

# Machine Learning Opportunities for EFT Analyses

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**EFT WG 16.11.23**

# How to perform a global effective field theory analysis of LHC data?

 Performing a global Effective Field Theory (EFT) analysis of LHC data involves:

- 1. Define EFT Framework:**
  - Choose an EFT framework relevant to the physics scenario of interest.
- 2. Generate Simulated Data:**
  - Simulate expected LHC data using Monte Carlo methods based on the chosen EFT framework.
- 3. Develop Analysis Pipeline:**
  - Create an analysis pipeline that includes event selection, background subtraction, and parameter estimation.
- 4. Implement Machine Learning:**
  - Integrate machine learning for event classification, anomaly detection, and optimization of experimental design.
- 5. Compare with Experimental Data:**
  - Compare the simulated data with actual LHC data, adjusting EFT parameters to match observed results.
- 6. Statistical Analysis:**
  - Perform statistical analyses to quantify the agreement between the EFT predictions and experimental data.
- 7. Iterative Refinement:**
  - Iterate through steps 2-6, refining the analysis based on feedback and incorporating additional data.
- 8. Collaboration and Peer Review:**
  - Engage in collaboration with experts, undergo peer review, and ensure the transparency and reproducibility of the analysis.
- 9. Publication and Interpretation:**
  - Publish the results, interpret the findings in the context of the chosen EFT, and contribute to the global understanding of particle physics.

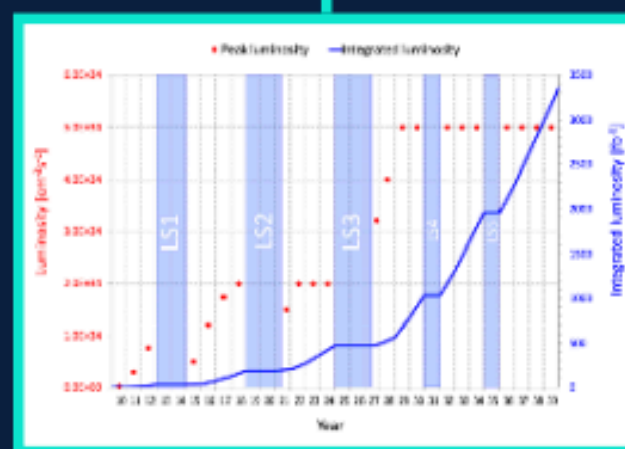
Global SMEFT analyses of LHC data can benefit from machine learning (according to machine learning)

# Why ML in HEP?

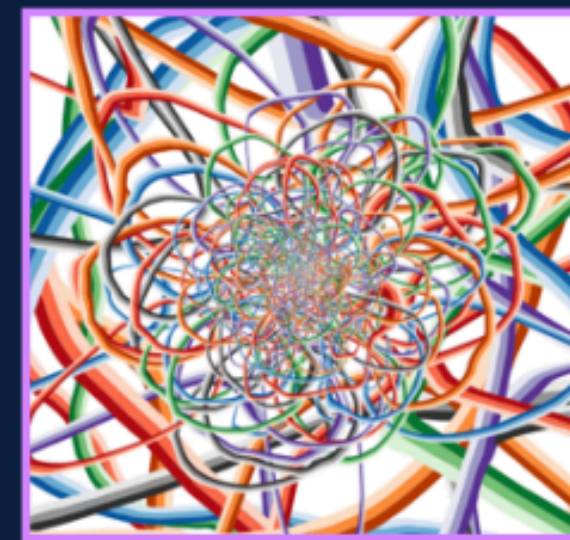
## Data volume

- Large amounts of data
- 1. labeled (Simulation)
- 2. unlabeled (Detector)

ML wants lots of data



1



2

## Complexity

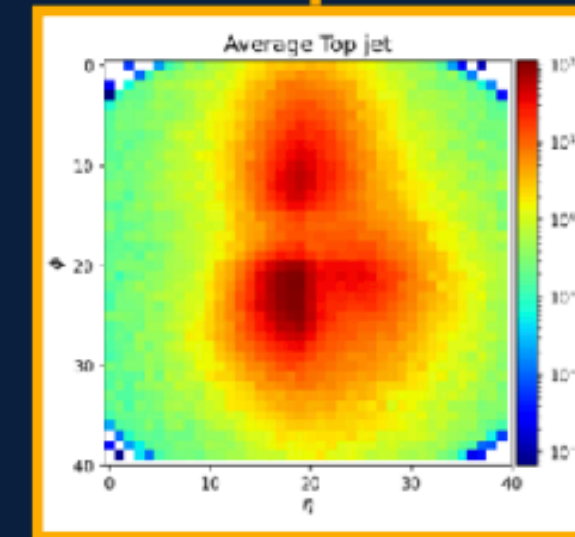
High-dimensional & highly correlated data structure

ML is expressive and interpretable

## Signal detection

Rare and elusive signals among large backgrounds

ML has high accuracy and sensitivity

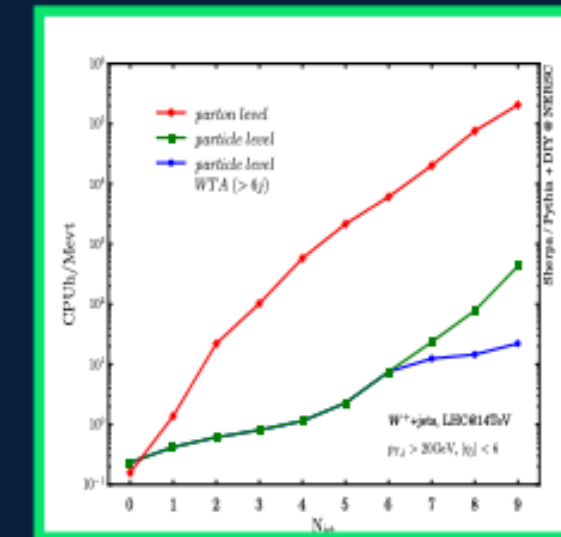


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## Computing Budget

Simulation & analysis is computationally expensive

ML is fast

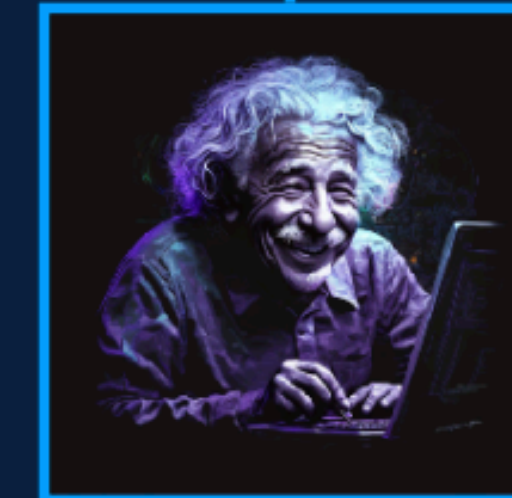


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## Increasing interest

> 150 paper/year  
Future of HEP?

ML is fun!



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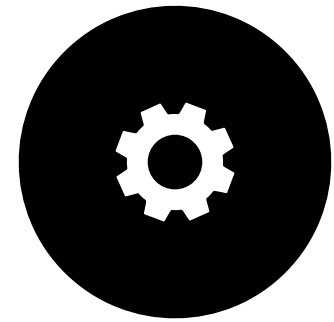
1987

2014

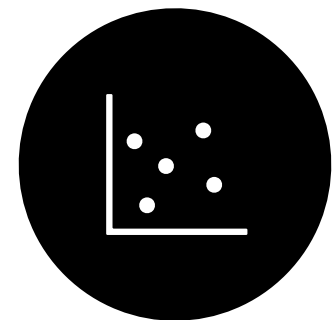
2022



# Machine Learning Opportunities for EFT Analyses : this talk



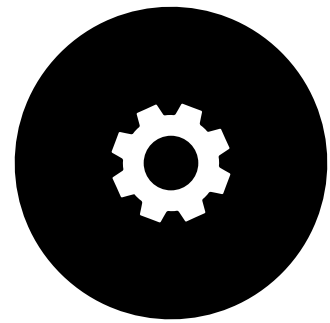
**Overview** of recent ML/EFT studies



**ML4EFT:** unbinned multivariate observables for global SMEFT fits



**SIMUnet:** global SMEFT and PDF determinations



## Overview of recent ML/EFT studies

LHC EFT WG - Area 3: Observables - Opportunities and Challenges of Machine Learning for EFT analyses

 Tuesday 24 Oct 2023, 15:00 → 16:40 Europe/Zurich

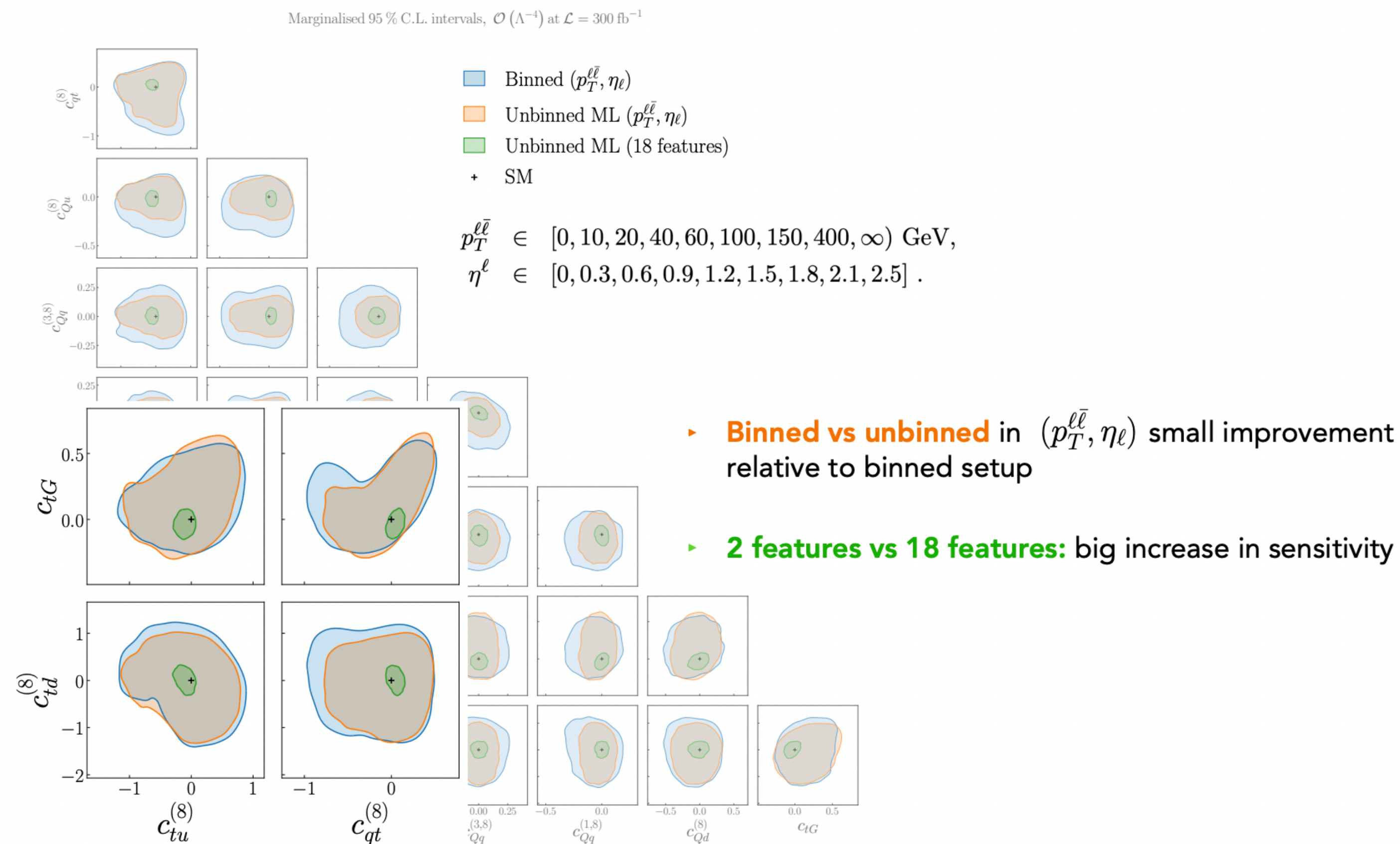
# Statistically optimal observables for global SMEFT fits

Jaco ter Hoeve

2211.02058

[Link to slides](#)

## Unbinned observables in the top sector



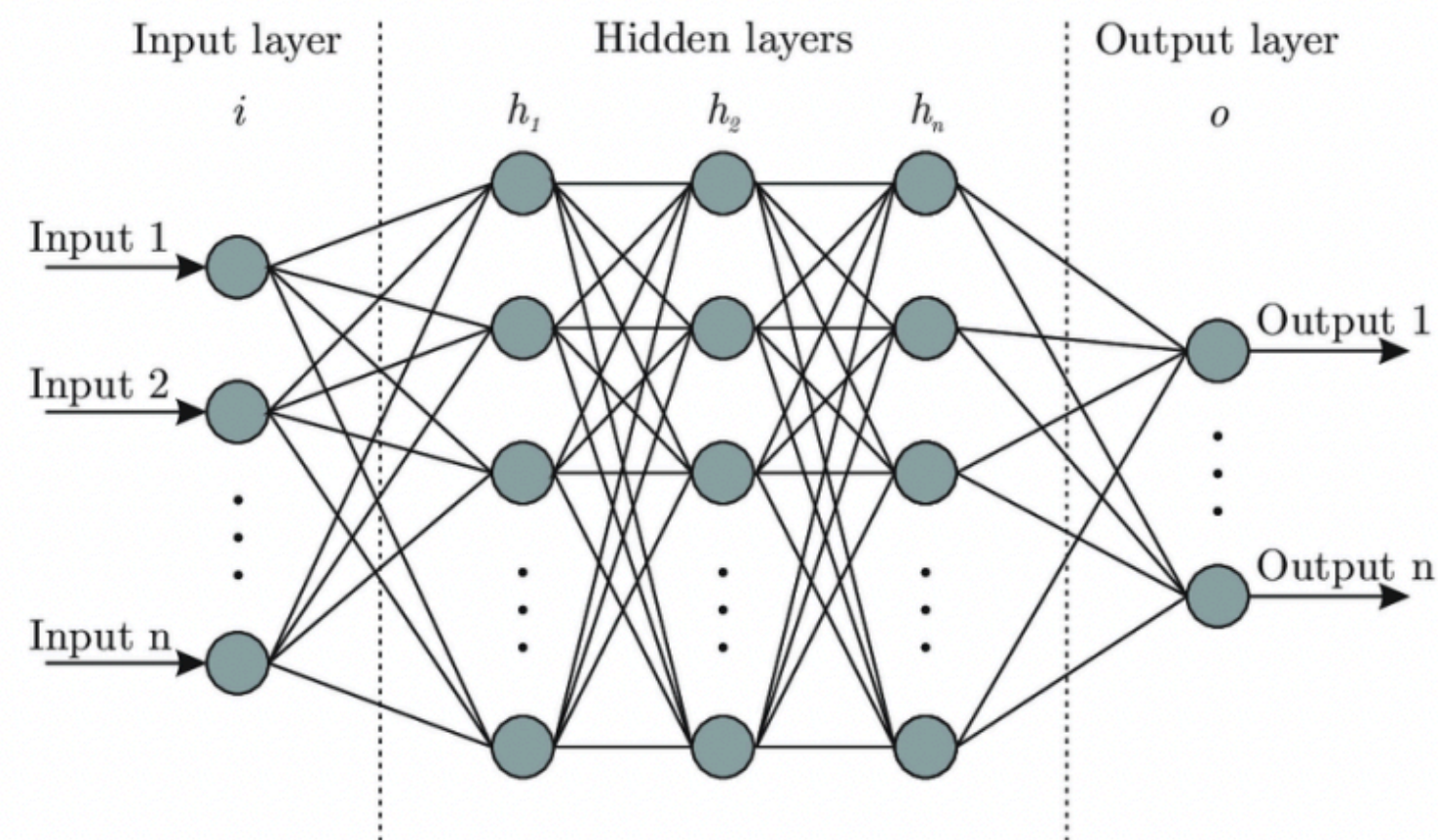
# Unbinned MVA techniques for EFT analyses

Alfredo Glioti

2007.10356, 2308.05704

[Link to slides](#)

**Basic idea:** approximate  $p(x|\theta)$  with **Neural Networks:**  $p(x|\theta) \leftrightarrow nn(x; w)$



The result will be **fully differential on all observables**, quick to evaluate and it can be obtained with a relatively small amount of Monte Carlo points.

**No transfer functions** modeling required.

**Universal and systematically improvable**

$$p(x|c) = \frac{1}{\sigma} \frac{d\sigma}{dx}(c)$$

- Multivariate in all features  $x$
- Extract full information on relation between data  $x$  and Wilson coefficients  $C$
- Optimal constraints on the EFT

# Unbinned MVA techniques for EFT analyses

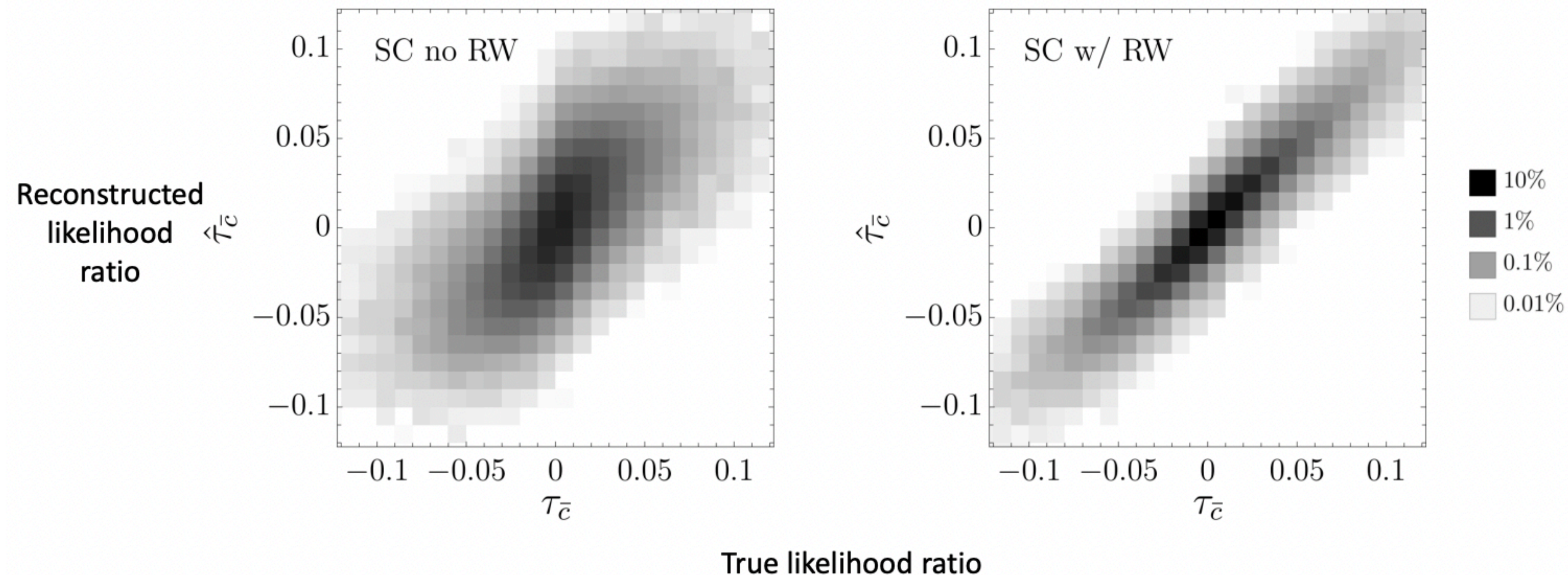
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Training on reweighted samples reduces number of training points needed and leads to a higher accuracy

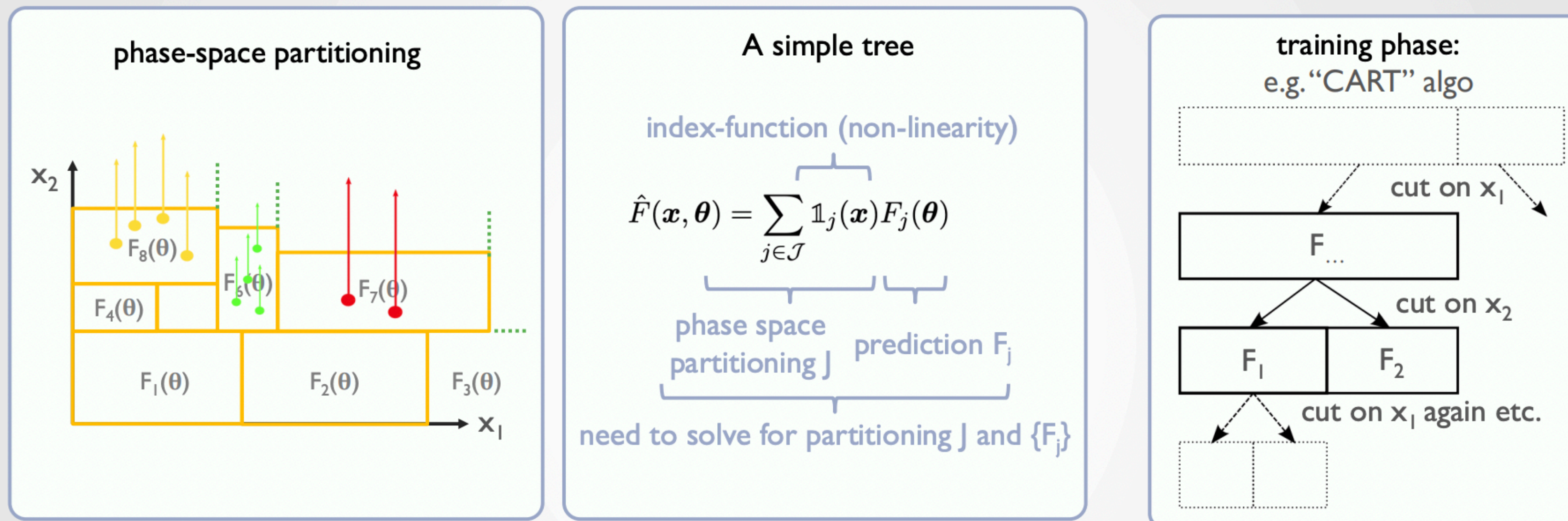
Without reweighting (left) vs with reweighting (right):





## A SIMPLE TREE ALGORITHM

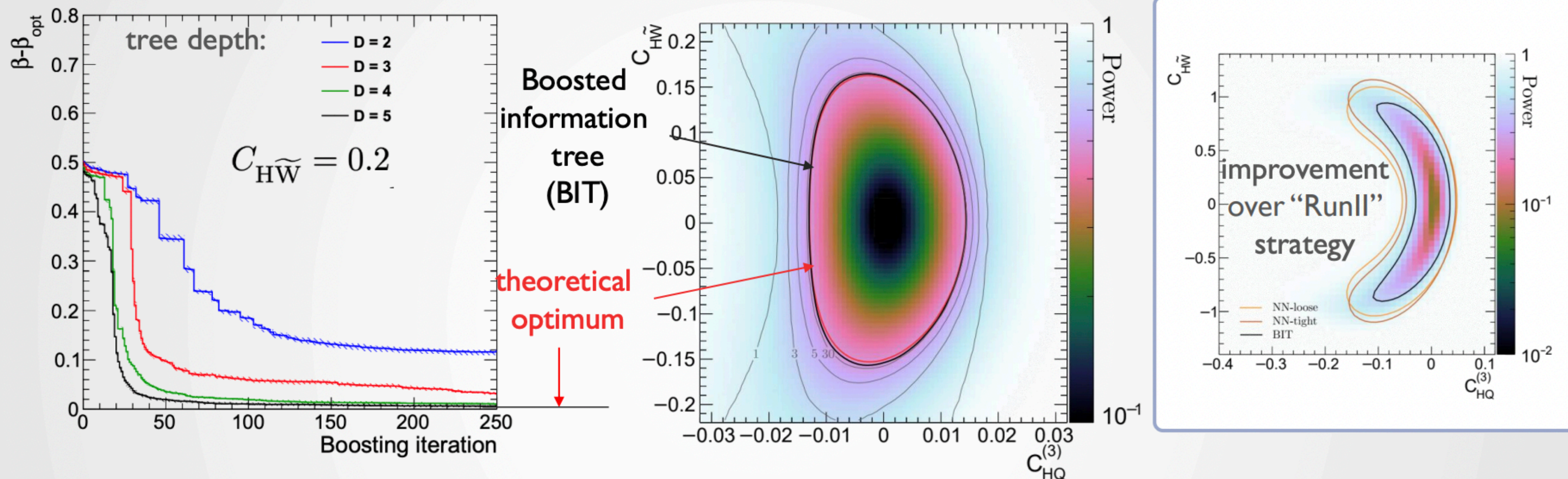
[[arXiv:2107.10859](#), [arXiv:2205.12976](#)]



- A tree is a hierarchical phase-space partitioning ( $\mathcal{J}$ )
  - the novelty in the Boosted Information Tree is that we associate each region  $j$  with a polynomial  $F_j(\boldsymbol{\theta})$
  - Note: A tree algorithm can have an arbitrarily complicated predictive function; here it is a SMEFT polynomial
- Fitting tree: Optimize "node split positions" on some loss. Trained (e.g. greedily) on the *ensemble*.

## OPTIMALITY IN TEST CASES

[[arXiv:2107.10859](#), [arXiv:2205.12976](#)]



- Obtain parametrized classifiers with 20-40% improvements in 2D toy cases (NOT marginalized!)
- No free lunch – Analysis dependent choices are needed
  - Binned analysis: variable binning  $\rightarrow$  background estimation is CPU intensive
  - Systematics treatment for unbinned analyses (beyond Higgs  $M_{4\ell}$ ) less far developed
- Is it all worth it in higher dimensions? Yes! More examples: [[ML<sub>4</sub>EFT](#)]; full list of references in backup

# Machine Learning for Higgs CP properties

António Jacques Costa

2112.05052

[Link to slides](#)

- CP violating operators affecting the Higgs/multiboson interactions can be probed via Effective Field Theory:

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)}$$

$c_i/\Lambda^2$  the Wilson coefficients,  
 $\Lambda$  the scale of new physics

$$\begin{aligned}\tilde{\mathcal{O}}_{\Phi\tilde{B}} &= \Phi^\dagger \Phi B^{\mu\nu} \tilde{B}_{\mu\nu}, \\ \tilde{\mathcal{O}}_{\Phi\tilde{W}} &= \Phi^\dagger \Phi W^{i\mu\nu} \tilde{W}_{\mu\nu}^i, \\ \tilde{\mathcal{O}}_{\Phi\tilde{W}B} &= \Phi^\dagger \sigma^i \tilde{W}^{i\mu\nu} B_{\mu\nu}.\end{aligned}$$

- Beyond-the-SM amplitude is then given by:  $|\mathcal{M}_{BSM}|^2 = |\mathcal{M}_{SM}|^2 + 2\text{Re}\{\mathcal{M}_{SM}\mathcal{M}_{d6}^*\} + |\mathcal{M}_{d6}|^2$
- Interference term leads to asymmetries in CP-odd observables

- Possible CP-odd observables

- Statistically optimal observables
- Angular observables – less sensitivity but easier implementation
- Machine-learning observables – attempt to recover best sensitivity while keeping feasibility



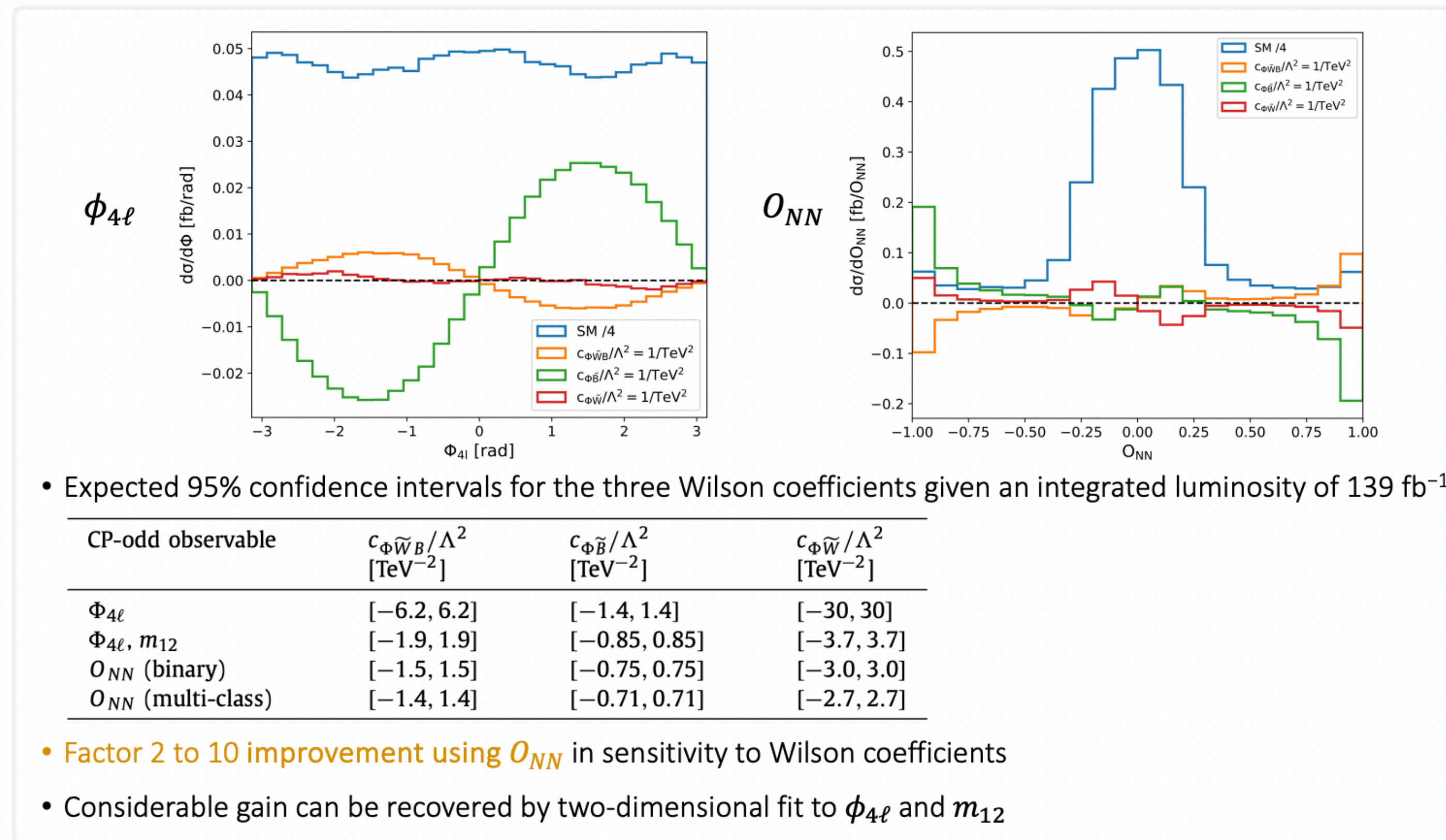
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## Neural network-based observable: $h \rightarrow 4\ell$ results



# Reusing Neural Networks: Experiences and suggestions for EFT cases

Tomasz Procter

[Link to slides](#)

## Initial experiences with reinterpretation



- So far: two publicly available LHC analysis NN networks - both from ATLAS SUSY:
  - [ANA-SUSY-2019-04](#) (RPV SUSY search in lep+jets final state)
  - [ANA-SUSY-2018-30](#) (gluino pair production in multi-b final states)
- This is a new type of experimental output - the experiments are still feeling out how to publicise this.
- ANA-SUSY-2018-30 has worked well in multiple frameworks (rivet, gambit, checkmate, ...).
- Several key features made this work:
  - Lots of extra info (ordering, units, usage example) - would have been impossible without SimpleAnalysis<sup>1</sup>.
  - All inputs are easily accessible to reinterpretation tools
    - Lepton/jet kinematics, MET, btag yes/no
    - No detector-level variables (including continuous btag score)

<u>Cut</u>	<u>Paper</u>	<u>Rivet</u>
0-lep	80.0	83.7
$\Delta\phi_{\min}^{4j} \geq 0.6$	52.5	54.6
2800-1400 NN Cut	21.7	23.9
$\Delta\phi_{\min}^{4j} \geq 0.6$	52.5	54.6
2300-1000 NN Cut	21.3	23.3
$\Delta\phi_{\min}^{4j} \geq 0.4$	61.1	63.8
2100-1600 NN Cut	6.20	6.50
$\Delta\phi_{\min}^{4j} \geq 0.4$	61.1	63.8
2000-1800 NN Cut	0.192	0.204

Tomasz Procter, LHC EFT WG Area 3, October 2023

<sup>1</sup>[Simplified ATLAS SUSY analysis framework](#)

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# Reusing Neural Networks: Experiences and suggestions for EFT cases

Tomasz Procter

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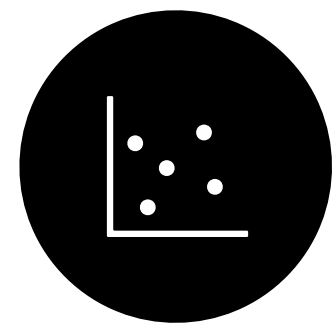
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- Key rule:  
***Reinterperatation is easiest when the analysis team think about it from the start***
  - Make sure models can be saved in a preservable format.
  - Example code snippets, metadata is very important.
  - Think about choice of inputs:
    - Do we need to use efficiencies/surrogates instead?



## **ML4EFT:** unbinned multivariate observables for global SMEFT fits

*2211.02058 Raquel Gomez Ambrosio, Jaco ter Hoeve, MM, Juan Rojo, Veronica Sanz*

# Why Unbinned Measurements?

*'Presenting Unbinned Differential Cross Section Results', Arratia et al, 2109.13243*

- 1. Inference-aware binning:**
  - optimal choice of binning can be made at the time of *each* statistical analysis or global fit

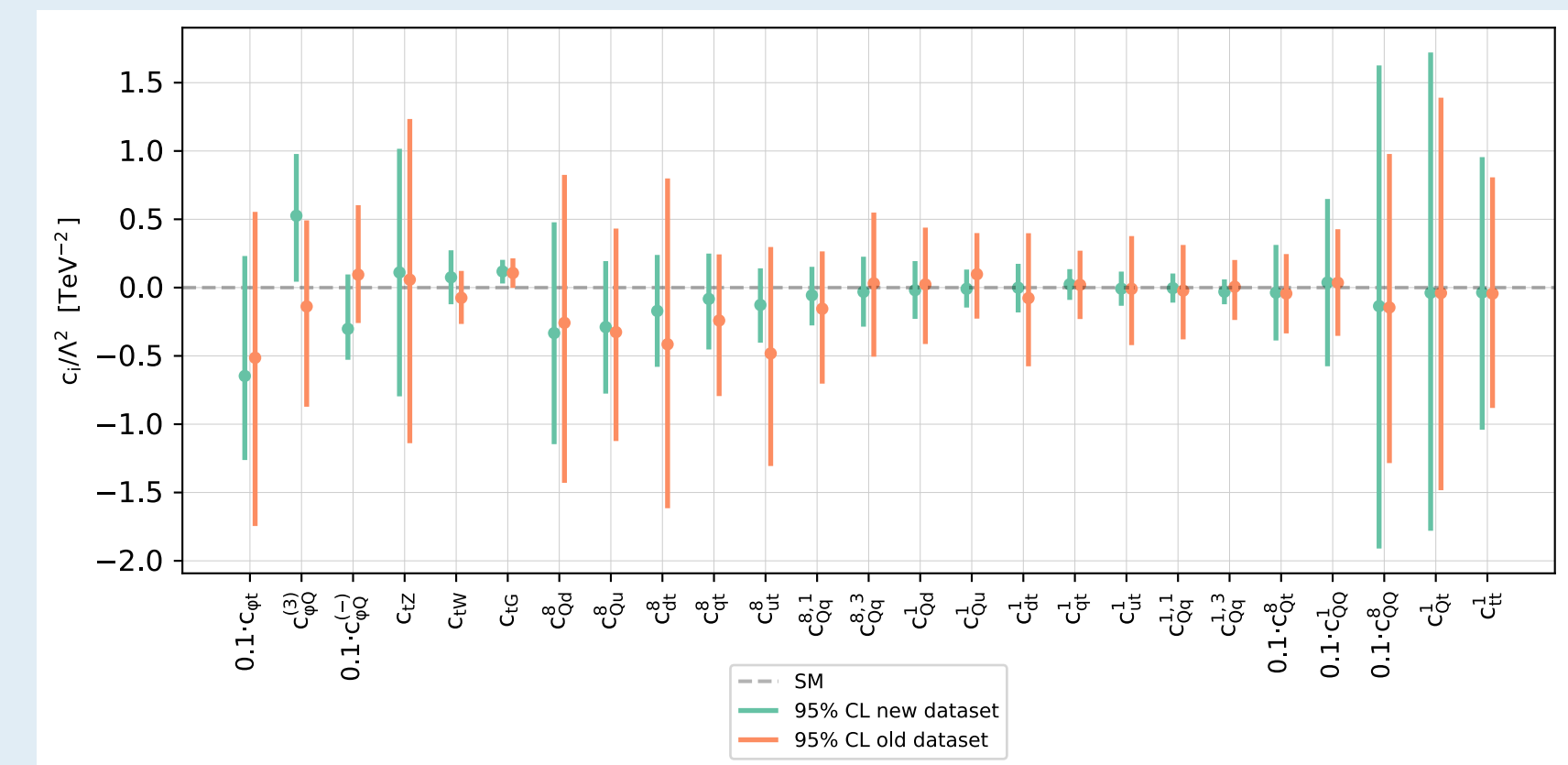
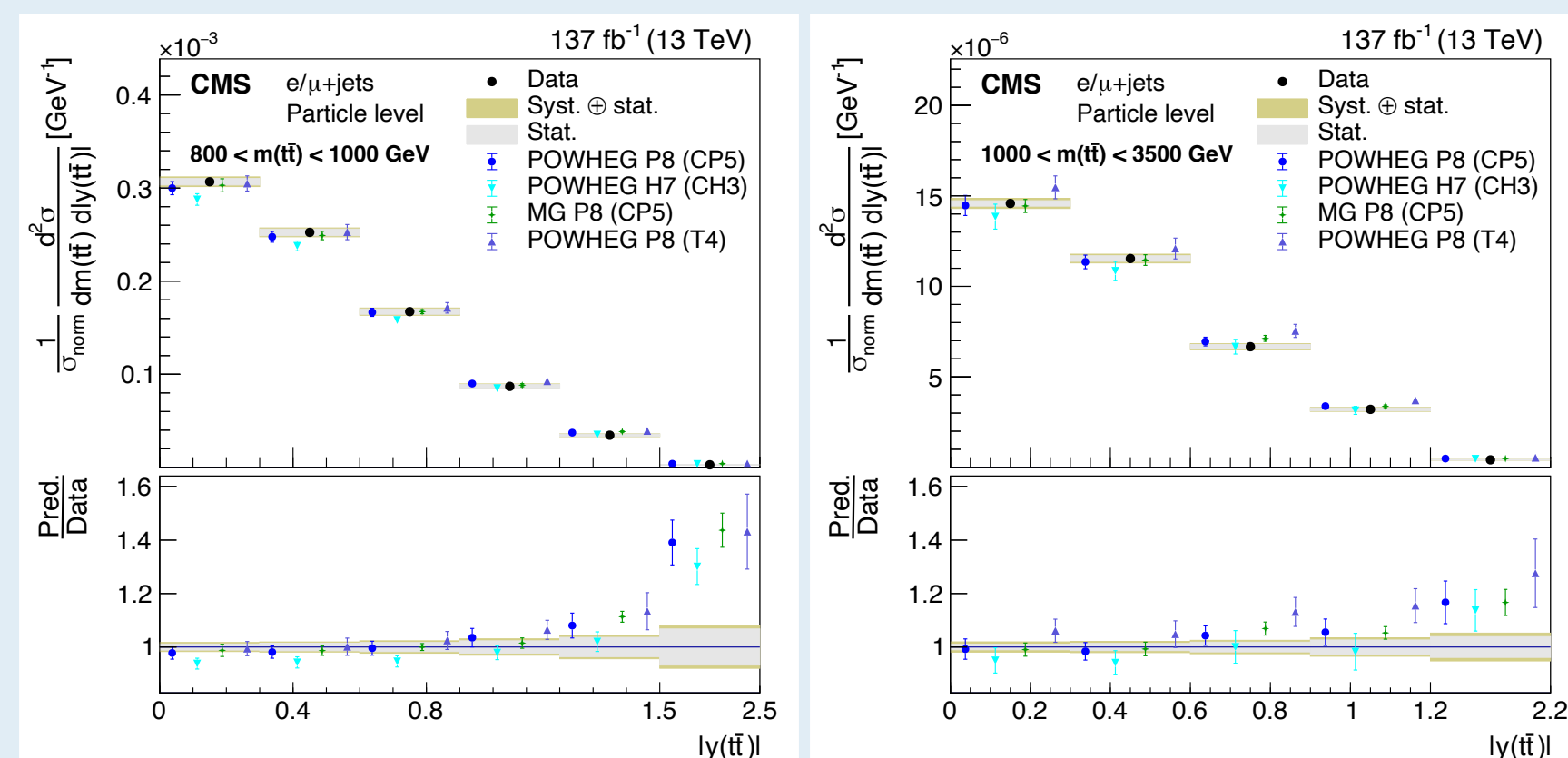


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Typically we **reinterpret** measurements optimised for SM measurements or NP resonance searches



*e.g. CMS measurement of top pair production in the l+jets channel 2108.02803*

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- 3. Extension to higher dimensions:** ▶ ML-based unbinned unfolding techniques well-suited to multiple features

# ML4EFT

<https://lhcfiteknikhef.github.io/ML4EFT/>

*2211.02058 Raquel Gomez Ambrosio, Jaco ter Hoeve, MM, Juan Rojo, Veronica Sanz*

Open-source NN-based python framework for the integration of **unbinned multivariate observables** into global SMEFT interpretations.

**Goal: to provide optimal constraints on the SMEFT**

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**Diagnostic tool:**

What is the information loss given a particular choice of bins?

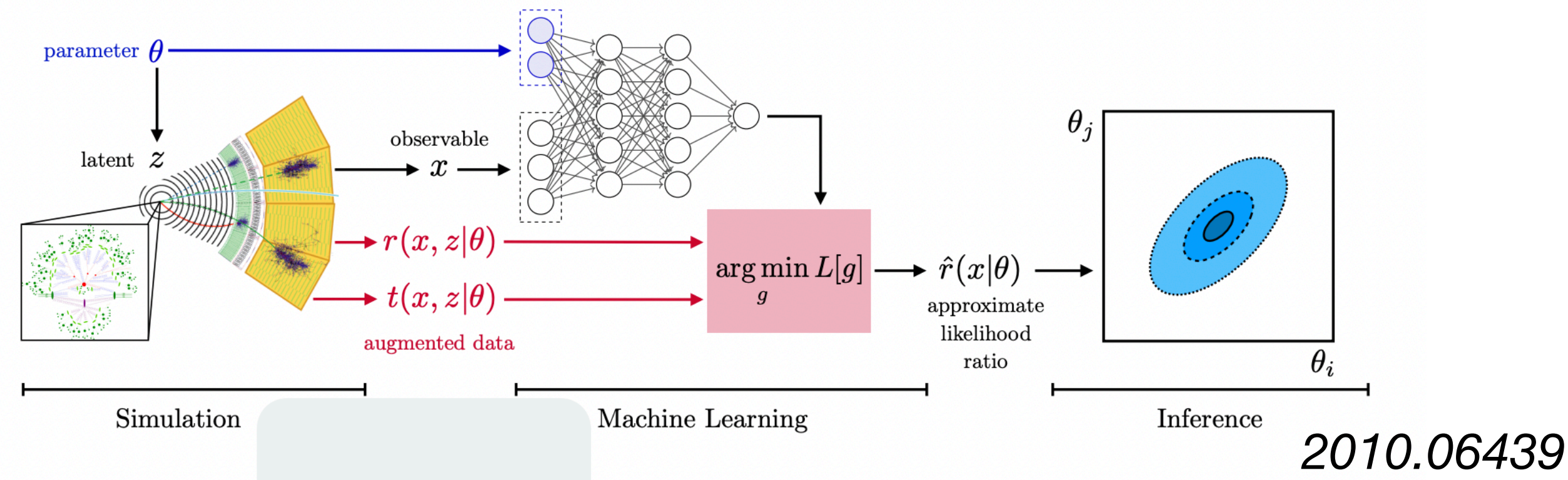
**Projections:**

If unbinned data are made available, how will SMEFT constraints improve?

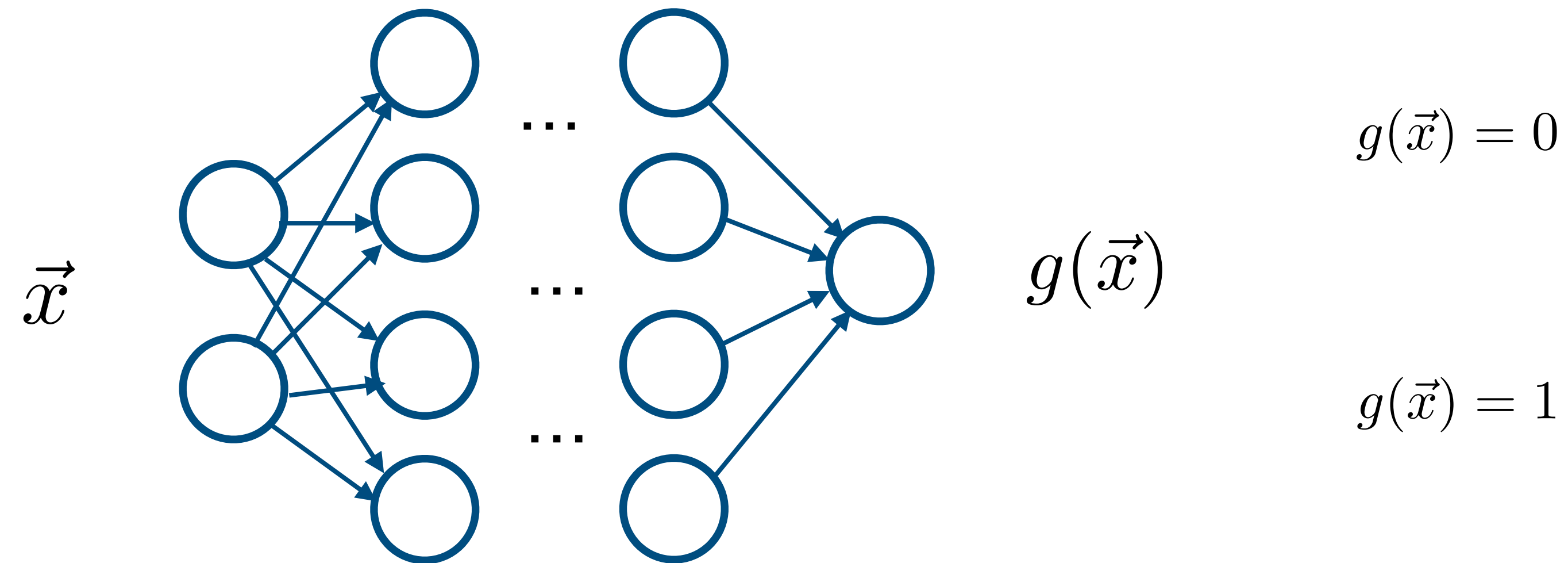
## Open-source NN-based python framework for the integration of **unbinned multivariate observables** into global SMEFT interpretations.

Related work:

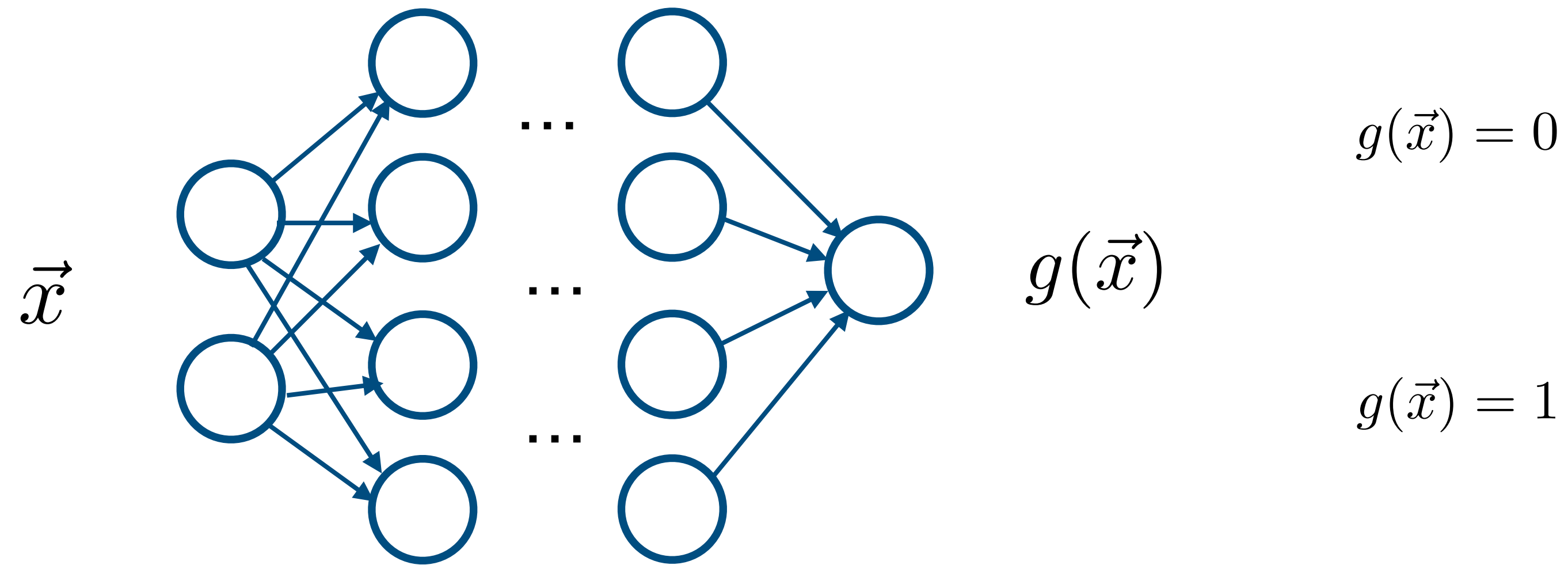
- 2007.10356 *Parameterized classifiers for SMEFT* A. Glioti et al.
- 2308.05704 *Boosted likelihood learning with event reweighting* A. Glioti et al
- 2205.12976 *Learning the EFT likelihood with tree boosting* R. Schöfbeck et al
- *MadMiner* (J.Brehmer, K.Cranmer, G.Louppe et al.) [1907.10621, 1805.00020, ...]
- + many others



# The 'classifier trick'



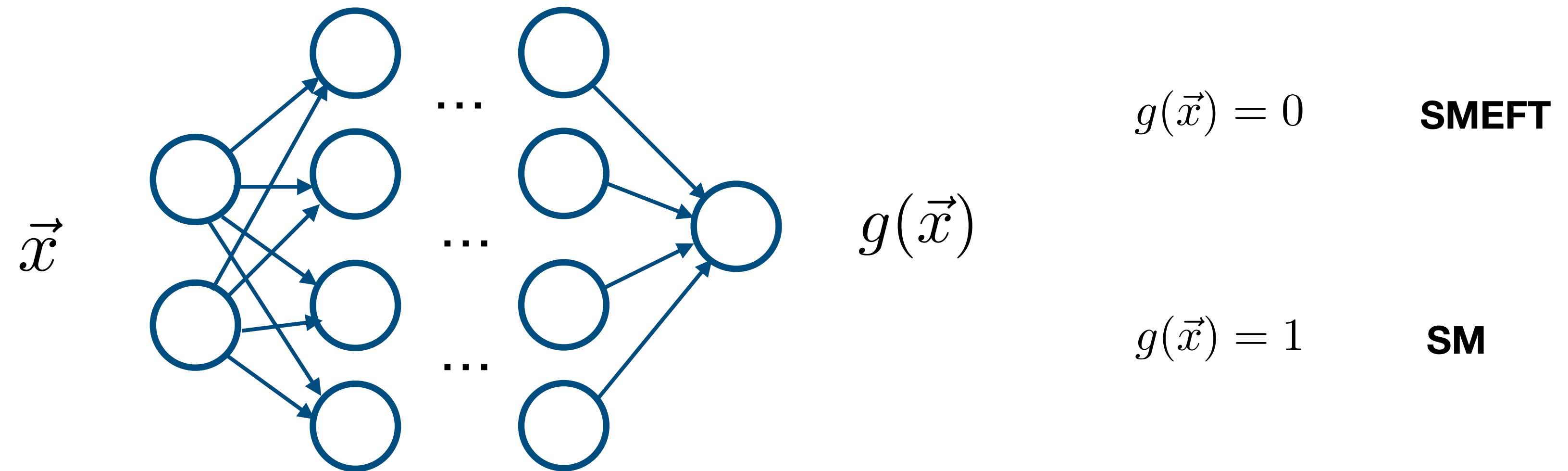
# The 'classifier trick'



$$\vec{x} = \{m_{t\bar{t}}, p_T^{\ell_1}, p_T^{\ell_2}, \Delta\eta_{\ell_1, \ell_2}, \Delta\phi_{\ell_1, \ell_2}, \dots\}$$



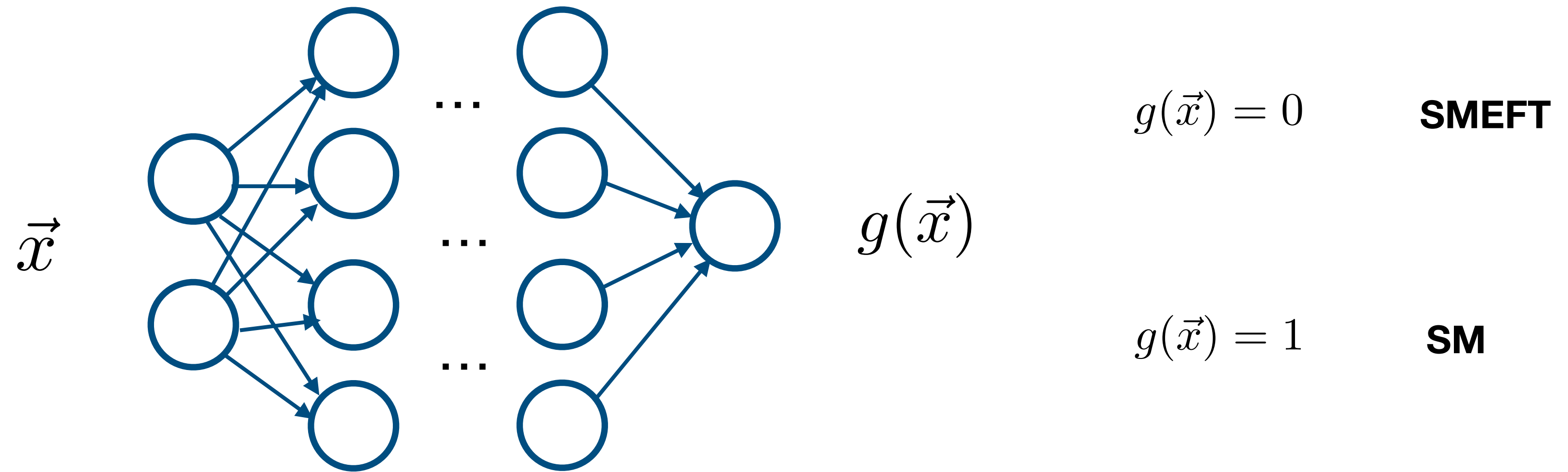
# The 'classifier trick'



$g(\vec{x}) = 0$  **SMEFT**

$g(\vec{x}) = 1$  **SM**

# The 'classifier trick'



Loss function:

$$L[g(\mathbf{x}, \mathbf{c})] = - \int d\mathbf{x} \frac{d\sigma(\mathbf{x}, \mathbf{c})}{d\mathbf{x}} \log(1 - g(\mathbf{x}, \mathbf{c})) - \int d\mathbf{x} \frac{d\sigma(\mathbf{x}, \mathbf{0})}{d\mathbf{x}} \log g(\mathbf{x}, \mathbf{c})$$

**SMEFT**
**SM**

# The 'classifier trick'

$$L[g(\mathbf{x}, \mathbf{c})] = - \int d\mathbf{x} \frac{d\sigma(\mathbf{x}, \mathbf{c})}{d\mathbf{x}} \log(1 - g(\mathbf{x}, \mathbf{c})) - \int d\mathbf{x} \frac{d\sigma(\mathbf{x}, \mathbf{0})}{d\mathbf{x}} \log g(\mathbf{x}, \mathbf{c})$$

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$$\frac{\delta L}{\delta g} = 0 \Rightarrow g(\mathbf{x}, \mathbf{c}) = \left( 1 + \frac{d\sigma(\mathbf{x}, \mathbf{c})}{d\mathbf{x}} / \frac{d\sigma(\mathbf{x}, \mathbf{0})}{d\mathbf{x}} \right)^{-1} \equiv \frac{1}{1 + r_\sigma(\mathbf{x}, \mathbf{c})}$$

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or: 
$$\frac{d\sigma(x, c)/dx}{d\sigma(x, 0)/dx} = \frac{1 - g(x, c)}{g(x, c)}$$

# Parametrised classifier

$$r_{\sigma}(x, c) = \frac{d\sigma(x, c)/dx}{d\sigma(x, 0)/dx} = \frac{1 - g(x, c)}{g(x, c)}$$

Exploit the polynomial structure of the SMEFT when defining the classifier  $g$ :

$$\hat{r}_{\sigma}(\mathbf{x}, \mathbf{c}) = 1 + \sum_{j=1}^{n_{\text{eft}}} \text{NN}^{(j)}(\mathbf{x}) c_j + \sum_{j=1}^{n_{\text{eft}}} \sum_{k \geq j}^{n_{\text{eft}}} \text{NN}_{\sigma}^{(j,k)}(\mathbf{x}) c_j c_k$$

$$\text{c.f. } k_i(C) = 1 + r_{\text{lin},i} C + r_{\text{quad},i} C^2 \quad i = 1, \dots, n_{\text{bins}}$$

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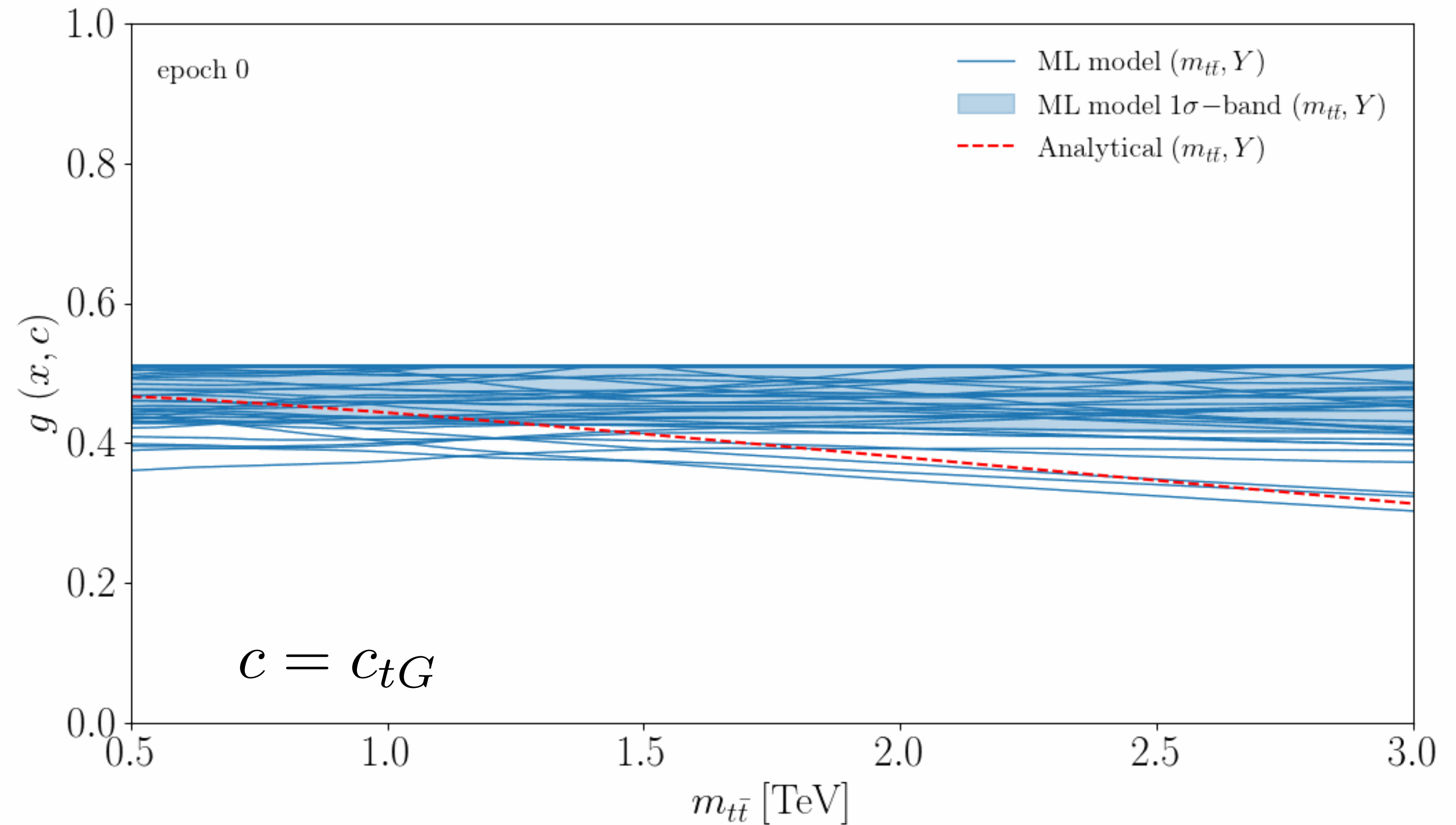
**Parallelisable:** generate a training sample with only  $c_i$  and learn only  $\text{NN}^i(\mathbf{x})$

well-suited to global fits of many SMEFT coefficients

# Proof of concept: top quark pair production

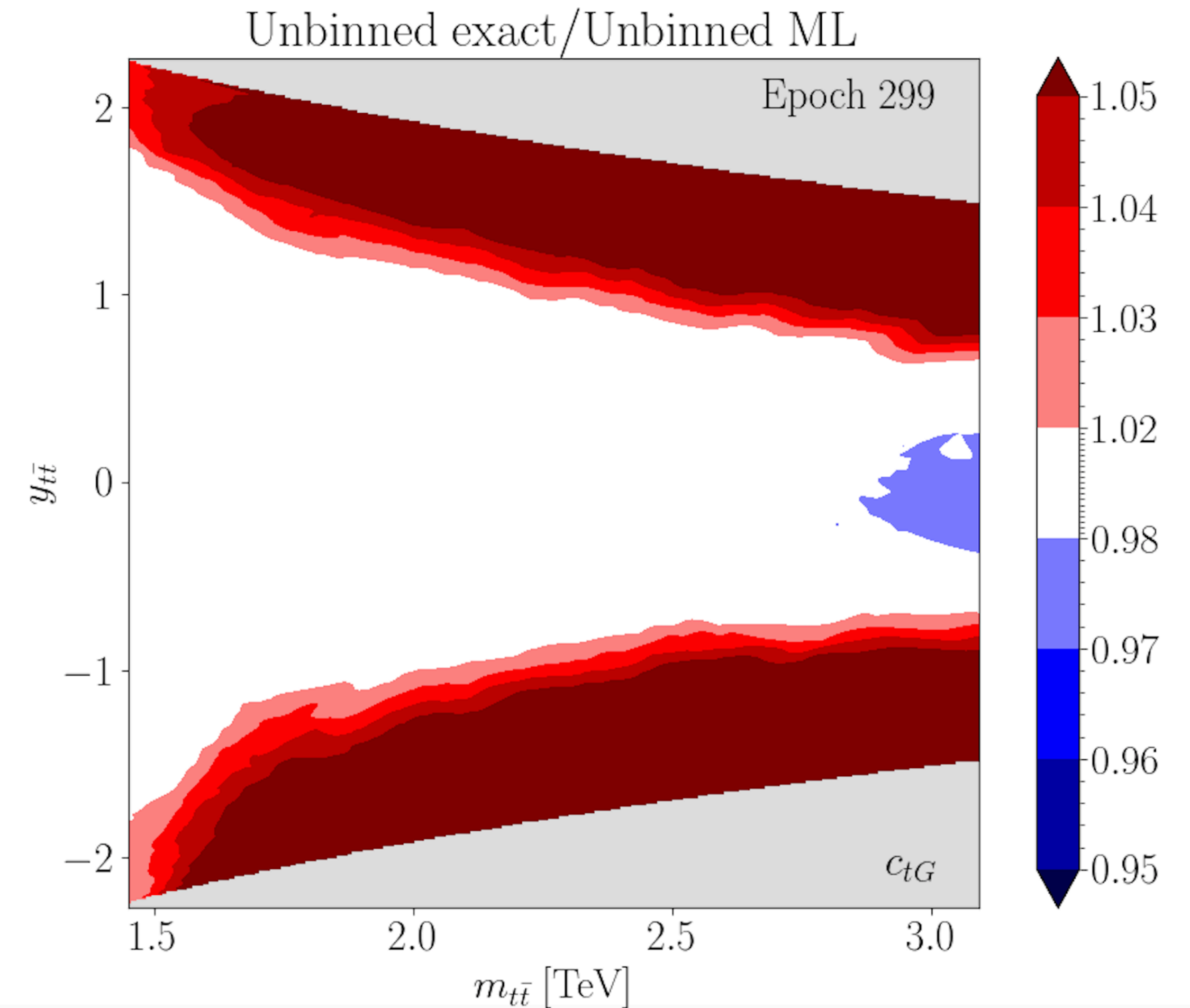
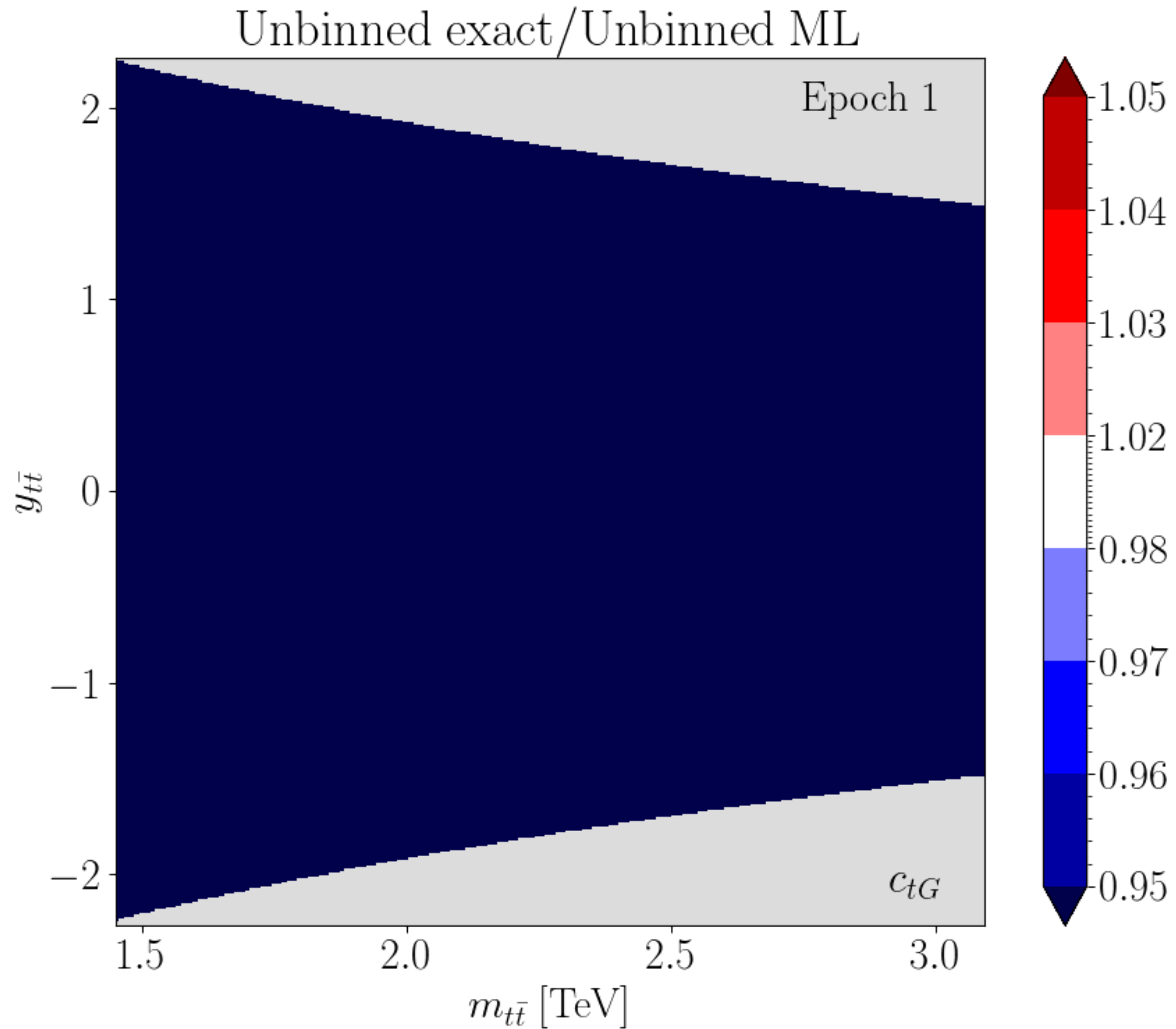
Validation against the analytical calculation of  $\frac{d^2\sigma}{dm_{t\bar{t}}dy_{t\bar{t}}}$  for **parton-level**  $t\bar{t}$  production

Train multiple instances of **g** to quantify the impact of finite training data samples





# Proof of concept: top quark pair production

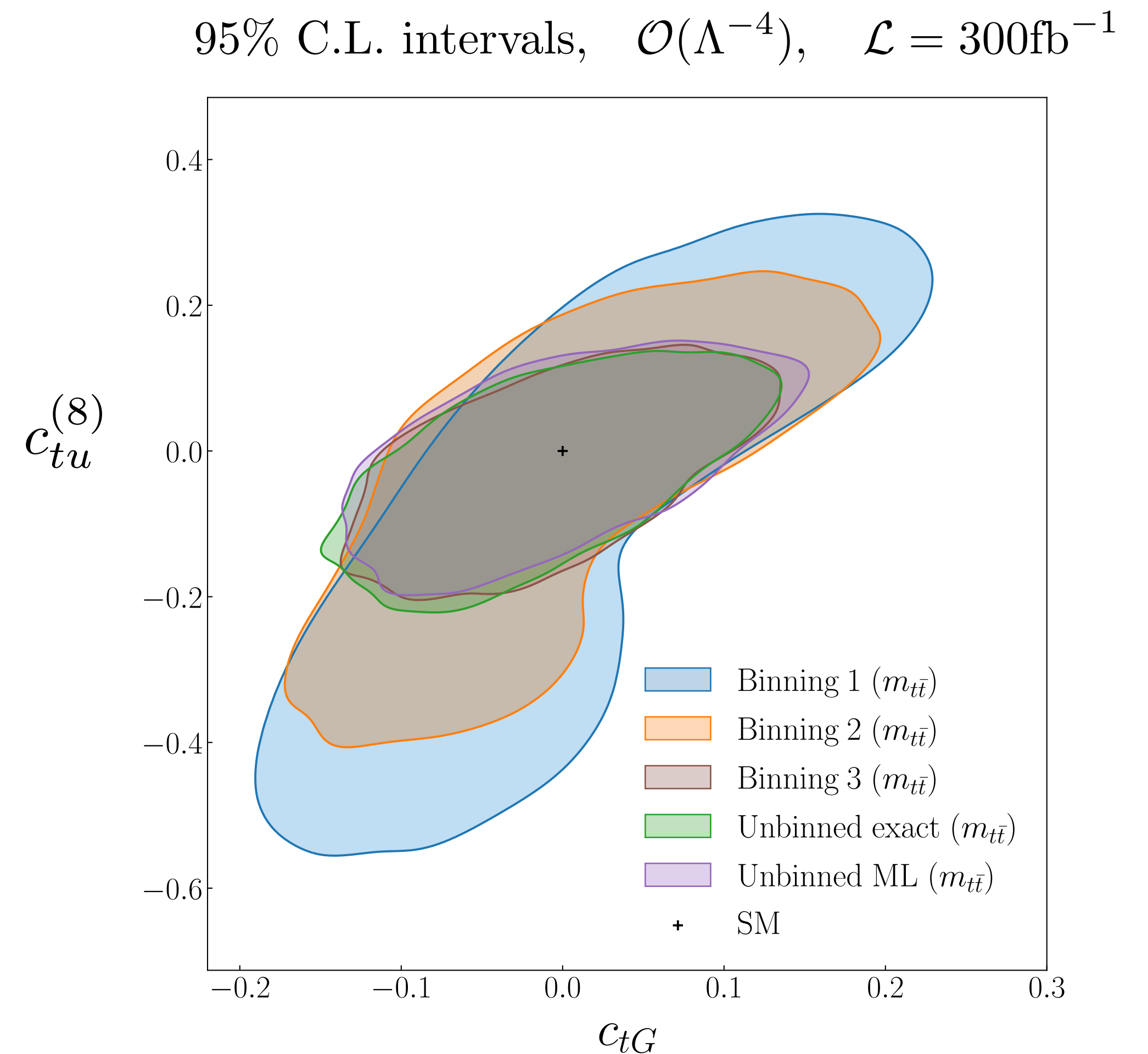


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- Unbinned **exact** and unbinned **ML** agree:

## ▸ validation of methodology



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▸ **validation of methodology**

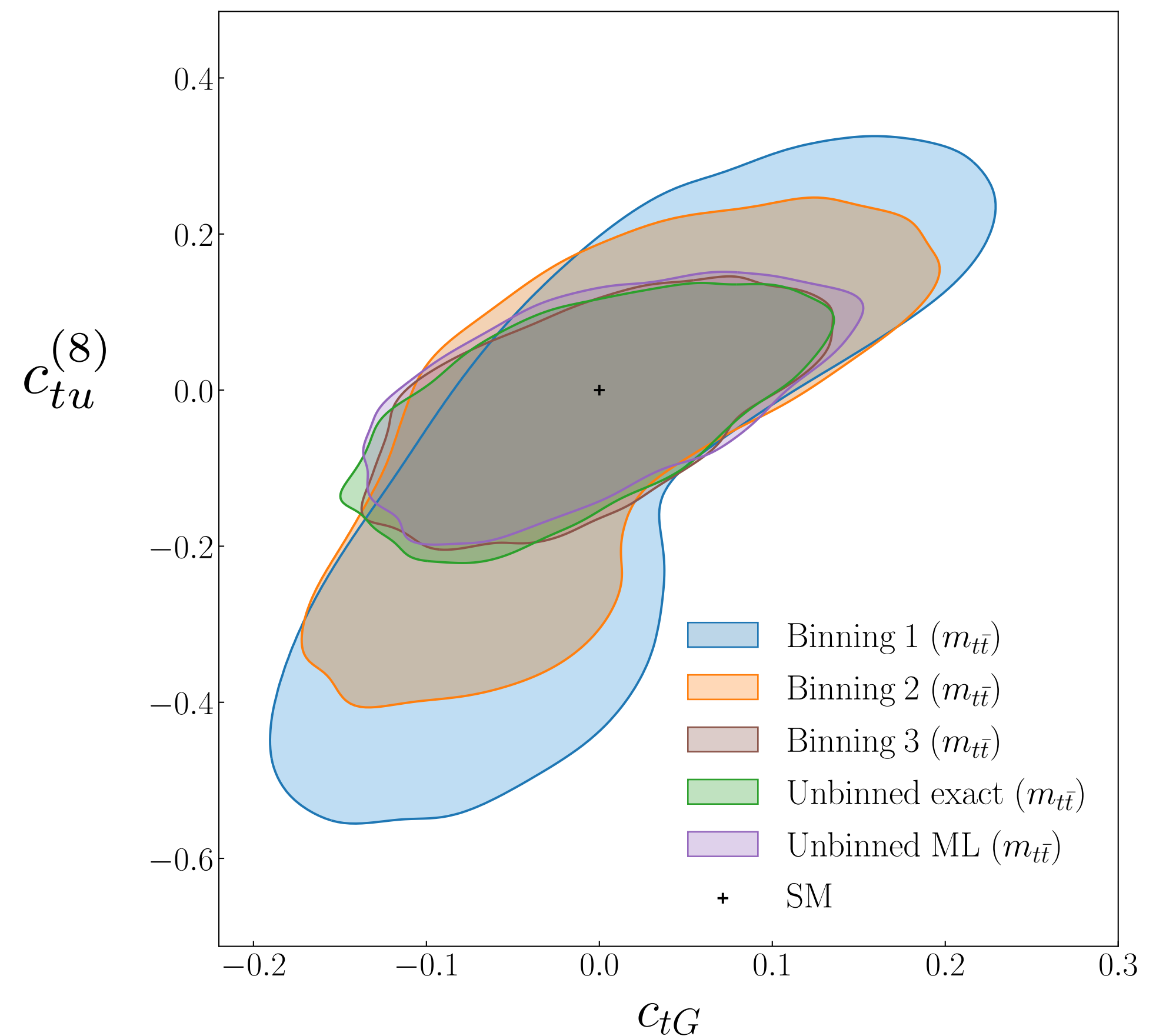
Binning 1 -> binning 2 -> binning 3

———— finer binning in  $m_{t\bar{t}}$  —————>

- Binning converges to unbinned constraints with finer binning:

**binning 3:**  $m_{t\bar{t}} \in [1.45, 1.5, 1.55, 1.6, 1.7, 1.8, 1.9, 2.0, 2.1, 2.2, 2.3, 2.4, 2.5, 2.6, 2.7, 2.8, 2.9, 3.0, 3.1, 3.2, 3.3, 3.4, 3.5, 3.6, 3.7, 3.8, 3.9, 4.0, 4.1, 4.2, 4.3, 4.4, 4.5, 4.6, 4.7, 4.8, 4.9, 5.0, \infty)$  TeV.

95% C.L. intervals,  $\mathcal{O}(\Lambda^{-4})$ ,  $\mathcal{L} = 300\text{fb}^{-1}$



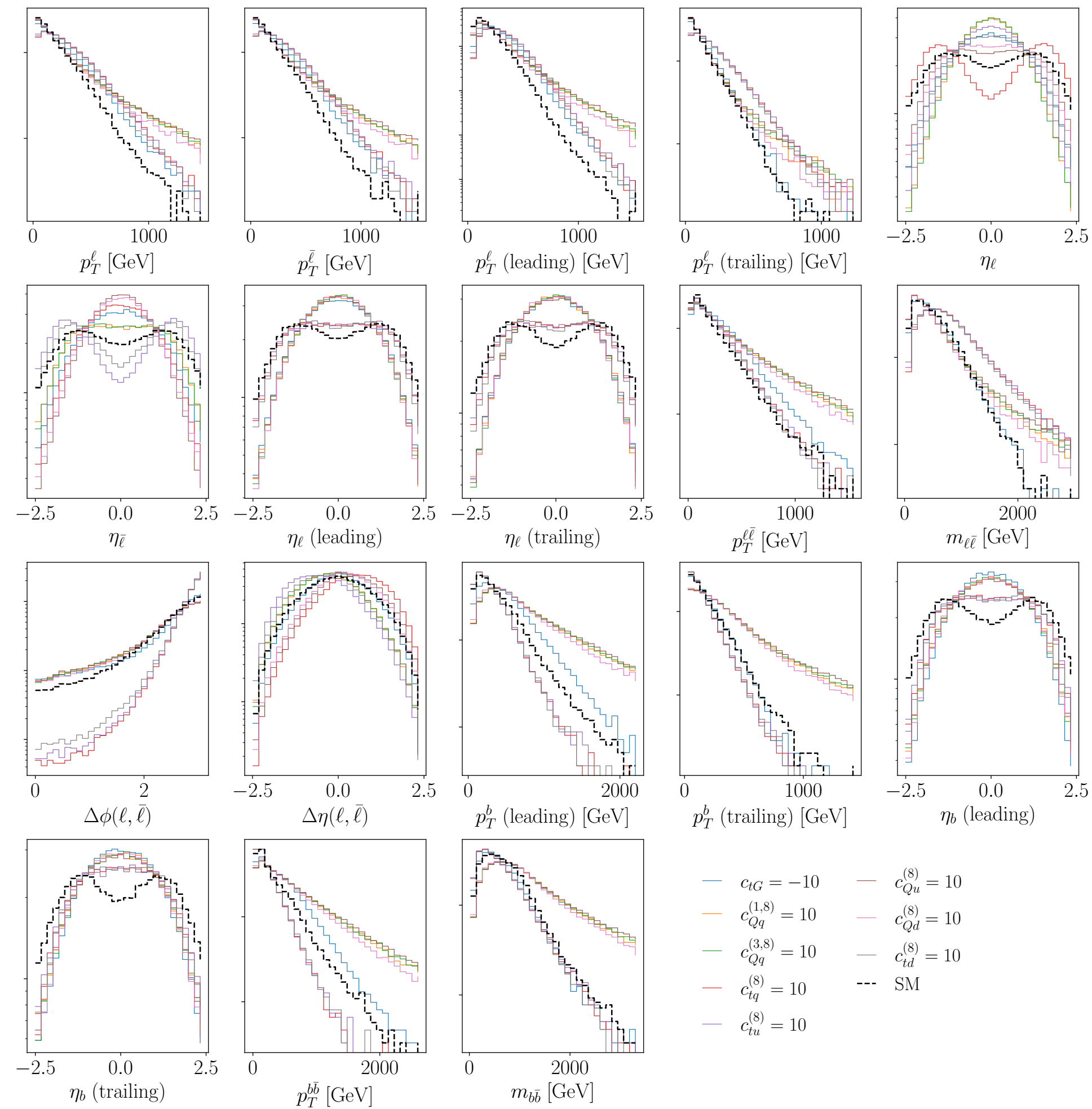
# Unbinned observables in the top sector

Particle-level top quark pair production in the dileptonic channel:

$$pp \rightarrow t\bar{t} \rightarrow \ell^+ \ell^- b\bar{b} \nu_\ell \bar{\nu}_\ell$$

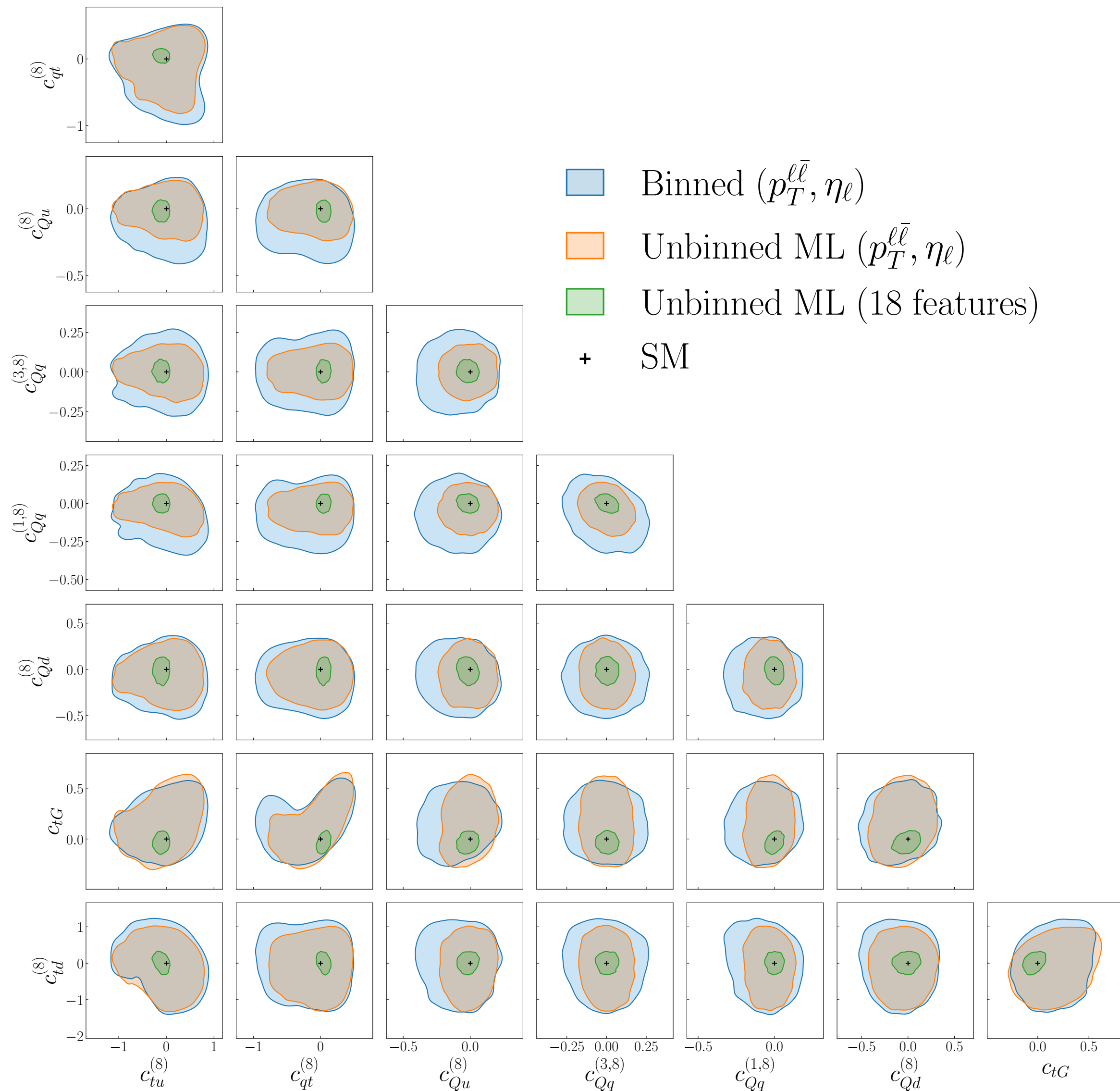
Constraints on 8 SMEFT operators:

$$O_{tG} + 4\text{-fermion operators}$$



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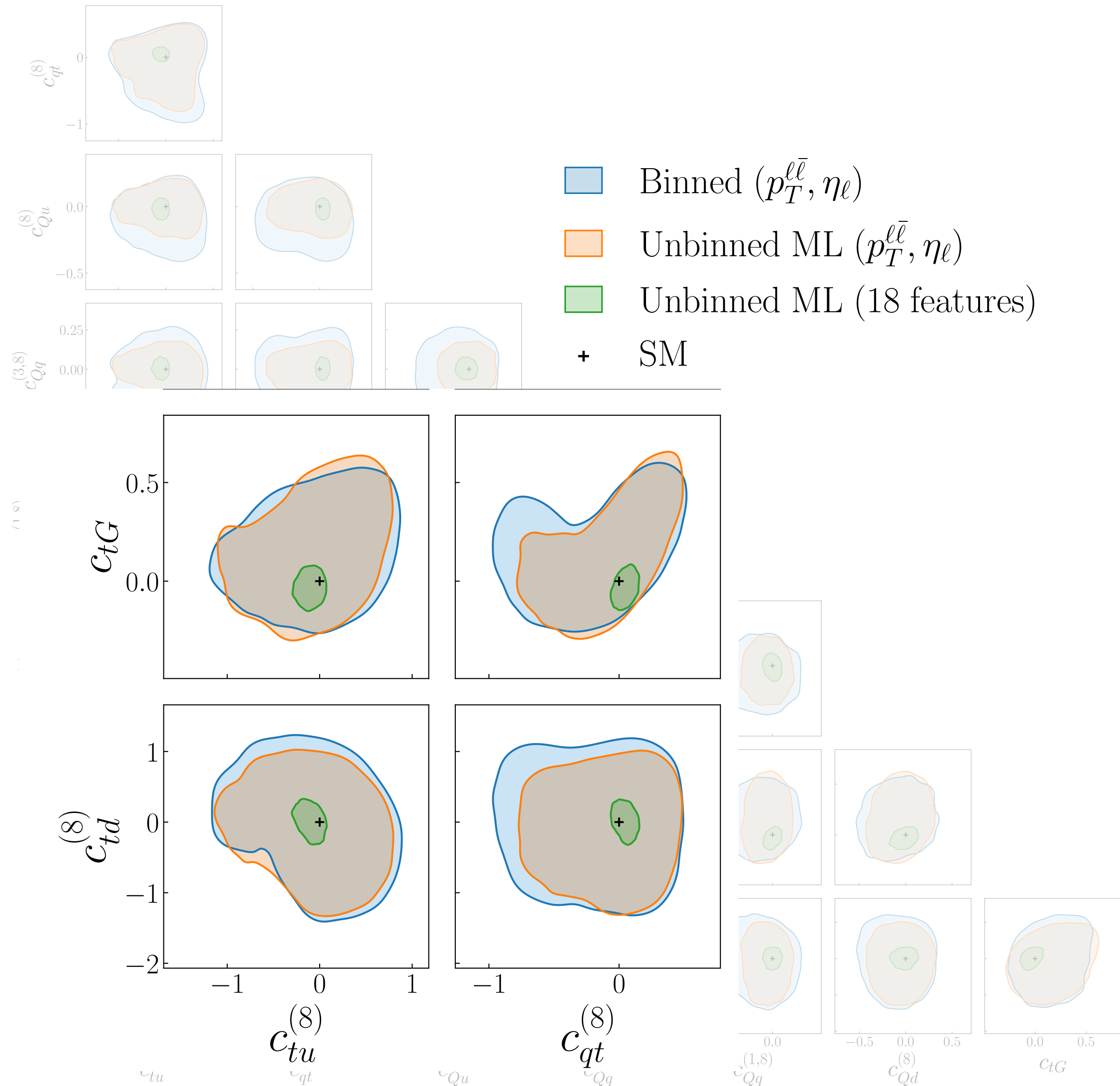
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# Unbinned observables in the top sector

Marginalised 95 % C.L. intervals,  $\mathcal{O} (\Lambda^{-4})$  at  $\mathcal{L} = 300 \text{ fb}^{-1}$

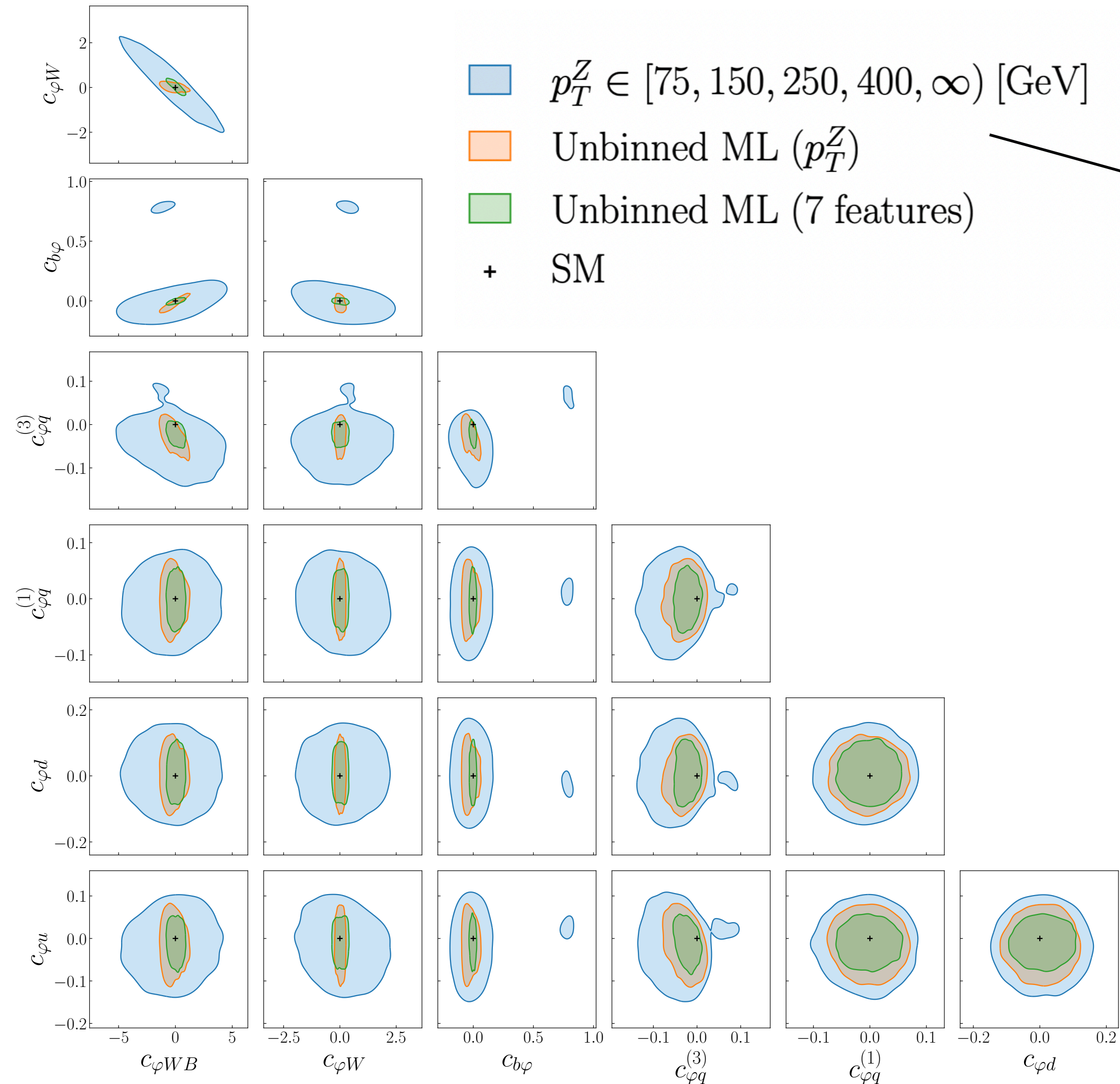


**Binned vs unbinned** in  $(p_T^{ll}, \eta_l)$  : small improvement from unbinned measurements, relative to nominal choice of bins

**2 features vs 18 features:** vast improvement in constraining power

# Unbinned observables in the Higgs sector

Marginalised 95 % C.L. intervals,  $\mathcal{O}(\Lambda^{-4})$  at  $\mathcal{L} = 300 \text{ fb}^{-1}$



$$pp \rightarrow hZ \rightarrow b\bar{b}l^+l^-$$

*STXS binning - see also 1908.06980, Brehmer et. al*

Constraints on 7 SMEFT coefficients:

$$c_{\varphi u}, c_{\varphi d}, c_{\varphi q}^{(1)}, c_{\varphi q}^{(3)}, c_{\varphi W}, c_{\varphi WB}, c_{b\varphi}$$

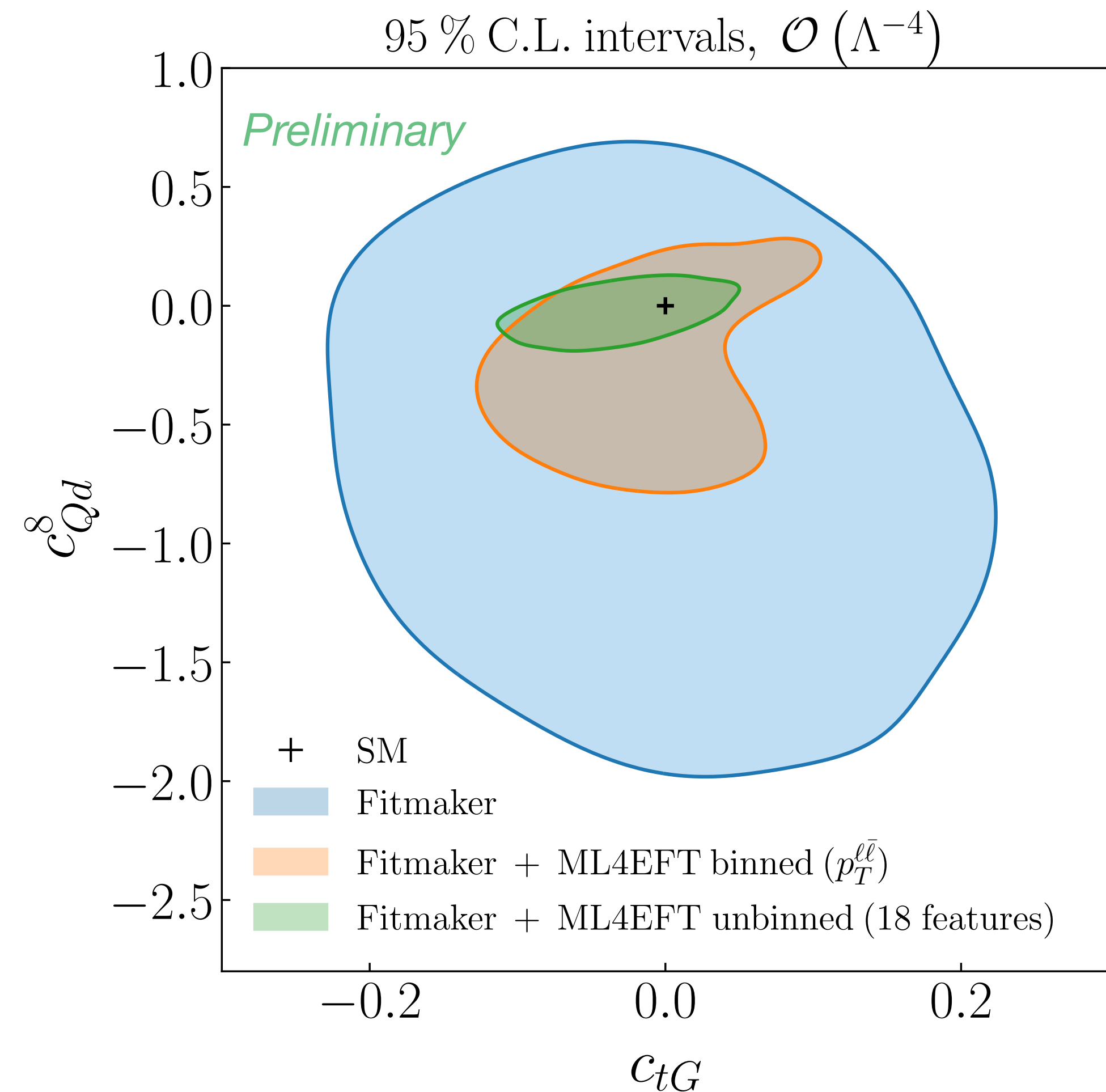
# Future directions

- treatment of systematic uncertainties
- New unbinned measurements can be combined alongside existing binned measurements:

$$\log \mathcal{L}(c) = \sum_{k=1}^{N_D^{(\text{unbinned})}} \log \mathcal{L}_k^{\text{unbinned}}(c) + \sum_{k=1}^{N_D^{(\text{binned})}} \log \mathcal{L}_k^{\text{binned}}(c)$$

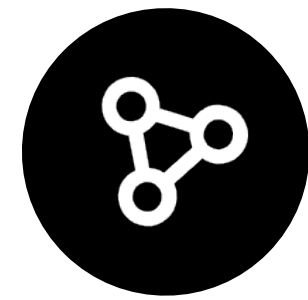
- incorporate parton-showered observables

*Work in progress, Pim Herbschleb, Jaco ter Hoeve*



*Work in progress, Jaco ter Hoeve, MM*



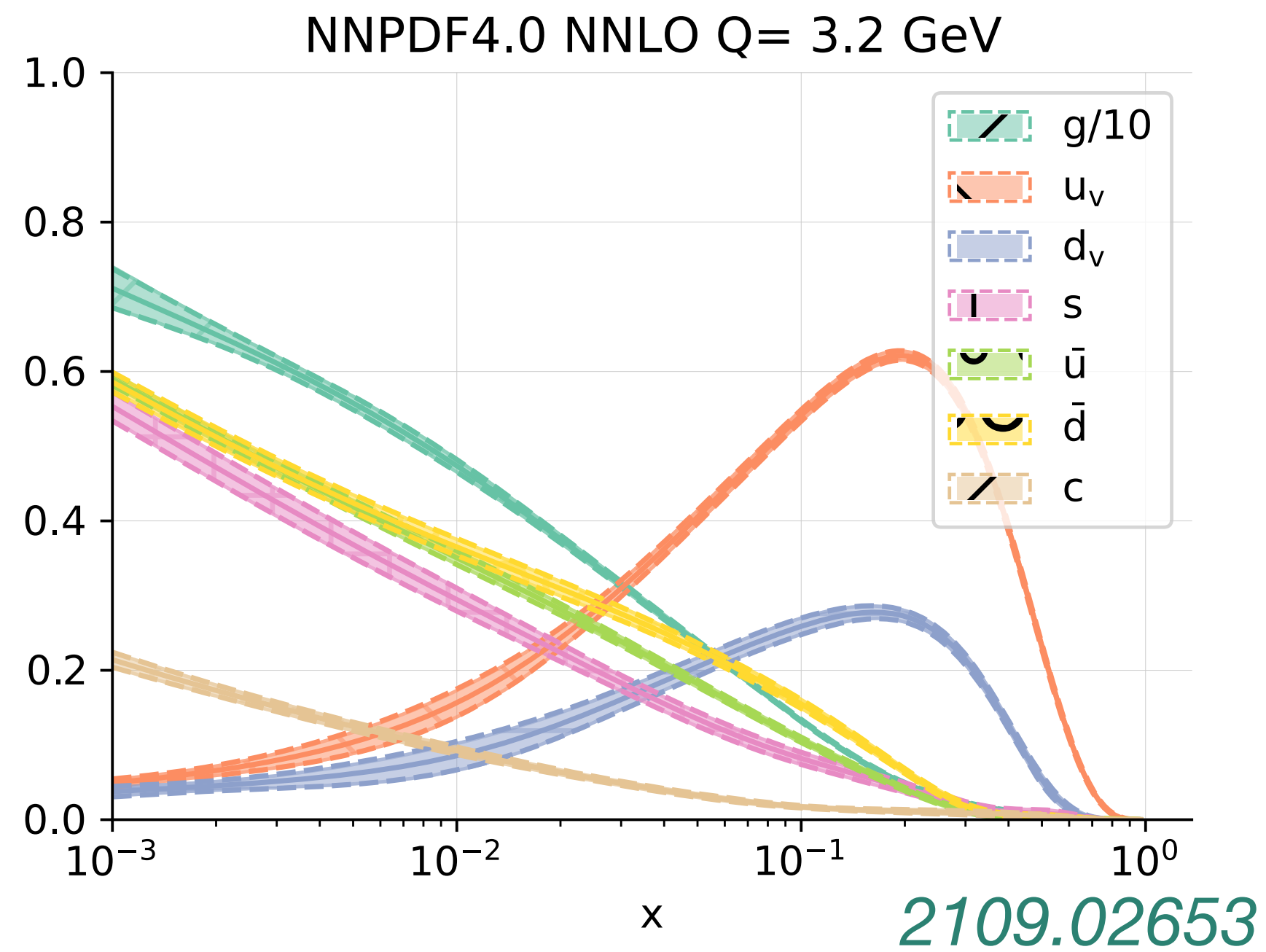


## **SIMUnet:** an open-source tool for the simultaneous fit of PDFs and SMEFT coefficients

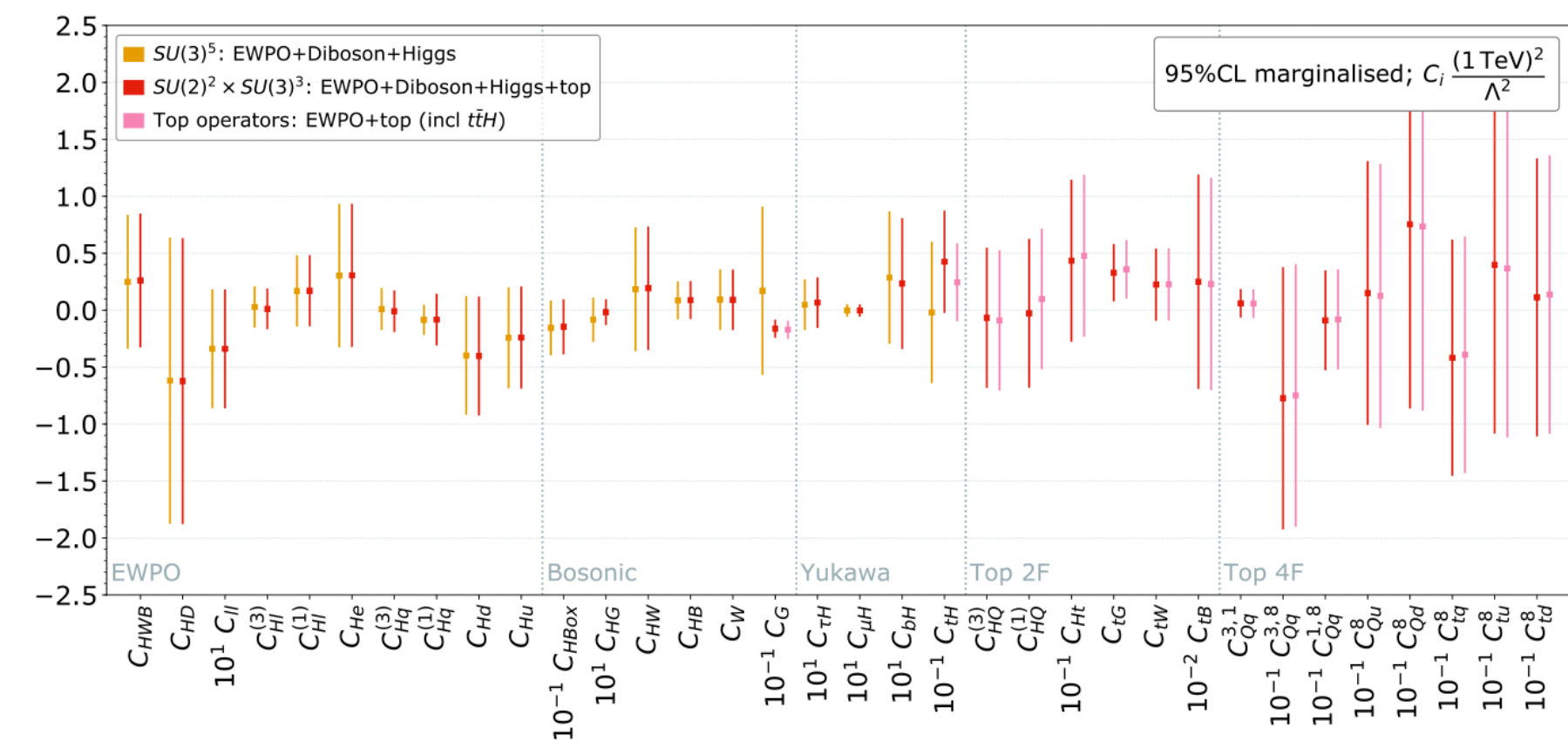
*Work in progress by Elie Hammou, Maeve Madigan, Luca Mantani, James Moore, Manuel Morales Alvarado, Mark Nestor Costantini, Maria Ubiali*

# PDF-EFT Interplay

## PDF fits



## SMEFT Fits and BSM searches



2012.02779

# PDF-EFT Interplay

## PDF fits

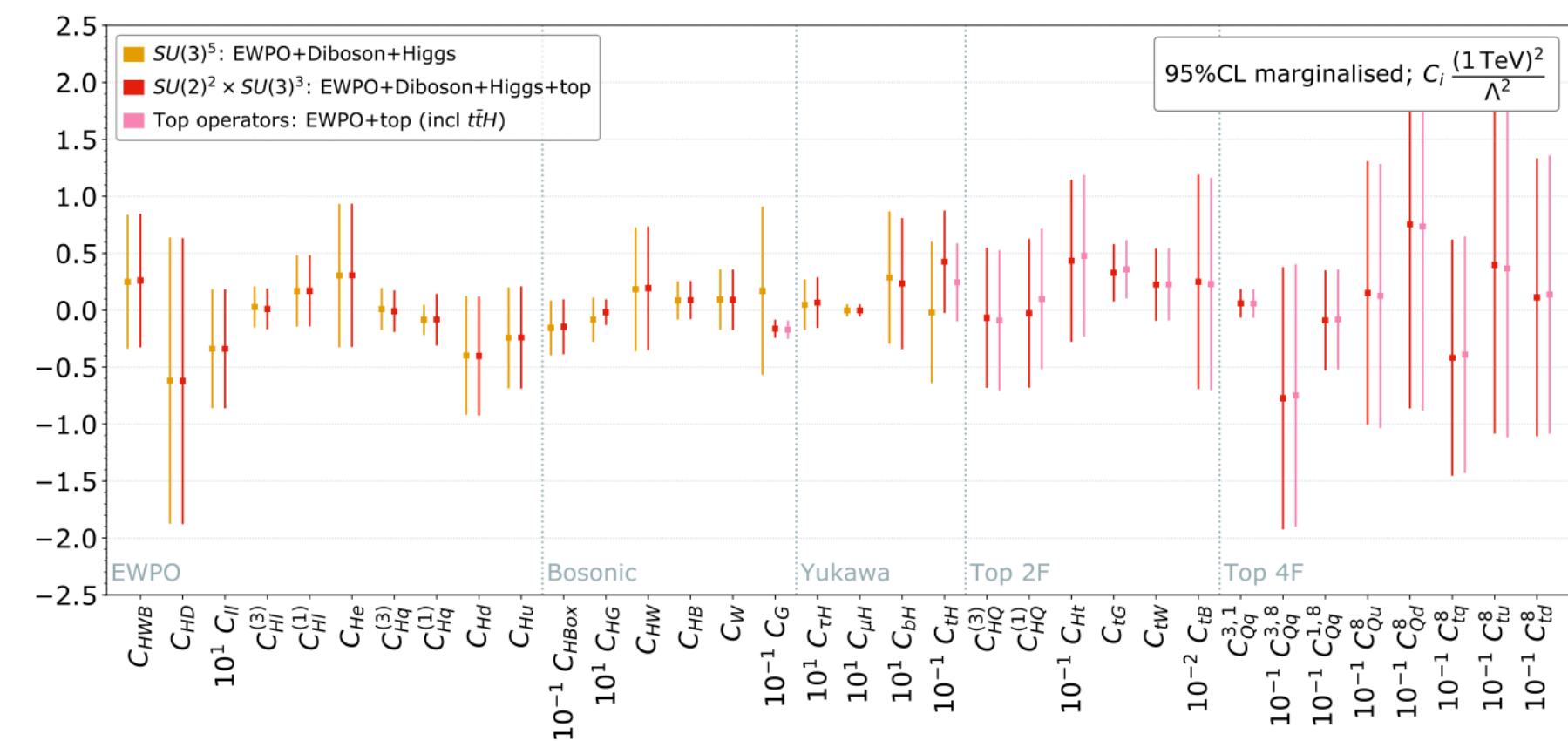
BSM parameters are kept fixed:

$$\sigma(\bar{c}, \theta) = f_1(\theta) \otimes f_2(\theta) \otimes \hat{\sigma}(\bar{c})$$

Typically PDF fits assume the SM:

$$\bar{c} = 0$$

## SMEFT Fits and BSM searches



2012.02779

# PDF-EFT Interplay

## PDF fits

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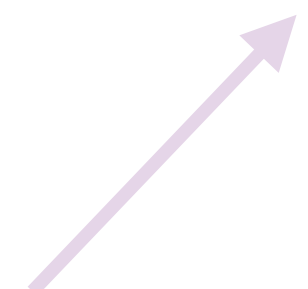
Typically PDF fits assume the SM:  
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## SMEFT Fits and BSM searches

PDF parameters are fixed:

$$\sigma(c, \bar{\theta}) = f_1(\bar{\theta}) \otimes f_2(\bar{\theta}) \otimes \hat{\sigma}(c)$$

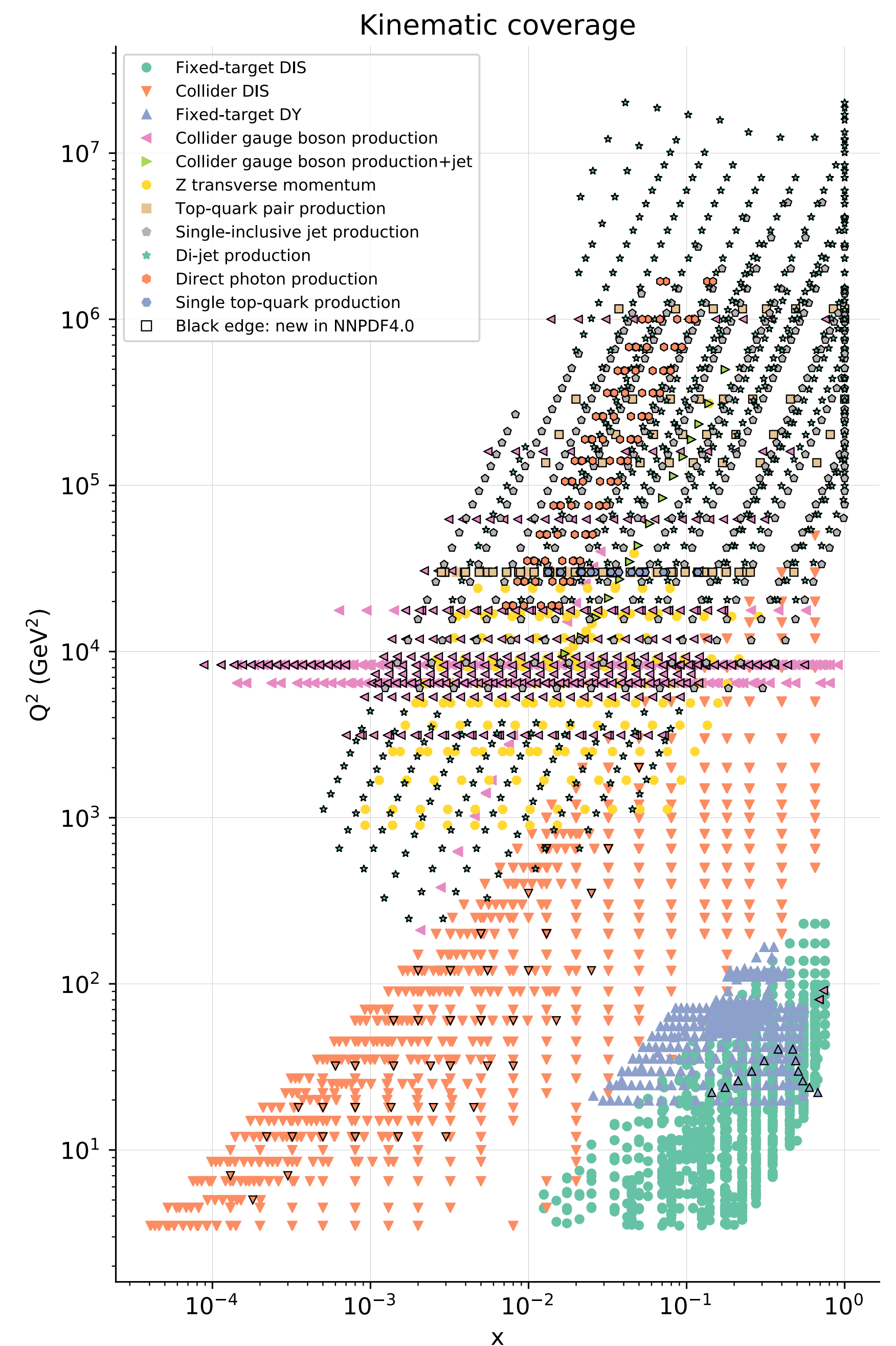
PDFs used in BSM searches rely on SM assumptions



# Data overlap

Often the data used in PDF fits are also used in EFT fits.

This overlap will grow as we take the global approach to constraining the SMEFT.

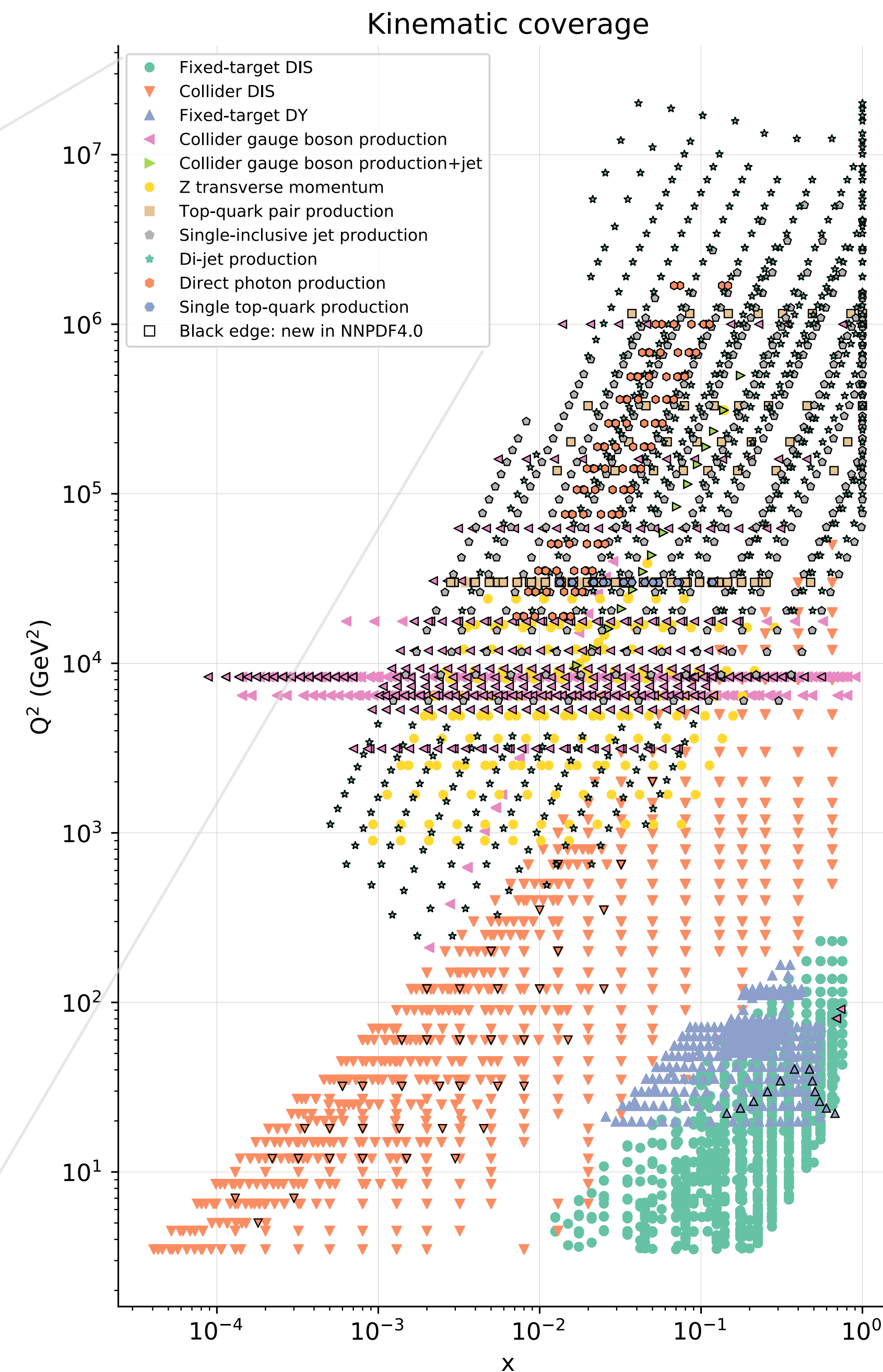
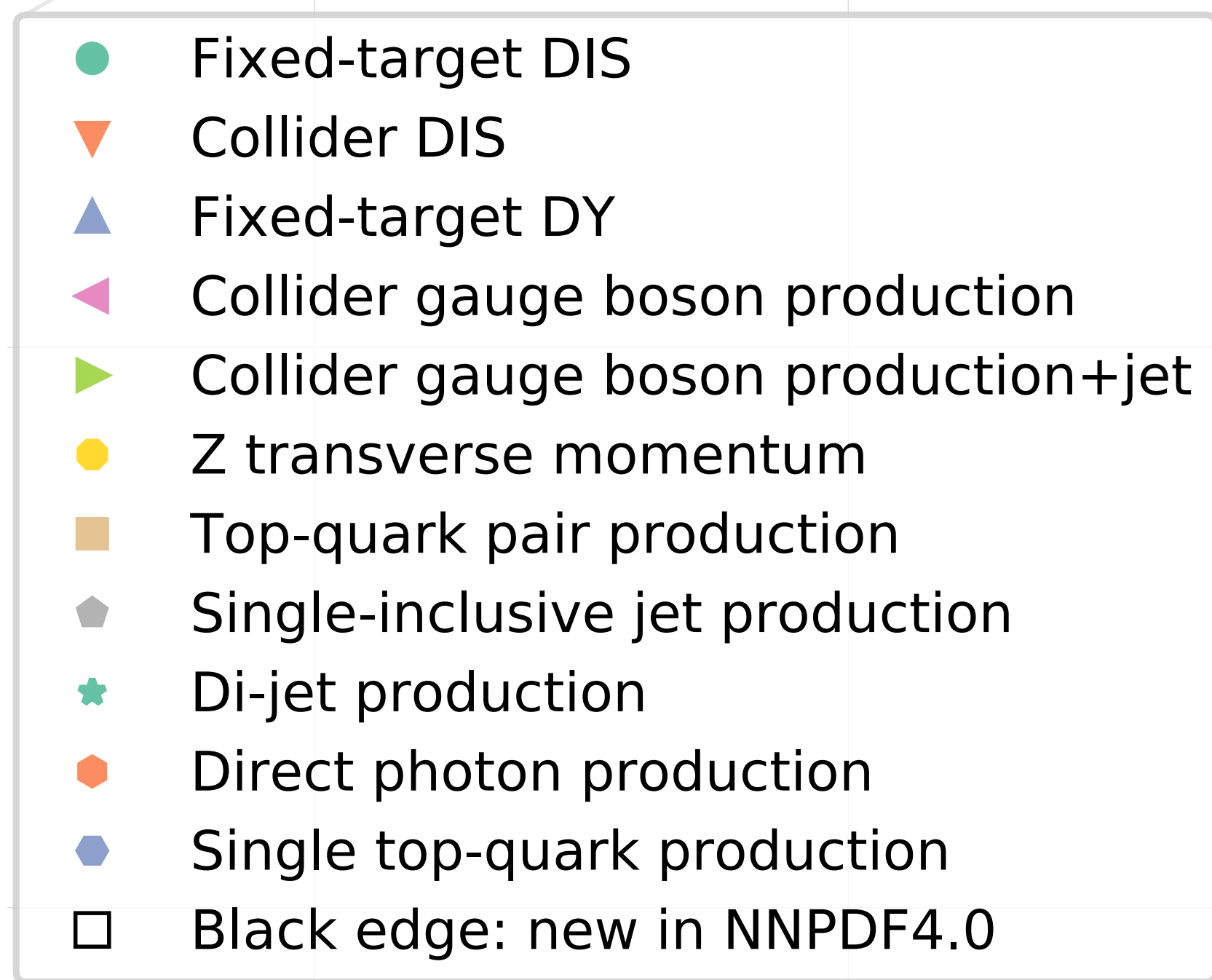


# Data overlap

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*Data included in NNPDF4.0, [2109.02653]:*

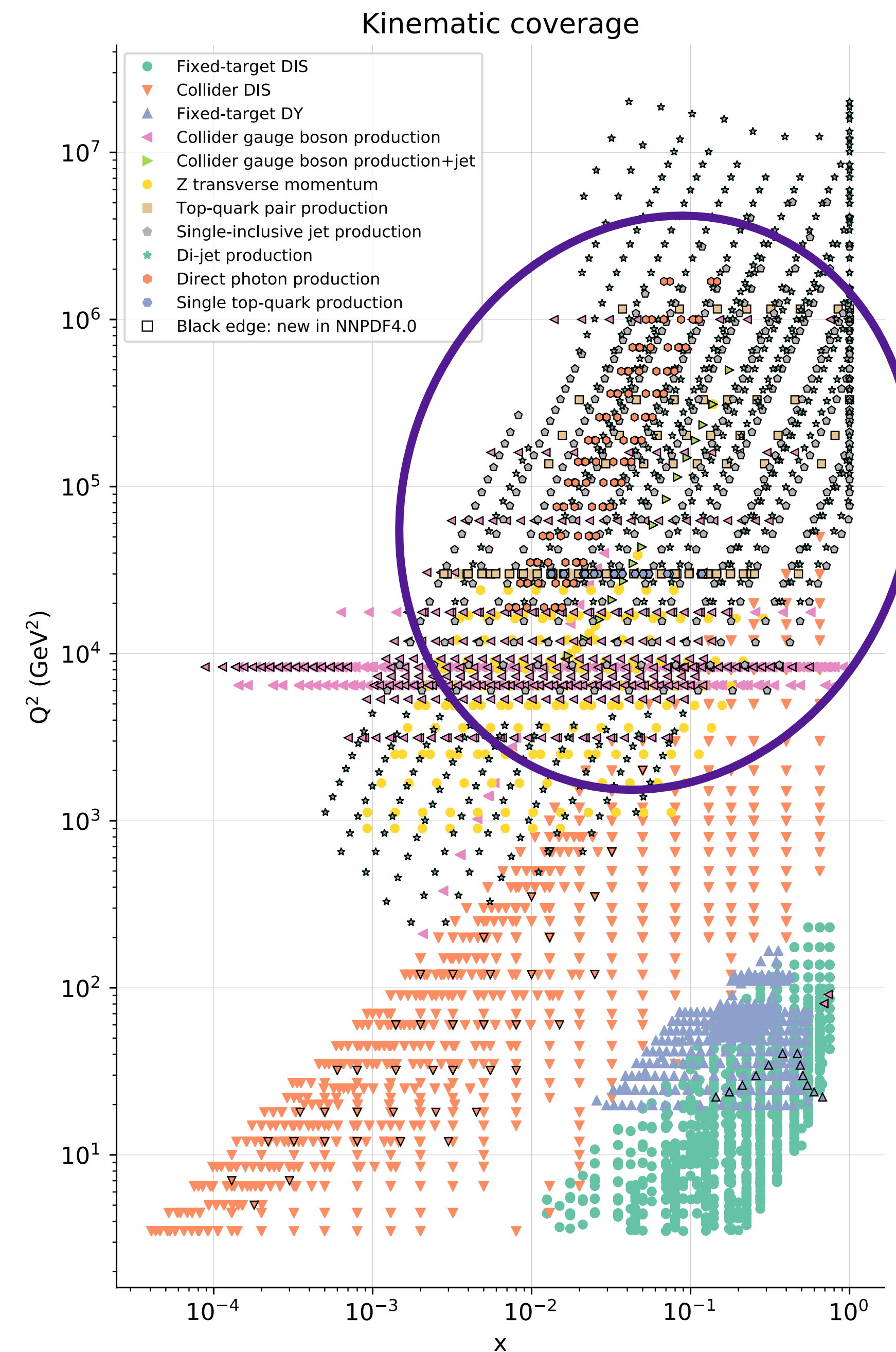
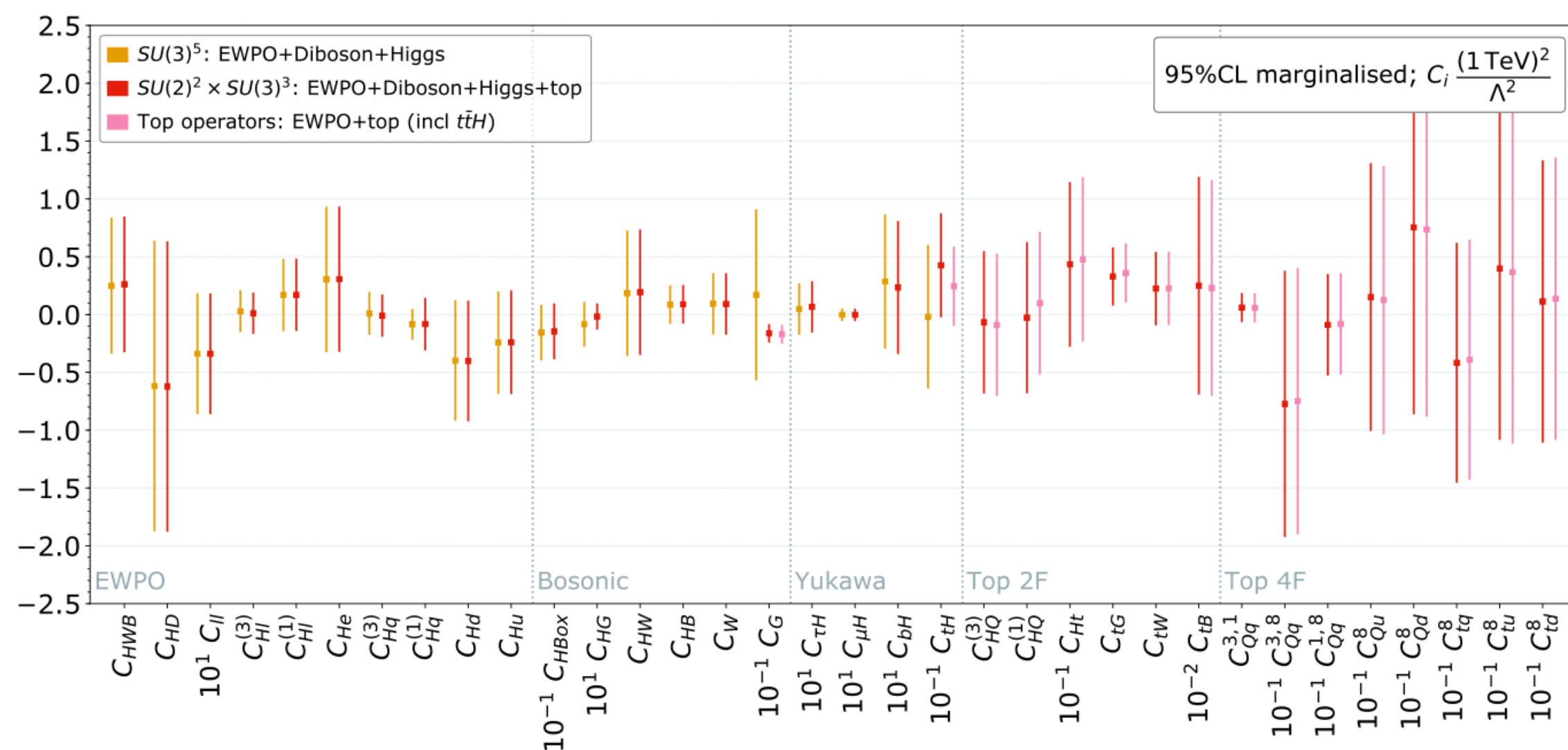


# Data overlap

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This overlap will grow as we take the global approach to constraining the SMEFT.

- ▶ e.g. Top quark data used to fit the SMEFT in the global fit of [2012.02779](#), *J. Ellis, MM, K. Mimasu, V. Sanz, T. You*

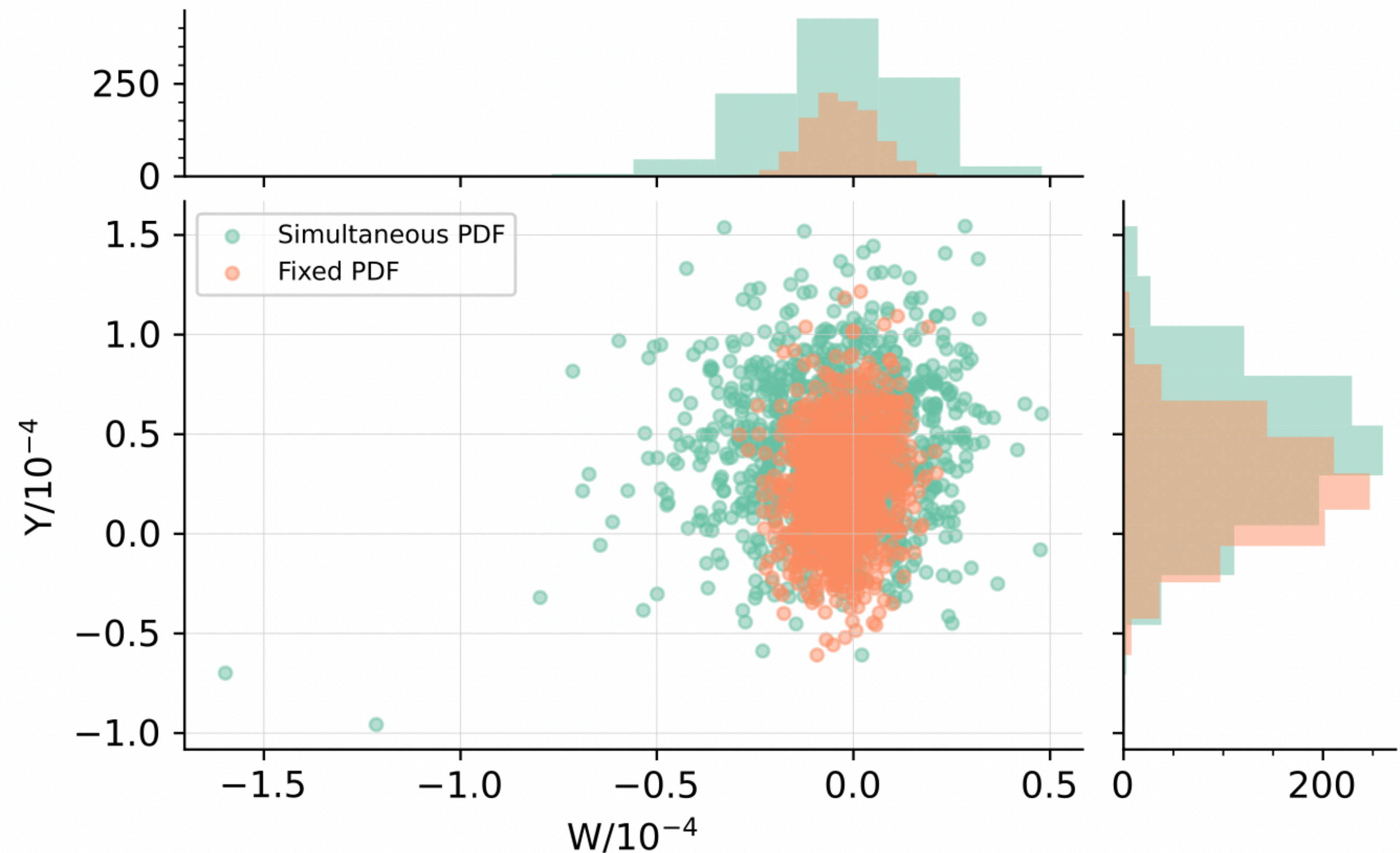


# Simultaneous PDF and SMEFT determinations

*S. Iranipour, M. Ubiali, 2201.07240*

## High-mass Drell-Yan

Neglecting PDF-EFT interplay at the **HL-LHC** leads to a significant overestimate of EFT constraints



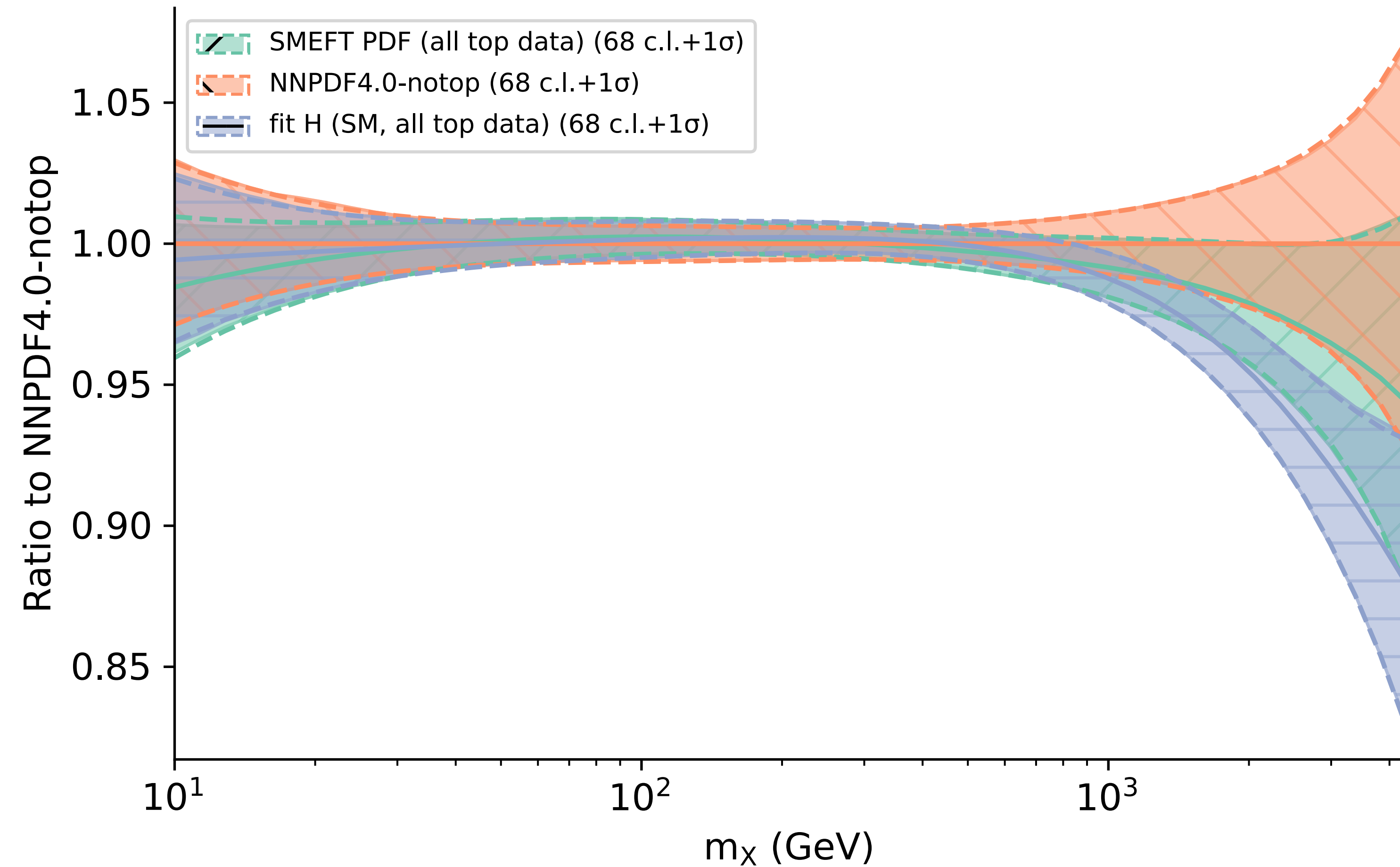


# Simultaneous PDF and SMEFT determinations

*Kassabov et. al: 2303.06159*

gg luminosity  
 $\sqrt{s} = 13$  TeV

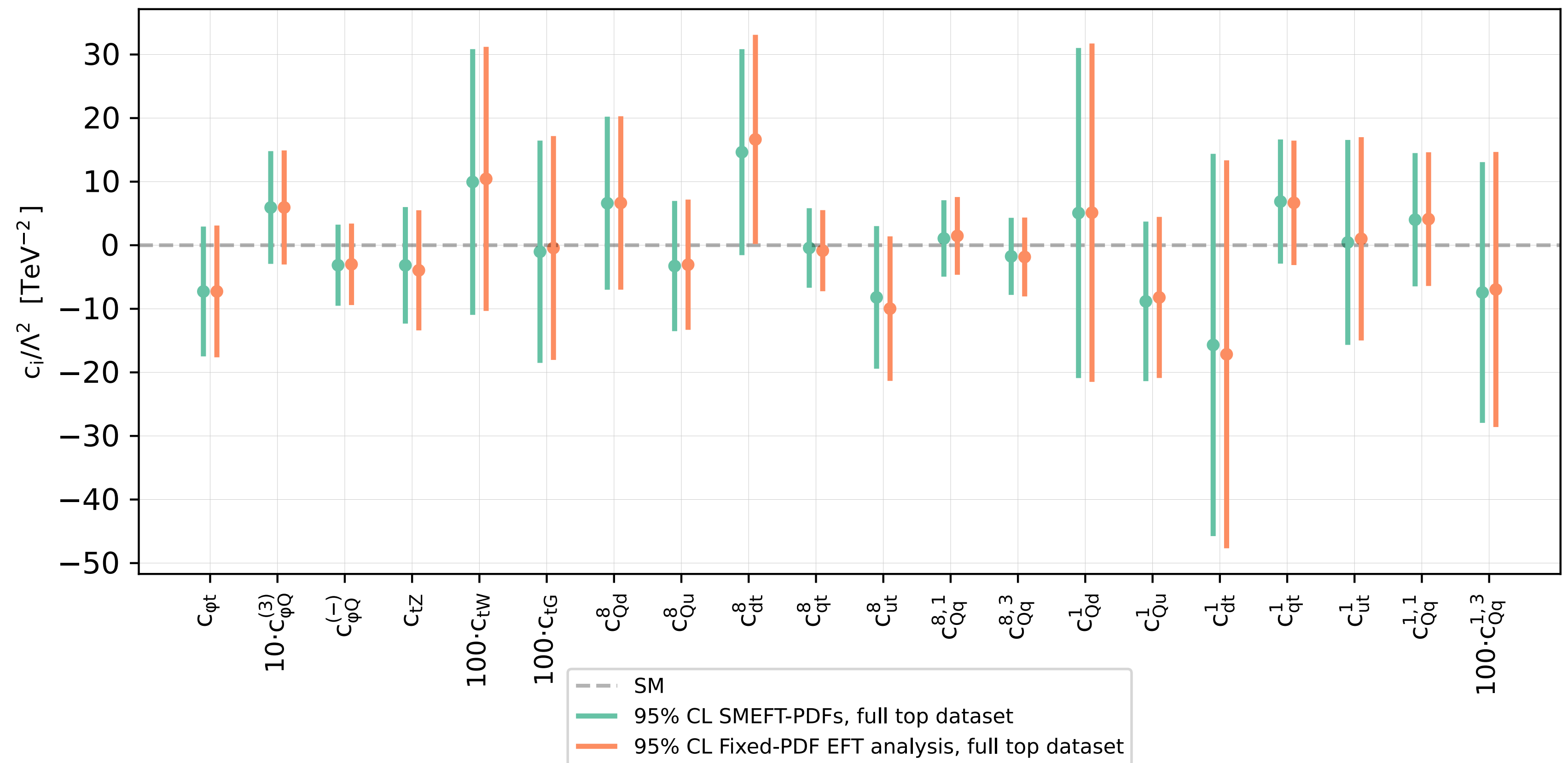
## Top quark data



# Simultaneous PDF and SMEFT determinations

*Kassabov et. al: 2303.06159*

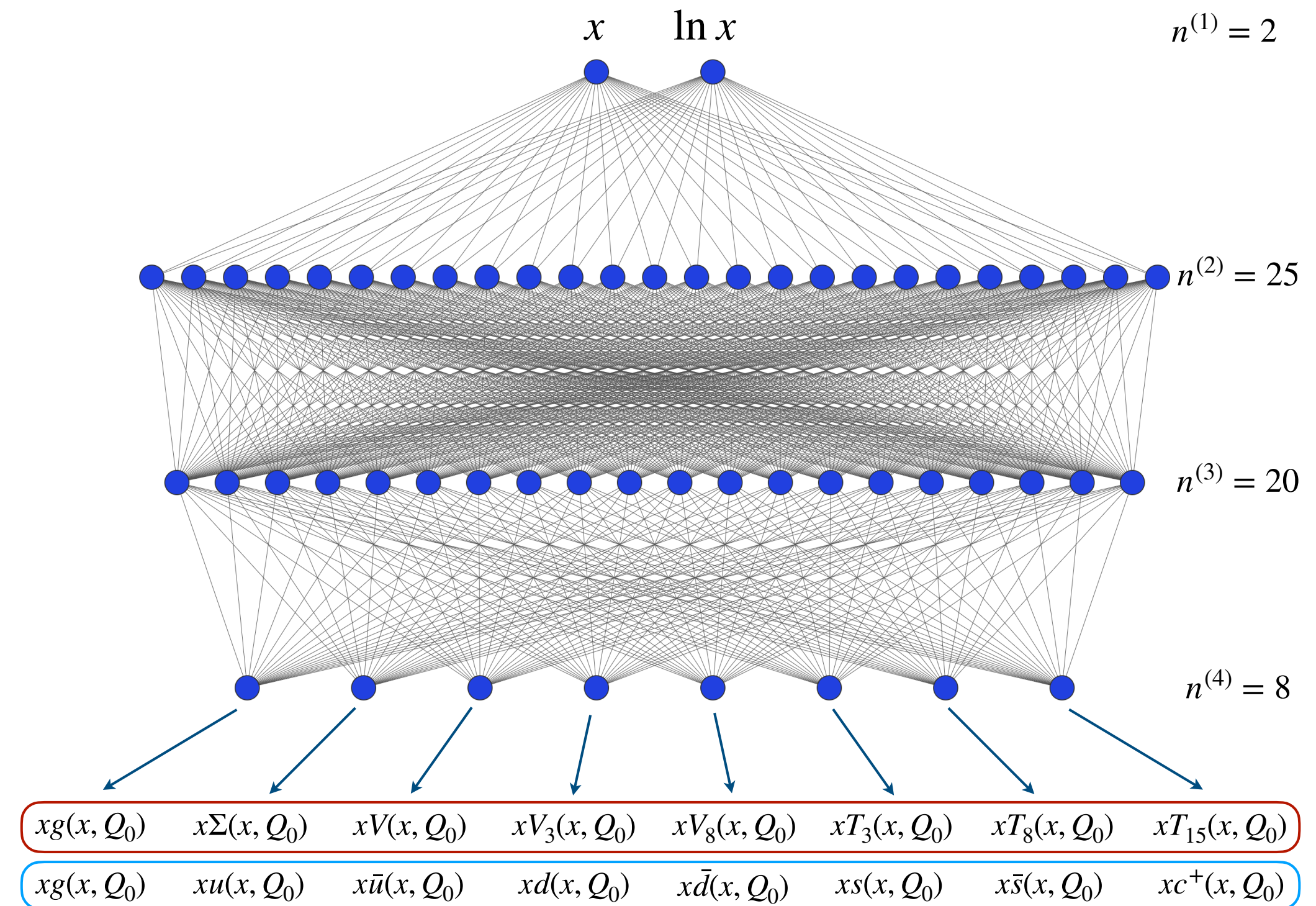
## Top quark data



# SIMUnet methodology

## An extension of the NNPDF framework

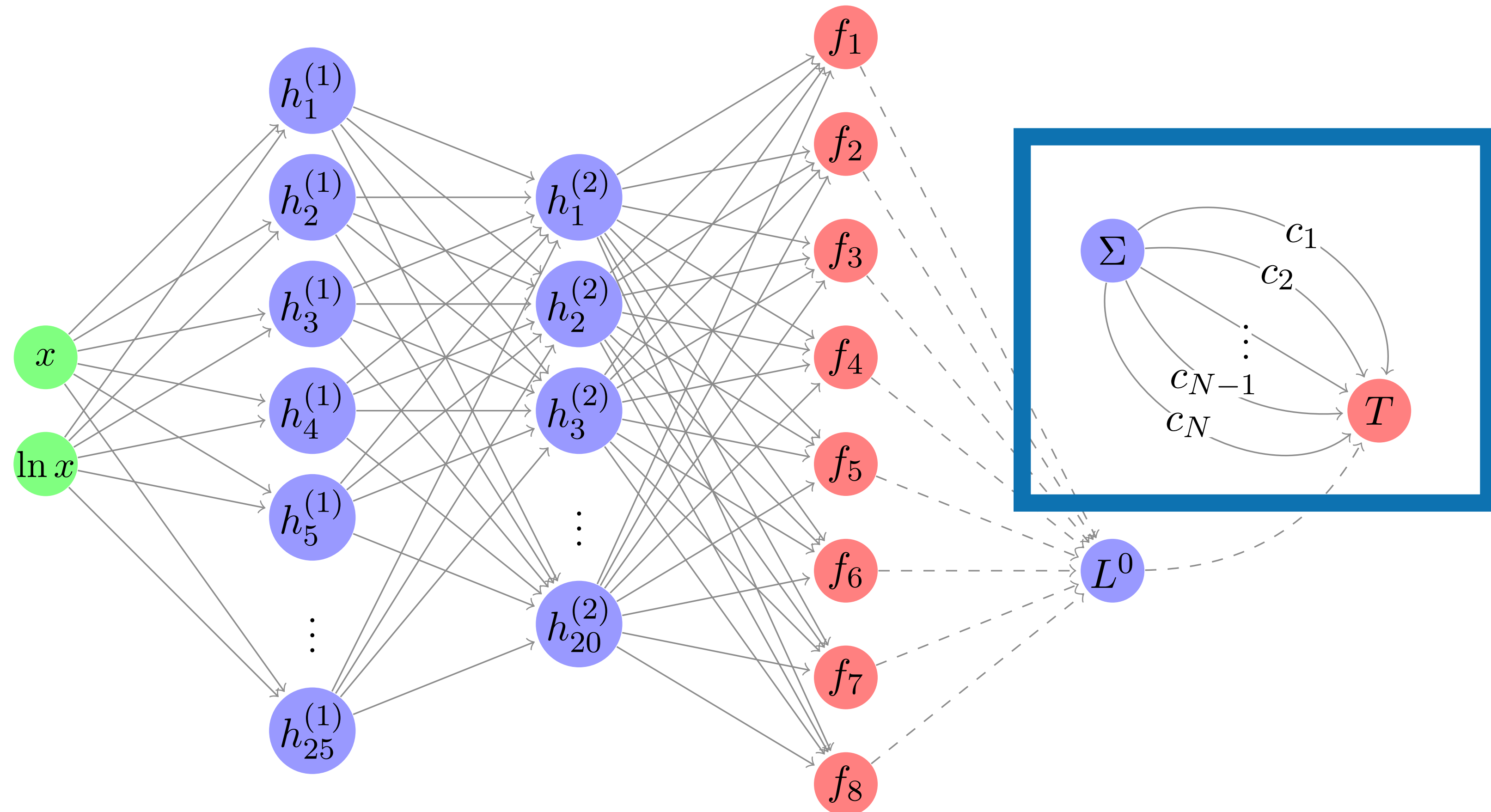
- PDFs parameterised by a neural network



*Ball et. al, NNPDF4.0, 2109.02653*

# SIMUnet methodology

Additional layer accounts for dependence of partonic cross section on Wilson coefficients via k-factor approximation

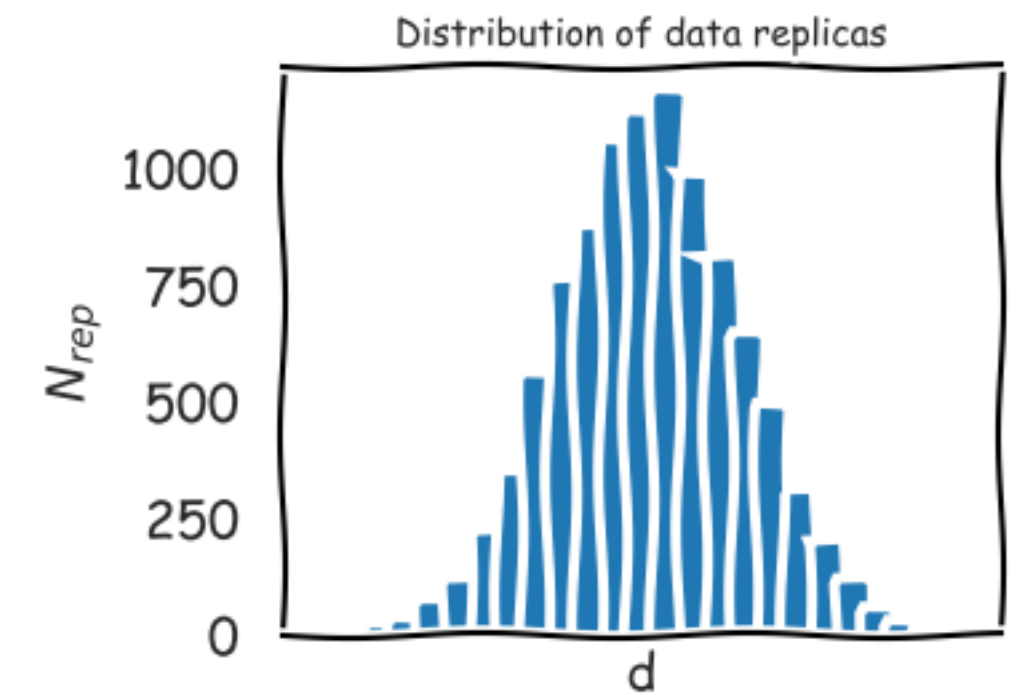


# SIMUnet methodology

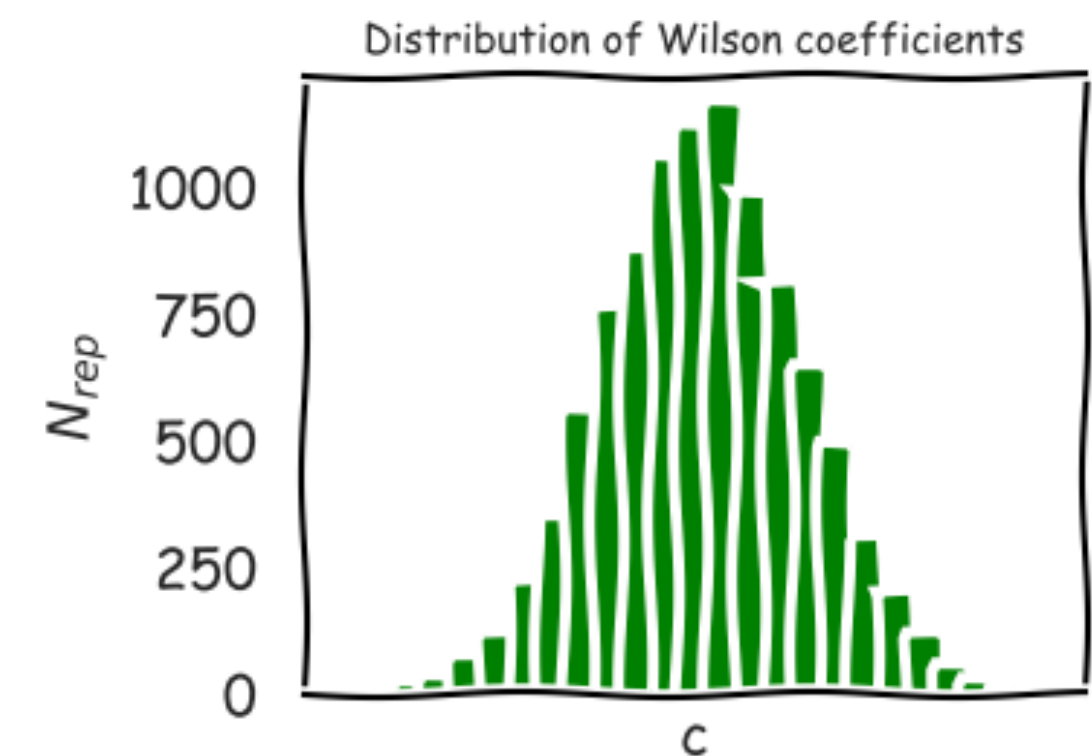
## An extension of the NNPDF framework

- PDFs parameterised by a neural network
- Propagates uncertainties from data to NN parameters using the Monte Carlo replica method

$$d_k \sim \mathcal{N}(d, \sigma)$$

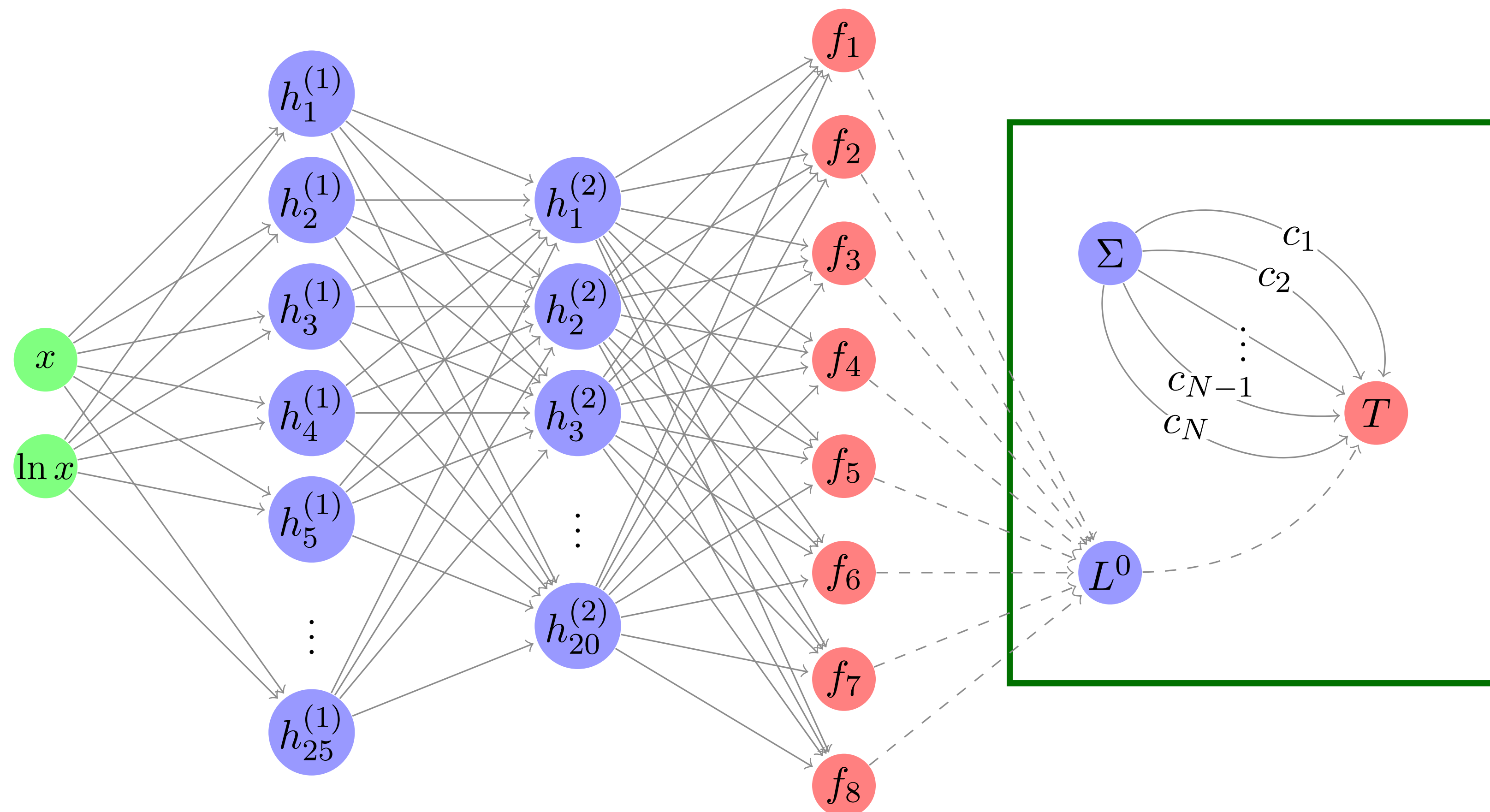


$$c_k = \arg \min_c \chi^2(c, d_k)$$



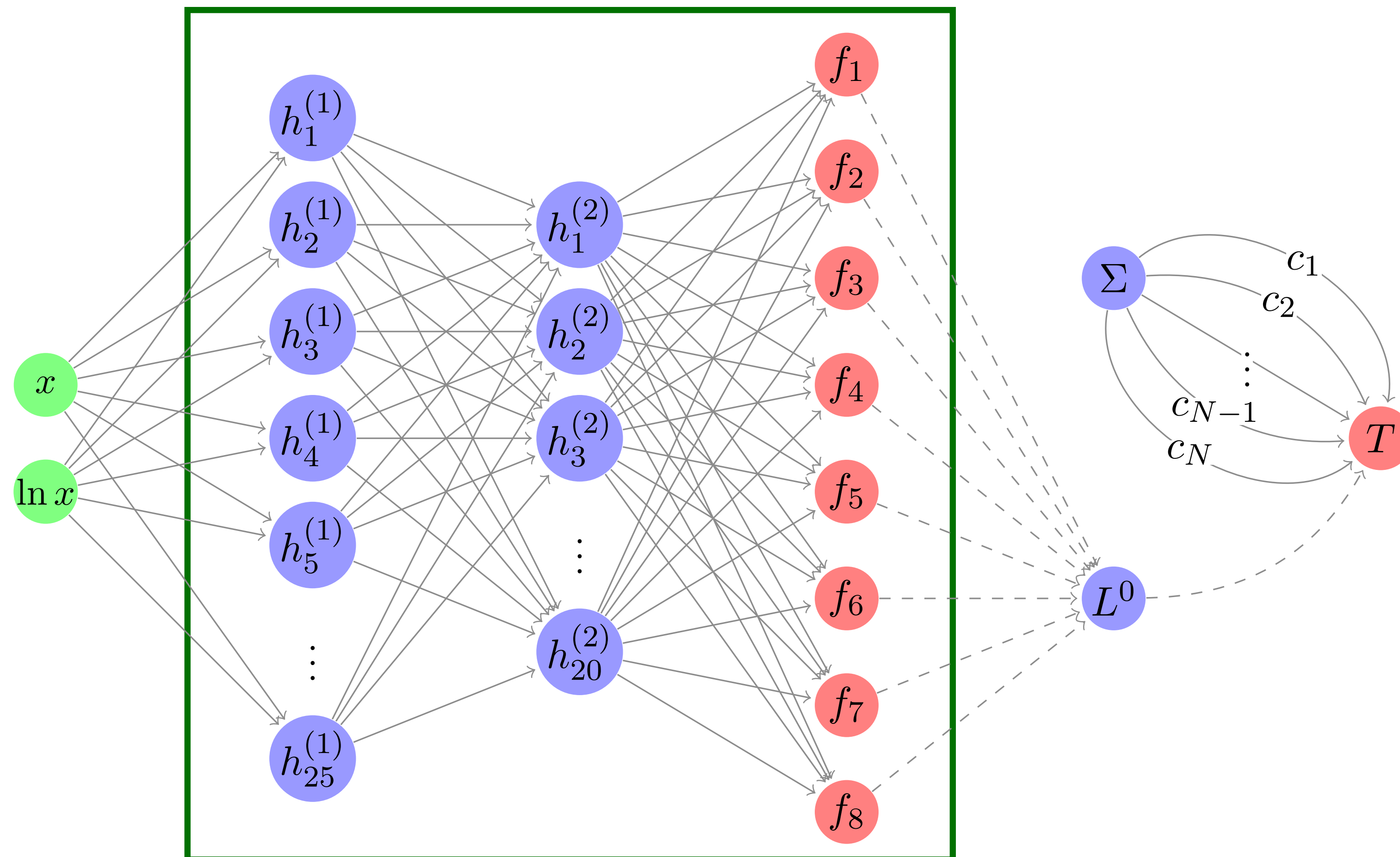
# SIMUnet fits: SMEFT

Train only the final layer: reproduce **SMEFT-only fits**



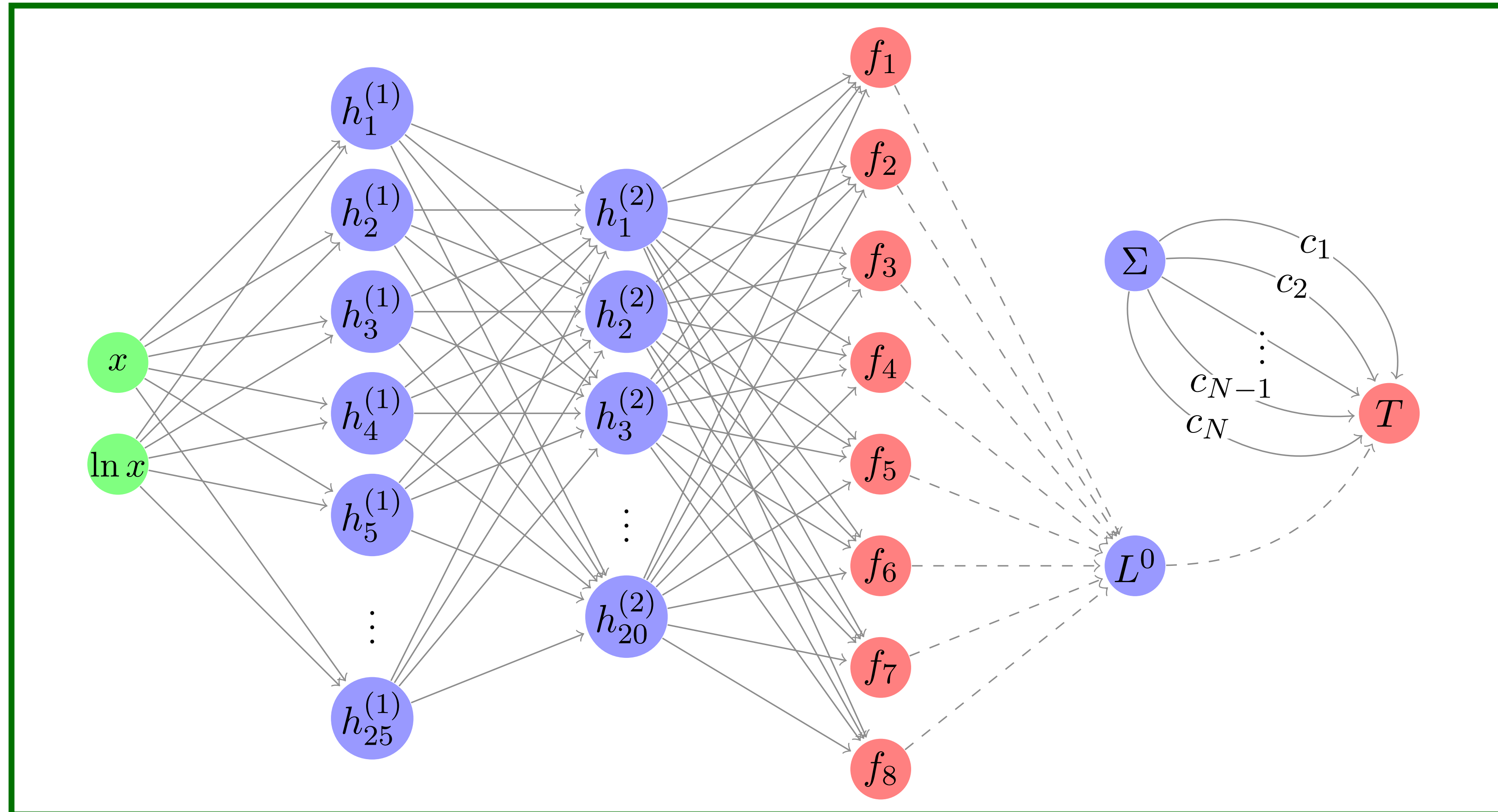
# SIMUnet fits: PDFs

Train only the PDF NN weights on all data: reproduce **PDF-only fits**



# SIMUnet fits: SMEFT and PDFs

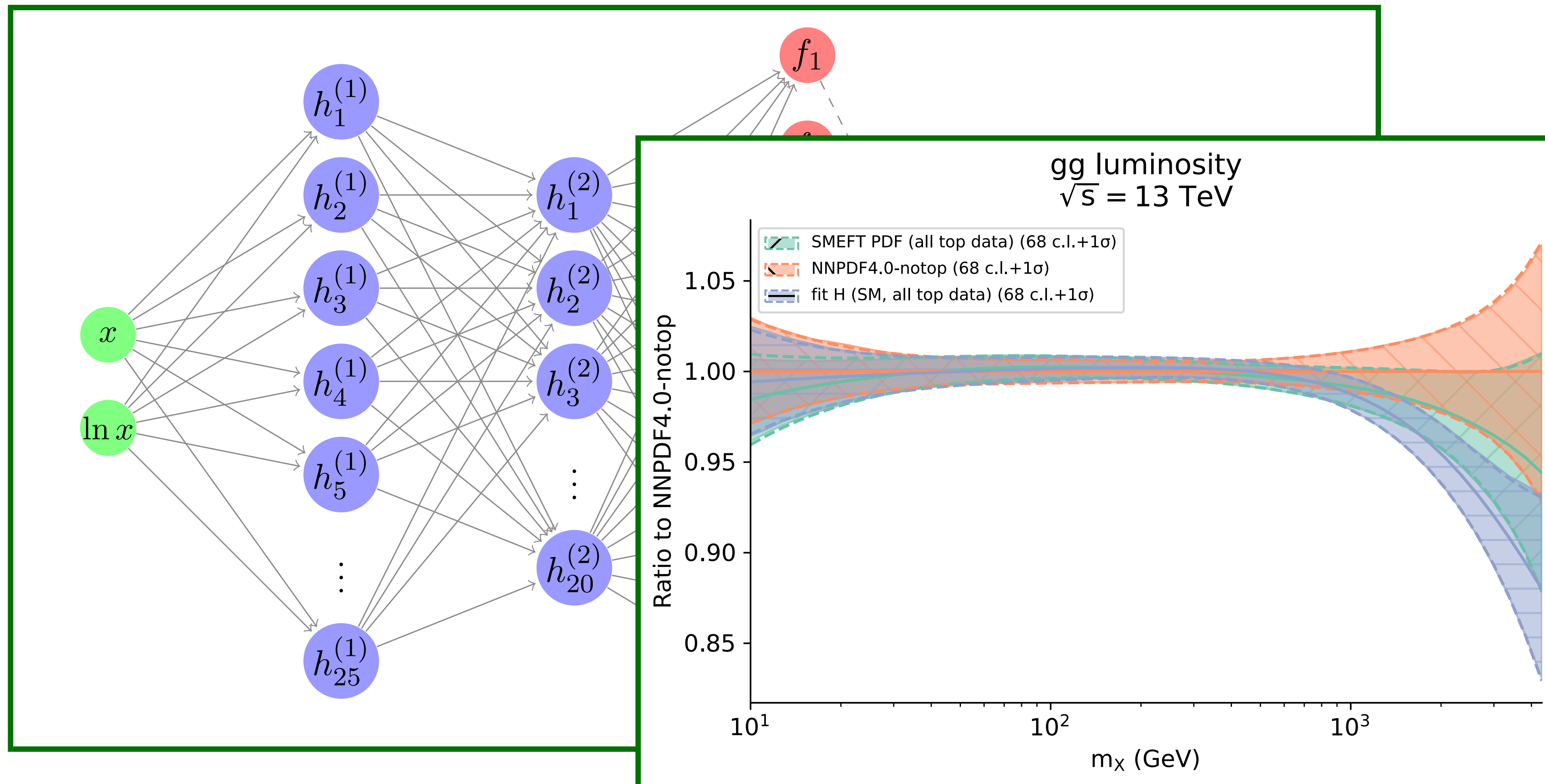
Train everything: **simultaneous fit**





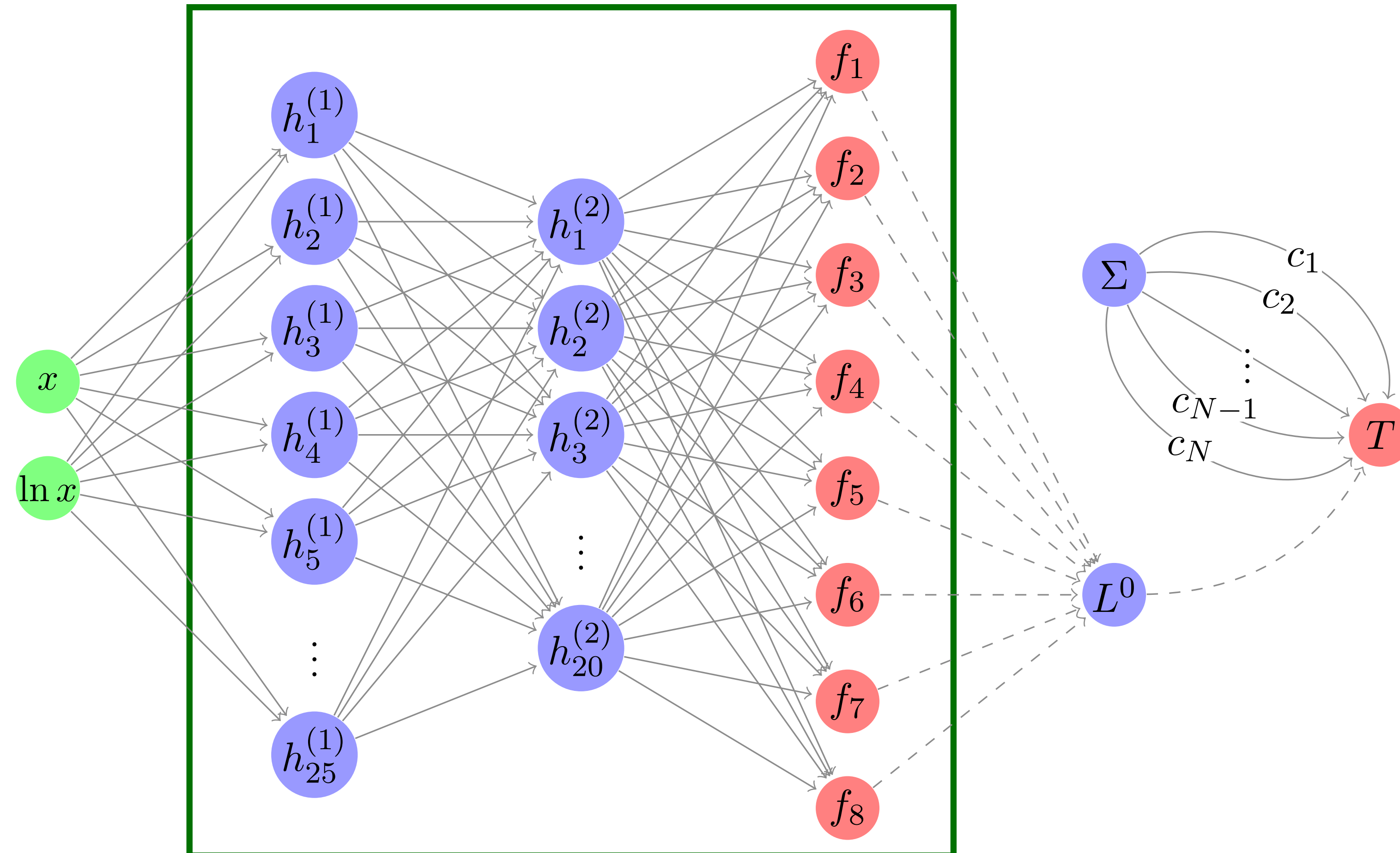
# SIMUnet fits: SMEFT and PDFs

Train everything: **simultaneous fit**



# SIMUnet fits: new physics contamination

Fit only the PDF to **pseudodata modified by new physics effects** and assess the fit quality: **is new physics absorbed?**

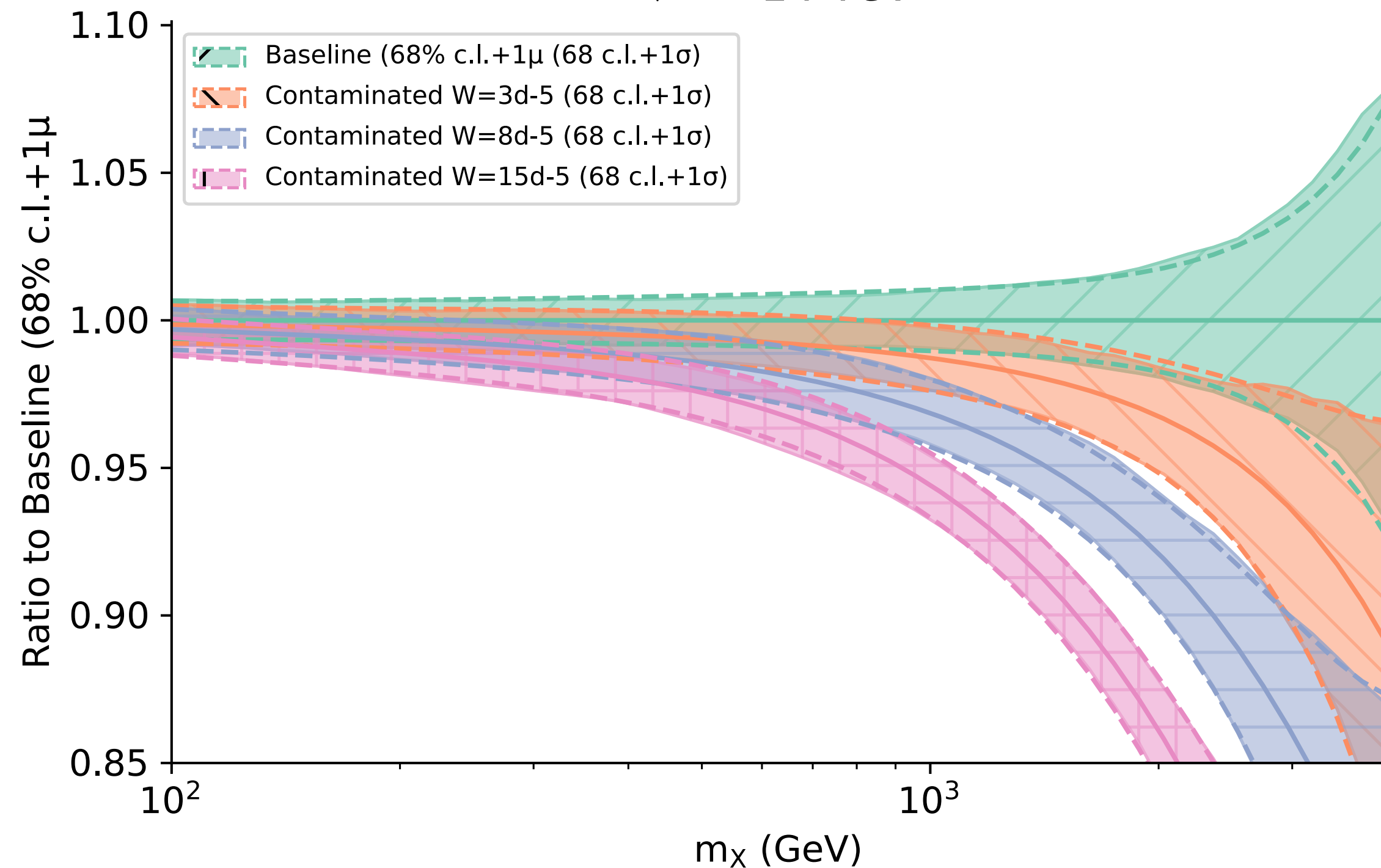


# SIMUnet fits: new physics contamination

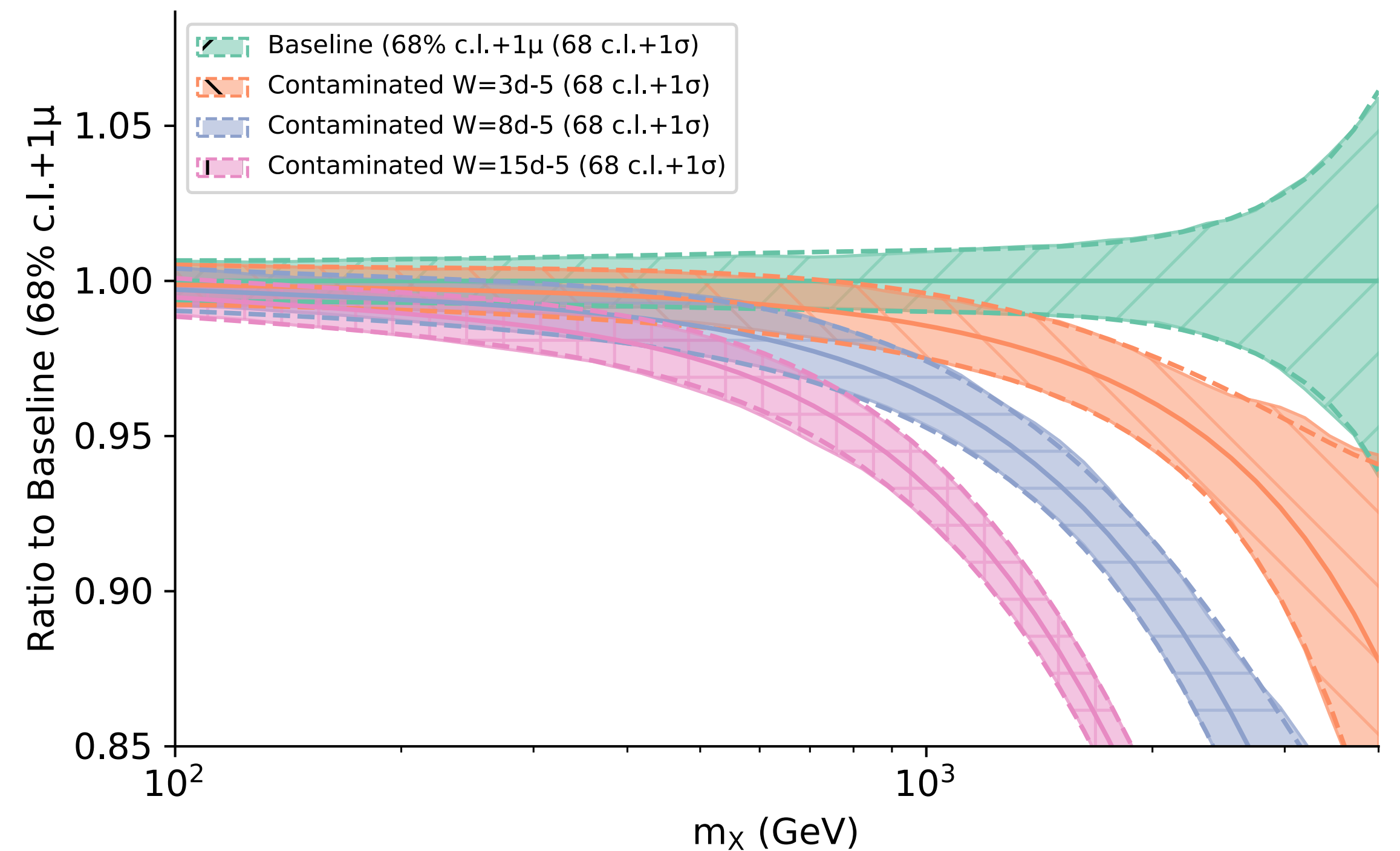
Fit only the PDF to **pseudodata modified by new physics effects** and assess the fit quality: **is new physics absorbed?**

e.g. HL-LHC high mass DY, *E. Hammou et. al 2307.10370*

$u\bar{u} + d\bar{d}$  luminosity  
 $\sqrt{s} = 14$  TeV



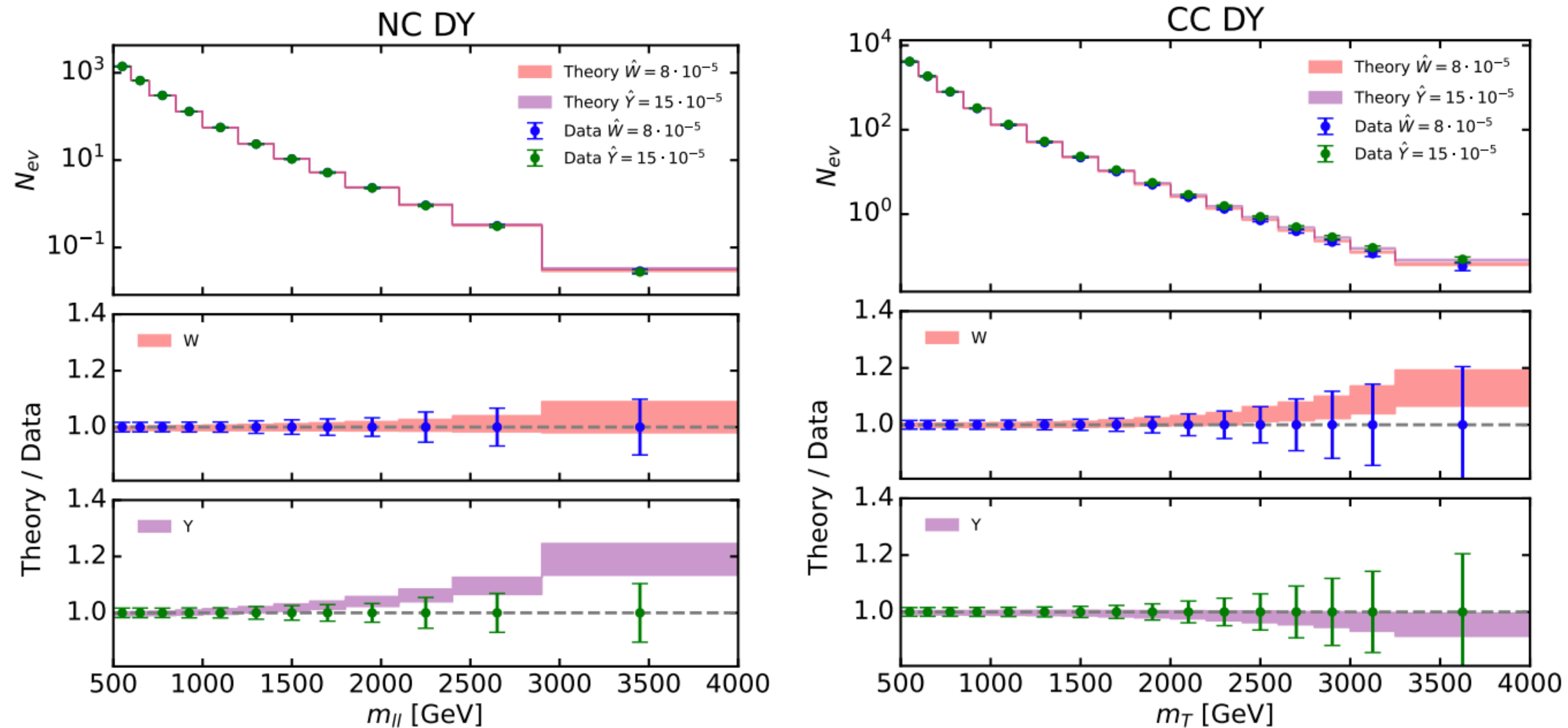
$u\bar{d} + d\bar{u}$  luminosity  
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# SIMUnet fits: new physics contamination

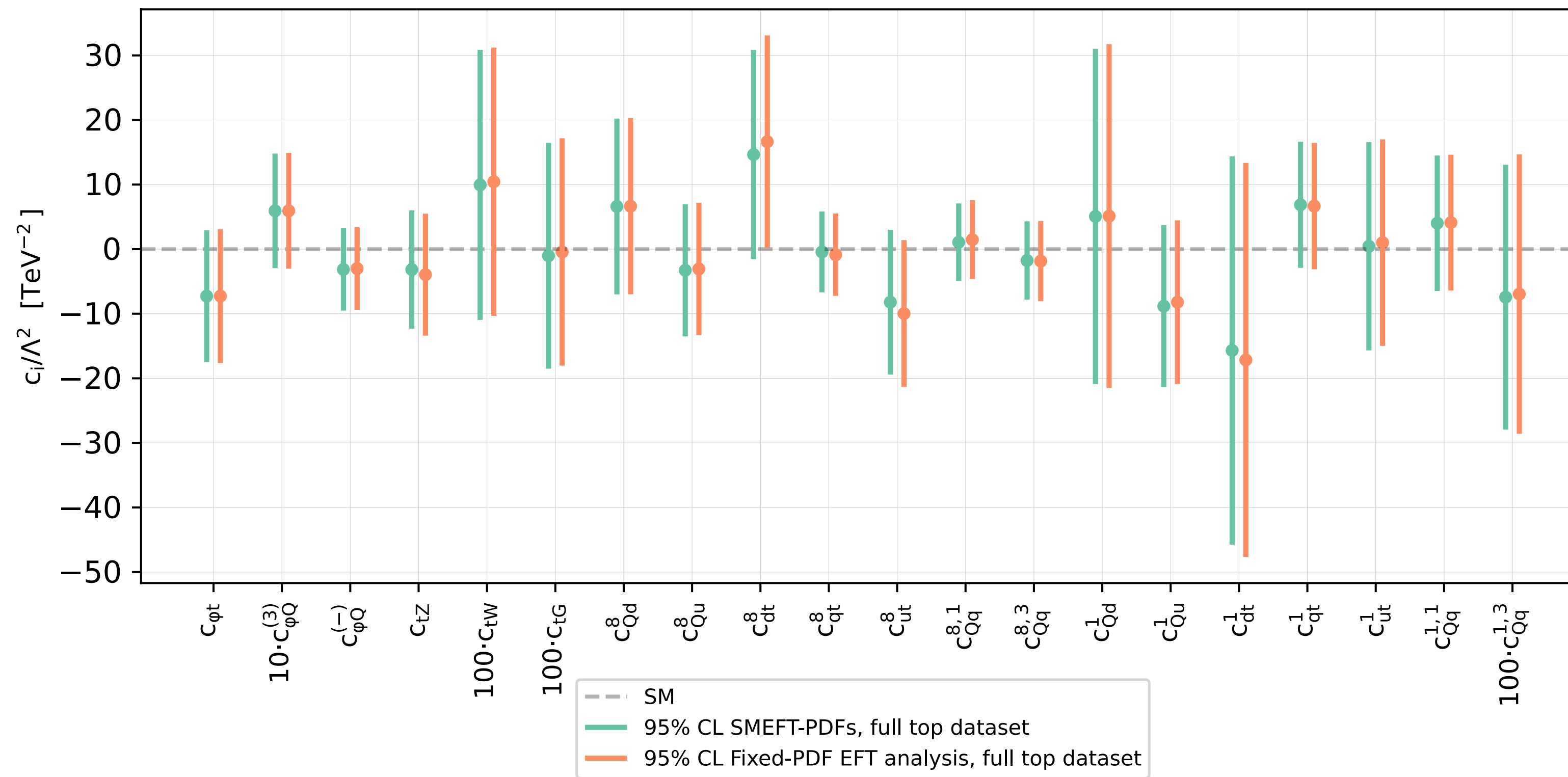
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# The SIMUnet release

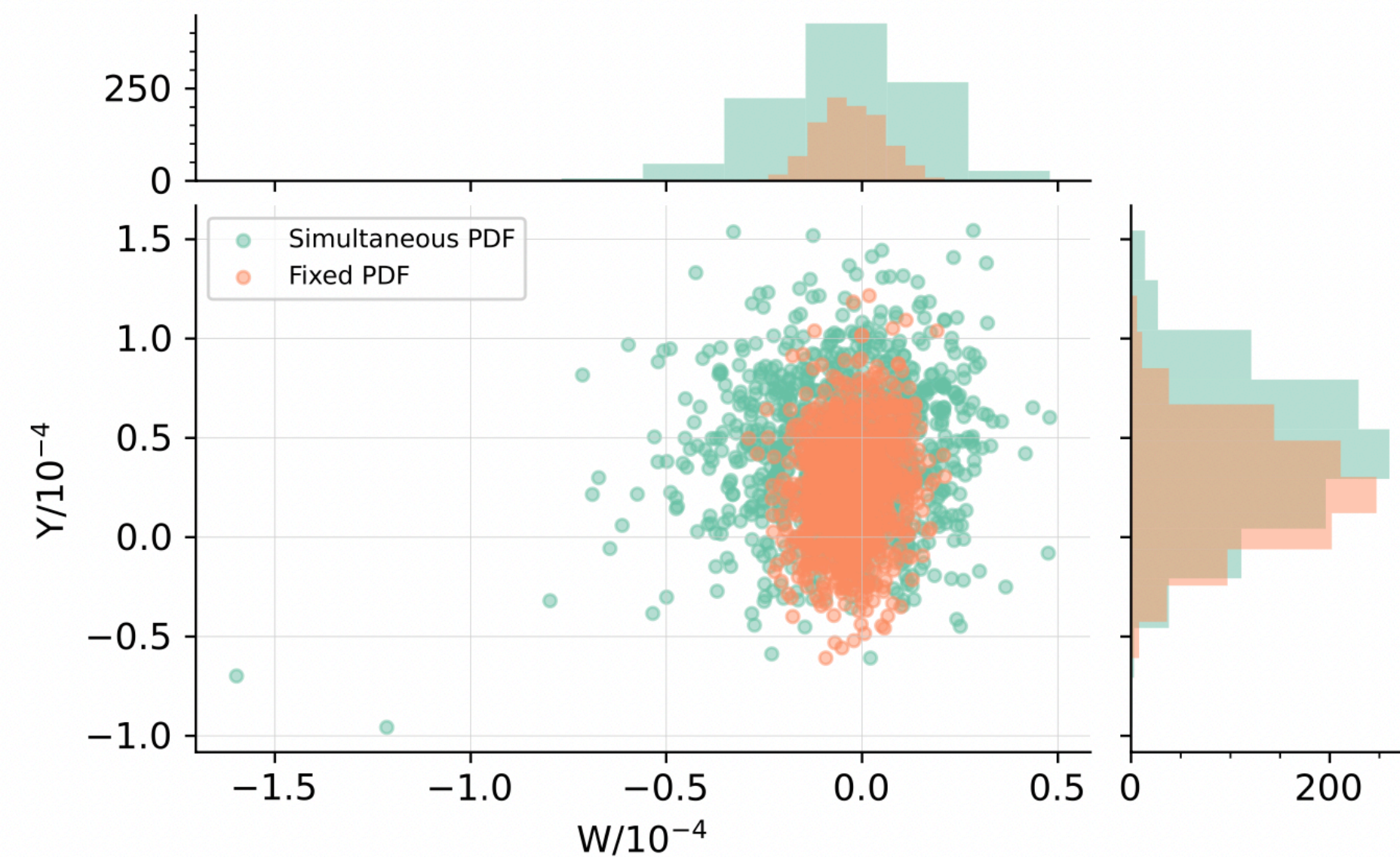
## Simultaneous fits of PDFs and linear SMEFT effects



# The SIMUnet release

- ⊗ Simultaneous fits of PDFs and linear SMEFT effects
- ⊗ + Fits of any linear combinations of Wilson coefficients

e.g. electroweak oblique parameters  $W, Y$



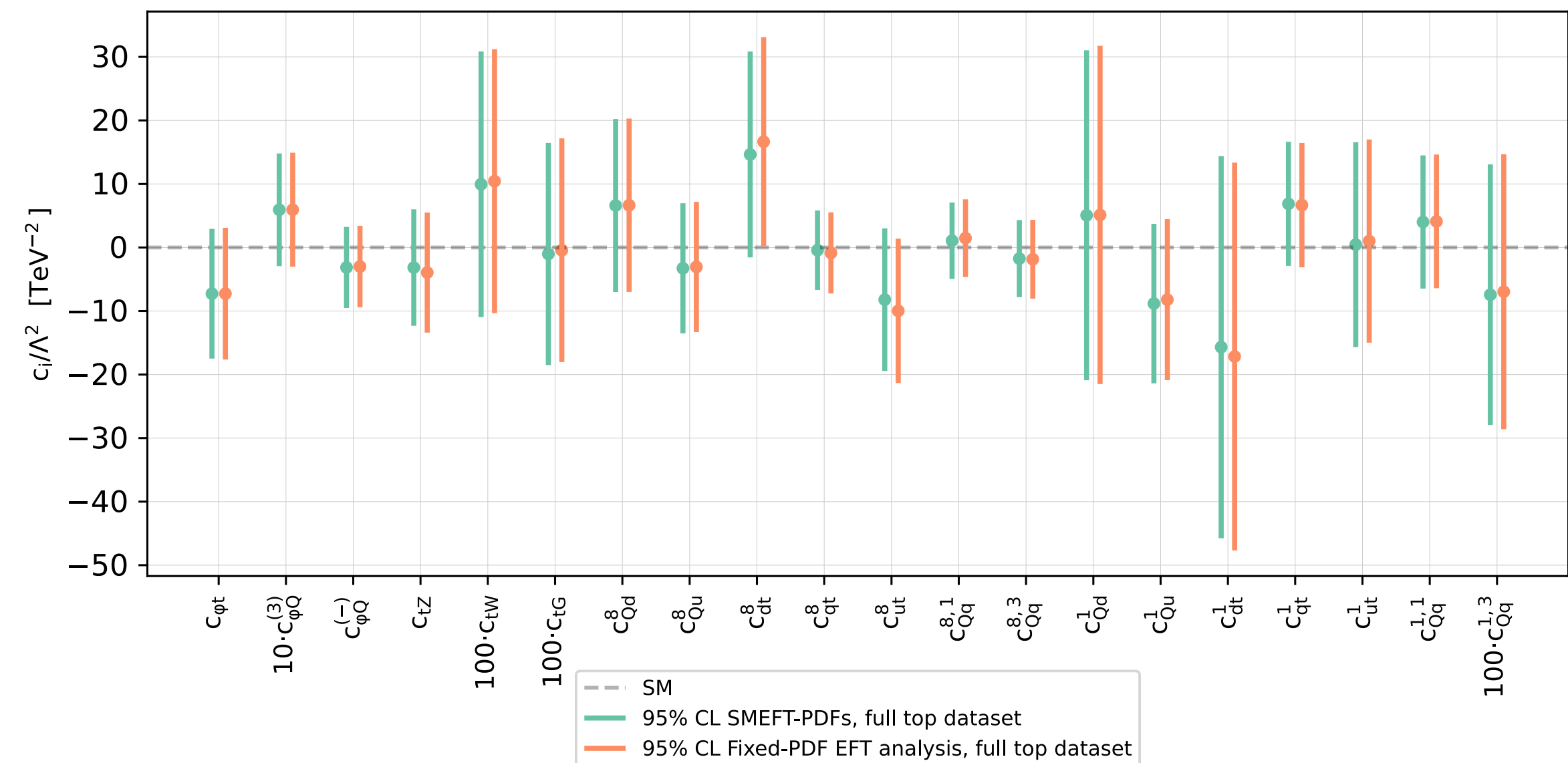
# The SIMUnet release

- Simultaneous fits of PDFs and linear SMEFT effects

- + Fits of linear combinations of Wilson coefficients

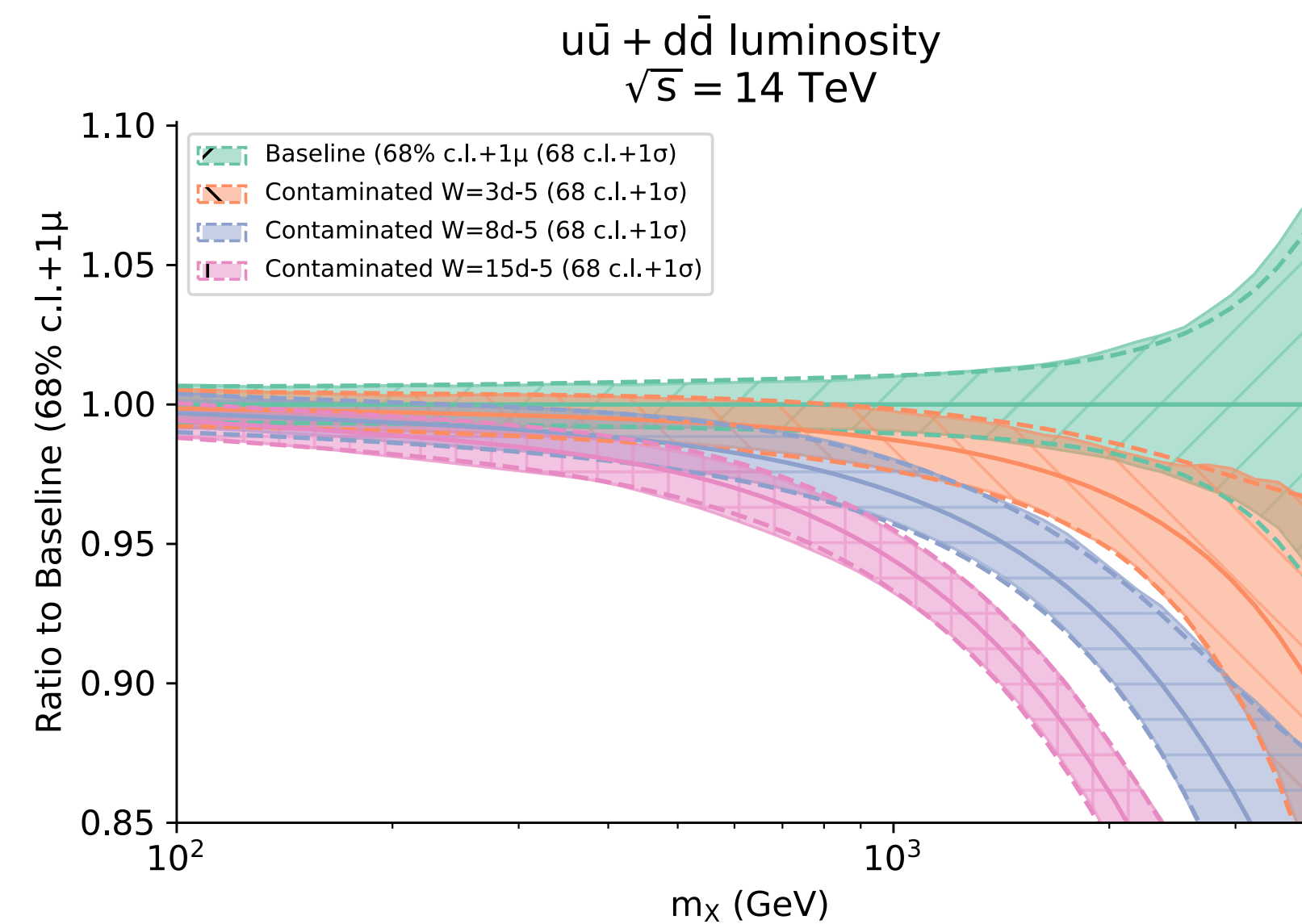
- PDF-independent observables**

E.g. measurements of  $W$  polarisations in top decay, electroweak precision observables



# The SIMUnet release

- Simultaneous fits of PDFs and linear SMEFT effects
- + Fits of linear combinations of Wilson coefficients
- PDF-independent observables
- Tests for new physics absorption**





# The SIMUnet release

⊗ Simultaneous fits of PDFs and linear SMEFT effects

⊗ + Fits of linear combinations of Wilson coefficients

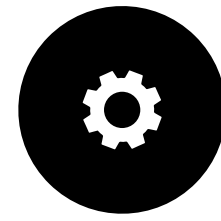
⊗ PDF-independent observables

⊗ Tests for new physics absorption

**+ new data from the Higgs, diboson, electroweak, Drell-Yan and top sectors**

**+ Tutorials, website and documentation**

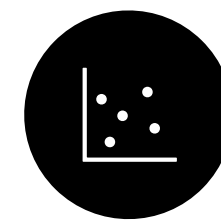
# Conclusions



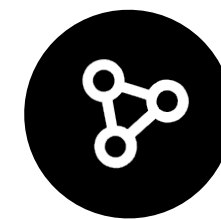
Many examples of the use of ML in EFT analyses

multivariate and unbinned analyses using parametrised classifiers and tree-boosting algorithms; NN-based observables; importance of re-interpretability

<https://indico.cern.ch/event/1331690/>



Unbinned multivariate observables for global SMEFT analyses from parametrised classifiers - optimal SMEFT constraints: **ML4EFT**

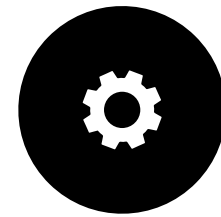


Simultaneous determinations of PDFs and the SMEFT made possible by **SIMUnet**

See also the HEP ML Living Review: <https://iml-wg.github.io/HEPML-LivingReview/>

# Conclusions

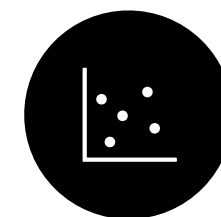
*Thank you for listening!*



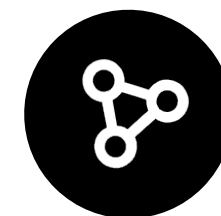
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# Backup

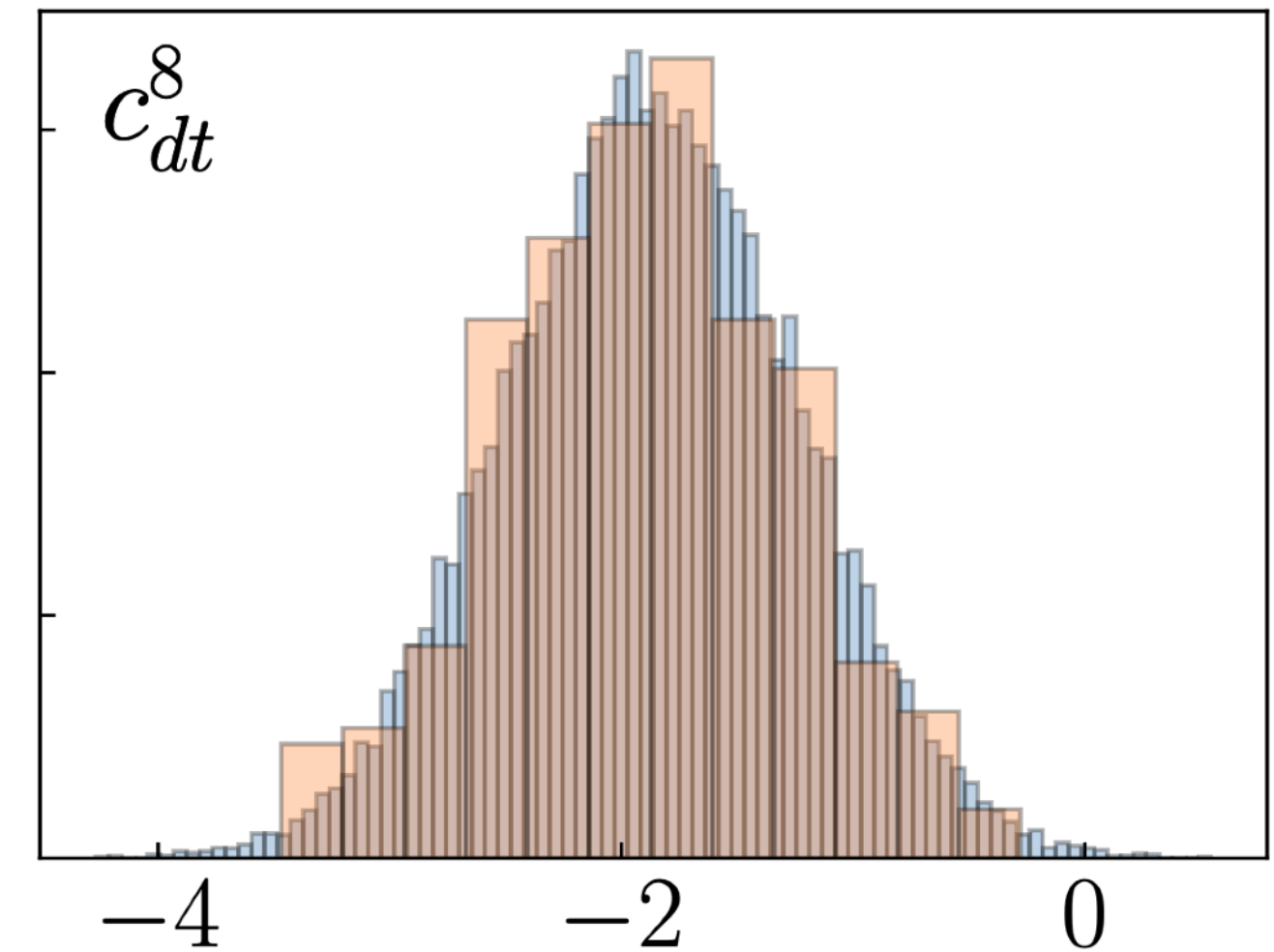
# The Monte Carlo Replica Method

Consider fitting 1 Wilson coefficient  $c$  to 1 datapoint  $\sigma_{\text{exp}}$ : define  $\chi^2 = \frac{(\sigma(c) - \sigma_{\text{exp}})^2}{\delta\sigma^2}$

1. Resample:  $\tilde{\sigma}_{\text{exp}} \sim \mathcal{N}(\sigma_{\text{exp}}, \Sigma)$

2. Minimise:  $\bar{c} = \arg \min_c \frac{(\sigma(c) - \tilde{\sigma}_{\text{exp}})^2}{\delta\sigma^2}$

3. Repeat, and treat the sample  $\{\bar{c}\}$  as a sample from the Bayesian posterior  $p(c|D)$

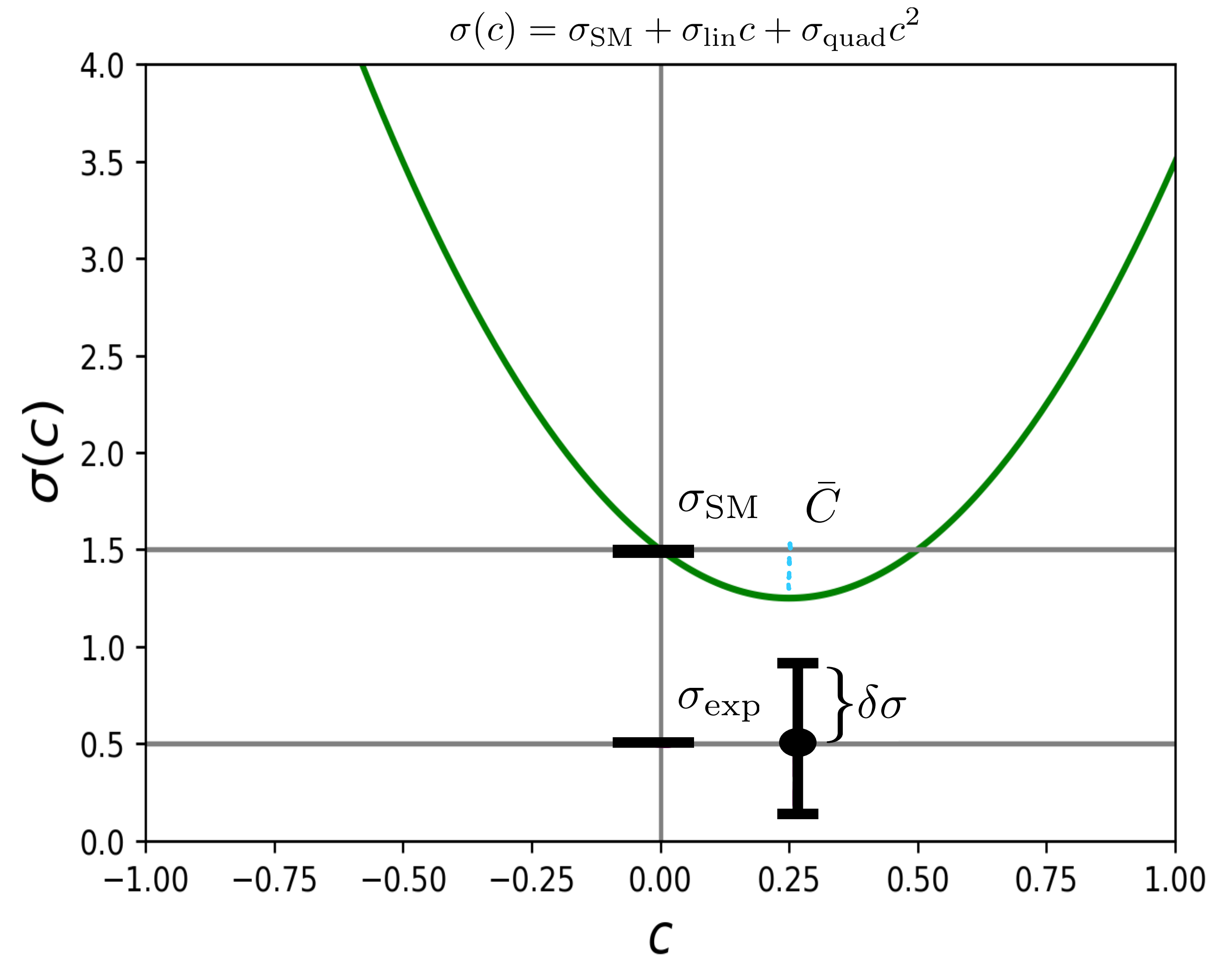


- Often used in the context of PDF fitting and SMEFT fitting, e.g. 2109.02653, 1901.05965

# The Monte Carlo Replica Method

**Problem:** in the presence of a quadratic theory, often the minimum  $\chi^2$  will be given by the same  $\bar{c}$ .

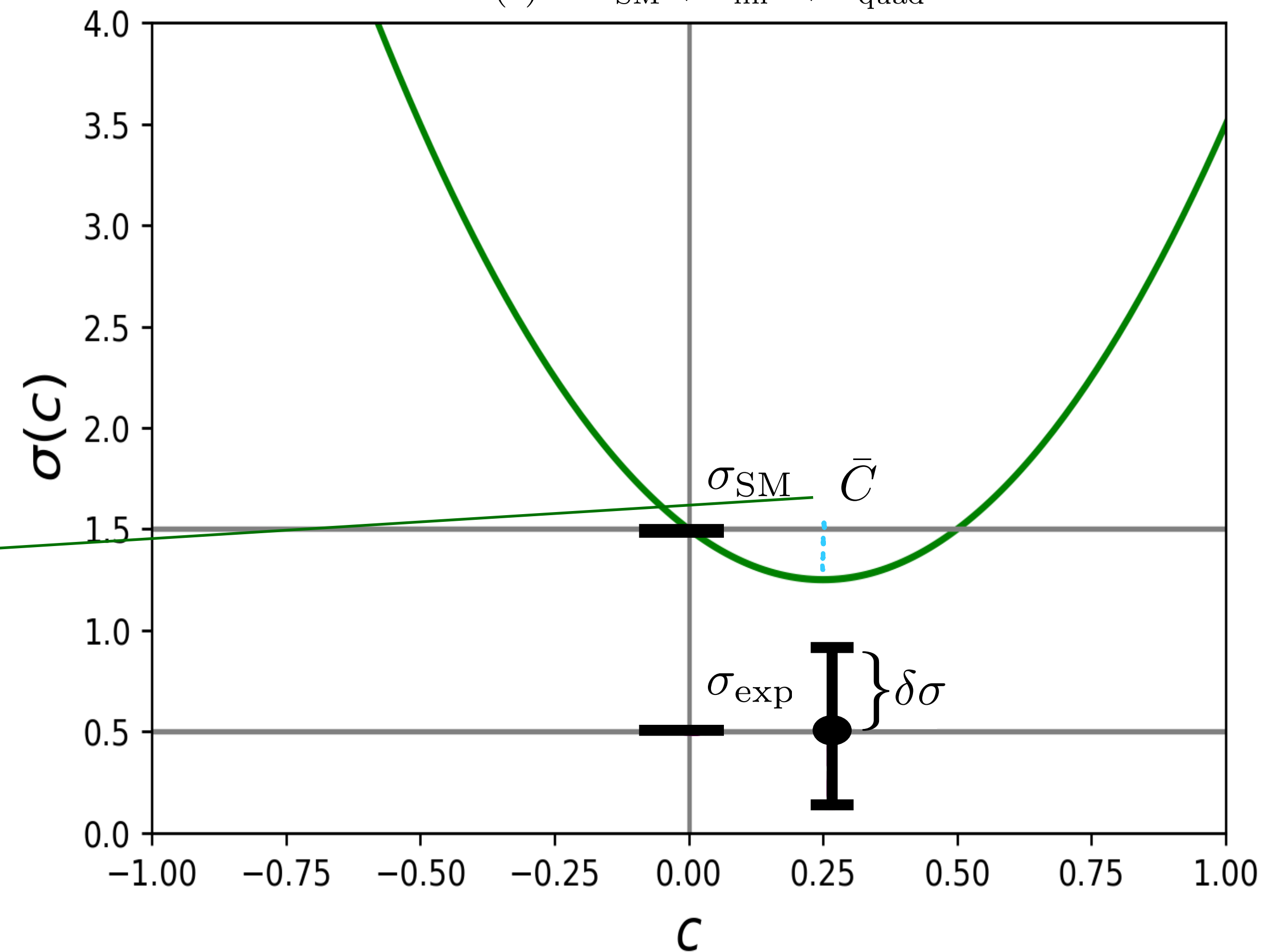
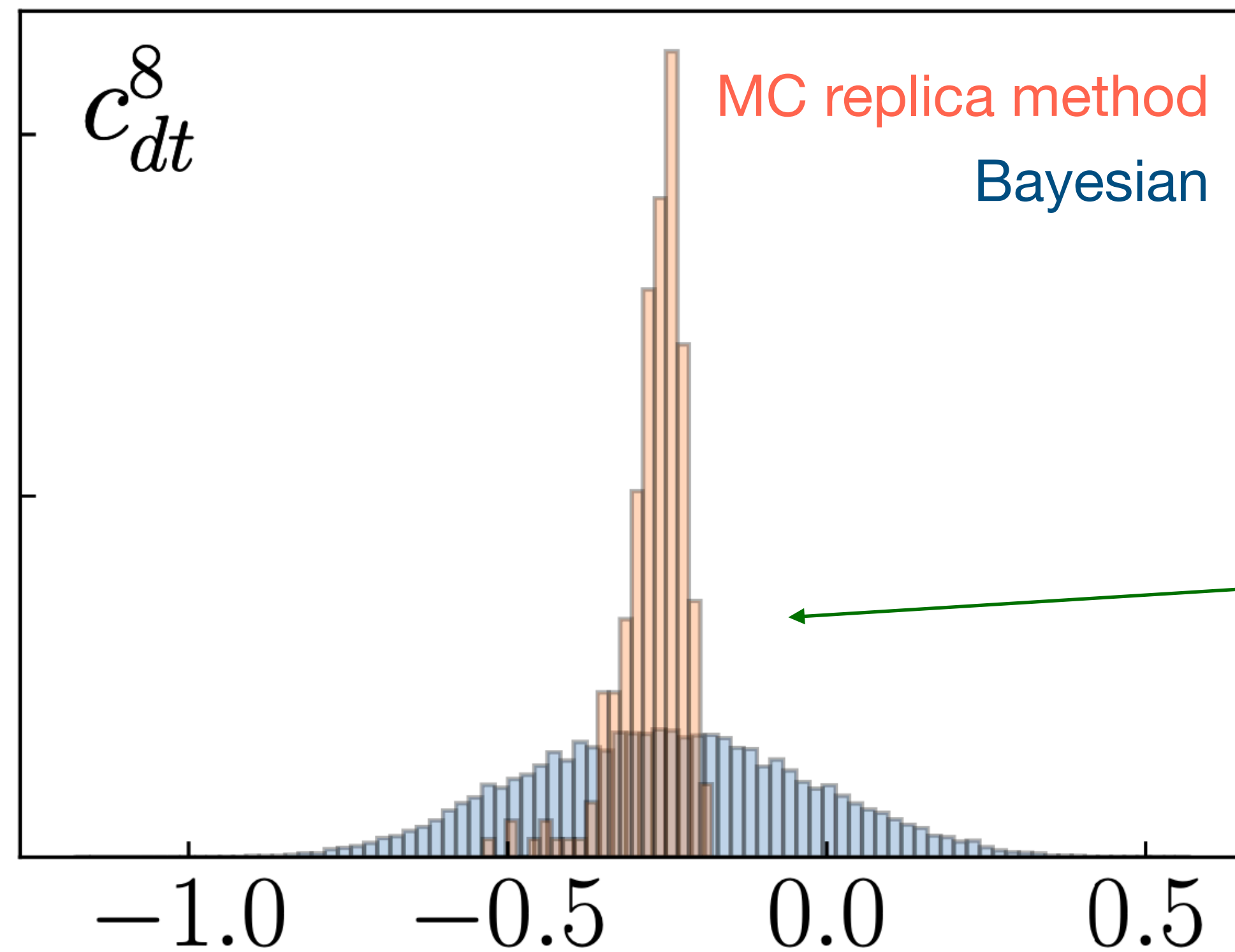
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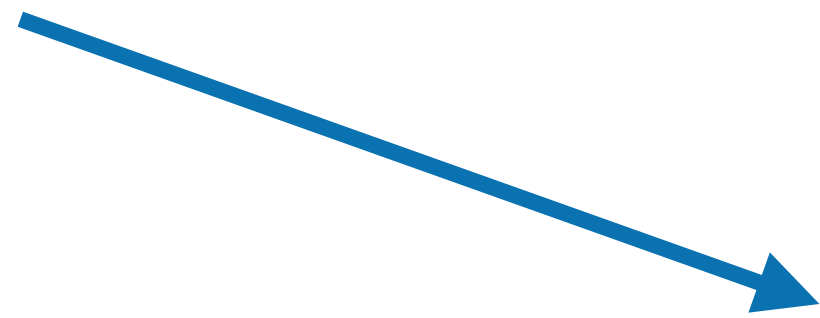
$$\sigma(c) = \sigma_{\text{SM}} + \sigma_{\text{lin}}c + \sigma_{\text{quad}}c^2$$



# ML4EFT

For parameter estimation, we would like to be able to calculate the **likelihood**:  $\mathcal{L}(D|\mathbf{c}) \propto \prod_{i=1}^{N_{ev}} f_{\sigma}(\mathbf{x}_i, \mathbf{c})$

where  $f_{\sigma}(\mathbf{x}, \mathbf{c}) = \frac{1}{\sigma(\mathbf{x}, \mathbf{c})} \frac{d\sigma(\mathbf{x}, \mathbf{c})}{d\mathbf{x}}$        $D = \{\mathbf{x}_i\}$        $\mathbf{x}_i = \{m_{t\bar{t}}, p_T^{\ell_1}, p_T^{\ell_2}, \Delta\eta_{\ell_1, \ell_2}, \Delta\phi_{\ell_1, \ell_2}, \dots\}$



multi-differential cross section in **all features**

**However:** analytical calculation of  $\mathcal{L}$  is intractable in most realistic cases.

**Instead:** approximate  $\mathcal{L}$  using neural networks



# ML4EFT

Train a classifier  $g$  to distinguish the SM from the SMEFT:

$$L[g(\mathbf{x}, \mathbf{c})] = - \sum_{i=1}^{N_{\text{ev}}^{\text{SMEFT}}} \frac{d\sigma(\mathbf{x}_i, \mathbf{c})}{dx} \log(1 - g(\mathbf{x}_i, \mathbf{c})) - \sum_{i=1}^{N_{\text{ev}}^{\text{SM}}} \frac{d\sigma(\mathbf{x}_i, \mathbf{0})}{dx} \log(g(\mathbf{x}_i, \mathbf{0}))$$

*SMEFT training pseudodata sample*

*SM training pseudo data sample*

# ML4EFT

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*SMEFT training pseudodata sample*                      *SM training pseudo data sample*

$$\frac{\delta L}{\delta g} = 0 \Rightarrow g(\mathbf{x}, \mathbf{c}) = \left( 1 + \frac{d\sigma(\mathbf{x}, \mathbf{c})}{dx} / \frac{d\sigma(\mathbf{x}, \mathbf{0})}{dx} \right)^{-1} \equiv \frac{1}{1 + r_{\sigma}(\mathbf{x}, \mathbf{c})}$$

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In the limit of infinite training samples, the decision boundary is 1:1 with the likelihood