

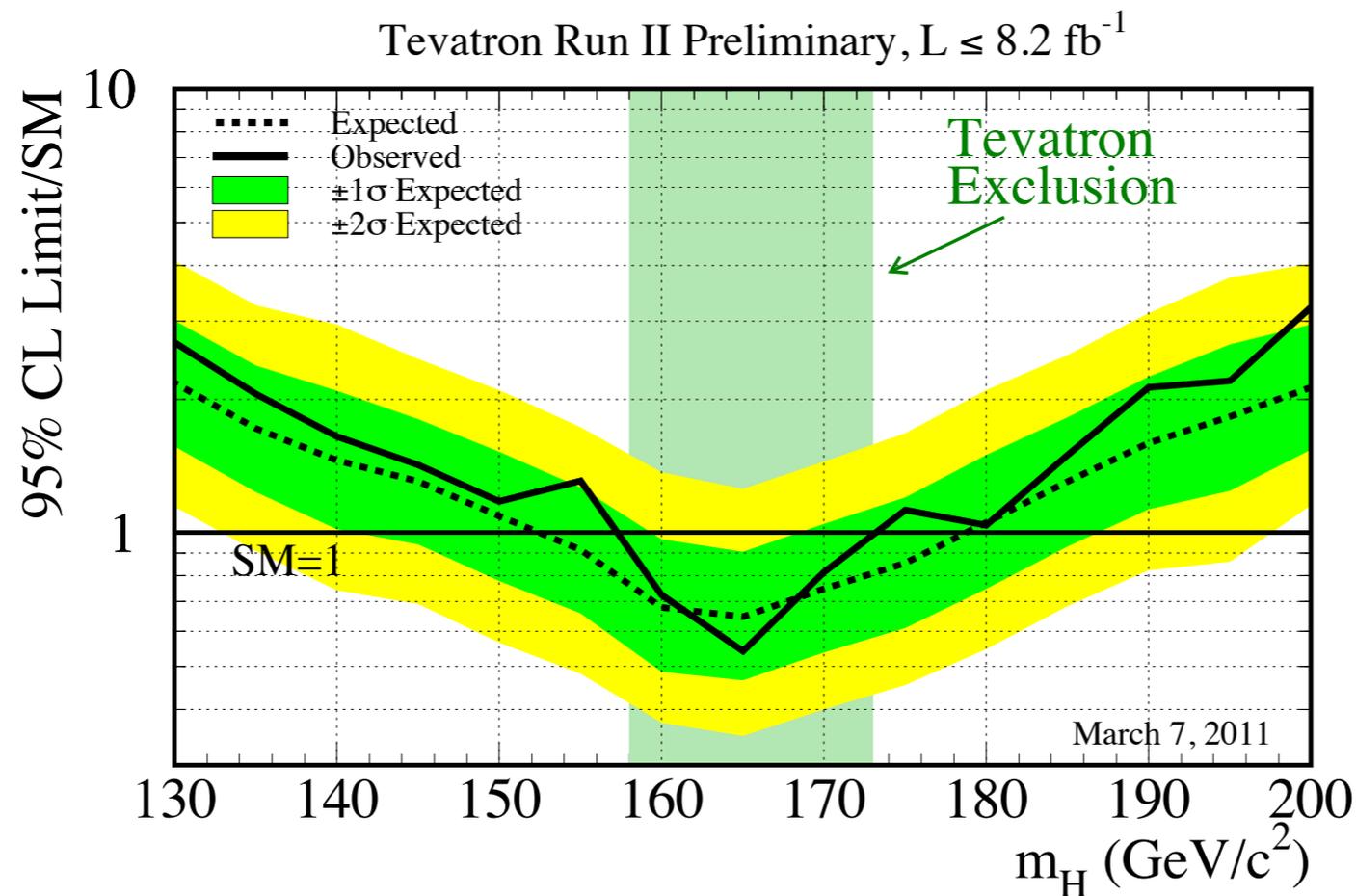
QCD Theory

DPF 2011

Christian Bauer

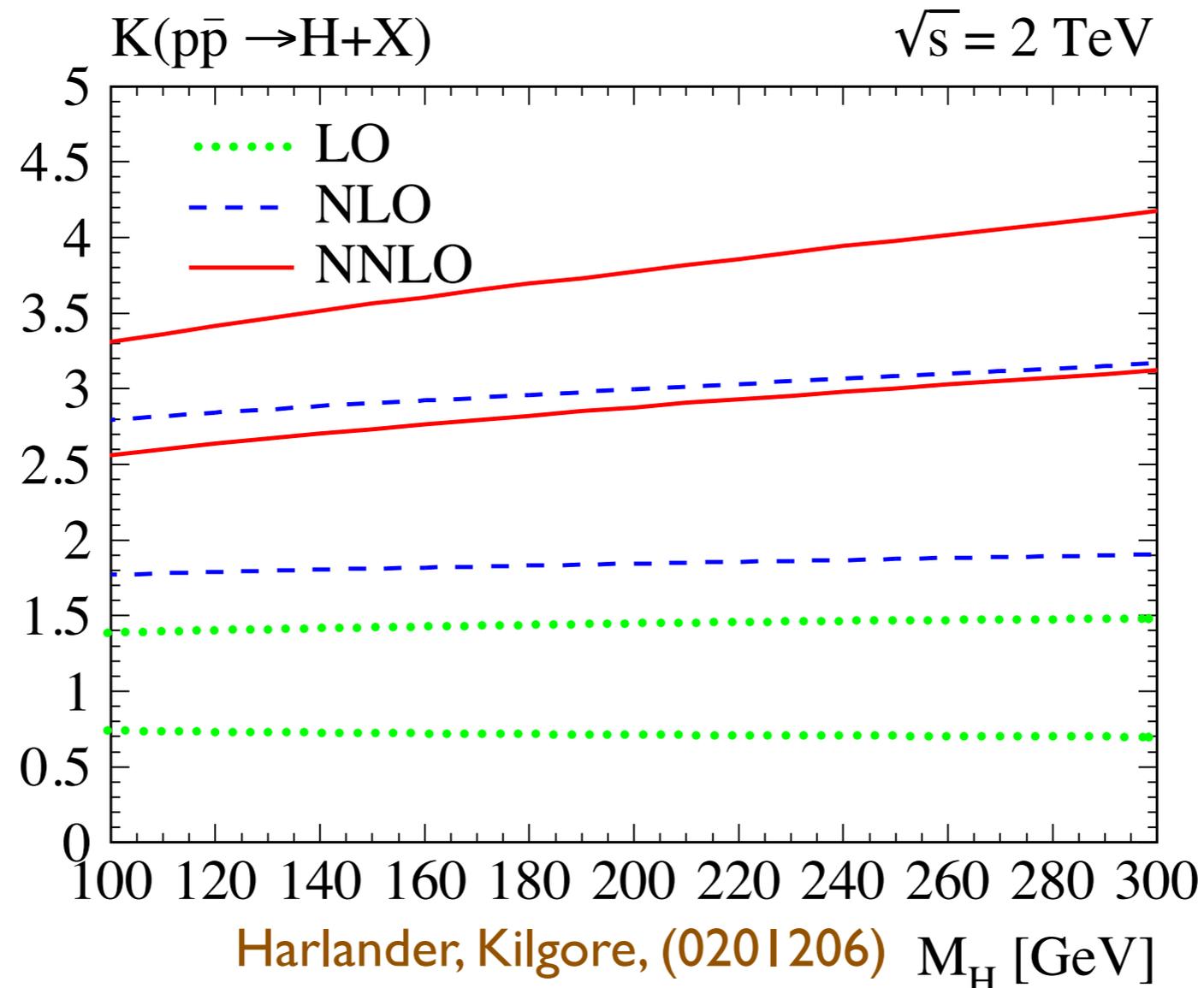
Some important lessons to remember...

... using Higgs search as example



Fixed order calculations

Cross-section for Higgs production known for long time



LO order result completely underestimates central value and uncertainties

Lesson 1:

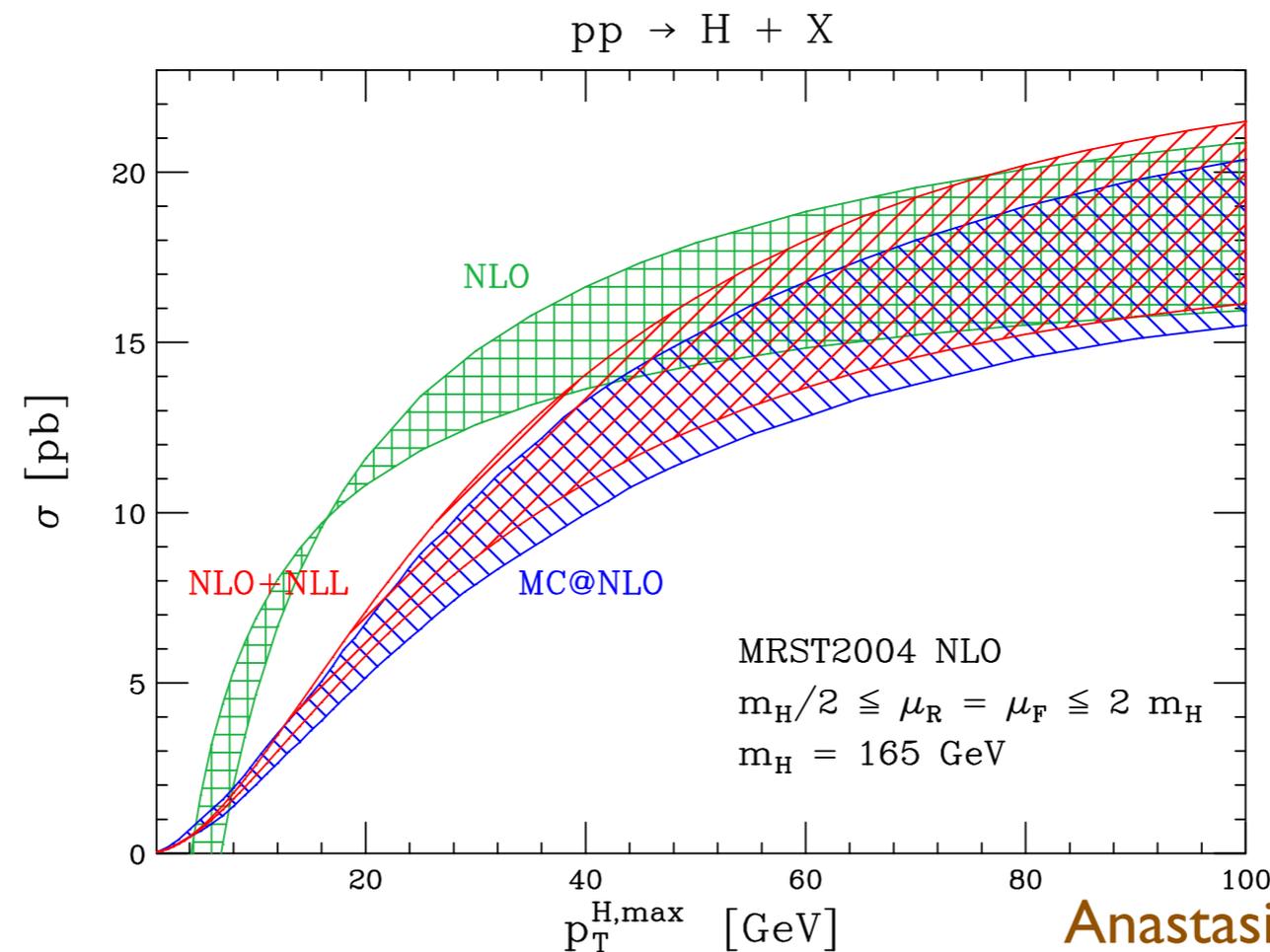
For reliable calculations, need
at least NLO results

Large logarithms

Ratio of scales appear as logarithms in perturbative results

For each power of α_s , typically get two powers of logarithms

If large ratios of scales present, should sum these logarithms for well behaved perturbation theory



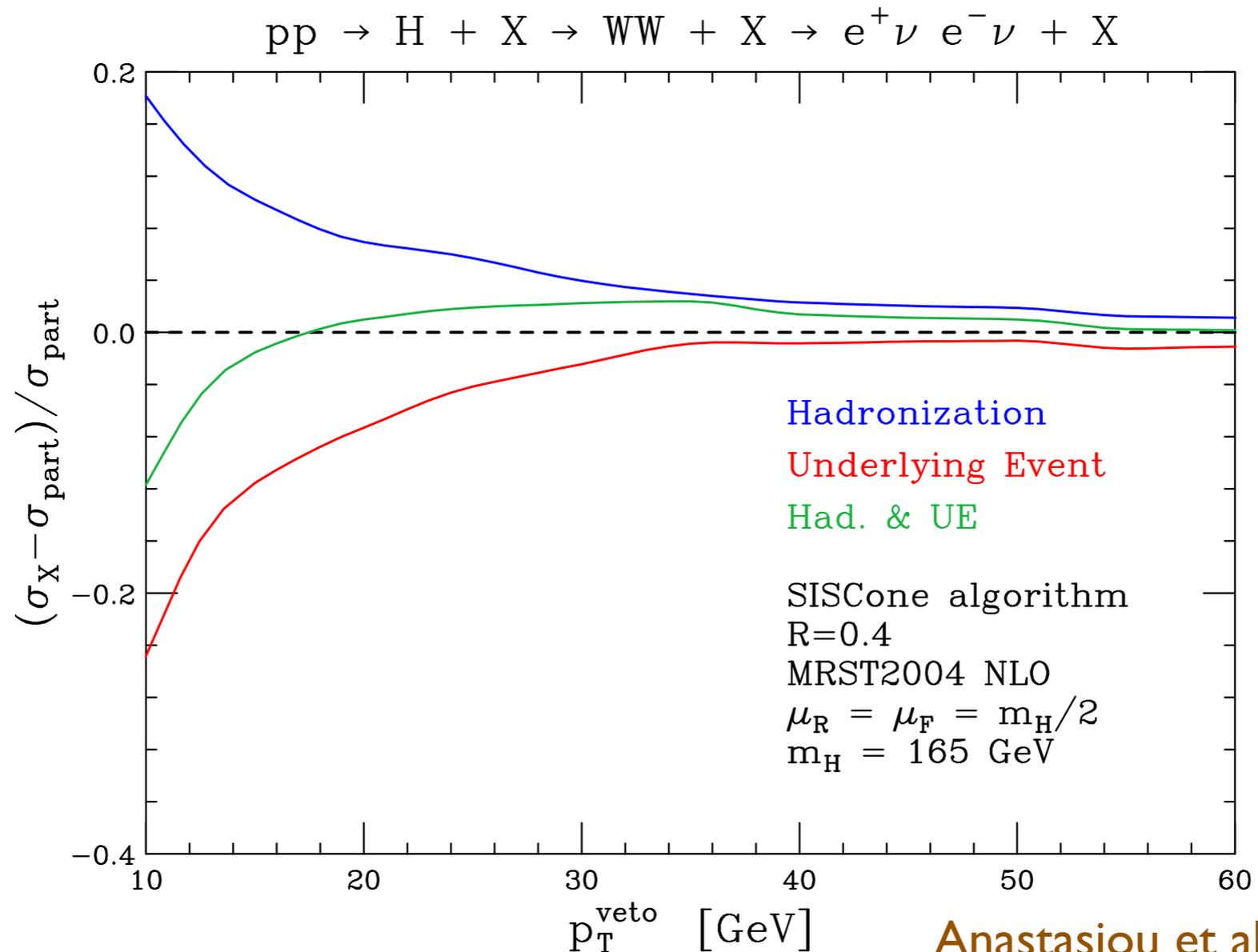
Anastasiou et al (0801.2682)

Lesson 2:

Large ratios of scales lead to
logarithms that should be
resummed

Hadronic effects

Hadronization and underlying event affects the cross sections



Can only be estimated using parton shower algorithms

Lesson 3:

Need parton shower
simulations to understand long
distance hadronic physics

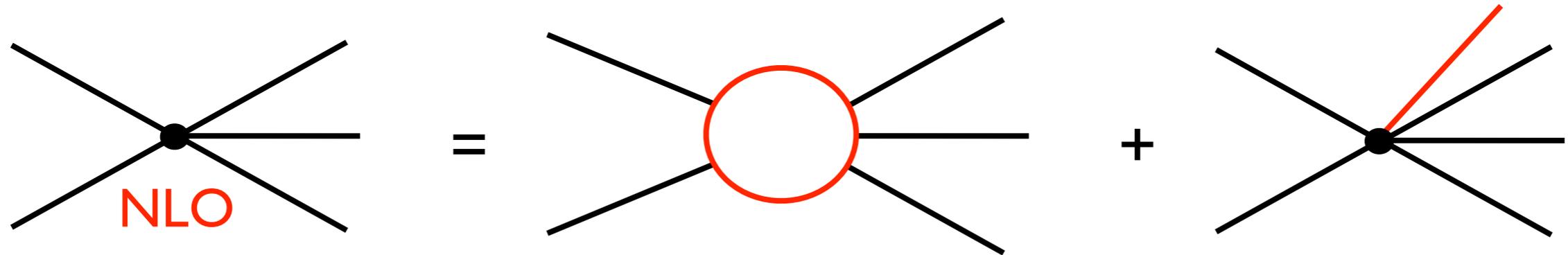
Outline

- Recent progress in NLO calculations
- Large logarithmic terms in perturbative expressions
- Merging fixed order calculations with parton showers

NLO calculations

General Idea

Need to combine one-loop with the real emission



Different complications for loop and real emissions diagrams:

- Loop diagrams require integrals over loop momenta
 - These integrals are divergent
- Real diagrams require integral over extra phase space
 - These integrals are divergent as well

At least part of the integrals have to be performed analytically

This issue has been solved for arbitrary real diagrams

Frixione, Kunzst, Signer, 9512328; Catani, Seymour, 9605323

Integrating loop diagrams

Generic loop integrals can get very complicated, so impossible to integrate by brute force

Strategy: Reduce to a set of master integrals

$$A = \sum d_i \text{[Square Diagram]} + \sum c_i \text{[Triangle Diagram]} + \sum b_i \text{[Bubble Diagram]} + \sum a_i \text{[Self-Energy Diagram]} + R$$

All master integrals have been calculated analytically

t'Hooft, Veltman (1979); Oldenborg, Vermaseren (19990); Beenakker, Denner, (1990);
Bern, Dixon, Kosower (9306240); Duplancic, Nizic, (0006249); ...

All results presented in Ellis, Zanderighi (0712.1851)
and available online at <http://qcdloop.fnal.gov>

Need to determine the coefficients d_i , c_i , b_i , a_i , and R

Obtaining the coefficients

Two different techniques available

Tensor reduction:

Decompose numerator into generic tensor structures, work out coefficients

$$C_1^\mu(p_1, p_2) = \int \frac{d^d k}{(2\pi)^d} \frac{k^\mu}{D_3(p_1, p_2)} = C_1^1 p_1^\mu + C_1^2 p_2^\mu$$

tHooft, Veltman, (1979); Passarino, Veltman, (1979);
Dittmaier, Denner, (0509141); ...

Gets complicated for high rank tensors,
since spurious divergences give
numerical instabilities

Tools using this method:

GOLEM, ...

Generalized unitarity

Use analytical structure on loop amplitudes to work out coefficients

$$\mathcal{A}_1|_{4-cut_i} = d_i D0_i|_{4-cut}$$

Bern, Dixon, Forde, Kosower, et al.; Ellis, Giele, Kuzst, Melnikov, Zanderighi; Ossala, Papadopolous, Pittau; Britto, Cachazo, Feng

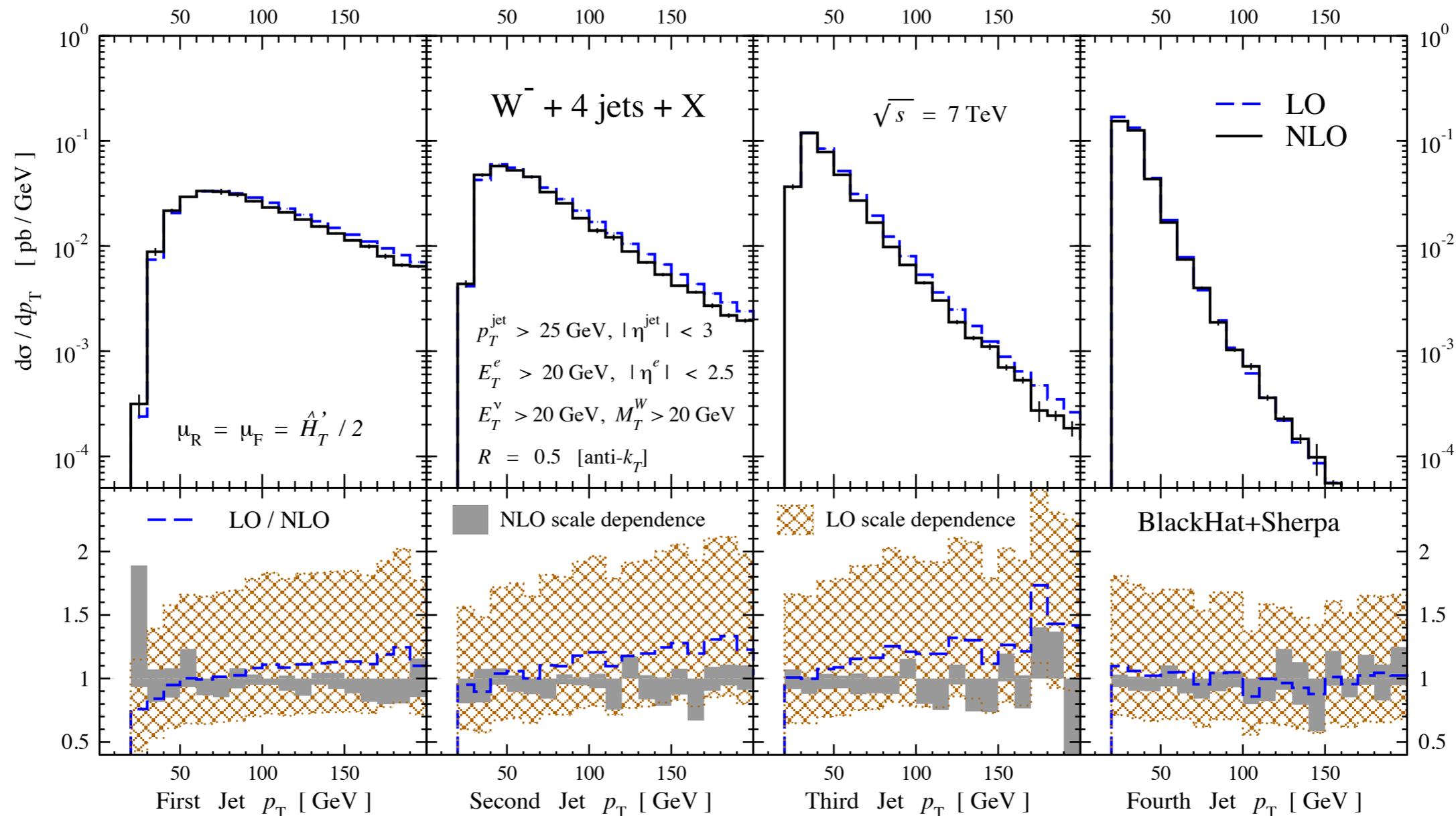
Lends itself to automatization, and this approach has allowed many of the recent new advances

Tools using this method:

Blackhat, Rocket, MadLoop, ...

Recent unitarity result

W+4 jet, first pp → 5 process



Blackhat,(1009.2338)

Nobody would have imagined that this is possible 5 years ago

Automatization of NLO

Madgraph is tool that creates code calculating arbitrary LO results

Alwall, Herquet, Maltoni, Mattelaer, Stelzer, ...

Madloop attempts same @ NLO

R. Frederix, S. Frixione, M.V. Garzelli, V. Hirschi, F. Maltoni, R. Pittau, (1103.0621)

Process	μ	n_{lf}	Cross section (pb)	
			LO	NLO
a.1 $pp \rightarrow t\bar{t}$	m_{top}	5	123.76 ± 0.05	162.08 ± 0.12
a.2 $pp \rightarrow tj$	m_{top}	5	34.78 ± 0.03	41.03 ± 0.07
a.3 $pp \rightarrow tjj$	m_{top}	5	11.851 ± 0.006	13.71 ± 0.02
a.4 $pp \rightarrow t\bar{b}j$	$m_{top}/4$	4	25.62 ± 0.01	30.96 ± 0.06
a.5 $pp \rightarrow t\bar{b}jj$	$m_{top}/4$	4	8.195 ± 0.002	8.91 ± 0.01
b.1 $pp \rightarrow (W^+ \rightarrow)e^+\nu_e$	m_W	5	5072.5 ± 2.9	6146.2 ± 9.8
b.2 $pp \rightarrow (W^+ \rightarrow)e^+\nu_e j$	m_W	5	828.4 ± 0.8	1065.3 ± 1.8
b.3 $pp \rightarrow (W^+ \rightarrow)e^+\nu_e jj$	m_W	5	298.8 ± 0.4	300.3 ± 0.6
b.4 $pp \rightarrow (\gamma^*/Z \rightarrow)e^+e^-$	m_Z	5	1007.0 ± 0.1	1170.0 ± 2.4
b.5 $pp \rightarrow (\gamma^*/Z \rightarrow)e^+e^- j$	m_Z	5	156.11 ± 0.03	203.0 ± 0.2
b.6 $pp \rightarrow (\gamma^*/Z \rightarrow)e^+e^- jj$	m_Z	5	54.24 ± 0.02	56.69 ± 0.07
c.1 $pp \rightarrow (W^+ \rightarrow)e^+\nu_e b\bar{b}$	$m_W + 2m_b$	4	11.557 ± 0.005	22.95 ± 0.07
c.2 $pp \rightarrow (W^+ \rightarrow)e^+\nu_e t\bar{t}$	$m_W + 2m_{top}$	5	0.009415 ± 0.000003	0.01159 ± 0.00001
c.3 $pp \rightarrow (\gamma^*/Z \rightarrow)e^+e^- b\bar{b}$	$m_Z + 2m_b$	4	9.459 ± 0.004	15.31 ± 0.03
c.4 $pp \rightarrow (\gamma^*/Z \rightarrow)e^+e^- t\bar{t}$	$m_Z + 2m_{top}$	5	0.0035131 ± 0.0000004	0.004876 ± 0.000002
c.5 $pp \rightarrow \gamma t\bar{t}$	$2m_{top}$	5	0.2906 ± 0.0001	0.4169 ± 0.0003
d.1 $pp \rightarrow W^+W^-$	$2m_W$	4	29.976 ± 0.004	43.92 ± 0.03
d.2 $pp \rightarrow W^+W^- j$	$2m_W$	4	11.613 ± 0.002	15.174 ± 0.008
d.3 $pp \rightarrow W^+W^- jj$	$2m_W$	4	0.07048 ± 0.00004	0.1377 ± 0.0005
e.1 $pp \rightarrow HW^+$	$m_W + m_H$	5	0.3428 ± 0.0003	0.4455 ± 0.0003
e.2 $pp \rightarrow HW^+ j$	$m_W + m_H$	5	0.1223 ± 0.0001	0.1501 ± 0.0002
e.3 $pp \rightarrow HZ$	$m_Z + m_H$	5	0.2781 ± 0.0001	0.3659 ± 0.0002
e.4 $pp \rightarrow HZ j$	$m_Z + m_H$	5	0.0988 ± 0.0001	0.1237 ± 0.0001
e.5 $pp \rightarrow Ht\bar{t}$	$m_{top} + m_H$	5	0.08896 ± 0.00001	0.09869 ± 0.00003
e.6 $pp \rightarrow Hb\bar{b}$	$m_b + m_H$	4	0.16510 ± 0.00009	0.2099 ± 0.0006
e.7 $pp \rightarrow Hjj$	m_H	5	1.104 ± 0.002	1.036 ± 0.002

- Attempts to bring same easy of use as Madgraph
- Combined with MadFKS to deal with IR divergences
- Currently limited to relatively low number of legs
- No public code available yet

NLO Conclusions

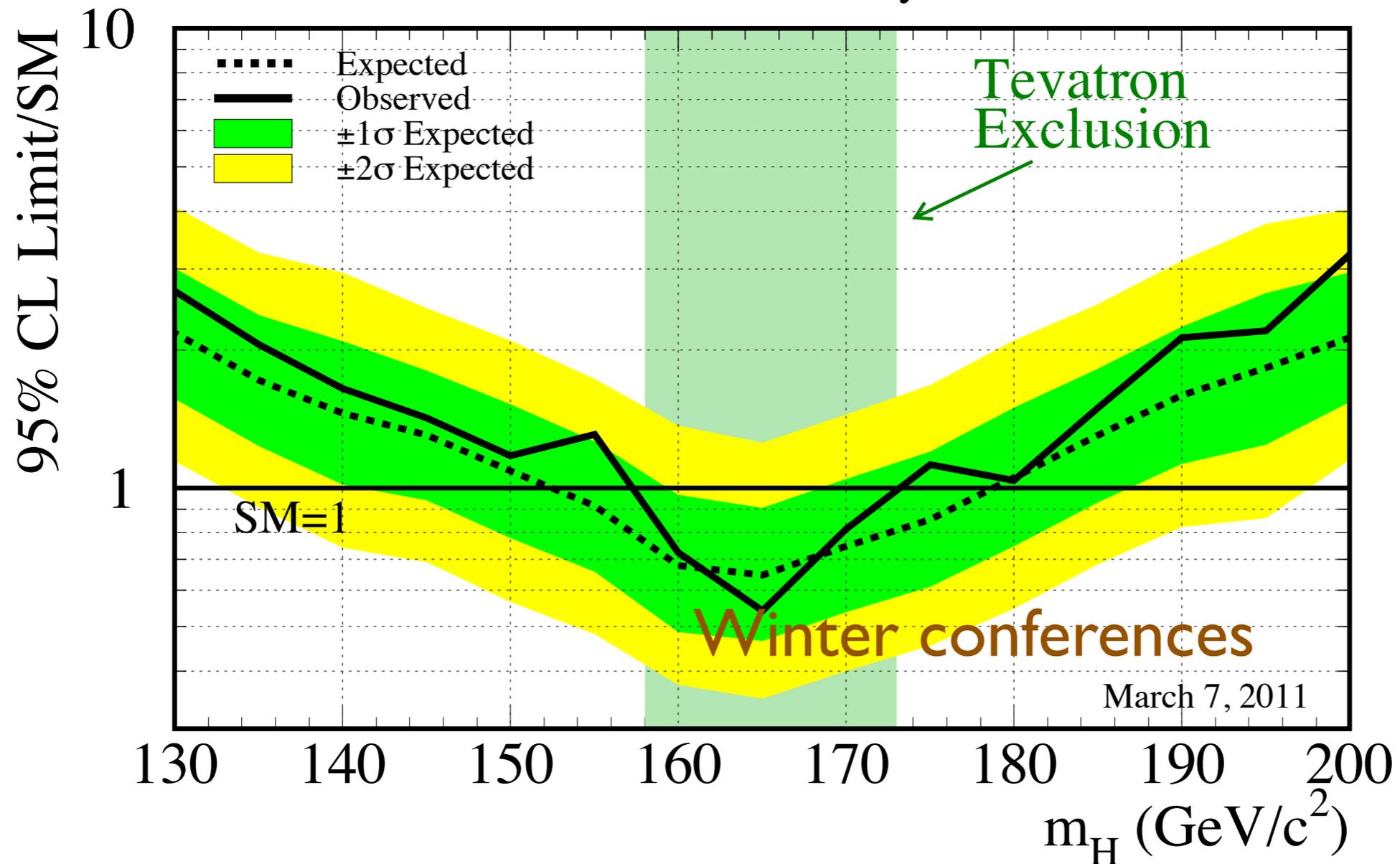
- For any reasonable precision, need at least NLO
- While procedure to calculate well known, we have hit the limit on analytical calculability
- Numerical calculations face issues of divergent integrals
- Tremendous progress over last few years for numerical techniques
- Has allowed results for previously unimaginable processes
- Much progress on complete automatization of NLO

Large logarithms

Searching for the Higgs

Tevatron has sensitivity to Higgs, mainly through $gg \rightarrow H \rightarrow WW$

Tevatron Run II Preliminary, $L \leq 8.2 \text{ fb}^{-1}$



Clearly, exclusion depends on the uncertainties assumed



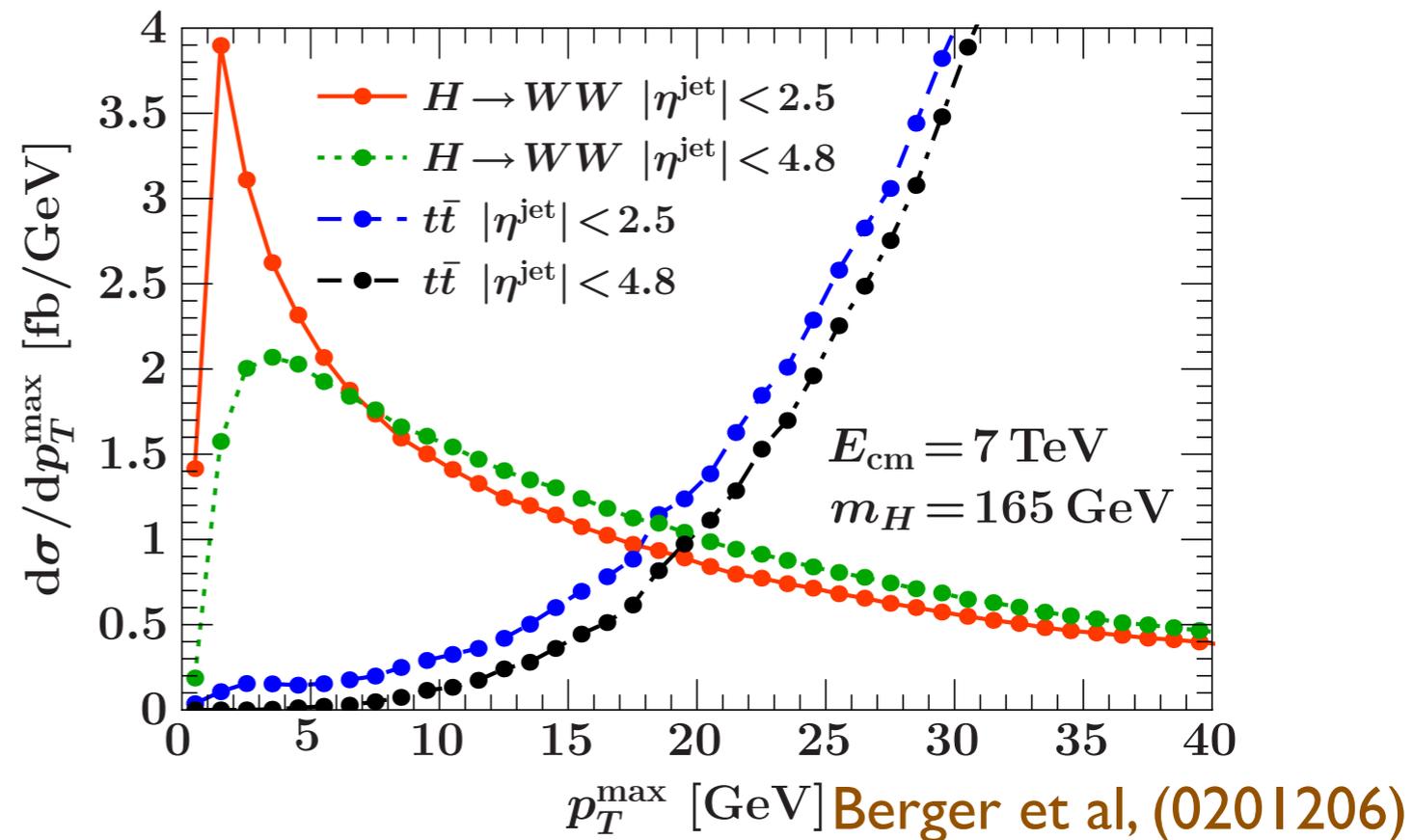
σ_{inc} [fb]	LO	NLO	NNLO	K^{NLO}	K^{NNLO}
$\mu = m_H/2$	1.998 ± 0.003	4.288 ± 0.004	5.252 ± 0.016	2.149 ± 0.008	2.629 ± 0.009
$\mu = m_H$	1.398 ± 0.001	3.366 ± 0.003	4.630 ± 0.010	2.412 ± 0.002	3.312 ± 0.008
$\mu = 2m_H$	1.004 ± 0.001	2.661 ± 0.002	4.012 ± 0.007	2.651 ± 0.008	3.996 ± 0.008

Anastasiou et al (0801.2682)

Scale variation at NNLO 14%

Large bkgr is tt production
 $pp \rightarrow tt \rightarrow bbWW \rightarrow bbl\nu\ell\nu$

Bottom quarks give rise to
 extra jet activity



Different amounts of jet activity should be treated differently

Fixed order predictions

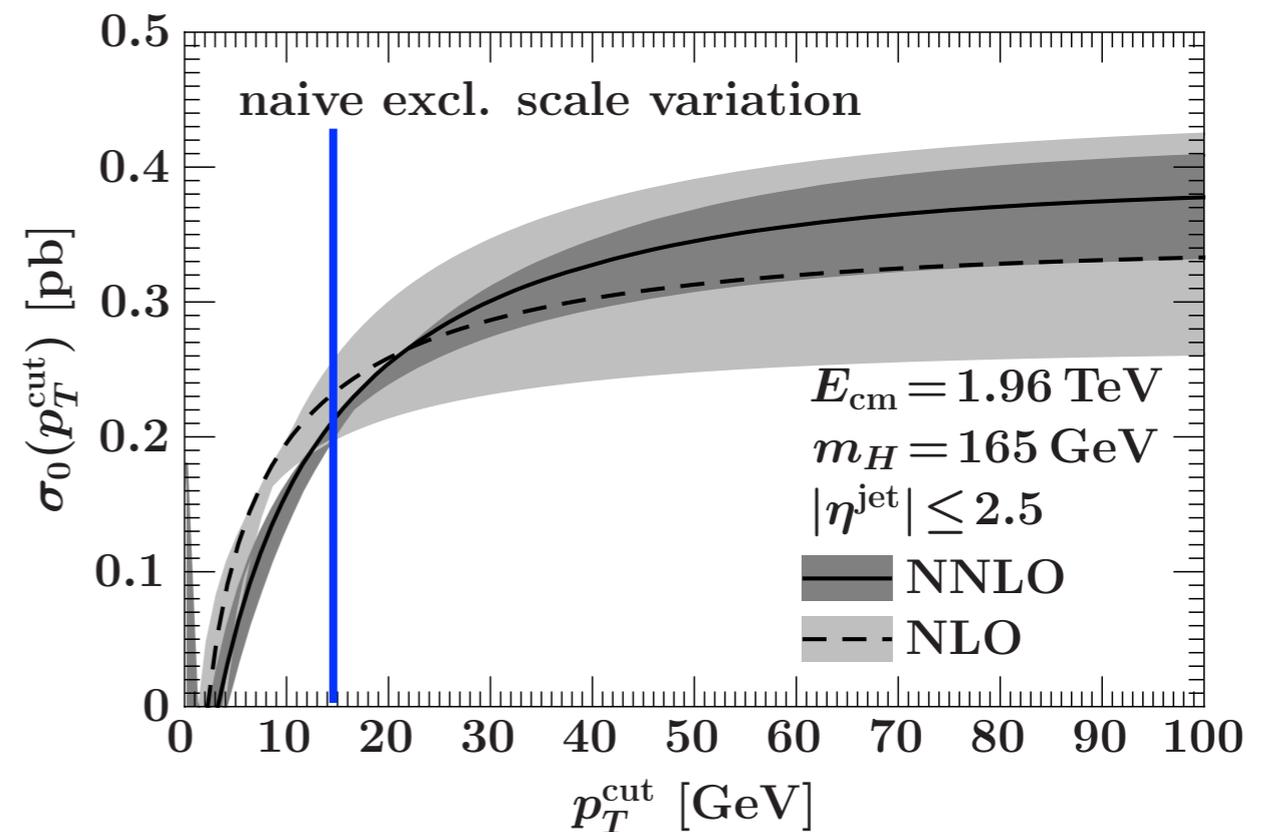
Results used by Tevatron winter analysis

σ [fb]	LO (pdfs, α_s)	NLO (pdfs, α_s)	NNLO (pdfs, α_s)
0-jets	$3.452^{+7\%}_{-10\%}$	$2.883^{+4\%}_{-9\%}$	$2.707^{+5\%}_{-9\%}$
1-jet	$1.752^{+30\%}_{-26\%}$	$1.280^{+24\%}_{-23\%}$	$1.165^{+24\%}_{-22\%}$
≥ 2 -jets	$0.336^{+91\%}_{-44\%}$	$0.221^{+81\%}_{-42\%}$	$0.196^{+78\%}_{-41\%}$

Anastasiou et al (0801.2682)

$$\frac{\Delta\sigma_{\text{total}}}{\sigma_{\text{total}}} = 66.5\% \times \begin{pmatrix} +5\% \\ -9\% \end{pmatrix} + 28.6\% \times \begin{pmatrix} +24\% \\ -22\% \end{pmatrix} + 4.9\% \times \begin{pmatrix} +78\% \\ -41\% \end{pmatrix} = \begin{pmatrix} +14\% \\ -14\% \end{pmatrix}.$$

How much can we trust uncertainties in individual channels?



Exclusive x-section has smaller uncertainties than inclusive one?

The perturbative expansion

Stewart, Tackmann, (1107.2117)

Inclusive cross section has form

$$\sigma_{\text{total}} = (0.15 \text{ pb}) [1 + 9.0 \alpha_s + 34 \alpha_s^2 + \mathcal{O}(\alpha_s^3)]$$

poorly behaved series

□ large K factor

Zero jet cross section has form

$$\sigma_0 (15 \text{ GeV}) = (0.15 \text{ pb}) [1 + 3.7 \alpha_s - 2 \alpha_s^2]$$

seems like much better behaved series

Should remember a simple relation

$$\sigma_0 (p^{\text{cut}}) = \sigma_{\text{total}} - \sigma_{\geq 1} (p^{\text{cut}})$$

$$\sigma_{\text{total}} \simeq \sigma_B [1 + \alpha_s + \alpha_s^2]$$

$$\sigma_{\geq 1} (p^{\text{cut}}) \simeq \sigma_B [\alpha_s (L^2 + L + 1) + \alpha_s^2 (L^4 + L^3 + L^2 + L + 1)]$$

$$L^2 = \ln^2(p^{\text{cut}}/Q)$$

Estimating uncertainties

Stewart, Tackmann, (1107.2117)

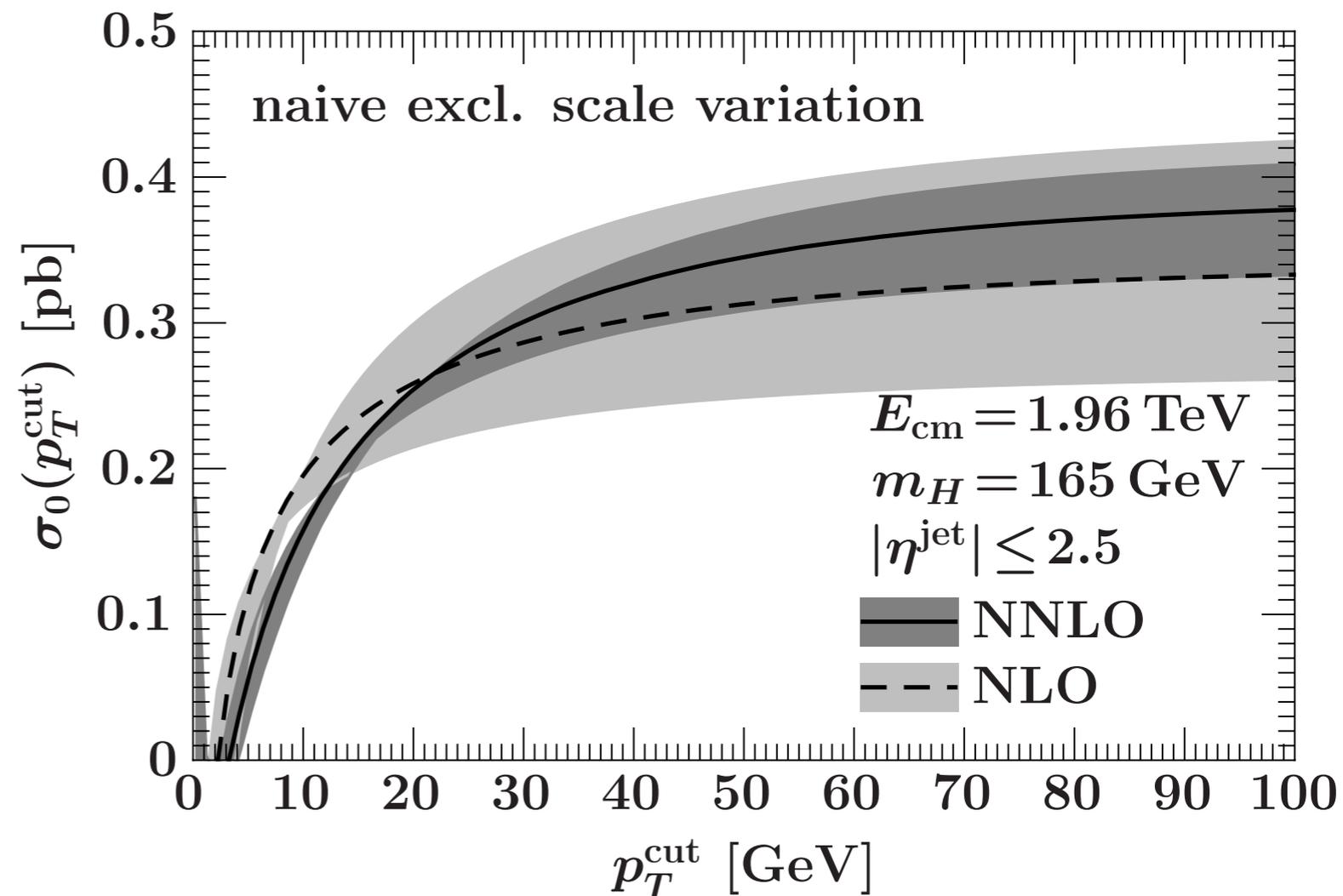
More conservative: Treat σ_{total} and $\sigma_{\geq 1}$ as uncorrelated, with separate uncertainties Δ_{total} and $\Delta_{\geq 1}$

Gives covariance matrix

$$\begin{matrix} \{ \sigma_{\text{total}}, \sigma_{\geq 1} \} \\ \begin{pmatrix} \Delta_{\text{total}}^2 & 0 \\ 0 & \Delta_{\geq 1}^2 \end{pmatrix} \end{matrix}$$

which implies that

$$\begin{matrix} \{ \sigma_{\text{total}}, \sigma_0, \sigma_{\geq 1} \} \\ \begin{pmatrix} \Delta_{\text{total}}^2 & \Delta_{\text{total}}^2 & 0 \\ \Delta_{\text{total}}^2 & \Delta_{\geq 1}^2 + \Delta_{\text{total}}^2 & -\Delta_{\geq 1}^2 \\ 0 & -\Delta_{\geq 1}^2 & \Delta_{\geq 1}^2 \end{pmatrix} \end{matrix}$$



Estimating uncertainties

Stewart, Tackmann, (1107.2117)

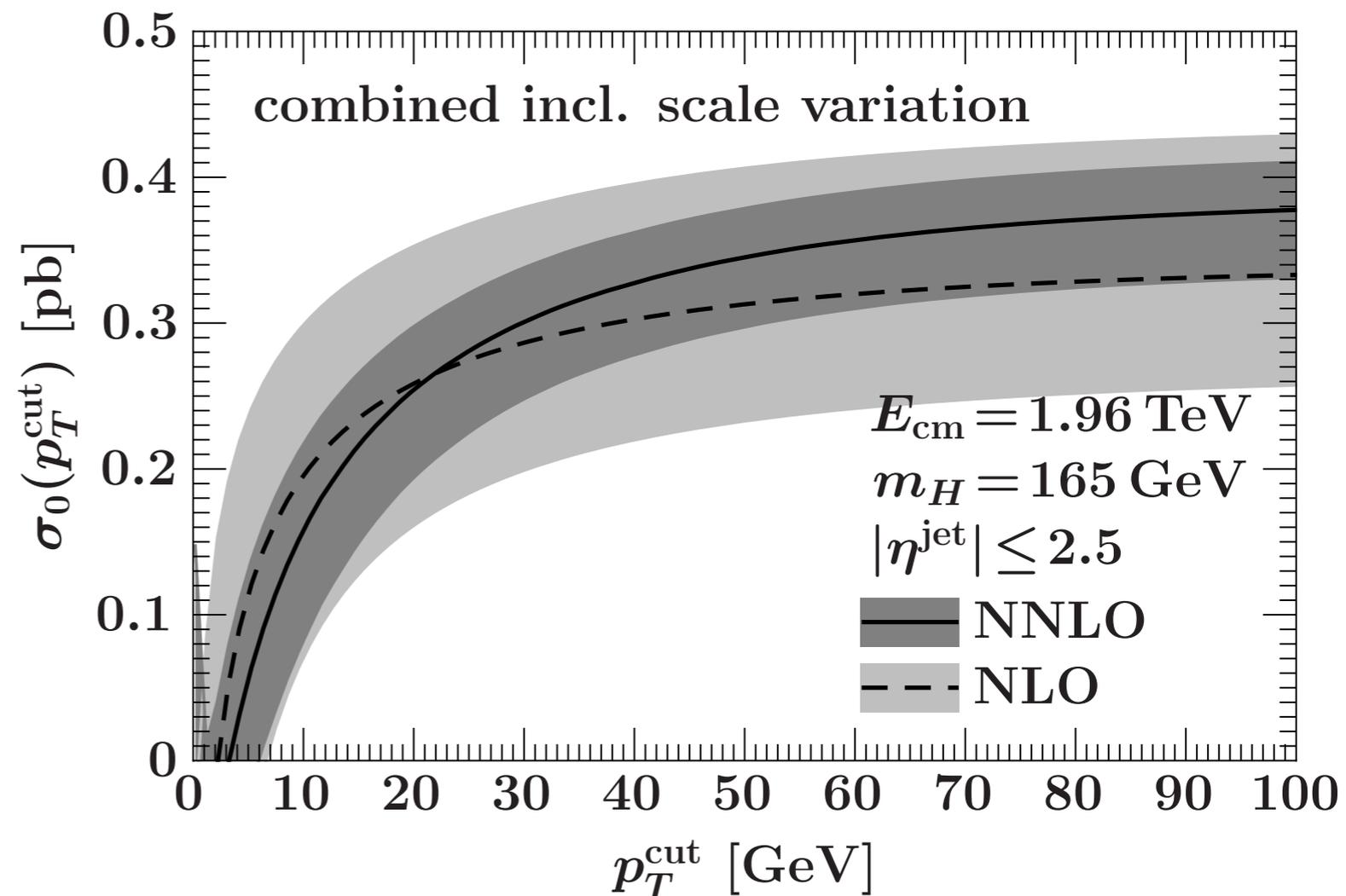
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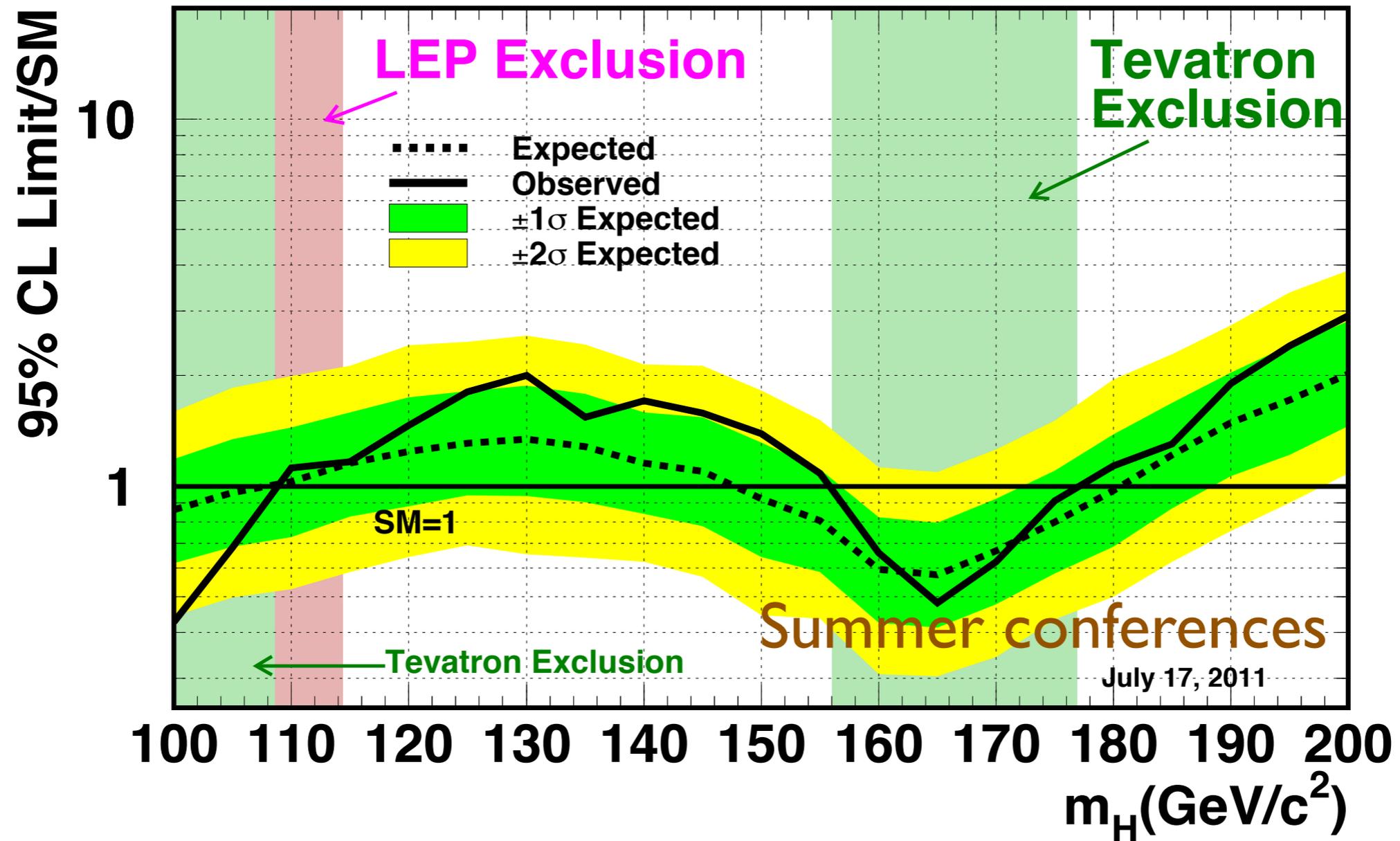
$$\begin{matrix} \{ \sigma_{\text{total}}, \sigma_0, \sigma_{\geq 1} \} \\ \begin{pmatrix} \Delta_{\text{total}}^2 & \Delta_{\text{total}}^2 & 0 \\ \Delta_{\text{total}}^2 & \Delta_{\geq 1}^2 + \Delta_{\text{total}}^2 & -\Delta_{\geq 1}^2 \\ 0 & -\Delta_{\geq 1}^2 & \Delta_{\geq 1}^2 \end{pmatrix} \end{matrix}$$



Effect on Tevatron analysis?

Procedure is used in latest Tevatron combinations

Tevatron Run II Preliminary, $L \leq 8.6 \text{ fb}^{-1}$



Used smaller overall uncertainties for inclusive cross-section

Log Conclusions

- Large ratios of scales \Rightarrow large logs in perturbative series
- Poorer convergence of series
- Gives rise to larger perturbative uncertainties
- Important to treat this effect carefully

Merging of fixed order calculations with parton showers

The need for merging

To combine predictions @ short distances with important long distance physics, need to merge fixed order calcs and parton showers

Problem is solved for LO accuracy ...

Catani, Krauss, Kuhn, Webber (0109231)

... and implemented in many computer codes

Alpgen, Madgraph, Sherpa, ...

Has been completely automated

For many analyses this is now standard theoretical tool

Idea behind merging

ϕ	
\vdots	
n+3-jets	
n+2-jets	
n+1-jets	
n-jets	

Idea behind merging

ϕ	
\vdots	
n+3-jets	
n+2-jets	
n+1-jets	
n-jets	FO

Idea behind merging

ϕ	
\vdots	\vdots
n+3-jets	PS
n+2-jets	PS
n+1-jets	PS
n-jets	FO

Idea behind merging

ϕ	S_I
\vdots	\vdots
n+3-jets	PS
n+2-jets	PS
n+1-jets	PS
n-jets	FO

Idea behind merging

ϕ	S_1	
\vdots	\vdots	
n+3-jets	PS	
n+2-jets	PS	
n+1-jets	PS	FO
n-jets	FO	

Idea behind merging

ϕ	S_1	
\vdots	\vdots	\vdots
n+3-jets	PS	PS
n+2-jets	PS	PS
n+1-jets	PS	FO
n-jets	FO	

Idea behind merging

ϕ	S_1	S_2
\vdots	\vdots	\vdots
n+3-jets	PS	PS
n+2-jets	PS	PS
n+1-jets	PS	FO
n-jets	FO	

Idea behind merging

ϕ	S_1	S_2	S_3
\vdots	\vdots	\vdots	\vdots
n+3-jets	PS	PS	PS
n+2-jets	PS	PS	FO
n+1-jets	PS	FO	
n-jets	FO		

Idea behind merging

ϕ	S_1	S_2	S_3	S_4
\vdots	\vdots	\vdots	\vdots	\vdots
n+3-jets	PS	PS	PS	FO
n+2-jets	PS	PS	FO	
n+1-jets	PS	FO		
n-jets	FO			

Idea behind merging

ϕ	S_1	S_2	S_3	S_4
\vdots	\vdots	\vdots	\vdots	\vdots
n+3-jets	PS	PS	PS	FO
n+2-jets	PS	PS	FO	
n+1-jets	PS	FO		
n-jets	FO			

Want to
combine all
samples
together

Idea behind merging

ϕ	S_1	S_2	S_3	S_4
\vdots	\vdots	\vdots	\vdots	\vdots
n+3-jets	PS	PS	PS	FO
n+2-jets	PS	PS	FO	
n+1-jets	PS	FO		
n-jets	FO			

Want to combine all samples together

Main Issue:

Avoid double counting:
Ensure that sum of all samples S_i still gives right answer for $m > n$ jets

Two different implementations possible...

Idea behind merging

ϕ	S_1	S_2	S_3	S_4
\vdots	\vdots	\vdots	\vdots	\vdots
n+3-jets	PS	PS	PS	FO
n+2-jets	PS	PS	FO	
n+1-jets	PS	FO		
n-jets	FO			

Want to combine all samples together

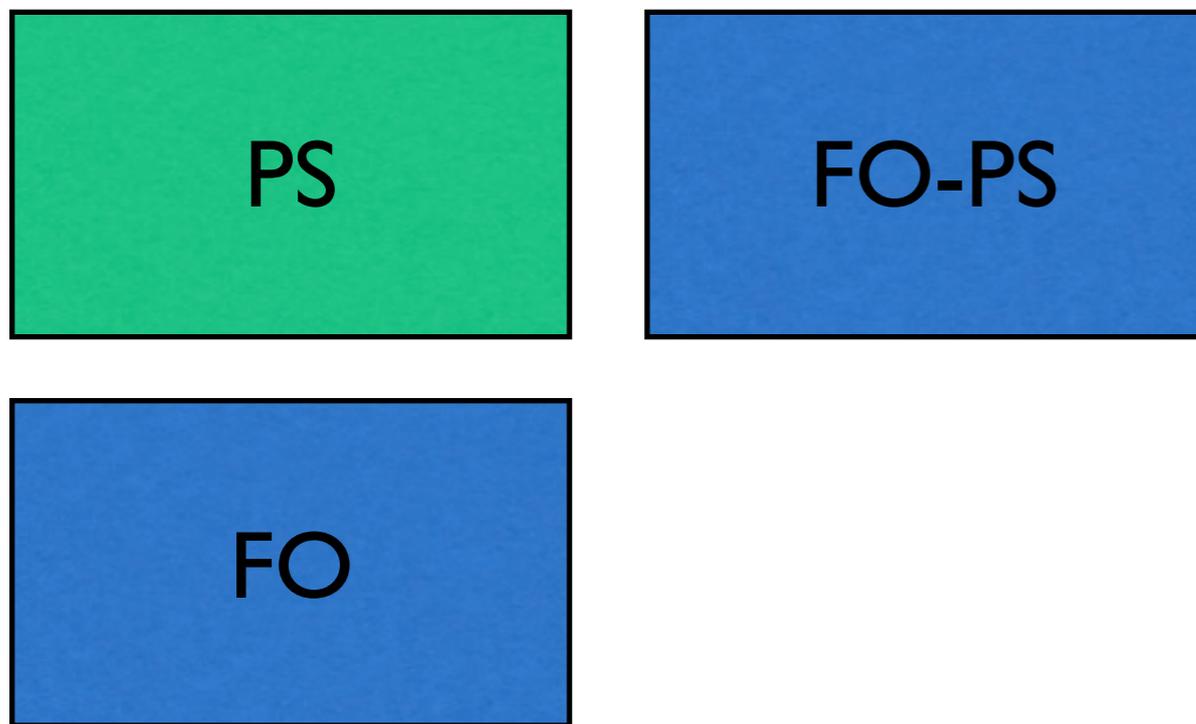
Main Issue:

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Two different implementations possible...

Two main choices

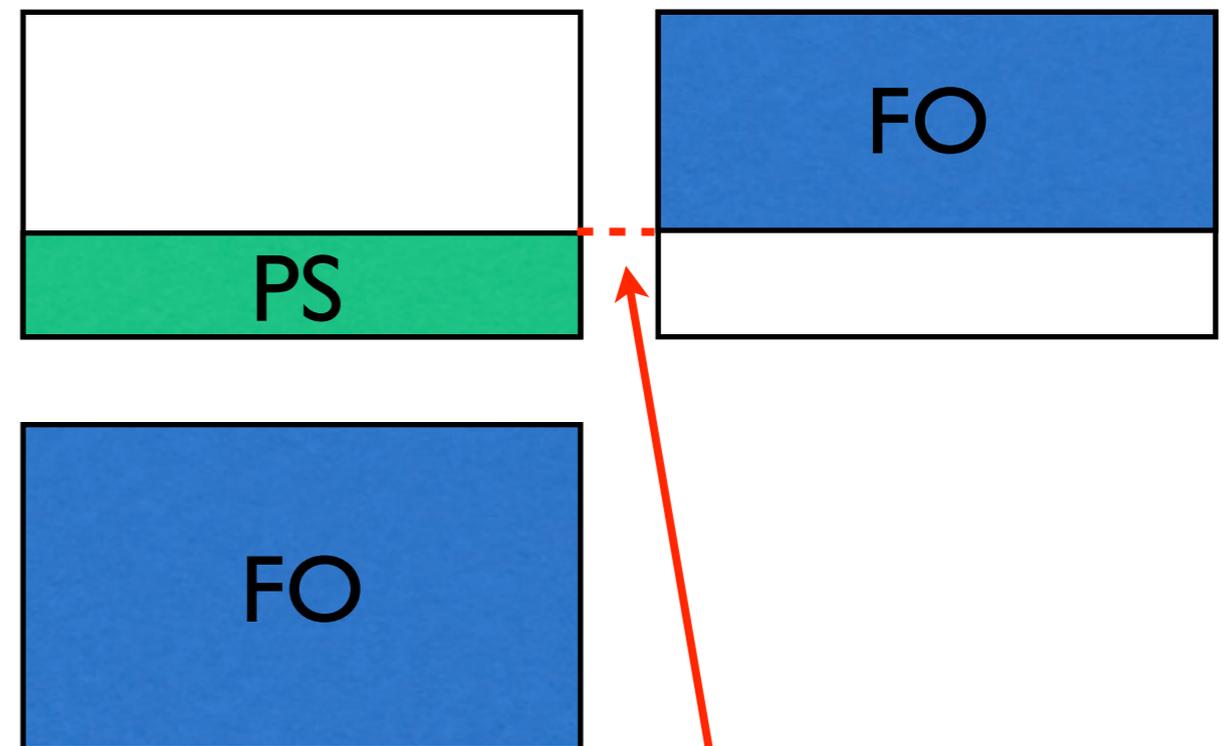
Subtraction



Need to know analytical expressions for parton shower

LO-PS can be negative!

Phase space separation



Need to be independent of arbitrary phase space separation

More difficult to sum correct logs

Merging at LO

As already mentioned, worked out at LO 10 years ago

Catani, Krauss, Kuhn, Webber (0109231); AlpGen, MadGraph, Sherpa,

ϕ	S_1	S_2	S_3	S_4
\vdots	\vdots	\vdots	\vdots	\vdots
n+3-jets	PS	PS	PS	LO
n+2-jets	PS	PS	LO	
n+1-jets	PS	LO		
n-jets	LO			

All approaches use phase space separation

Merging at NLO

NLO methods available since

ϕ	S_1	S_2	
\vdots	\vdots	\vdots	
n+3-jets	PS	PS	<u>Subtraction:</u> MC@NLO
n+2-jets	PS	PS	
n+1-jets	PS	LO	<u>Phase space separation:</u> POWHEG
n-jets	NLO		

Both Subtraction and phase space separation used

General framework to implement new processes (POWHEG Box)

Ongoing work to completely automate process (aMC@NLO)

Merging at LO

Recent ideas allow to combine this with LO at higher multiplicity

ϕ	S_1	S_2	S_3	S_4
\vdots	\vdots	\vdots	\vdots	\vdots
n+3-jets	PS	PS	PS	LO
n+2-jets	PS	PS	LO	
n+1-jets	PS	LO		
n-jets	NLO			

CWB, Tackmann, Thaler (0801.4026),
 Hamilton, Nason (2010)
 Höche, Krauss, Schönherr, Siegert (2010)

Merging at LO

Ongoing work to have NLO for all multiplicities

CWB, with Geneva collaboration

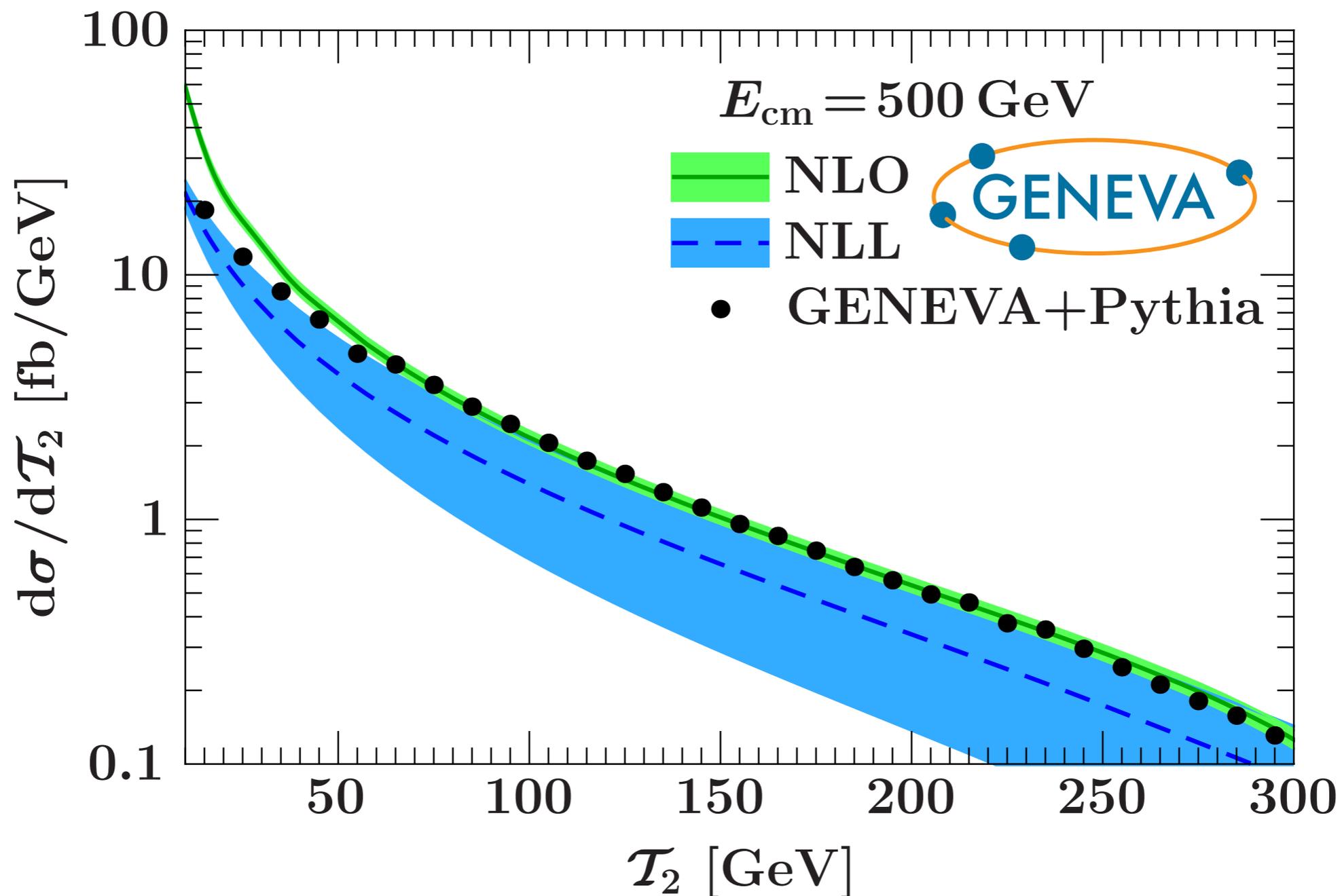
ϕ	S_1	S_2	S_3	S_4
\vdots	\vdots	\vdots	\vdots	\vdots
n+3-jets	PS	PS	PS	NLO
n+2-jets	PS	PS	NLO	
n+1-jets	PS	NLO		
n-jets	NLO			

Made possible
through applications
of SCET to event generation

CWB, Fleming, Luke, Pirjol Stewart;
CWB, Schwartz (0607296)

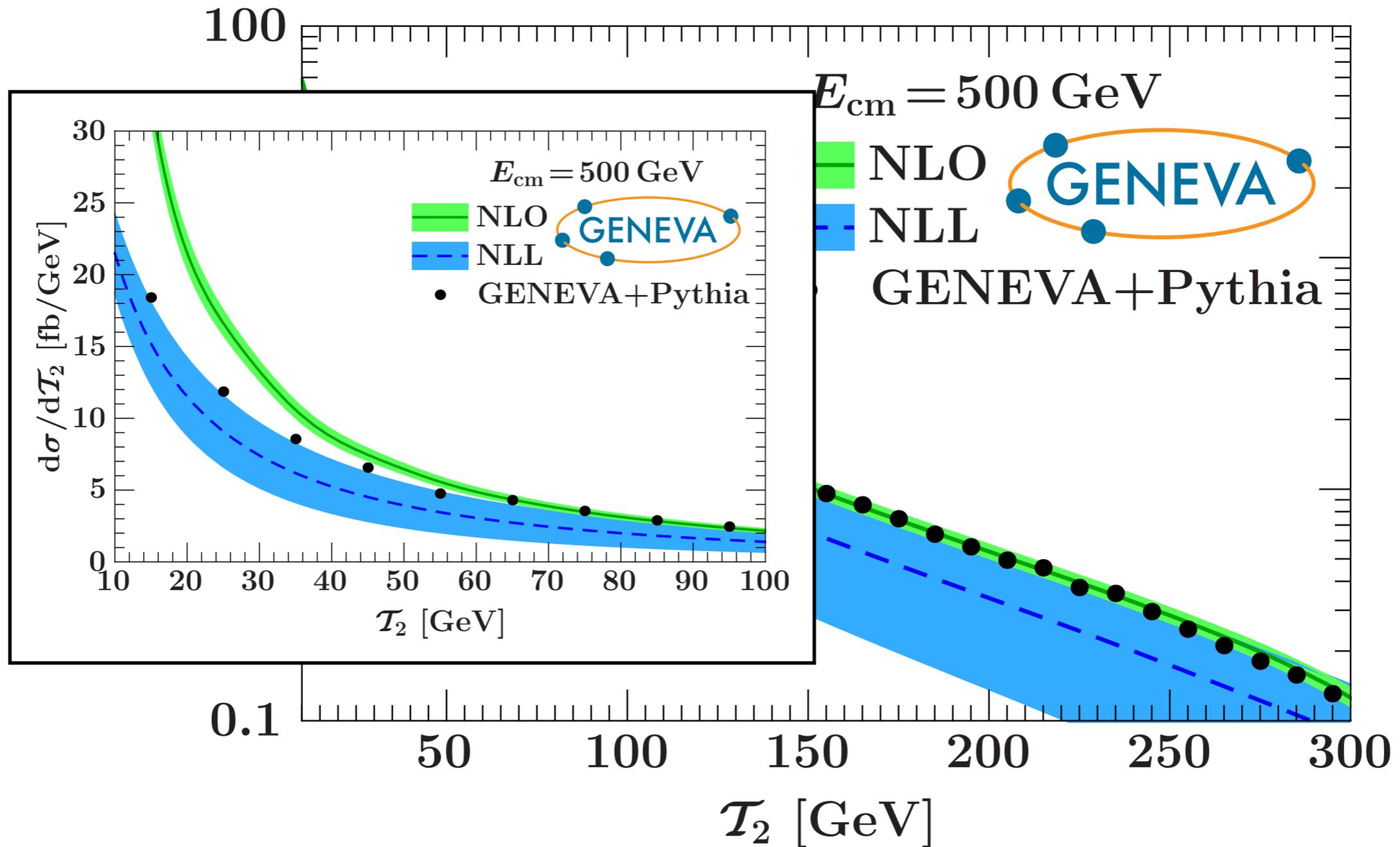
Preliminary Geneva

Event shape in $e^+e^- \rightarrow \text{jets}$, with both 2j and 3j final states correct to NLO



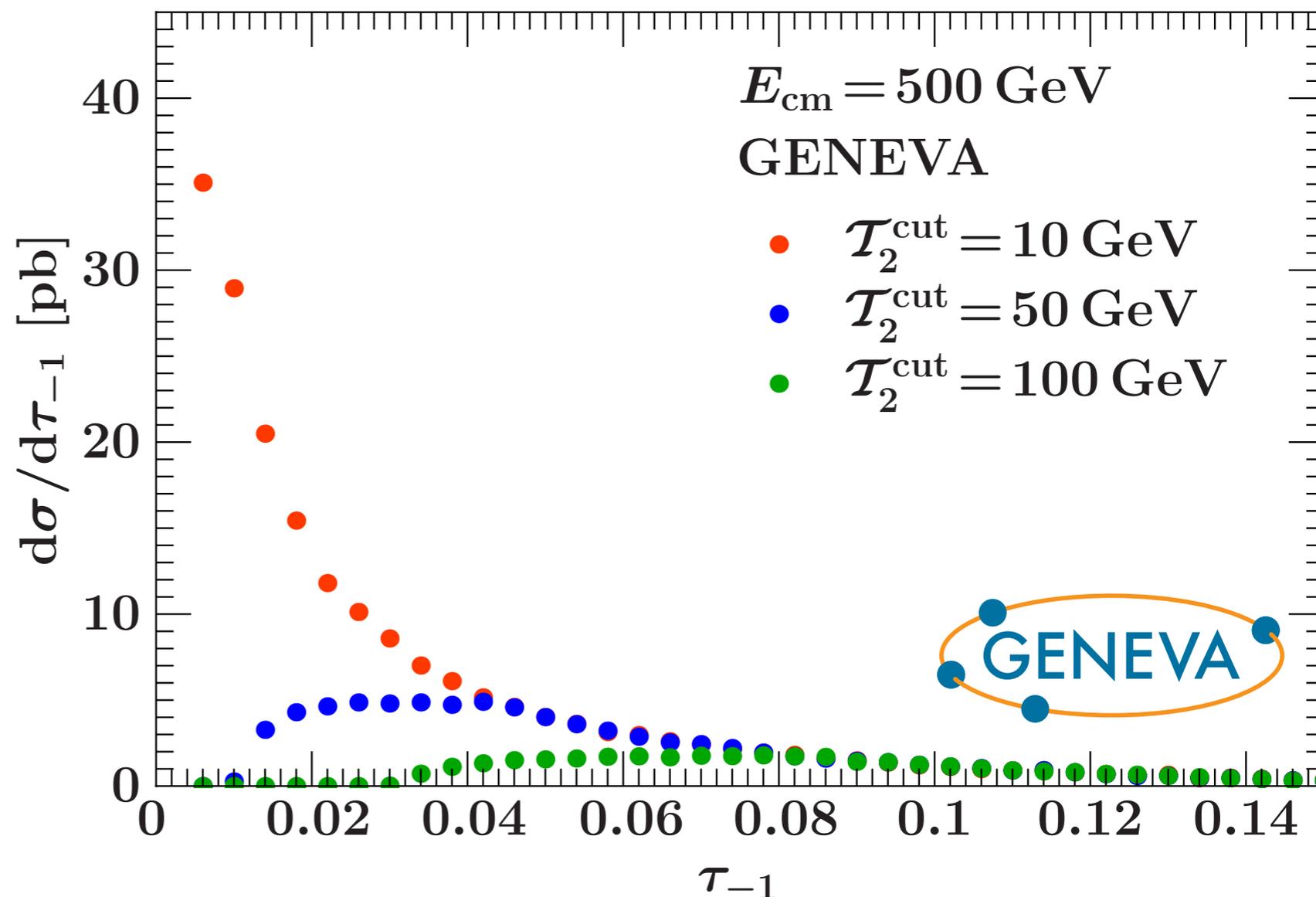
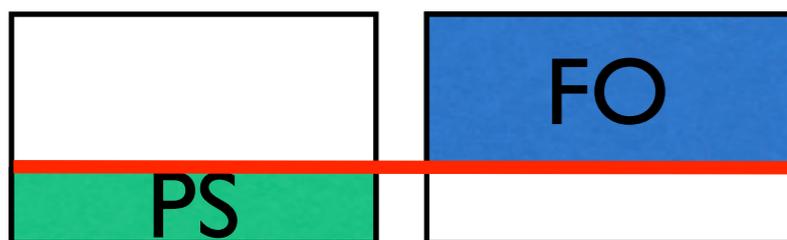
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Event shape in $e^+e^- \rightarrow \text{jets}$, with both 2j and 3j final states correct to NLO



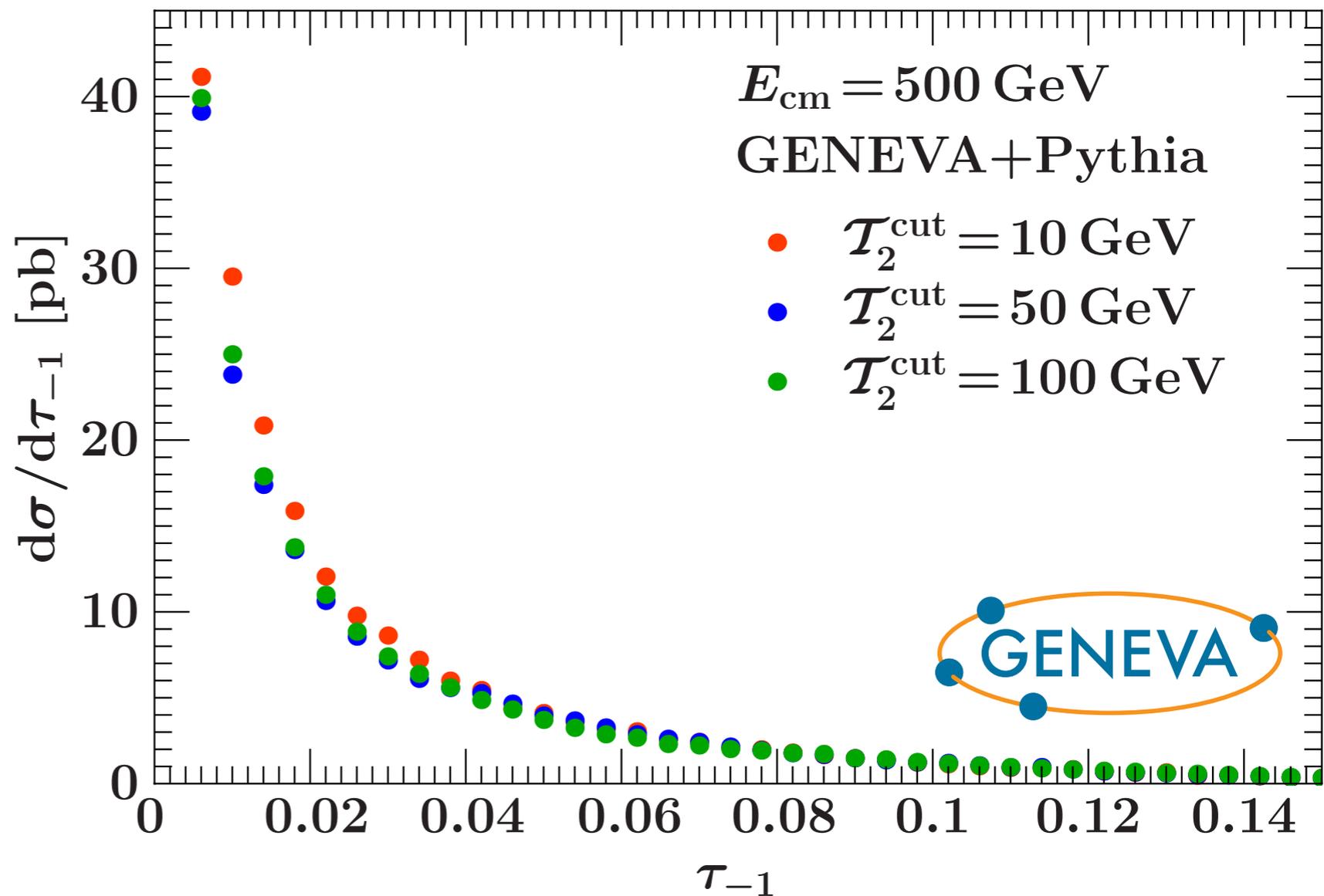
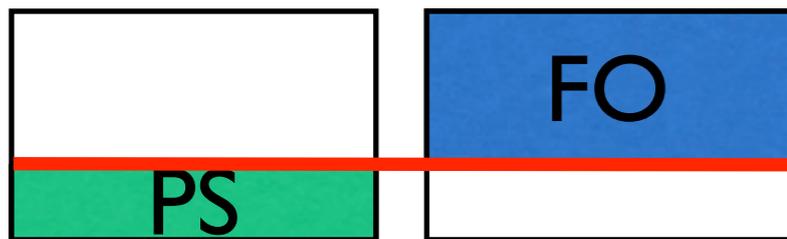
Preliminary Geneva

Can also study the cancellation of the dependence on merging scale



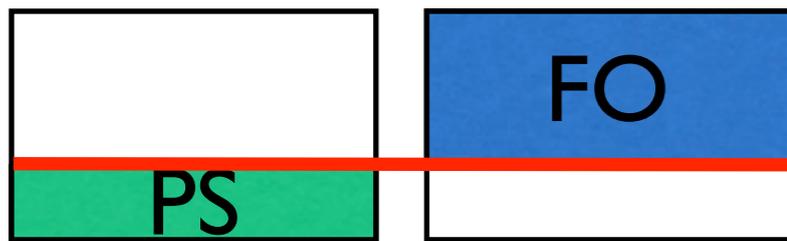
Preliminary Geneva

Can also study the cancellation of the dependence on merging scale

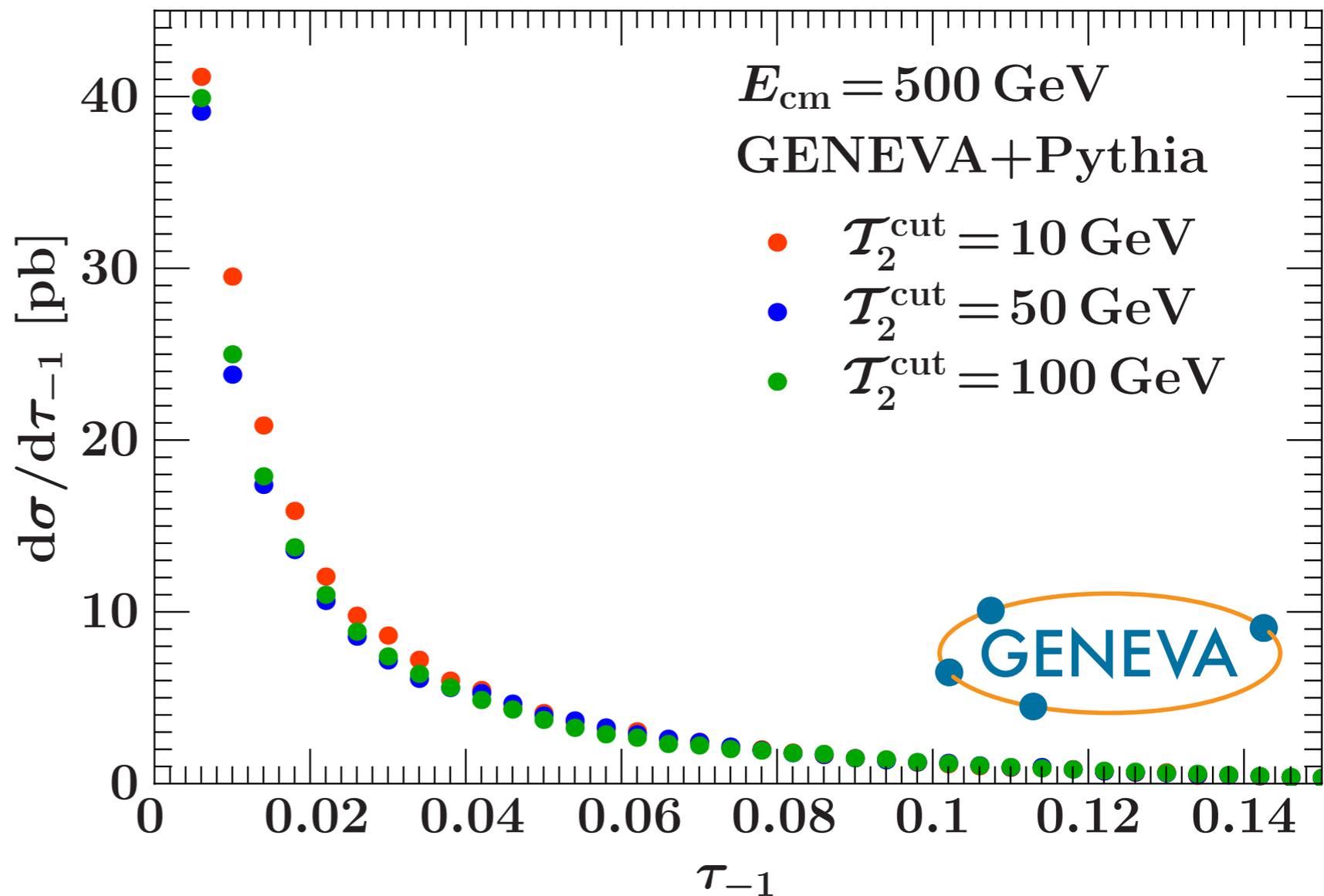


Preliminary Geneva

Can also study the cancellation of the dependence on merging scale

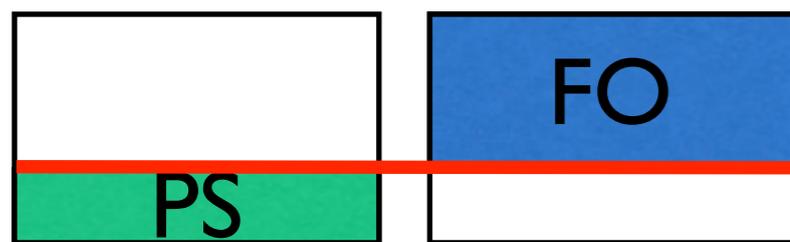


Proper cancellation once PS is added

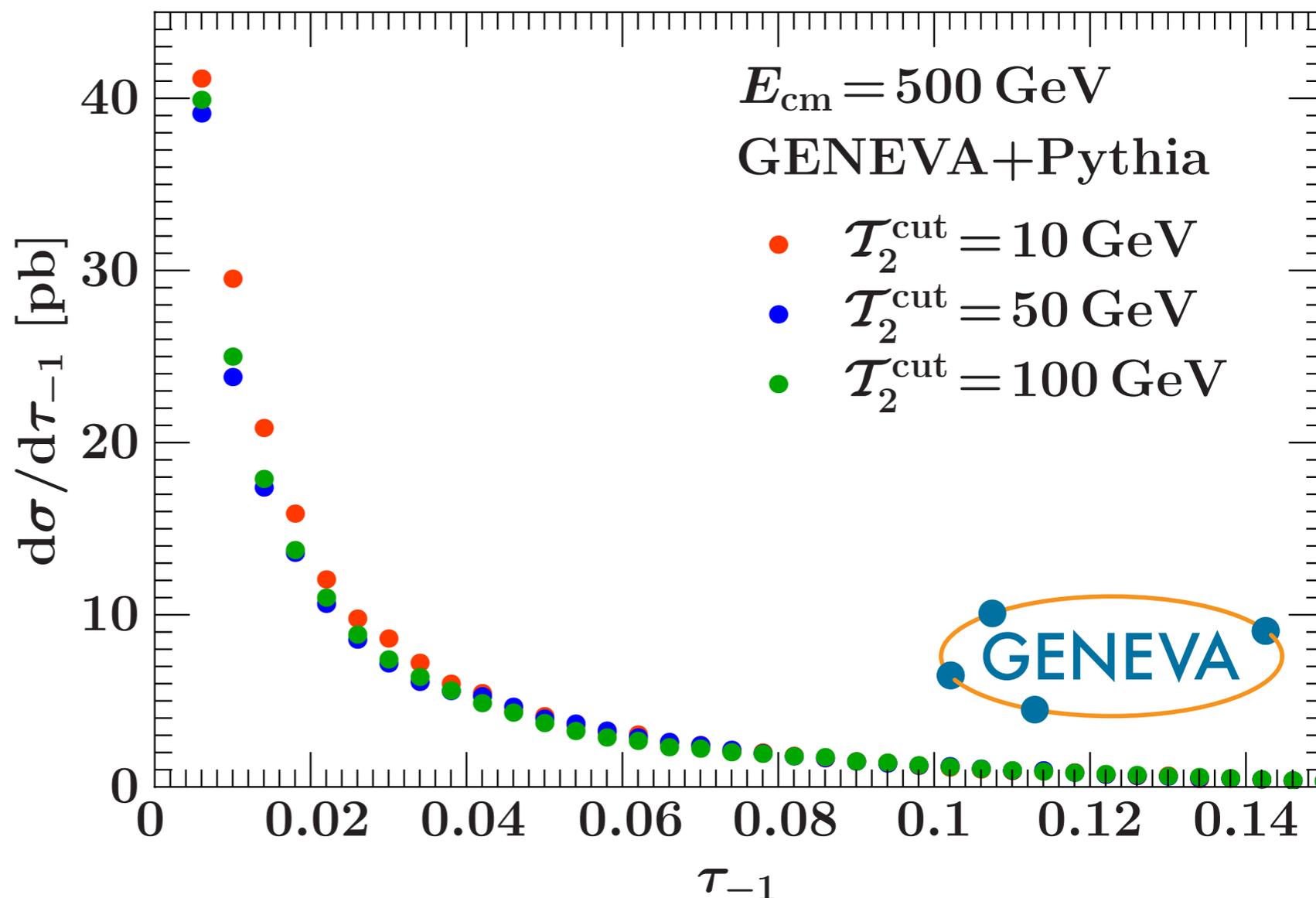


Preliminary Geneva

Can also study the cancellation of the dependence on merging scale



Proper cancellation once PS is added



All above easily generalized to pp collisions, in particular $W+0$ and $W+1$ jet production

Merging Conclusions

- Much progress in recent years on merging fixed order calculations with parton shower algorithms
- Merging of LO calculations completely standard
- Merging of NLO calculation getting quite mature
- Next goal is to find general way to combine multiple NLO calculations

Final Thoughts

QCD is a fascinating theory, which needs many different theoretical approaches

The LHC has given theorists a new problems to think about

Have only been able to show you a small selection of progress made in past years