

Probing Transverse Momentum Distributions with the Drell-Yan process

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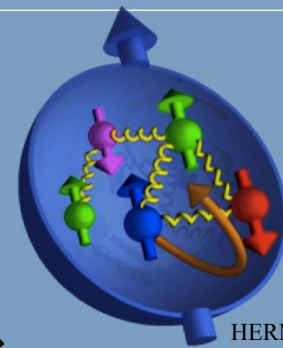
1. The Drell-Yan Process: History and Future
2. Unpolarized Transverse Momentum Measurements

$$\frac{d\sigma}{d\Omega} \propto 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi$$

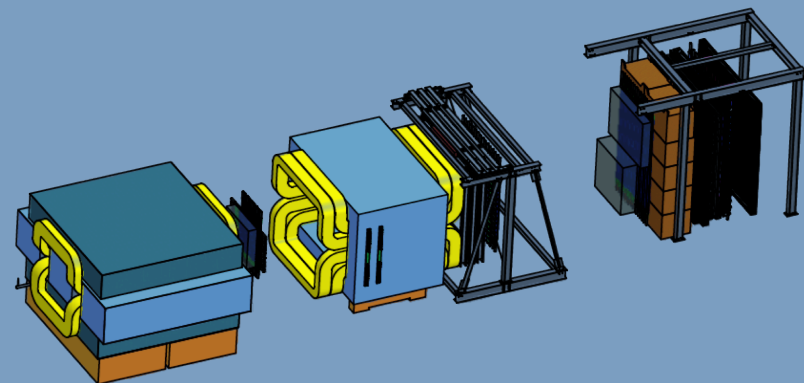
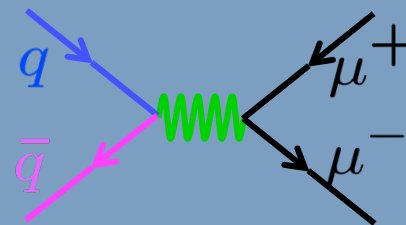
3. Polarized Measurements

$$f_{1T}^{\perp q} \Big|_{DIS} = - f_{1T}^{\perp q} \Big|_{D-Y}$$

4. What can the future hold?



HERMES
U. Elschenbroich



This work is supported in part by the U.S. Department of Energy, Office of Nuclear Physics, under Contract No. DE-AC02-06CH11357.

The journey begins



Observation of Massive Muon Pairs in Hadron Collisions*

J. H. Christenson, G. S. Hicks, L. M. Lederman, P. J. Limon, and B. G. Pope

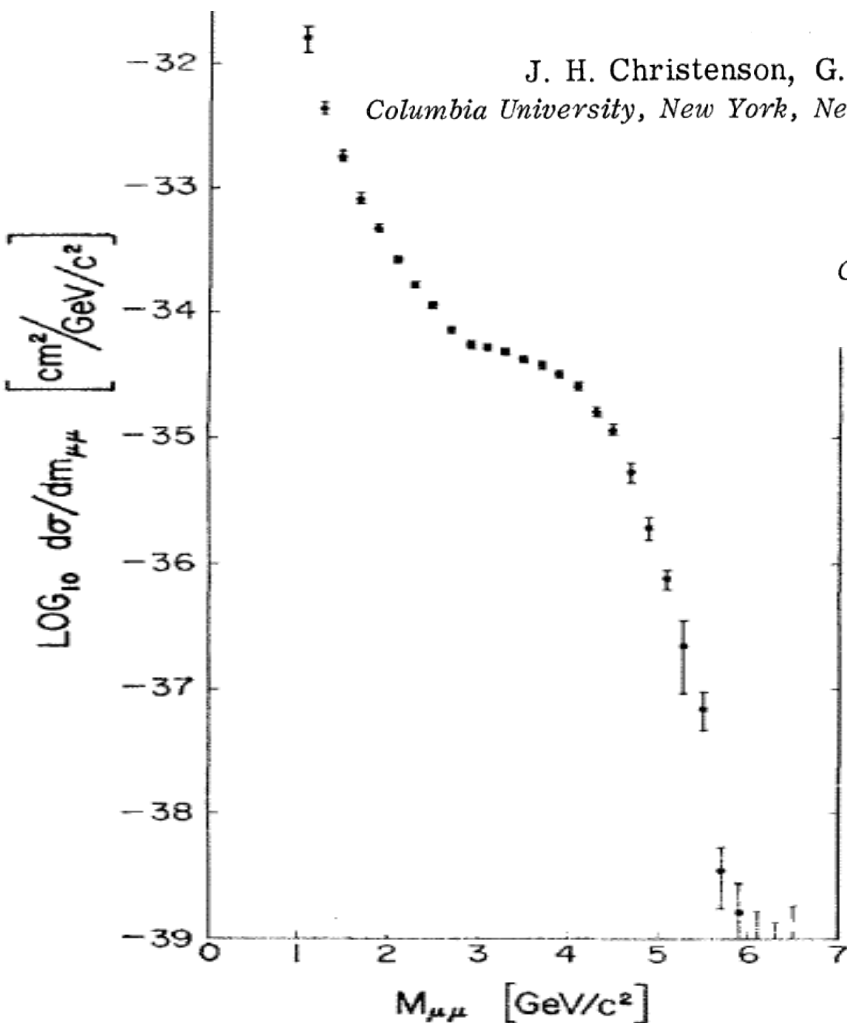
Columbia University, New York, New York 10027, and Brookhaven National Laboratory, Upton, New York 11973

and

E. Zavattini

CERN Laboratory, Geneva, Switzerland

(Received 8 September 1970)



Muon Pairs in the mass range $1 < m_{\mu\mu} < 6.7 \text{ GeV}/c^2$ have been observed in collisions of high-energy protons with uranium nuclei. At an incident energy of 29 GeV, **the cross section varies smoothly as $d\sigma/dm_{\mu\mu} \approx 10^{-32} / m_{\mu\mu}^5 \text{ cm}^2 (\text{GeV}/c)^{-2}$ and exhibits no resonant structure.** The total cross section increases by a factor of 5 as the proton energy rises from 22 to 29.5 GeV.

Early Muon Pair Data—soon to be called Drell-Yan

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PHYSICAL REVIEW LETTERS

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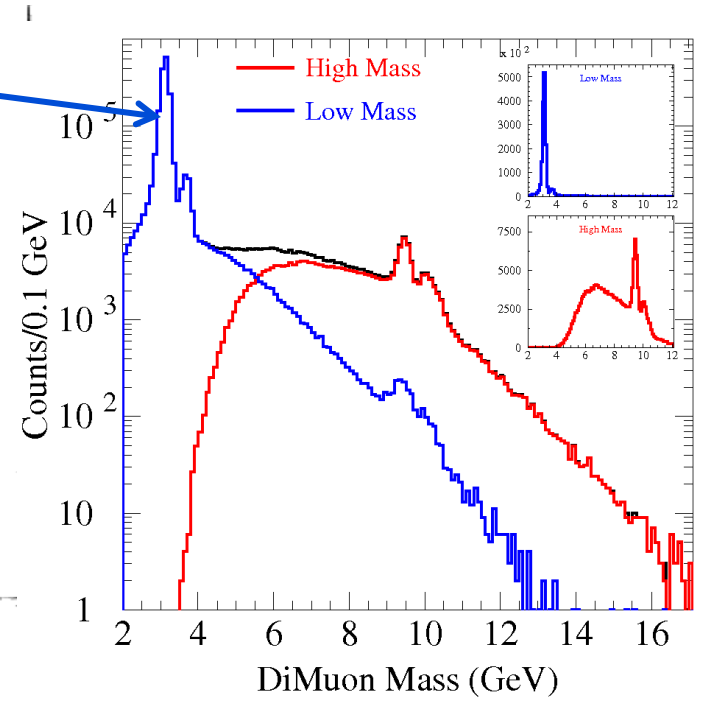
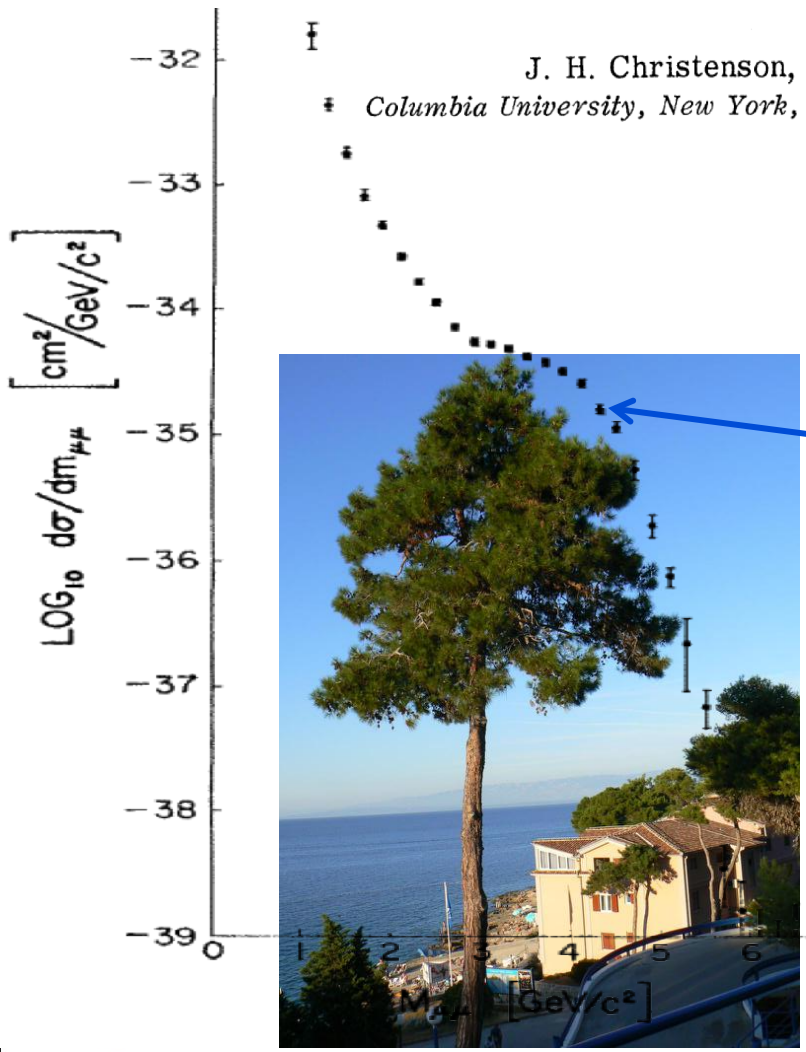
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Don't miss the tree when looking at the beautiful scenery

Paul E. Reimer Transversity 2011

MASSIVE LEPTON-PAIR PRODUCTION IN HADRON-HADRON COLLISIONS AT HIGH ENERGIES*

Sidney D. Drell and Tung-Mow Yan

Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305

(Received 25 May 1970)

On the basis of a parton model studied earlier we consider the production process of large-mass lepton pairs from hadron-hadron inelastic collisions in the limiting region, $s \rightarrow \infty$, Q^2/s finite, Q^2 and s being the squared invariant masses of the lepton pair and the two initial hadrons, respectively. General scaling properties and connections with deep inelastic electron scattering are discussed. In particular, a rapidly decreasing cross section as $Q^2/s \rightarrow 1$ is predicted as a consequence of the observed rapid falloff of the inelastic scattering structure function νW_2 near threshold.

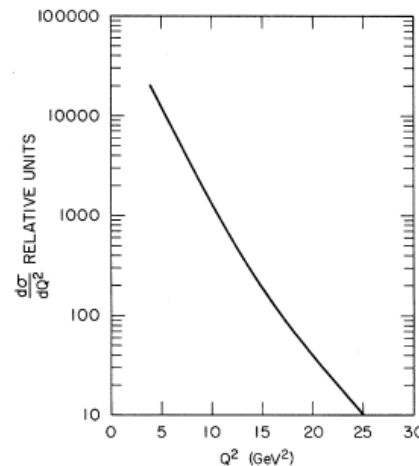
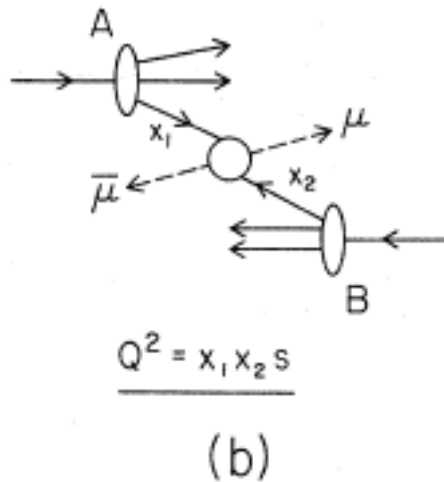
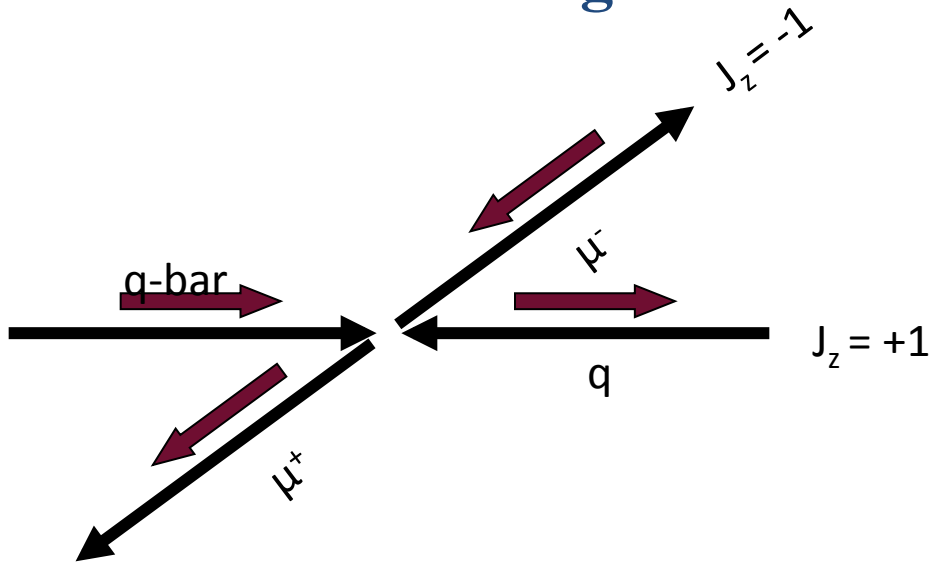


FIG. 2. $d\sigma/dQ^2$ computed from Eq. (10) assuming identical parton and antiparton momentum distributions and with relative normalization.

Also predicted
 $\lambda(1+\cos^2\theta)$
 angular distributions.

LO Drell-Yan Angular Distributions: $d\sigma/d\Omega / (1+\cos^2\theta)$

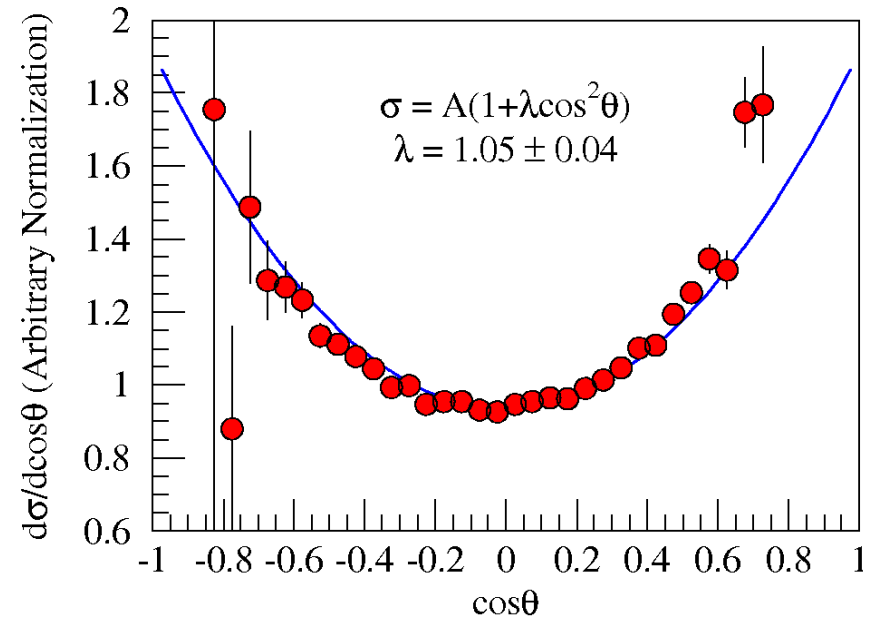


$$\mathcal{M} \propto d_{\lambda', \lambda}^j(\theta) = \langle j \lambda' | e^{-i\theta J_y} | j \lambda \rangle$$

$$d_{1,1}^1 = d_{-1,-1}^1 = \frac{1}{2} (1 + \cos \theta)$$

$$d_{-1,1}^1 = d_{1,-1}^1 = \frac{1}{2} (1 - \cos \theta)$$

$$\frac{d\sigma}{d\Omega} \propto \mathcal{M}^2 \propto (1 + \cos^2 \theta)$$



Fermilab E-866/NuSea unpublished

Helped to validate the Drell-Yan picture of quark-antiquark annihilation for lepton pair production

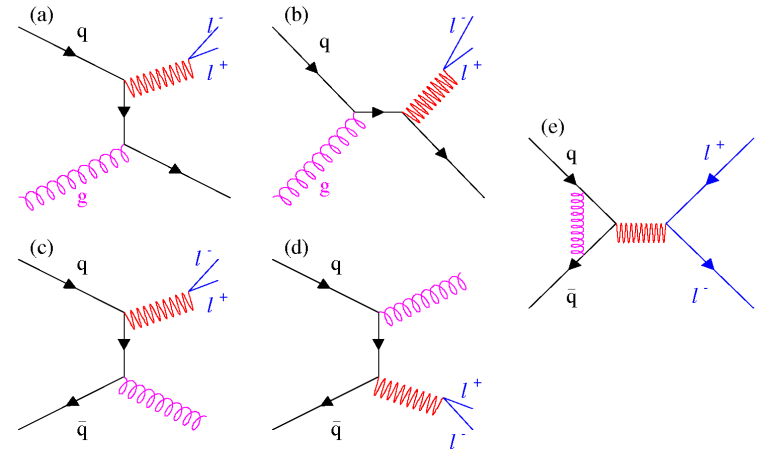
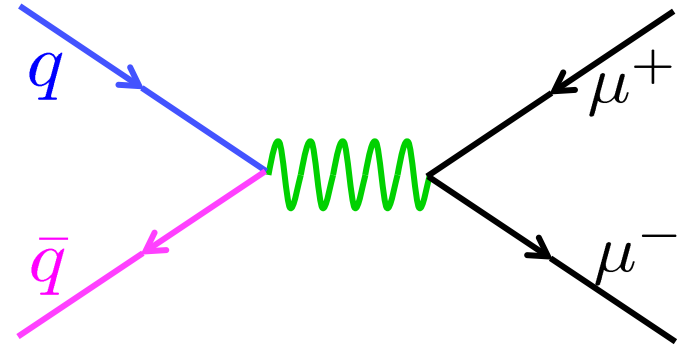
Drell-Yan Cross Section

- Measured cross section is a convolution of beam and target parton distributions

$$\frac{d^2\sigma}{dx_b dx_t} = \frac{4\pi\alpha^2}{x_b x_t s} \sum_{q \in \{u, d, s, \dots\}} e_q^2 [\bar{q}_t(x_t) q_b(x_b) + \bar{q}_b(x_b) q_t(x_t)]$$

- u-quark dominance
 - $(2/3)^2$ vs. $(1/3)^2$

- Next-to-leading order diagrams complicate the picture and must be considered
- These diagrams are responsible for up to **50% of the measured cross section**
- Intrinsic transverse momentum of quarks (although a small effect, $\lambda > 0.8$)
- Soft gluon resummation at all orders**



Drell-Yan Cross Section

Measured cross section is a convolution of beam and target parton distributions

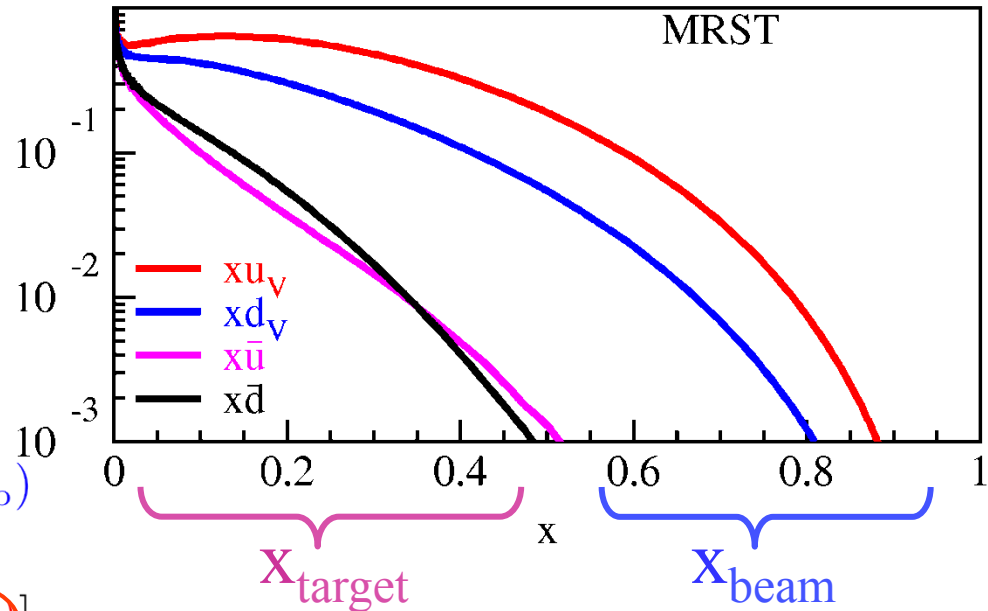
Proton Beam

– Target antiquarks and beam

$$\frac{d^2\sigma}{dx_b dx_t} = \frac{4\pi\alpha^2}{x_b x_t s} \sum_{q \in \{u, d, s, \dots\}} e_q^2 [\bar{q}_t(x_t) q_b(x_b) + \cancel{q_b(x_b) \bar{q}_t(x_t)}]$$

~~$q_b(x_b) \bar{q}_t(x_t)$~~

Acceptance limited



π beam

– Valence beam anti-u quark and u target quark

$$\left. \frac{d^2\sigma}{dx_\pi dx_N} \right|_{\pi^- N} = \frac{4\pi\alpha^2}{x_\pi x_N s} \left[\begin{aligned} & \frac{4}{9} \bar{u}_\pi(x_\pi) u_N(x_N) \\ & + \frac{1}{9} d_\pi(x_\pi) \bar{d}_N(x_N) \\ & + \frac{4}{9} u_\pi(x_\pi) \bar{u}_N(x_N) \\ & + \frac{1}{9} \bar{d}_\pi(x_\pi) d_N(x_N) \end{aligned} \right]$$

Valence × Valence → $\frac{4}{9} \bar{u}_\pi(x_\pi) u_N(x_N)$

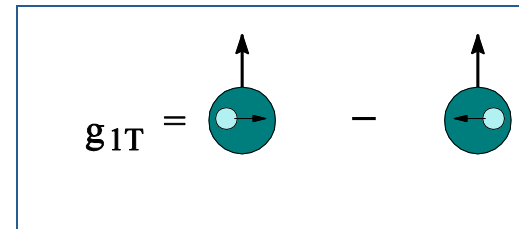
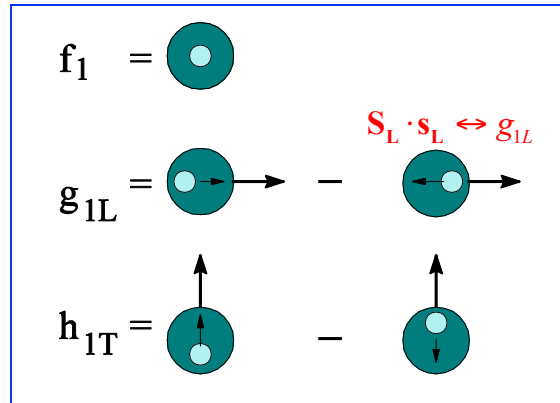
Valence-sea × 1/4 → $\frac{1}{9} d_\pi(x_\pi) \bar{d}_N(x_N)$

Sea-Sea → $\frac{1}{9} \bar{d}_\pi(x_\pi) d_N(x_N)$

| Beam | Target | Experiment |
|-------------|---|---|
| Hadron | Beam valence quarks target antiquarks | Fermilab E-906, RHIC (forward acpt.) J-PARC |
| Anti-Hadron | Beam val. antiquarks Target valence quarks | GSI-FAIR Fermilab Collider |
| Meson | Beam val. antiquarks Target valence quarks | COMPASS |

Transverse Momentum Distributions: Introduction

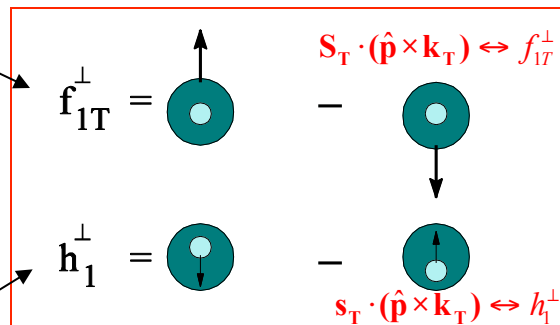
Survive k_T integration



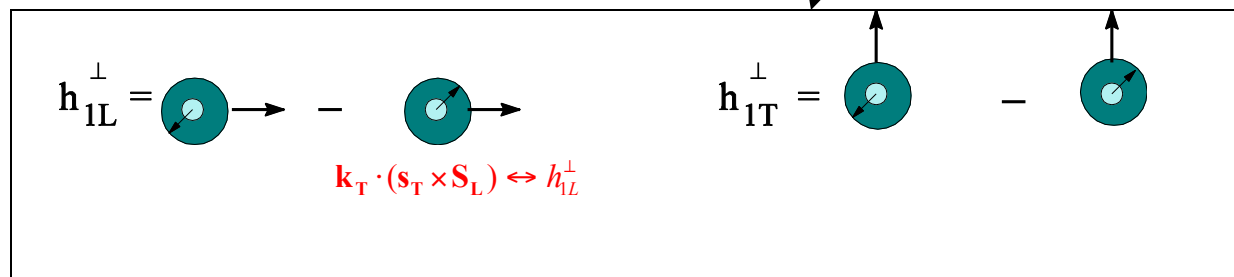
k_T - dependent, T-even

Sivers Function

k_T - dependent, Naive T-odd



Boer-Mulders Function



Transverse Momentum Distributions: Polarized Drell-Yan

1. What is the Lam-Tung relation? How does it relate to transverse structure?

- Boer-Mulders Function

$$h_1^\perp = \text{[Diagram: circle with dot and arrow pointing down]} - \text{[Diagram: circle with dot and arrow pointing up]}$$

$s_T \cdot (\hat{p} \times \mathbf{k}_T) \leftrightarrow h_1^\perp$

k_T - dependent,
Naïve T-odd

2. What do the available data teach us?

- πN (CERN NA10, Fermilab E-615)
- pp and pd (Fermilab E866)

3. Is it consistent?



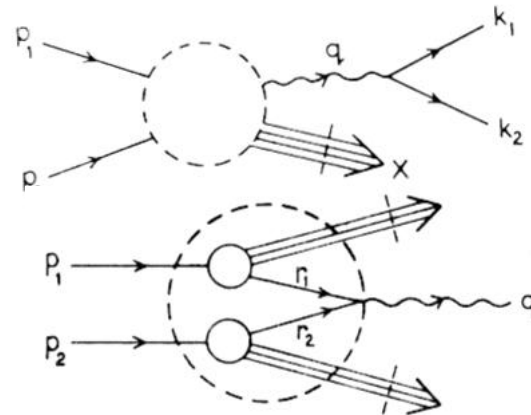
Generalized Angular Distributions & Lam Tung relation

Chi-Sing Lam and Wu-Ki Tung—basic formula for lepton pair production angular distributions PRD 18 2447 (1978)

$$\frac{d\sigma}{d^4q d\Omega_k^*} = \frac{1}{2} \frac{1}{(2\pi)^4} \left(\frac{\alpha}{M_S}\right)^2 \left[W_T (1 + \cos^2 \theta) + W_L (1 - \cos^2 \theta) \right. \\ \left. + W_{\Delta} \sin 2\theta \cos \phi + W_{\Delta\Delta} \sin^2 \theta \cos 2\phi \right]$$

Structure function formalism

- Derived in analogy to DIS
- Independent of Drell-Yan and parton “models”
- Showed same relations follow as a general consequence of the quark-parton model



Lam-Tung relation

- Derived in analogy to Colin-Gross relation of DIS

$$W_T = 2W_{\Delta\Delta}$$

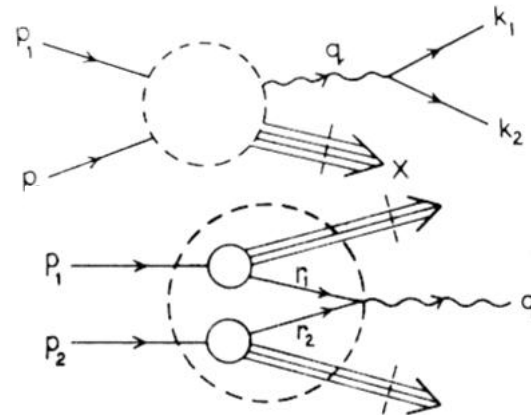
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$$\frac{d\sigma}{d\Omega} \propto 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi$$

Structure function formalism

- Derived in analogy to DIS
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Lam-Tung relation

- Derived in analogy to Colin-Gross relation of DIS

$$1 - \lambda = 2\nu$$

Why is Lam-Tung relation important? What does it have to do with Transverse Momentum?

■ Lam-Tug Relation is theoretically robust?

- Unaffected by $O(\alpha_s)$ (NLO) corrections
- NNLO [$O(\alpha_s^2)$] corrections also small Mirkes and Ohnemus, PRD 51 4891 (1995)
- Soft Gluon Resummation–Berger, Qiu and Rodrigues-Pedraza showed that **the Lam-Tung relation is preserved** under resummation. arXiv:0707.3150, and PRD 76 074006 (2007)

■ Factorization breaking QCD Vacuum

■ k_T dependent transverse momentum distribution (Boer Mulders h_1^\perp)

■ Alternatives:

- Nuclear effects
- Higher-Twist effects from quark-antiquark binding in pion

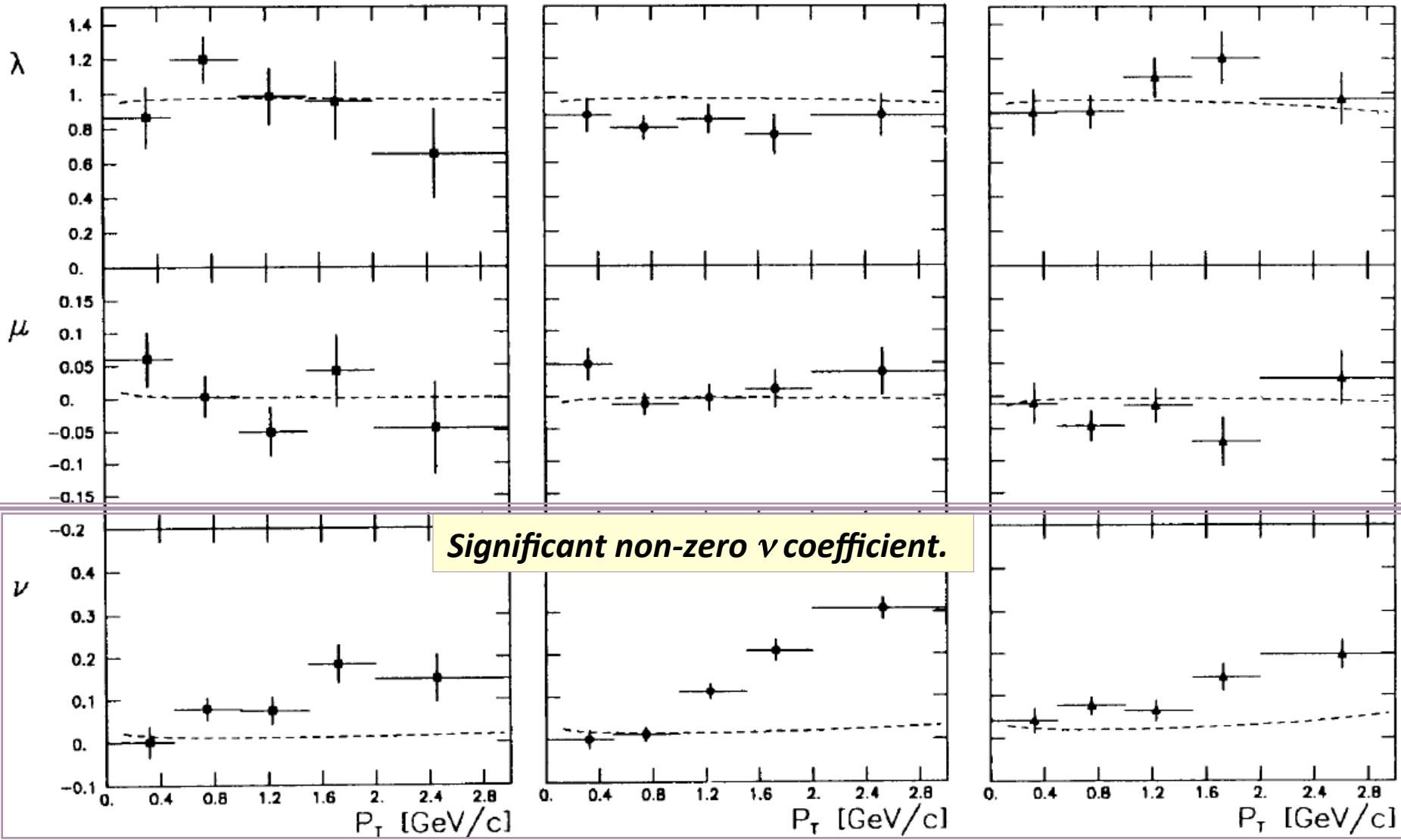
- Neither seen in data

NA10 angular distributions vs. p_T

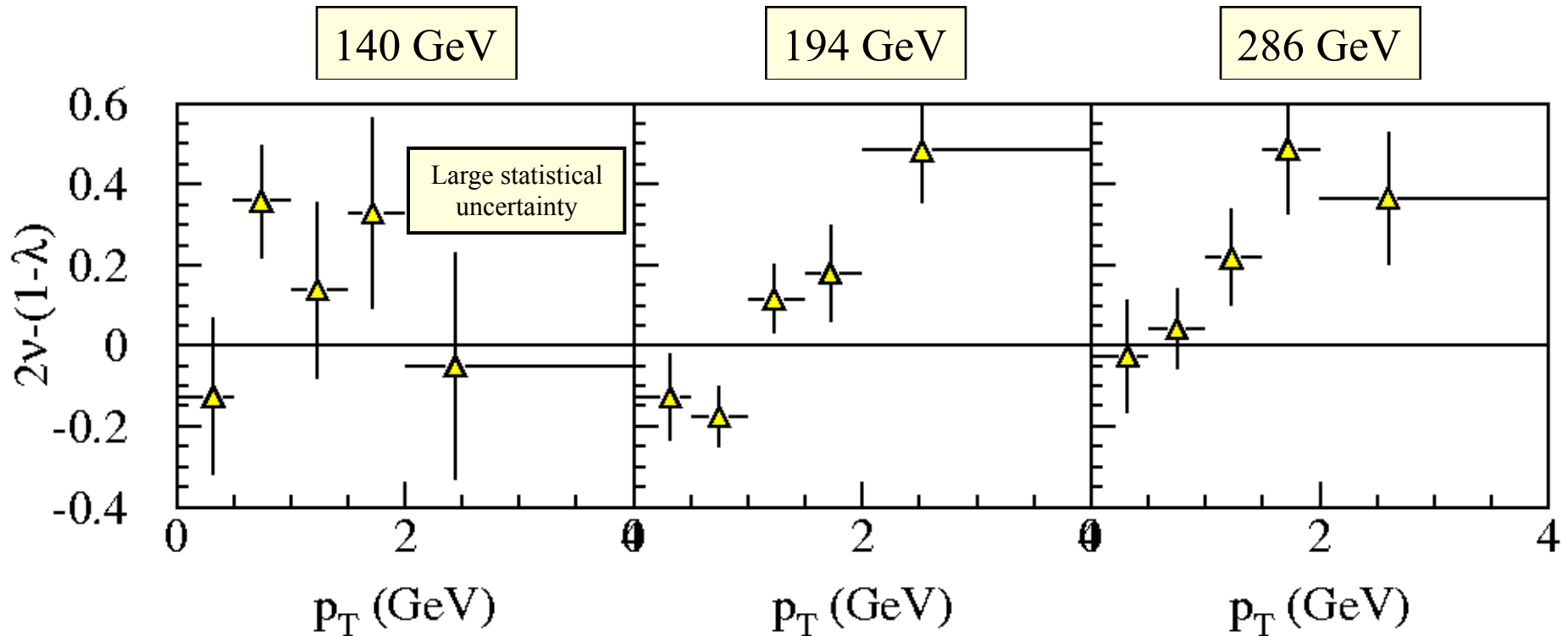
140 GeV/c

194 GeV/c

286 GeV/c



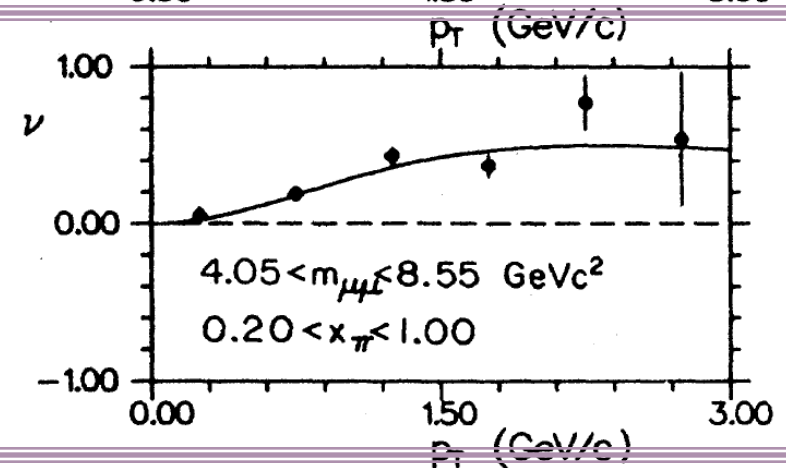
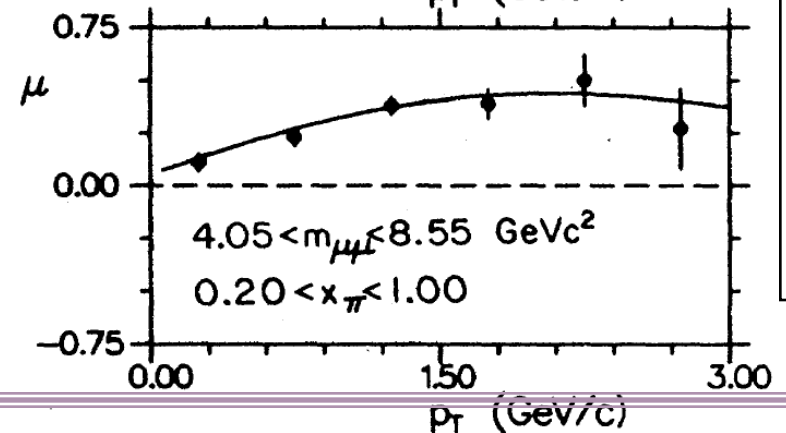
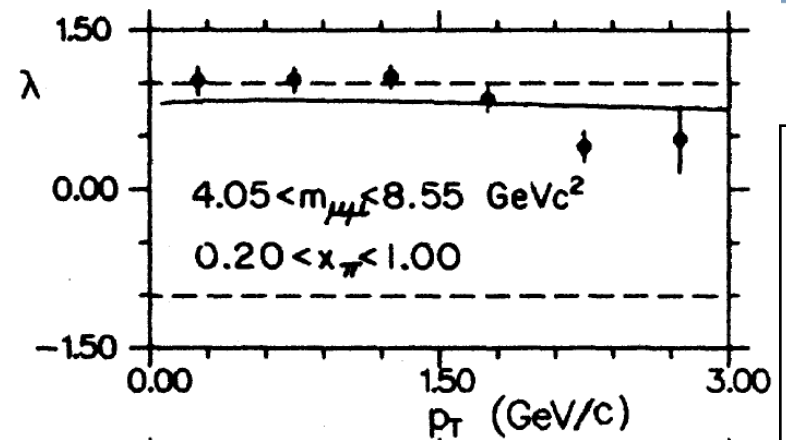
NA10 Lam-Tung Relation vs. p_T



- Violation of Lam-Tung relation as p_T increases in higher momentum data. Statistics poor in 140 GeV data.
- Note: Correlation between λ and v uncertainties not considered.
- Since most data is at low p_T , *on average* the Lam-Tung relationship holds

Pionic Data Fermilab E615

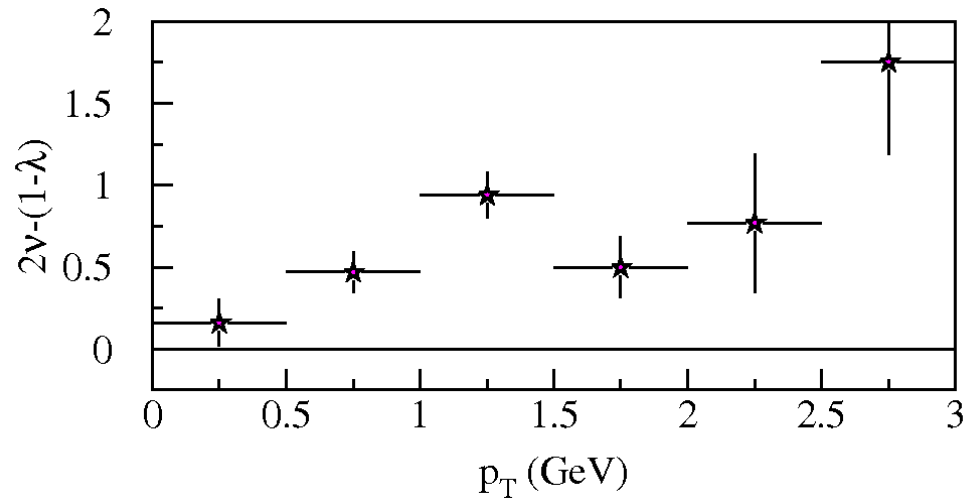
- Significant non-zero ν coefficient
- Shows other kinematic dependencies



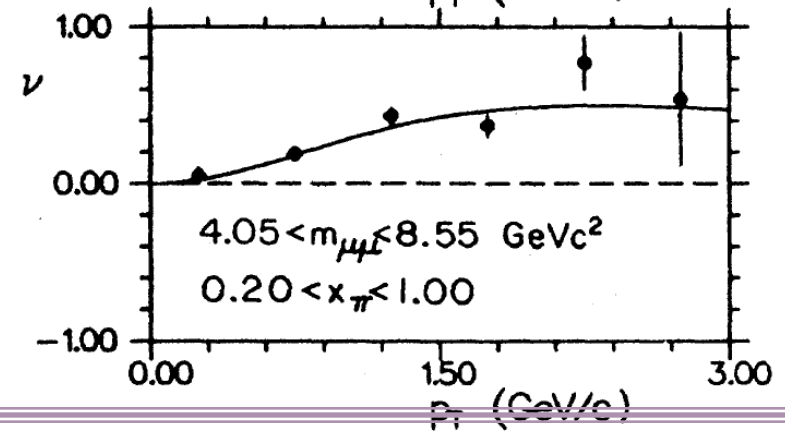
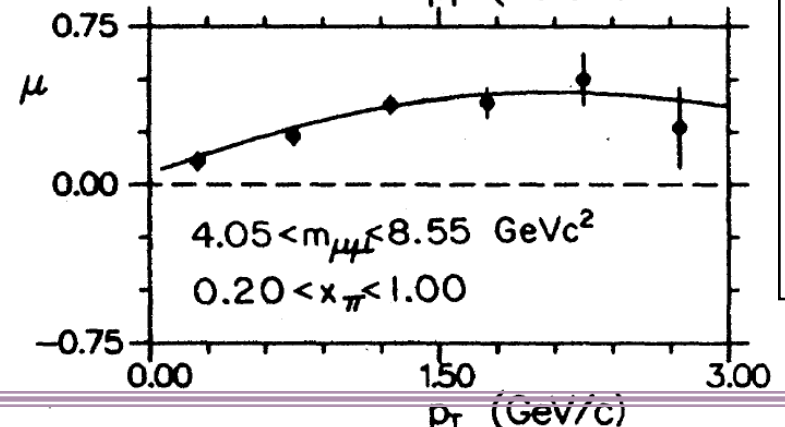
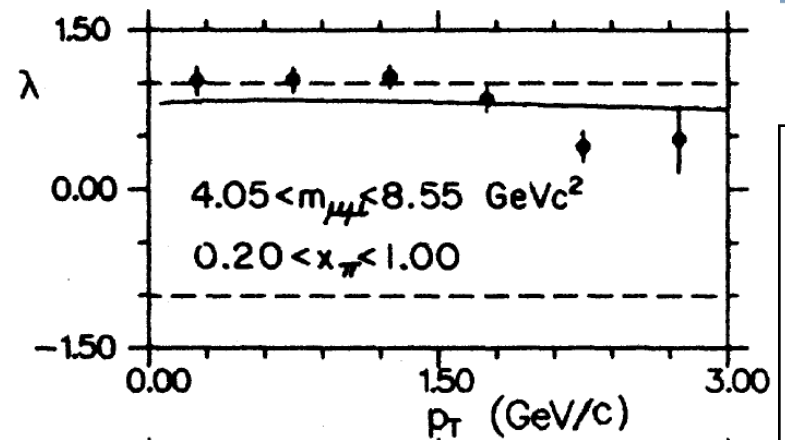
Conway et al. PRD 39 92 (1989)

Pionic Data Fermilab E615

- Significant non-zero ν coefficient
- Shows other kinematic dependencies



- Clear violation of Lam-Tung Relation vs. p_T .
- Violation larger than NA10



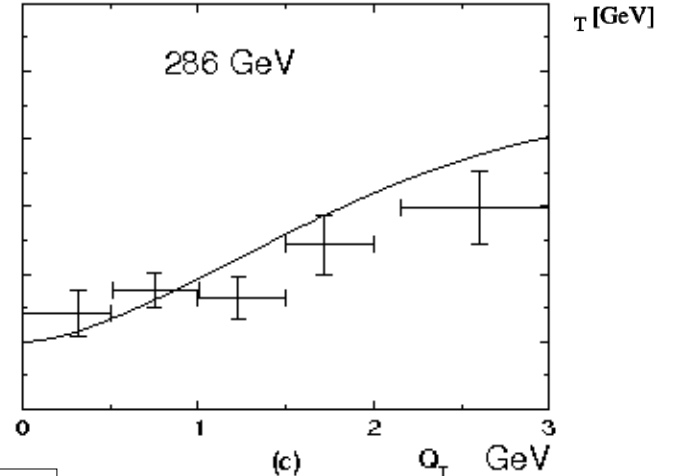
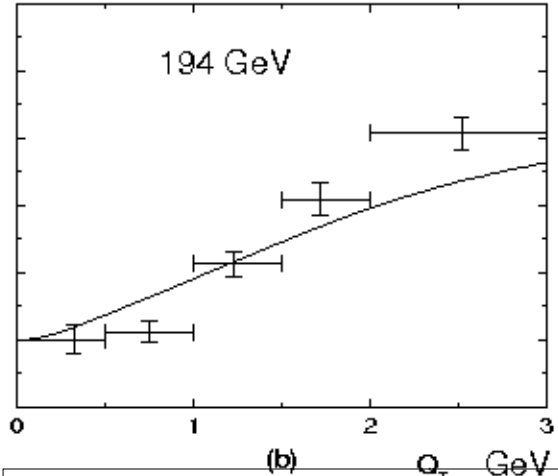
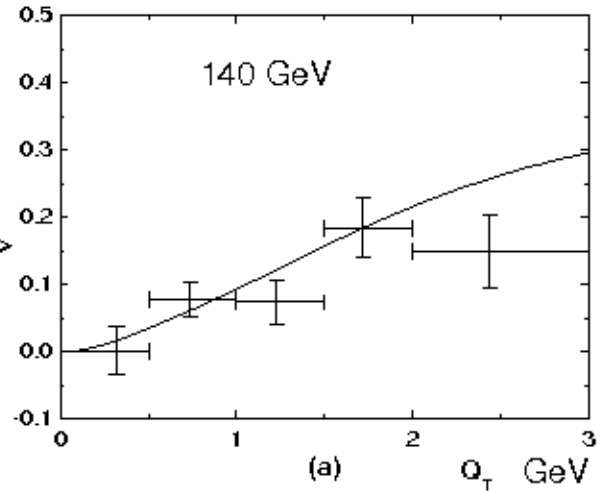
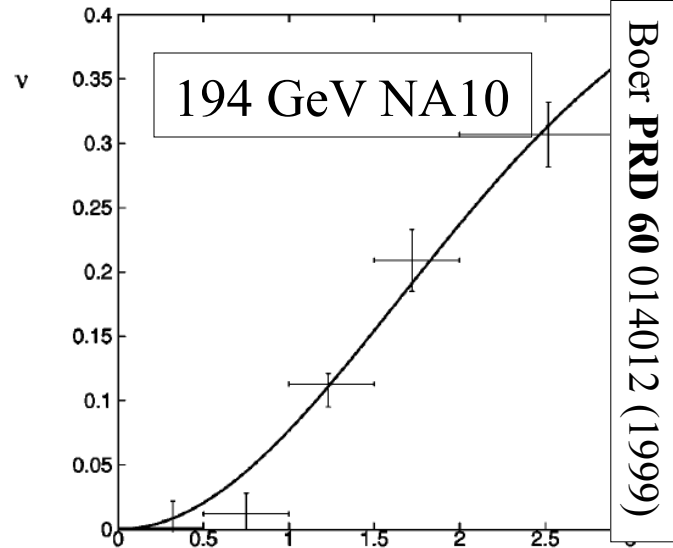
Conway et al. PRD 39 92 (1989)

Boer-Mulders Structure Function

- Relates transverse spin and transverse momentum (k_T) in an unpolarized nucleon.
- Presence in both quark and antiquark in annihilation could form correlation contributing to $\cos(2\phi)$ distribution

$$\nu \propto h_{1q}^\perp(x_1) h_{1\bar{q}}^\perp(x_2)$$

- Reasonable fits to pionic data



Lu, Ma, PLB 615, 200 (2005)

Lam-Tung Relation—Alternative View: QCD effects

- Factorization breaking Brandenburg, Nachtmann and Mirkes, ZPC **60**, 679 (1993).
 - QCD Vacuum *may* correlate the spins and momenta of incoming partons
 - Effect could be instanton-induced Boer, Brandenburg, Nachtmann, Utermann, EPJC 40 55 (2005), Brandenburg, Ringwald, Utermann NPB 754, 107 (2006).

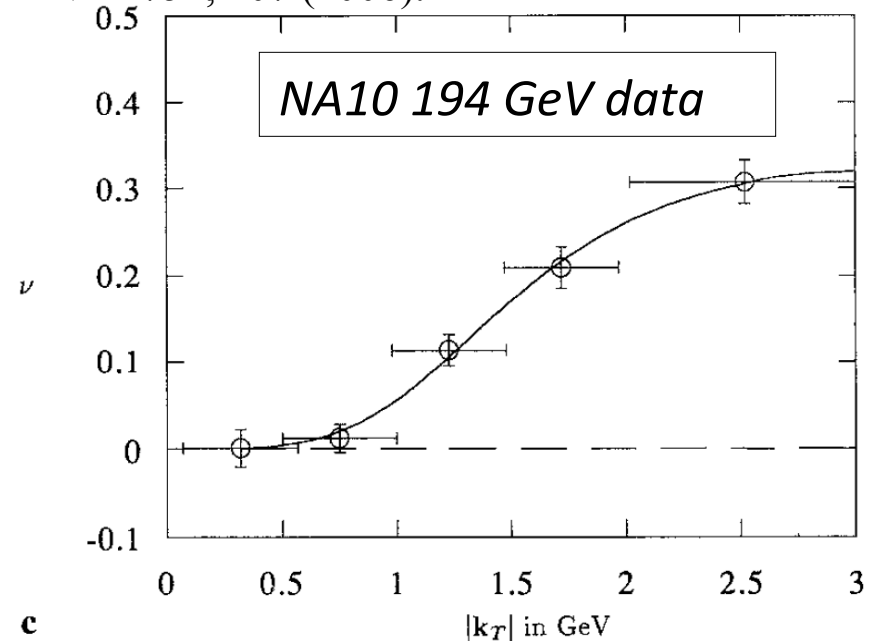
$$\nu \approx 2\mathcal{K} = 2\mathcal{K}_0 \frac{p_T^4}{p_T^4 + m_T^4}$$

$$\lambda \approx 1 \quad \mu \approx 0$$

- Fit NA10:

$$\kappa = 0.17$$

$$m_T = 1.5$$



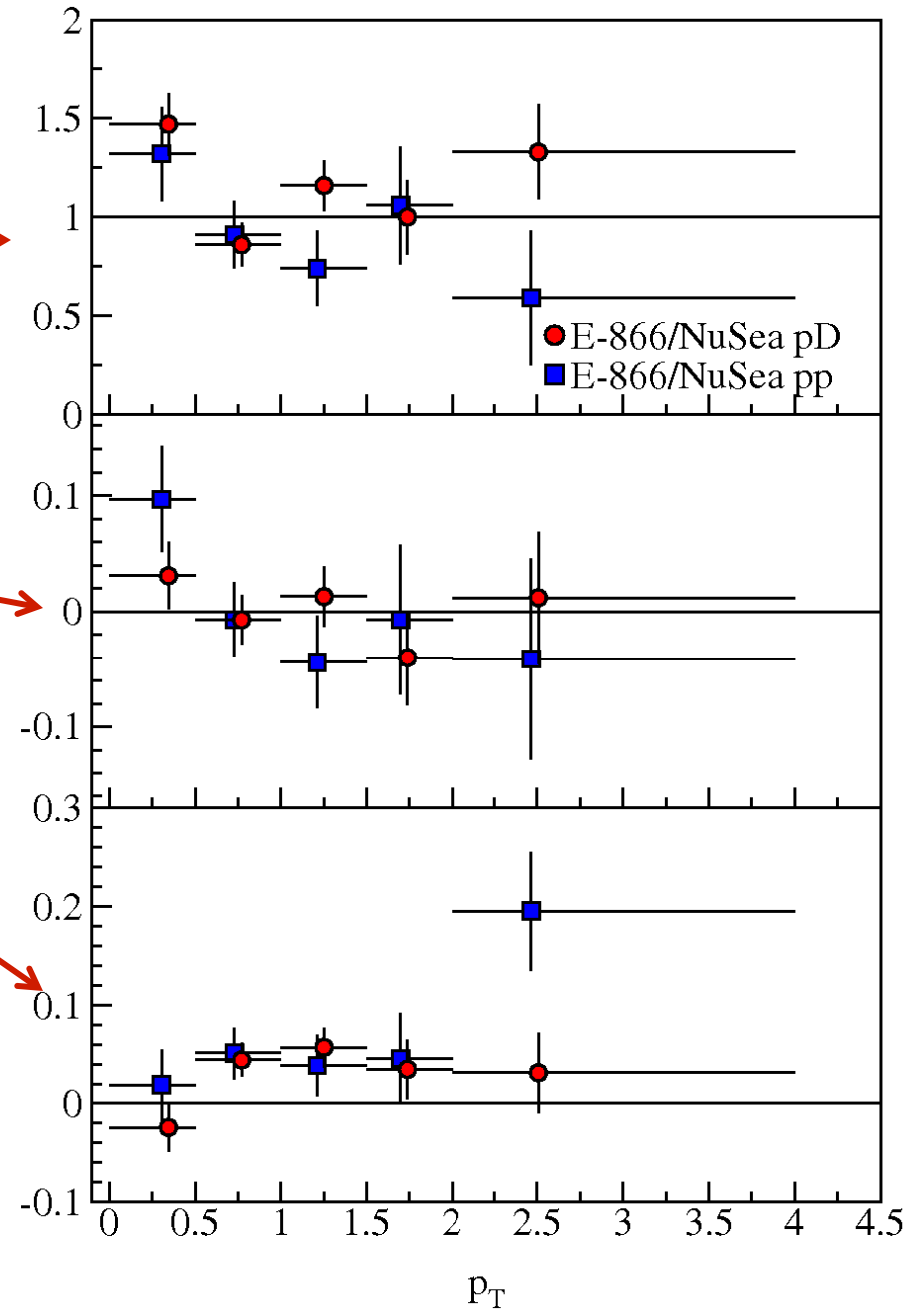
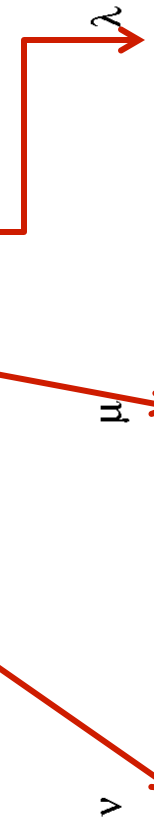
Should be flavor blind and seen in both sea and valence distributions

Protonic Drell-Yan

E-866/NuSea: Angular distrib.

$$\frac{d\sigma}{d\Omega} \propto 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + (\nu/2) \sin^2 \theta \cos 2\phi$$

- λ consistent with 1
- μ consistent with 0
- ν consistent with 0 (or slightly positive)

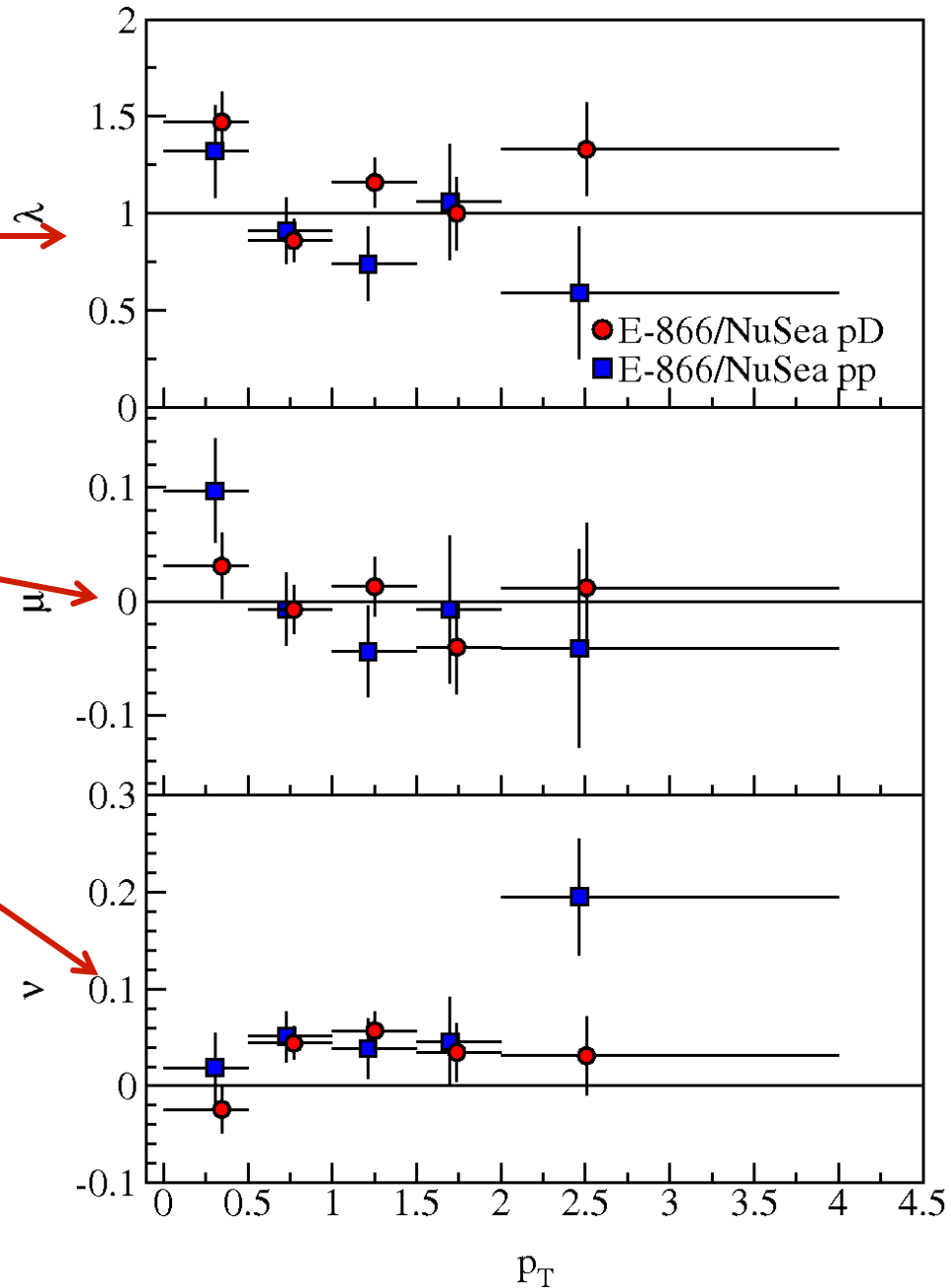
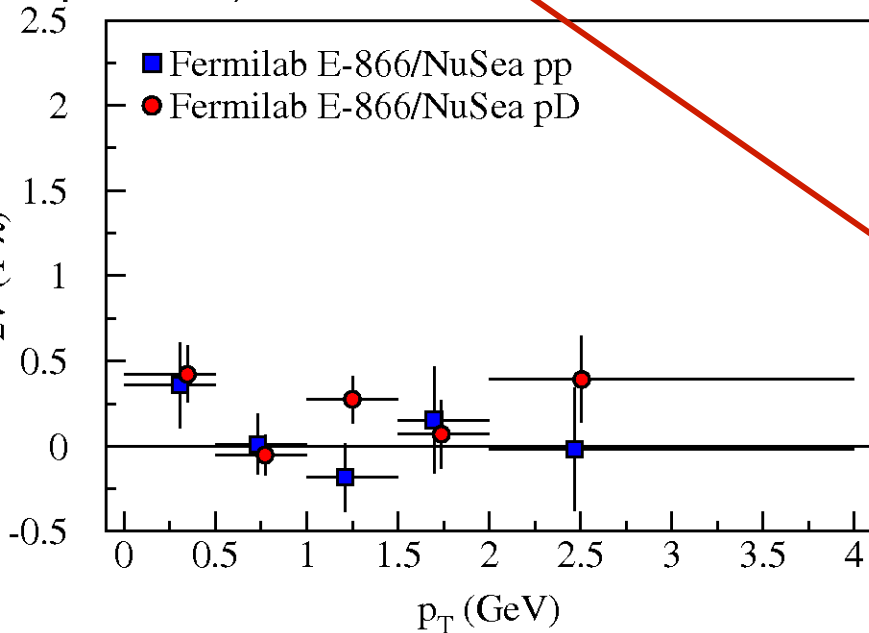


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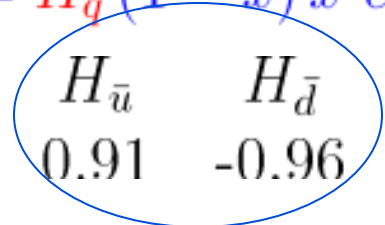


Extraction of Boer-Mulders function from pD Drell-Yan

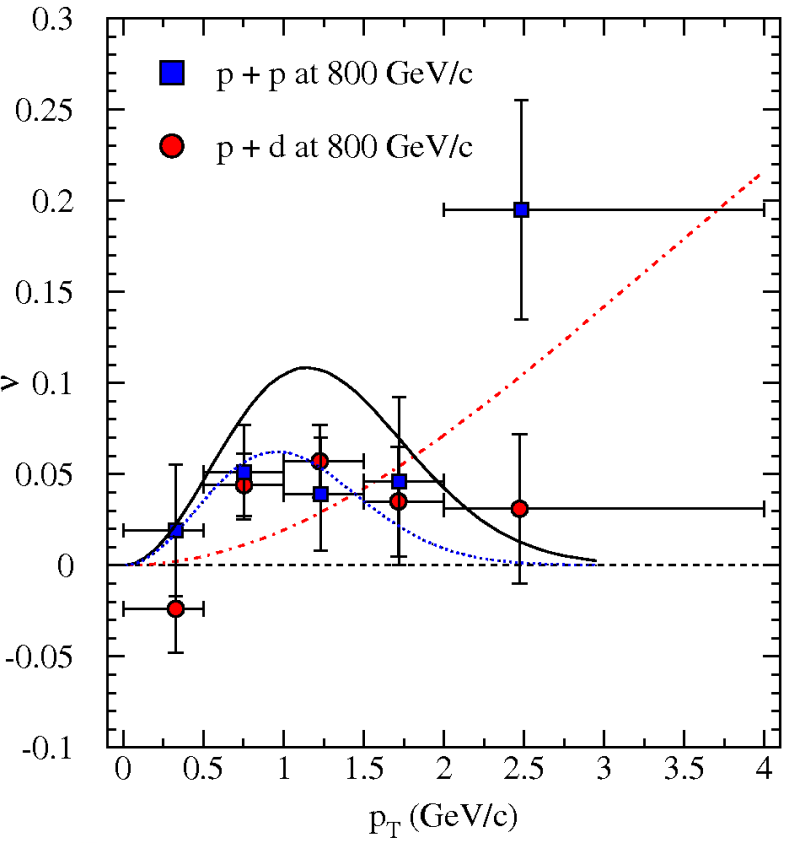
- Zhang, Lu, Ma, Schmidt, Phys.Rev.**D77**:054011,2008.
- Fit pD data in p_T , x_1 and x_2
- Extract $h_1^{\perp,q}$. (flavor separation)
- Prediction for pp Drell-Yan

$$h_1^{\perp,q}(x, p_{\perp}^2) = H_q (1-x) x^c e^{-p_{\perp}^2/p_{B-M}^2} f_1^q(x)$$

| | | | | | |
|-------|-------|---------------|---------------|------|-----------|
| H_u | H_d | $H_{\bar{u}}$ | $H_{\bar{d}}$ | c | p_{B-M} |
| 3.99 | 3.83 | 0.91 | -0.96 | 0.16 | 0.45 |



Small sea quark $h_1^?$



QCD effects in Drell-Yan

$\nu(pp) = \nu(pD)$ because of same kinematic coverage

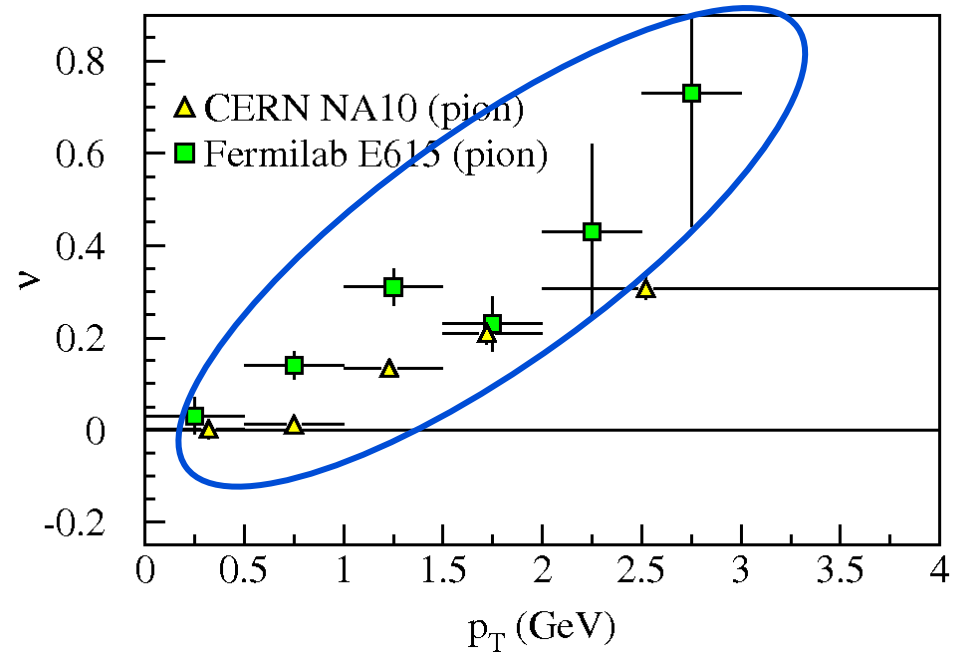
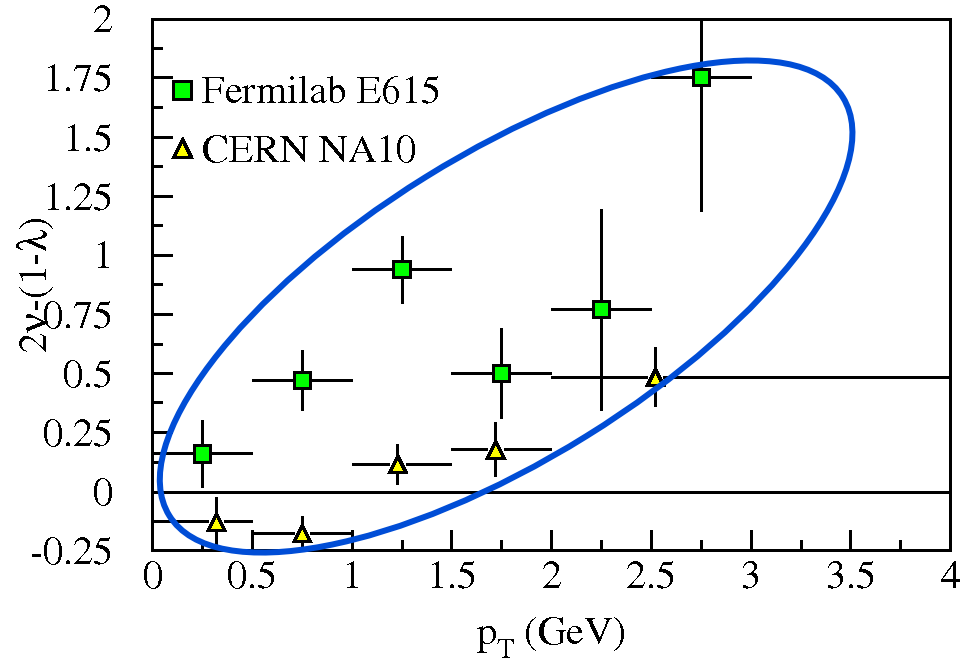
$$\nu = \frac{Q_{\perp}^2/Q^2}{1 + \frac{3}{2}Q_{\perp}^2/Q^2}$$

Boer *et al* EPJ **C 40**, 55 (2005) and
 Berger *et al* PR **D 76**, 074006 (2007).

Lam-Tung Relation

■ π^- Drell-Yan

- Violates L-T relation
- Large ν ($\cos 2\phi$) dependence
- Strong with p_T



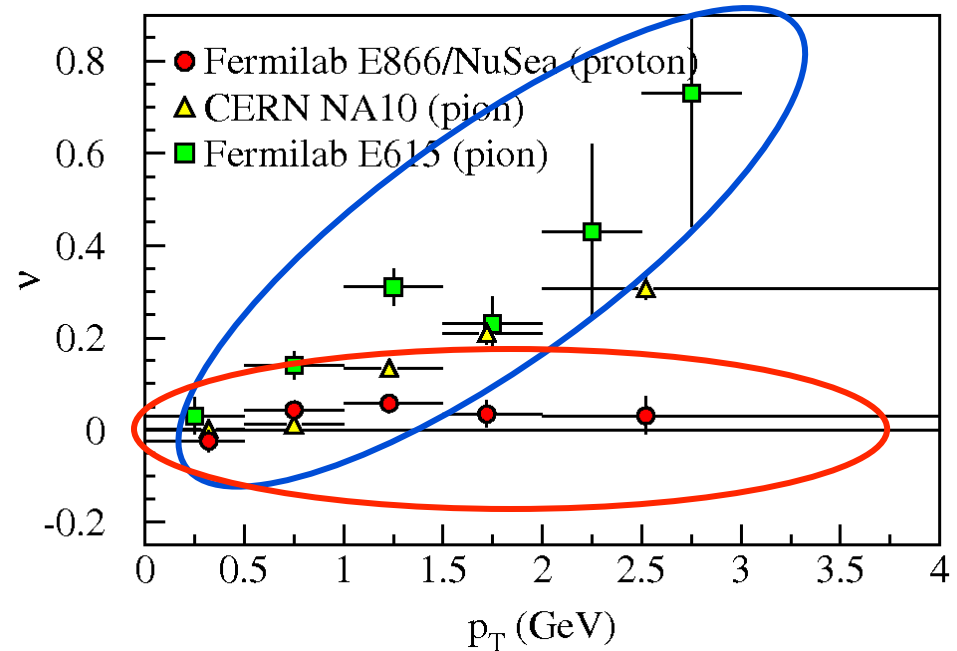
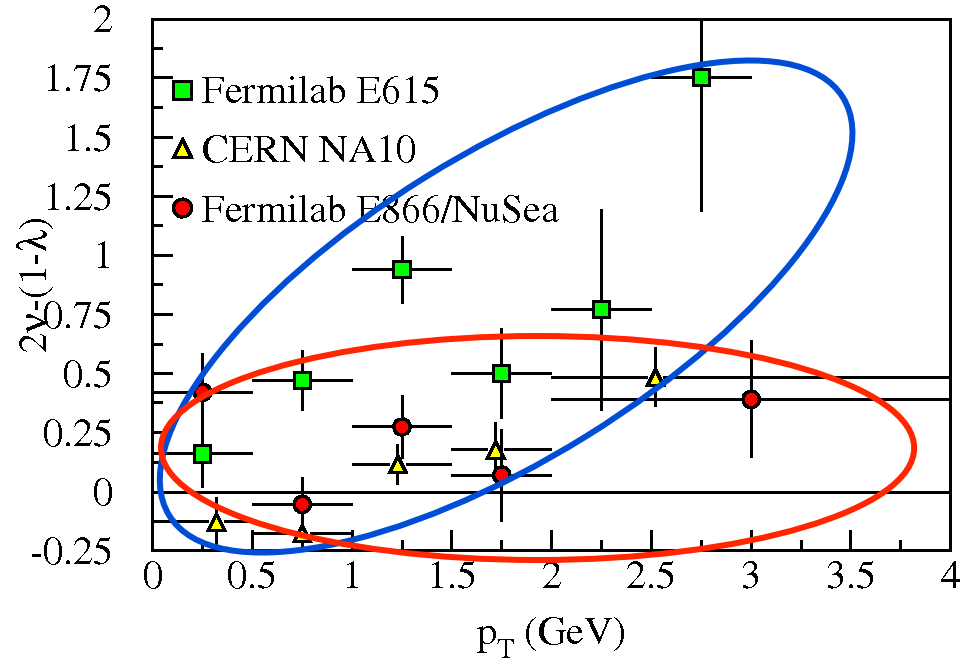
Lam-Tung Relation

■ π^- Drell-Yan

- Violates L-T relation
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- Strong with p_T

■ Proton Drell-Yan

- Consistent with L-T relation
- No ν ($\cos 2\phi$) dependence
- No p_T dependence



Lam-Tung Relation

■ π^- Drell-Yan

- Violates L-T relation
- Large v ($\cos 2\phi$) dependence
- Strong with p_T

■ Proton Drell-Yan

- Consistent with L-T relation
- No v ($\cos 2\phi$) dependence
- No p_T dependence

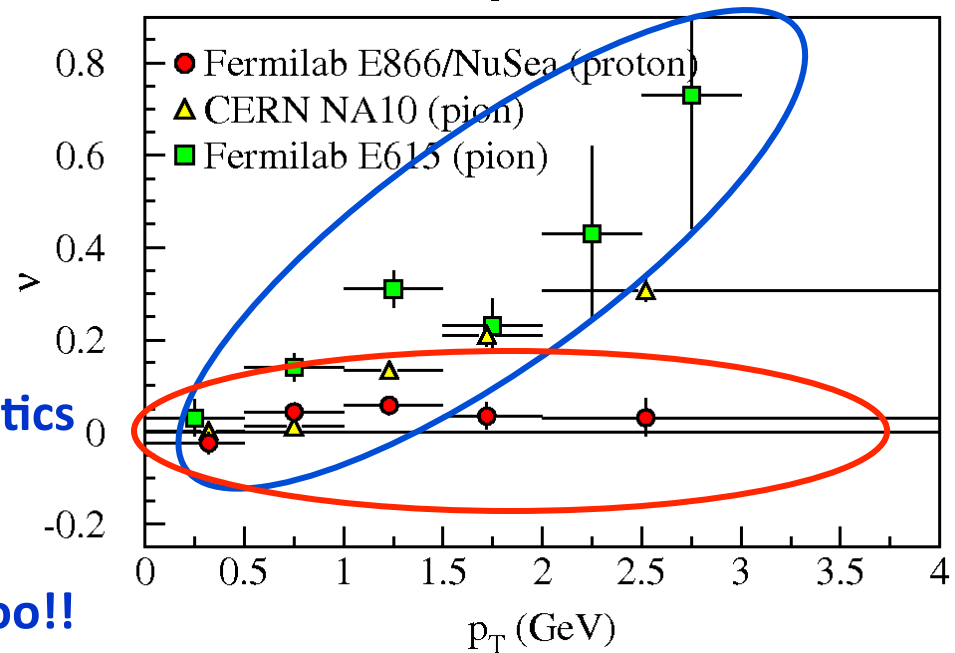
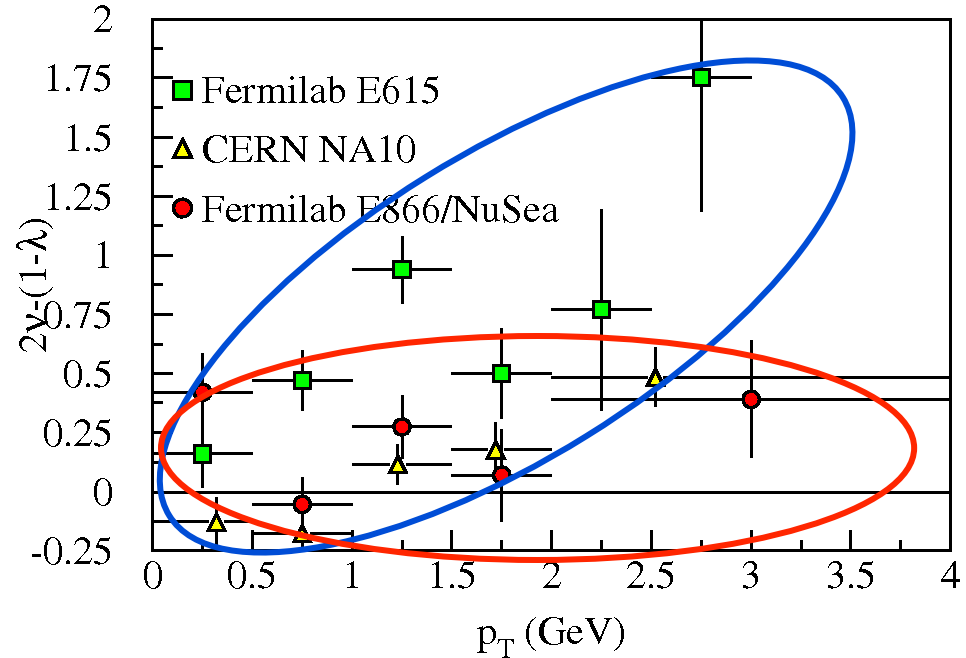
■ With Boer-Mulders function h_1^{\perp} :

- $v \pi^- W \rightarrow \mu^+ \mu^- X$
 $valence h_1^{\perp}(\pi) * valence h_1^{\perp}(p)$
- $v pd \rightarrow \mu^+ \mu^- X$
 $valence h_1^{\perp}(p) * sea h_1^{\perp}(p)$

■ Consistent story but need better statistics

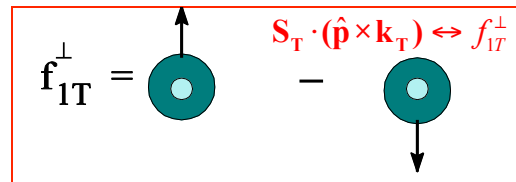
- $\pi^- p \rightarrow \mu^+ \mu^- X$ —CERN Compass
- $pd \rightarrow \mu^+ \mu^- X$ —Fermilab E906/SeaQuest

– Remember QCD effects are important too!!



Transverse Momentum Distributions: The Sivers function

1. What is the Sivers function?



k_T - dependent,
Naïve T-odd

2. What do the available Drell-Yan data teach us?
 - Nothing. . . . Yet
3. Is it consistent?
 - The big question



$$f_1^\perp(x, k_T) \Big|_{\text{DIS}} \stackrel{?}{=} -f_1^\perp(x, k_T) \Big|_{\text{D-Y}}$$

$$h_1^\perp(x, k_T) \Big|_{\text{DIS}} \stackrel{?}{=} -h_1^\perp(x, k_T) \Big|_{\text{D-Y}}$$

Single Spin Leading Order Drell-Yan Cross Section

$$\frac{d\sigma^{\text{LO}}}{d^4q d\Omega} = \frac{\alpha^2}{Fq^2} \hat{\sigma}_U^{\text{LO}} \left[1 + D_{\sin^2 \theta}^{\text{LO}} A_U^{\cos 2\phi} \cos 2\phi \right. \\ \left. + S_L D_{\sin^2 \theta}^{\text{LO}} A_L^{\sin 2\phi} \sin 2\phi \right. \\ \left. + \left| \vec{S}_T \right| A_T^{\sin \phi_S} \sin \phi_S \right. \\ \left. + \left| \vec{S}_T \right| D_{\sin^2 \theta}^{\text{LO}} A_T^{\sin(2\phi + \phi_S)} \sin(2\phi + \phi_S) \right. \\ \left. + \left| \vec{S}_T \right| D_{\sin^2 \theta}^{\text{LO}} A_T^{\sin(2\phi - \phi_S)} \sin(2\phi - \phi_S) \right]$$

Formula: from Aram Kotzinian

Slide: content from Oleg Denisov's talk

$A_U^{\cos 2\phi}$ Boer-Mulders of beam hadron

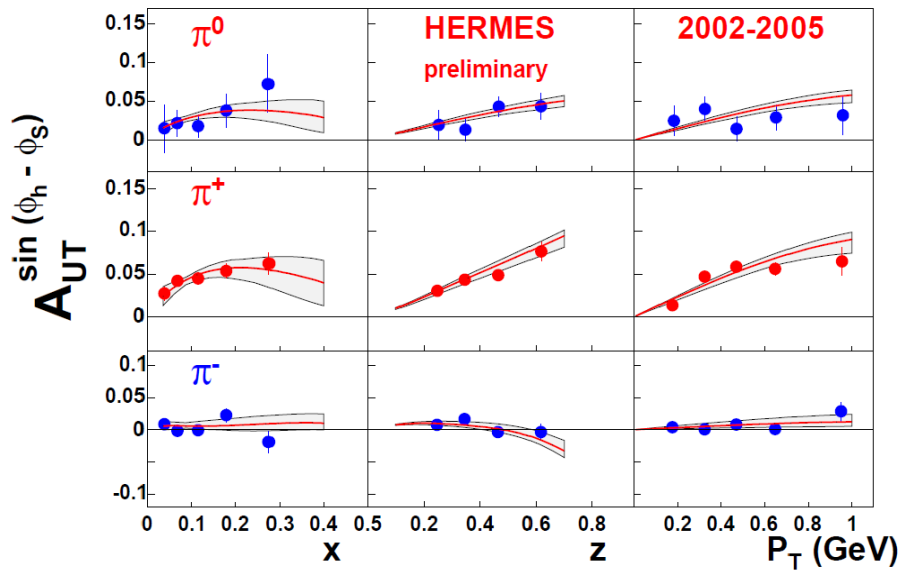
$A_T^{\sin \phi_S}$ Sivers for target nucleon

$A_T^{\sin(2\phi + \phi_S)}$ Boer-Mulders of beam hadron and h_1^\perp and pretzelosity of target nucleon

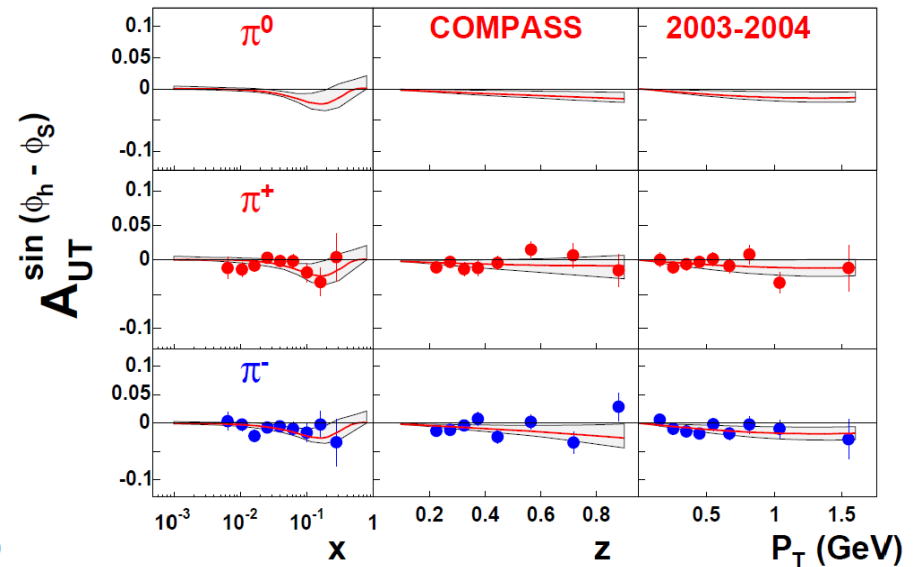
$A_T^{\sin(2\phi - \phi_S)}$ Boer-Mulders of beam hadron and h_1 and transversity of target nucleon

SIDIS Sivers measurements

HERMES

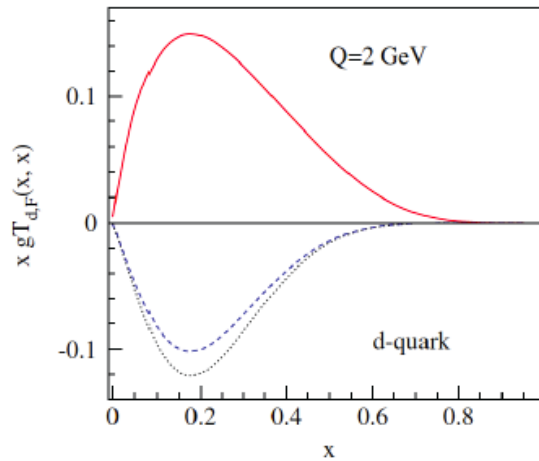
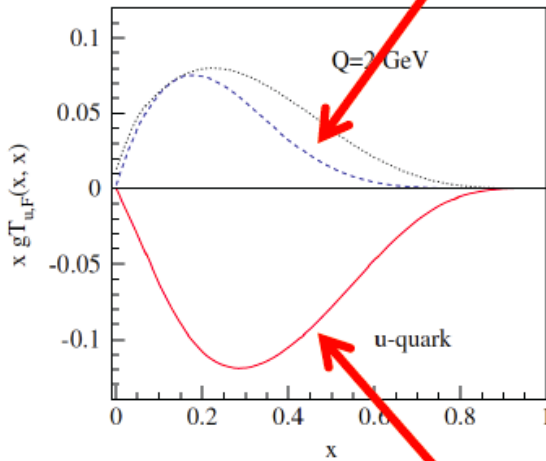


COMPASS



- Global fit to $\sin(\phi_h - \phi_S)$ asymmetry in SIDIS (**HERMES, COMPASS**)
- Comparable measurements needed for single spin asymmetries in Drell-Yan process
 - Caution: Nothing excludes a node
 - Must cover same kinematics with Drell-Yan and SIDIS

from $-\int \frac{d^2 \vec{k}_\perp}{2\pi} \frac{\vec{k}_\perp^2}{M^2} (f_{1T}^\perp(x, k_\perp))_{\text{DIS}}$



- basic observation:

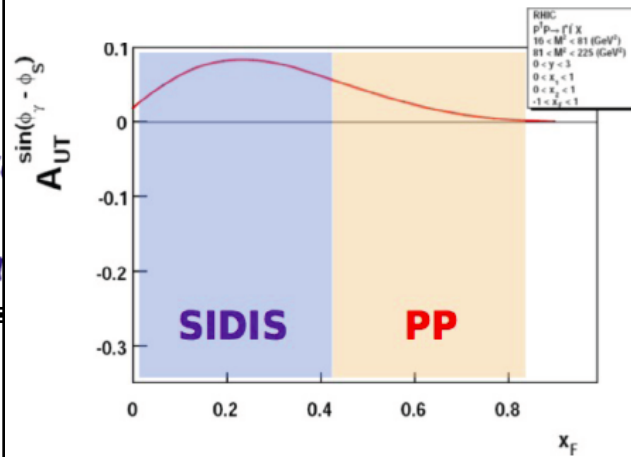
SIDIS π^+ : *final-state*

$pp \rightarrow \pi^+ X$: *initial-state*

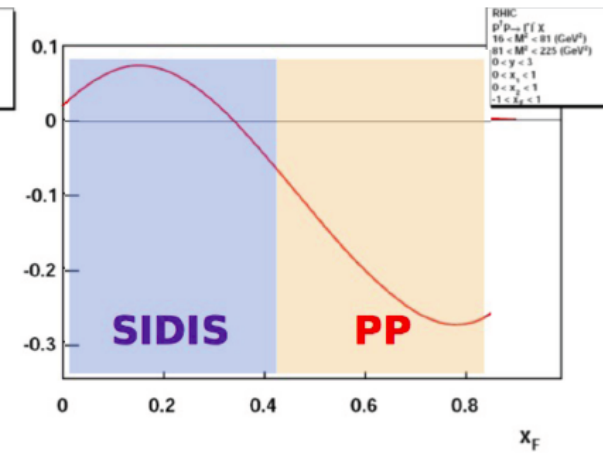
From V. Vogelsang's talk

Has ramifications for DY spin asymmetry:

Kang, Prokudin



w/o node



with node

Strengthens case for study of DY "sign change" !



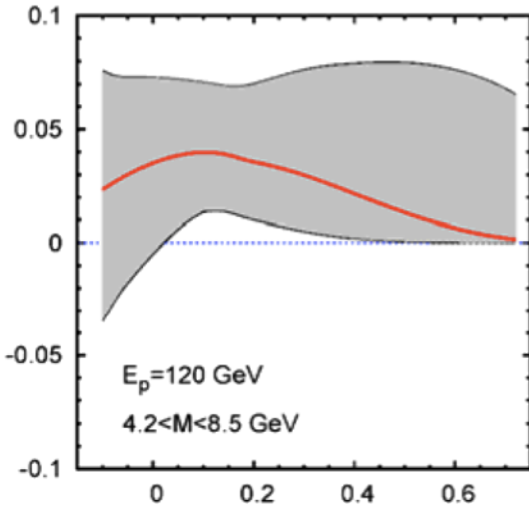
Future Polarized Drell-Yan Measurements

| experiment | particles | energy | x_1 or x_2 | luminosity | timeline |
|-------------------------------------|-------------------------------|---|---|---|-----------------|
| COMPASS (CERN) | $\pi^\pm + p^\uparrow$ | 160 GeV $\sqrt{s} = 17.4$ GeV | $x_2 = 0.2 - 0.3$ $x_2 \sim 0.05$ (low mass) | $2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$ | >2012 |
| PAX (GSI) | $p^\uparrow + p_{\text{par}}$ | collider $\sqrt{s} = 14$ GeV | $x_1 = 0.1 - 0.9$ | $2 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$ | >2017 |
| PANDA (GSI) | $p_{\text{par}} + p^\uparrow$ | 15 GeV $\sqrt{s} = 5.5$ GeV | $x_2 = 0.2 - 0.4$ | $2 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$ | >2016 |
| J-PARC | $p^\uparrow + p$ | 50 GeV $\sqrt{s} = 10$ GeV | $x_1 = 0.5 - 0.9$ | $1 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ | >2015 ?? |
| NICA (JINR) | $p^\uparrow + p$ | collider $\sqrt{s} = 20$ GeV | $x_1 = 0.1 - 0.8$ | $1 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$ | >2014 |
| PHENIX (RHIC) | $p^\uparrow + p$ | collider $\sqrt{s} = 500$ GeV | $x_1 = 0.05 - 0.1$ | $2 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$ | >2018 |
| RHIC internal target phase-1 | $p^\uparrow + p$ | 250 GeV $\sqrt{s} = 22$ GeV | $x_1 = 0.25 - 0.4$ | $2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$ | >2018 |
| RHIC internal target phase-1 | $p^\uparrow + p$ | 250 GeV $\sqrt{s} = 22$ GeV | $x_1 = 0.25 - 0.4$ | $6 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ | >2018 |
| A_n DY RHIC (IP-2) | $p^\uparrow + p$ | 500 GeV $\sqrt{s} = 32$ GeV | $x_1 = ??$ | $?? \text{ cm}^{-2} \text{ s}^{-1}$ | >2015 |
| SeaQuest (unpol.) (FNAL) | $p + p$ | 120 GeV $\sqrt{s} = 15$ GeV | $x_1 = 0.3 - 0.9$ | $3.4 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ | 2011 |
| pol. SeaQuest (FNAL) | $p^\uparrow + p$ | 120 GeV $\sqrt{s} = 15$ GeV | $x_1 = 0.3 - 0.9$ | $1 \times 10^{36} \text{ cm}^{-2} \text{ s}^{-1}$ | >2014 |

From W. Lorenzon's talk

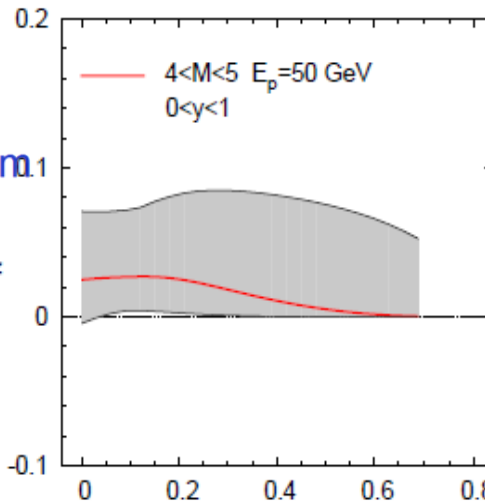
$A_N^{\sin(\phi_T - \phi_S)}$ prediction Anselmino *et al*

Fermilab: $p \uparrow p$



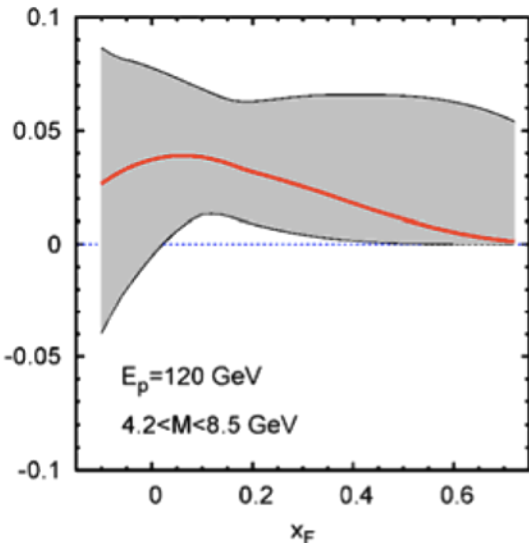
FNAL
120 GeV
polarized beam
 $\sqrt{s} \sim 15$ GeV
(hydrogen)

JPARC: $p \uparrow p$



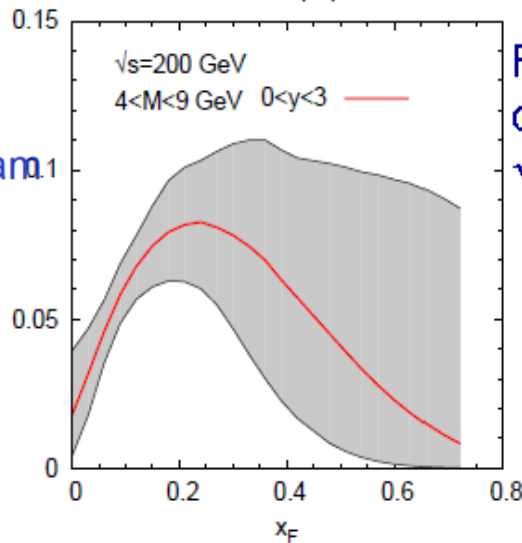
J-PARC
50-GeV
polarized beam
 $\sqrt{s} \sim 10$ GeV

Fermilab: $p \uparrow d$



FNAL
120 GeV
polarized beam
 $\sqrt{s} \sim 15$ GeV
(deuterium)

RHIC: $p \uparrow p$



RHIC
collider
 $\sqrt{s} = 200$ GeV

- $A_N \approx 0.04$
- Uncertainty dominated by lack of SIDIS data

Anselmino *et al.* priv. comm. 2010

Anselmino *et al.* PRD79, 054010 (2009)

Drell-Yan Measurements of the Sivvers Distribution

- Brookhaven/RHIC ANDY pp^{\rightarrow}
 - Possibility for both beams to be polarized?
 - See talk by Les Bland

- CERN Compass $\pi^- p^{\rightarrow}$
 - Polarized valence u-quark distributions at small-x
 - See talks by Oleg Denisov and by Oleg Teryaev Wolfgang Lorenzon

- Fermilab $p\bar{p}^{\rightarrow}$
 - Polarized Valence Distributions at large-x
 - See talk by Wolfgang Lorenzon

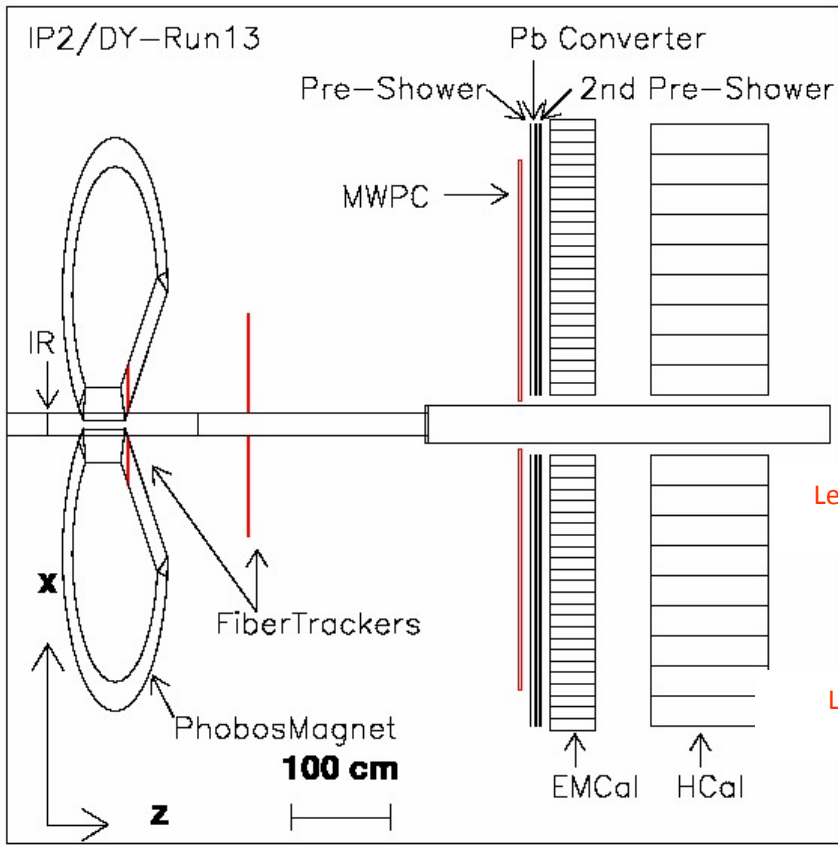
Remember 9(beginning of the journey) different beams and targets probe different distribs.

| Beam | Target | Experiment |
|-------------|--|---|
| Hadron | Beam valence quarks target antiquarks | Fermilab E-906, RHIC (forward acpt.) J-PARC |
| Anti-Hadron | Beam val. antiquarks Target val. quarks | GSF-FAIR Fermilab Collider |
| Meson | Beam val. antiquarks Target val. quarks | COMPASS |



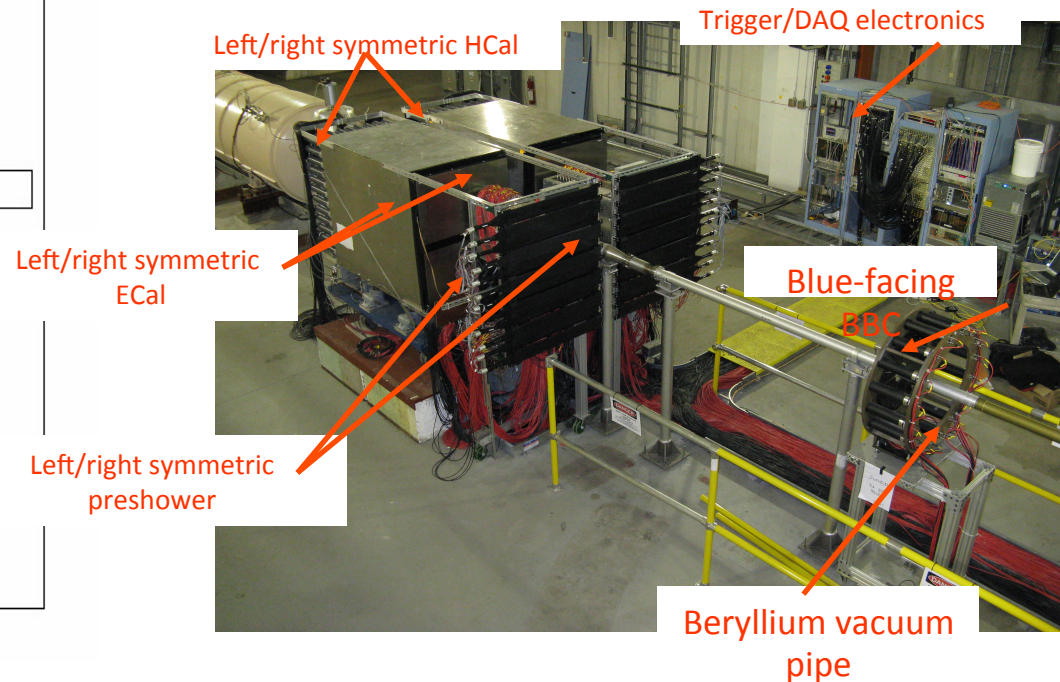
A_NDY@RHIC

- See talk by Les Bland
- Calorimetry required for experiment is at BNL. It is on loan from JLab until July 2014.



http://www.star.bnl.gov/~akio/ip2/topview_run13.jpeg

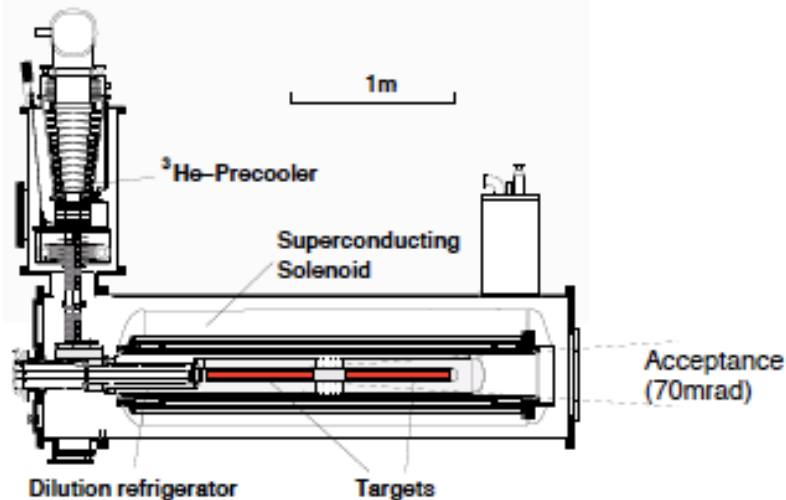
- Layout for second stage (no magnet) is complete, with specification of azimuthally complete HCal, ECal and segmented preshower detector.
- Existing split dipole does require modification to match acceptance of final calorimetry plan.
- Staging of apparatus awaits funding review.
- At present, first attempt at transverse spin DY measurement will be in RHIC run 13



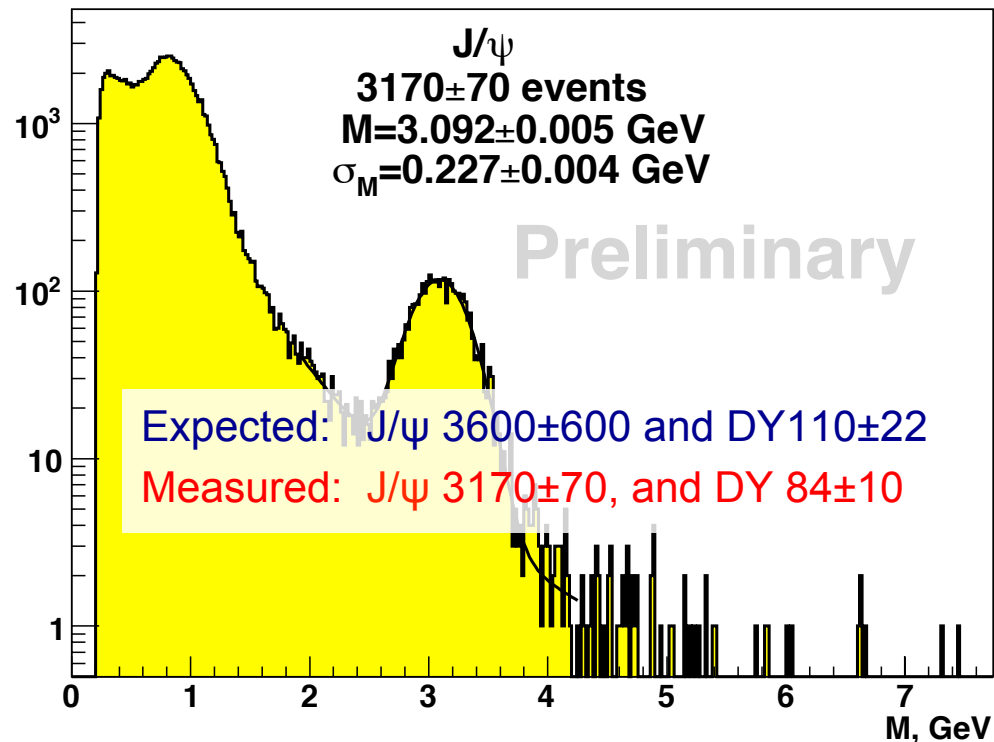
COMPASS@CERN

- See talk by Oleg Denisov
- Very well understood apparatus
 - A_N angular measurements already completed for SIDIS
 - MC and test data match
- Polarized target
- Data in 2014

Polarized target: ${}^6\text{LiD}$ or NH_3

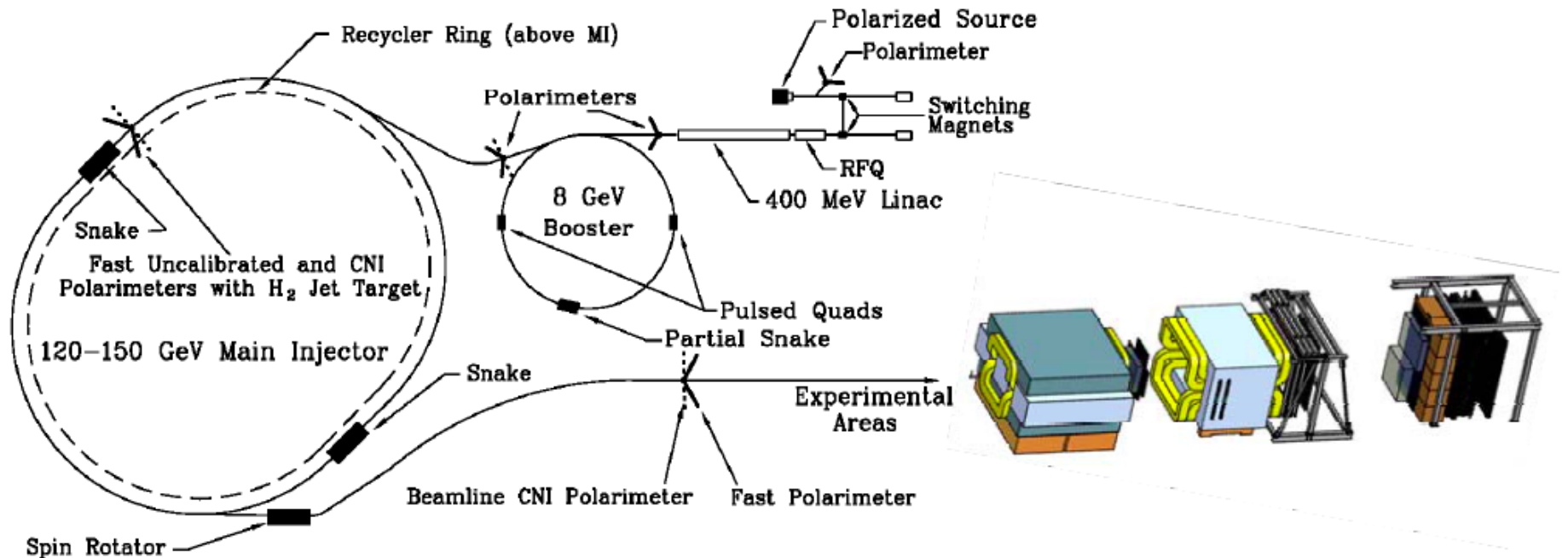


COMPASS DY beam test 2009



Fermilab Polarized Main Injector

- See talks by W. Lorenzon, A. Krisch, D. Sivers
- Polarized beam
 - Major advantage—the beam is a blow torch—Luminosity
 - Major disadvantage—the beam polarization is presently virtual—only a proposal
 - **Relatively inexpensive!!**
- By 2014, spectrometer will be well understood, including angular acceptance





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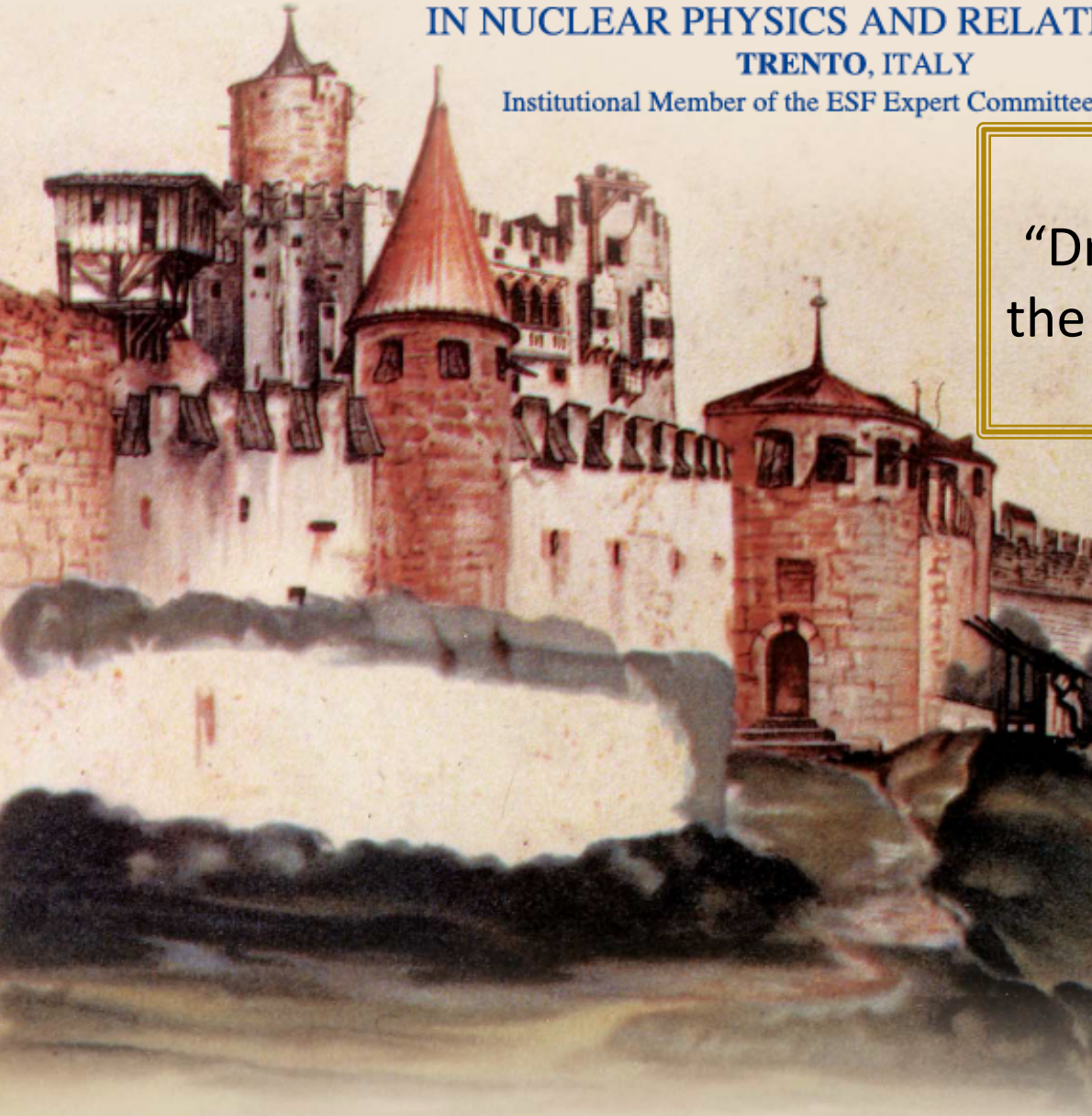


EUROPEAN CENTRE FOR THEORETICAL STUDIES
IN NUCLEAR PHYSICS AND RELATED AREAS

TRENTO, ITALY

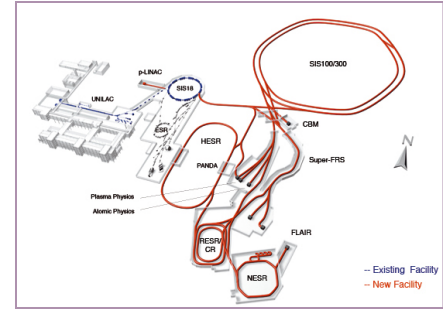
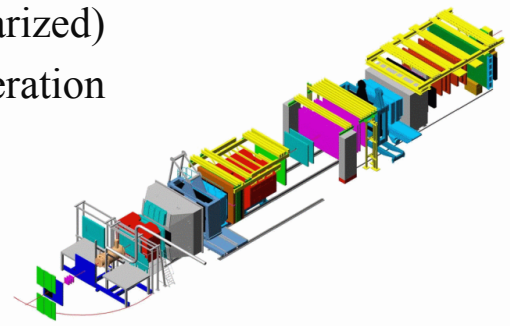
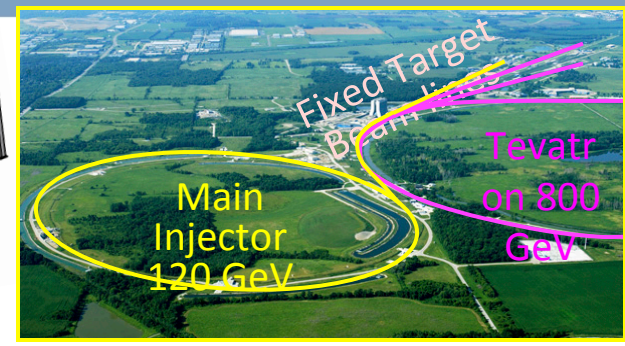
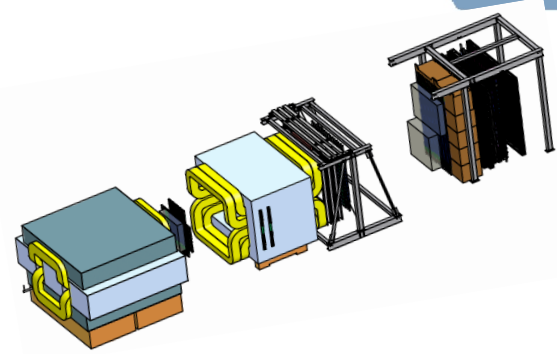
Institutional Member of the ESF Expert Committee NuPECC

ECT* workshop on
“Drell-Yan Scattering and
the Structure of Hadrons”
21-25 May 2012

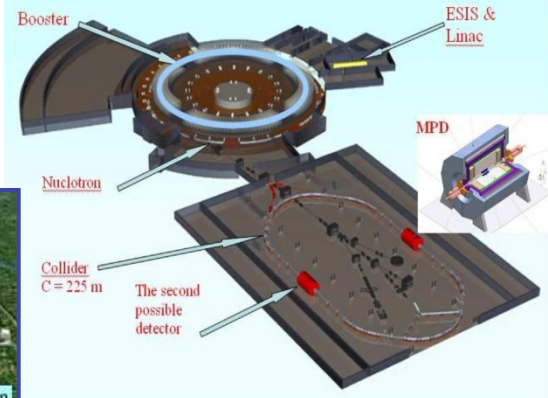
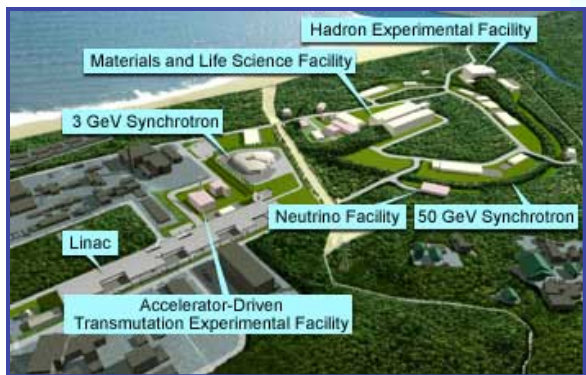


Future Drell-Yan Experiments

- Fermilab E-906/Drell-Yan
 - Better statistical precision (unpolarized)
 - Polarized extension under consideration
- COMPASS
 - Pion beam—valence distributions
- GSI FAIR—PAX experiment
 - Antiproton beam will sample valence distributions of targets



- JINR Dubna-NICA
- J-PARC
- RHIC



It is a long trip, but in the end, it's worth it!

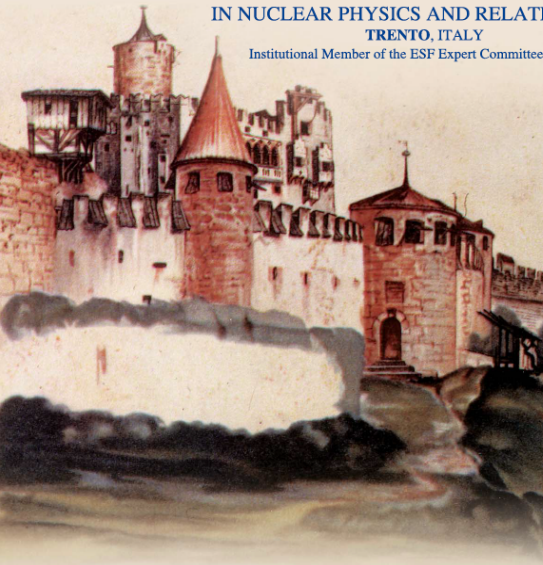


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Castello di Trento ("Tirolo", Wagnersaal, 13.8 x 27.2, painted by A. Diener on his way back from Venice (1495)

Additional Material

Boer-Mulders Structure Function

- Relates parton's transverse spin and transverse momentum (k_T) in an unpolarized nucleon.
- Presence in both quark and antiquark in annihilation could form correlation contributing to $\cos(2\phi)$ distribution

$$\nu \propto h_{1q}^\perp(x_1) h_{1\bar{q}}^\perp(x_2)$$

$$h_1^\perp(x, k_T^2) = C_H \frac{\alpha_T}{\pi} \frac{1}{k_T^2 + M_C^2} e^{-\alpha_T k_T^2} f_1(x)$$

$$\nu = 16C_1C_2 \frac{p_T^2 M_C^2}{(p_T^2 + 4M_C^2)^2}$$

$$M_C = 1.2 \pm 0.5 \text{ GeV}$$

$$16C_1C_2 = 7 \pm 2$$

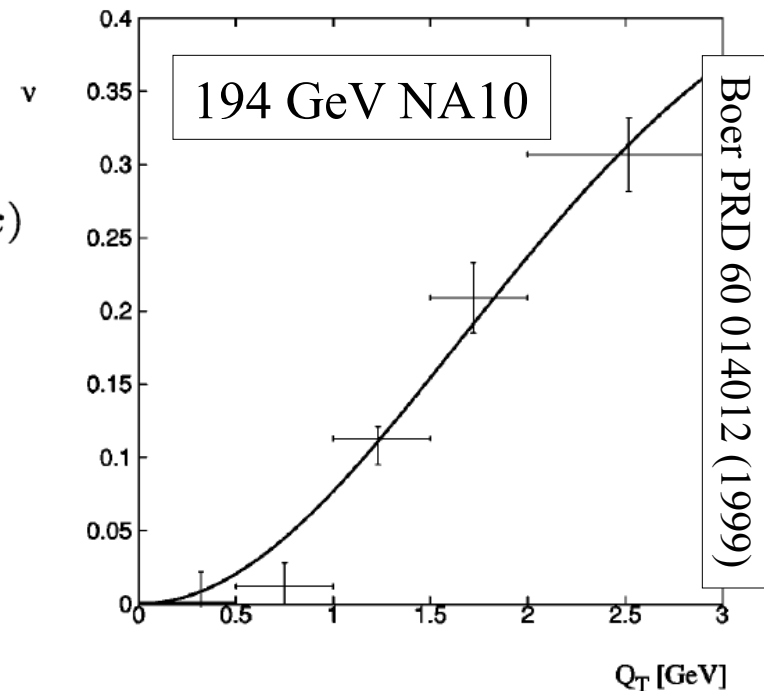


FIG. 4. Data from [3] at 194 GeV and fit [using Eq. (49)] to $\nu = 2\kappa$ as a function of the transverse momentum Q_T of the lepton pair. The fitted parameters are $M_C = 2.3 \pm 0.5$ GeV and $16\kappa_1 = 7 \pm 2$.

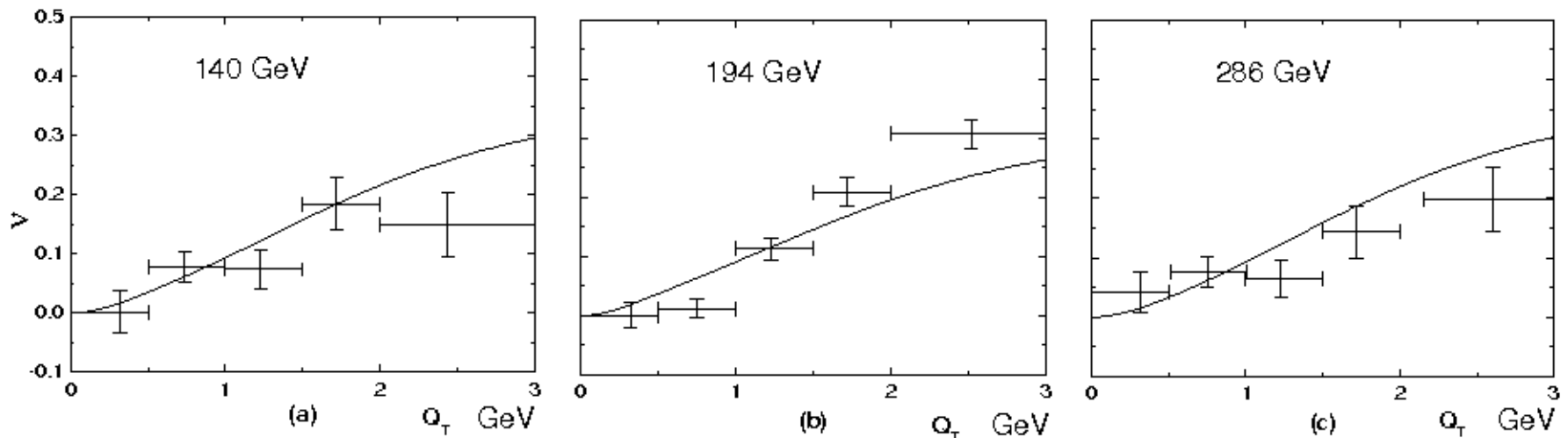
Boer-Mulders Structure Function

- Lu and Ma—quark-spectator-antiquark model

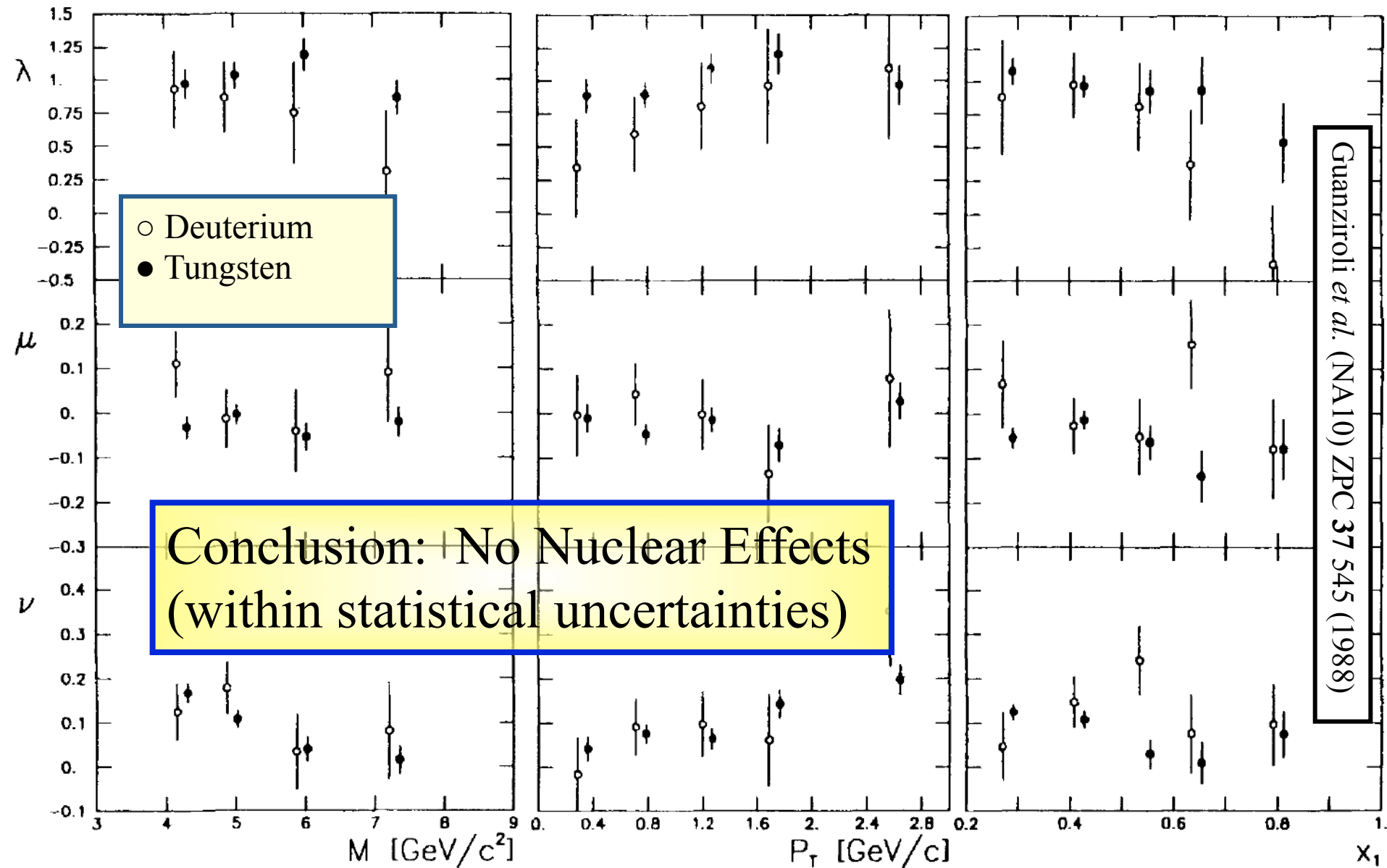
$$h_{1\pi}^{\perp} = \frac{A_{\pi}(x)}{k_{\perp}^2 [k_{\perp}^2 + B + \pi(x)]} \ln \left[\frac{k_{\perp}^2 + B_{\pi}(x)}{B_{\pi}(x)} \right]$$

- Fit all three NA10 energies

Lu, Ma, PLB **615**, 200 (2005)

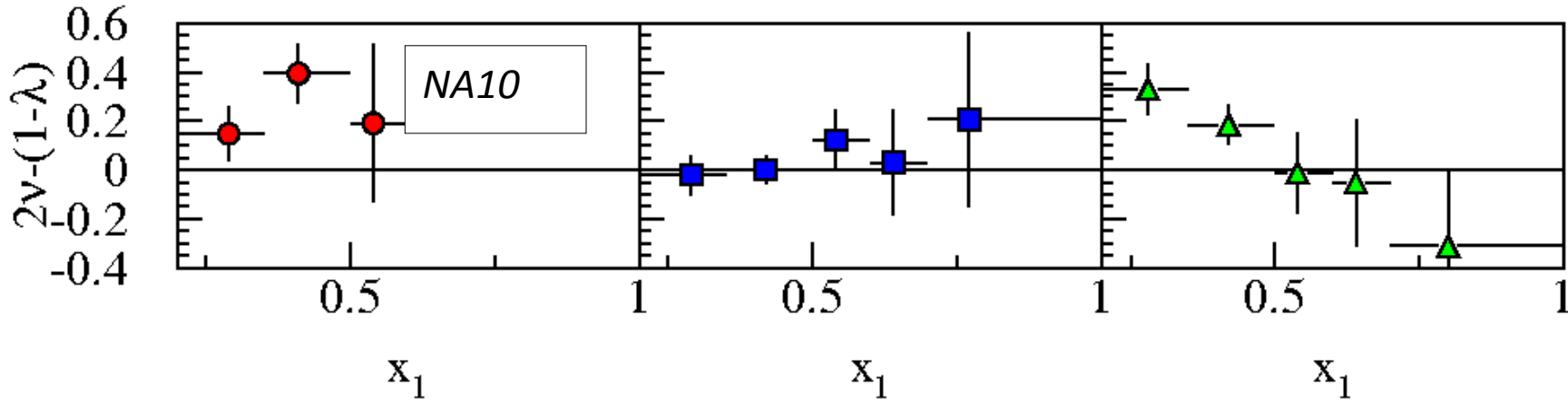
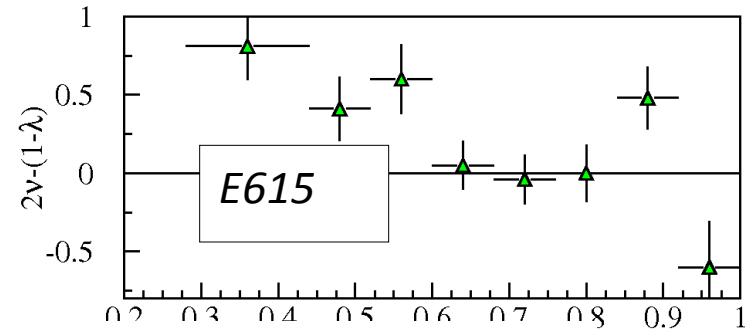


Nuclear Effect? Compare NA10 Deuterium and Tungsten



Higher twist/QCD effects in pion

- Expected only as $x_1 \rightarrow 1$



- QCD effects in pion—model only applicable for $x_1 > 0.6$ Brandenburg, Brodsky, Khoze and Muller Phys.Rev.Lett.73:939-942,1994

- Effect not large enough

