Femtoscopy for exotic hadrons and nuclei





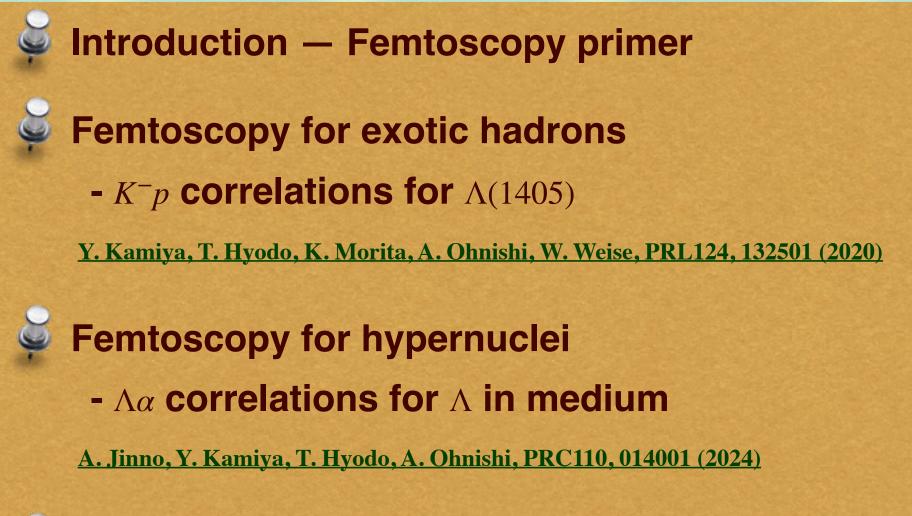
Tetsuo Hyodo

Tokyo Metropolitan Univ.





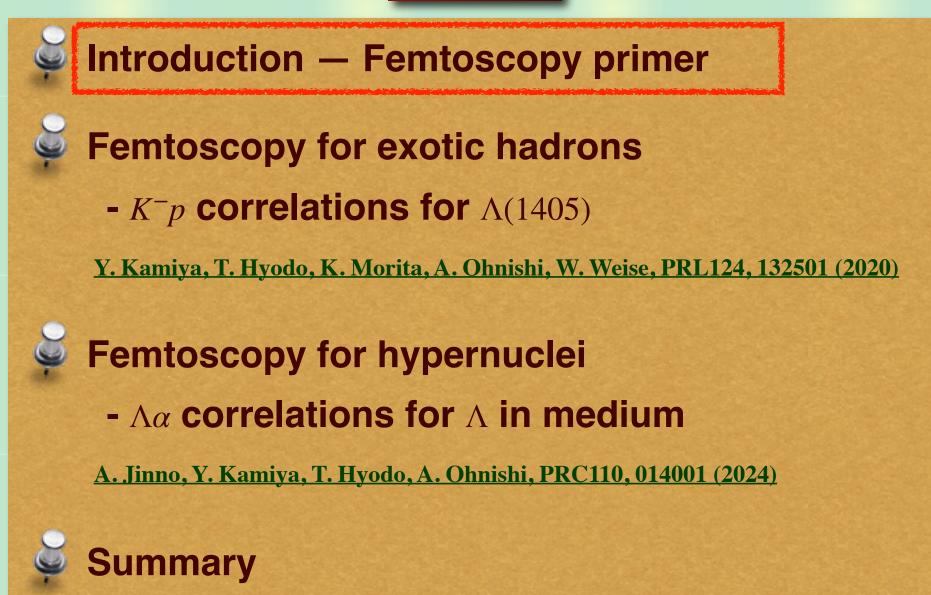
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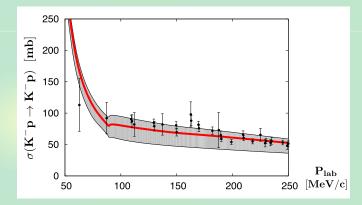


Scattering experiments and femtoscopy

Traditional methods: scattering experiments

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011)

- Limited channels: NN, YN, πN , KN, $\bar{K}N$, ...
- Heavy (c, b) hadrons: impossible
- Limited statistics (low-energy)



Scattering experiments and femtoscopy

Traditional methods: scattering experiments

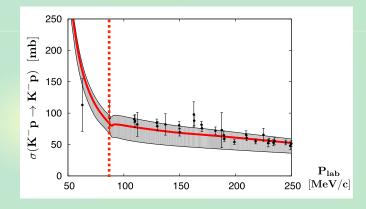
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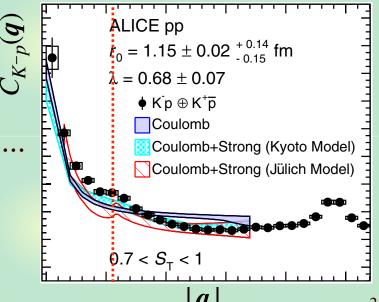
- Limited channels: NN, YN, πN , KN, $\bar{K}N$, ...
- Heavy (c, b) hadrons: impossible
- Limited statistics (low-energy)

Femtoscopy: correlation function

ALICE collaboration, PRL 124, 092301 (2020)

- Various systems: $\Lambda\Lambda$, $N\Omega$, ϕN , $\bar{K}\Lambda$, DN, ...
- Heavy hadrons: possible!
- Excellent precision (\bar{K}^0n cusp)



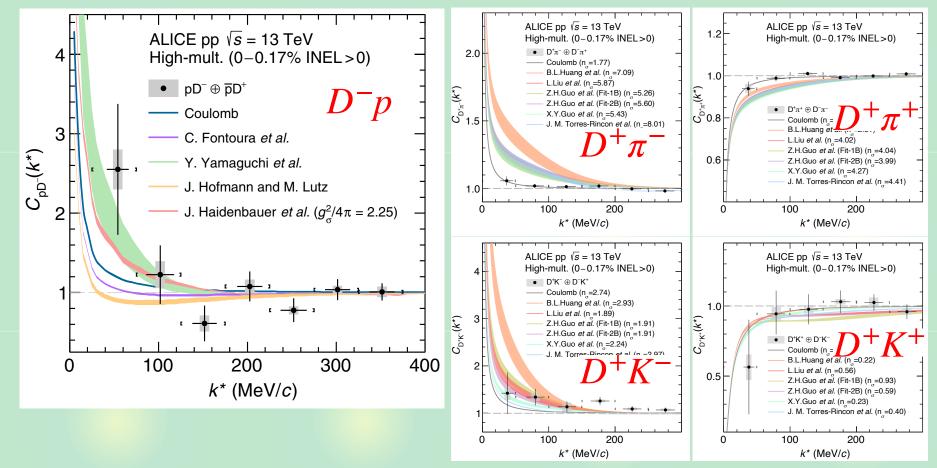


Experimental data in charm sector

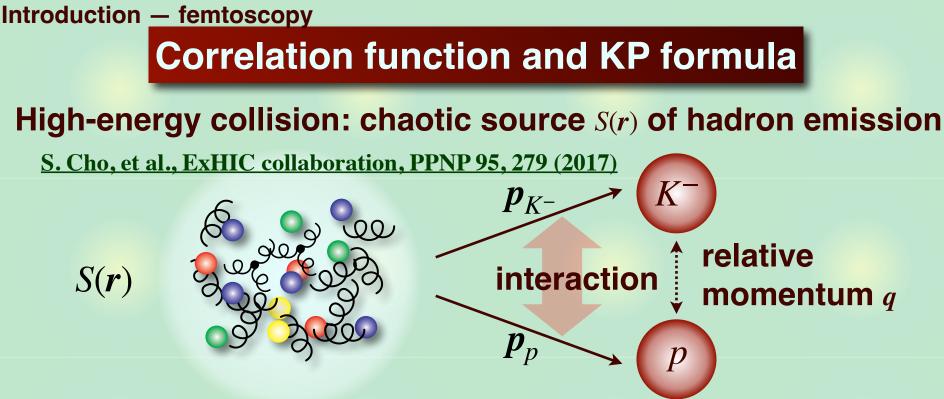
Observed correlation functions with charm: DN, $D\pi$, DK

ALICE collaboration, PRD 106, 052010 (2022);

ALICE collaboration, PRD 110, 032004 (2024)



Unique way to obtain data in charm sector (yet low statistics),



- Definition

$$C(\boldsymbol{q}) = \frac{N_{K^-p}(\boldsymbol{p}_{K^-}, \boldsymbol{p}_p)}{N_{K^-}(\boldsymbol{p}_{K^-})N_p(\boldsymbol{p}_p)} \quad \text{(= 1 in the absence of FSI/QS)}$$

Introduction – femtoscopy **Correlation function and KP formula High-energy collision: chaotic source** *S*(*r*) of hadron emission <u>S. Cho, et al., ExHIC collaboration, PPNP 95, 279 (2017)</u> *S*(*r*)

- Definition

$$C(\boldsymbol{q}) = \frac{N_{K^-p}(\boldsymbol{p}_{K^-}, \boldsymbol{p}_p)}{N_{K^-}(\boldsymbol{p}_{K^-})N_p(\boldsymbol{p}_p)} \quad \text{(= 1 in the absence of FSI/QS)}$$

- Theory (Koonin-Pratt formula) incoming + outgoing S.E. Koonin, PLB 70, 43 (1977); S. Pratt, PRD 33, 1314 (1986) $C(q) \simeq \int d^3r S(r) |\Psi_q^{(-)}(r)|^2, \quad \Psi_q^{(-)}(r) \propto S^{\dagger} e^{-iqr} - e^{+iqr} \quad (r \to \infty)$ Source function S(r) < -> wave function $\Psi_q^{(-)}(r)$ (interaction)

 \boldsymbol{p}_p

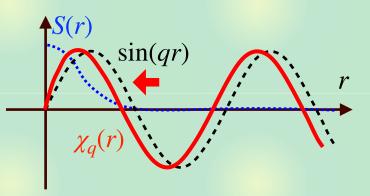
Introduction – <u>femtoscopy</u>

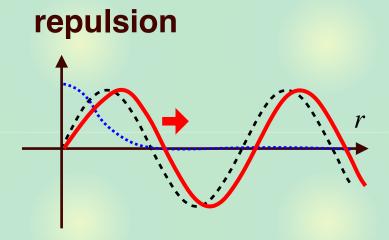
Wave functions and correlations

Spherical source with s-wave interaction dominance

$$C(q) \simeq 1 + \int_0^\infty dr \, S(r) \{ |\chi_q(r)|^2 - \sin^2(qr) \}$$

attraction

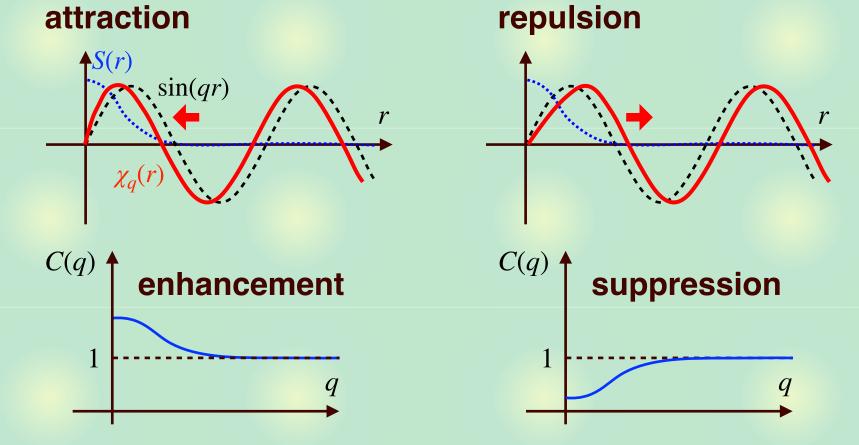




Wave functions and correlations

Spherical source with s-wave interaction dominance

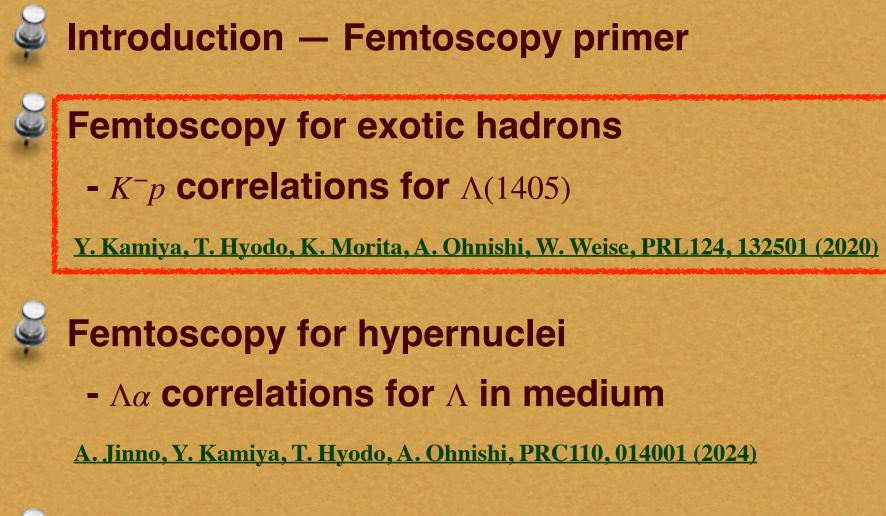
$$C(q) \simeq 1 + \int_0^\infty dr \, S(r) \{ |\chi_q(r)|^2 - \sin^2(qr) \}$$



Correlation function <-> nature of interaction

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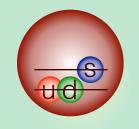


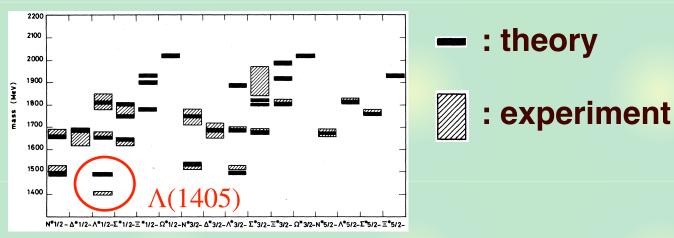


$\Lambda(1405)$ and $\bar{K}N$ scattering

$\Lambda(1405)$ does not fit in standard picture —> exotic candidate

N. Isgur and G. Karl, PRD18, 4187 (1978)

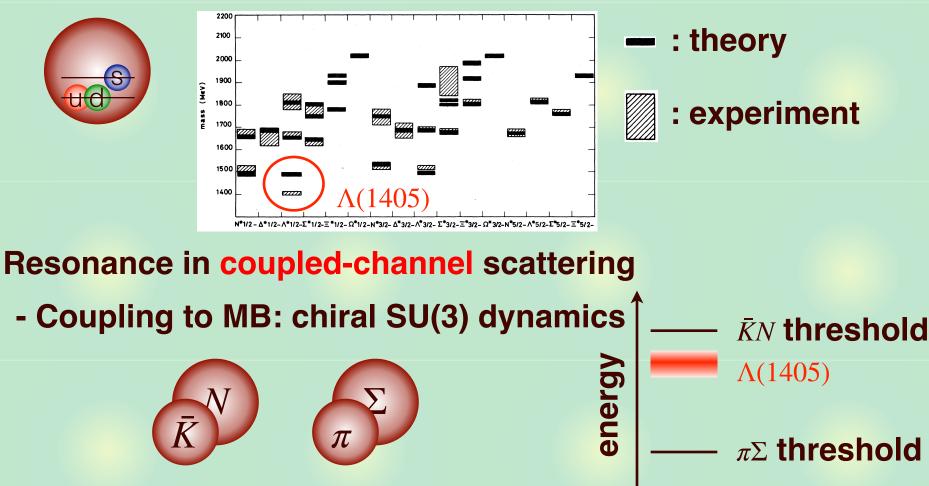




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T. Hyodo, W. Weise, arXiv:2202.06181 [nucl-th] (Handbook of Nuclear Physics)

Coupled-channel effects

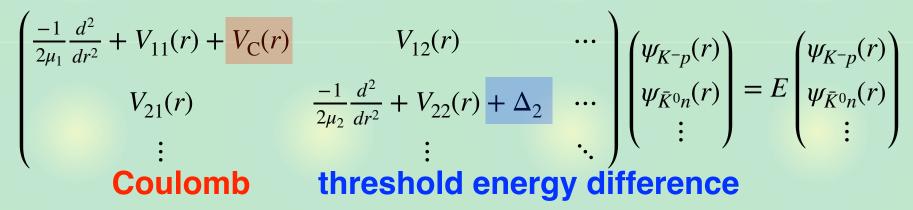
Schrödinger equation (s-wave)

$$\begin{pmatrix} \frac{-1}{2\mu_1} \frac{d^2}{dr^2} + V_{11}(r) + V_{\mathbf{C}}(r) & V_{12}(r) & \cdots \\ V_{21}(r) & \frac{-1}{2\mu_2} \frac{d^2}{dr^2} + V_{22}(r) + \Delta_2 & \cdots \\ \vdots & \ddots \end{pmatrix} \begin{pmatrix} \psi_{K^-p}(r) \\ \psi_{\bar{K}^0n}(r) \\ \vdots \end{pmatrix} = E \begin{pmatrix} \psi_{K^-p}(r) \\ \psi_{\bar{K}^0n}(r) \\ \vdots \end{pmatrix}$$

Coulomb threshold energy difference

Coupled-channel effects

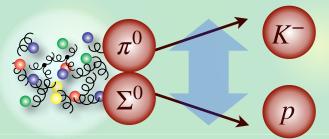
Schrödinger equation (s-wave)



Asymptotic ($r \rightarrow \infty$) wave function (incoming + outgoing)

$$\begin{pmatrix} \psi_{K^-p}(r) \\ \psi_{\bar{K}^0n}(r) \\ \vdots \end{pmatrix} \propto \begin{pmatrix} S_{11}^{\dagger}e^{-iqr} - e^{iqr} \\ S_{12}^{\dagger}e^{-iq_2r} - 0 \times e^{iq_2r} \\ \vdots \end{pmatrix} \quad (r \to \infty)$$

- Transition from $\bar{K}^0 n, \pi^+ \Sigma^-, \pi^0 \Sigma^0, \pi^- \Sigma^+, \pi^0 \Lambda$ is in $\psi_i(r)$ with $i \neq K^- p$



Coupled-channel correlation function

Coupled-channel Koonin-Pratt formula

R. Lednicky, V.V. Lyuboshitz, V.L. Lyuboshitz, Phys. Atom. Nucl. 61, 2950 (1998); J. Haidenbauer, NPA 981, 1 (2019);

Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020)

$$C_{K^{-p}}(\boldsymbol{q}) \simeq \int d^3 \boldsymbol{r} \, S_{K^{-p}}(\boldsymbol{r}) \left| \Psi_{K^{-p},\boldsymbol{q}}^{(-)}(\boldsymbol{r}) \right|^2$$

Coupled-channel correlation function

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$$C_{K^{-p}}(\boldsymbol{q}) \simeq \int d^3 \boldsymbol{r} \, S_{K^{-p}}(\boldsymbol{r}) \, |\Psi_{K^{-p},\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2 + \sum_{i \neq K^{-p}} \boldsymbol{\omega}_i \int d^3 \boldsymbol{r} \, S_i(\boldsymbol{r}) \, |\Psi_{i,\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2$$

- Transition from $\bar{K}^0 n, \pi^+ \Sigma^-, \pi^0 \Sigma^0, \pi^- \Sigma^+, \pi^0 \Lambda$
- ω_i : weight of channel *i* source relative to K^-p

Coupled-channel correlation function

Coupled-channel Koonin-Pratt formula

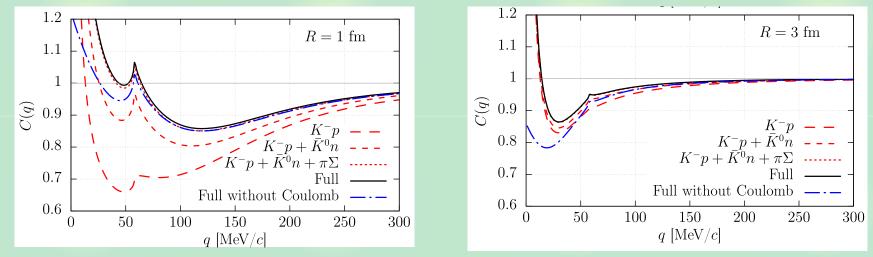
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Francition from $\bar{k}^0 r \, \sigma^+ \Sigma^- \sigma^0 \Sigma^0 \, \sigma^- \Sigma^+ \, \sigma^0 \Lambda$

- Transition from $\bar{K}^0 n, \pi^+ \Sigma^-, \pi^0 \Sigma^0, \pi^- \Sigma^+, \pi^0 \Lambda$

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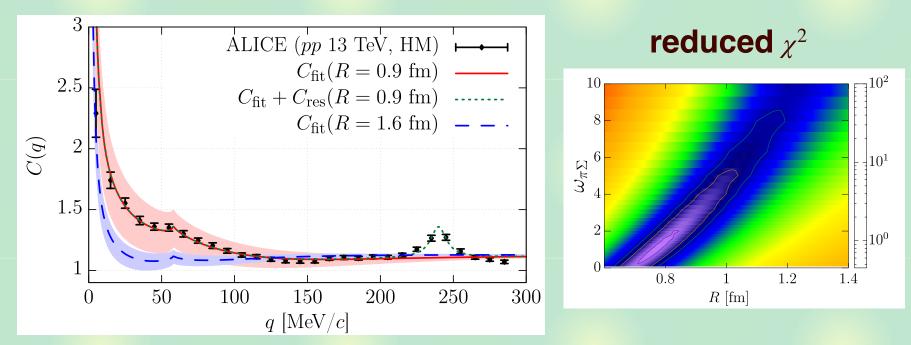
Coupled-channel effect is enhanced for small sources

Correlation from chiral SU(3) dynamics

Wave function $\Psi_{i,g}^{(-)}(r)$: Kyoto $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$ potential

K. Miyahara, T. Hyodo, W. Weise, PRC98, 025201 (2018)

- Source function S(r): gaussian, $R \sim 1$ fm from K^+p data
- Source weight $\omega_{\pi\Sigma} \sim 2$ by simple statistical model estimate



Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020)

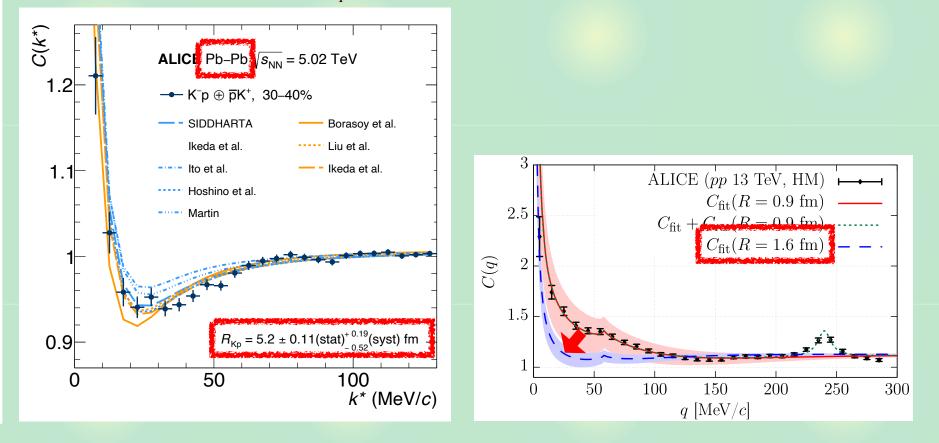
Correlation is well reproduced by chiral SU(3) potential

Large source case

New data with Pb-Pb collisions at 5.02 TeV

ALICE collaboration, PLB 822, 136708 (2021)

- Scattering length $a_{K^-p} = -0.91 + 0.92i$ fm



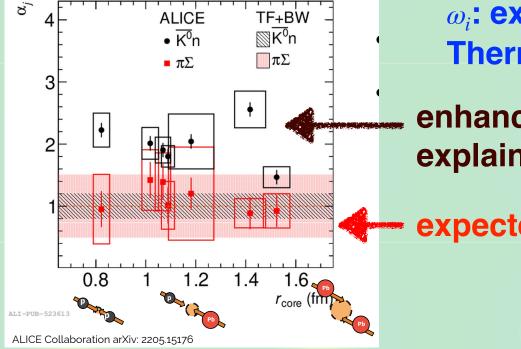
Correlation is suppressed at larger *R***, as predicted**

Systematic study of source size dependence

Correlations in *pp*, *p*-Pb, Pb-Pb **by Kyoto** $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$ **potential**

ALICE collaboration, EPJC 83, 340 (2023)

$$C_{K^{-}p}(\boldsymbol{q}) \simeq \int d^3 \boldsymbol{r} \, S_{K^{-}p}(\boldsymbol{r}) \, |\Psi_{K^{-}p,\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2 + \sum_{i \neq K^{-}p} \alpha_i \, \omega_i \int d^3 \boldsymbol{r} \, S_i(\boldsymbol{r}) \, |\Psi_{i,\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2$$



ω_i: expected weight by Thermal Fist + Blast Wave

enhancement needed to explain data

expected weight is OK

More strength is needed in the $\bar{K}^0 n$ channel



Contents





- K^-p correlations for $\Lambda(1405)$

Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020)

Femtoscopy for hypernuclei - $\Lambda \alpha$ correlations for Λ in medium

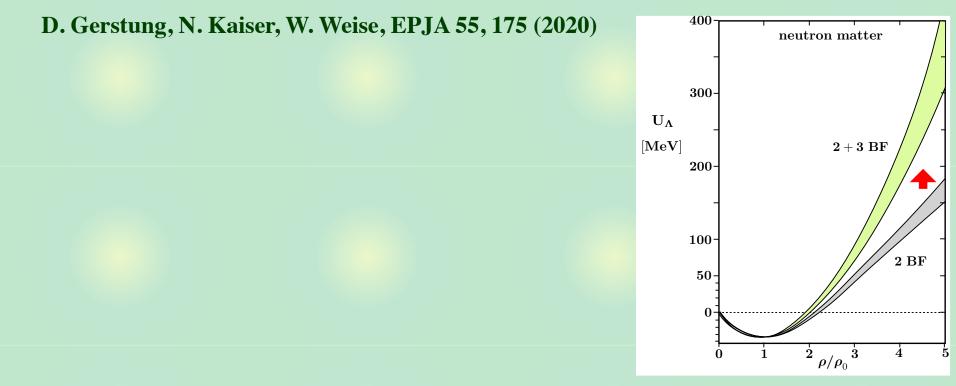
A. Jinno, Y. Kamiya, T. Hyodo, A. Ohnishi, PRC110, 014001 (2024)



Motivation

A solution to hyperon puzzle in neutron stars

- ANN three-body force for repulsion at high density



Motivation

A solution to hyperon puzzle in neutron stars

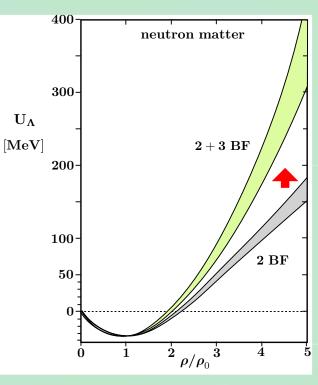
- ANN three-body force for repulsion at high density

D. Gerstung, N. Kaiser, W. Weise, EPJA 55, 175 (2020)

How to verify this in experiments?

- Λ directed flow in heavy ion collisions

Y. Nara, A. Jinno, K. Murase, A. Ohnishi, PRC 106, 044902 (2022)



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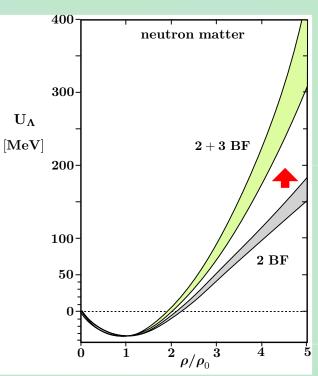
- Λ directed flow in heavy ion collisions

Y. Nara, A. Jinno, K. Murase, A. Ohnishi, PRC 106, 044902 (2022)

A-nucleus correlation function?

- Heavy nuclei are difficult to produce
- Strong binding of *α*: two-body treatment justified

$\Lambda \alpha$ correlation function —> nature of $\Lambda \alpha$ potential?

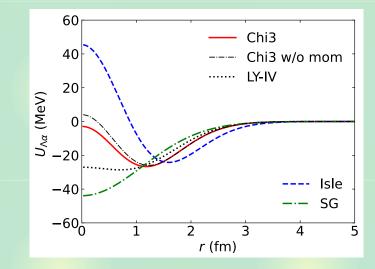


$\Lambda \alpha$ potentials

Phenomenological $\Lambda \alpha$ **potentials** (⁵ He binding energy)

I. Kumagai-Fuse, S. Okabe, Y. Akaishi, PLB 345, 386 (1997)

- SG: single gaussian
- Isle: two gaussians (with core)



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- Isle: two gaussians (with core)
- **Skyrme-Hartree Fock methods**
 - LY4: phenomenorogical

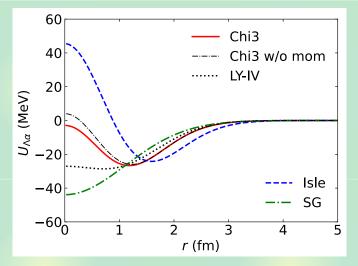
D.E. Lanskoy, Y. Yamamoto, PRC 55, 2330 (1997)

- Chi3: based on chiral EFT with ANN force

A. Jinno, K. Murase, Y. Nara, A. Ohnishi, PRC 108, 065803 (2023)

- Both potentials reproduce hypernuclear data from C to Pb
- α density distribution —> $\Lambda \alpha$ potentials

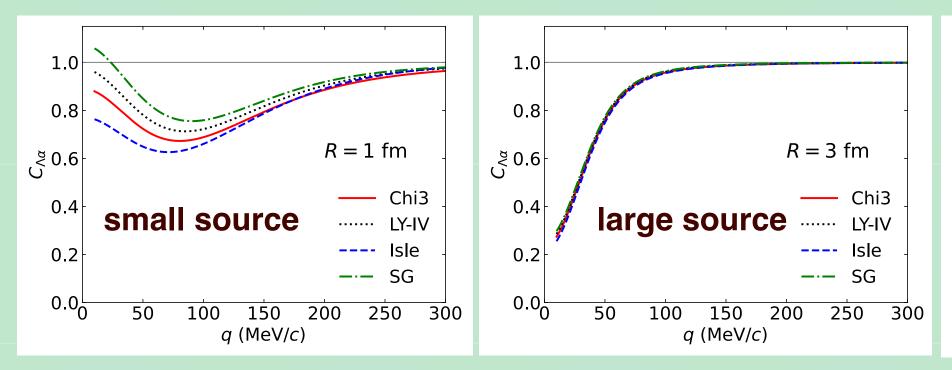
Effect of repulsive core —> correlation function?



$\Lambda \alpha$ correlation functions: source size dependence

Correlation functions from small and large sources

A. Jinno, Y. Kamiya, T. Hyodo, A. Ohnishi, PRC110, 014001 (2024)

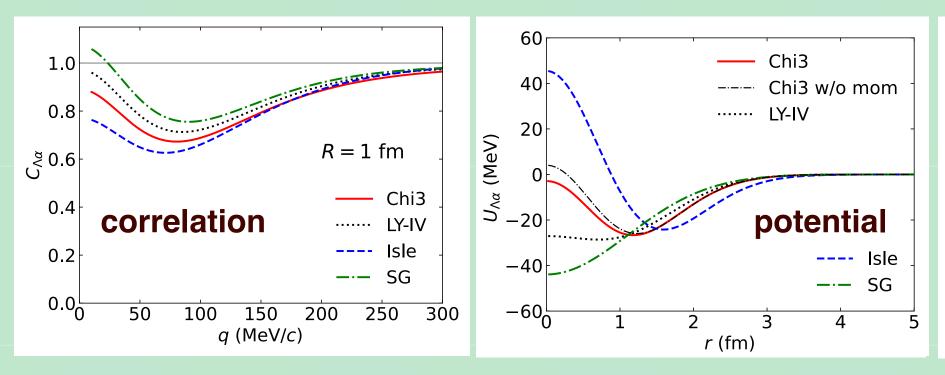


- Bound state signature (dip at low q in small source)
- No difference in large source ($R \sim 3 \text{ fm}$)
- Potential dependence in small source ($R \sim 1 \text{ fm}$)

$\Lambda \alpha$ correlation functions: potential dependence

Correlation functions and $\Lambda \alpha$ **potentials**

A. Jinno, Y. Kamiya, T. Hyodo, A. Ohnishi, PRC110, 014001 (2024)

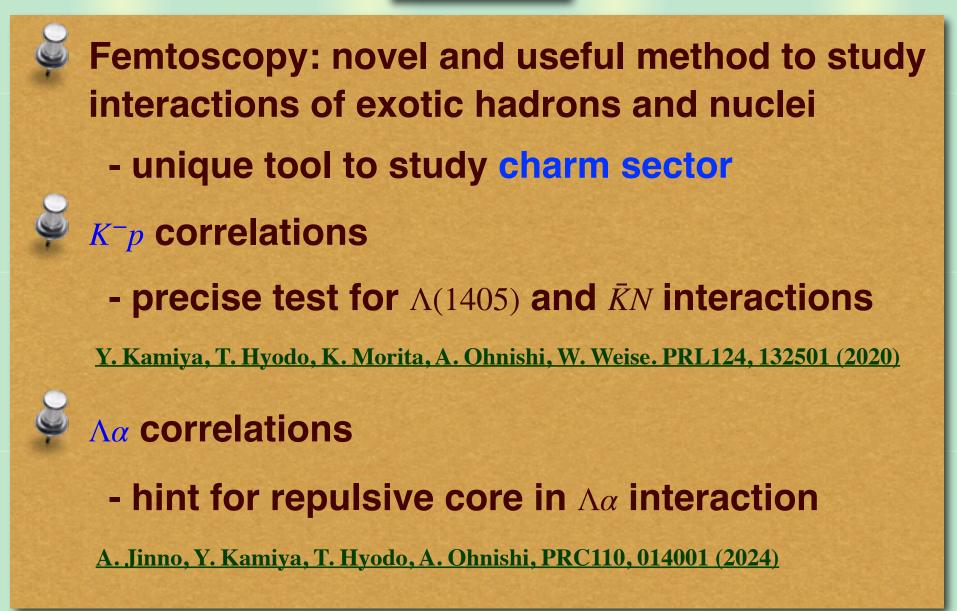


- $U_{\Lambda\alpha}(r=0)$: Isle > LY-IV > Chi3 > SG
- $C_{\Lambda\alpha}(q=0)$: Isle < LY-IV < Chi3 < SG
- Central repulsion suppresses correlation at low q

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Summary and future prospects

Summary



Hadron2025

Hadron2025

- Hadron 2025 conference will be held in Osaka
 - March 27-31, 2025
 - Registration will be open soon

https://hadron2025.rcnp.osaka-u.ac.jp/





LL formula

Correlation function <-> **observables** ($a_0, r_e, f(q)$)

R. Lednicky, V.L. Lyuboshits, Yad. Fiz. 35, 1316 (1981)

- Gaussian (relative) source $S(r) = \exp(-r^2/4R^2)/(4\pi R^2)^{3/2}$
- R : source size (gaussian width is $\sqrt{2R}$)

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- Gaussian (relative) source $S(r) = \exp(-r^2/4R^2)/(4\pi R^2)^{3/2}$
- R : source size (gaussian width is $\sqrt{2R}$)
- s-wave interaction only
- zero-range interaction : $R \gg R_{int}$ (use asymptotic w.f.)

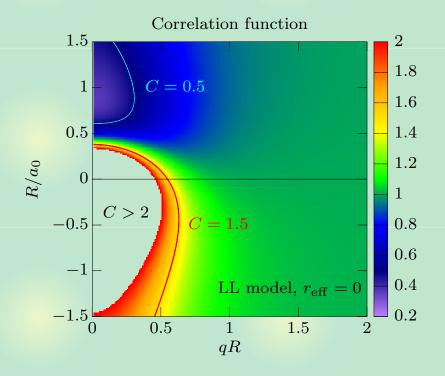
$$C(q) = 1 + \frac{|f(q)|^2}{2R^2} F_3(r_e/R) + \frac{2\text{Re } f(q)}{\sqrt{\pi R}} F_1(2qR) - \frac{\text{Im } f(q)}{R} F_2(2qR)$$

F_i(x) : known functions, *f(q)* : s-wave scattering amplitude
 <u>S. Cho, et al., ExHIC collaboration, PPNP 95, 279 (2017)</u>

$$f(q) = \frac{1}{q \cot \delta - iq} \simeq \frac{1}{-\frac{1}{a_0} + \frac{r_e}{2}q^2 - iq}$$

LL formula and correlations

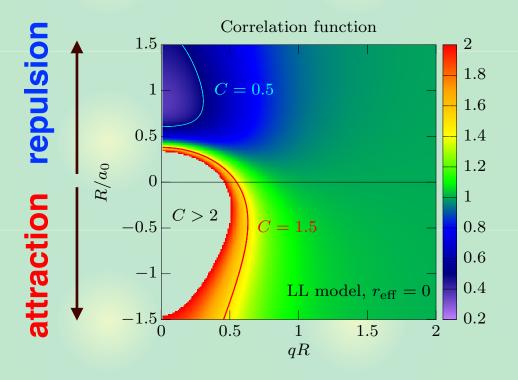
LL formula with $r_e = 0$



LL formula and correlations

LL formula with $r_e = 0$

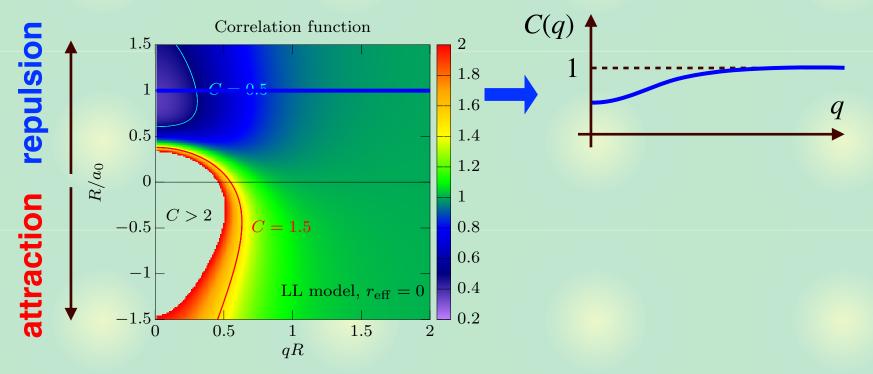
- fixed R > 0
- repulsion: $R/a_0 > 0$, attraction: $R/a_0 < 0$



LL formula and correlations

LL formula with $r_e = 0$

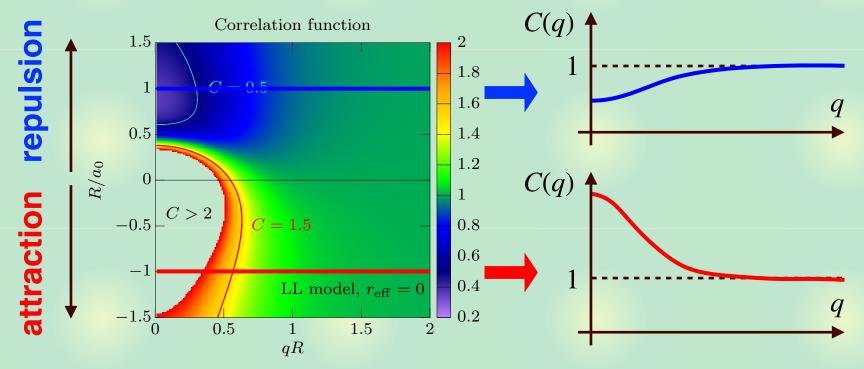
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LL formula and correlations

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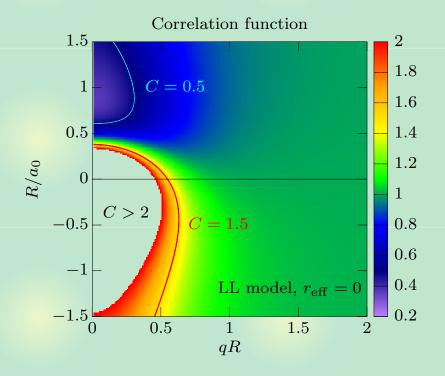
- fixed R > 0
- repulsion: $R/a_0 > 0$, attraction: $R/a_0 < 0$



Consistent with KP formula

Shallow bound state case

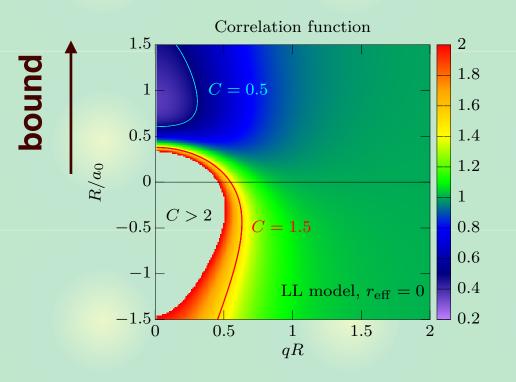
LL formula with $r_e = 0$



Shallow bound state case

LL formula with $r_e = 0$

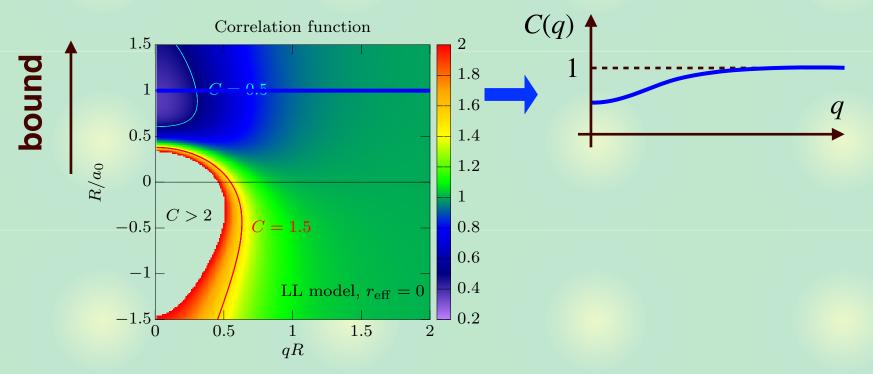
- shallow bound state: fixed $a_0 > 0$, $|a_0| \gg R_{int}$
- large source: $R/a_0 \sim 1$, small source: $R/a_0 \ll 1$



Shallow bound state case

LL formula with $r_e = 0$

- shallow bound state: fixed $a_0 > 0$, $|a_0| \gg R_{int}$
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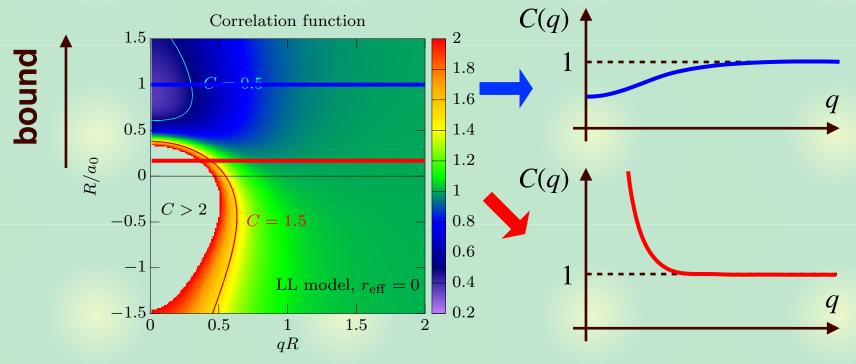


Shallow bound state case

LL formula with $r_e = 0$

<u>Y. Kamiya, K. Sasaki, T. Fukui, T. Hyodo, K. Morita, K. Ogata, A. Ohnishi, T. Hatsuda,</u> <u>PRC 105, 014915 (2022)</u>

- shallow bound state: fixed $a_0 > 0$, $|a_0| \gg R_{int}$
- large source: $R/a_0 \sim 1$, small source: $R/a_0 \ll 1$



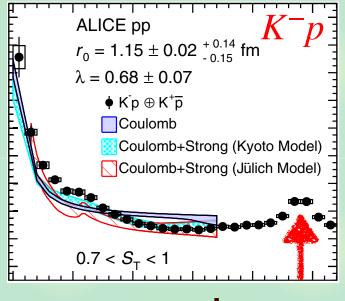
-> qualitative difference in large/small source

Resonance contributions

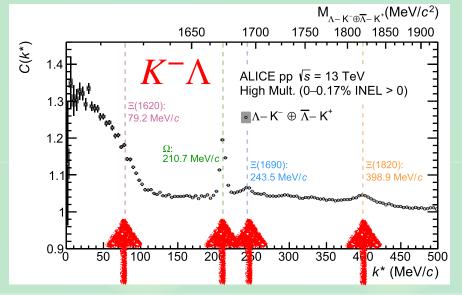
Higher partial waves and resonance contributions

Resonances in $\ell = 0$ and in $\ell \neq 0$ are seen

- Simple Breit-Wigner function has been used



Λ(1520) : **d-wave**



E(1620), E(1690): s-wave Ω : p-wave (weak decay) E(1820): d-wave

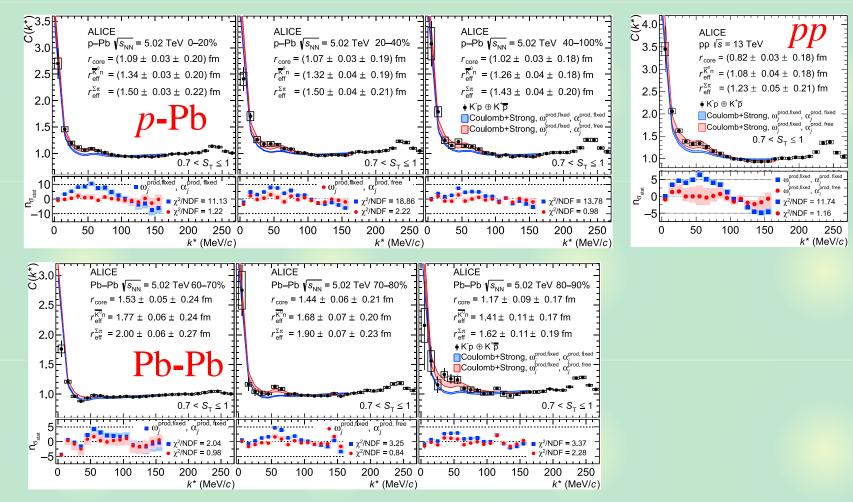
Questions

- Contribution from higher partial waves?
- Is Breit-Wigner function fine for resonance?

Systematic study of source size dependence

Correlations in *pp*, *p*-Pb, Pb-Pb **by Kyoto** $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$ **potential**

ALICE collaboration, EPJC 83, 340 (2023)



Correlations with different source size