

**New Mathematics  
for**

**Old Scattering Amplitudes**



**M. Spradlin, Brown University**

**DPF 2011**

Many have contributed significantly to these developments including

L. Alday  
N. Arkani-Hamed  
T. Bargheer  
N. Beisert  
Z. Bern  
R. Boels  
J. Bourjaily  
N.E. Bjerrum-Bohr  
A. Brandhuber  
M. Bullimore  
R. Britto  
F. Cachazo  
S. Caron-Huot  
J.J. Carrasco  
V. Del Duca  
L. Dixon  
L. Dolan  
C. Duhr  
J. Drummond

B. Eden  
H. Elvang  
B. Feng  
D. Freedman  
D. Gaiotto  
P. Goddard  
A. Goncharov  
J. Henn  
P. Heslop  
A. Hodges  
Y. Huang  
H. Ita  
H. Johansson  
V. Khoze  
M. Kiermaier  
G. Korchemsky  
D. Kosower  
L. Lipatov  
F. Loebbert

D. Maitre  
J. Maldacena  
L. Mason  
J. Plefka  
R. Roiban  
A. Sabio Vera  
A. Sever  
D. Skinner  
A. Smirnov  
E. Sokatchev  
W. Spence  
M. Staudacher  
G. Travaglini  
J. Trnka  
C. Vergu  
P. Vieira  
A. Volovich  
C. Wen  
E. Witten

The goals of this research program are:

- to **explore** the rich mathematical structure of scattering amplitudes
- to **exploit** that structure to make previously impossible calculations trivial

# The Appeal of Long-Lasting Vitality

From 2004 through 2011 (and beyond?) the  
impossible → hard → trivial  
phase transition often occurred on the  
timescale of months, not years.

Communal pride trumps (occasional)  
individual shame!

# Why $N=4$ Super-Yang Mills Theory?

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- Moreover it = quantum gravity in AdS space

# Why $N=4$ Super-Yang Mills Theory?

- A nontrivial but apparently solvable four-dimensional QFT.
- But not just some random QFT, it is a gauge theory, and hence a cousin of QCD.
- Moreover it = quantum gravity in AdS space
- Most importantly, because it is there and

YES WE CAN!



## S-Matrix Theory Reloaded?

The spirit behind much recent progress is similar to that of "S-Matrix Theory".

If we have seen further than the giants of the past, it is because

# S-Matrix Theory Reloaded?

The spirit behind much recent progress is similar to that of "S-Matrix Theory".

If we have seen further than the giants of the past, it is because we're focused squarely on the well-manicured lawn beneath our feet ( $N=4$ ) rather than the dark jungle looming ahead (QCD).

# A Completely Different Philosophy

Then, the idea was

- (1) enumerate the principles  
(locality, analyticity, etc.)
- (2) solve!

# A Completely Different Philosophy

Then, the idea was

(1) enumerate the principles  
(locality, analyticity, etc.)

(2) solve!

Now, our aim is to

(1) write down the solution **BY ANY**  
**MEANS NECESSARY**

(2) deduce the principles from understanding  
its structure

BUCKLE

YOUR

SAFETY BELTS

I'll focus almost exclusively on planar, color-stripped partial amplitudes in  $\mathcal{N}=4$  Yang-Mills theory.

$$\mathcal{A}(1, \dots, n) = \sum_{\sigma \in S_n} \mathcal{A}(\sigma(1), \dots, \sigma(n)) \text{Tr} [T^{a_{\sigma(n)}} \dots T^{a_{\sigma(1)}}]$$

$\sigma \in S_n \rightarrow$

color-ordered partial amplitude

# Amplitude Representation

In massless  $\phi^3$  theory, two representations of the tree-level four-particle amplitude are

$$A = -\frac{1}{s} - \frac{1}{t} - \frac{1}{u} \quad \text{and}$$

$$A = \frac{10s^2}{tu} + \frac{20s}{t} - \frac{9s}{tu} + \frac{32s}{u} + \frac{10u}{t} - \frac{15}{t} + \frac{34t}{u} - \frac{15}{u} + 42 + \frac{22t}{s} - \frac{5u}{ts} + \frac{10u}{s} + \frac{12t^2}{us} - \frac{5t}{us} - \frac{11}{s}$$

They are equal, if you use  
 $s+t+u=0$

A linear constraint  $\Rightarrow$  switch to

*free variables*

Amplitude  
Representation  
Spinor Helicity

For particles with spin we need spinor  
helicity variables.

$$\text{Solve } k^2=0 \text{ via } k^M = \sigma^M_{\alpha\dot{\alpha}} \lambda^\alpha \bar{\lambda}^{\dot{\alpha}}$$

2 2-component  
objects

$$\text{Notation } \langle ij \rangle = \epsilon_{ab} \lambda_i^a \lambda_j^b \quad [ij] = \epsilon_{\dot{a}\dot{b}} \lambda_i^{\dot{a}} \lambda_j^{\dot{b}}$$

$$\langle ij+k|e \rangle = \langle ij \rangle [je] + \langle ik \rangle [ke]$$



Amplitude  
Representation  
Spinor Helicity  
MHV etc.

A spin-1 particle has 2 physical states  
helicity +1 or -1

The  $n$ -particle MHV amplitude  
involves  $n-2$  + helicity &  $2$  - helicity particles

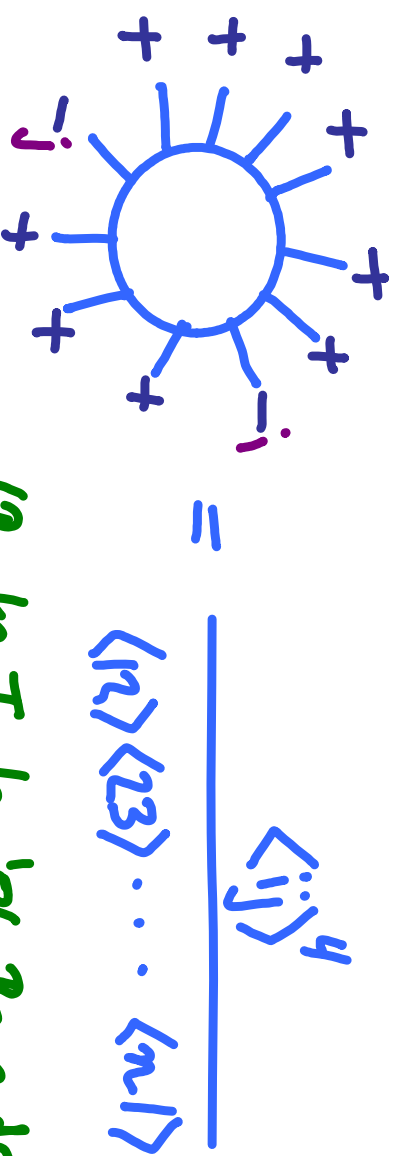
NMHV:  $(n-3, 3)$

$N^2$ -MHV:  $(n-4, 4)$

etc.

Amplitude  
Representation  
Spinor Helicity  
MHV etc.  
Parke-Taylor

The tree-level MHV amplitude is


$$= \frac{\langle ij \rangle^4}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle}$$

(Parke Taylor '86, Berends Giele '88)

This is, psychologically, the most important formula in our field.

It's message is ...

# The Philosophy of Amplification

1. Simplifications do not happen by accident.

Amplitude  
 Representation  
 Spinor Helicity  
 MHV etc.  
 Parke-Taylor  
 Superamplitude

Working in  $\mathcal{N}=4$  superspace we write

$$A_{\text{MHV}} = \frac{\delta^8 \left( \sum_{i=1}^n \lambda_i \eta_i^A \right)}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle} \quad \leftarrow \text{(Nair '88)}$$

supermomentum conservation

$$A = A_{\text{MHV}} [1 + \mathcal{P}_{\text{NMHV}} + \mathcal{P}_{\text{NNMHV}} + \dots]$$

$$\begin{array}{ccc} \uparrow & \uparrow & \uparrow \\ \mathcal{O}(\eta^8) & \mathcal{O}(\eta^4) & \mathcal{O}(\eta^8) \text{ etc.} \end{array}$$

Amplitude  
Representation  
Spinor Helicity  
MHV etc.  
Parke-Taylor  
Superamplitude  
BCFW

BCFW is a quadratic on-shell  
recursion relation; schematically

$$\text{Diagram with } n \text{ legs} = \sum_{k=2}^{n-2} \text{Diagram with } k \text{ and } n-k \text{ legs}$$

It is important both conceptually and  
practically — for rapidly generating **data**.

There is no single "BCFW representation",  
different discrete choices give apparently  
different (but secretly  $\Rightarrow$ ) results.

Example: The amplitude  $A(1^-, 2^-, 3^-, 4^+, 5^+, 6^+)$  may be expressed as either

$$\frac{[4|5+6|1\rangle^3}{s_{123}} + \frac{[6|1+2|3\rangle^3}{s_{612}} \quad \text{or}$$

$$\frac{[12][23]\langle 45 \rangle [1|2+3|4 \rangle [3|4+5|6 \rangle}{s_{123}^2} + \frac{[34]\langle 61 \rangle [3|4+5|6 \rangle [5|6+1|2 \rangle s_{612}}{\langle 12 \rangle^2 [45]^3} + \frac{\langle 34 \rangle [61] [1|2+3|4 \rangle [5|6+1|2 \rangle s_{234}}{\langle 23 \rangle^2 [56]^3}$$

Equality of these two expressions can (essentially) only be checked by **numerical experimentation**.

[Richardson's theorem: there is no algorithm for determining if two mathematical expressions are equal.]

A. Hodges followed his nose and found that there is a reason why in this case  $(2 \text{ terms}) = (3 \text{ terms})$ .

Each expression computes the volume of a polytope in complex projective space, triangulated in different ways!

This is one of many examples showing that...

# The Philosophy of AmplITUDEology

1. Simplifications do not happen by accident.
2. This is an experimental science.  
(Get the answer first, by any means necessary, then analyze it.)



Amplitude  
Representation  
Spinor Helicity  
MHV etc.  
Parke-Taylor  
Superamplitude  
BCFW  
A=Volumes  
Twistor Space

“We can go to twistor space by carrying out a  $1/2$ -fourier transform

$$A(\lambda_a, \mu_a) = \int d^2 \bar{\chi}_a e^{i \bar{\chi}_a \mu_a} A(\lambda_a, \bar{\chi}_a) \quad \gg$$

Since an overall rescaling of the 4 variables  $(\lambda, \mu)$  leaves  $k^\mu$  unchanged,

they parameterize  $\mathbb{CP}^3$ .

points in  $\mathbb{CP}^3 \Leftrightarrow$  null rays in spacetime  
lines in  $\mathbb{CP}^3 \Leftrightarrow$  points in spacetime

Amplitude  
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 Superamplitude  
 BCFW  
 A=Volumes  
 Twistor Space  
 Twistor String

Nothing that

momentum conserving delta function

$$A_{\text{MHV}} = \int d^4x \underbrace{\exp\left(i x_a \sum_{i=1}^n \lambda_i^a \bar{\lambda}_i^a\right)}_{\langle 12 \rangle \dots \langle n1 \rangle}$$

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Nothing that

momentum conserving delta function

$$A_{\text{MHV}} = \int_{\mathcal{P}} d^4x \exp(i x_a \sum_{i=1}^n \lambda_i^a \bar{\lambda}_i^{\dot{a}}) \frac{1}{\langle 12 \rangle \dots \langle n1 \rangle}$$

measure on the moduli space of lines in  $\mathbb{P}^3$

Witten invented a topological string

theory on  $\mathbb{P}^3$  which computes the

tree-level  $N^k_{\text{MHV}}$  amplitude as an integral

over curves of degree  $k+1$ !

MOVING ON

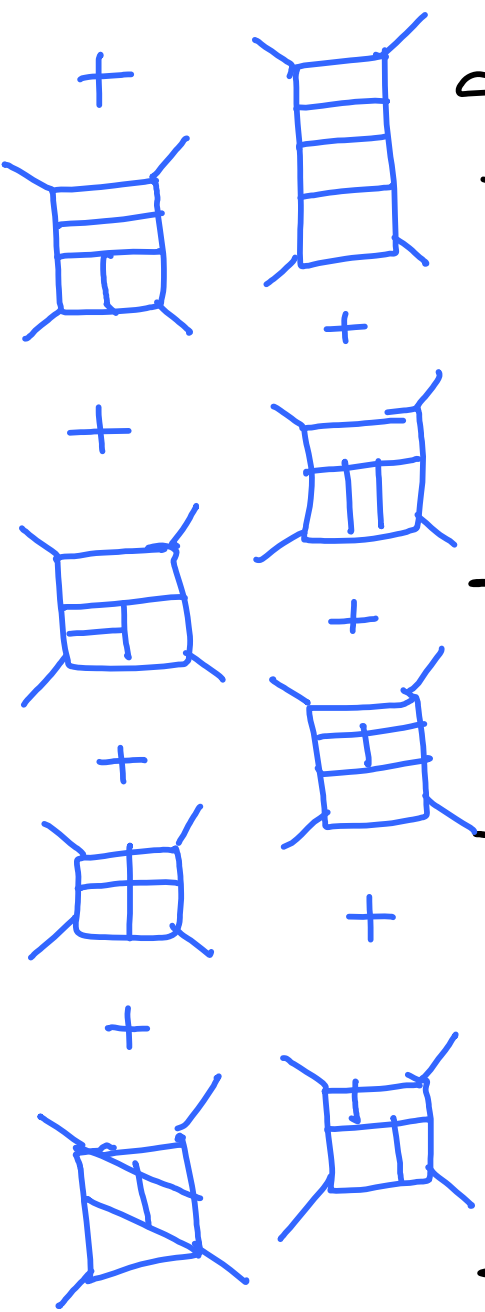
TO

LOOP LEVEL

Amplitude  
Representation  
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Parke-Taylor  
Superamplitude  
BCFW  
A=Volumes  
Twistor Space  
Twistor String  
Gen. Unitarity

Union of the unitarity-based method  
(Bem, Dixon, Dunbar, Kosower) with more  
recent advances (Britto, Cachazo, Feng)  
allowed much more **data** to be collected

eg the four-loop four-particle amplitude



(Bem, Zakon, Dixon, Kosower, Smirnov)

A pattern in this data reveals a surprise...

# What is Dual Conformal Symmetry?

A strange non-local symmetry under which scattering amplitudes are **co-variant**

- free level gluon amplitudes in QCD
- all orders in SYM

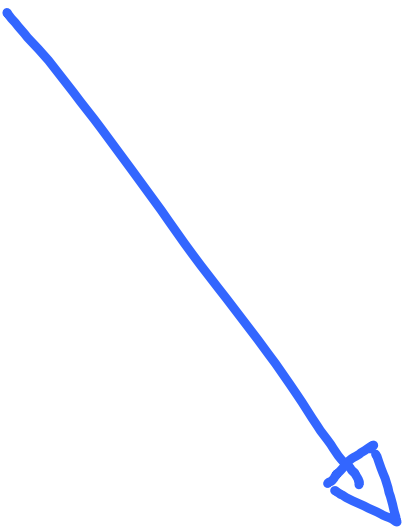
}  $N_{\text{color}} = \infty$

No understanding at the level of the Lagrangian.

# What is Dual Conformal Symmetry?

A strange non-local symmetry under which scattering amplitudes are **co-variant**

$A_{n\text{-particle}}$   
 $A_{\text{MHV}}$   
 $A_{n\text{-particle}}$   
is **invariant**



# Dual conformal symmetry

(Drummond, Henn, Korchemsky, Sokatchev)

Parameterize

$$\left. \begin{aligned} K_1 &= x_1 - x_2 \\ K_2 &= x_2 - x_3 \\ &\vdots \\ K_n &= x_n - x_1 \end{aligned} \right\} \sum K_i = 0$$

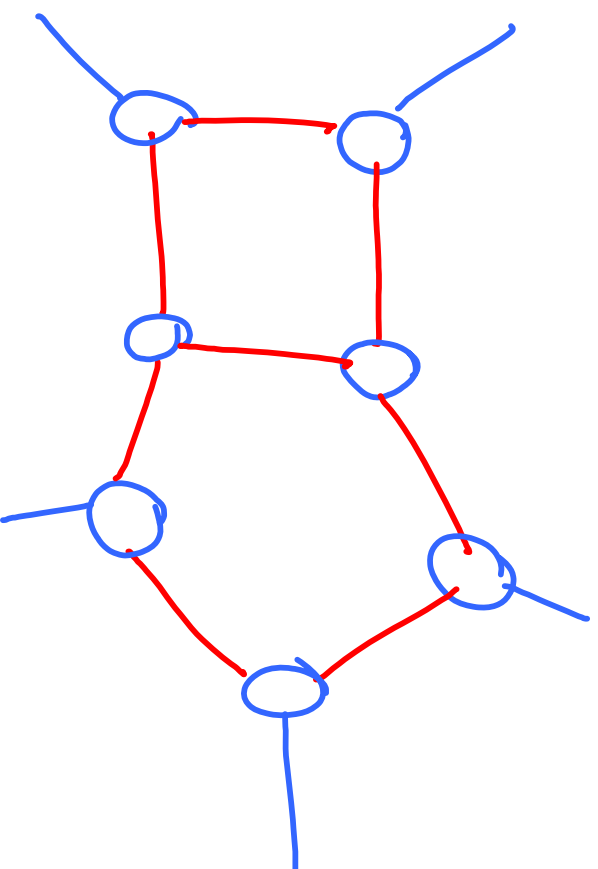
$\Rightarrow$  Then amplitudes are invariant under conformal transformations on the  $x_i$ .



Amplitude  
 Representation  
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 Parke-Taylor  
 Superamplitude  
 BCFW  
 A=Volumes  
 Twistor Space  
 Twistor String  
 Dual Conformal Inv.  
 The Integrand  
 Leading Sing.

The  $L$ -loop integrand is a rational function with (by definition) only simple poles in  $\mathbb{C}^{4L}$  (the loop momentum variables).

A **leading singularity** is a residue of the integrand @ one of these poles.



has a pole when all 8 propagators go on-shell

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Dual Conformal Inv.  
The Integrand  
Leading Sings.  
Yangian Symmetry

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## Yangian Symmetry

= the closure under commutation of  
{ ordinary conformal transformations,  
dual conformal transformations }

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Leading Sings.  
Yangian Symmetry  
R-Invariants

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R-invariants = a natural collection of Yangian invariant objects which constitute the building blocks for e.g. tree-level BCFW representations & leading singularities of loop integrands  
(=  $\delta$ -functions constraining five points in  $\mathbb{CP}^{3|4}$  to lie inside a common  $\mathbb{CP}^3$  subspace.)

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 Yangian Symmetry  
 R-Invariants  
 Grassmannian Int.

Grassmannian Integral = A contour integral over the space of  $K \times n$  matrices

$$\int [dC]_{K \times n} \frac{\delta^{4|4}(C_{ai} Z_i)}{(|\dots k| |\dots k+1| \dots |n+k-1|)} \quad Z_i = (\lambda_i, \mu_i)$$

minors of the matrix  $C$

(Arkani-Hamed, Cachazo, Cheung, Kaplan)

which is a generating function for all  $n$ -particle  $\mathcal{N}^k_{\text{MHV}}$  Yangian invariants.

(It's residues are  $R$ -invariants.)

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Leading Sings.  
Yangian Symmetry  
R-Invariants  
Grassmannian Int.  
Strong Coupling

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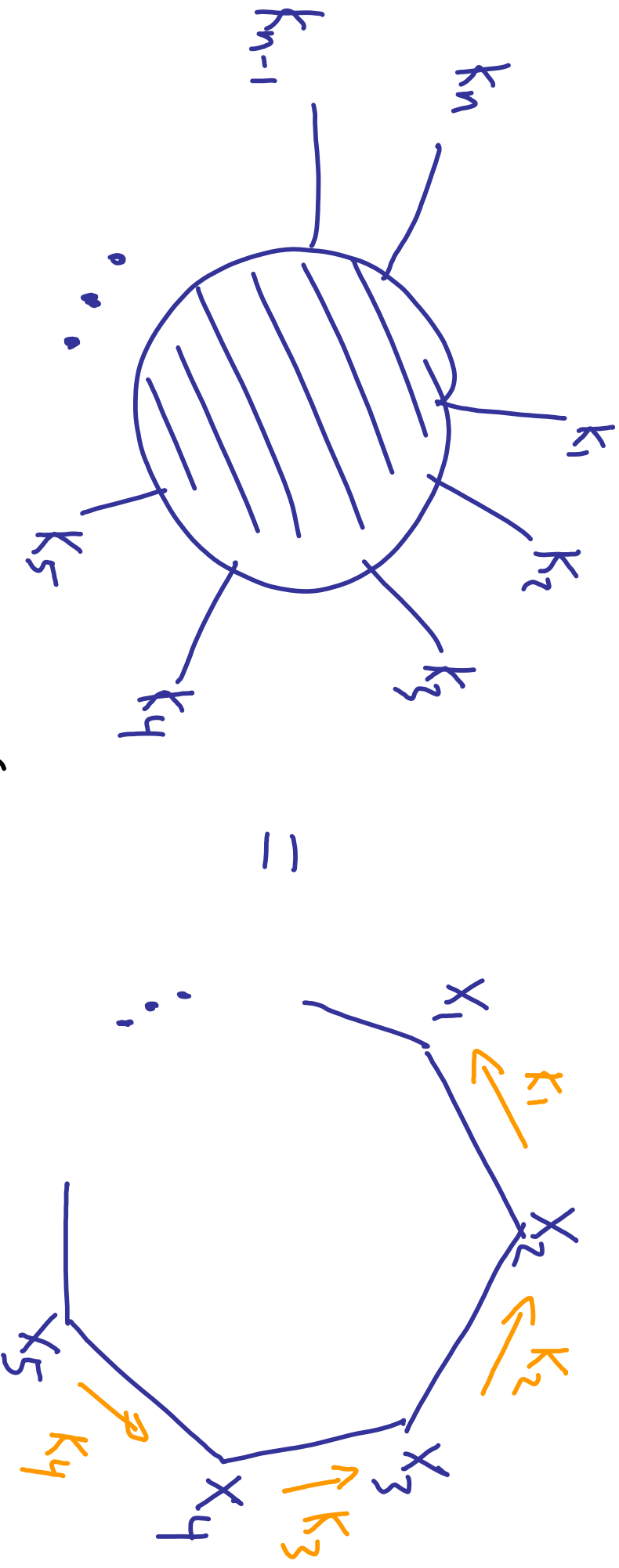
Using AdS/CFT, Alday & Maldacena showed that at strong coupling, amplitudes are computed as certain Wilson loops = minimal area surfaces in  $AdS_5$ .

Hence integrability methods can be brought to bear on the problem.

(Alday, Gaiotto, Maldacena)

# Amplitudes and Wilson Loops

Another remarkable feature of  $\mathcal{N}=4$  supersymmetric Yang-Mills is



The MHV scattering amplitude, for all  $n$  and to all orders in perturbation theory, = the expectation value of a Wilson loop on an  $n$ -sided polygon with lightlike edges.

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Yangian Symmetry  
R-Invariants  
Grassmannian Int.  
Strong Coupling  
A = Wilson Loop  
A = Correlation Fn.

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The integrand may in fact be computed  
yet **another** way, by correlation functions  
of the form

$\langle \underbrace{\mathcal{O}(x_1) \dots \mathcal{O}(x_n)}_{\text{light like separated } 1/2 \text{ BPS operators}} \underbrace{\mathcal{L}(y_1) \dots \mathcal{L}(y_L)}_{\text{Lagrangian insertions}} \rangle$

Eden, Heslop, Korchemsky, Sokatchev  
Adamo, Bullimore, Mason, Skinner

# You Say You Want a Revolution

One revolution has come to fruition

$$A = \int d^4 p_1 \dots d^4 p_L \underbrace{\sum \text{Feynman diagrams}}$$

↓ exciting new  
technology

$$= \int d^4 p_1 \dots d^4 p_L \text{ (relatively simple integrand)}$$



# You Say You Want a Revolution

One revolution has come to fruition

$$A = \int d^4 p_1 \dots d^4 p_L \underbrace{\sum \text{Feynman diagrams}}$$

↓ exciting new  
technology

=  $\int d^4 p_1 \dots d^4 p_L$  (relatively simple integrand)

still very hard to evaluate!

But another is desperately needed,

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R-Invariants  
Grassmannian Int.  
Strong Coupling  
A = Wilson Loop  
A = Correlation Fn.  
BDS Remainder

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The natural "observable" for loop-level  
MHV amplitudes is the  
BDS remainder function  
which is what is left after  
appropriate subtraction of IR  
divergences, leaving a finite &  
dual conformal invariant quantity.

The first nontrivial loop amplitude  
in  $\mathcal{N}=4$  SYM is the

2-loop (1-loop is absorbed into BDS  
subtraction)

6-particle (~~≠~~ any cross-ratios for  $n < 6$ )

MHV remainder function.

Computed analytically in late 2009 by  
Del Duca, Duhr & Smirnov















$$\begin{aligned}
& \frac{3}{4}\mathcal{G}\left(\frac{1}{v_{231}}, 1, \frac{1}{1-u_2}; 1\right) H(0; u_3) + \frac{3}{4}\mathcal{G}\left(v_{231}, \frac{1}{1-u_2}, 1; 1\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{v_{312}}, 1, \frac{1}{1-u_3}; 1\right) H(0; u_3) + \frac{1}{4}\mathcal{G}\left(v_{312}, \frac{1}{1-u_3}, 1; 1\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(v_{231}, 1, \frac{1}{1-u_3}; 1\right) H(0; u_3) + \frac{1}{4}\mathcal{G}\left(v_{231}, \frac{1}{1-u_3}, 1; 1\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(0; u_1) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, \frac{u_3-1}{u_2+u_3-1}; 1\right) H(0; u_1) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right) H(0; u_1) H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_1) H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_1) H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_1) H(0; u_3) + \\
& \frac{5}{24}\pi^2 H(0; u_1) H(0; u_3) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}; 1\right) H(0; u_2) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{u_2}, \frac{1}{u_2+u_3}; 1\right) H(0; u_2) H(0; u_3) + \frac{1}{4}\mathcal{G}\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right) H(0; u_2) H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right) H(0; u_2) H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right) H(0; u_2) H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right) H(0; u_2) H(0; u_3) + \frac{5}{24}\pi^2 H(0; u_2) H(0; u_3) + \\
& \frac{3}{4}H(0; u_2) H(0, 0; u_1) H(0; u_3) + 3H(0; u_1) H(0; u_2) H(0; u_3) + \\
& \frac{1}{4}H(0; u_2) H\left(0, 1; \frac{u_2+u_3-1}{u_1+u_2-1}\right) H(0; u_3) + \frac{1}{2}H(0; u_3) H(0, 1; (u_1+u_3)) H(0; u_3) + \\
& \frac{1}{4}H(0; u_1) H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right) H(0; u_3) + \frac{1}{2}H(0; u_2) H(0, 1; (u_2+u_3)) H(0; u_3) + \\
& \frac{3}{4}H(0; u_2) H(1, 0; u_1) H(0; u_3) + \frac{3}{4}H(0; u_1) H(1, 0; u_2) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; 0; u_1) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; 0; u_1) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{312}; 1\right) H(0; 0; u_1) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{321}; 1\right) H(0; 0; u_1) - \frac{23}{24}\pi^2 H(0, 0; u_1) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right) H(0; 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right) H(0; 0; u_2) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right) H(0; 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{321}; 1\right) H(0, 0; u_2) - \\
& \frac{25}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{312}; 1\right) H(0; 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{321}; 1\right) H(0, 0; u_2) + \\
& \frac{25}{4}H(0, 0; u_1) H(0, 0; u_2) - \frac{23}{24}\pi^2 H(0, 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right) H(0, 0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right) H(0, 0; u_3) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0, 0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0, 0; u_3) + 3H(0; u_1) H(0; u_2) H(0, 0; u_3) - \\
& \frac{25}{4}H(0, 0; u_1) H(0, 0; u_3) - \frac{25}{4}H(0, 0; u_2) H(0, 0; u_3) - \frac{23}{24}\pi^2 H(0, 0; u_3) + \frac{1}{12}\pi^2 H(0, 1; u_1) + \\
& \frac{1}{12}\pi^2 H(0, 1; u_2) - \frac{1}{24}\pi^2 H\left(0, 1; \frac{u_2+u_3-1}{u_2-1}\right) + \frac{1}{2}H(0; u_1) H(0; u_2) H(0, 1; (u_1+u_2)) +
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{12}\pi^2 H(0, 1; (u_1+u_2)) + \frac{1}{12}\pi^2 H(0, 1; u_3) + \frac{1}{4}H(0; u_1) H(0; u_2) H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) - \\
& \frac{1}{24}\pi^2 H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) + \frac{1}{12}\pi^2 H(0, 1; (u_1+u_3)) - \frac{1}{24}\pi^2 H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right) + \\
& \frac{1}{12}\pi^2 H(0, 1; (u_2+u_3)) - \frac{1}{2}G\left(0, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_1) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_1) + \frac{1}{4}G\left(\frac{u_2}{u_3}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_3}, \frac{u_1-1}{u_1+u_3-1}; 1\right) H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_1) - \\
& \frac{1}{4}G\left(\frac{1}{1-u_3}, v_{312}; 1\right) H(1, 0; u_1) - \frac{3}{4}H(0, 0; u_2) H(1, 0; u_1) - \frac{3}{4}H(0, 0; u_3) H(1, 0; u_1) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) H(1, 0; u_1) - \frac{1}{3}\pi^2 H(1, 0; u_1) - \frac{1}{2}G\left(0, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_2) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}; 1\right) H(1, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_2) + \frac{1}{4}G\left(\frac{u_2}{u_3}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_2) - \\
& \frac{1}{4}G\left(\frac{1}{1-u_1}, u_{123}; 1\right) H(1, 0; u_2) - \frac{3}{4}H(0, 0; u_1) H(1, 0; u_2) - \frac{3}{4}H(0, 0; u_3) H(1, 0; u_2) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_1+u_2-1}{u_2-1}\right) H(1, 0; u_2) - \frac{1}{4}H(1, 0; u_1) H(1, 0; u_2) - \frac{1}{3}\pi^2 H(1, 0; u_2) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) - \frac{1}{2}G\left(0, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{1-u_2}, \frac{u_3-1}{u_2+u_3-1}; 1\right) H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) - \frac{1}{3}\pi^2 H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_2}, \frac{u_2+u_3}{u_3}; 1\right) H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_3}, \frac{u_2+u_3}{u_3}; 1\right) H(1, 0; u_3) - \frac{1}{4}G\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(1, 0; u_3) + \\
& \frac{3}{4}H(0; u_1) H(0; u_2) H(1, 0; u_3) - \frac{3}{4}H(0, 0; u_1) H(1, 0; u_3) - \frac{3}{4}H(0, 0; u_2) H(1, 0; u_3) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right) H(1, 0; u_3) - \frac{1}{4}H(1, 0; u_1) H(1, 0; u_3) - \frac{1}{4}H(1, 0; u_2) H(1, 0; u_3) + \\
& \frac{1}{24}\pi^2 H(1, 1; u_1) + \frac{1}{24}\pi^2 H(1, 1; u_2) + \frac{1}{24}\pi^2 H(1, 1; u_3) + \frac{1}{2}H(0; u_2) H(0, 0; u_1) + \\
& \frac{1}{2}H(0; u_3) H(0, 0; u_2) + \frac{1}{2}H(0; u_1) H(0, 0; u_3) - \frac{1}{2}H(0; u_2) H\left(0, 0, 1; \frac{u_1+u_2-1}{u_2-1}\right) - \\
& \frac{1}{2}H(0; u_3) H\left(0, 0, 1; \frac{u_1+u_2-1}{u_2-1}\right) - H(0; u_1) H(0, 0, 1; (u_1+u_2)) - \\
& \frac{1}{2}H(0; u_2) H(0, 0, 1; (u_1+u_2)) - \frac{1}{2}H(0; u_1) H\left(0, 0, 1; \frac{u_1+u_3-1}{u_1-1}\right) -
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{2}H(0;u_2)H\left(0,0,1;\frac{u_1+u_3-1}{u_1-1}\right)-H(0;u_1)H(0,0,1;(u_1+u_3))- \\
& \frac{1}{2}H(0;u_3)H(0,0,1;(u_1+u_3))-\frac{1}{2}H(0;u_1)H\left(0,0,1;\frac{u_2+u_3-1}{u_3-1}\right)- \\
& \frac{1}{2}H(0;u_3)H\left(0,0,1;\frac{u_2+u_3-1}{u_3-1}\right)-H(0;u_2)H(0,0,1;(u_2+u_3))- \\
& H(0;u_3)H(0,0,1;(u_2+u_3))-\frac{1}{2}H(0;u_2)H(0,1,0;u_1)-\frac{1}{2}H(0;u_3)H(0,1,0;u_2)- \\
& \frac{1}{2}H(0;u_1)H(0,1,0;u_3)+\frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{u_1+u_2-1}{u_2-1}\right)- \\
& \frac{1}{4}H(0;u_3)H\left(0,1,1;\frac{u_1+u_2-1}{u_2-1}\right)+\frac{1}{4}H(0;u_1)H\left(0,1,1;\frac{u_1+u_3-1}{u_1-1}\right)- \\
& \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{u_1+u_3-1}{u_1-1}\right)-\frac{1}{4}H(0;u_1)H\left(0,1,1;\frac{u_2+u_3-1}{u_3-1}\right)+ \\
& \frac{1}{4}H(0;u_3)H\left(0,1,1;\frac{u_2+u_3-1}{u_3-1}\right)+\frac{1}{2}H(0;u_2)H(1,0,0;u_1)-\frac{1}{2}H(0;u_3)H(1,0,0;u_1)- \\
& \frac{1}{2}H(0;u_1)H(1,0,0;u_2)+\frac{1}{2}H(0;u_3)H(1,0,0;u_2)+\frac{1}{2}H(0;u_1)H(1,0,0;u_3)- \\
& \frac{1}{2}H(0;u_2)H(1,0,0;u_3)-\frac{1}{4}H(0;u_3)H\left(1,0,1;\frac{u_1+u_2-1}{u_2-1}\right)- \\
& \frac{1}{4}H(0;u_2)H\left(1,0,1;\frac{u_1+u_3-1}{u_1-1}\right)-\frac{1}{4}H(0;u_1)H\left(1,0,1;\frac{u_2+u_3-1}{u_3-1}\right)- \\
& \frac{1}{4}H(0,0,0;0;u_1)-7H(0,0,0,0;0;u_2)-7H(0,0,0,0;0;u_3)+\frac{3}{2}H\left(0,0,0,1;\frac{u_1+u_2-1}{u_2-1}\right)+ \\
& 7H(0,0,0,0;0;u_1)+\frac{3}{2}H\left(0,0,0,1;\frac{u_1+u_3-1}{u_1-1}\right)+3H(0,0,0,1;(u_1+u_3))+ \\
& 3H(0,0,0,1;(u_1+u_2))+\frac{3}{2}H\left(0,0,0,1;\frac{u_1+u_3-1}{u_1-1}\right)+3H(0,0,0,1;(u_1+u_3))+ \\
& \frac{3}{2}H\left(0,0,0,1;\frac{u_2+u_3-1}{u_3-1}\right)+3H(0,0,0,1;(u_2+u_3))+\frac{4}{4}H(0,0,1,0;u_1)+ \\
& \frac{9}{4}H(0,0,1,0;u_2)+\frac{9}{4}H(0,0,1,0;u_3)-\frac{1}{2}H(0,1,0,0;u_1)-\frac{1}{2}H(0,1,0,0;u_2)- \\
& \frac{1}{2}H(0,1,0,0;u_3)+\frac{1}{2}H\left(0,1,0,1;\frac{u_1+u_2-1}{u_2-1}\right)+\frac{1}{2}H\left(0,1,0,1;\frac{u_1+u_3-1}{u_1-1}\right)+ \\
& \frac{1}{2}H\left(0,1,0,1;\frac{u_2+u_3-1}{u_3-1}\right)+H(0,1,1,0;u_1)+H(0,1,1,0;u_2)+H(0,1,1,0;u_3)- \\
& \frac{1}{4}H\left(0,1,1,1;\frac{u_1+u_2-1}{u_2-1}\right)-\frac{1}{4}H\left(0,1,1,1;\frac{u_1+u_3-1}{u_1-1}\right)- \\
& \frac{1}{4}H\left(0,1,1,1;\frac{u_2+u_3-1}{u_3-1}\right)+H(1,0,0,1;\frac{u_1+u_2-1}{u_2-1})+H(1,0,0,1;\frac{u_1+u_3-1}{u_1-1})+ \\
& H(1,0,0,1;\frac{u_2+u_3-1}{u_3-1})+2H(1,0,1,0;u_1)+2H(1,0,1,0;u_2)+2H(1,0,1,0;u_3)+ \\
& \frac{1}{4}H(1,1,0,1;\frac{u_1+u_2-1}{u_2-1})+\frac{1}{4}H(1,1,0,1;\frac{u_1+u_3-1}{u_1-1})+ \\
& \frac{1}{4}H(1,1,0,1;\frac{u_2+u_3-1}{u_3-1})+\frac{1}{2}H(1,1,1,0;u_1)+\frac{1}{2}H(1,1,1,0;u_2)+\frac{1}{2}H(1,1,1,0;u_3)- \\
& \frac{1}{24}\pi^2H(0;u_3)\mathcal{H}\left(1;\frac{1}{u_{123}}\right)-\frac{1}{24}\pi^2H(0;u_1)\mathcal{H}\left(1;\frac{1}{u_{231}}\right)-\frac{1}{24}\pi^2H(0;u_2)\mathcal{H}\left(1;\frac{1}{u_{312}}\right)+ \\
& \frac{1}{8}\pi^2H(0;u_2)\mathcal{H}\left(1;\frac{1}{v_{123}}\right)-\frac{1}{8}\pi^2H(0;u_3)\mathcal{H}\left(1;\frac{1}{v_{123}}\right)+\frac{1}{24}\pi^2H(0;u_2)\mathcal{H}\left(1;\frac{1}{v_{132}}\right)-
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{24}\pi^2H(0;u_3)\mathcal{H}\left(1;\frac{1}{v_{132}}\right)-\frac{1}{24}\pi^2H(0;u_1)\mathcal{H}\left(1;\frac{1}{v_{213}}\right)+\frac{1}{24}\pi^2H(0;u_3)\mathcal{H}\left(1;\frac{1}{v_{213}}\right)- \\
& \frac{1}{8}\pi^2H(0;u_1)\mathcal{H}\left(1;\frac{1}{v_{231}}\right)+\frac{1}{8}\pi^2H(0;u_3)\mathcal{H}\left(1;\frac{1}{v_{231}}\right)+\frac{1}{8}\pi^2H(0;u_1)\mathcal{H}\left(1;\frac{1}{v_{312}}\right)- \\
& \frac{1}{8}\pi^2H(0;u_2)\mathcal{H}\left(1;\frac{1}{v_{312}}\right)+\frac{1}{24}\pi^2H(0;u_1)\mathcal{H}\left(1;\frac{1}{v_{321}}\right)-\frac{1}{24}\pi^2H(0;u_2)\mathcal{H}\left(1;\frac{1}{v_{321}}\right)- \\
& \frac{1}{4}H(0;u_2)H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{u_{123}}\right)-\frac{1}{4}H(1,0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{u_{123}}\right)+\frac{1}{24}\pi^2H\left(0,1,1;\frac{1}{u_{123}}\right)+ \\
& \frac{1}{24}\pi^2\mathcal{H}\left(0,1,1;\frac{1}{u_{231}}\right)-\frac{1}{4}H(0;u_1)H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{u_{231}}\right)-\frac{1}{4}H(1,0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{u_{231}}\right)- \\
& \frac{1}{4}H(0;u_1)H(0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{u_{312}}\right)-\frac{1}{4}H(1,0;u_1)\mathcal{H}\left(0,1,1;\frac{1}{u_{312}}\right)+\frac{1}{24}\pi^2H\left(0,1,1;\frac{1}{u_{312}}\right)- \\
& \frac{1}{4}H(0;u_2)H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{u_{123}}\right)+\frac{1}{4}H(0,0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{v_{123}}\right)+ \\
& \frac{1}{4}H(0,0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{123}}\right)+\frac{1}{6}\pi^2H\left(0,1,1;\frac{1}{v_{123}}\right)-\frac{1}{4}H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{132}}\right)+ \\
& \frac{1}{4}H(0,0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{v_{132}}\right)+\frac{1}{4}H(0,0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{132}}\right)+\frac{1}{6}\pi^2H\left(0,1,1;\frac{1}{v_{132}}\right)- \\
& \frac{1}{4}H(0;u_1)H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{213}}\right)+\frac{1}{4}H(0,0;u_1)\mathcal{H}\left(0,1,1;\frac{1}{v_{213}}\right)+\frac{1}{6}\pi^2H\left(0,1,1;\frac{1}{v_{213}}\right)+ \\
& \frac{1}{4}H(0,0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{213}}\right)+\frac{1}{6}\pi^2H\left(0,1,1;\frac{1}{v_{213}}\right)-\frac{1}{4}H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{231}}\right)+ \\
& \frac{1}{4}H(0,0;u_1)\mathcal{H}\left(0,1,1;\frac{1}{v_{231}}\right)+\frac{1}{4}H(0,0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{231}}\right)+\frac{1}{6}\pi^2H\left(0,1,1;\frac{1}{v_{231}}\right)- \\
& \frac{1}{4}H(0;u_1)H(0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{v_{312}}\right)+\frac{1}{4}H(0,0;u_1)\mathcal{H}\left(0,1,1;\frac{1}{v_{312}}\right)+\frac{1}{6}\pi^2H\left(0,1,1;\frac{1}{v_{312}}\right)+ \\
& \frac{1}{4}H(0,0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{v_{312}}\right)+\frac{1}{6}\pi^2H\left(0,1,1;\frac{1}{v_{312}}\right)-\frac{1}{4}H(0;u_1)H(0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{v_{321}}\right)+ \\
& \frac{1}{4}H(0,0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{321}}\right)+\frac{1}{6}\pi^2H\left(0,1,1;\frac{1}{v_{321}}\right)-\frac{1}{4}H(0;u_2)H(0;u_3)\mathcal{H}\left(1,1,1;\frac{1}{v_{321}}\right)+ \\
& \frac{1}{2}H(0,0;u_3)\mathcal{H}\left(1,1,1;\frac{1}{v_{123}}\right)+\frac{11}{24}\pi^2\mathcal{H}\left(1,1,1;\frac{1}{v_{123}}\right)-\frac{1}{24}\pi^2\mathcal{H}\left(1,1,1;\frac{1}{v_{132}}\right)- \\
& \frac{1}{24}\pi^2\mathcal{H}\left(1,1,1;\frac{1}{v_{213}}\right)-\frac{1}{2}H(0;u_1)H(0;u_3)\mathcal{H}\left(1,1,1;\frac{1}{v_{231}}\right)+\frac{1}{2}H(0,0;u_1)\mathcal{H}\left(1,1,1;\frac{1}{v_{231}}\right)+ \\
& \frac{1}{2}H(0,0;u_3)\mathcal{H}\left(1,1,1;\frac{1}{v_{231}}\right)+\frac{11}{24}\pi^2\mathcal{H}\left(1,1,1;\frac{1}{v_{231}}\right)-\frac{1}{2}H(0;u_1)H(0;u_2)\mathcal{H}\left(1,1,1;\frac{1}{v_{312}}\right)+ \\
& \frac{1}{2}H(0,0;u_1)\mathcal{H}\left(1,1,1;\frac{1}{v_{312}}\right)+\frac{1}{2}H(0,0;u_2)\mathcal{H}\left(1,1,1;\frac{1}{v_{312}}\right)+\frac{11}{24}\pi^2\mathcal{H}\left(1,1,1;\frac{1}{v_{312}}\right)- \\
& \frac{1}{24}\pi^2\mathcal{H}\left(1,1,1;\frac{1}{v_{321}}\right)+\frac{1}{2}H(0;u_2)\mathcal{H}\left(0,0,1;\frac{1}{u_{123}}\right)+\frac{1}{2}H(0;u_3)\mathcal{H}\left(0,0,1;\frac{1}{u_{123}}\right)+ \\
& \frac{1}{2}H(0;u_1)\mathcal{H}\left(0,0,1;\frac{1}{u_{231}}\right)+\frac{1}{2}H(0;u_3)\mathcal{H}\left(0,0,1;\frac{1}{u_{231}}\right)+\frac{1}{2}H(0;u_1)\mathcal{H}\left(0,0,1;\frac{1}{u_{312}}\right)+ \\
& \frac{1}{2}H(0;u_2)\mathcal{H}\left(0,0,1;\frac{1}{u_{312}}\right)+\frac{1}{2}H(0;u_3)\mathcal{H}\left(0,0,1;\frac{1}{u_{312}}\right)+\frac{1}{4}H(0;u_1)\mathcal{H}\left(0,1,1;\frac{1}{u_{123}}\right)+ \\
& \frac{1}{4}H(0;u_2)\mathcal{H}\left(0,0,1;\frac{1}{u_{312}}\right)+\frac{1}{4}H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{u_{123}}\right)+\frac{1}{4}H(0;u_1)\mathcal{H}\left(0,1,1;\frac{1}{u_{231}}\right)+
\end{aligned}$$



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Goncharov, Vergu, Volovich and I used some mathematical magic called the **symbol** of a **transcendental function** to simplify this expression down to a single line, involving nothing more complicated than the "standard" **Li** polylogarithm functions.

The worth of our formula is NOT  
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Nobody cares about the value of this amplitude!

Rather, it is important because it provides hope, where previously was none, that we might really unlock the secrets to all orders...

# The Philosophy of AmplITUDEology

1. Simplifications do not happen by accident.
2. This is an experimental science.  
(Get the answer first, by any means necessary, then analyze it.)
3. Simplicity has to be believed to be seen.