

Light Scalars at FASER

FPF Theory Workshop

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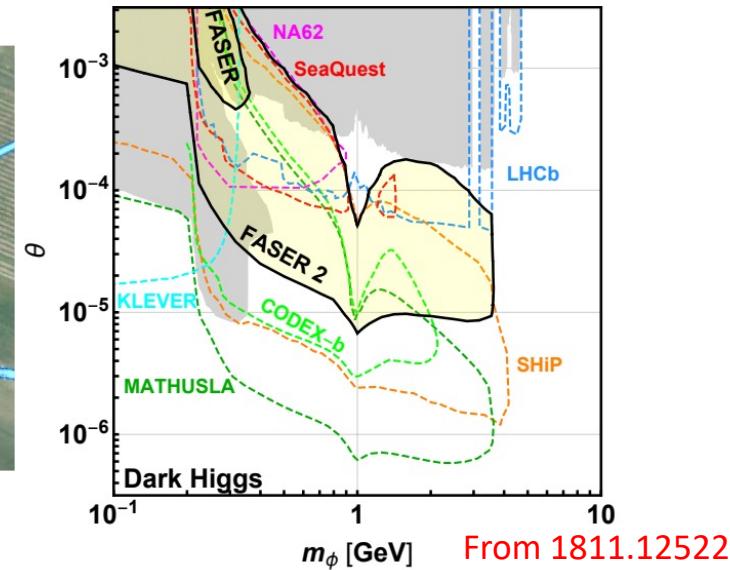
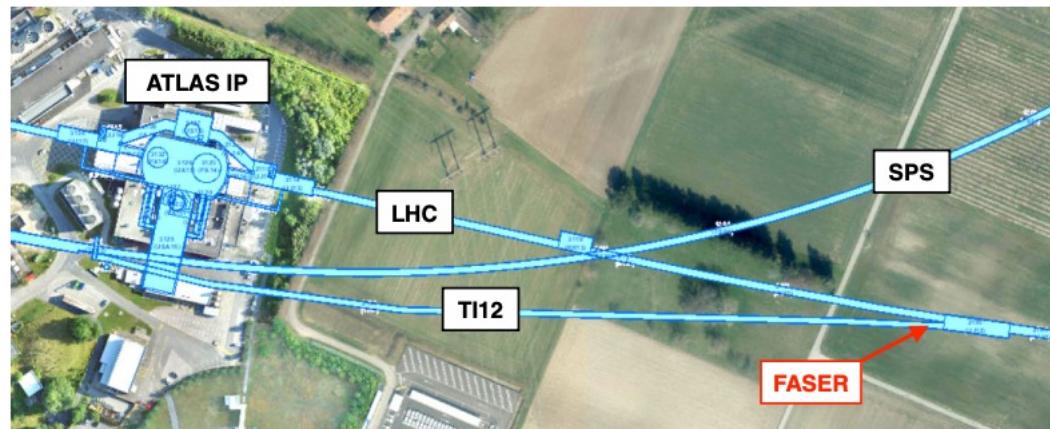


Light Scalars @ FASER

Many Beyond Standard Models including extended Higgs sector permit the light and weakly coupled scalars, such as Dark Higgs (SM+Singlet), 2HDM, 2HDM+(P)S, NMSSM,

Simplest prototype model:
Dark Higgs

$$\mathcal{L} = -m_\phi^2 \phi^2 - \sin \theta \frac{m_f}{v} \phi \bar{f} f - \lambda v h \phi \phi + \dots$$



Light Scalars @ FASER

**Model-independent framework with the most general interactions
for CP-even and CP-odd scalar under EFT/coupling modifier.**

- developed general formalism for scalar production and decay
- CP-odd A mix with light meson states
- developed a program to calculate scalar decay, can be used for other new physics models
- more complicated comparing to the simplest scenario
- case study of 2HDM.

Light cP-even Scalar

**Effective
Lagrangian**

$$\mathcal{L} = -\frac{1}{2}m_\phi^2\phi^2 - \sum_f \xi_\phi^f \frac{m_f}{v} \phi \bar{f} f + \xi_\phi^W \frac{2m_W^2}{v} \phi W^\mu{}^+ W_\mu^- + \xi_\phi^Z \frac{m_Z^2}{v} \phi Z^\mu Z_\mu \\ + \xi_{\phi\phi}^W \frac{g^2}{4} \phi \phi W^\mu{}^+ W_\mu^- + \xi_{\phi\phi}^Z \frac{g^2}{8 \cos^2 \theta_W} \phi \phi Z^\mu Z_\mu + \xi_\phi^g \frac{\alpha_s}{12\pi v} \phi G_{\mu\nu}^a G^{a\mu\nu} + \xi_\phi^\gamma \frac{\alpha_{ew}}{4\pi v} \phi F_{\mu\nu} F^{\mu\nu}$$

coupling modifiers

loop generated

Production

- decay of mesons, hadrons, radiative bottomonium
- Bremsstrahlung: small for high beam energies
- photon/gluon fusion: smaller, small in forward region
- h/Z/W decay: small in forward region

$$\mathcal{L}_{eff} = \frac{\phi}{v} \sum \xi_\phi^{ij} m_{f_j} \bar{f}_i P_R f_j + h.c. \quad \mathcal{L} \supset \xi_{\phi\phi}^{ij} \frac{\phi^2}{v^2} m_j \bar{f}_i P_R f_j + h.c.$$

effective coupling for flavor changing quark interactions

ϕ Production

- Heavy B meson decay

$$B \rightarrow X_s \phi \quad \xi_{\phi}^{bs}$$

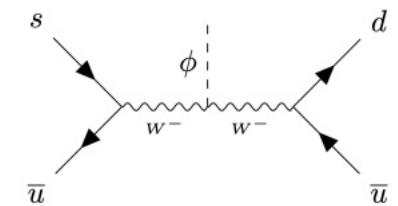
- Semileptonic decay of mesons

$$X \rightarrow \phi e \nu \quad \xi_{\phi}^W$$

- Kaon decay

$$K \rightarrow \pi \phi$$

$$\xi_{\phi}^{ds} \quad \xi_{\phi}^W$$



- $\eta^{(\prime)}$ decay

$$\eta^{(\prime)} \rightarrow \pi \phi$$

$$g_{\phi \eta^{(\prime)} \pi}$$

- Radiative bottomonium decay

$$\Upsilon \rightarrow \gamma \phi \quad \xi_{\phi}^b$$

- Double scalar production

$$B \rightarrow X_s \phi \phi \quad K \rightarrow \pi \phi \phi \quad \xi_{\phi \phi}^{ij}$$

ϕ Decay

Decay into a pair of photons, leptons, pair of quarks (gluons)/multiple hadrons

- Decay into diphoton

$$\Gamma_{\gamma\gamma} = \frac{G_F \alpha_{ew}^2 m_\phi^3}{32\sqrt{2}\pi^3} |\xi_\phi^\gamma|^2$$

- Decay into dilepton

$$\Gamma_{\ell^+\ell^-} = \frac{G_F m_\phi m_\ell^2 \beta_\ell^3}{4\sqrt{2}\pi} |\xi_\phi^\ell|^2$$

$m_\phi > 2$ GeV: perturbative spectator model

- Decay into diquark

$$\Gamma_{\ell^+\ell^-} : \Gamma_{s\bar{s}} : \Gamma_{c\bar{c}} : \Gamma_{b\bar{b}} = |\xi_\phi^\ell|^2 m_\ell^2 \beta_\ell^3 : 3|\xi_\phi^s|^2 m_s^2 \beta_K^3 : 3|\xi_\phi^c|^2 m_c^2 \beta_D^3 : 3|\xi_\phi^b|^2 m_b^2 \beta_B^3$$

- Decay into digluon

$$\Gamma_{gg} = \frac{G_F \alpha_s^2 m_\phi^3}{36\sqrt{2}\pi^3} |\xi_\phi^g|^2$$

$m_\phi < 2$ GeV: dispersive analyses

- Hadronic decay into pions and kaons

$$\Gamma_\pi \quad \Delta_\pi \quad \Theta_\pi$$

$$\xi_\phi^u \quad \xi_\phi^d \quad \xi_\phi^s \quad \xi_\phi^d$$

- Further hadronic decays

$$\phi \rightarrow 4\pi, \eta\eta, KK\pi\pi, \rho\rho \dots$$

$$\Gamma_{4\pi, \eta\eta, \rho\rho, \dots} = C |\xi_\phi^g|^2 m_\phi^3 \beta_{2\pi}$$

Light cP-odd Scalar

Effective Lagrangian

$$\mathcal{L}_A = -\frac{1}{2}m_A^2 A^2 + \sum_{f=u,d,e} \xi_A^f \frac{im_f}{v} \bar{f} \gamma_5 f A + \xi_{AA}^W \frac{g^2}{4} A A W^{\mu+} W_\mu^- + \xi_{AA}^Z \frac{g^2}{8 \cos^2 \theta_W} A A Z^\mu Z_\mu$$

$$+ \xi_A^g \frac{\alpha_s}{4\pi v} A G_{\mu\nu}^a \tilde{G}^{a\mu\nu} + \xi_A^\gamma \frac{\alpha_{ew}}{4\pi v} A F_{\mu\nu} \tilde{F}^{\mu\nu}$$

loop generated coupling modifiers

Mixing

$$A \approx O_{A\pi^0}\pi^0 + O_{A\eta}\eta + O_{A\eta'}\eta' + O_{AAA}A_{CP-odd}$$

typically small except in the resonant region $m_A \sim m_i$

Production

- Production via meson mixing $\sigma_A \approx |O_{A\pi^0}|^2 \sigma_{\pi^0} + |O_{A\eta}|^2 \sigma_\eta + |O_{A\eta'}|^2 \sigma_{\eta'}$
- B meson and Kaon decay $K \rightarrow \pi A \quad B \rightarrow X_s A \quad \xi_A^{ij}$
- Bottomonium decay $\Upsilon \rightarrow \gamma A \quad J/\psi \rightarrow \gamma A \quad \xi_A^f$
- Double pseudoscalar production $B \rightarrow X_s A A \quad K \rightarrow \pi A A \quad \xi_{AA}^{ij}$

A Decay

Decay into a pair of photons, leptons, pair of quarks (gluons)/multiple hadrons

- **Decay into diphoton**

$$\Gamma(A \rightarrow \gamma\gamma) = \frac{\alpha_{\text{ew}}^2 m_A^3}{64\pi^3} \left| O_{AA} C_A^\gamma + O_{A\pi^0} C_{\pi^0}^\gamma + O_{A\eta} C_\eta^\gamma + O_{A\eta'} C_{\eta'}^\gamma \right|^2$$

- **Decay into dilepton**

$$\Gamma(A \rightarrow \ell^+ \ell^-) = \frac{G_F m_A m_\ell^2 \beta_\ell}{4\sqrt{2}\pi} |\xi_A^\ell|^2$$

$m_A > 3 \text{ GeV}$: perturbative spectator model

- **Decay into diquark**

$$\Gamma_{\bar{\ell}\ell} : \Gamma_{\bar{s}s} : \Gamma_{\bar{c}c} : \Gamma_{\bar{b}b} = (\xi_A^\ell)^2 m_\ell^2 \beta_\ell : 3(\xi_A^s)^2 m_s^2 \beta_s : 3(\xi_A^c)^2 m_c^2 \beta_c : 3(\xi_A^b)^2 m_b^2 \beta_b$$

- **Decay into digluon**

$$\Gamma(A \rightarrow gg) = \frac{G_F \alpha_s^2 m_A^3}{4\sqrt{2}\pi^3} |\xi_A^g|^2$$

$1.3 \text{ GeV} < m_A < 3 \text{ GeV}$: spectator model with partonic dynamic and hadronic kinematics

- **Hadronic decay**

$$\mathcal{L}_{\text{spect.}} = \frac{i}{\sqrt{2}} A_1 (\mathcal{Y}_u^A \bar{u} \gamma_5 u + \mathcal{Y}_d^A \bar{d} \gamma_5 d + \mathcal{Y}_s^A \bar{s} \gamma_5 s)$$

$$\mathcal{Y}_u^A \approx \frac{\sqrt{2}B}{\sqrt{3}v f_\pi^2} m_u \xi_A^u$$

A Decay continued

$m_A < 1.3 \text{ GeV}$: chiral perturbation theory

- Hadronic decay into tri-meson

$$\Gamma(A \rightarrow \Pi_i \Pi_j \Pi_k) = \frac{1}{256 S_{ijk} \pi^3 m_A} \int_{(m_j+m_k)^2}^{(m_A-m_i)^2} ds |\mathcal{M}_A^{ijk}|^2$$

$$\sqrt{1 - \frac{2(m_j^2 + m_k^2)}{s} + \frac{(m_j^2 - m_k^2)^2}{s^2}} \times \sqrt{\left(1 + \frac{s - m_i^2}{m_A^2}\right)^2 - \frac{4s}{m_A^2}}$$

$$\mathcal{M}_A^{ijk} \propto O_{AA} \mathcal{A}_A^{ijk} + \sum_l O_{Al} \mathcal{A}^{ijkl}$$

- Radiative hadronic decay

$$A \rightarrow \pi^+ \pi^- \gamma$$

$$\Gamma(A \rightarrow \pi^+ \pi^- \gamma) = \int_{4m_\pi^2}^{m_A^2} ds \Gamma_0(s) |O_{A\eta} B_\eta(s) + O_{A\eta'} B_{\eta'}(s)|^2$$

2HDM

Two Higgs Doublet Model (CP-conserving): $\phi_{1,2}$

After EWSB, 5 physical Higgses:

CP-even Higgses: h , H , CP-odd Higgs: A , charged Higgses: H^\pm

parameters (CP-conserving, flavor limit, Z_2 symmetry)

$v, \tan \beta, \alpha, m_h, m_H, m_A, m_{H^\pm}$ soft Z_2 breaking: m_{12}^2

Alignment limit: h is 125 GeV Higgs, $\cos(\beta - \alpha) \sim 0$

- Type I: ϕ_1 couples quarks and leptons
all fermion couplings suppressed at large $\tan \beta \Rightarrow$ LLP
- Type II, L, F: $\phi_{1,2}$ couples to at least one type of quarks or leptons
unsuppressed couplings of scalars to at least one type of fermions for the entire region of $\tan \beta \Rightarrow$ difficult to realize very weakly coupled long-lived scalars

constraints

- Theoretical constraints: unitarity, perturbativity, vacuum stability
- EW precision constraints
- Flavor constraints
- Invisible Higgs decay
- LEP & LHC H^\pm search

Two benchmark scenarios

Light H : $\cos(\beta - \alpha) = \frac{1}{\tan \beta}$, $m_A = m_{H^\pm} = 600$ GeV, $\lambda v^2 = 0$

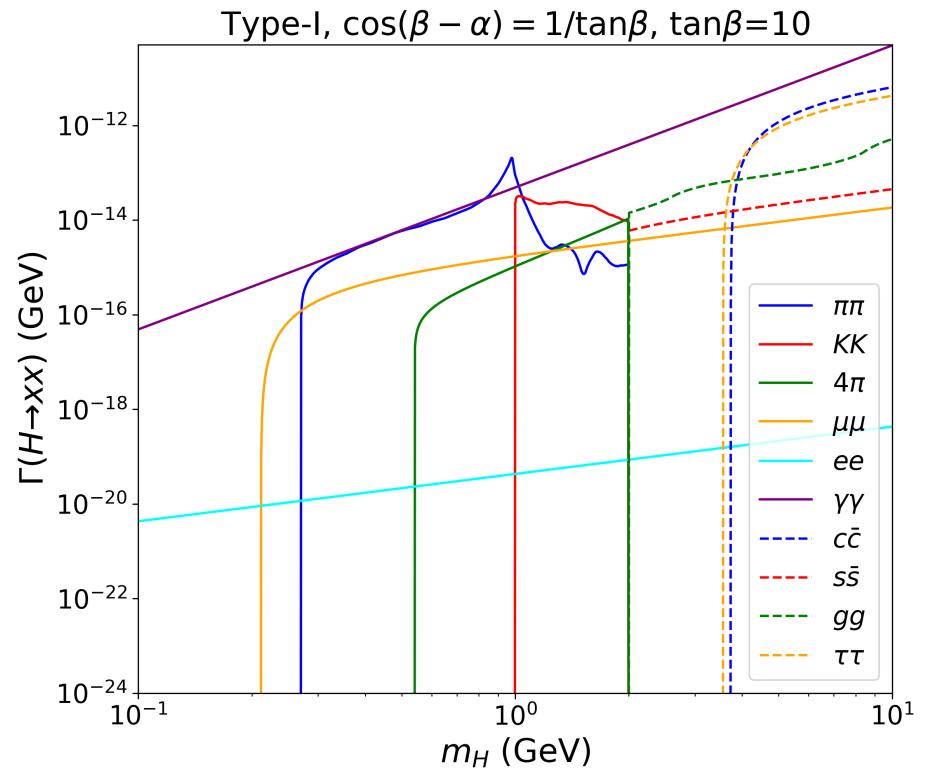
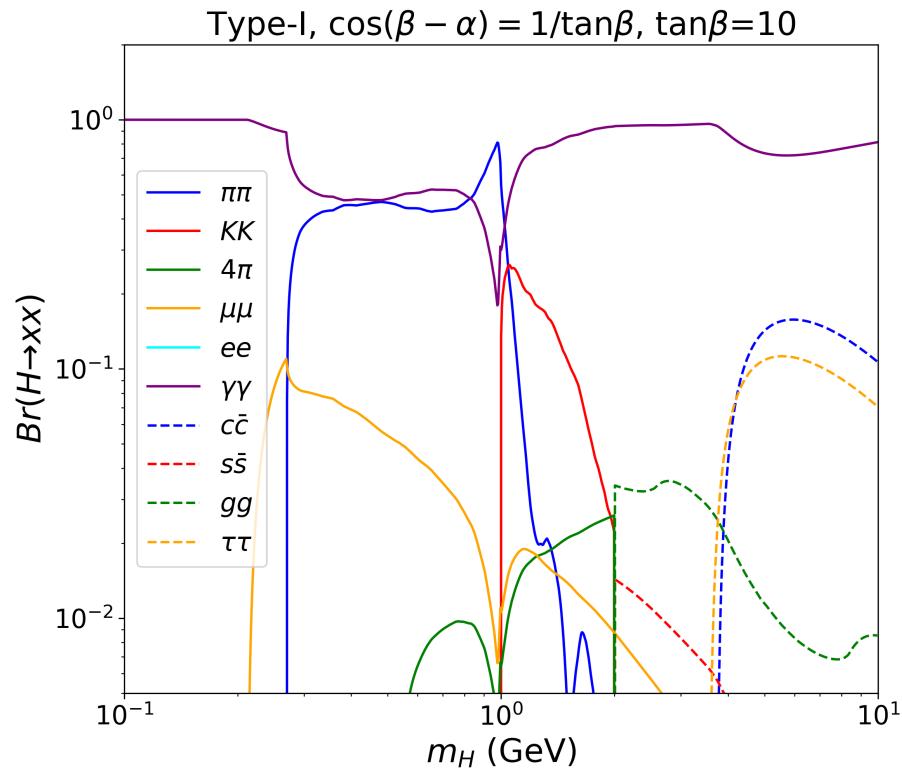
Light A : $\cos(\beta - \alpha) = 0$, $m_H = m_{H^\pm} = 90$ GeV, $\lambda v^2 = 0$,

$$\xi_A^f|_{\cos(\beta-\alpha)=0} = 1/\tan \beta,$$

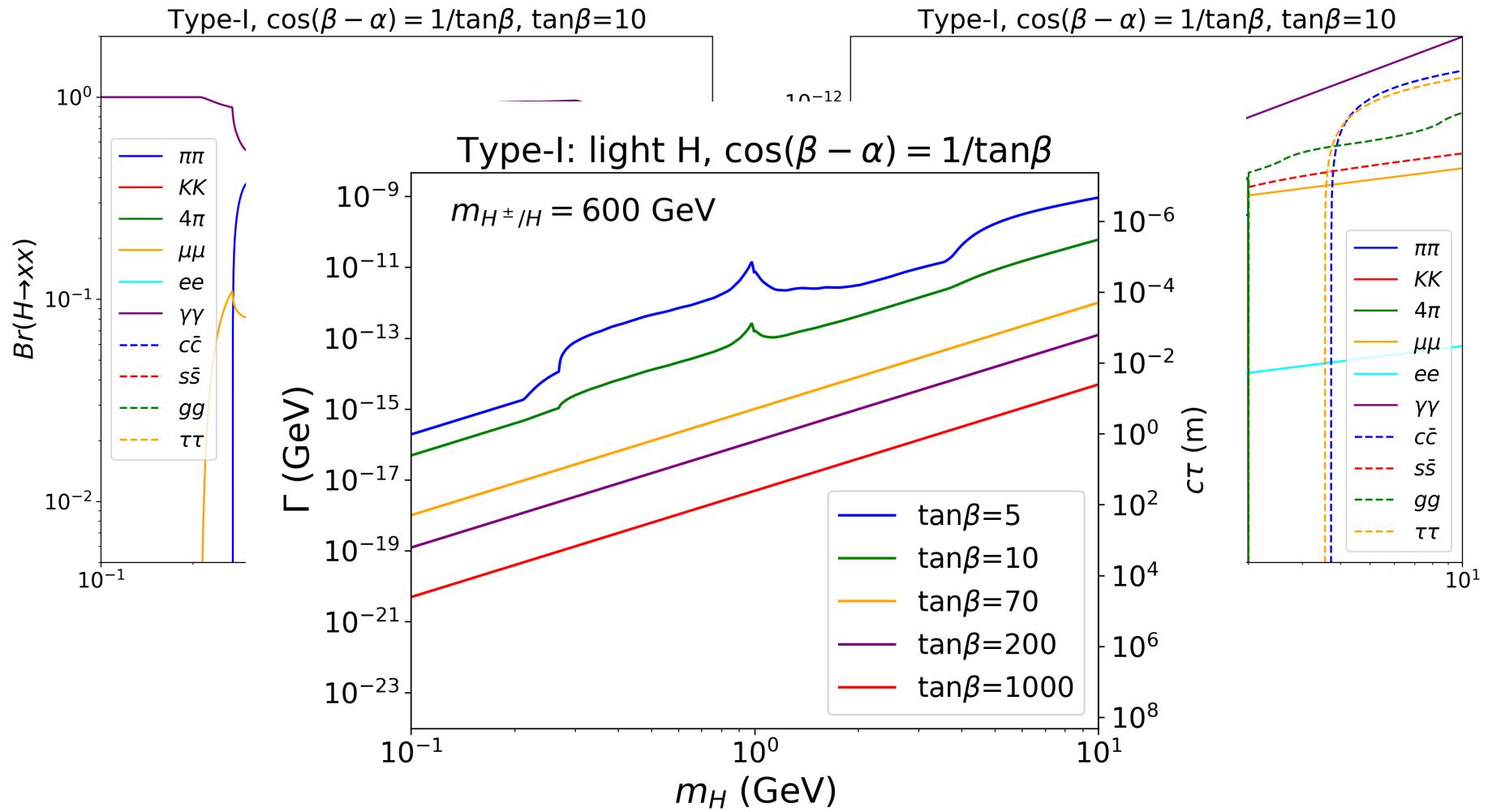
$$\xi_H^V = c_{\beta-\alpha} = 1/\tan \beta,$$

$$\xi_H^f = c_{\beta-\alpha}(1 - s_{\beta-\alpha}) \approx 1/(2 \tan^3 \beta) + \mathcal{O}(c_{\beta-\alpha}^5)$$

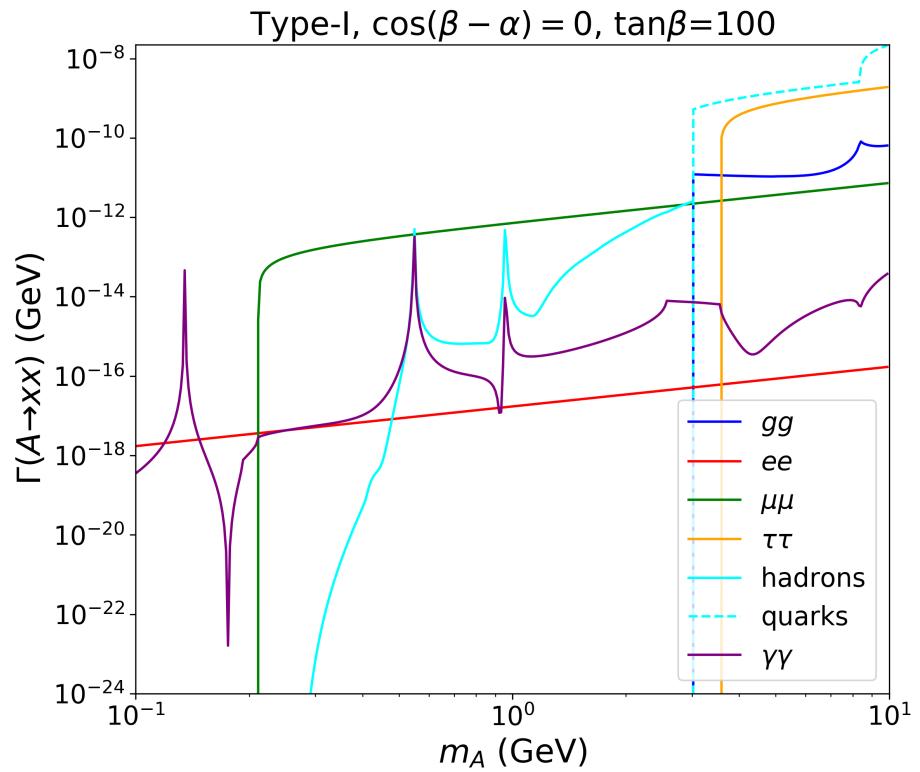
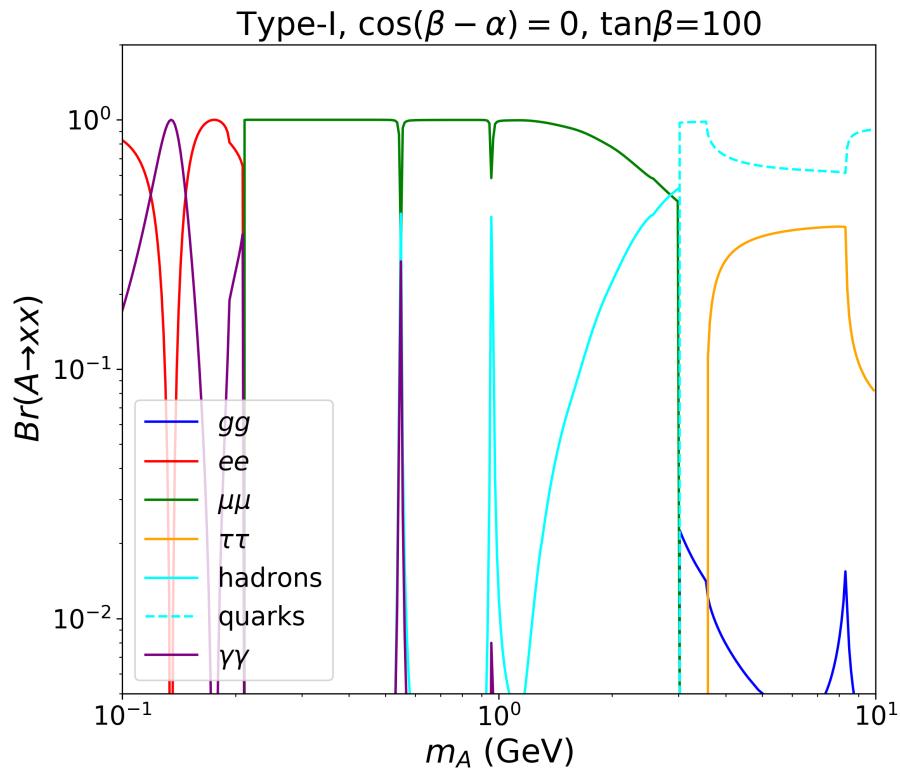
Light cP-even Scalar



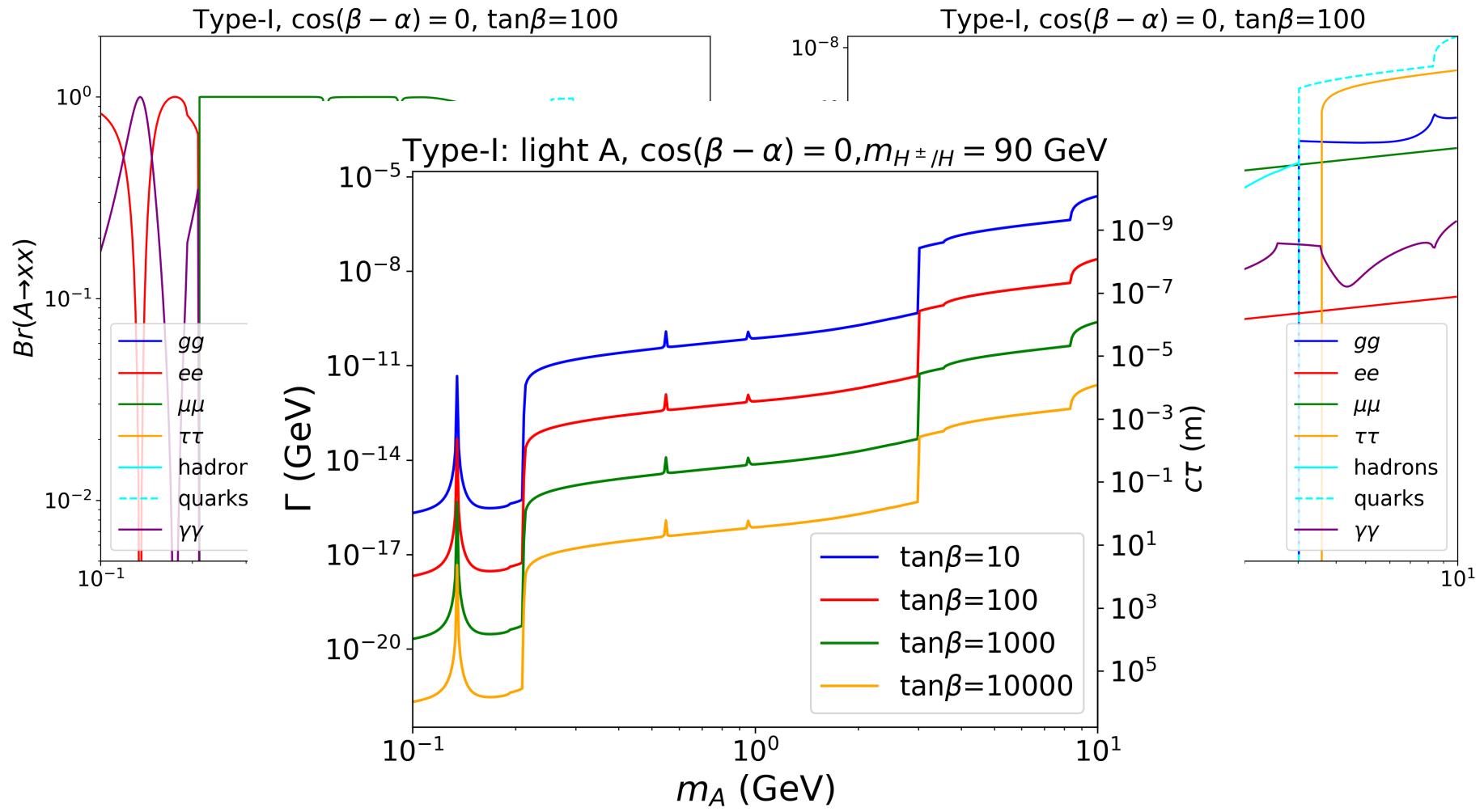
Light cP-even Scalar



Light cP-odd Scalar



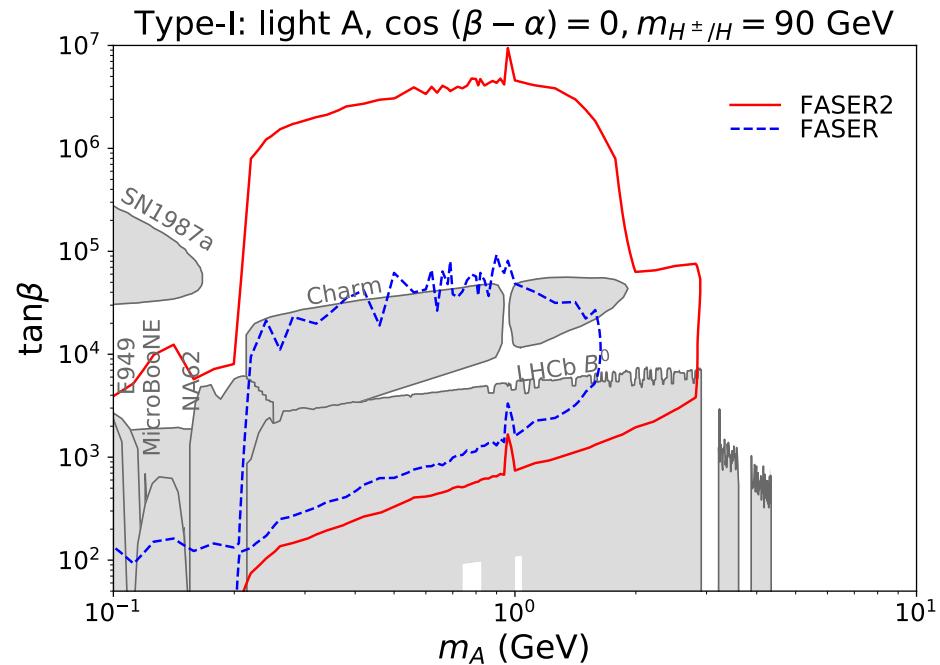
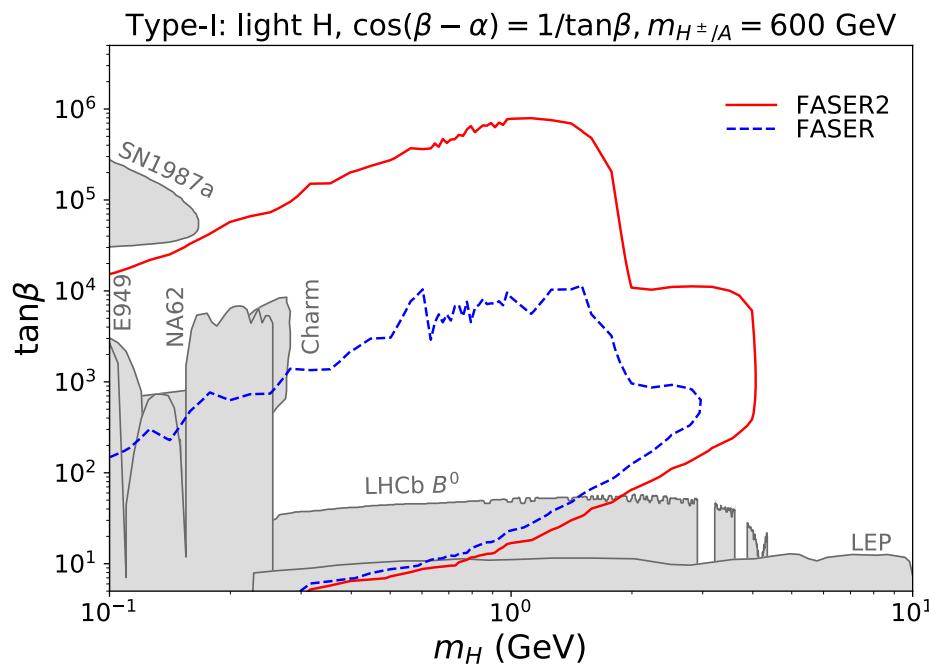
Light cP-odd Scalar



Other constraints on Light Scalar Searches

- **CHARM bounds: light ALP** CHARM, PLB 157 (1985) 458
- **Supernova: $NN \rightarrow NNS(A)$** Turner, PRL 60 (1988) 1797
- **B meson decays: $B \rightarrow K^* \phi$ (LHCb)** LHCb, 1508.04094, 1612.07818
- **D meson decays: $D^+ \rightarrow \pi^+ \phi$ (LHCb)** PDG, LHCb, 2011.00217
NA62, 2103.15389
- **Kaon decays: $K^+ \rightarrow \pi^+ \phi$ (NA62, MicroBooNE, E949)** MicroBooNE, 2106.00568
BNL-E949, 0903.0030
- **LEP: $e^- e^+ \rightarrow Z^* \phi$** Winkler, 1809.01876
Clarke, Foot and Volkas, 1310.8042

Light Scalars Reaches at FASER

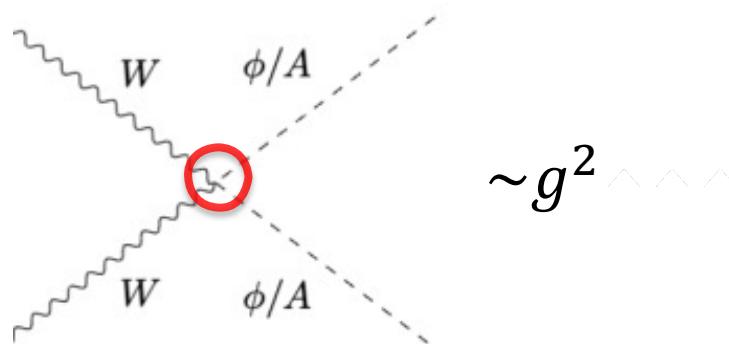


$$\xi_A^f|_{\cos(\beta-\alpha)=0} = 1/\tan\beta,$$

$$\xi_H^V = c_{\beta-\alpha} = 1/\tan\beta,$$

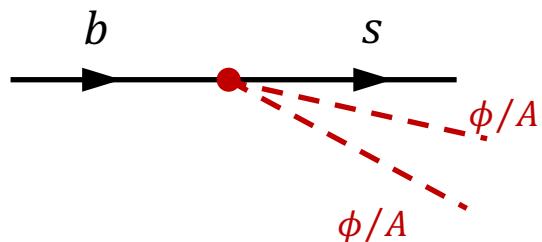
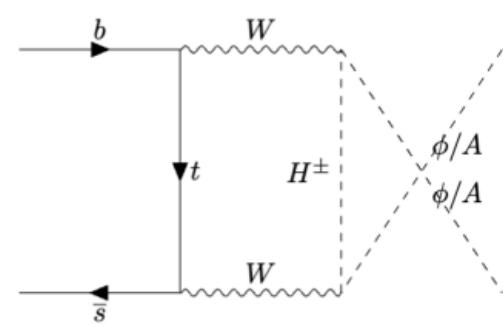
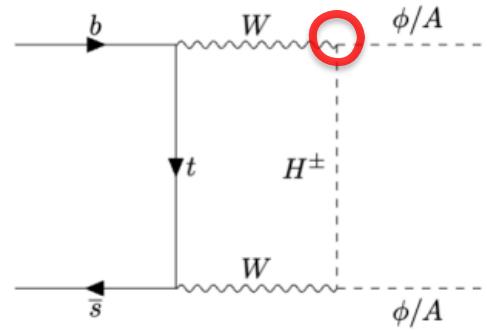
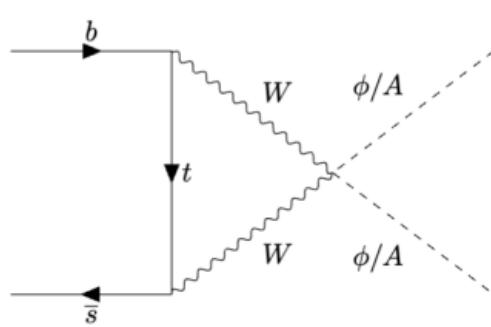
$$\xi_H^f = c_{\beta-\alpha}(1 - s_{\beta-\alpha}) \approx 1/(2\tan^3\beta) + \mathcal{O}(c_{\beta-\alpha}^5)$$

Double Scalar Production



$\sim g^2$

Governed by gauge symmetry
and not suppressed



$$\mathcal{L} \supset \xi_{\phi\phi}^{ij} \frac{\phi^2}{v^2} m_j \bar{f}_i P_R f_j + \xi_{AA}^{ij} \frac{A^2}{v^2} m_j \bar{f}_i P_R f_j + h.c.$$

conclusion

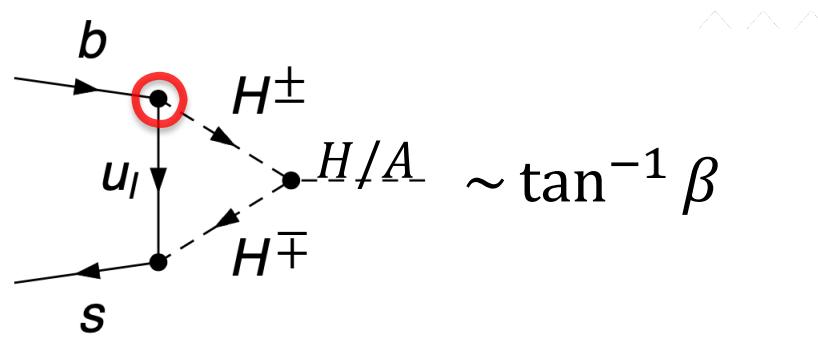
- ❖ Light LLP appear in many new physics scenarios
- ❖ Light particle copiously produced in the forward region of LHC, and FASER/FASER2 (FPF): new experiments to detect light LLP
- ❖ Light (pseudo)scalar
 - Model-independent framework, coupling modified in EFT
 - Scalar production and decay (hadronic)
 - Public code to calculate decay
(https://github.com/shiggs90/Light_scalar_decay.git)
- ❖ 2HDM case study: large $\tan\beta$ region of Type-I 2HDM
 - decay length: 10^{-8} to 10^5 m, probe very large $\tan\beta$
 - FASER2 vs. FASER: higher Lum, larger detector
- ❖ Complementary to LHC prompt search, LLP search in transverse region, and fixed target exp at low energies, or other astrophysical processes (e.g. supernova)



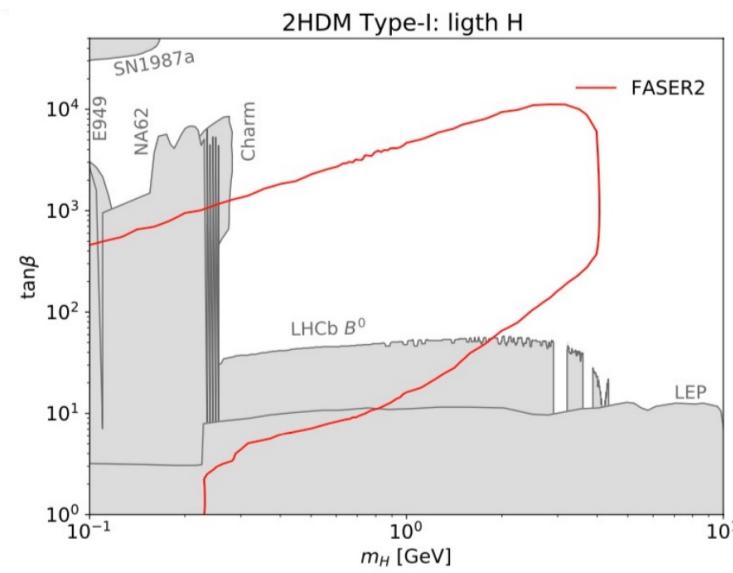
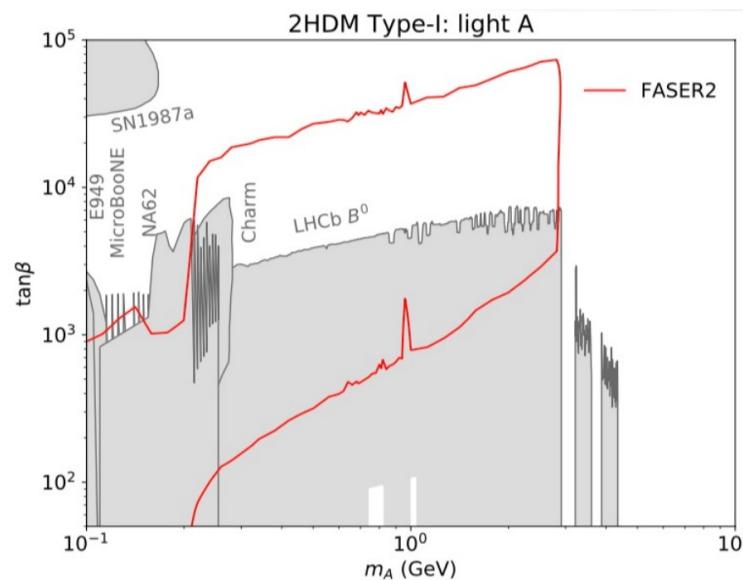
Backup Slides

without Double Scalar Production

In Type-I 2HDM

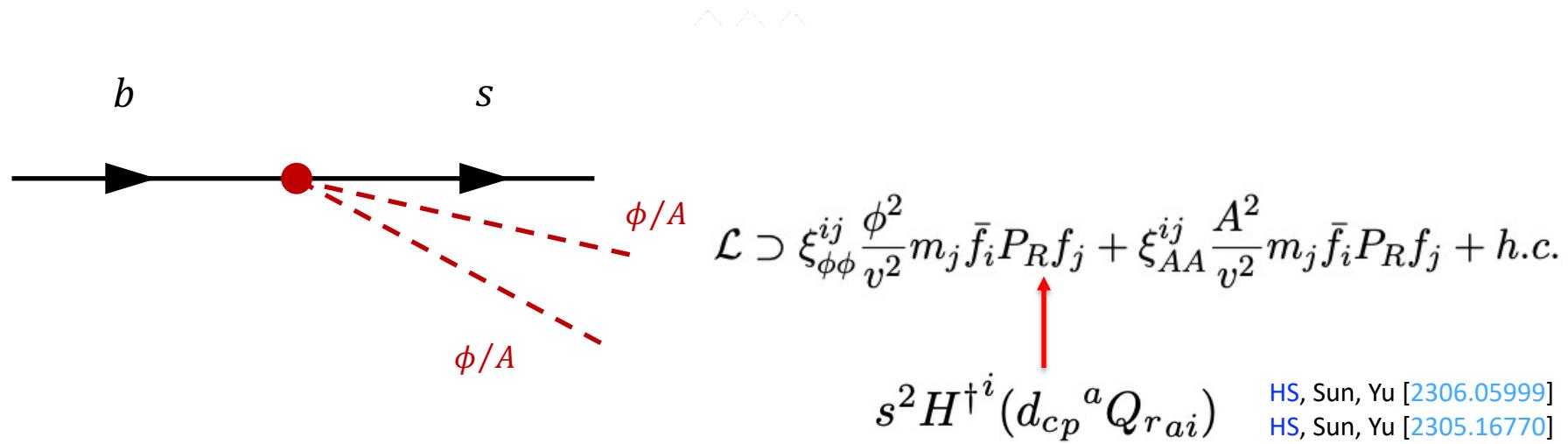


Similar to minimal model
(Dark Higgs)



Effective couplings

In Type-I 2HDM

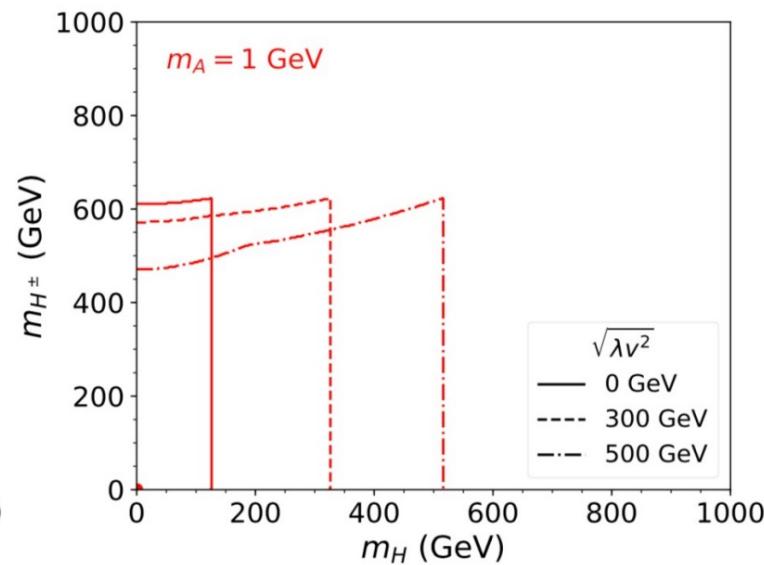
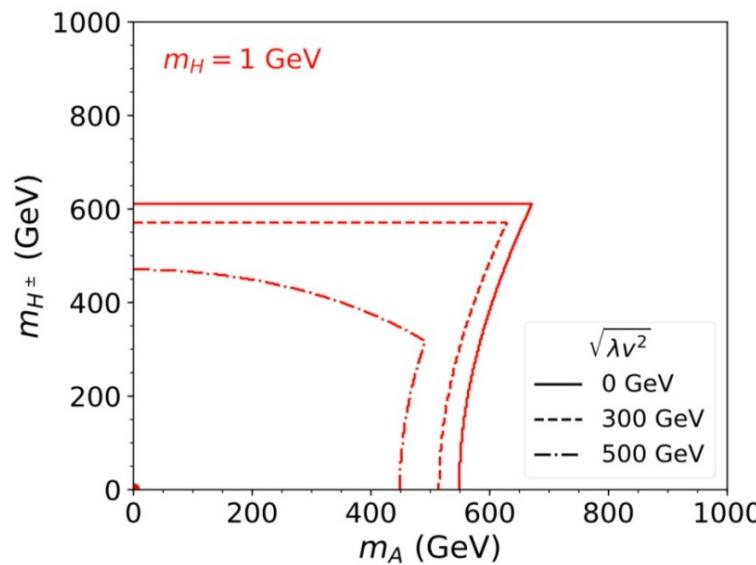


$$\begin{aligned} \xi_{\phi\phi}^{ij} \simeq \xi_{AA}^{ij} \simeq \frac{g^2}{64\pi^2} \sum_k V_{ki}^* & [f_0(x_k, x_{H^\pm}) + f_1(x_k, x_{H^\pm}) \log x_k \\ & + f_2(x_k, x_{H^\pm}) \log x_{H^\pm}] V_{kj} + \mathcal{O}(\cos(\beta - \alpha), 1/\tan\beta) \end{aligned}$$

constraints

Theoretical constraints

$$\lambda v^2 \equiv m_H^2 - m_{12}^2 / s_\beta c_\beta = 0$$



$$m_H \sim 0 : \quad m_{A/H^\pm} \lesssim 600 \text{ GeV}$$

$$m_A \sim 0 : \quad m_{H^\pm} \lesssim 600 \text{ GeV}, \quad m_H \lesssim m_h$$

constraints

Invisible Higgs decays

$$\text{Br}(h \rightarrow \phi\phi) = \frac{\Gamma(h \rightarrow \phi\phi)}{\Gamma_h} \approx \frac{1}{\Gamma_h^{\text{SM}}} \frac{g_{h\phi\phi}^2}{8\pi m_h^2} \left(1 - \frac{4m_H^2}{m_h^2}\right)^{1/2} \simeq 4700 \cdot \left(\frac{g_{h\phi\phi}}{v}\right)^2 < 0.24$$

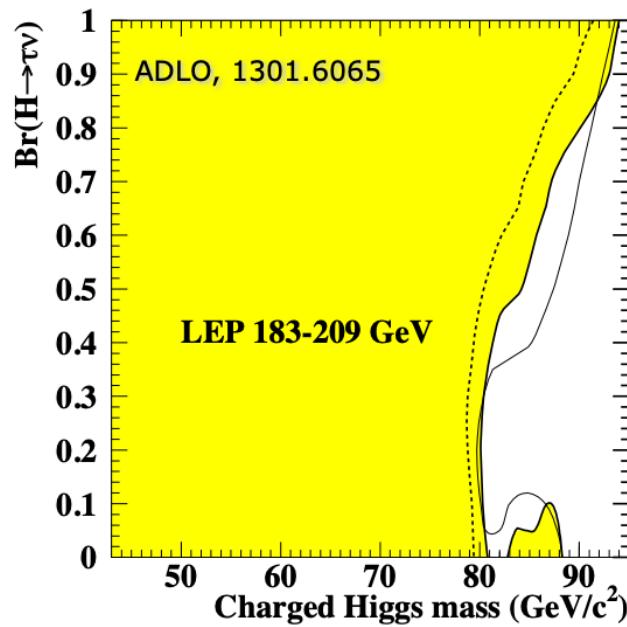
$$\text{Br}(h \rightarrow \phi\phi) = 0$$

$$\text{Light } H : \cos(\beta - \alpha) = \tan 2\beta \frac{2\lambda v^2 + m_h^2}{2(m_H^2 - 3\lambda v^2 - m_h^2)} \approx \frac{1}{\tan \beta},$$

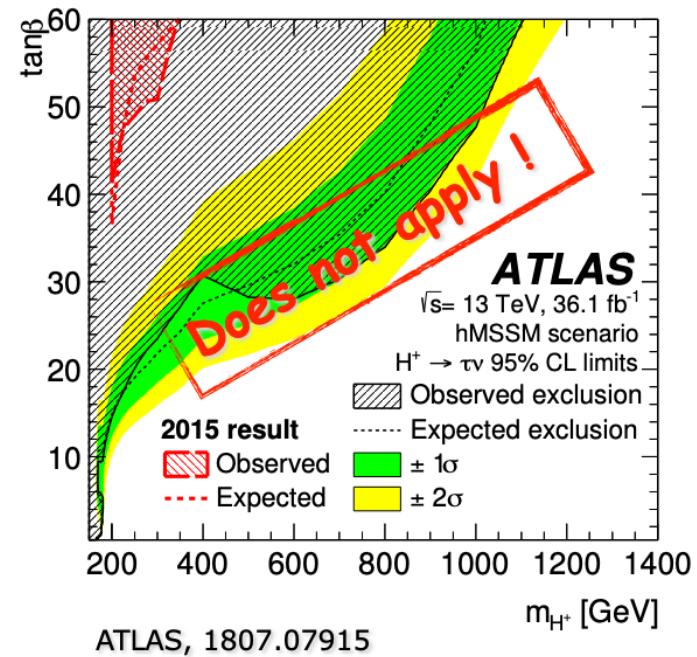
$$\text{Light } A : \cos(\beta - \alpha) = \tan 2\beta \frac{2\lambda v^2 + m_h^2 + 2m_A^2 - 2m_H^2}{2(m_H^2 - \lambda v^2 - m_h^2)} \approx \frac{1}{\tan \beta} \frac{2m_H^2 - m_h^2}{m_H^2 - m_h^2},$$

constraints

LEP H \pm search:
 $m_{H^\pm} > 85 \text{ GeV}$ viable

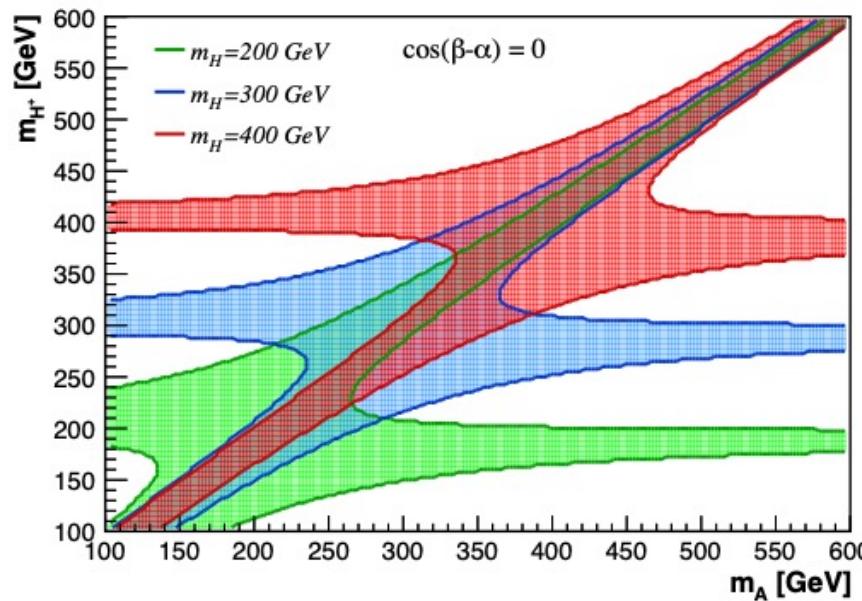


LHC H \pm search



constraints

- EW precision constraints:
 $m_{H^\pm} \sim m_H$ or m_A



$m_H \sim 0 : m_A \sim m_{H^\pm} \lesssim 600$ GeV
 $m_A \sim 0 : m_{H^\pm} \sim m_H \lesssim m_h,$
 $\lambda v^2 \approx 0 \quad |\cos(\beta - \alpha)| \sim 0.$