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Motivation

- Emission times of carriers are crucial for defect characterization
 - at close to RT conventional methods (TSC,DLTS) are difficult to operate particularly at high fluences
- Direct observation of de-trapped charge previously trapped during the drift of carriers (prompt generation of non-equilibrium carriers)
- De-trapping can play some role in CCE for longer shaping times and low bias voltages (not really LHC case)



Conventional TCT was used with small modifications:

- Signals taken on the time scale of few μs
- Bias-T with different bandwidth was used
- Self-triggering of oscilloscope to avoid small DC offset

Samples

- MCz-n type samples (WODEAN), 1 kΩcm, V_{fd}~250 V
- Samples irradiated with neutrons to 10¹⁴ cm⁻² and 3x10¹⁴ cm⁻²
- Measurements shown are after annealing of 80 min at 60°C



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Method (II) – non irradiated sample



As there is no charge trapping the evolution of the signal in time is the consequence electronics (amplifier, Bias-T):

$$Q(t) = \int_{0}^{t} I(x) * h(t - x) dx \approx Q(t_d) h(t) \quad t \gg t_d$$
$$h(t) = \sum_{i=0}^{2} a_i \exp(-t/\tau_i) \qquad \begin{aligned} \mathbf{a}_1 = 0.245, \, \tau_1 = 342 \, \mathrm{ns} \\ \mathbf{a}_2 = 0.755, \, \tau_1 = 68000 \, \mathrm{ns} \end{aligned}$$

Fit is shown in the insets!

Method (III) – diode irradiated to 10¹⁴ cm⁻²



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Method (IV) – diode irradiated to 10¹⁴ cm⁻²



The shape of the Q(t) changes with voltage:

- □ It increases with time for lower voltages
- □ It has the shape similar as non-irradiated sample for very high voltages

The reason for the increase is the de-trapping of the charge trapped during the drift (note that at low voltages a lot of charge is trapped)

Method (V) – diode irradiated to 10¹⁴ cm⁻²

The difference between two plots can be expressed:

$$- Q_2(t) - Q_1(t) = \int_0^t [I_2(x) - I_1(x)] * h(t - x) \, dx =$$
$$= \int_0^{t_d} [I_2(x) - I_1(x)] * h(t - x) \, dx + \int_{t_d}^t [I_2(x) - I_1(x)] * h(t - x) \, dx$$

$$T(t > t_d) = \sum_i \frac{Q_{t_i}}{\tau_{d_i}} \exp(-t/\tau_{d_i})$$

$$\Delta Q(t) - \Delta Q(t_d) h(t) = \sum_{j=1}^{2} \sum_{i=1}^{2} \left(\Delta \mathbf{Q}_{ii} \left\{ \frac{a_j \tau_j}{\tau_{d_i} - \tau_j} \left(\exp\left(\frac{t}{\tau_{d_i}}\right) - \exp\left(-\frac{t}{\tau_j}\right) \right) \right] \right)$$

This equation is fit to the measured data! Two traps were assumed (index *i*).

- Unknowns: τ_{d1} , τ_{d2} , $\Delta Q_{t1}/\Delta Q_{t2}$
- Constrain: $\sum_{i=1}^{2} \Delta Q_{ti} = Q_2(t_d) Q_1(t_d)$

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- Fits are reasonably good, however at T<=-10°C the measurements can be fit with single exponential (only one free parameter)
- De-trapping times are in the range from 1-10 µs, the long term dominates (~80% of the de-trapping amplitude)

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Interpretation of the measured time constants

- Trapped charge can be trapped again which influences measurements, hence measured de-trapping times \(\tau_d\) should be corrected – scaled to get proper detrapping times.
- MC-simulation was used assuming constant trapping distance over the entire volume of the detector (determines CCE=Q/Q₀). The Q(t) obtained from fit to the simulated values was compared to input values.
- The correction factor depends only on trapping times (probability of being re-trapped).



Determination of the traps



- Short component can only be established reliably up to 0°C:
 - 20% of the de-trapping amplitude (trapped charge)
 - longer time constants and smaller amount of trapped charge make determination very unreliable
- The introduction rate for the defect responsible for short component should be around 0.1 cm⁻¹ (to explain trapping times/amount of trapped charge)
- Contrary introduction rates should be very high ~100 cm⁻¹ for the defect responsible for long component?

Can they be correlated with DLTS/TSC measurements?

Microscopic measurements:

I. Pintilie, 3rd MC-PAD training event, Ljubljana, 2010



TCT measurements:

Fast (could be H(152)?): $E_{t1}-E_v=0.42 \text{ eV}$ $\sigma_p=8e-14 \text{ cm}^{-2}$

Slow (could be H(140) or/and H(116)?): E_{t2} - E_v =0.31 eV σ_p =1e-16 cm⁻²



What about de-trapping times of electrons?



Sensitivity is lower due to less charge trapped (mainly the effect of velocity)

 Emission times of electrons are larger than the time scale investigated – de-trapping times of electrons are larger than ~10 µs

Conclusions & outlooks

- A method for measurement of de-trapping times using TCT was developed.
- It was found that de-trapping times of holes are of order of μs
 - □ effective traps with $E_t=0.31 \text{ eV}$, $\sigma_h=1e-16 \text{ cm}^{-2}$ (20% of the trapped charge) and $E_t=0.42 \text{ eV}$, $\sigma_h=8e-14 \text{ cm}^{-2}$ (80% of the trapped charge) result in good fit to the data
- Electron de-trapping times are longer that few tens μs
- Future work
 - □ annealing studies
 - more temperature points
 - □ improved modeling (bias dependent de-trapping?)
 - □ material dependence?

But as usually there is a lack of manpower