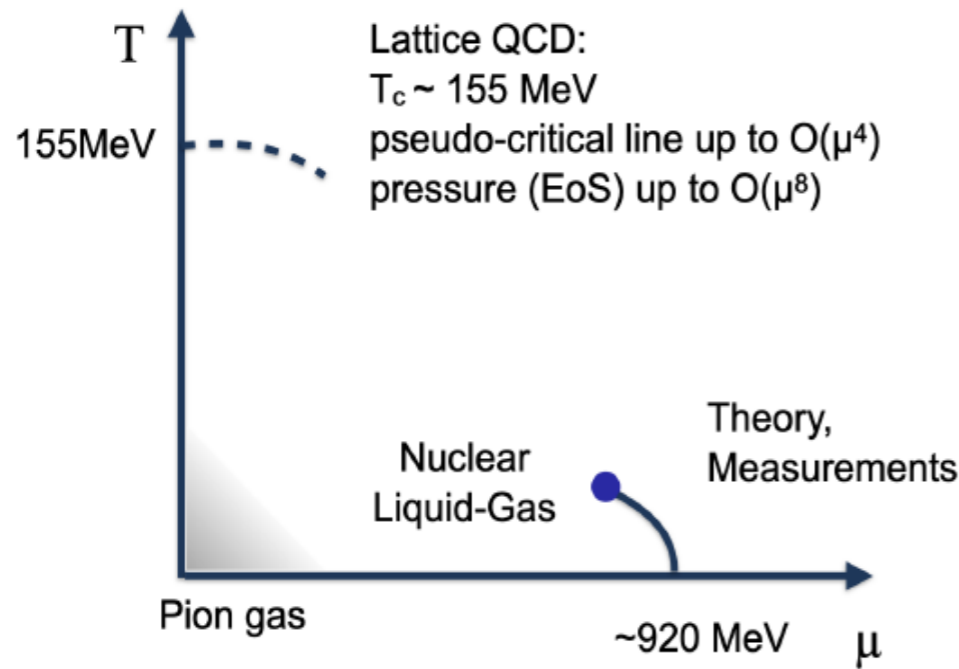


Figure courtesy of V. Koch

What we hope to know

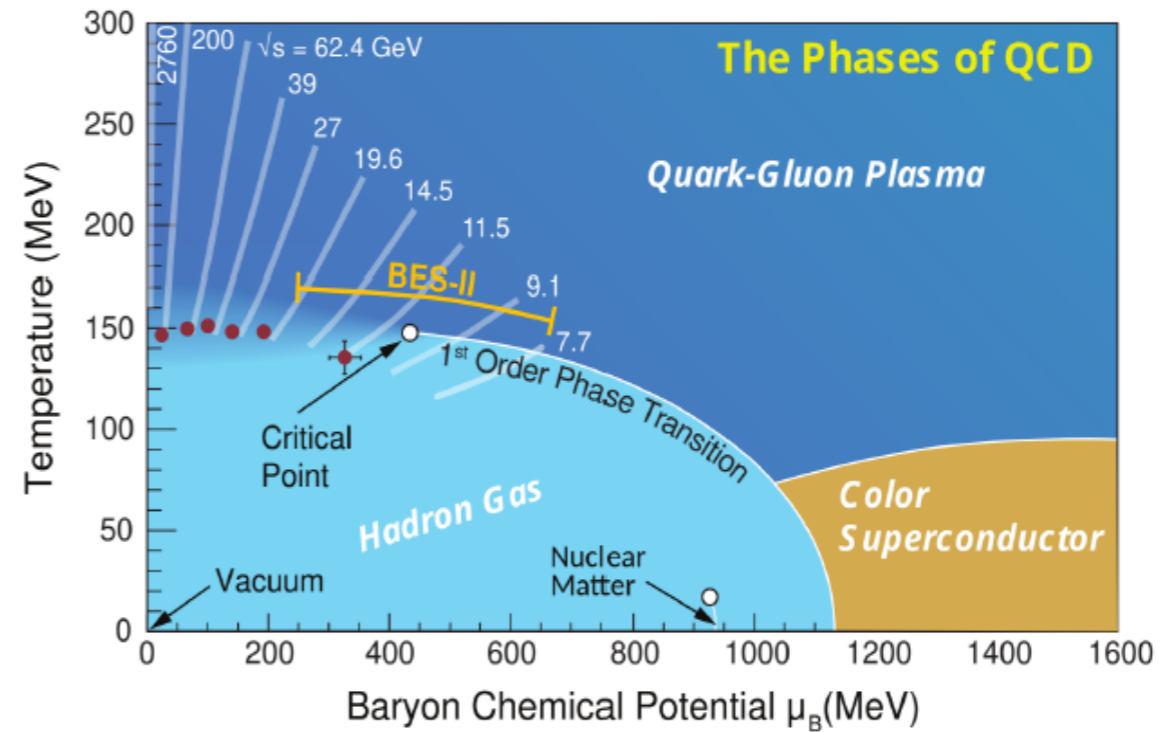
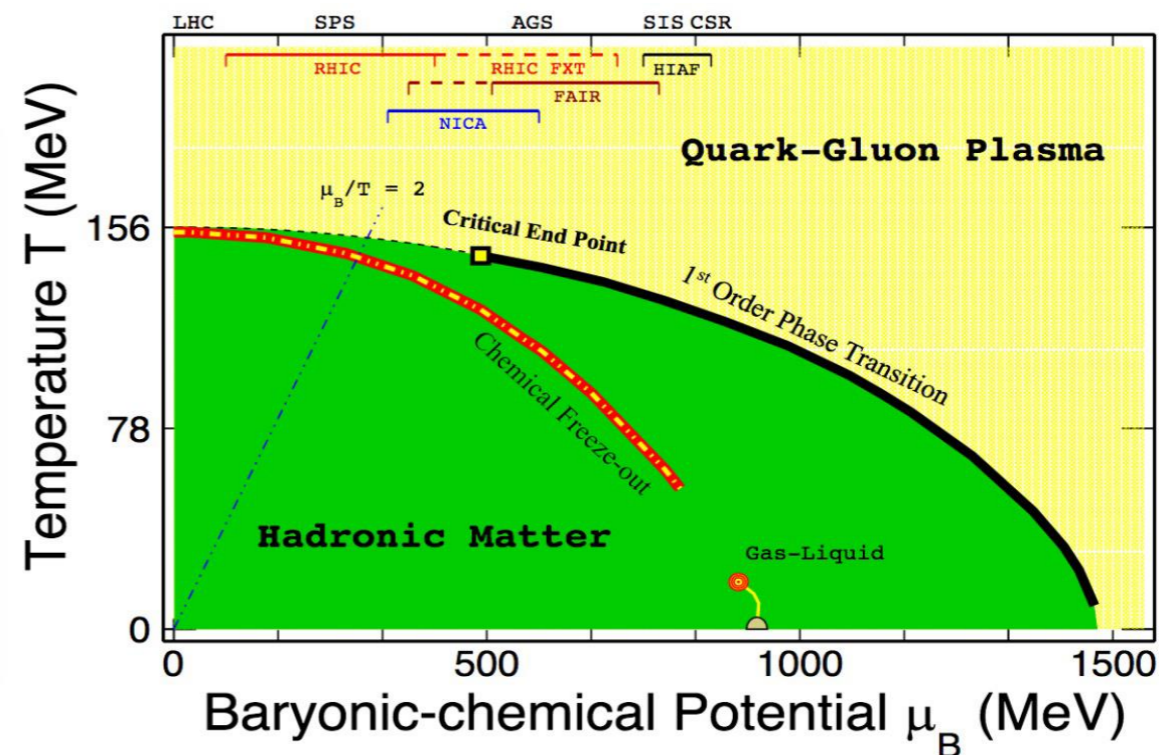


Figure from Bzdak et al., Phys. Rept. '20 & 2015 Long Range Plan



# Experimental signatures potentially based on critical multiplicity fluctuations

- Critical point in QCD phase diagram at finite  $\mu_B$ .  
Experimental conditions reproduced at RHIC and GSI.
- Chiral transition at  $\mu_B = 0$  (which temperature, which order). Experimental conditions reproduced at the LHC.
- Evidence for strong magnetic field at the onset of the deconfined quark gluon phase. LHC measurements.

# The Basics: Relation between susceptibilities, event by event fluctuations and statistical mechanics

Cumulant generating function

$$K_N(t) = \ln \langle e^{tN} \rangle = \sum_{n=1}^{\infty} \kappa_n \frac{t^n}{n!}$$

$$\kappa_n \propto \frac{\partial^n (\ln Z^{\text{gce}})}{\partial \mu^n}$$

Grand partition function

$$\ln Z^{\text{gce}}(T, V, \mu) = \ln \left[ \sum_N e^{\mu N/T} Z^{\text{ce}}(T, V, N) \right]$$

Cumulants measure chemical potential derivatives of the (QCD) equation of state

- Cumulants of conserved quantities

Net-baryon (B) (net-proton as proxy)

Net-electric charge (Q)

Net-strangeness (S) (net-kaon as proxy)

$$\begin{aligned} \delta N &= N - \langle N \rangle & \frac{C_2}{C_1} &= \frac{\sigma^2}{M}, & \frac{C_3}{C_2} &= S\sigma \\ C_1 &= \langle N \rangle = M & & & & \\ C_2 &= \langle (\delta N)^2 \rangle = \sigma^2 & \frac{C_4}{C_2} &= \kappa \sigma^2 & & \\ C_3 &= \langle (\delta N)^3 \rangle & & & & \\ C_4 &= \langle (\delta N)^4 \rangle - 3 \langle (\delta N)^2 \rangle^2 & & & & \end{aligned}$$

- Sensitive to correlation length

$$C_3 = \langle (\delta N)^3 \rangle \sim \xi^{4.5}$$

$$C_4 = \langle (\delta N)^4 \rangle - 3 \langle (\delta N)^2 \rangle^2 \sim \xi^7$$

- Related to susceptibility

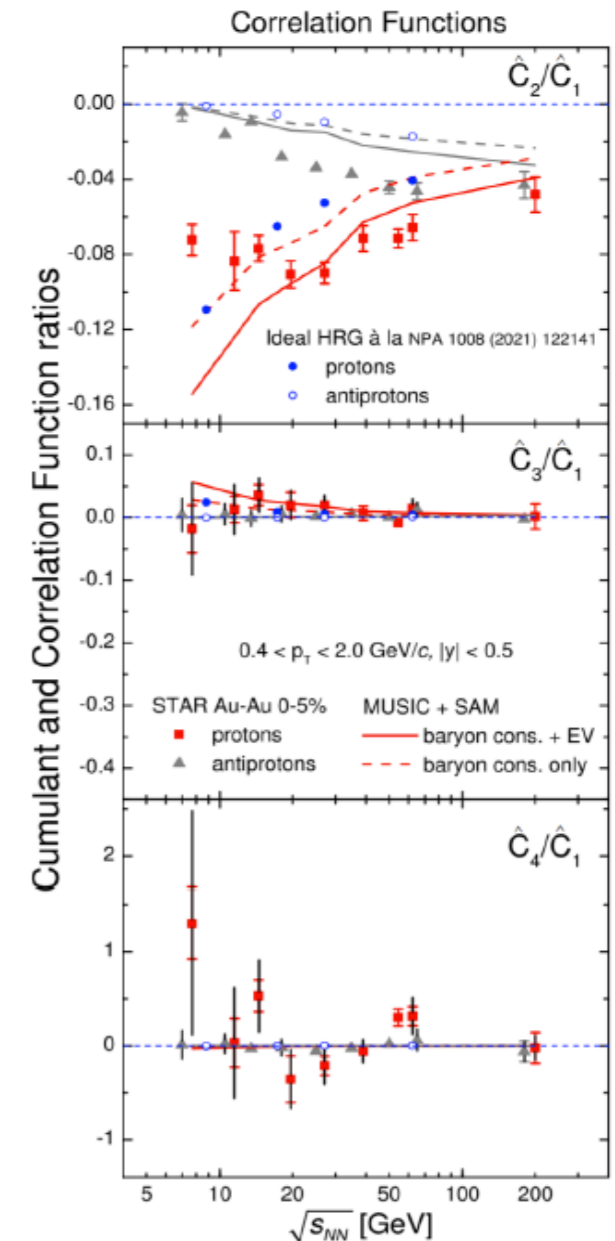
$$\frac{\chi_4^q}{\chi_2^q} = \kappa \sigma^2 = \frac{C_4^q}{C_2^q}, \quad \frac{\chi_3^q}{\chi_2^q} = S\sigma = \frac{C_3^q}{C_2^q}$$

$$\chi_n^q = \frac{1}{VT^3} \cdot C_n^q = \frac{\partial^n (p/T^4)}{\partial (\mu^q)^n}, \quad q = B, Q, S$$

# The experimental improvement: factorial moments

- **Factorial cumulants  $\hat{C}_n$**  [Bzdak, Koch, Strodthoff, PRC 95, 054906 (2017)]
  - Remove the Poisson contribution and probe genuine correlations
 
$$\hat{C}_1 = \kappa_1, \quad \hat{C}_3 = 2\kappa_1 - 3\kappa_2 + \kappa_3,$$

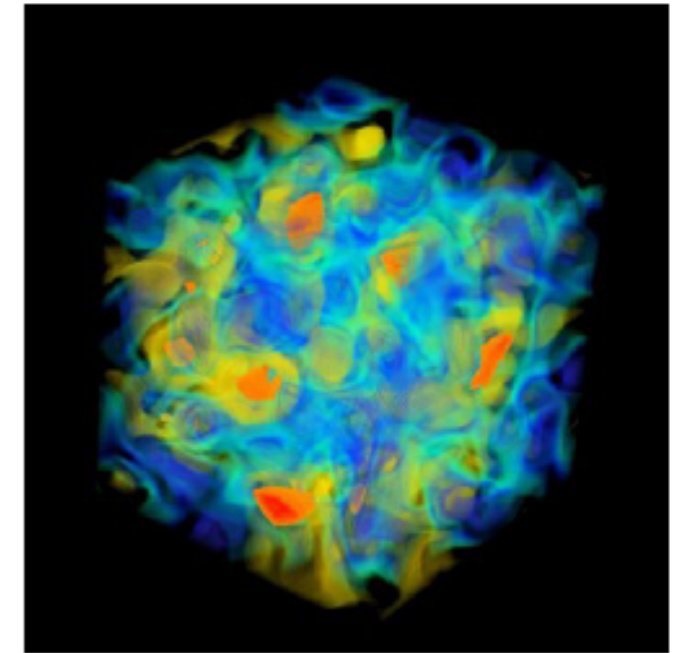
$$\hat{C}_2 = -\kappa_1 + \kappa_2, \quad \hat{C}_4 = -6\kappa_1 + 11\kappa_2 - 6\kappa_3 + \kappa_4.$$
- **Expectation:** High-order ( $n > 3$ ) factorial cumulants
  - have small contributions from non-critical effects [Bzdak, Koch, Skokov, EPJC '17; VV et al, PLB '17]
  - are as singular as ordinary cumulants near the critical point [Ling, Stephanov, PRC '16]
- **Observations from STAR data:**
  - $\hat{C}_3$  &  $\hat{C}_4$  are largely consistent with zero within errors
    - Reanalyze (non-)monotonic energy dependence for  $\hat{C}_4/\hat{C}_1$  instead of  $\kappa_4/\kappa_2$ ?
  - Statistically significant effects appear to be driven by two-proton correlations in  $\hat{C}_2$



⚠ **Notation:** We use  $\kappa_n$  for cumulants and  $\hat{C}_n$  for factorial cumulants, STAR uses the opposite

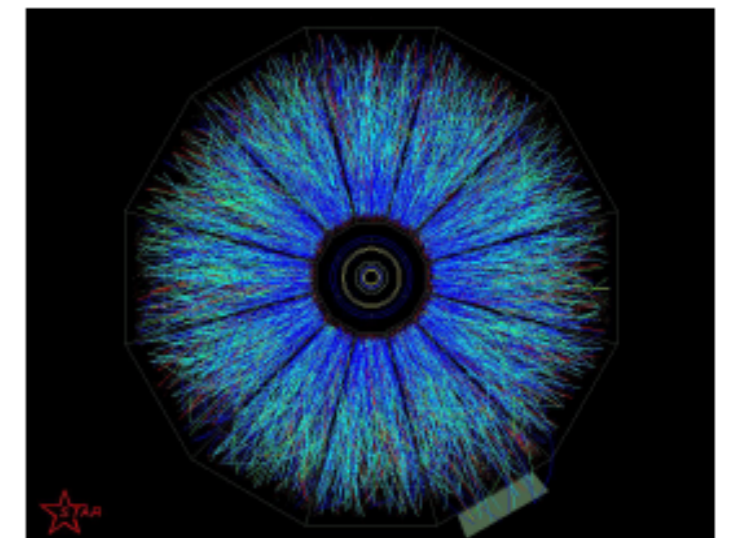
# Theory vs. experiment: the caveats

## Theory



© Lattice QCD@BNL

## Experiment



STAR event display

- **Grand-canonical heat bath vs expansion in vacuum**

- Vovchenko et al., PLB 811, 135868 (2020), PRC 105, 014903 (2022)
- Bzdak et al., PRC 87, 014901 (2013)

- **Coordinate vs. momentum space**

- Ling et al., PRC93, 034915 (2016)

- **Conserved charges (net-baryon) vs. proxies (net-proton)**

- Kitazawa et al., PRC 85, 021901 (2012)
- Koch et al., PRC 103, 044903 (2021)

- **Volume, initial state, baryon stopping: fixed vs. fluctuating**

- Gorenstein et al., PRC 84, 014904 (2011)
- Skokov et al., PRC 88, 034911 (2013)

- **Hadronic phase**

- Steinheimer et al., PLB 776, 32 (2018)
- Savchuk et al. PLB 827, 136983 (2022)

- **Non-equilibrium (memory) effects**

- Mukherjee et al., PRC92, 034912 (2015)
- Asakawa et al., PRC 101, 034913 (2020)

# Topic 1: Critical Point

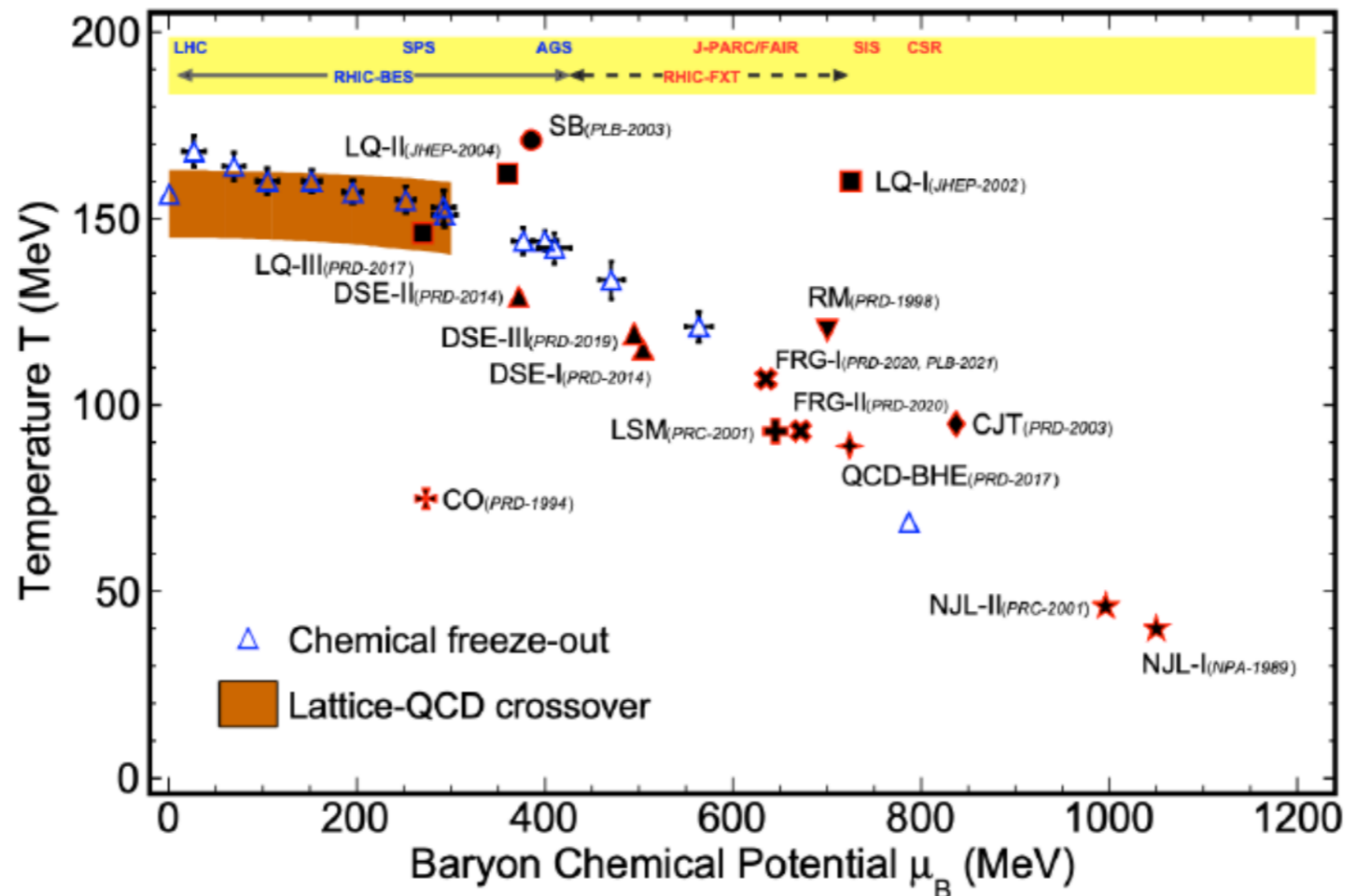


Figure adapted from A. Pandav, D. Mallick, B. Mohanty, Prog. Part. Nucl. Phys. 125 (2022)

Until recently: no viable critical point calculation based on lattice QCD alone

Including the possibility that the QCD critical point does not exist at all

de Forcrand, Philipsen, JHEP 01, 077 (2007); VV, Steinheimer, Philipsen, Stoecker, PRD 97, 114030 (2018)

# Sign problem overcome by Standard Taylor Expansion

## Taylor Expansion around $\mu_B = 0$

$$\frac{n_B(T, \mu_B)}{T^3} = \sum_{n=0}^{\infty} \frac{1}{(2n-1)!} \chi_{2n}(T, \mu_B = 0) \left(\frac{\mu_B}{T}\right)^{2n-1}$$

[Borsanyi, S. et al High Energy Physics.9(8), 1-16.(2012)]

[Bazavov, A et al PhysRevD.95, 054504 (2017)]

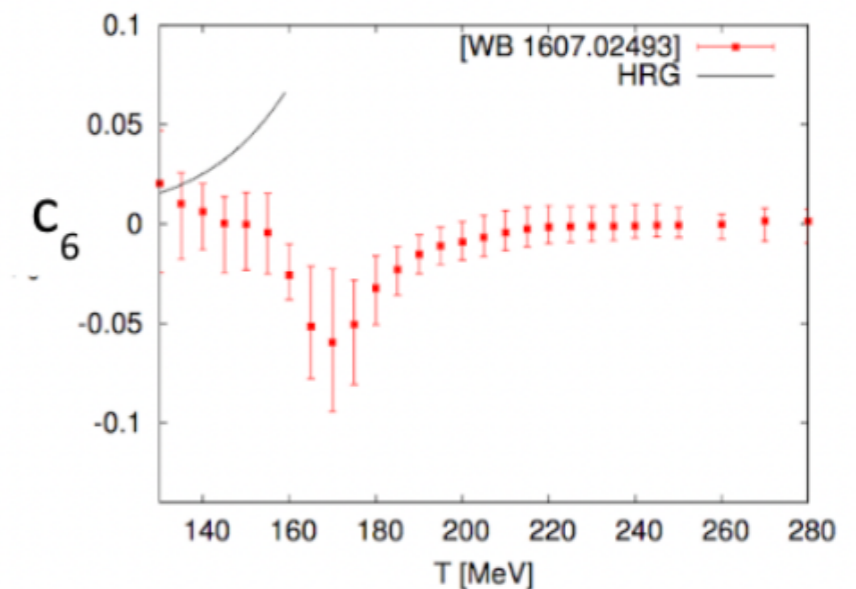
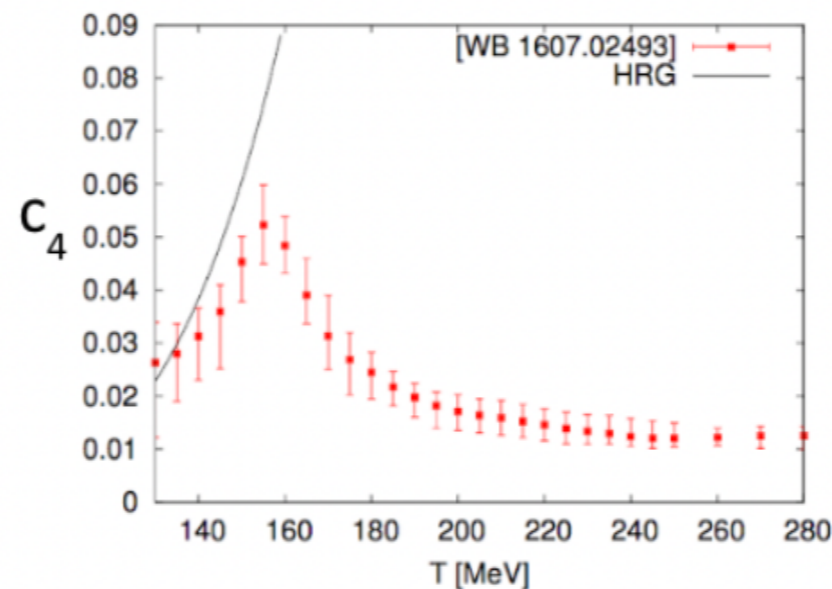
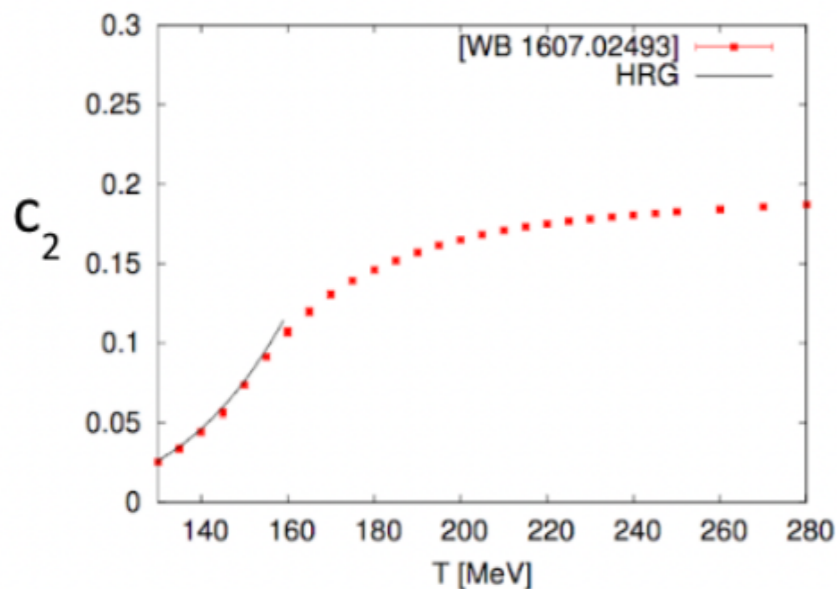
$$c_n(T) = \frac{\chi_n^B(T, \mu_B = 0)}{n!} = \frac{1}{n!} \left( \frac{\partial}{\partial(\mu_B/T)} \right)^n (P/T^4) \Big|_{\mu_B=0}$$

## Limitations

- Currently limited to  $\frac{\mu_B}{T} \leq 3$  despite great computational power
- Including one more higher-order term does not remove unphysical behavior due to truncation of Taylor series

[Bollweg, D. et al Phys.Rev.D 108 (2023) 1, 014510]

[Borsanyi, S et al arXiv:2312.07528v1. (2023)]



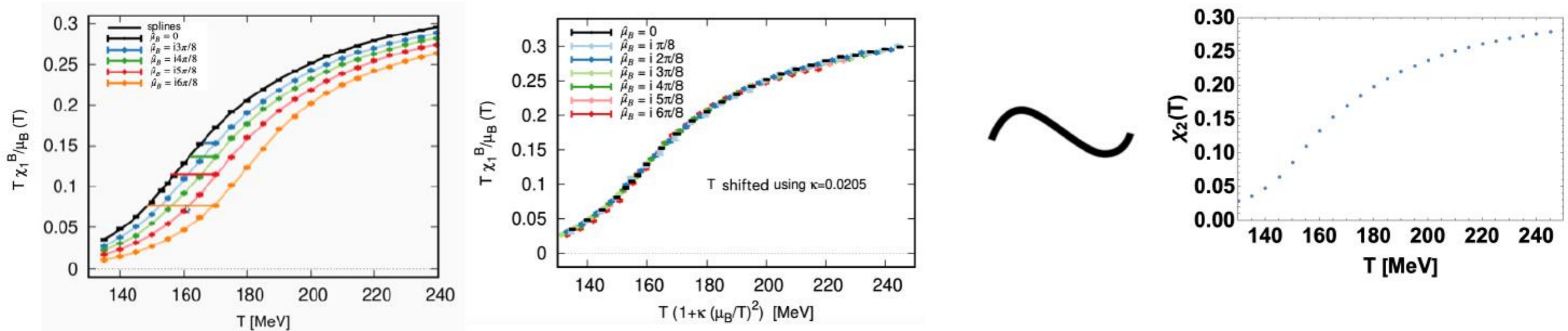
[WB Lattice QCD Collaboration]



# Promising new expansion scheme

## T' Expansion scheme (T ExS)

Simulating at Imaginary  $\mu_B$



[Borsányi, S et al PRL. 108(1), 101.034901(2021)]

$$T \frac{\chi_1^B(T, \mu_B)}{\mu_B} = \chi_2^B(T', 0)$$

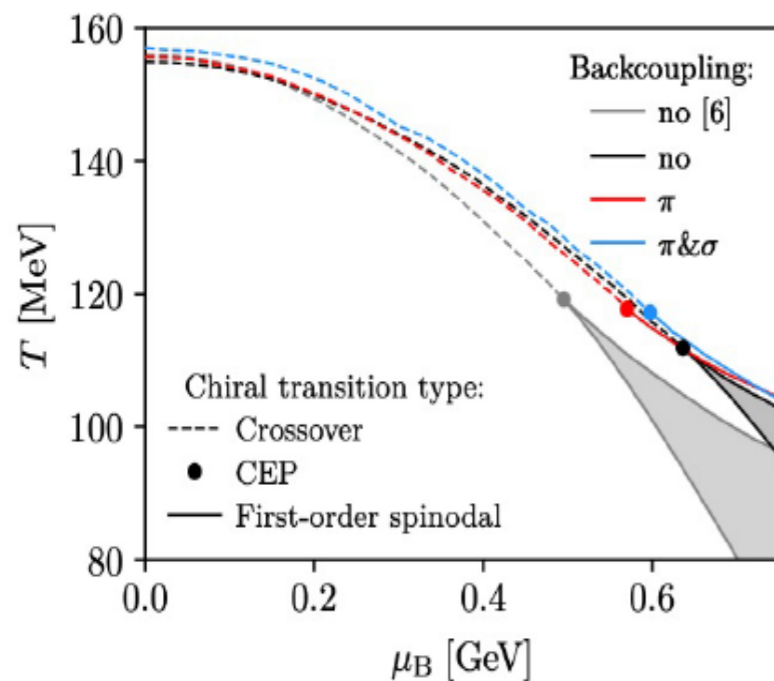
$$\chi_n^B(T, \mu_B) = \frac{1}{n!} \left( \frac{\partial}{\partial(\mu_B/T)} \right)^n (P/T^4)$$

$$T'(T, \mu_B) = T \left[ 1 + \kappa_2^{BB}(T) \left( \frac{\mu_B}{T} \right)^2 + \kappa_4^{BB}(T) \left( \frac{\mu_B}{T} \right)^4 + \mathcal{O} \left( \frac{\mu_B}{T} \right)^6 \right]$$

- $\mu_B$  dependence is captured in T-rescaling.
- Trusted up to  $\frac{\mu_B}{T} = 3.5$

# Updated predictions (anchored to lattice QCD at $\mu_B=0$ )

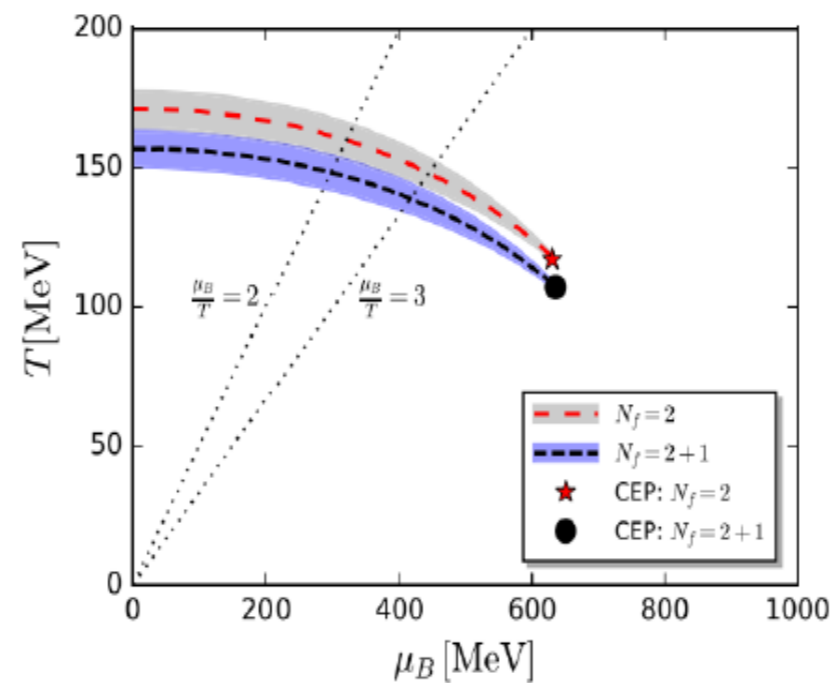
Dyson-Schwinger equations



Gunkel, Fischer, PRD 104, 054202 (2021)

$T \sim 120$  MeV  $\mu_B \sim 600$  MeV

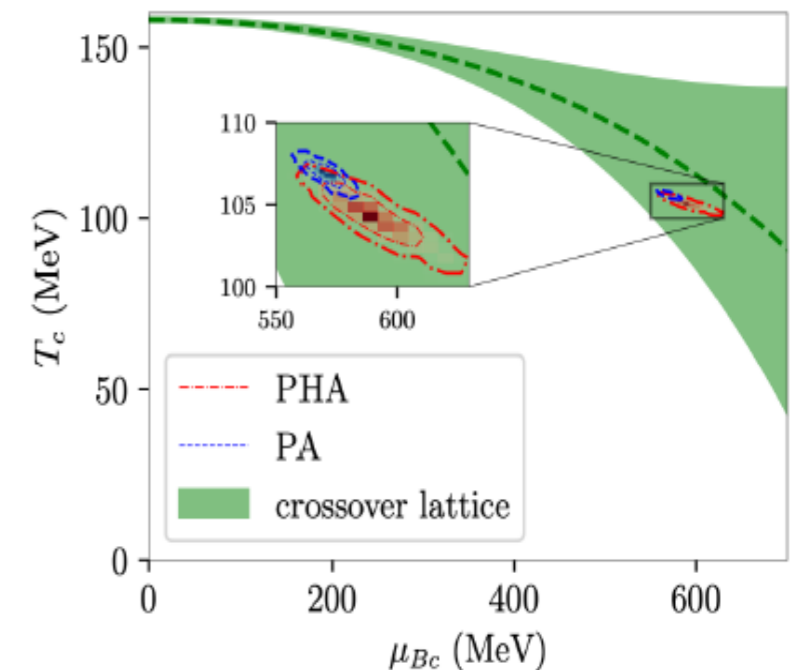
Functional renormalization group



Fu, Pawłowski, Rennecke, PRD 101, 053032 (2020)

$T \sim 100$  MeV  $\mu_B \sim 600 - 650$  MeV

Black-hole engineering



Hippert et al., arXiv:2309.00579

$T \sim 105$  MeV  $\mu_B \sim 580$  MeV

All in excellent agreement with lattice QCD at  $\mu_B = 0$   
and predict QCD critical point in a similar ballpark of  $\mu_B/T \sim 5-6$

If true, reachable in heavy-ion collisions at  $\sqrt{s_{NN}} \sim 3 - 5$  GeV

# Interesting new lattice results using Lee-Yang edge singularities

D.A. Clarke et al., 2405.10196

Similar results: G. Basar, 2312.06952

Multi-point Pade approximation

AIC - Akaike information criterion

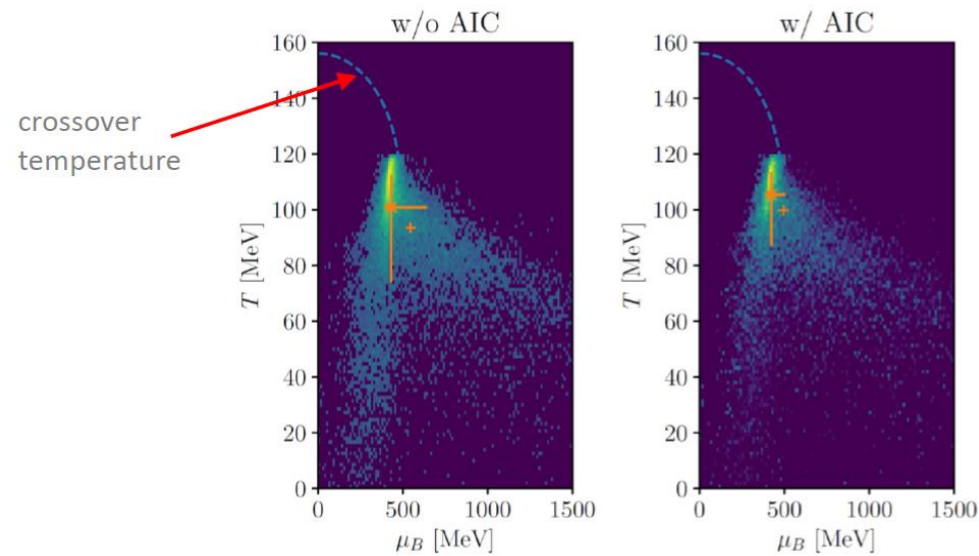
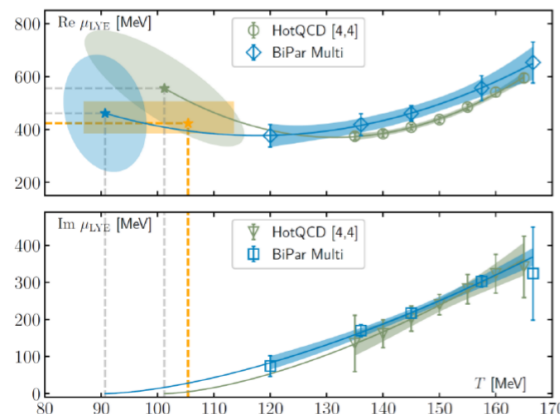


FIG. 1. Probability distribution of the QCD critical point from extrapolating Lee-Yang singularities to the real domain using universal scaling. For a detailed description see the text.

We analyze the trajectory of the Lee-Yang edge singularities of the QCD equation of state in the complex baryon chemical potential ( $\mu_B$ ) plane... By extrapolating from this information, we estimate for the location of the QCD critical point,  $T_c \approx 100 \text{ MeV}$ ,  $\mu_c \approx 580 \text{ MeV}$ .

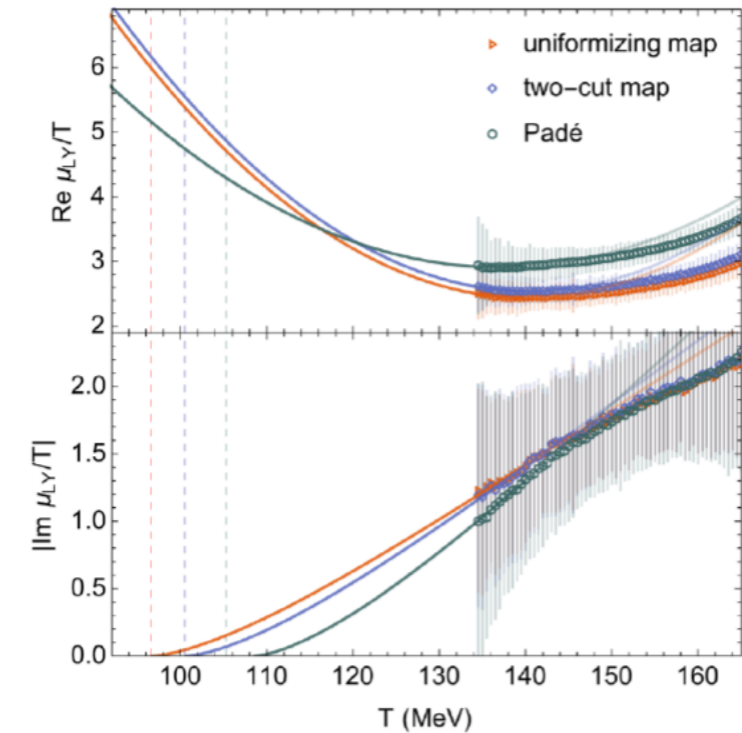
Using the multi-point Pade approach, the Lee-Yang edge singularities are located in the QCD pressure in the complex baryon chemical potential.

$$(T^{\text{CEP}}, \mu_B^{\text{CEP}}) = (105_{-18}^{+8}, 422_{-35}^{+80}) \text{ MeV}$$



“An important limitation of the estimate presented here is that they are not fully extrapolated to the continuum limit”

Can we conclude that there is the QCD critical point?



# A more dynamical phenomenological approach

## 1. Dynamical model calculations of critical fluctuations [X. An et al., Nucl. Phys. A 1017, 122343 (2022)]

- Fluctuating hydrodynamics (hydro+) and (non-equilibrium) evolution of fluctuations
- Equation of state with a tunable critical point [P. Parotto et al, PRC 101, 034901 (2020); J. Karthein et al., EPJ Plus 136, 621 (2021)]
- Generalized Cooper-Frye particlization [M. Pradeep, et al., PRD 106, 036017 (2022); PRL 130, 162301 (2023)]

Alternatives at high  $\mu_B$ : hadronic transport/molecular dynamics with a critical point

[A. Sorensen, V. Koch, PRC 104, 034904 (2021); V. Kuznietsov et al., PRC 105, 044903 (2022)]

## 2. Deviations from precision calculations of non-critical fluctuations

- Non-critical baseline is not flat [Braun-Munzinger, Rustamov, Stachel, NPA 1008, 122141 (2021)]
- Include essential non-critical contributions to (net-)proton number cumulants
- Exact **baryon conservation** + **hadronic interactions** (hard core repulsion)
- Based on realistic hydrodynamic simulations tuned to bulk data

[VV, C. Shen, V. Koch, Phys. Rev. C 105, 014904 (2022)]

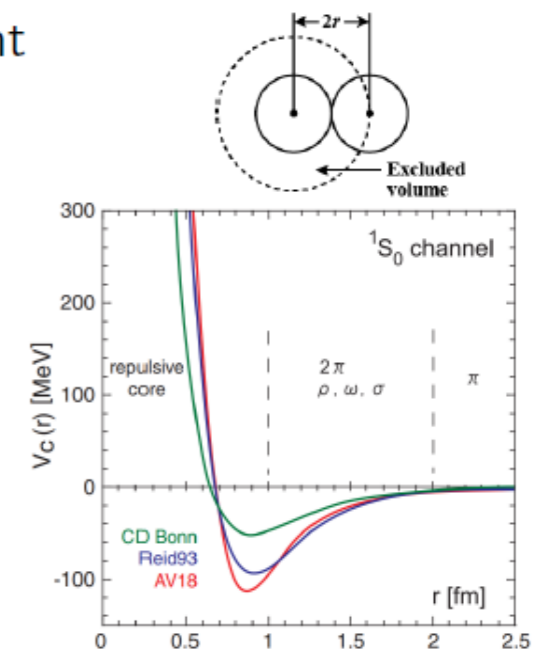
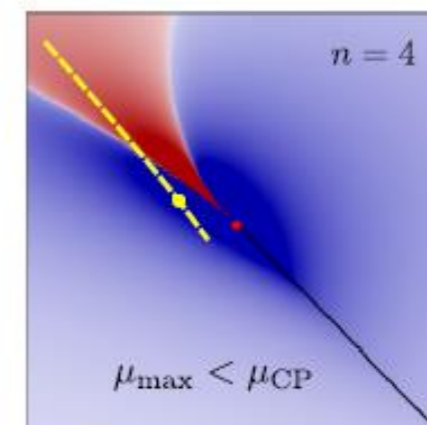
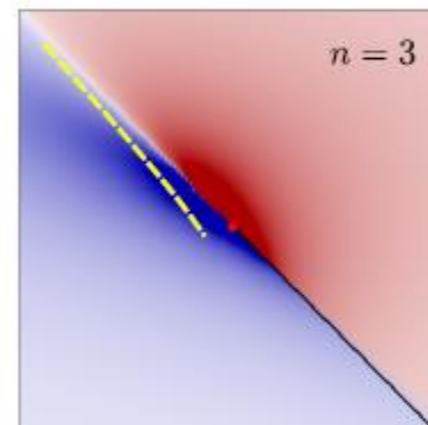
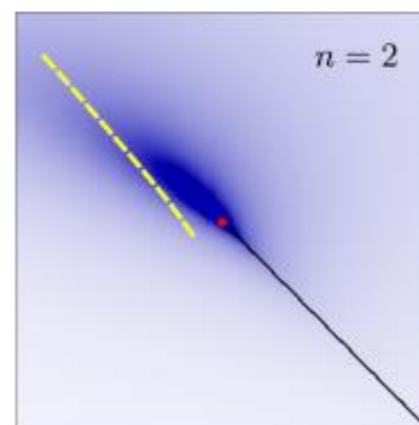


Figure from Ishii et al., PRL '07

Review by A. Bzdak  
(1906.00936)

(universal EOS) critical  $\chi_n$ :



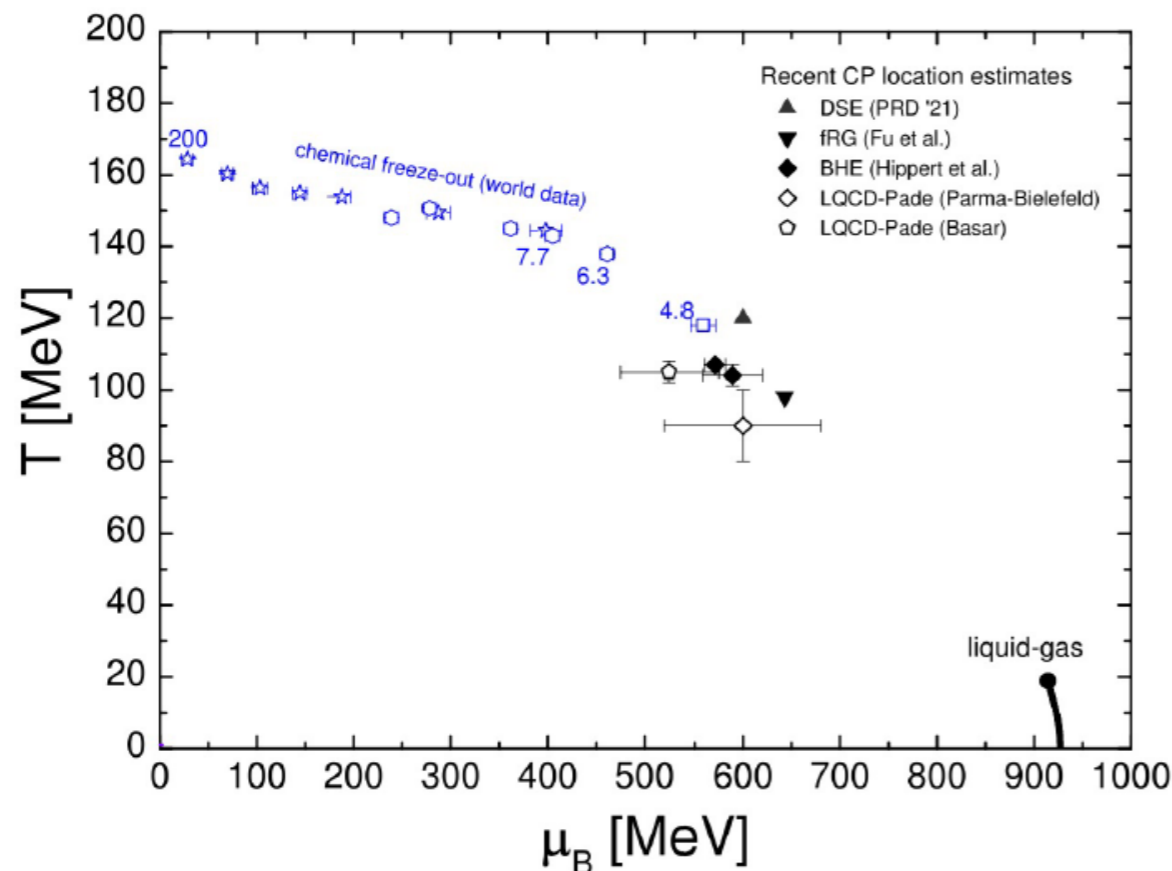
# Experiment: Heavy Ion Collisions

## Control parameters

- Collision energy  $\sqrt{s_{NN}} = 2.4 - 5020$  GeV
  - Scan the QCD phase diagram
- Size of the collision region
  - Expect stronger signal in larger systems

## Measurements

- Final hadron abundances and momentum distributions **event-by-event**



STAR BES-I Program: Au+Au collisions

$\sqrt{s_{NN}}$ (GeV)	Events / $10^6$	$\mu_B$ (MeV)
200	220	25
62.4	43	75
54.4	550	85
39	92	112
27	31	156
19.6	14	206
14.5	14	264
11.5	7	315
7.7	3	420
3.0	140	750

$\mu_B$  from J. Cleymans et al, PRC 73,034905(2006)

STAR BES Program:  
Au+Au collisions

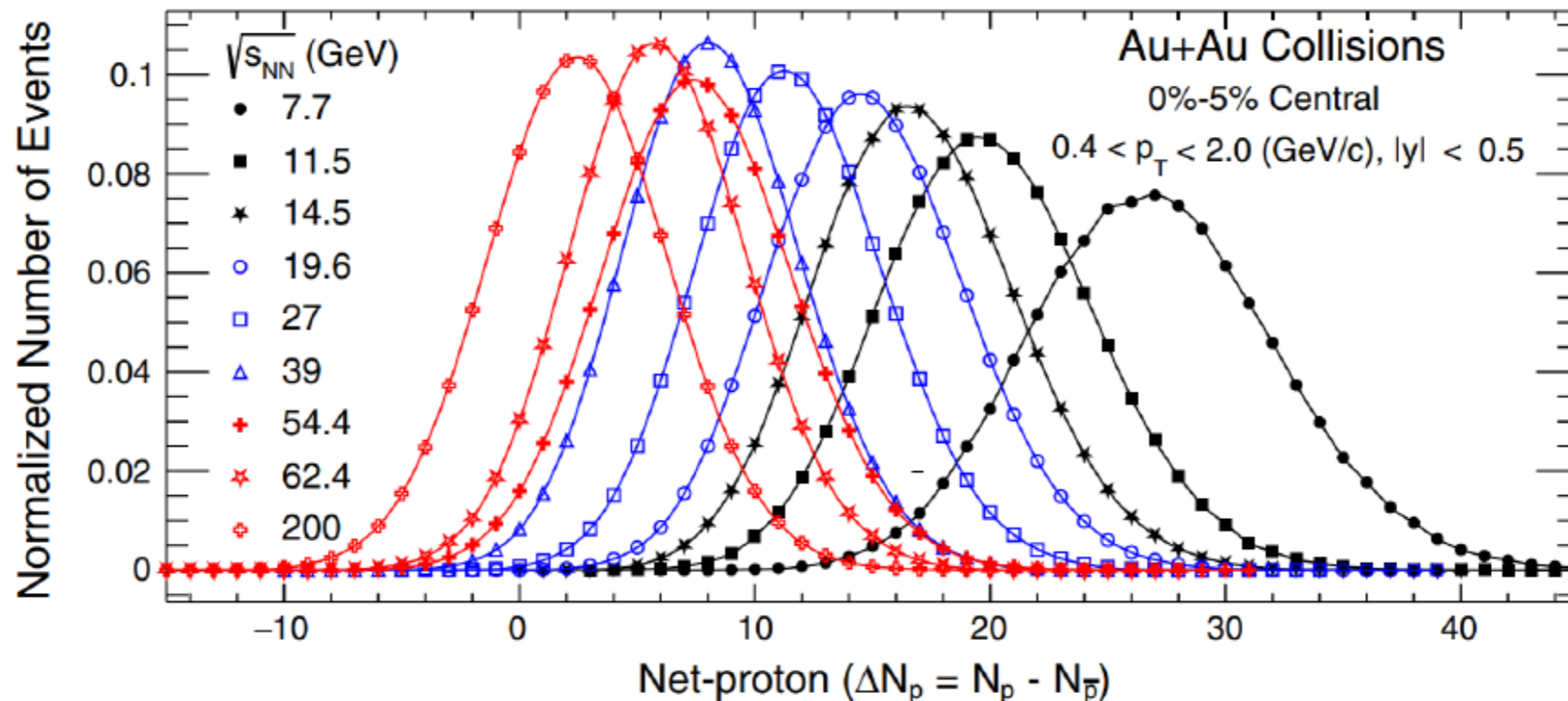
$\sqrt{s_{NN}}$ (GeV)	Events BES-I ( $10^6$ )	Events BES-II ( $10^6$ )
7.7	3	45
9.2	-	78
11.5	7	110
14.6	20	178
17.3	-	116
19.6	15	270
27	30	220

# Measuring cumulants in event by event multiplicity distributions

(particles are proxies for conserved quantum numbers)

Count the number of events with given number of e.g. (net) protons  $P(\Delta N_p) \sim \frac{N_{\text{events}}(\Delta N_p)}{N_{\text{events}}^{\text{total}}}$

STAR Collaboration, Phys. Rev. Lett. 126, 092301 (2021)



Cumulants are extensive,  $\kappa_n \sim V$ , use ratios to cancel out the volume

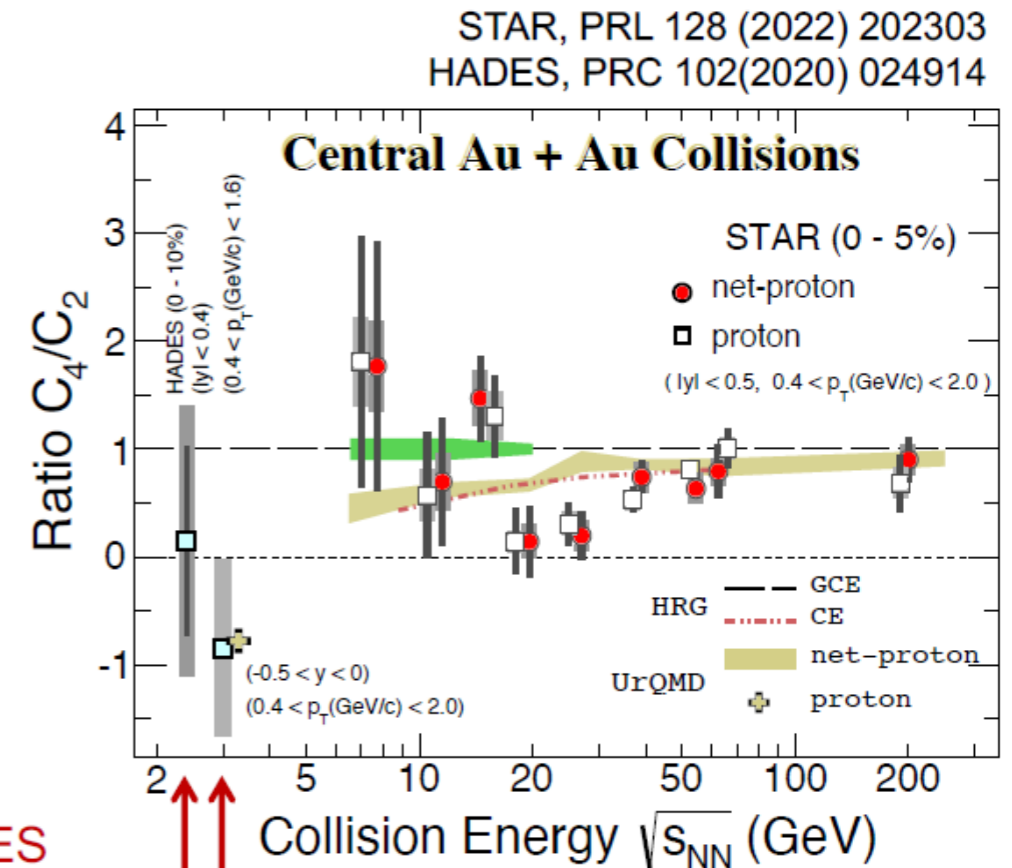
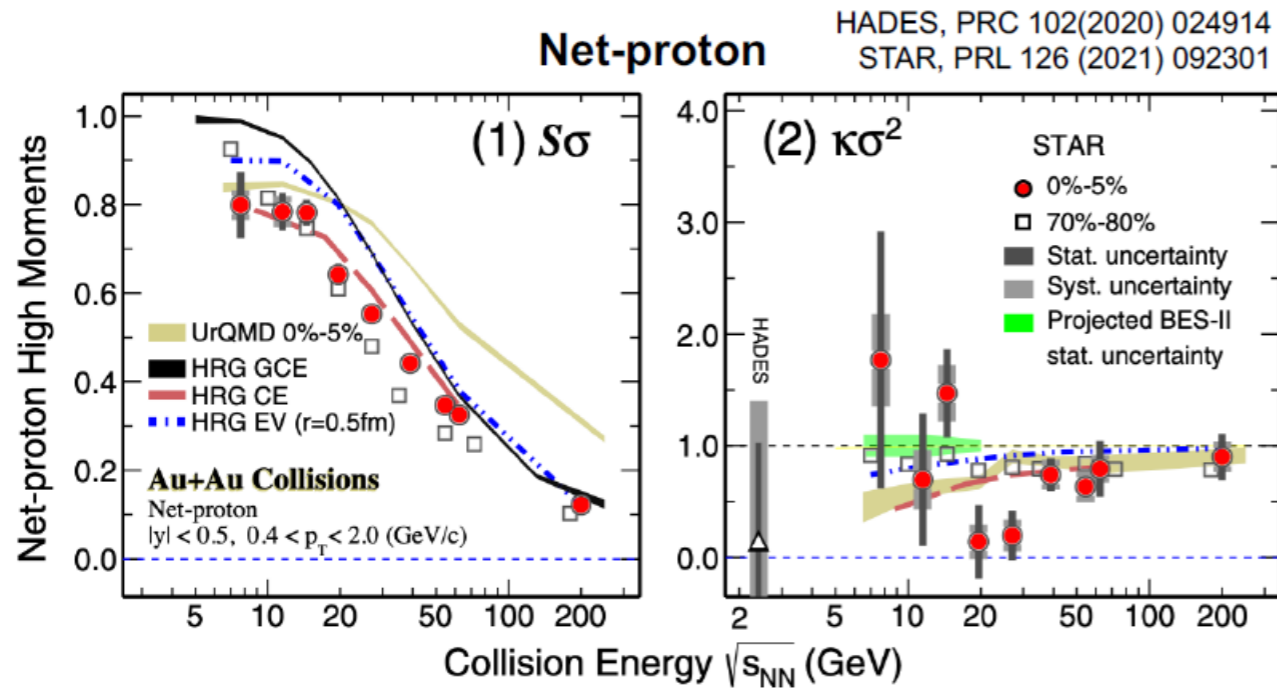
$$\frac{\kappa_2}{\langle N \rangle}, \quad \frac{\kappa_3}{\kappa_2}, \quad \frac{\kappa_4}{\kappa_2}$$

Look for subtle critical point signals (tails of the distribution)

# Experimental results: cumulant ratios

HADES, STAR BES-I (2020, 2021)

HADES, STAR BES-II FXT (2020,2022)



1. Full measurement on BES-I datasets
2. With TOF detector,  $p_T$  coverage is extended to 2.0 GeV/c
3. **Non-monotonic energy dependence trend is observed with  $3.1\sigma$  significance**

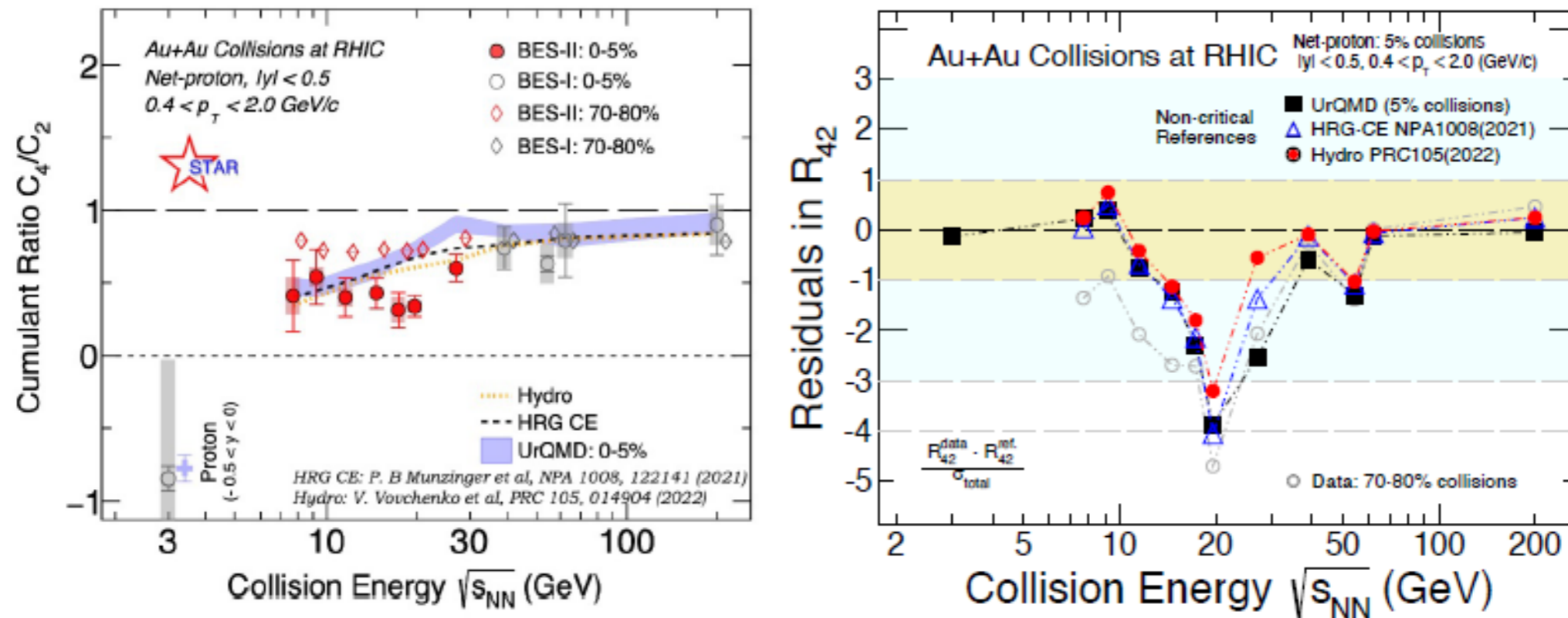
HADES  
 $\sqrt{s_{NN}} = 2.4$  GeV  
 Au+Au collisions

STAR  
 $\sqrt{s_{NN}} = 3.0$  GeV  
 Au+Au collisions

STAR 3 GeV and HADES data indicate a purely hadronic system

# New STAR results (CPOD 2024, SQM2024)

## Precision Measurement on BES-II

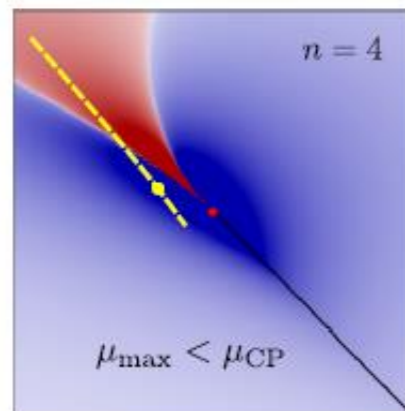
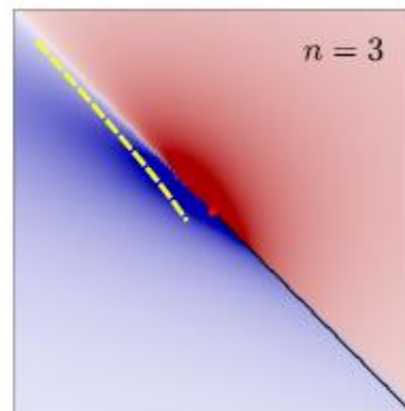
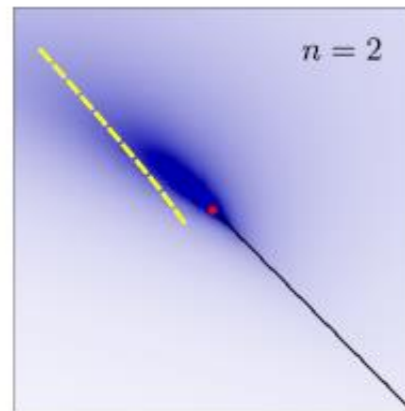


1. Net-proton  $C_4/C_2$  in 0-5% shows a minimum around ~20 GeV comparing to non-CP models (Hydro, HRG, UrQMD) and 70-80% data.
2. Maximum deviation:  $3.2-4.7\sigma$  at 20 GeV ( $1.3-2\sigma$  at BES-I)
3. Overall deviation from 7.7-27 GeV:  $1.9-5.4\sigma$  ( $1.4-2.2\sigma$  at BES-I)

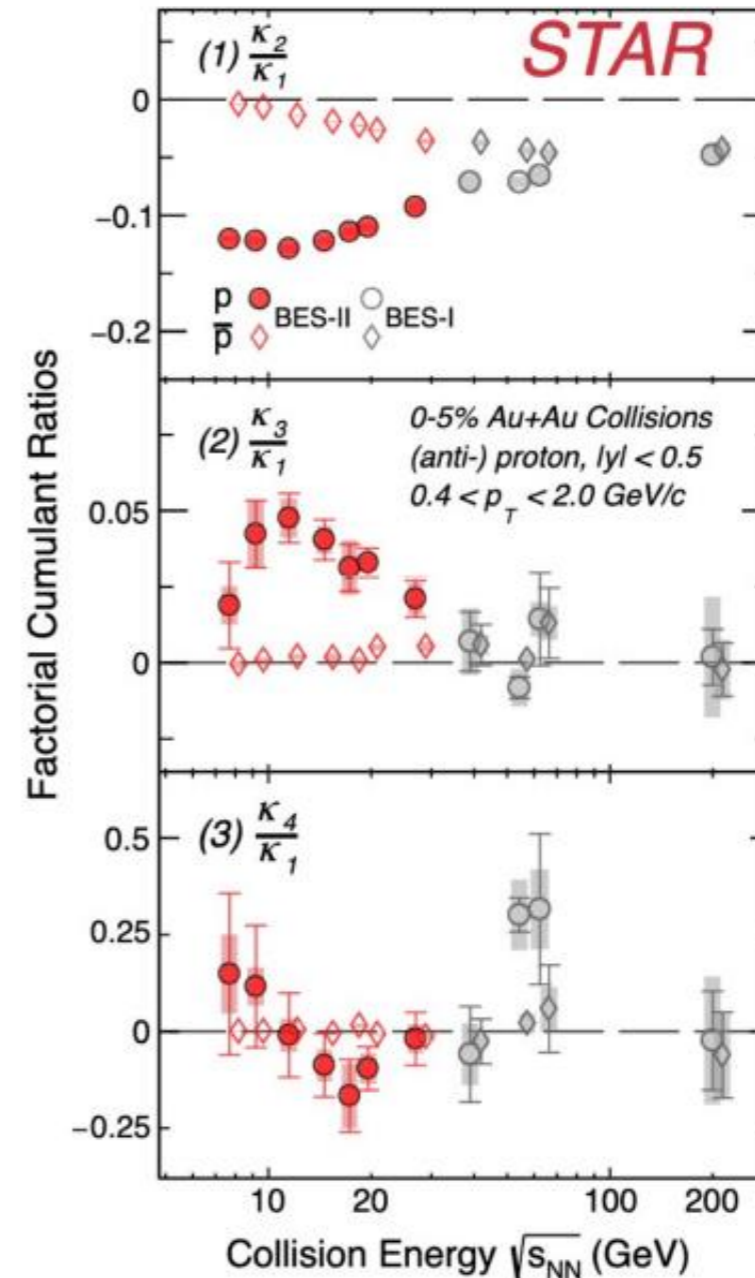
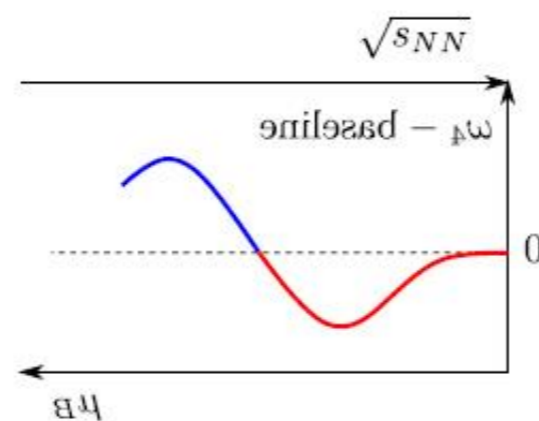
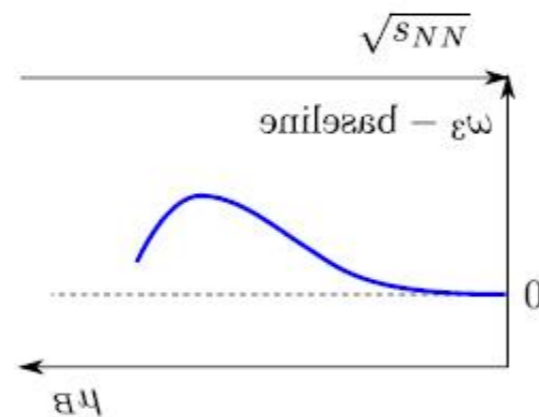
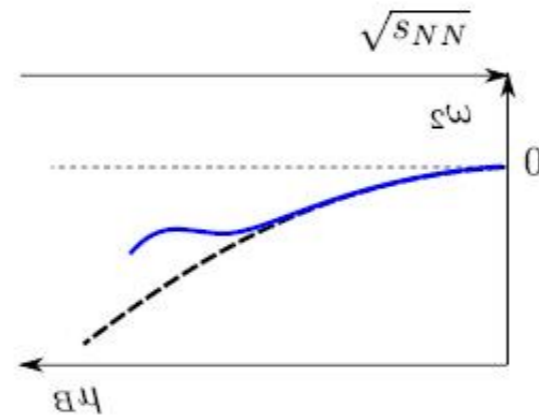


# Comparison to theory: no smoking gun but still intriguing structures in all measured cumulants

(universal EOS) critical  $\chi_n$ :



(irreducible correlations)  $FC_n [N_p] \sim \chi_n$  (Pradeep, MS 2211.09142),  $\omega_n \equiv FC_n / FC_1$

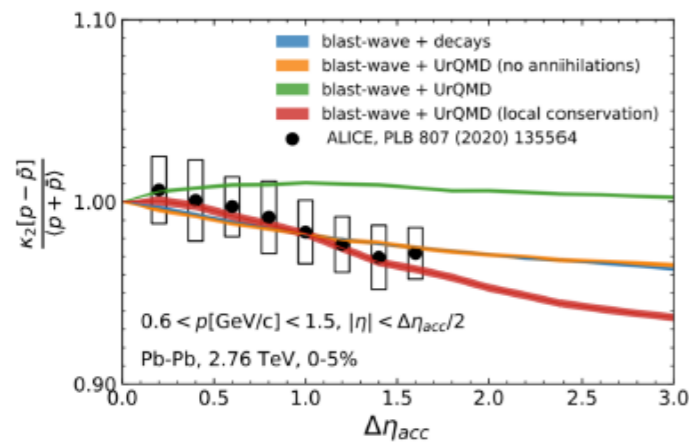


Bzdak et al review 1906.00936

Expected signatures: **bump** in  $\omega_2$  and  $\omega_3$ , **dip** then **bump** in  $\omega_4$  for CP at  $\mu_B > 420$  MeV

# Topic 2: Chiral transition at $\mu_B = 0$

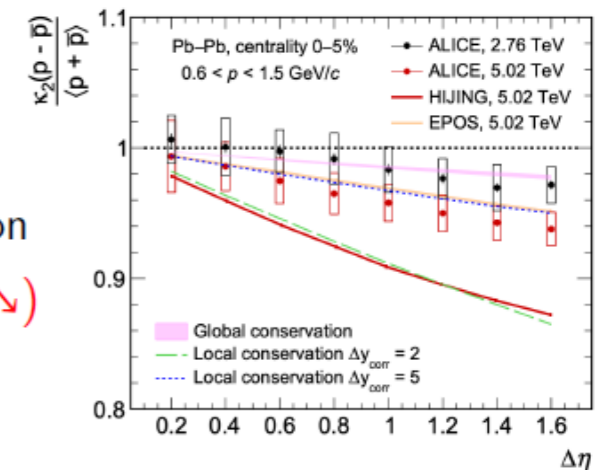
At LHC energies  $\mu_B=0$ , the lower cumulants are well understood by baryon number conservation



$$\kappa_2[p - \bar{p}] / \langle p + \bar{p} \rangle:$$

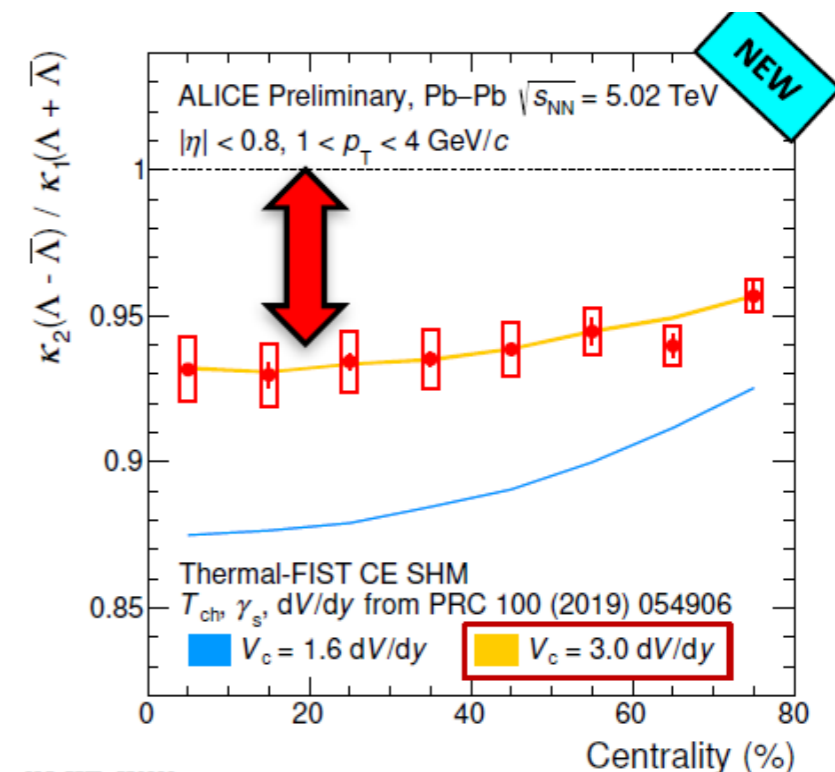
- Largely understood as (global) baryon conservation
  - Larger suppression at 5 TeV contrary to naïve expectation
- Interplay: baryon annihilation(↗) vs local conservation(↘)
  - Additional measurement of  $\kappa_2[p + \bar{p}]$  can resolve it

O. Savchuk et al., PLB 827, 136983 (2022)



Deviation from Skellam allows you to calculate correlation volume. Large volume = early correlations

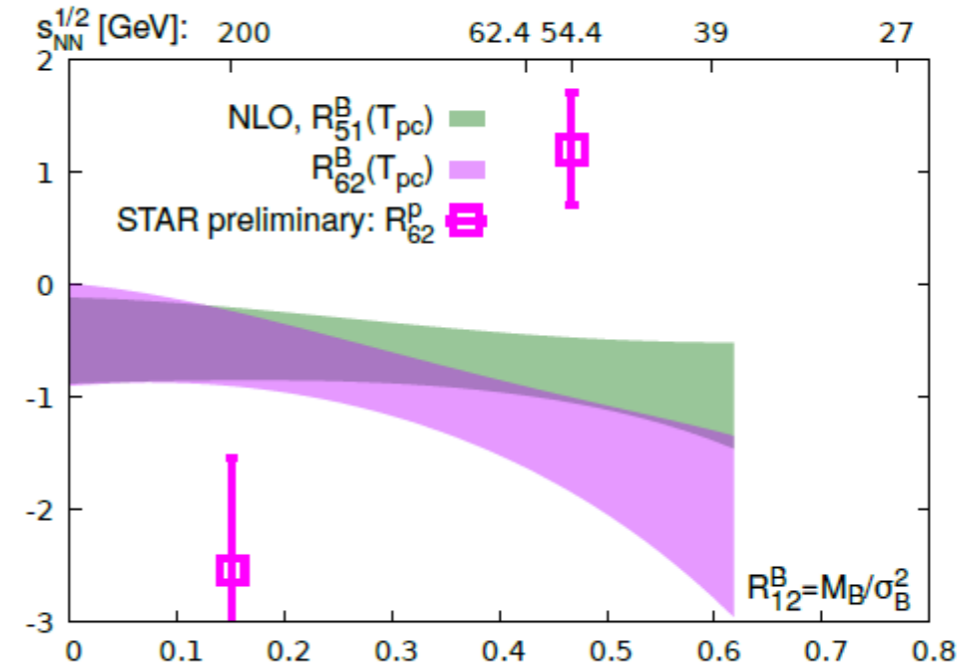
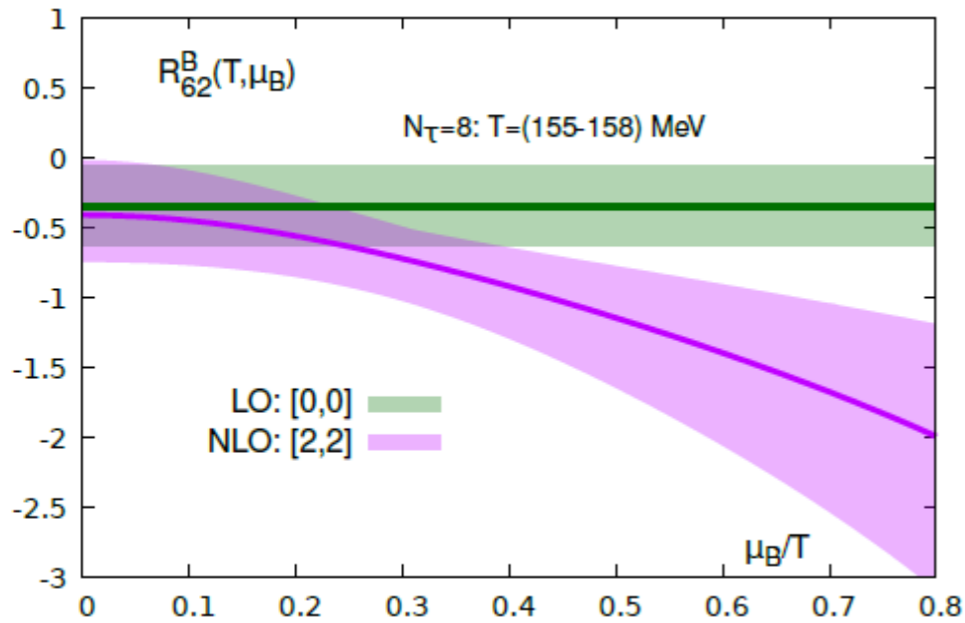
Strong early baryon correlations ( $\Lambda, p$ ) (ALICE, SQM 2024)



ALI-PREL-570339

# Lattice QCD predictions on chiral criticality in higher order cumulants

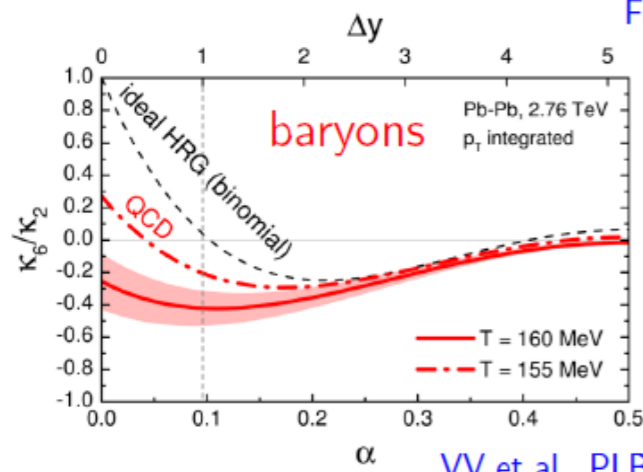
A. Bazavov et al. (HotQCD), Phys.Rev.D 101 (2020) 7, 074502: [c6/c2 goes negative](#)



ALICE Collaboration, PLB 844, 137545 (2023)

High-order cumulants: probe remnants of chiral criticality

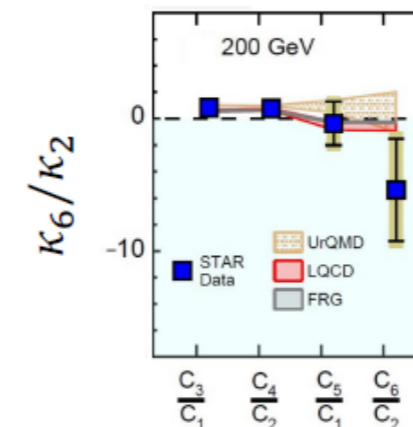
Friman et al., EPJC 71, 1694 (2011)



- negative  $\kappa_6$  of baryons

VV et al., PLB 811, 135868 (2020)

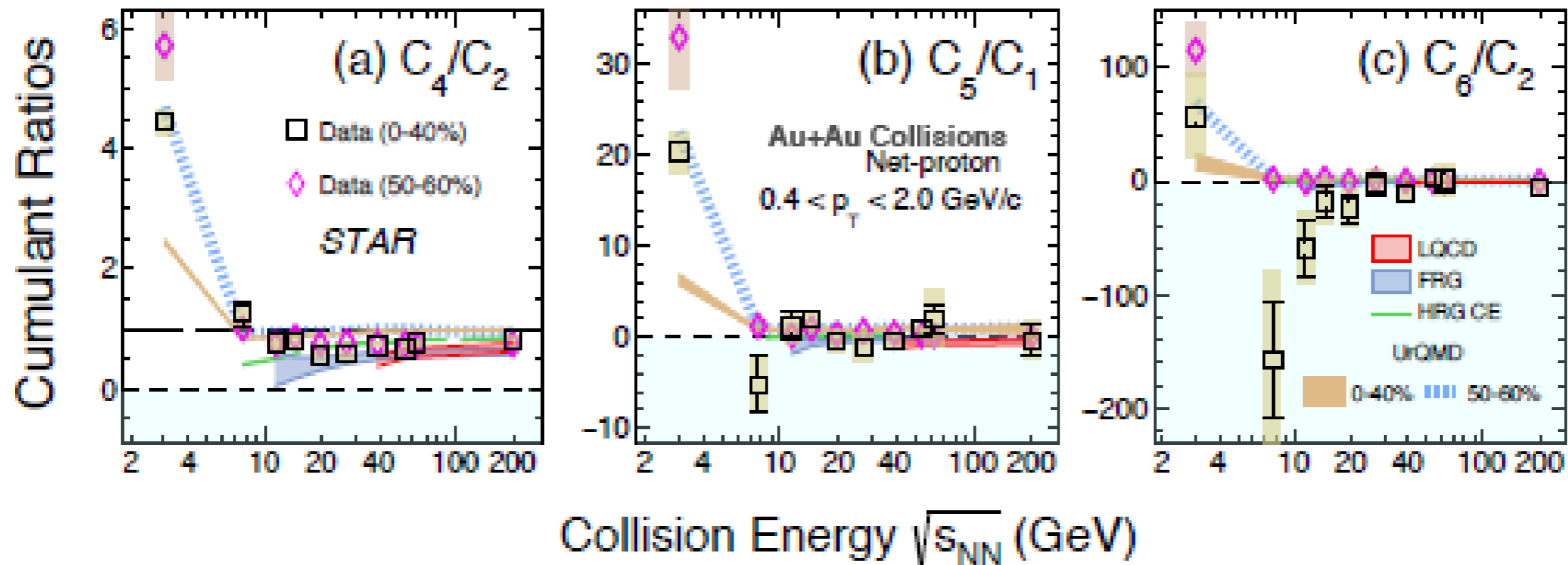
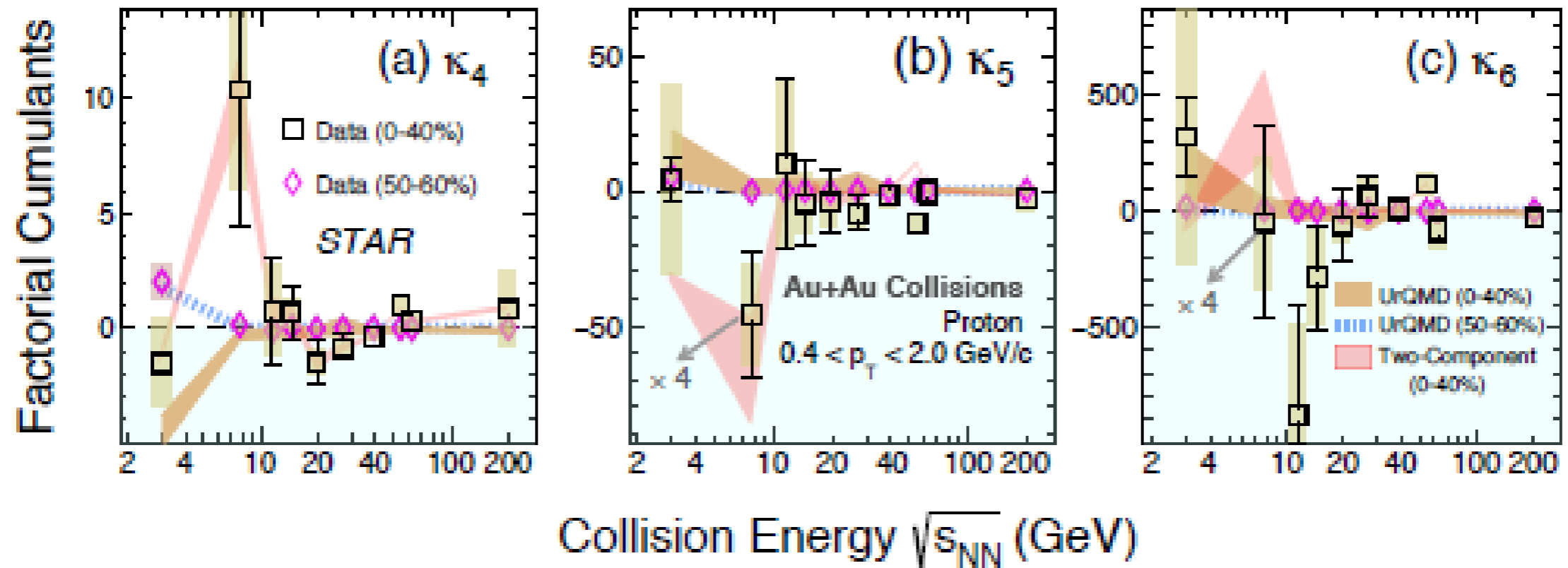
RHIC 200 GeV: hints of negative  $\kappa_6 < 0$  (protons)



- are baryons even more negative?

STAR Collaboration, PRL 130, 082301 (2023)

# Detailed experimental results (STAR, 2207.09837)



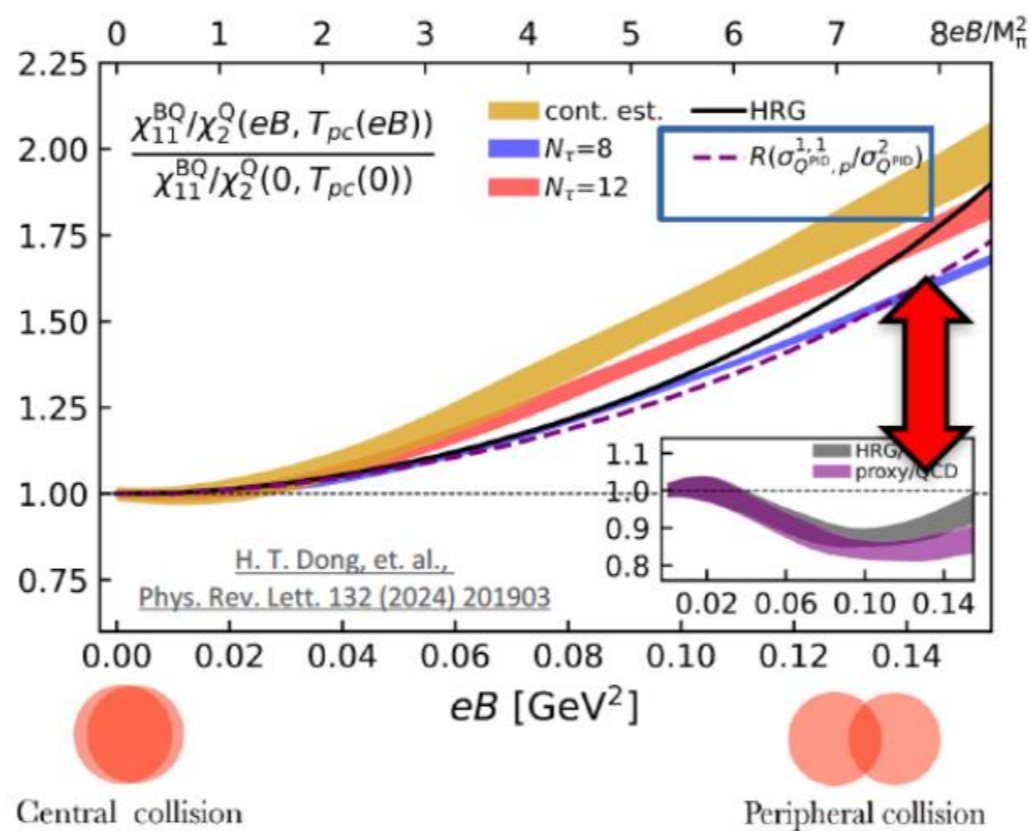
# Topic 3: Conserved charge fluctuations to probe early magnetic field

H.-T. Ding et al., PRL 132, 201903 (2024)

Magnetic field + LQCD EOS → Any modification in chiral susceptibilities?

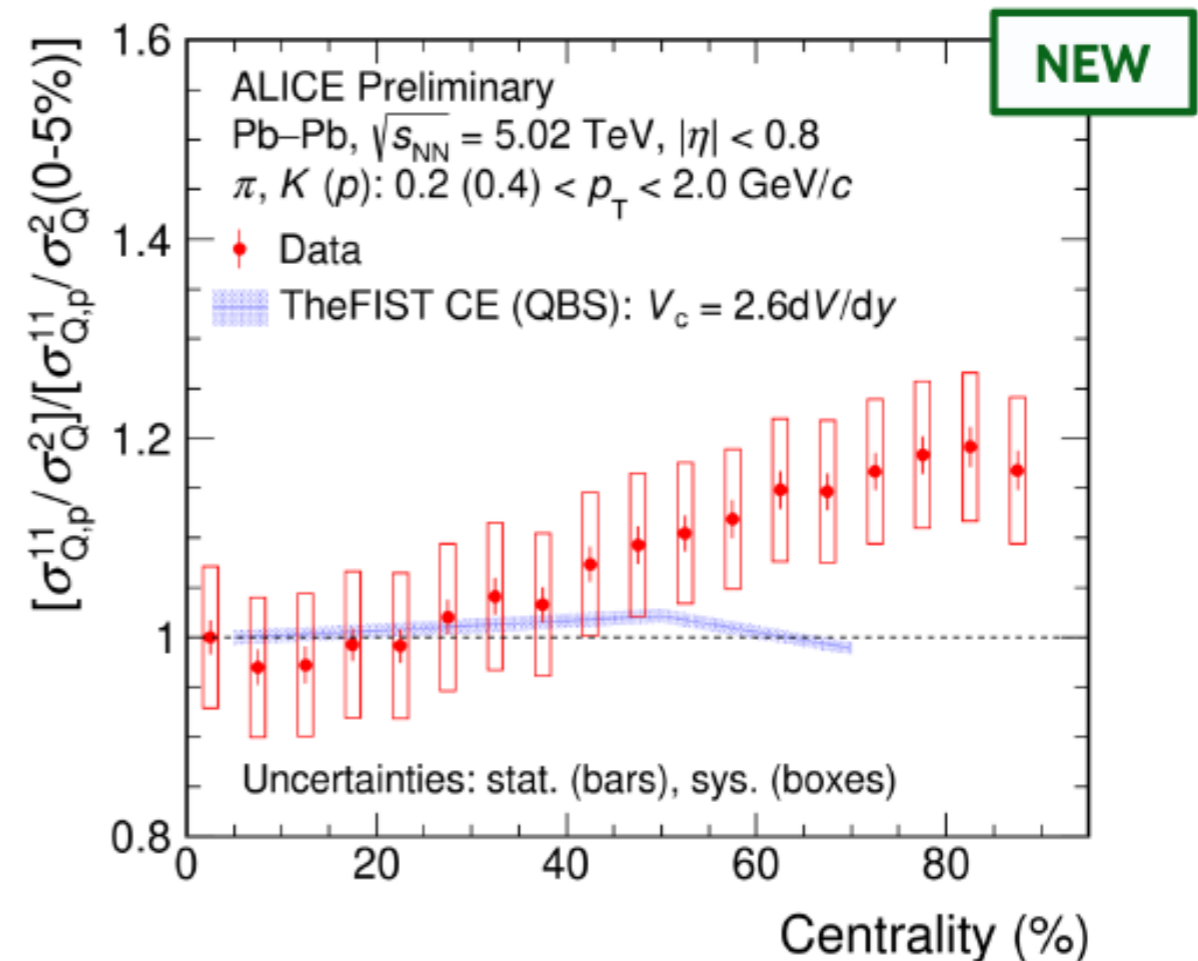
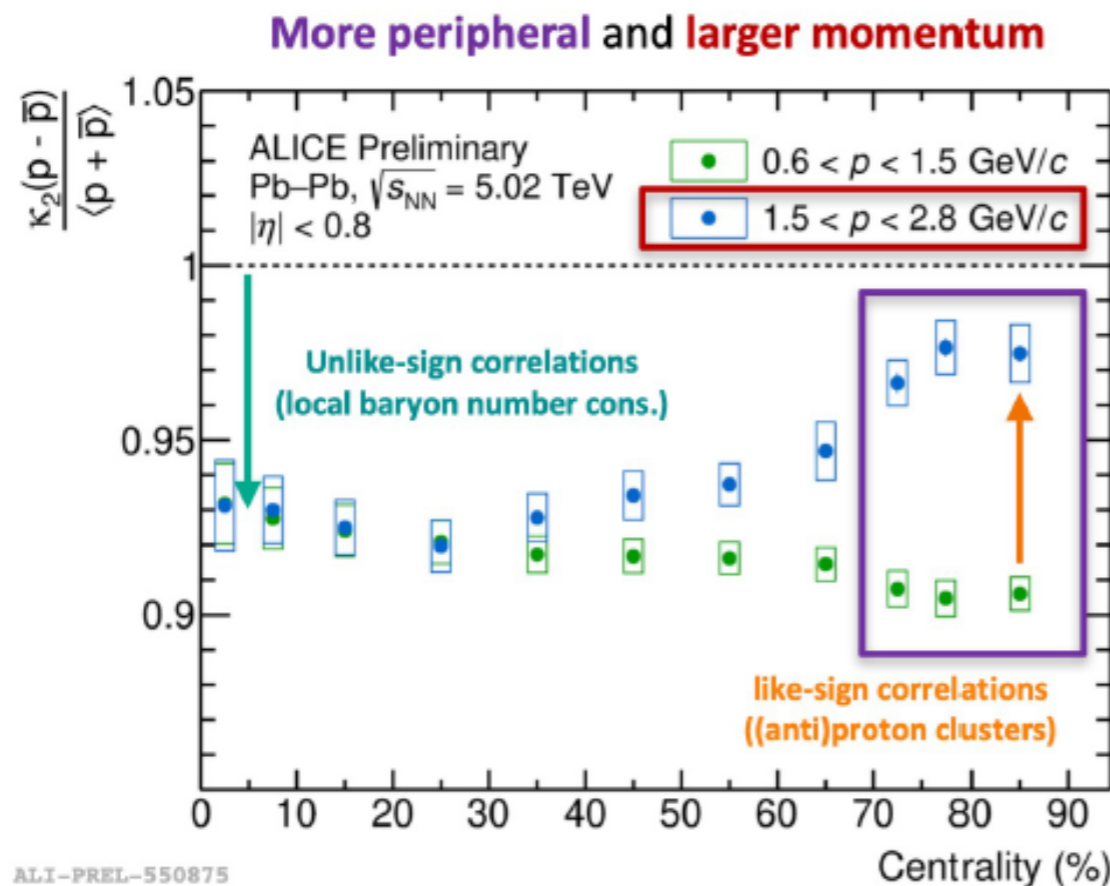
$$\chi_{B,S,Q}^{lmn} = \frac{1}{VT^3} \sigma_{B,S,Q}^{lmn}$$

**Proxies:**  
 Charge: K,  $\pi$ , p  
 Baryon: p  
 Strangeness: K



1) Increase towards peripheral collisions

# New preliminary ALICE data (SQM 2024)



1. A rise of  $\chi_2^B$ ,  $\chi_{11}^{BQ}$  in peripheral collisions due to magnetic field suggested by LQCD
2. **Hint of magnetic field in peripheral collisions in data?**

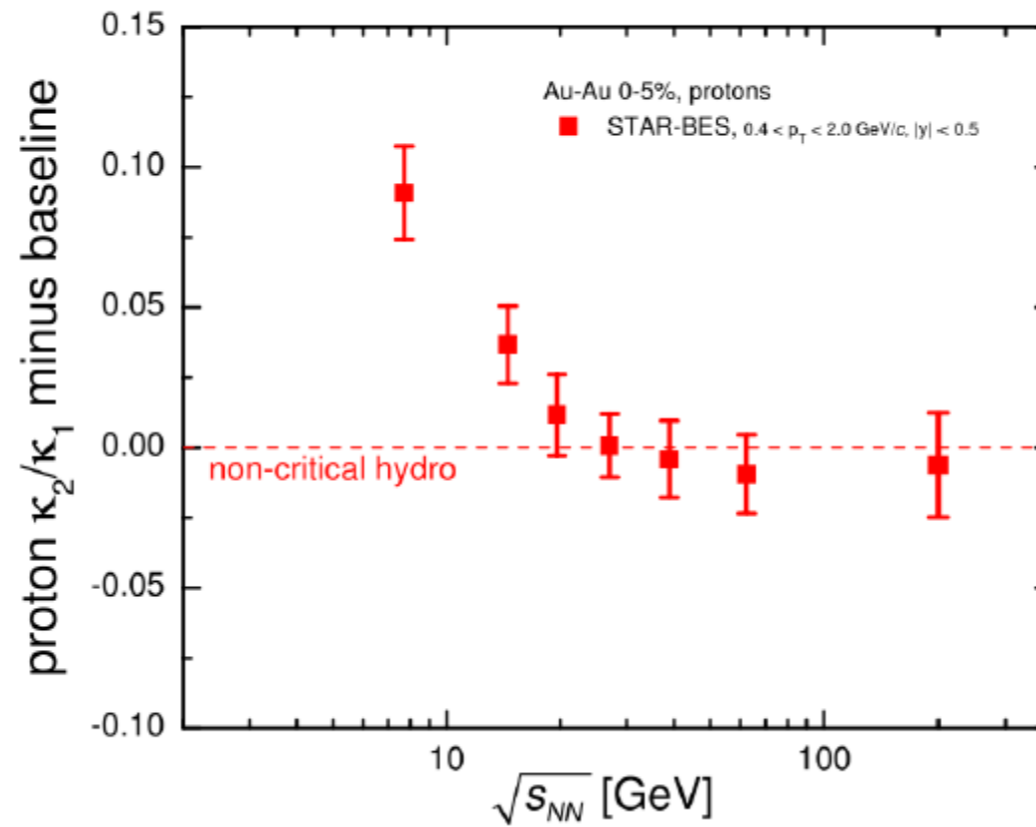
- 1) Increase towards peripheral collisions
- 2) **Similar behavior** is also observed in **net-p fluctuations** for the larger momentum range  
⇒ Magnetic field or ?

# Conclusions and Summary

- New expansion schemes for lattice QCD extend the region of validity to values of  $\mu_B/T = 3.5$
- Indication of critical point in Lee-Yang singularities.
- Location of point in agreement with DSE, FRG, and holography approaches that use lattice QCD as anchor.
- Location in the range of collision energies around 3-5 A GeV (potentially measurable at RHIC, FAIR, NICA, HPARC). Data still inconclusive, but more to come.
- Indication of chiral criticality at  $\mu_B=0$  and strong early magnetic fields from higher order fluctuations

# Backup: Factorial results

Proton  $\kappa_2/\kappa_1$  excess over baseline



- Intriguing hints from HADES@2.4 GeV and STAR-FXT@3GeV: huge excess of two-proton correlations!

[HADES Collaboration, Phys. Rev. C 102, 024914 (2020)]

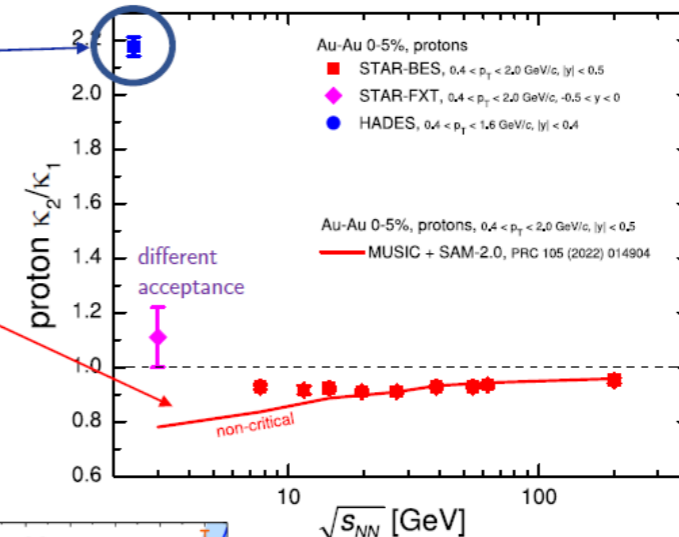
[STAR Collaboration, Phys. Rev. Lett. 128, 202303 (2022)]

- No change of trend in the non-critical reference
- Additional mechanisms:
  - Nuclear liquid-gas transition (the other QCD critical point)
  - Light nuclei formation/fragmentation
  - Stronger initial state, volume, and baryon stopping fluctuations

Talk by A. Bzdak, Wed 14:20; Poster by A. Rustamov

- Difference in acceptance ( $-0.5 < y < 0$  vs  $|y| < 0.5$ )

- Improved modeling of lower energies required



VV, Phys. Rev. C 106, 064906 (2022)

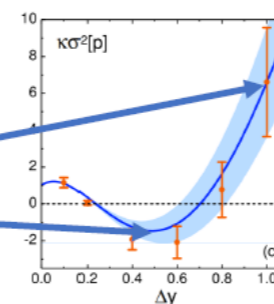


Figure from O. Savchuk et al., PLB 835, 137540 (2022)