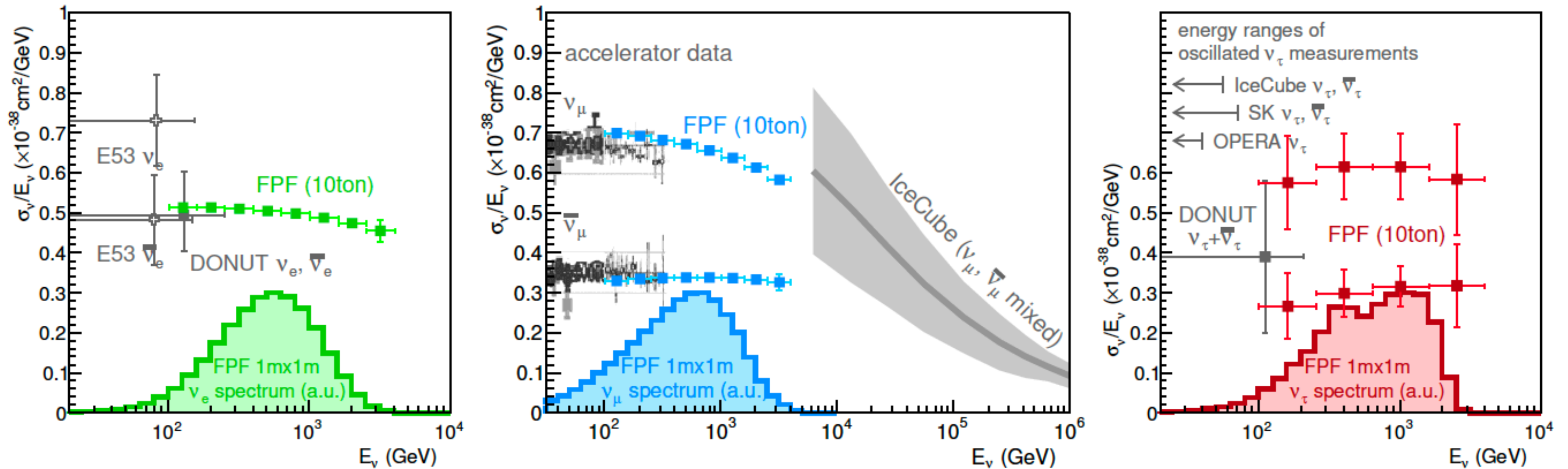


# Neutrino cross sections for shallow- and deep-inelastic scattering

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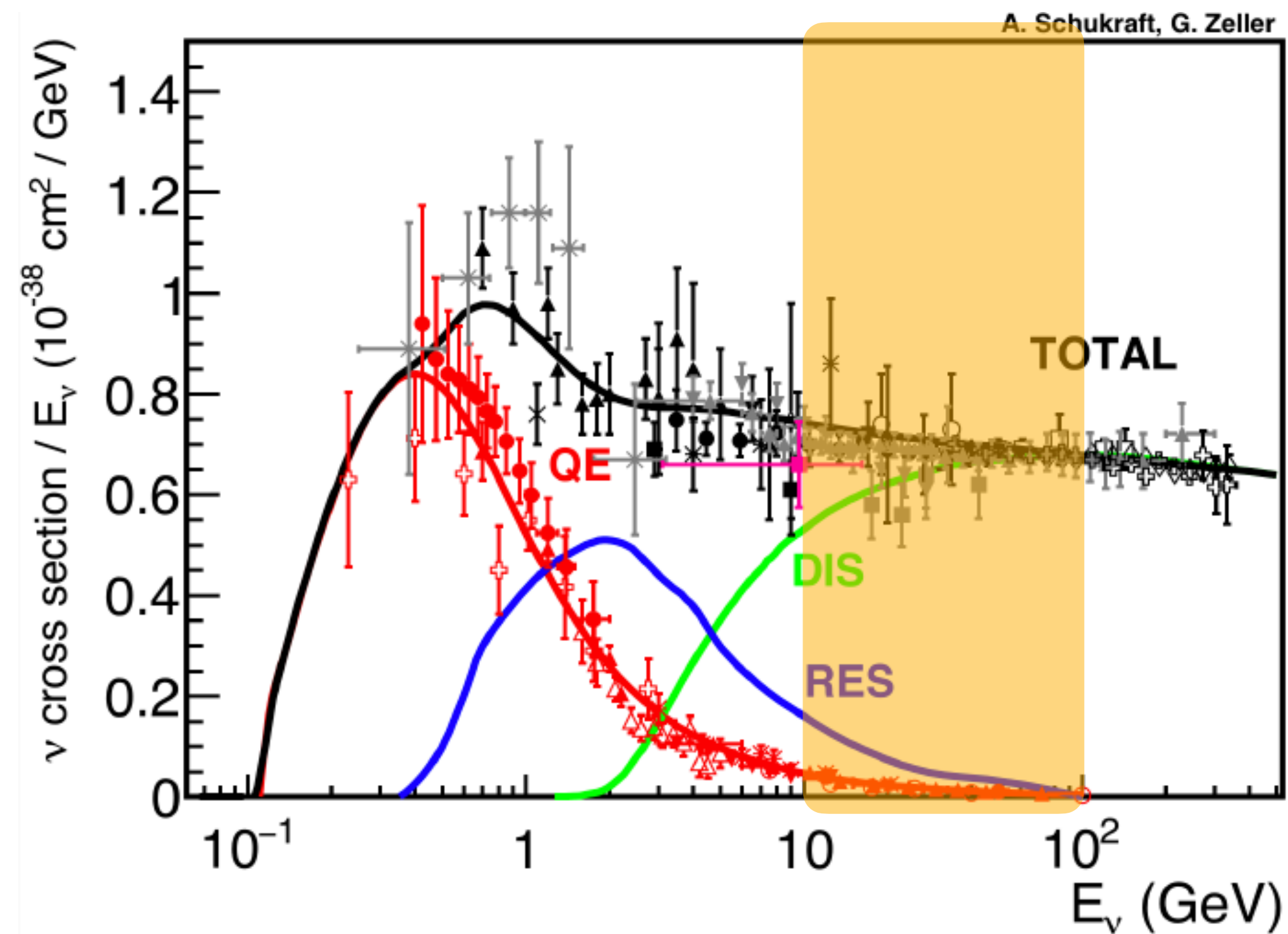
Work with M. H. Reno (Univ. of Iowa) (arXiv: 2307.09241)

# Neutrinos at the FPF



- The LHC produces a large number of neutrinos in the forward direction for all three flavors.
- The neutrinos that reach to the FPF can have energies up to a few TeV.
  - mainly distributed above 100 GeV, but there are considerable amount in 10s GeV region.

# Neutrino interactions



- For  $E > 100$  GeV, neutrinos interact through deep inelastic scattering.
- In the energies of  $10 - 100$  GeV, there exist different types of interactions, so there should be the transition regions.
- It is important to understand the cross sections in the transition regions to avoid double counting.

Detector	CCQE					CCRES					NCEL all	NCRES all
	$\nu_e$	$\bar{\nu}_e$	$\nu_\mu$	$\bar{\nu}_\mu$	$\nu_\tau + \bar{\nu}_\tau$	$\nu_e$	$\bar{\nu}_e$	$\nu_\mu$	$\bar{\nu}_\mu$	$\nu_\tau + \bar{\nu}_\tau$		
FASER $\nu$ 2	60	50	570	350	3.5	170	180	1.6k	1.1k	10	170	1.3k
FLArE	40	40	420	260	3.5	120	140	1.2k	860	10	130	940

Table from FPF White paper: *J.Phys.G* 50 (2023) 3, 030501

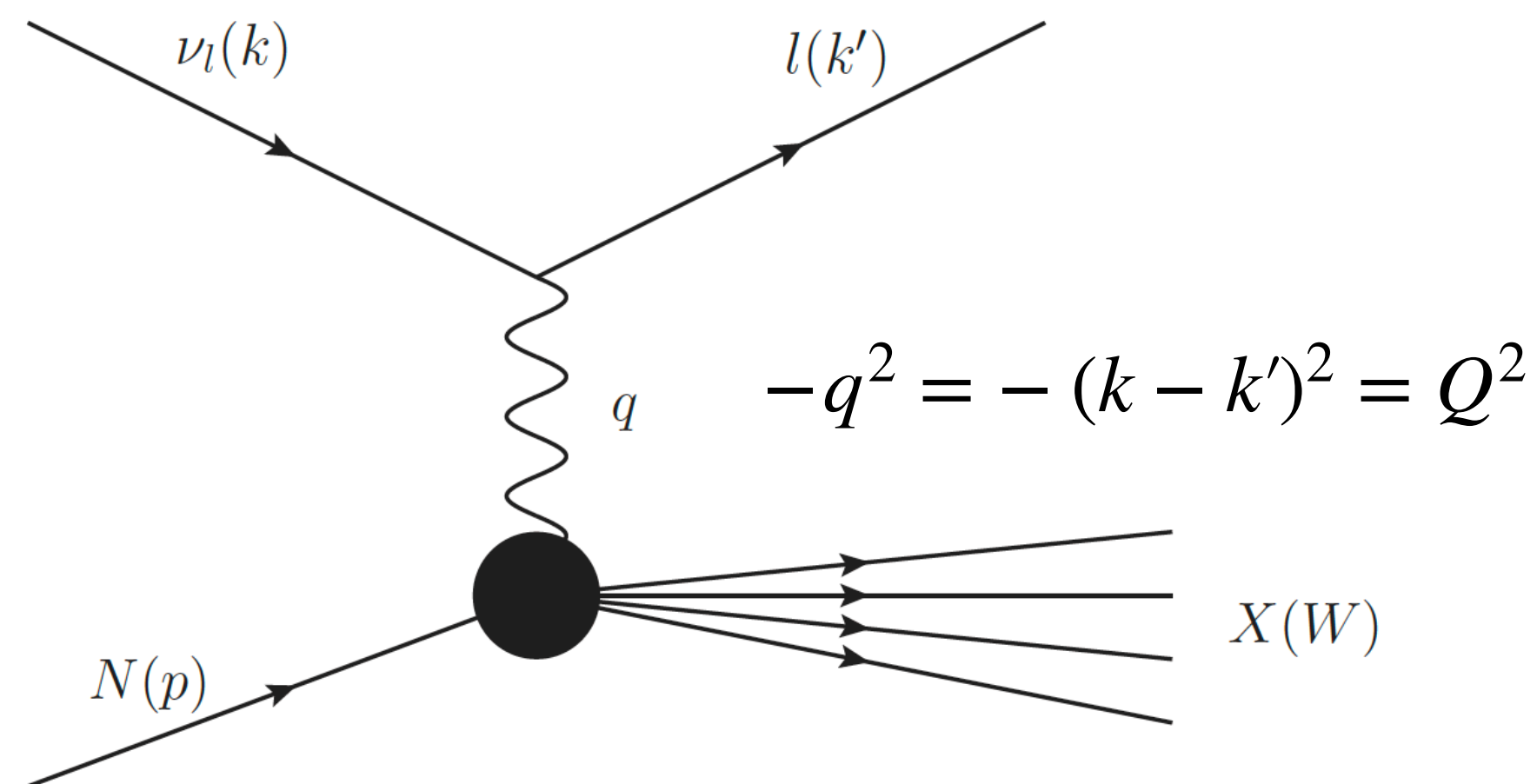
# Deep inelastic scattering cross sections

## ■ Neutrino-nucleon charged-current (CC) cross section for deep inelastic scattering

$$\frac{d^2\sigma^{\nu(\bar{\nu})}}{dx dy} = \frac{G_F^2 m_N E_\nu}{\pi(1 + Q^2/M_W^2)^2} \left( \left( y^2 x + \frac{m_\ell^2 y}{2E_\nu m_N} \right) F_{1,CC} + \left[ \left( 1 - \frac{m_\ell^2}{4E_\nu^2} \right) - \left( 1 + \frac{m_N x}{2E_\nu} \right) y \right] F_{2,CC} \right. \\ \left. \pm \left[ xy \left( 1 - \frac{y}{2} \right) - \frac{m_\ell^2 y}{4E_\nu m_N} \right] F_{3,CC} + \frac{m_\ell^2 (m_\ell^2 + Q^2)}{4E_\nu^2 m_N^2 x} F_{4,CC} - \frac{m_\ell^2}{E_\nu m_N} F_{5,CC} \right)$$

$y = (E - E')/E$

Y. S. Jeong and M. H. Reno, *Phys. Rev. D* 81 (2010) 114012



$$W^2 = Q^2 \left( \frac{1}{x} - 1 \right) + m_N^2$$

## ■ Kinematic region

- DIS:  $W > 2 \text{ GeV}$  and  $Q^2 > 1 \text{ GeV}^2$
- SIS:  $m_N + m_\pi < W \lesssim 2 \text{ GeV}$ , all  $Q^2$



# Structure functions $F_i(x, Q^2)$

## ■ Structure functions are:

- essential components in evaluating the DIS cross section.
- expressed in terms of Bjorken- $x$  and the parton distribution functions (PDF)  $q(x, Q^2)$

$$\text{e.g. } F_{2,CC}(x, Q^2) = \sum_{q,q'} 2x(q(x, Q^2) + \bar{q}'(x, Q^2)) \quad (\text{at LO in pQCD})$$

- not reliable for  $Q^2 < 1 \text{ GeV}^2$ , where the perturbative QCD is not applicable.

## ■ Non-perturbative structure functions at low $Q$ ( $Q^2 < 1 \text{ GeV}^2$ )

- phenomenologically constructed by fitting to the data.
- (e.g.) Bodek-Yang, CKMT, NNSFV

■ Parameterization from fits to the electromagnetic structure function data.

$$F_2^{\text{CKMT}}(x, Q^2) = Ax^{-\Delta(Q^2)}(1-x)^{n(Q^2)+4} \left( \frac{Q^2}{Q^2+a} \right)^{1+\Delta(Q^2)} + Bx^{1-\alpha_R}(1-x)^{n(Q^2)} \left( \frac{Q^2}{Q^2+b} \right)^{\alpha_R} (1+f(1-x))$$

$$n(Q^2) = \frac{3}{2} \left( 1 + \frac{Q^2}{Q^2+c} \right)$$

$$\Delta(Q^2) = \Delta_0 \left( 1 + \frac{2Q^2}{Q^2+d} \right)$$

■ Using the same functional form, normalization parameters are modified for neutrino-nucleon charged-current scattering.

$\Delta_0$	$\alpha_R$	$a$ [GeV <sup>2</sup> ]	$b$ [GeV <sup>2</sup> ]	$c$ [GeV <sup>2</sup> ]	$d$ [GeV <sup>2</sup> ]
0.07684	0.4150	0.2631	0.6452	3.5489	1.1170
Process	$A$	$B$	$f$		
EM $F_2$	0.1502	1.2064	0.15		
$\nu N$ $F_2$	0.5967	2.7145	0.5962		
$\nu N$ $xF_3$	$9.3955 \times 10^{-3}$	2.4677	0.5962		
$\bar{\nu} N$ $xF_3$	$9.3955 \times 10^{-3}$	-2.4677	0.5962		

M. H. Reno, *Phys. Rev. D*74 (2006) 033001

Y. S. Jeong and M. H. Reno, *arXiv:2307.09241*

# CKMT + PCAC

- Partially conserved axial-vector current (PCAC) gives non-zero structure function at low  $Q^2$  limit through  $F_L$ .
- According to Kulagin and Petti for the PCAC corrections to  $F_L$  at low  $Q^2$ :

S. A. Kulagin, and R. Petti, *Phys. Rev. D* 76, 094023 (2007)

$$F_L^{\text{PCAC}} = \frac{f_\pi^2 \sigma_\pi(W^2)}{\pi} f_{\text{PCAC}}(Q^2)$$

$$\sigma_\pi \simeq X(W^2)^\epsilon + Y(W^2)^{-\eta_1}, \quad f_{\text{PCAC}}(Q^2) = \left(1 + \frac{Q^2}{M_{\text{PCAC}}^2}\right)^{-2} \quad \text{with } M_{\text{PCAC}} = 0.8 \text{ GeV}$$

- We further implemented PCAC correction approximately to CKMT parameterizations.

# Other structure functions at low Q

## ■ Bodek–Yang

A. Bodek, I. Park, U. K. Yang,  
*Nucl. Phys. B Proc. Suppl.* 139, 113 (2005)

- provides the effective PDFs at low  $Q^2$  to construct structure functions.
- GRV98 LO PDF + charged lepton DIS data

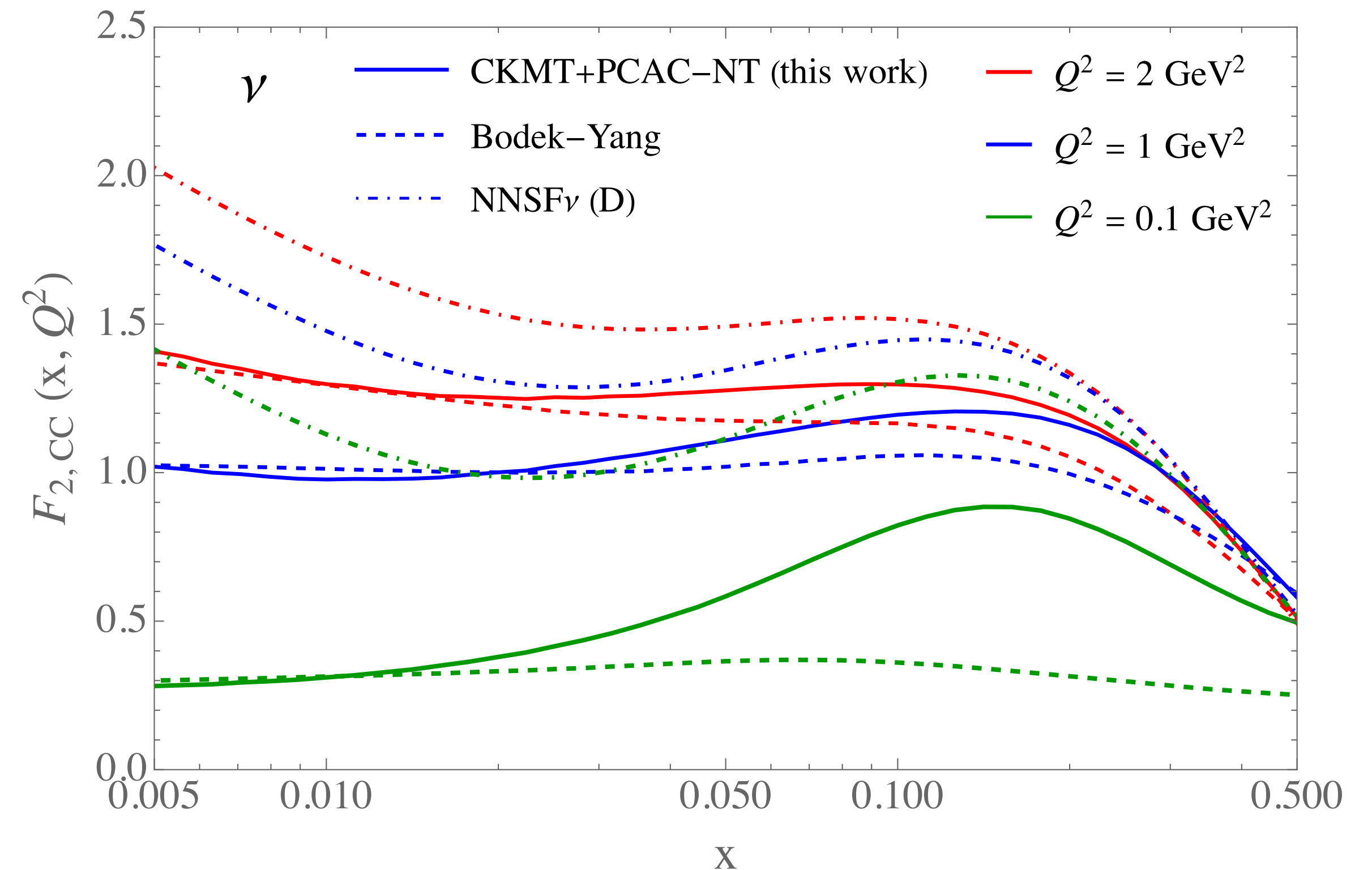
$$F_2(x, Q^2 < 0.8 \text{ GeV}^2) = K(Q^2) \times F_2(\xi_\omega, Q^2 = 0.8 \text{ GeV}^2)$$

## ■ NNSFv

A. Candido, A. Garcia, G. Magni, T. Rabemananjara,  
J. Rojo, and R. Stegeman, *JHEP* 05 (2023) 149

- machine learning parameterization fitted to neutrino scattering data.

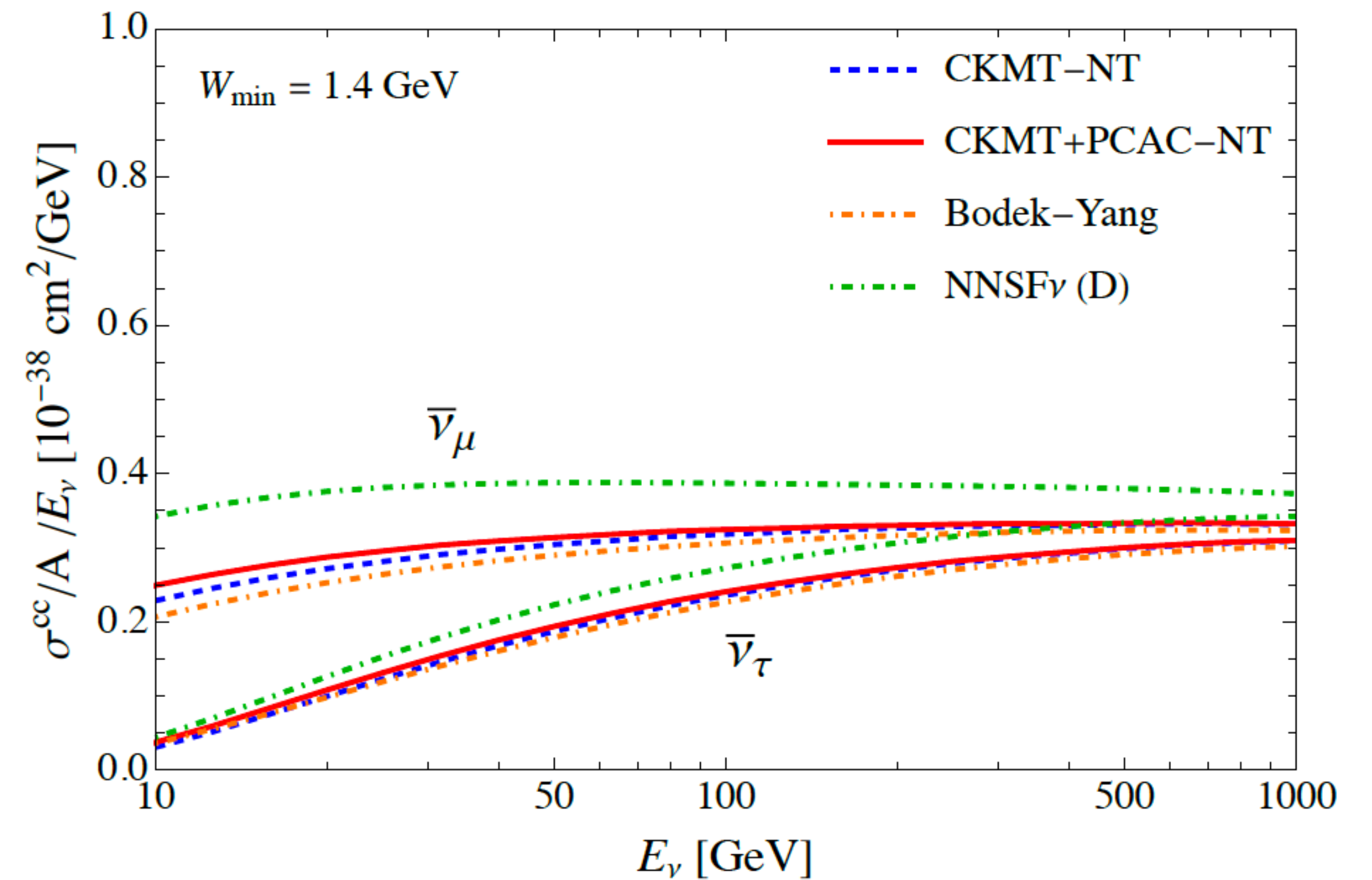
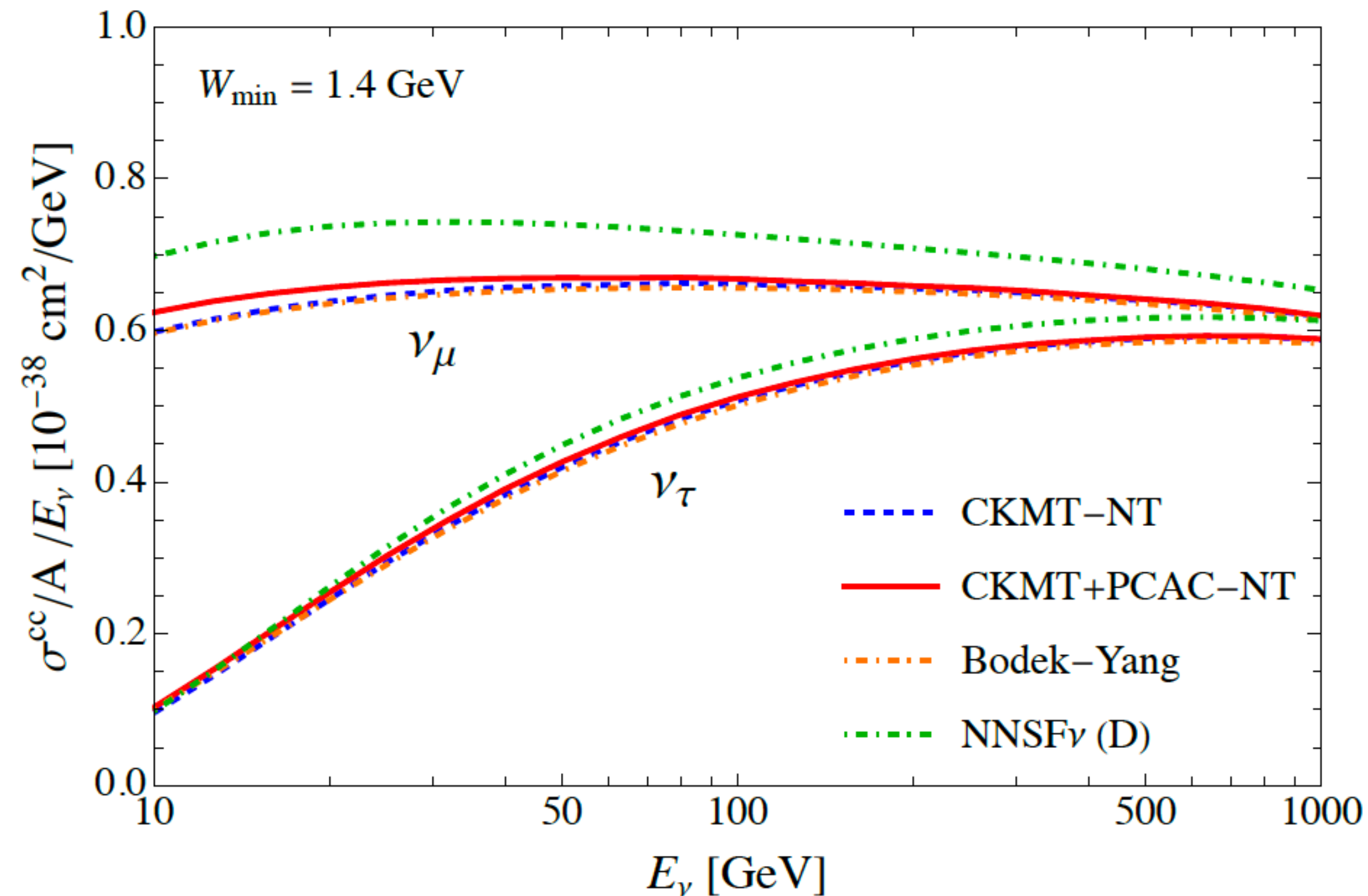
See Tanjona's talk for NNSFv



\* NT: NLO+TMC

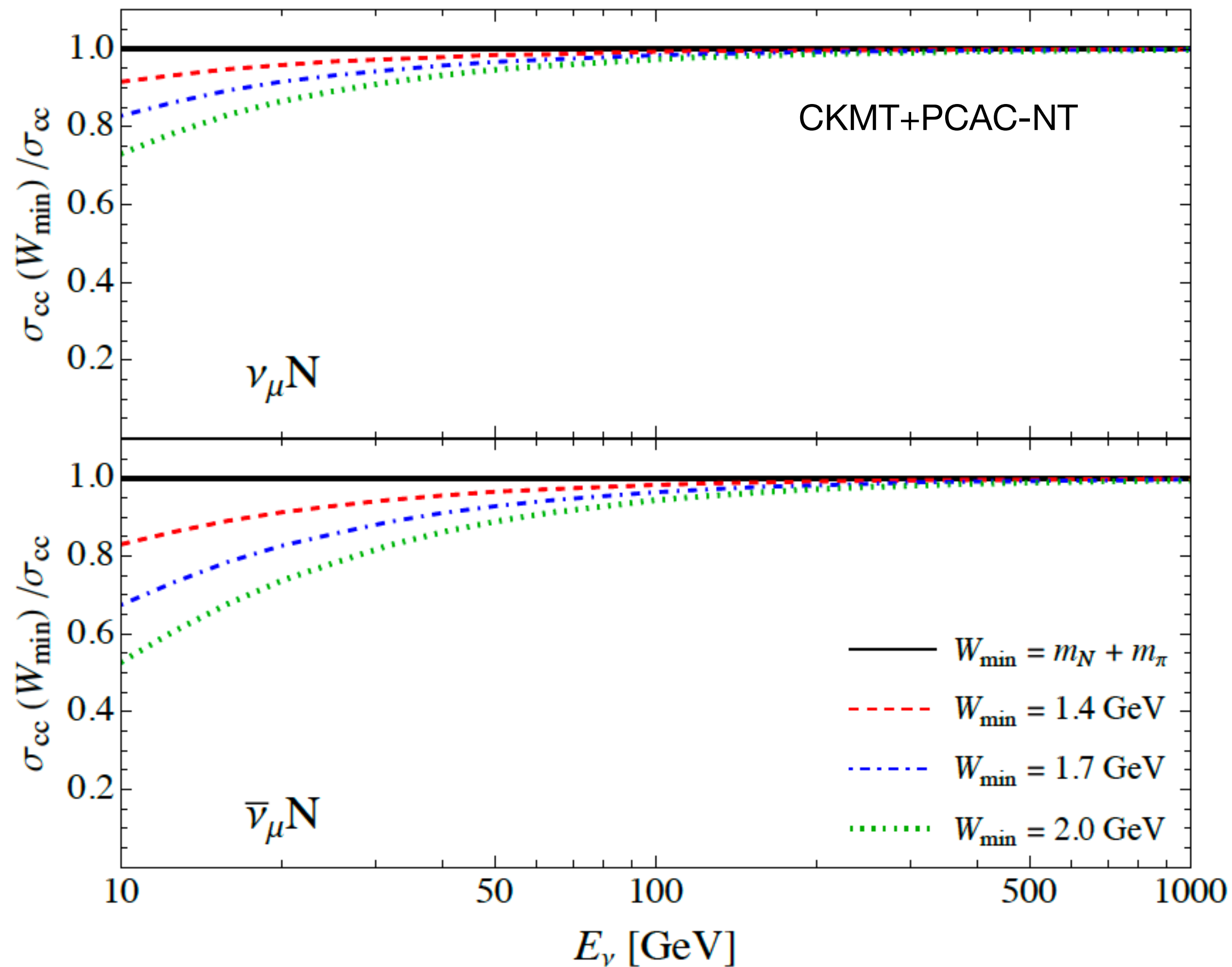


# The neutrino-nucleon DIS cross sections



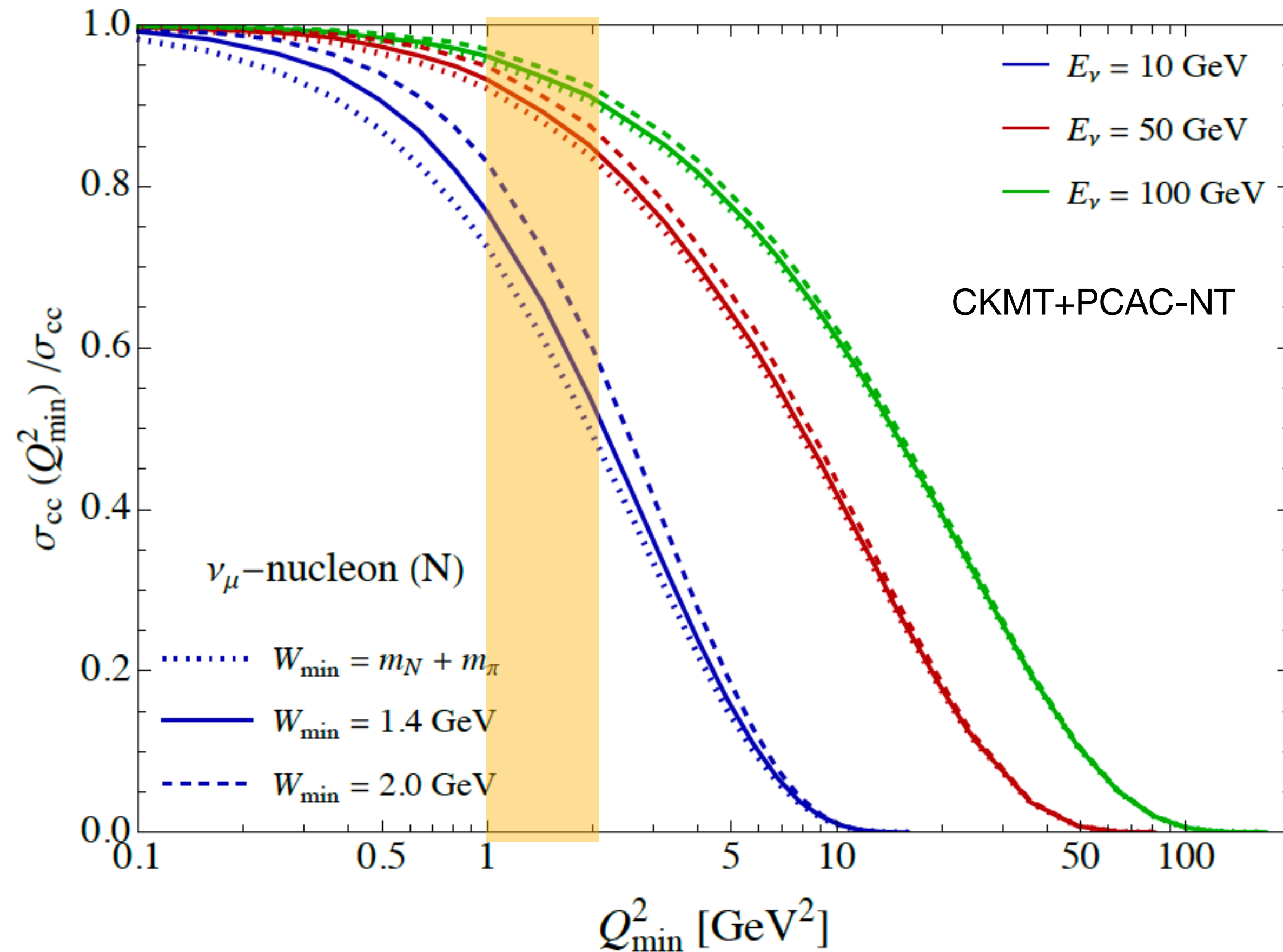
- Cross sections with CKMT+PCAC structure functions are similar to the Bodek-Yang results with about 10% difference at 10 GeV and closer at higher energies.
- NNSFnu results are much larger than the CKMT and BY across the energies.

# Impact of the $W_{\min}$



- The impact of  $W_{\min}$  appears below a few hundred GeV, and more significant at lower energies.
- The hadronic final states between RES and DIS ( $m_N + m_{\pi} < W_{\min} < 2$  GeV) can affect the cross sections by a few % at 100 GeV and about 30 % at 10 GeV and have larger impact on antineutrinos.
- These results with  $\sigma$ -ratio are approximately the same for nuclear targets.

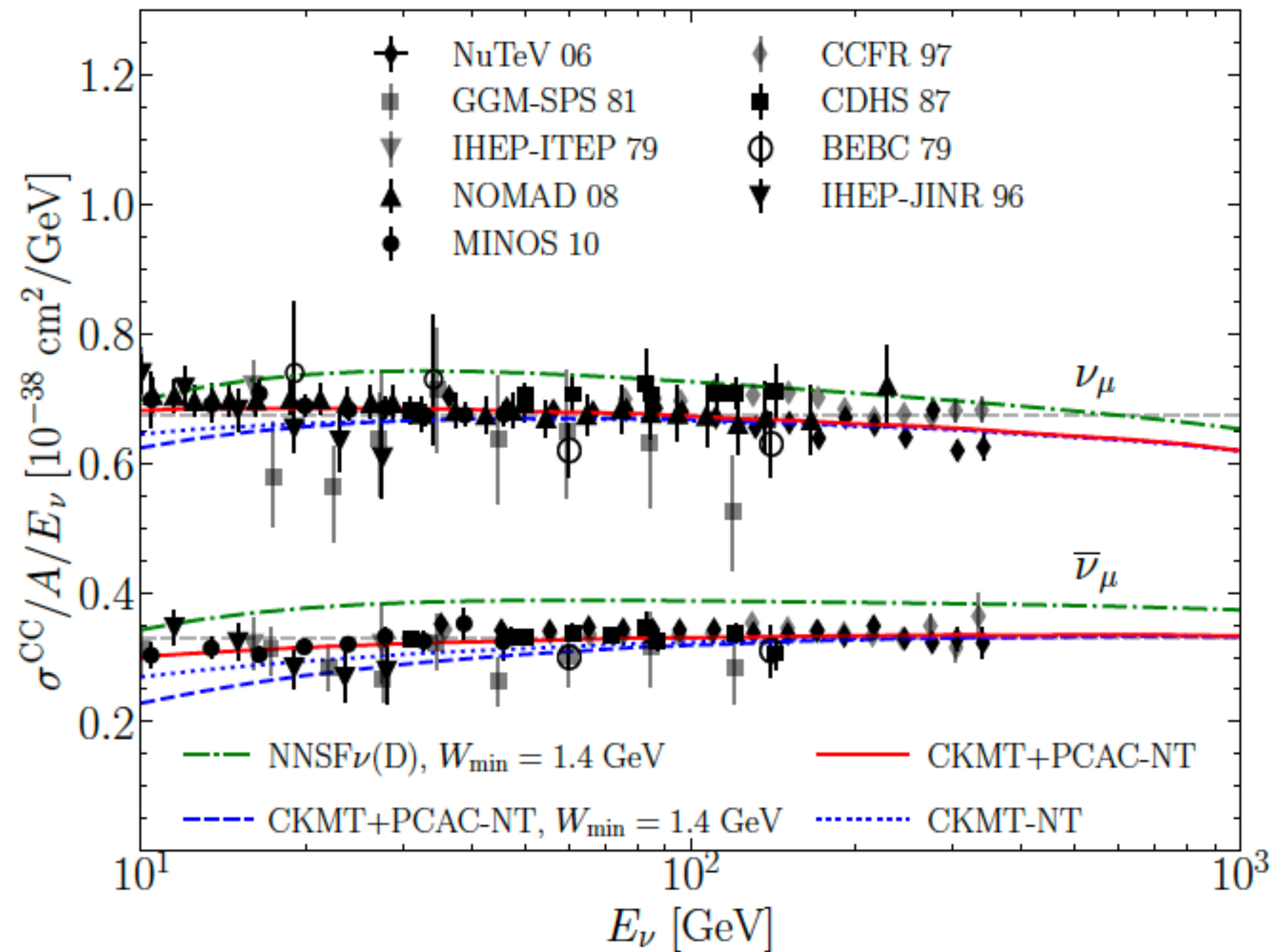
# Impact of the $Q_{min}^2$ and $W_{min}$



- The impact of  $Q_{min}^2$  and  $W_{min}$  is more evident at low energies.
- (e.g.) For  $W_{min} = 1.4$  GeV, the contribution from  $Q_{min}^2 < 1$  GeV<sup>2</sup> to the CC cross section for muon neutrino-nucleon scattering:
  - is less than 5 % at  $E_{\nu} = 100$  GeV
  - increases to about 25 % at  $E_{\nu} = 10$  GeV.



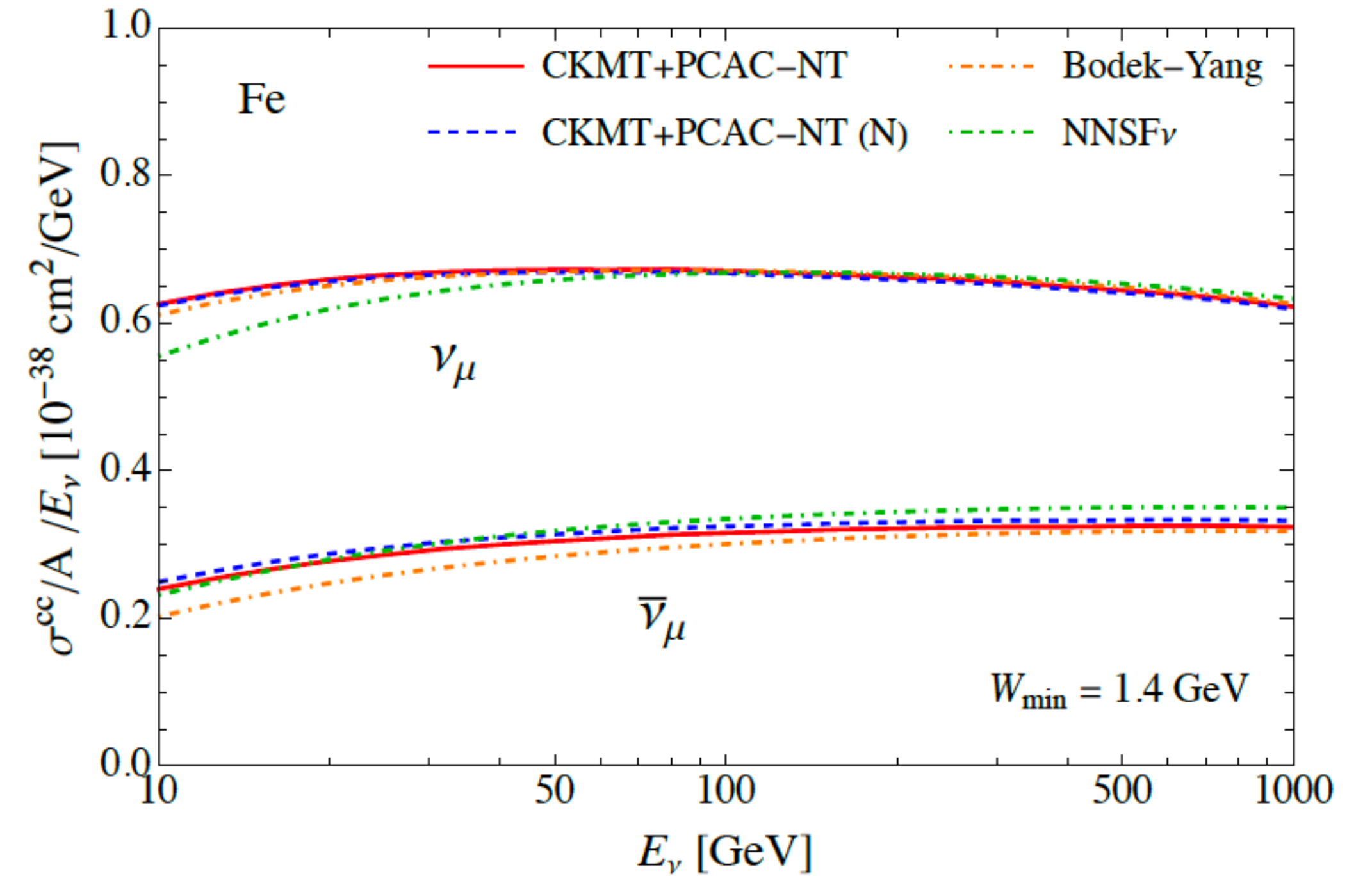
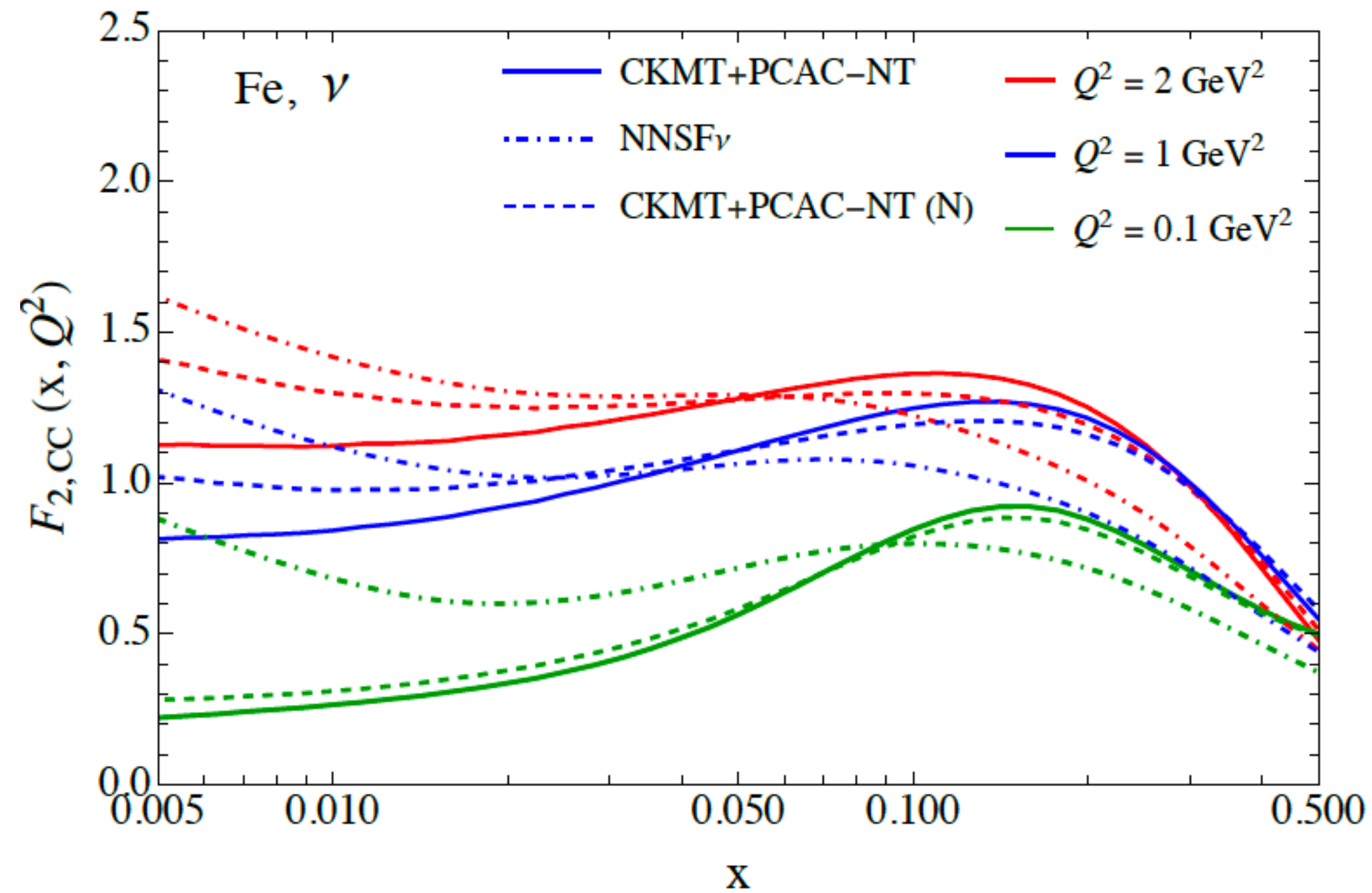
# Neutrino cross sections for $\nu_\mu / \bar{\nu}_\mu$ - nucleon scattering



- Cross sections per nucleon normalized with  $E_\nu$
- Red: CKMT+PCAC-NT with  $W_{\min} = m_N + m_\pi$ 
  - Account for all CC cross sections for inelastic collisions beyond quasi-elastic scattering.
  - Very well matched to the experimental data.
- PCAC improves the prediction to have a better agreement with the data.

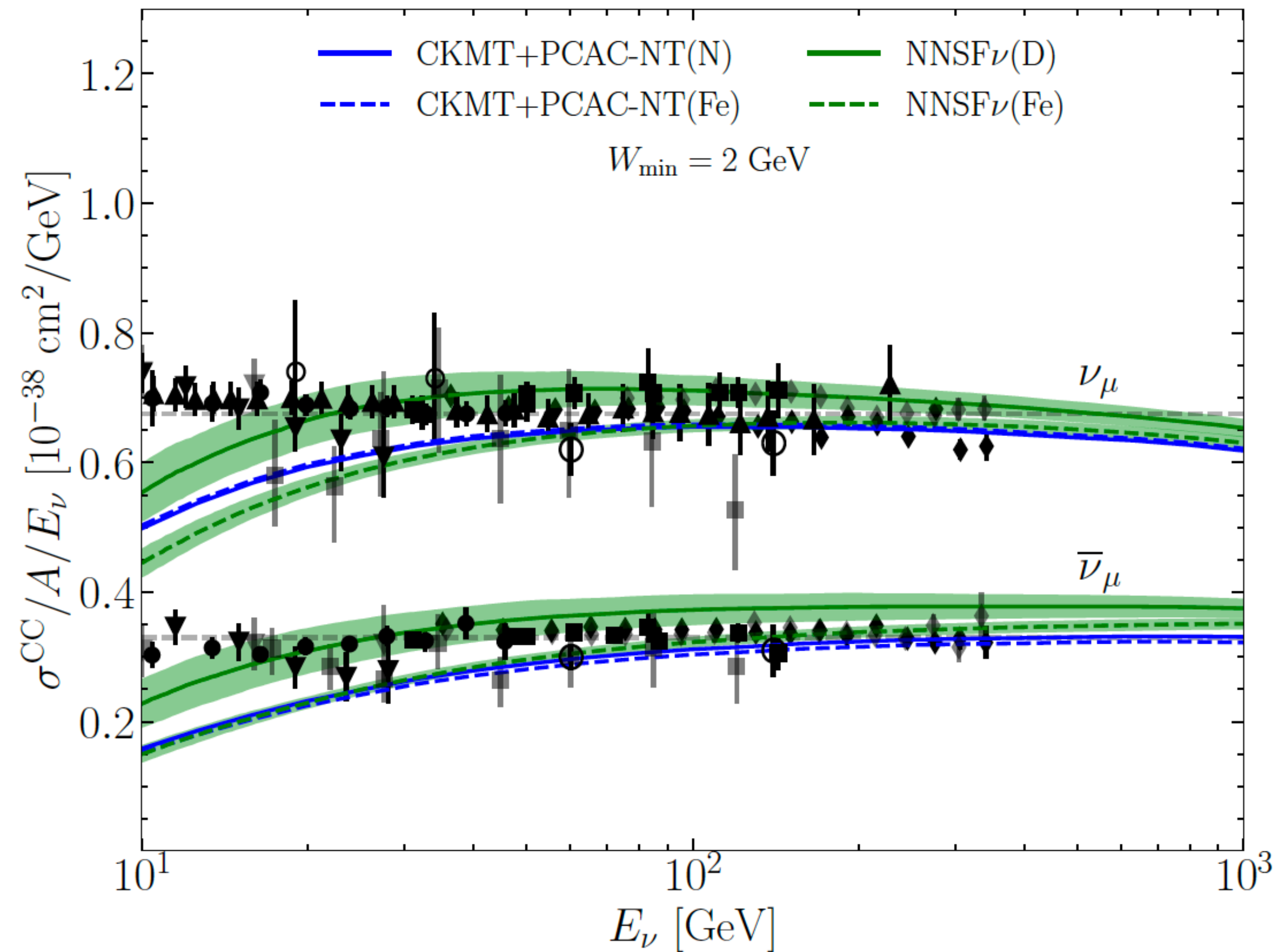


# More comparison with NNSF $\nu$ for Fe target



- For heavier nuclear target, disagreement between NNSF $\nu$  and CKMT based result is relieved compared to the nucleon target.

# More comparison with NNSF $\nu$ ( $W_{\min} = 2$ GeV)



- There is still disagreement in the CC cross sections for neutrino scattering with isoscalar nucleon with  $W_{\min} = 2$  GeV.

# Summary

- We have evaluated neutrino CC DIS cross sections using the structure functions at low  $Q^2$  and investigated the impact of low  $Q^2$  and low  $W$  outside DIS interaction region.
- Although the impact of low  $W$  and  $Q^2$  is negligible for the main energy region of the FPF, it appears below 100 GeV, where the FPF could obtain sufficient events for interaction study.
- There is discrepancy between the low  $Q$  structure functions and cross sections (esp. NNSFV).
- More theoretical study on structure functions for low  $Q$ /cross sections for low  $E$  will be helpful. Also, the FPF could provide the useful data for interaction study in the transition region.

*Thank you for your attention*