

*Search for New Physics at  $\sqrt{s} = 7 \text{ TeV}$  in Hadronic  
Final States with  $\cancel{E}_T$  and Heavy Flavor*

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DPF 2011

August 9th, 2011



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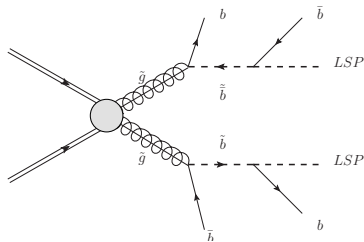
# Introduction

- Supersymmetry (SUSY) is one of the most popular and compelling theoretical possibilities for physics beyond the standard model. In the R-parity conserving, minimal SUSY extension to the Standard Model, sparticles are produced in pairs and decay to Standard Model particles and the stable lightest supersymmetric particle ( $\tilde{\chi}^0$ ).
- Mixing between left and right-handed squarks ( $\tilde{q}$ ) can lead to heavy flavor squark mass eigenstates which are significantly lighter than the first and second-generation quark super-partners, thereby enhancing the gluino ( $\tilde{g}$ ) decay branching ratio to heavy flavor.
- This analysis is an update and expansion of a [previous ATLAS analysis](#) done with  $35 \text{ pb}^{-1}$  of integrated luminosity.



# The General Case: $\tilde{g} \rightarrow \tilde{b}, \tilde{b} \rightarrow b + \tilde{\chi}^0$

The general SUSY process we are looking for is gluino pair production and subsequent decay via sbottom to 4  $b$ -jets and 2  $\tilde{\chi}^0$  (the LSP). The  $\cancel{E}_T$  in these events arises from the momentum mismatch between the two outgoing  $\tilde{\chi}^0$ , resulting in a  $4b + \cancel{E}_T$  final state signature. In this phenomenological model, there are 3 parameters:  $m_{\tilde{g}}$ ,  $m_{\tilde{b}}$ , and  $m_{\tilde{\chi}^0}$ .

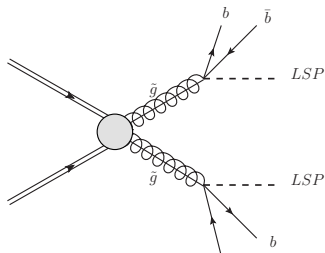


As we are in the 0 lepton ( $e, \mu$ ) channel, we are not very sensitive to  $\tilde{t}$ -mediated decays. These decays also tend to have lower  $\cancel{E}_T$  due to the reduced phase space afforded the LSP.

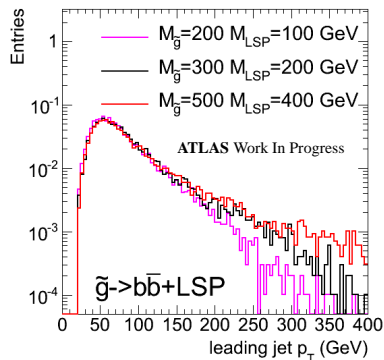


# Simplified Models: A Special Case

- $m_{\tilde{b}} \gg m_{\tilde{g}} > m_{\tilde{\chi}^0}$
- $m_{\tilde{q}_{1,2}}, m_{\tilde{t}} \gg m_{\tilde{b}}$
- $m_{\tilde{g}}$  and  $m_{\tilde{\chi}^0}$  varied
- $\tilde{b}$  is off-shell, decay chain is effectively  $\tilde{g} \rightarrow b\bar{b} + \tilde{\chi}^0$



In this limit, the squark mass spectrum is decoupled and the model kinematics are, to first order, determined by a single parameter:  $m_{\tilde{g}} - m_{\tilde{\chi}^0}$ .



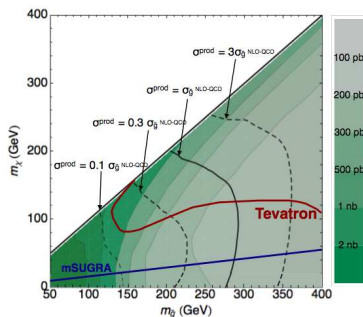
## More on Simplified Models

Simplified models have many advantages as signal samples when attempting to design analyses sensitive to as broad a slice of phase space as is possible and practical.

- Well-motivated, SM-like production and decay processes common to many Beyond the Standard Model (BSM) scenarios
- Kinematics determined by few parameters ( $m_{\tilde{g}} - m_{\tilde{\chi}^0}$ )
- Single final state and final state signature

Moreover, final states are often insensitive to model details, so we pick the simplest to characterize.

These features together enable cross section-independent signal region optimization and results with broad applicability to BSM models with  $\tilde{g}$ -like decays to heavy flavor.



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# Baseline Selections

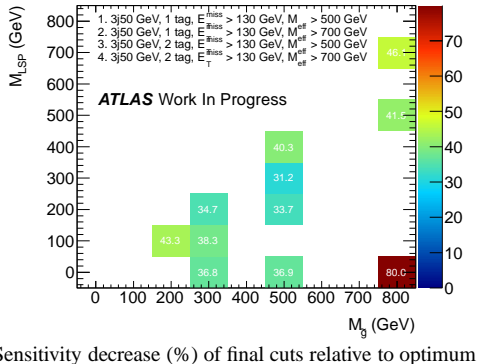
This search was conducted using  $0.83 \text{ fb}^{-1}$  of data collected with the ATLAS detector and a single jet +  $\cancel{E}_T$  trigger. The baseline selections are driven by data quality, trigger requirements, QCD rejection, and heavy flavor signal enhancement.

Motivation	Cut	Details
Data quality	Detector flags and jet cleaning At least 1 good vertex	Detector status, electronics noise Beam background/cosmic rejection
Trigger	1 jet with $p_T > 130 \text{ GeV}$ $\cancel{E}_T > 130 \text{ GeV}$	Turn-on plateau Turn-on plateau
QCD rejection	$\Delta\phi(J1/2/3, \cancel{E}_T) > 0.4$ $\cancel{E}_T/M_{\text{eff}} > 0.25$	Fake $\cancel{E}_T$ due to mis-measured jets
Signal enhancement	Lepton vetos ( $e, \mu$ ) 2 additional jets with $p_T > 50 \text{ GeV}$ 1 or more SV0 $b$ -tagged jets	No loose $e$ or $\mu$  Decay-length significance



# Optimization procedure

- 1 Evaluate kinematic variables for signal/background separation  
power:  $\cancel{E}_T$ ,  $H_T$ ,  $M_{eff}$  ( $\cancel{E}_T + H_T$ ),  
jet multiplicity, b-tagged jet multiplicity, etc.
- 2 Choose best variable set, create n-dimensional cut grid
- 3 Evaluate systematic-corrected significance at each cut grid point for each signal grid point to create a set of "optimal" kinematic cuts
- 4 Reduce number of selections to a manageable number while ensuring each signal point retains a signal region reasonably (50% of optimum) sensitive to it.

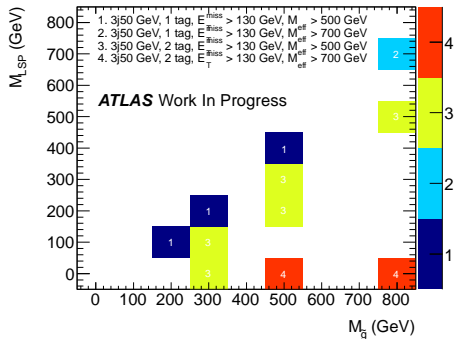




# Signal Regions

4 signal regions were chosen to represent good compromises among mass plane coverage, sensitivity, and practical concerns. These regions form a 2 by 2 grid in  $b$ -tagged jet multiplicity and  $M_{eff}$ .

Strong dependence on  $m_{\tilde{g}} - m_{\tilde{\chi}^0}$ , trends confirmed with full grid.



Most sensitive cut as a function of mass

	$M_{eff} > 500$ GeV	$M_{eff} > 700$ GeV
$\geq 1$ $b$ -tag	Region 1	Region 2
$\geq 2$ $b$ -tag	Region 3	Region 4



# Monte Carlo Background Estimation

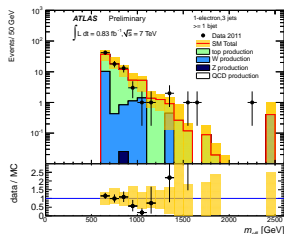
Non-QCD background estimates were taken from Monte Carlo simulation, and in the  $t\bar{t}$  case validated using a 1-lepton ( $e$  or  $\mu$ ) control region.  $t\bar{t}$  makes up the largest background component in all signal regions.

The 1-lepton control region uses the same kinematic cuts as the 0-lepton signal region except that 1 lepton with  $p_T > 20$  GeV,  $M_{eff} > 600$  GeV and  $40 \text{ GeV} < M_T < 100 \text{ GeV}$  are required.

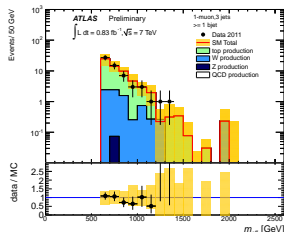
$M_T$  = transverse mass of lepton and  $\cancel{E}_T$

Physics process	$\sigma \cdot \text{BR}$ [nb]
$W \rightarrow l\nu$ (+jets)	$31.4 \pm 1.6$
$Z/\gamma^* \rightarrow ll$ (+jets)	$3.20 \pm 0.16$
$Z \rightarrow \nu\nu$ (+jets)	$5.82 \pm 0.29$
$t\bar{t}$	$0.165^{+0.011}_{-0.016}$
Single top	$0.085 \pm 0.003$

## 1e $t\bar{t}$ estimate validation



## 1 $\mu$ $t\bar{t}$ estimate validation

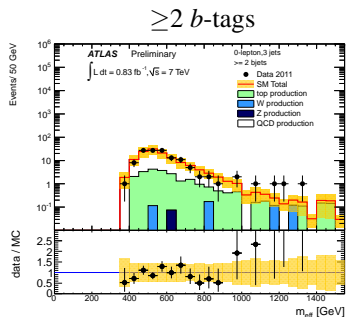
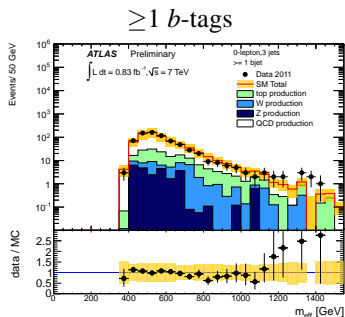


# Data-driven QCD Estimation

The QCD background estimate in the signal region was estimated from data, as leading-order Monte Carlo was not sufficient to provide a reliable estimate.

**Fundamental Assumption:**  $\cancel{E}_T$  in QCD multi-jet background due to mis-measured jets

- Jet momenta in data events with low  $\cancel{E}_T$  significance smeared with a jet response function to generate pseudoevents with large  $\cancel{E}_T$ .
- Validated by comparing data and pseudoevent distributions in QCD-enriched ( $\Delta\phi(\text{jet}, \cancel{E}_T) < 0.4$ ) control regions.



$M_{\text{eff}}$  distributions in QCD control regions



# Systematic Uncertainties

Non-QCD background - 30-35%

- Jet energy scale: 20-25%
- Theoretical: 25-30%
- $W/Z$  + heavy flavor: 50-100%
- $b$ -tagging efficiency: 10-30%
- Integrated luminosity: 4.5%

QCD background - 50%

- Smearing function dependency on flavor composition of the low- $\cancel{E}_T$  unsmear sample

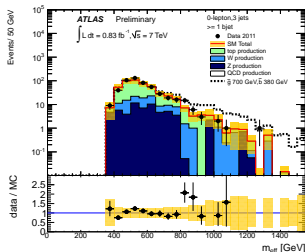
SUSY signal processes

- $b$ -tagging and jet energy scale: 10-35%
- Renormalization/factorization scales: 16-30%
- PDF: 7-25%
- Integrated luminosity: 4.5%



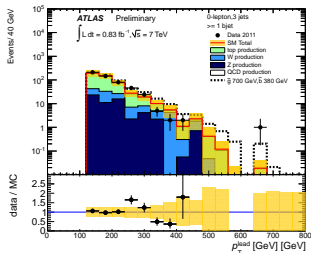
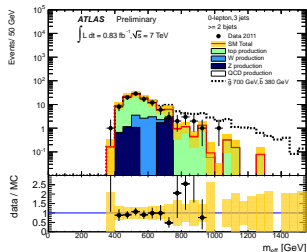
# Signal Region Data/MC Comparisons

$\geq 1$   $b$ -tag

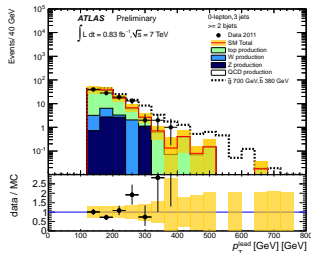


$M_{\text{eff}}$

$\geq 2$   $b$ -tags



$p_{\text{T}}^{\text{lead}}$



# Signal Region Data/MC Comparisons

Good agreement is observed between the SM expectation and data for all signal regions.

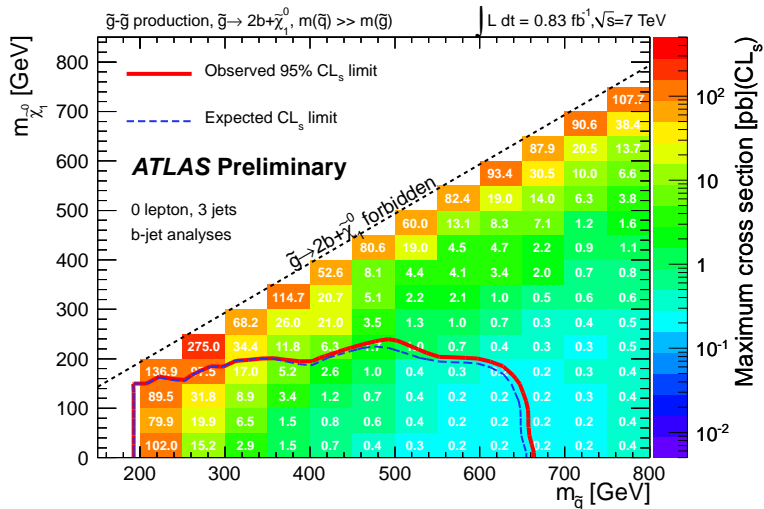
Sig. Reg.	Data (0.83 fb <sup>-1</sup> )	Top	W/Z	QCD	Total
1 (1 <i>b</i> -tag $M_{eff} > 500$ GeV)	361	221 <sup>+82</sup> <sub>-68</sub>	121 ± 61	15 ± 7	356 <sup>+103</sup> <sub>-92</sub>
2 (1 <i>b</i> -tag $M_{eff} > 700$ GeV)	63	37 <sup>+15</sup> <sub>-12</sub>	31 ± 19	1.9 ± 0.9	70 <sup>+24</sup> <sub>-22</sub>
3 (2 <i>b</i> -tag $M_{eff} > 500$ GeV)	76	55 <sup>+25</sup> <sub>-22</sub>	20 ± 12	3.6 ± 1.8	79 <sup>+28</sup> <sub>-25</sub>
4 (2 <i>b</i> -tag $M_{eff} > 700$ GeV)	12	7.8 <sup>+3.5</sup> <sub>-2.9</sub>	5 ± 4	0.5 ± 0.3	13.0 <sup>+5.6</sup> <sub>-5.2</sub>

Signal-independent upper limits on the non-SM contributions to each signal region can be defined in terms of event yield and effective cross section ( $\sigma_{eff} = \sigma \times BR \times \epsilon$ ).

Sig. Reg.	95% C.L. N events	95% C.L. $\sigma_{eff}$ (pb)
	$CL_s$ (PCL)	$CL_s$ (PCL)
1 (1 <i>b</i> -tag $M_{eff} > 500$ GeV)	240 (206)	0.288 (0.247)
1 (1 <i>b</i> -tag $M_{eff} > 700$ GeV)	51 (40)	0.061 (0.048)
2 (2 <i>b</i> -tag $M_{eff} > 500$ GeV)	65 (53)	0.078 (0.064)
3 (2 <i>b</i> -tag $M_{eff} > 700$ GeV)	14 (11)	0.017 (0.014)



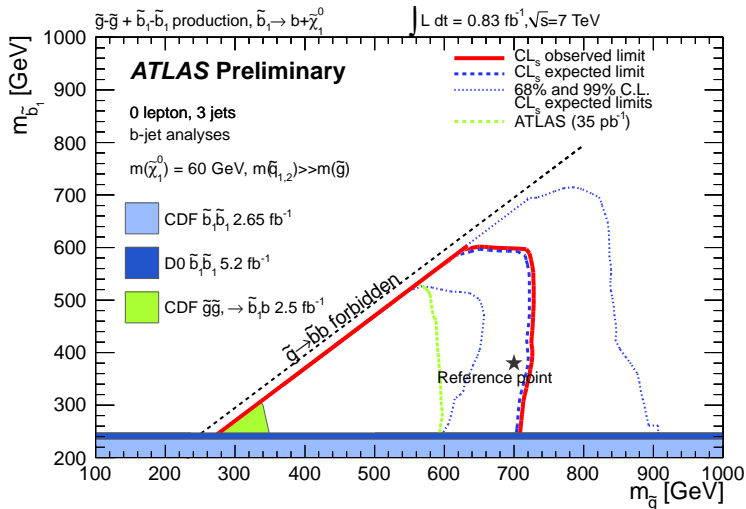
# Limits on $\tilde{g} \rightarrow b\bar{b} + \tilde{\chi}^0$ Simplified Models



LSP masses below 200-250 GeV are excluded for gluino masses in the range 200-660 GeV, if  $m_{\tilde{g}} - m_{\tilde{\chi}^0} > 100$  GeV



# Limits on $\tilde{g} \rightarrow \tilde{b}, \tilde{b} \rightarrow b + \tilde{\chi}^0$ Phenomenological Models



Glauino masses below 720 GeV are excluded for sbottom masses up to 600 GeV and  $m_{\tilde{\chi}^0} = 60 \text{ GeV}$ .





# Summary

A search for new physics with signatures characterized by heavy flavor,  $\cancel{E}_T$ , and no leptons was conducted on  $0.83 \text{ fb}^{-1}$  of data collected with the ATLAS detector at the LHC.

- For  $\tilde{g} \rightarrow b\bar{b} + \tilde{\chi}^0$ , LSP masses below 200-250 GeV are excluded for gluino masses in the range 200-660 GeV, if  $m_{\tilde{g}} - m_{\tilde{\chi}^0} > 100 \text{ GeV}$
- For  $\tilde{g} \rightarrow \tilde{b}$  general models, gluino masses below 720 GeV are excluded for sbottom masses up to 600 GeV and  $m_{\tilde{\chi}^0} = 60 \text{ GeV}$ . This extends the previous ATLAS limit by 130 GeV.
- Upper limits on the  $\tilde{g} \rightarrow b\bar{b} + \tilde{\chi}^0$  cross section between 0.2 and 275 pb are set for gluino masses between 200 and 750 GeV and neutralino masses between 1 and 700 GeV. These limits apply to in some measure to any BSM model where gluino-like particles are pair-produced and decay to  $b\bar{b} + \tilde{\chi}^0$ .



# Backup

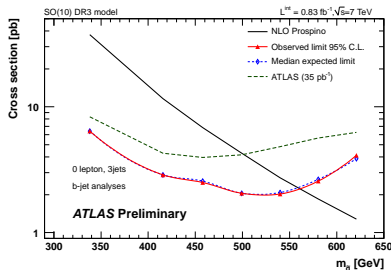


# SO10 Results

## DR3

$m_{\tilde{g}} < 570$  GeV is excluded.

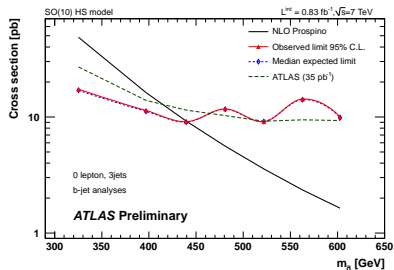
$\tilde{g} \rightarrow b\bar{b} + \tilde{\chi}^0$  dominates up to  $m_{\tilde{g}} = 550$  GeV, above which other decay channels lower the sensitivity of this analysis.



## HS

$m_{\tilde{g}} < 450$  GeV is excluded.

$\tilde{g} \rightarrow b\bar{b} + \chi_0^2$  has a large branching ratio and  $m_{\tilde{\chi}^0}$  is  $\approx 2X$  higher than the DR3 case.



# ATLAS Detector

The ATLAS detector is one of two general-purpose detectors built to collect pp collision data from the Large Hadron Collider at CERN. It consists of 4 major components: the inner detector (tracking), the calorimeters (energy measurements), the muon spectrometer, and the magnet and cooling systems.

