

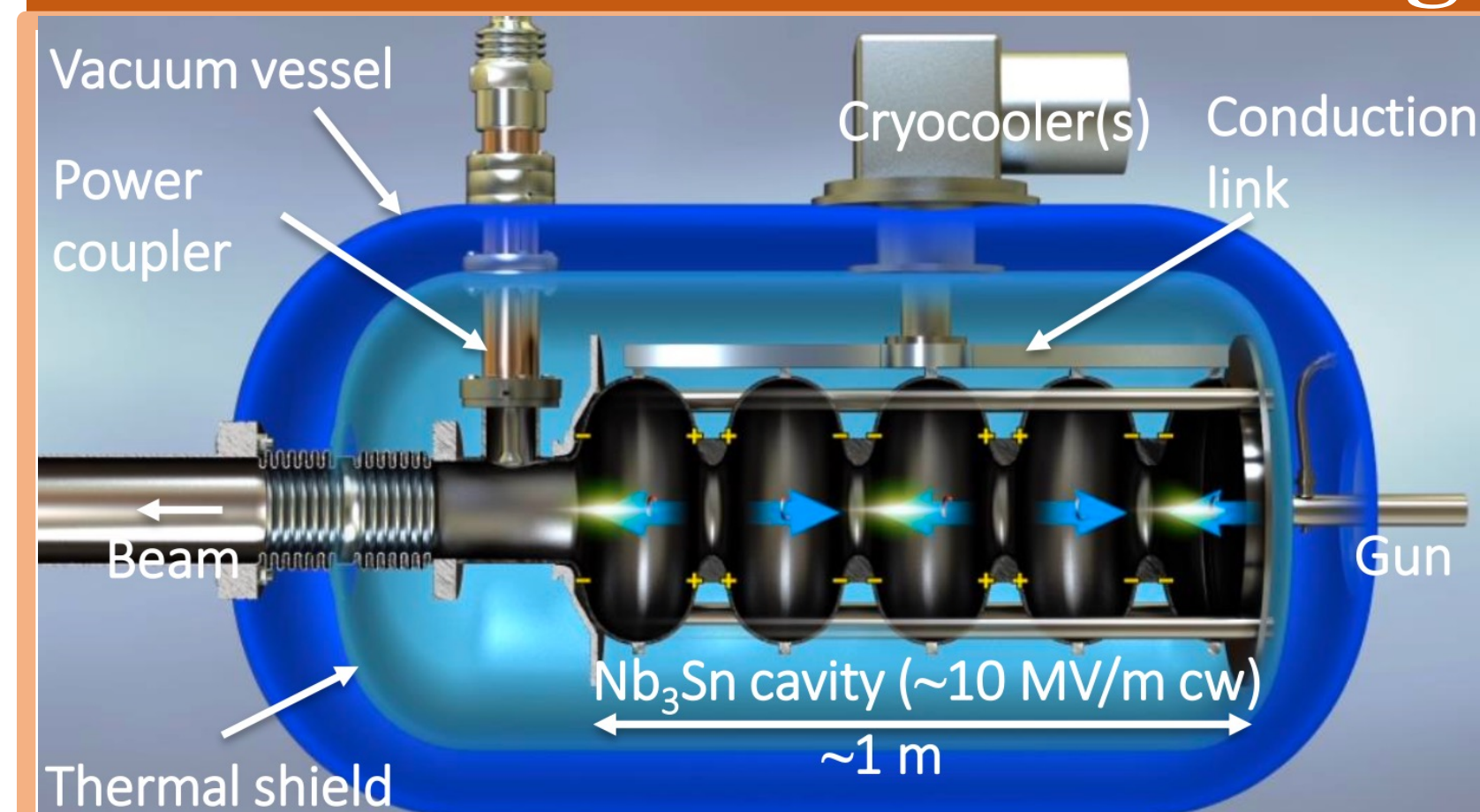
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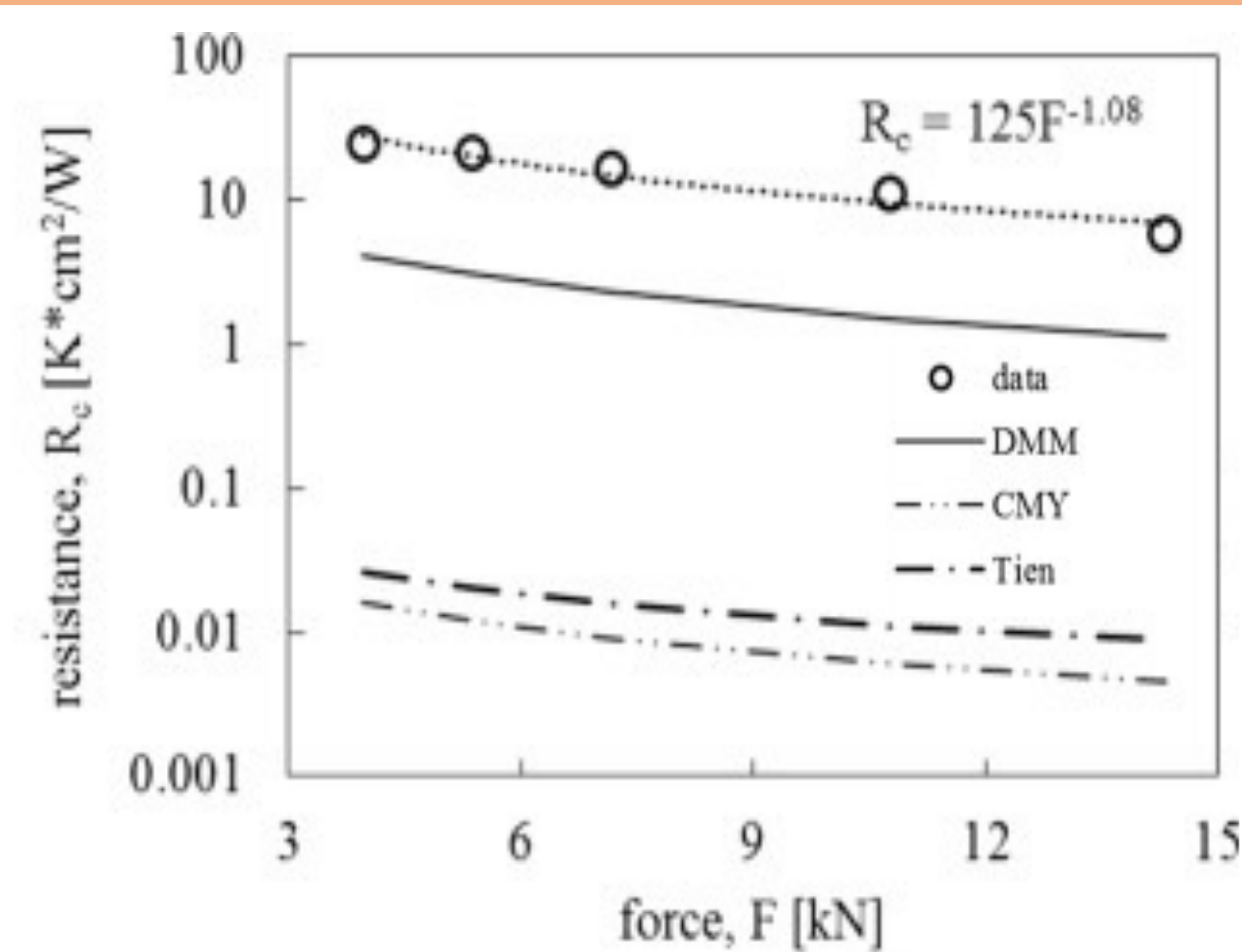
Abstract

Cryocooler-based conduction cooling has gained traction in replacing liquid He baths for cooling in applications such as superconducting radiofrequency cavities. One of the primary governing factors in conduction cooling is thermal contact resistance (R_c) between two different metals, which, for example, varies from 10^{-4} to 10^{-3} m²K/W near 4 K for a bolting force of 3 to 10 kN between high-purity aluminum and niobium. However, the measured R_c is far from the intrinsic R_c expected of metal contact, thus leading to losses at the contact. With the present study, we aim to measure the intrinsic R_c between two distinct materials using time-resolved EUV diffraction measurements. The intrinsic value will serve as a benchmark for evaluating the thermal contact's effectiveness, which is currently lacking. The sensitivity of our setup is 10^{-9} m²K/W and surpasses the existing reported values by order of magnitudes.

Need for conduction cooling and why measure contact resistance



Conduction-cooled SRF accelerator [1]

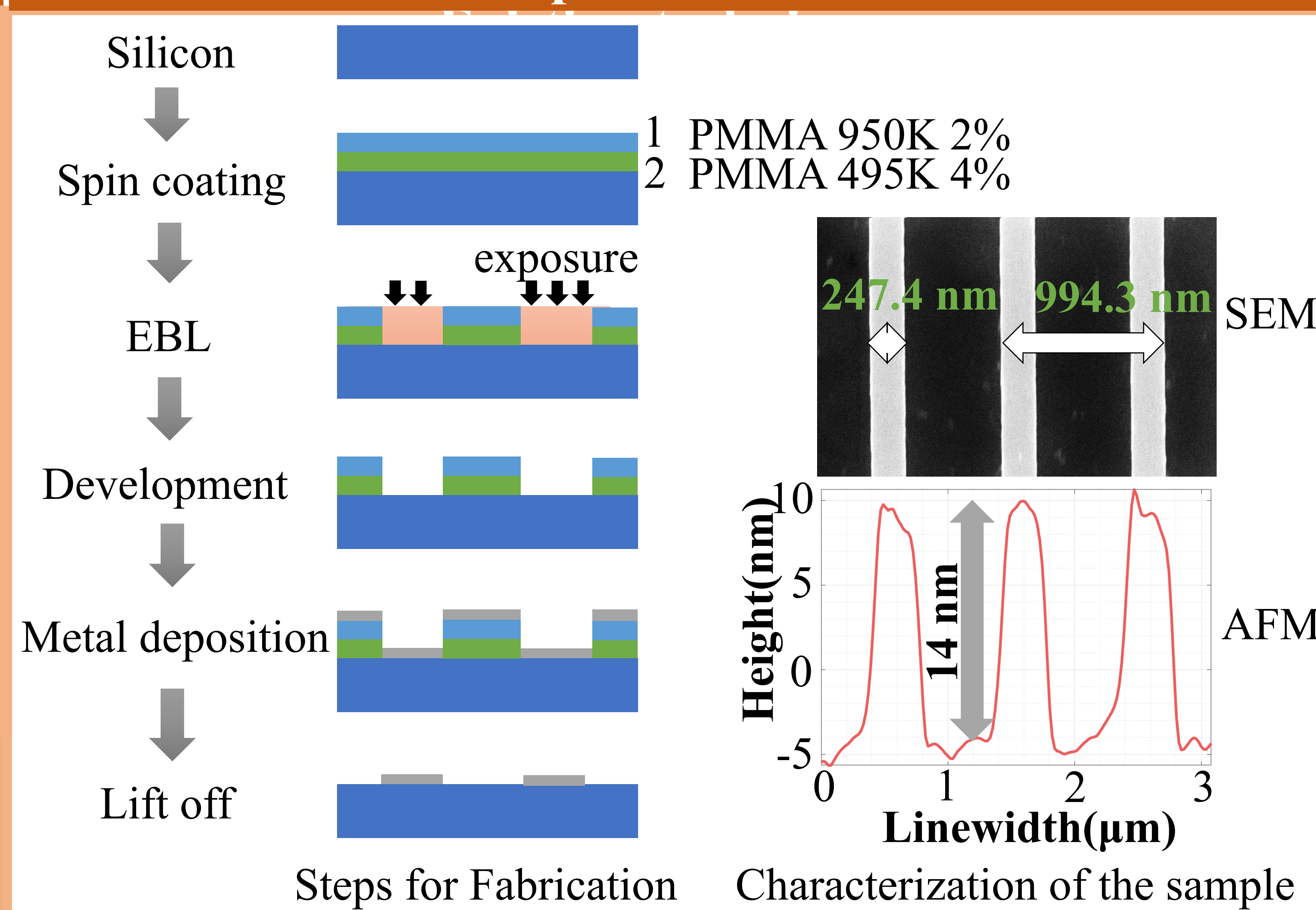


R_c values between Al and Nb [2]

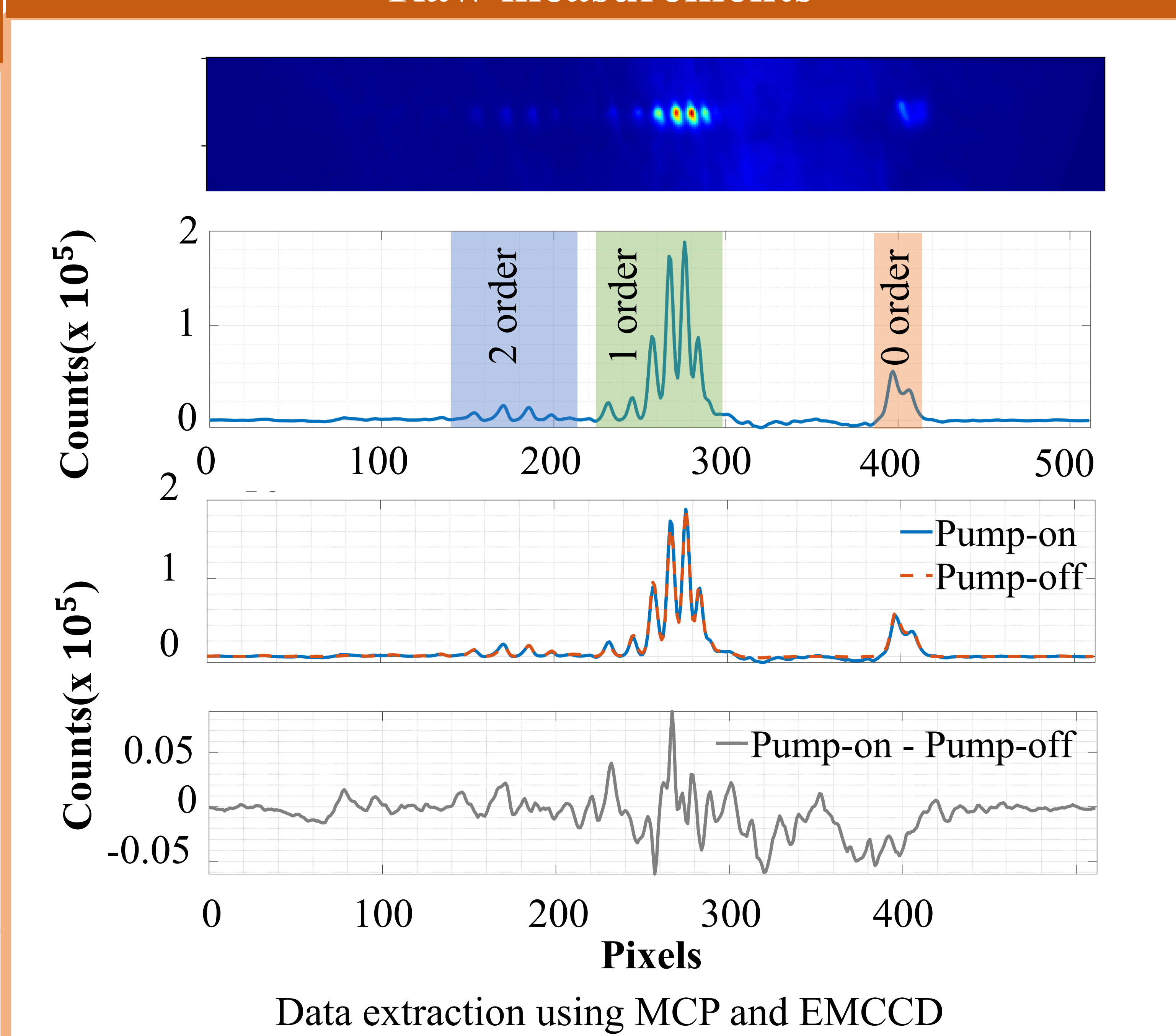
- Intrinsic R_c significantly lower than measured R_c
- Need to determine the effectiveness of contact to identify scope of improvement

R_c measurement between Al and Nb [2]

Sample fabrication



Raw measurements



Conclusion

- Our experiments provide unprecedented high sensitivity in R_c ($\sim 10^{-9}$ m²K/W) that can serve to classify effectiveness of contacts
- Potential strategies to decrease R_c can further improve the efficiency of conduction-cooled SRF cavity

Coupled thermo-mechanical and Fresnel diffraction modeling

$$\nabla \cdot (c : \nabla(u - \alpha \Delta T_p)) = \rho \frac{\partial^2 u}{\partial t^2}$$

$$\rho C_p \frac{\partial T}{\partial t} + \rho C_p \mathbf{u} \cdot \nabla T = \nabla \cdot (K_{Ni} \nabla T) + Q$$

$$\nabla \cdot (c : \nabla(u - \alpha \Delta T_s)) = \rho \frac{\partial^2 u}{\partial t^2} \quad C_s \frac{\partial T_s}{\partial t} + \rho C_s \mathbf{u} \cdot \nabla T_s = \nabla \cdot q$$

KCM [3]:

$$\tau \frac{\partial q}{\partial t} + q + k_{Si} \nabla T_s = l^2 (\nabla^2 q + \alpha_1 \nabla \nabla \cdot q)$$

$$(T_p - T_s) = -R_c q \cdot n + \frac{1}{\gamma_{Si}} (\beta \nabla \cdot q - \nabla q : X)$$

Effective resistivity [4]: $q + k_{Si} \nabla T_s = 0 \rightarrow \frac{T_p - T_s}{R_{c,eff}} = q$

Viscous heat equations [5]:

$$C \frac{\partial T}{\partial t} + \sum_{i,j=1}^3 W_{0j}^1 \sqrt{T} A^j C \frac{\partial u_d^j}{\partial r^i} - \sum_{i,j=1}^3 k^{i,j} \frac{\partial^2 T}{\partial r^i \partial r^j} = 0$$

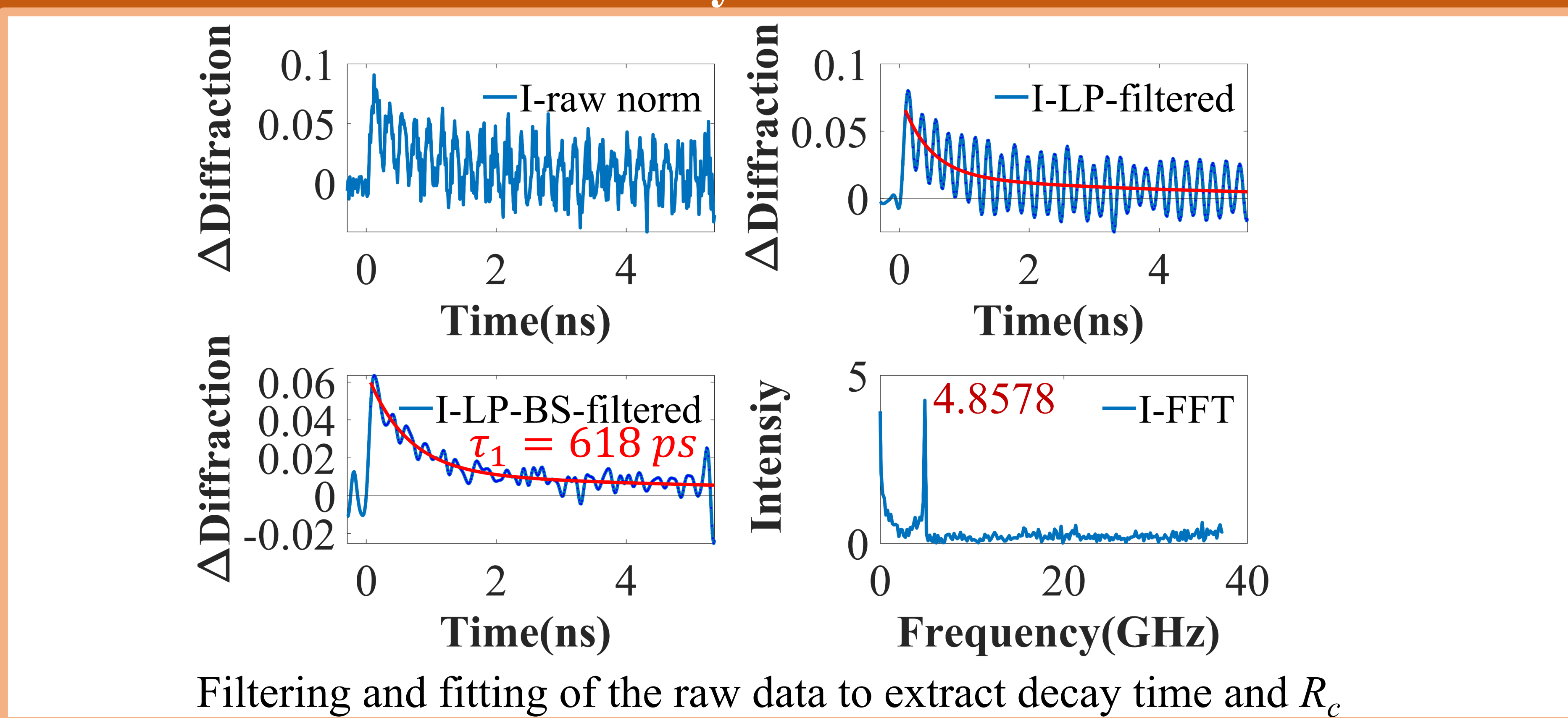
$$A^i \frac{\partial u_d^i}{\partial t} + \sqrt{\frac{CA^i}{T}} \sum_{j=1}^3 W_{i0}^j \frac{\partial T}{\partial r^j} - \sum_{j,k,l=1}^3 \mu^{ijkl} \frac{\partial^2 u_d^k}{\partial r^j \partial r^i} = - \sum_{j=1}^3 \sqrt{A^i A^j} D_U^{ij} u_d^j$$

where, T_p, T_s, T : phononic temperature, u : displacement, c : elasticity tensor, k : bulk thermal conductivity, Q : heat source term, C_p, C_s, C : heat capacity, α : thermal expansion coefficient, α_1 : constant, τ : decay constant, l : non-local length, $R_c, R_{c,eff}$: contact resistivity, W_{0j}^1 : velocity tensor, u_d^j : drift velocity, \bar{T} : reference temperature, A^j : specific momentum, D_U^{ij} : momentum dissipation rate, μ^{ijkl} : thermal viscosity tensor

Fresnel diffraction

$$E(x, y, z) = \frac{e^{ikz}}{i\lambda z} \int \int_{-\infty}^{\infty} E(x', y', 0) e^{\frac{ik}{2z} [(x-x')^2 + (y-y')^2]} dx' dy'$$

Data analysis and Results



References

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- Dhuley *et al.* Cryogenics 93, 86-93, (2018)
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Experimental setup: NIR pump-XUV probe

