

# Transversity 2011

3<sup>rd</sup> international workshop on

Transverse Polarization Phenomena in Hard Scattering

Veli Lošinj (Croatia), Aug. 29th – Sept. 2nd 2011

# Unpolarized and Polarized Fragmentation Functions (only for light quarks in vacuum)

for a review see also  
*Parton fragmentation in the vacuum  
and in the medium*

Mini-workshop ECT\*, 25–28 Feb. 2008  
arXiv:0804.2021 [hep-ph]

**Marco Radici**





# Outline

- Unpol. 1-hadron Fragm. Functions (1h FF)  
status of “collinear” parametrizations  
what do we know about 1h “TMD” FF ?
- Pol. 1h FF: the Collins function
- Models of 1h FF
- 2h FF (or Dihadron Fragm. Functions – DiFF)  
BELLE (+BaBar?) data and parametrizations (next 2 talks)  
BELLE+HERMES (+COMPASS) data and extraction of  $h_1$  (Braun)  
extraction of  $e$  and  $h_L$  with DiFF at JLab (Avakian)
- Conclusions and Outlooks



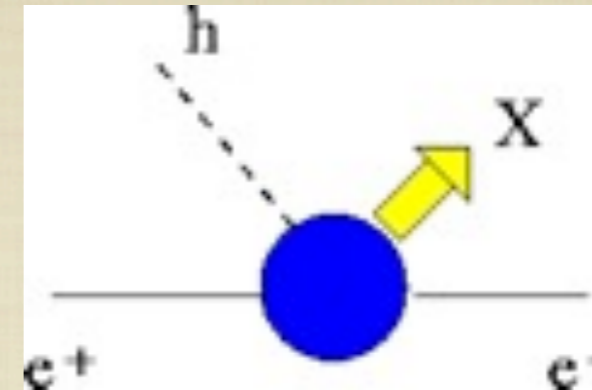
unpolarized 1h FF



# 1h FF main source of data

$$e^+e^- \rightarrow hX$$

$$h = \pi^\pm, K^\pm, K_s^0, p, \bar{p}, \Lambda, \bar{\Lambda}$$

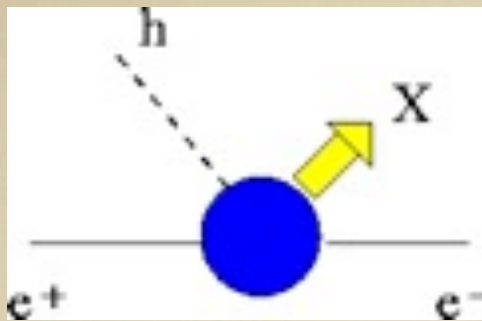


## Energy range

- $\sqrt{s} = 12-36$  GeV at DESY (ARGO, JADE, CELLO, TASSO)
  - $\sqrt{s} = 29$  at SLAC (HRS, MARK II, TPC)
  - $\sqrt{s} = 58$  at KEK (TOPAZ)
  - $\sqrt{s} = 91.2$  (Z0) at LEP-1 (ALEPH, DELPHI, OPAL)
  - “ “ “ at SLAC (SLD)
  - $\sqrt{s} = 133-209$  at LEP-2 (DELPHI, L3, OPAL)
  - $\sqrt{s} = 10.58$  ( $Y_{4S}$ ) at B-factories (BaBar, BELLE, CLEO)
- } 80's
- } '95-'06
- $5 \times 10^{-3} \leq z \leq 0.8$



# $e^+e^- \rightarrow hX$



$$\frac{1}{\sigma_{\text{tot}}} \frac{d\sigma^h}{dz} \equiv F^h(z, Q^2)$$

$$= \sum_{i=q, \bar{q}} D_i^h(z, Q^2) \quad \text{at LO}$$

$$= \sum_{i=q, \bar{q}, g} C_i(z, Q^2) \otimes D_i^h(z, Q^2) \quad \text{beyond}$$

- direct connection (at LO) to parton-to-hadron FF
- $C_i$  known up to NNLO in  $\overline{\text{MS}}$  (Mitov & Moch (2006))
- flavor analysis  $\sim \{u, d, s\} + c + b$   
except OPAL (full separation)

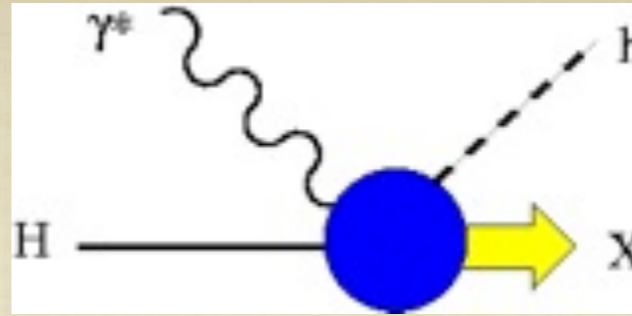
**but**

- \*  $D_g^h$  less constrained
- \* access only to  $D_q^h + D_{\bar{q}}^h = D_q^{h/\bar{h}}$  (at LO)
- \* virtuality fixed by c.m. energy  $Q = \sqrt{s}/2$



$$e^{\pm} p \rightarrow e^{\pm} h X$$

$$h = \pi^{\pm}, K^{\pm}, h^{\pm}, \Lambda, \bar{\Lambda}$$



## Energy range

- $1 \leq Q \leq 200$  GeV at HERA (H1, ZEUS, HERMES)
- $1 \leq Q \leq 5$  at CERN (COMPASS)
- $1 \leq Q \leq 10$  also at NOMAD with  $\nu_{\mu}$  probes
  
- $0.1 \leq z < 1$

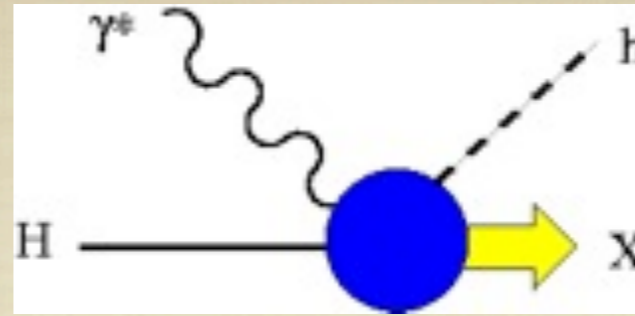
– larger phase space in  $\{z, Q^2\}$  than in  $e^+e^-$

– separate  $D_q^h$  from  $D_{\bar{q}}^h$  (at least for  $x_B \geq 0.1$ )



$$e^{\pm} p \rightarrow e^{\pm} h X$$

$$h = \pi^{\pm}, K^{\pm}, h^{\pm}, \Lambda, \bar{\Lambda}$$



- SIDIS in Breit frame

$$x_p = P_h / Q/2$$

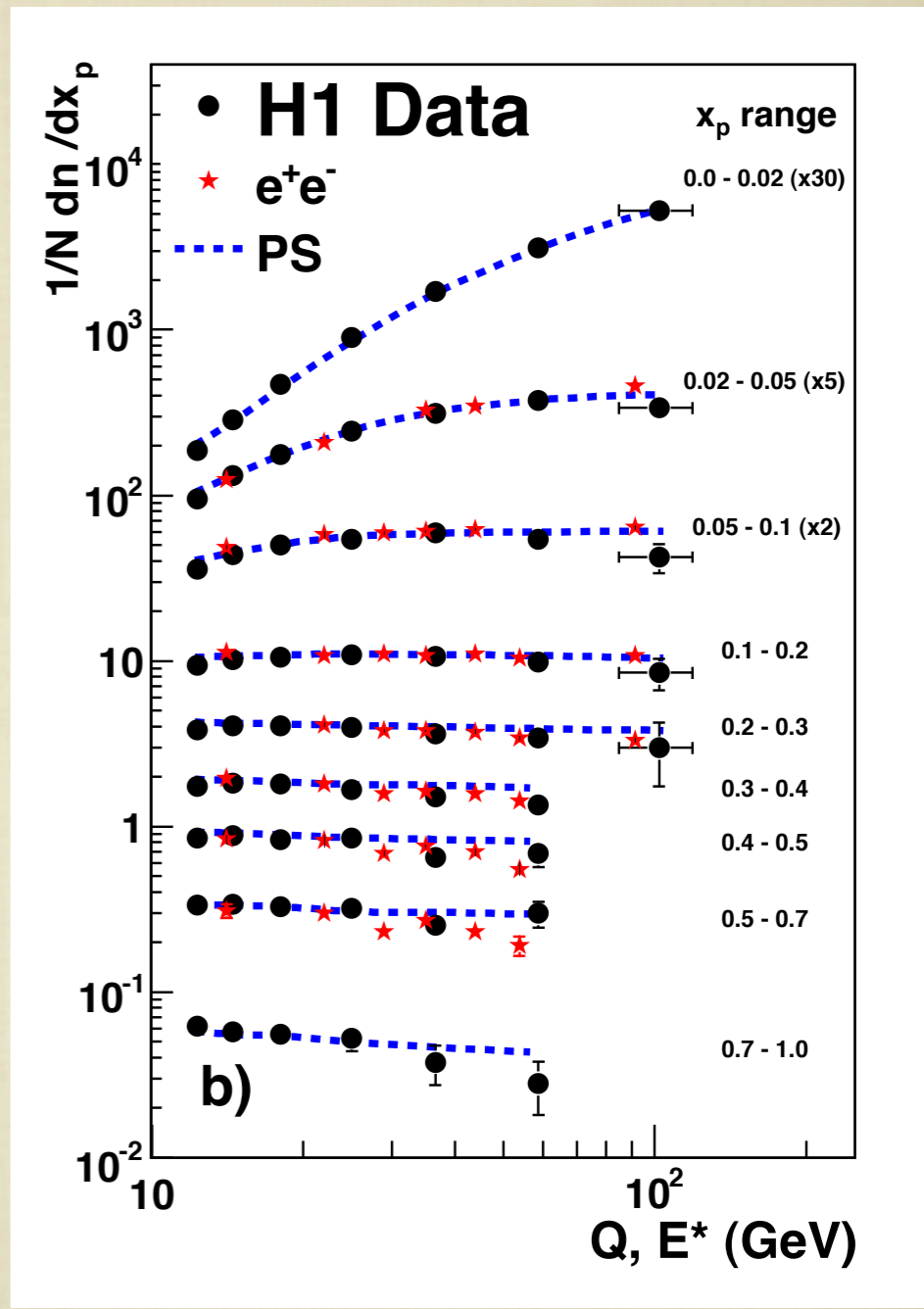
$h^{\pm}$  scaled mom. distr.

$$1/N \, dn/dx_p$$

- compare with  $e^+e^-$

$$\text{at } E^* \equiv Q = \sqrt{s}/2$$

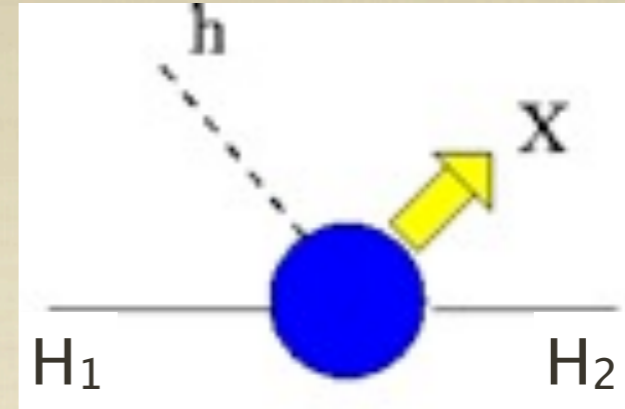
➤ universality test





$$p\bar{p} \rightarrow hX$$

$$h = \pi^{\pm,0}, K^{\pm}, K_s^0, p, \bar{p}, \Lambda, \bar{\Lambda}$$



## Energy range

- mid  $\eta$ ,  $1 \leq P_{\perp}(\pi^0) \leq 20$  GeV at RHIC (PHENIX)
  - large  $\eta > 0$ ,  $1 \leq P_{\perp}(\pi^0 - \pi^{\pm}, K^{\pm}) \leq 10$  at RHIC (STAR — BRAHMS)
  - mid  $\eta$ ,  $1 \leq P_{\perp}(K_s^0, p, \bar{p}, \Lambda, \bar{\Lambda}) \leq 10$  at RHIC (STAR)
  - $80 \leq M_{jj} \leq 600$ ,  $1 \leq P_{\perp}(h^{\pm}) \leq 20$  at CDF
- }  $\sqrt{s} = 200$   
pp  
p $\bar{p}$

– constrain  $D_g^h$ , especially at  $x_B \ll 1$

– probe FF at large  $z$  (complementary to  $e^+e^-$ )

–  $1/N \, dn/dx_p$  test universality with  $e^+e^-$  and SIDIS



# status of parametrizations

## before 2007

- **AKK** Albino, Kniehl, Kramer, 2005
- **BKK** Binnewies, Kniehl, Kramer, 1995
- **BFG** Bourhis, Fontannaz, Guillet, 1998
- **BFGW** Bourhis, Fontannaz, Guillet, Werlen, 2001
- **CGRW** Chiappetta, Greco, Guillet, Rolli, Werlen, 1994
- **GRV** Glück, Reya, Vogt, 1993
- **KKP** Kniehl, Kramer, Potter, 2000
- **Kr** Kretzer, 2000

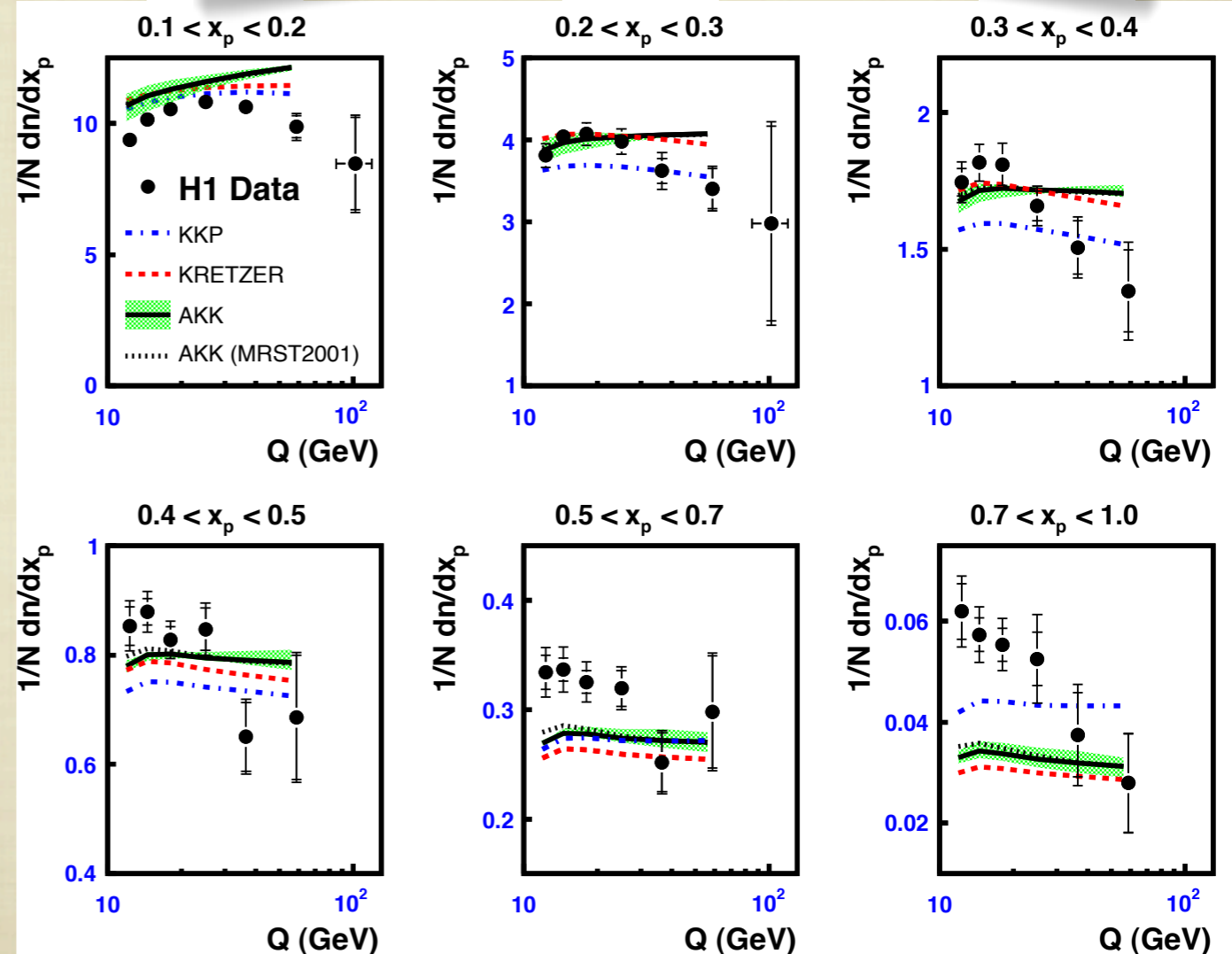


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- Kr Kretzer, 2000

fail to reproduce  
scaling violations  
of recent H1 data



H1 Coll., P.L. B654 (2007) 148



# status of parametrizations

after 2007

- AKK08 Albino, Kniehl, Kramer, 2008
- DSS De Florian, Sassot, Stratmann, 2007
- HKNS Hirai, Kumano, Nagai, Sudoh, 2007

## main ingredients

DSS	AKK08	HKNS
$e^+e^-$ SIDIS pp	$e^+e^-$ pp $p\bar{p}$	$e^+e^-$
$\pi^\pm, K^\pm, p, \bar{p}, h^\pm (\Lambda)$	$\pi^\pm, K^\pm, K_s^0, p, \bar{p}, \Lambda, \bar{\Lambda}$	$\pi^\pm, \pi^0, K^\pm, K^0+\bar{K}^0, n, p+\bar{p}$
$0.05-0.1 \leq z \leq 1 \leq Q^2 \leq 10^5 \text{ GeV}^2$	$0.05 \leq z \leq 1 \leq Q^2 \leq 4 \times 10^4 \text{ GeV}^2$	$0.01 \leq z \leq 1 \leq Q^2 \leq 10^8 \text{ GeV}^2$
NLO DGLAP in Mellin space $D(z, Q_0) = N z^a (1-z)^b [1 - c(1-z)^d]$ $N$ fixed by $\sum_h \int dz z D_i^h(z, Q^2) = 1$	NLO DGLAP in Mellin space + resum $\log^n(1-z)/1-z$ at NLL $D(z, Q_0)$ and $N$ fixed as DSS	NLO DGLAP direct integration $D(z, Q_0) = N z^a (1-z)^b$ $N$ fixed as DSS
SU(2) symmetric unfavoured $d+\bar{d} \propto u+\bar{u}$	SU(2) symmetric favoured ( $\pi$ ) and unfavoured build $D_i^{h++h-}, D_i^{h+-h-}$	SU(2) symmetric favoured and unfavoured $s =$ unfavoured
Lagrange multipliers	no error analysis	Hessian errors
	$m_h \neq 0$ effect $\rightarrow z \neq x_p$ resum log's at NLL also in $C_i$	

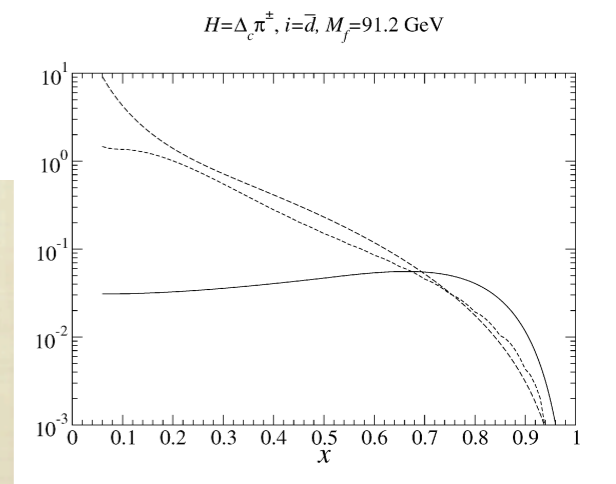
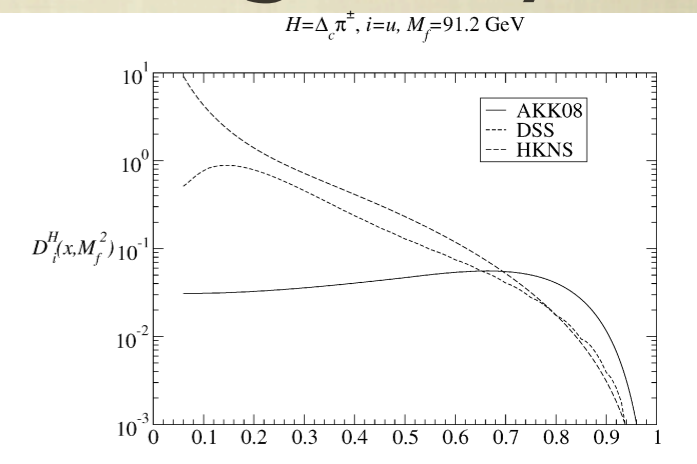
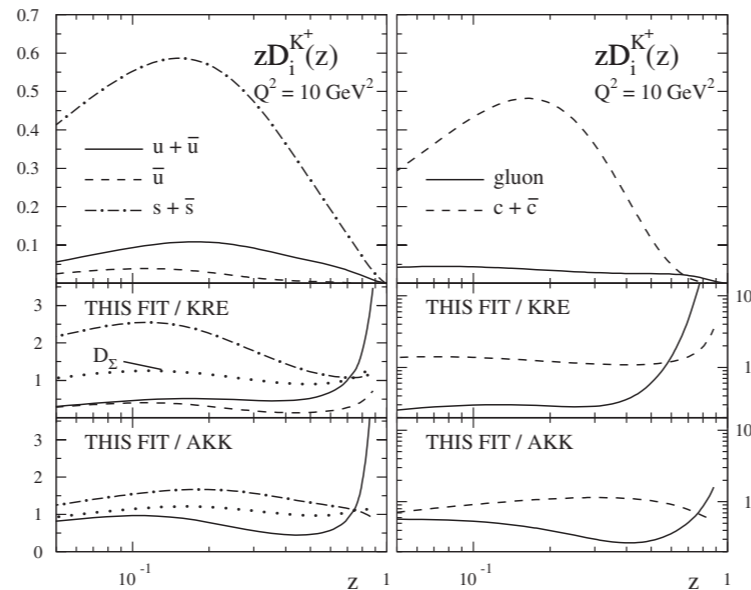
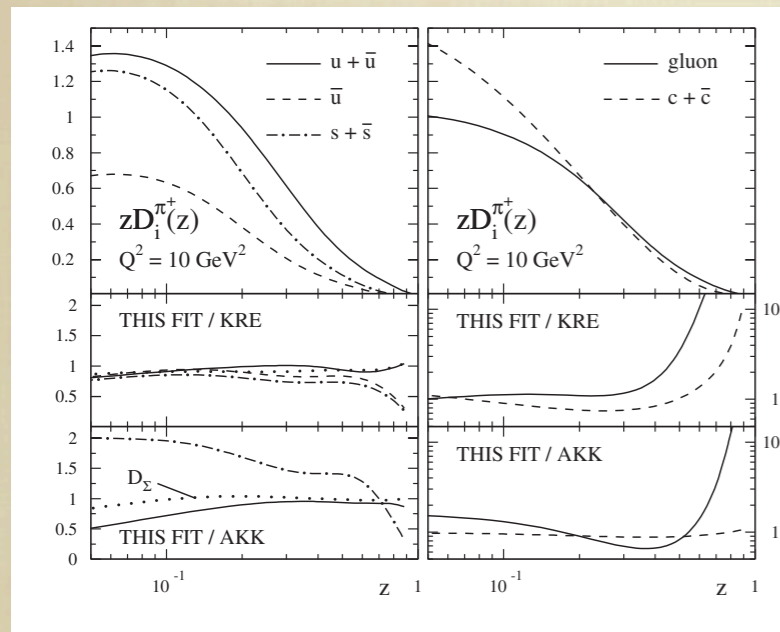
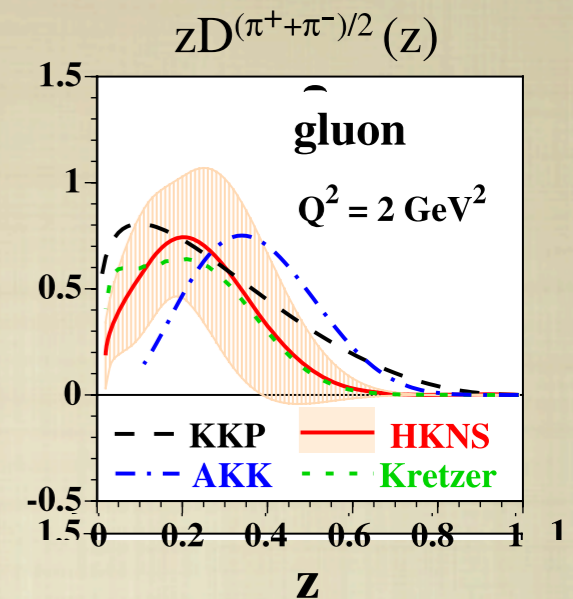


# main differences

- HKNS: no constrain on  $D_g^h$  from pp data, reliable at LHC ?

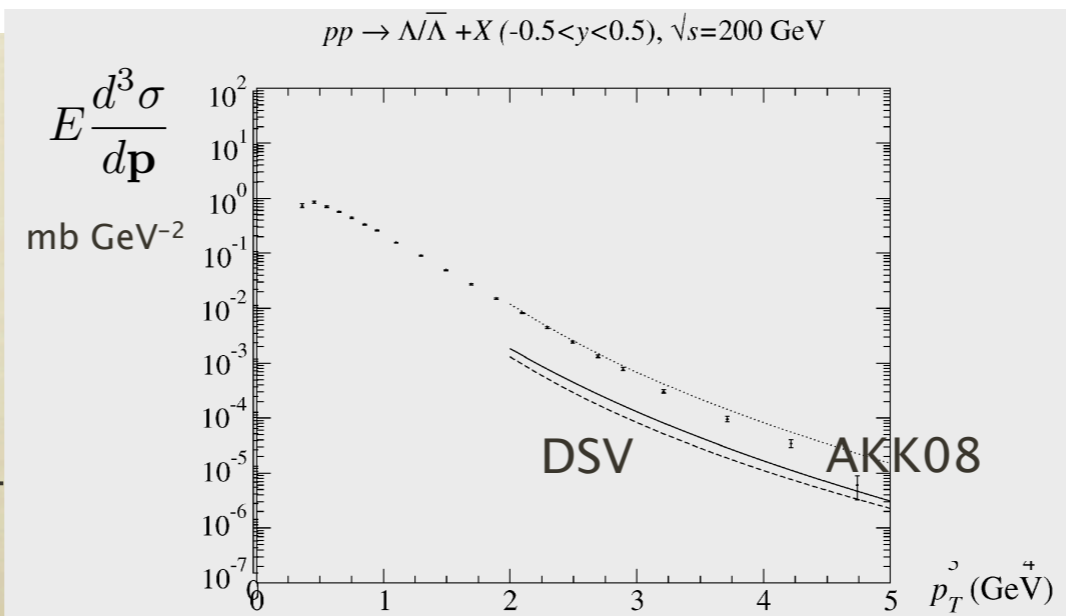
- AKK08-DSS discrepancies at large  $z$  and

charge asym.



- the puzzle of STAR  $\Lambda, \bar{\Lambda}$  data

STAR Coll., P.R.C 75 (07) 064901



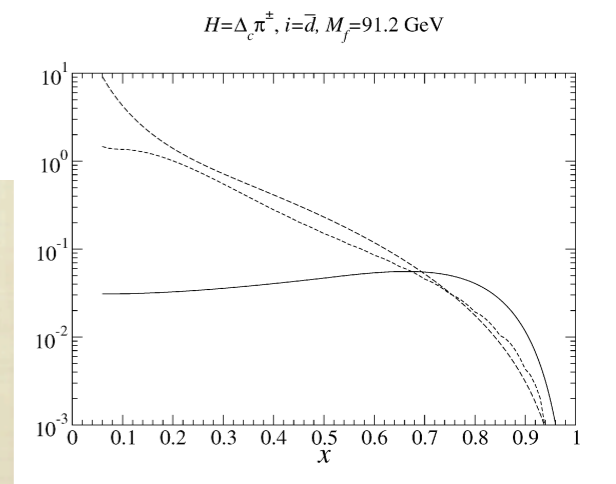
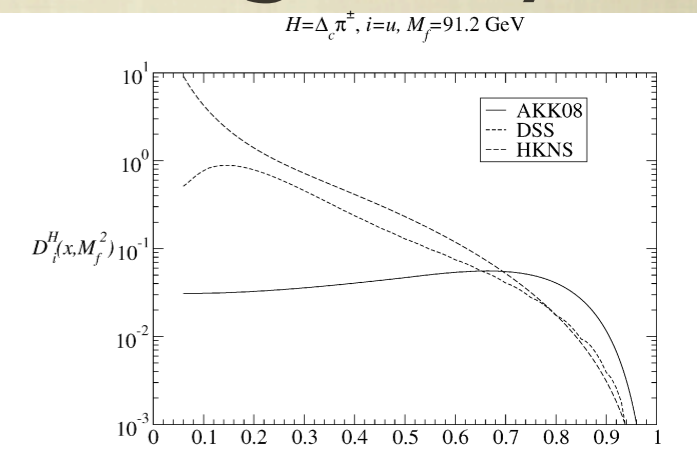
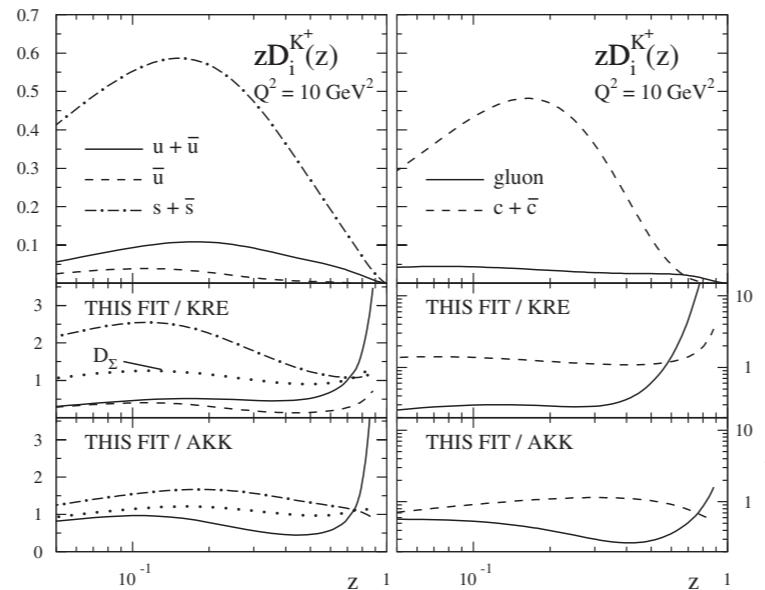
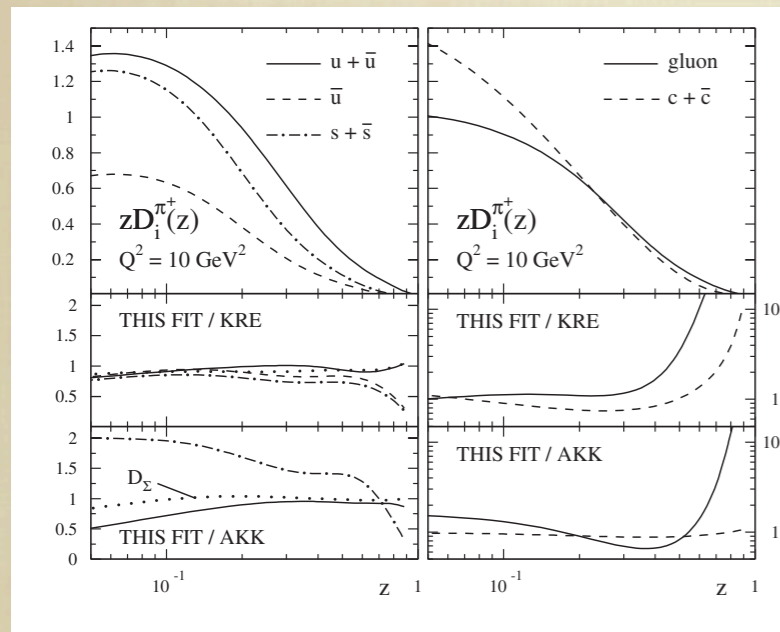
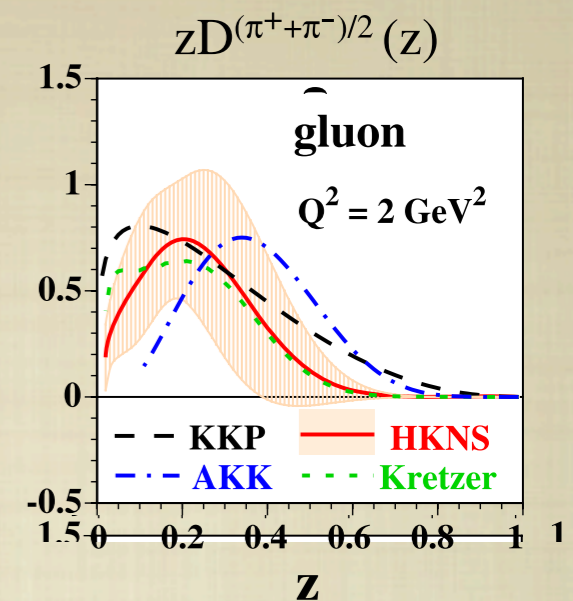


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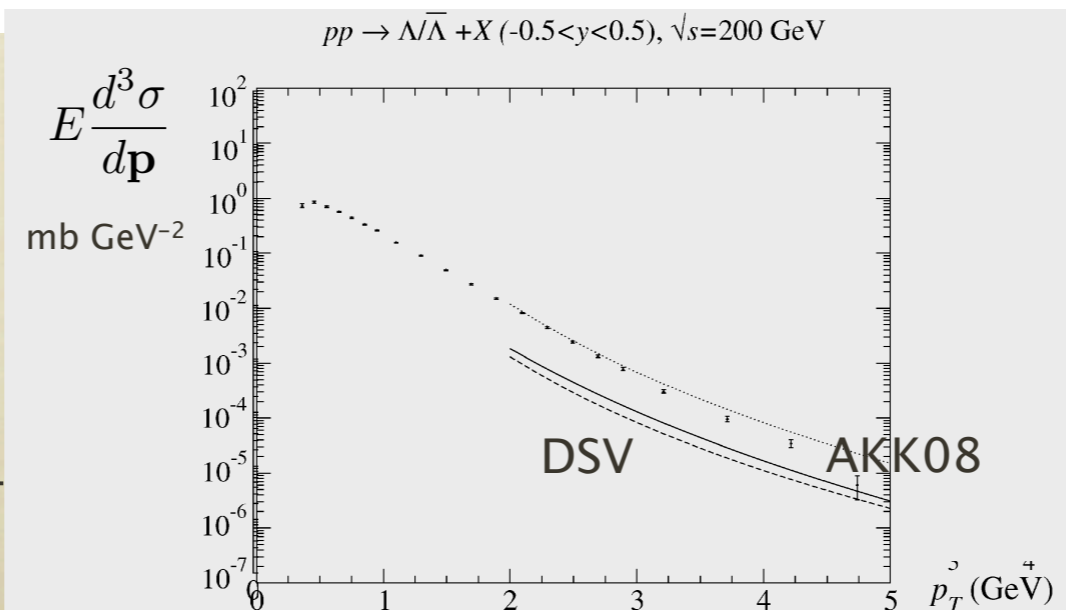
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STAR Coll., P.R.C 75 (07) 064901



# future of parametrizations

## – towards NNLO analysis

Almasy, Moch, Vogt, arXiv:1107.2263 [hep-ph]  
Albino et al., arXiv:1108.3948 [hep-ph]

$$\frac{dD_i^h(z, Q^2)}{d \ln Q^2} = \sum_{i=q, \bar{q}, g} P_{ji}(z, Q^2) \otimes D_j^h(z, Q^2) \quad \text{non-singlet } o(\alpha_s^3)$$



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$$\sigma^{K^\pm} - 2\sigma^{K_s^0} = [C_u - C_d] \otimes D_{u-d}^{K^\pm} \quad \text{at any order for SU(2) sym.}$$

NS  $K^\pm$  FF directly from data with NNLO  $C_i$

but data put not enough constrains yet

Albino, Christova, P.R.D81 (10) 094031



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NS  $K^\pm$  FF directly from data with NNLO  $C_i$

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Albino, Christova, P.R.D81 (10) 094031

– determine “non-perturbative” error from FF

➤ need a common interface like LHAPDF

at present only <http://www.pv.infn.it/~radici/FFdatabase>



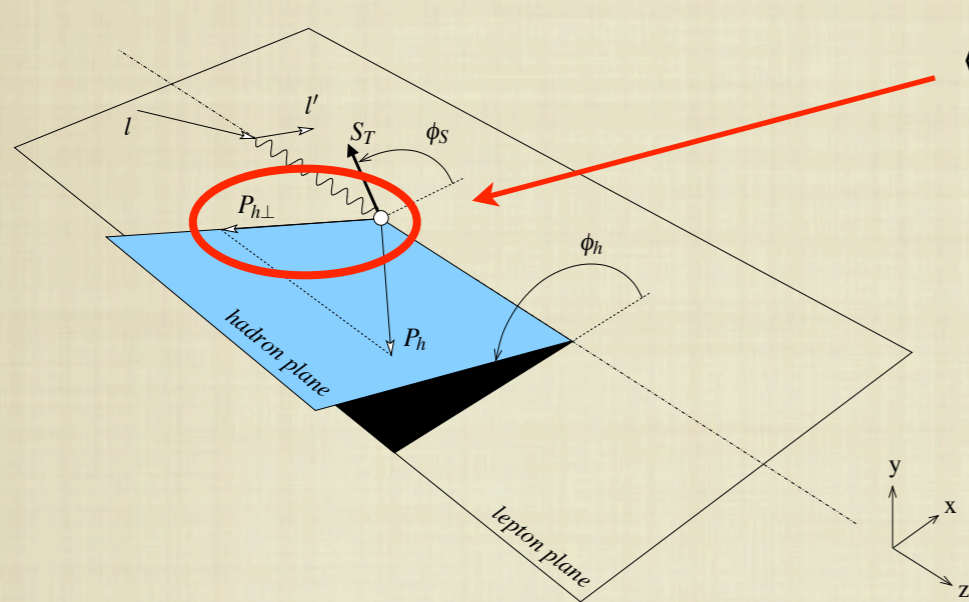
# what about 1h TMD FF ?

## Gaussian ansatz for SIDIS $d\sigma_{UU}$

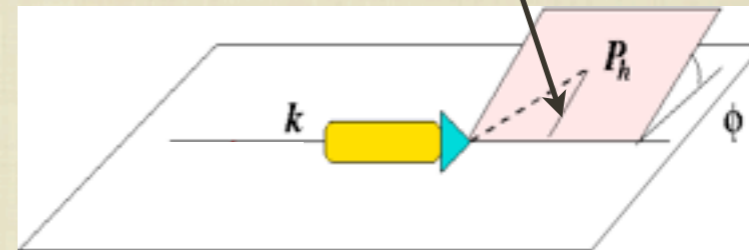
$$f_1^q(x, \mathbf{p}_T) = f_1^q(x) \frac{\exp[-\mathbf{p}_T^2 / \langle \mathbf{p}_T^2 \rangle]}{\pi \langle \mathbf{p}_T^2 \rangle}$$

$$D_1^q(z, \mathbf{K}_T) = D_1^q(z) \frac{\exp[-\mathbf{K}_T^2 / \langle \mathbf{K}_T^2 \rangle]}{\pi \langle \mathbf{K}_T^2 \rangle}$$

$$\longrightarrow \frac{d\sigma_{UU}(P_{h\perp})}{dz d\mathbf{P}_{h\perp}^2} = \frac{d\sigma_{UU}(0)}{dz d\mathbf{P}_{h\perp}^2} \exp[-\mathbf{P}_{h\perp}^2 / \langle \mathbf{P}_{h\perp}^2 \rangle]$$



$$\langle \mathbf{P}_{h\perp}^2 \rangle = z^2 \langle \mathbf{p}_T^2 \rangle + \langle \mathbf{K}_T^2 \rangle \quad \mathbf{K}_T = -z \mathbf{k}_T$$



$\langle \mathbf{p}_T^2 \rangle = 0.25$  ,  $\langle \mathbf{K}_T^2 \rangle = 0.20 \text{ GeV}^2$  by fitting Cahn effect in EMC data ('83)

(Anselmino et al., P.R.D71 (05) 074006)

$\langle \mathbf{p}_T^2 \rangle = 0.33$  ,  $\langle \mathbf{K}_T^2 \rangle = 0.16 \text{ GeV}^2$  by reproducing HERMES  $\langle P_{h\perp} \rangle$  data ('98-'00)

(Collins et al., P.R.D73 (06) 014021)

used in many phenomenological studies, but...



$$\langle \mathbf{p}_T^2 \rangle = 0.25, \quad \langle \mathbf{K}_T^2 \rangle = 0.20 \text{ GeV}^2$$

$A_{UU}^{\cos \phi}$  in EMC not only from Cahn effect

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HERMES data not corrected for acceptance effects



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Since 2007, new data (including  $\cos \varphi$  and  $\cos 2\varphi$ ) from JLab, HERMES, COMPASS

combined analysis of SIDIS and (old+new) DY data

(Schweitzer, Teckentrup, Metz, P.R.D81 (10) 094019)

➤ new parameters      $\langle \mathbf{p}_T^2 \rangle = 0.38 \pm 0.06$  ,  $\langle \mathbf{K}_T^2 \rangle = 0.16 \pm 0.01 \text{ GeV}^2$

➤ various tests of Gaussian ansatz

➤  $p_T$  and  $K_T$  broadening with  $s$  ➤  $\langle \mathbf{p}_T^2(s) \rangle = 0.3 + C_h s$

$$C_p = 7 \times 10^{-4}$$

$$C_\pi = 2.1 \times 10^{-3}$$



$\langle \mathbf{p}_T^2 \rangle = 0.25$  ,  $\langle \mathbf{K}_T^2 \rangle = 0.20 \text{ GeV}^2$       $A_{UU}^{\cos \phi}$  in EMC not only from Cahn effect

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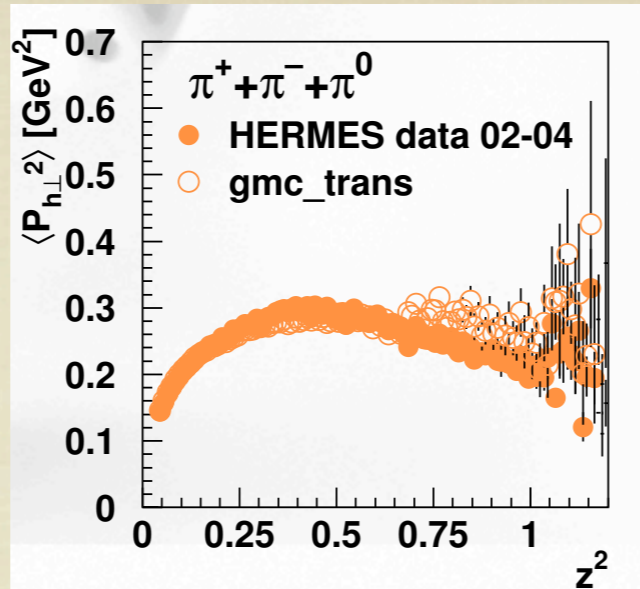
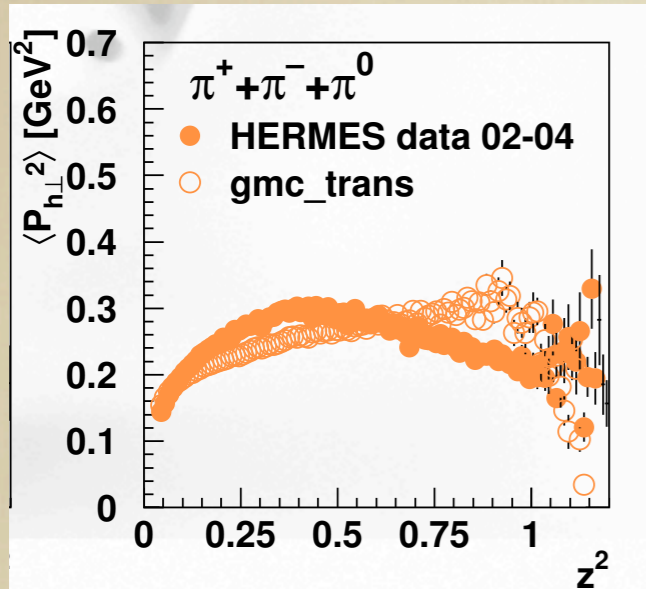
BUT...



# gmc\_trans MC (Schnell, ECT\* '07)

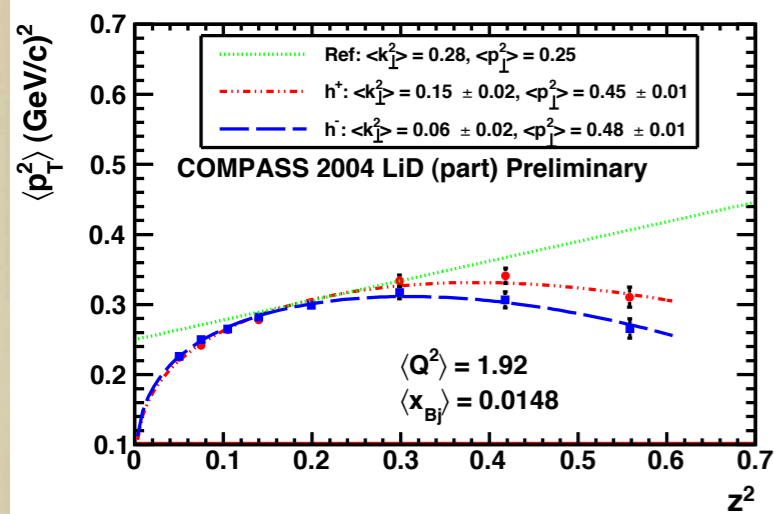
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$$\langle \mathbf{P}_{h\perp}^2 \rangle = z^2 \langle \mathbf{p}_T^2 \rangle + \langle \mathbf{K}_T^2(z) \rangle$$



$$\langle \mathbf{p}_T^2 \rangle = 0.14 \text{ GeV}^2$$

$$\langle \mathbf{K}_T^2 \rangle = 0.42 z^{0.54} (1 - z)^{0.37} \text{ GeV}^2$$



## similarly COMPASS (Rajotte, arXiv:1008.5125)

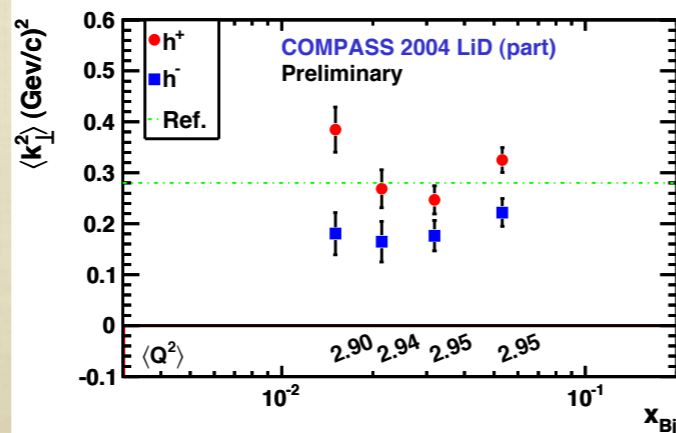
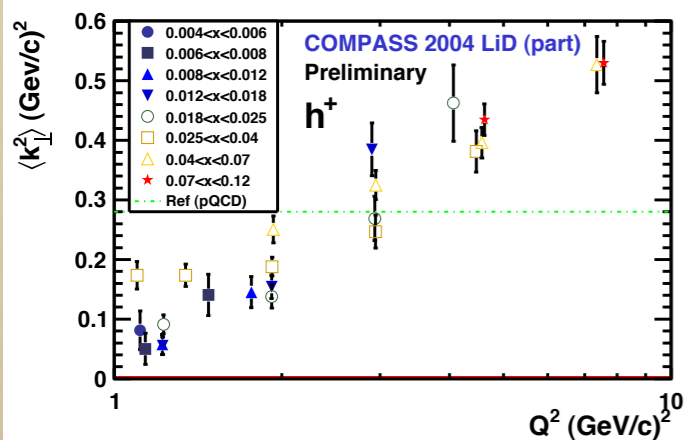
$$\langle \mathbf{p}_T^2 \rangle_{h^+} = 0.15 \text{ GeV}^2$$

$$\langle \mathbf{K}_T^2 \rangle_{h^+} = 0.45 z^{0.5} (1 - z)^{1.5} \text{ GeV}^2$$

$$\langle \mathbf{p}_T^2 \rangle_{h^-} = 0.06 \text{ GeV}^2$$

$$\langle \mathbf{K}_T^2 \rangle_{h^-} = 0.48 z^{0.5} (1 - z)^{1.5} \text{ GeV}^2$$

— fit with constant  
 $\langle \mathbf{p}_T^2 \rangle, \langle \mathbf{K}_T^2 \rangle$



also  
 $\langle \mathbf{p}_T^2(x, Q^2) \rangle_{h^+} \neq \langle \mathbf{p}_T^2(x, Q^2) \rangle_{h^-}$

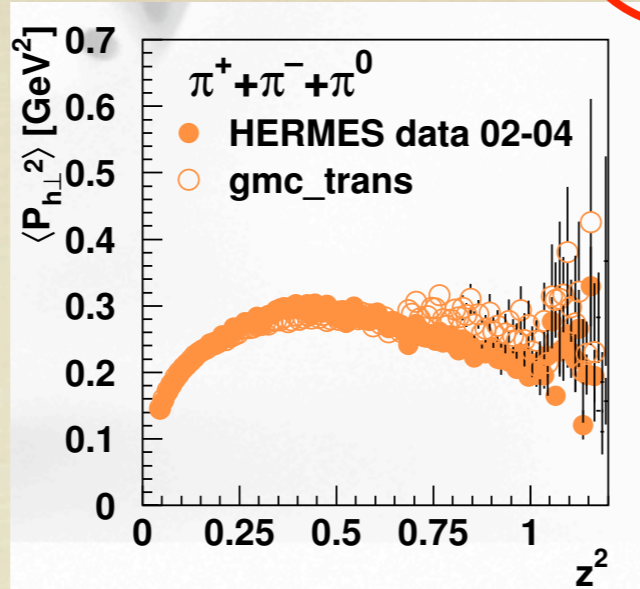
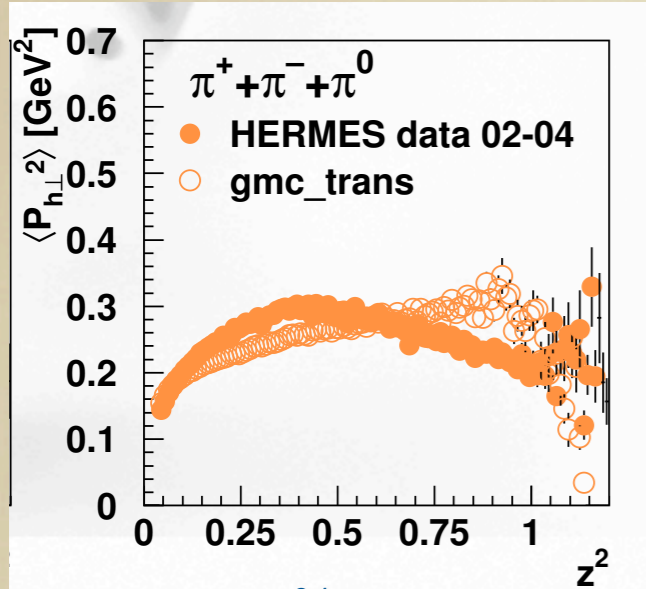
Moreover,...



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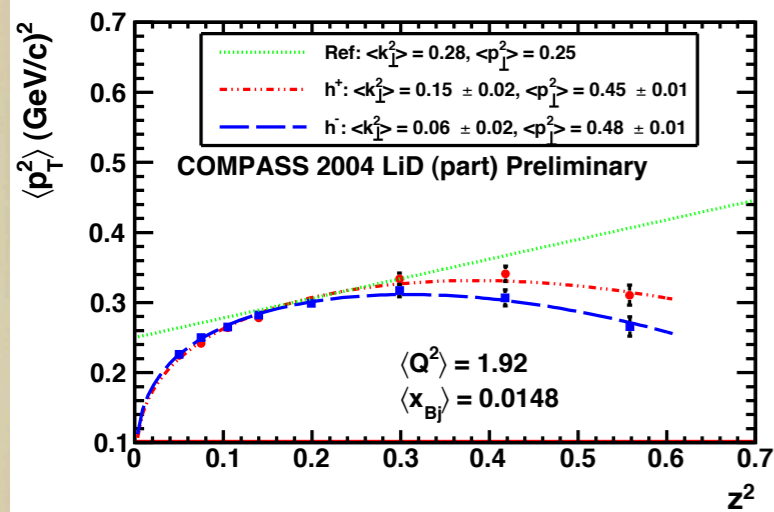
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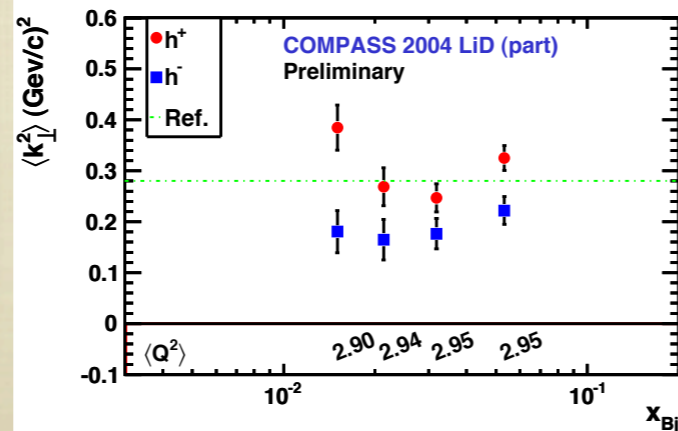
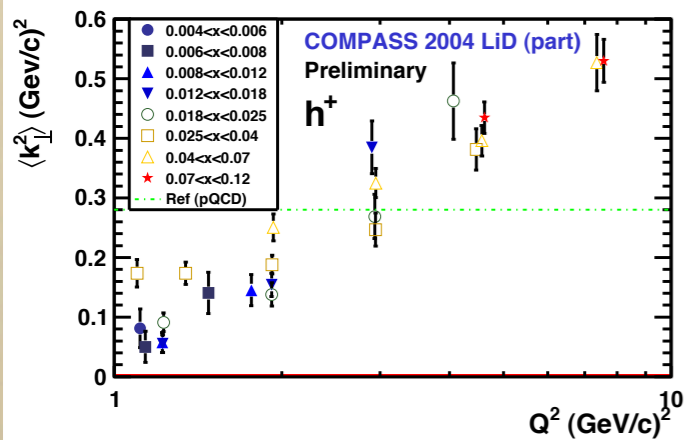
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$$\langle \mathbf{K}_T^2 \rangle_{h^-} = 0.48 z^{0.5} (1-z)^{1.5} \text{ GeV}^2$$

— fit with constant  
 $\langle p_T^2 \rangle, \langle K_T^2 \rangle$



also

$$\langle p_T^2(x, Q^2) \rangle_{h^+} \neq \langle p_T^2(x, Q^2) \rangle_{h^-}$$

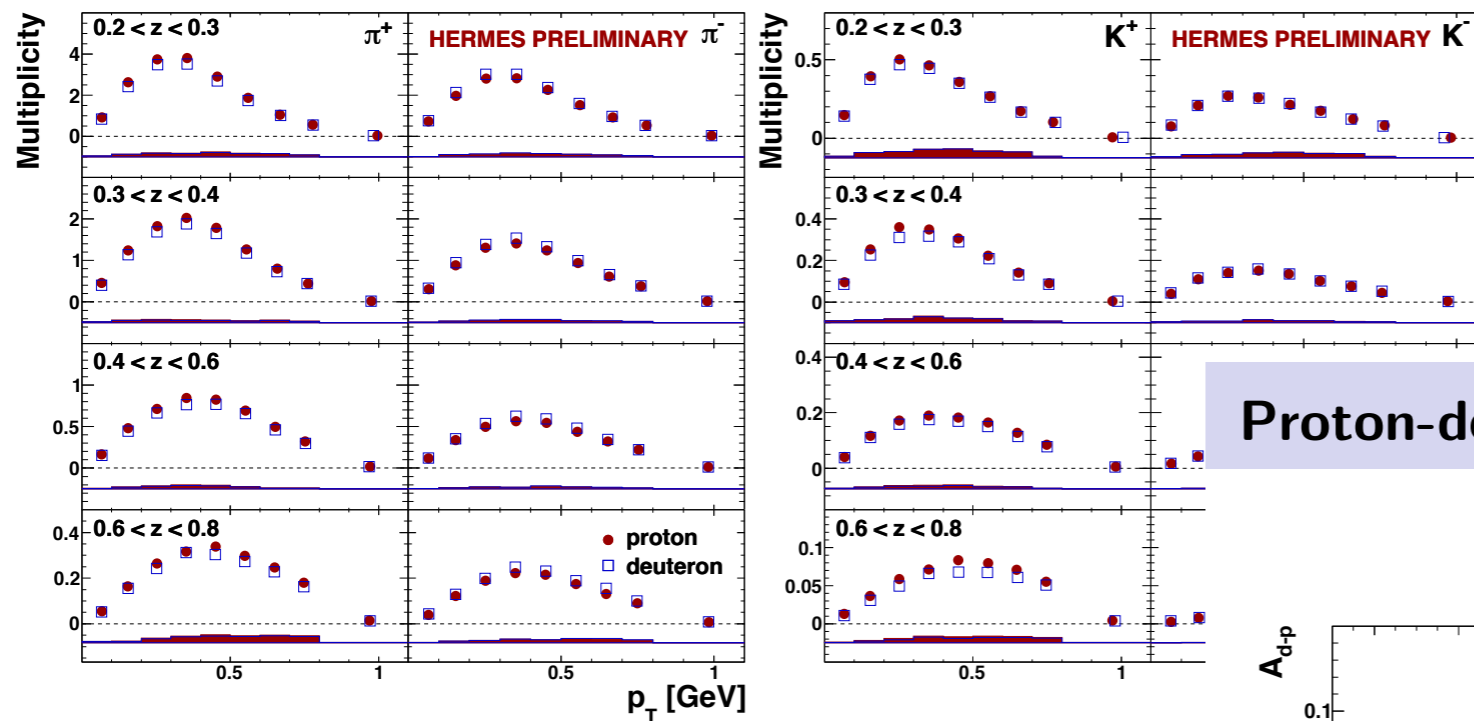
Moreover, ...



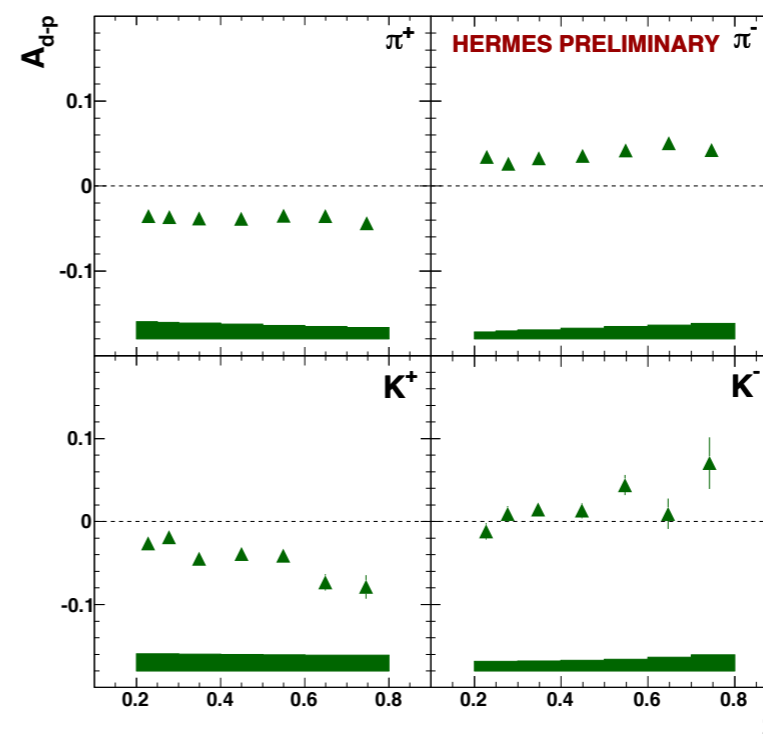
# ► HERMES multiplicity (Joosten, DIS 2011)

## Results: Projections vs $z p_T$

- Disentanglement of  $z$  and  $p_T$
- Access to the transverse intrinsic quark  $p_T$  and fragmentation  $k_T$ .



## Proton-deuteron multiplicity asymmetry



definition:

$$A_{d-p}^h \equiv \frac{\mathcal{M}_d^h - \mathcal{M}_p^h}{\mathcal{M}_d^h + \mathcal{M}_p^h}$$

- Reflects different valence quark content
- Improved precision by cancellations in the systematic uncertainty

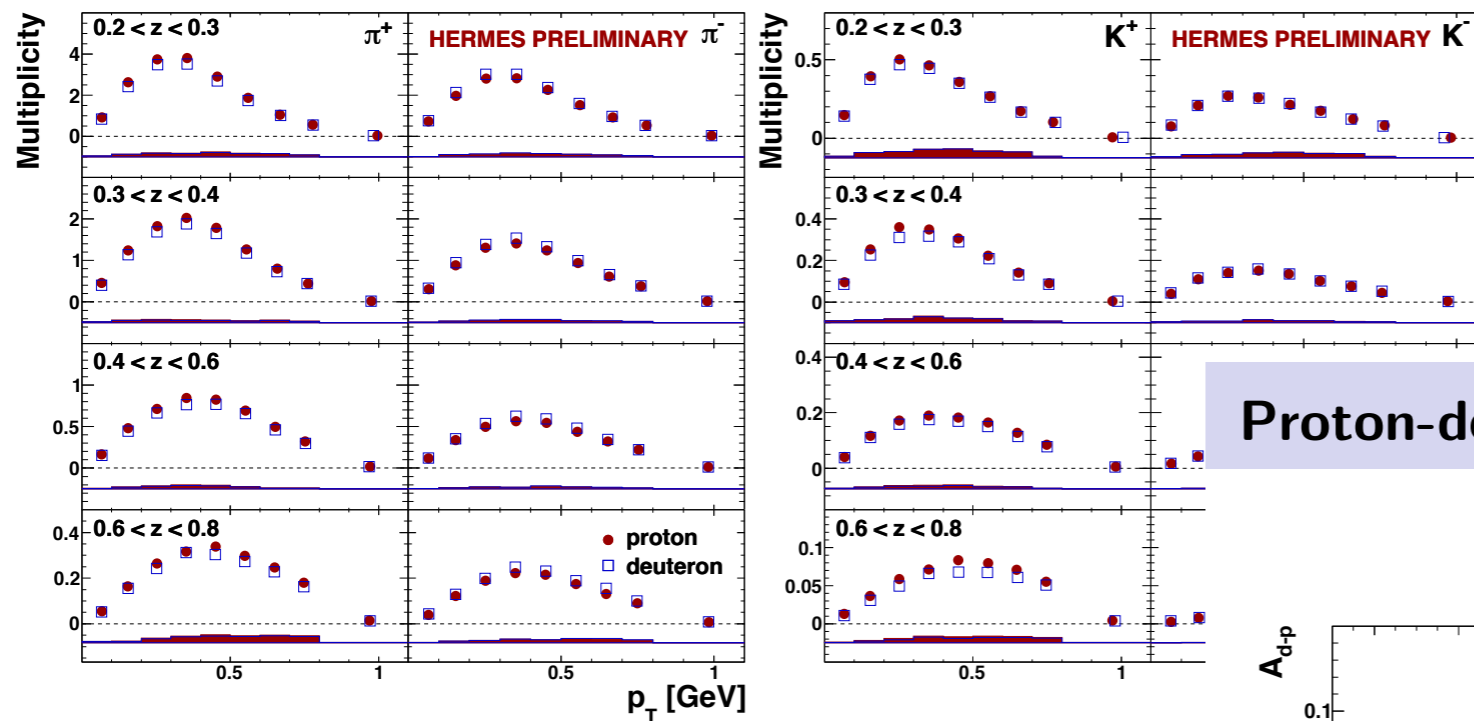
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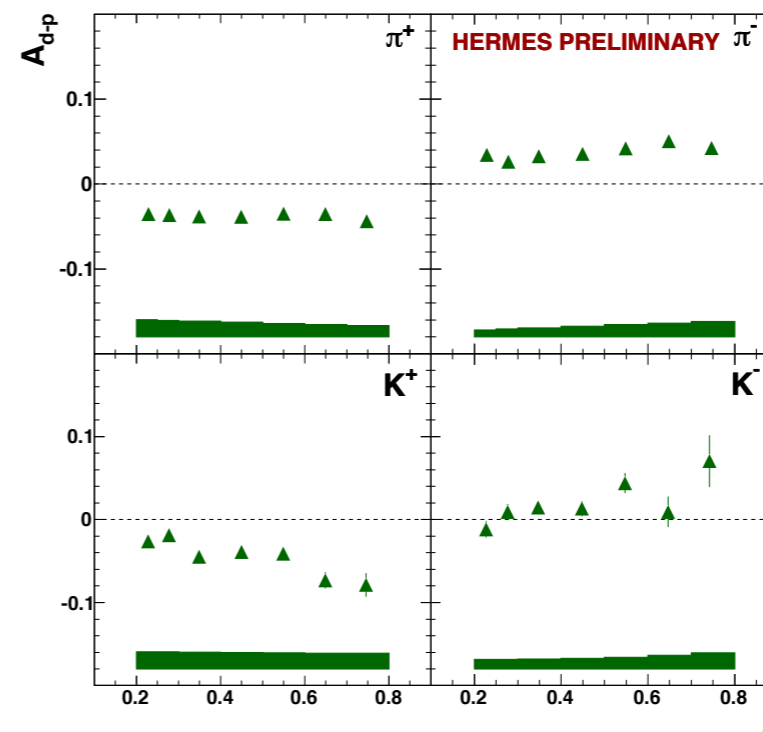
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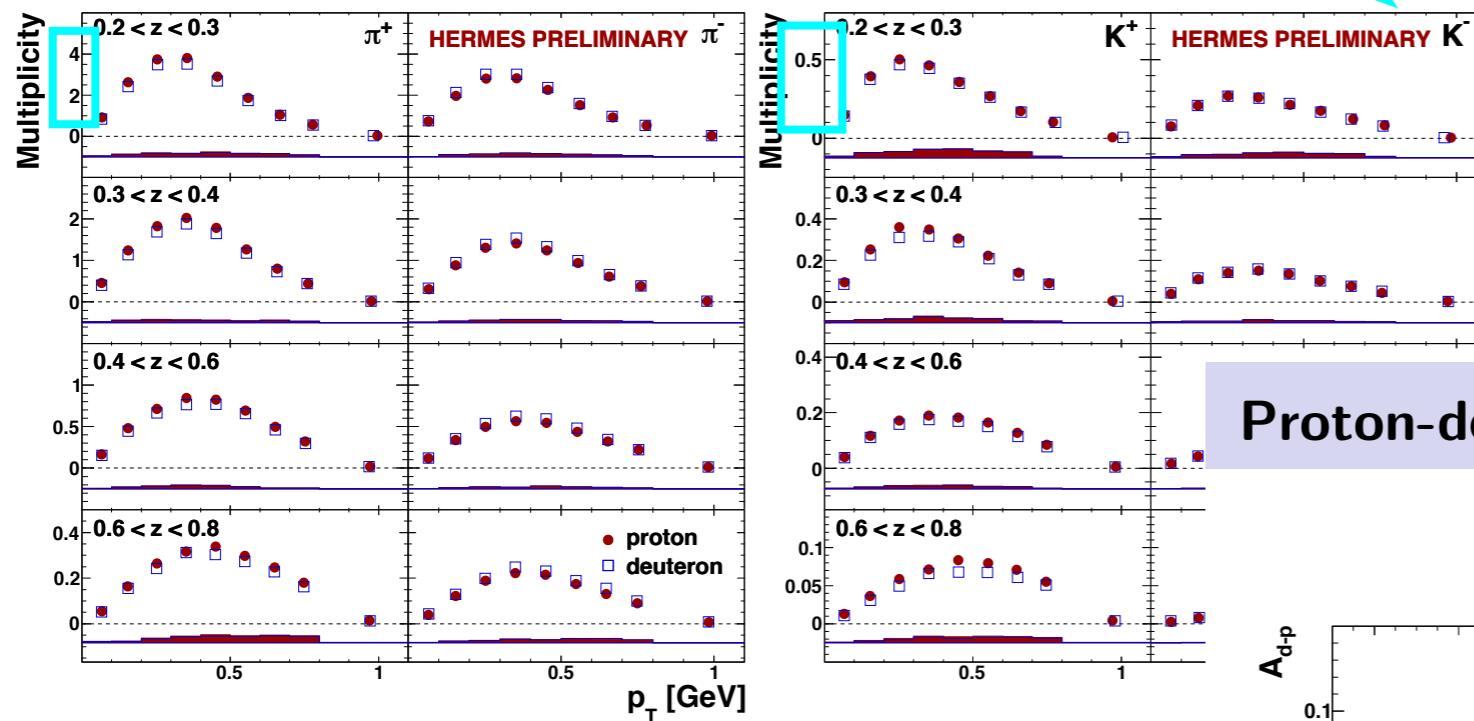


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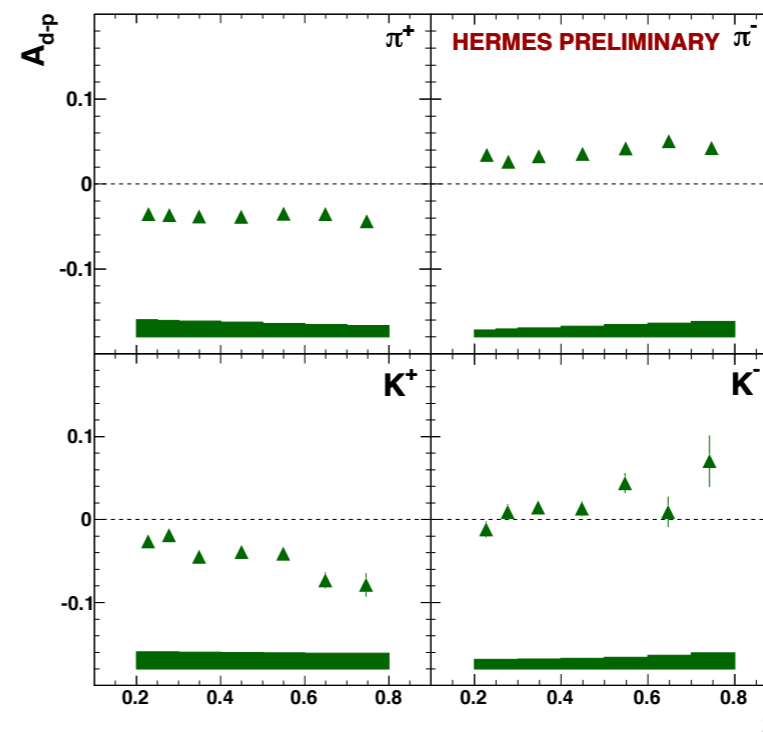
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- Access to the transverse intrinsic quark  $p_T$  and fragmentation  $k_T$ .

evidence for flavor dependence



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..., moreover, ...



# 1h TMD FF evolution

(Rogers & Aybat, P.R.D83 (11) 114042)

in config. space

$$D_i^h(z, \mathbf{b}_T; Q, \zeta) = A \times B \times C$$



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$$\sum_j C_{ij} \otimes D_j^h(z)$$

at  $b_T \ll 1/\Lambda_{\text{QCD}}$



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The diagram shows the factorization of the TMD FF into three parts: A, B, and C. Part A is represented by the sum  $\sum_j C_{ij} \otimes D_j^h(z)$  at the scale  $b_T \ll 1/\Lambda_{\text{QCD}}$ . Part B is shown to be related to an anomalous dimension via a double-headed arrow. Part C is the remaining factor in the product.







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$$D_i^h(z, \mathbf{b}_T; Q, \zeta) = A \times B \times C$$

$$\sum_j C_{ij} \otimes D_j^h(z) \quad \sim \exp[\text{anom. dim.}] \quad e^{-g(Q) \mathbf{b}_T^2}$$

at  $b_T \ll 1/\Lambda_{\text{QCD}}$       dim.]

nonperturb.  
g universal  
**scale dep.**  
**at low  $K_T$**

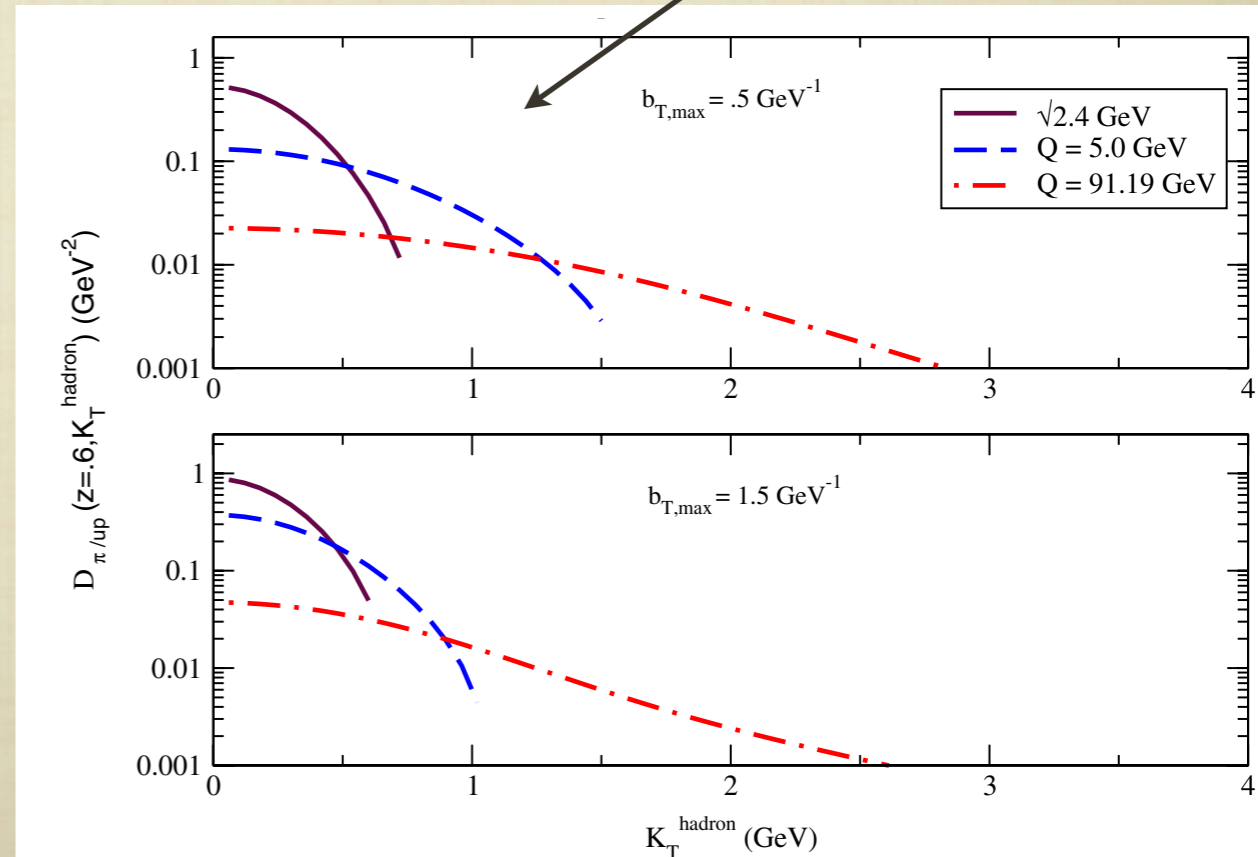
$$o(\alpha_s^0) \approx D_i^h(z) \exp[-g(Q) \mathbf{b}_T^2]$$

Gaussian ← TMD FF → BLNY fit ⇒ fix  $g(Q)$

$$Q_0^2 = 2.4$$

$$Q^2 = M_Z^2$$

$u \rightarrow \pi^+$





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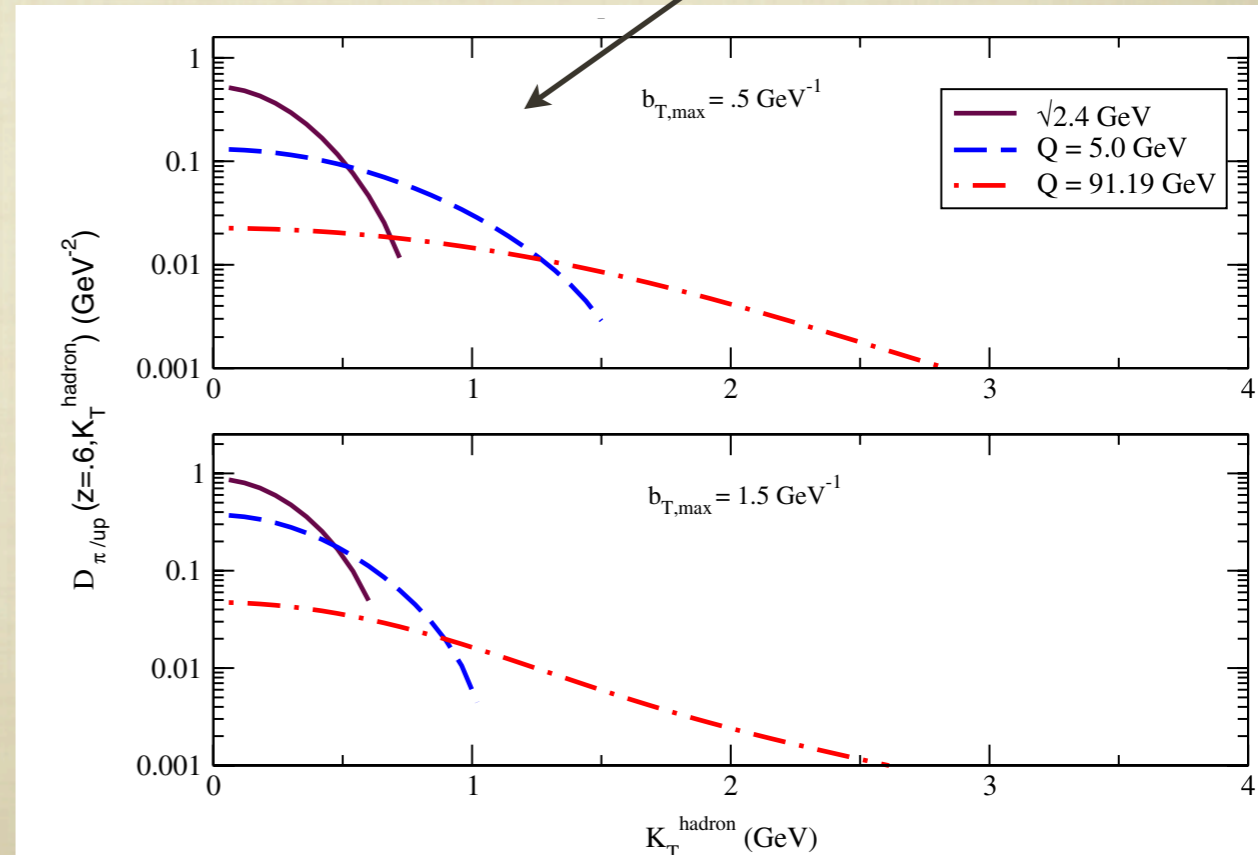
$Q^2 = M_Z^2$

**strong evolution effects**

$\langle k_T^2 \rangle^{1/2} (M_Z^2)$

$b_{T\text{max}}$ GeV <sup>-1</sup>	Gauss GeV	TMD FF GeV
0.5	1.74	2.15
1.5	1.06	1.85

$u \rightarrow \pi^+$





gaussian ansatz: too narrow point of view ?



It's amazing...



gaussian ansatz: too narrow point of view ?



It's amazing...



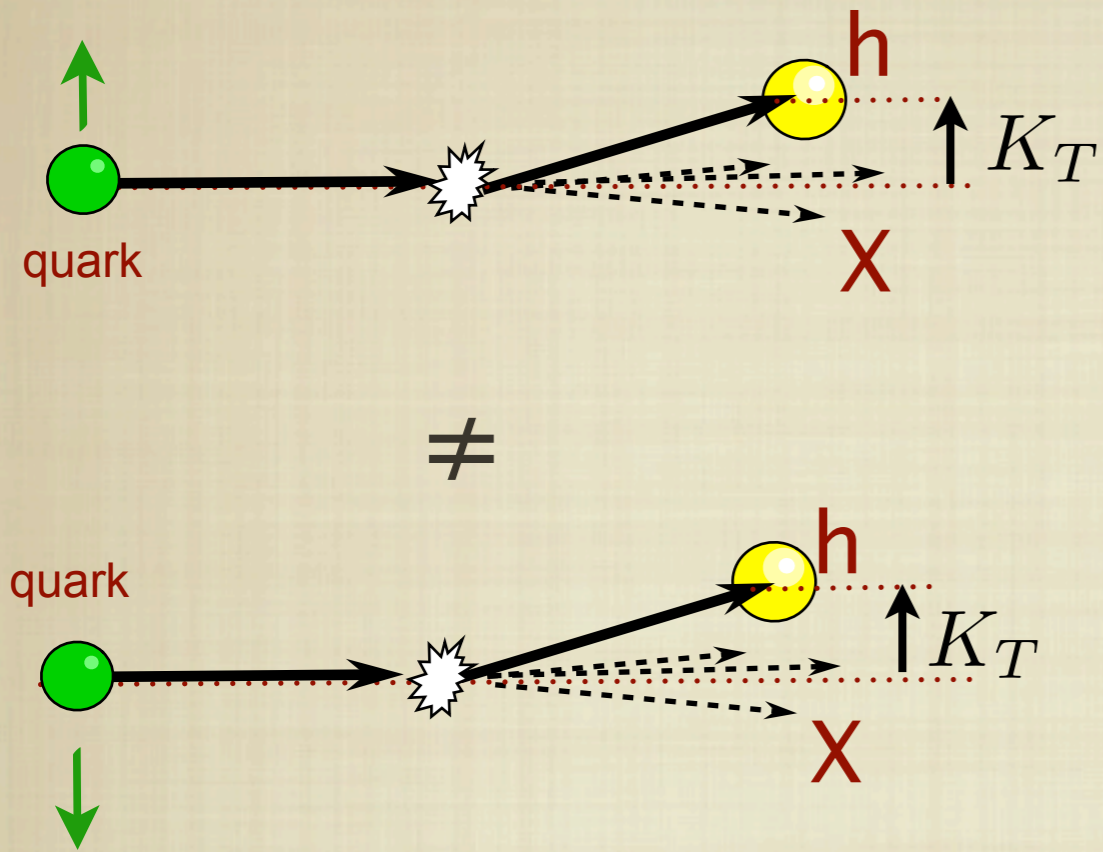
... how the earth  
looks so **FLAT** ...



polarized 1h TMD FF



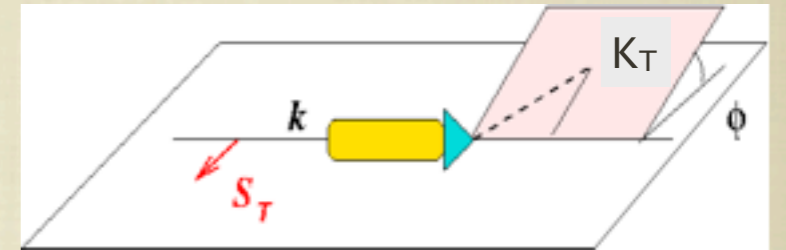
# the Collins function



# density of h produced by  $q^\uparrow$

$$D_{h/q^\uparrow}(z, \mathbf{K}_T) = D_1^q(z, \mathbf{K}_T^2) + H_1^{\perp q}(z, \mathbf{K}_T^2) \frac{|\mathbf{K}_T|}{zM_h} \mathbf{S}_\perp^q \cdot \underbrace{(\hat{\mathbf{k}}_q \times \hat{\mathbf{K}}_T)}_{\sim \sin \phi}$$

transv. motion  
of h  
=  
spin analyzer  
of fragm. q



positivity bound

$$|H_1^{\perp q}(z, \mathbf{K}_T^2)| \frac{|\mathbf{K}_T|}{zM_h} \leq D_1^q(z, \mathbf{K}_T^2)$$

Schäfer–Teryaev sum rule

$$\sum_{h, S_h} \int_0^1 dz z M_h H_1^{\perp (1)q}(z) = 0$$

Schäfer & Teryaev, P.R.D61 (00) 077903  
Meissner, Metz, Pitonyak, P.L.B690 (10) 296

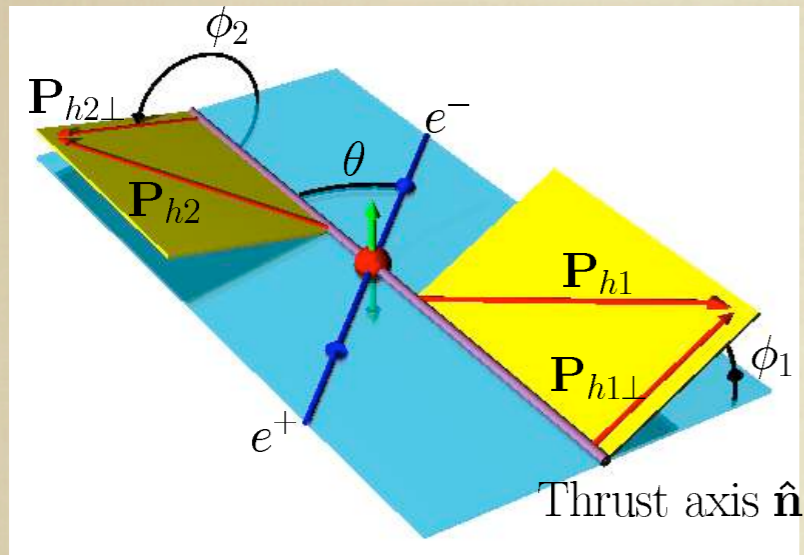
$$H_1^{\perp (n)q}(z) = \int d\mathbf{K}_T \frac{1}{2} \left( \frac{\mathbf{K}_T^2}{z^2 M_h^2} \right)^n H_1^{\perp q}(z, \mathbf{K}_T^2)$$



# extraction of Collins function

$$e^+e^- \rightarrow \pi^+\pi^-X$$

$$A^{\cos(\phi_1+\phi_2)}(\cos\theta, z, \bar{z}) = \frac{\sin^2\theta}{1+\cos^2\theta} \frac{\sum_q e_q^2 H_{1,q\rightarrow h_1}^{\perp(1/2)}(z) H_{1,\bar{q}\rightarrow h_2}^{\perp(1/2)}(\bar{z})}{\sum_q e_q^2 D_{1,q\rightarrow h_1}(z) D_{1,\bar{q}\rightarrow h_2}(\bar{z})}$$



old data:

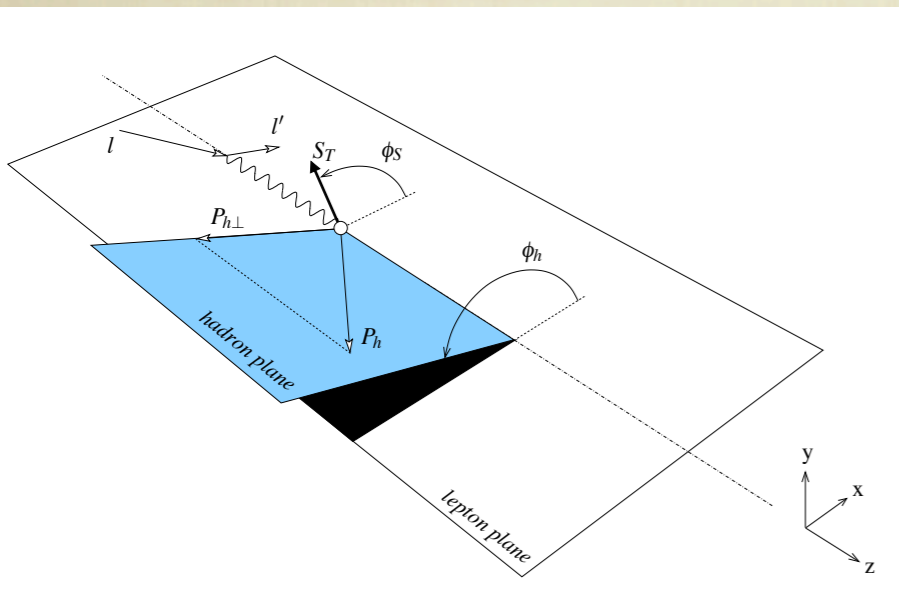
Abe et al. (Belle), P.R.L.96 (06) 232002

new data:

Seidl et al. (Belle), P.R.D78 (08) 032011

“thrust axis” method, or Collins–Soper frame  
also “ $\cos(2\phi_0)$ ” method, or Gottfried–Jackson frame

## in combination with SIDIS



$$A_{UT}^{\sin(\phi_h+\phi_S)} \propto - \frac{\sum_q e_q^2 \left[ h_1^q \otimes H_{1,q\rightarrow h}^{\perp} \right] (x, z, P_{h\perp}^2)}{\sum_q e_q^2 f_1^q(x) D_{1,q\rightarrow\pi}(z)}$$

old data:

Airapetian et al. (HERMES), P.R.L.94 (05) 012002

Ageev et al. (COMPASS), N.P.B765 (07) 31

new data:

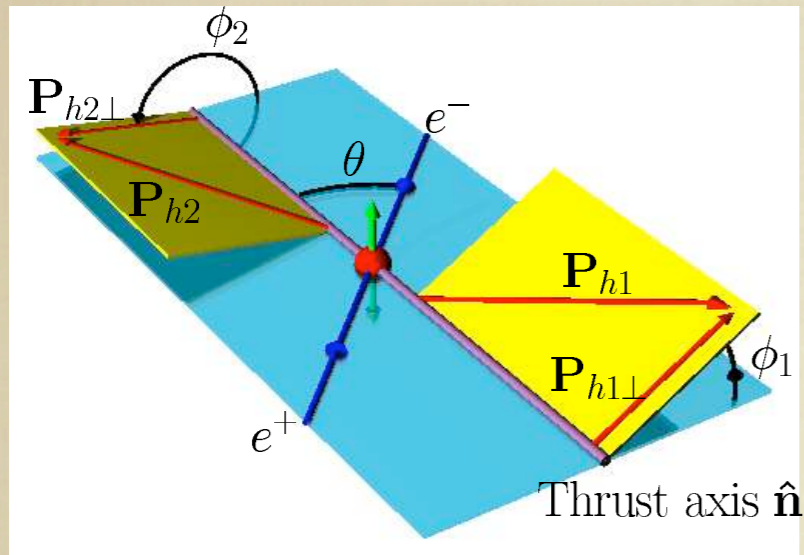
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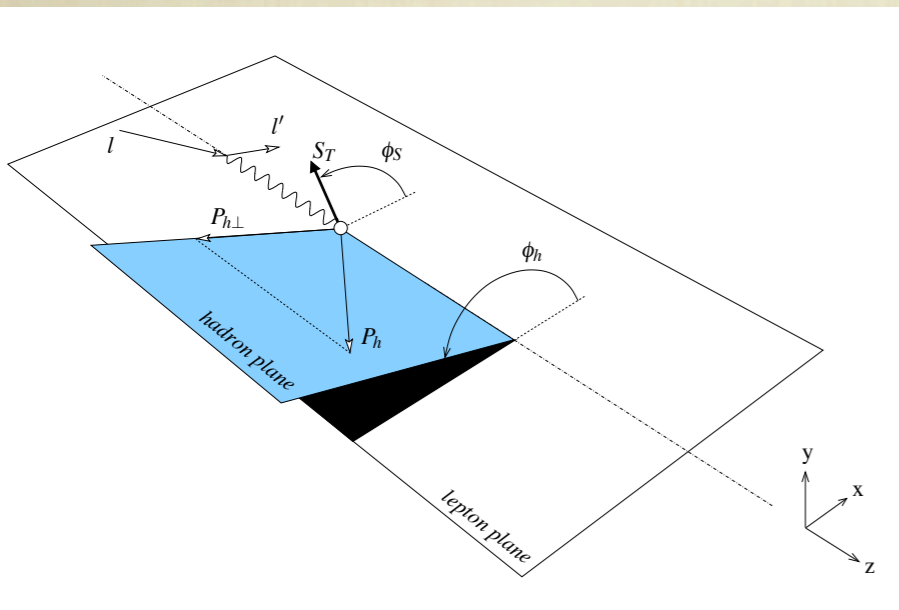
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need

- factoriz. th.
- universality
- evolution

“thrust axis” method, or Collins–Soper frame  
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# extraction of Collins function

$$\begin{array}{c}
 H_{1q \rightarrow h_1}^{\perp(1/2)} \quad H_{1\bar{q} \rightarrow h_2}^{\perp(1/2)} \\
 \updownarrow \\
 h_1^q \otimes H_{1q \rightarrow h}^{\perp}
 \end{array}$$

Gaussian ansatz:

$$D_1^q(z, \mathbf{K}_T) = D_1^q(z) \frac{\exp[-\mathbf{K}_T^2 / \langle \mathbf{K}_T^2 \rangle]}{\pi \langle \mathbf{K}_T^2 \rangle}$$

$$H_1^{\perp q}(z, \mathbf{K}_T) = F_{N^q, \gamma, \delta}(z) D_1^q(z) \frac{\exp[-\mathbf{K}_T^2 / \langle \mathbf{K}_T^2 \rangle_C]}{\pi \langle \mathbf{K}_T^2 \rangle_C}$$



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2 channels: favoured , unfavoured

5 params:  $N^{\text{fav}}$ ,  $N^{\text{unfav}}$ ,  $\gamma$ ,  $\delta$ ,  $M_H$



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- factoriz. th.
- universality
- evolution

↔ LO DGLAP for  $D_1(z)$  and  $H_1^{\perp(n)}(z) \sim D_1(z)$



# extraction of Collins function

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$$\updownarrow$$

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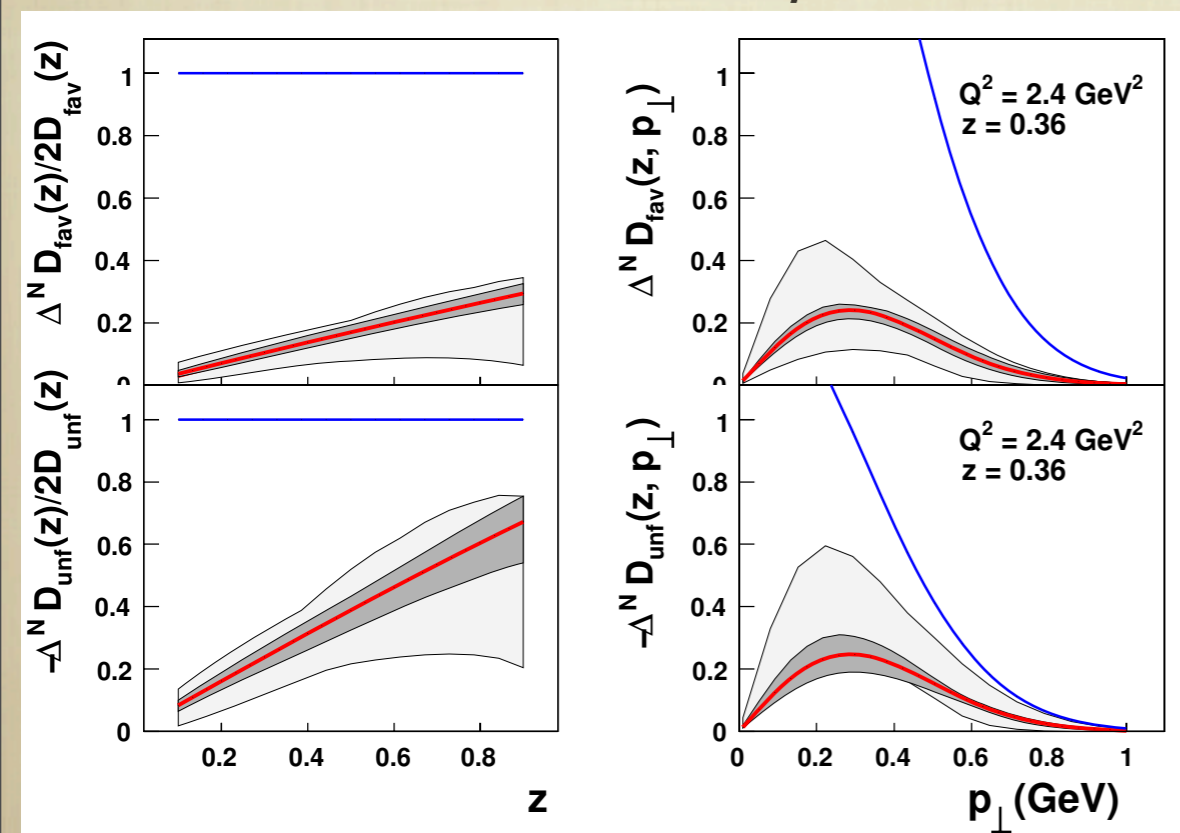
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↔ LO DGLAP for  $D_1(z)$  and  $H_1^{\perp(n)}(z) \sim D_1(z)$

2  $H_1^{\perp(1/2)} / D_1$



old data

Anselmino et al., P.R.D75 (07) 054032

error band  
 $\Delta\chi^2 \approx 17$

new data

Anselmino et al., N.P.B191(Pr.Sup.) (09) 98

positivity bound

see also Vogelsang & Yuan, P.R.D72 (05) 054028

Efremov, Goeke, Schweitzer, P.R.D73 (06) 094025



# extraction of Collins function

$$H_{1q \rightarrow h_1}^{\perp(1/2)} H_{1\bar{q} \rightarrow h_2}^{\perp(1/2)} \longleftrightarrow h_1^q \otimes H_{1q \rightarrow h}^{\perp}$$

Gaussian ansatz:

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- universality
- evolution

$\longleftrightarrow$  LO DGLAP for  $D_1(z)$  and  $H_1^{\perp(n)}(z) \sim D_1(z)$

2  $H_1^{\perp(1/2)} / D_1$

► unfav.  $\approx$  - fav.

old data

Anselmino et al., P.R.D75 (07) 054032

error band

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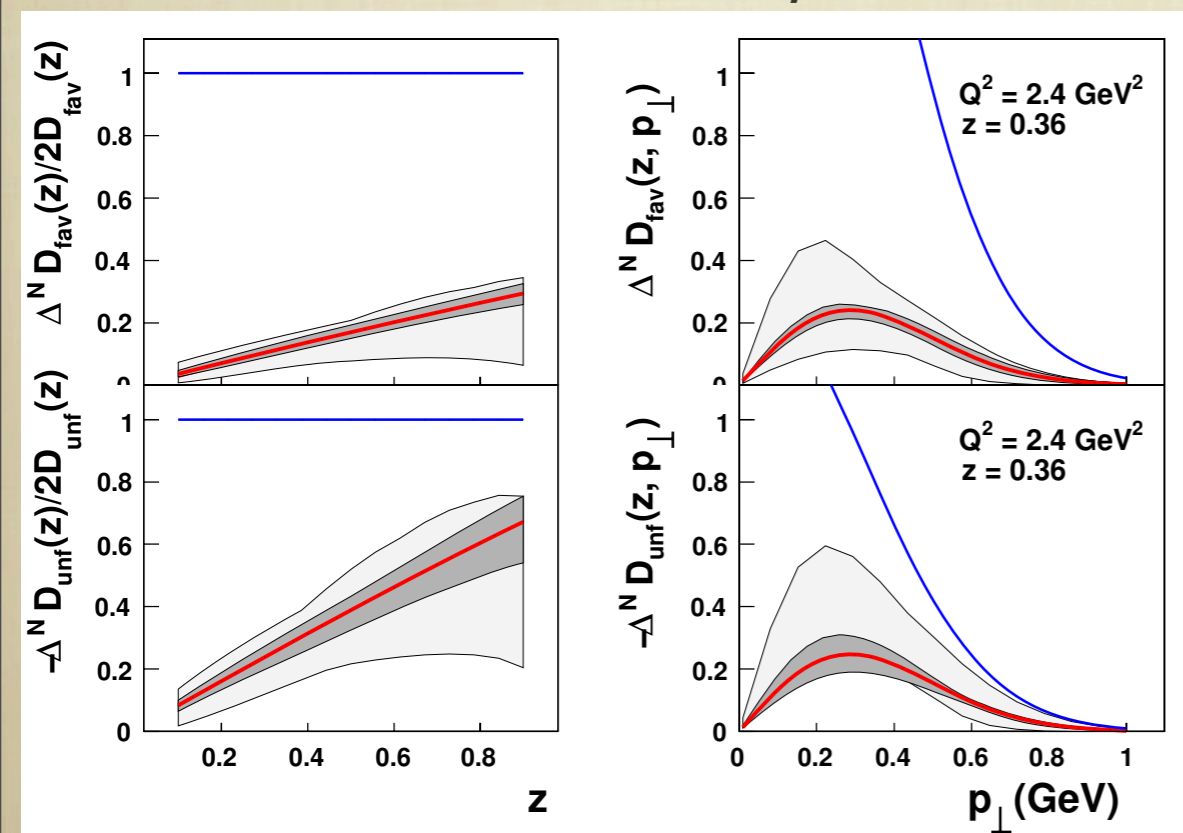
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positivity bound

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Efremov, Goeke, Schweitzer, P.R.D73 (06) 094025





# But...

- access only to  $H_{1\perp}^{\perp(n)}(z) \Rightarrow K_T$  dep. unconstrained  
 $\langle K_T^2 \rangle_C \neq \langle K_T^2 \rangle$  but flavor- /  $z$ - /  $Q^2$ -independent
- SIDIS kin.:  $x \leq 0.3$ ,  $0.2 \leq z \leq 0.7$ ,  $Q^2 = 2.5$  (need EIC)
- only fav./unfav. flavors (u & d)
- LO DGLAP evolution of  $H_{1\perp}^{\perp(n)}(z) \sim D_1(z)$   
but the chiral-odd kernel of  $H_{1\perp}^{\perp(1)} \sim h_1$  (Kang, P.R.D83 (11) 036006)
- full TMD evolution missing [ $Q_{\text{Belle}}^2 \sim 100 \leftrightarrow Q_{\text{SIDIS}}^2 \sim 2.5$ ]  
 $H_{1\perp}^{\perp(1)}$  kernel: diag. piece ( $\sim h_1$ ) + off-diag. piece (small ?)  
 $D_{1T\perp}^{\perp(1)}$  kernel: diag. piece ( $\sim D_1$ ) + off-diag. piece

(Kang, P.R.D83 (11) 036006; see also Meissner, Metz, P.R.L.102 (09) 172003; Yuan, Zhou, P.R.L.103 (09) 052001)  
Boer et al., P.R.L.105 (10) 202001; Gamberg et al., P.R.D83 (11) 071503(R)



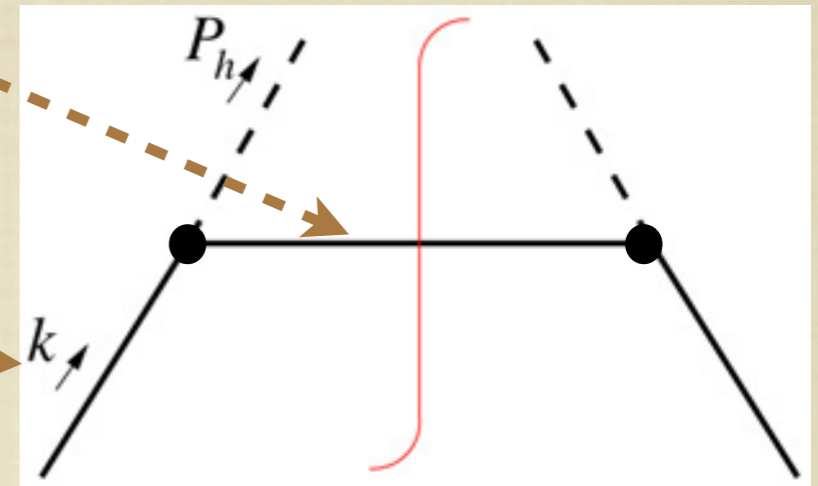
models of 1h (TMD) FF



# 1<sup>st</sup> category: the spectator approximation

on-shell spectator

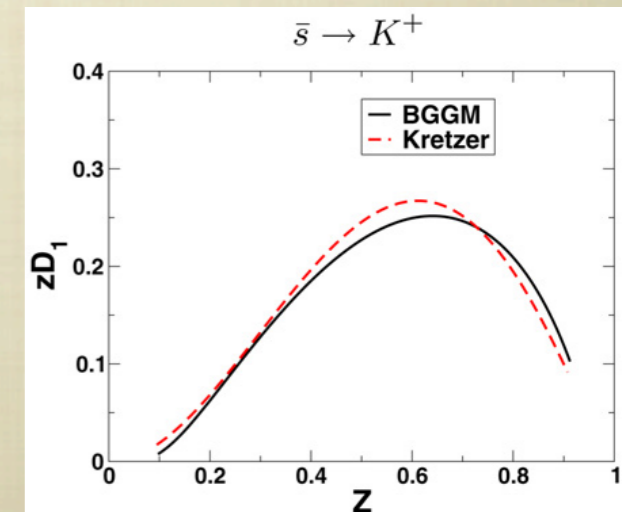
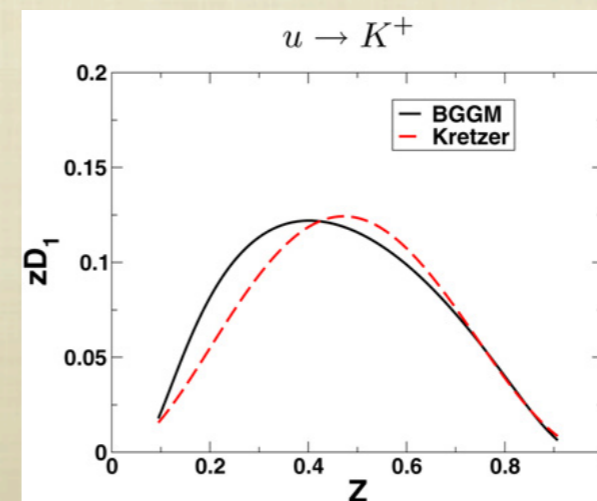
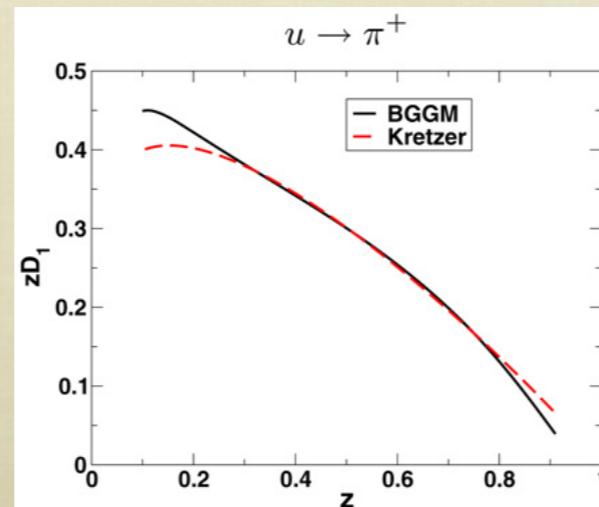
- $\delta$  funct.  $\Rightarrow$   $q$ - $q$  correlator analytic
- off-shell  $k^2(z)$  analytic
- ✗ only favoured channel



- $q\pi$  vertex: PS  $g_{\pi q} \gamma_5 \tau_i$  (Jakob et al., N.P.A626 (97); Bacchetta et al., P.L.B506 (01), B659 (08); Gamberg et al., P.R.D68 (03); Amrath et al., P.R.D71 (05) )
- PV  $g_{\pi q} \gamma_5 \not{P}_h$  (Bacchetta et al., P.R.D65 (02), P.L.B574 (03); Amrath et al., P.R.D71 (05) )

$$g_{\pi q}(z, k^2) \sim \exp[-k^2/\Lambda^2(z)] \quad (\text{Gamberg et al., P.R.D68 (03); Bacchetta et al., P.L.B659 (08) )$$

fit  $D_1^q$  to Kretzer  
@  $Q^2 = 0.4 \text{ GeV}^2$

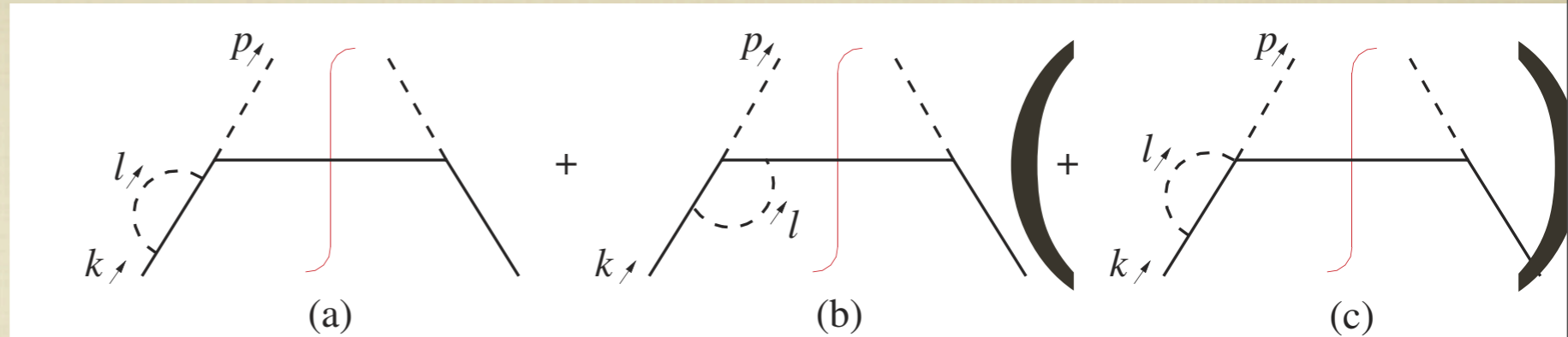




# the spectator approximation: the Collins funct.

interference from:

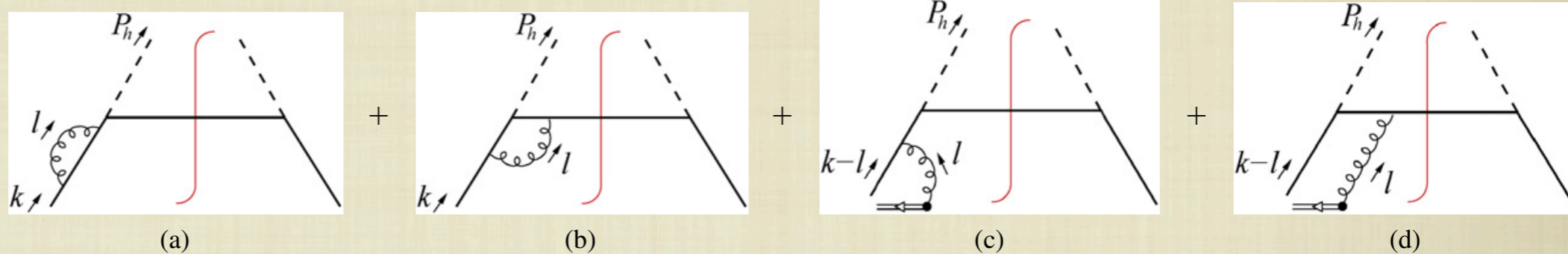
$\pi$  loops



needed in PV coupl.

and / or

g loops

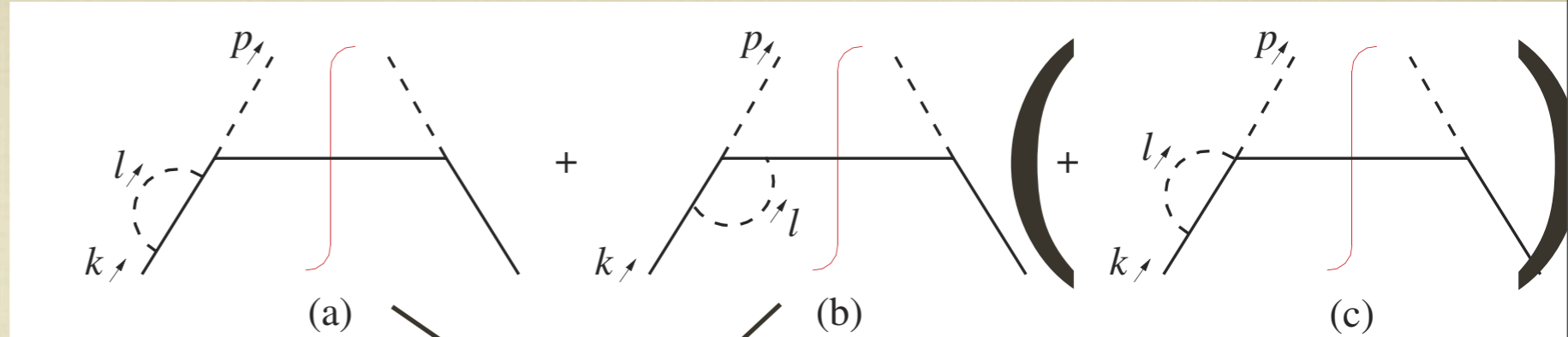




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interference from:

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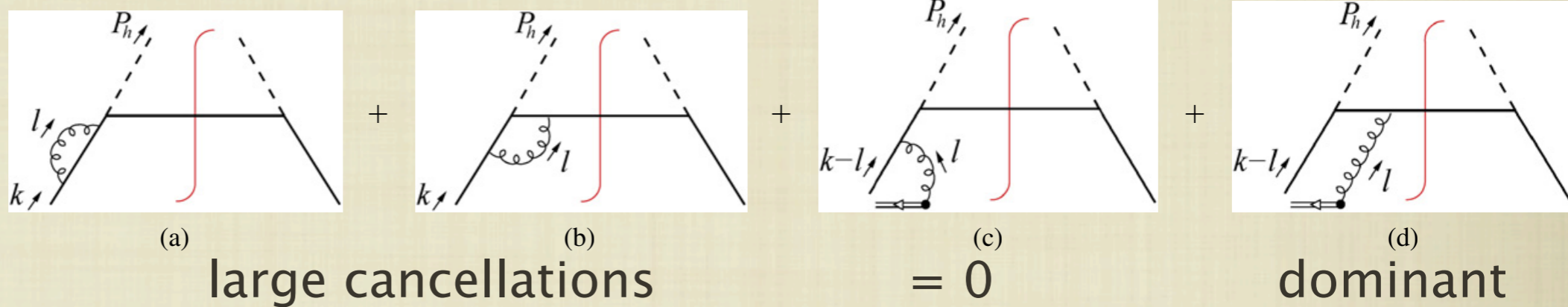
and / or

large cancellations

needed in PV coupl.

$\Rightarrow$  c) dominant

$g$  loops

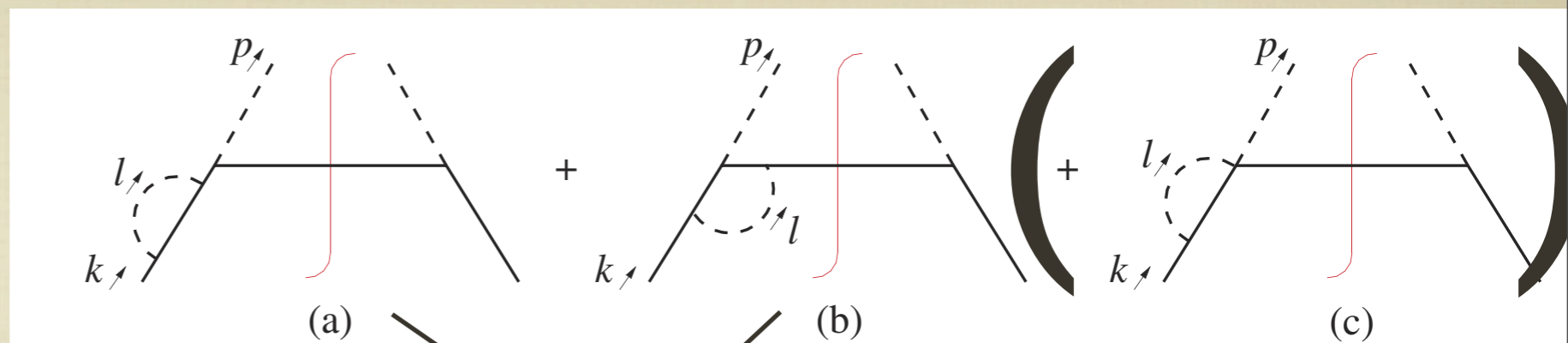




# the spectator approximation: the Collins funct.

interference from:

**$\pi$  loops**

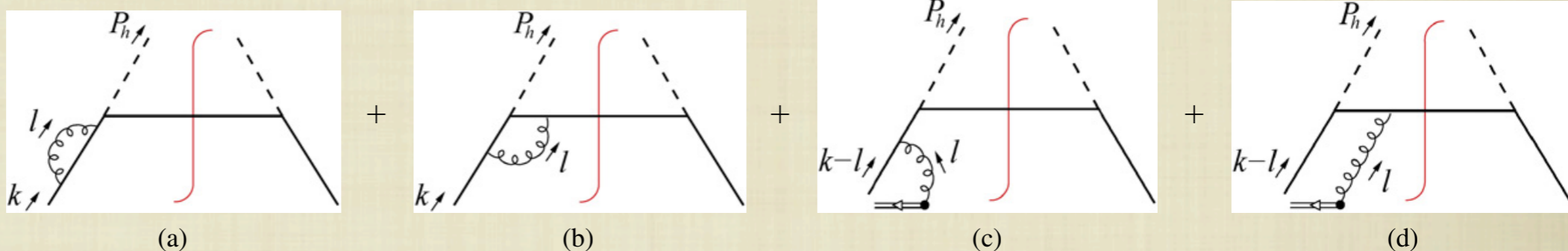


and / or

large cancellations

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**g loops**



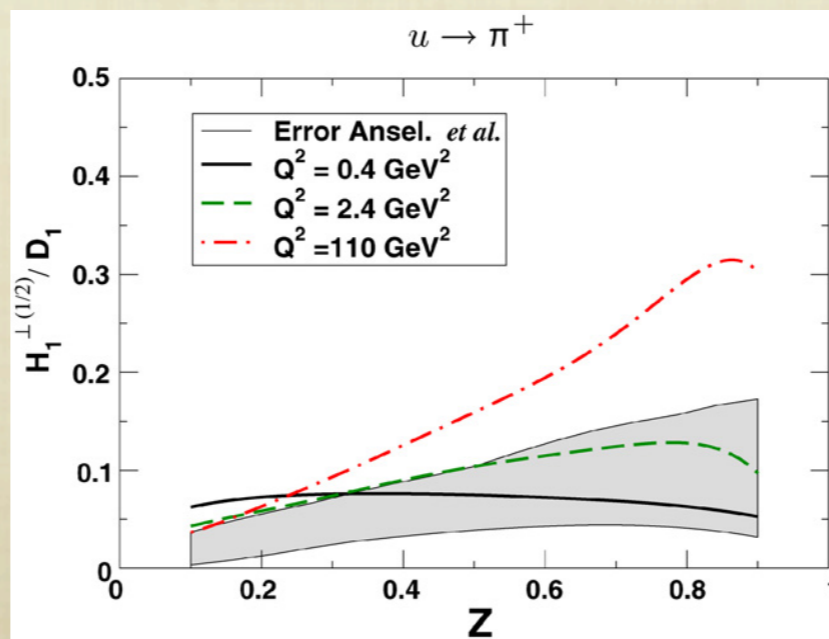
large cancellations

= 0

dominant

PS coupling  
g loops  
params fitted to  $D_1$

Bacchetta et al., P.L.B659 (08)



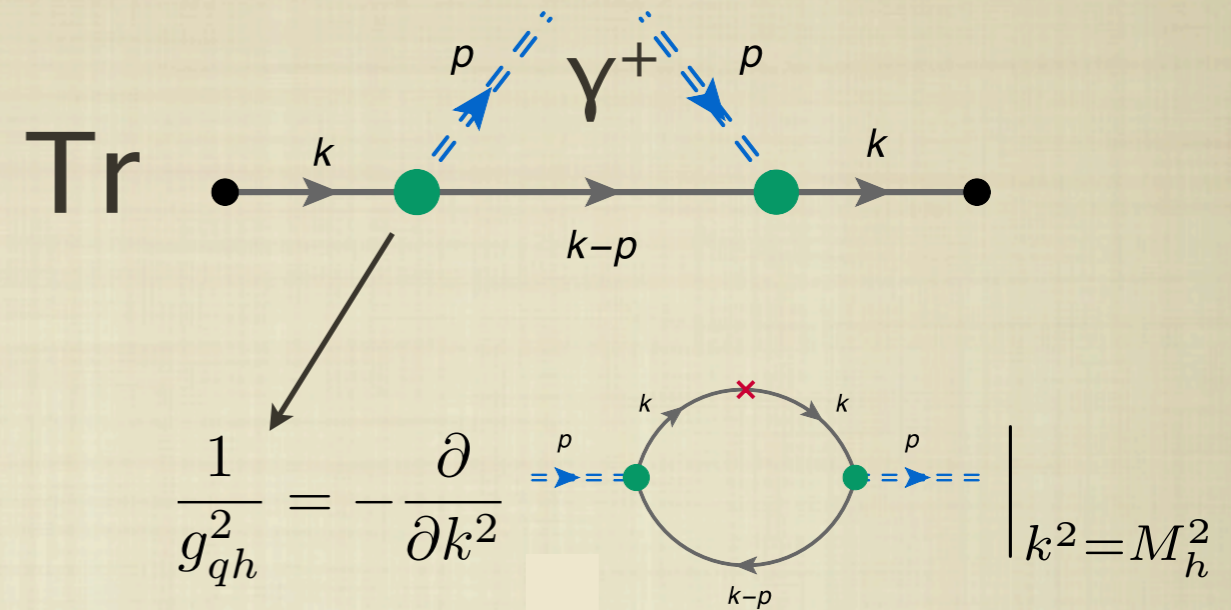
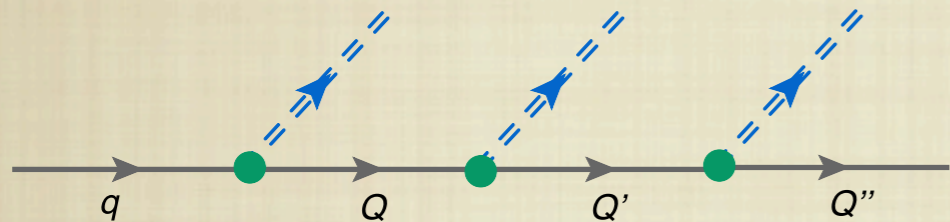


# 2<sup>nd</sup> category: the NJL-jet model

( Ito et al., P.R.D80 (09) 074008 )

- elementary fragm.  $d_q^h(z)$  from
- multiplicative ansatz

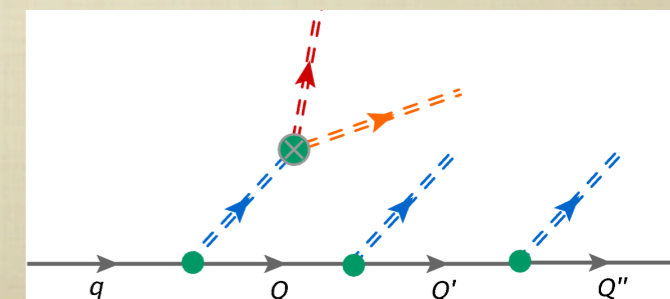
$$D_q^h(z) = d_q^h(z) + \sum_Q [d_q^Q \otimes D_Q^h](z)$$



- mom. sum rule satisfied in Bjorken limit ( $\#h's \rightarrow \infty$ )
- probabilistic interpretation  $\rightarrow$  Monte Carlo (sample based on  $d_q^h$ )

$$D_q^h(z) \Delta z = 1/N \sum_N N_q^h(z, z + \Delta z) \quad \text{MC for } N \rightarrow \infty \rightarrow \text{ansatz}$$

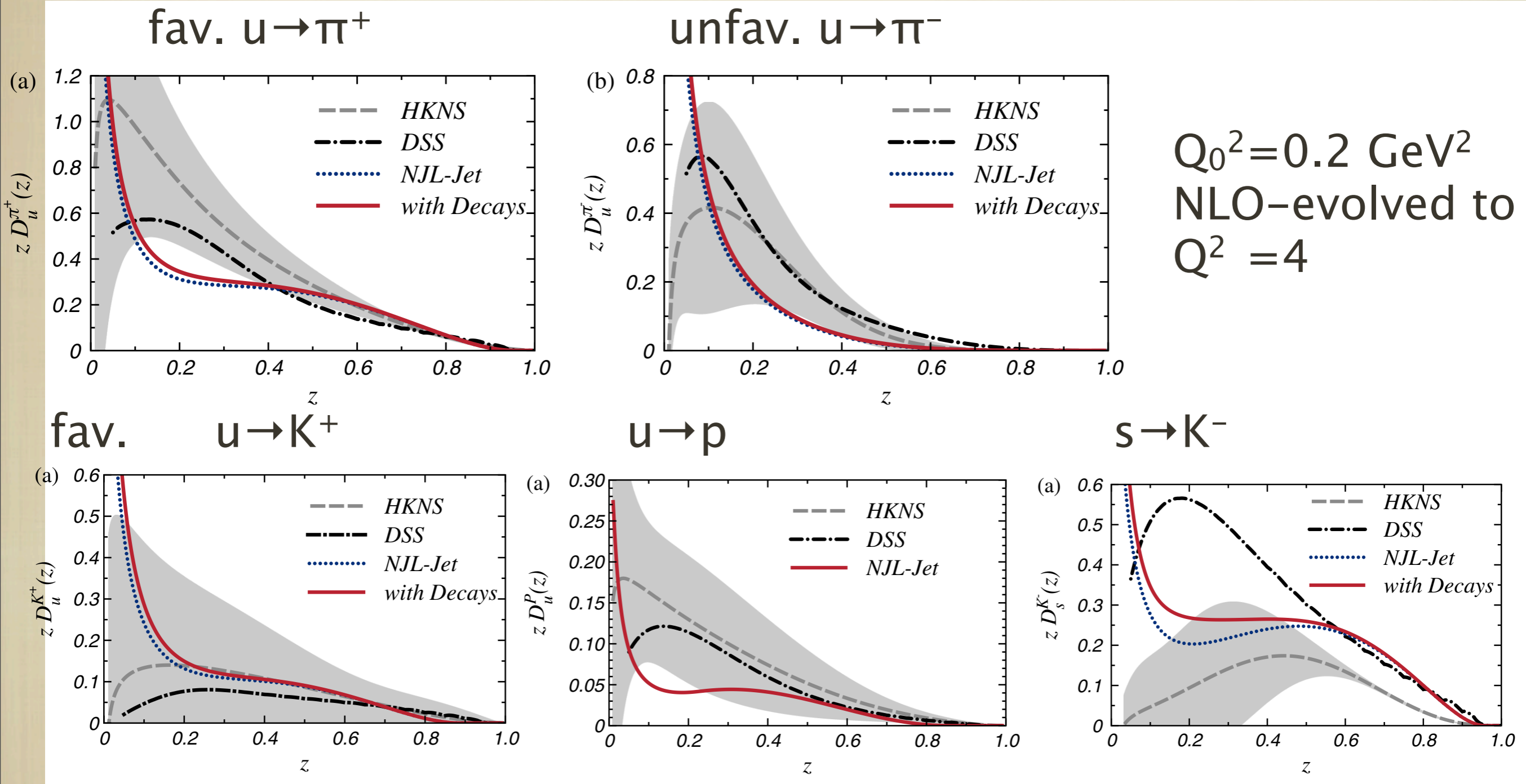
- $h = \pi, K, \rho, K^*, \phi, \rho, n \rightarrow$  spect. diquark model  
**only scalar**





# results for the MC ~ NJL-jet model

( Matevosyan et al., P.R.D83 (11) 114010 )

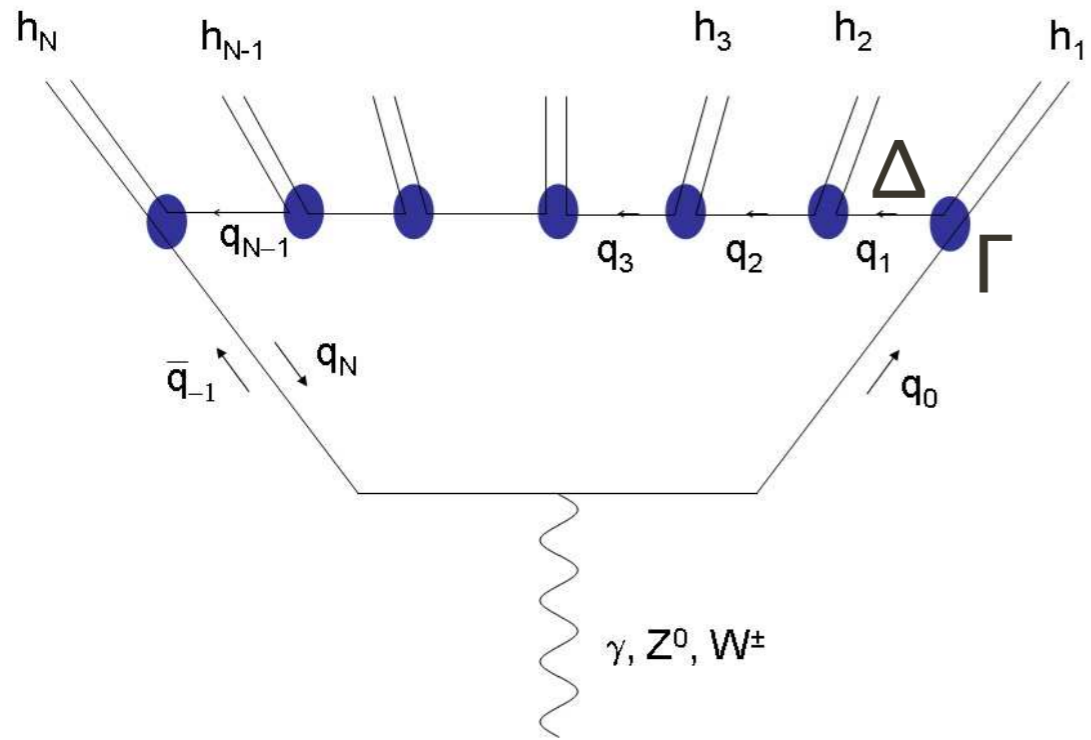


- ✘  $zD(z, Q_0^2) \rightarrow \text{const}$  for  $z \rightarrow 0$  (mult.  $\rightarrow \infty$ ), larger effect at  $Q^2 = 4$
- ✘ LB regular. scheme  $\Rightarrow z_{\min}(h) \leq z \leq z_{\max}(h)$



# 3<sup>rd</sup> category: recursive model with spin

( Artru, arXiv:1001.1061 )



$$e^+ e^- \rightarrow q_0 \bar{q}_{-1} \rightarrow h_1 + h_2 + \dots + h_N$$

$$p_{q_0} = p_{h_1} + p_{q_1}$$

$$p_{q_1} = p_{h_2} + p_{q_2}$$

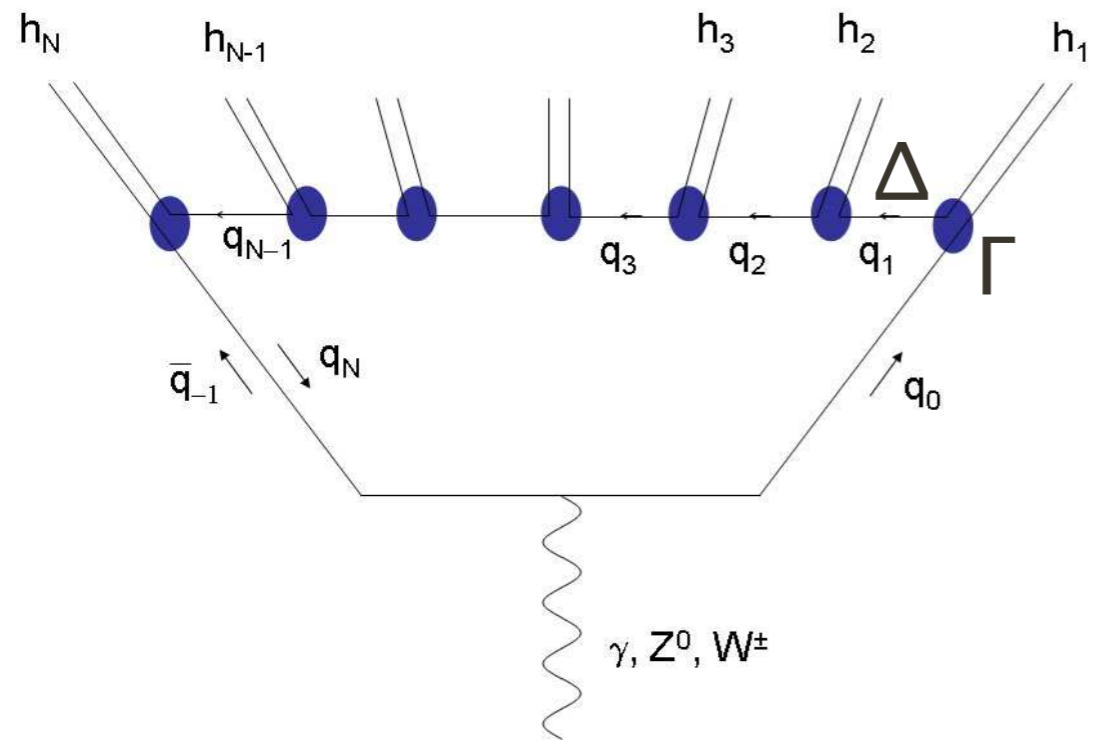
.....

$$\mathcal{M} = \bar{v}(-1) \Gamma(N) \Delta(N-1) \dots \Delta(1) \Gamma(1) u(0)$$



# 3<sup>rd</sup> category: recursive model with spin

( Artru, arXiv:1001.1061 )



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Simplifications :

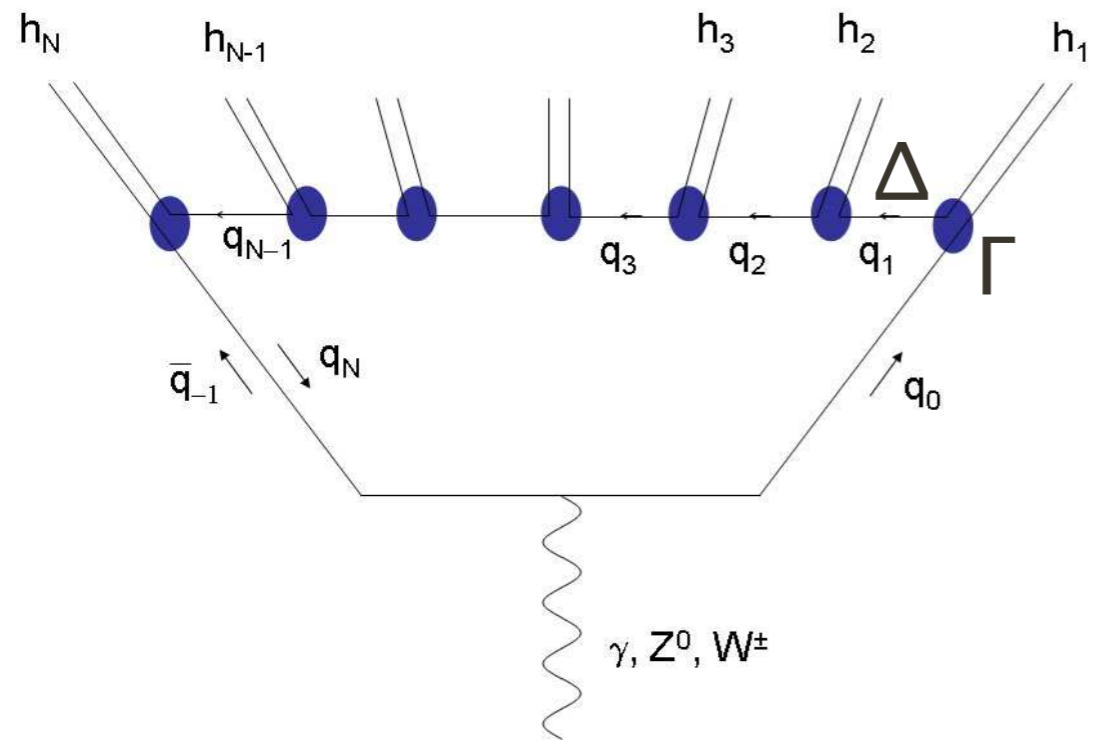
1-  $\Gamma = \text{const.}$

2-  $\Delta(p_q) \approx \exp[-b \mathbf{p}_{qT}^2 / 2] [ \mu(\mathbf{p}_{qT}^2) + i \boldsymbol{\sigma} \cdot \check{\mathbf{z}} \times \mathbf{p}_{qT} ]$  with  $b$  some parameter



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Simplifications :

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$$\mathcal{M} \mathcal{M}^\dagger \approx \exp[-b \mathbf{p}_{h_1 T}^2 \dots - b \mathbf{p}_{h_N T}^2] \text{Tr} \{ M_1 \dots M_N (1 + \mathbf{S}_0 \cdot \boldsymbol{\sigma}) M_N^\dagger \dots M_1^\dagger \}$$

$$M_i = [ \mu(\mathbf{p}_{h_i T}^2) + i \boldsymbol{\sigma} \cdot \check{\mathbf{z}} \times \mathbf{p}_{h_i T} ] \sigma_z$$

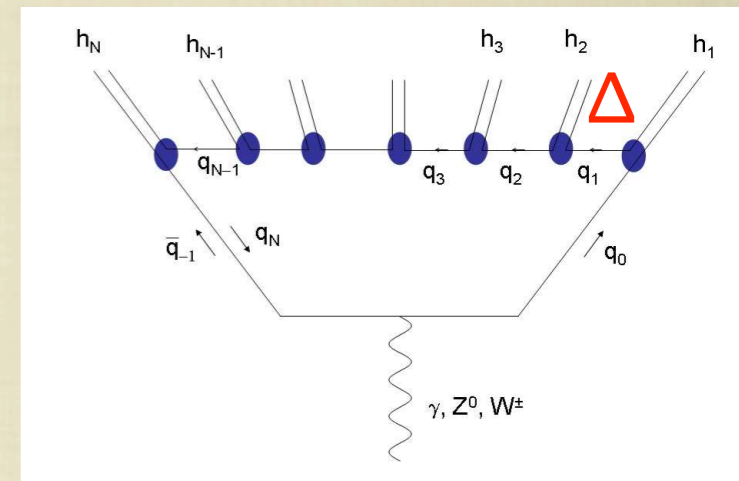


# recursive model with spin: Collins and jet handedness

( Artru, arXiv:1001.1061 )

## N=1: the Collins effect

$$\begin{aligned} \mathcal{M}\mathcal{M}^\dagger &\approx \exp[-b\mathbf{p}_{h1T}^2] \text{Tr} \{M_1 (1 + \mathbf{S}_0 \cdot \boldsymbol{\sigma}) M_1^\dagger\} \\ &= \exp[-b\mathbf{p}_{h1T}^2] [\sigma^0(\mathbf{p}_{h1T}^2) + \text{Im}(\mu) \mathbf{S}_0 \cdot \check{\mathbf{z}} \times \mathbf{p}_{h1T}] \end{aligned}$$



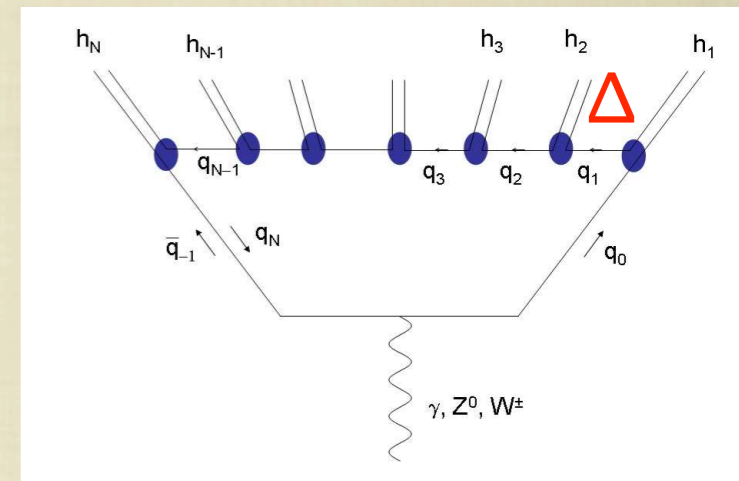


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N=2: iterated **Collins** effect + **jet handedness**

$$\begin{aligned} \mathcal{M}\mathcal{M}^\dagger &\approx \exp[-b\mathbf{p}_{h1T}^2 - b\mathbf{p}_{h2T}^2] \text{Tr} \{M_1 M_2 (1 + \mathbf{S}_0 \cdot \boldsymbol{\sigma}) M_2^\dagger M_1^\dagger\} \\ &= \dots + A(\mathbf{p}_{h2T}^2) \text{Im}(\mu) \mathbf{S} \cdot \check{\mathbf{z}} \times \mathbf{p}_{h1T} + A'(\mathbf{p}_{h1T}^2) \text{Im}(\mu) \mathbf{S} \cdot \check{\mathbf{z}} \times \mathbf{p}_{h2T} \\ &\quad - 2 \text{Im}(\mu^2) S_z \check{\mathbf{z}} \cdot \mathbf{p}_{h1T} \times \mathbf{p}_{h2T} \end{aligned}$$

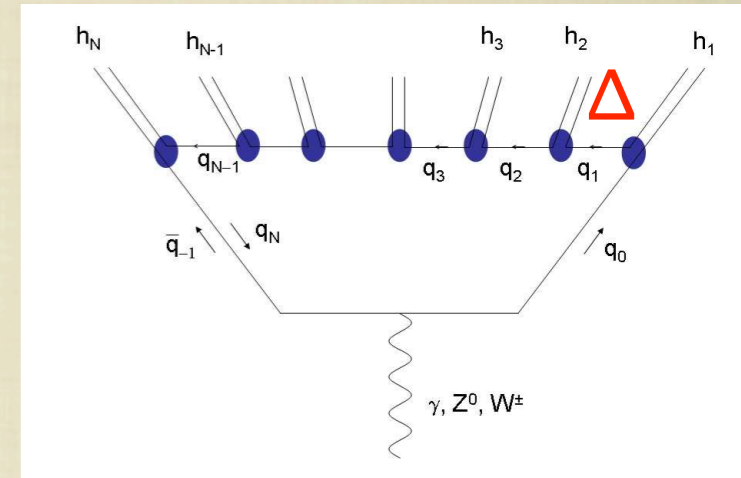


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N=2: iterated **Collins** effect + **jet handedness**

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why ?



# recursive model with spin: Collins and jet handedness

( Artru, arXiv:1001.1061 )

- define recursive property
- $$R_N = M_1 \dots M_N (1 + \mathbf{S}_0 \cdot \boldsymbol{\sigma}) M_N^\dagger \dots M_1^\dagger$$
- $$R_N = M_N R_{N-1} M_N^\dagger$$





# recursive model with spin: Collins and jet handedness

( Artru, arXiv:1001.1061 )

- define recursive property  $R_N = M_1..M_N (1 + \mathbf{S}_0 \cdot \boldsymbol{\sigma}) M_N^\dagger .. M_1^\dagger$   
 $R_N = M_N R_{N-1} M_N^\dagger$
- implies  $\mathbf{S}_N = 1/\text{Tr}\{R_N\} [ \text{Im}(\mu) \check{\mathbf{z}} \times \mathbf{p}_{qNT} + \mathcal{R}(\check{\mathbf{z}}; \mu, \mathbf{p}_{qT}^2) \mathbf{S}_{N-1} ]$ 
  - $\text{Im}(\mu) \neq 0 \Rightarrow \mathbf{S}_{NT} \neq 0$  even if  $\mathbf{S}_{N-1} = 0$

**helicity  $\mathbf{S}_{N-1z} \leftrightarrow$  transversity  $\mathbf{S}_{NT}$**



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**helicity  $S_{N-1z} \leftrightarrow$  transversity  $S_{NT}$**

- **jet handedness** = 1)  $S_{0z} \rightarrow S_{1T} \parallel \mathbf{p}_{h1T} \neq 0$   
2) Collins effect  $\check{\mathbf{z}} \cdot \mathbf{p}_{h2T} \times \mathbf{S}_{1T} \Rightarrow \check{\mathbf{z}} \cdot \mathbf{p}_{h2T} \times \mathbf{p}_{h1T}$



# recursive model with spin: Collins and jet handedness

( Artru, arXiv:1001.1061 )

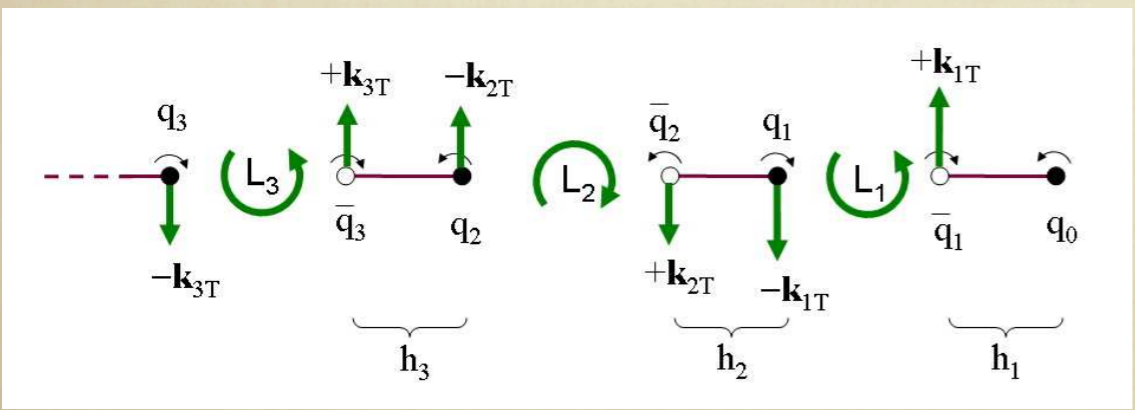
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 2) Collins effect  $\check{\mathbf{z}} \cdot \mathbf{p}_{h2T} \times \mathbf{S}_{1T} \Rightarrow \check{\mathbf{z}} \cdot \mathbf{p}_{h2T} \times \mathbf{p}_{h1T}$

- implies  $S_{Nz} = D_{LL}(|\mu^2|) S_{N-1z}$  ;  $\mathbf{S}_{NT} = D_{TT}(|\mu^2|) \mathbf{S}_{N-1T}$   $2|D_{TT}| \leq 1 + D_{LL}$
- $D_{TT} < 0 \Rightarrow$  **alternate Collins effects** on  $h_1, h_2..$  as in Lund  $^3P_0$  model

**unfav.  $\sim$  - fav.**





# recursive model with spin: Collins and jet handedness

( Artru, arXiv:1001.1061 )

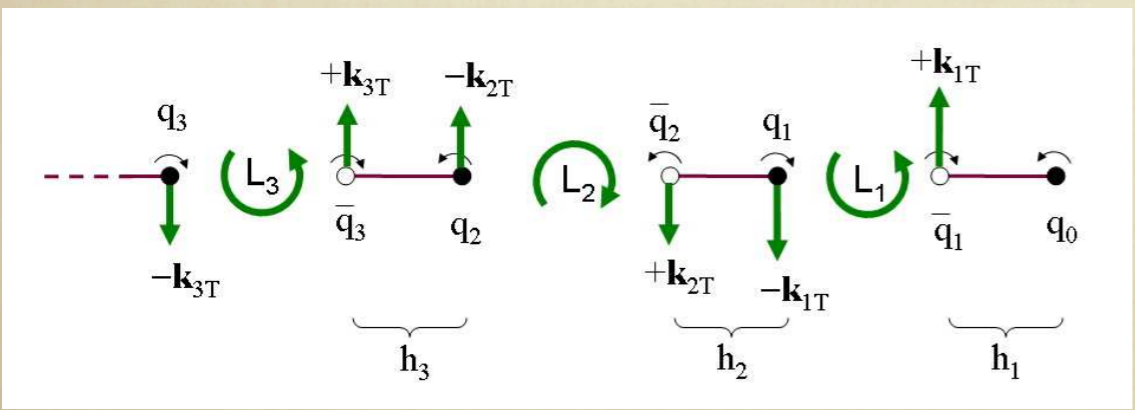
- define recursive property  $R_N = M_1..M_N (1 + \mathbf{S}_0 \cdot \boldsymbol{\sigma}) M_N^\dagger .. M_1^\dagger$   
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**unfav. ~ - fav.**



**BUT**  $A_{UT}^{\text{Coll}}(K^-) \sim 0$  at HERMES  
 different trend at COMPASS  
 large  $A_{UU}^{\cos^2\phi}(K^-)$  at HERMES

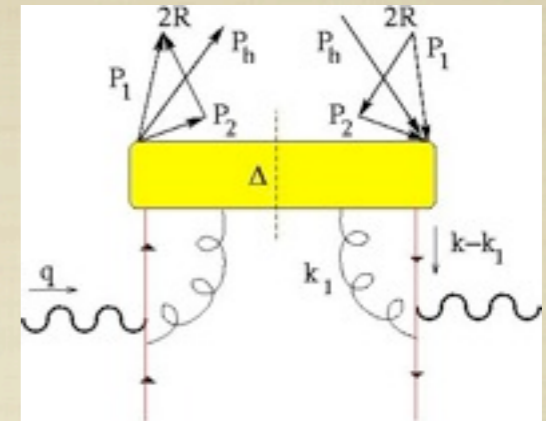


2h FF (DiFF)



# Di-hadron Fragn. Functions (DiFF)

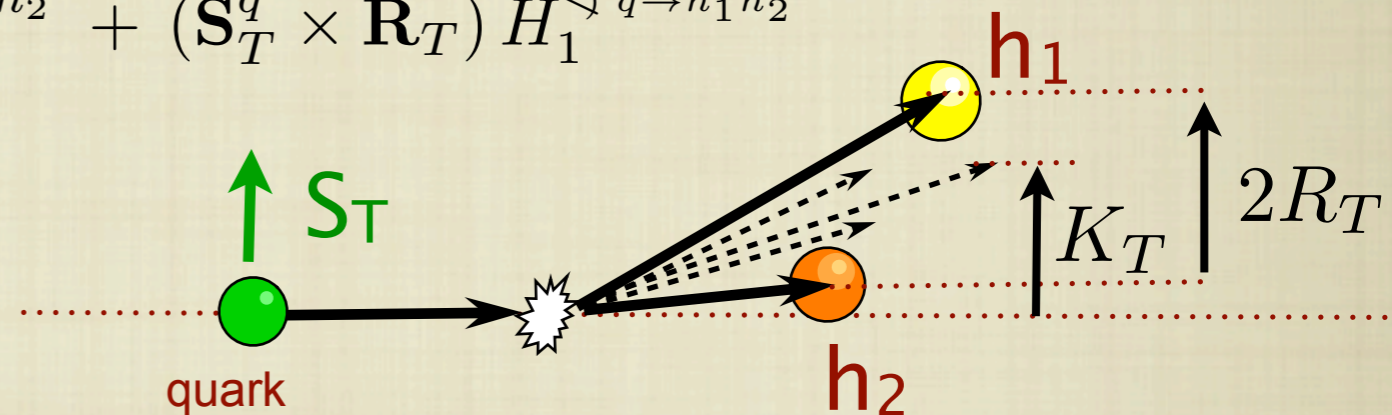
from q-q correlator  $\Delta(z_1, z_2, \mathbf{K}_T, \mathbf{R}_T)$   
project out (at leading twist):



$$\text{Tr} [\Delta \gamma^-] \rightarrow D_1^{q \rightarrow h_1 h_2}(z_1, z_2, K_T^2, R_T^2, \mathbf{K}_T \cdot \mathbf{R}_T)$$

$$\text{Tr} [\Delta \gamma^- \gamma_5] \rightarrow (\mathbf{R}_T \times \mathbf{K}_T) G_1^\perp{}^{q \rightarrow h_1 h_2}$$

$$\text{Tr} [\Delta i\sigma^{i-} \gamma_5] \rightarrow (\mathbf{S}_T^q \times \mathbf{K}_T) H_1^\perp{}^{q \rightarrow h_1 h_2} + (\mathbf{S}_T^q \times \mathbf{R}_T) H_1^{\triangleleft}{}^{q \rightarrow h_1 h_2}$$



First suggested in Konishi et al., P.L.B78 (78)

Polarized DiFF in Collins et al., N.P.B420 (94); Jaffe et al., P.R.L.80 (98); Artru & Collins, Z.Ph.C69 (96)

Jet handedness in Efremov et al., P.L.B284 (92); Stratmann & Vogelsang, P.L.B295 (92); Boer et al., P.R.D67 (03)

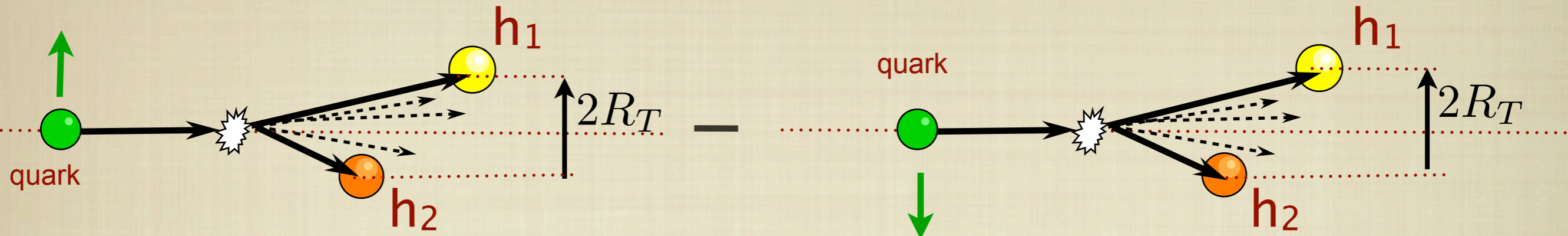
full analysis at twist 2 Bianconi et al., P.R.D62 (00); at twist 3 Bacchetta & Radici, P.R.D69 (04)

LO evolution eqs. Ceccopieri et al., P.L.B650 (07)



chiral-odd  $H_1^{\triangleleft} q \rightarrow h_1 h_2$  survives  $\int d\mathbf{K}_T$  ( $H_1^{\perp} q \rightarrow h$  doesn't)

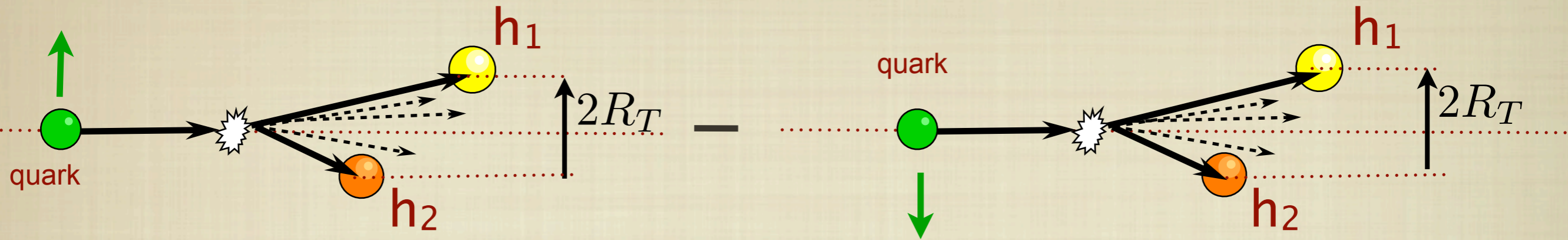
(memo:  $h_1, h_2$  must be distinguishable!)





# chiral-odd $H_1^{\triangleleft q \rightarrow h_1 h_2}$ survives $\int d\mathbf{K}_T$ ( $H_1^{\perp q \rightarrow h}$ doesn't)

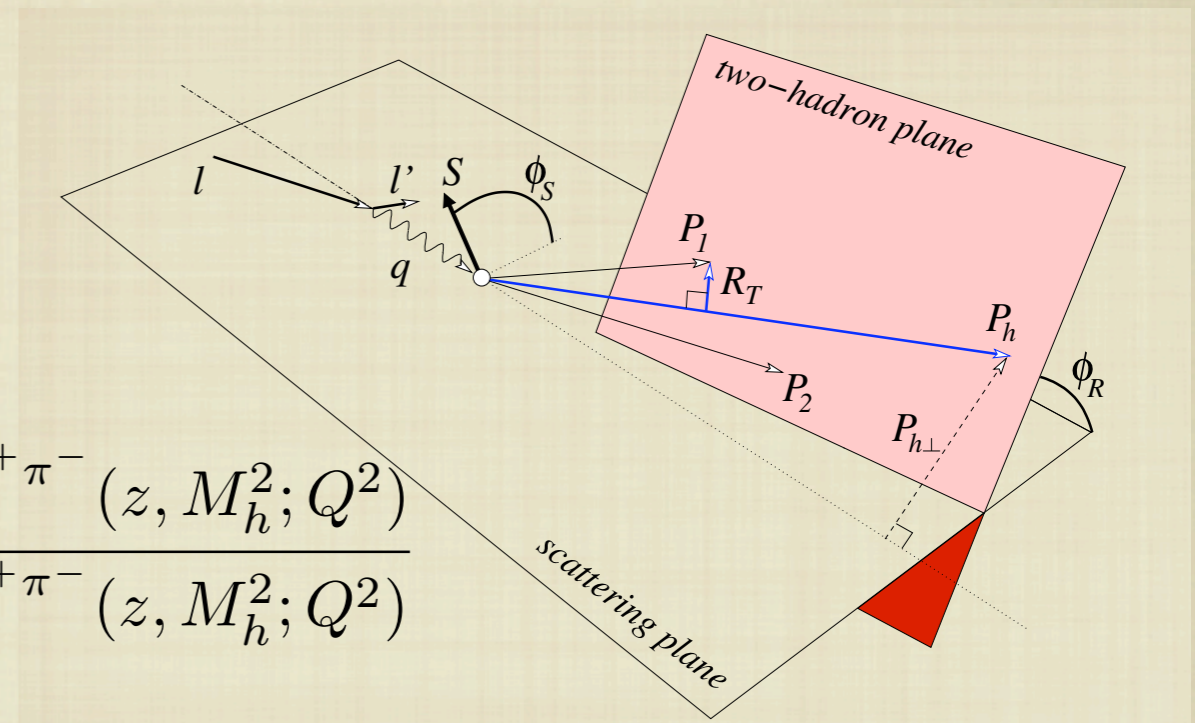
(memo:  $h_1, h_2$  must be distinguishable!)



partner of transversity

$$A_{UT}^{\sin(\phi_R + \phi_S) \sin \theta}(x, z, M_h^2; Q^2) =$$

$$- C_y \frac{|\mathbf{R}|}{M_h} \frac{\sum_q e_q^2 h_1^q(x, Q^2) H_1^{\triangleleft q \rightarrow \pi^+ \pi^-}(z, M_h^2; Q^2)}{\sum_q e_q^2 f_1^q(x, Q^2) D_1^{\triangleleft q \rightarrow \pi^+ \pi^-}(z, M_h^2; Q^2)}$$



Radici et al., PR D65 (02)  
Bacchetta & Radici, PR D67 (03)

- coll. fact.  $\rightarrow$  simple product (no  $\otimes$ )
- DGLAP (LO) evolution
- universality
- cleaner  $e^+e^- \rightarrow (\pi^+\pi^-)(\pi^+\pi^-)X$  (expect less background)



2007: preliminary SIDIS data on D from COMPASS:  $A_{UT} \sim 0$

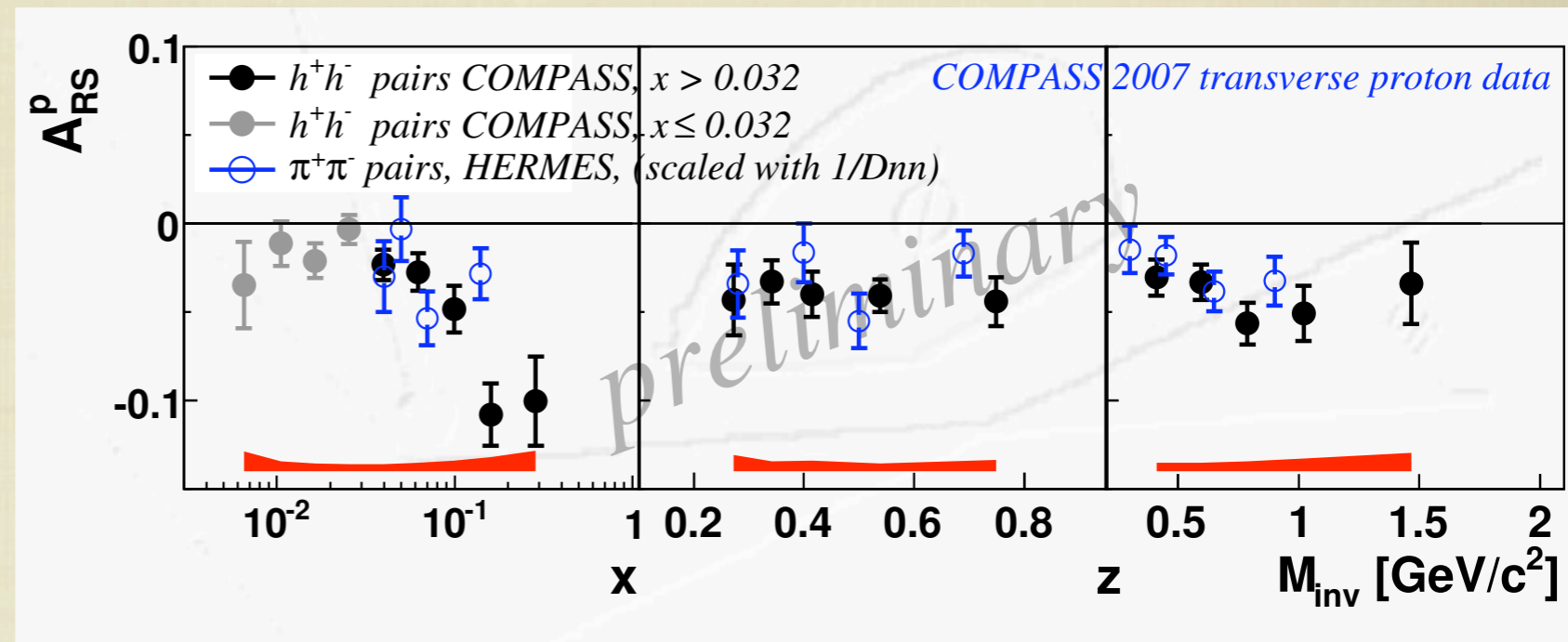
Martin, hep-ex/0702002

2008: first SIDIS data on  $p^\uparrow$  from HERMES

Airapetian et al. (HERMES), JHEP 06 (08)

2009: preliminary data on  $p^\uparrow$  from COMPASS

Wollny (COMPASS), DIS 2009, arXiv:0907.0961





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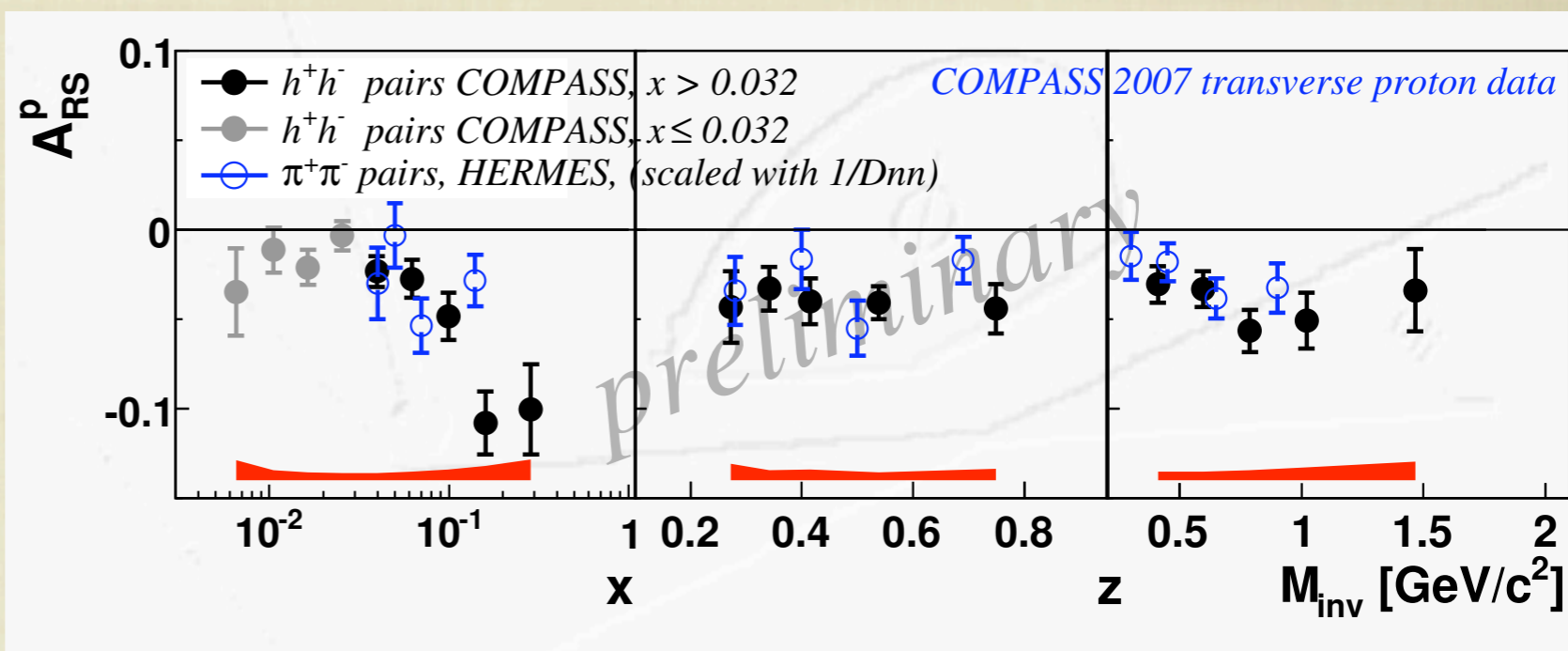
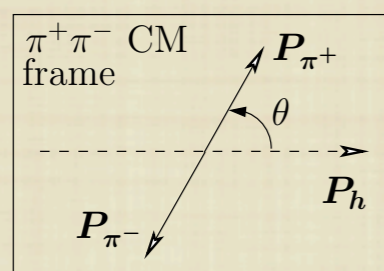
Airapetian et al. (HERMES), JHEP 06 (08)

2009: preliminary data on  $p^\uparrow$  from COMPASS

Wolny (COMPASS), DIS 2009, arXiv:0907.0961

observable is

$$A_{UT}^{\sin(\phi_R + \phi_S) \sin \theta}$$



partial wave expansion

Bacchetta & Radici, P.R.D67 (03)

$$H_1^{\triangleleft} = H_{1,sp}^{\triangleleft}(z, M_h^2) + \cos \theta H_{1,pp}^{\triangleleft}(z, M_h^2)$$

access to interference  $(\pi^+\pi^-)_s \leftrightarrow (\pi^+\pi^-)_p$

model prediction

Bacchetta & Radici, P.R.D74 (06)

model analysis

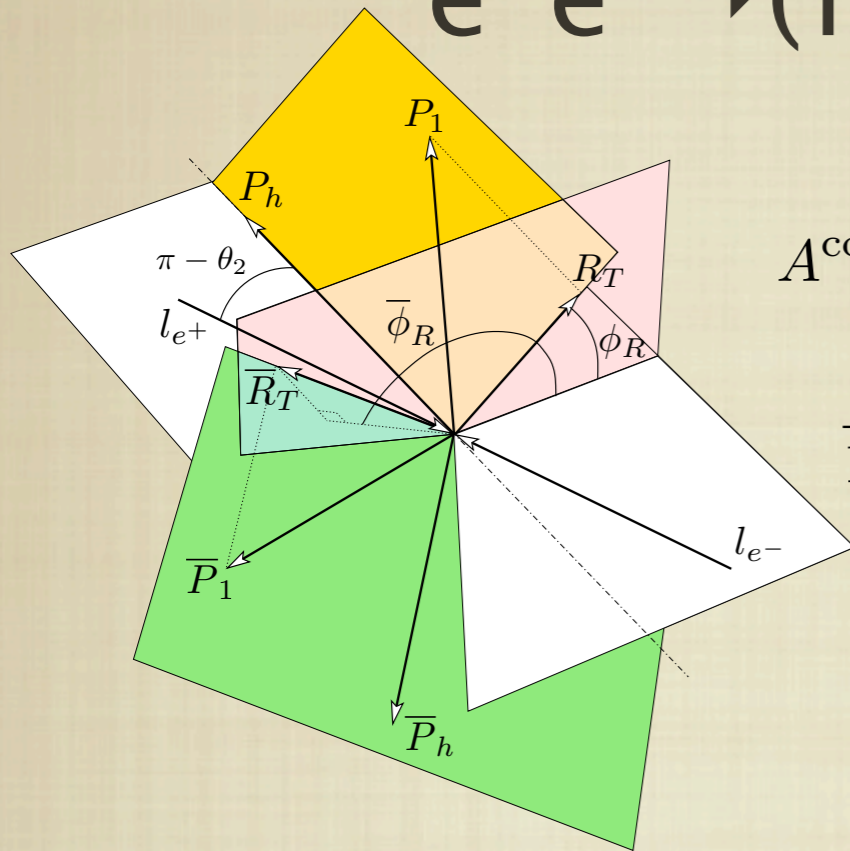
Bacchetta et al., P.R.D79 (09); She et al., P.R.D77 (08)



# 2011: the BELLE data for $a_{12R}$

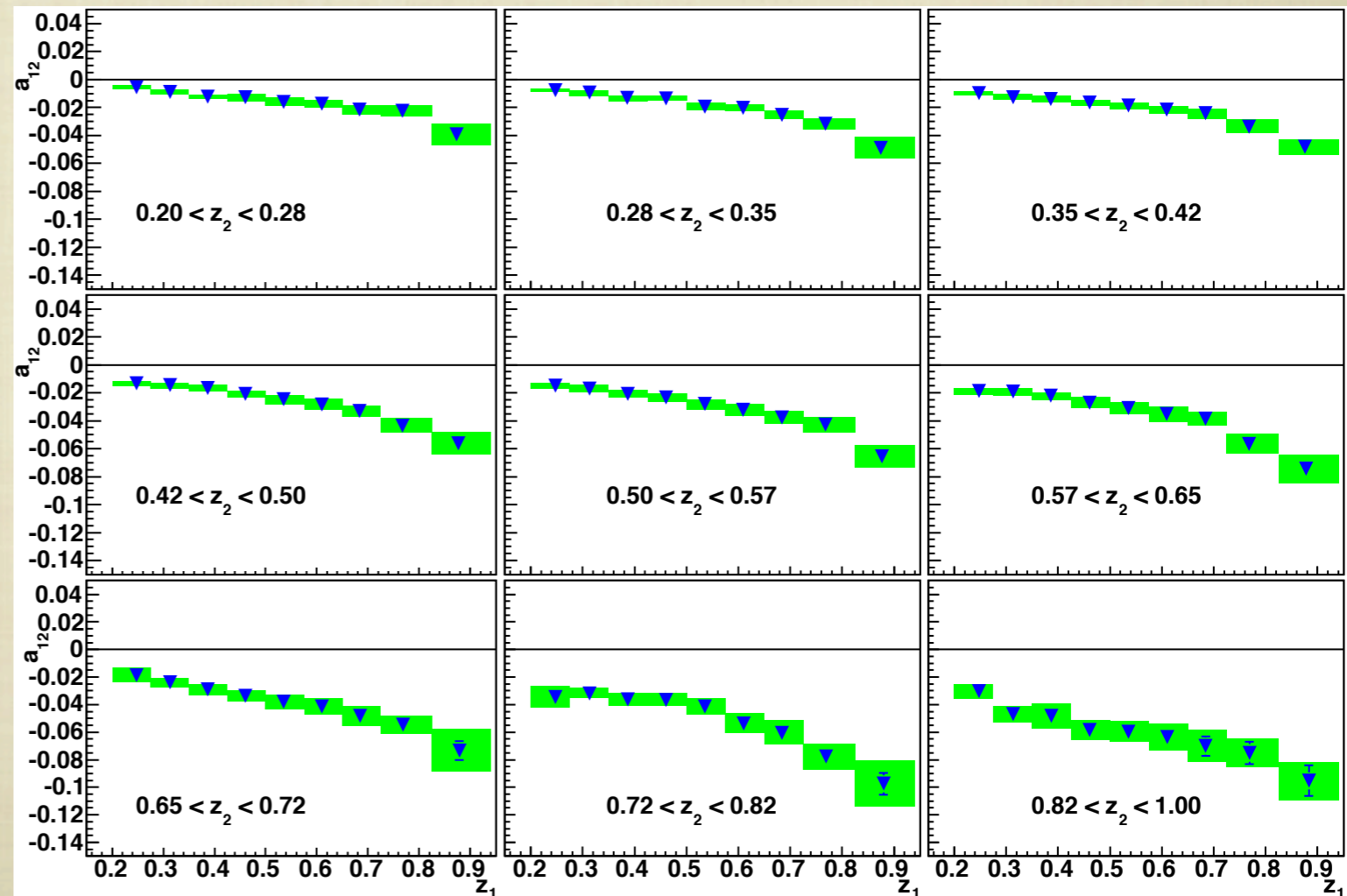
$$e^+e^- \rightarrow (\pi^+\pi^-)(\pi^+\pi^-)X$$

Artru & Collins, Z.Ph.C69 (96)  
Boer et al., P.R.D67 (03)



$$A^{\cos(\phi_R + \bar{\phi}_R)}(\cos\theta_2, z, M_h^2, \bar{z}, \bar{M}_h^2) \equiv a_{12R} \propto$$

$$\frac{\sin^2\theta_2}{1 + \cos^2\theta_2} \frac{|\mathbf{R}_T|}{M_h} \frac{|\bar{\mathbf{R}}_T|}{\bar{M}_h} \frac{\sum_q e_q^2 H_{1,q \rightarrow \pi^+\pi^-}^{\triangleleft}(z, M_h^2) H_{1,\bar{q} \rightarrow \pi^+\pi^-}^{\triangleleft}(\bar{z}, \bar{M}_h^2)}{\sum_q e_q^2 D_{1,q \rightarrow \pi^+\pi^-}(z, M_h^2) D_{1,\bar{q} \rightarrow \pi^+\pi^-}(\bar{z}, \bar{M}_h^2)}$$



Vossen et al. (BELLE),  
arXiv:1104.2425 [hep-ex]



# parametrizing DiFF: fitting BELLE data

$$a_{12R} = \frac{\langle \sin^2 \theta_2 \rangle}{\langle 1 + \cos^2 \theta_2 \rangle} \frac{\langle \sin \theta \rangle |\mathbf{R}|}{M_h} \frac{\langle \sin \bar{\theta} \rangle |\bar{\mathbf{R}}|}{\bar{M}_h} \frac{\sum_q e_q^2 H_{1,q \rightarrow \pi^+ \pi^-}^{\triangleleft}(z, M_h^2) H_{1,\bar{q} \rightarrow \pi^+ \pi^-}^{\triangleleft}(\bar{z}, \bar{M}_h^2)}{\sum_q e_q^2 D_{1,q \rightarrow \pi^+ \pi^-}(z, M_h^2) D_{1,\bar{q} \rightarrow \pi^+ \pi^-}(\bar{z}, \bar{M}_h^2)}$$

$$|\mathbf{R}| = \frac{M_h}{2} \sqrt{1 - \frac{4m_\pi^2}{M_h^2}}$$

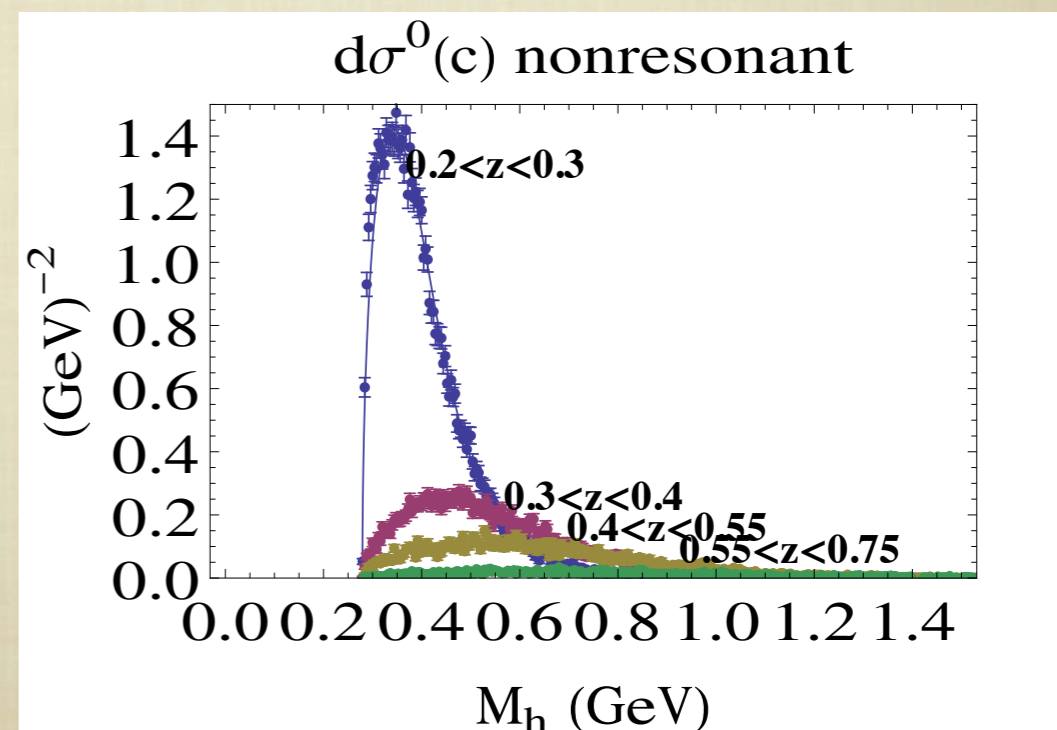
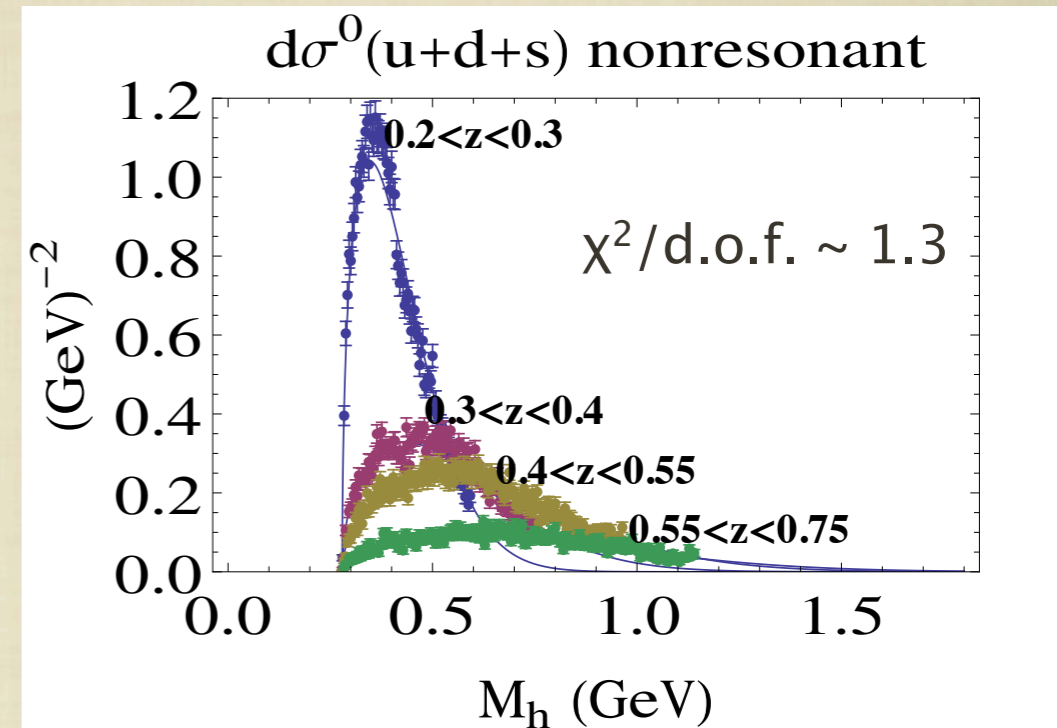
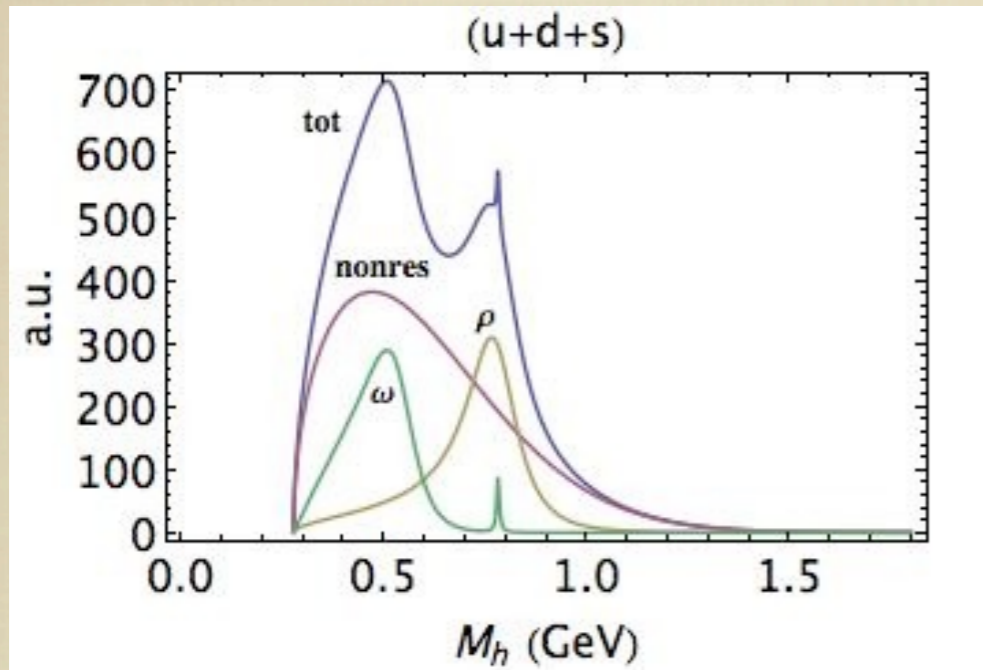
## strategy

1. fit the denominator using the unpolarized cross section generated by PYTHIA MC adapted to BELLE
2. fit the asymmetry  $a_{12R}$  multiplied by denominator ( $\approx [\text{statistical error}]^{-1}$ )  
→ get the numerator, bin by bin



# 1. fitting the BELLE (MC) $d\sigma^0 \rightarrow D_1^q \rightarrow \pi^+\pi^-$

1. flavor decomposition: {uds} – charm
2. resonant ( $\rho, \omega$ ; only {uds}) and nonresonant contributions

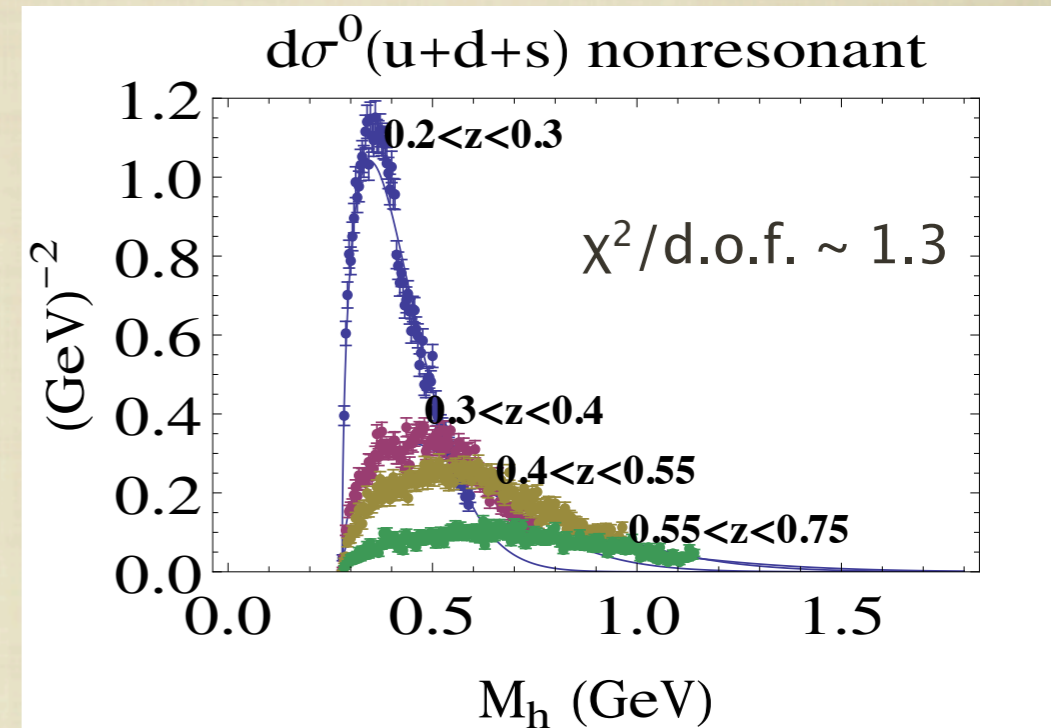
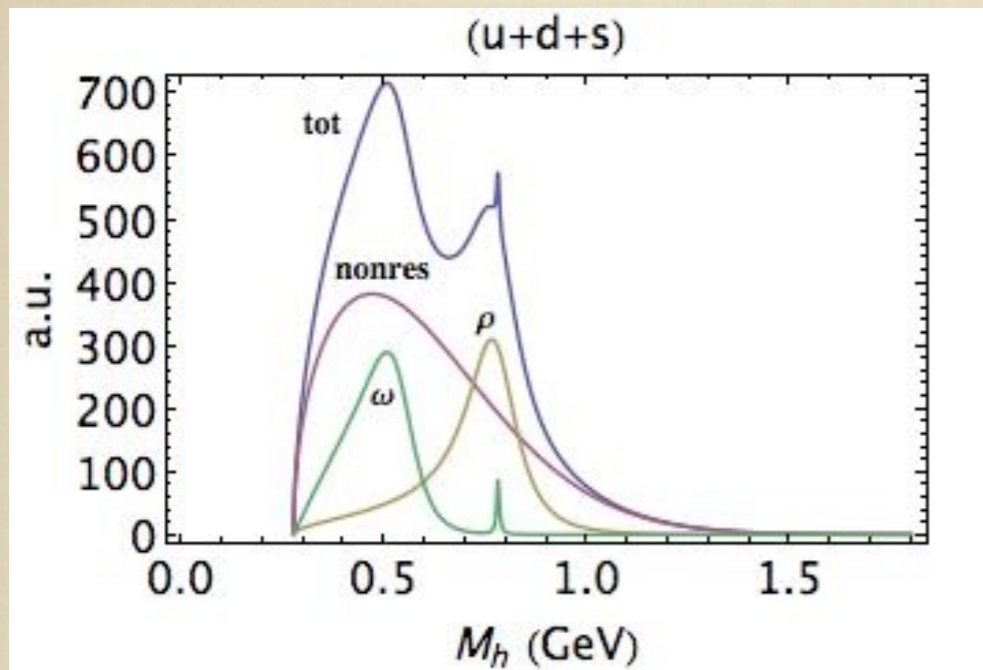


Courtoy et al., arXiv:1012.0054 [hep-ph]



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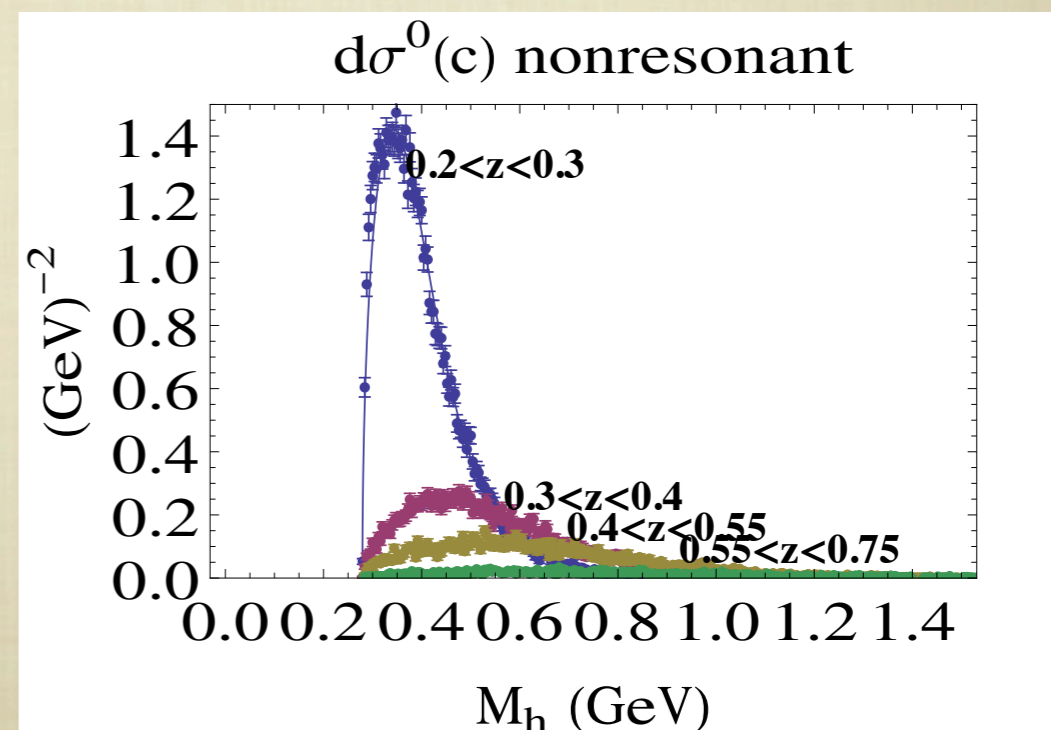


- big effect from charm
- no factorization of  $(z, M_h)$  depend.

work in progress for  $d\sigma^0 \times a_{12R} \dots$

... but ...

Courtoy et al., arXiv:1012.0054 [hep-ph]





# 1<sup>st</sup> extraction of transversity in coll. framework

Bacchetta et al., P.R.L. 107 (11)

$$A_{UT}^{\text{SIDIS}}(x, z, M_h^2; Q^2) = -C_y \frac{|\mathbf{R}|}{M_h} \frac{\sum_q e_q^2 h_1^q(x, Q^2) H_1^{\triangleleft q \rightarrow \pi^+ \pi^-}(z, M_h^2; Q^2)}{\sum_q e_q^2 f_1^q(x, Q^2) D_1^{\triangleleft q \rightarrow \pi^+ \pi^-}(z, M_h^2; Q^2)}$$

$$\int dz \int dM_h^2$$

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symmetry

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$$x h_1^{u_v}(x, Q^2) - \frac{1}{4} x h_1^{d_v}(x, Q^2)$$

$$\Rightarrow - \frac{A_{UT}^{\text{SIDIS}}(x, Q^2)}{C_y} \frac{n_u(Q^2)}{n_u^\uparrow(Q^2)} \sum_{q=u,d,s} \frac{e_q^2 N_q}{e_u^2} x f_1^{q+\bar{q}}(x, Q^2)$$

HERMES (+ COMPASS)

MSTW08LO



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goal  $\longrightarrow xh_1^{uv}(x, Q^2) - \frac{1}{4}xh_1^{dv}(x, Q^2)$

BELLE

$$\xrightarrow{=} -\frac{A_{UT}^{\text{SIDIS}}(x, Q^2)}{C_y} \frac{n_u(Q^2)}{n_u^\uparrow(Q^2)} \sum_{q=u,d,s} \frac{e_q^2 N_q}{e_u^2} x f_1^{q+\bar{q}}(x, Q^2)$$

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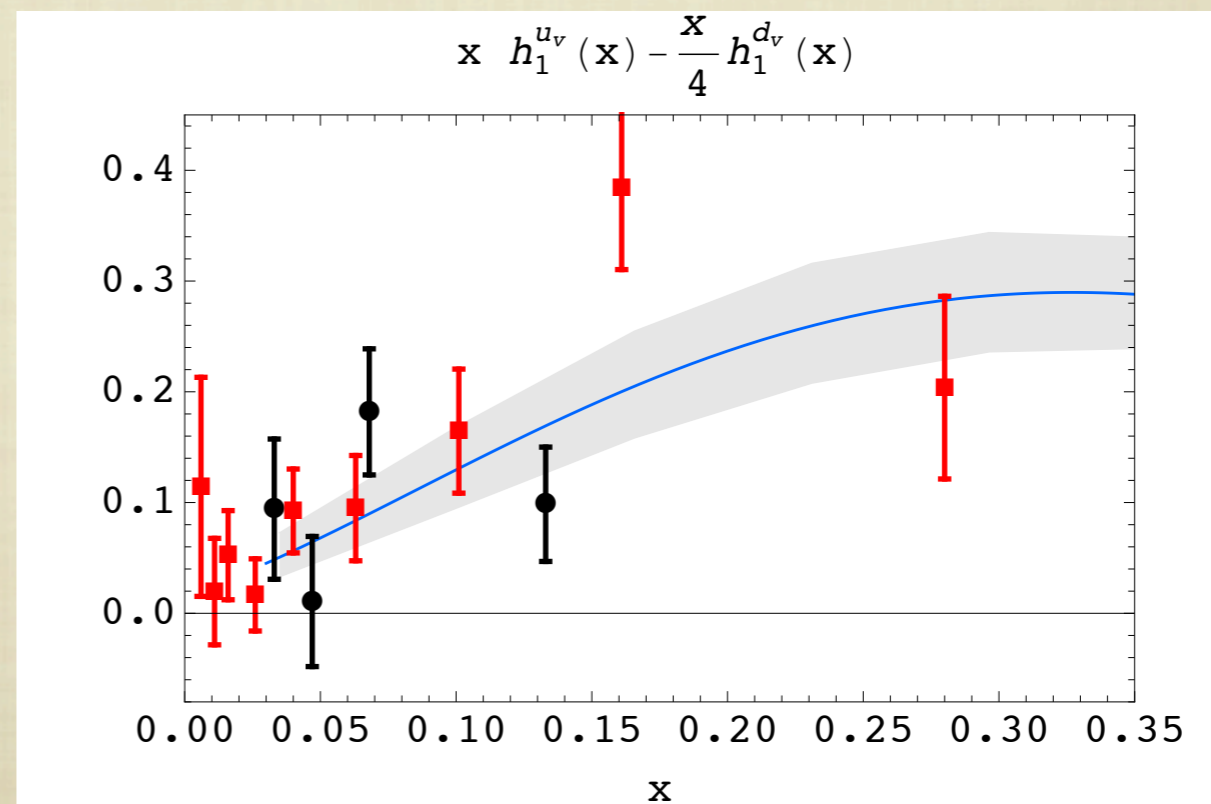
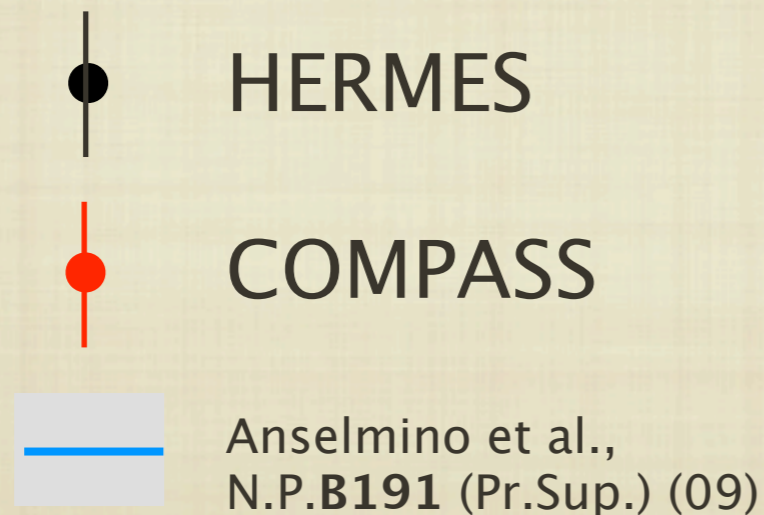
1. start from  $D_1^{q=u,s,c}(z, M_h; Q_0^2=1)$ ,  $H_1^{<)u}(z, M_h; Q_0^2=1)$  Bacchetta & Radici, P.R.D74 (06)  
resonant + nonresonant channel inspired by spect. model
2. evolve at LO with HOPPET (updating with chiral-odd kernel)
3. fit  $d\sigma^0$  from PYTHIA (adapted to BELLE) and  $d\sigma^0 \times a_{12R}$  bin by bin
4. integrate  $D_1^q$  and  $H_1^{<)u}$  in HERMES range  $0.5 \leq M_h \leq 1$ ,  $0.2 \leq z \leq 0.7$
5. get  $n_u^\uparrow(Q^2)/n_u(Q^2)$ :  $Q^2=2.5 \text{ GeV}^2$   $n_u^\uparrow/n_u = -0.251 \pm 0.006_{\text{ex}} \pm 0.023_{\text{th}}$   
 $[n_u^\uparrow/n_u(2.5)] / [n_u^\uparrow/n_u(100)] \sim 92\%(\pm 8\%)$



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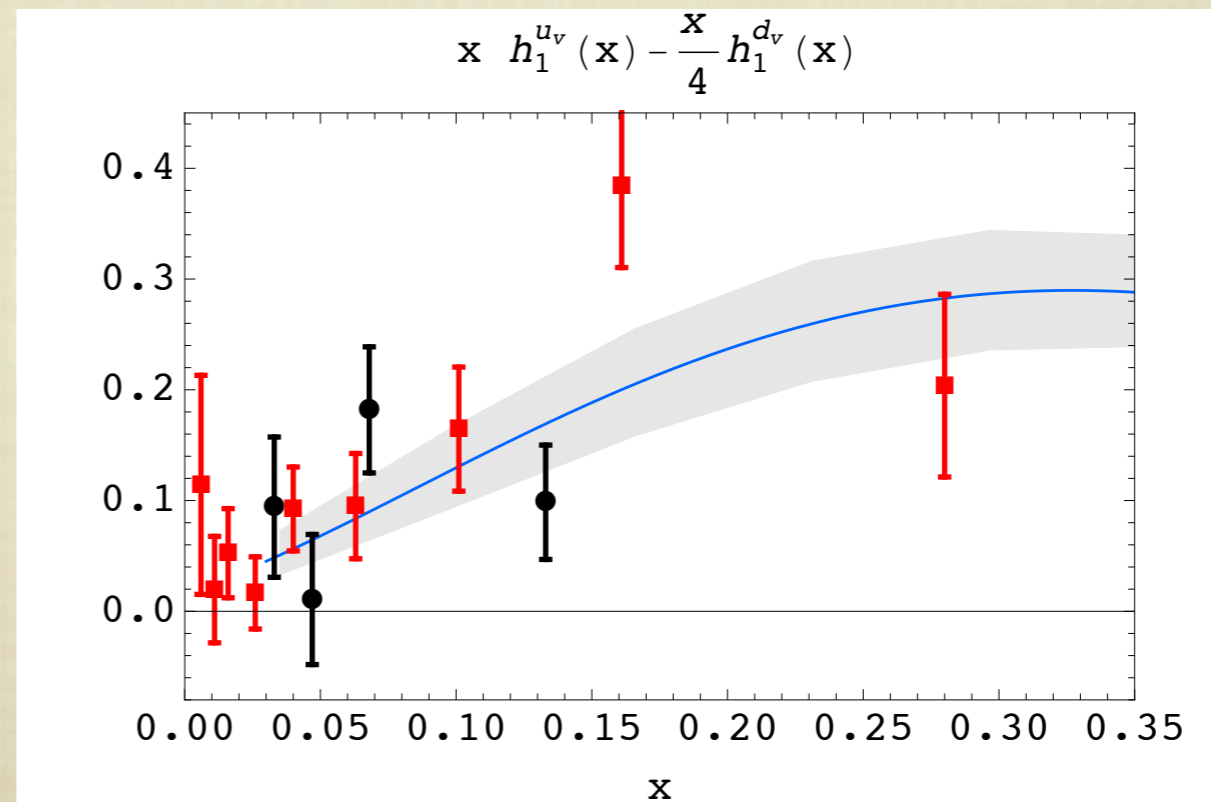




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several “BUT..”  
work in progress  
stay tuned..