

# Effect of metamaterial engineering on the superconductive properties of ultrathin layers of NbTiN

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# **Engineered superconductors: historic overview**

#### - Proposals of 1D and 2D superconductors:

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Possibility of Synthesizing an Organic Superconductor\*

W. A. LITTLE Department of Physics, Stanford University, Stanford, California (Received 13 November 1963; revised manuscript received 27 January 1964)

V. L. Ginzburg, "On surface superconductivity", Phys. Lett. 13. 101 (1964).

## - Electromagnetic approach to superconductivity:

 $V(\vec{q},\omega) = \frac{4\pi e^2}{q^2 \varepsilon_{off}(\vec{q},\omega)} \quad \text{$\epsilon$ should be negative and small!}$ 

D.A. Kirzhnits, E.G. Maksimov, D.I. Khomskii, "The description of superconductivity in terms of dielectric response function", J. Low Temp. Phys. 10, 79 (1973).

#### New solution: artificial high polarizability metamaterials

I.I. Smolyaninov and V.N. Smolyaninova, "Metamaterial superconductors", Phys. Rev. B 91, 094501

Nanofabrication has reached a few nanometers scale – ability to engineer artificial material, which has enhanced superconducting critical temperature

FIG. 1. Proposed

model of a superconducting organic molecule. The molecule A is a long un-

saturated polvene chain called the "spine." The mole-

cules B are side chains attached to

the spine at points P, P', · · · .



**Recent developments in** metamaterials

Artificial "metamaterials" may be created from much larger building blocks than atoms, and the electromagnetic properties of these fundamental building blocks ("metaatoms") may be engineered at will



# Engineering small and negative dielectric function $\varepsilon(q,\omega)$



with higher  $T_c \rightarrow NbTiN$  hyperbolic metamaterials

# Hyperbolic metamaterial scenario



$$\frac{k_x^2 + k_y^2}{\varepsilon_{zz}} + \frac{k_z^2}{\varepsilon_{xx}} = \frac{\omega^2}{c^2}$$
$$\varepsilon_{yy} = \varepsilon_{xx} < 0; \ \varepsilon_{zz} > 0$$

#### **Effective Coulomb potential:**

$$V(\vec{q},\omega) = \frac{4\pi e^2}{q_z^2 \varepsilon_{zz}(\vec{q},\omega) + (q_x^2 + q_y^2) \varepsilon_{xx}(\vec{q},\omega)}$$

diverges at: 
$$q_z^2 \mathcal{E}_{zz}(\vec{q},\omega) + (q_x^2 + q_y^2) \mathcal{E}_{xx}(\vec{q},\omega) \approx 0$$

- Electron-electron interaction is strongly enhanced in hyperbolic metamaterials
- The best choice of geometry appears to be metal/dielectric layered structure

#### Smolyaninov, Smolyaninova, PRB 91, 094501

## Hyperbolic metamaterial

BiO

SrO

CuO<sub>2</sub>

Ca CuO

SrO BiO

BiO SrO

CuO<sub>2</sub>

Ca

CuO<sub>2</sub>

SrO



# Effect of metamaterial engineering on the $T_c$ and $H_c$ of ultrathin layers of NbTiN







## **Samples**

In ultrathin layers of NbTiN T<sub>c</sub> depends on film thickness

NbTiN films were grown using reactive high-power impulse magnetron sputtering (R-HiPIMS) 2 nm < t  $_{\rm NbTiN}$ < 3.6 nm 0.8 nm < t  $_{\rm AIN}$ < 2 nm

## T<sub>c</sub> dependence on number of layers



The  $T_c$  increases with the number of layers, with largest increase of the  $T_c$  by 3.6 K (by 32%) for 17 NbTiN/AlN layers, as expected for a superconducting hyperbolic metamaterial. Indeed, such a material with a large number of building NbTiN/AlN blocks can be considered to be a "metamaterial".

J. Appl. Phys. 130, 073901

## Polarization reflectometry: anisotropy of the dielectric function



Hyperbolic properties of multilayers have been achieved:  $\varepsilon_{in plane} < 0$ ;  $\varepsilon_{out of plane} > 0$ 

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# Variation of volume fraction of superconductor

The magnitude of the electron-electron interaction depends on an anisotropic dielectric function, which, in turn is a function of the volume fraction of the metal, *n*. In the case of a layered hyperbolic metamaterial, the dielectric function dependence on *n* can be found from the Maxwell-Garnett approximation:

$$\varepsilon_1 = n\varepsilon_m + (1-n)\varepsilon_d$$
  $\varepsilon_2 = \frac{\varepsilon_m \varepsilon_d}{(1-n)\varepsilon_m + n\varepsilon_d}$ 

 $T_{\rm c}$  should depend on the volume fraction of metal in the metamaterial. To test this assumption, the volume fraction of the dielectric in the multilayers was varied, while the thickness of the metal layer remained constant, since the  $T_{\rm c}$  of a single layer depends on its thickness.

 $T_{\rm c}$  of these superconducting hyperbolic metamaterials decreases with increasing volume fraction of the metal. The results are consistent with those for aluminium and tin based hyperbolic metamaterials.

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## Anomalous behavior of upper critical field in NbTiN-based multilayers



• While  $H_{c2\perp}$  of the single NbTiN layer follows the Werthamer–Hefand–Hohenberg (WHH) model, the multilayers exhibit a higher  $H_{c2\perp}$  with an anomalous linear temperature dependence, or a slight positive curvature.

• Extrapolated  $H_{c\parallel}(0) \sim 420 \ kOe$ 

Ultrathin superconducting films are known to exhibit extremely high values of the critical magnetic field  $H_c$ . However, the  $T_c$  of a single layer of ultrathin films is significantly reduced.

We have demonstrated the ability to maintain the  $T_c$  of a superconducting coating while keeping (or even increasing) its critical magnetic field at the same level as for the ultrathin films. This result can lead to many advances in technological applications of superconductors.

# **Conclusions**

## NbTiN/AIN multilayered metamaterials with ultrathin layers

- exhibit up to a 32% enhancement of  $T_c$  with respect to the  $T_c$  of a single ultrathin NbTiN layer.
- This  $T_c$  increase can be attributed to an enhanced electron-electron interaction in superconducting hyperbolic metamaterials.
- The critical fields in these multilayers are high and have anomalous linear temperature dependence in the perpendicular to the magnetic field direction.

These results demonstrate that the metamaterial approach to superconductor engineering can enable the increase of  $H_{c2}$  as well as  $T_{c}$ .

# **Anisotropy of coherence length**

