

Forward D meson production

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2108.03741, 2109.10905, 2203.05090

HF at hadron colliders

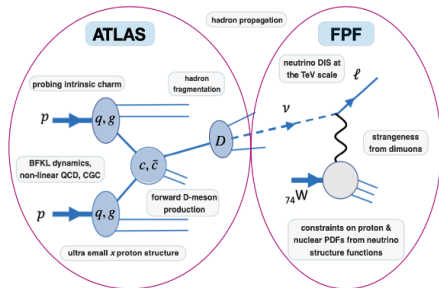
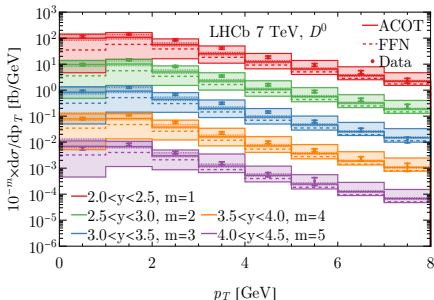
Data:

- Heavy-flavor hadron (D, B -meson) production, especially LHCb
- $Z + b/c$ production [See Boettcher's talk]
- Neutrino resource measured at the FASER as well as other FPFs

Theoretical interests:

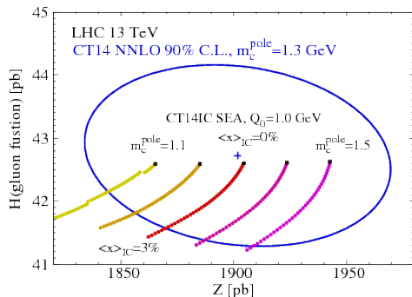
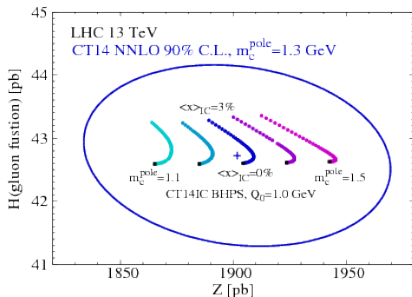
- pQCD: factorization theorem, scale uncertainty, fragmentation, etc.
- PDF: Forward heavy flavor production probes gluon PDF at small x .

$$x \sim \frac{m}{\sqrt{s}} e^{-y} \sim 10^{-6}$$



The importance of HF PDFs

- Intrinsic vs extrinsic, *i.e.*, perturbative vs non-perturbative (fitted)? [See Hobbs' talk]
- Can data tell the difference?
- Heavy flavor mass: dynamics (ME) and kinematics (phase space or threshold)
- Multiple scales: $Q(p_T)$ vs m_Q . PDF resums large logarithms $\alpha_s \log(Q^2/m_Q^2)$

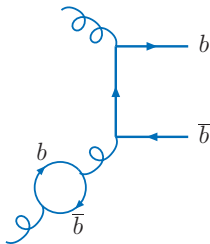


[1707.00657]

Massive Fixed-Flavor-Number (FFN) scheme

For consistency, we should take $N_f = 3(4)$ for charm(bottom) flavor production, in both α_s and PDF running.

- The heavy-quark running in the virtual loops is missing.
- No Flavor Excitation (FE) contributions as no heavy-flavor PDF.



Inconsistency when using $N_f = 5$ PDF in MCFM, MadGraph_aMC@NLO, POWHEG,

- $N_f = 5$ in the α_s running, e.g. reading directly from LHAPDF;
- No FE contributions, equivalent to $N_f = 3(4)$ in the PDFs.

We need treat heavy flavor consistently.

Theory for heavy-flavor production

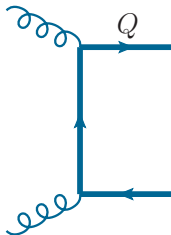
Energy scale Q , such as invariant mass M_{QQ} or p_T

- $Q \lesssim m$ (low energy), Flavor Creation (FC), massive FFN scheme (N_f)
- $Q \gg m$ (high energy), Flavor Excitation (FE), Zero-mass (ZM) scheme ($N_f + 1$), resum $\alpha_s^m \log^n(Q^2/m^2)$ as heavy-flavor PDF (massless)
- $Q \sim m$, General-mass (GM) variable flavor number (VFN) scheme

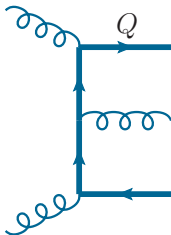
matching: subtracting the double-counted terms

$$\text{VFN} = \text{FC} + \text{FE} - \text{SB}$$

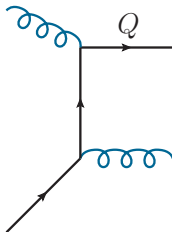
- $Q \lesssim m$, $\text{FE} \simeq \text{SB}$, $\text{VFN} \rightarrow \text{FC}$ FFN scheme
- $Q \gg m$, $\text{FC} \simeq \text{SB}$, $\text{VFN} \rightarrow \text{FC}$ ZM scheme



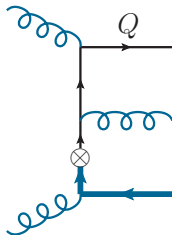
(a) FC LO



(b) FC NLO

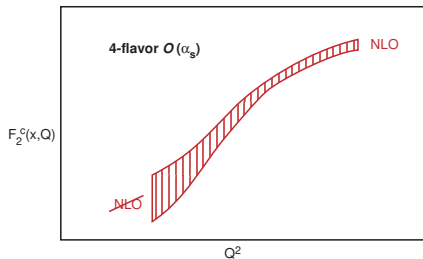
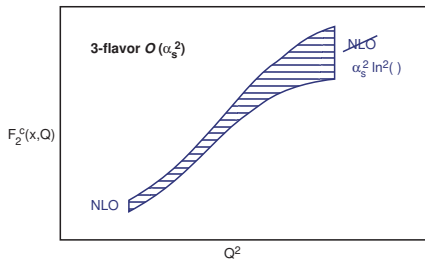
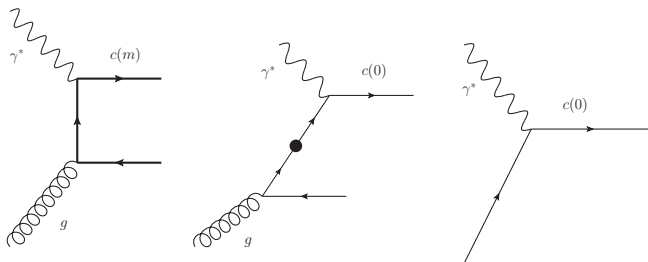


(c) FE



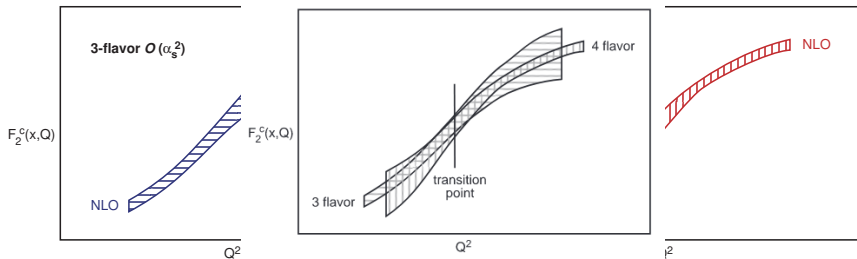
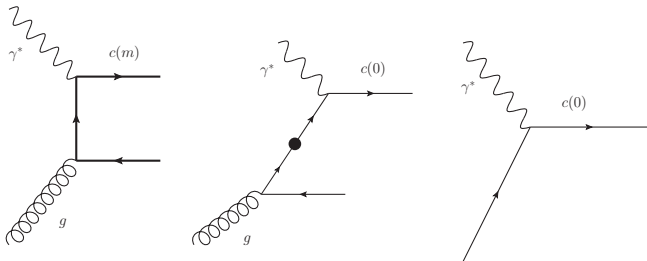
(d) SB

ACOT scheme



[W. Tung, et al., 0110247]

ACOT scheme



[W. Tung, et al., 0110247]

- $Q \gtrsim m_Q$, m_Q matters, $f_Q(x, \mu) \approx 0$, Flavor Creation (FFN 3-flv).
- $Q \gg m_Q$, $m_Q \approx 0$, $f_Q(x, \mu)$ matters, Flavor Excitation (ZM 4-flv).

ACOT series

- Aivazis-Collins-Olness-Tung [PRD1994] introduce an asymptotic subtraction (SB) term to get rid of the double-counting between Flavor Creation (FC) and Flavor Excitation (FE), which switches from N_f to $N_f + 1$ scheme (Variable Flavor Number Scheme).

$$\text{ACOT} = \text{FC} - \text{SB} + \text{FE}$$

- $Q \gtrsim m_Q$, $\text{SB} \simeq \text{FE}$, $\text{ACOT} \rightarrow \text{FFN 3-flv scheme}$;
- $Q \gg m_Q$, $\text{SB} \simeq \text{FC}$, $\text{ACOT} \rightarrow \text{ZM 4-flv scheme}$.
- Simplified-ACOT scheme [J. Collins PRD1998, M. Kramer et al., PRD2000] treats heavy-quark as massless in Flavor Excitation. Warning: instability in the cancellation between SB and FE around the switching point.
- The S-ACOT- χ scheme [W. Tung et al., 0110247] introduces rescaling variable $\chi = x(1 + 4m_Q^2/Q^2)$ to capture the mass threshold effect. It stabilizes the perturbative convergence near the switching point by enforcing energy-momentum conservation in all scattering contributions.
- The S-ACOT- m_T scheme [I. Helenius et al., 1804.03557]
- The S-ACOT-MPS [K. Xie et al., 2108.03741] scheme extends the S-ACOT- χ method to hadron-hadron collisions.

Formulation of the S-ACOT-MPS scheme

- FC+FE-SB

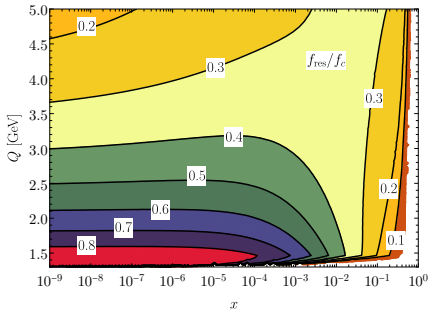
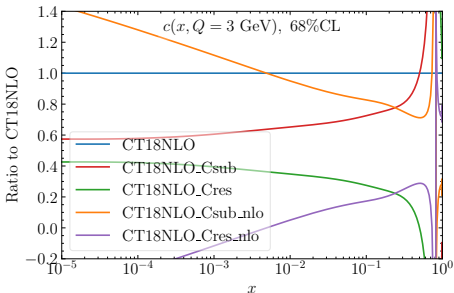
$$\sigma_{\text{FC}} = \sum_{i,j} f_i(x_i, \mu^2) f_j(x_j, \mu^2) \hat{\sigma}_{ij \rightarrow QX},$$

$$\sigma_{\text{FE}} = \sum_i f_i(x_i, \mu^2) f_Q(x_Q, \mu^2) \hat{\sigma}_{iQ \rightarrow QX} + (i \leftrightarrow Q),$$

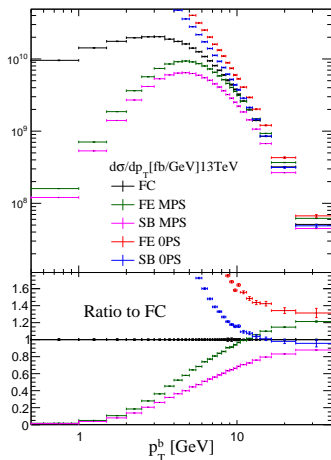
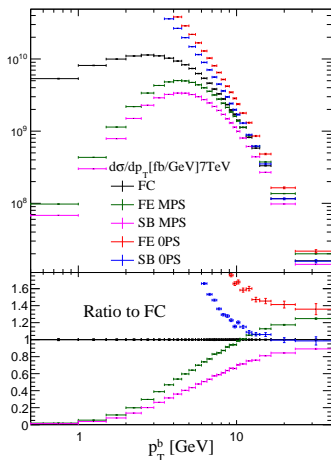
$$\sigma_{\text{SB}} = \sum_{i,j} f_i(x_i, \mu^2) [P_{Qj} \otimes f_j](x_Q, \mu^2) \hat{\sigma}_{iQ \rightarrow QX} + (i \leftrightarrow Q).$$

- We can define the subtracted and residual PDFs

$$\tilde{f}_Q(x, \mu^2) = \sum_j [P_{Qj} \otimes f_j](x, \mu^2), \quad \delta f_Q(x, \mu^2) = [f_Q - \tilde{f}_Q](x, \mu^2)$$



The massive phase space



Caveat: The Lorentz violation for the heavy parton

$$p_b = xp_{\text{proton}} : p_{\text{proton}}^2 = 0 \leftrightarrow p_b^2 = m_b^2.$$

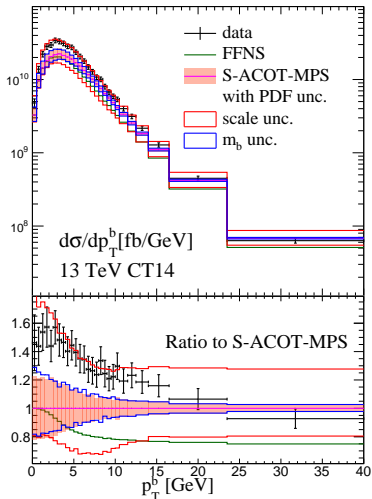
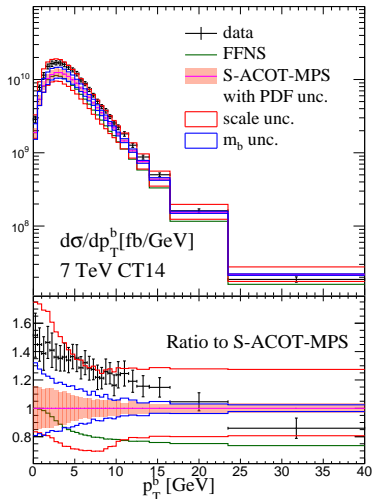
We enforce $E_b = xE_{\text{beam}} > m_b$.

A correction term $\mathcal{O}(m_b^2/Q^2)$ needs to be got back order by order.

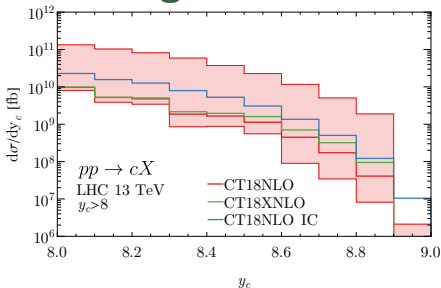
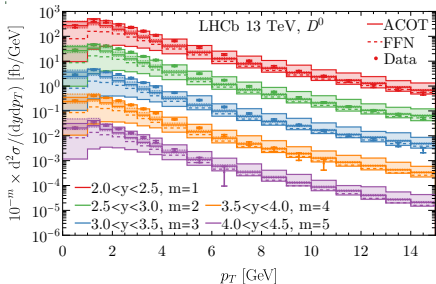
Bottom production at LHCb

Scale $(\mu_R, \mu_F) = (1/2, 1, 2)\sqrt{p_{T,b}^2 + m_b^2}$ uncertainty is large:

- $\alpha_s(\mu_R)$ is large and varies drastically around $\mu_R \sim m_Q$,
- Heavy-flavor PDF $f_Q(x, \mu_F)$ starts to be generated perturbatively at $\mu_F = m_Q$.

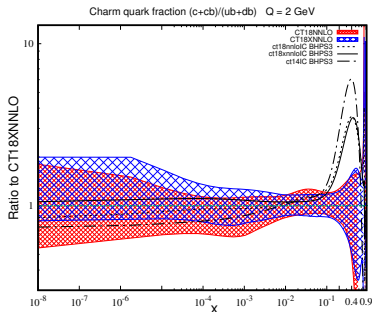


Charm production in the forward region



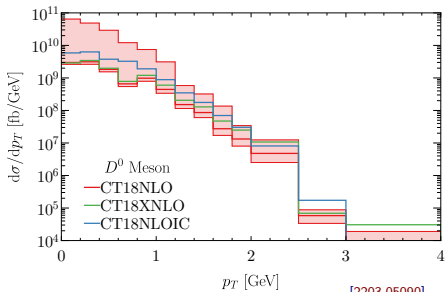
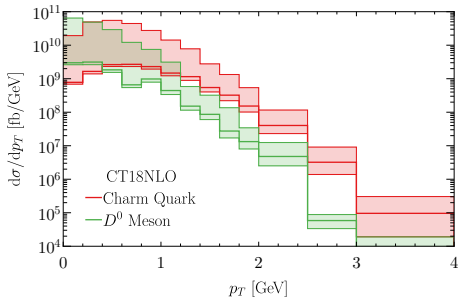
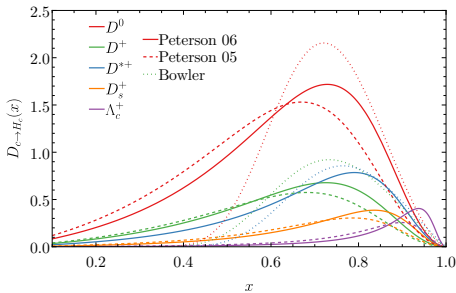
[2108.03741]

[2109.10905]



- Charm production in the forward region are sensitive to both small and large x charm and gluon PDFs.
- Intrinsic charm can potentially show up in the large x region.
- Both the LHCb and the FASER measurement can provide probe to the gluon at small x and intrinsic charm at large x .

Final-state fragmentation



[2203.05090]

- The overall size gets reduced roughly by a factor of fragmentation fraction $\mathcal{B}(c \rightarrow D)$
- Fragmentation shift the p_T spectrum to a lower value.
- Forward D production probes both large and small x PDFs, involving no-pert. charm and gluon components.

Summary

- We develop S-ACOT-MPS scheme for the heavy flavor production at hadron colliders
 - Inclusive heavy quark production from both Flavor Creation and Flavor Excitation;
 - The double-counted term from gluon splitting is subtracted;
 - We introduce massive phase space to capture the threshold effect.
- We obtain good cancellations behaviors in both asymptotic limits:
 - $p_T \ll m_Q$, the SB cancels the FE term, FFN scheme,
 - $p_T \gg m_Q$, the SB cancels the FC term, ZM scheme.
- Our calculations agree well with the LHCb D measurements.
- Our theory directly applies the forward D production, which provide the neutrino source measured at FPFs.
- Implementation in MCFM can be easily extended to NNLO.
- We have obtained the subtraction $\tilde{f}_Q = P_{Qg} \otimes g$ and residual $\delta f_Q = f_Q - \tilde{f}_Q$ PDF, which can be easily applied to other heavy-flavor process, such as $H/V + Q$. Available on [HEPForge](#).
- Fast computation tables are generated.