Minimal Holography Higher spin gravity from solveable CFTs

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Based on: 1101.2910, 1106.1897, 1107.5098 and papers in progress

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Outline

Motivation

Higher Spin Holography

Applications/Examples

$$\Lambda < 0$$

AdS/CFT provides a non-perturbative definition of quantum gravity in certain circumstances (string theory, supersymmetry, $\Lambda < 0$, etc.)

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1. How does a local extra dimension emerge from CFT?

And when does it work?

- 2. When does CFT = AdS quantum gravity?
- 3. How does holography apply in other spacetimes?
- 4. Can time be an emergent dimension?

Simplified Holography

A possible approach:

• Find holographic dualities simple enough to solve, but complicated enough to look like gravity in d>2.

This talk: A summary of higher spin dualities

- Simple bulk: Higher spin gravity
- Simple boundary: Free CFTs and other exactly solveable models

(This is a toy model; higher spin gravity has no Einstein-gravity-like limit.)

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Higher Spin Gravity

Large (or infinite) number of massless fields, with spins

$$s = 0, 1, 2, 3, 4, \cdots, N$$

Spin-2 = graviton. Massless higher spin fields mean very large gauge symmetry extending diffeomorphism invariance.

Consistent interacting theory exists for $\Lambda \neq 0$

Fradkin and Vasiliev, 1987 Vasiliev, 1990

Toy model for string field theory in the stringy limit

AdS₄: Dual to Free Fields

Free fields have an infinite number of conserved currents of all spins:

$$J_{\mu_1 \cdots \mu_s} \sim \phi^a \partial_{(\mu_1} \cdots \partial_{\mu_s)} \phi^a$$

These are dual to bulk higher spin gauge fields:

$$J_{\mu_1\cdots\mu_s} \leftrightarrow A_{\mu_1\cdots\mu_s}$$

eg: Fronsdal 1979; Witten 02; Sezgin, Sundell 02; Mikhailov 02; Klebanov, Polyakov 02.

Roughly, the large gauge symmetry of the bulk theory guarantees that the boundary CFT₃ is exactly solveable (in this case, noninteracting).

In CFT₂ the result is solveable but interacting.

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Applications/Examples

- 1. Free Fields
- 2. Minimal Models
- 3. de Sitter Space

Example 1: The O(N) Model

Bulk

• Vasiliev's higher spin gravity in AdS₄

Boundary CFT₃

- The free O(N) model
- This CFT has an interacting IR fixed point that is dual to higher spin gravity with different boundary conditions.

Klebanov, Polyakov 02; Giombi, Yin 09.

Evidence

- Symmetries
- RG flows
- 3-point functions

Example 2: Minimal Model CFTs

Bulk

- Vasiliev's higher spin gravity in AdS₃
 - Higher spin sector has no propagating degrees of freedom
 - Dynamics from boundary d.o.f., scalar field, and black holes.

Boundary

- 2d W_N minimal model CFTs
- Higher spin generalizations of the c < 1 Virasoro Minimal Models (eg Ising model)
- Exactly solveable
- Admit a classical limit in the bulk

Campoleoni et al '10; Henneaux, Rey '10; Gaberdiel, Gopakumar '10; Gaberdiel, Hartman '11; Gaberdiel et al '11

Minimal Model Results

Classical Limit

- Bulk spectrum = boundary spectrum at large N
- Matching known RG flows
 Gaberdiel, Gopakumar, Hartman, Raju 'l l
- New RG flows from black hole solutions that interpolate between different W_N CFTs

Ammon, Gutperle, Kraus, Perlmutter 'I I

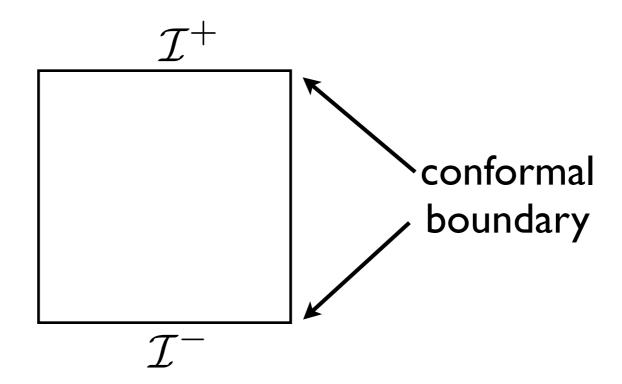
Quantum Limit

• In certain isolated cases, the full quantum theory has a bulk dual.

Castro, Gaberdiel, Hartman, Maloney, Volpato (to appear)

Example 3: dS/CFT

The dS/CFT correspondence, if it exists, is very different from AdS/CFT because time is emergent:



A general dS/CFT dictionary has been proposed, but an explicit example is lacking.

Witten '01; Strominger '01; Maldacena '02.

Higher Spin dS/CFT

Anninos, Hartman, Strominger (to appear)

Bulk

Vasiliev higher spin gravity in dS₄

Boundary

- The free Sp(N) CFT
- Anticommuting scalars; nonunitary

Motivation

$$\Lambda \sim N$$

This is, roughly, the O(-N) model.

Conclusion

Higher spin gravity is dual to solveable (sometimes free) CFTs

Some open questions

- How does a local extra dimension emerge from CFT?
 Jevicki, Jin, Ye '10,'11; Douglas, Mazzucato, Razamat '10.
- Quantum effects and black hole thermodynamics
- RG flows, cosmological evolution, and the big bang