Accelerating the search for mass bumps using the Data-Directed Paradigm

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Introduction

Introduction

Motivation

What? \rightarrow Train a neural network to **identify mass bumps in real data** without the need of simulation or analytical fit to estimate the background

- Rapidly scan many different regions of the observable-space
- O Complementary to the standard analysis approach
- Why? \rightarrow Exploit the discovery potential of the data
 - O Impossible to cover all possible searches with the traditional analysis
 - Many possible resonances in unexplored final states

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			-	ala	h.	1	~	Z/W = H		$BSM \rightarrow SM_1 \times SM_1$		$BSM \rightarrow SM_1 \times SM_2$				$BSM \rightarrow complex$				
	-	"	· ·	47.9	0		· · ·	10,10		q/g	$\gamma/\pi^{0}s$	b		tZ/H	bH		$\tau = \tau q q$	eqq'	$\mu q q'$	
e	[37, 38]	[39, 40]	[39]	ø	ø	ø	[41]	[42]	ø	ø	ø	ø		ø	ø	ø	ø	[43, 44]	ø	
μ		[37, 38]	[39]	ø	ø	ø	[41]	[42]	ø	ø	ø	ø		ø	ø	ø	ø	ø	43,44	1
τ			[45, 46]	ø	[47]	ø	ø	ø	ø	ø	ø	ø		ø	ø	ø	[48,	[9] Ø	ø	
q/g				[29, 30, 50, 51]	[52]	ø	[53, 54]	[55]	ø	ø	ø	ø		ø	ø	ø	ø	ø	ø	
b					[29, 52, 56]	[57]	[54]	[58]	[59]	ø	ø	ø		[60]	ø	ø	ø	ø	ø	
t						[61]	ø	[62]	[63]	ø	ø	ø		[64]	[60]	ø	ø	ø	ø	
γ							[65, 66]	[67-69]	[68, 70]	ø	ø	ø		ø	ø	ø	ø	ø	ø	
Z/W								[71]	[71]	ø	ø	ø		ø	ø	ø	ø	ø	ø	
H									[72, 73]	[74]	ø	ø		ø	ø	ø	ø	ø	ø	
_ 9/9										ø	ø	ø		ø	ø	ø	ø	ø	ø	
$\tilde{S} \gamma/\pi^{0}s$											[75]	ø		ø	ø	ø	ø	ø	ø	
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Existing searches for two-body resonances^[1]

[1] J. H. Kim et al., J. High Energ. Phys. 2020, 30 (2020), arXiv:1907.06659 [hep-ph]

• The Data Directed Paradigm (DDP) is a search strategy to efficiently identify regions of interest in the data. It requires two ingredients:



• Proof of concept performed for symmetries^[2] and bump searches^[3]

^[2] S. Volkovich et al., Eur. Phys. J. C 82, 265 (2022), arXiv:2107.11573 [hep-ex]

^[3] M. Birman et al., Eur. Phys. J. C 82, 508 (2022), arXiv:2203.07529 [hep-ph]

- Bump search performed with a network mapping invariant mass distribution to statistical significance
 - Input: vector of bin entries from invariant mass histogram
 - Target: vector of statistical significance Z from likelihood-ratio test
 - Representative of an ideal analysis without modelling uncertainties



- Multiple combinations of objects at the LHC
- Selections on variables such as H_T, E_{Tmiss}, leading object p_T, etc



Electron	Leptonic Z
Muon	Boosted hadronic W/Z
Photon	Boosted top
Jet	High mass jet (m > 200 GeV)

Neural network and synthetic data generation

Implementation

Architecture:

- O Use of 1D convolution layers followed by a dense layer
- Intuitive and agnostic to the number of bins in the histogram



Synthetic data generation

Data generation workflow:

1. Obtain background shapes

Analytical functions

$$\begin{array}{l} be^{-ax}, \ \ ax+b, \ \ \frac{1}{ax}+b, \ \ \frac{1}{ax^2}+b, \ \ \frac{1}{ax^3}+b, \\ \frac{1}{ax^4}+b, \ \ a\left(x-x_2\right)^2+y_2, \ \ -a\cdot\ln\left(x\right)+b, \\ (y_1-y_2)\cos\left(a\left(x-b\right)\right)+y_2, \ \ \cosh\left(a\left(x-x_2\right)\right)+b \right) \end{array}$$

Fits to simulation data



2. Inject signal

- Select background
- Ø Generate Gaussian signal
- Ocombine both into observed histogram
- Poisson fluctuate the histogram
- O Calculate true significance with likelihood-ratio test



• Training data:

- Histograms with 30 to 100 bins
- Direction Broad dynamic range, from 10 to 100k entries per bin
- ig> Signals injected from broad significance range, from 1 to 20 σ

Bruna Pascua

Accelerating the search for mass bumps using the Data-Directed Paradigm

Dark Machines sample description

Using the Dark Machines dataset^[4]

- Designed to test anomaly detection techniques
- Contains all of the highest cross-section processes at the LHC
- Generation with Madgraph and Pythia, including fast detector simulation using Delphes
- Events divided into signal regions/channels e.g. channel 3, which is more inclusive with cuts on E_{Tmiss} > 100 GeV and H_T ≥ 600 GeV
- Dataset equivalent to 10 fb⁻¹

SM processes							
Physics process	Process ID	σ (pb)	$N_{\text{tot}} (N_{10 \text{ fb}^{-1}})$				
$pp \rightarrow jj(+2j)$	njets	$19718_{H_T > 600 \text{GeV}}$	415331302 (197179140)				
$pp \rightarrow l^{\pm}\nu_l(+2j)$	w_jets	$10537_{H_T>100 \text{GeV}}$	135692164 (105366237)				
$pp \rightarrow \gamma j(+2j)$	gam_jets	$7927_{H_T} > 100 \text{GeV}$	123709226 (79268824)				
$pp \rightarrow l^+l^-(+2j)$	z_jets	$3753_{H_T>100GeV}$	60076409 (37529592)				
$pp \rightarrow t\bar{t}(+2j)$	ttbar	541	13590811 (5412187)				
$pp \rightarrow t + jets(+2j)$	single_top	130	7223883 (1297142)				
$pp \rightarrow \bar{t} + \text{jets}(+2j)$	single_topbar	112	7179922 (1116396)				
$pp \rightarrow W^+W^-(+2j)$	ww	82.1	17740278 (821354)				
$pp \rightarrow W^{\pm}t(+2j)$	wtop	57.8	5252172 (577541)				
$pp \rightarrow W^{\pm}\bar{t}(+2j)$	wtopbar	57.8	4723206 (577541)				
$pp \rightarrow \gamma \gamma (+2j)$	2gam	47.1	17464818 (470656)				
$pp \rightarrow W^{\pm}\gamma(+2j)$	Wgam	45.1	18633683 (450672)				
$pp \rightarrow ZW^{\pm}(+2j)$	zw	31.6	13847321 (315781)				
$pp \rightarrow Z\gamma(+2j)$	Zgam	29.9	15909980 (299439)				
$pp \rightarrow ZZ(+2j)$	ZZ	9.91	7118820 (99092)				
$pp \rightarrow h(+2j)$	single_higgs	1.94	2596158 (19383)				
$pp \rightarrow t\bar{t}\gamma(+2j)$	ttbarGam	1.55	95217 (15471)				
$pp \rightarrow t\bar{t}Z$	ttbarZ	0.59	300000 (5874)				
$pp \rightarrow t\bar{t}h(+1j)$	ttbarHiggs	0.46	200476 (4568)				
$pp \rightarrow \gamma t(+2j)$	atop	0.39	2776166 (3947)				
$pp \rightarrow t\bar{t}W^{\pm}$	ttbarW	0.35	279365 (3495)				
$pp \rightarrow \gamma \bar{t}(+2j)$	atopbar	0.27	4770857 (2707)				
$pp \rightarrow Zt(+2j)$	ztop	0.26	3213475 (2554)				
$pp \rightarrow Z\bar{t}(+2j)$	ztopbar	0.15	2741276 (1524)				
$pp \rightarrow t\bar{t}t\bar{t}$	4top	0.0097	399999 (96)				
$pp \rightarrow t\bar{t}W^+W^-$	ttbarWW	0.0085	150000 (85)				

^[4] T. Aarrestad et al., SciPost Phys. 12, 043 (2022), arXiv:2105.14027 [hep-ph]

- Consider all possible combinations of objects and selections
 - With 0 to 4 objects per type
 - Additional kinematic cuts:
 - $E_{Tmiss} > 200, 500$ GeV; leading object $p_T > 100, 200, 400$ GeV, ...
- Split the sub-dataset according to jet multiplicity
 - $igodoldsymbol{0}$ 0 jet, 1 jet, 2 jets, ... , \geq 6 jets (depends on the available stat)
 - Should allow to improve S/B
 - Help reducing look-elsewhere effect (since bump should appear at the same place in neighboring jet multiplicities)
- **D** Build variables from available objects
 - Mass distributions of the objects and their combinations
 - Transverse mass distributions including E_{Tmiss}
 - So For jets, use only the 4 leading jets using b-tagging information
- Find the maximum of the histogram and start from there

 $\begin{array}{l} \blacksquare \quad 1\mu + 3j + E_{Tmiss} > 200 \text{ GeV} \\ + \text{ 0e, } 0\gamma, \text{ 0T, } 0Z, \text{ 0Wh, } 0HM \end{array}$

 $1W + 1Z + 3j + p_T(Z) > 100 \text{ GeV}$ $+ 0e, 0\gamma, 0T, 0HM$

...

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 $\begin{array}{|c|c|c|c|c|c|} \hline & 1\mu + 3j + E_{Tmiss} > 200 \ \text{GeV} \\ + 0e, 0\gamma, 0T, 0Z, 0Wh, 0HM \\ \hline & \underline{lal} & m(\mu j_1), m(\mu j_1 E_{Tmiss}), \dots \\ \hline & \underline{lal} & m(\mu j_1 j_2), m(\mu j_1 j_2 E_{Tmiss}), \dots \\ \hline & \underline{lal} & \dots \end{array}$

 $\begin{array}{c|c} \hline & 1W + 1Z + 3j + p_T(Z) > 100 \text{ GeV} \\ &+ 0e, 0\gamma, 0T, 0HM \\ \hline & \underline{lal} & m(Wh, j_1), m(Wh, j_1, E_{Tmiss})... \\ \hline & \underline{lal} & m(j_1j_2), m(j_1j_2, E_{Tmiss})... \\ \hline & \underline{lal} & \dots \end{array}$

...

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 $1W + 1Z + 3j + p_T(Z) > 100 \text{ GeV}$ $+ 0e, 0\gamma, 0T, 0HM$ $\underline{Int} m(Wh, j_1), m(Wh, j_1, E_{Tmiss})...$ $\underline{Int} m(j_1j_2), m(j_1j_2, E_{Tmiss})...$

<u>dd</u> ...

hh

...

Total of 30 000 mass histograms

Processing and calibration

- In real data, signal width determined by detector resolution
 - Different for different final states and/or mass
 - Narrow signal should produce a bump in few bins
- Rebin histogram to reflect the detector resolution we would see in actual experimental data
 - $oldsymbol{0}$ e,μ,γ : Delphes Card formulas, depends on p_{T} and η
 - \circ jets: ATLAS report on jet resolution, depends on p_T



larger for small masses

- HM/T/Wh accounts for 3/3/2 jets
- Using $p_T \approx m/2$ approximation



Binning reflects this with larger bin width when resolution is smaller

Dark Machines datasets

Background-only datasets:

- Obtain shapes to be included in the training sample
- Assess false positive rate

BSM signal datasets:

- Simulated signal data added on top of the backgrounds
- Test the network in a more realistic scenario
- Different levels of difficulty (e.g. cross-section, mass values, etc)

Some of the new physics models we have available include:

- **Q** RPV stop $\rightarrow b\ell$
- **Q** $W' \rightarrow WZ \rightarrow \ell \nu q q, q q \nu \nu$
- **Q** $LQ \rightarrow beb\mu$, bebe, $t\nu t\nu$
- **Q** $Z' \rightarrow 3l$





Performance and finding BSM signals

Performance over synthetic data

- Performance quantified in terms of the difference between predicted and true maximum significance
 - Z^{max}_{true}: maximal significance calculated via the likelihood ratio test
 - Z^{max}_{pred}: maximal predicted significance
- Majority of entries should have $Z_{pred}^{max} Z_{true}^{max}$ close to 0 with the smallest variance possible
- Non-zero significance for background-only histograms
 - Due to look-elsewhere-effect
 - O Could artificially bias the performance at low significance



Performance and finding BSM signals

Performance on the testing sample

• Trained model accurately predicts maximum significance with no bias and a variance of ±0.64



• Excellent discriminating performance of signal and background with an AUC of 0.900



Performance stability

- **o** Good agreement between Z_{pred}^{max} and Z_{true}^{max} is **stable**
 - For all Z^{max}_{true}
 - Over all mass range
 - For all dynamic ranges
 - Sor linear combinations of the functions



Mean & sigma as a function of signal width



Training with fixed signal width has increased bias as signal width increases

- Same behaviour regardless of number of bins and dynamic range
- Accurate calibration takes time (and the DDP should work fast)
- Room for improvement, yet well defined behaviour

Finding BSM signals

Promising results when finding the Higgs bump

Data sampled from the ATLAS plot^[5] using a digitizer



- Predicted resonance at the correct mass
- ${\small 0} \hspace{0.1 cm} \mbox{Predicted significance of 4.6} \sigma$ whereas the ATLAS significance is 3.7 σ

^[5] ATLAS Collaboration, Physics Letters B 716, 1–29 (2012), arXiv:1207.7214 [hep-ex]

Finding BSM signals

• Tested over simulated BSM signals added to the Dark Machines background

Q RPV stop $\rightarrow b\ell$



 $\mathbf{Q} W' \to WZ$



- Successfully finds an excess at the expected mass of the stop at 1 TeV
- Successfully finds bump for W' (with a boosted Z in the final state)
- \odot Other signals tested and successfully found include $LQ \to beb\mu, bebe$ and $Z' \to 3\ell$
- Section 2012 False-positive rate of 0.1% when tested over background-only sample

Conclusion

• Data directed paradigm bump search to scan unexplored final states in search for resonances

- Target smoothly falling invariant mass spectra across a variety of final states
- Exploit full potential of data without the need of simulation or analytical fit

Network implementation and performance

- O Using Dark Machines datasets with highest cross-section processes at the LHC
- Produce mass histograms with binning that reflects detector resolution
- Successfully finds Higgs bump and BSM signals on top of the Dark Machines background, such as RPV stop and W

Future developements

- Application to real experimental data, focusing on Run 2
- O Use full MC simulation data with basic selections
- Sirst iteration using single-lepton trigger and same objects
- > Eventually adding more objects, such as large-R jets

Appendix

Calculation of target significance

For each bin *i*, we perform a **hypothesis test** and obtain a significance value z_i (see arXiv:1007.1727 \mathbb{C})

- Step 1: Define the shape of the signal S for which we are looking, e.g. a Gaussian centered on bin *i*
 - We are currently using the mass values on the x-axis

$$\mathbf{s}_i = \int_{\mathrm{bin}i} f(\mathbf{x}; \mathbf{x}_o, \sigma) \, \mathrm{d}\mathbf{x} \approx f(\mathbf{x}; \mathbf{x}_o, \sigma) \Delta \mathbf{x}$$

$$\begin{cases} x_o = \text{center of bin } i \\ \sigma = \text{width of bin } i \times \text{given number of bins} \end{cases}$$

- **Step 2:** Obtain the **maximum likelihood estimator** $\hat{\mu}$ for this signal S
 - Solution Expected value is $N_{exp} = b_i + \mu s_i$

$$\hat{\mu} = \arg \min - \ln[L(\mu)]$$

$$-\ln[L(\mu)] = -\sum_{i} N_{\rm obs} \ln(N_{\rm exp}) - N_{\rm exp} + \ln\left(\frac{1}{N_{\rm obs}!}\right)$$



- Observed data (Nobs) is the synthetic data
- Unfluctuated background (b) is our null hypothesis
- Injected signal is the one added into the observed data
- Test signal is the signal that maximizes the likelihood for each bin

Calculation of target significance

For each bin *i*, we perform a **hypothesis test** and obtain a significance value z_i (see arXiv:1007.1727 \mathbb{C})

 Step 3: Obtain the likelihood for both the background-only hypothesis and the maximum-likelihood estimator (MLE) signal hypothesis

$$oldsymbol{\partial}$$
 $-\ln[L(\mu=0)]$ $ightarrow$ Background-only ($N_{
m exp}=b_i$)

- \bullet $-\ln[L(\mu = \hat{\mu})] \rightarrow$ Signal MLE ($N_{exp} = b_i + \hat{\mu}s_i$)
- Step 4: Obtain the significance using the profile likelihood ratio test

If $\hat{\mu} \geq 0$:

$$z_i = \sqrt{q_0}$$
 with $q_0 = -2 \ln \left[\frac{L(0)}{L(\hat{\mu})} \right]$

If $\hat{\mu} < 0$:

$$z_i = -\sqrt{-q_0} \quad \text{with} \quad q_0 = \left\{ \begin{array}{l} -2 \ln \left[\frac{L(\hat{\mu})}{L(0)} \right], \text{if two-sided} \\ 0, \text{otherwise} \end{array} \right.$$



- Observed data (Nobs) is the synthetic data
- **O** Unfluctuated background (b) is our null hypothesis
- Injected signal is the one added into the observed data
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Background functions in the framework

ax + b	<pre>def linear(x, p0, p1): return p0 * x + p1</pre>	ae ^{-bx}	<pre>def exponential(x, p0, p1): return p0 * np.exp(-p1 * x)</pre>
$\frac{1}{ax} + b$	<pre>def one_over_x(x, p0, p1): return 1/(p0*x) + p1</pre>	$a(x-x_{max})^2+y_{min}$	<pre>def parabola_half(x, p0, p1, x2, y2): return p0*(x - x2)**2 + y2</pre>
$\frac{1}{ax^2} + b$	<pre>def one_over_x_squared(x, p0, p1): return 1/(p0*x**2) + p1</pre>	$\Delta y \cos[a(x-b)] + y_{max}$	<pre>def cos_quarter(x, p0, p1, y1, y2): return (y1 - y2)*np.cos(p0*(x - p1)) + y1</pre>
$\frac{1}{ax^3} + b$	<pre>def one_over_x_cubed(x, p0, p1): return 1/(p0*x**3) + p1</pre>	$-a\ln(x)+b$	<pre>def ln_negative(x, p0, p1): return -p0*np.log(x) + p1</pre>
$\frac{1}{ax^4} + b$	<pre>def one_over_x_to_4th(x, p0, p1): return 1/(p0*x**4) + p1</pre>	$\cosh[a(x - x_{max})] + b$	<pre>def cosh_half(x, p0, p1, x2): return np.cosh(p0*(x - x2)) + p1</pre>
$\frac{1}{ax^n} + b$	<pre>def one_over_x_to_nth(x,p0,p1,n): return 1/(p0*x**n) + p1</pre>	\rightarrow where <i>n</i> is	taken randomly from the range [0.01, 10]