

Physics Cases For Muon Colliders

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(Brown University, Aug. 11, 2011)

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It is these features: heavy mass, short lifetime that dictate the physics.

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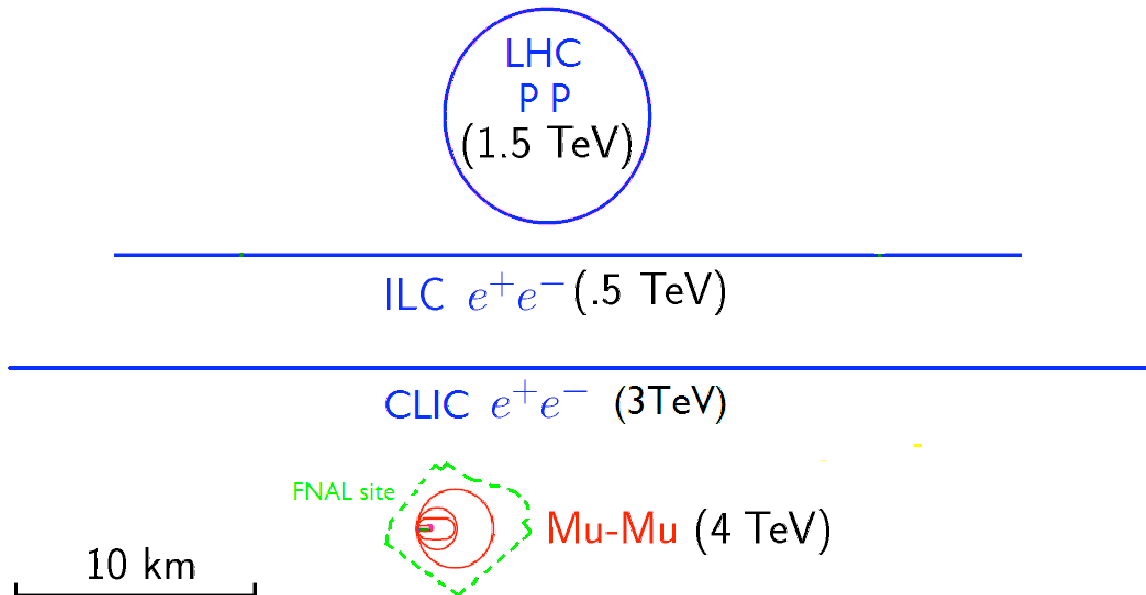
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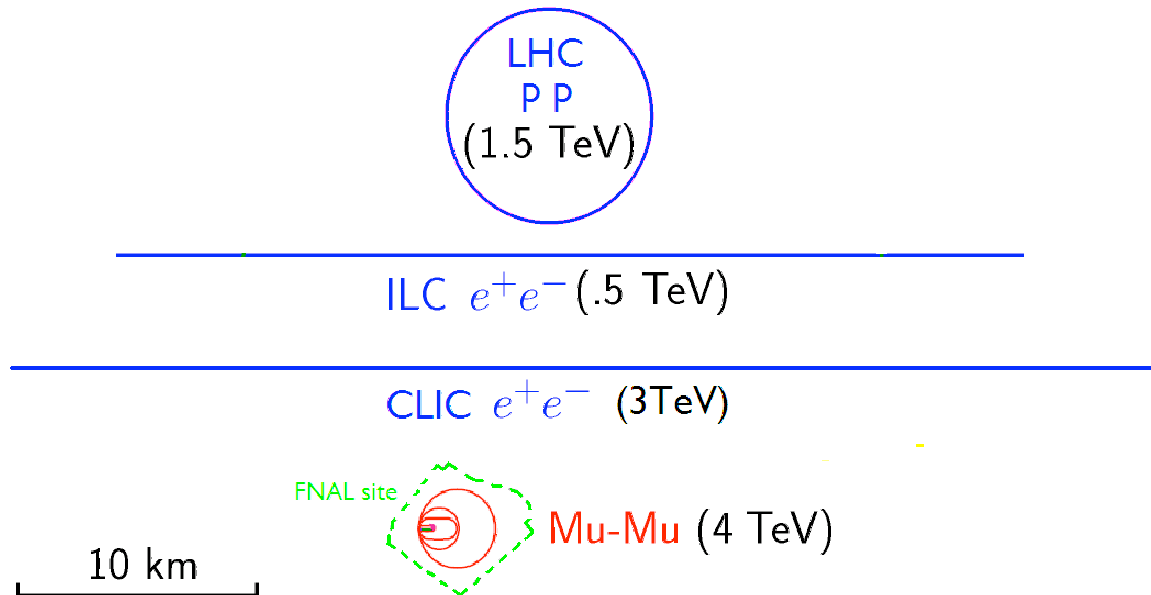
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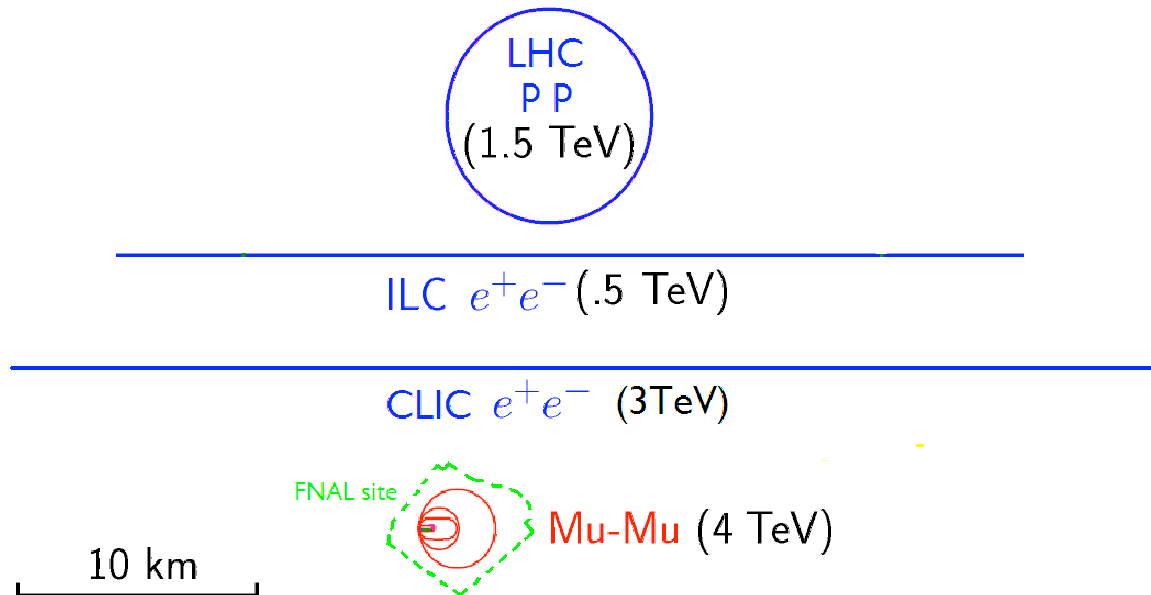
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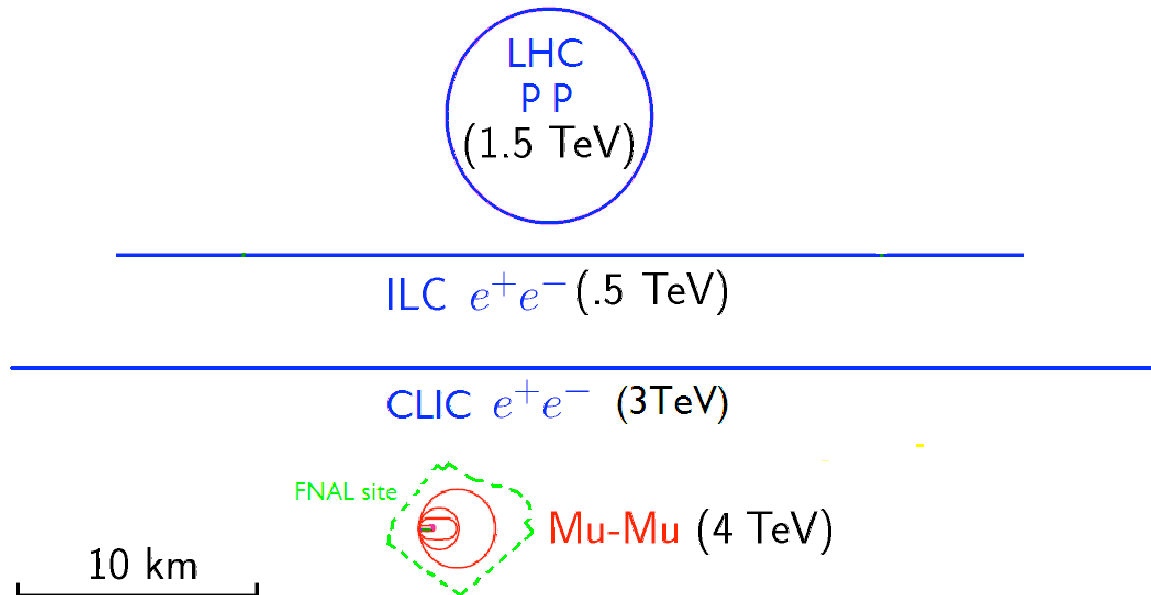
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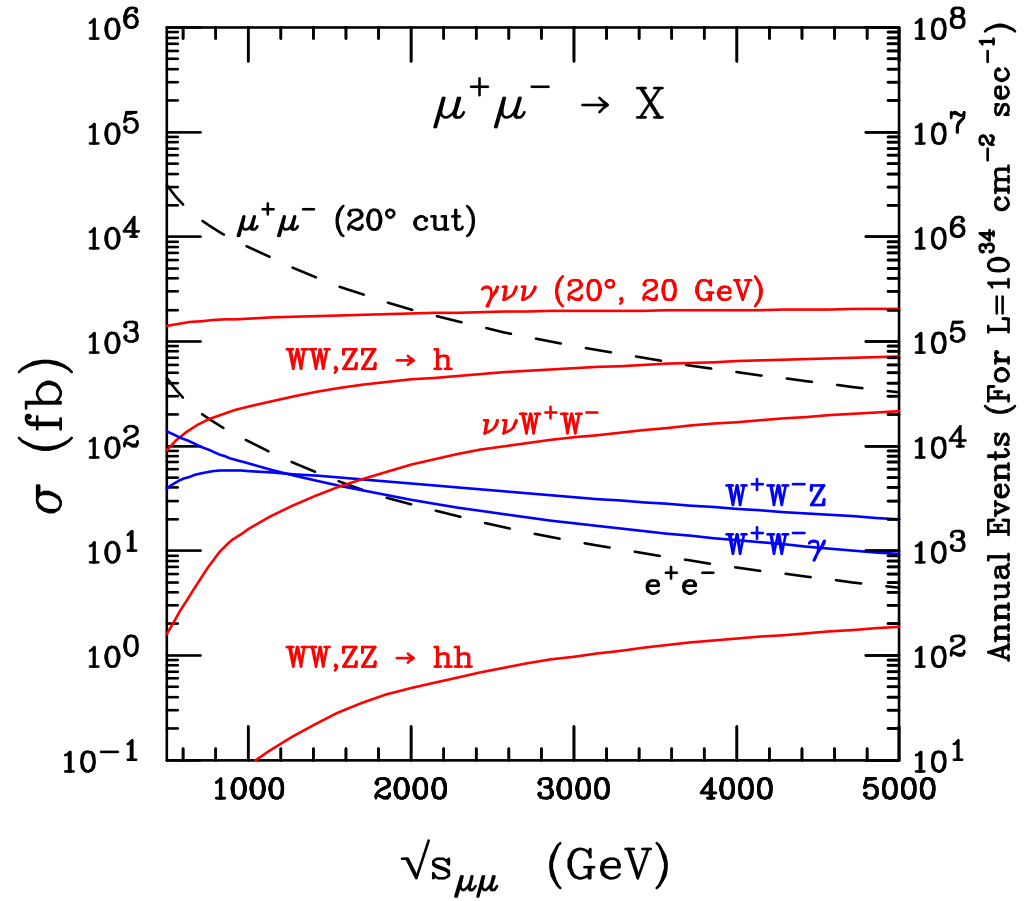
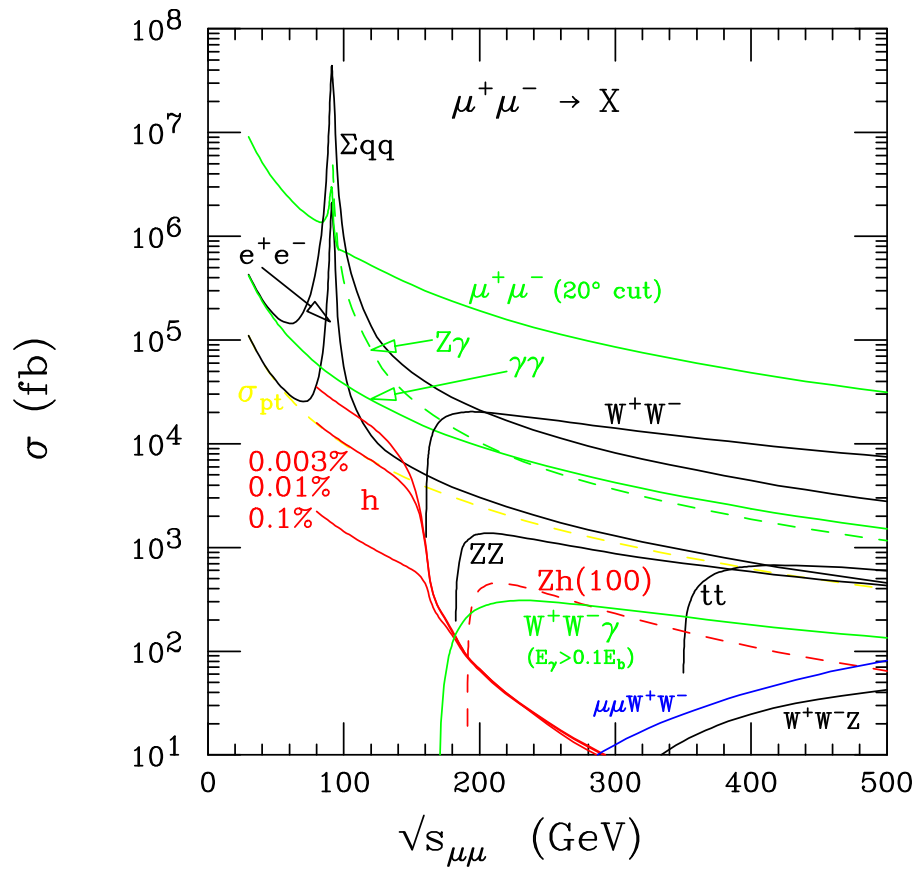
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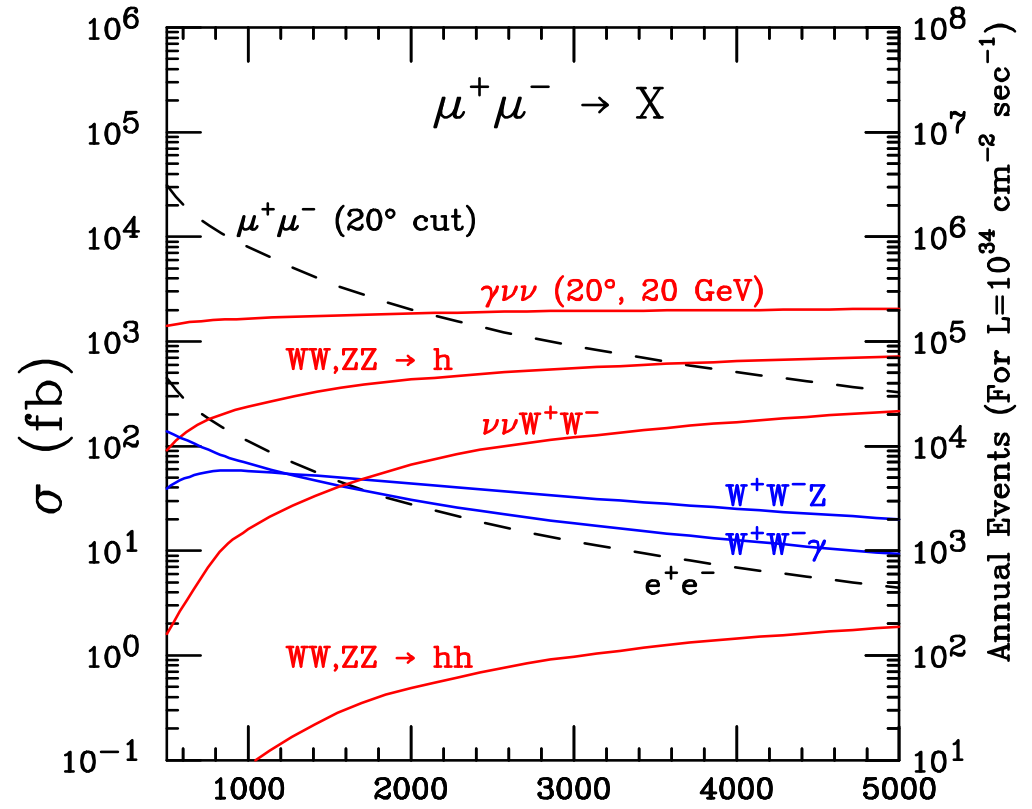
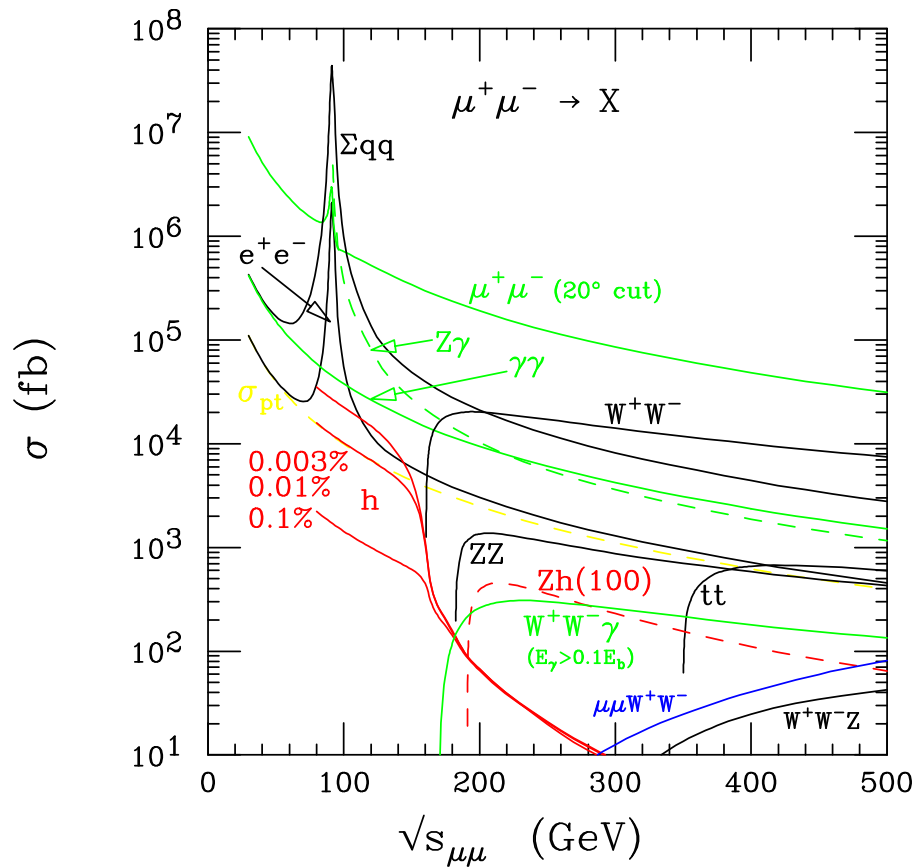
(3). Challenges: “Never play with an unstable thing!”
(see Ron Lipton’s talk next)

*Palmer

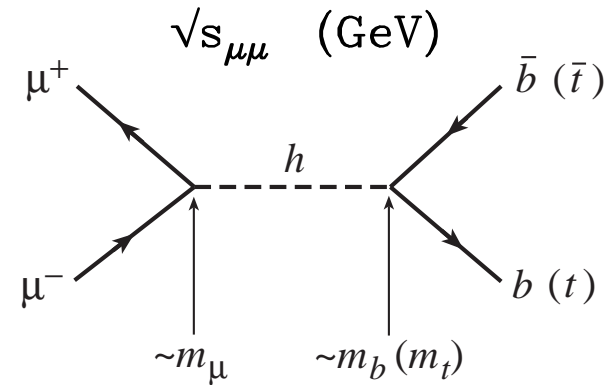
Physics at Muon Colliders



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Most unique of all at a muon collider:
 the s -channel scalar resonance.[†]



[†]Barger, Berger, Gunion, Han

A Higgs Factory

The s -channel resonant production:

$$\begin{aligned}\sigma(\mu^+\mu^- \rightarrow H, A \rightarrow X) &= \frac{4\pi\Gamma(H, A \rightarrow \mu^+\mu^-) \Gamma(H, A \rightarrow X)}{(s - M_H^2)^2 + \Gamma_H^2 M_H^2} \\ \bar{\sigma}(s) &= \int d\sqrt{s} \sigma(\mu^+\mu^- \rightarrow H, A \rightarrow X) \frac{dL}{d\sqrt{s}} \\ \delta E \ll \Gamma_H &\Rightarrow \frac{4\pi\Gamma(H, A \rightarrow \mu^+\mu^-) \Gamma(H, A \rightarrow X)}{\Gamma_H^2 M_H^2}.\end{aligned}$$

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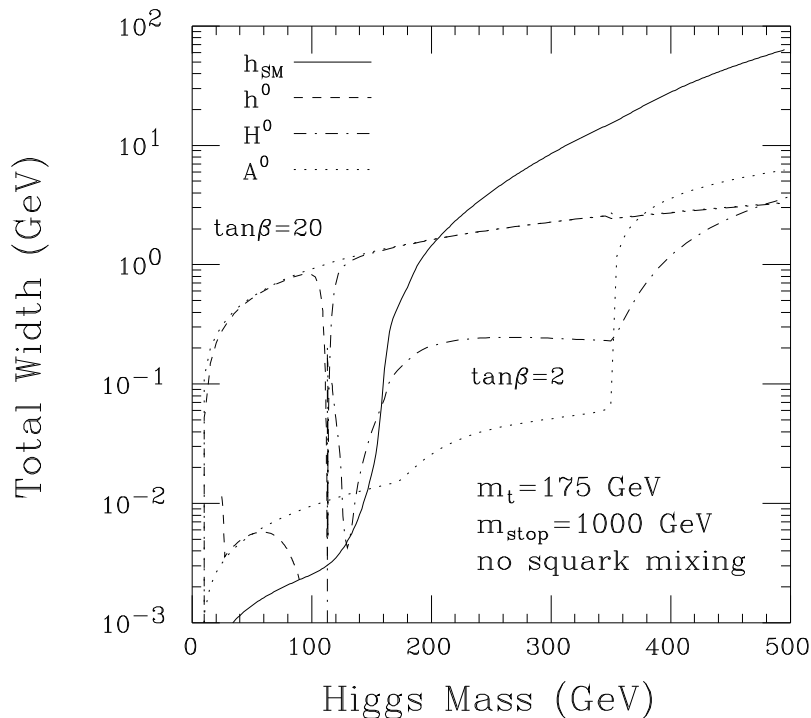
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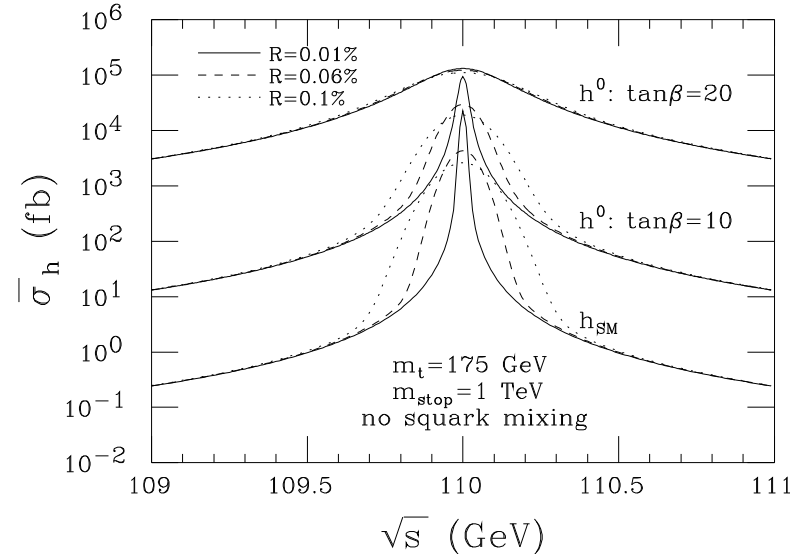
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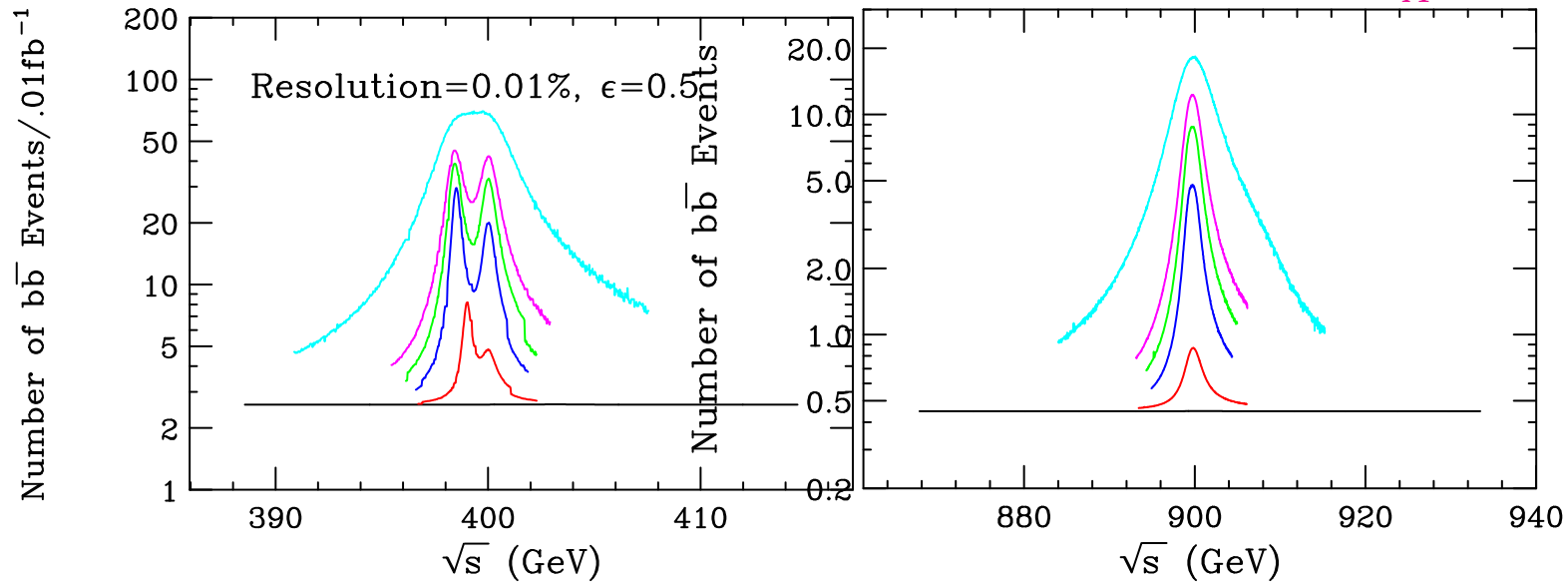
Higgs Total Widths



Effective Cross Sections: $m_h=110$ GeV



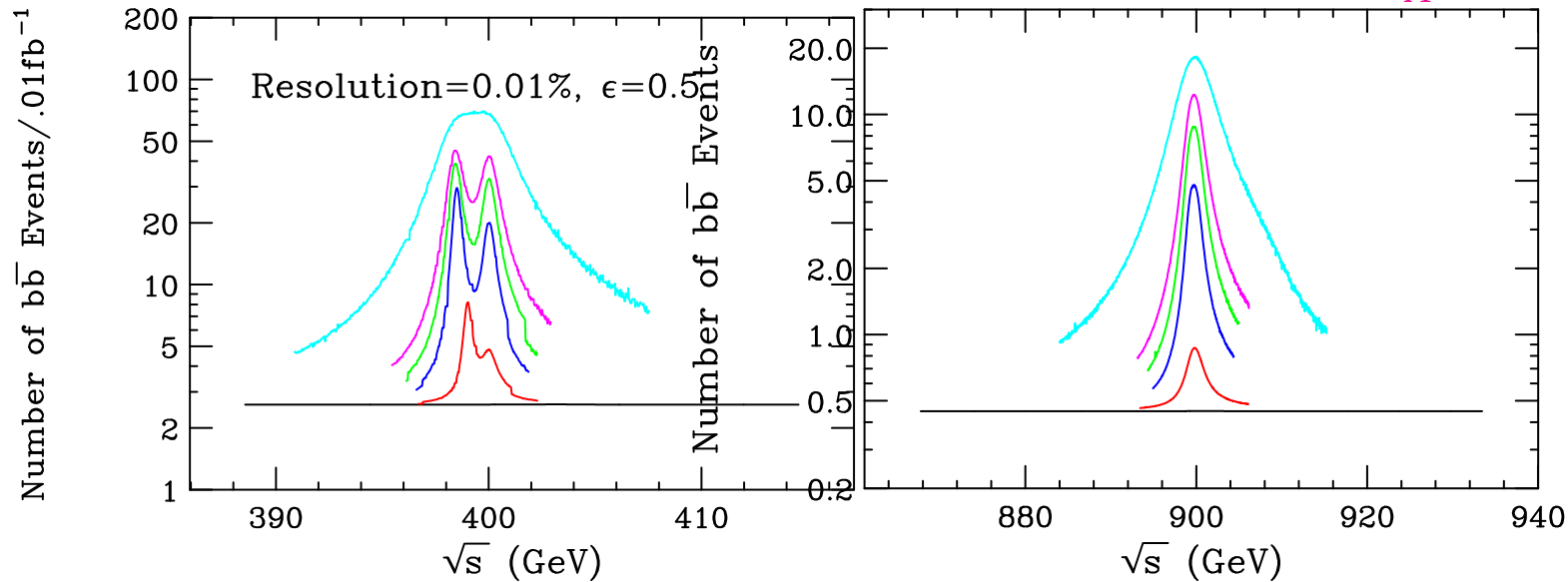
Heavy Higgs degenerate as M_A large: $\delta M \approx \frac{M_Z^2}{2M_A} \sin^2 2\beta$.



400 GeV Higgses resolved!

900 GeV Higgses not resolvable.‡

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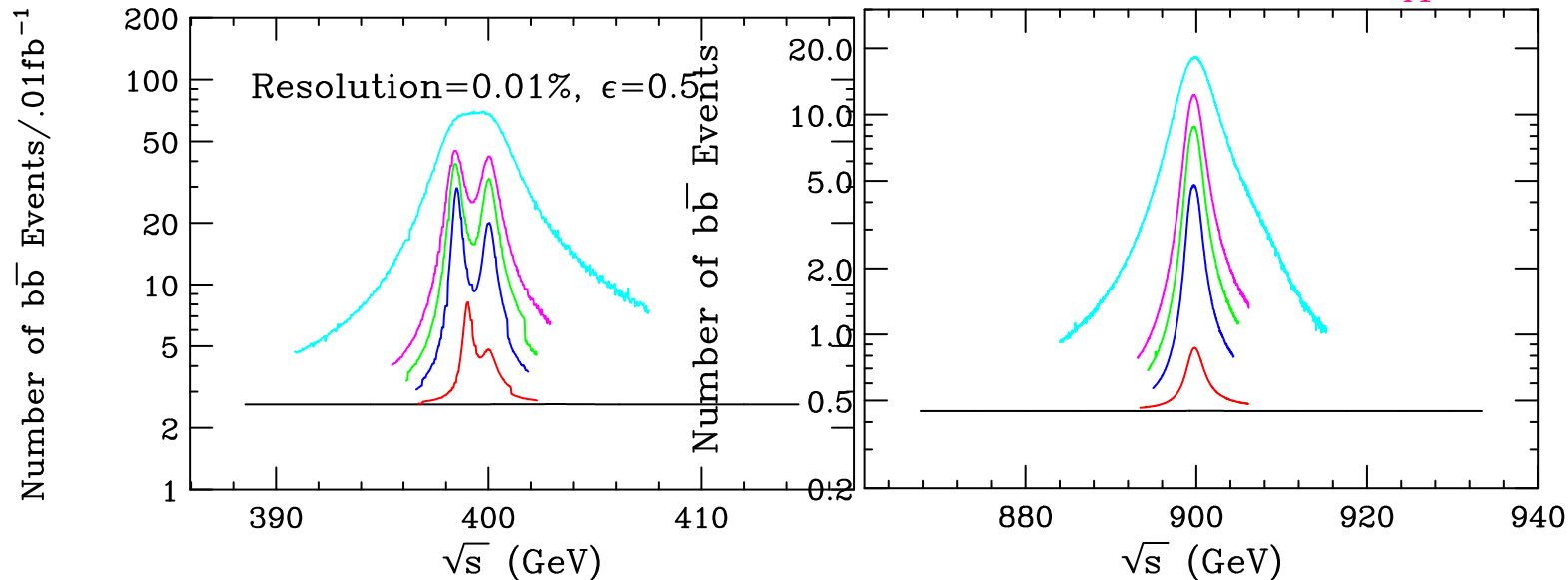
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Even at high mass, there is sufficient info understanding the Higgs sector:

$$\bar{\sigma}^{measured}(b\bar{b}, t\bar{t}, \tau\tau) \Rightarrow \frac{4\pi\Gamma(H, A \rightarrow \mu^+\mu^-) \Gamma(H, A \rightarrow X)}{\Gamma_{tot}^2 M_H^2}.$$

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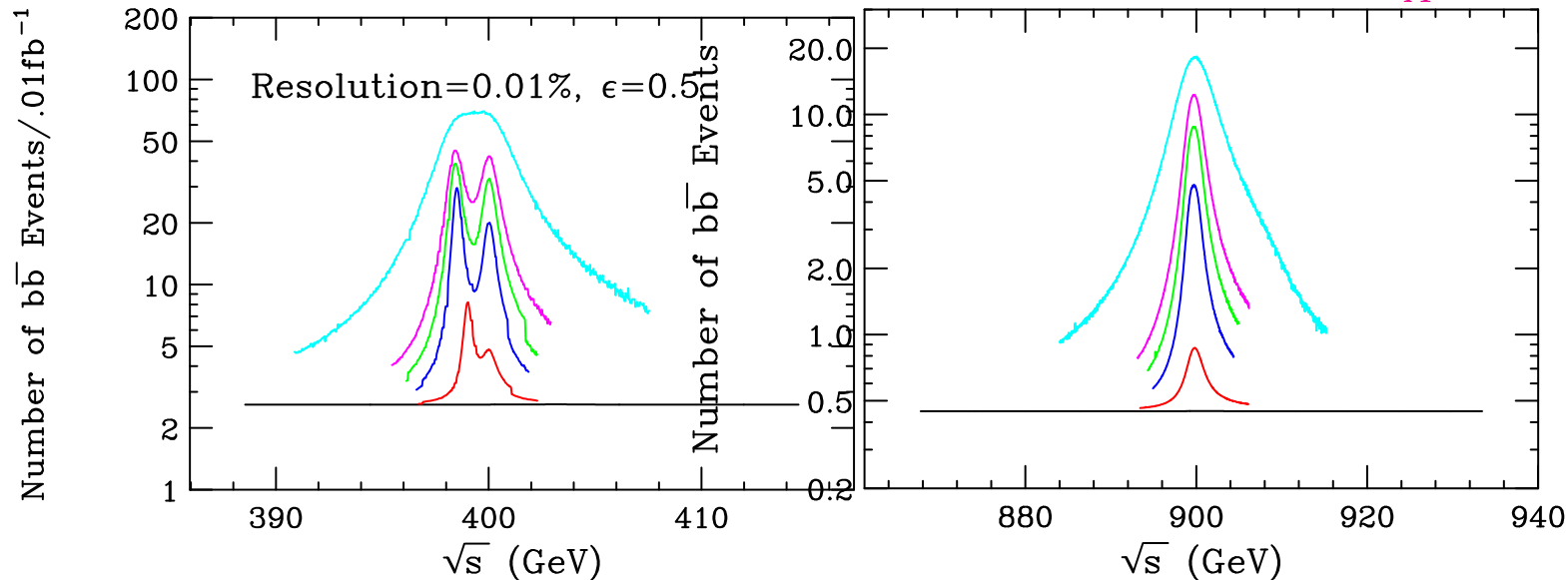
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- M_H : peak, accurate!
- Γ_{tot} : profile, accurate by scanning!
- $\bar{\sigma}^{measured}$: $(b\bar{b})/(t\bar{t}) \approx (m_b^2/m_t^2) \tan^4 \beta$, $(b\bar{b})/(\tau\tau) \approx 3m_b^2/m_\tau^2$ upto radiative corrections.
- $\bar{\sigma}^{tot} = (b\bar{b}) + (t\bar{t}) + (\text{smaller ones}) \Rightarrow \Gamma(\mu^+\mu^-)$ upto missing channels.

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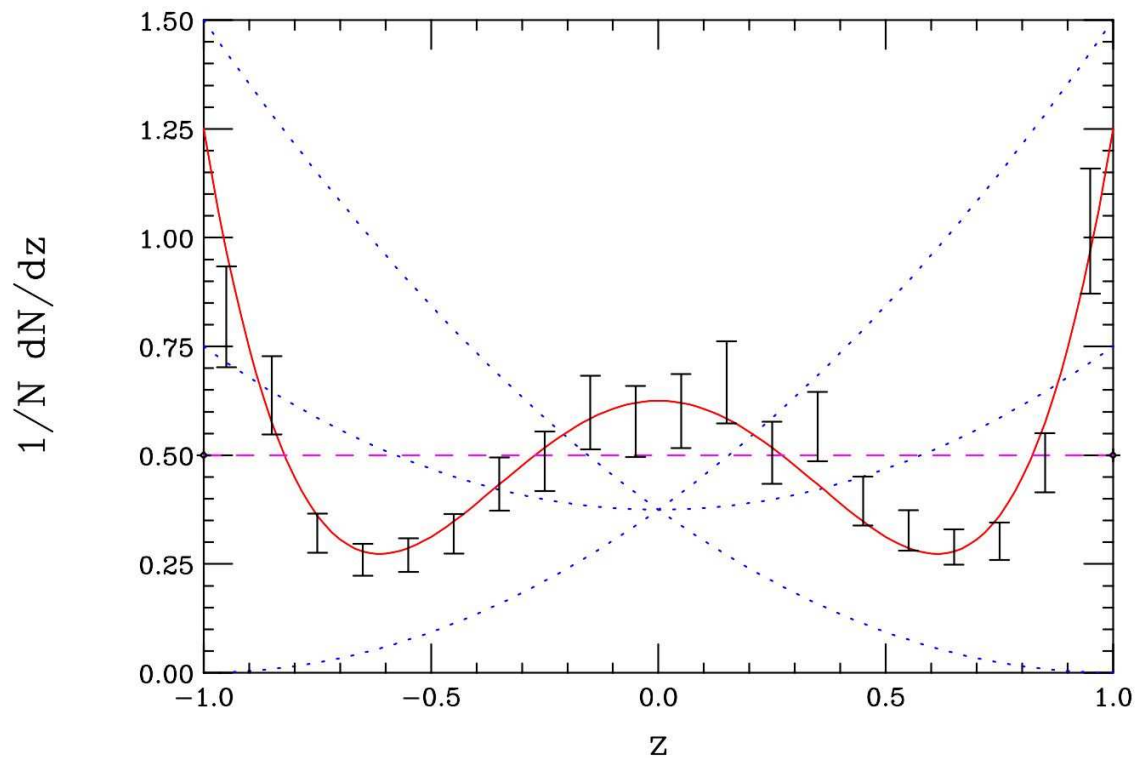
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- $\bar{\sigma}^{tot} = (b\bar{b}) + (t\bar{t}) + (\text{smaller ones}) \Rightarrow \Gamma(\mu^+\mu^-)$ upto missing channels.
- Compare with theory: $\Gamma(H, A \rightarrow \mu^+\mu^-)$, learn how many H, A 's contributing.
- If $t\bar{t}, \tau\tau$ decay kinematics reconstructed, hope to see CP violation!

‡Gunion, Han

A Z' Factory

Or any resonance R , that couples to a muon

$$\sigma(\mu^+\mu^- \rightarrow R \rightarrow X) \approx \frac{4(2J+1)\pi\Gamma(R \rightarrow \text{initial})\Gamma(R \rightarrow X)}{\Gamma^2 M^2}.$$



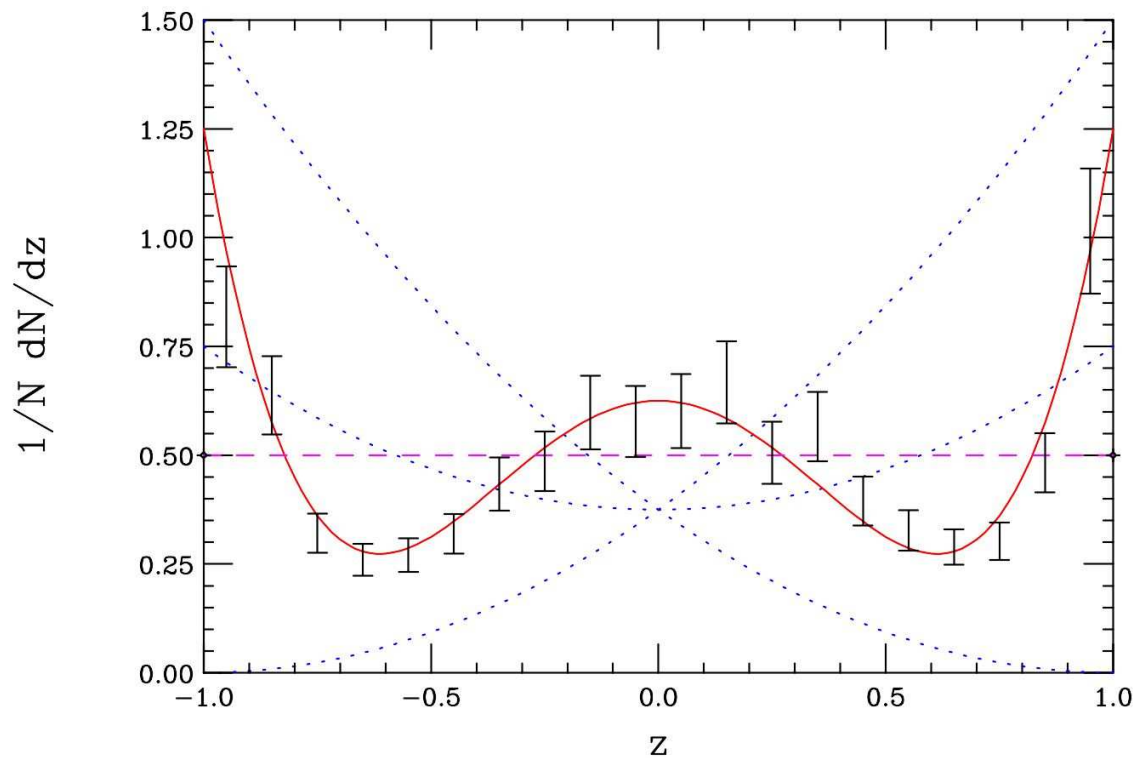
for various couplings and spins (1, 2).[§]

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Comment: The LHC will have the full coverage upto $M_{Z'} \sim 4 - 6$ TeV, which will soon motivate/define the machine needs (or otherwise).

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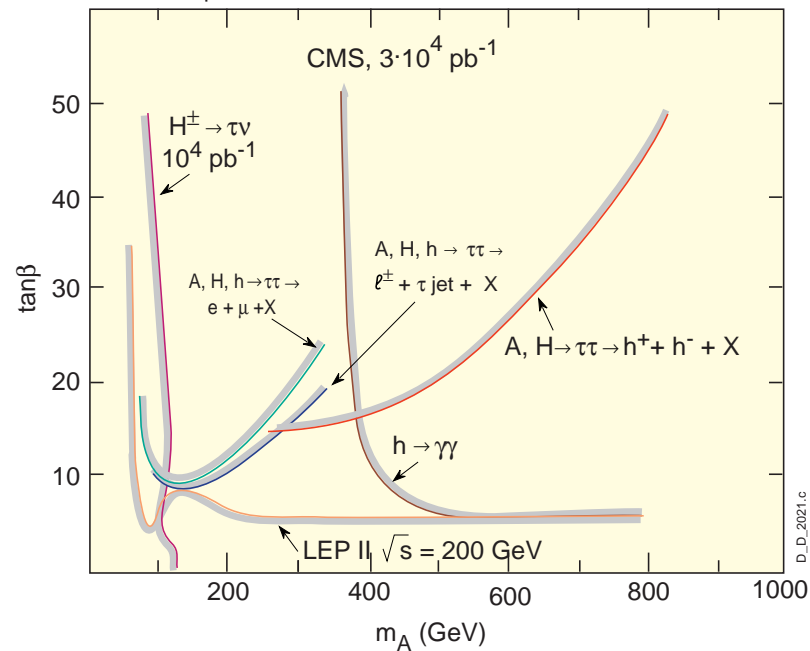
Electro-Weak Particle Pairs

(1). At LHC, h_{SM} fully covered, but H , A , H^\pm may not.
 At $\sqrt{s} = 14$ GeV, still a large hole, especially $M_{H,A} > 500$ GeV:¶

Significance contours for SUSY Higgses

Regions of the MSSM parameter space (m_A , $\tan\beta$)
 explorable through various SUSY Higgs channels

- 5σ significance contours
- two-loop / RGE-improved radiative corrections
- $m_{top} = 175$ GeV, $m_{SUSY} = 1$ TeV, no stop mixing ;



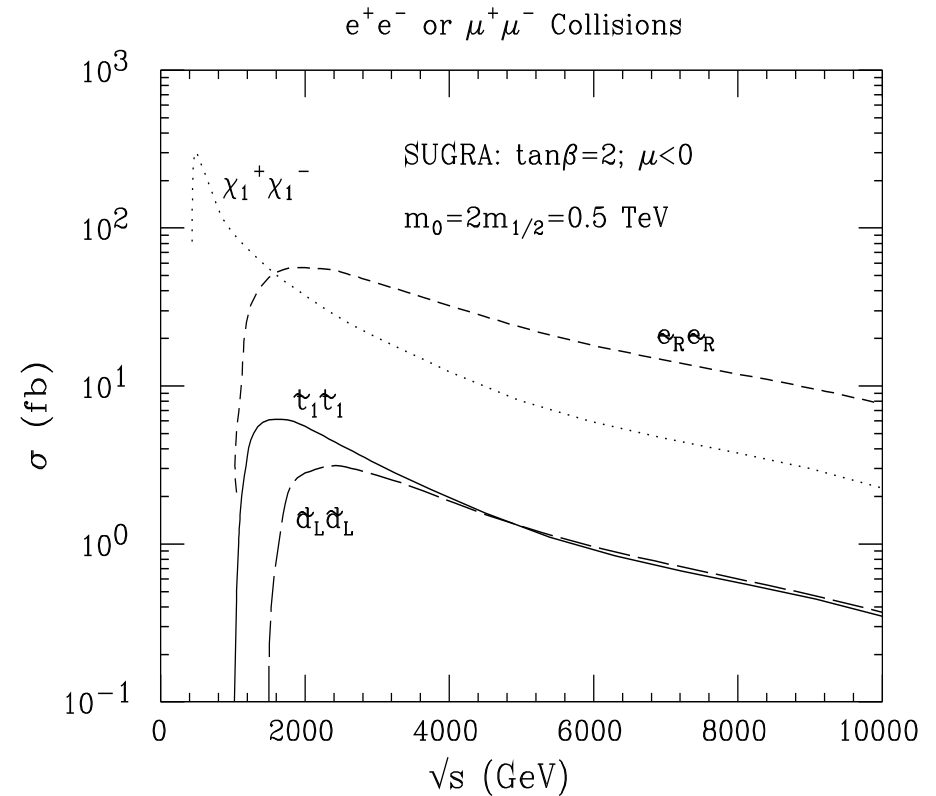
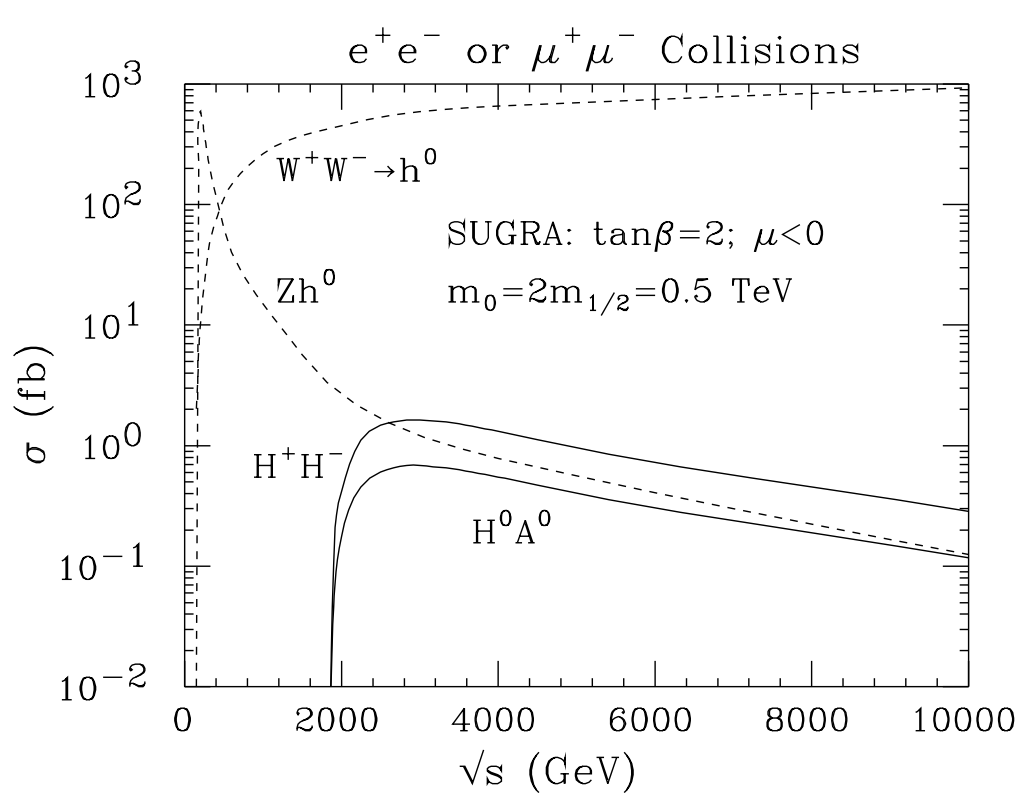
(2). At LHC, heavy EW pairs are difficult to search for
 $\mu^+\mu^- \rightarrow H_1H_2, \tilde{H}^+\tilde{H}^-, \tilde{H}^0\tilde{H}^0, \tilde{\ell}\tilde{\ell}.$

IF no help from colored states $\tilde{g} \rightarrow q\tilde{q} \rightarrow qq' \tilde{\chi}^{0,\pm} \dots$

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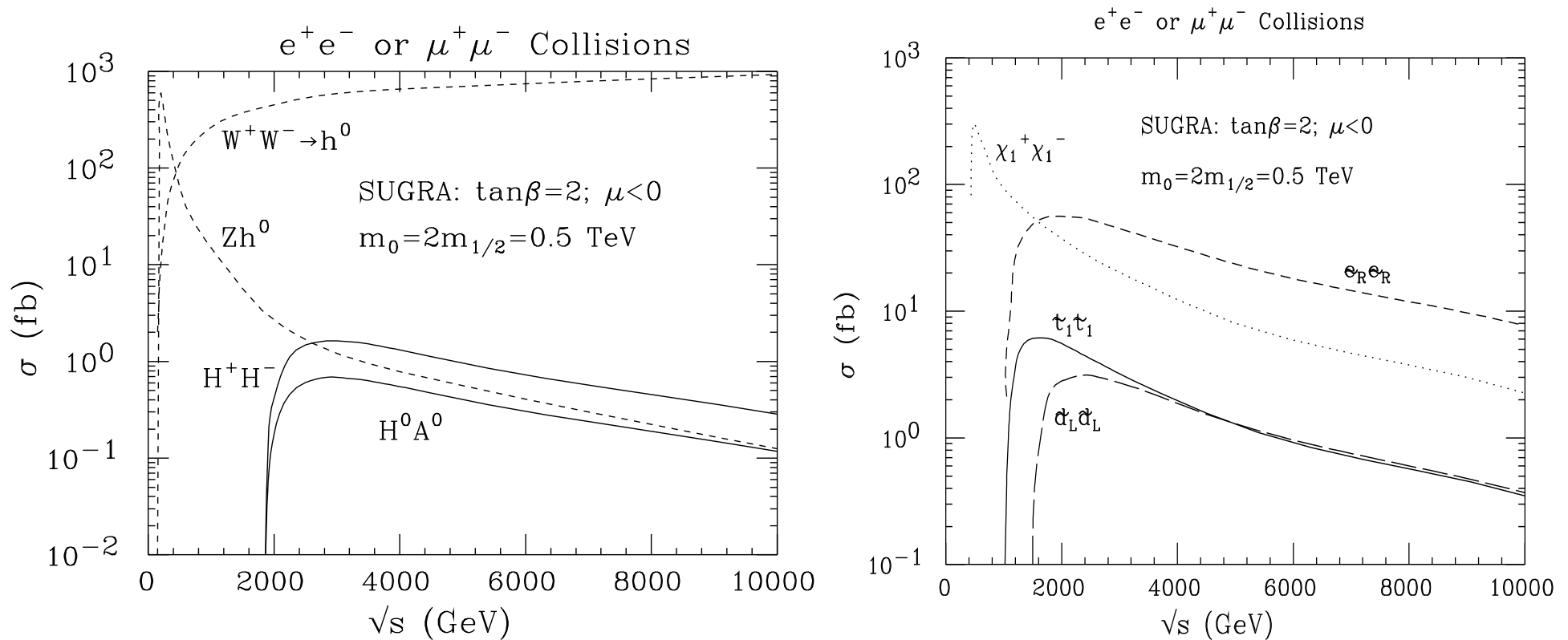
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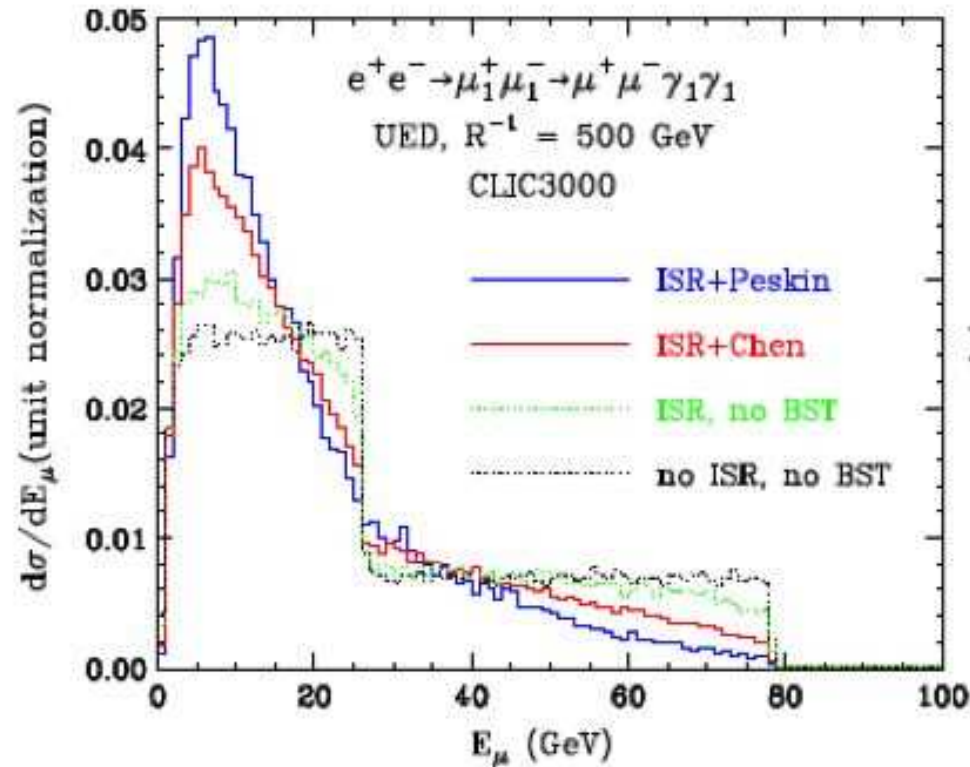
Once crossing the pair threshold, observation straightforward.
 (rather model-independent, like in Two-Higgs Doublet model etc.)

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Energy edges in chain decays:^{||}

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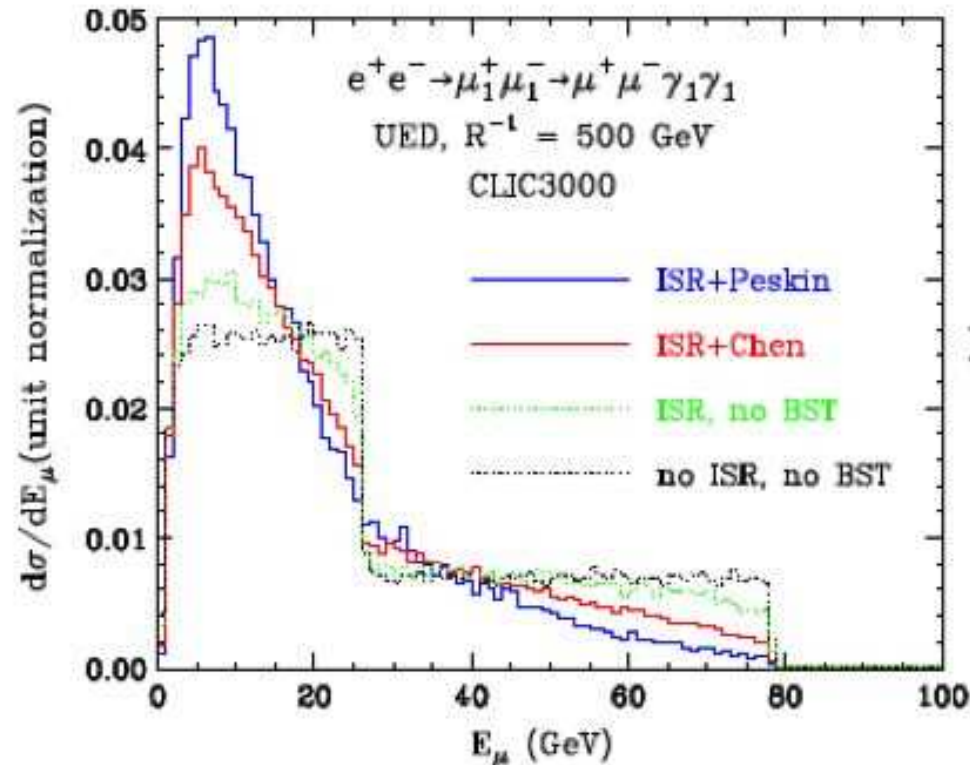
$$E_{max, min} = \frac{\sqrt{s}}{4} \left(1 - \frac{M_{\tilde{\chi}}^2}{M_{\tilde{\mu}}^2}\right) (1 \pm \beta), \quad \beta = \left(1 - \frac{4M_{\tilde{\mu}}^2}{s}\right)^{1/2}.$$

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Comment: very difficult at LHC due to under-constrained kinematics.

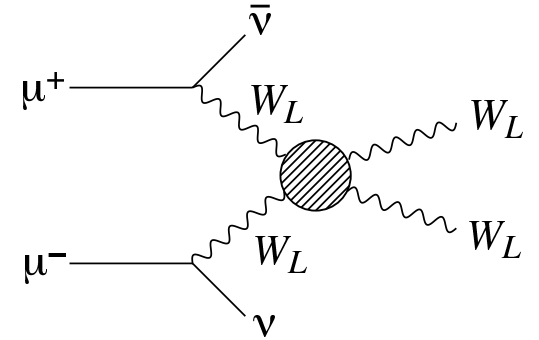
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Strong Electroweak Dynamics

$W_L W_L$ scattering

The scattering amplitude behaves as

$$A \sim \begin{cases} m_H^2/v^2 & \text{if light Higgs,} \\ s_{WW}/v^2 & \text{if no light Higgs.} \end{cases}$$



Partial wave unitarity implies: m_H or $\sqrt{s_{WW}} \leq 1.2$ TeV.

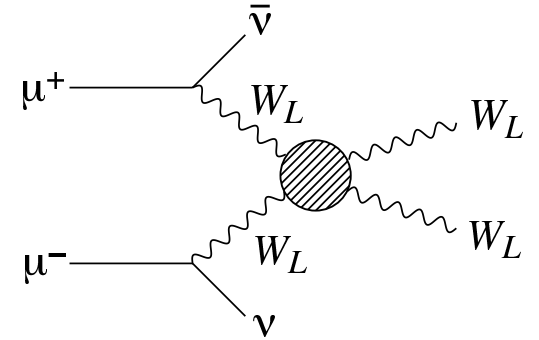
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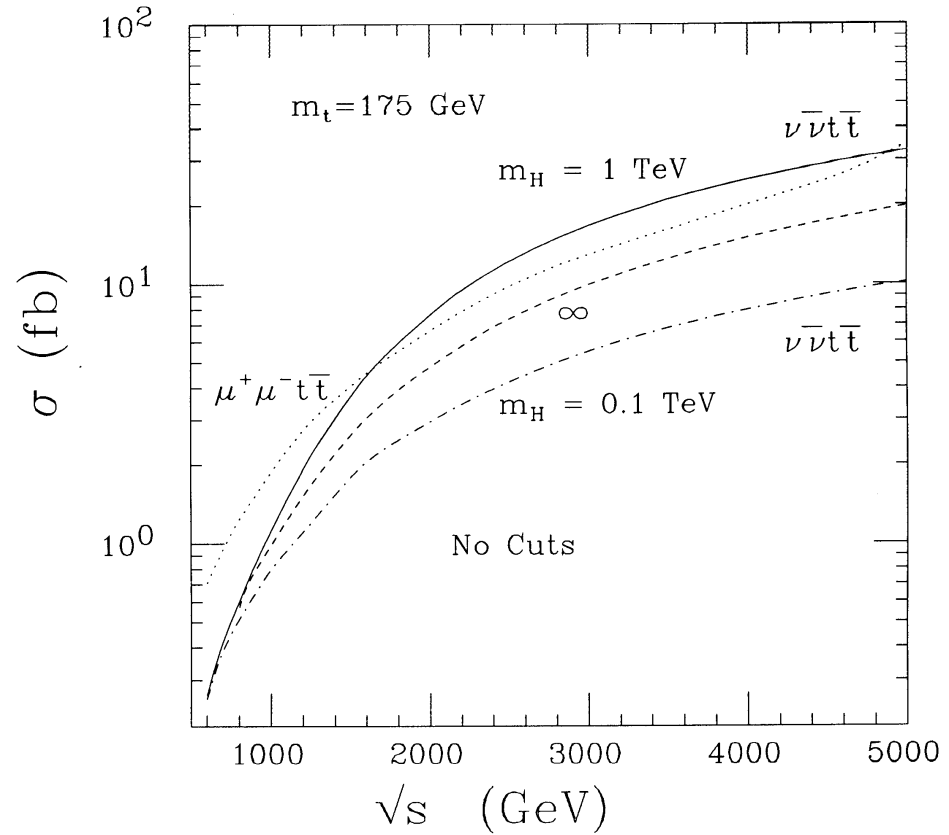
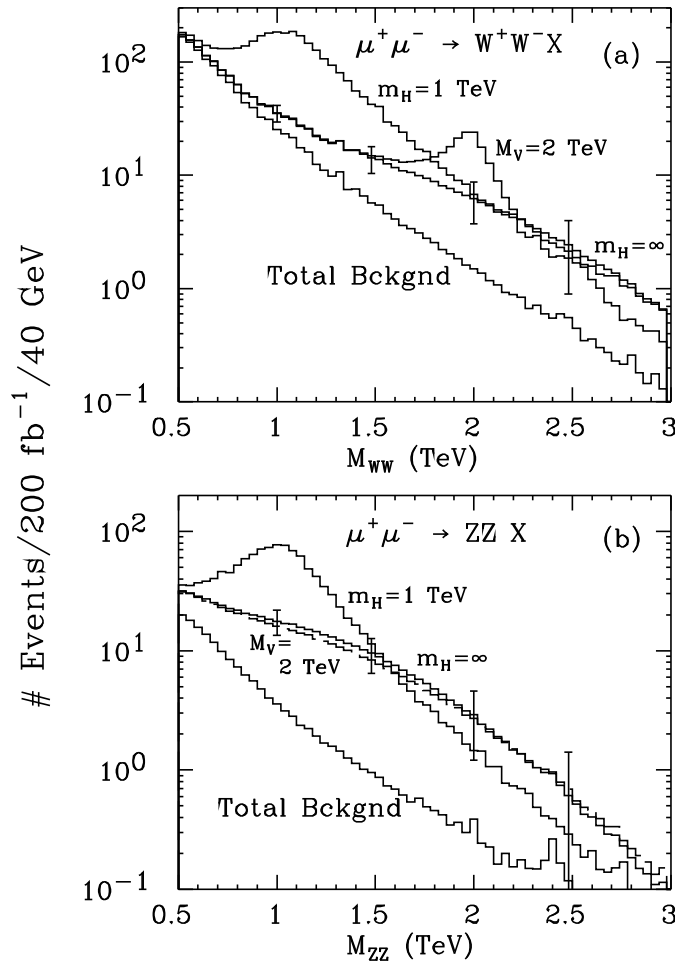
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For model connections:

$$\frac{\sigma(W_L^+ W_L^- \rightarrow W_L^+ W_L^-)}{\sigma(W_L^+ W_L^- \rightarrow Z_L Z_L)} \begin{cases} \sim 2 & \text{scalar } H^0, \\ \gg 1 & \text{vector } \rho_{TC}^0, \\ \sim 2/3 & \text{LET } \sqrt{s} \ll M. \end{cases}$$

Consider $\mu^+\mu^- \rightarrow \nu\nu W^+W^-$, $\nu\nu ZZ$ and $\nu\nu t\bar{t}$ via H , V or non-resonance at $\sqrt{s} = 4$ TeV.



Comment: This would be very challenging to test at the LHC.**

**Bagger et al.

Benchmark Processes

Recent meeting at Telluride, CO, *Muon Collider 2011*:
<http://conferences.fnal.gov/muon11>

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- (6) Contact interactions

Comparative Remarks:

Representative Physics Reach:

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CLIC $(1 - 2) \times 10^{34}$	$\checkmark \checkmark$ H potential	\checkmark	\checkmark	\checkmark e flavor	\checkmark CP-V
μ -Collider	$\checkmark \checkmark \checkmark$ H resonances CP-V	\checkmark	\checkmark	\checkmark μ flavor	\checkmark CP-V

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2. LHC: much larger SM backgrounds.
3. LHC: less constrained kinematics.

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Need to take advantage of the complementarity!