

QCD & New Physics in Extreme Astrophysical Environments: Neutron Stars & their Mergers

Chair: Nicole Bell

Panel members: Susan Gardner
Anthony Thomas
Steven Harris

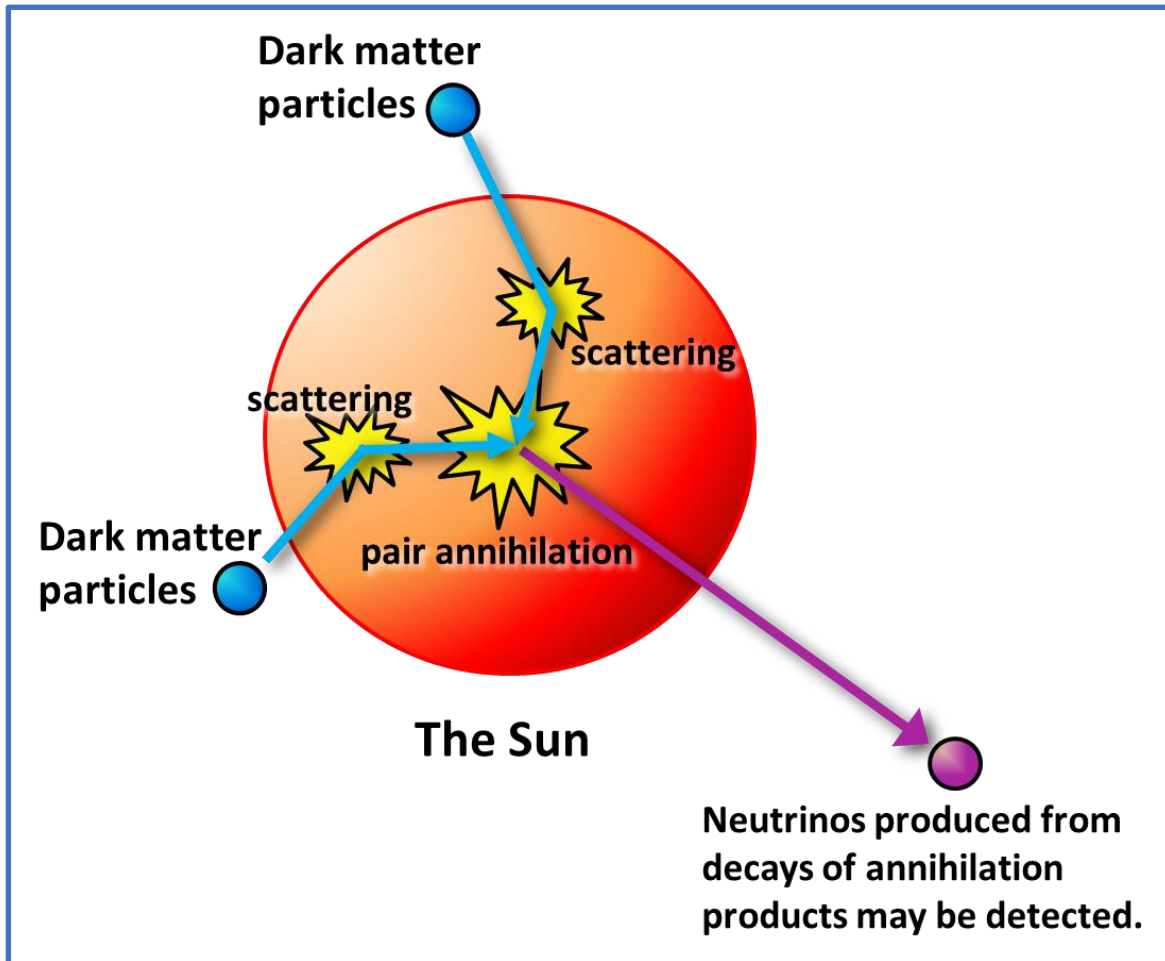
QCD & New Physics in Extreme Astrophysical Environments: Neutron Stars & their Mergers

- What lurks in the heart of a neutron star? Can they contain dark matter?
- What effects would “dark” neutron decays have? What can we learn about such DM ?
- How can neutron star observations constrain QCD-like hidden sectors?
- What new physics can we learn from gravitational wave detections (in NS mergers)?
- What new constraints emerge on axions from the study of neutron star mergers?

- What is the maximum NS mass?
- To what extent do NS possess strangeness? Can we generate heavy NS with hyperons?
- Is there a transition to deconfined quark matter?
- Can we observe evidence of QCD phase transition, under what conditions?
- How do GW observations constrain the EoS?
- Can we observe prompt black hole formation?

Dark Matter Capture in Stars

→ an alternative approach to Dark Matter Direct Detection experiments



Due to their extreme density, **neutron stars** capture dark matter **very** efficiently.

Capture probability saturates at order unity when the cross section satisfies the **geometric limit**

$$\sigma_{th} \sim \pi R^2 \frac{m_n}{M_*} \sim 10^{-45} \text{cm}^2$$

Neutron Stars → Black holes?

Kouvaris; Kouvaris & Tinyakov; McDermott, Yu & Zurek; Bramante, Fukushima & Kumar; NFB, Petraki & Melatos; Bertone, Nelson & Reddy; and others.

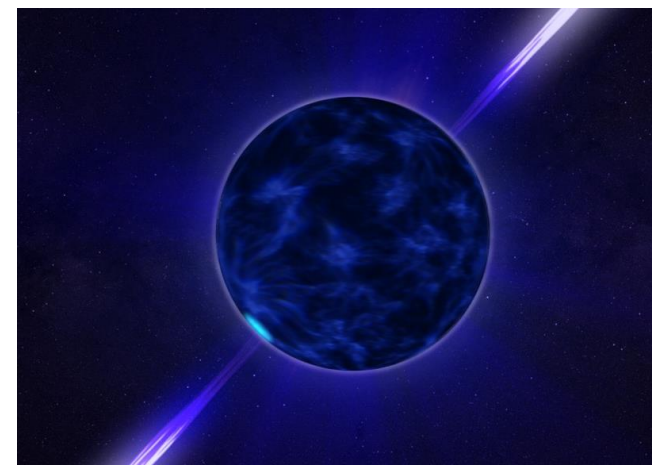
Can neutron stars accumulate so much dark matter they would collapse to black holes?

Yes, provided that:

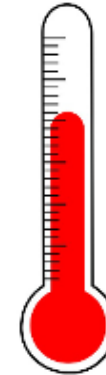
- No annihilation (e.g. asymmetric dark matter, with DM-antiDM asymmetry))
- DM is bosonic (and condenses to a small self gravitating BEC), or
- DM is fermionic with attractive self-interactions, and
- No repulsive-self interactions that prevent collapse (even very very tiny self-interaction is enough) [NFB, Petraki & Melatos, PRD 2013](#)

→ Black hole formation is possible for non-annihilating dark matter, while unlikely for typical WIMP-like dark matter

Neutron star heating



- **Capture of dark matter** (plus subsequent energy loss)
 - DM *kinetic energy* heats neutron star ~ **1700K** (Baryakhtar et al)
- **Annihilation** of thermalized dark matter
 - DM *rest mass energy* heats neutron star ~ **additional 700K**

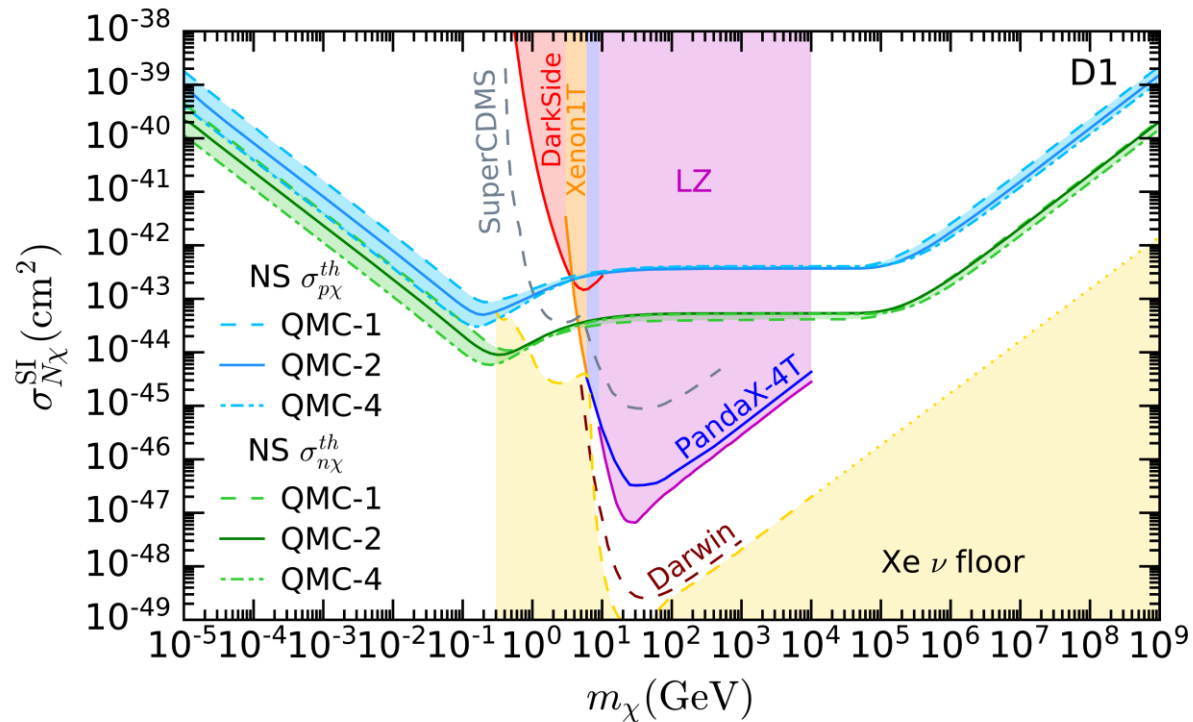


Coollest known neutron star (PSR J2144-3933) has a temperature of $\sim 4.2 \times 10^4$ K.

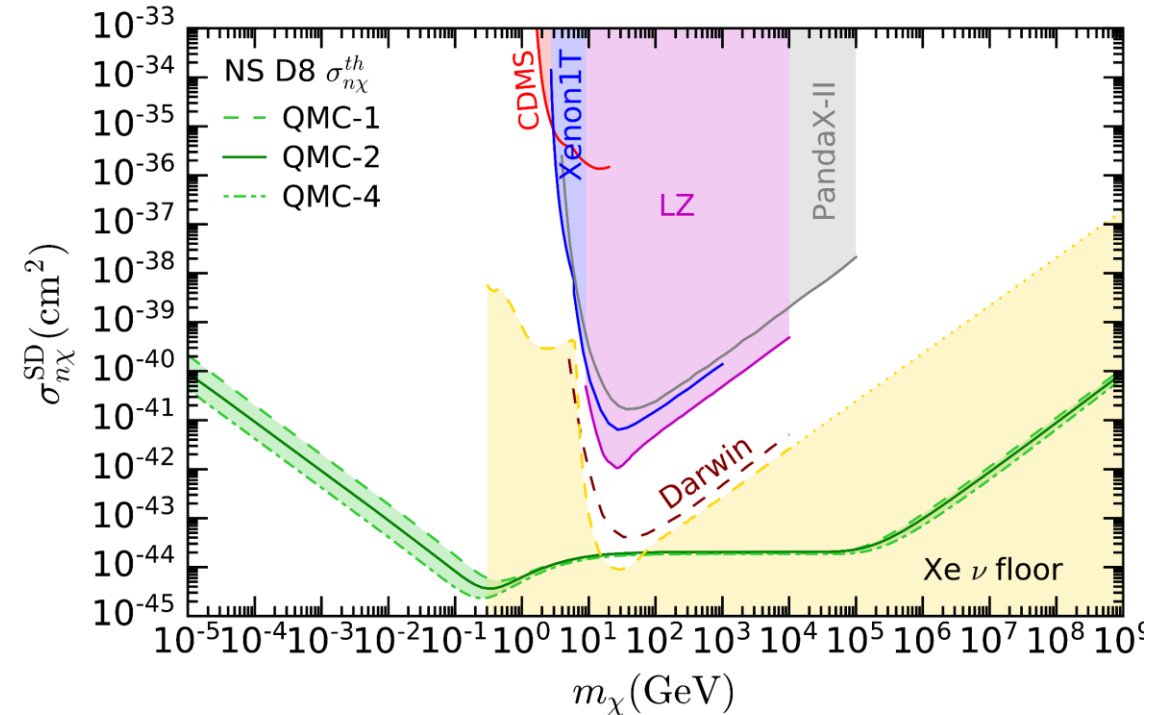
Old isolated neutron stars should cool to: 1000 K after ~ 10 Myr
100 K after ~ 1 Gyr

NS Heating Sensitivity:

Spin-independent DM-nucleon scattering



Spin-dependent DM-nucleon scattering

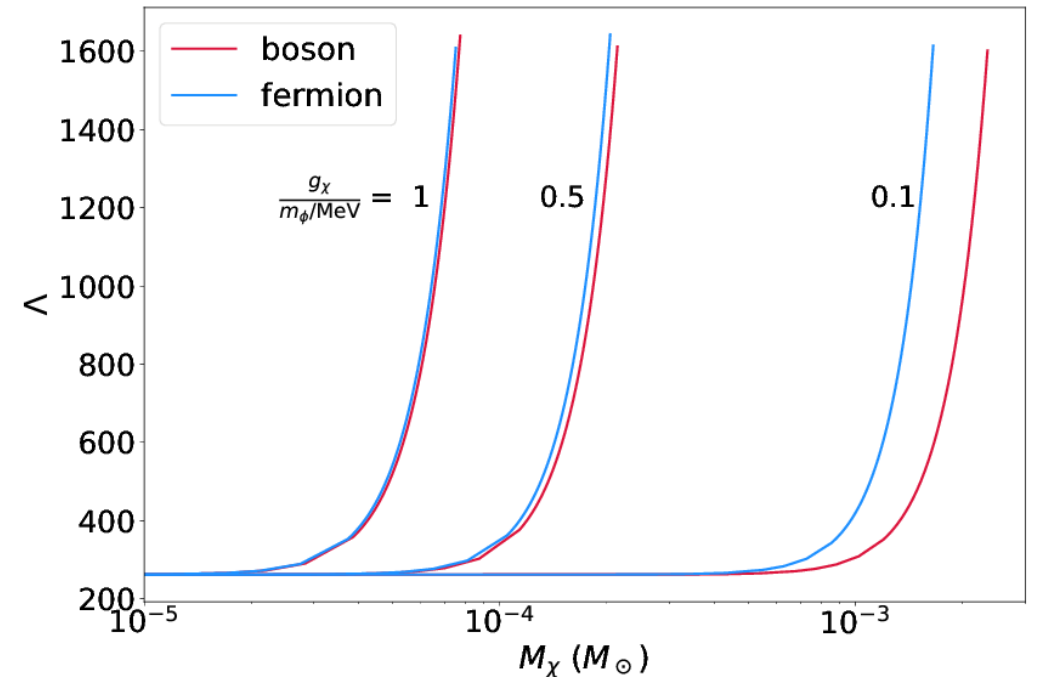
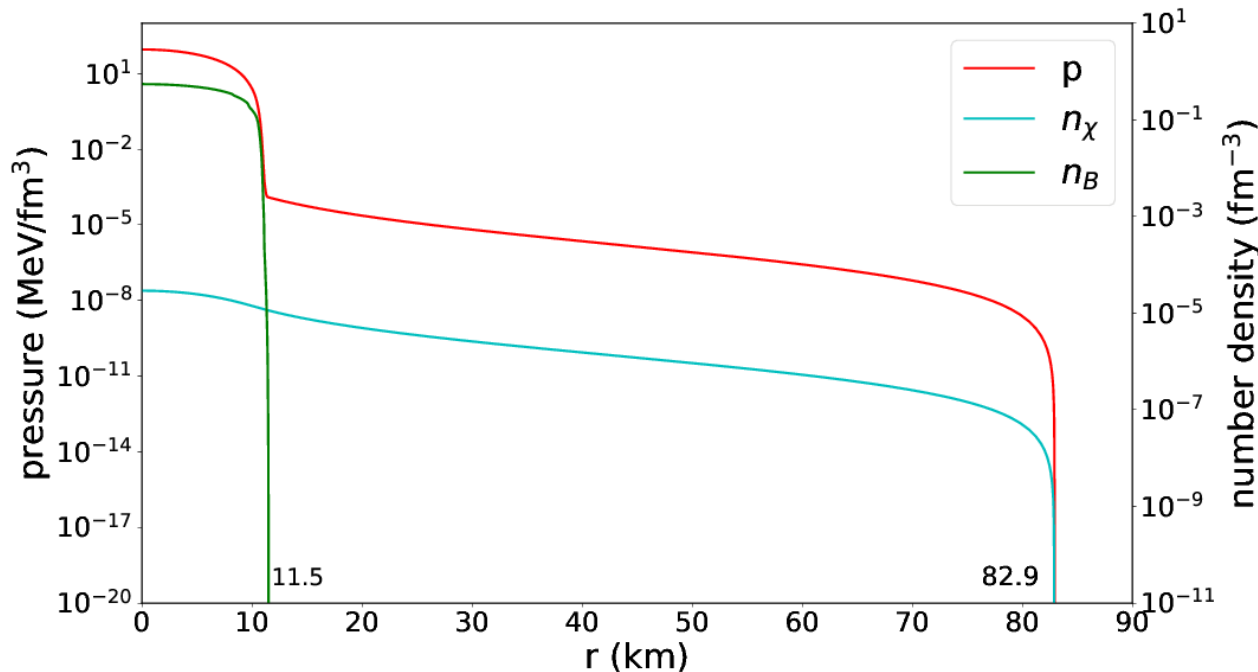


Anzuni, NFB, Busoni, Motta, Robles, Thomas and Virgato, JCAP 11, 056 (2021)

Neutron star mergers → gravitational waves

Nelson, Reddy & Zhou,
JCAP 07, 012 (2019)

- Light DM + light mediators (MeV scale)
- Either fermionic DM (degeneracy pressure) or bosonic with repulsive self-interactions
 - **DM component extends to large radii → NS dark matter halo**
- Increases the NS tidal deformability, Λ .
 - LIGO observation of NS-NS merger, GW170817, constrains $\Lambda < 800$
 - strong bounds, even for small DM component $\sim 10^{-4} M_\odot$



Light Dark Matter

Recently there was a very interesting proposal from Fornal and Grinstein (1801.01124).

Originated in long-standing puzzle concerning free neutron lifetime:

- Measurement for trapped n's: 879.6 ± 0.6 sec
- Measurement in beam decay : 888.0 ± 2.0 sec

This 3.5σ discrepancy solved by existence of new decay mode, which would not be seen in the beam decay experiment

$$n \rightarrow \text{Dark Matter } (\chi) + \text{something}$$

“Something” not a photon : Tang *et al.*, Los Alamos 1802.01595

Light Dark Matter (cont.)

- There are very strict limits on the mass of the new DM particle. It should be within an MeV or so of the neutron mass. For the case

$$n \rightarrow \chi + \varphi$$

with χ carrying baryon number and φ also dark,

$$937.9 \text{ MeV} < m_\chi + m_\varphi < 939.565 \text{ MeV}$$

- Serebrov *et al.*, 1802.06277 also claim this particle would resolve a reactor anti-neutrino anomaly
- Also nice discussion of tension with best value of neutron axial charge by [Czarnecki et al.](#), 1802.01804 (constrains but does not rule out the explanation)

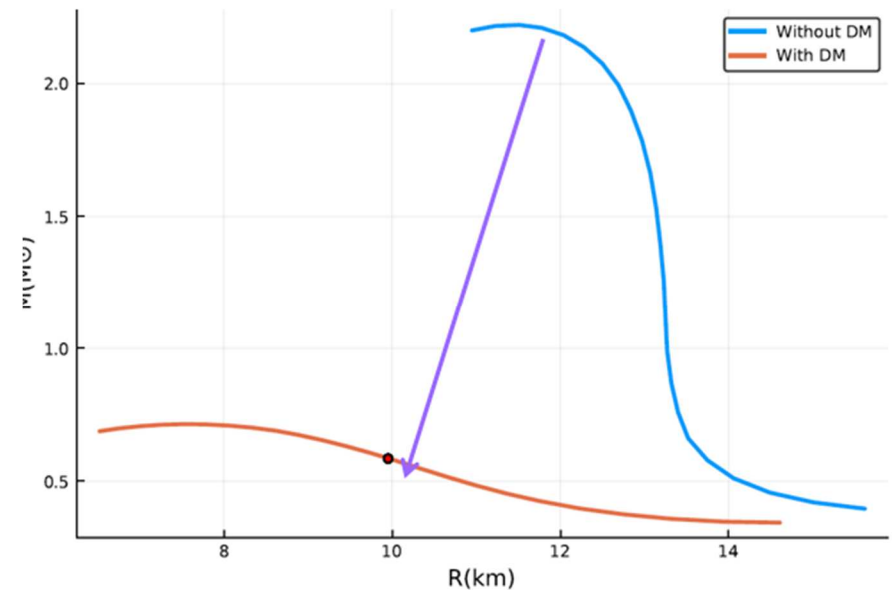
Compatibility of Fornal-Grinstein Hypothesis with Neutron Star Properties?

- A rush of papers
 - McKeen *et al.*, 1802.08244
 - Motta *et al.*, 1802.08427
 - Baym *et al.*, 1802.08282
- All reach a similar conclusion
- I follow the work of Motta, Guichon and Thomas (1802.08427)
- If such a dark matter particle exists neutrons high in the Fermi sea of a neutron star will be unstable
- This will replace high energy/pressure neutrons with lower energy/pressure dark matter particles: Consequences?

Dark matter particles must have strong repulsive interaction: like omega in NN force

Solve Tolman-Oppenheimer-Volkoff Equations

- Maximum allowed mass for stable neutron star drops from $2.21 M_{\odot}$ to $0.7 M_{\odot}$
- But cannot even get that as maximum stable star goes to just $0.58 M_{\odot}$

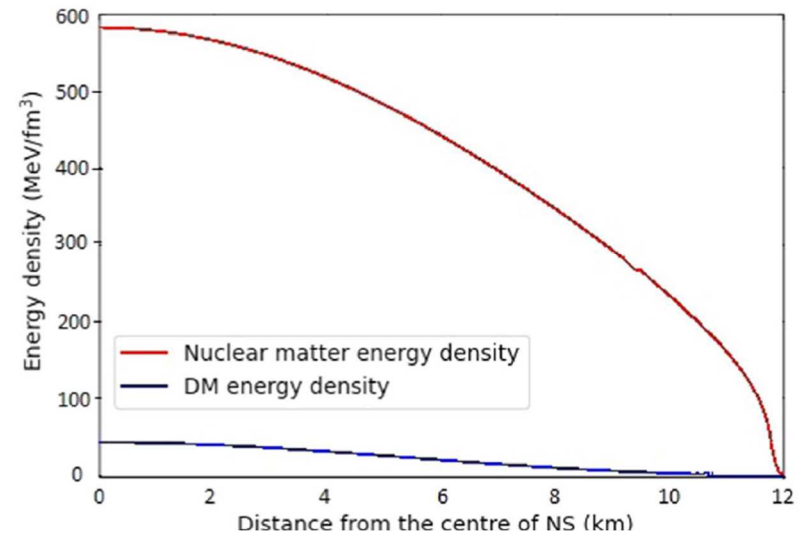
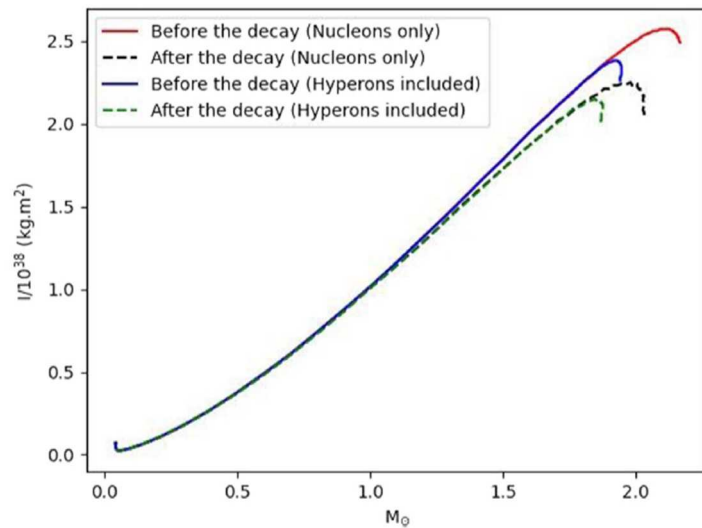


Motta *et al.*, 1802.08427

Alternative due to Strumia (JHEP 2022)

$n \rightarrow \chi\chi\chi$ with m_χ very close to $m_n/3$
No repulsive interaction needed.

**Moment of inertia decreases after decay:
hence star spins up and gets hotter!**



Dark Decay Models [Fornal & Grinstein, 2018]

Minimal ingredients, considered broadly

At lower energies...

$$\mathcal{O}_{abc} = u_a d_b d_c \chi$$

[Alonso-Alvarez et al., 2022]

— to induce visible-dark baryon mixing

Dark Decays of Hadrons

Neutron decay anomaly

$$\mathcal{O} = u d d \chi \quad m_{\text{DS}} \lesssim m_n$$

Hyperon dark decays (this work)

$$\mathcal{O} = u d s \chi \quad m_{\text{DS}} \lesssim m_\Lambda$$

B-Mesogenesis

$$\mathcal{O} = u d b \chi \quad m_{\text{DS}} \lesssim m_B$$

**CLAS, BESIII,
SN1987A**

[Elor, Escudero, Nelson, 2019;
Elor & McGehee, 2021;...]

$$\mathcal{L}_1^{\text{eff}} = \bar{n} \left(i\not{\partial} - m_n + \frac{g_n e}{2m_n} \sigma^{\mu\nu} F_{\mu\nu} \right) n$$

$$+ \bar{\chi} (i\not{\partial} - m_\chi) \chi + \varepsilon (\bar{n} \chi + \bar{\chi} n)$$

largest
dark sector
mass

mediates $n \rightarrow \chi\gamma$ (or $\Lambda \rightarrow \chi\gamma$) and more

Energy Loss Observables

Can be particularly sensitive to new physics,
such as dark matter capture, BSM decays

Observables include...

Pulsar spin down measurements...



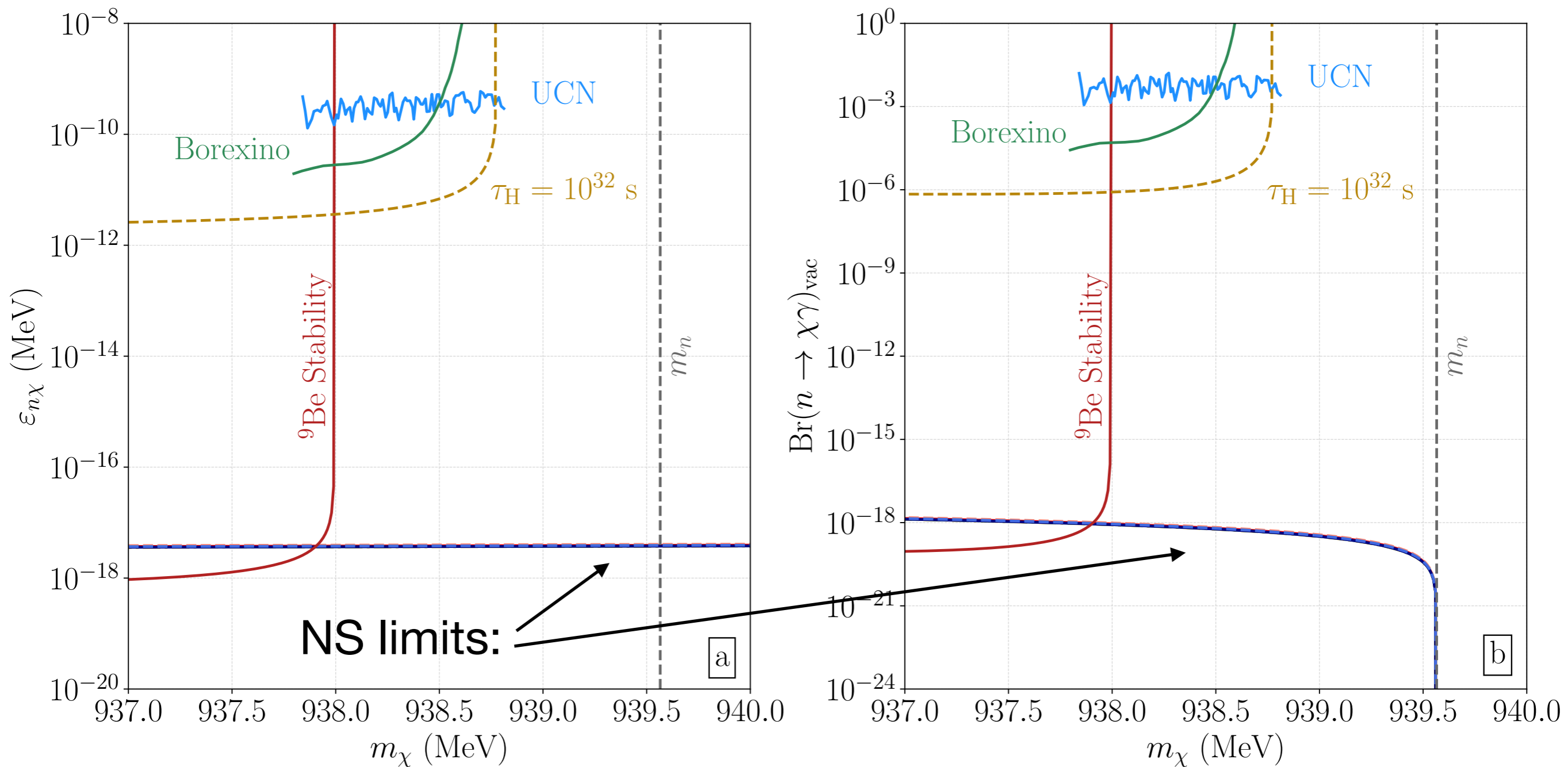
Measurements of pulsar binary
orbital period decay rates
(Recall Hulse-Taylor binary pulsar, GR test!)

Temperature measurements of the cooling history of a neutron star (or a white dwarf)...

G. Bussoni, talk in F

These effects *constrain* the BSM possibilities...

Exclusion Limits (at 2σ)



N.B. dark sector choices

[J. Berryman, SG, M. Zakeri, PRD 109, 023021 (2024)]

The Theory of Neutron Stars is not settled

Much ongoing discussion, also at QCHSC 24!

QCD @ $T = 0$; $\rho_c \simeq (1 - \text{few}) \rho_0$

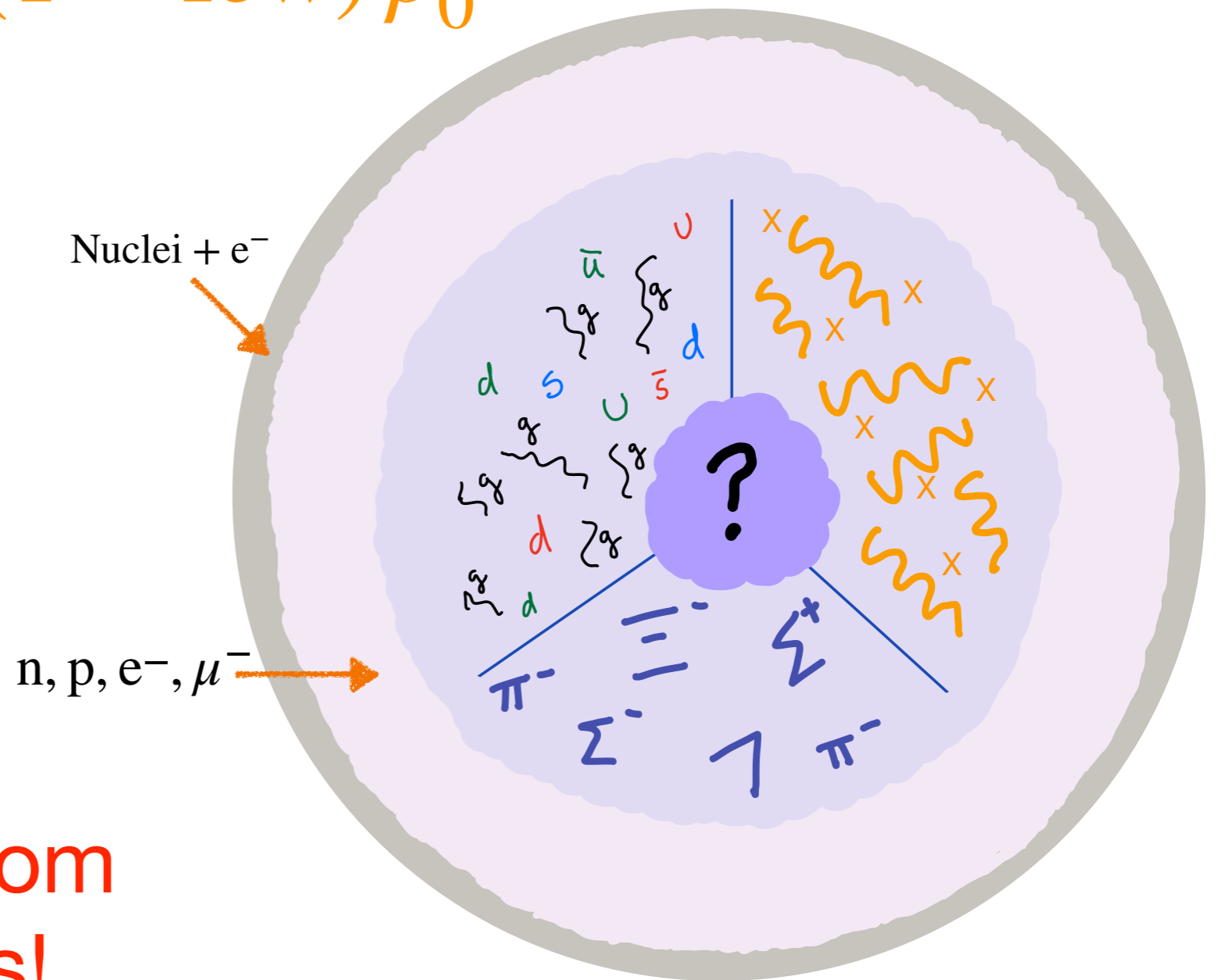
Description?

EoS?

Role of Hyperons?

Observed ns's
already constrain
dark matter, dark
sector models!

Further constraints from
studies of energy loss!



[Berryman, SG, & Zakeri, 2022;
after Baym & Pethick, 1975]

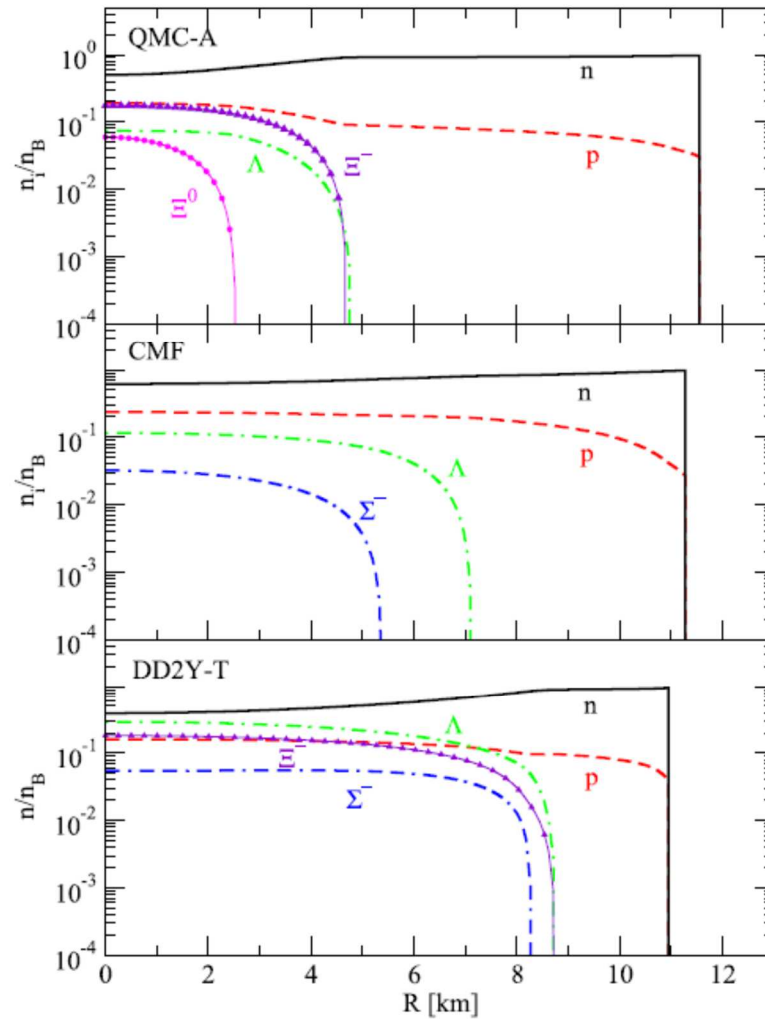
Hyperons in Neutron Star Cores?

- Nuclear matter traditionally n, p and e
- Charge neutral so $\mu_n = \mu_p + \mu_e$
- But density in typical NS reaches 2-3 times nuclear matter density (i.e. 0.32 to 0.48 fm⁻³)
- Fermi level of n's higher than m_Λ AND time to form core in supernova explosion ~seconds
- Whereas WEAK interactions, which can change strangeness take nano-seconds:
hence hyperons CAN appear
- They are stable because decays are forbidden by the Pauli Principle

The Hyperon Puzzle

- Energetically hyperons must appear
- But when new species appears the momentum of lowest energy level ~zero
- Hence the pressure is lower than nucleons alone
- A high mass NS requires high pressure in the core and this appears impossible if hyperons are present
- How are 2 M_o NS possible?

Radial Distribution of Hyperons In Heavy NS



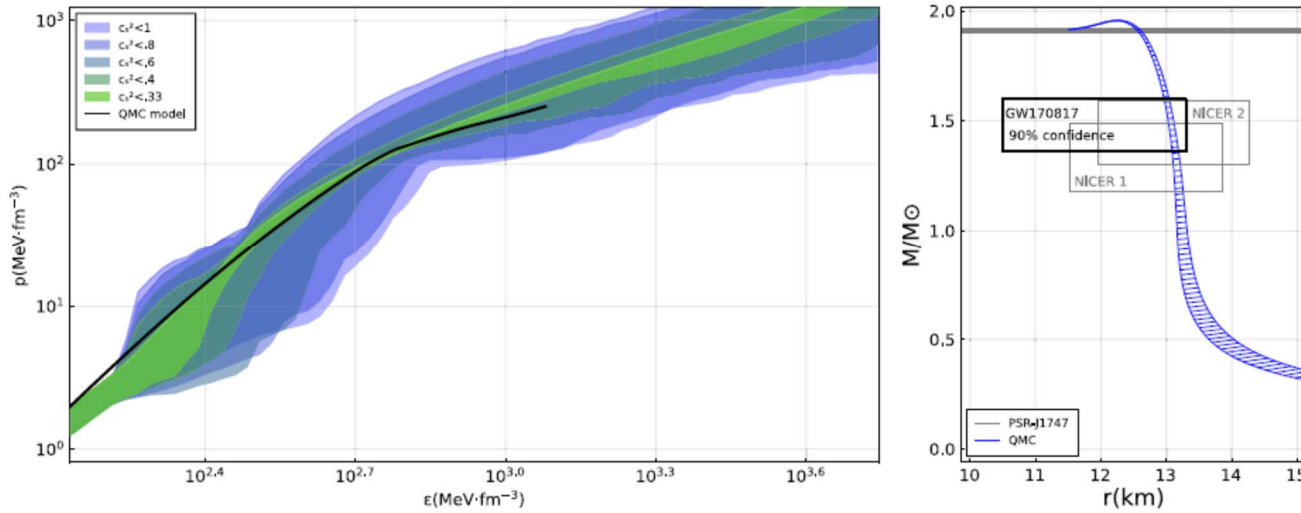
CMF: Papazoglou et al. (1999)

DD2Y-T : Typel et al. (2010)

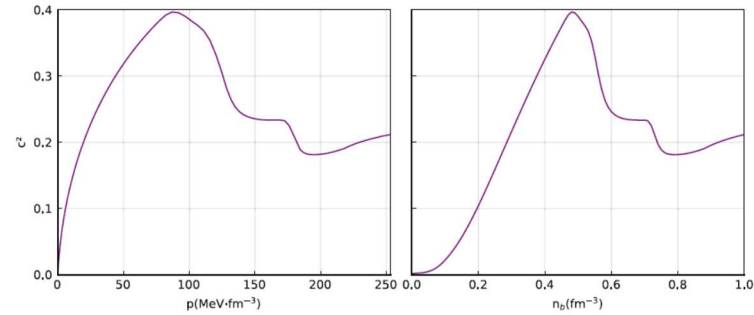
MNRAS 502, 3476–3490 (2021)

On the sound speed in hyperonic stars

T.F. Motta ^{a,*}, P.A.M. Guichon ^b, A.W. Thomas ^a



**Follow up on Annala et al.,
Nature Physics (2020) model
independent EoS based on
speed of sound interpolation
between low and high density
- claim low value implies quark matter**



Get to know GW230529

Full name GW230529_181500

Discovered on 29 May 2023 at 18h15 UTC

most likely a merger between a
Neutron Star & Black Hole (NSBH)



~1.4 M_{\odot}



~3.6 M_{\odot}

Most symmetric NSBH event so far

more likely than prior GW NSBHs to have the neutron star
ripped apart by the black hole

@astronerdika

~ 650 million light years away

Detectors



- Offline OR not operational
- Online BUT not used for analysis*
- Online AND used for analysis

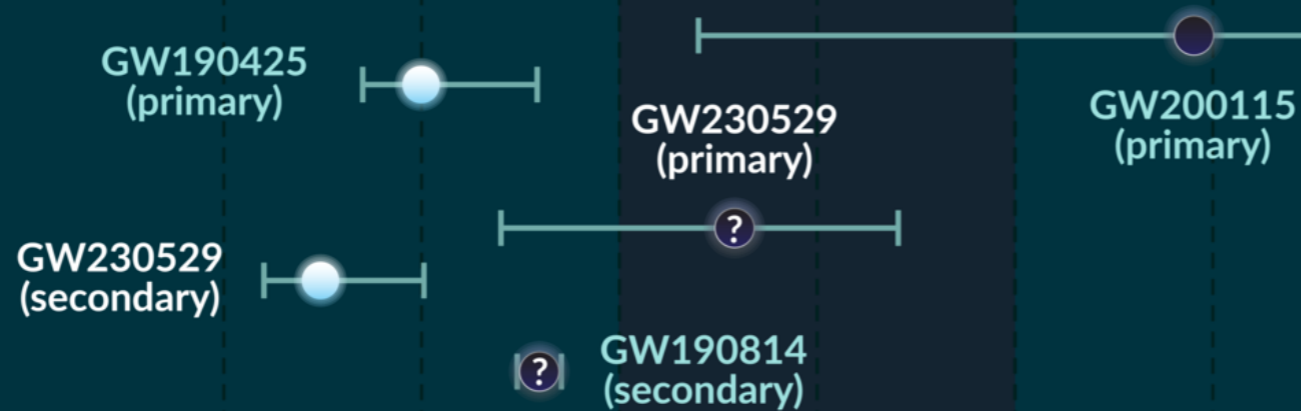
Primary object in lower mass gap
further supports that this region is not empty



* Although the KAGRA detector was in observing mode, its sensitivity was insufficient to impact the analysis of GW230529

FILLING THE MASS ↔ GAP

with observations of compact binaries from gravitational waves



Mass of compact object (M_{\odot}) 1 2 3 4 5 6

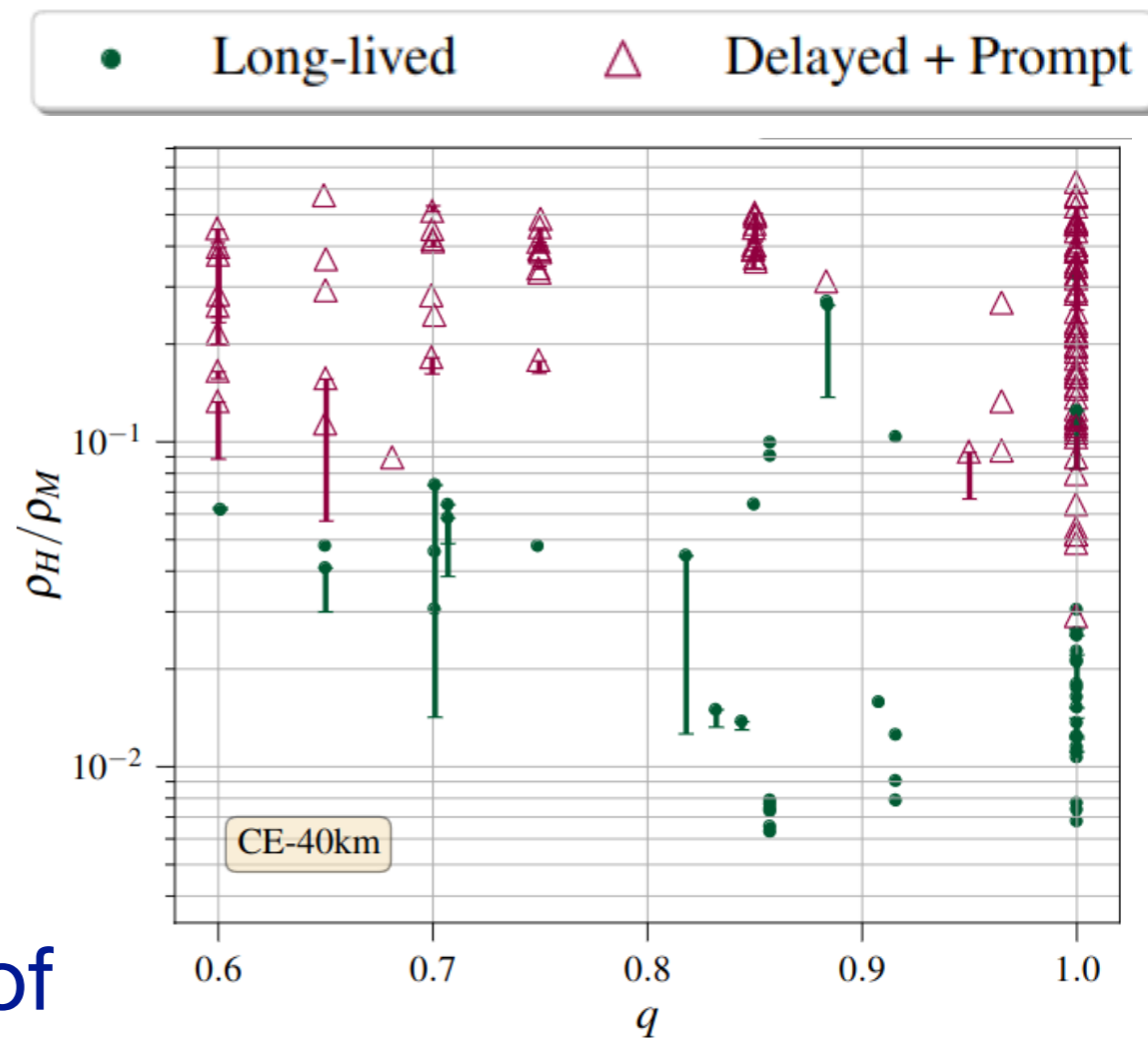
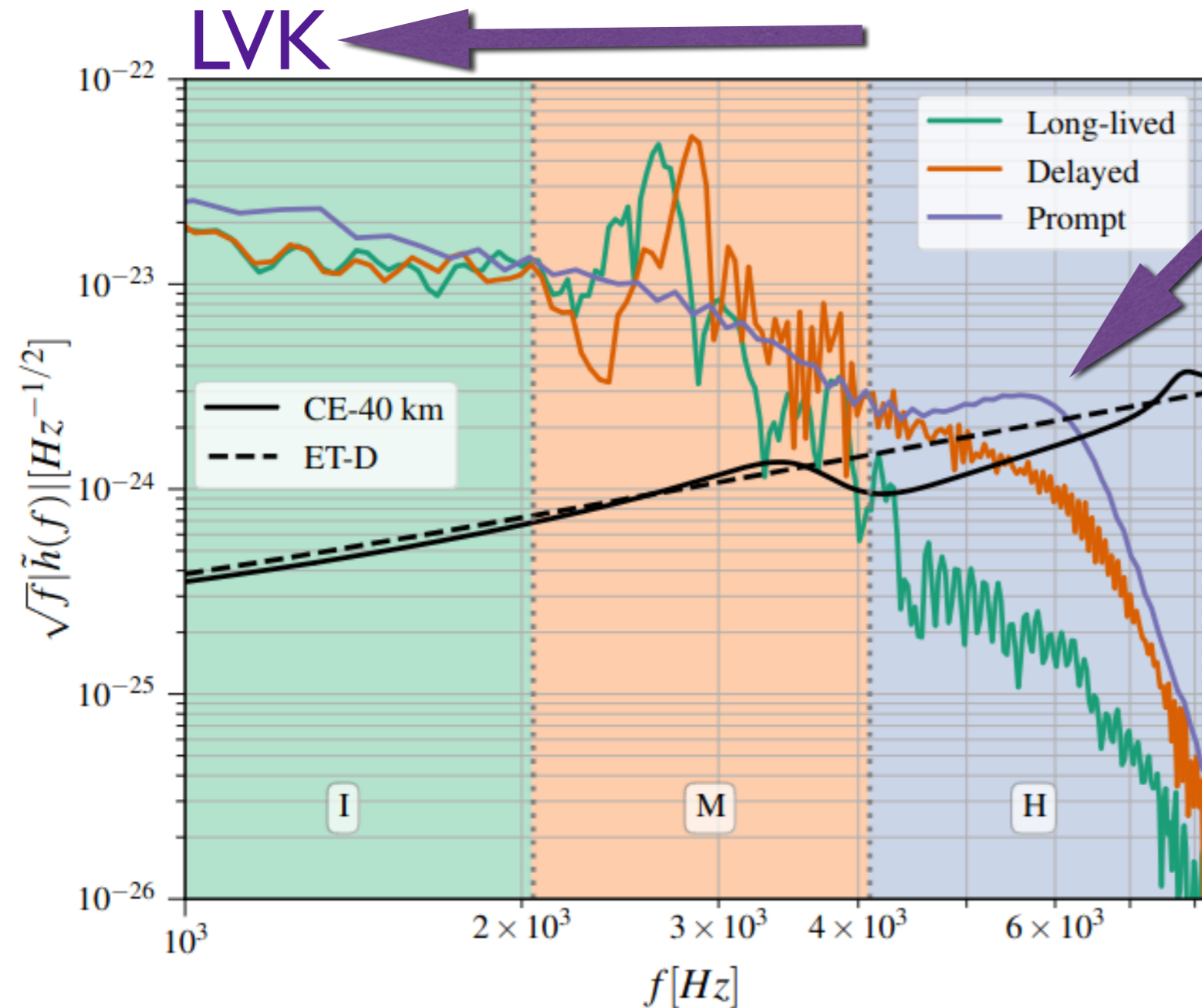
Includes components of compact binary mergers detected with a False Alarm Rate (FAR) of less than 0.25 per year

Neutron star or black hole?

NS/BH transition speaks to EoS, rotation, ...

Improvements in sensitivity opens new windows on new physics

Direct Detection of Black Hole Formation



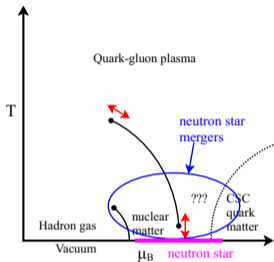
[Arnab Dhani, Radice,
Schuette-Engel, SG+, PRD 109
044071 (2024)]

Detectability of
QCD phase transitions also possible

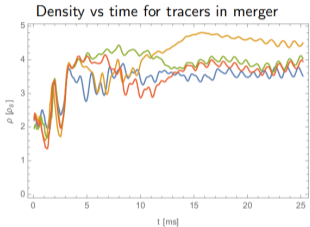
[Aviral Prakash ..., PRD 109 103008 (2024)]

What do neutron star mergers offer that neutron stars don't?

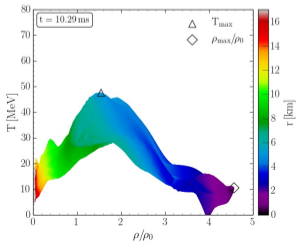
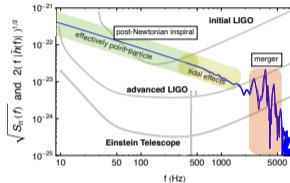
1) Probe interior of QCD phase diagram



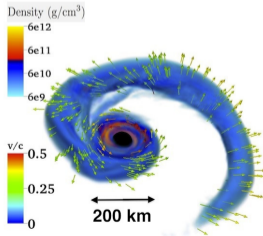
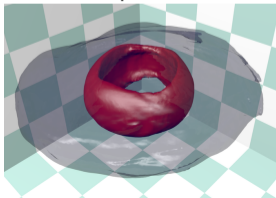
2) Dramatic variation of thermodynamic properties



3) Unique multimessenger observables



Temperature



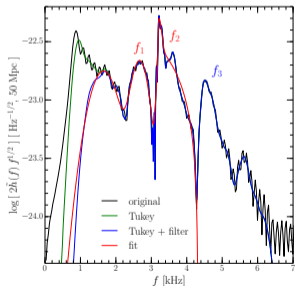
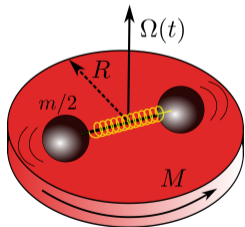
How can BSM physics show up in NS mergers? (1/3)

1) BSM particles that are trapped/accumulated in the merger remnant

- a) Oscillation modes
- b) Transport properties

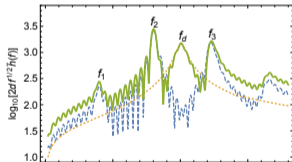
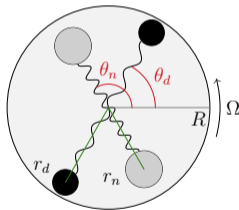
2) BSM particles that escape the merger remnant

GW spectrum without DM



Takami, Rezzolla, Baiotti arXiv:1412.3240

GW spectrum with DM



Ellis, Hektor, Hütsi, Kannike, Marzola, Raidal, Vaskonen, arXiv:1710.05540

Dark matter fluid changes properties of GW spectrum from NS merger

How can BSM physics show up in NS mergers? (2/3)

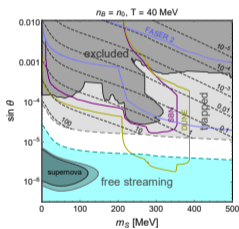
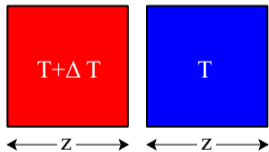
1) BSM particles that are trapped/accumulated in the merger remnant

- a) Oscillation modes
- b) Transport properties

2) BSM particles that escape the merger remnant

Thermal conductivity (& shear viscosity)

Energy transfer via particles with intermediate MFPs



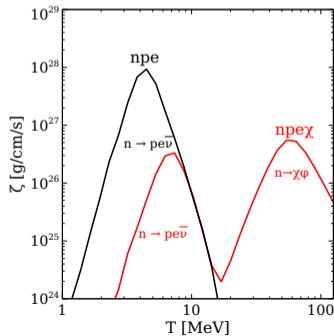
Trapped BSM particles could enhance thermal conductivity or shear viscosity.

Effects: Slow down differential rotation, and ???

Energy dissipation from compression and resultant reequilibration.

Bulk viscosity

Neutron lifetime anomaly: To reconcile, maybe neutron decays into dark sector (e.g. $n \rightarrow \chi\phi$)



Novel neutron decay yields high bulk viscosity at high T .

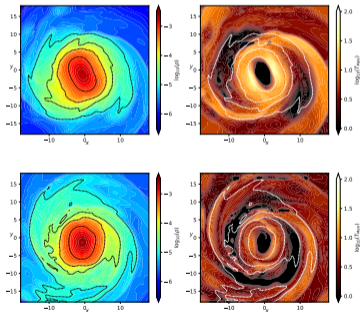
How can BSM physics show up in NS mergers? (3/3)

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- b) Transport properties

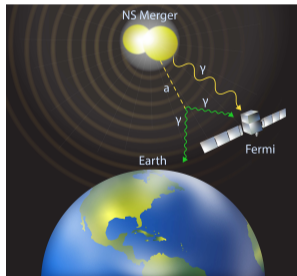
2) BSM particles that escape the merger remnant

Cooling from axion emission



Remnant cools, but no change in GW signal.
Cooling could effect ν emission and thus, kilonova

Decay of emitted axions



Emitted axions can decay to photons, generating an additional EM signal in NS mergers.
Fermi-LAT constrains axions from GW170817.

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