

# COMPASS results on unpolarised azimuthal asymmetries and distributions

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PHENOMENA IN  
HARD SCATTERING**

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# OUTLINE

- COMPASS experiment
- results
  - azimuthal asymmetries
  - hadron multiplicities

**COM**mon  
**MU**on and  
**PR**oton

29 Institutes, ~230 physicists

**Apparatus for  
Structure and  
Spectroscopy**

**NA58**



wide physics program carried on using both muon and hadron beam

<b>muon beam</b>	<b>deuteron</b> ( ${}^6\text{LiD}$ )	2002	L/T target polarization
	polarized target	2003	L/T target polarization
		2004	L/T target polarization
	<i>spectrometer upgrade</i>	2005	<i>SPS shutdown</i>
		2006	L target polarization
	<b>proton</b> ( $\text{NH}_3$ )	2007	L/T target polarization
	polarized target	2010	T target polarization
		2011	L target polarization

hadron beam	LH target	2008
		2009

### **longitudinally polarised muon beam**

beam intensity:  $2 \cdot 10^8 \mu^+/\text{spill}$  (4.8s/16.2s)

beam momentum: **160 GeV/c**

luminosity:  $\sim 5 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$

- muon beam:
  - nucleon spin structure
    - $\Delta G/G$
    - helicity distributions
    - transverse spin effects
    - $\Lambda$  physics
    - $\rho^0$  production
    - ...

## ← SIDIS

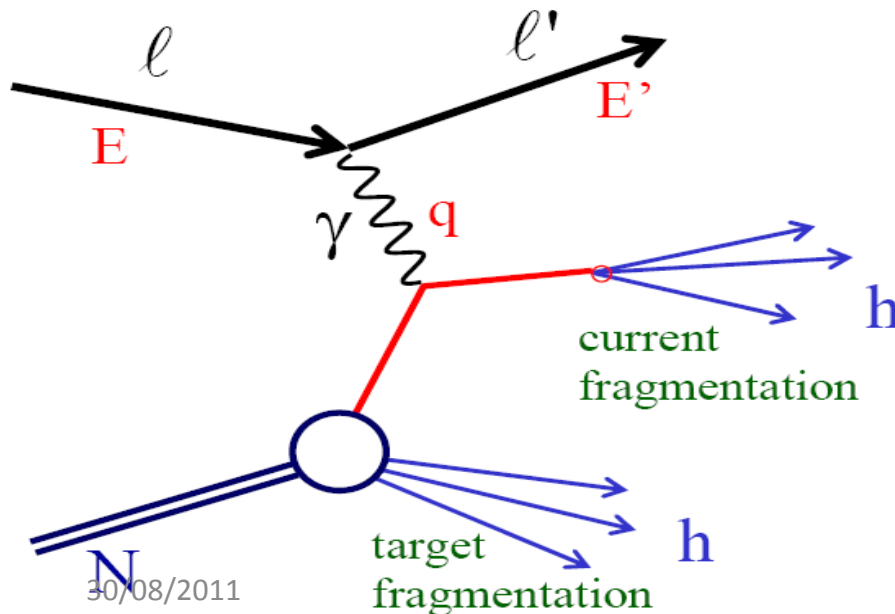
lepton interacts with a **single constituent** of the nucleon ( $Q^2 > 1 \text{ GeV}^2/c^2$ )

$$q = \ell - \ell' \quad Q^2 = -q^2$$

$$W^2 = (P + q)^2$$

$$x = \frac{Q^2}{2P \cdot q} \quad \text{Bjorken scaling variable}$$

$$y = \frac{P \cdot q}{P \cdot \ell} =_{\text{LAB}} \frac{E - E'}{E}$$



at least one hadron is detected in the final state  
(information on the **struck quark**)

$$z = \frac{P \cdot P_h}{P \cdot q} =_{\text{LAB}} \frac{E_h}{E - E'}$$

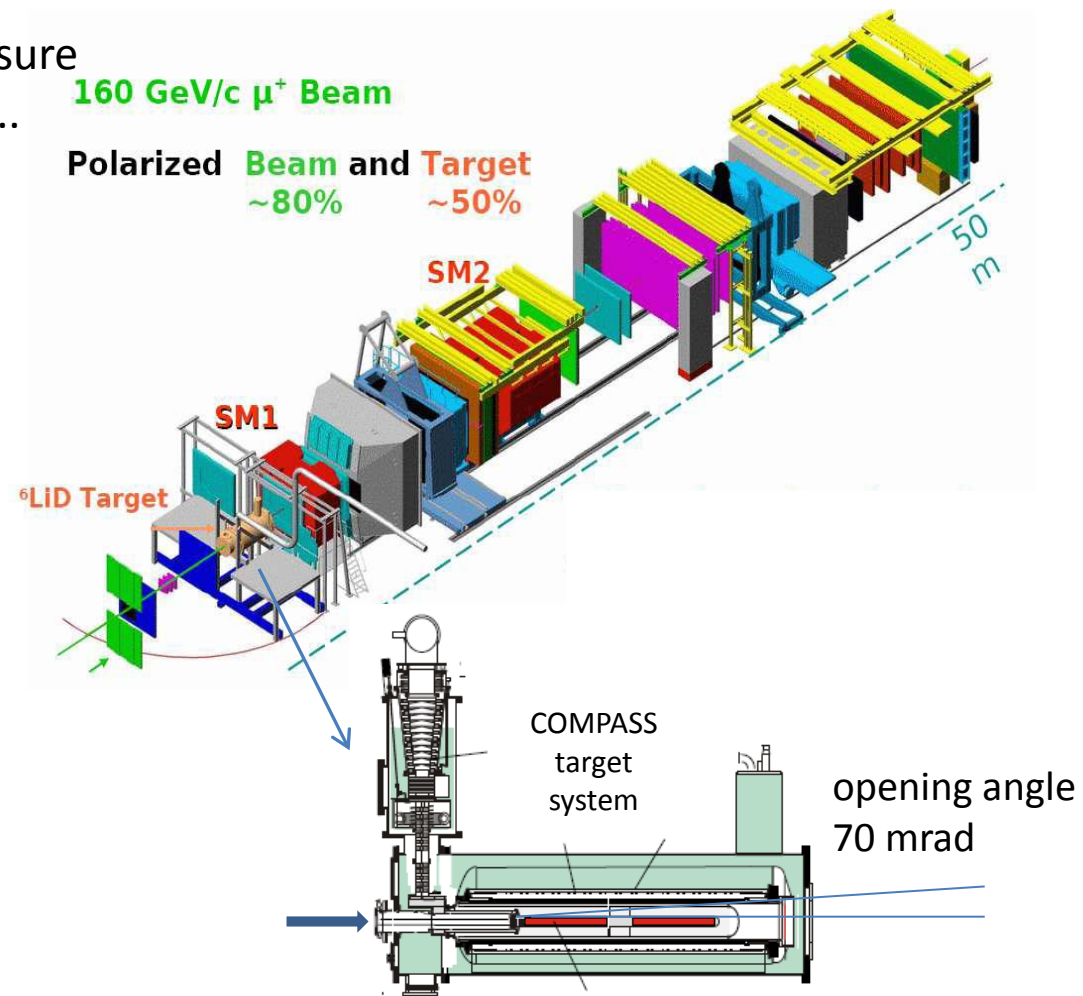
**complex apparatus** optimized to measure **spin dependent asymmetries**, but also ...

160 GeV/c  $\mu^+$  Beam

Polarized Beam and Target  
~80% ~50%

spin averaged physics  
*combining + and - pol.*

**MC simulation** needed to **correct the real data distributions** for the acceptance of the apparatus (geometrical acceptance, detector and trigger efficiencies,...)



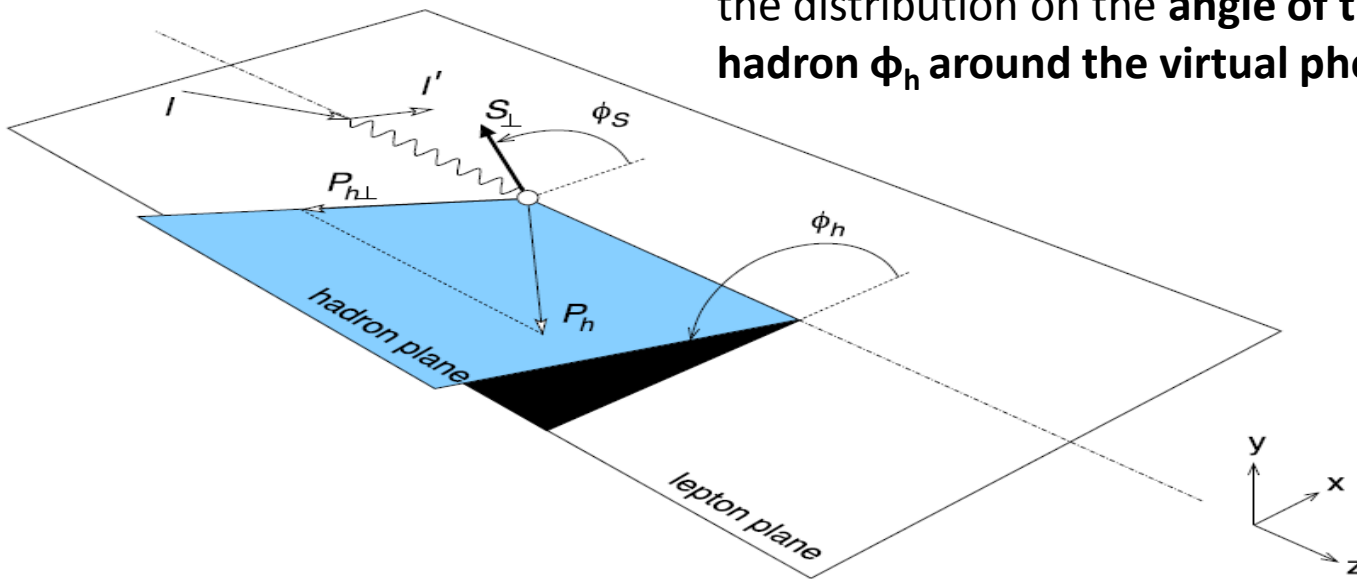
*simulation of the 2004 data taking widely studied and the apparatus behaviour well understood already a large amount of data to measure unpolarized distributions*

Results from the analyses of the  
**unpolarized distributions** at COMPASS

# 1. azimuthal asymmetries

$$N(\phi_h) \propto N_0 \cdot (1 + \varepsilon_1 A_{\cos \phi_h}^{UU} \cos \phi_h + \varepsilon_2 A_{\cos 2\phi_h}^{UU} \cos 2\phi_h + \lambda_l \varepsilon_3 A_{\sin \phi_h}^{LU} \sin \phi_h)$$

3 independent azimuthal modulations on the distribution on the **angle of the hadron  $\phi_h$**  around the **virtual photon direction**



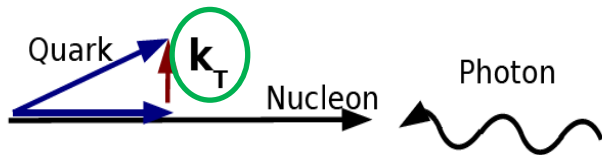


**QPM** convolution on the **TM of the quark** between different PDFs and a FFs

$$\sum_q e_q^2 f_q(x, k_\perp) \otimes D_q^h(z, p_\perp)$$

$$N(\phi_h) \propto N_0 \cdot (1 + \varepsilon_1 A_{\cos \phi_h}^{UU} \cos \phi_h + \varepsilon_2 A_{\cos 2\phi_h}^{UU} \cos 2\phi_h + \lambda_l \varepsilon_3 A_{\sin \phi_h}^{LU} \sin \phi_h)$$

mainly **Cahn** effect: **kinematical effect** proportional to the **quark transverse momentum**



$$d\sigma^{lq \rightarrow lq} \propto \hat{s}^2 + \hat{u}^2 \propto \left( 1 + \varepsilon_1 \frac{k_\perp}{Q} \cos \phi \right)$$

**Boer-Mulders** (*T-odd*!) function, one of the most famous **TMD PDF**, convoluted with the **Collins FF**



the **Boer-Mulders** function correlates the **quark transverse momentum** and the **quark spin** in an **unpolarized nucleon**

higher twist effect proportional to beam polarization

no clear interpretation in terms of PM

## Basic idea of the method

for each bin (k) in  $x, z$  and  $P_T^h$  (transv. mom. of the hadron w.r.t. the virtual photon)

The **measured azimuthal distributions** have been **corrected** for the **apparatus acceptance** which depends on  $\phi_h$

Azimuthal acceptance calculated from as the **ratio** between **Reconstructed** and **Generated** distributions from **MC**

*acceptance extracted adding one more dimension (x in the example)*

$$N_k^{corr}(\phi_h, x) = \frac{N_k(\phi_h, x)}{Acc_k(\phi_h, x)}$$

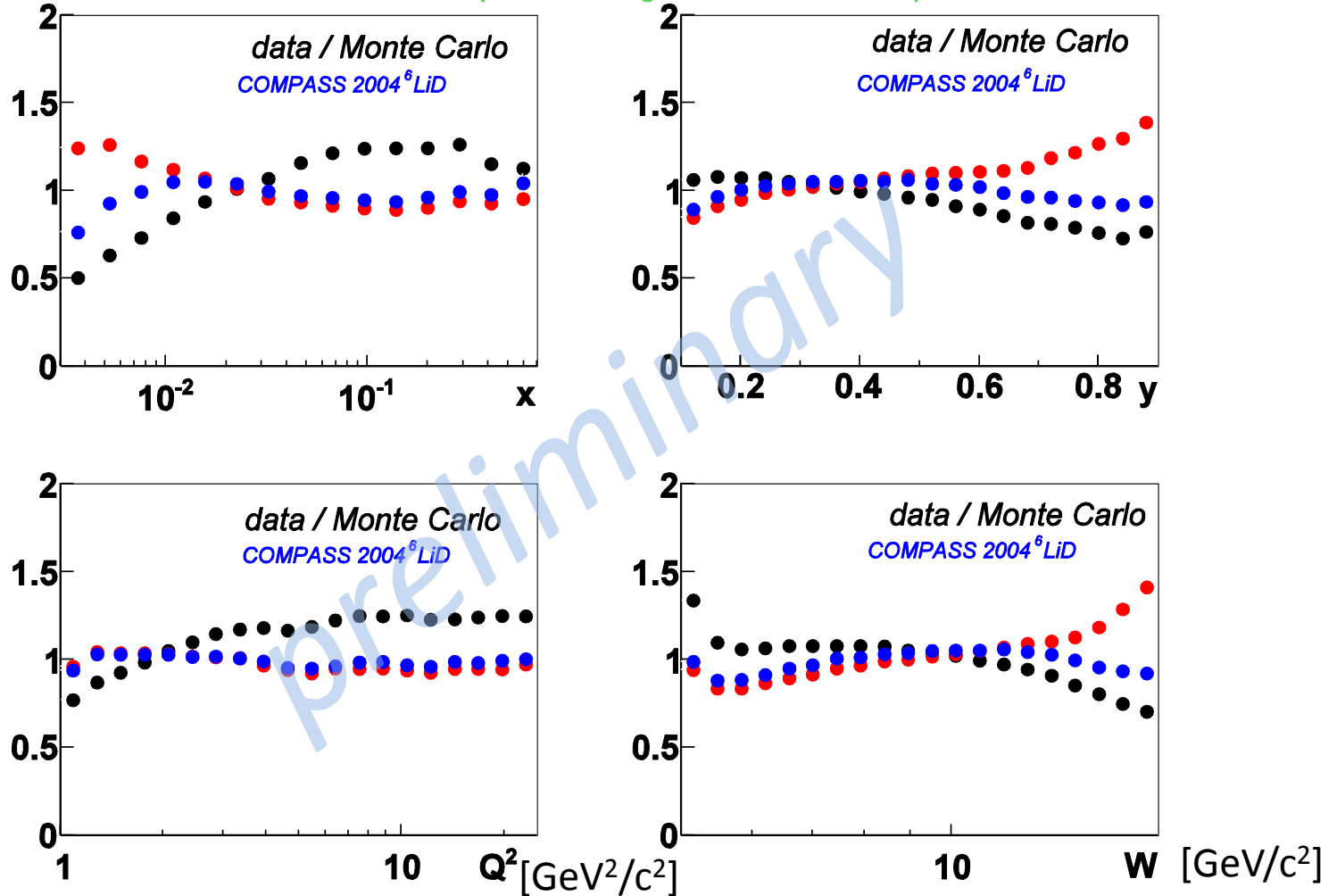
$$Acc_k(\phi_h, x) = \frac{R_k^{mc}(\phi_h, x)}{G_k^{mc}(\phi_h, x)}$$

amplitudes extracted with a fit

$$p_0 \cdot (1 + p_1 \cdot \cos \phi_h + p_2 \cdot \cos 2\phi_h + p_3 \cdot \sin \phi_h)$$

it has been checked that the extracted **azimuthal amplitudes are ~the same** by **using 3 different MC simulations** describing equally well the apparatus

data / MC comparison using 3 different MC samples



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**The apparatus azimuthal acceptance as a function of the event kinematics has been *studied* at length**

*the regions giving rise to*  
**large azimuthal modulations in the acceptance**  
*(above 50% for some bin)*  
have been **excluded**

$$Q^2 > 1 \text{ ( GeV/c )}^2$$

$$\theta_y^{\text{lab}} < 0.06$$

$$0.003 < x < 0.13$$

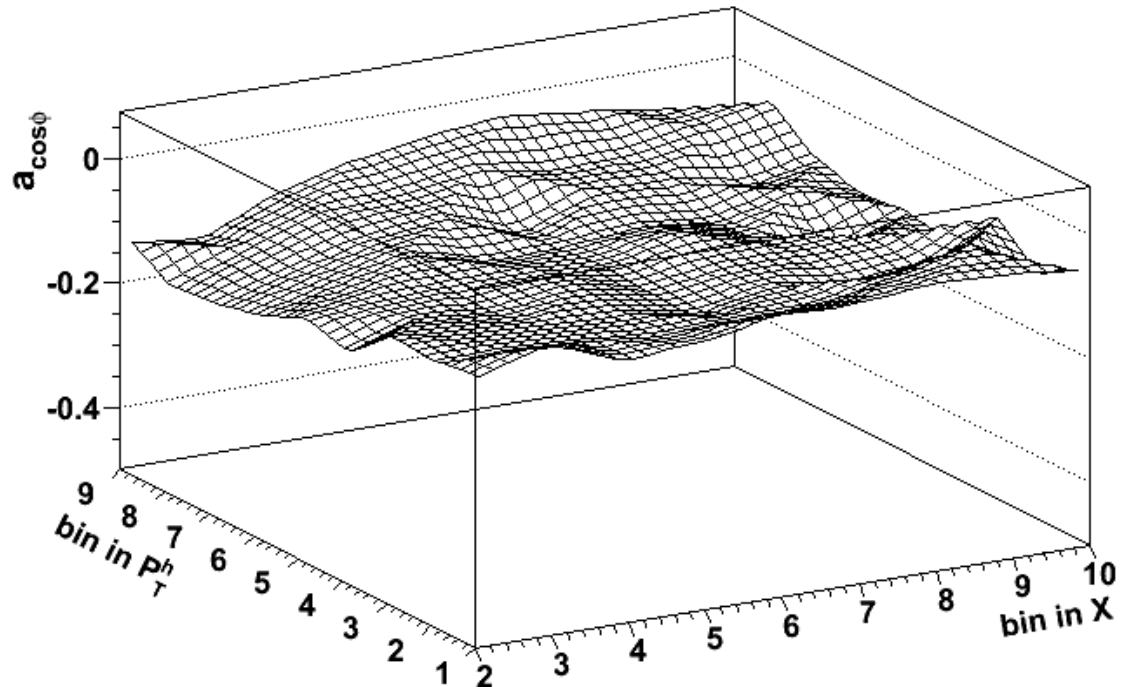
$$0.2 < y < 0.9$$

$$W > 5 \text{ GeV/c}^2$$

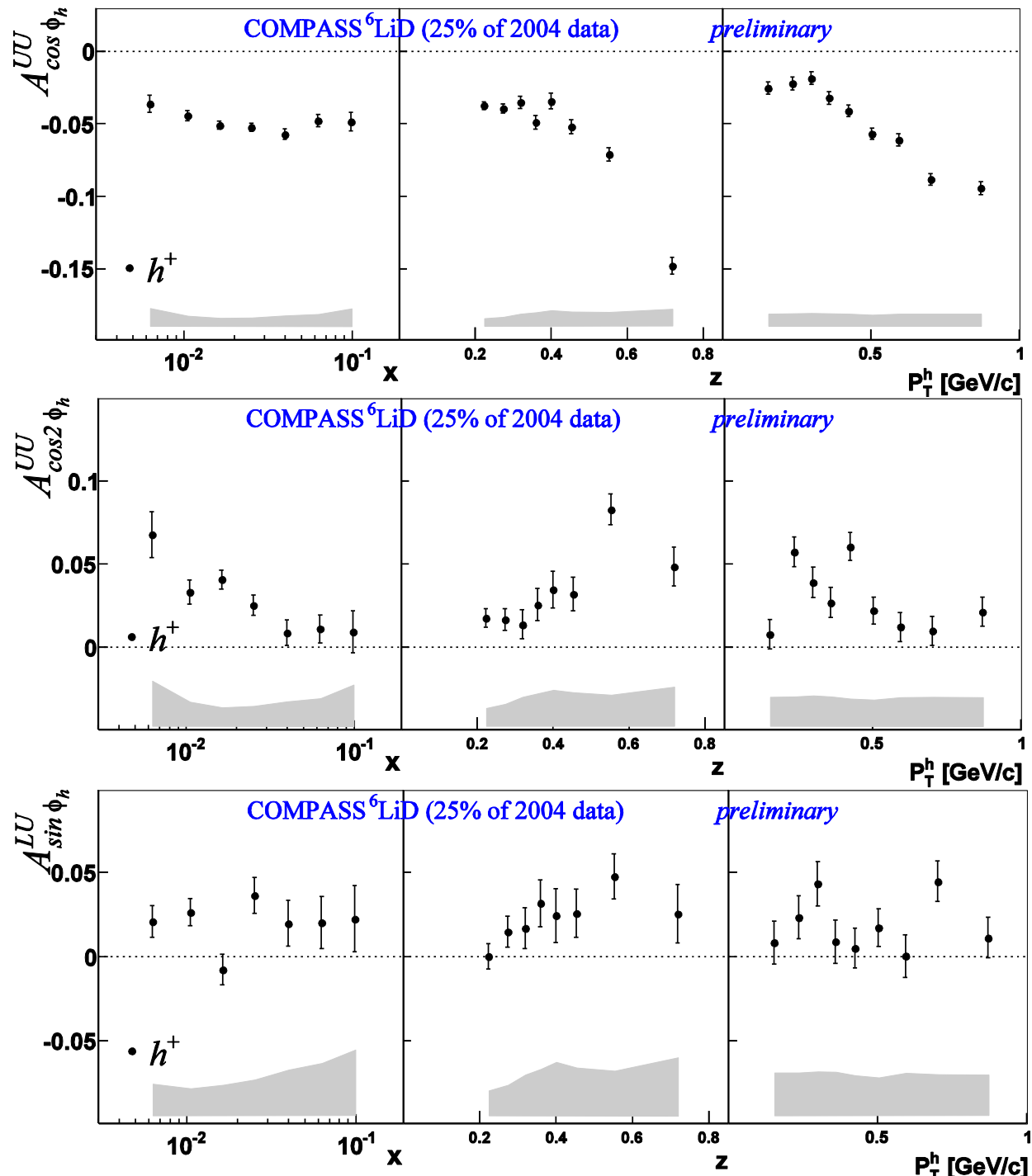
$$0.2 < z < 0.85$$

$$0.1 < P_h^T < 1 \text{ GeV/c}$$

amplitude of the  $\cos\phi$  acceptance modulation ( $a_{\cos\phi}$ )

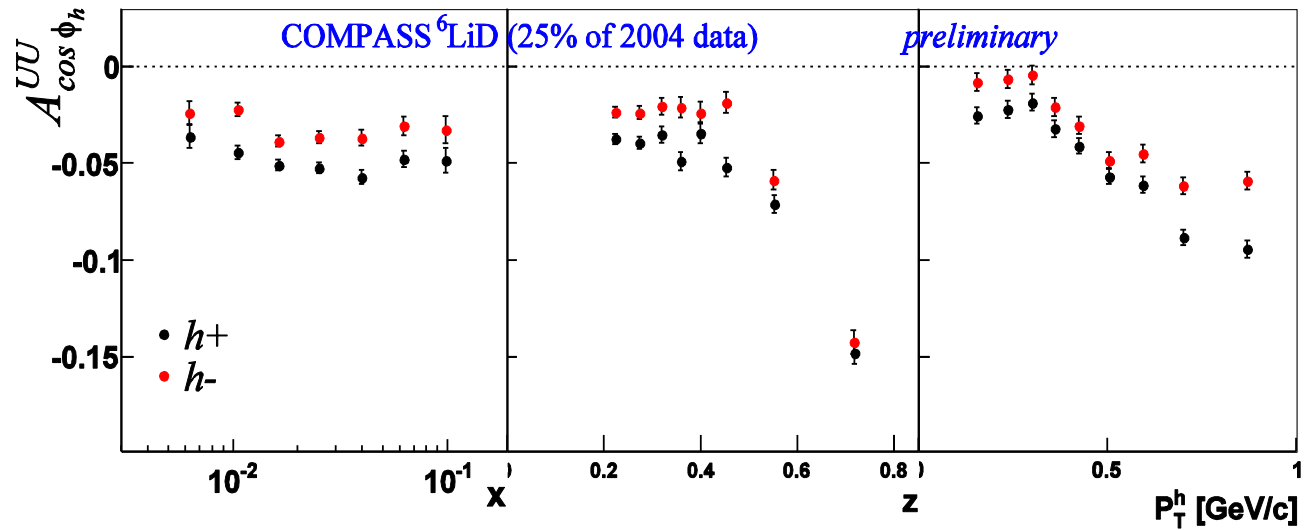


The amplitudes of the 3 azimuthal modulations have been measured separately for positive and negative hadrons, as functions of the kinematical variables  $x, z$  and  $P_T^h$

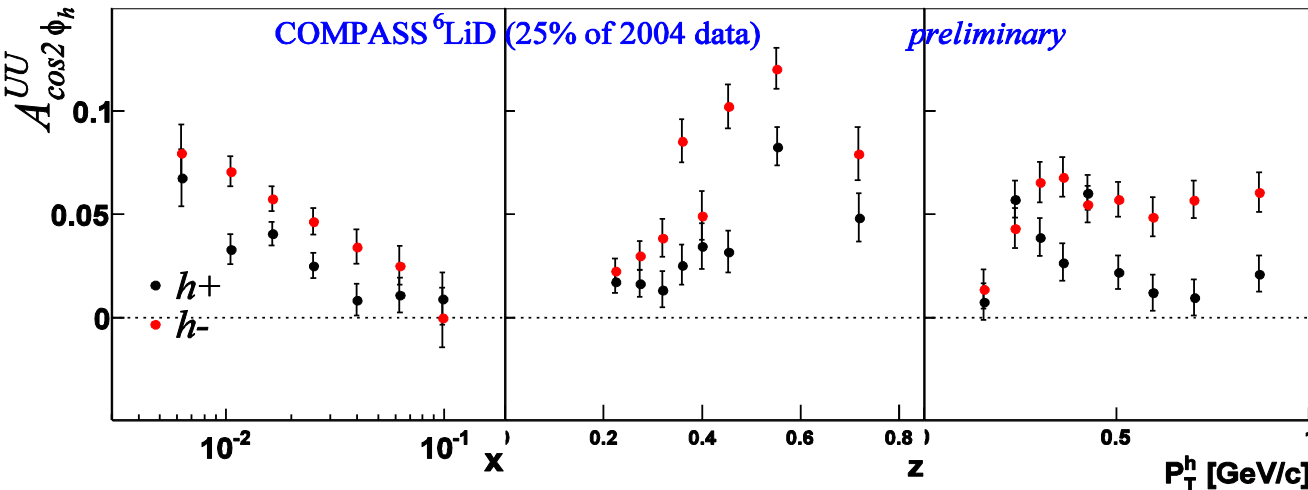


*presented at spin2010*

comparison between results from positive and **negative** hadrons



strong kinematical trends  
in  $z$  and  $P_T^h$   
(difficult to describe)



different results  
for positive and  
**negative** hadrons :  
possible signature of the  
Boer-Mulders TMD PDF

prediction for the Cahn effect can be calculated assuming

$$f_q(x, k_\perp) = f_q(x) \frac{1}{\pi \langle k_\perp^2 \rangle} e^{-k_\perp^2 / \langle k_\perp^2 \rangle}$$

$$D_q^h(z, \vec{p}_\perp) = D_q^h(z) \frac{1}{\pi \langle p_\perp^2 \rangle} e^{-p_\perp^2 / \langle p_\perp^2 \rangle}$$

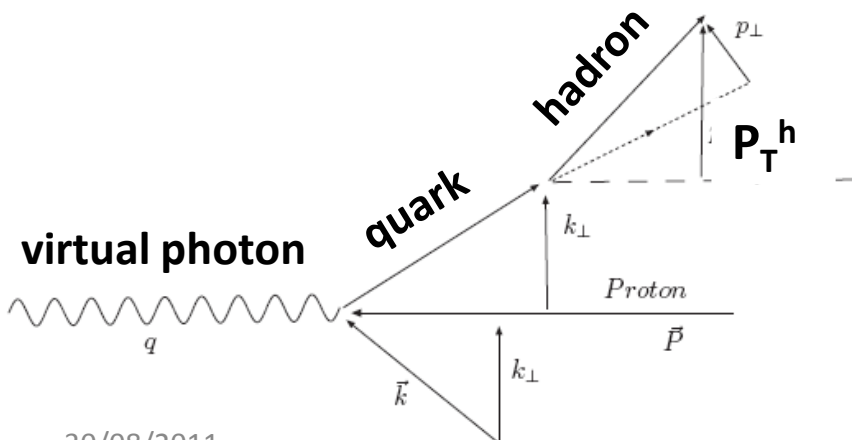


$$\sum_q e_q^2 f_q(x, k_\perp) \otimes D_q^h(z, p_\perp)$$



and the  $\cos \phi_h$  amplitude is expected to be

$$A_{\cos \phi_h}^{UU}(z) = \frac{z \langle k_\perp^2 \rangle \sqrt{\pi}}{2 \langle Q \rangle \sqrt{\langle P_T^{h2} \rangle}}$$



$$\langle P_T^{h2} \rangle = z^2 \langle k_\perp^2 \rangle + \langle p_\perp^2 \rangle$$

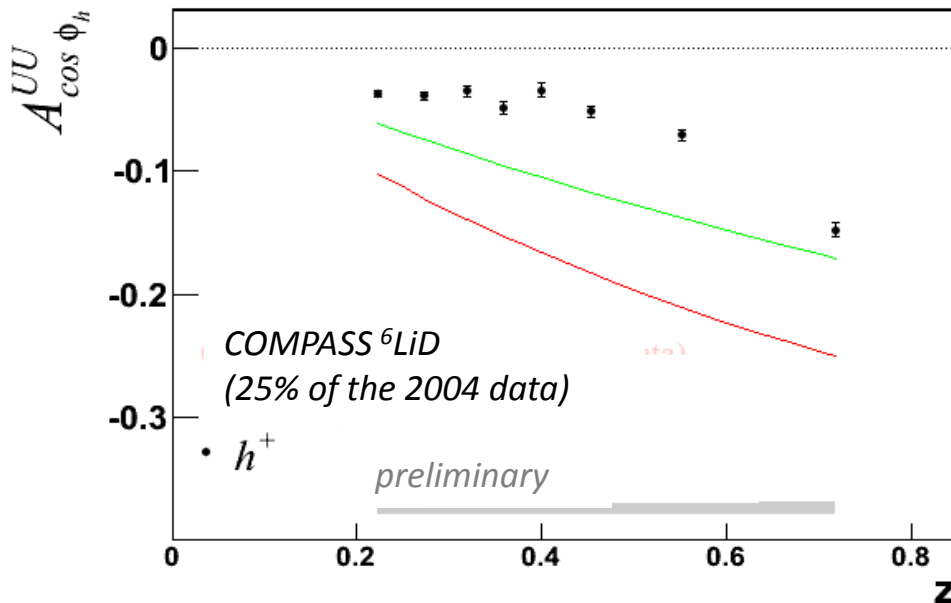
$$A_{\cos \phi_h}^{UU}(z) = \frac{z \langle k_{\perp}^2 \rangle \sqrt{\pi}}{2 \langle Q \rangle \sqrt{z^2 \langle k_{\perp}^2 \rangle + \langle p_{\perp}^2 \rangle}}$$

results are not well reproduced by calculations of this expression using:

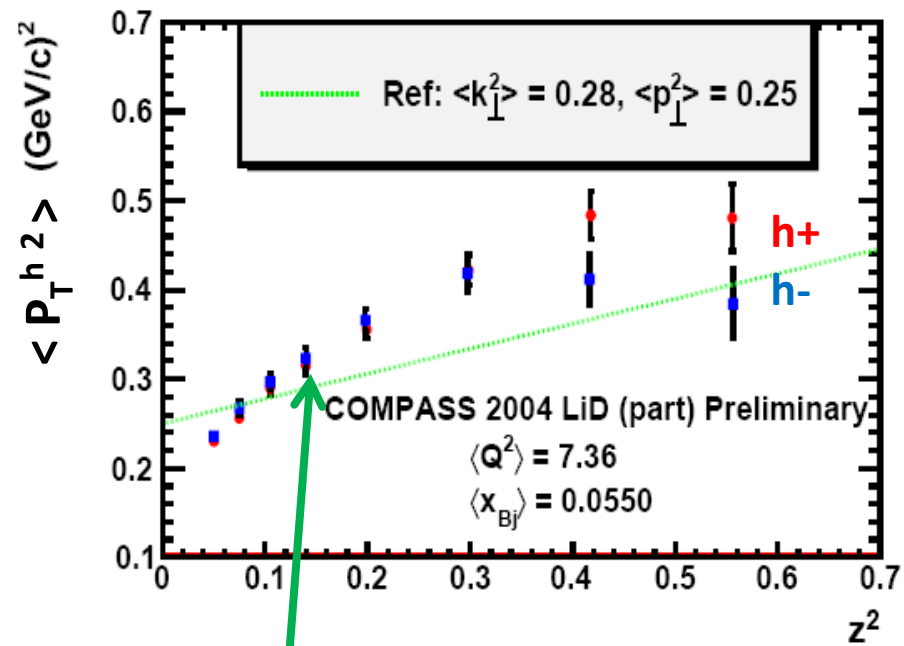
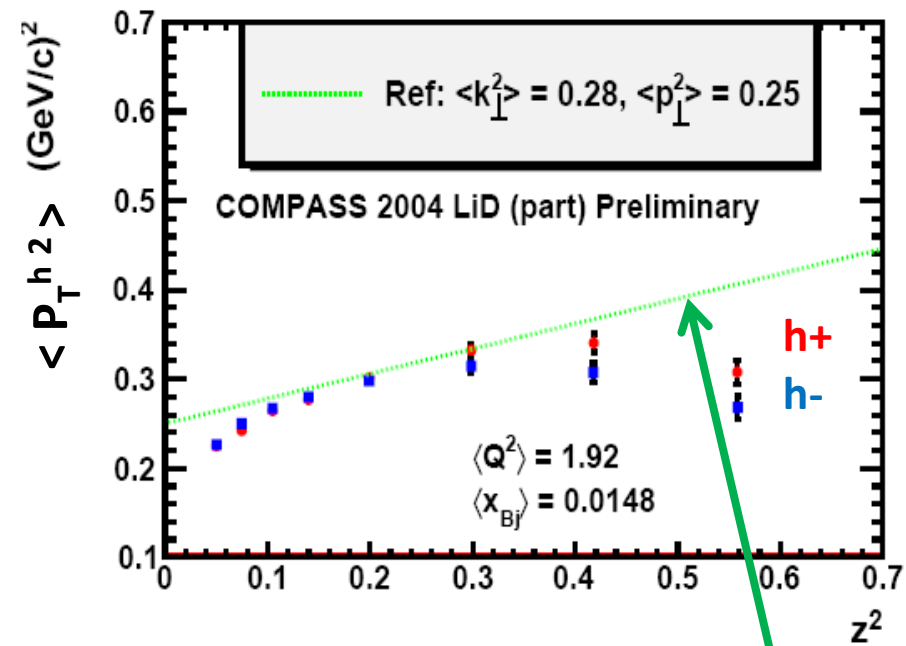
$\langle k_{\perp}^2 \rangle, \langle p_{\perp}^2 \rangle$  (GeV/c)<sup>2</sup>

**0.25, 0.20** Anselmino et al. (PRD71 2005)  
combined analysis of EMC data

**0.38, 0.16** Schweitzer et al. (PRD81 2010)  
from  $\langle P_T^{h2}(z) \rangle$  measured by HERMES







moreover the relation

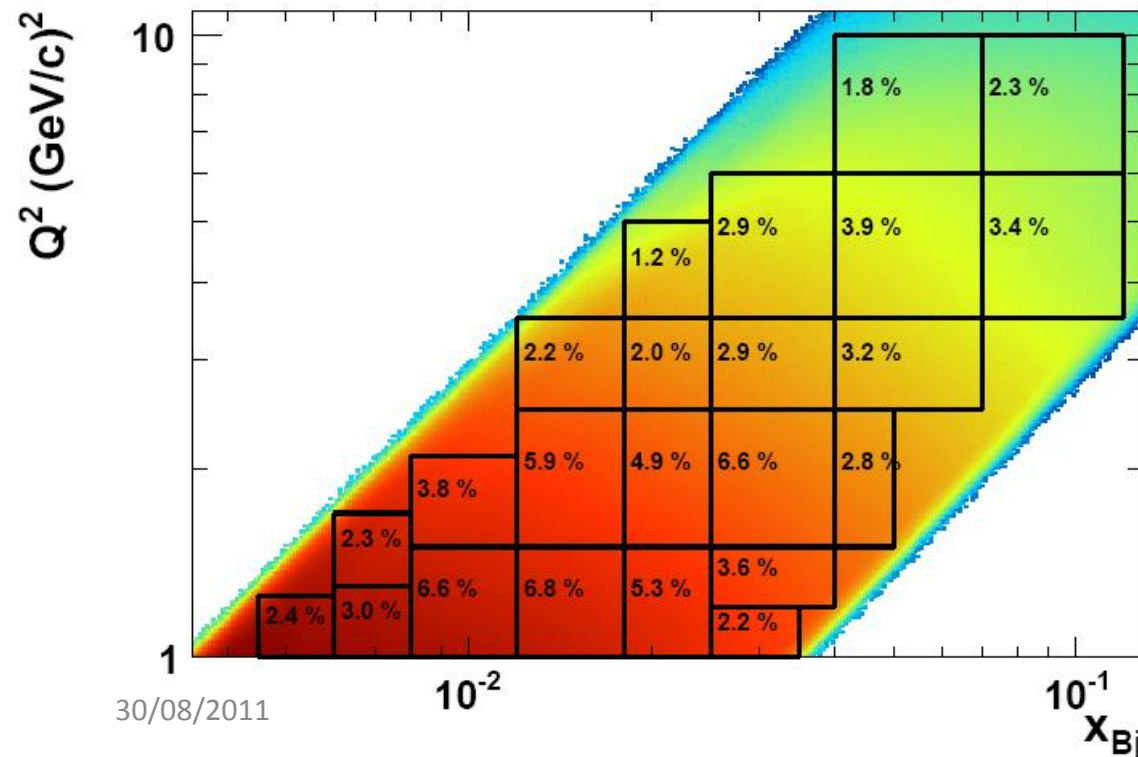
$$\langle P_T^{h2} \rangle = z^2 \langle k_{\perp}^2 \rangle + \langle p_{\perp}^2 \rangle$$

seems not to hold for COMPASS data ...

## 2. *transverse momentum dependent hadron multiplicities*

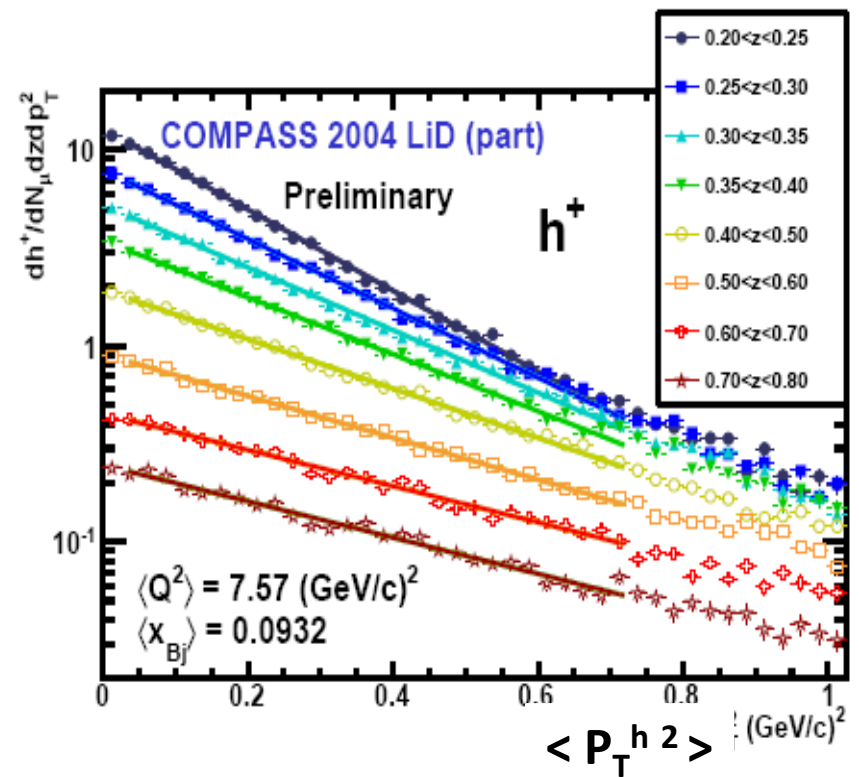
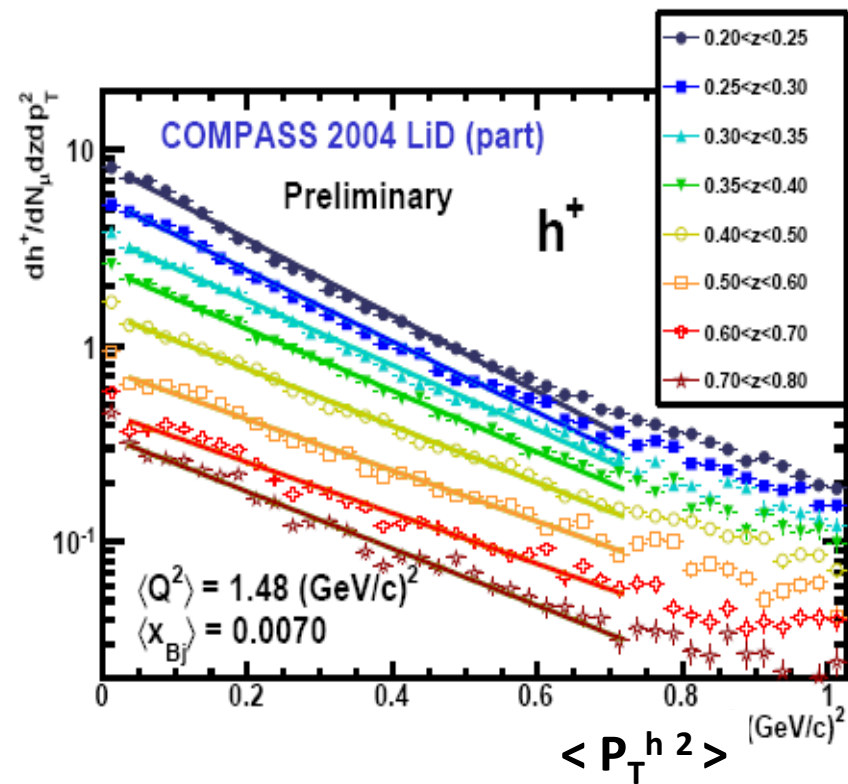
$P_T^h$  dependent distributions measured from the unpolarized hadron multiplicities  $\propto e^{-\frac{P_T^{h2}}{\langle P_T^{h2} \rangle}}$

**fits to the acceptance corrected hadron multiplicities** in **different kinematical bins ( $Q^2, \mathbf{x}$ )** to get  $\langle P_T^{h2} \rangle$  as function of  $z$  separately for positive and negative hadrons



$Q^2 > 1 ; 0.1 < y < 0.9$

$0.2 < z < 0.8$

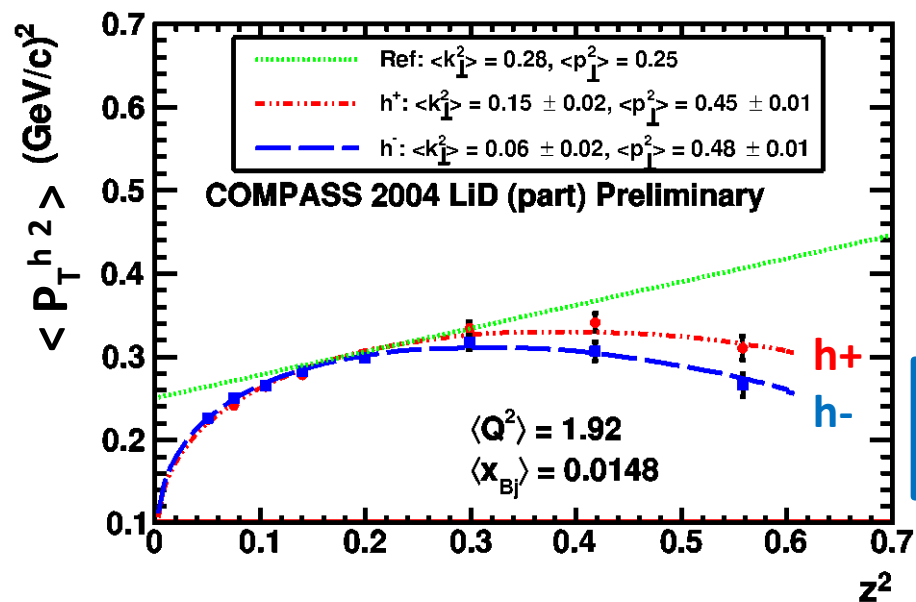


from the fits to the hadron multiplicities it seems that the dependence on  $z^2$  is not linear...

$$\langle P_T^{h2} \rangle = z^2 \langle k_\perp^2 \rangle + \langle p_\perp^2 \rangle$$



$$\langle P_T^{h2} \rangle = z^2 \langle k_\perp^2 \rangle + z^{1/2} (1 - z)^{3/2} \langle p_\perp^2 \rangle$$

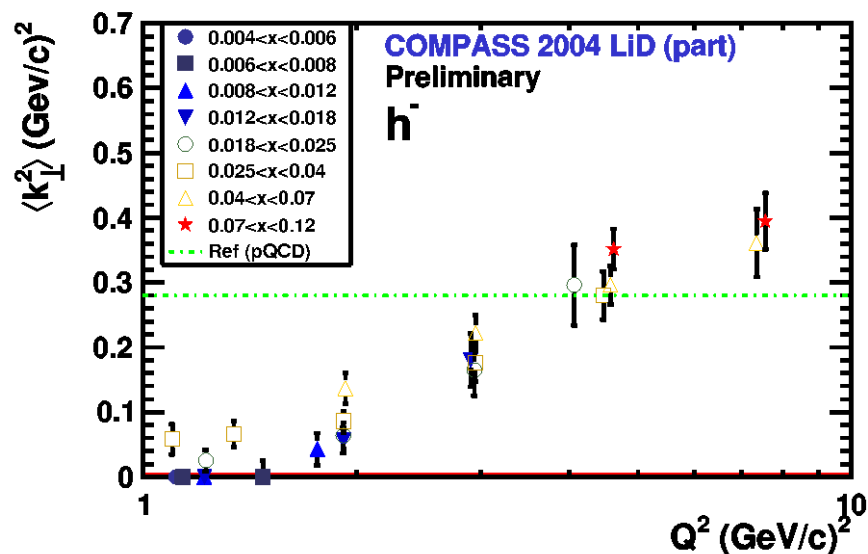
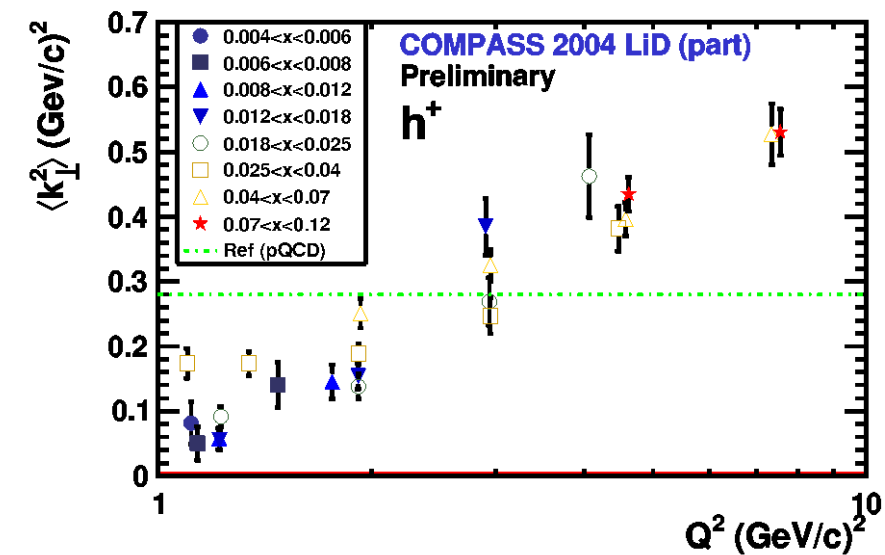


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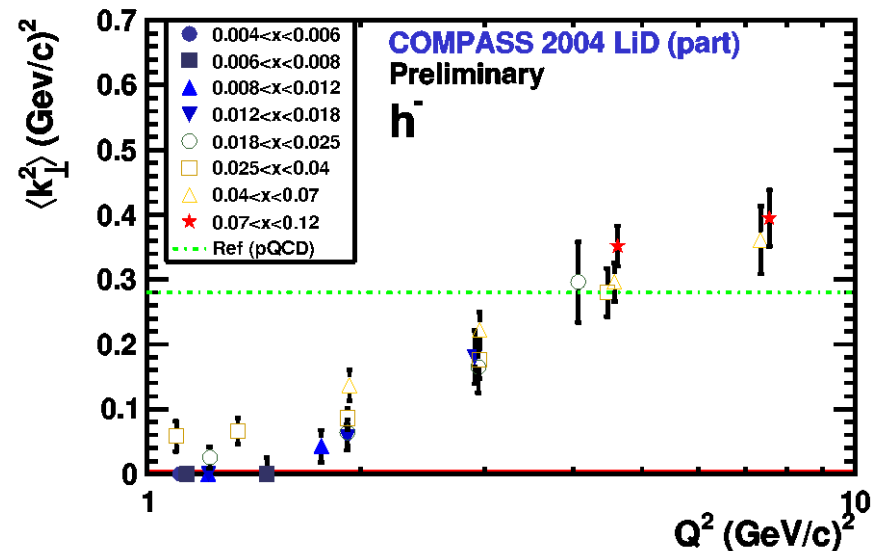
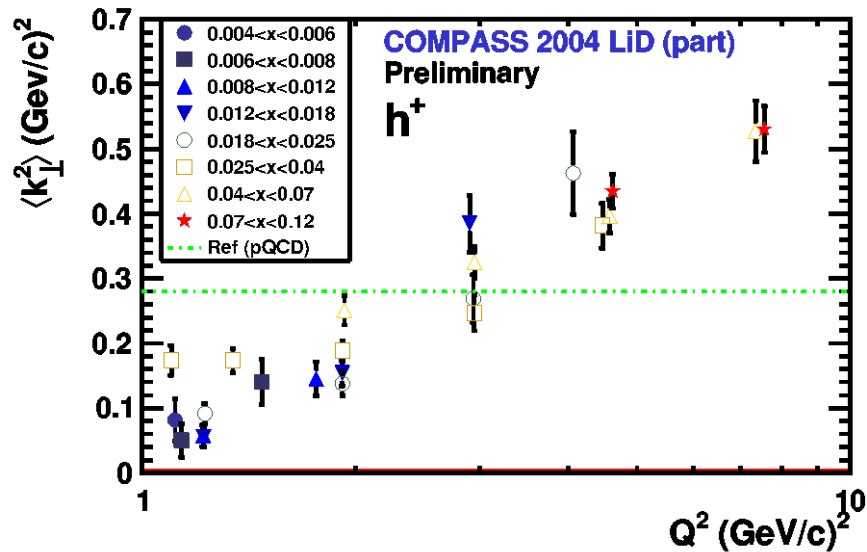


... extracting  $k_{\perp}$  using this parameterization ...



# intriguing results !

... extracting  $k_{\perp}$  using this parameterization ....



### *3. hadron multiplicities vs x and z*

QPM  
collinear

$$\frac{dM^h(x, z)}{dz} = \frac{\sum_q e_q^2 f_q(x) \cdot D_q^h(z)}{\sum_q e_q^2 f_q(x)}$$

 ***allow to extract FFs***

analysis performed in bins  
of  $(x, z)$  and  $(Q^2, z)$

$Q^2 > 1$  ;  $0.1 < y < 0.9$

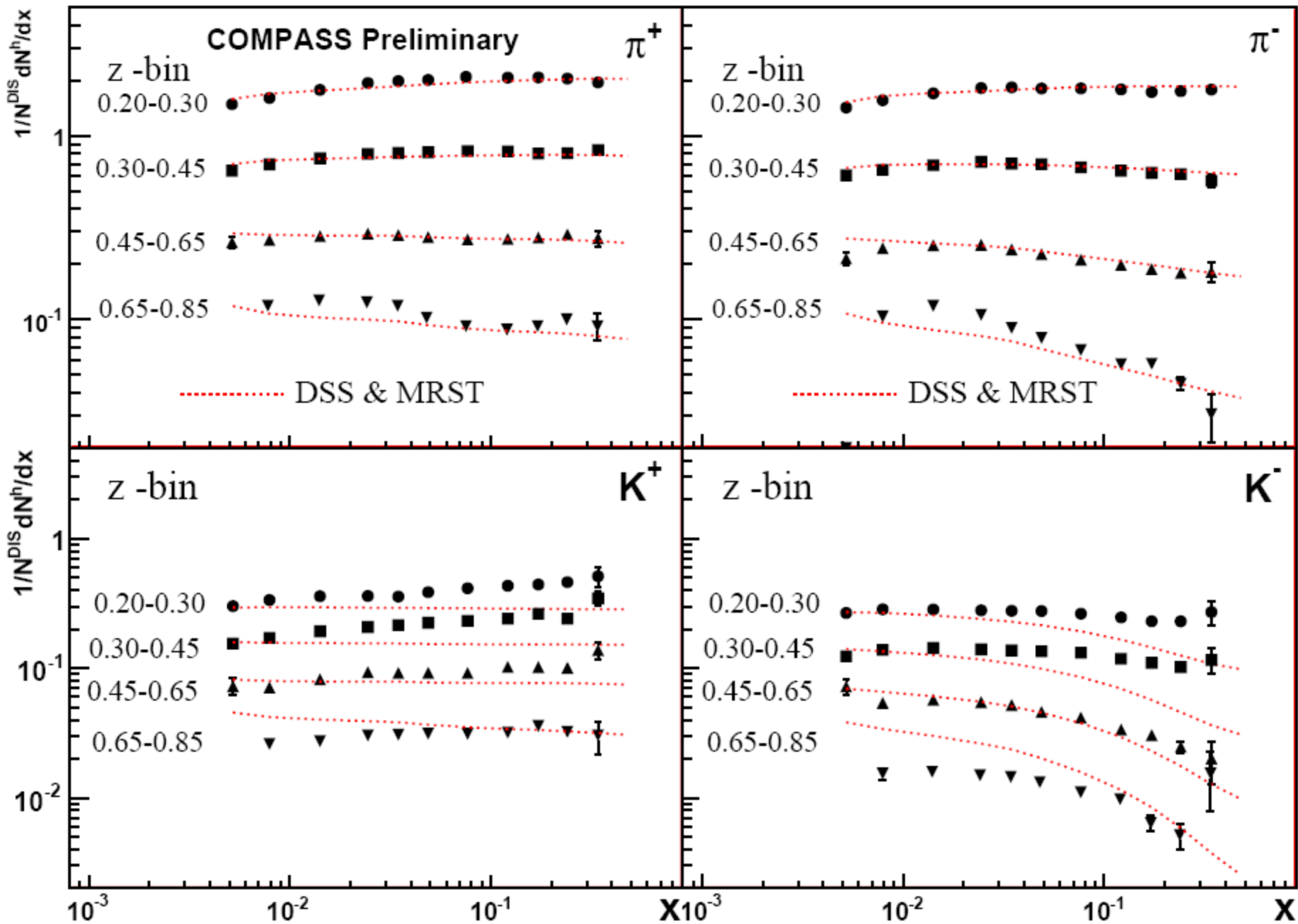
$W > 7$  (to avoid large acceptance corrections )

$0.2 < z < 0.85$

real data multiplicities corrected for the acceptance

$\pi^+$ ,  $\pi^-$ ,  $K^+$ ,  $K^-$  identification performed





## ***CONCLUSIONS and OUTLOOK***

- COMPASS provides interesting results from the unpolarized measurements on deuterium
- TM effects have been studied in the COMPASS kinematics both from the azimuthal asymmetries and the hadron multiplicities
- interesting kinematical dependencies in  $z$  and  $P_T^h$  from both azimuthal asymmetries and hadron multiplicities
- interesting inputs for theory and global analyses

## ***CONCLUSIONS and OUTLOOK***



### **COMPASS2**

SIDIS program in parallel with DVCS  
on proton (LH<sub>2</sub> target)  
apparatus more suitable for  
cross section measurement  
***more precise measurements***  
on the ***unpolarized distributions***  
and for the ***extraction of FFs***

backup

# COMPASS MC chain

- **generation**

Lepto

DIS events simulation

- **propagation**

COMGEANT

simulates the interaction and the propagation of the particles inside the spectrometer

magnets, materials, detectors, triggers, ...

*different setups are individually described*

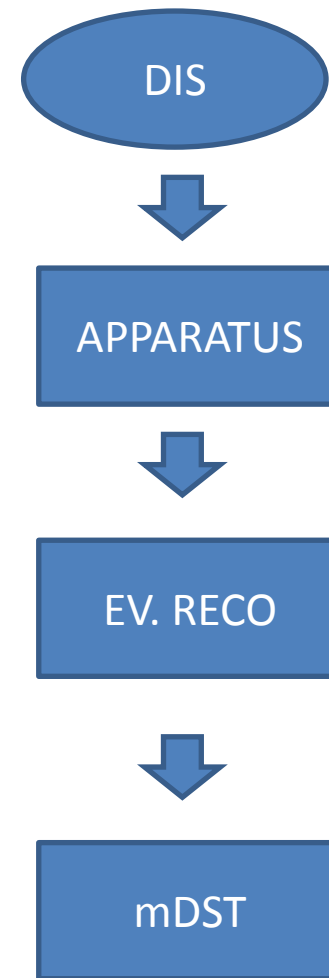
- **reconstruction**

CORAL (program for the data reconstruction)

vertices, tracking, momentum,

*the same program used in MC and real data*

- files with the reconstructed quantities are produced in the same format as for the real data



# COMPASS MC full chain simulation

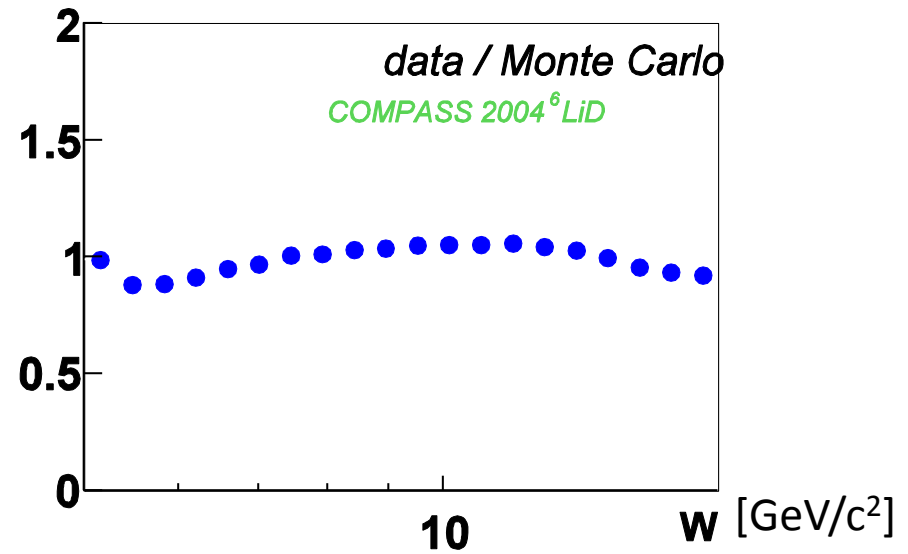
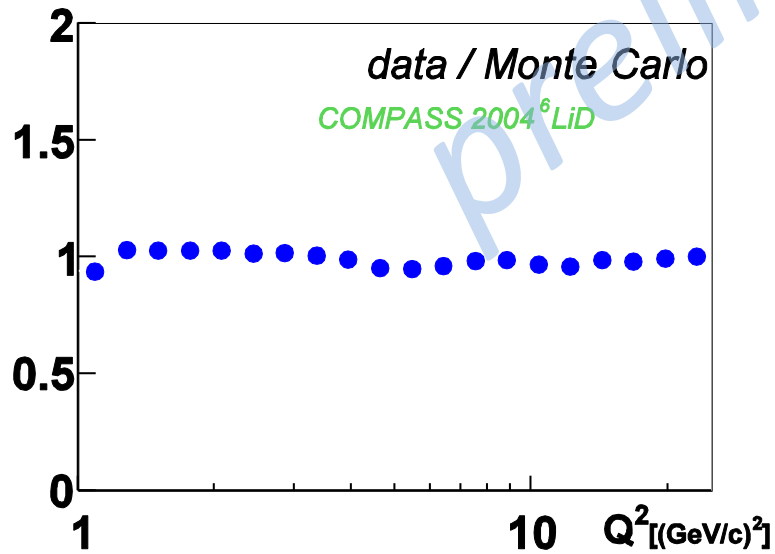
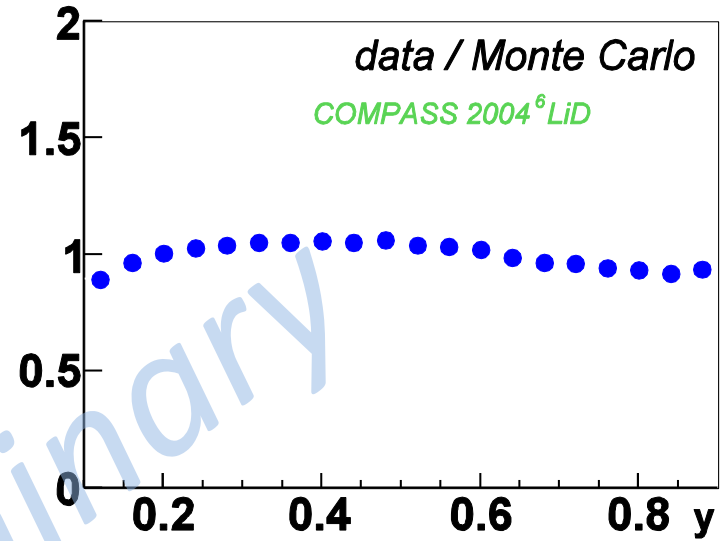
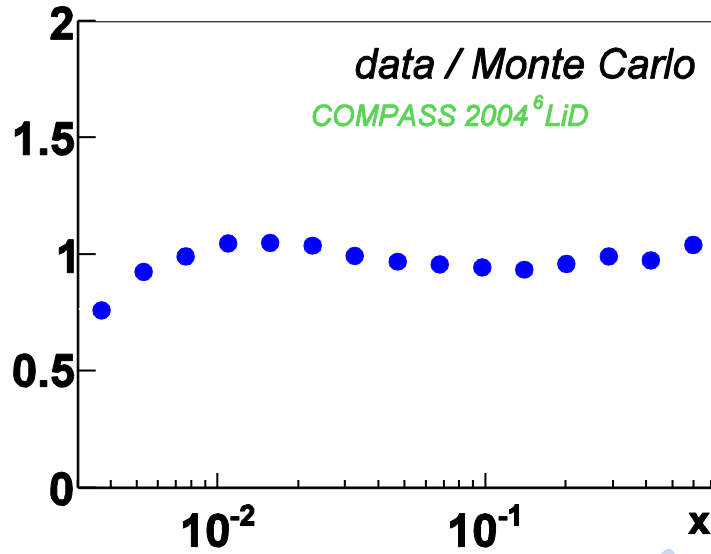
- **generation**
  - Lepto
  - DIS events simulation
- **propagation**
  - COMGEANT
  - simulates the interaction and the propagation of the particles inside the spectrometer
  - magnets, materials, detectors, triggers, ...
  - different setups are individually described*
- **reconstruction**
  - CORAL (program for the data reconstruction)
  - tracking , vertices, momentum,
  - the same program used in MC and real data*
- files with the reconstructed quantities are produced in the same format as for the real data

***fine tuning can be different for each analysis***

indeed huge work has been done in order to ***optimize the description of the real data conditions***

*an example of Monte Carlo – Real Data comparison in the next 2 slides*

# Events RD/MC ratios



# Hadrons RD/MC ratios

