

The Role of Effective Field Theories in Precision Physics

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DK-PI Meeting

28-29th September 2023



HELMHOLTZ
RESEARCH FOR GRAND CHALLENGES

Why Precision Physics?

New Physics can manifest itself as:

signals in processes where the SM predictions are exceedingly small
small corrections to Standard Model (SM) predictions

LHC and other **modern experiments** now often reach the **percent level precision** in their measurements
A high luminosity run of the LHC is expected to deliver integrated luminosity reaching inverse attobarns

To do list for theorists:

Established theories must be understood better (in particular strong and weak interactions)
Higher precision calculations (incl. resummation, control non-perturbative physics, improve MC-generators, ...)
Precise measurements of fundamental parameters (coupling constants, particle masses, CKM, ...)

In this context Effective Field Theories (EFT) play a key role

Effective Field Theory

Effective Field Theory

Effective Field Theory is a **field theory**, very much like QED or QCD

Provides a coherent description of a system using only the **relevant degrees of freedom**

The **power counting and symmetry** arguments tells you which contributions are relevant

Example: Hydrogen Atom

Obviously, SM of particle physics contains everything (electromagnetic, strong and weak interaction)

In practice the SM-action is way too complicated and much of it is entirely irrelevant for describing the H-atom (quark structure of the proton, weak interaction, Higgs)

H-atom Hamiltonian contains only relevant dof's

Leading order:

- Non-relativistic electron
- Non-relativistic proton (the only property of the nucleus we need is the electric charge)
- Interacting via Coulomb potential

} $\mathcal{O}(\alpha^2)$

Higher order corrections

- Fine structure (m_e, m_p): relativistic correction, darwin correction, spin-orbit interaction
- Lamb shift (QED effects)
- ...

$\mathcal{O}(\alpha^4)$

$\mathcal{O}(\alpha^5)$

The effective theory for H-atom is non-relativistic QED


How to construct an EFT ?

- ✓ Identify the relevant degrees of freedom
- ✓ Identify the symmetries that constraint the interaction
- ✓ Identify the power counting (e.g. $\lambda = \Lambda_{\text{QCD}}/E_{\text{cm}}$)
- ✓ Do the expansion correctly (integrate out all the off-shell dof's)

$$L = L^{(0)} + \lambda L^{(1)} + \lambda^2 L^{(2)} + \lambda^3 L^{(3)} + \dots$$

The key concept is **locality** which separates the high energy and the low energy dynamics (factorization)

The EFT Lagrangian is a sum of local, gauge and Lorentz invariant operators

$$L^{(n)} = \sum_{m=0}^N C_m O^{(m)}$$


Short distance coefficients accounts for off-shell dof's

Operators constructed from relevant dof's

Why are EFTs very powerful tools in QCD?

- ✓ Simplifies the calculation by only including relevant interactions and dealing with **one scale at a time**.
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- ✓ **Sum large logarithms** of the ratios of scales which could spoil the perturbative expansions

$$\sigma = \sum_{n \geq k} a_{nk} \alpha_s^n L^k \quad \text{with} \quad L = \log \left(\frac{M_w}{m_b} \right) \quad \text{or} \quad L = \log \left(\frac{\Lambda_{\text{QCD}}}{m_b} \right)$$

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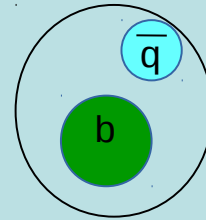
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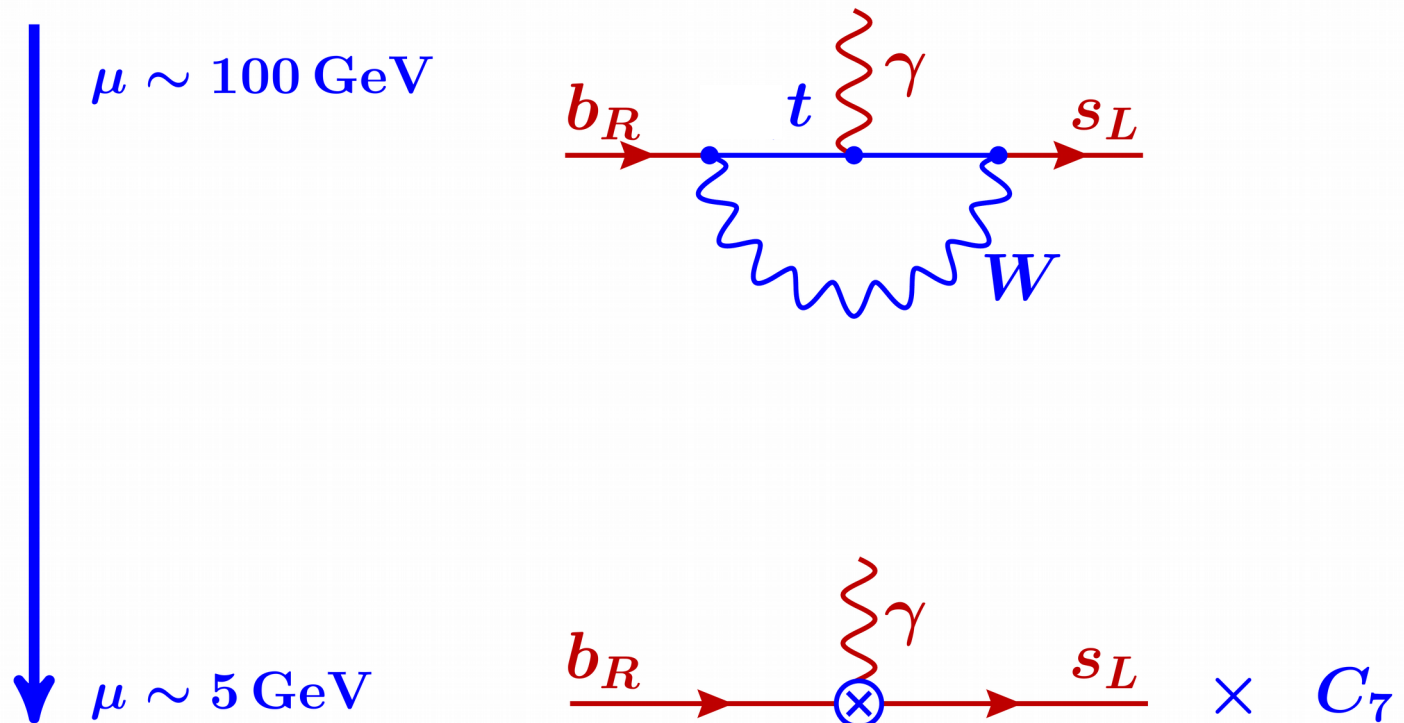
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- ✓ One can strictly prove **factorization theorems** where the dynamics of different dof's are completely decoupled
- ✓ Makes the **relevant physics** and **symmetries manifest**

Inclusive decays of B-mesons



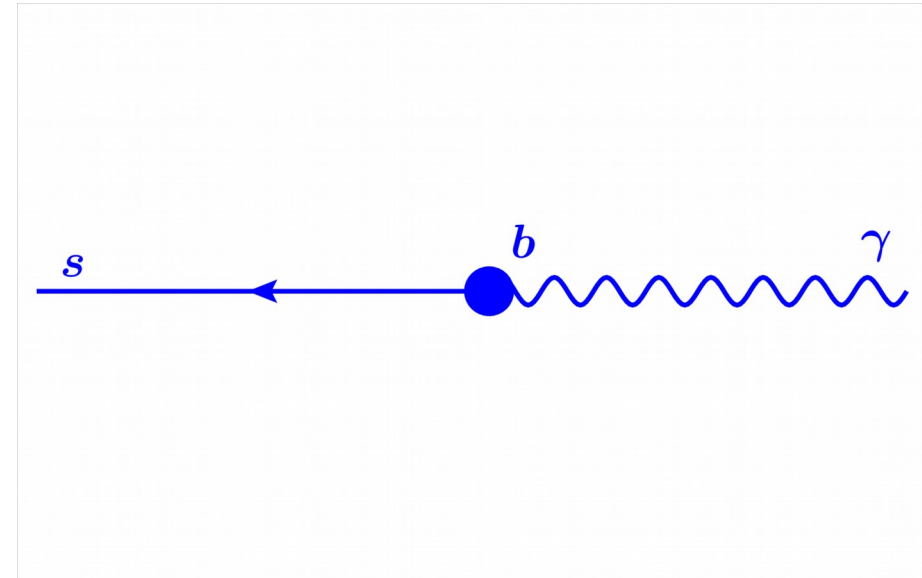
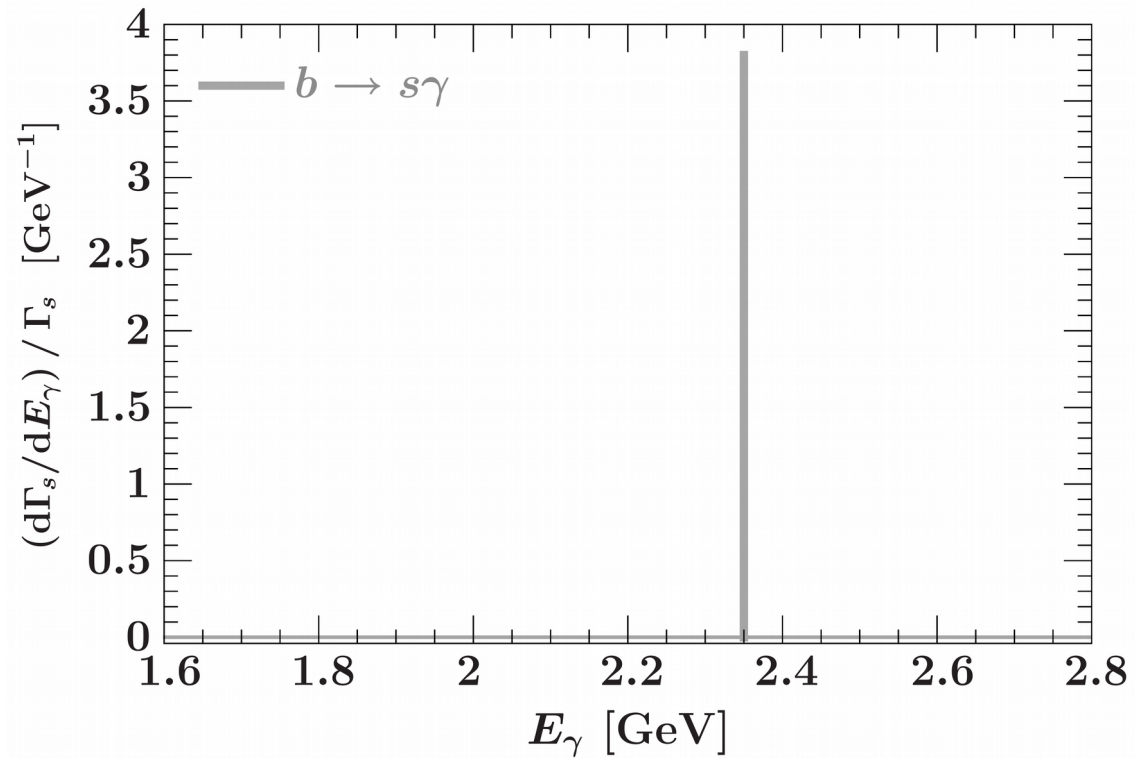
EFT setup in inclusive B-meson decay

We wish to describe the photon energy spectrum in $B \rightarrow X_s \gamma$



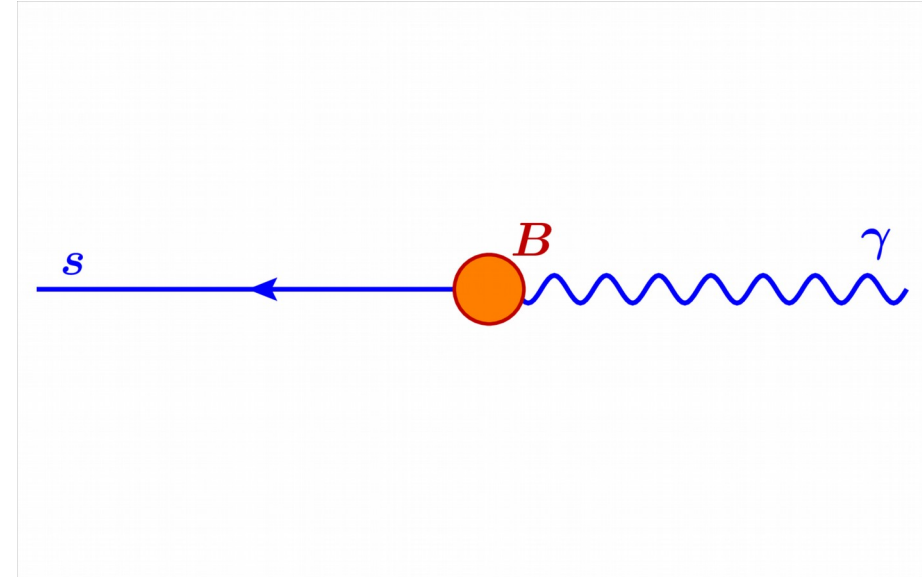
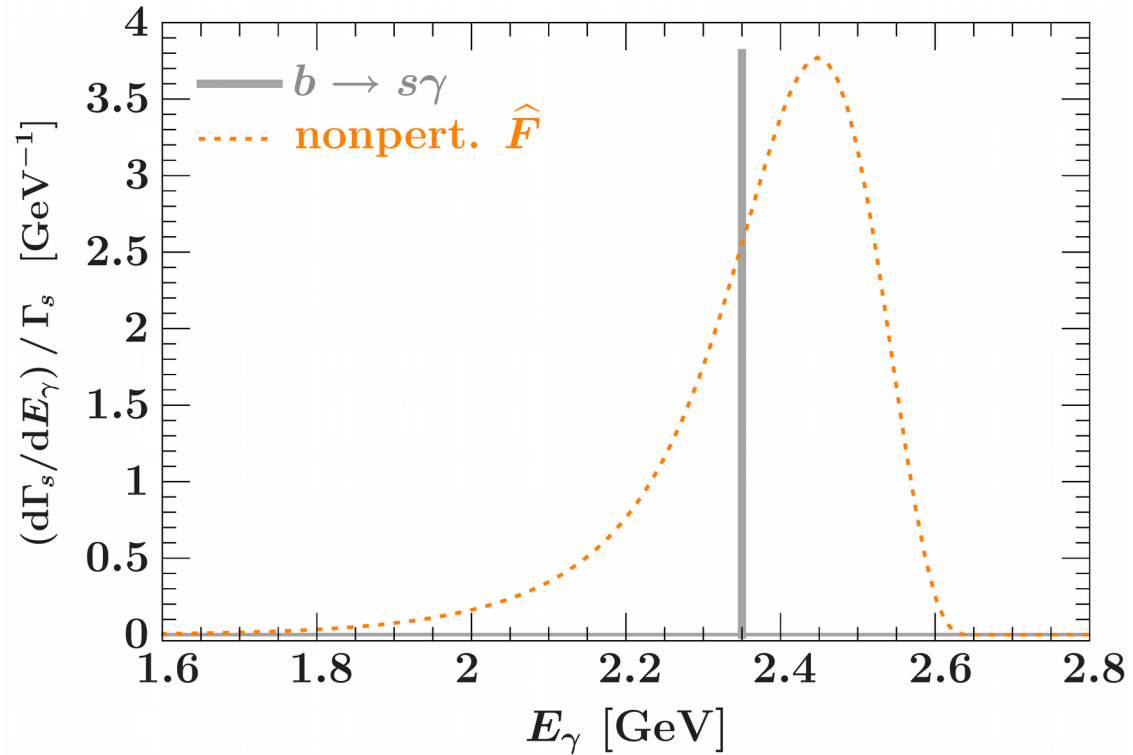
The bottom & strange quark masses and their momenta are much smaller than the W mass
Integrate out all the heavy particles i.e. W, Z, H, and top quark

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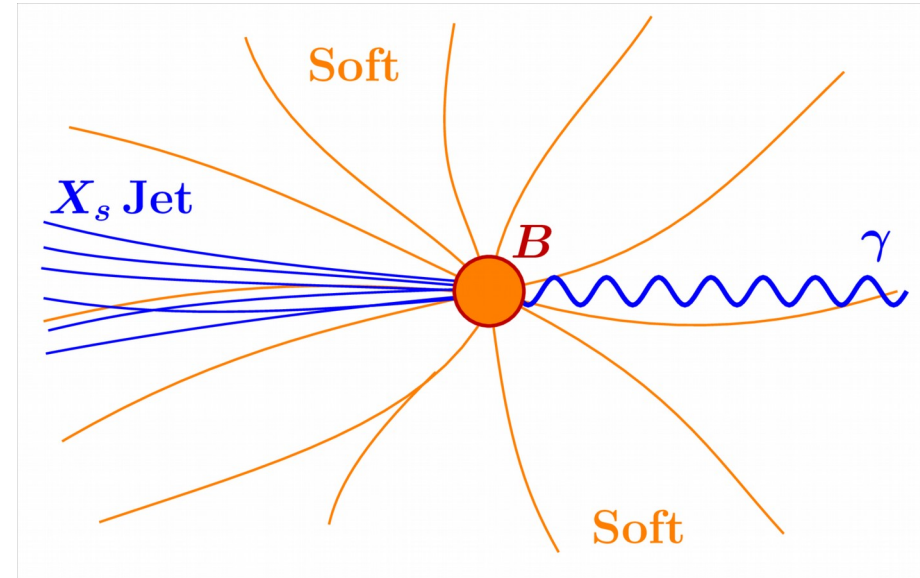
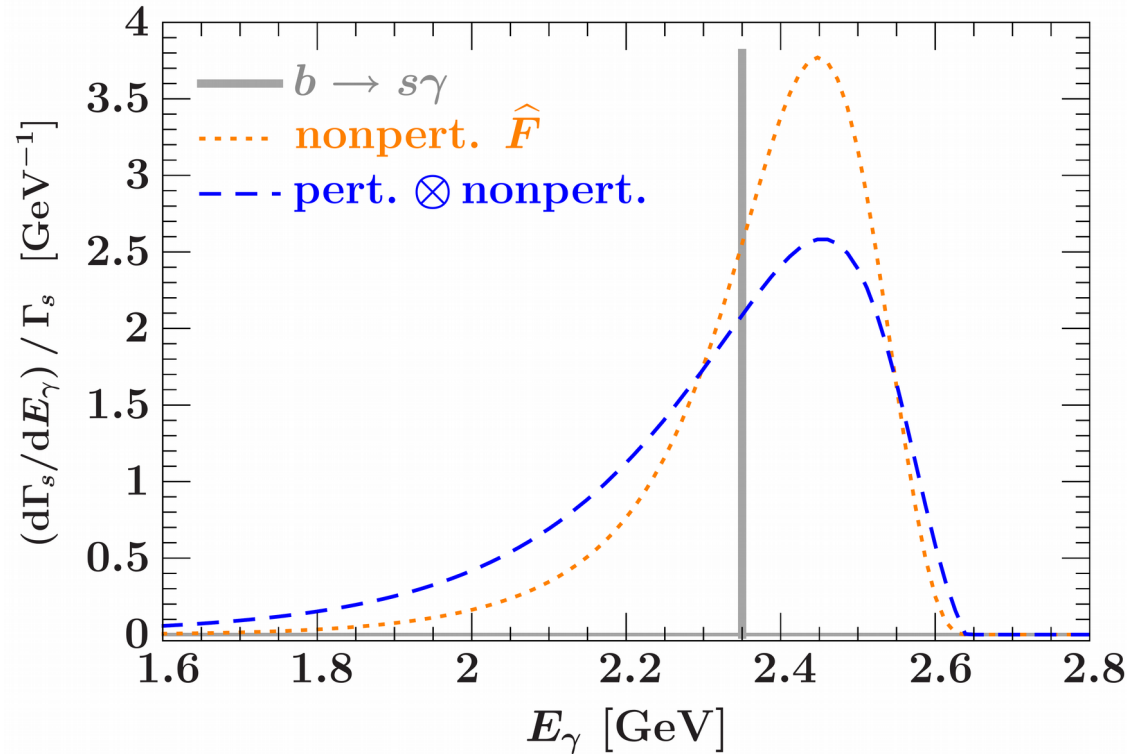
- Start with a stationary b-quark, the transition rate at tree level is: $d\Gamma/dE_\gamma = |C_7|^2 \delta(E_\gamma - m_b/2)$

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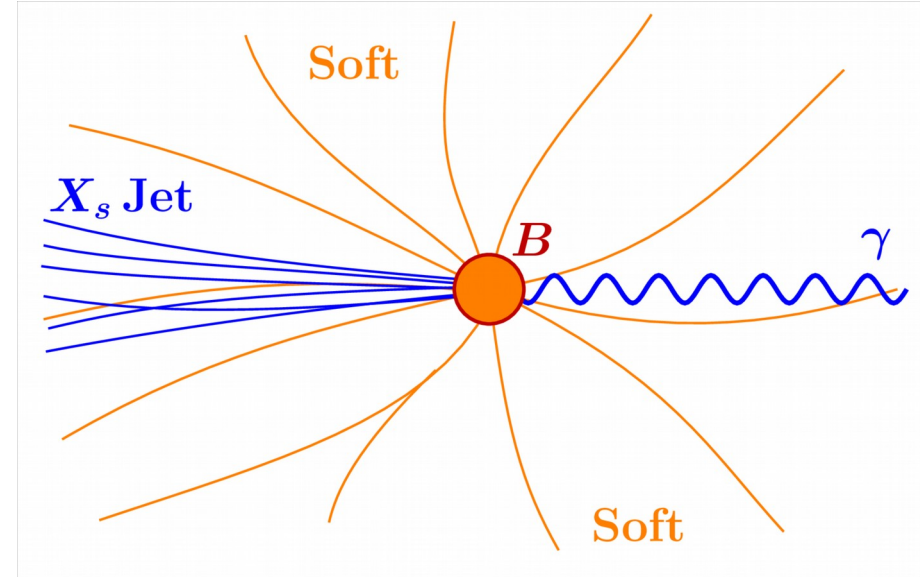
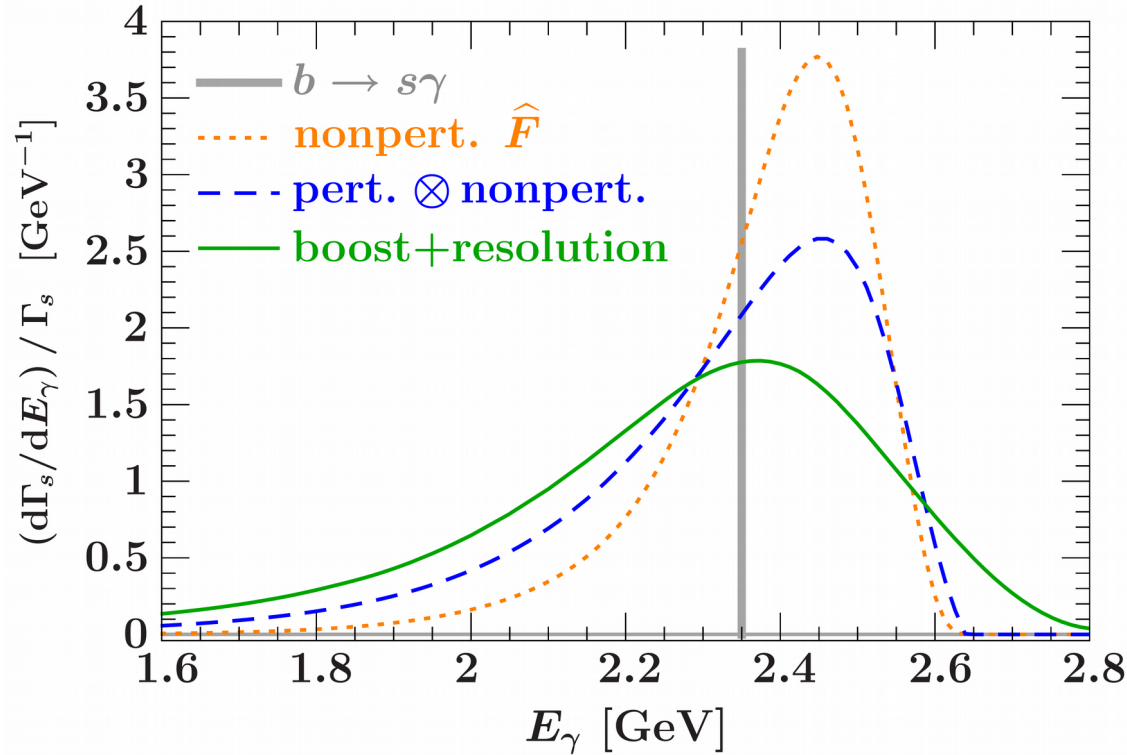
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- Account for the motion of b-quark in the B-meson (non-perturbative): $d\Gamma/dE_\gamma = |C_7|^2 F_B(E_\gamma - m_B/2)$

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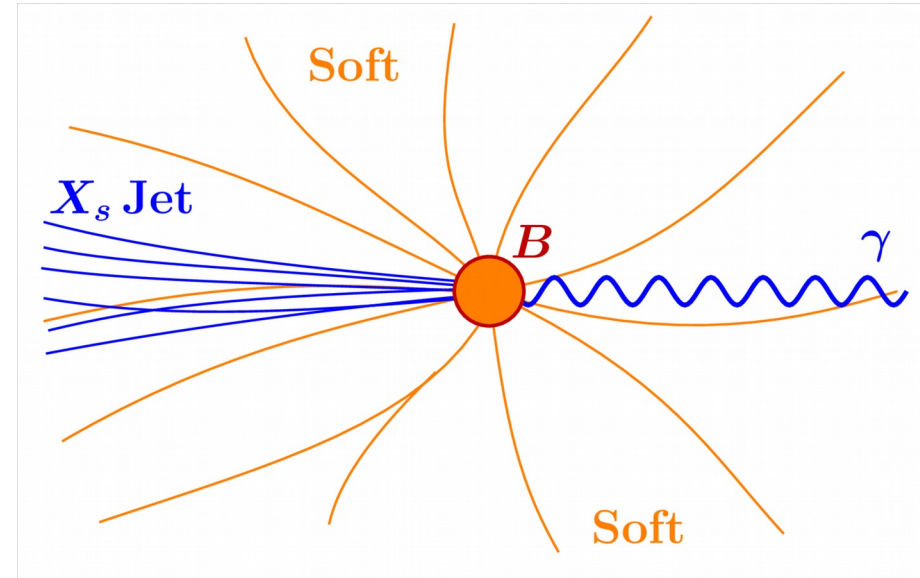
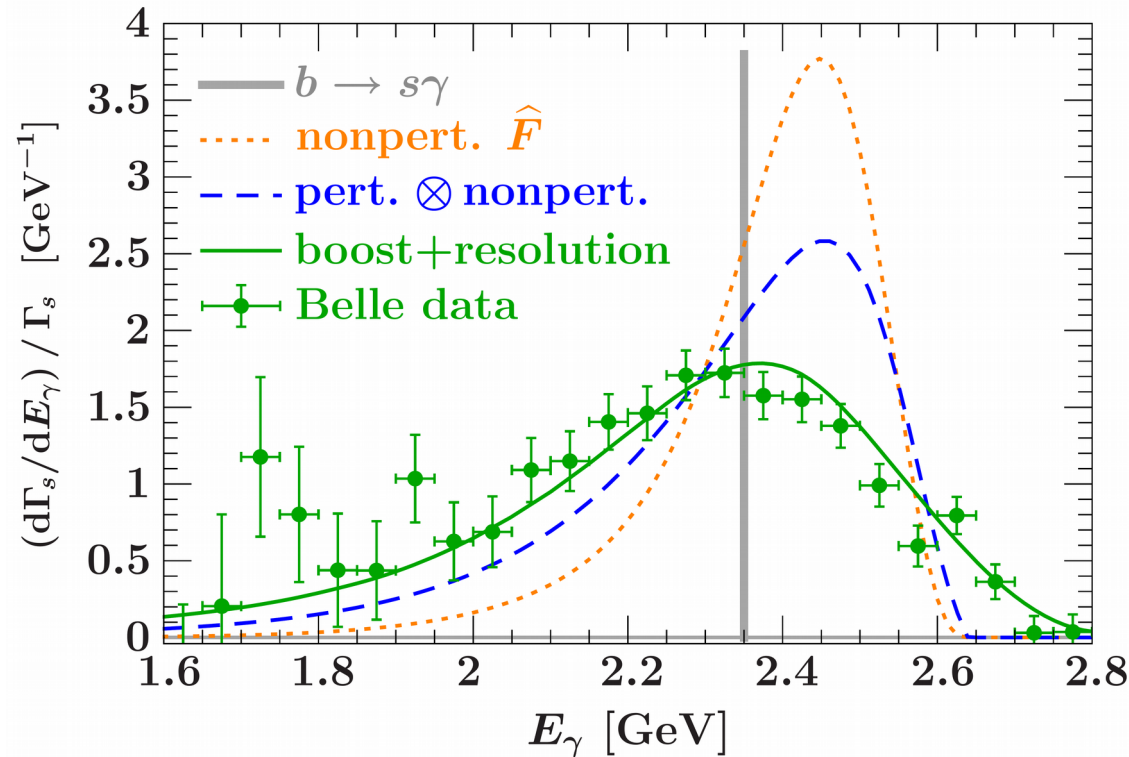
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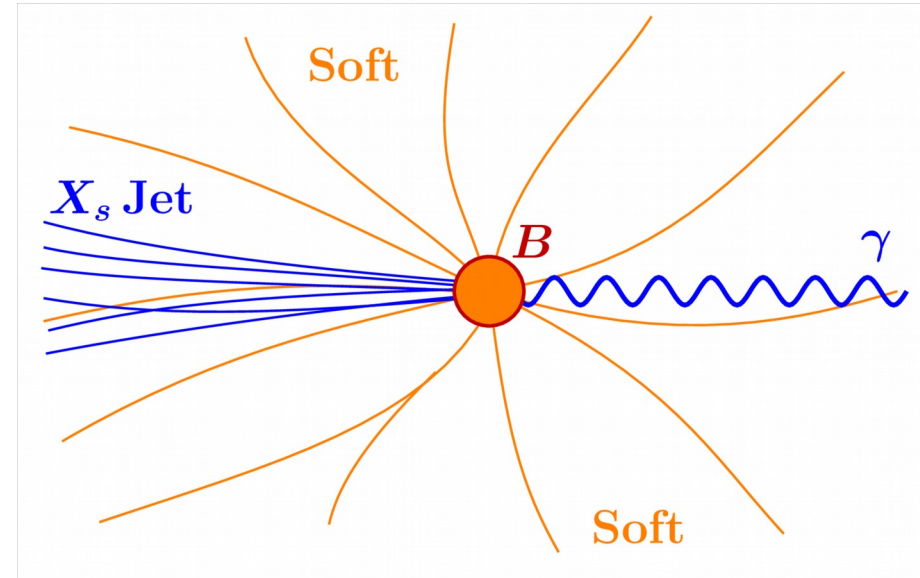
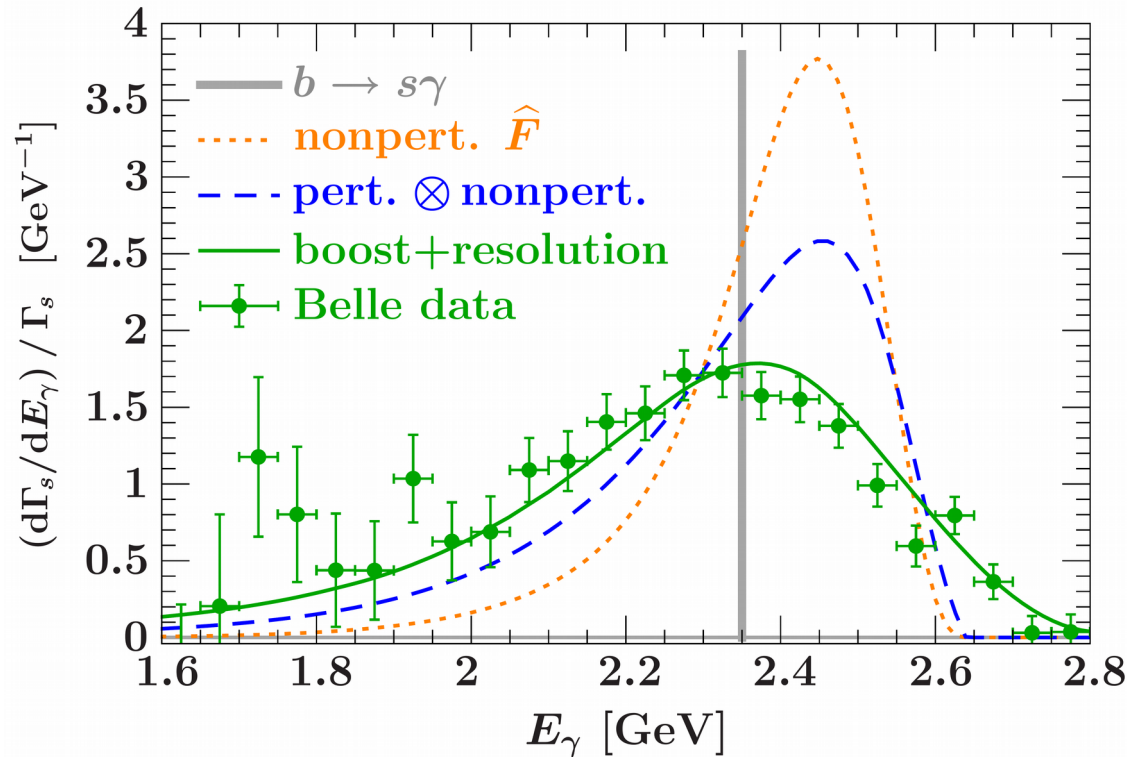
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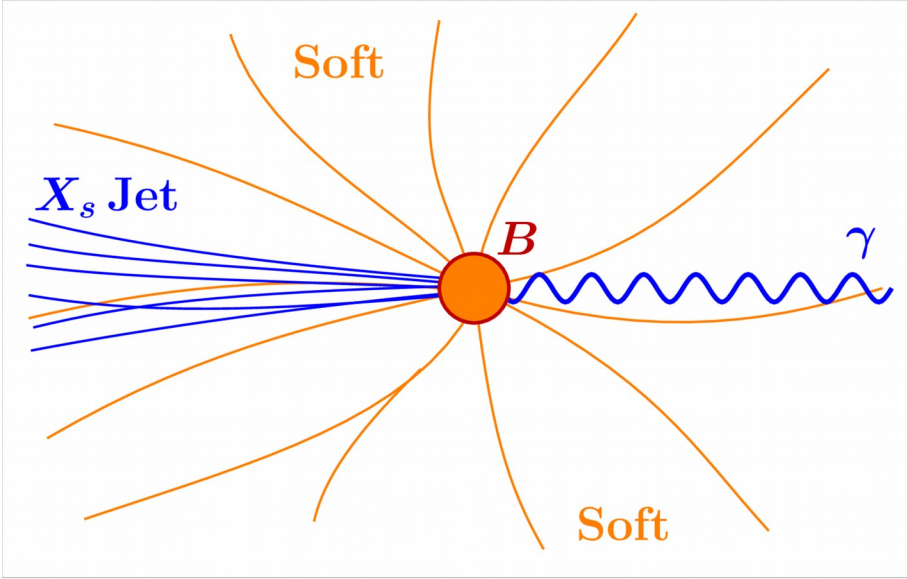
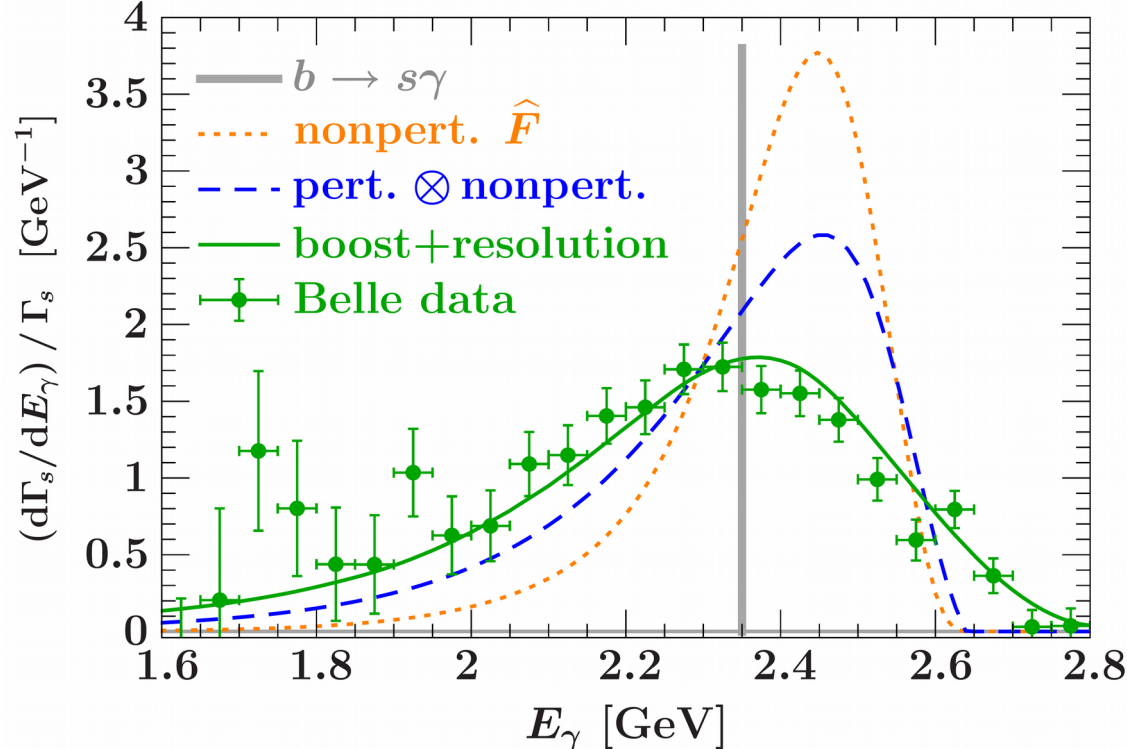
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- Analogous story for other inclusive B-meson decays: $B \rightarrow X_u l \nu$ & $B \rightarrow X_s l^+ l^-$

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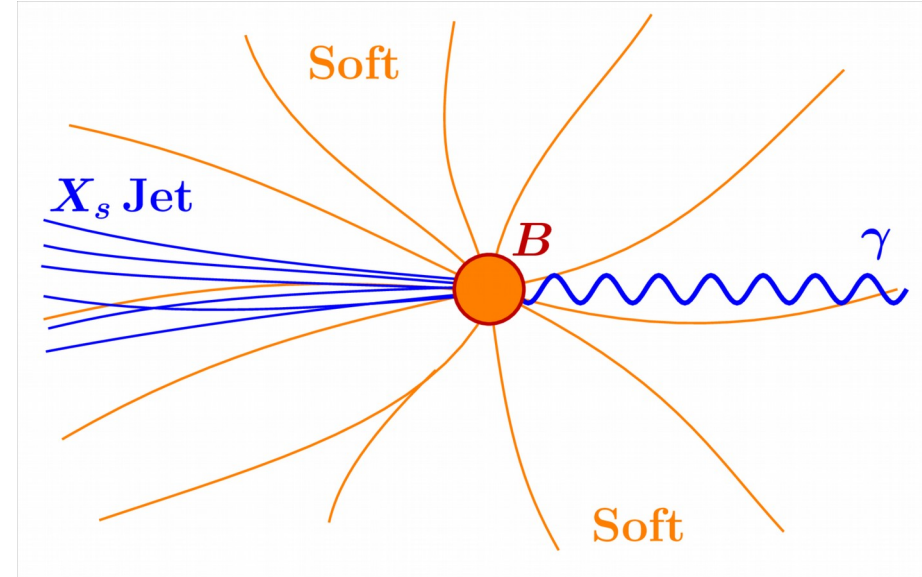
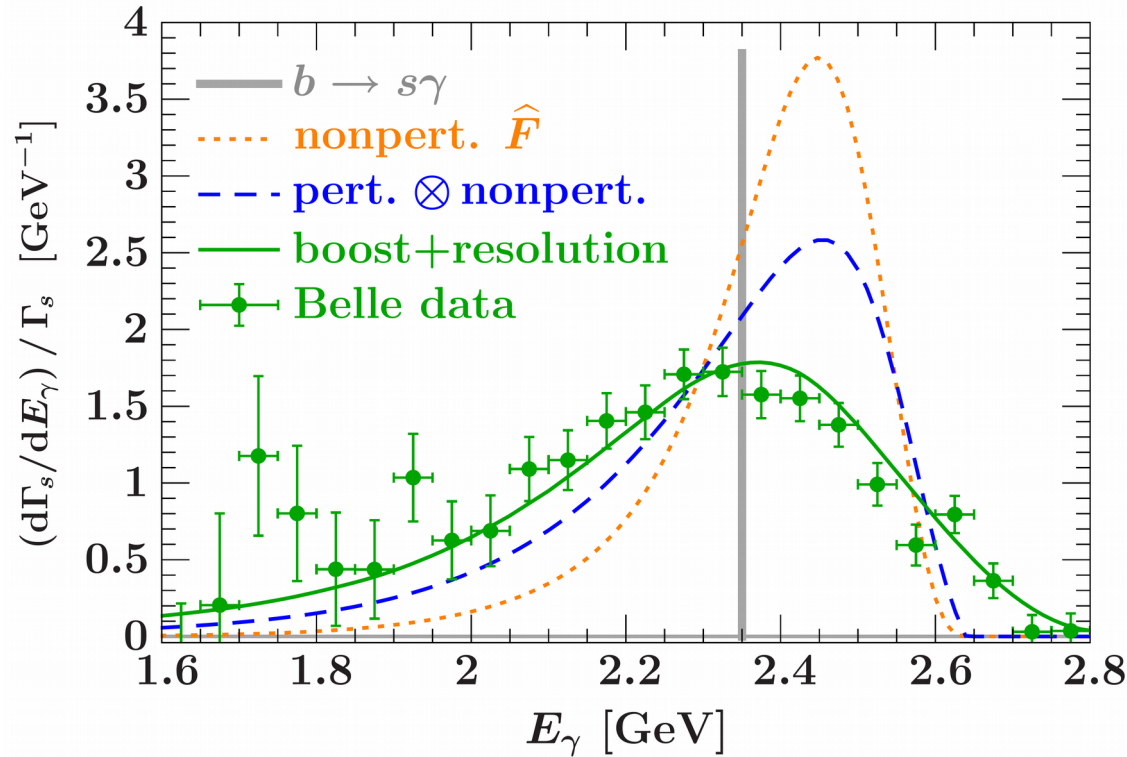
Effective theory setup:

Effective weak theory

Heavy quark effective theory (HQET)

Soft & collinear effective theory (SCET)

EFT setup in inclusive B-meson decay



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Factorization Theorem

$$d\Gamma/dE_\gamma = |C_7|^2 H \times J \otimes S \otimes F_B$$

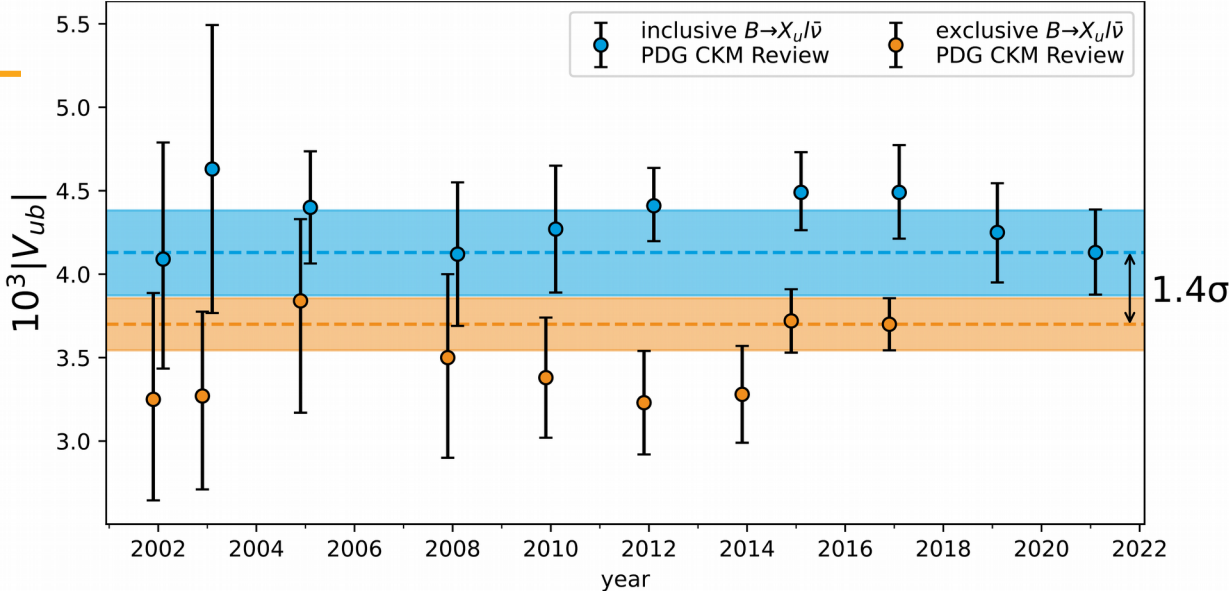
Applications of EFTs in Precision Physics

Inclusive B-meson decays

Problem: persistent tension between the **inclusive** and **exclusive** determinations of $|V_{ub}|$

$$\sum_{X_u} d\Gamma(B \rightarrow X_u l \bar{\nu})$$

$$\begin{aligned} & d\Gamma(B \rightarrow \pi l \bar{\nu}) \\ & d\Gamma(B \rightarrow \rho l \bar{\nu}) \\ & d\Gamma(B_s \rightarrow K l \bar{\nu}) \end{aligned}$$

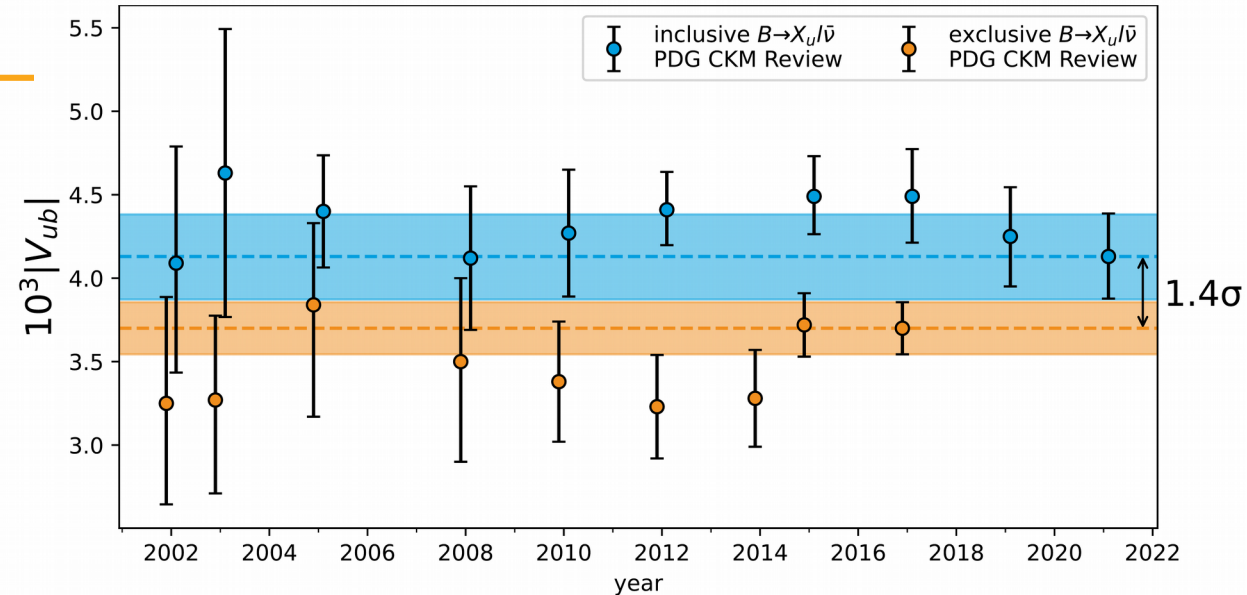


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$$\begin{aligned} d\Gamma(B \rightarrow X_s \gamma) / dE_\gamma &= H^{(B \rightarrow X_s \gamma)} \times J \otimes S \otimes F_B \\ d\Gamma(B \rightarrow X_u l \nu) / dp_X^+ &= H^{(B \rightarrow X_u l \nu)} \times J \otimes S \otimes F_B \end{aligned}$$

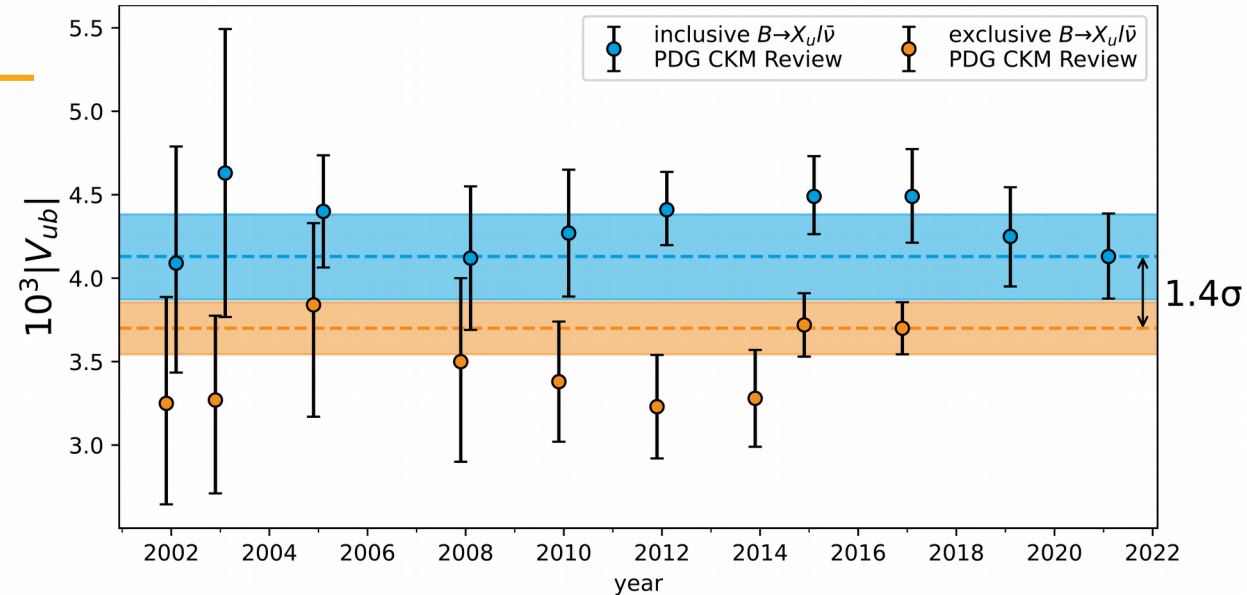
- ✓ Factorization theorem separates the short-distance from the long distance physics
- ✓ $B \rightarrow X_s \gamma$ can be used to extract the **universal non-perturbative shape function**, then use it as an input for the other inclusive B-meson decay modes

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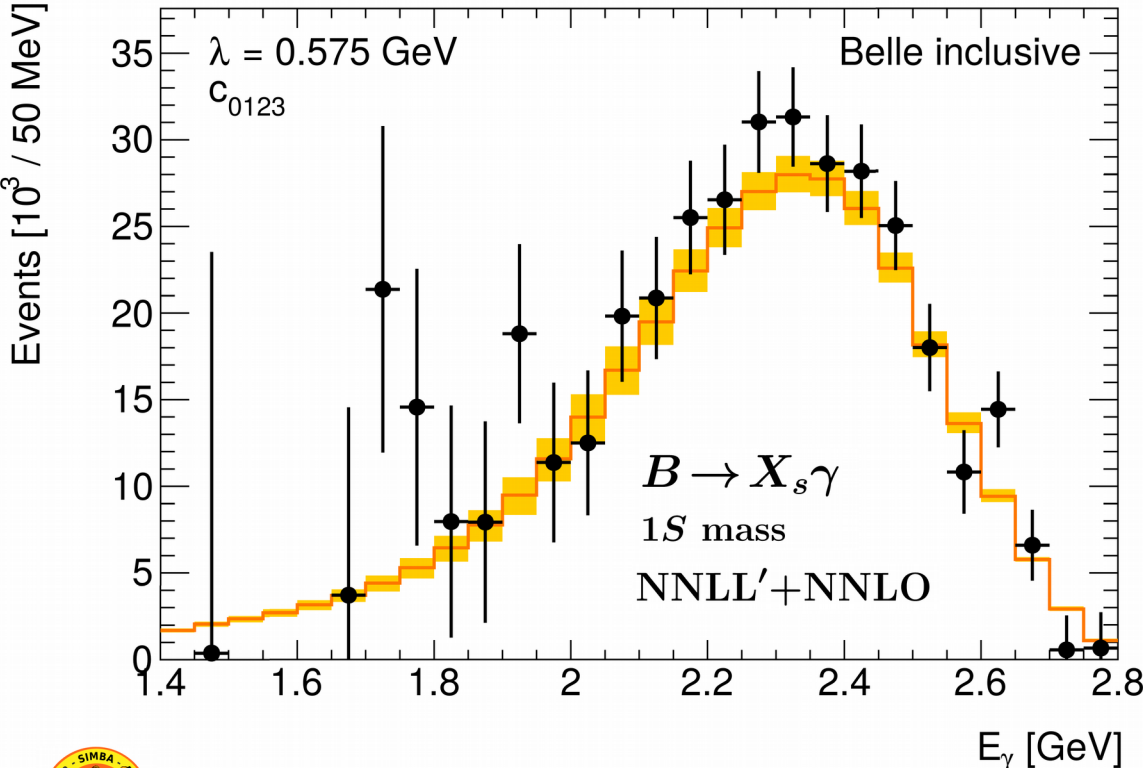
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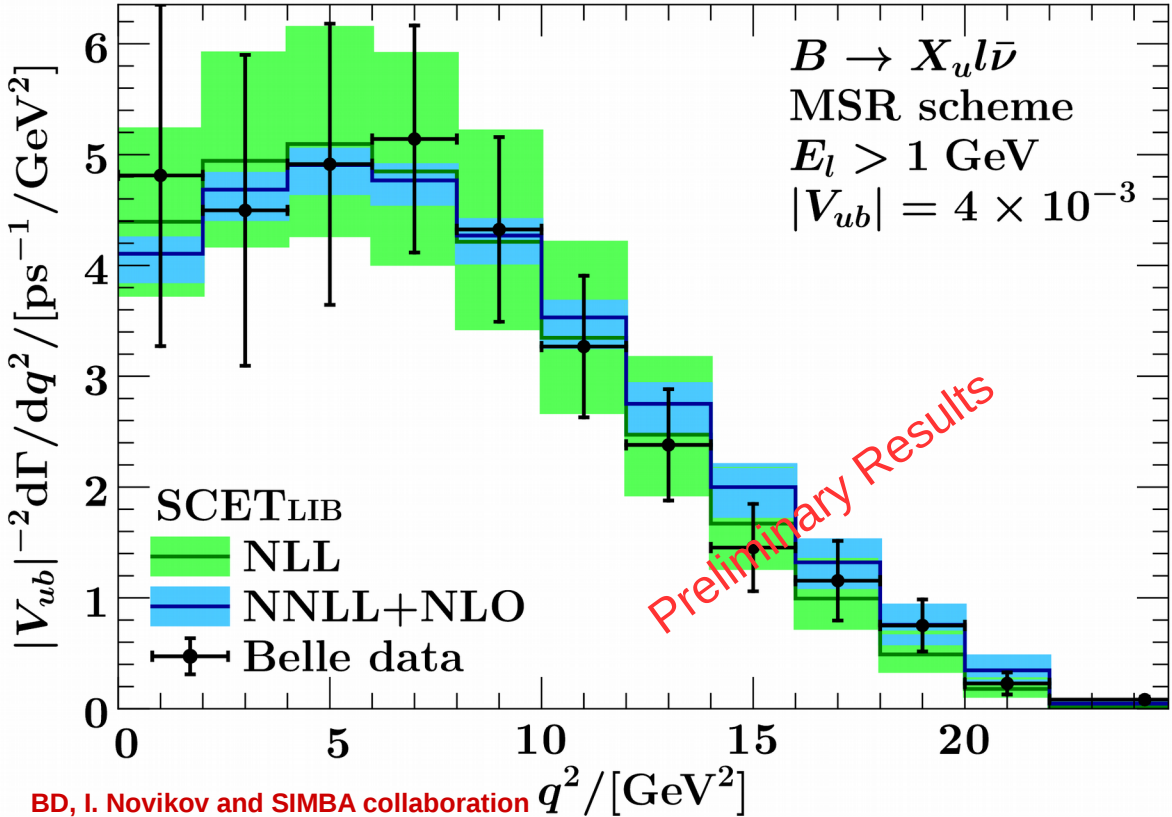
Aim: determine $|V_{ub}|$ from a global fit to inclusive B-meson decay

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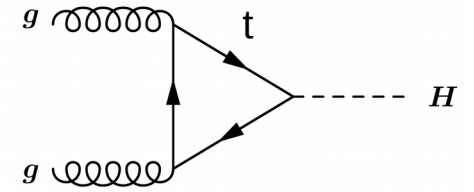


SIMBA collaboration, F. U. Bernlochner, H. Lacker, Z. Ligeti, I. W. Stewart, F. J. Tackmann and K. Tackmann (2021)

BD, I. Novikov and SIMBA collaboration $q^2 / [\text{GeV}^2]$

Total Higgs fiducial cross section

QCD correction to the total inclusive cross sections for $gg \rightarrow H$ is known up to N³LO



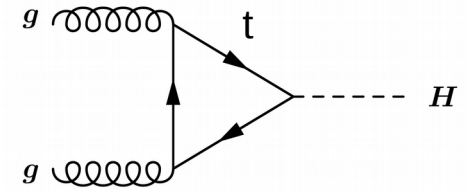
Experimental measurements at LHC apply kinematic selection cuts on Higgs decay products

For example ATLAS fiducial cuts for $H \rightarrow \gamma\gamma$ \rightarrow $p_T^{\gamma 1} \geq 0.35 m_H$, $p_T^{\gamma 2} \geq 0.25 m_H$
 $|\eta^\gamma| \leq 1.37$ or $1.52 \leq |\eta^\gamma| \leq 2.37$

Problem: fiducial cuts break the azimuthal symmetry and introduce linear power corrections which increase the IR sensitivity of the fiducial obs.

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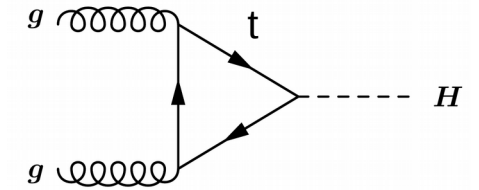
rEFT : integrate out the top quark

Soft & collinear effective theory (SCET)

$$d\sigma^{(\text{f.d.})}/dq_T = \int dY A(q_T, Y; \Theta) H \times B \otimes B \otimes S(q_T)$$

✓ Factorization theorem allows us to account for the **leading and sub-leading power corrections** systematically and **sum them to all orders**

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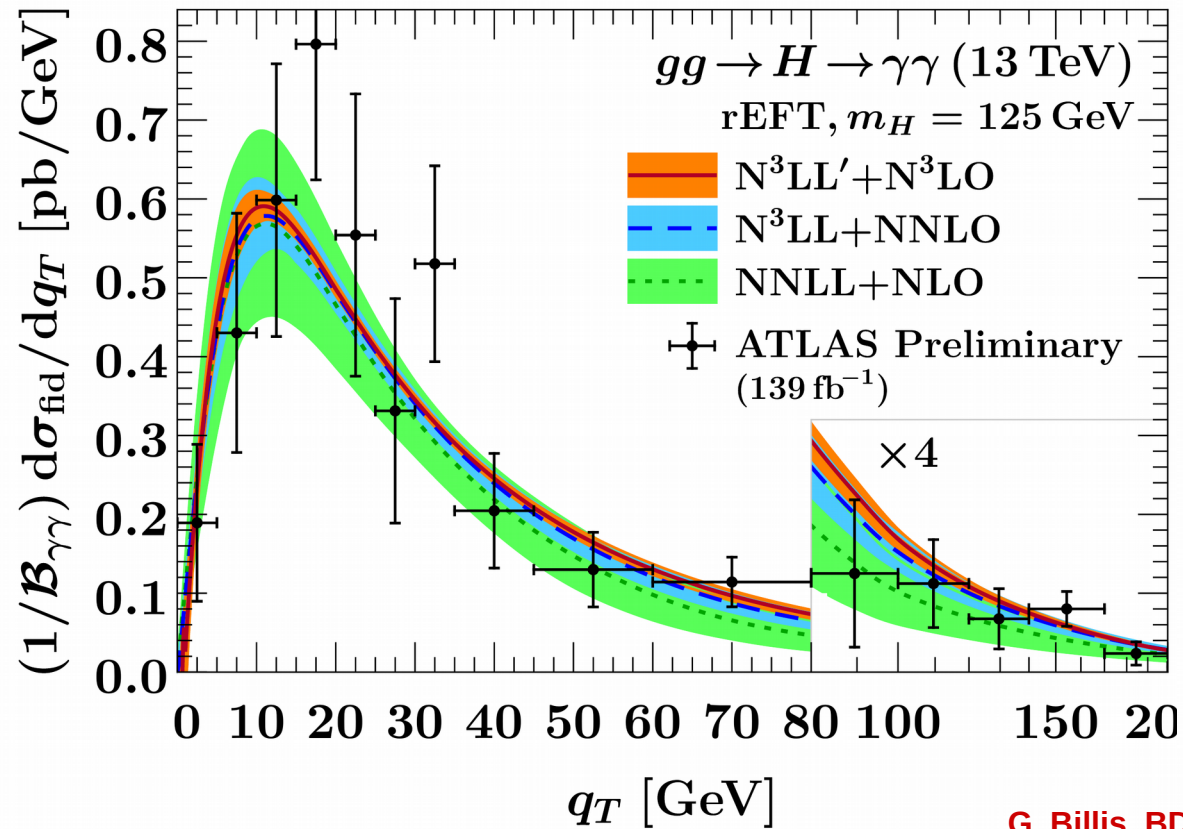
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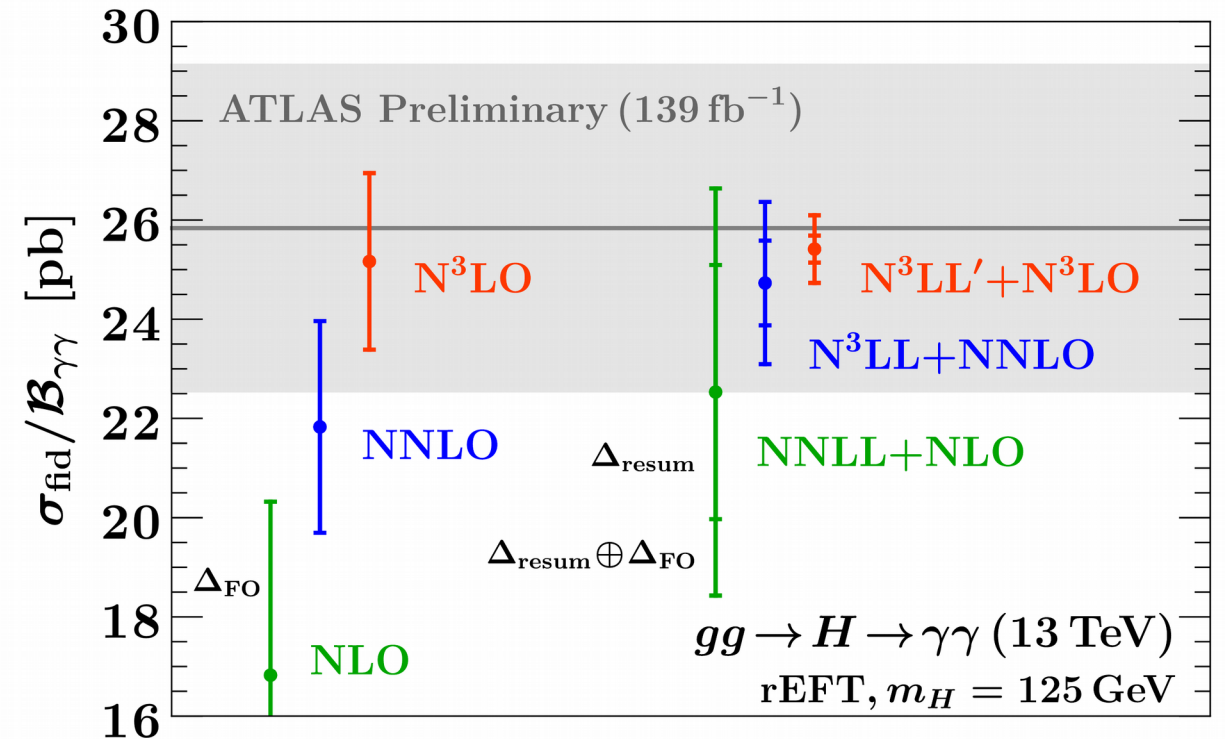
Aim: compute the total (and differential) Higgs cross section with fiducial cut including resummations

Total Higgs fiducial cross section

Higgs transverse momentum spectrum



total fiducial cross section



G. Billis, BD, M A. Ebert, J. K.L. Michel, F. J. Tackmann (2021)

Summary

- ✓ **Effective Field Theories** are very **powerful tools** in **precision physics**
- ✓ **Effective Field Theories** provide **coherent description** of a system using only the **relevant degrees of freedom**
- ✓ **Effective Field Theories** have often **more predictive powers** which can be **systematically improved**
- ✓ **Effective Field Theories** equip us with **factorization theorems** which allow us to systematically understand and control the physics at lower energies (and non-perturbative dynamics)

Acknowledgment

Thank you all,
for the nice
memories!

