

Transverse-Spin Physics: Highlights and News

Werner Vogelsang
Tübingen University

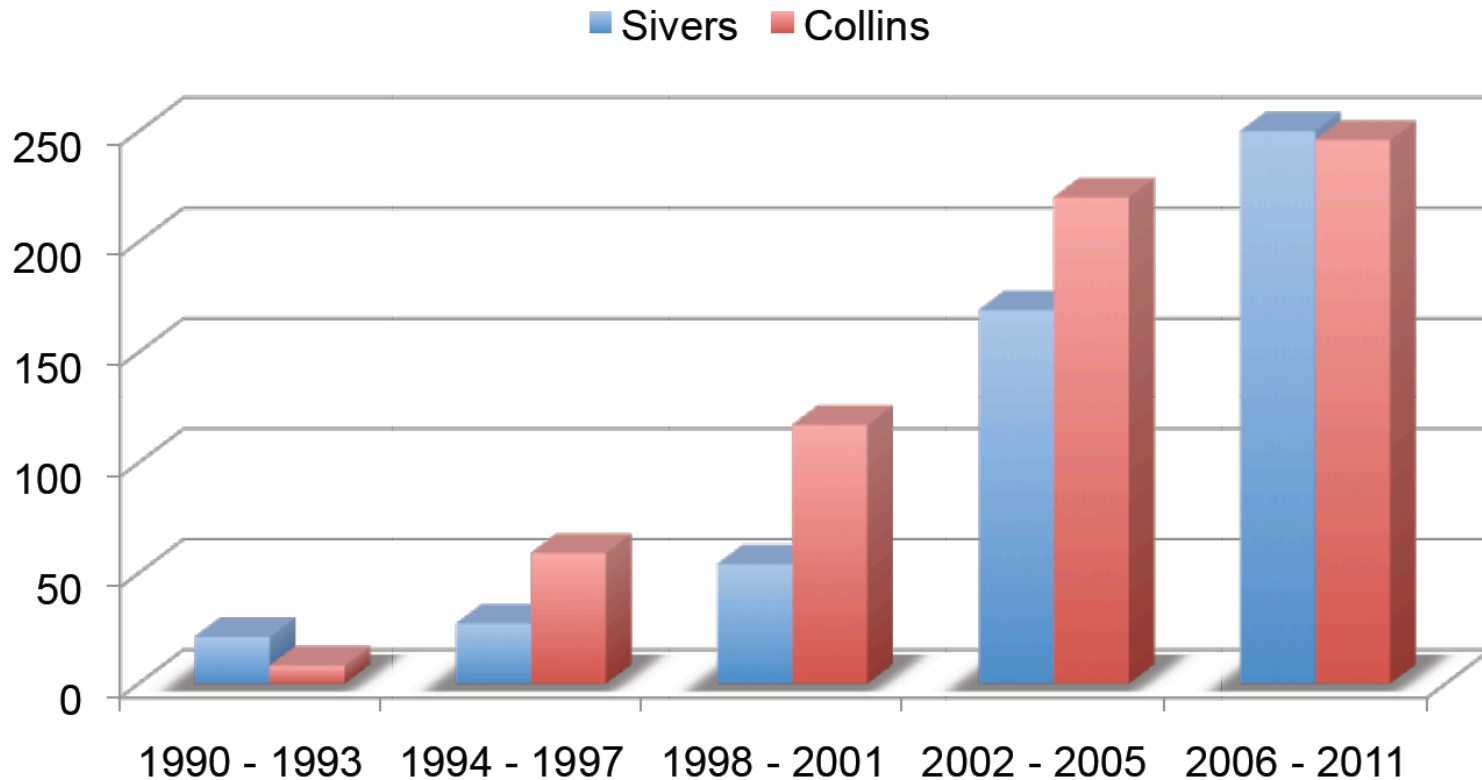
Veli Losinj, 29.08.2011

- Remarkable development of this field:
from

The sidelines in strong interaction physics

to

Center stage in our efforts to figure out QCD



Courtesy of Z. Kang, shown at 2011 RHIC Users meeting

- Numerous exciting new developments over past ~3 years
- Continue to get new insights and face new puzzles
- A lot of progress in advancing theory framework
- Have uncovered connections to other areas of QCD

Will touch on a few aspects of this
(with a strong theory and personal bias)

- Particular boost by INT program on the **Electron Ion Collider (EIC)**

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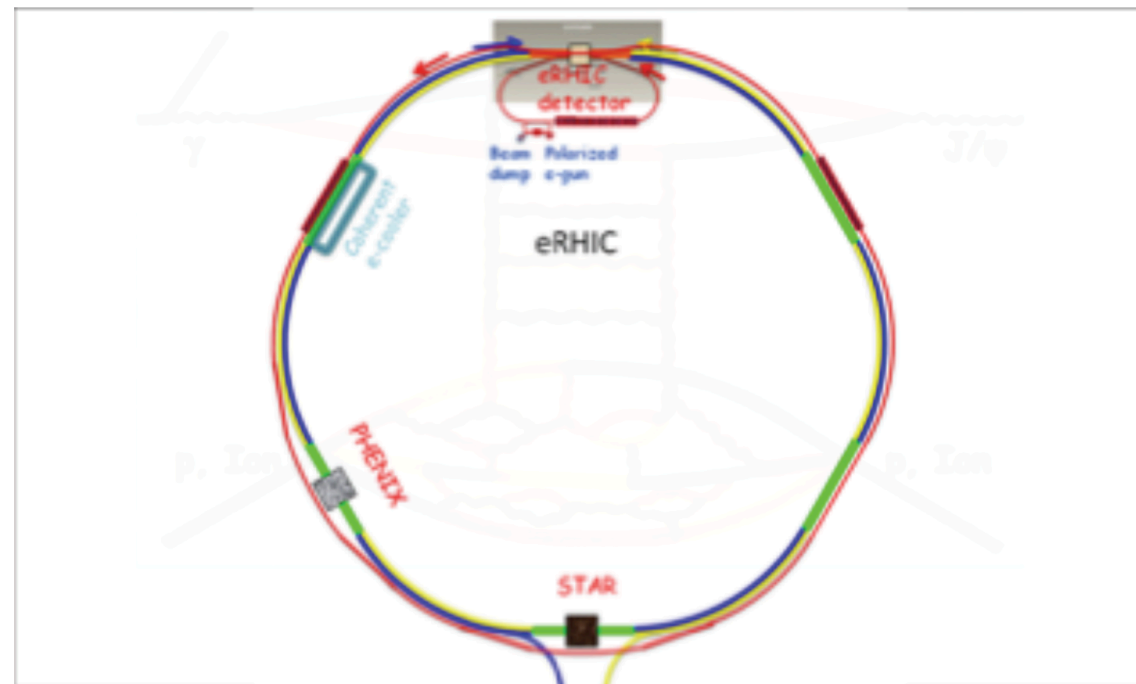
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Gluons and the quark sea at high energies: distributions, polarization, tomography

September 13 to November 19, 2010

Report from the INT program "[Gluons and the quark sea at high energies: distributions, polarization, tomography](#)"



The EIC Science case: a report on the joint BNL/INT/JLab program

Gluons and the quark sea at high energies:
distributions, polarization, tomography

arXiv:1108.1713v1 [nucl-th]

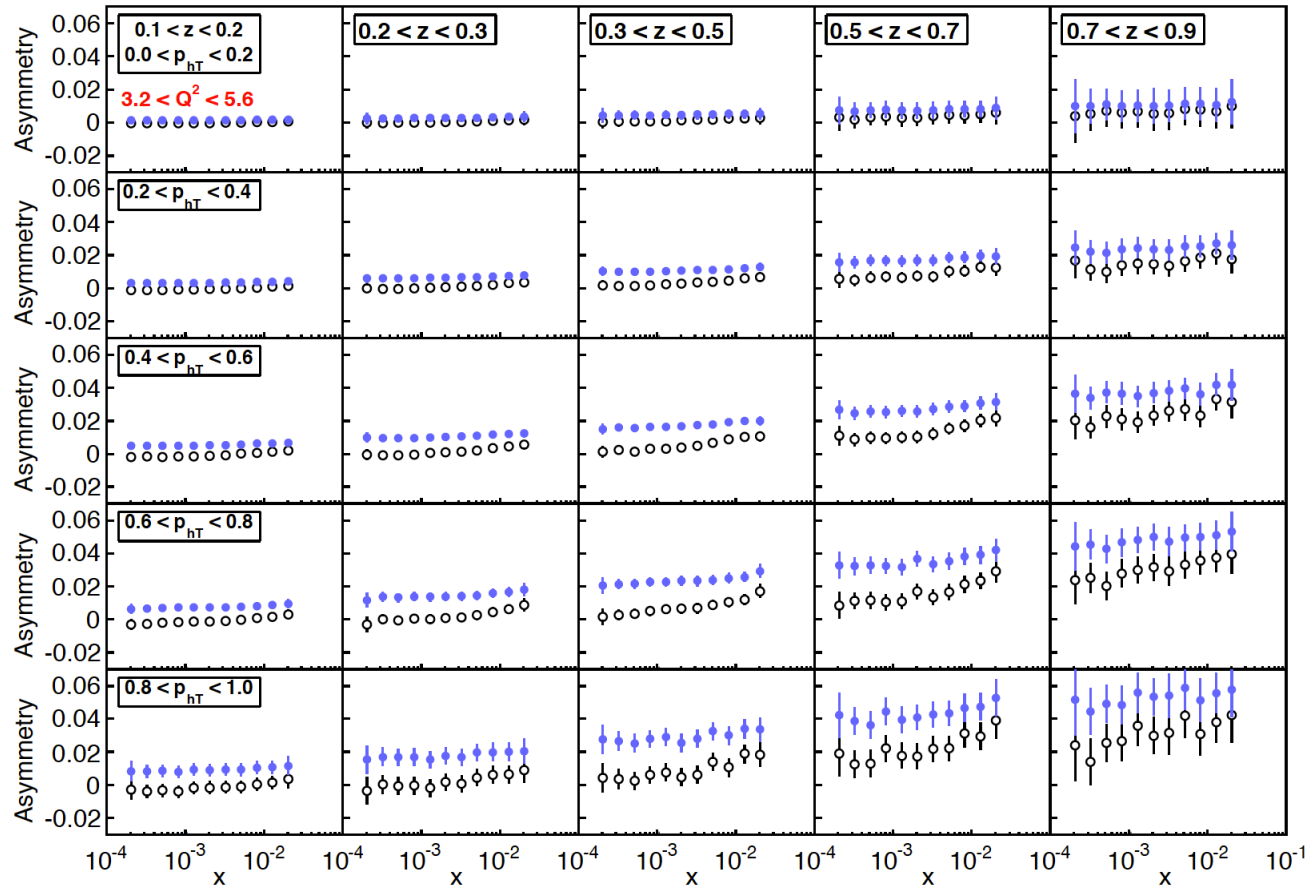
Institute for Nuclear Theory, University of Washington, USA
September 13 to November 19, 2010

Editors:

D. Boer, Universiteit Groningen, The Netherlands
M. Diehl, Deutsches Elektronen-Synchrotron DESY, Germany
R. Milner, Massachusetts Institute of Technology, USA
R. Venugopalan, Brookhaven National Laboratory, USA
W. Vogelsang, Universität Tübingen, Germany

- 1 **The spin and flavor structure of the proton** **Stratmann**
- 2 **Three-dimensional structure of the proton and nuclei: transverse momentum** **Hasch, Yuan**
- 3 **Three-dimensional structure of the proton and nuclei: spatial imaging** **Burkardt, Guzey, Sabatie**
- 4 **Input from lattice QCD** **Hasch, Yuan**
- 5 **QCD matter under extreme conditions** **Accardi, Lamont, Marquet**
- 6 **Electroweak physics** **Kumar, Li, Marciano**
- 7 **Experimental aspects** **Aschenauer, Ent**

Sivers projections for EIC:

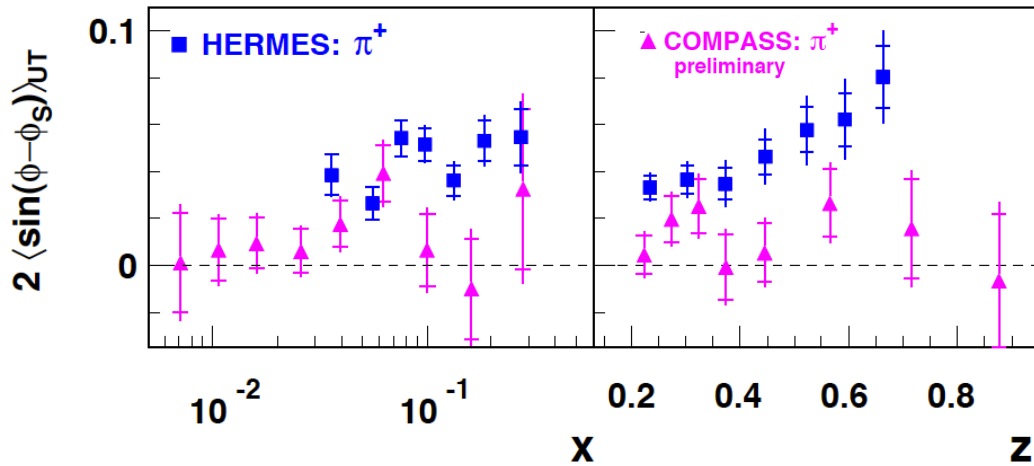
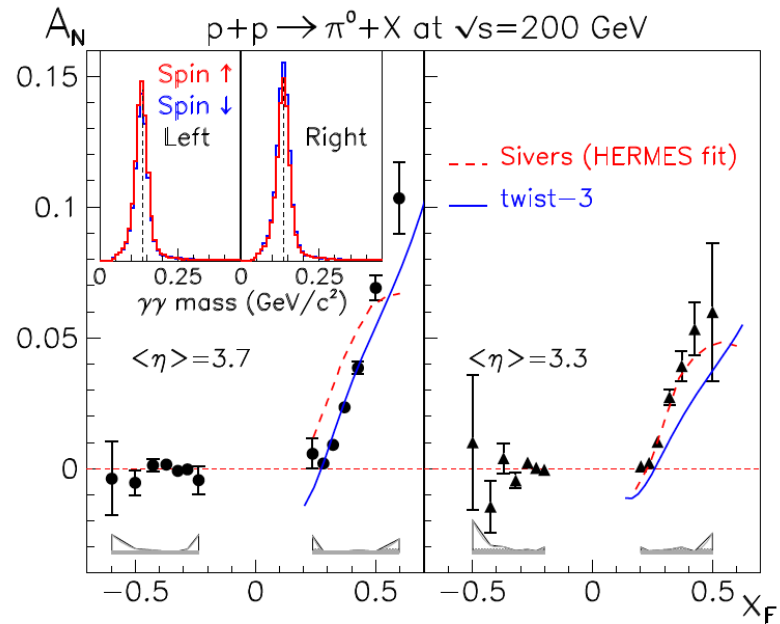
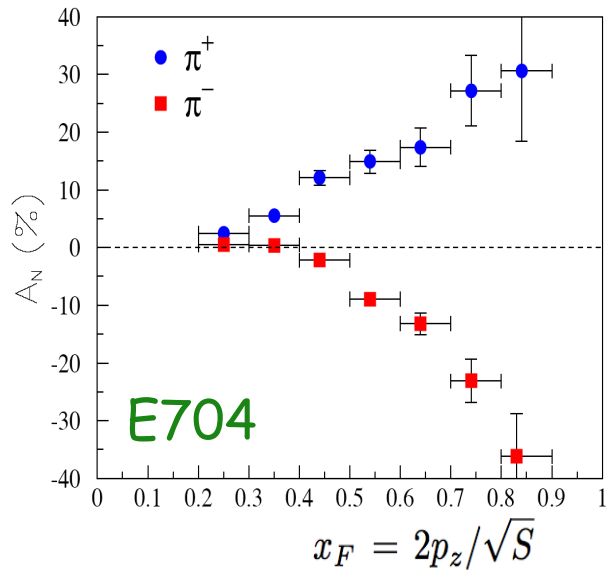


$$\pi^+, \sqrt{s} = 140 \text{ GeV}, 4/\text{fb}$$

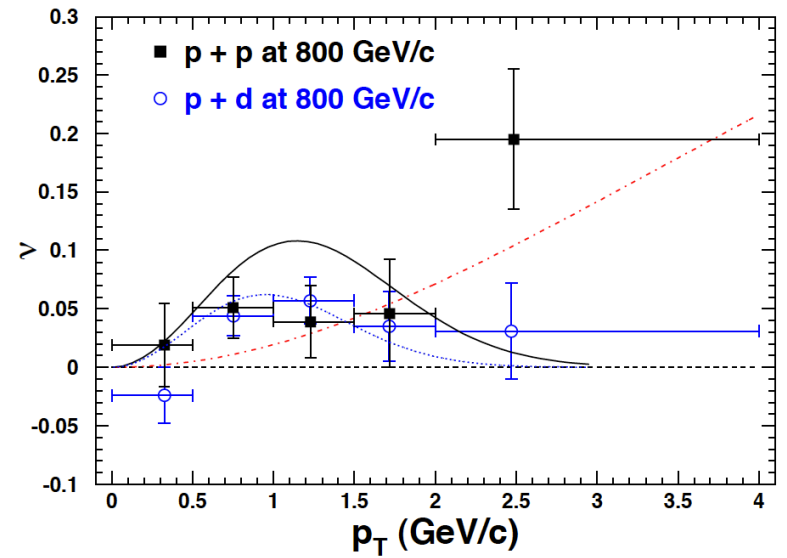
Outline:

- **Single-spin and azimuthal asymmetries**
 - TMDs, twist-3 correlations
 - factorization, evolution, resummation
 - role of gluon distributions
 - ramifications for other areas
- **Other things I am looking forward to at Transversity 2011**
- **Conclusions**

Single-spin & azimuthal asymmetries



HERMES, COMPASS,
JLab

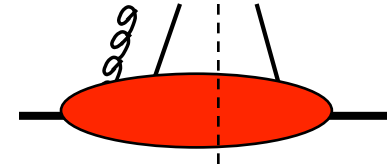
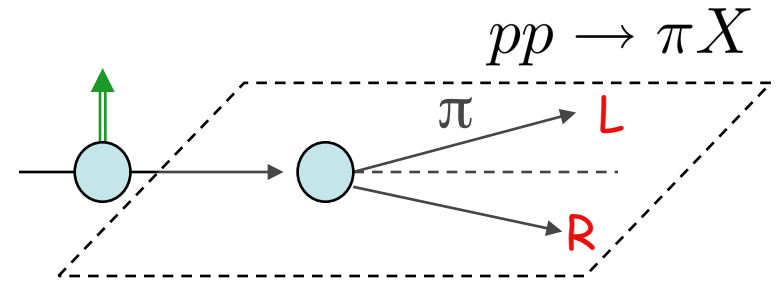


E866/NuSea

Distinguish:

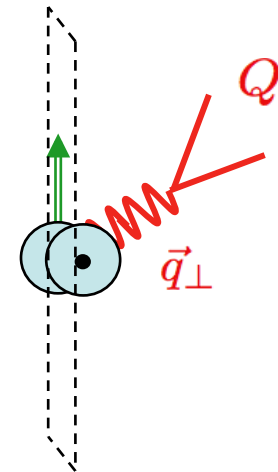
- *single-inclusive processes*

- single large scale p_{\perp}
- power-suppressed $\sim 1/p_{\perp}$ in QCD
- collinear factorization, e.g. $T_F(x,y)$
- phase in hard scattering



- *two-scale processes: small & measured $q_{\perp} \ll Q$*

- SIDIS at HERMES, COMPASS
- TMD factorization holds only for simplest observables
- crucial role of gauge links



- quark TMD correlator

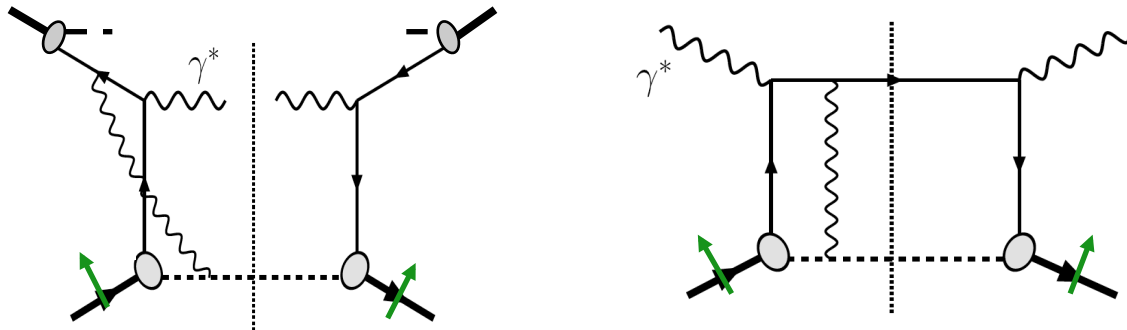
$$\Phi_{ij}^q(x, \mathbf{k}_\perp, \mathbf{S})_\eta = \int \frac{dz^- d^2 z_\perp}{(2\pi)^3} e^{ik \cdot z} \langle \mathbf{P}, \mathbf{S} | \bar{\psi}_j^q(0) \mathcal{W}_\eta(0, z) \psi_i^q(z) | \mathbf{P}, \mathbf{S} \rangle \Big|_{z^+=0}$$

- leading-power TMDs

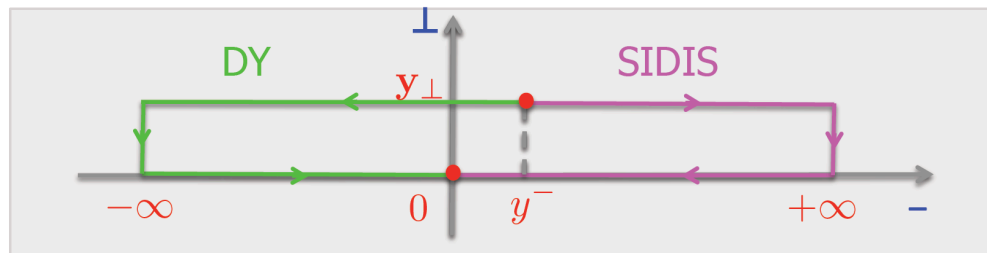
$$\begin{aligned} \frac{1}{2} \text{tr} [\gamma^+ \Phi^q(x, \mathbf{k}_\perp, \mathbf{S})] &= f_1^q(x, k_\perp) - \frac{\varepsilon^{jk} k_\perp^j S_T^k}{M} f_{1T}^{\perp q}(x, k_\perp) \\ \frac{1}{2} \text{tr} [\gamma^+ \gamma_5 \Phi^q(x, \mathbf{k}_\perp, \mathbf{S})] &= S_L g_{1L}^q(x, k_\perp) + \frac{\mathbf{k}_\perp \cdot \mathbf{S}_T}{M} g_{1T}^q(x, k_\perp) \\ \frac{1}{2} \text{tr} [i\sigma^{j+} \gamma_5 \Phi^q(x, \mathbf{k}_\perp, \mathbf{S})] &= S_T^j h_1^q(x, k_\perp) + S_L \frac{k_\perp^j}{M} h_{1L}^{\perp q}(x, k_\perp) \\ &+ \frac{(k_\perp^j k_\perp^k - \frac{1}{2} \mathbf{k}_\perp^2 \delta^{jk}) S_T^k}{M^2} h_{1T}^{\perp q}(x, k_\perp) + \frac{\varepsilon^{jk} k_\perp^k}{M} h_1^{\perp q}(x, k_\perp) \end{aligned}$$

- extensive phenomenology Anselmino, D'Alesio, Murgia, Prokudin, Schlegel, ...
- most of the TMDs involve orbital angular momentum:
Spin-orbit corr., "lensing", 3D imaging (-> models) Burkardt
- $f_{1T}^{\perp q}$, $h_1^{\perp q}$ have particularly interesting physics

- gauge link determined by hard process and its **color flow**



Brodsky, Hwang,
Schmidt
Belitsky, Ji, Yuan
Collins
Boer, Mulders, Pijlman
...



- central quest for our field: verify consequence

$$\left(f_{1T}^{\perp q} \right)_{\text{DY}} = - \left(f_{1T}^{\perp q} \right)_{\text{DIS}}$$

- goes beyond "just" check of TMD factorization
- strong motivation for Drell-Yan-type experiments

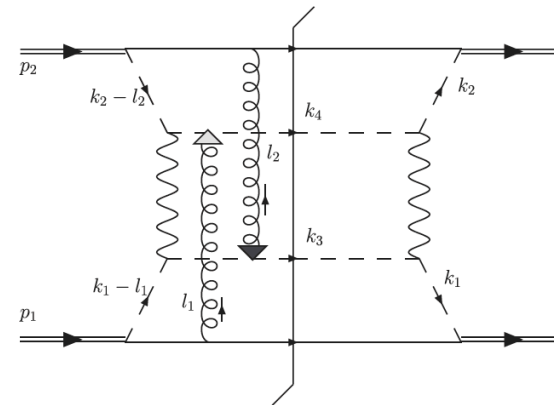
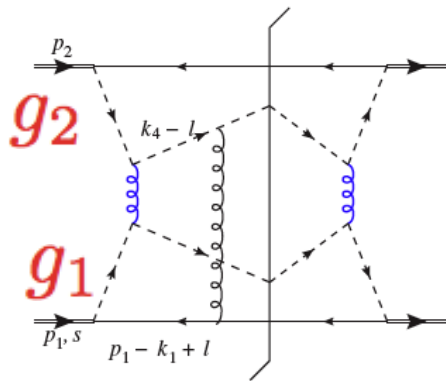
AnDY, COMPASS, E906, W bosons at RHIC

Barone et al., Anselmino et al., Yuan, WV, Schlegel et al., Kang, Qiu, Metz, Zhou



No TMD factorization for general QCD hard scattering !

Bomhof, Mulders, Pijlman;
Collins, Qiu; Mulders, Rogers



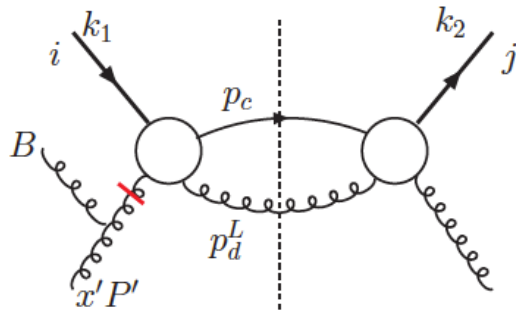
- gauge links in parton distributions “know” about full hard process
→ loss of universality
- even “generalized factorization” does not hold Mulders, Rogers
- **Dynamics!** QM interplay of probe and structure
- a lot to learn from phenomenological studies of fact. breaking
- wider ramifications for strong-interaction physics

- Twist-3 collinear correlator, e.g.

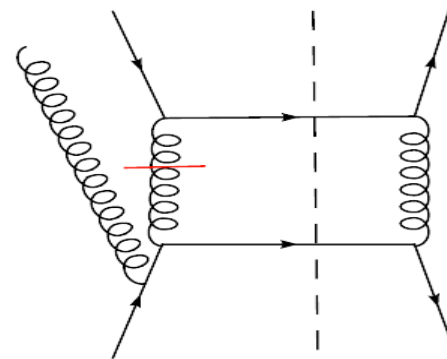
$$\begin{aligned} \tilde{\Phi}_{\alpha\beta}^{\mu}(x, x_1) &= \int \frac{dy^-}{2\pi} \frac{dy_1^-}{2\pi} e^{-ixP^+y^-} e^{i(x_1-x)P^+y_1^-} \langle P, S | \bar{\psi}_{\beta}(y^-) gF_+^{\mu}(y_1^-) \psi_{\alpha}(0) | P, S \rangle \\ &= \frac{M}{2} \left\{ \underbrace{T_F(x, x_1)}_{\text{Efremov, Teryaev, Qiu, Serman, Kanazawa, Koike, ...}} \epsilon_{\perp}^{\nu\mu} S_{\perp\nu} \not{p} + \underbrace{T_F^{(\sigma)}(x, x_1)}_{\text{Efremov, Teryaev, Qiu, Serman, Kanazawa, Koike, ...}} i\gamma_{\perp}^{\mu} \not{p} + \dots \right\}_{\alpha\beta} \end{aligned}$$

Efremov, Teryaev, Qiu, Serman, Kanazawa, Koike, ...

- accessible in A_N for inclusive processes such as $pp \rightarrow \pi X$



soft-gluon pole
 $T_F(x, x)$



soft-fermion pole
 $T_F(x, 0)$

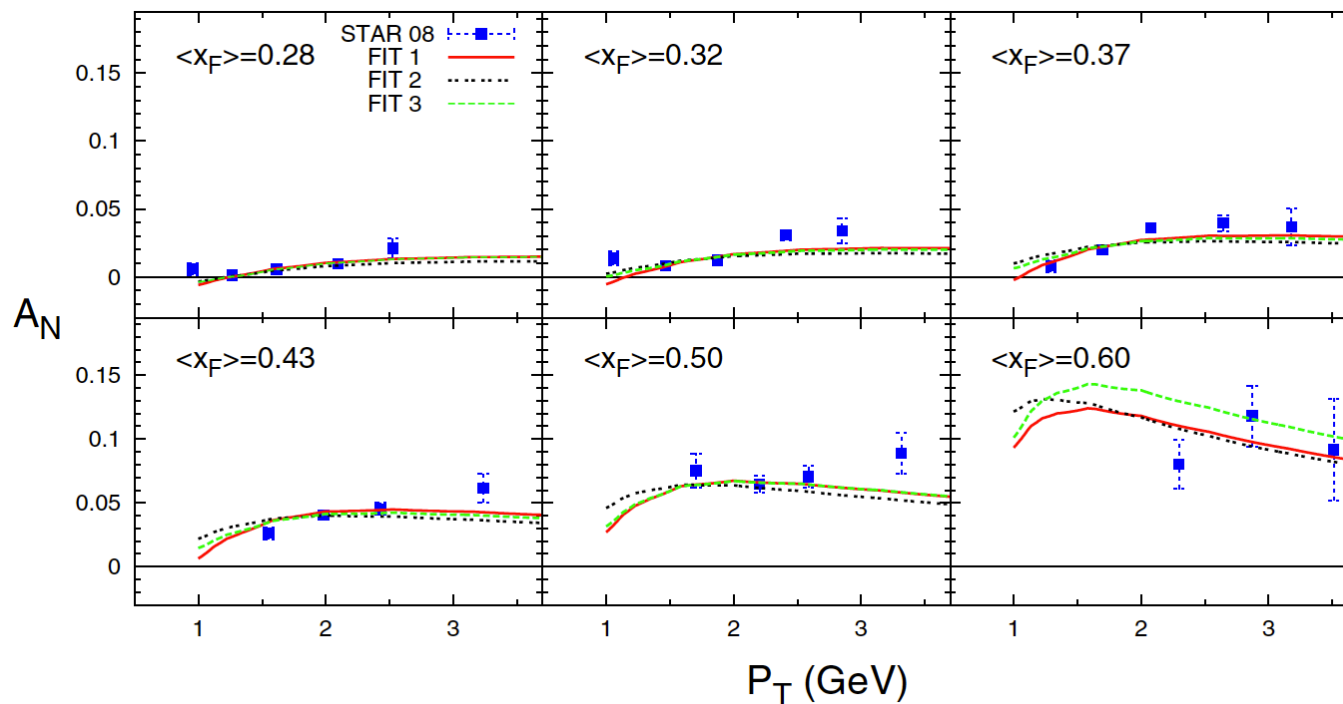
- factorization and universality



Most complete (to date) analysis of $pp \rightarrow \pi X$ data

Kanazawa, Koike

- better description of p_T dependence through interplay of SGP and SFP contributions



We have learned that TMDs and twist-3 functions are closely related

- at operator level:

Boer, Mulders, Pijlman
Ma, Wang

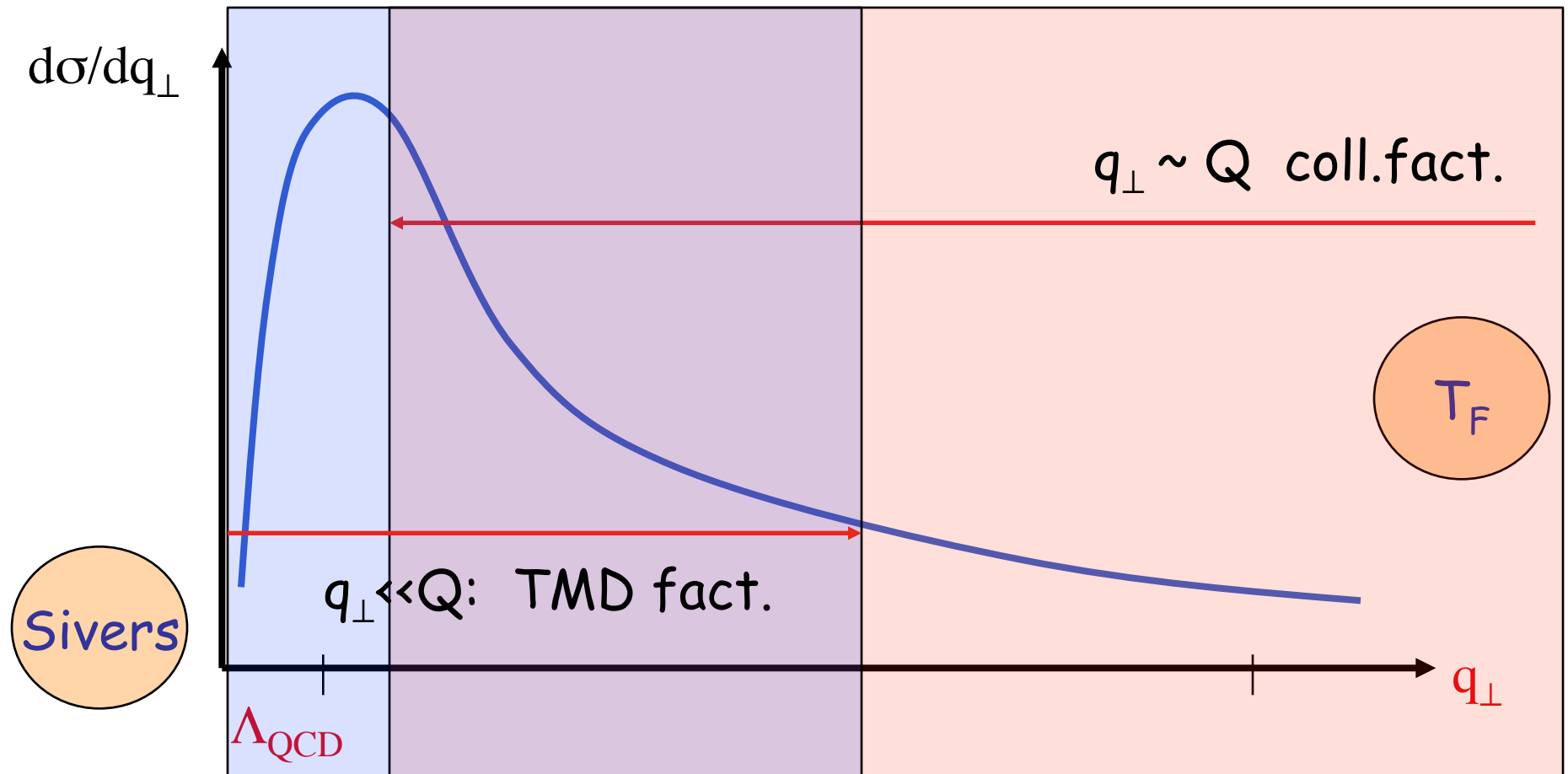
$$T_F(x, x) = - \int \frac{d^2 \vec{k}_\perp}{2\pi} \frac{\vec{k}_\perp^2}{M^2} (f_{1T}^\perp(x, k_\perp))_{\text{DIS}}$$

$$T_F^{(\sigma)}(x, x) = - \int \frac{d^2 \vec{k}_\perp}{2\pi} \frac{\vec{k}_\perp^2}{M^2} (h_1^\perp(x, k_\perp))_{\text{DIS}}$$

(caveat: renormalization)

- r.h.s. realized by q_T moments of cross sections
- it means we can use SIDIS and RHIC data jointly

- at level of a physical process:



$\Lambda_{\text{QCD}} \ll q_{\perp} \ll Q$ same physics

Verified to 1-loop

Ji, Qiu, WV, Yuan; Koike, WV, Yuan
 Zhou, Yuan, Liang
 Bacchetta, Boer, Diehl, Mulders



A sign puzzle ?

Kang, Qiu, Yuan, WV
Gamberg, Kang
Boer
Kang, Prokudin

- previous twist-3 calculations of pion A_N had incorrect sign
- as a result, signs of extracted T_F functions **to be reversed**

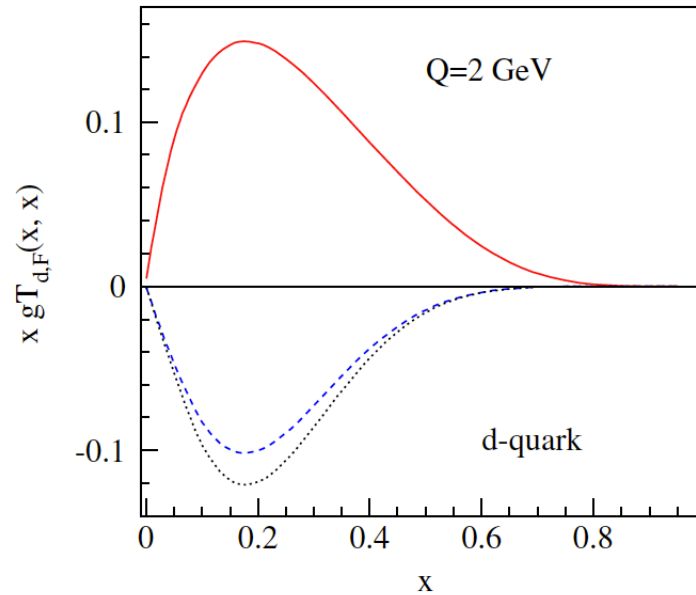
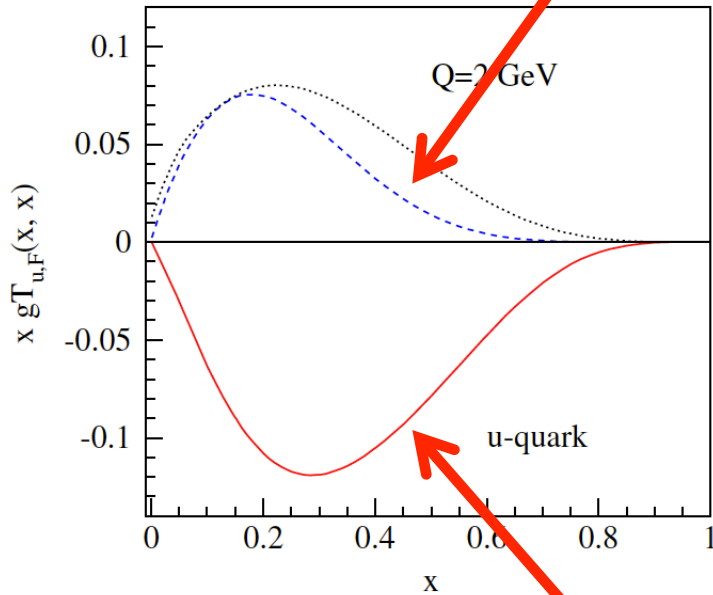
- use
$$T_F(x, x) = - \int \frac{d^2 \vec{k}_\perp}{2\pi} \frac{\vec{k}_\perp^2}{M^2} (f_{1T}^\perp(x, k_\perp))_{\text{DIS}}$$

with Sivers functions extracted from SIDIS

Anselmino et al.

- and compare...

from $-\int \frac{d^2\vec{k}_\perp}{2\pi} \frac{\vec{k}_\perp^2}{M^2} (f_{1T}^\perp(x, k_\perp))_{\text{DIS}}$



Kang, Qiu, Yuan, WV

$T_F(x, x)$ from pp

- basic observation:

SIDIS π^+ : *final-state* interaction

$pp \rightarrow \pi^+ X$: *initial-state* interaction dominates in $ug \rightarrow ug$

A puzzle? What are the implications?

Possibly: inconsistency in QCD formalism for single-spin

More likely, one of the following:

- Collins-type effect dominates in $pp \rightarrow \pi X$
- Sivers fct. mostly constrained at $k_{\perp} \sim \Lambda_{\text{QCD}}$
maybe "Gaussian" large- k_{\perp} produces different sign?

$$- \int \frac{d^2 \vec{k}_{\perp}}{2\pi} \frac{\vec{k}_{\perp}^2}{M^2} (f_{1T}^{\perp}(x, k_{\perp}))_{\text{DIS}}$$

Also: UV renormalization

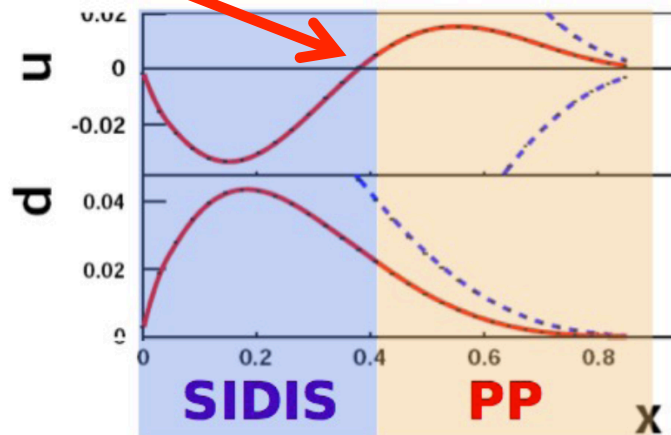
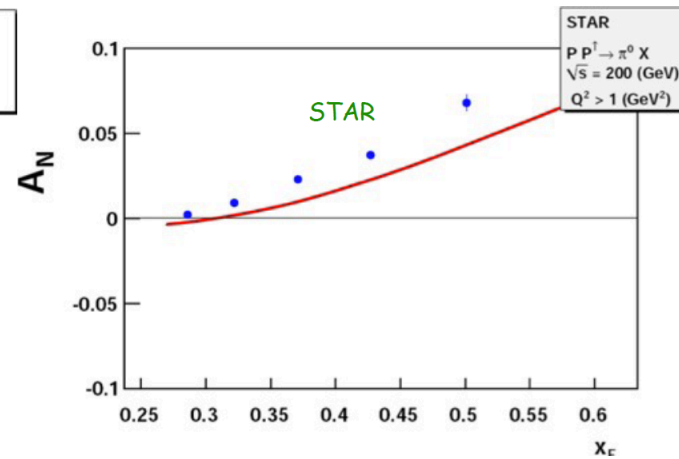
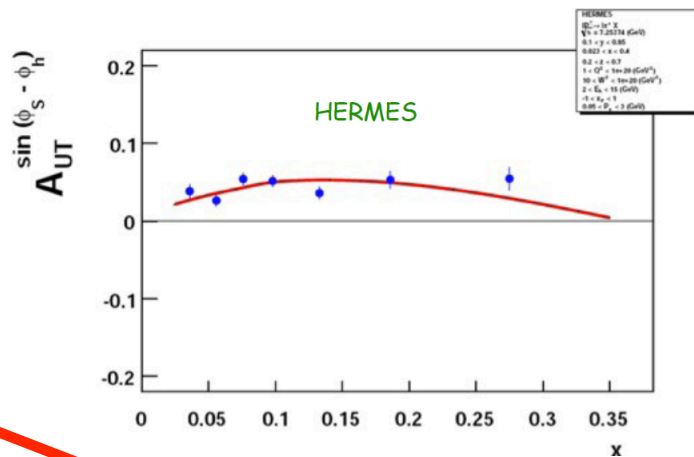
-> map out k_{\perp} dependence in experiment!

- Perhaps $T_F(x,x)$ has node in x ?

Boer
Kang, Prokudin

joint fit to SIDIS and pp data:

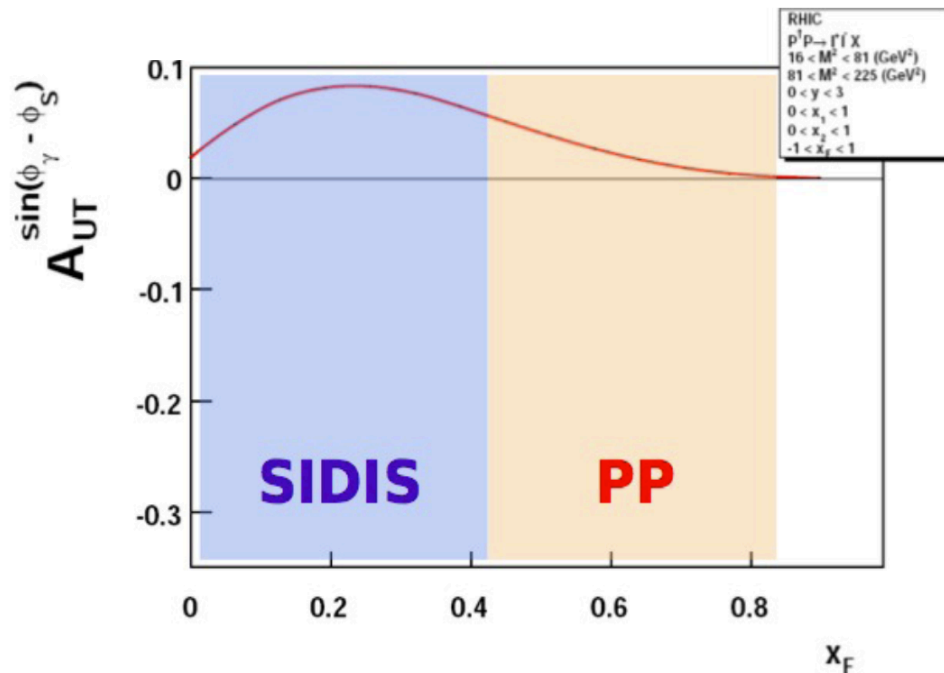
Kang, shown at RHIC Users meeting 2011



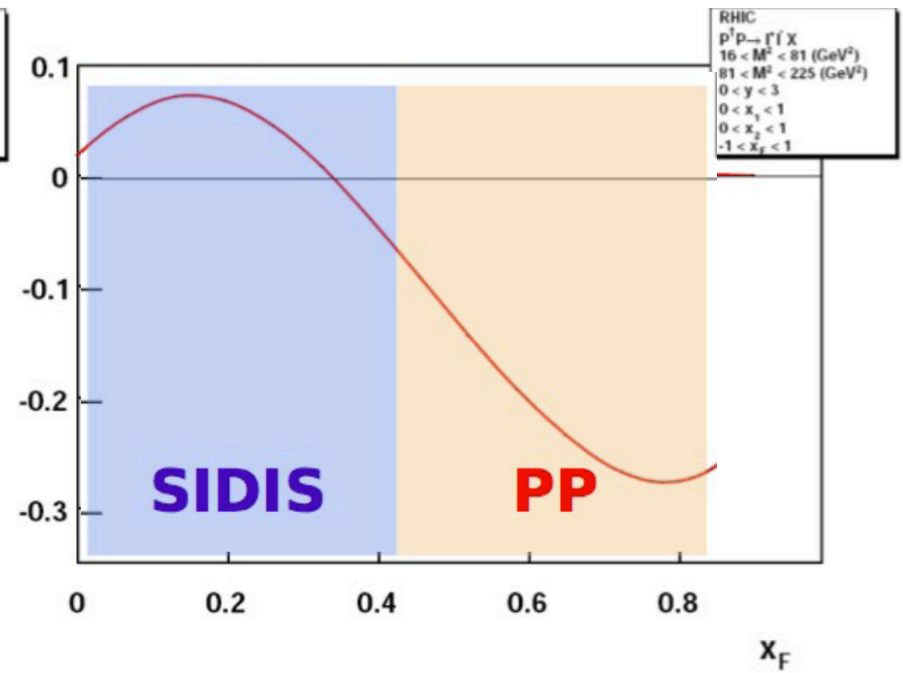
(conflict with BRAHMS data?)

Has ramifications for DY spin asymmetry:

Kang, Prokudin



w/o node



with node

Strengthens case for study of DY "sign change" !

Twist-3 calculation is the only known consistent framework for describing A_N for $pp \rightarrow \pi X$ in pQCD

- a "hybrid" TMD/Twist-3 model: Gamberg, Kang
D'Alesio, Gamberg, Kang,
Murgia, Pisano

Use TMDs (Sivers) and factorization, but with proper IS and FS interactions included.

Reproduces most of Twist-3 terms, but has more.

- how much to be learned from it ?

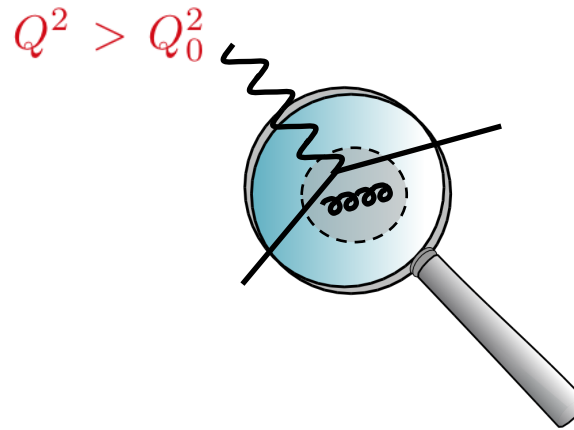
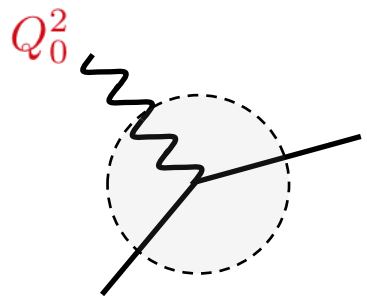


QCD corrections to single-spin observables

Mert Aybat, Rogers
Kang, Xiao, Yuan
Kang, Qiu
Yuan, WV
Braun, Manashov, Pirnay
Zhou, Yuan, Liang
Kang

- takes this field to new level
- crucial for future physics at an EIC
- closely tied to definitions of TMDs and twist-3 fcts., evolution, factorization, universality

- Like ordinary pdfs, twist-3 correlation functions are scale dependent: $T_F(x, x, Q^2)$



$$\left(\alpha_s \log \frac{Q}{Q_0} \right)^k$$

- DGLAP type evolution:

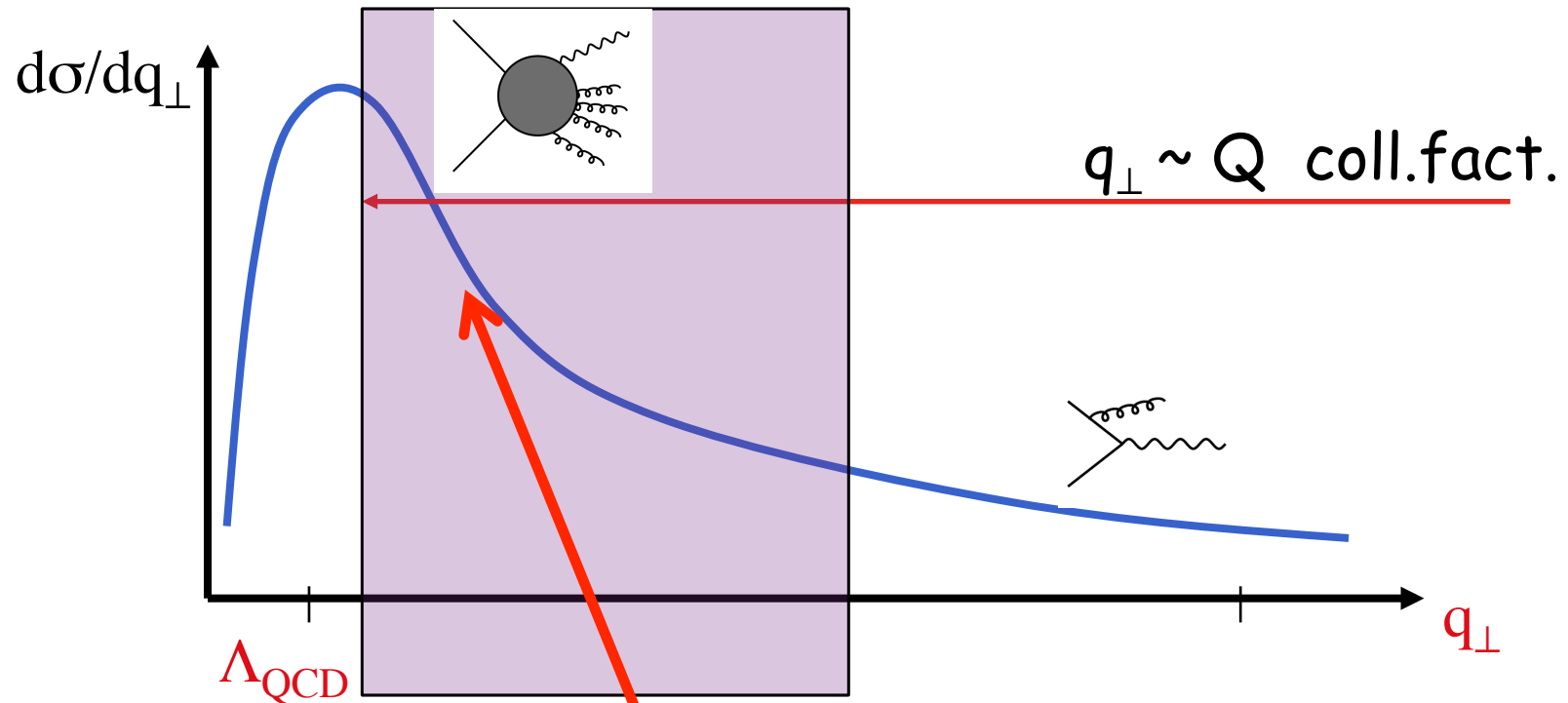
$$\frac{\partial T_F(x, x, \mu_F)}{\partial \ln \mu_F^2} = \frac{\alpha_s}{2\pi} \int_x^1 \frac{d\xi}{\xi} \left\{ P_{qq}(z) T_F(\xi, \xi, \mu_F) + \frac{C_A}{2} \left[\frac{1+z^2}{1-z} [T_{q,F}(\xi, x, \mu_F) - T_{q,F}(\xi, \xi, \mu_F)] + z T_{q,F}(\xi, x, \mu_F) \right] + \dots \right\}$$

Kang, Qiu
Yuan, WV
Braun, Manashov, Pirnay

- N.B.:** various calculations not in complete agreement
- Full NLO of Drell-Yan single-spin asymmetry (would be great to have for $pp \rightarrow \pi X \dots$)

Yuan, WV

- Evolution for TMDs ?



well-known feature: emergence of Sudakov logarithms

$$\alpha_s^k \frac{\log^{2k-1} \left(\frac{Q^2}{q_\perp^2} \right)}{q_\perp^2}$$

- can be resummed to all orders in strong coupling

(e.g. Drell-Yan, simplified)

Collins, Soper, Sterman;...

$$\frac{d\sigma}{d^2q_\perp} \sim \sigma_0 \int d^2b e^{-i\vec{b}\cdot\vec{q}_\perp} q(x_1, 1/b) \otimes \bar{q}(x_2, 1/b) e^{-\underbrace{\frac{C_F}{\pi} \int_{1/b^2}^{Q^2} \frac{dk_\perp^2}{k_\perp^2} \left(\alpha_s(k_\perp^2) \log \frac{Q^2}{k_\perp^2} + \dots \right)}_{\text{Sudakov exponent}}}$$

- idea: re-organize in terms of simple TMD-like formula

$$\frac{d\sigma}{dq_\perp^2} \sim \sigma_0 \int d^2k_{\perp,1} \int d^2k_{\perp,2} F(x_1, k_{\perp,1}, Q) \bar{F}(x_2, k_{\perp,2}, Q) \delta^{(2)}(\vec{k}_{\perp,1} + \vec{k}_{\perp,2} - \vec{q}_\perp)$$

Mert Aybat, Rogers

- this gives the evolution of TMDs

- requires a number of developments, such as
convenient definition of TMDS
treatment of light-cone divergences

Collins
(Cherednikov,
Stefanis)

- final result:

$$\begin{aligned}
 \tilde{F}_{f/P}(x, \mathbf{b}_T; \mu, \zeta_F) = & \underbrace{\sum_j \int_x^1 \frac{d\hat{x}}{\hat{x}} \tilde{C}_{f/j}(x/\hat{x}, b_*; \mu_b^2, \mu_b, g(\mu_b)) f_{j/P}(\hat{x}, \mu_b)}_A \text{ collinear piece} \\
 & \times \underbrace{\exp\left\{ \ln \frac{\sqrt{\zeta_F}}{\mu_b} \tilde{K}(b_*; \mu_b) + \int_{\mu_b}^{\mu} \frac{d\mu'}{\mu'} \left[\gamma_F(g(\mu'); 1) - \ln \frac{\sqrt{\zeta_F}}{\mu'} \gamma_K(g(\mu')) \right] \right\}}_B \sim \text{Sudakov} \\
 & \times \underbrace{\exp\left\{ g_{j/P}(x, b_T) + g_K(b_T) \ln \frac{\sqrt{\zeta_F}}{\sqrt{\zeta_{F,0}}} \right\}}_C \text{ non-perturbative piece}
 \end{aligned}$$

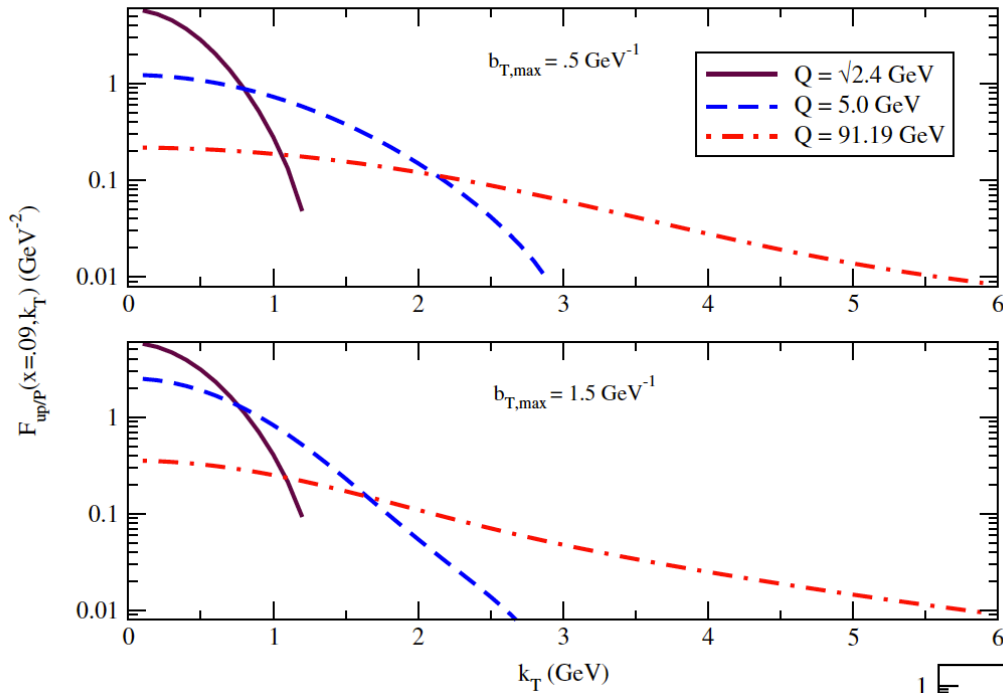
Mert Aybat, Rogers

- Note, separation of non-perturbative piece into "intrinsic" and "extrinsic" not so clear. Depends on large-b prescription.
- Appealing! (at the end of the day can use original CSS)

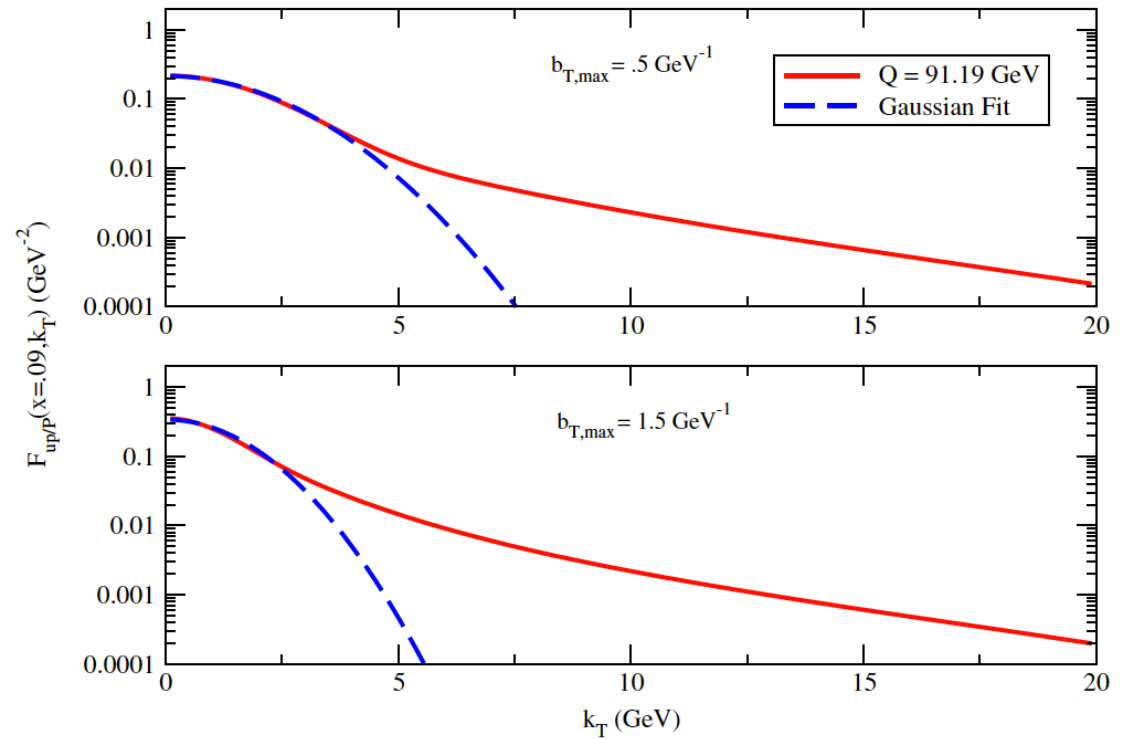
Koike, Nagashima, WV

Mert Aybat, Rogers

evolution of unpolarized TMDs



comparison to Gaussian:



The latest:

- TMD resummation for single-spin asymmetries

Kang, Xiao, Yuan

- uses CSS formalism, starting from explicit NLO calc.
- same Sudakov exponent as for spin-averaged case
- result involves $T_F(x, x)$
- important applications at EIC



Study of gluon distributions

Boer, Brodsky, Mulders, Pisano
Metz, Zhou
Qiu, Schlegel, WV
Beppu, Koike, Tanaka, Yoshida

- again, a topic also for the EIC
- have wider relevance: small- x , Higgs

- gluons & spin: focus is presently on helicity $\Delta g(x)$
- for inclusive single-spin asym.: triple-gluon correlations
- complete classification: Beppu, Koike, Tanaka, Yoshida

$$O^{\alpha\beta\gamma}(x_1, x_2) = -g(i)^3 \int \frac{d\lambda}{2\pi} \int \frac{d\mu}{2\pi} e^{i\lambda x_1} e^{i\mu(x_2-x_1)} \langle pS | d_{bca} F_b^{\beta n}(0) F_c^{\gamma n}(\mu n) F_a^{\alpha n}(\lambda n) | pS \rangle$$

$$N^{\alpha\beta\gamma}(x_1, x_2) = -g(i)^3 \int \frac{d\lambda}{2\pi} \int \frac{d\mu}{2\pi} e^{i\lambda x_1} e^{i\mu(x_2-x_1)} \langle pS | i f_{bca} F_b^{\beta n}(0) F_c^{\gamma n}(\mu n) F_a^{\alpha n}(\lambda n) | pS \rangle$$

- for example, photon production in pp:

$$\begin{aligned}
E_\gamma \frac{d\sigma}{d^3q} &= \frac{4\alpha_{em}\alpha_s M_N \pi}{S} \sum_a \int \frac{dx'}{x'} f_a(x') \int \frac{dx}{x} \delta(\hat{s} + \hat{t} + \hat{u}) \epsilon^{qpnS_\perp} \frac{1}{\hat{u}} \\
&\times \left[\delta_a \left(\frac{d}{dx} O(x,x) - \frac{2O(x,x)}{x} + \frac{d}{dx} O(x,0) - \frac{2O(x,0)}{x} \right) \right. \\
&\left. - \frac{d}{dx} N(x,x) + \frac{2N(x,x)}{x} + \frac{d}{dx} N(x,0) - \frac{2N(x,0)}{x} \right] \left(\frac{1}{N} \frac{\hat{s}^2 + \hat{u}^2}{\hat{s}\hat{u}} \right)
\end{aligned}$$

- contribution to $pp \rightarrow \pi X$ would be very interesting !

- gluon TMDs originally studied in Mulders, Rodrigues '01
- correlator

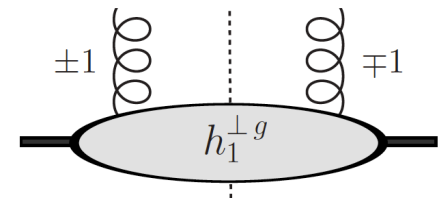
$$\Gamma_{\mu\nu;\lambda\eta}(x, \vec{k}_T) = \int \frac{dz^- d^2 z_T}{(2\pi)^3 x P^+} e^{ik \cdot z} \langle P, S | F_{\mu\nu}^\alpha(0) \mathcal{W}^{\alpha\beta}[0; z] F_{\lambda\eta}^\beta(z) | P, S \rangle \Big|_{z^+=0}$$

- leading terms

$$\Gamma_U^{+i;+j}(x, \vec{k}_T) = \frac{\delta^{ij}}{2} f_1^g + \frac{k_T^i k_T^j - \frac{1}{2} \vec{k}_T^2 \delta^{ij}}{2M^2} \underbrace{h_1^{\perp g}}_{\text{"Boer Mulders"}}$$

$$\Gamma_T^{+i;+j}(x, \vec{k}_T) = -\frac{\delta^{ij}}{2} \frac{\epsilon_T^{rs} k_T^r S_T^s}{M} \underbrace{f_{1T}^{\perp g}}_{\text{Sivers}} + \frac{i\epsilon_T^{ij}}{2} \frac{\vec{k}_T \cdot \vec{S}_T}{M} g_{1T}^{\perp g} \\ + \frac{S_T^{\{i} \epsilon_T^{j\}r} k_T^r + k_T^{\{i} \epsilon_T^{j\}r} S_T^r}{8M} h_{1T}^g + \frac{k_T^{\{i} \epsilon_T^{j\}r} k_T^r}{4M^2} \frac{\vec{k}_T \cdot \vec{S}_T}{M} h_{1T}^{\perp g}$$

- $h_1^{\perp g}$: gluons with *linear* polarization in unpolarized hadron (T-even, $\Delta L=2$)



Possibilities for measurement of $h_1^{\perp g}$:

- heavy-flavor leptonproduction $ep \rightarrow eQ\bar{Q}X$

factorization should hold \rightarrow EIC

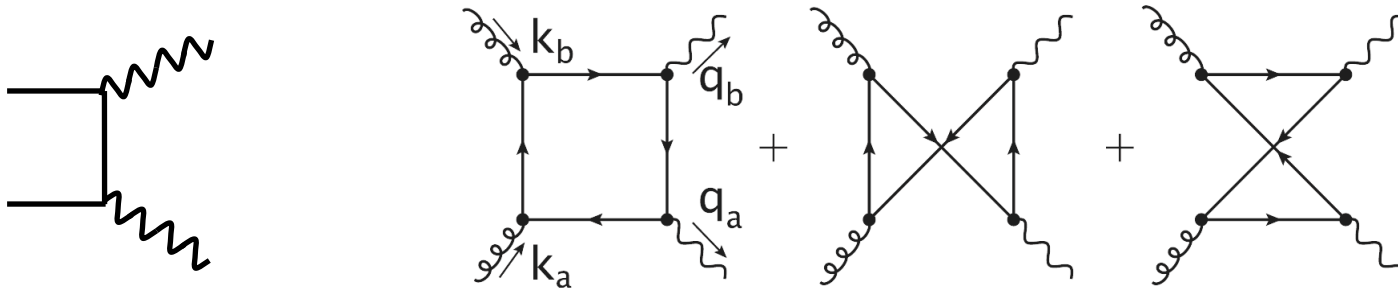
Boer, Brodsky,
Mulders, Pisano

- heavy-flavor hadroproduction $pp \rightarrow Q\bar{Q}X$

factorization breaking expected from IS, FS interactions

- double-photon production $pp \rightarrow \gamma\gamma X$

Qiu, Schlegel, WV

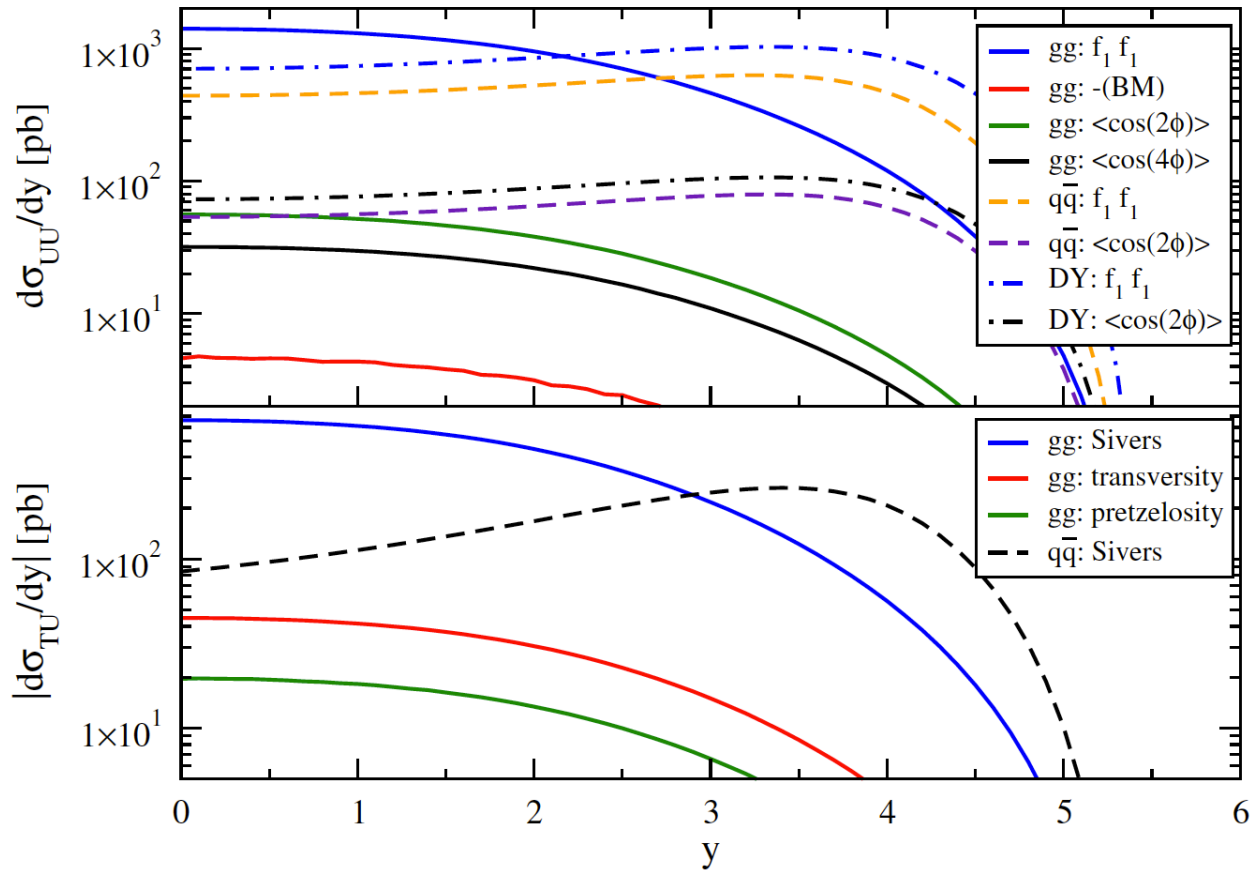


"Drell-Yan like" gauge links \rightarrow factorizes (?)

RHIC

$$\sqrt{s} = 500 \text{ GeV}$$

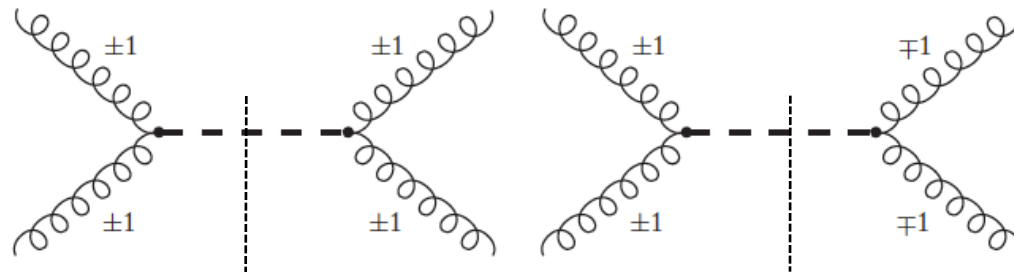
(assuming Gaussians with saturated positivity bounds)



Qiu, Schlegel, WV

- $h_1^{\perp g}$ may contribute to Higgs production !

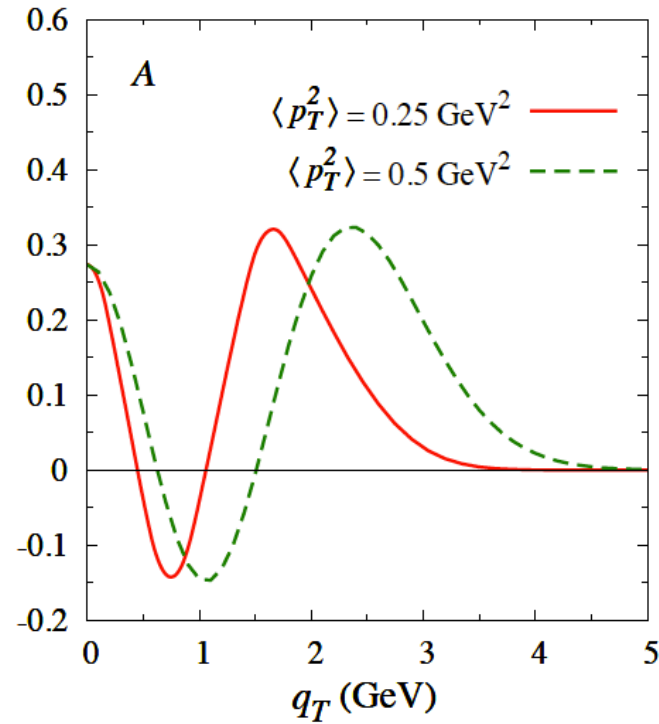
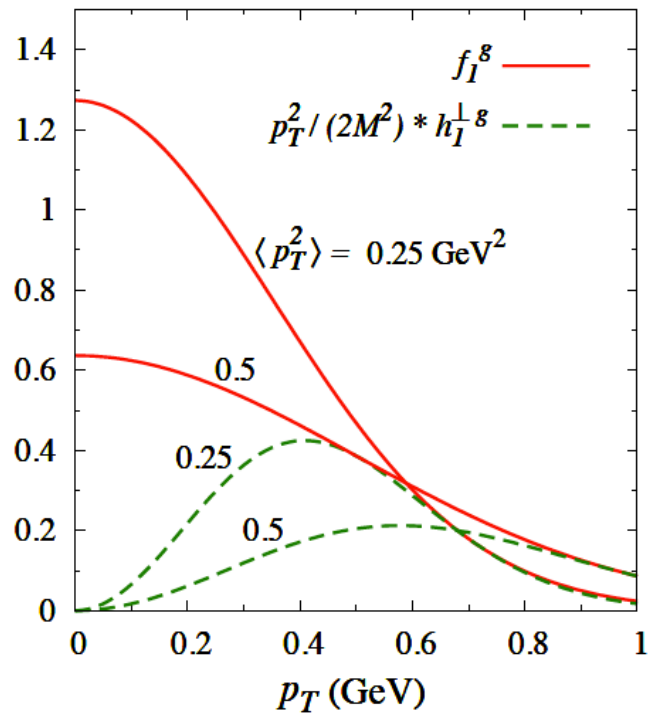
Boer, den Dunnen,
Pisano, Schlegel, WV
(in prep.)



distinguishes parity of Higgs:

Higgs = scalar:
$$\sigma \propto \mathcal{C} [f_1^g f_1^g] + \mathcal{C} [w_H h_1^{\perp g} h_1^{\perp g}]$$

Higgs = pseudoscalar:
$$\sigma \propto \mathcal{C} [f_1^g f_1^g] - \mathcal{C} [w_H h_1^{\perp g} h_1^{\perp g}]$$



$$A = \frac{\mathcal{C}[w_H h_1^{\perp g} h_1^{\perp g}]}{\mathcal{C}[f_1^g f_1^g]}$$

- could make significant non-perturbative contribution to perturbative resummation

Sun, Xiao, Yuan (in prep.)
 Also: Catani, Grazzini
 Nadolsky, Balazs, Berger, Yuan

- relevance reduced when decay & background taken into acc.

Other things I am looking forward
to at Transversity 2011

News

- spin sum rule / orbital angular momentum: a long-standing debate

$$\mathbf{J}_{QCD} = \mathbf{S}^q + \mathbf{L}^q + \mathbf{S}^g + \mathbf{L}^g$$

$$\mathbf{S}^q = \int \psi^\dagger \frac{1}{2} \boldsymbol{\Sigma} \psi d^3x,$$

$$\mathbf{L}^q = \int \psi \mathbf{x} \times (\mathbf{p} - g \mathbf{A}) \psi d^3x,$$

$$\mathbf{S}^g = \int \mathbf{E}^a \times \mathbf{A}_{phys}^a d^3x,$$

$$\mathbf{L}^g = \int E^{aj} (\mathbf{x} \times \nabla) A_{phys}^{aj} d^3x + g \int \psi^\dagger \mathbf{x} \times \mathbf{A}_{phys} \psi d^3x$$

$$A_{phys}^\mu(x) \rightarrow U(x) A_{phys}^\mu(x) U^{-1}(x),$$

$$A_{pure}^\mu(x) \rightarrow U(x) \left(A_{pure}^\mu(x) - \frac{i}{g} \partial^\mu \right) U^{-1}(x)$$

Recently:
Wakamatsu
Chen et al.
Leader

There are many observables that are sensitive to OAM.
 Question is connection to spin sum rule



Sivers \leftrightarrow OAM in a quantitative way ?

Bacchetta, Radici

- combines:
 - ◆ Ji's expression for J_q in terms of GPDs
 - ◆ connection between moment of f_{1T}^\perp and GPD \mathcal{E} (Burkardt's lensing idea)
 - ◆ joint fit of SIDIS Sivers asymmetries and magnetic moments

- OAM from Wigner distributions ?

Lorce, Pasquini

$$\rho^{[\Gamma]}(\vec{b}_\perp, \vec{k}_\perp, x, \vec{S}) \equiv \int \frac{d^2\Delta_\perp}{(2\pi)^2} \left\langle p^+, \frac{\vec{\Delta}_\perp}{2}, \vec{S} \left| \hat{W}^{[\Gamma]}(\vec{b}_\perp, \vec{k}_\perp, x) \right| p^+, -\frac{\vec{\Delta}_\perp}{2}, \vec{S} \right\rangle$$

$$\ell_z^q = \int dx d^2k_\perp d^2b_\perp (\vec{b}_\perp \times \vec{k}_\perp)_z \rho^{[\gamma^+]q}(\vec{b}_\perp, \vec{k}_\perp, x, \vec{e}_z)$$

- the latest on transversity (dihadron fragmentation)

Bacchetta, Courtoy, Radici

- experimental data

- and a lot more !

Enjoy
Transversity 2011 !