

SUSY and Electroweak Precision Tests at Tera-Z Colliders

Kevin Langhoff

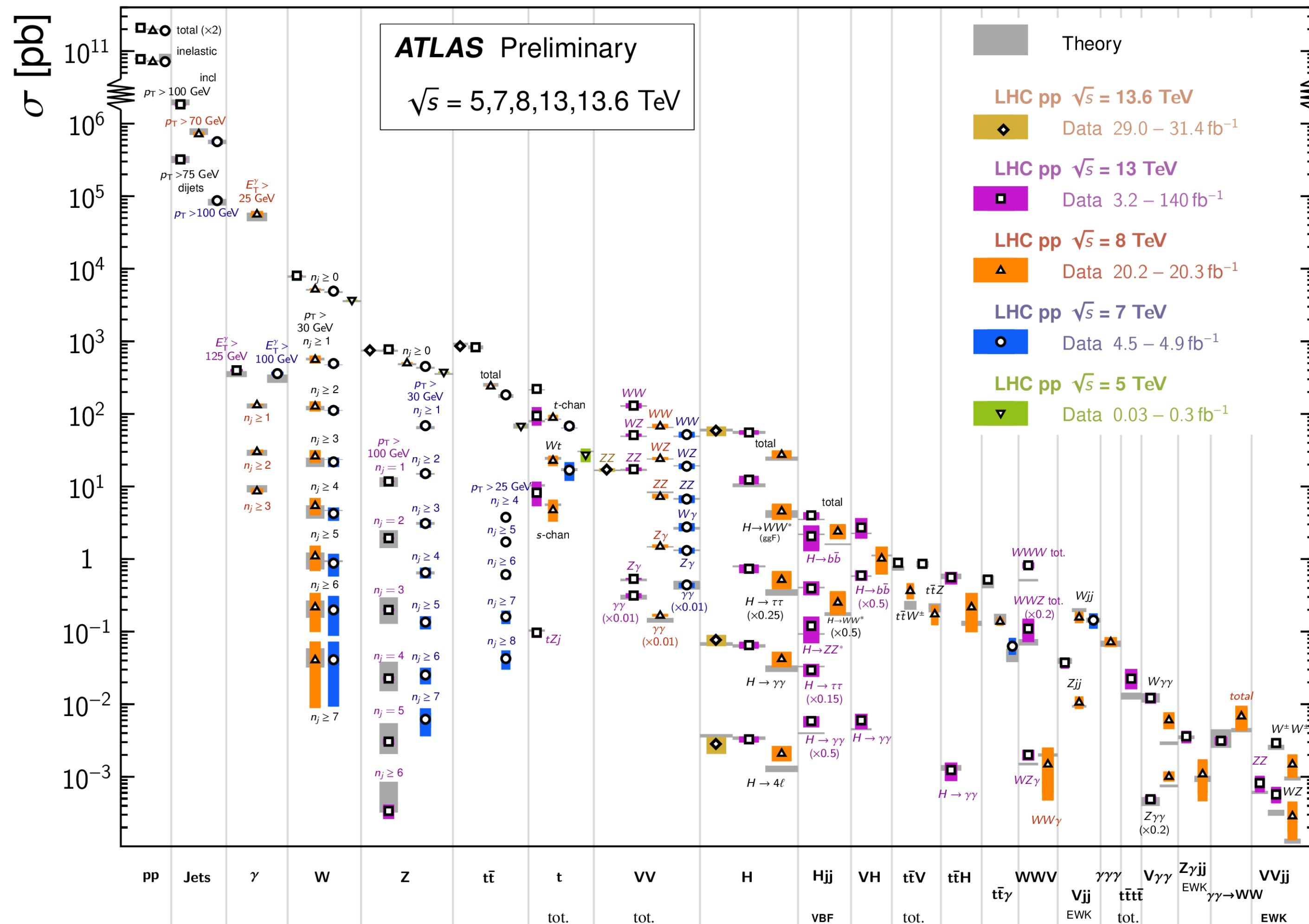
(Work in progress with Simon Knapen and Zoltan Ligeti)

FCC-Week (June 11th, 2024)

The Standard Model is extremely predictive!

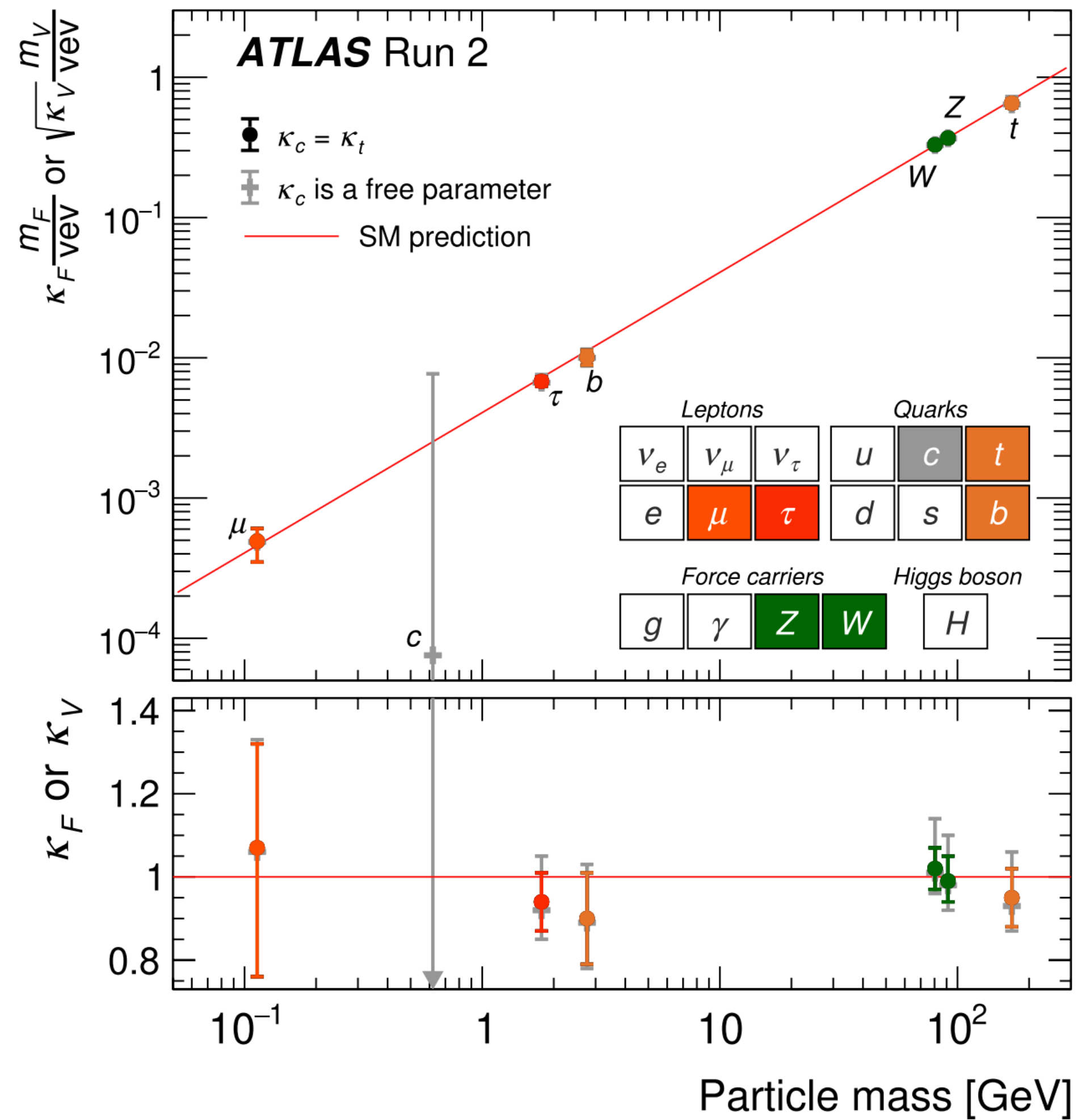
Standard Model Production Cross Section Measurements

Status: October 2023



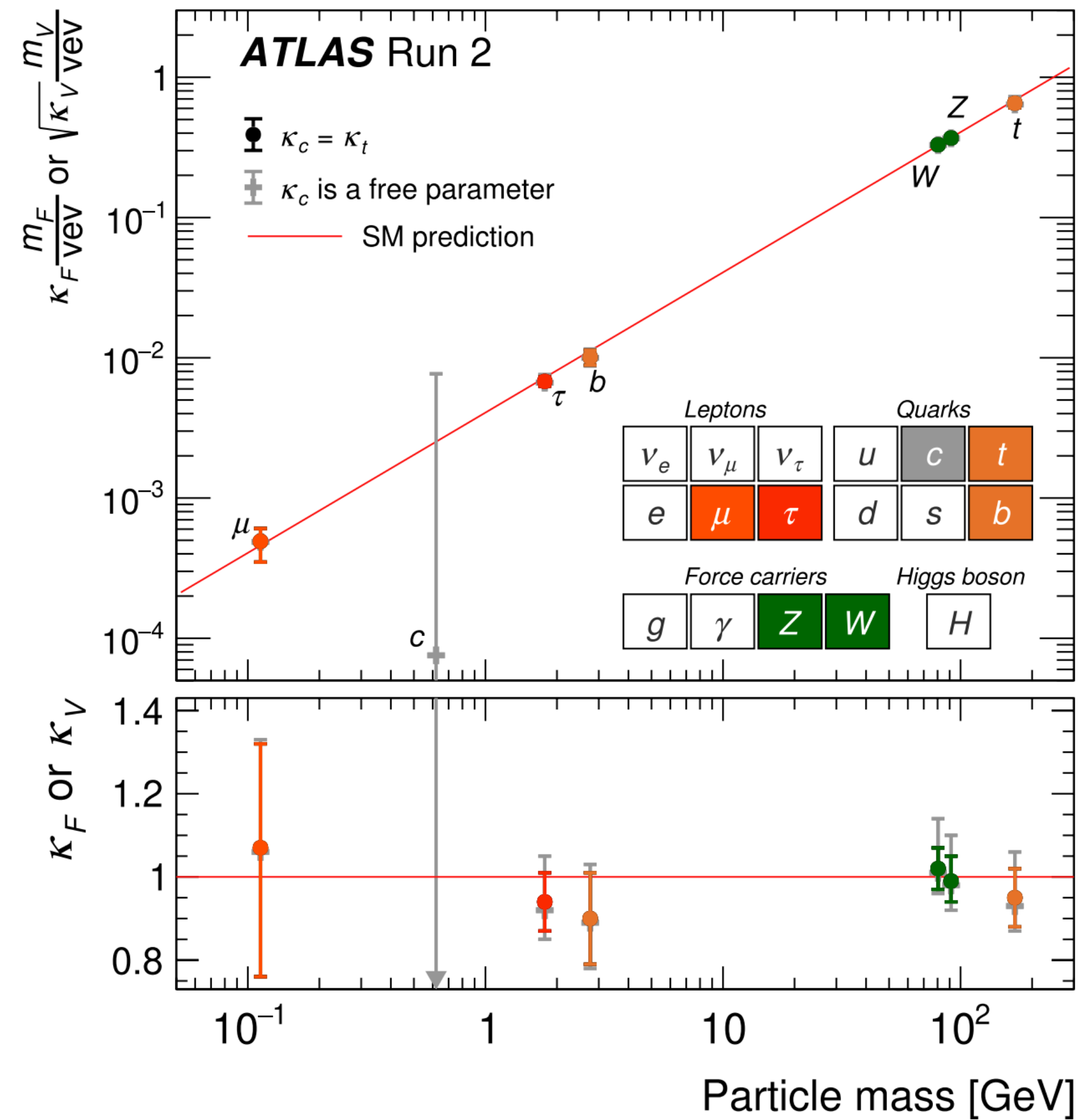
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Higgs couplings

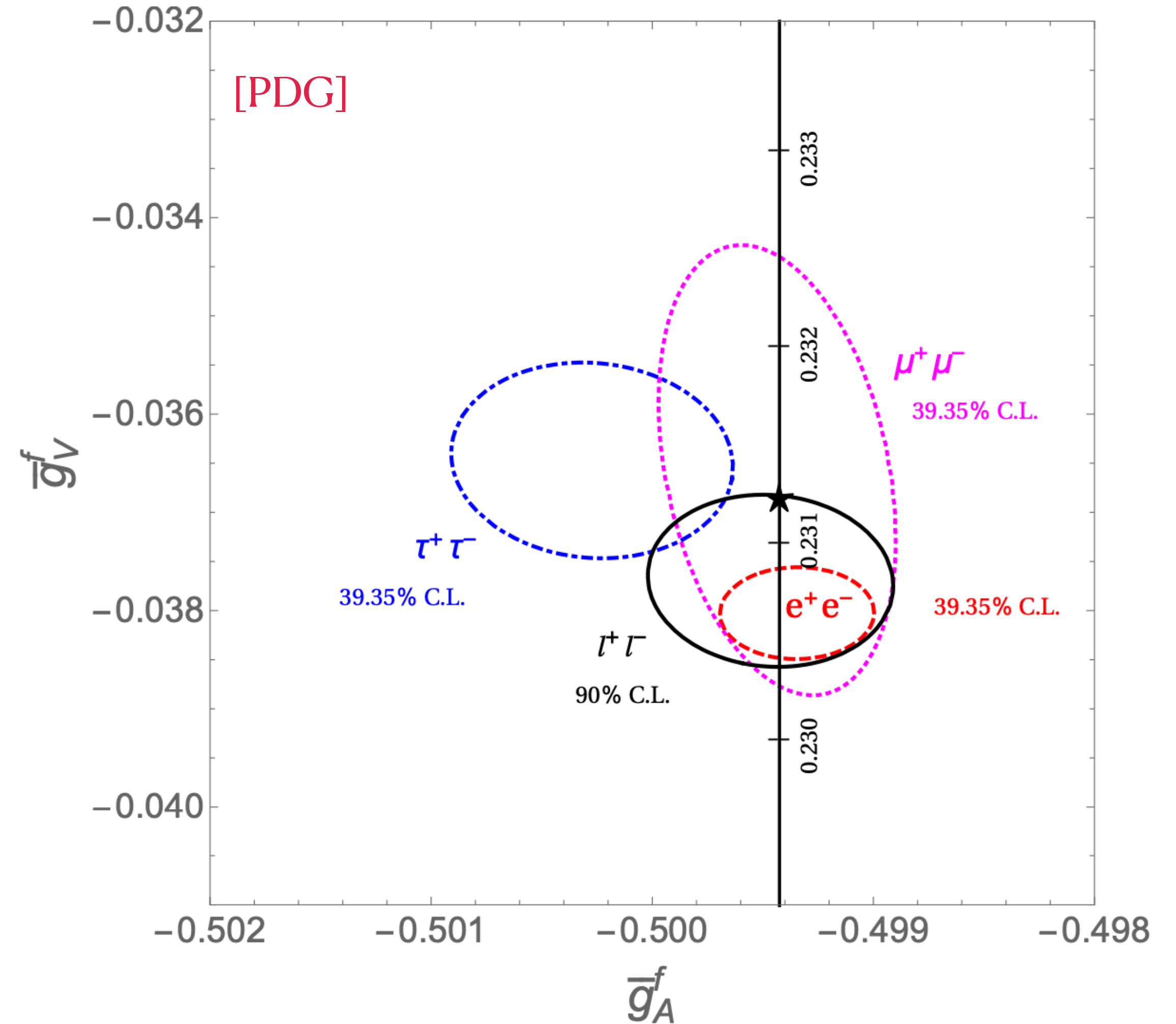


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Z boson couplings



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Observables

1.	$G_F = \left(\sqrt{2} v^2 \right)^{-1}$	3.	$m_Z = \frac{1}{2} \sqrt{g^2 + g'^2} v$
2.	$m_W = \frac{1}{2} g v$	4.	$\sin \theta_W = \frac{g'}{\sqrt{g^2 + g'^2}}$

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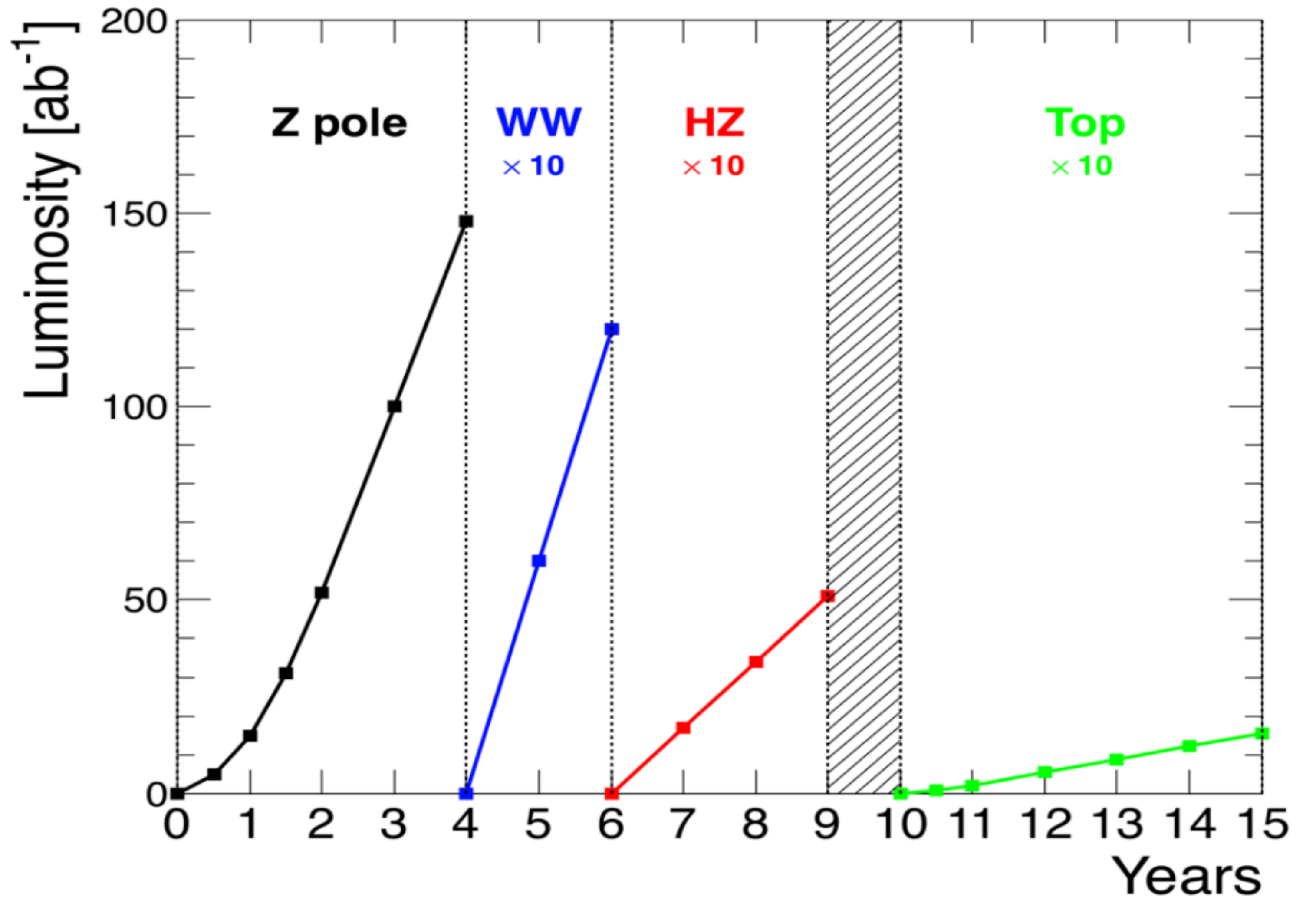


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- Checks like this give us a method of indirectly discovering new physics!

The FCC-ee is an incredible precision device!

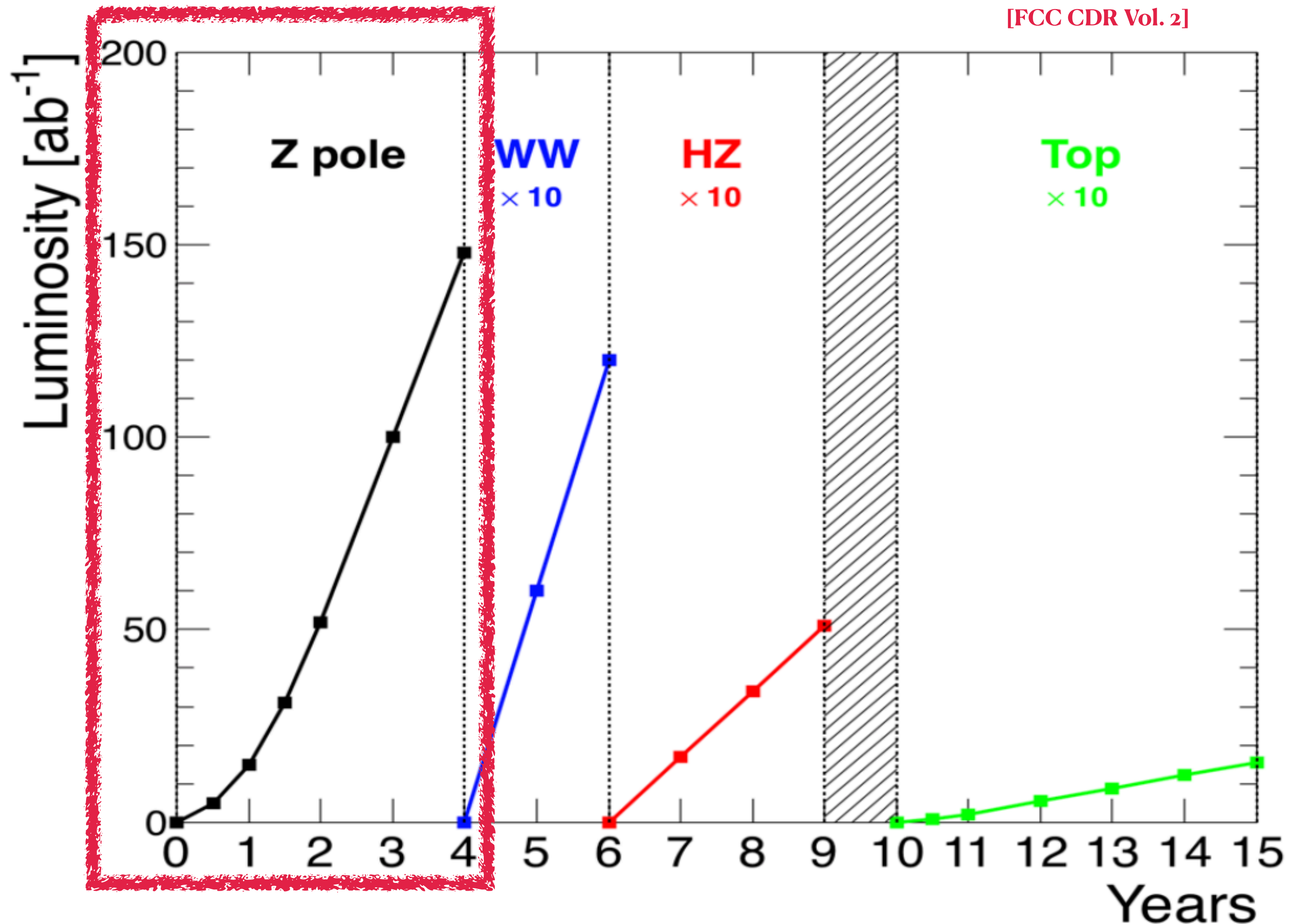
[FCC CDR Vol. 2]



The FCC-ee is an incredible precision device!

The FCC-ee will produce roughly 5×10^{12} Z-bosons!

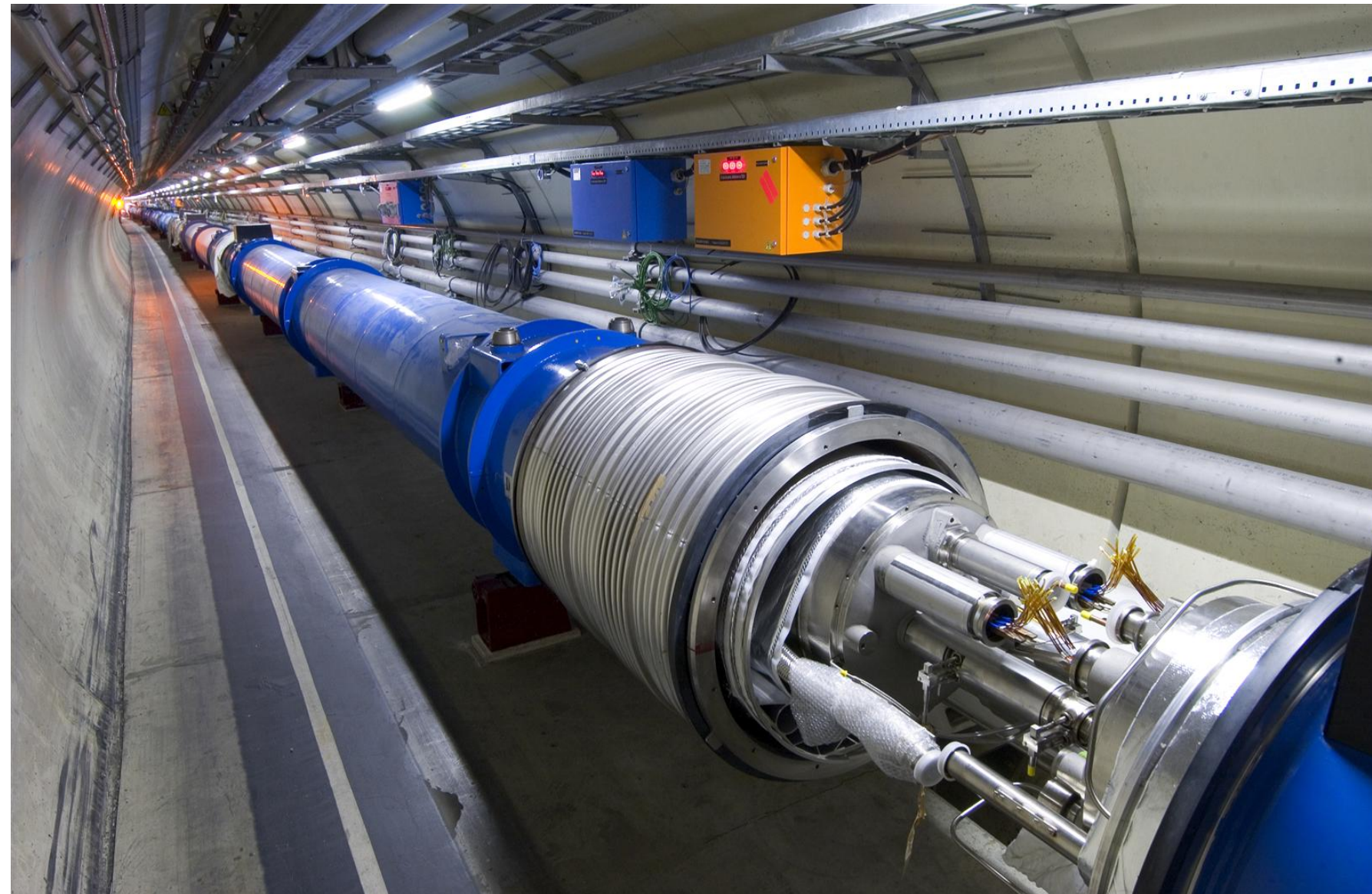
(Roughly a factor of 10^5 more than produced at LEP)



What experiment explores the highest energy scales?

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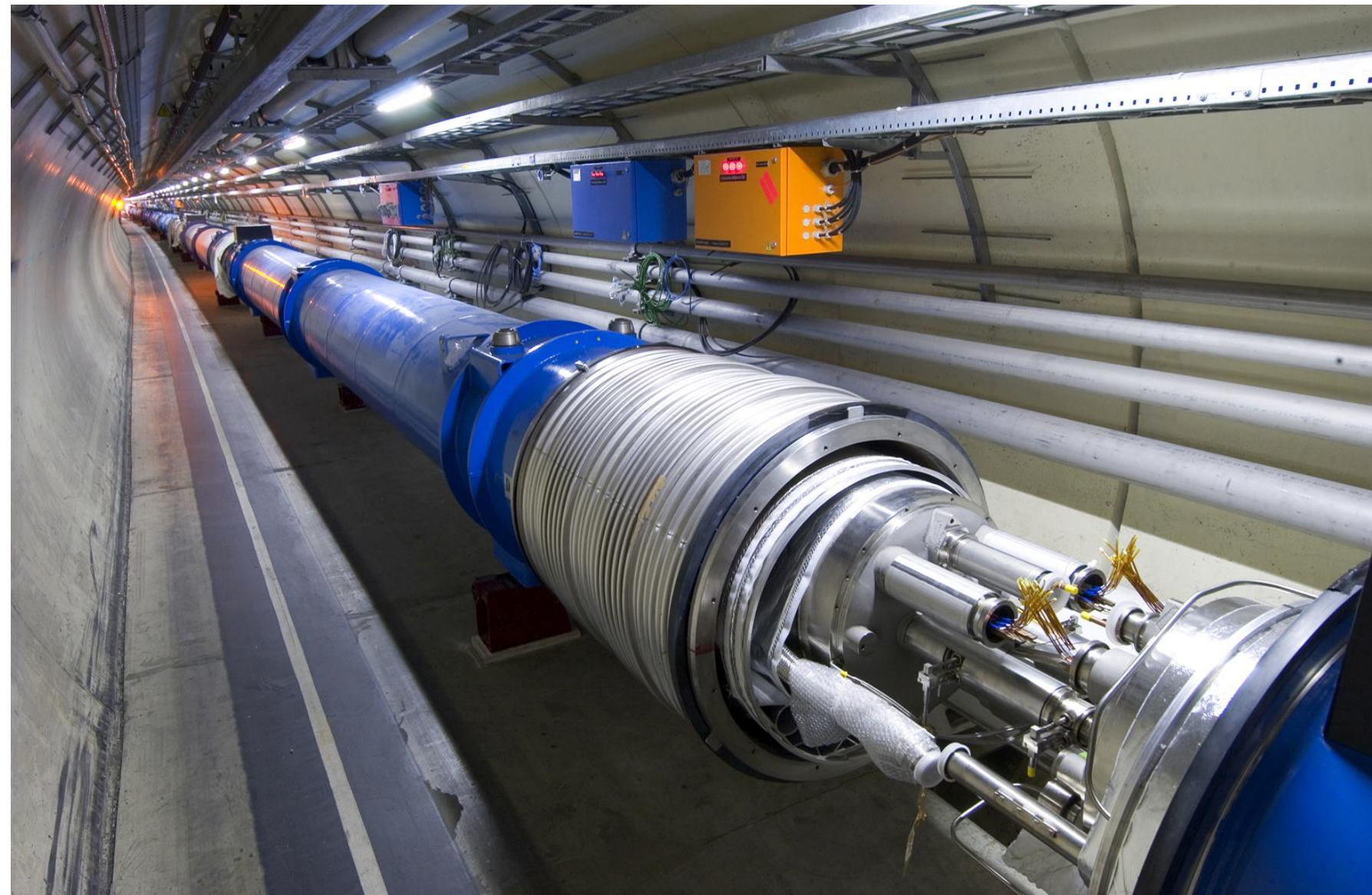
LHC?



Directly explores energy scales $\Lambda \sim 10^3$ GeV.

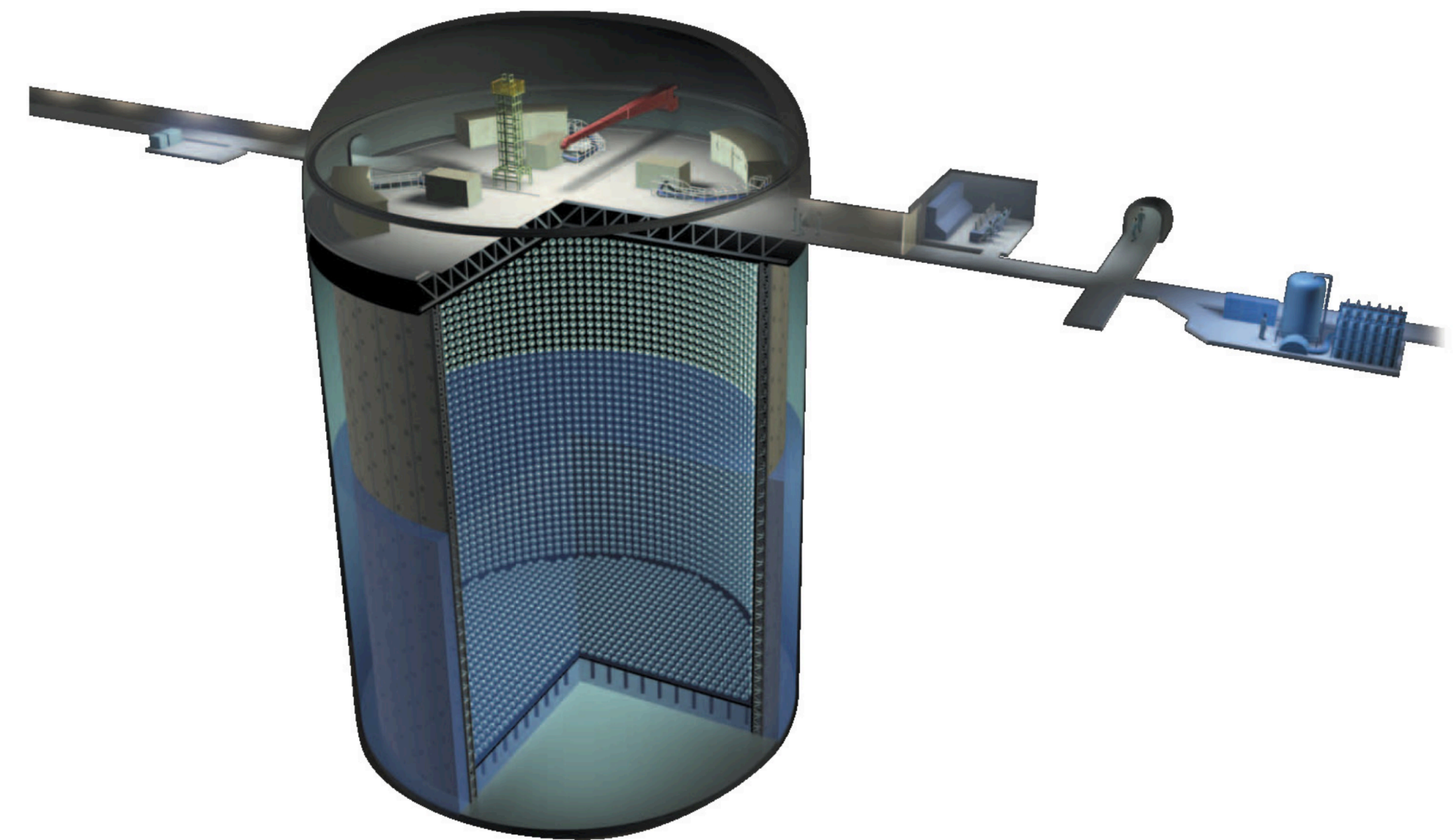
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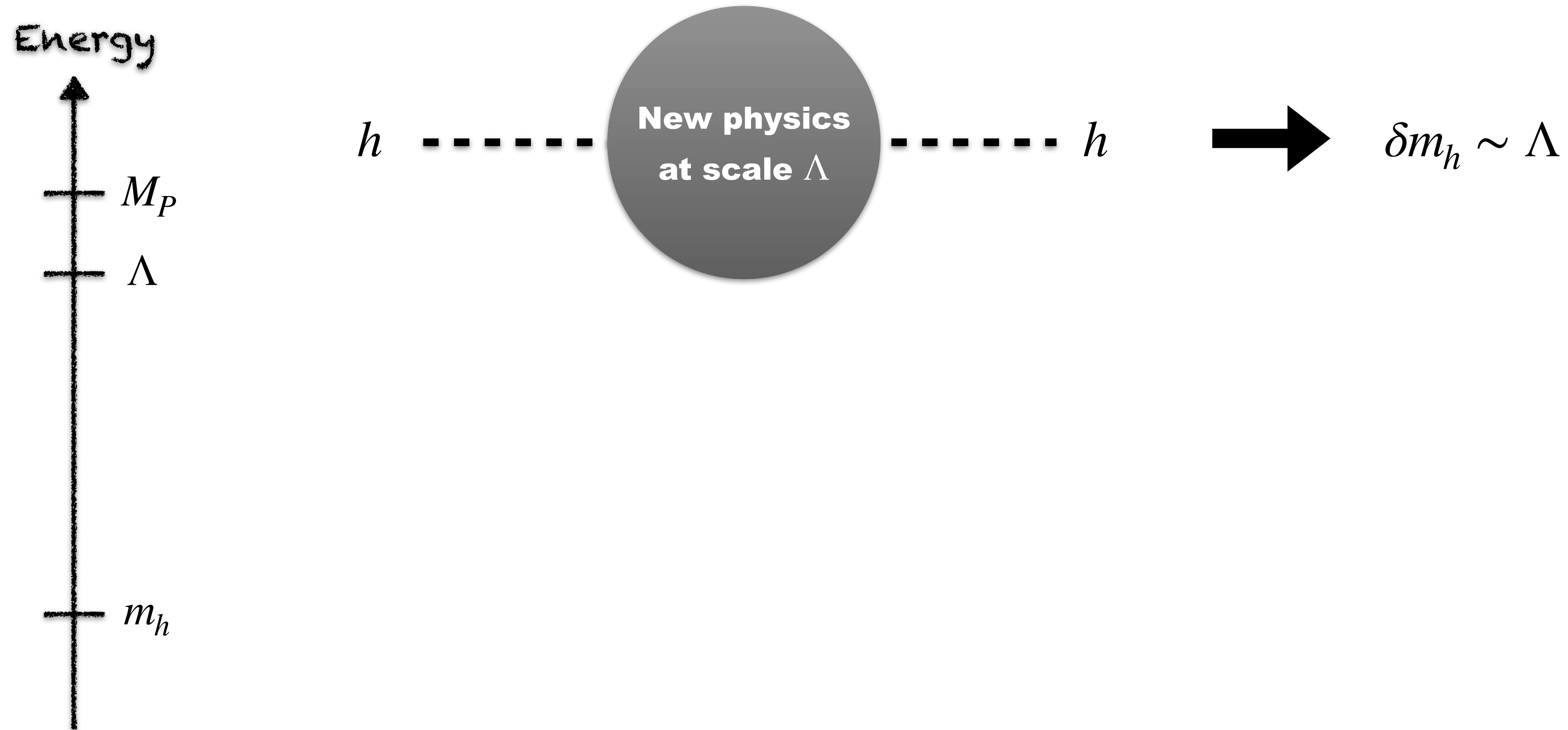
Super-Kamiokande



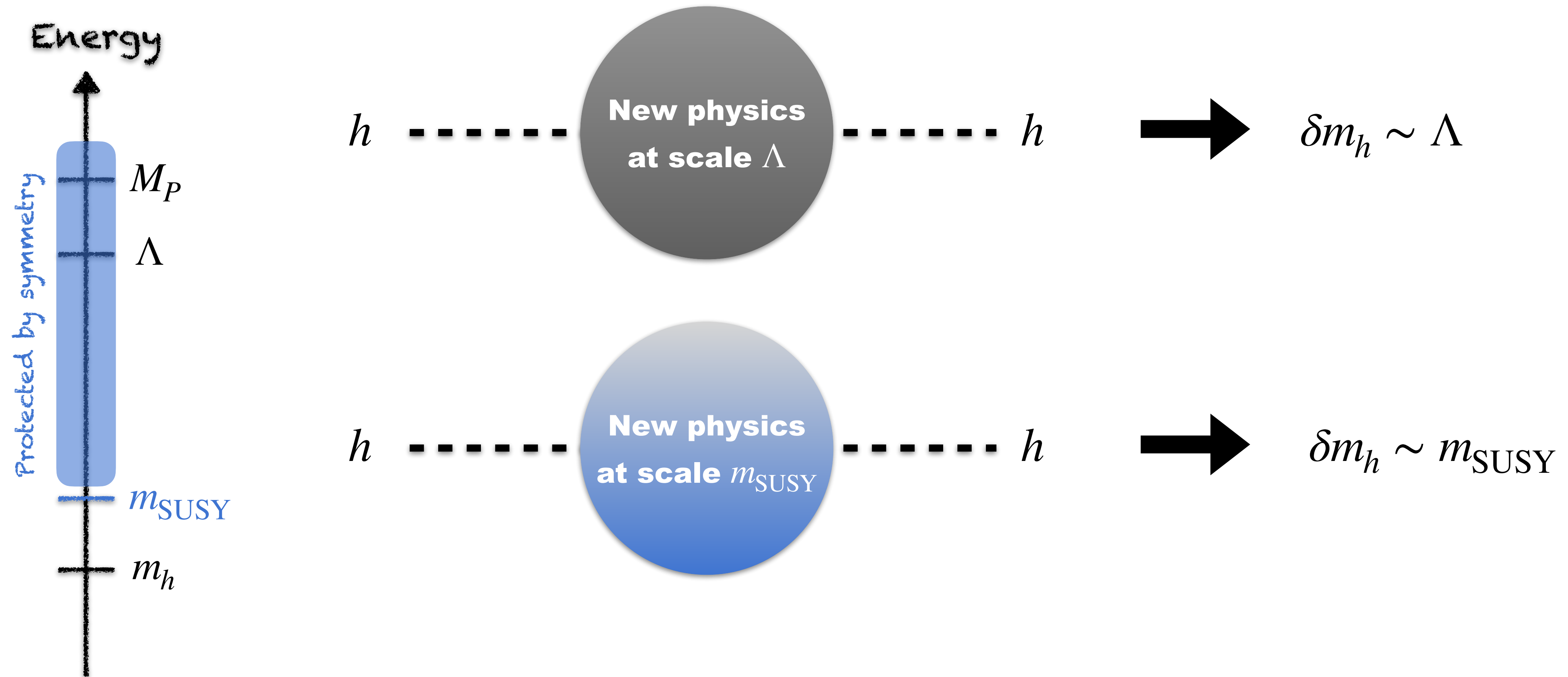
Using 10^{34} protons, indirectly explores baryon violating Dim-6 operators at scales $\Lambda \sim 10^{16}$ GeV.

What mysteries can we explore using these methods?

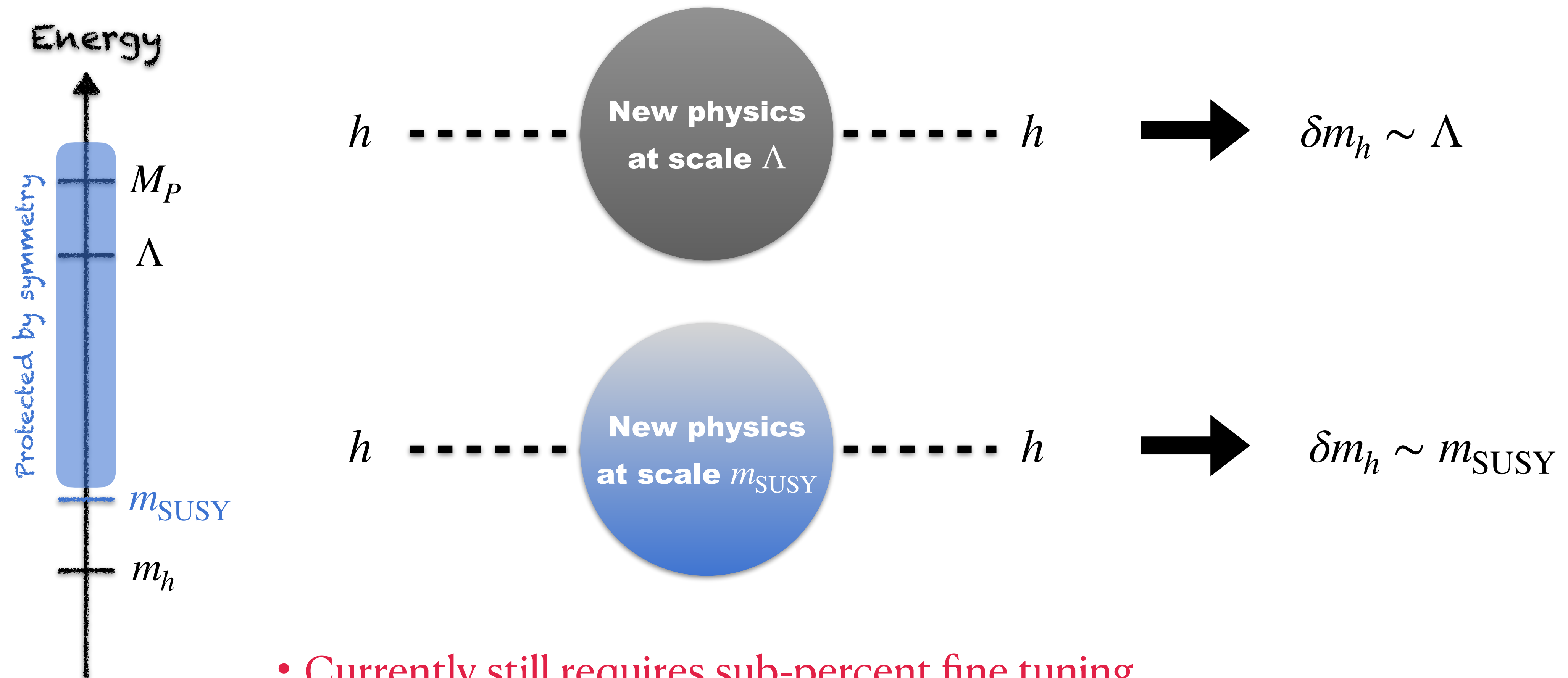
The Higgs Hierarchy Problem



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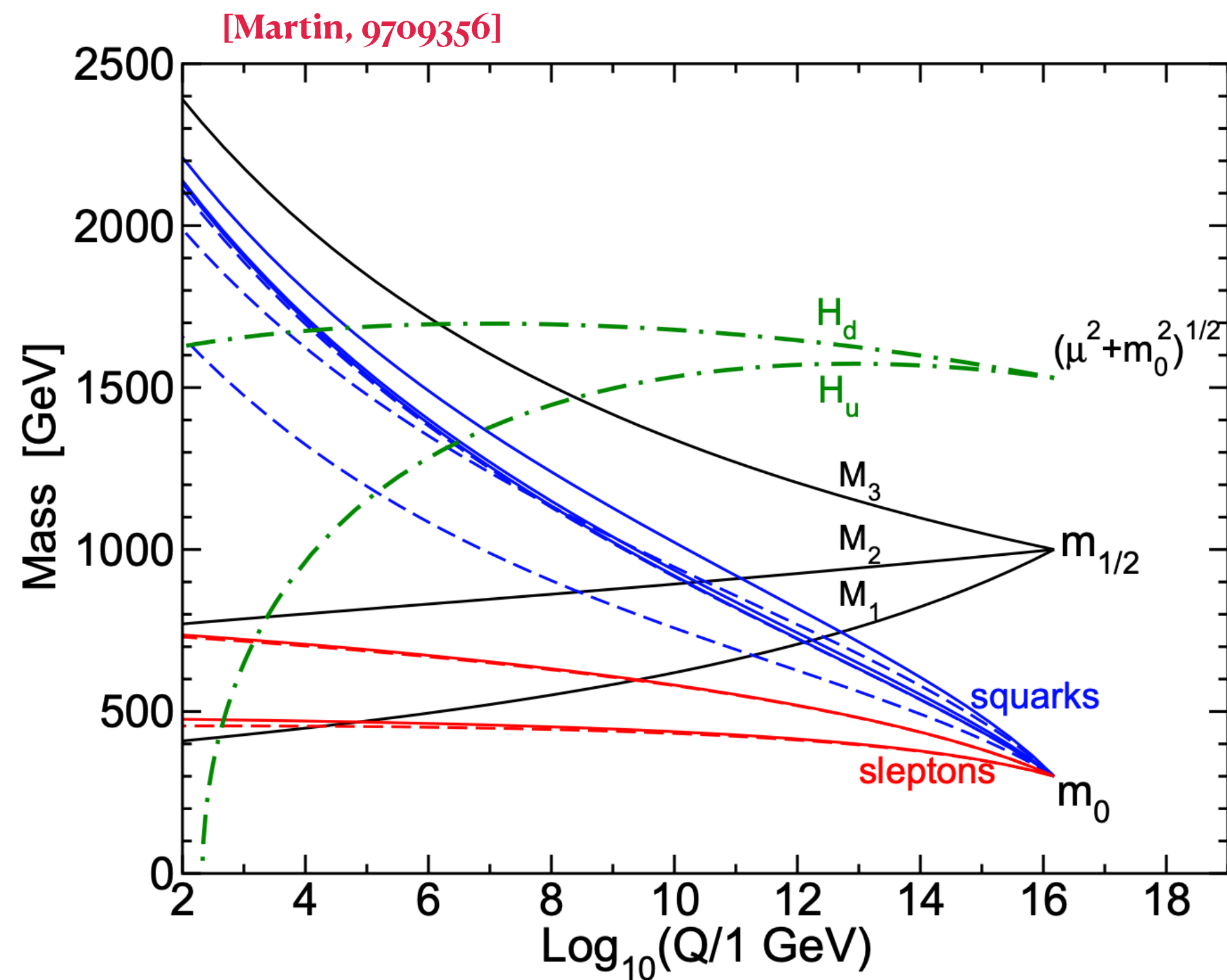


The Higgs Hierarchy Problem



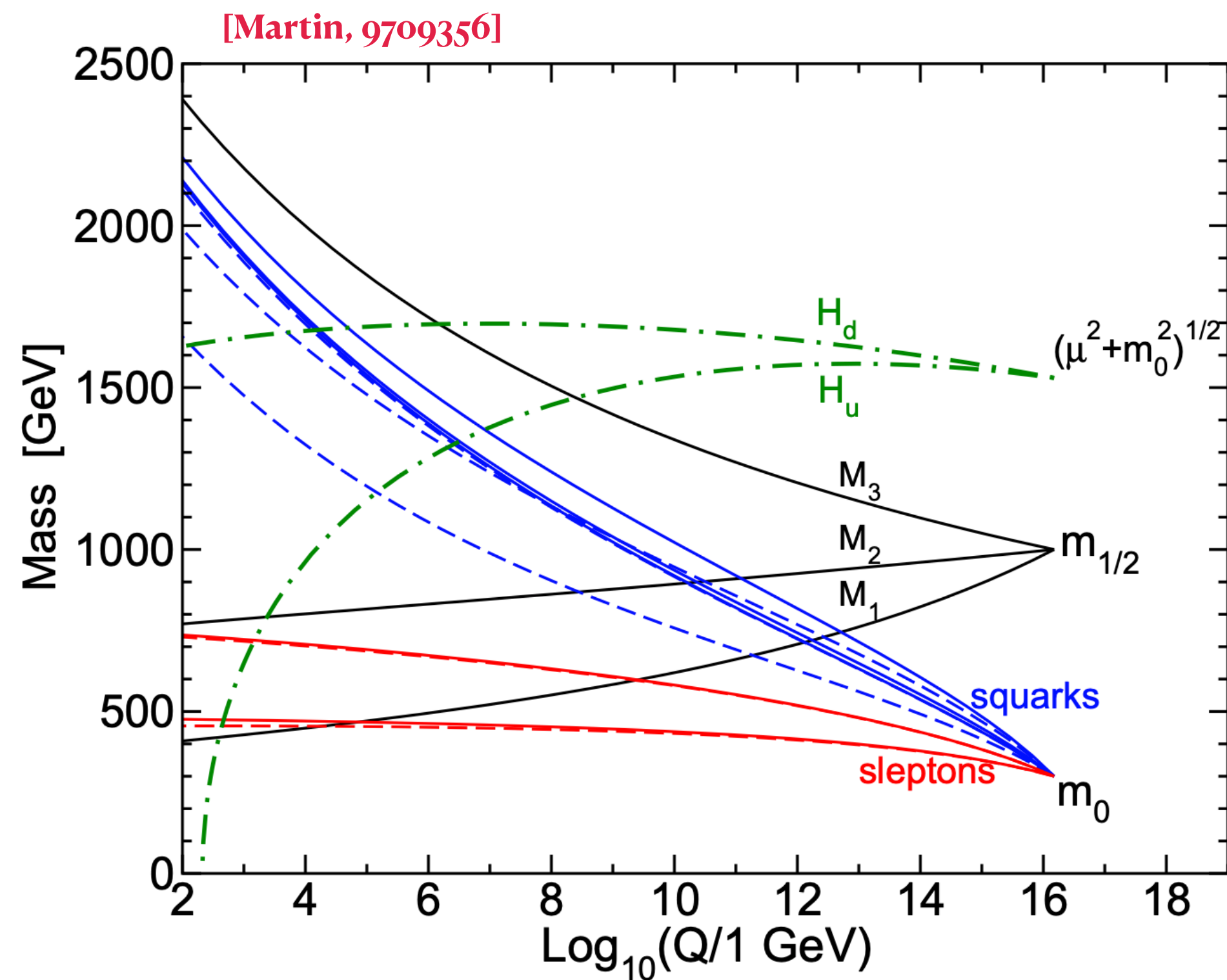
- Currently still requires sub-percent fine tuning.
- Strongest constraints are on colored sparticles (e.g. gluinos and squarks).

Can the FCC-ee See What The LHC Can't?



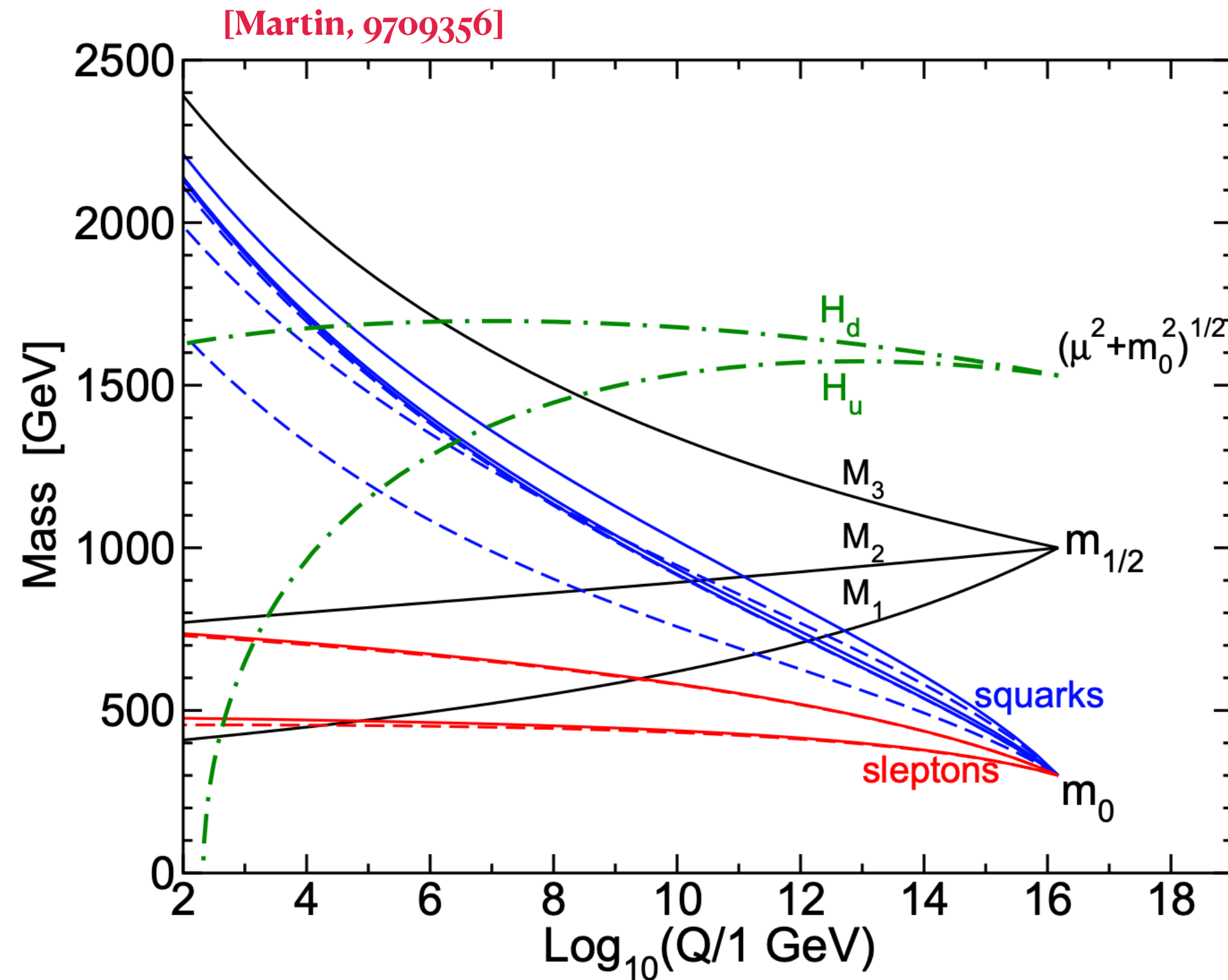
- Running motivates $m_{\text{colored}} > m_{\text{uncolored}}$.

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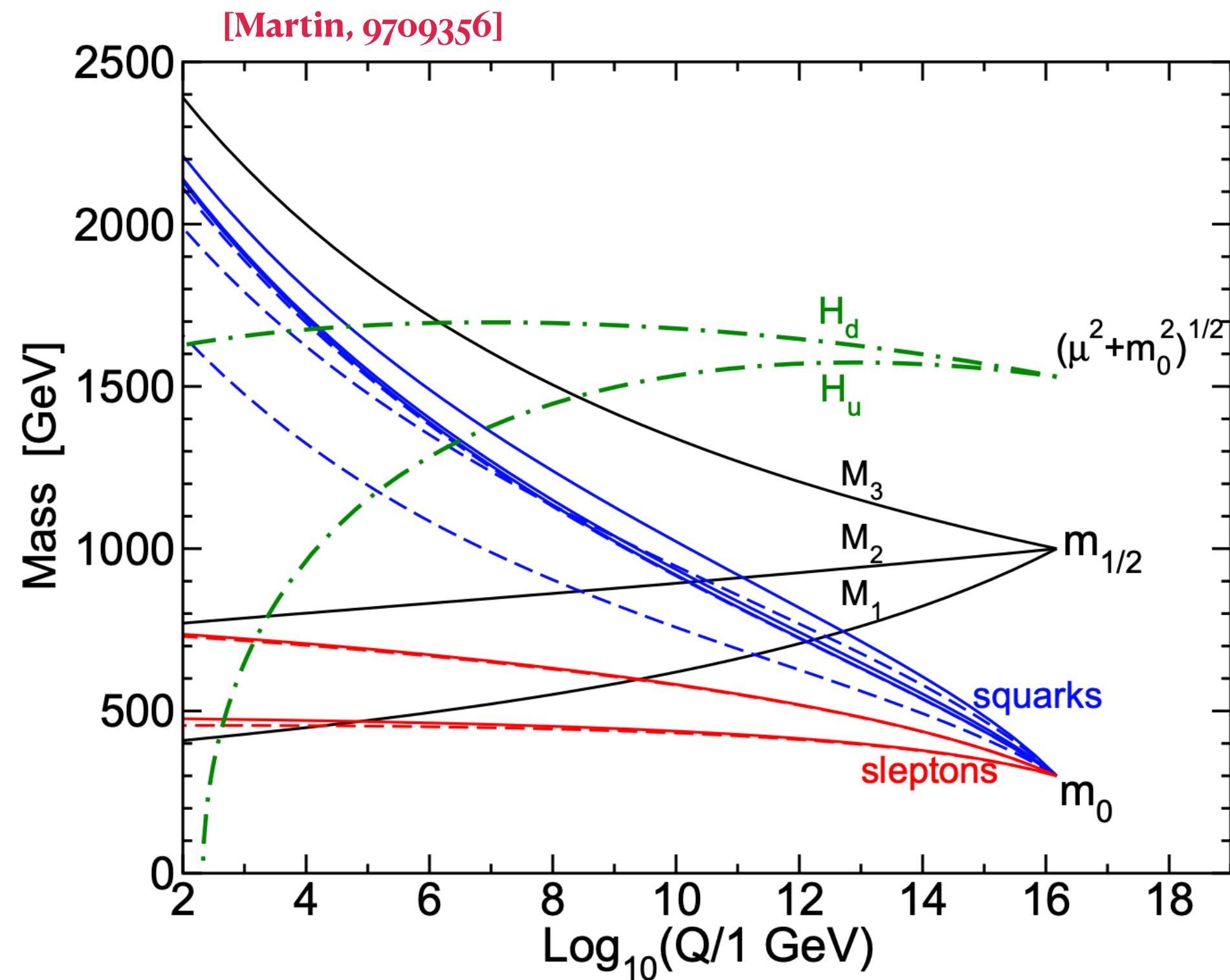
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- Dominant effect on EWPTs may come from color neutral sparticles.

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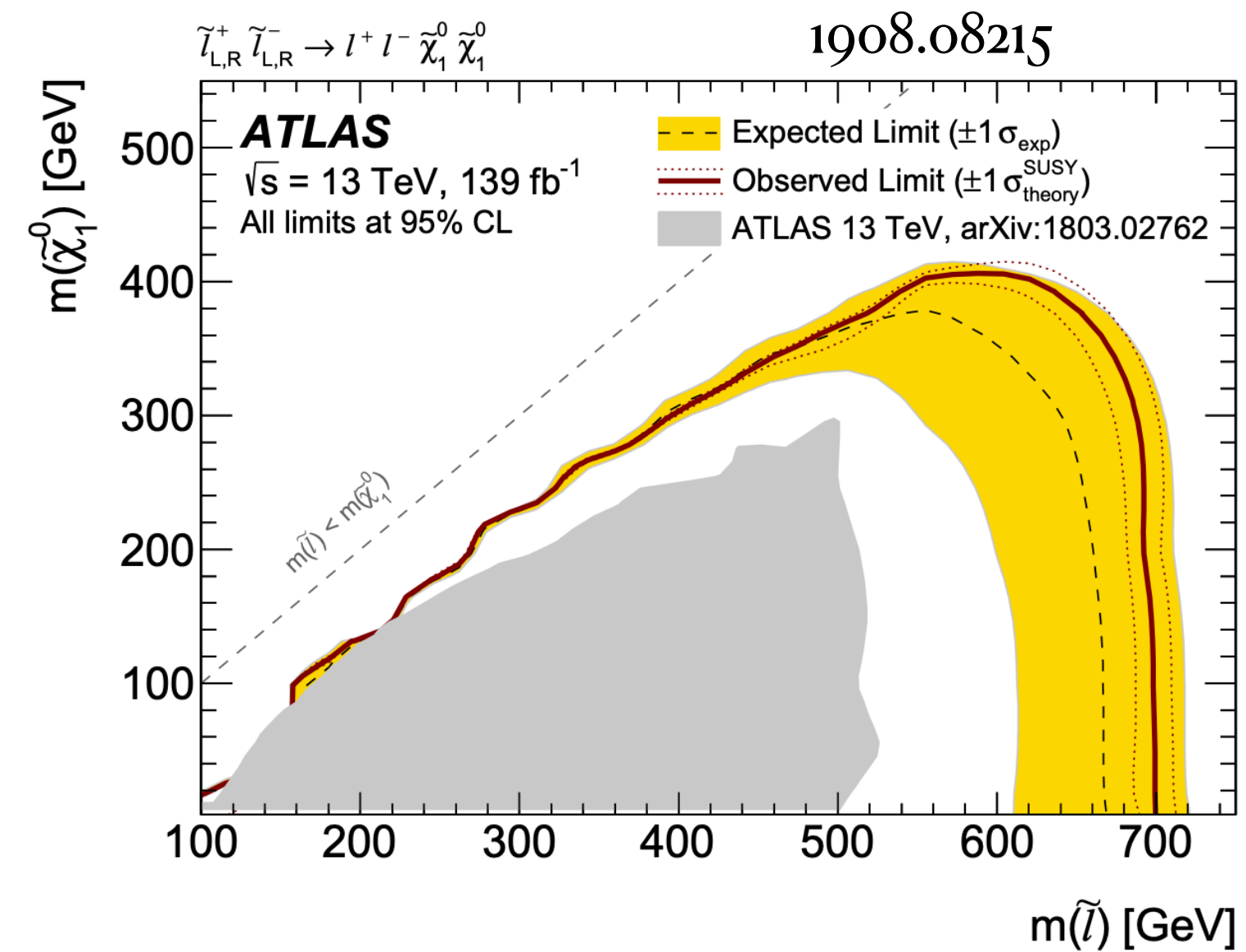
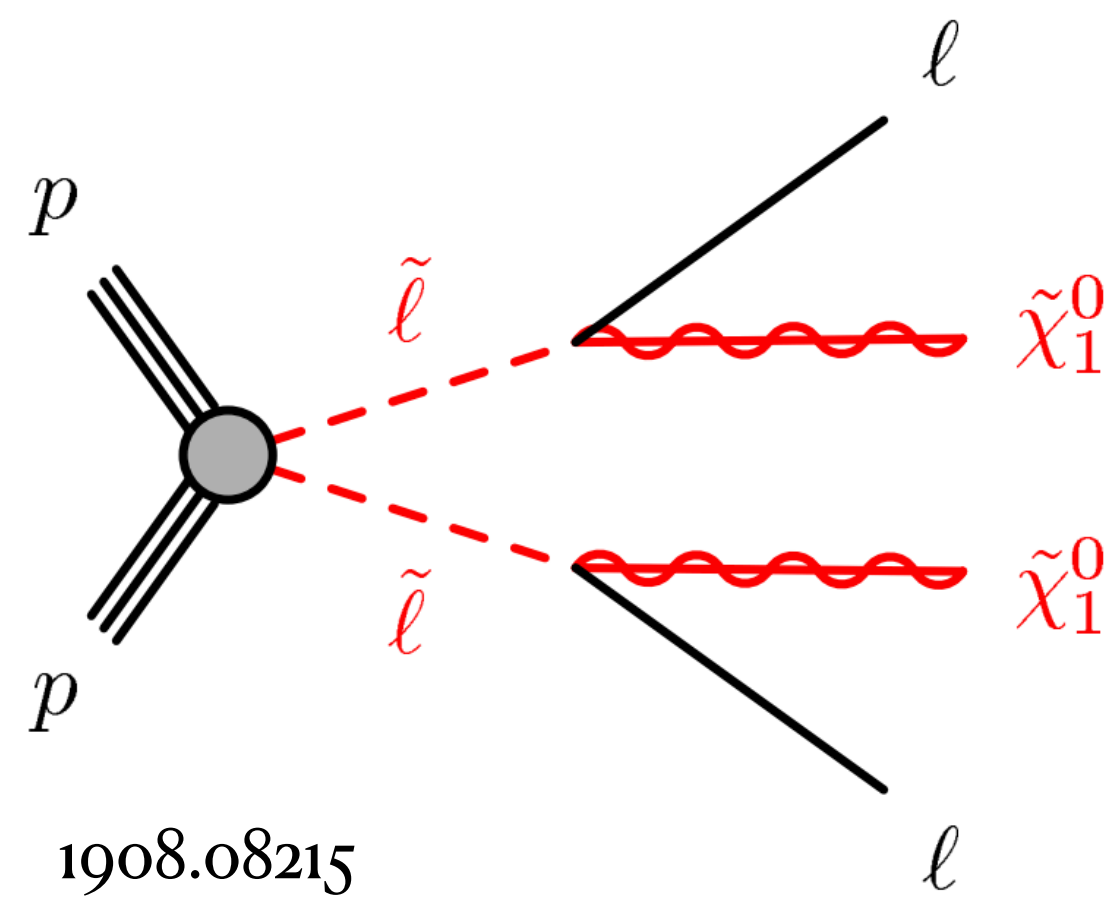


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⇒ FCC-ee may indirectly see SUSY using EWPTs even if the LHC sees nothing.

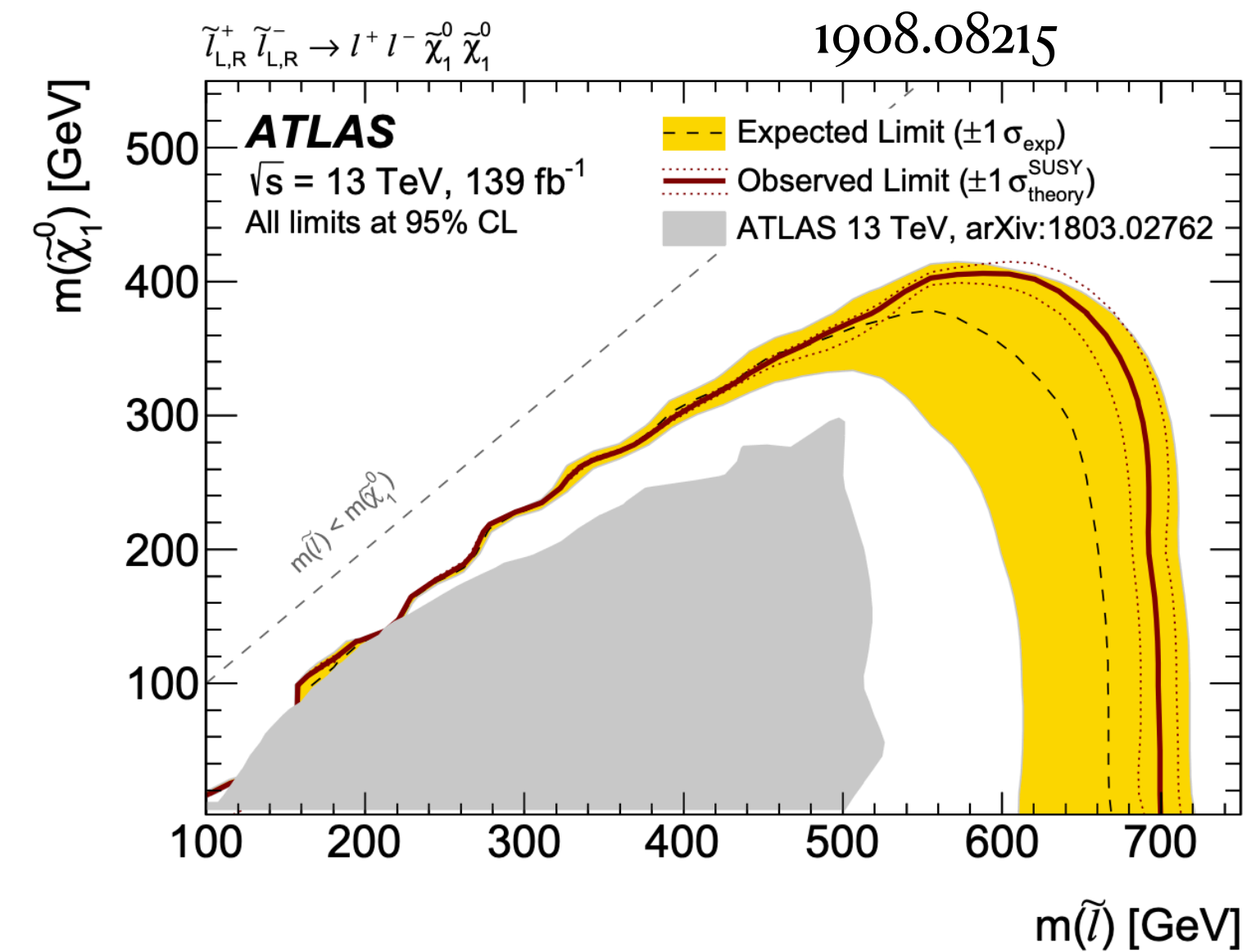
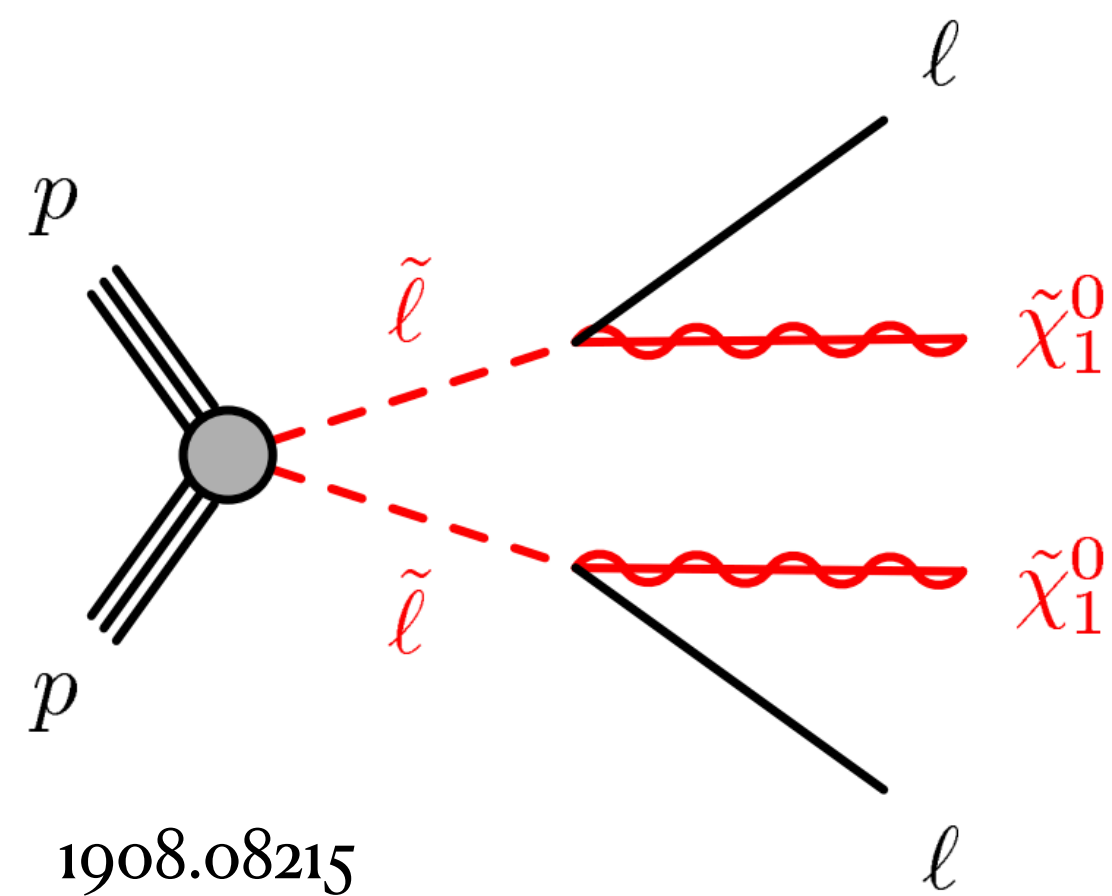
Simplified SUSY Models

LHC SUSY searches often consider representative simplified models.



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For simplicity we will do the same:

1. Pure Bino + RH Slepton model $\supset (\tilde{B}, \tilde{e})$
2. Pure Wino + LH Slepton model $\supset (\tilde{W}, \tilde{L})$.

How might SUSY show up in EWPTs?

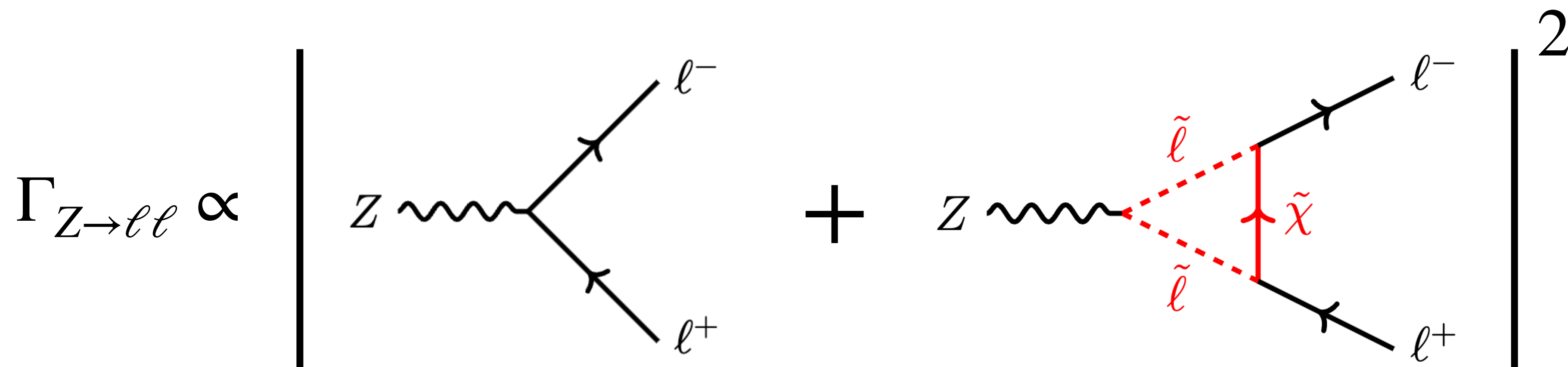
Let $\tilde{\chi} = (\tilde{W}, \tilde{B})$ and $\tilde{\ell} = (\tilde{L}, \tilde{e})$.

$$\Gamma_{Z \rightarrow \ell\ell} \propto \left| \begin{array}{c} Z \text{ wavy line} \\ \swarrow \nearrow \\ l^- \\ l^+ \end{array} \right. + \left. \begin{array}{c} Z \text{ wavy line} \\ \swarrow \nearrow \\ \tilde{\ell} \text{ (dashed)} \\ \tilde{\chi} \text{ (solid)} \\ \tilde{\ell} \text{ (dashed)} \\ \swarrow \nearrow \\ l^- \\ l^+ \end{array} \right|^2$$

(Just one of several diagrams)

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$$\frac{\Gamma_{Z \rightarrow \ell\ell} - \Gamma_{Z \rightarrow \ell\ell}^{(SM)}}{\Gamma_{Z \rightarrow \ell\ell}^{(SM)}} \propto \frac{g^2}{16\pi^2} \left(\frac{m_Z}{M_{SUSY}} \right)^2 \longrightarrow M_{SUSY}^{\text{probed}} \sim 1 \text{ TeV} \times \left(\frac{\delta\Gamma/\Gamma}{10^{-5}} \right)$$

Some details

$\Gamma(Z \rightarrow \ell \bar{\ell})$ is not the best observable. Instead we use

$$R_\ell \equiv \frac{\Gamma(Z \rightarrow \text{hadrons})}{\Gamma(Z \rightarrow \ell \bar{\ell})}$$

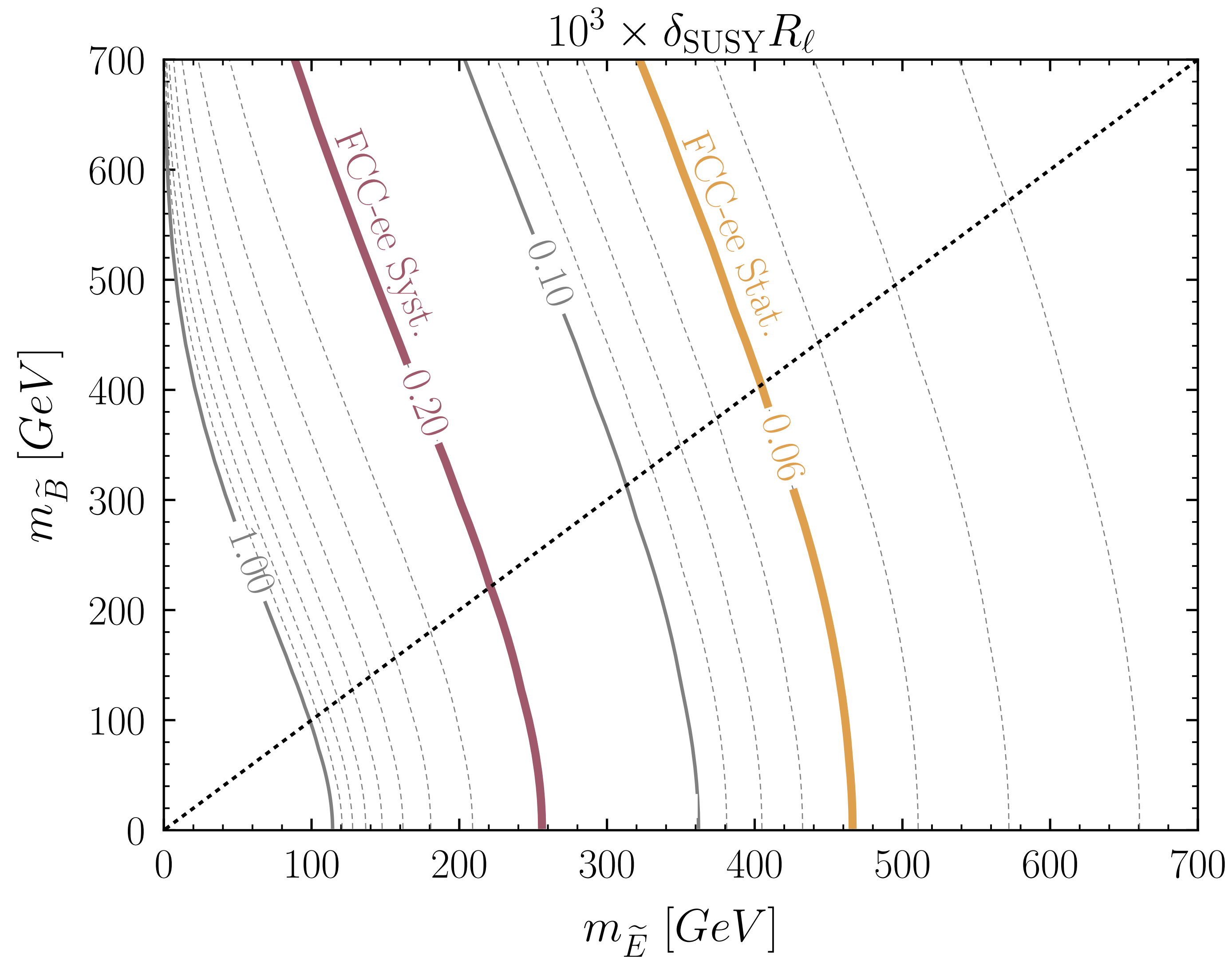
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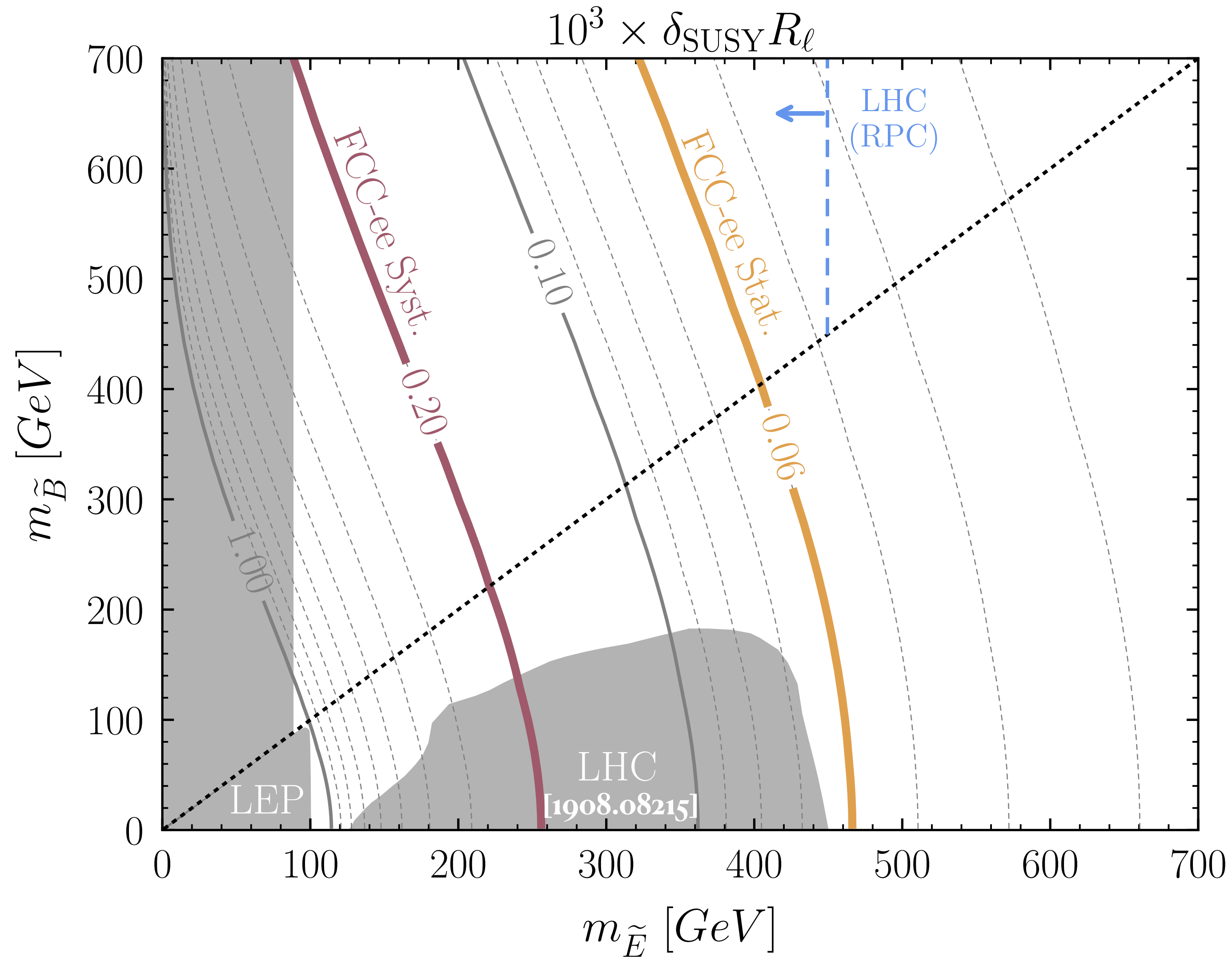
Still depends on θ_W and $\alpha_s(M_Z)$. These must be determined by other measurements.

Results



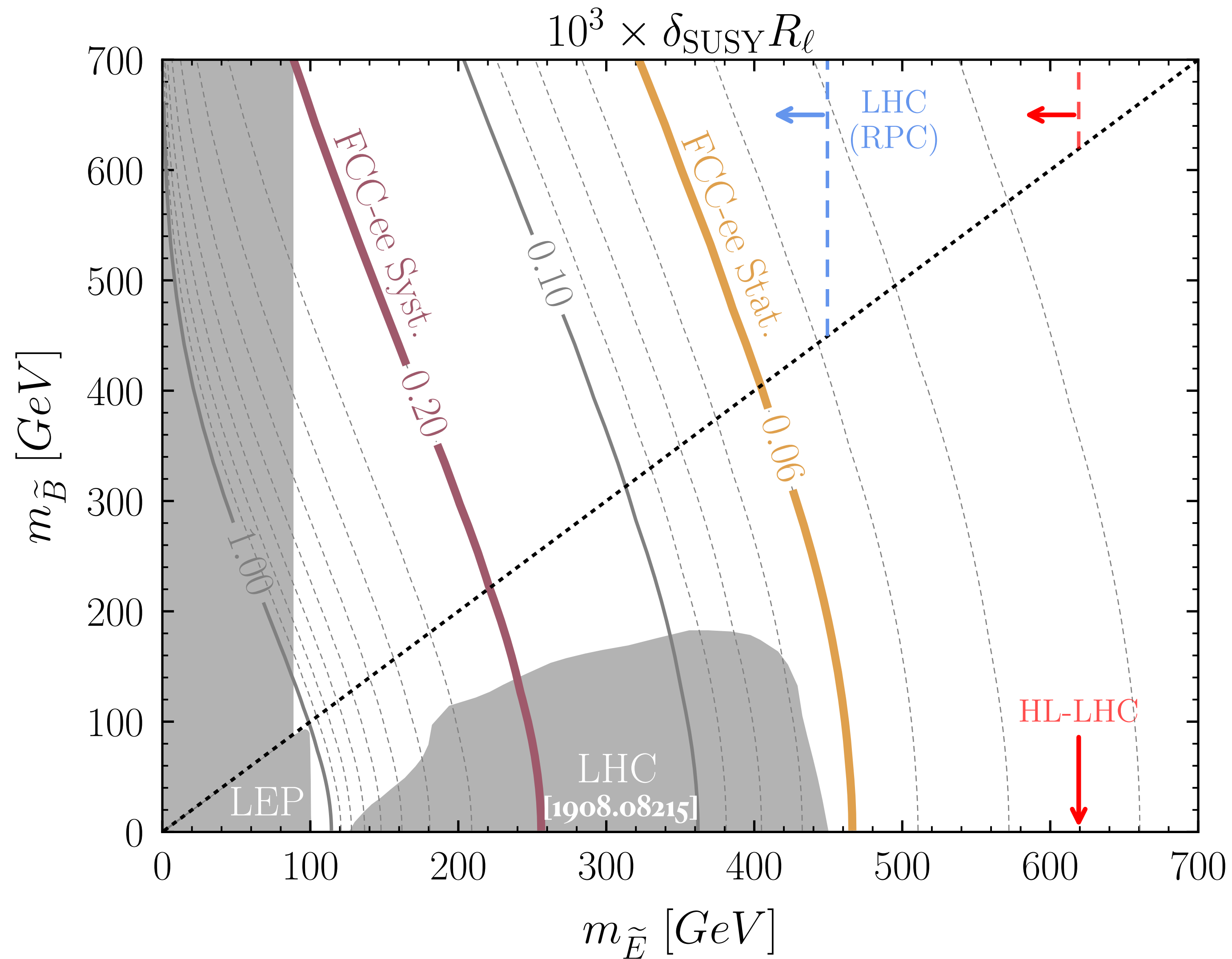
**Bino + RH Slepton
(Preliminary)**

Results



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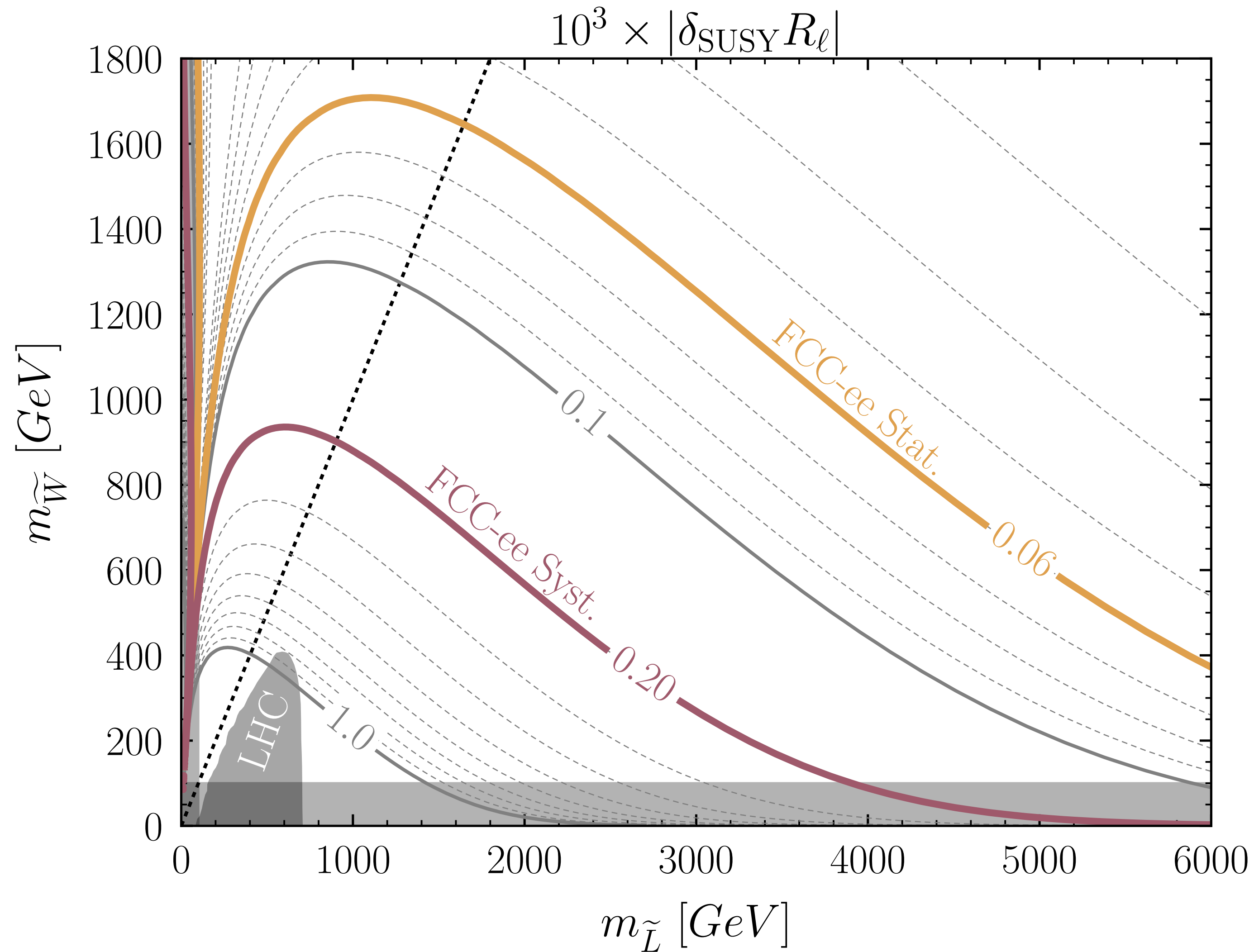
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$$\cos(2\beta) = 0$$

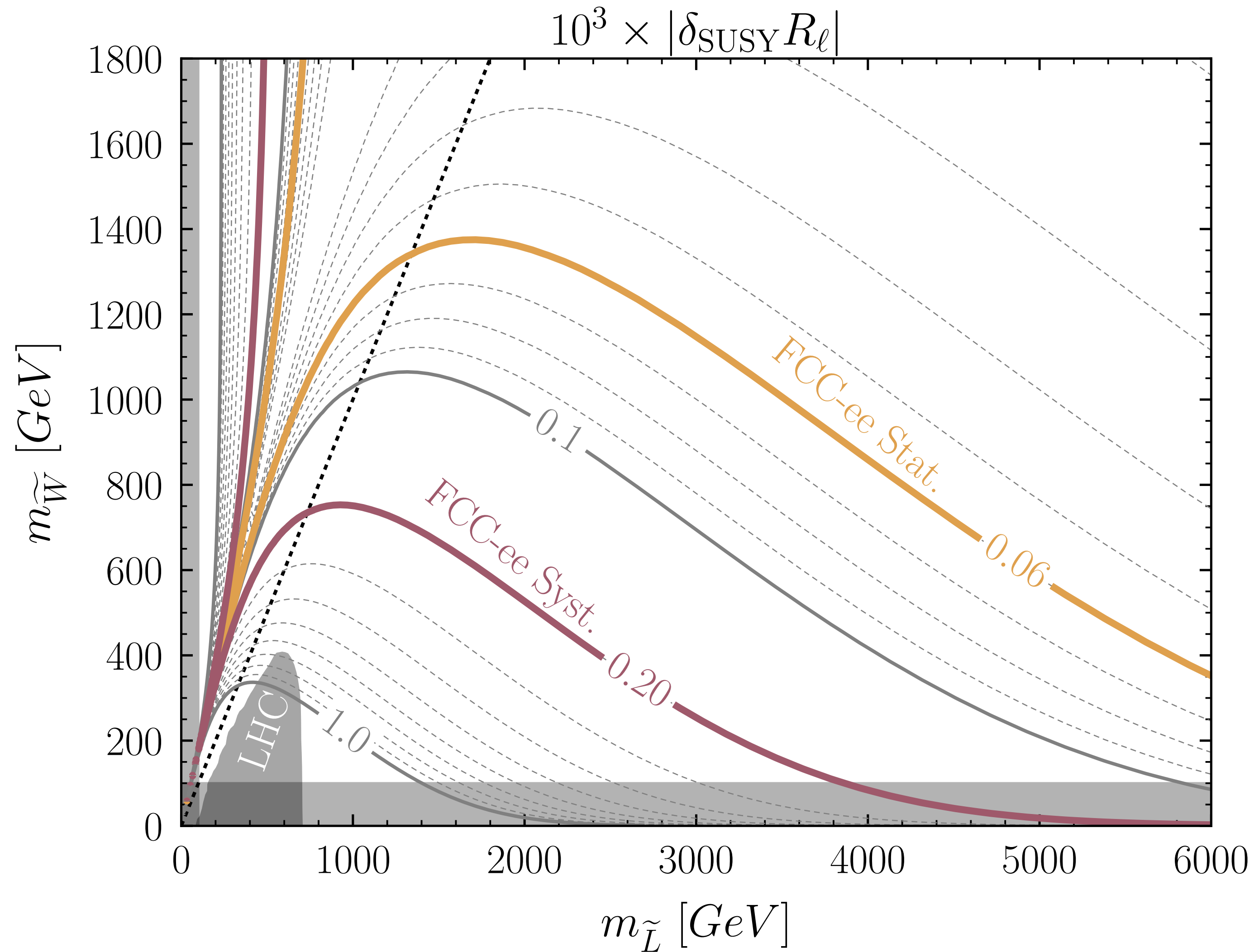


**Wino + LH Slepton
(Preliminary)**

[1908.08215]

Results

$$\cos(2\beta) = -1$$

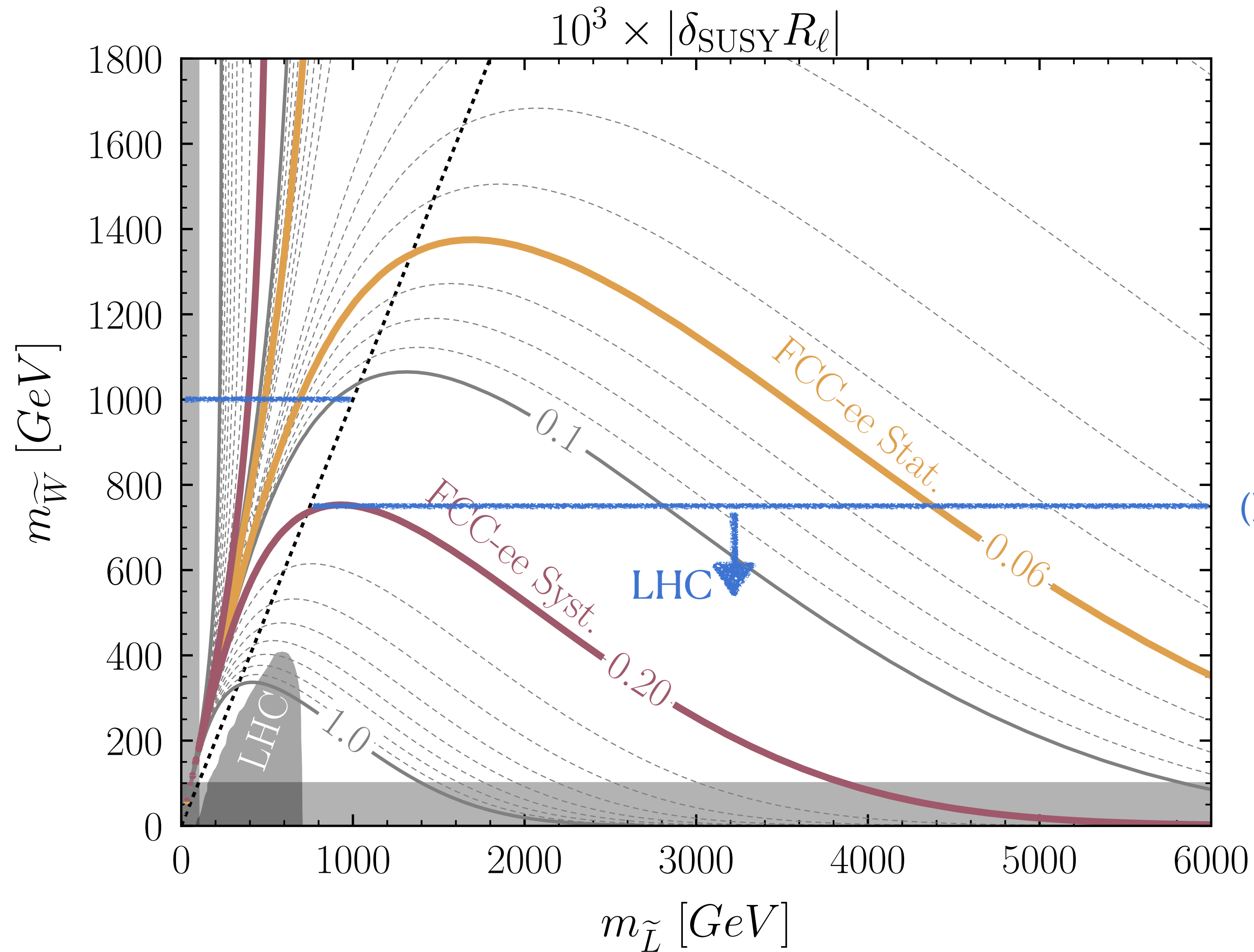


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[1908.08215]

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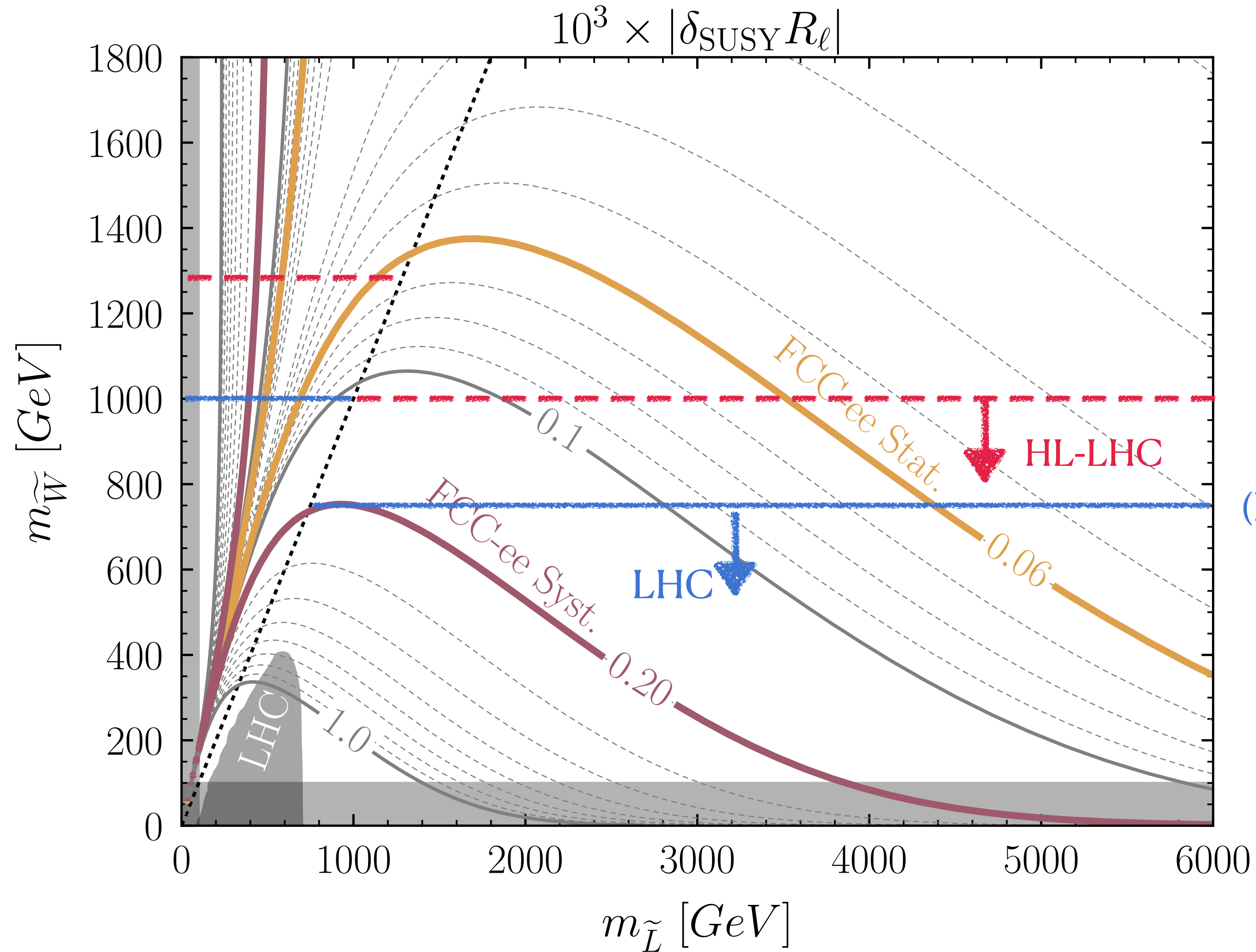
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(Depends on $m_{\chi_1^0}$!)

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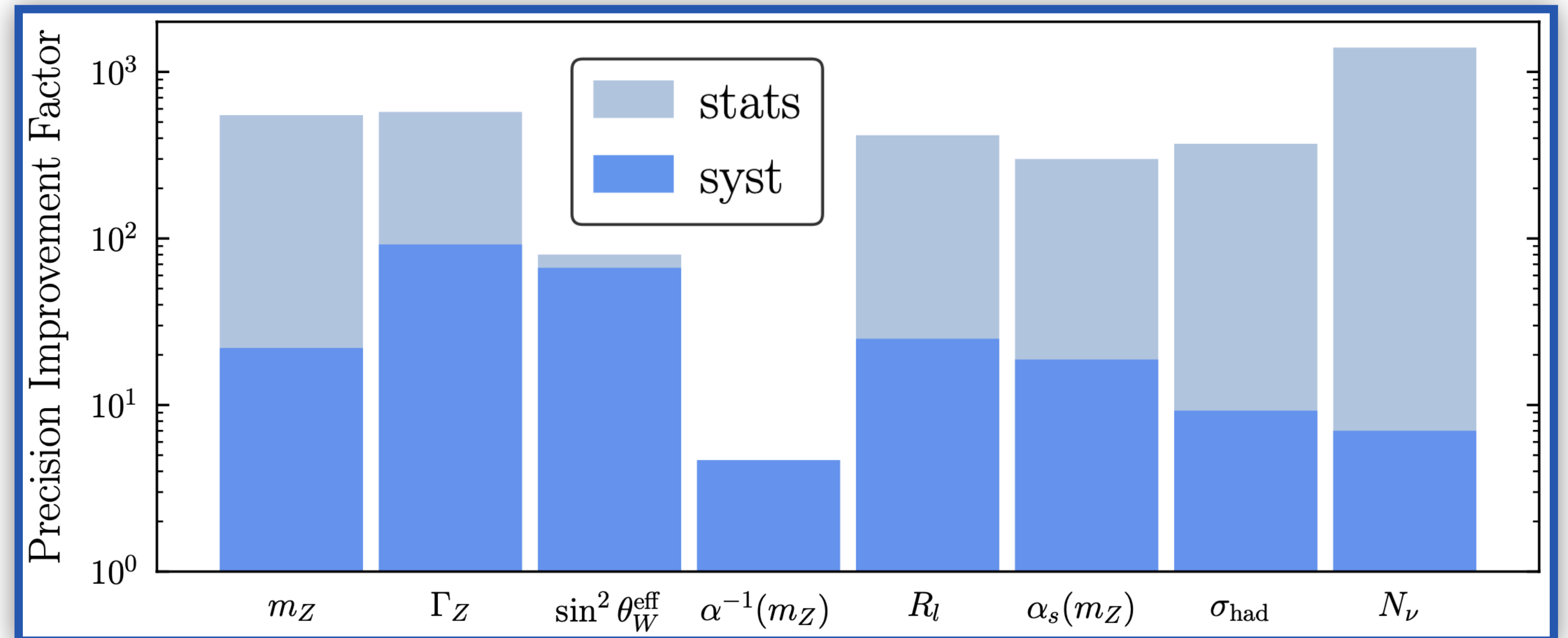
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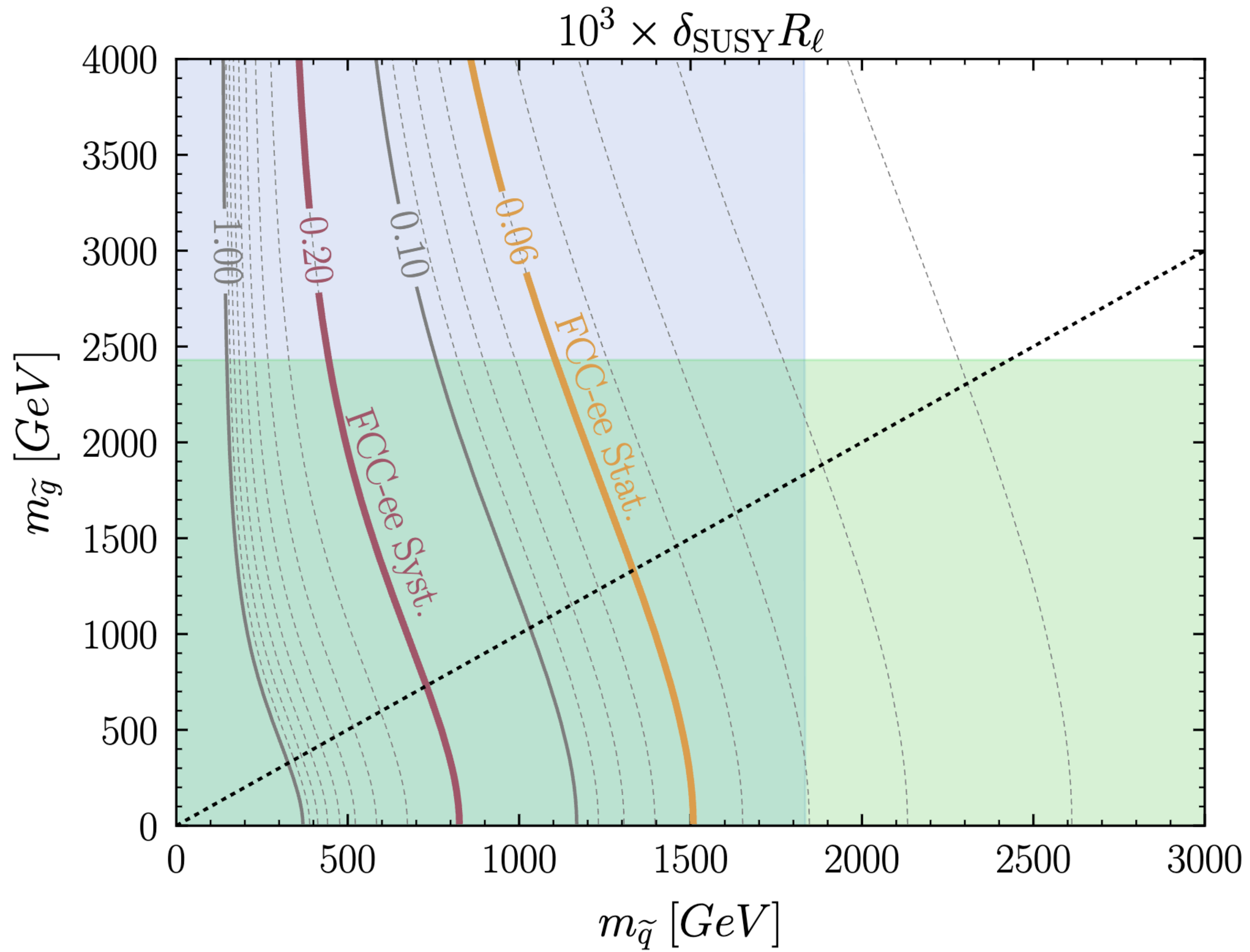
Conclusion

- EWPTs at the FCC-ee are interesting/complimentary ways to search for new physics.
- There exists some SUSY parameter space which may be explored at the FCC-ee.
- A more thorough investigation into which observables/signatures are motivated by various models may motivate further dedication to the reduction of certain systematics.



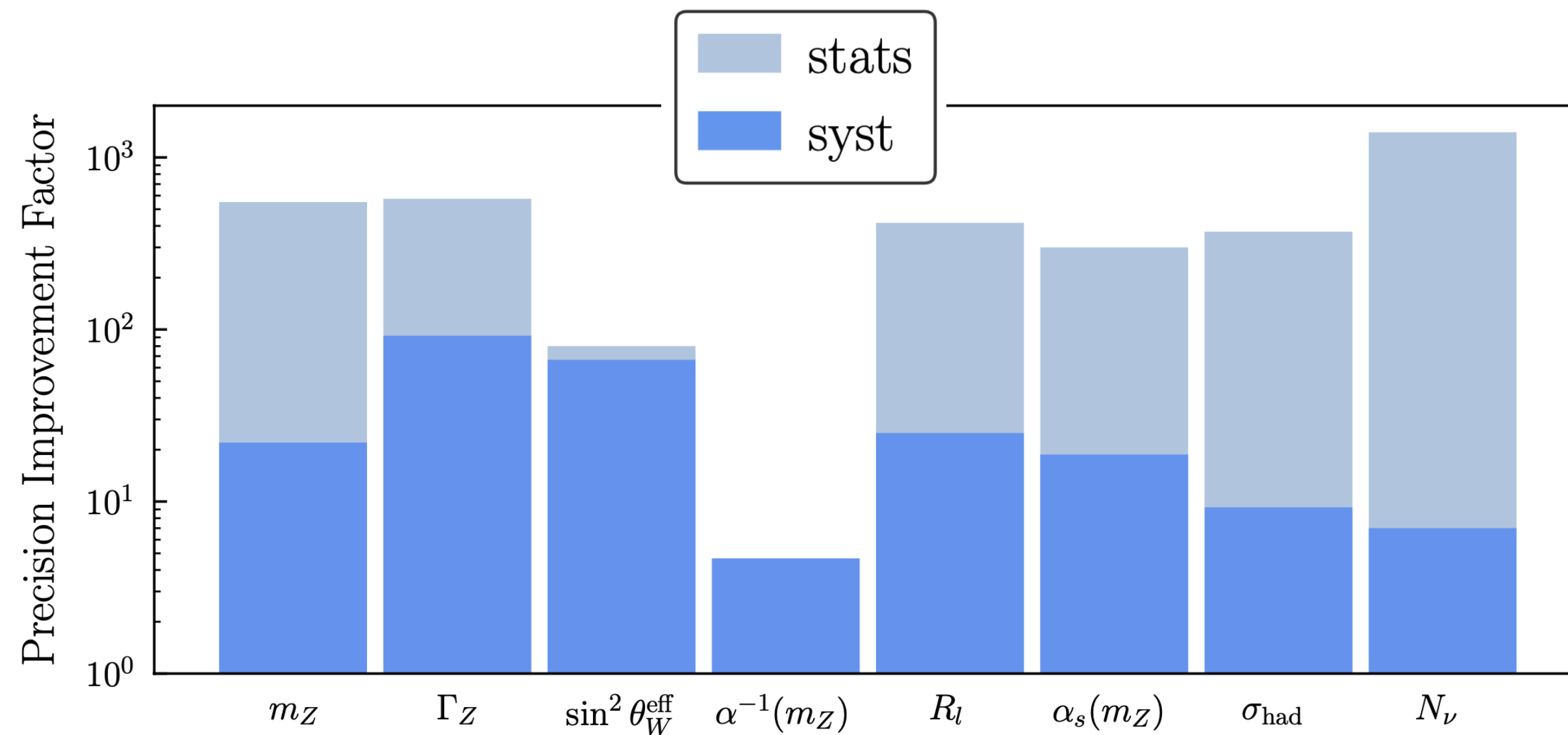
[FCC CDR Vol. 2]

Backup Slides



Electroweak Precision Tests at the Z-pole

There are many measurements which can be performed at the Z-pole.



Many measurements are systematics limited!

Which systematics should we prioritize reducing?

Observable	Present value \pm error	FCC-ee Stat.	FCC-ee Syst.
m_Z (keV)	$91,186,700 \pm 2200$	5	100
Γ_Z (keV)	$2,495,200 \pm 2300$	8	100
R_ℓ^Z ($\times 10^3$)	$20,767 \pm 25$	0.06	0.2–1.0
$\alpha_s(m_Z)$ ($\times 10^4$)	1196 ± 30	0.1	0.4–1.6
R_b ($\times 10^6$)	$216,290 \pm 660$	0.3	< 60
σ_{had}^0 ($\times 10^3$) (nb)	$41,541 \pm 37$	0.1	4
N_ν ($\times 10^3$)	2991 ± 7	0.005	1
$\sin^2 \theta_W^{\text{eff}}$ ($\times 10^6$)	$231,480 \pm 160$	3	2–5
$1/\alpha_{\text{QED}}(m_Z)$ ($\times 10^3$)	$128,952 \pm 14$	4	Small
$A_{\text{FB}}^{b,0}$ ($\times 10^4$)	992 ± 16	0.02	1–3
$A_{\text{FB}}^{\text{pol},\tau}$ ($\times 10^4$)	1498 ± 49	0.15	< 2
m_W (MeV)	$80,350 \pm 15$	0.5	0.3
Γ_W (MeV)	2085 ± 42	1.2	0.3
$\alpha_s(m_W)$ ($\times 10^4$)	1170 ± 420	3	Small
N_ν ($\times 10^3$)	2920 ± 50	0.8	Small
m_{top} (MeV)	$172,740 \pm 500$	17	Small
Γ_{top} (MeV)	1410 ± 190	45	Small
$\lambda_{\text{top}}/\lambda_{\text{top}}^{\text{SM}}$	1.2 ± 0.3	0.1	Small
ttZ couplings	$\pm 30\%$	0.5–1.5%	Small

[FCC CDR]

ATLAS SUSY Searches* - 95% CL Lower Limits

July 2020

ATLAS Preliminary

$\sqrt{s} = 13$ TeV

Model	Signature	$\int \mathcal{L} dt$ [fb ⁻¹]	Mass limit	Reference		
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0 e, μ mono-jet	2-6 jets 1-3 jets E_T^{miss} 139 36.1	\tilde{q} [10x Degen.] 1.9 \tilde{q} [1x, 8x Degen.] 0.43 0.71	$m(\tilde{\chi}_1^0) < 400$ GeV $m(\tilde{q}) - m(\tilde{\chi}_1^0) = 5$ GeV	ATLAS-CONF-2019-040 1711.03301
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0 e, μ	2-6 jets E_T^{miss} 139	\tilde{g} 2.35 \tilde{g} Forbidden 1.15-1.95	$m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{\chi}_1^0) = 1000$ GeV	ATLAS-CONF-2019-040 ATLAS-CONF-2019-040
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}W\tilde{\chi}_1^0$	1 e, μ	2-6 jets	\tilde{g} 2.2	$m(\tilde{\chi}_1^0) < 600$ GeV	ATLAS-CONF-2020-047
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	$ee, \mu\mu$	2 jets E_T^{miss} 36.1	\tilde{g} 1.2	$m(\tilde{g}) - m(\tilde{\chi}_1^0) = 50$ GeV	1805.11381
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}WZ\tilde{\chi}_1^0$	0 e, μ	7-11 jets E_T^{miss} 139	\tilde{g} 1.97	$m(\tilde{\chi}_1^0) < 600$ GeV	ATLAS-CONF-2020-002
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}WZ\tilde{\chi}_1^0$	SS e, μ	6 jets E_T^{miss} 139	\tilde{g} 1.15	$m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200$ GeV	1909.08457
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{\chi}_1^0$	0-1 e, μ SS e, μ	3 b 6 jets E_T^{miss} 79.8 139	\tilde{g} 2.25 \tilde{g} 1.25	$m(\tilde{\chi}_1^0) < 200$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300$ GeV	ATLAS-CONF-2018-041 1909.08457
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/\tilde{u}_1^\pm$	Multiple Multiple	36.1 139	\tilde{b}_1 Forbidden 0.9 \tilde{b}_1 Forbidden 0.74	$m(\tilde{\chi}_1^0) = 300$ GeV, $\text{BR}(b\tilde{\chi}_1^0) = 1$ $m(\tilde{\chi}_1^0) = 200$ GeV, $m(\tilde{\chi}_1^\pm) = 300$ GeV, $\text{BR}(t\tilde{\chi}_1^\pm) = 1$	1708.09266, 1711.03301 1909.08457
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_2^0 \rightarrow bh\tilde{\chi}_1^0$	0 e, μ 2 τ	6 b 2 b E_T^{miss} 139 139	\tilde{b}_1 Forbidden 0.23-1.35 \tilde{b}_1 0.13-0.85	$\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 100$ GeV $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 0$ GeV	1908.03122 ATLAS-CONF-2020-031
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	0-1 e, μ	≥ 1 jet E_T^{miss} 139	\tilde{t}_1 1.25	$m(\tilde{\chi}_1^0) = 1$ GeV	ATLAS-CONF-2020-003, 2004.14060
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$	1 e, μ	3 jets/1 b E_T^{miss} 139	\tilde{t}_1 0.44-0.59	$m(\tilde{\chi}_1^0) = 400$ GeV	ATLAS-CONF-2019-017
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b\nu, \tilde{\tau}_1 \rightarrow \tau\tilde{G}$	1 $\tau + 1 e, \mu, \tau$	2 jets/1 b E_T^{miss} 36.1	\tilde{t}_1 1.16	$m(\tilde{\tau}_1) = 800$ GeV	1803.10178
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c} \rightarrow c\tilde{\chi}_1^0$	0 e, μ	2 c E_T^{miss} 36.1	\tilde{t}_1 0.46 0.85 \tilde{t}_1 0.43	$m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 50$ GeV $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 5$ GeV	1805.01649 1805.01649 1711.03301
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z/h\tilde{\chi}_1^0$	1-2 e, μ	1-4 b E_T^{miss} 139	\tilde{t}_1 0.067-1.18	$m(\tilde{\chi}_2^0) = 500$ GeV	SUSY-2018-09
$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 e, μ	1 b E_T^{miss} 139	\tilde{t}_2 Forbidden 0.86	$m(\tilde{\chi}_1^0) = 360$ GeV, $m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 40$ GeV	SUSY-2018-09	
EW direct	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via WZ	3 e, μ $ee, \mu\mu$	≥ 1 jet E_T^{miss} 139 139	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.64 $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.205	$m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5$ GeV	ATLAS-CONF-2020-015 1911.12606
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via WW	2 e, μ	E_T^{miss} 139	$\tilde{\chi}_1^\pm$ 0.42	$m(\tilde{\chi}_1^0) = 0$	1908.08215
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via Wh	0-1 e, μ	2 $b/2 \gamma$ E_T^{miss} 139	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ Forbidden 0.74	$m(\tilde{\chi}_1^0) = 70$ GeV	2004.10894, 1909.09226
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via $\tilde{\ell}_L/\tilde{\nu}$	2 e, μ	E_T^{miss} 139	$\tilde{\chi}_1^\pm$ 1.0	$m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$	1908.08215
	$\tilde{\tau}\tilde{\tau}, \tilde{\tau} \rightarrow \tau\tilde{\chi}_1^0$	2 τ	E_T^{miss} 139	$\tilde{\tau}$ [$\tilde{\tau}_L, \tilde{\tau}_{R,L}$] 0.16-0.3 0.12-0.39	$m(\tilde{\chi}_1^0) = 0$	1911.06660
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ $ee, \mu\mu$	0 jets ≥ 1 jet E_T^{miss} 139 139	$\tilde{\ell}$ 0.7 $\tilde{\ell}$ 0.256	$m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\ell}) - m(\tilde{\chi}_1^0) = 10$ GeV	1908.08215 1911.12606
	$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	0 e, μ 4 e, μ	$\geq 3 b$ 0 jets E_T^{miss} 36.1 139	\tilde{H} 0.13-0.23 0.29-0.88 \tilde{H} 0.55	$\text{BR}(\tilde{\chi}_1^0 \rightarrow h\tilde{G}) = 1$ $\text{BR}(\tilde{\chi}_1^0 \rightarrow Z\tilde{G}) = 1$	1806.04030 ATLAS-CONF-2020-040
Long-lived particles	Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet E_T^{miss} 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_1^\mp$ 0.46 $\tilde{\chi}_1^\pm$ 0.15	Pure Wino Pure higgsino	1712.02118 ATL-PHYS-PUB-2017-019
	Stable \tilde{g} R-hadron	Multiple	36.1	\tilde{g} 2.0		1902.01636, 1808.04095
	Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	Multiple	36.1	\tilde{g} [$\tau(\tilde{g}) = 10$ ns, 0.2 ns] 2.05 2.4	$m(\tilde{\chi}_1^0) = 100$ GeV	1710.04901, 1808.04095
RPV	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_1^0, \tilde{\chi}_1^\pm \rightarrow Z\ell \rightarrow \ell\ell\ell$	3 e, μ	139	$\tilde{\chi}_1^\pm/\tilde{\chi}_1^0$ [BR(Z τ)=1, BR(Z e)=1] 0.625 1.05	Pure Wino	ATLAS-CONF-2020-009
	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/\tau\mu$	$e\mu, e\tau, \mu\tau$	3.2	$\tilde{\nu}_\tau$ 1.9	$\lambda'_{311} = 0.11, \lambda'_{132/133/233} = 0.07$	1607.08079
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0 \rightarrow WW/Z\ell\ell\nu\nu$	4 e, μ	0 jets E_T^{miss} 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ [$\lambda'_{133} \neq 0, \lambda'_{12k} \neq 0$] 0.82 1.33	$m(\tilde{\chi}_1^0) = 100$ GeV	1804.03602
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$	4-5 large-R jets Multiple	36.1 36.1	\tilde{g} [$m(\tilde{\chi}_1^0) = 200$ GeV, 1100 GeV] 1.3 1.9 \tilde{g} [$\lambda'_{112} = 2e-4, 2e-5$] 1.05 2.0	Large λ'_{12} $m(\tilde{\chi}_1^0) = 200$ GeV, bino-like	1804.03568 ATLAS-CONF-2018-003
	$\tilde{t}\tilde{t}, \tilde{t} \rightarrow t\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow tbs$	Multiple	36.1	\tilde{t} [$\lambda'_{323} = 2e-4, 1e-2$] 0.55 1.05	$m(\tilde{\chi}_1^0) = 200$ GeV, bino-like	ATLAS-CONF-2018-003
	$\tilde{t}\tilde{t}, \tilde{t} \rightarrow b\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow bbs$	$\geq 4b$	139	\tilde{t} Forbidden 0.95	$m(\tilde{\chi}_1^\pm) = 500$ GeV	ATLAS-CONF-2020-016
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$	2 jets + 2 b	36.7	\tilde{t}_1 [qq, bs] 0.42 0.61		1710.07171
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$	2 e, μ 1 μ	2 b DV 36.1 136	\tilde{t}_1 0.4-1.45 \tilde{t}_1 [$1e-10 < \lambda'_{23k} < 1e-8, 3e-10 < \lambda'_{23k} < 3e-9$] 1.0 1.6	$\text{BR}(\tilde{t}_1 \rightarrow b\ell) > 20\%$ $\text{BR}(\tilde{t}_1 \rightarrow q\mu) = 100\%, \cos\theta_t = 1$	1710.05544 2003.11956	

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

10⁻¹ 1 Mass scale [TeV]

Oblique Corrections

- Assuming heavy new physics dominantly modifies SM gauge boson propagators, corrections from heavy new physics can be quantified by a set of oblique parameters.



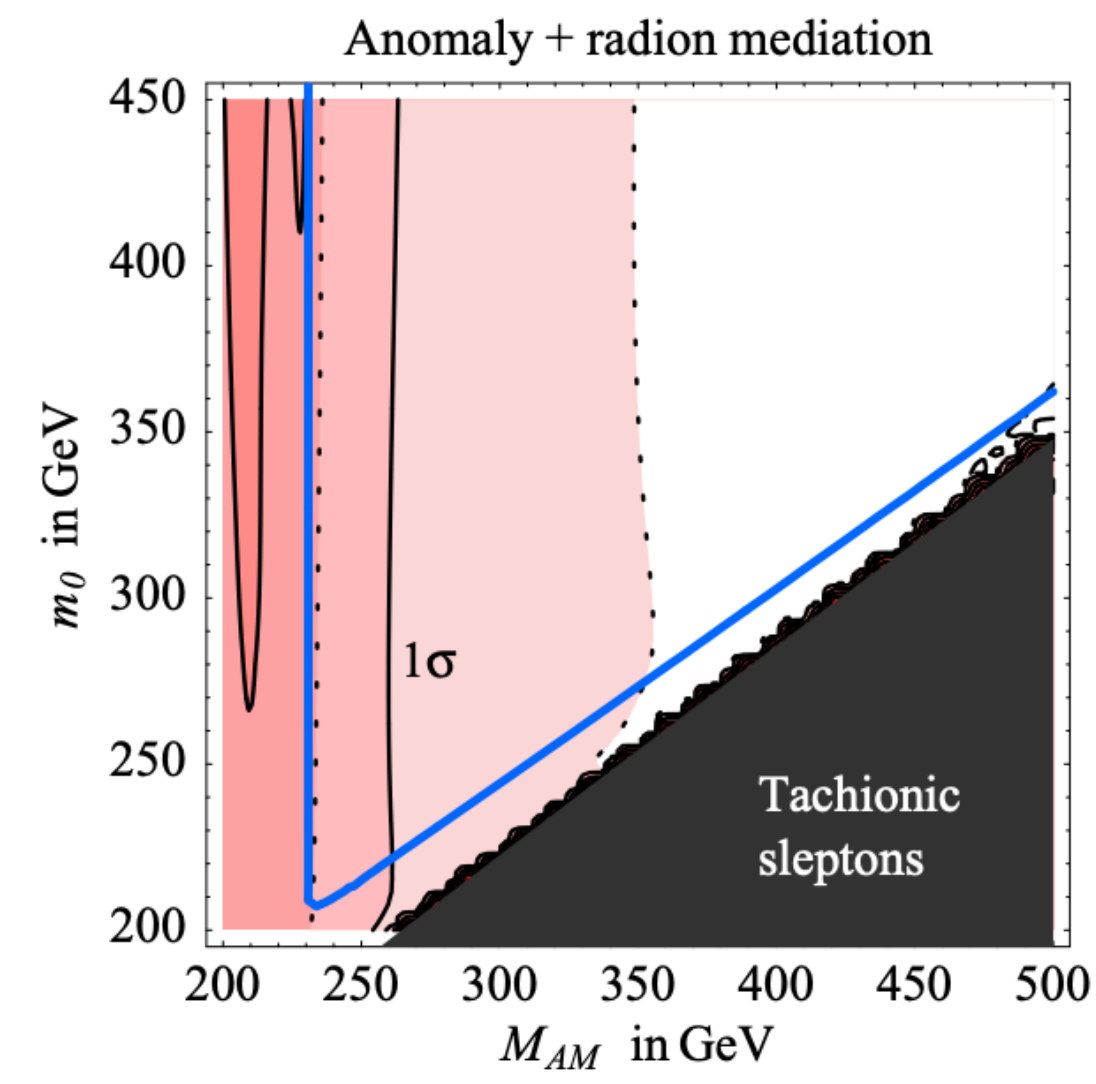
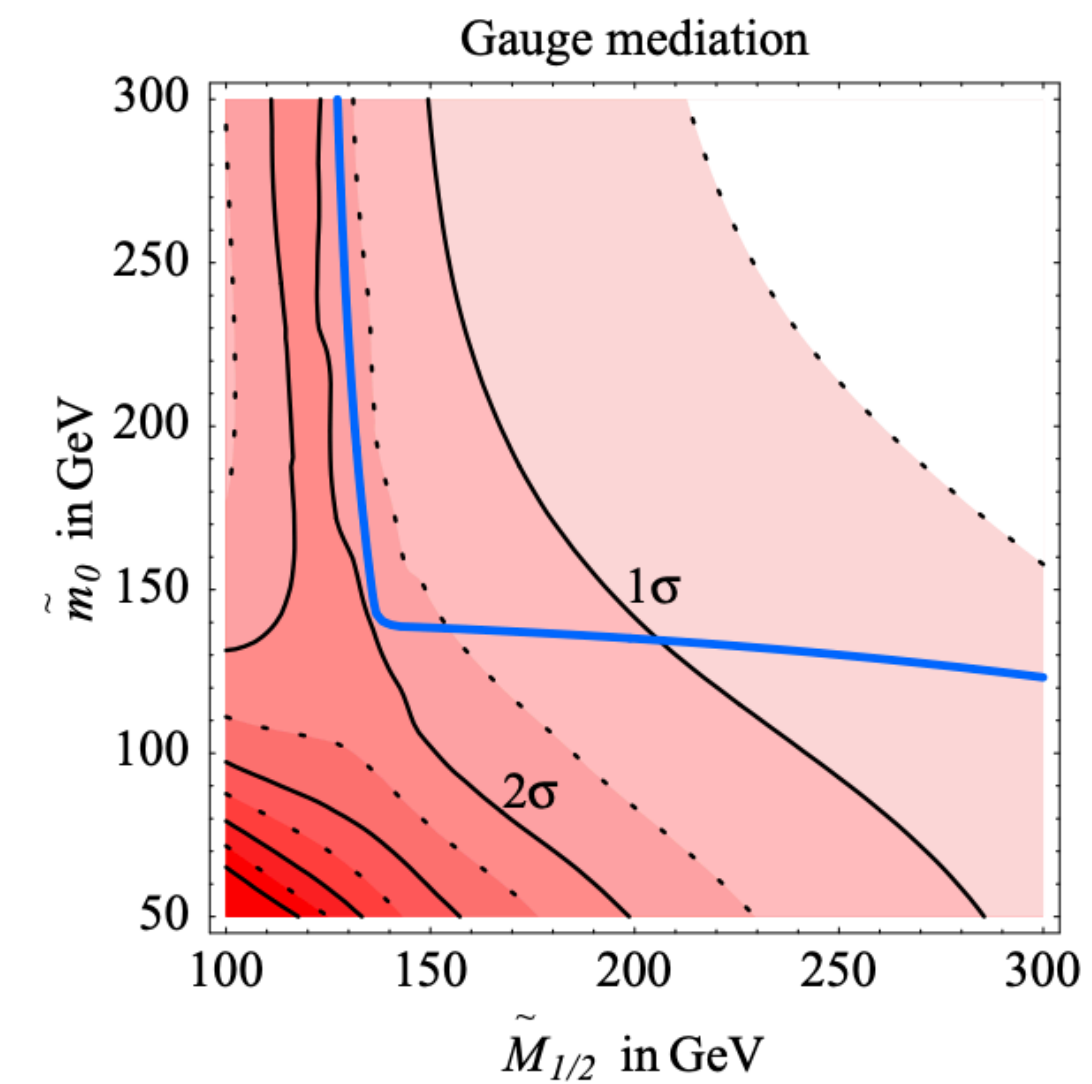
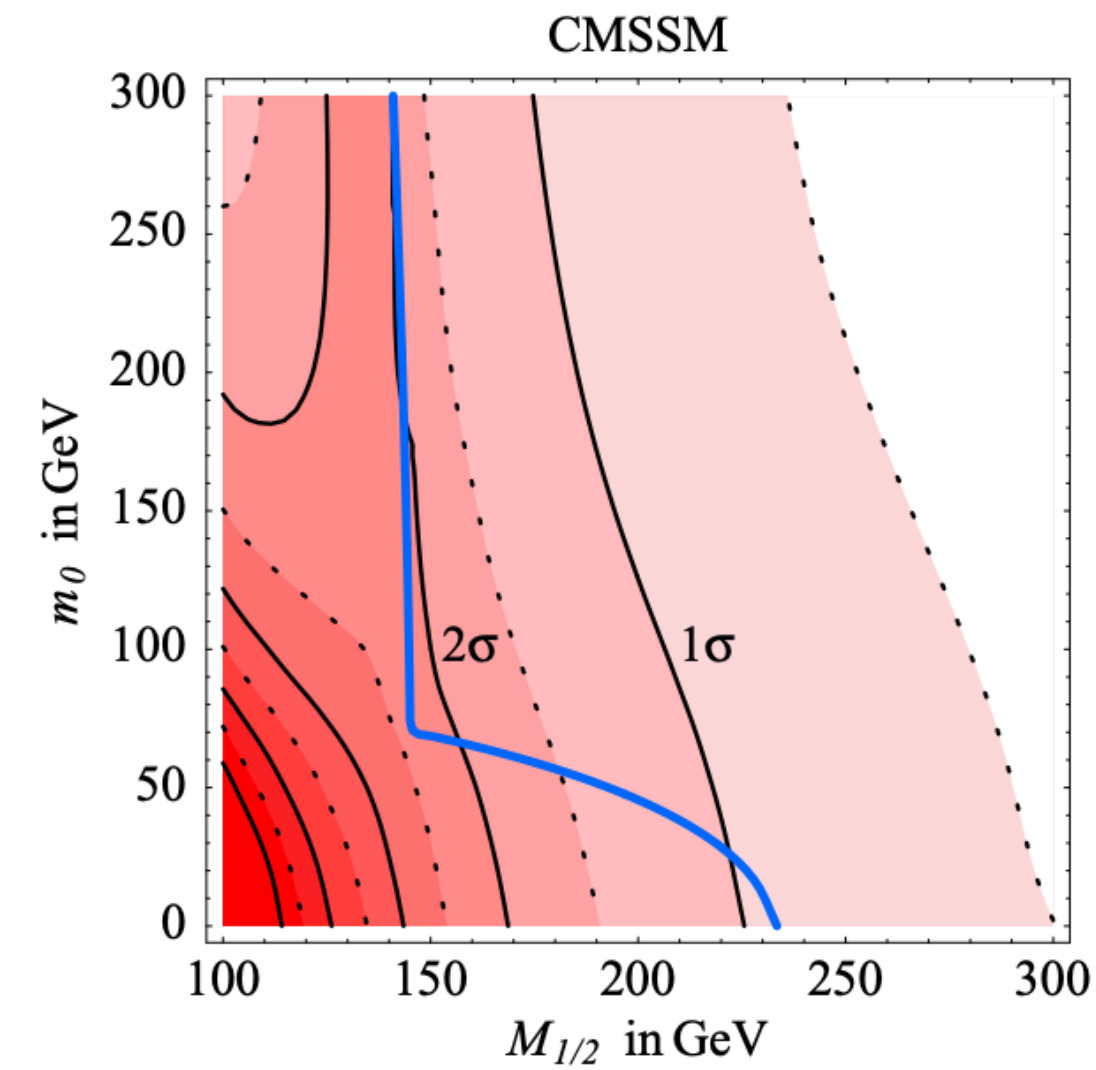
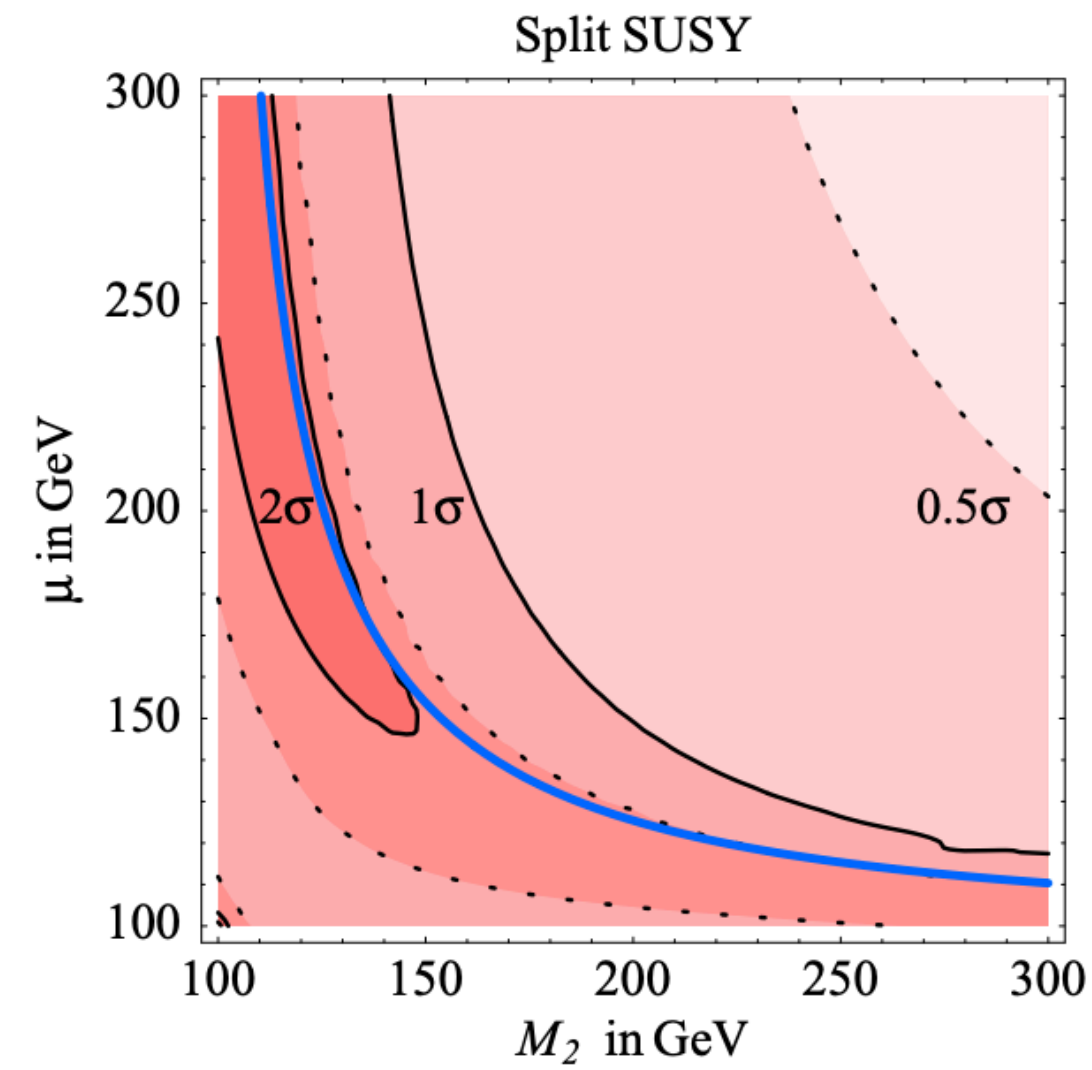
Oblique parameter	Corresponding dim-6 operator
$\hat{S} = \frac{g}{g'} \Pi'_{W^3 B}(0)$	$\mathcal{O}_S = H^\dagger \tau^a H W_{\mu\nu}^a B^{\mu\nu}$
$\hat{T} = m_W^{-2} [\Pi_{W^3 W^3}(0) - \Pi_{W^+ W^-}(0)]$	$\mathcal{O}_T = H^\dagger D_\mu H ^2$
$W = \frac{1}{2} m_W^2 \Pi''_{W^3 W^3}(0)$	$\mathcal{O}_W = \frac{1}{2} (D_\rho W_{\mu\nu}^a)^2$
$Y = \frac{1}{2} m_W^2 \Pi''_{BB}(0)$	$\mathcal{O}_Y = \frac{1}{2} (\partial_\rho B_{\mu\nu})^2$

- Simple, but not completely general.

X	$\hat{S} \times \left(\frac{m_X^2}{m_W^2}\right)$	$\hat{T} \times \left(\frac{m_X^2}{m_W^2}\right)$	$W \times \left(\frac{m_X^2}{m_W^2}\right)$	$Y \times \left(\frac{m_X^2}{m_W^2}\right)$
\tilde{E}	0	0	0	$\frac{\alpha_Y}{40\pi}$
\tilde{L}	$-\frac{\alpha_W c_{2\beta}}{16\pi}$	$\frac{\alpha_W c_{2\beta}^2}{16\pi}$	$\frac{\alpha_W}{80\pi}$	$\frac{\alpha_Y}{80\pi}$
\tilde{B}	0	0	0	0
\tilde{W}	0	0	$\frac{\alpha_W}{15\pi}$	0

Supersymmetry and precision data after LEP2 [0502095]

Guido Marandella^a, Christian Schappacher^b, Alessandro Strumia^c



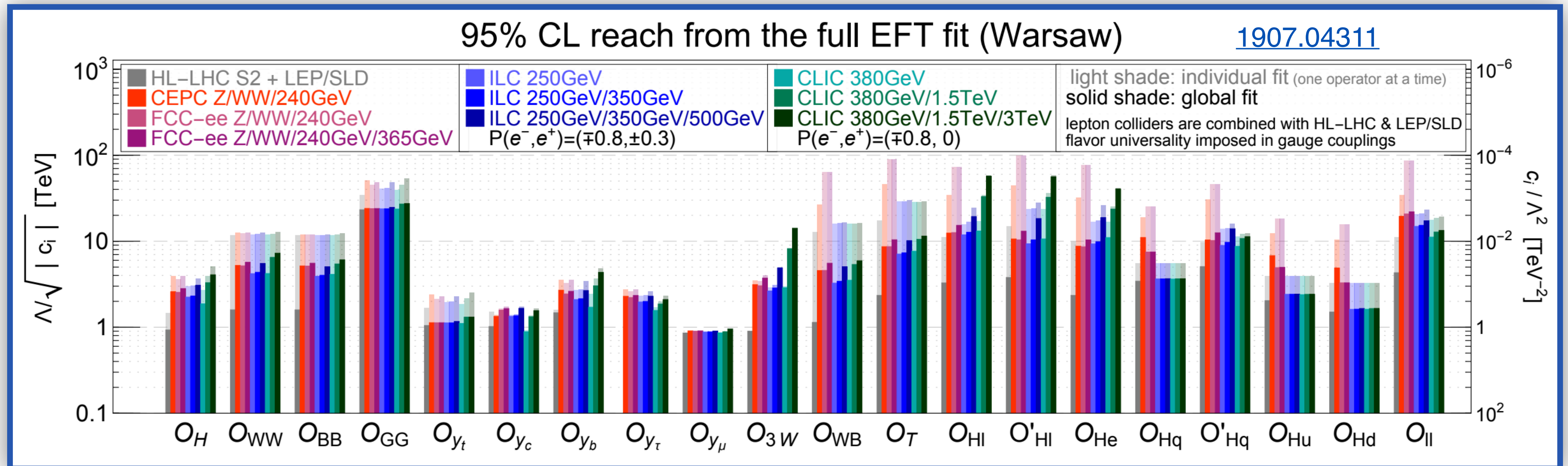
Focuses on
oblique corrections

Electroweak Precision Tests

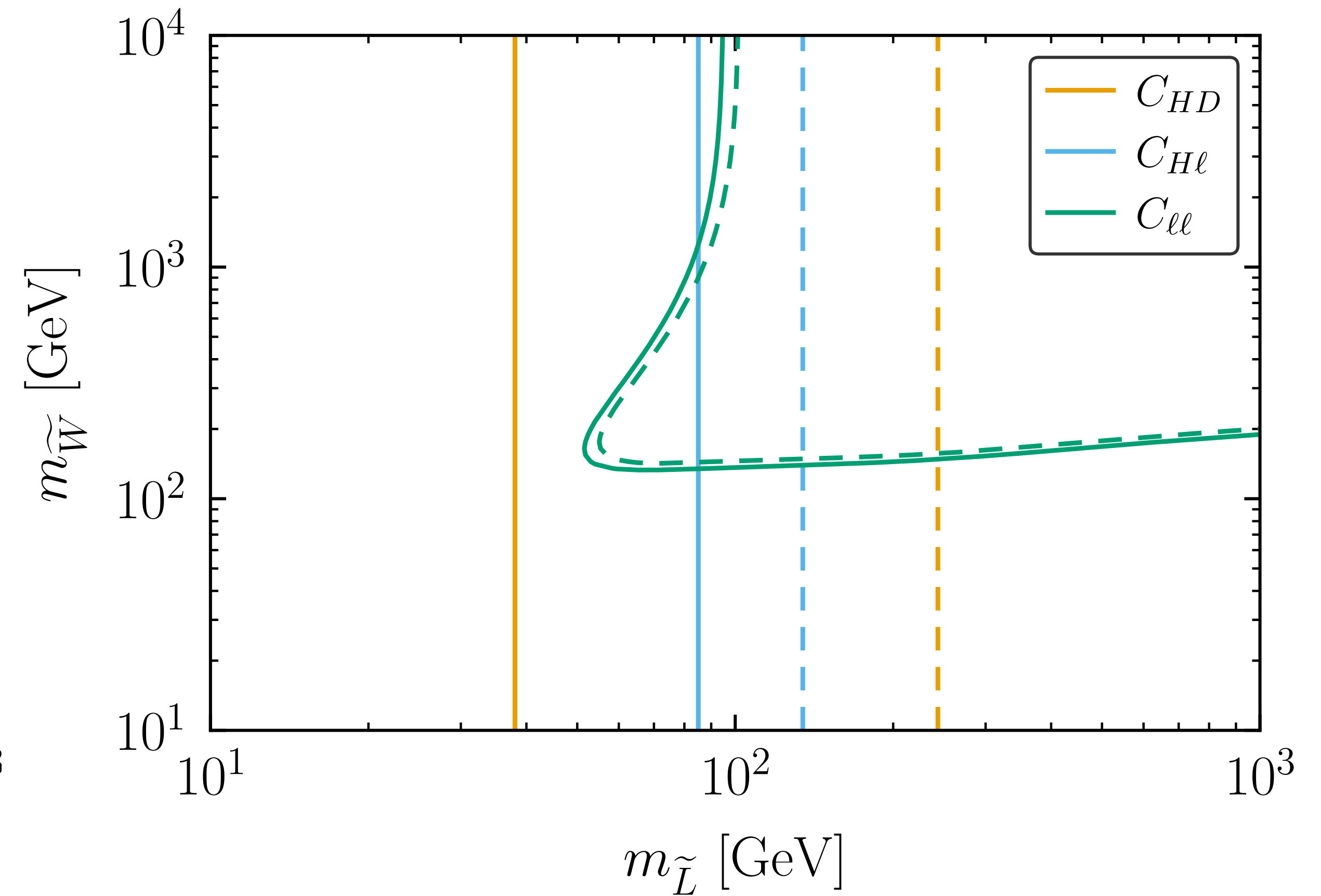
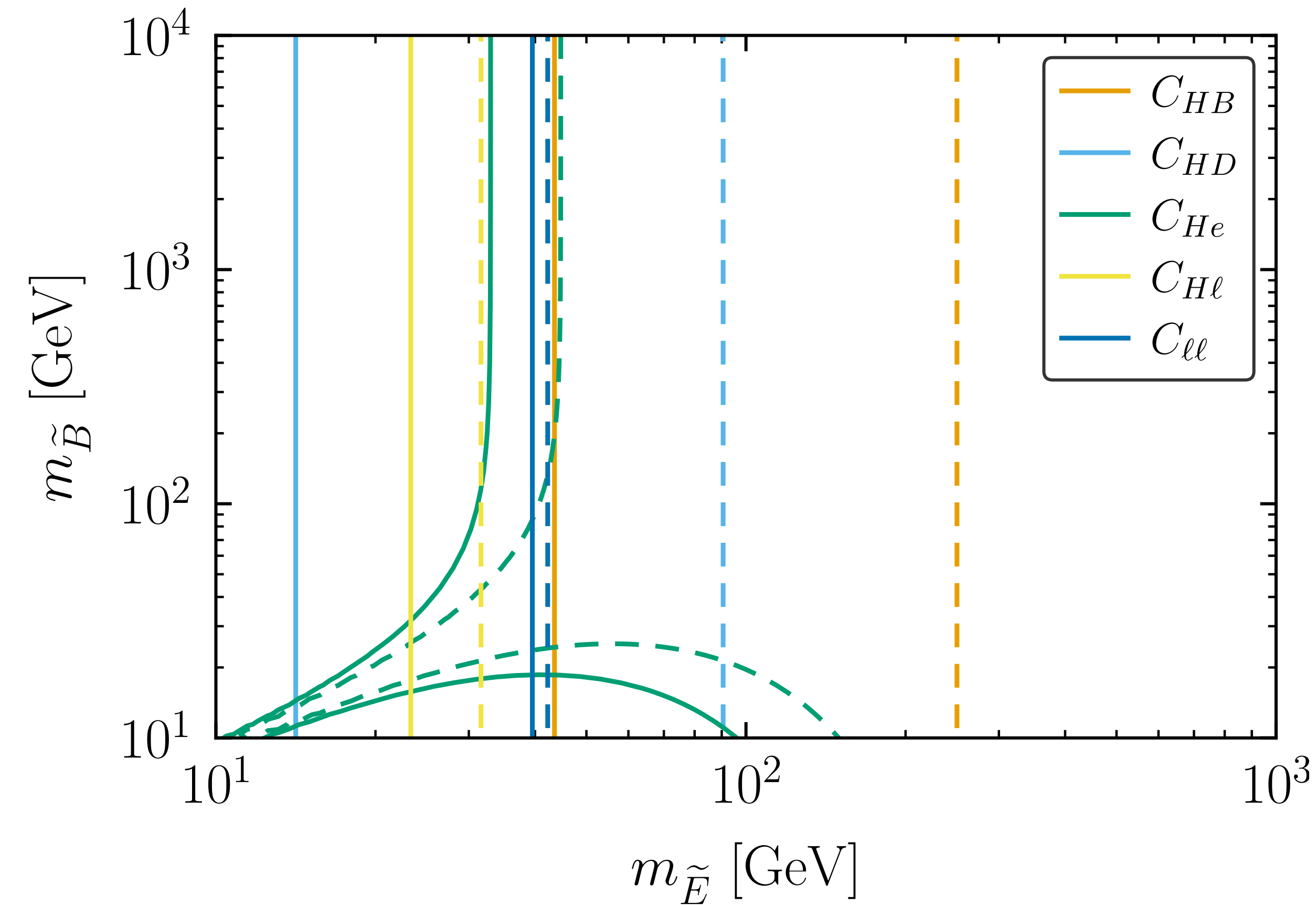
- The most general method of indirectly searching for heavy new physics is SMEFT.

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{n=5}^{\infty} \sum_i \frac{c_i^{(n)}}{\Lambda^{n-4}} \mathcal{O}_i^{(n)}$$

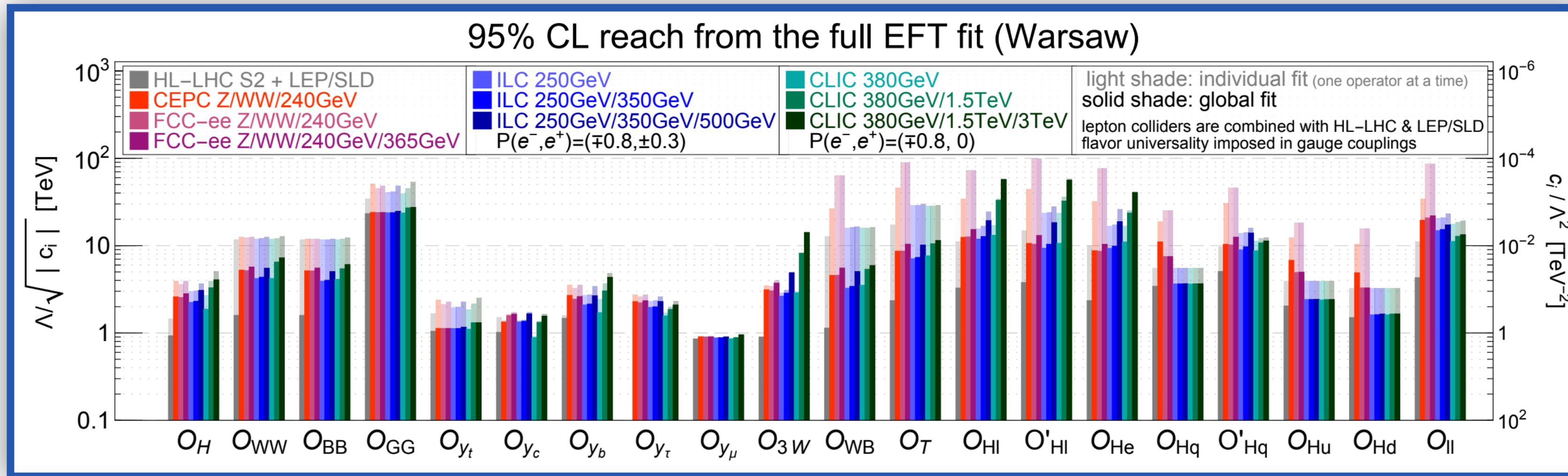
- Assuming CP conservation and MFV, about 20 operators are relevant for EWPTs.



Constraints from SMEFT



A Tale of Two Bar Plots



How do we interpret these?

Questions

1. Which SMEFT operators are most interesting?
2. Which systematics should experimentalists be most motivated to decrease?

