
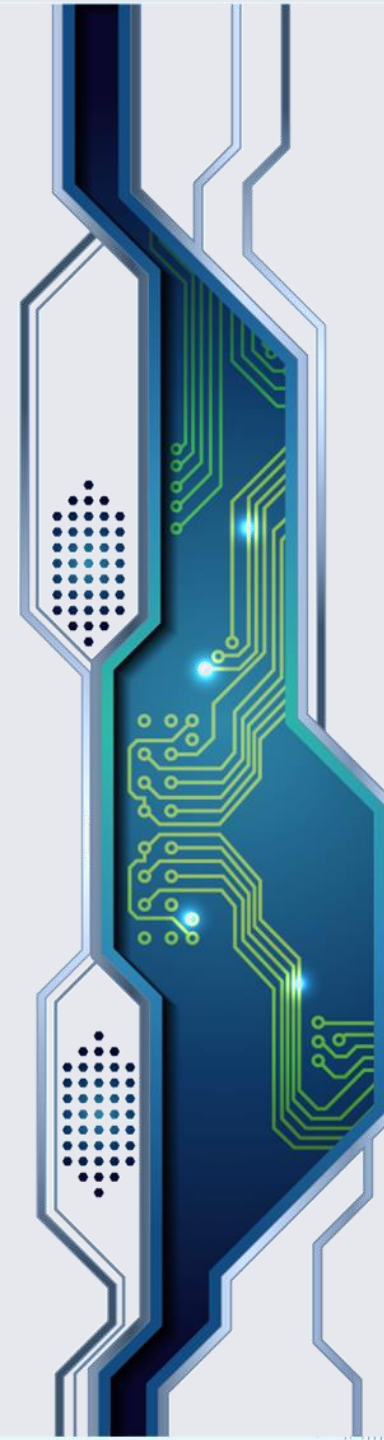


COFI Advanced Instrumentation and Analysis Techniques  
School

December 9-17, 2023



Principles (and future) of  
**Quantum Computing**

Jens Koch | Northwestern University

**logistics**

## Quantum Simulation for High-Energy Physics

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# HHL algorithm & ML relevance

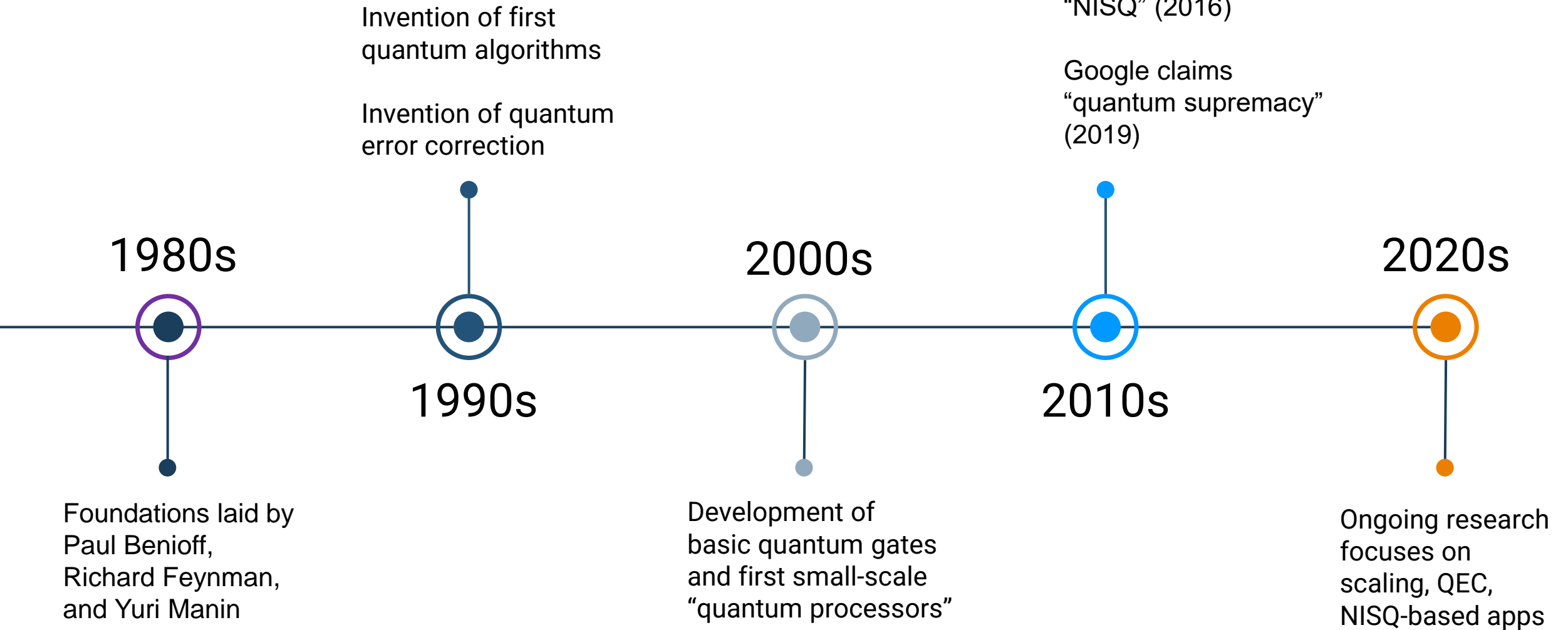
Harrow, Hassadim & Lloyd 2008

$$A\vec{x} = \vec{b}$$

$$\vec{x} = A^{-1}\vec{b}$$

Use of HHL in training of deep neural networks offers exponential speedup over classical training

# quantum computing history



# computer

INPUT



OUTPUT



# computer

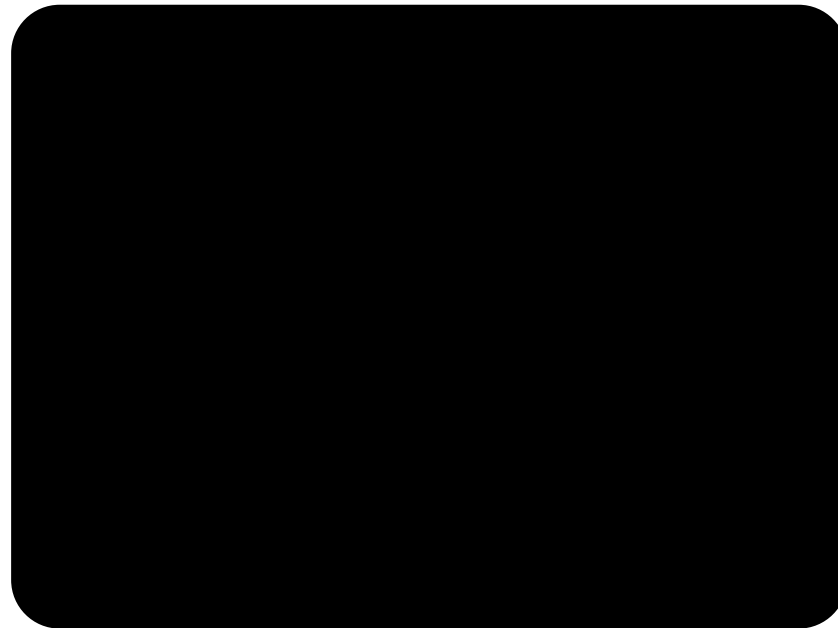
INPUT →  
011010100...101101010




→ OUTPUT  
0101010110...1001

# computer

INPUT   
011010100...101101010  
  
*n* binary digits



 OUTPUT  
0101010110...1001  
  
*m* binary digits




# computer

INPUT   
011010100...101101010

$$\vec{x} \in \mathbb{B}^n$$



 OUTPUT  
0101010110...1001

$$\vec{y} \in \mathbb{B}^m$$

## Notation

$$\mathbb{B} = \{0, 1\}, \quad n, m \in \mathbb{N}$$

# computer

INPUT

$$\vec{x} \in \mathbb{B}^n$$



$$f : \mathbb{B}^n \rightarrow \mathbb{B}^m$$



OUTPUT

$$\vec{y} \in \mathbb{B}^m$$

## Notation

$$\mathbb{B} = \{0, 1\}, \quad n, m \in \mathbb{N}$$

# computer

INPUT

$$\vec{x} \in \mathbb{B}^n$$



$$f : \mathbb{B}^n \times \Omega \rightarrow \mathbb{B}^m$$



OUTPUT

$$\vec{y} \in \mathbb{B}^m$$

## Notation

$$\mathbb{B} = \{0, 1\}, \quad n, m \in \mathbb{N}$$

$\Omega$  sample space  
for random variable(s)

# computer

a black box\* implementing map  $f$

## Examples



$$f : \mathbb{B}^n \times \Omega \rightarrow \mathbb{B}^m$$

\*a "device" obeying the laws of physics

# computer

a black box\* implementing map  $f$

## Examples



$$f: \mathbb{B}^n \times \Omega \rightarrow \mathbb{B}^m$$

\*a “device” obeying the laws of physics

# computer

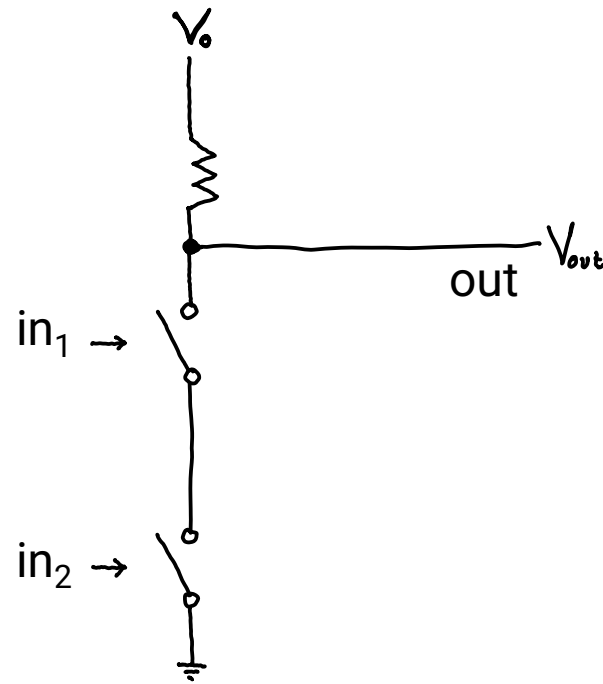
a black box\* implementing map  $f$

## Examples

NAND gate



in <sub>1</sub>	in <sub>2</sub>	out
0	0	1
0	1	0
1	0	0
1	1	0



$$f : \mathbb{B}^n \times \Omega \rightarrow \mathbb{B}^m$$

\*a "device" obeying the laws of physics

# computer

a black box\* implementing map  $f$

## Examples



$$f : \mathbb{B}^n \times \Omega \rightarrow \mathbb{B}^m$$

\*a “device” obeying the laws of physics

# programmable computer

a black box\* that can implement\*\*  
any\*\*\* given  $f$

$$\text{NAND: } \mathbb{B}^2 \rightarrow \mathbb{B}^1$$

$$\text{random: } \Omega \rightarrow \mathbb{B}^1$$

$$\text{bubble sort: } \mathbb{B}^{64 \cdot 100} \rightarrow \mathbb{B}^{64 \cdot 100}$$

$$\text{PFD: } \mathbb{B}^{1024} \rightarrow \mathbb{B}^{1024 \cdot 2}$$

prime factor decomposition

...

\* a “device” obeying the laws of physics

\*\* with input & output sizes  $n, m \leq N_0$

\*\*\* output is provided in finite time



# classical vs. quantum computer

NAND:  $\mathbb{B}^2 \rightarrow \mathbb{B}^1$

random:  $\Omega \rightarrow \mathbb{B}^1$

bubble sort:  $\mathbb{B}^{64 \cdot 100} \rightarrow \mathbb{B}^{64 \cdot 100}$

PFD:  $\mathbb{B}^{1024} \rightarrow \mathbb{B}^{1024 \cdot 2}$

prime factor decomposition

...

# classical

“uses purely classical physics paradigms for computing”



# quantum

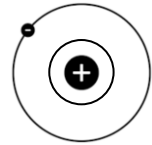
“uses concepts specific to quantum physics”

# Jozsa & Linden (2003)

For any quantum algorithm operating on pure states, we prove that the presence of **multi-partite entanglement**, with a number of parties that **increases unboundedly with input size**, is **necessary** if the quantum algorithm is to offer an **exponential speed-up over classical computation**.

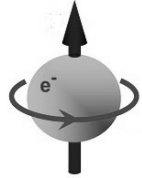
**entanglement**

# quantum computing hardware



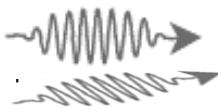
$|1s\rangle, |2s\rangle$

atoms  
ions  
NV centers



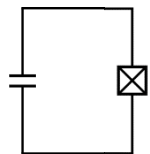
$|\uparrow\rangle, |\downarrow\rangle$

spins

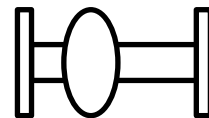


$|H\rangle, |V\rangle$

photons





$|0\rangle, |1\rangle$



superconducting  
circuits and cavities

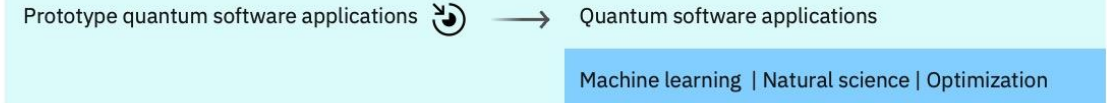
# Development Roadmap

Executed by IBM   
On target 

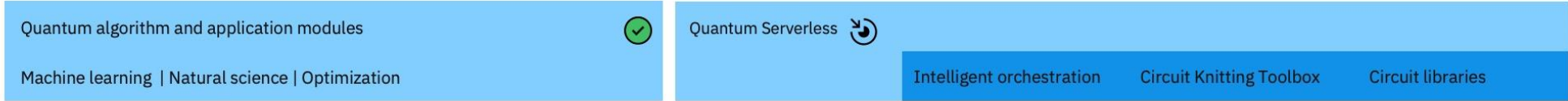
IBM Quantum

2019 	2020 	2021 	2022 	2023	2024	2025	2026+
Run quantum circuits on the IBM cloud	Demonstrate and prototype quantum algorithms and applications	Run quantum programs 100x faster with Qiskit Runtime	Bring dynamic circuits to Qiskit Runtime to unlock more computations	Enhancing applications with elastic computing and parallelization of Qiskit Runtime	Improve accuracy of Qiskit Runtime with scalable error mitigation	Scale quantum applications with circuit knitting toolbox controlling Qiskit Runtime	Increase accuracy and speed of quantum workflows with integration of error correction into Qiskit Runtime

Model Developers



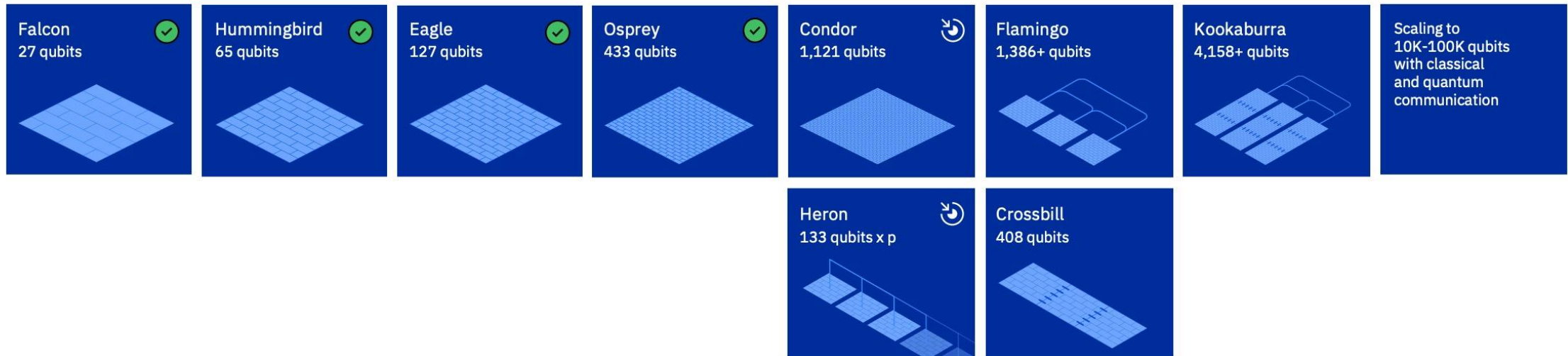
Algorithm Developers



Kernel Developers



System Modularity



# Quantum Computing Market Map

Non exhaustive and in no particular order. Excludes details on control systems, assembly languages, circuit design, etc.

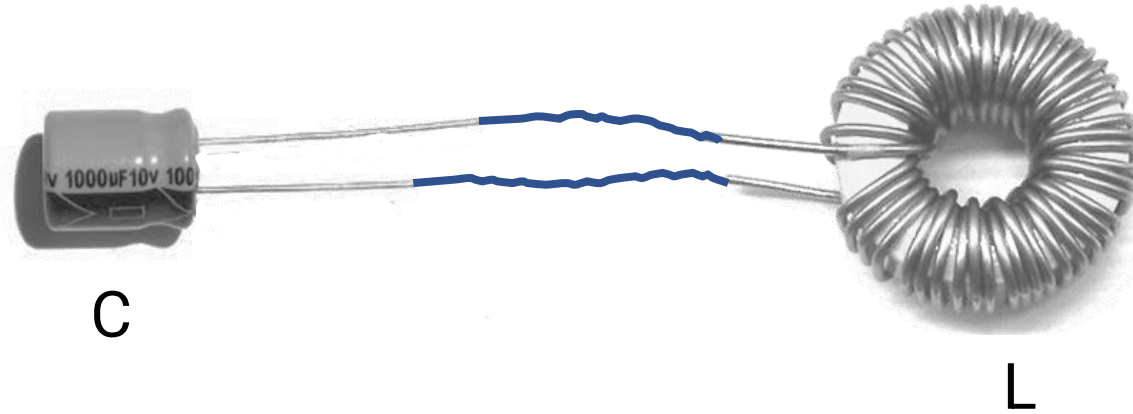
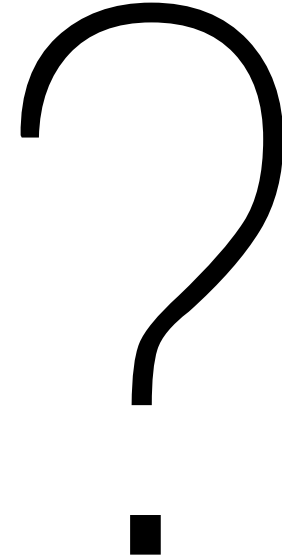
Users <i>Select examples</i>	Applications <i>Not mapped to verticals</i>	Software offerings <i>Includes control software</i>	QPUs <sup>2</sup>		Hardware / components <i>Select examples only – not representative of entire ecosystem</i>
Material Science	Not strictly categorized given diversity of operations <sup>1</sup>		Superconducting		Cryogenics (includes testing)
Finance			<p>Ion Trap</p>		<p>Neutral Atoms</p>
	<p>Life Sciences</p>		<p>Lights and lasers</p>		
	<p>Cloud access to QPUs</p>		<p>Silicon</p>		<p>Other componentry (examples)</p>
Other	<p>Simulators / q-inspired / etc</p>		<p>Photonics</p>		<p>Other</p>

<sup>1</sup> Software offerings can be further classified into SDKs, firmware / enablers, algorithms / applications, simulators etc. but many companies are offering a mixture across the stack

<sup>2</sup> Many QPU providers are offering full stack services (e.g. Pasqal acquired Qu&Co, Quantinuum was originally CQC prior to merger with HQS, etc.)



# Building quantum bits from **circuits**

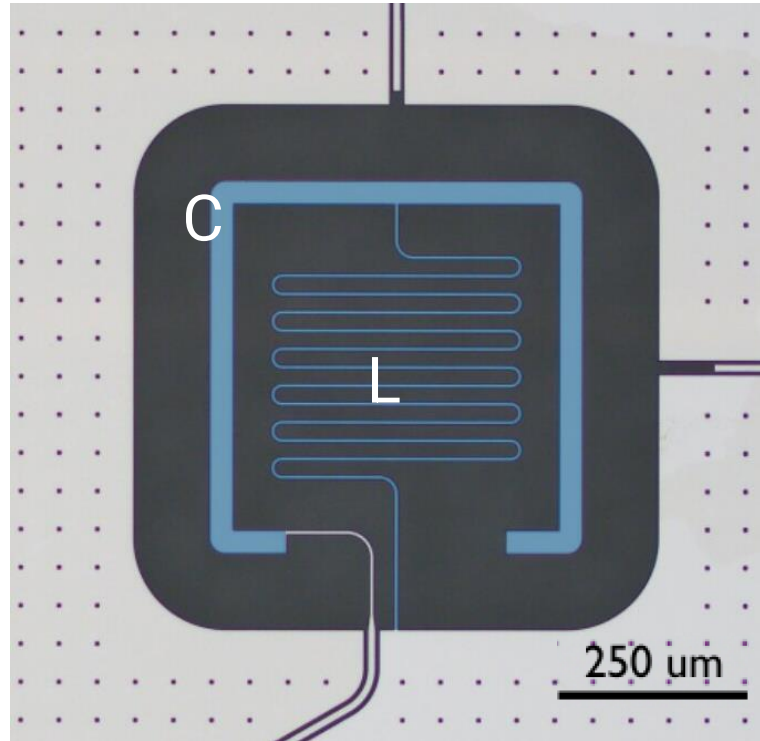
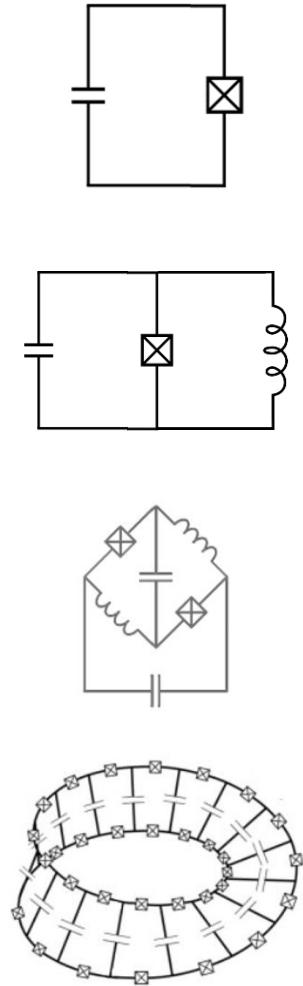


Ohmic losses  
energy dissipation,  $V = RI$

interaction with the environment **dramatically reduces quantum effects**



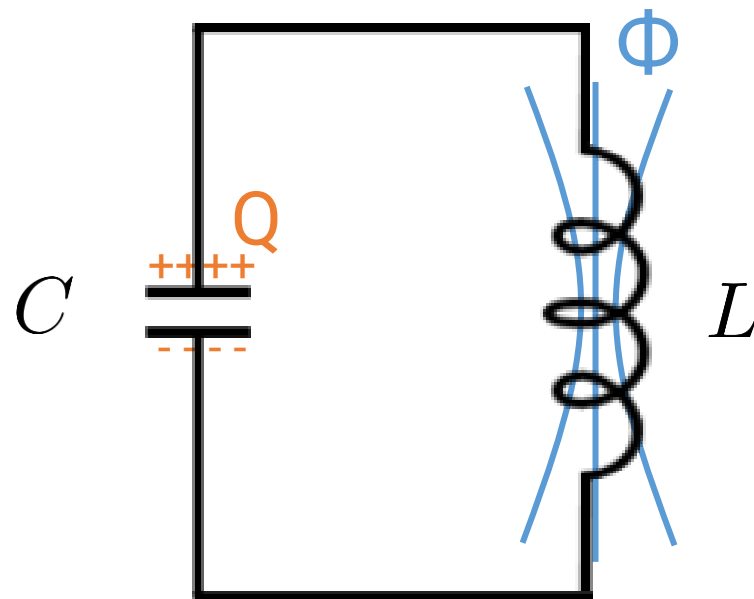
# Building quantum bits from superconducting circuits



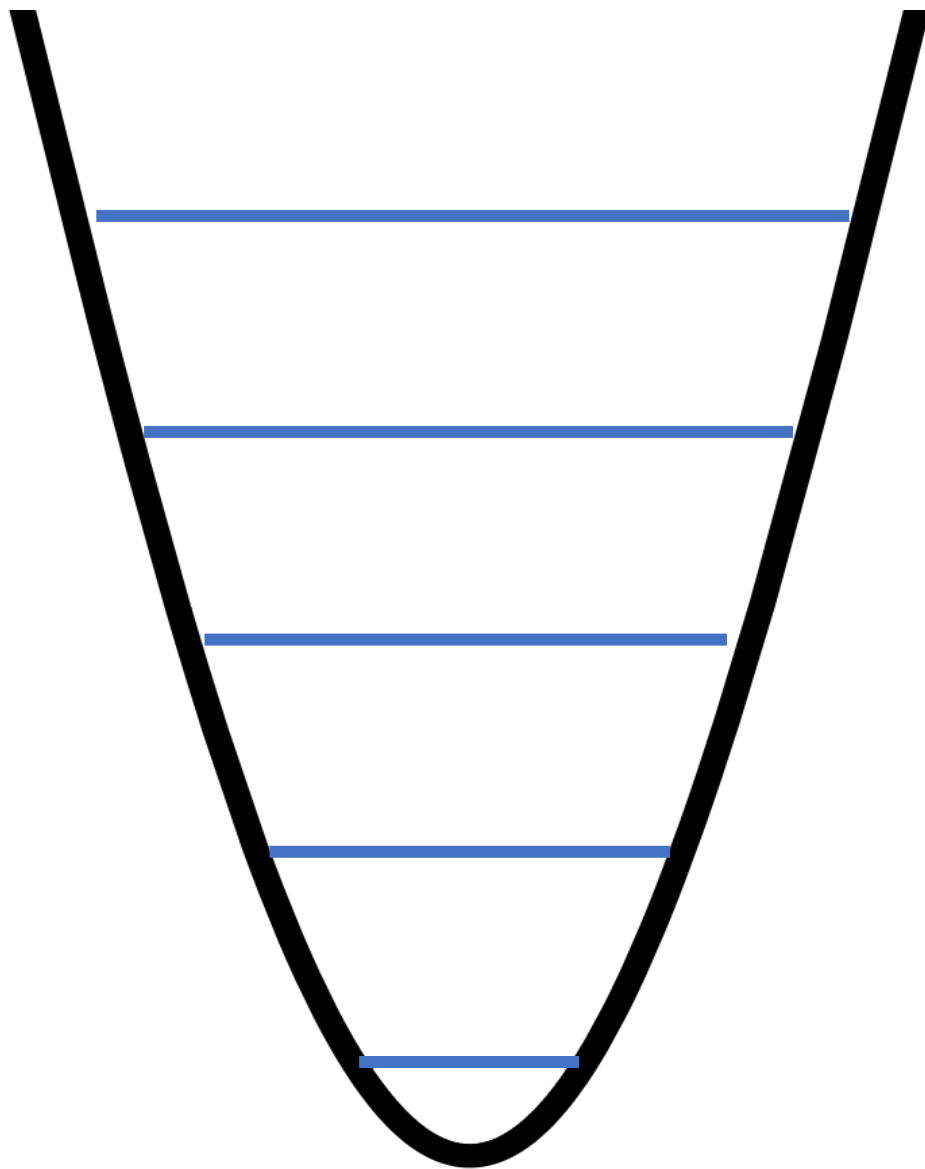
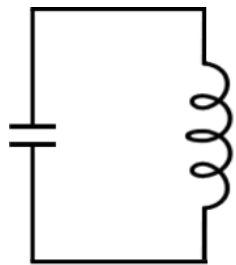
$$H = \frac{Q^2}{2C} + \frac{\Phi^2}{2L}$$

charging energy    inductive energy

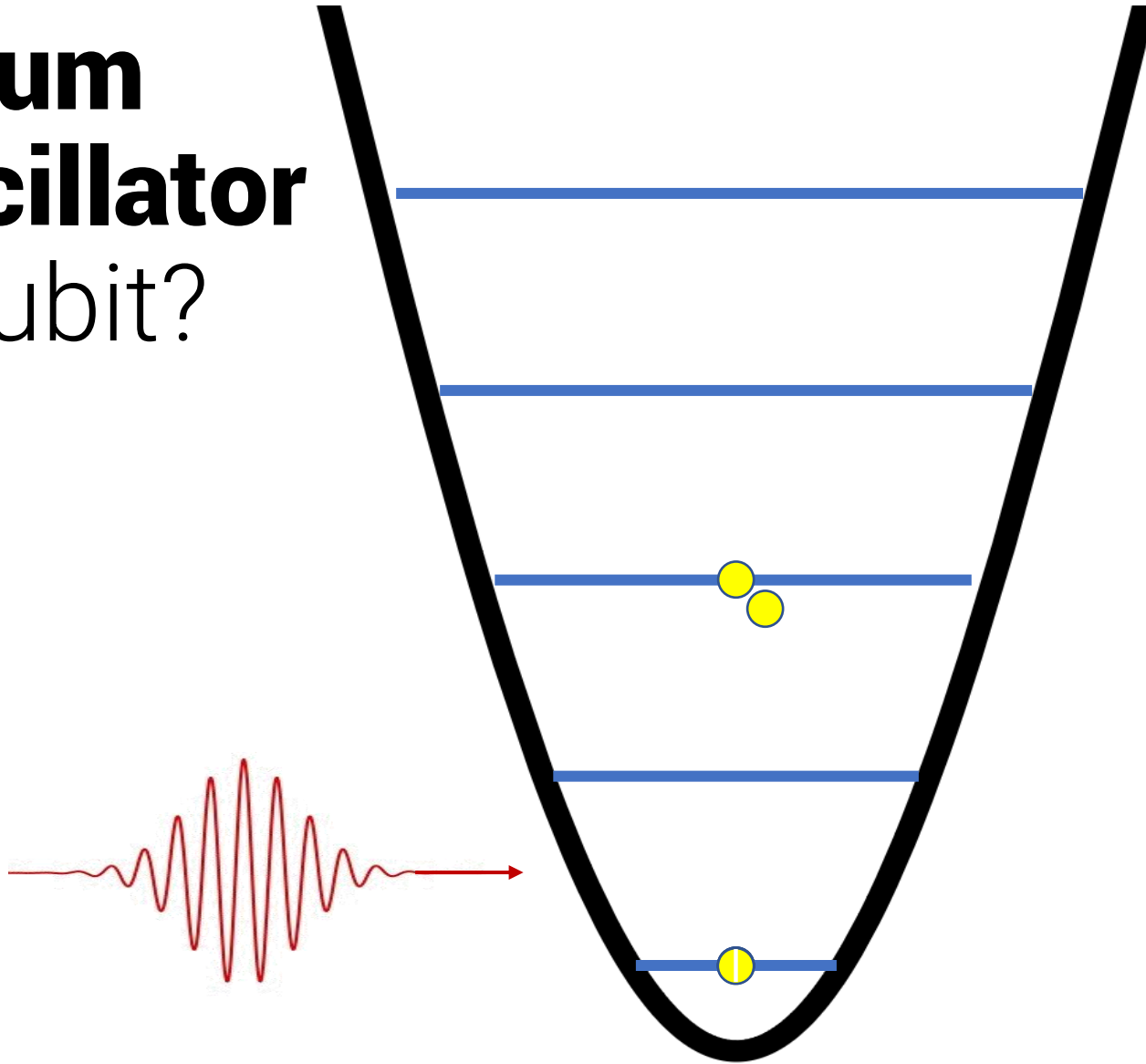
$$= \hbar\omega \left( a^\dagger a + \frac{1}{2} \right)$$



$$[\Phi, Q] = i\hbar$$

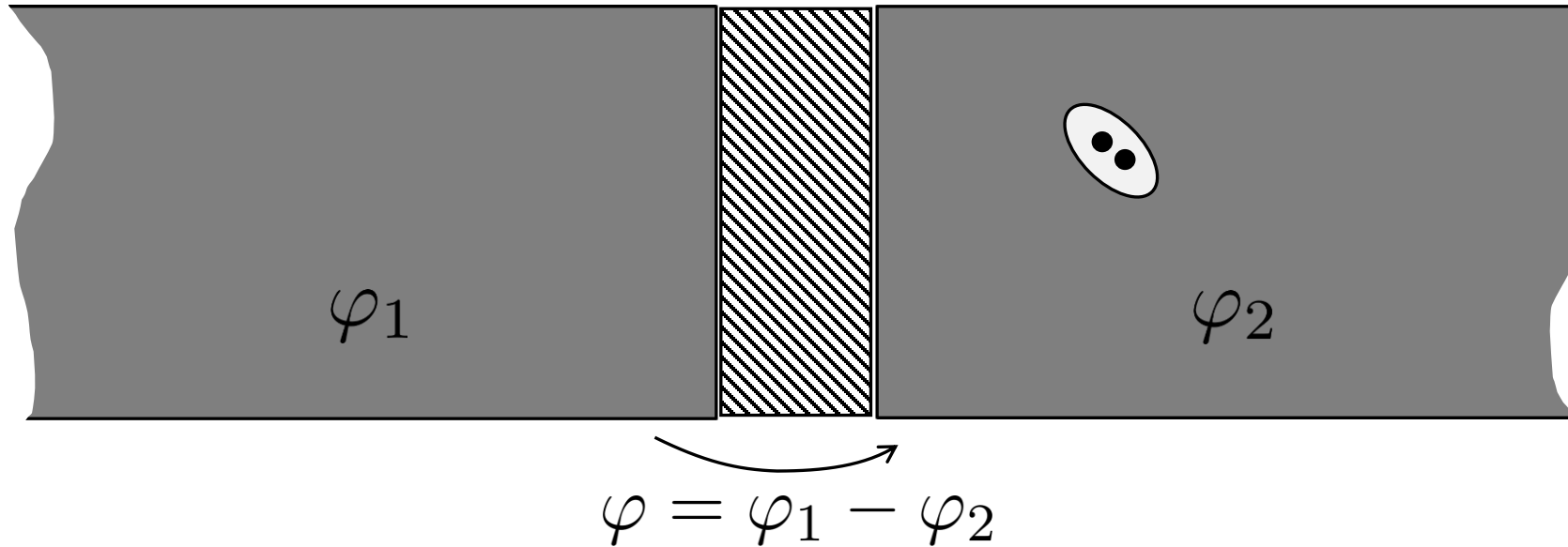


# Quantum LC oscillator as a qubit?



**not** a good qubit...  
need **nonlinearity**

# From LC oscillator to transmon



Brian Josephson  
Nobel Prize 1973

## Josephson effect

current-phase relation:  $I = I_0 \sin \varphi$

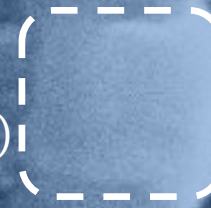
phase evolution equation:  $\frac{d\varphi}{dt} = V \cdot \frac{2\pi}{\Phi_0}$

# Josephson junction

V 1 = 141.4 nm



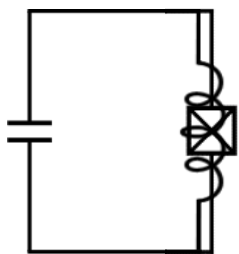
Bottom  
SC electrode 1 (Al)



Top  
SC electrode 1 (Al)



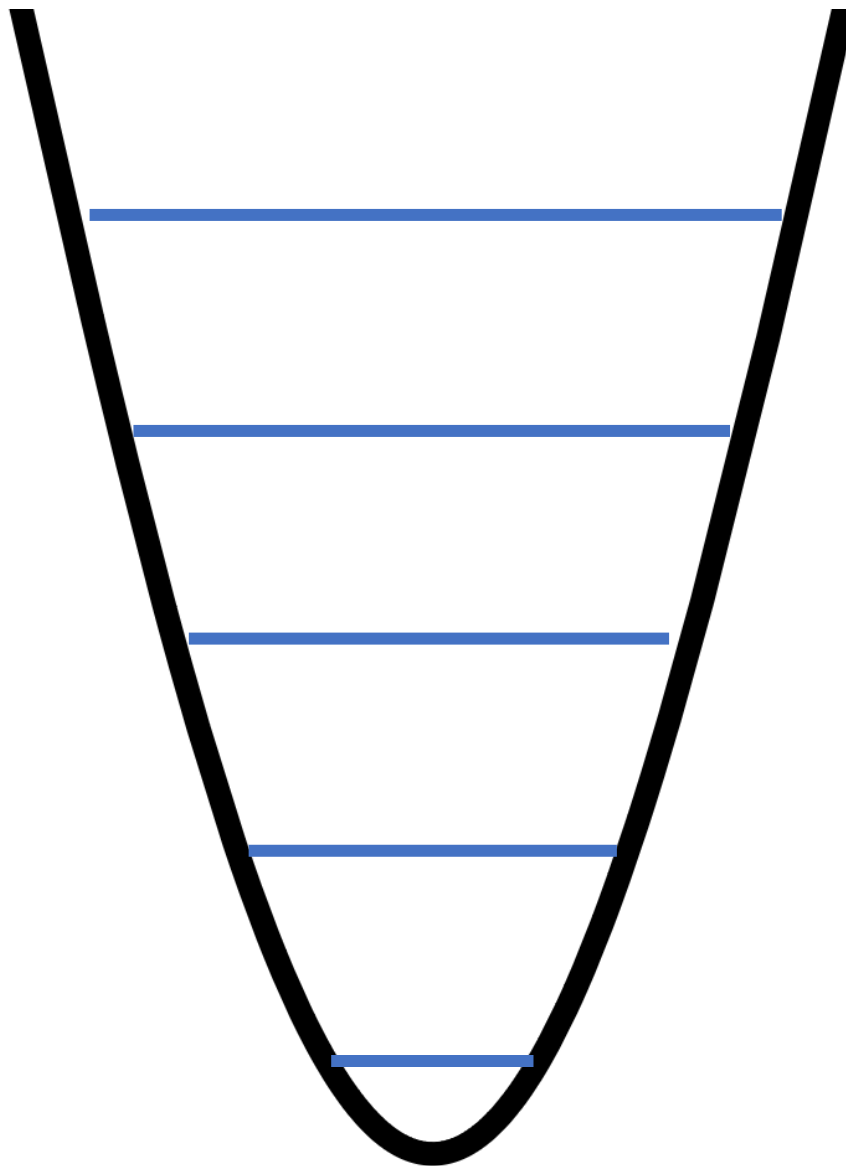
H 1 = 121.0 nm

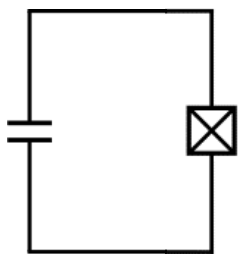


$$\frac{\Phi^2}{2L}$$



$$-E_J \cos(2\pi\Phi/\Phi_0)$$

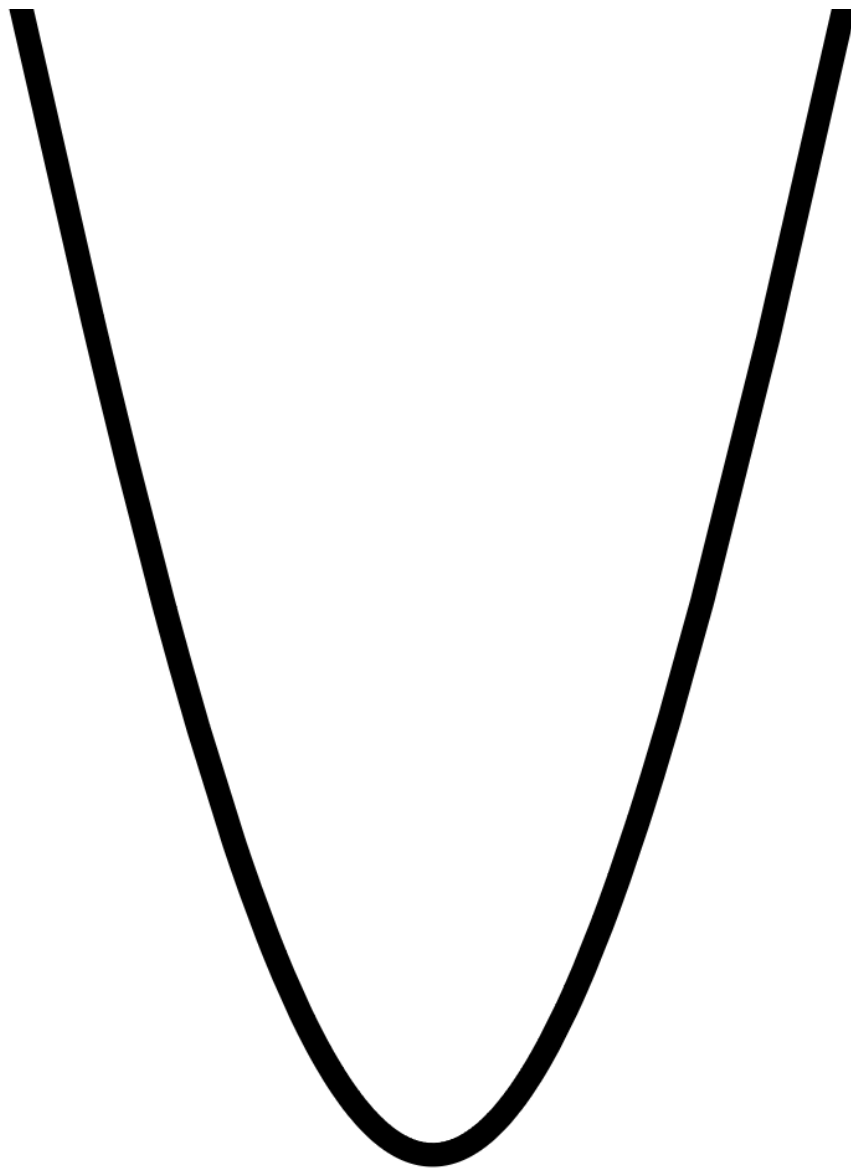




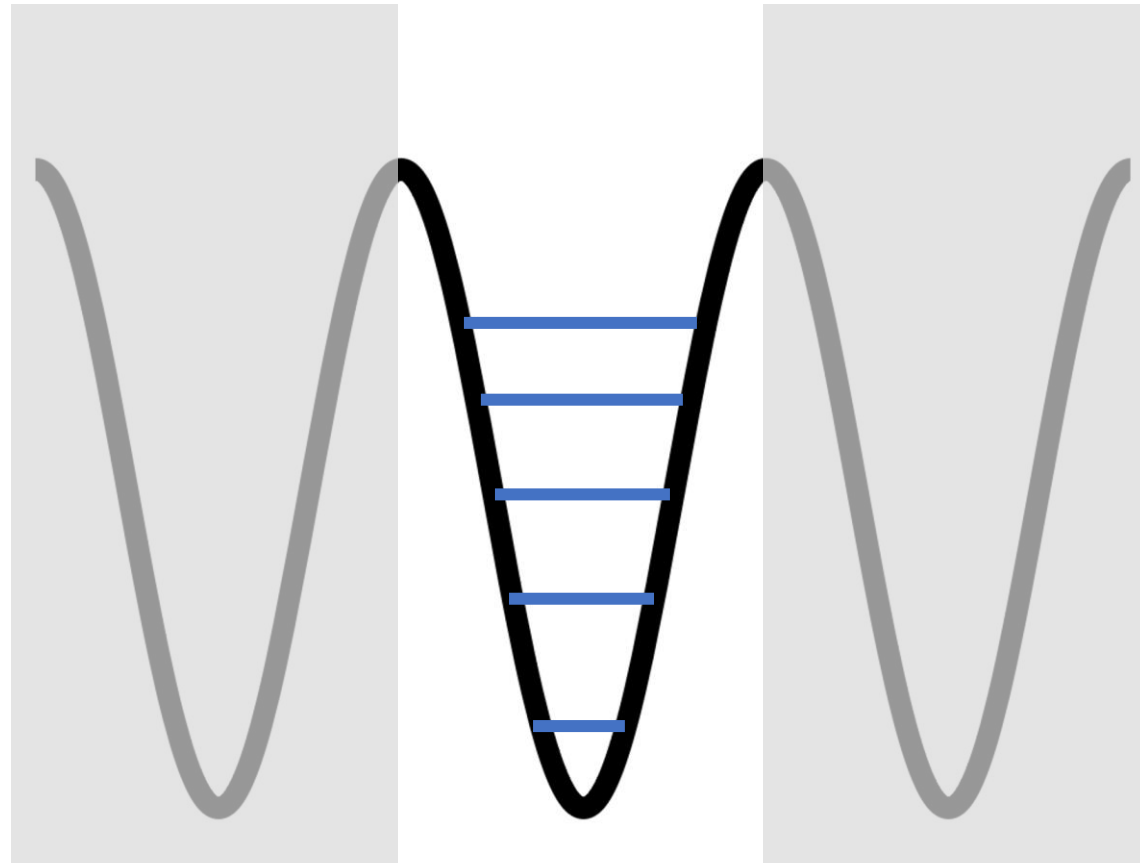
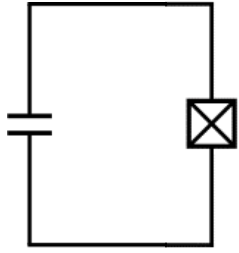
$$\frac{\Phi^2}{2L}$$



$$-E_J \cos(2\pi\Phi/\Phi_0)$$









quantum computing:

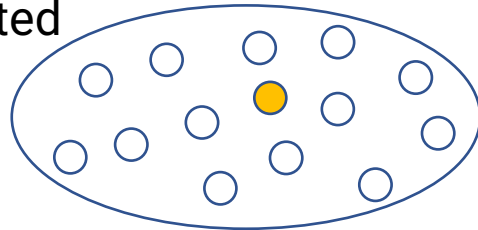
use quantum parallelism  
and entanglement  
to beat classical  
computers

# Shor's algorithm

$N = p q$  ( $p, q$  prime)  
period-finding of modular functions

# Grover's algorithm

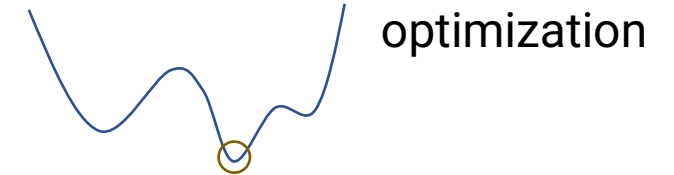
search unsorted  
database



# Harrow-Hassidim-Lloyd (HHL) algorithm

matrix inversion\*  $\vec{x} = \mathcal{A}^{-1}\vec{b}$

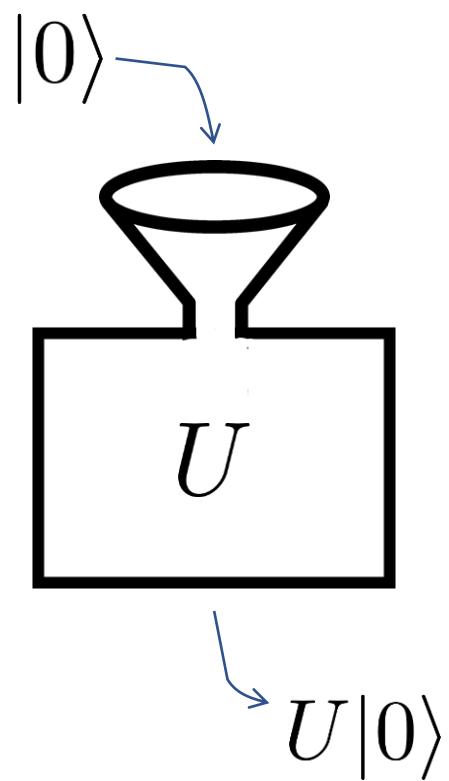
# quantum adiabatic algorithm



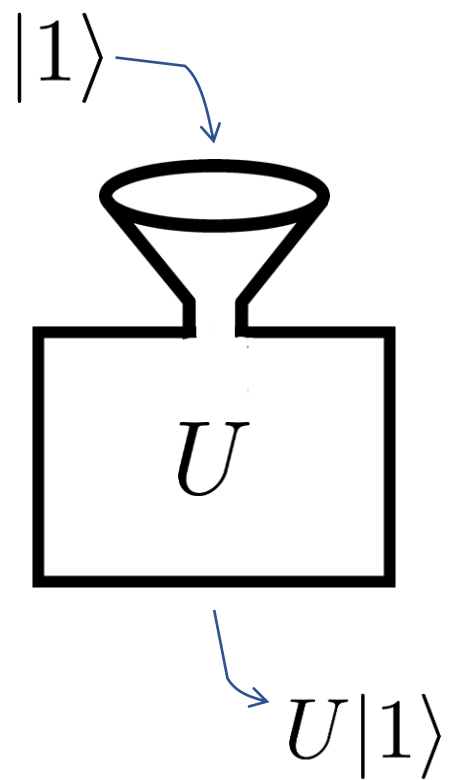
# digital quantum simulation

$$i\hbar \frac{d}{dt} |\Psi(t)\rangle = H |\Psi(t)\rangle$$

# quantum parallelism

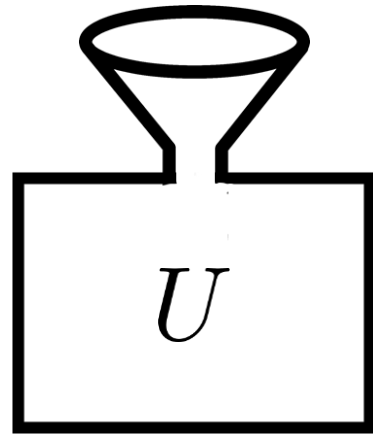


# quantum parallelism



# What makes a quantum computer tick?

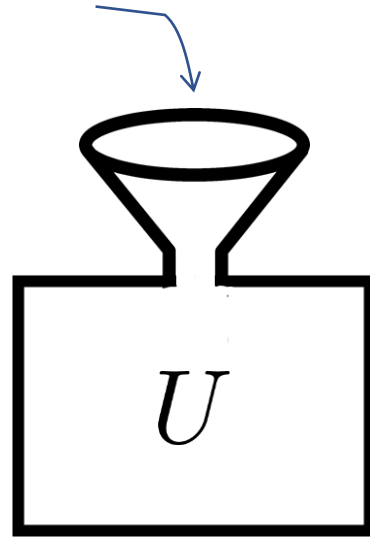
$|01010010100\rangle$



$U|01010010100\rangle$

# quantum parallelism

$$|00000 \dots 0\rangle + |00000 \dots 1\rangle + \dots + |11111 \dots 1\rangle$$



$$U|00000 \dots 0\rangle + U|00000 \dots 1\rangle + \dots + U|11111 \dots 1\rangle$$



# storing quantum states

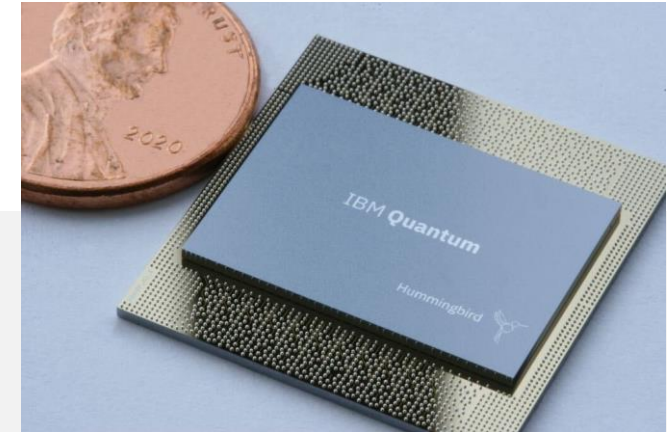
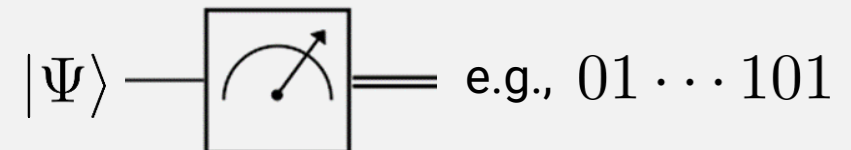


$$\begin{aligned}
 |\Psi\rangle = & \alpha_0 |00 \dots 000\rangle \\
 & + \alpha_1 |00 \dots 001\rangle \\
 & + \alpha_2 |00 \dots 010\rangle \\
 & + \alpha_3 |00 \dots 011\rangle \\
 & \vdots \\
 & + \alpha_{2^N} |1 \dots 111\rangle
 \end{aligned}$$

- $2^N$  complex amplitudes
- storage:  $2^N \times 128$  bits

$N = 65$  qubits  $\longrightarrow$  memory/disk space: 590 billion GB

- **but:** output only  $N$  classical bits

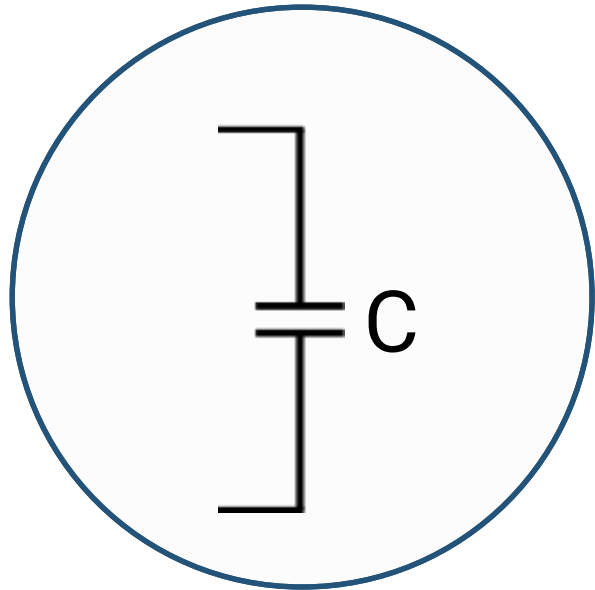


## quantum parallelism

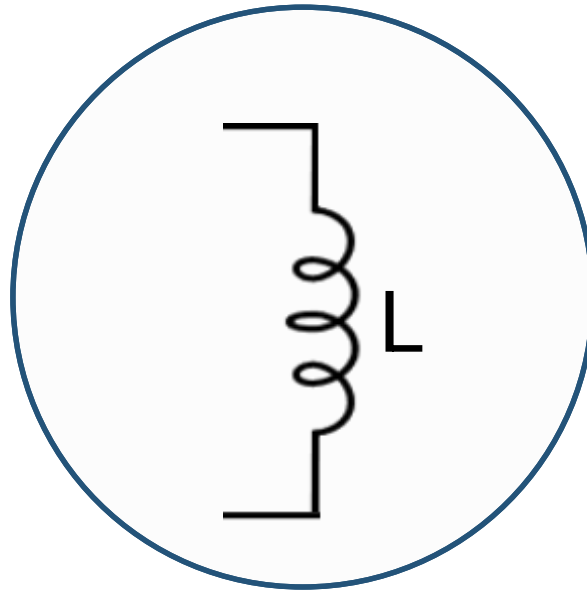
- can process all possible inputs at once
- **not sufficient** for quantum speedup
  - measurement only accesses one output state

# Circuit elements

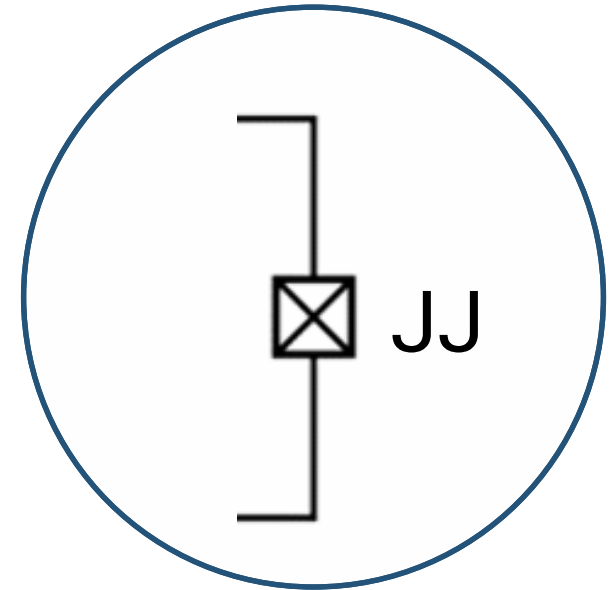
Building blocks for superconducting qubits



Capacitance

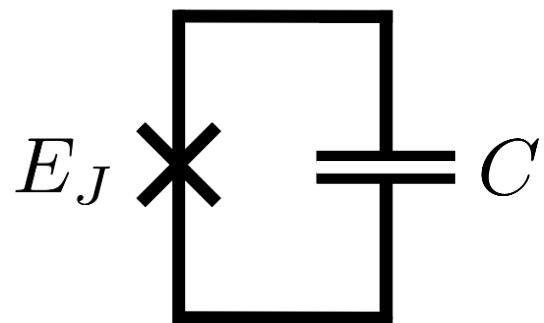


Inductance



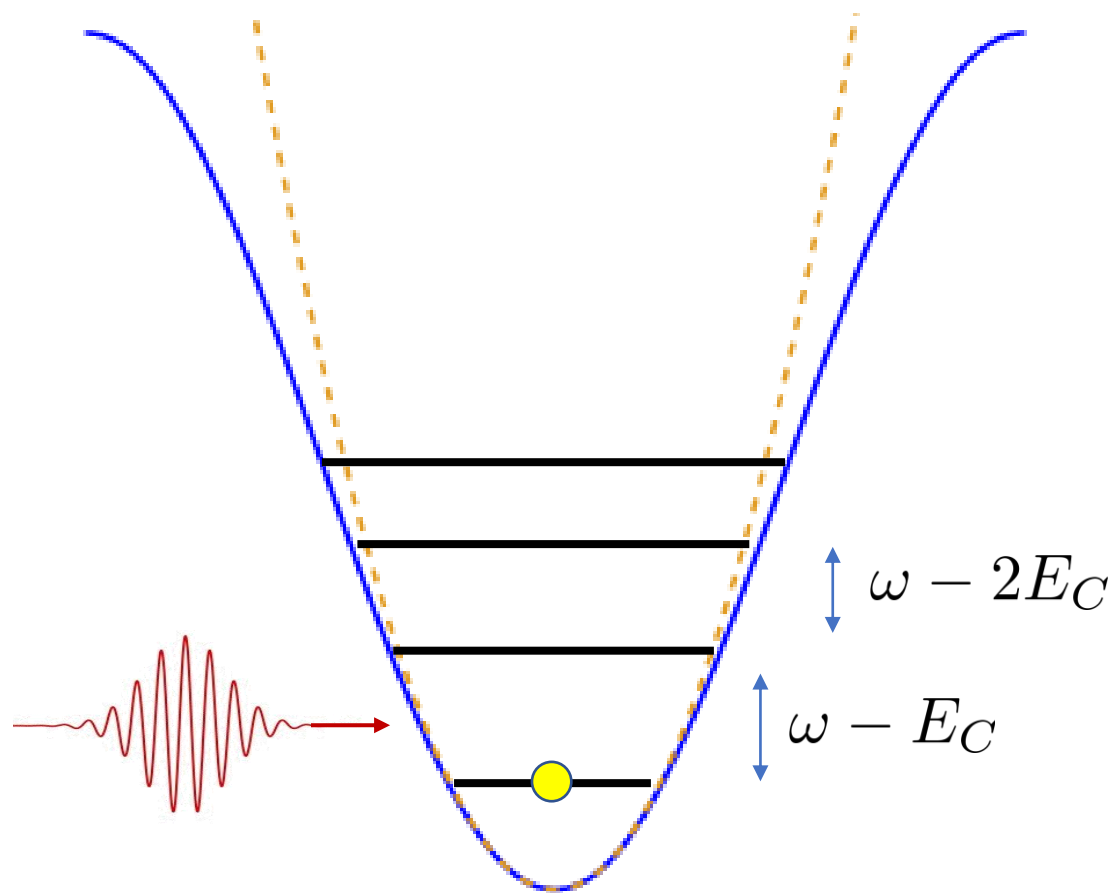
Josephson junction

# transmon qubit



**Introduce nonlinearity:**  
replace inductor with  
Josephson junction

$$H = \frac{Q^2}{2C} - E_J \cos(\overbrace{2\pi\Phi/\Phi_0}^{\varphi})$$



# transmon qubit

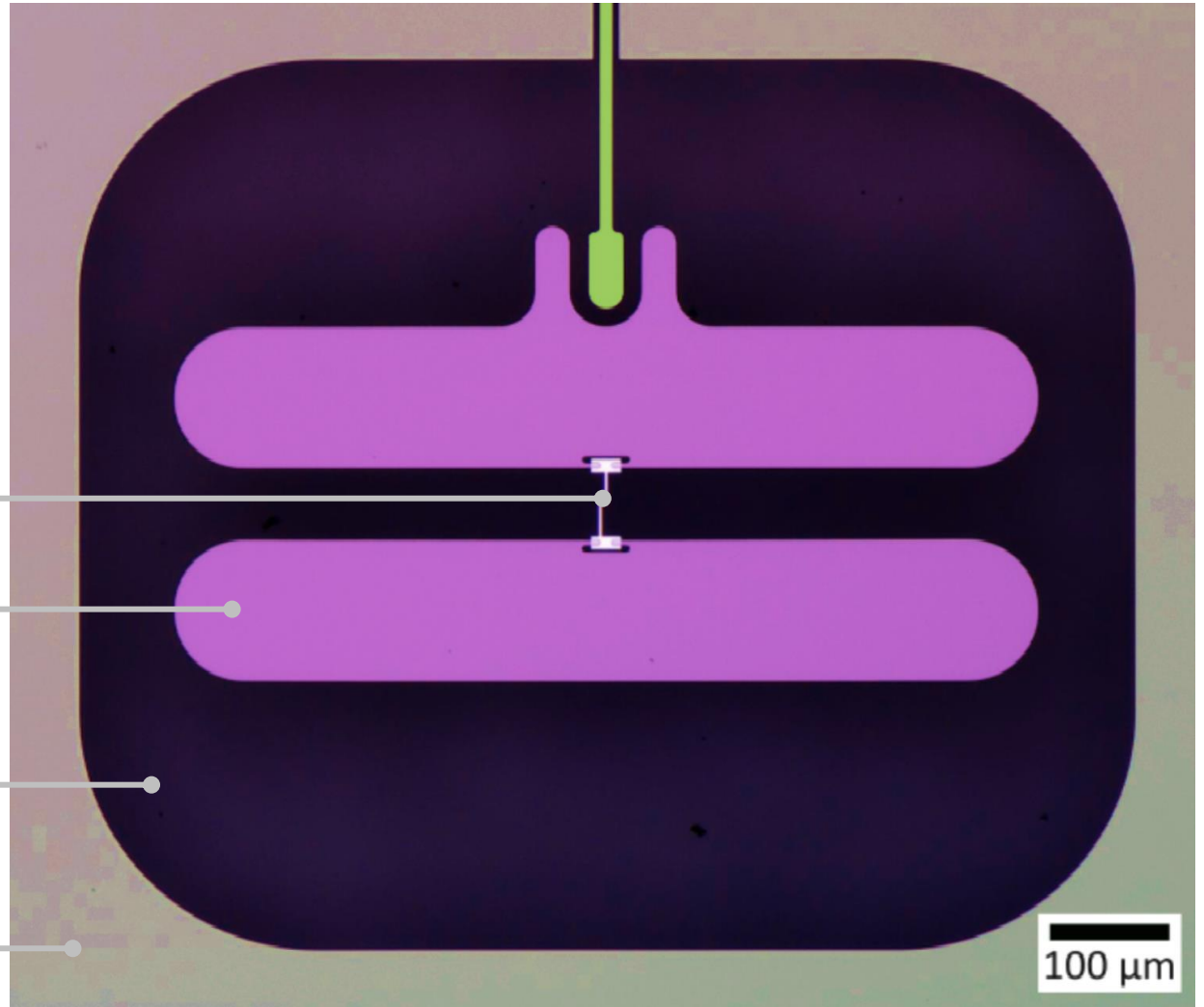
JJ  
(Al-AlOx)

capacitor pads  
(Nb, Al, or Ta)

substrate  
(Si or sapphire)

ground plane  
(Nb, Al, or Ta)

Houck lab



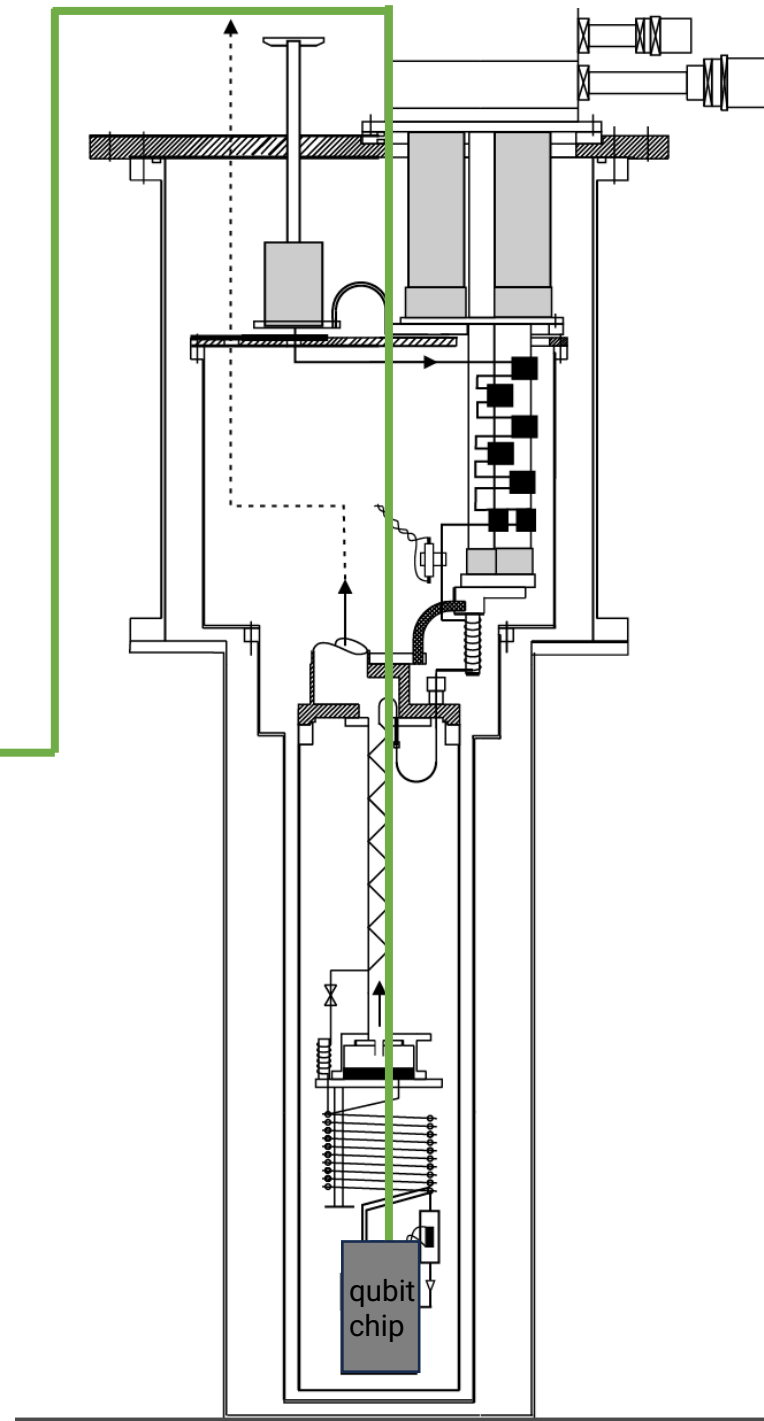
100 μm

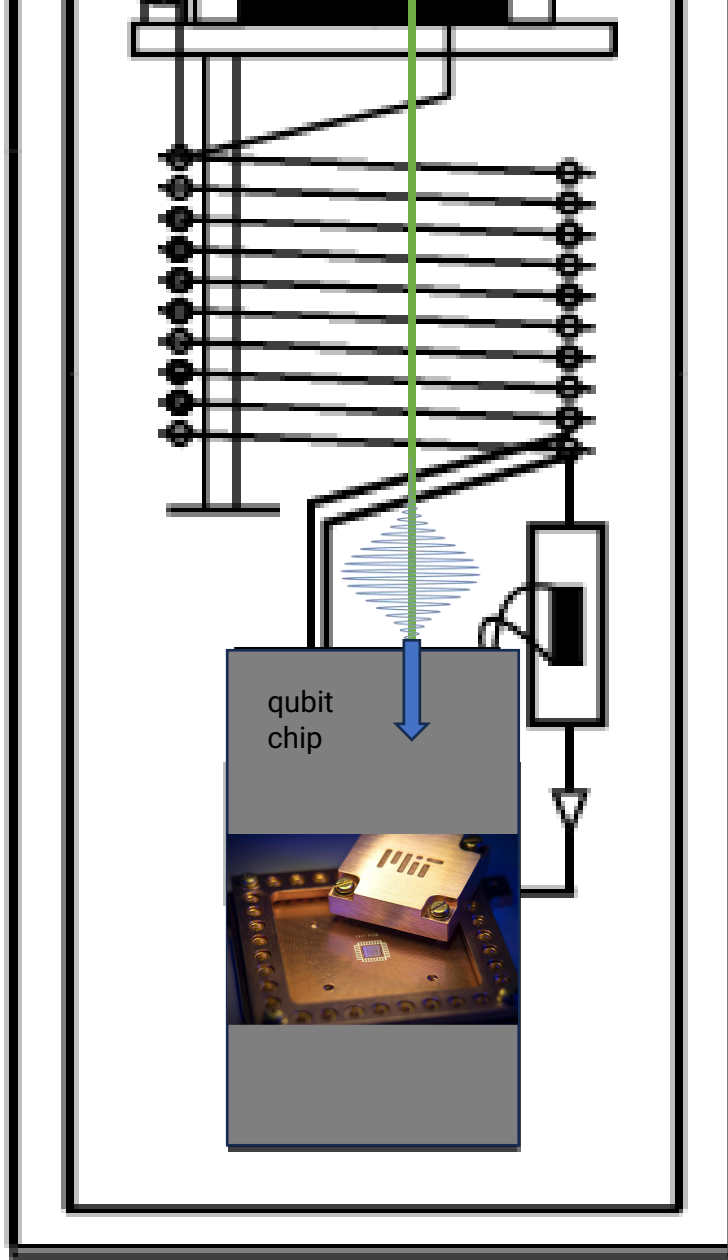
# single-qubit gates



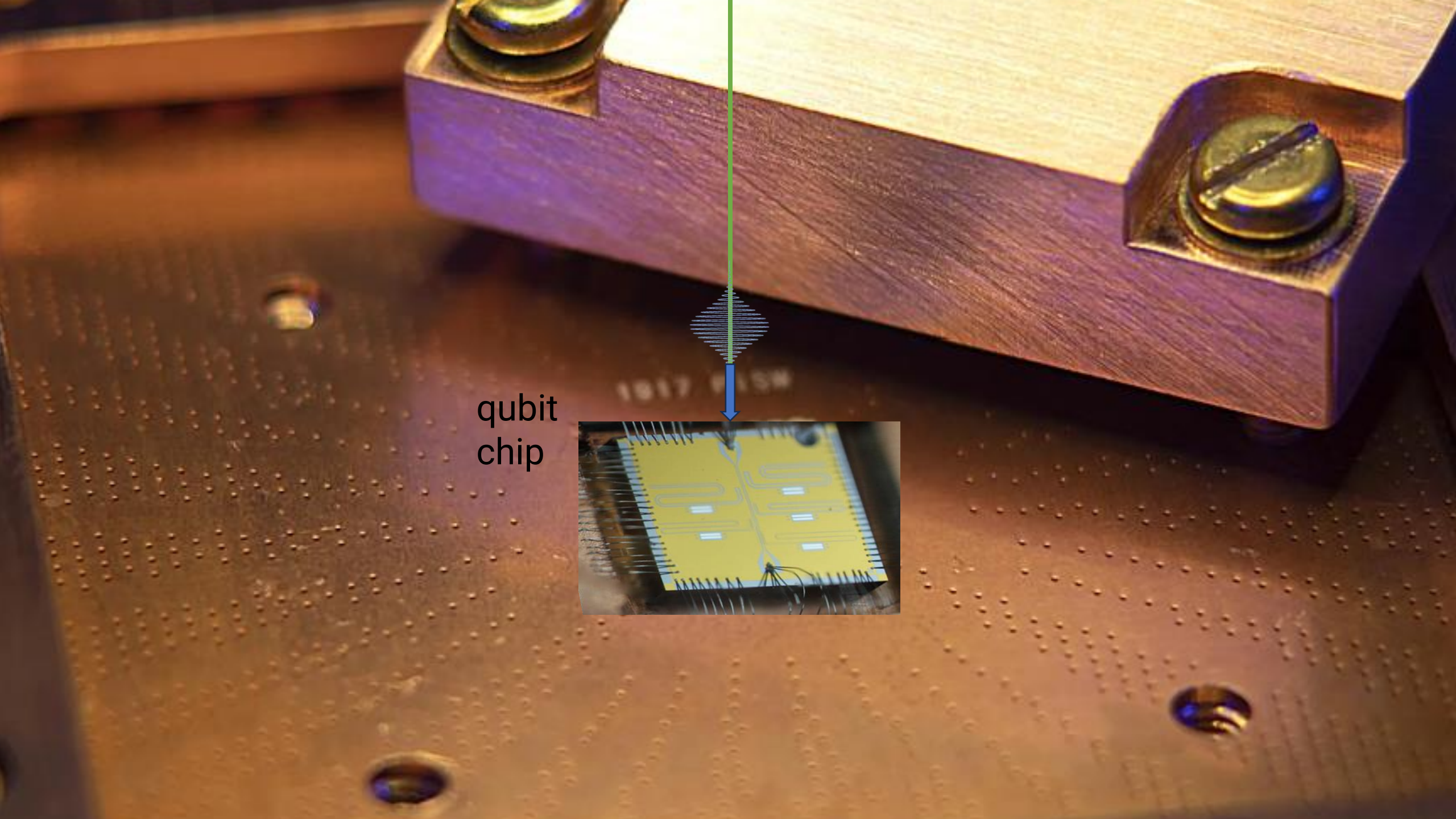
Microwave generator  
(AWG or FPGA / DAC)

$\sim$  GHz

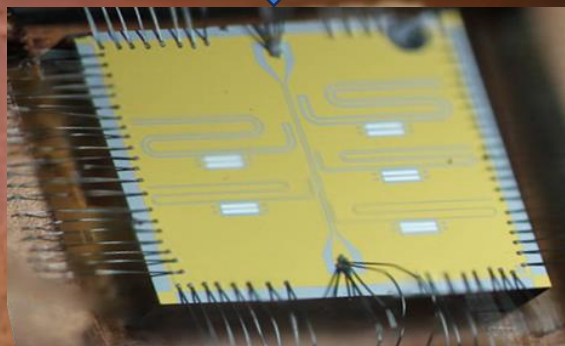




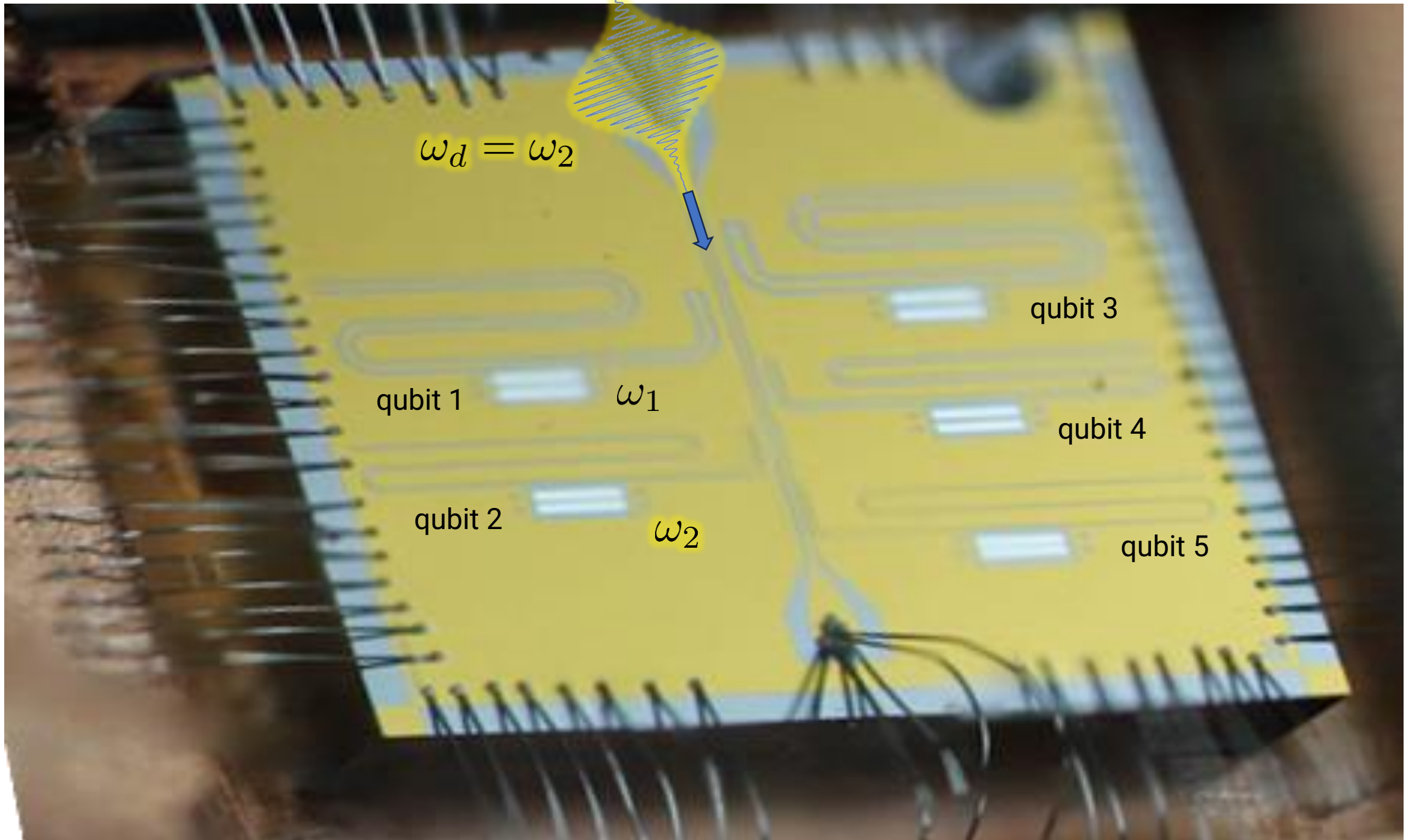




qubit  
chip







$$\omega_d = \omega_2$$

qubit 1

$$\omega_1$$

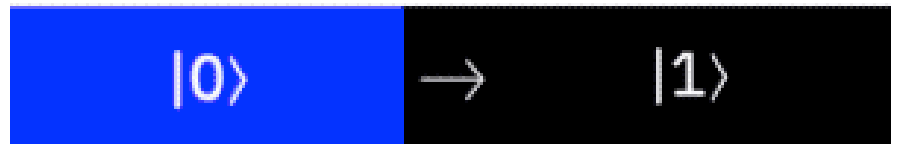
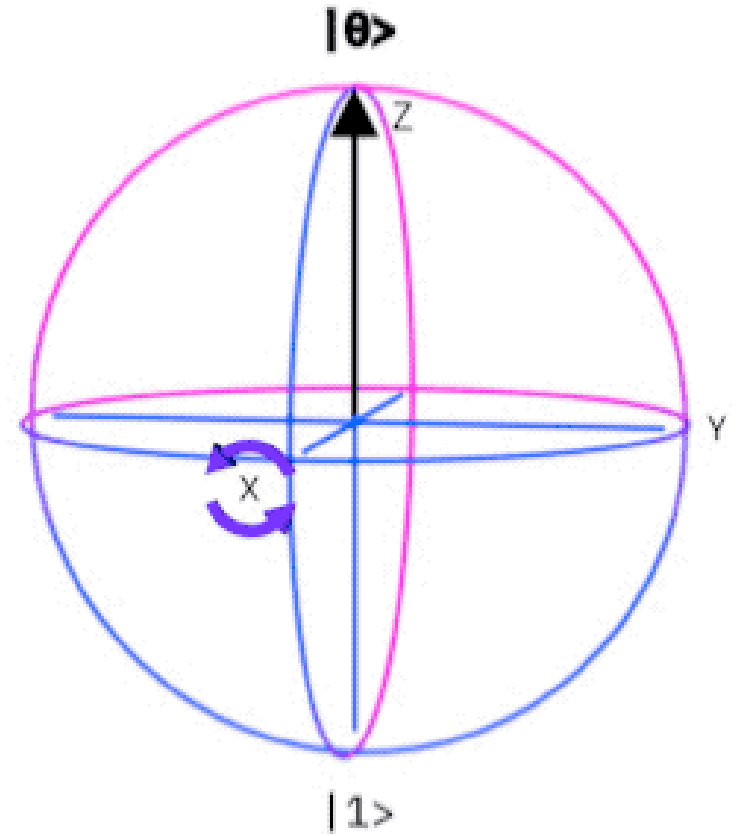
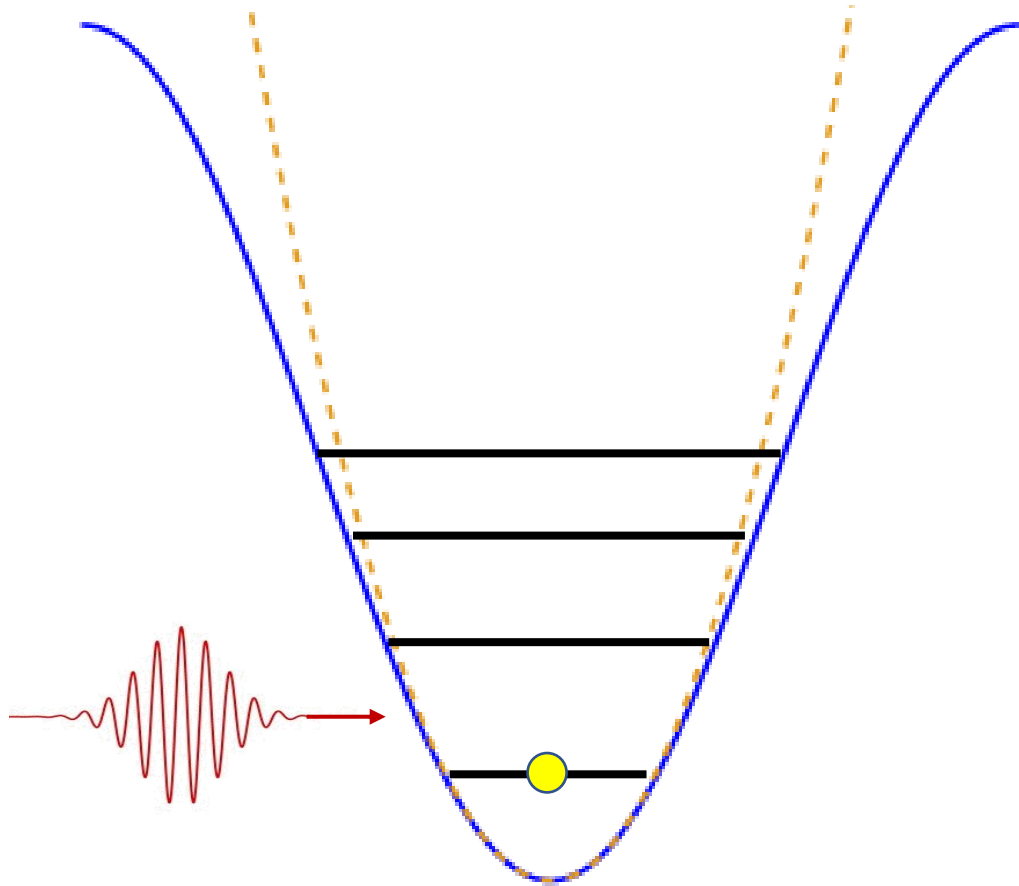
qubit 3

qubit 2

$$\omega_2$$

qubit 4

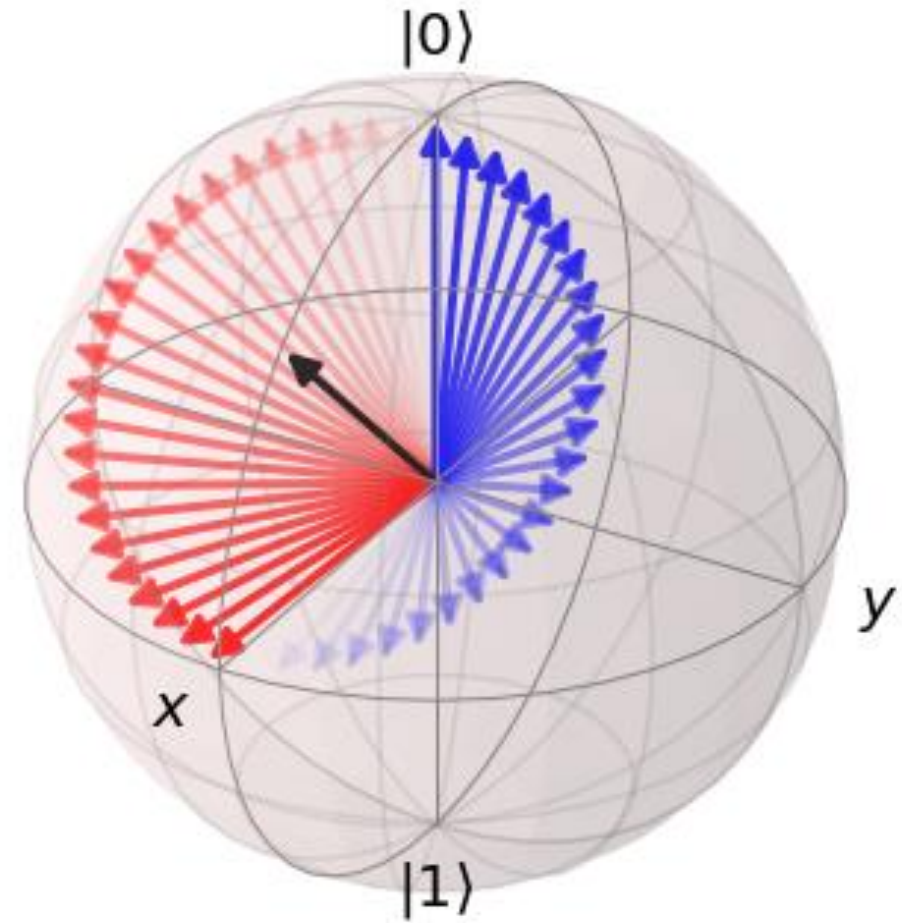
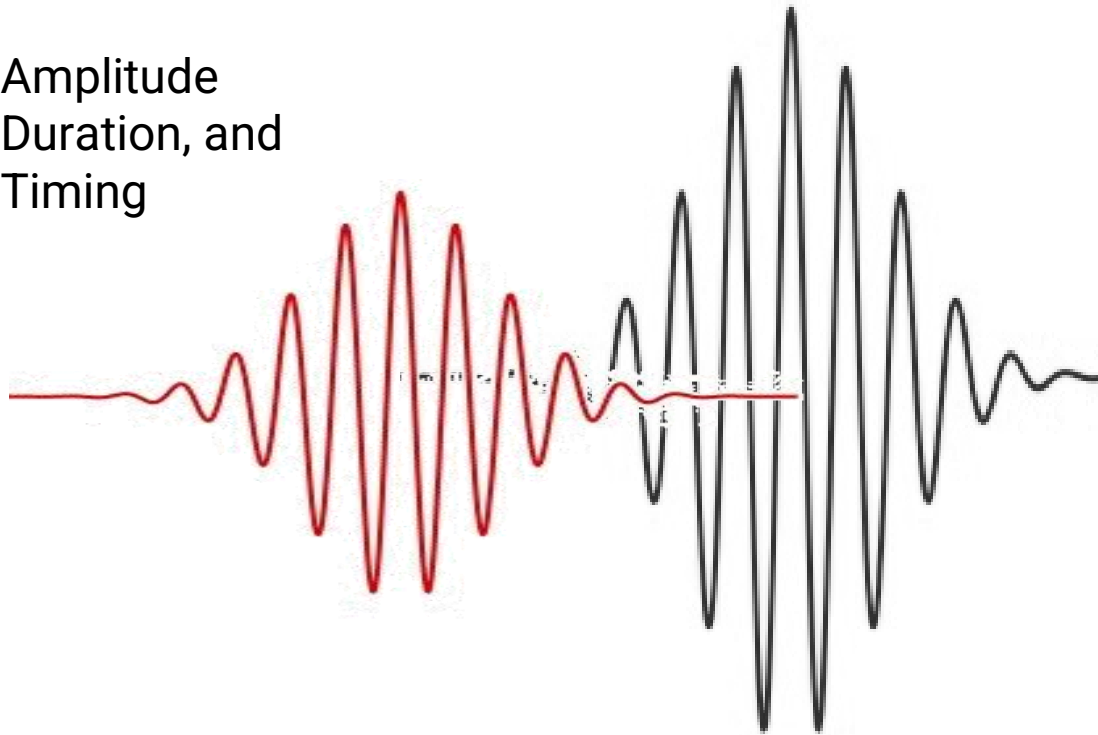
qubit 5



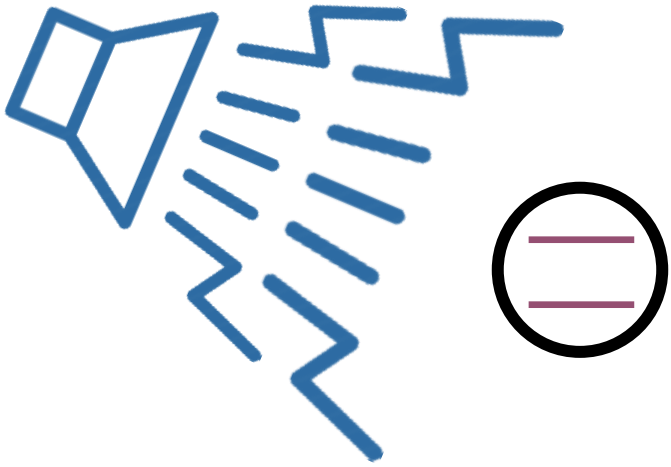
# arbitrary single-qubit gates

Vary pulse

- Amplitude
- Duration, and
- Timing



# noise



## QUBIT

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$$

1/f charge noise  
1/f flux noise  
critical current noise  
dielectric loss  
photon shot noise  
Purcell decay  
nonequilibrium quasiparticles  
...

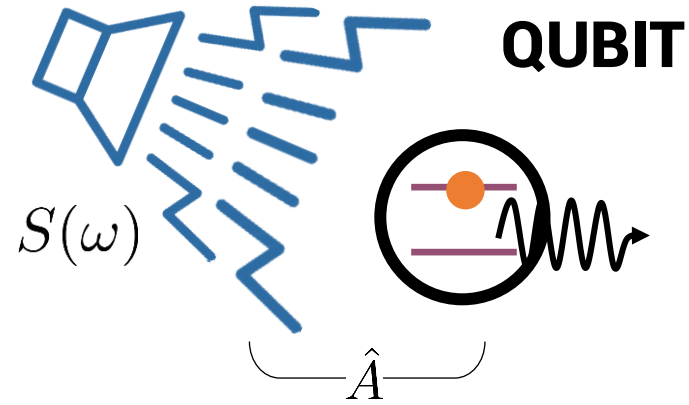
## DECOHERENCE / ERRORS



Your qubit ran into a problem that it couldn't handle, and now it needs to restart.

You can search for the error online: [HBAR\\_DIED\\_A\\_HORRIBLE\\_DEATH](#)

# noise



**QUBIT**

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$$



$|0\rangle$

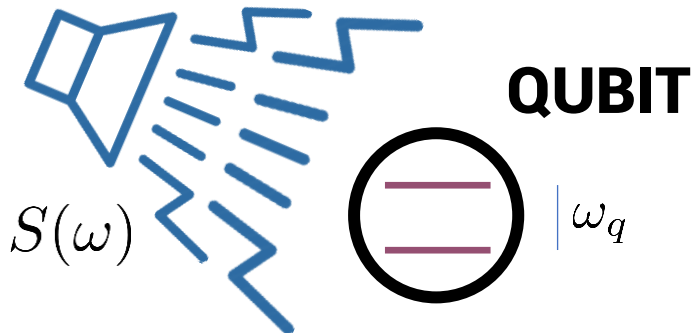
## **DECOHERENCE**

**depolarization**  
( $T_1$ )

rate: (Fermi's golden rule)

$$1/T_1 \sim \underbrace{|\langle 0|\hat{A}|1\rangle|^2}_{\text{transition matrix element}} \underbrace{S(\omega_q)}_{\text{noise spectral density}}$$

# noise



$$\rho = |\psi\rangle\langle\psi| = \begin{pmatrix} |\alpha|^2 & \alpha\beta^* e^{-i\omega_q t} \\ \alpha^* \beta e^{i\omega_q t} & |\beta|^2 \end{pmatrix}$$

$$\begin{pmatrix} |\alpha|^2 & 0 \\ 0 & |\beta|^2 \end{pmatrix}$$

## DECOHERENCE

**depolarization**  
( $T_1$ )

rate: (Fermi's golden rule)

$$1/T_1 \sim |\langle 0|\hat{A}|1\rangle|^2 S(\omega_q)$$

**pure dephasing**  
( $T_\varphi$ )

$$1/T_\varphi \sim \left( \frac{\partial \omega_q}{\partial \lambda} \right)^2 S(0)$$

linear sensitivity  
of qubit freq. to noise

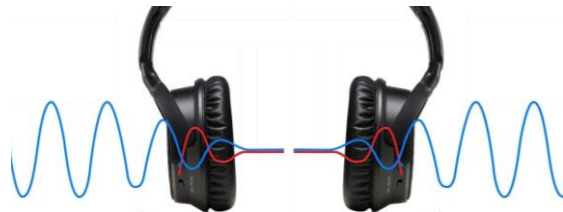
# error mitigation strategies



- 1 reduce noise  
(materials science, fabrication,  
microwave engineering, ...)



- 2 intrinsic noise protection  
(develop circuits that are insensitive to noise)



- 3 active quantum error correction  
(monitor for errors, apply correction steps  
when required in feedback loop)

# 1/f noise



**charge noise**  
**flux noise**

Cooper pair box

rf SQUID qubit, 3-junction flux qubit



# Cooper pair box

**letters to nature**

---

## **Coherent control of macroscopic quantum states in a single-Cooper-pair box**

**Y. Nakamura\*, Yu. A. Pashkin† & J. S. Tsai\***

*\* NEC Fundamental Research Laboratories, Tsukuba, Ibaraki 305-8051, Japan*

*† CREST, Japan Science and Technology Corporation (JST), Kawaguchi,  
Saitama 332-0012, Japan*

NATURE | VOL 398 | 29 APRIL 1999 | [www.nature.com](http://www.nature.com)

Physica Scripta. Vol. T76, 165–170, 1998

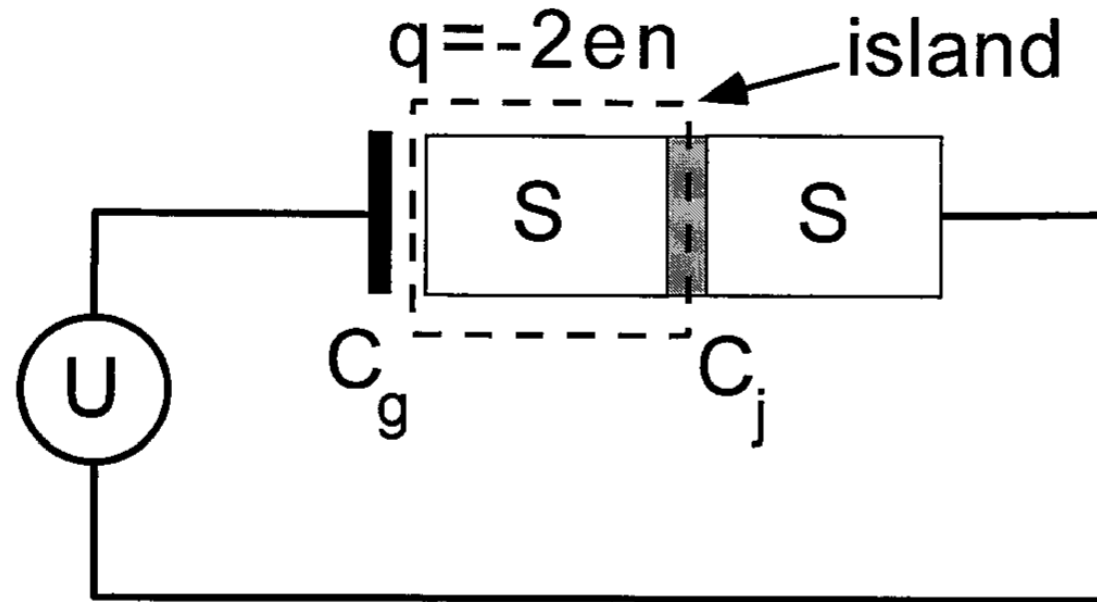
## **Quantum Coherence with a Single Cooper Pair**

V. Bouchiat,\* D. Vion, P. Joyez, D. Esteve and M. H. Devoret

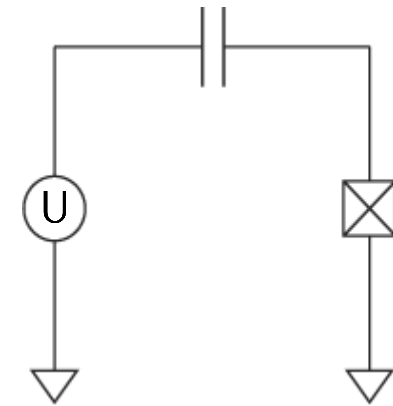
Quantronics group, Service de Physique de l'Etat Condensé CEA-Saclay, F-91191 Gif-sur-Yvette, France

*Received October 27, 1997; revised version received January 15, 1998; accepted January 23, 1998*

# Cooper pair box

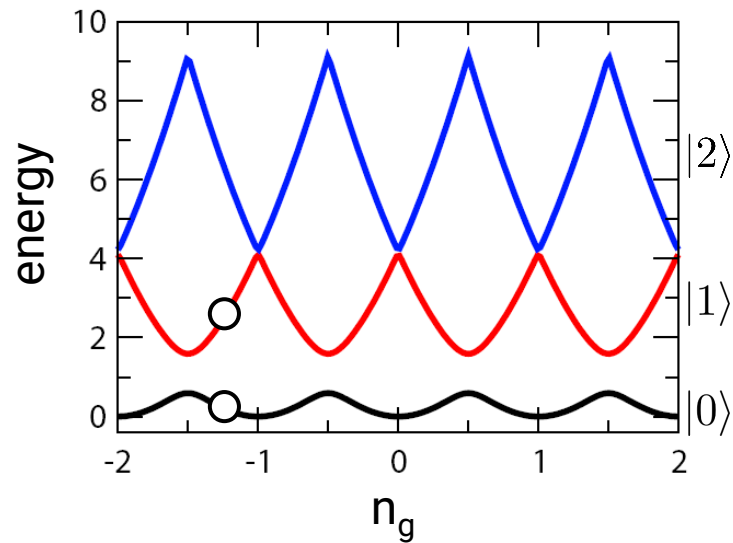
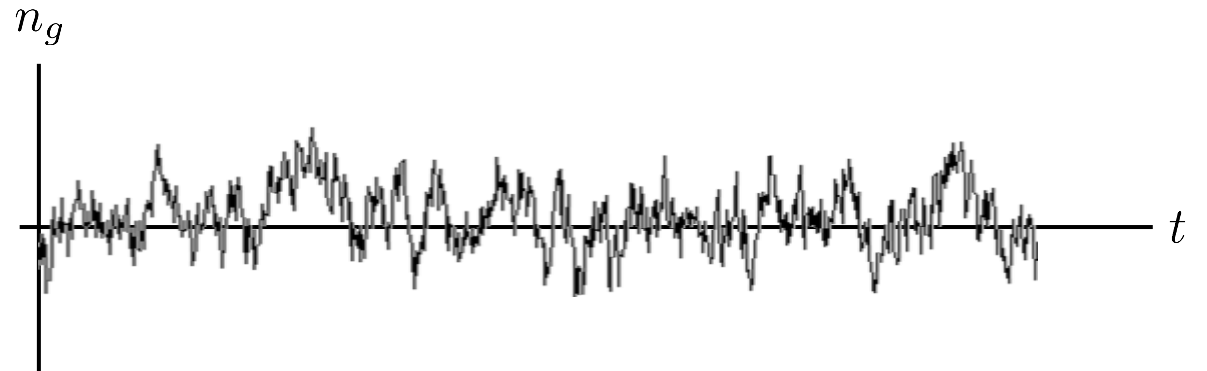
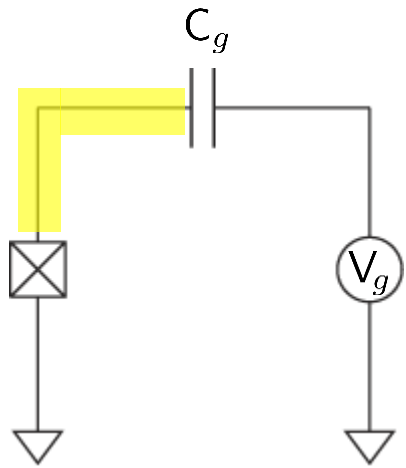


*Fig. 1.* Schematics of the single Cooper box: a superconducting electrode (island) is in contact with a superconducting reservoir through a tunnel junction (grey zone) with capacitance  $C_j$ . Excess Cooper pairs tunnel onto the island in response to an electric field applied by means of the gate capacitance  $C_g$  and voltage  $U$ .



$$T_\varphi \approx 1 \text{ ns}$$

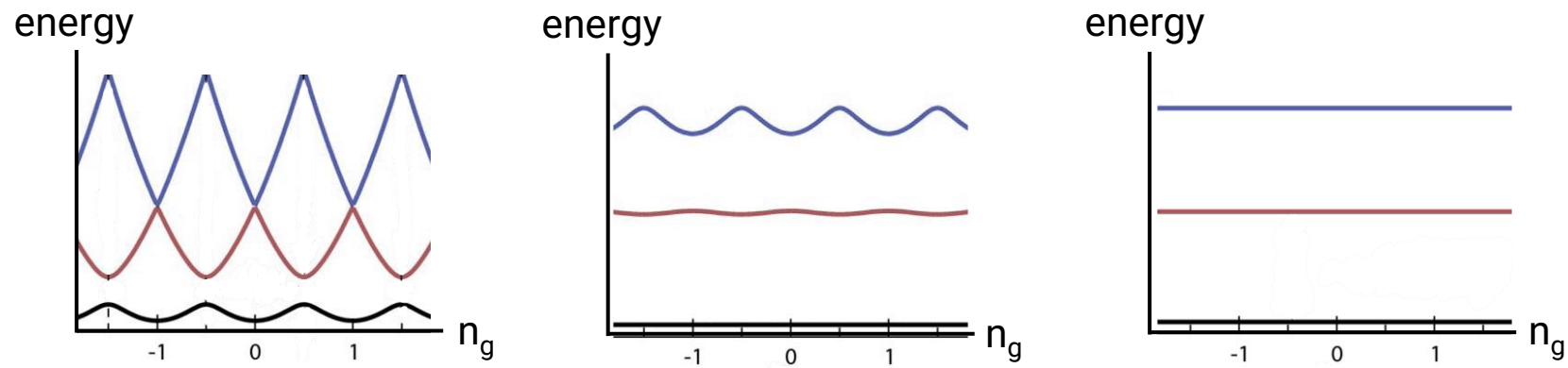
# charge noise dephasing



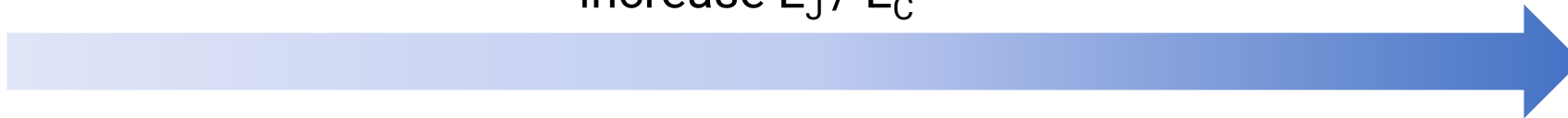
► dephasing

$$1/T_\varphi \sim \left( \frac{\partial \omega_q}{\partial \lambda} \right)^2 S(0)$$

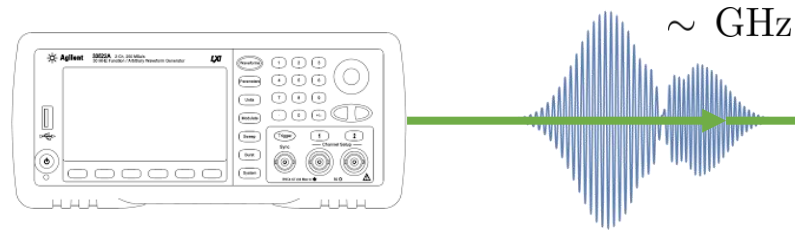
# Flattening the spectrum for charge noise protection



increase  $E_J / E_C$



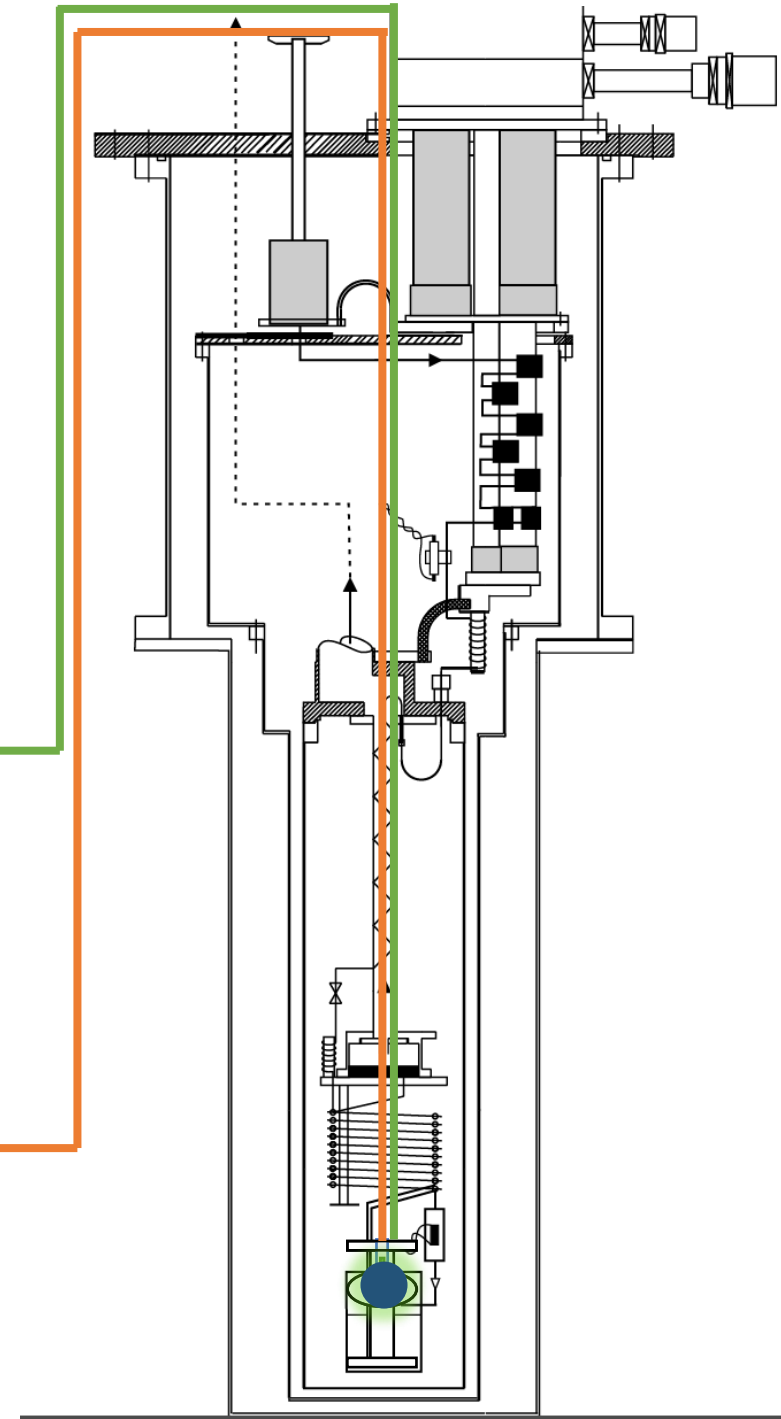
# basics of qubit control & readout



Microwave generator  
(AWG or FPGA / DAC)



Microwave measurement  
(e.g., homodyne measurement)



# Qubit operations and readout

$$\hat{H}_{\text{JC}} = \frac{1}{2}\hbar\omega_q\hat{\sigma}_z + \hbar\omega_r\hat{a}^\dagger\hat{a} + \hbar g(\hat{a}\hat{\sigma}_+ + \hat{a}^\dagger\hat{\sigma}_-)$$

**Perturbation theory in  $g/\Delta \ll 1$**

(Schrieffer-Wolff transf. / adiabatic elimination)

$$\hat{H}_{\text{JC}}^{\text{eff}} = \hbar\omega_r\hat{a}^\dagger\hat{a} + \frac{\hbar\omega'_q}{2}\hat{\sigma}_z + \hbar\chi\hat{a}^\dagger\hat{a}\hat{\sigma}_z$$

Lamb shift:  $\hbar\omega'_q = \hbar\omega_q + \frac{g^2}{\Delta}$

ac Stark shift:  $\chi = \frac{g^2}{\Delta}$

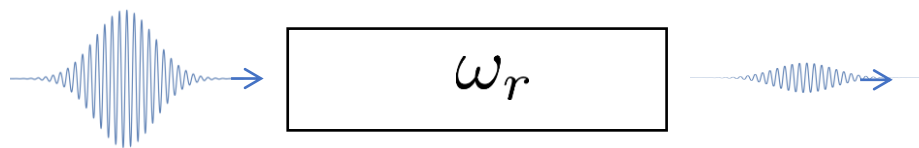
# dispersive regime

$$\hat{H}_{\text{JC}}^{\text{eff}} = \hbar\omega_r \hat{a}^\dagger \hat{a} + \frac{\hbar\omega'_q}{2} \hat{\sigma}_z + \hbar\chi \hat{a}^\dagger \hat{a} \hat{\sigma}_z$$

$$= \hbar(\omega_r + \chi \hat{\sigma}_z) \hat{a}^\dagger \hat{a} + \frac{\hbar\omega'_q}{2} \hat{\sigma}_z$$

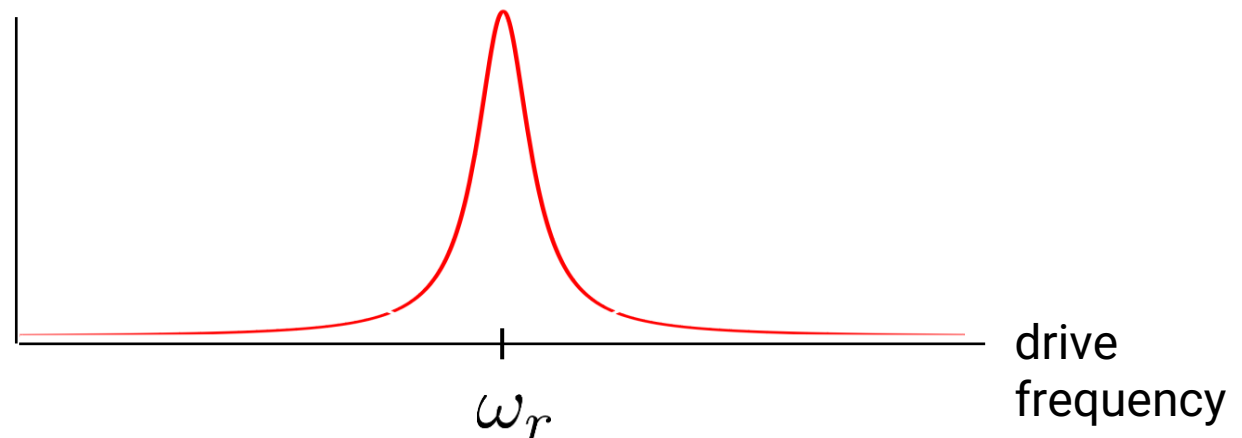
Resonator frequency  
depends  
on qubit state!

# dispersive readout



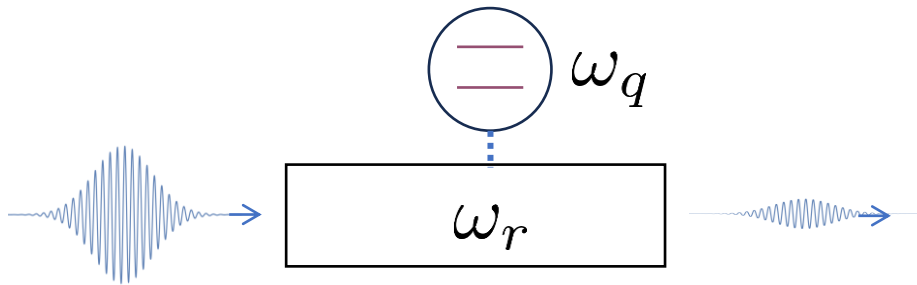
$$\hat{H} = \hbar\omega_r\hat{a}^\dagger\hat{a} + \hat{H}_{\text{drive}}$$

trans-  
mission  
 $S_{21}$

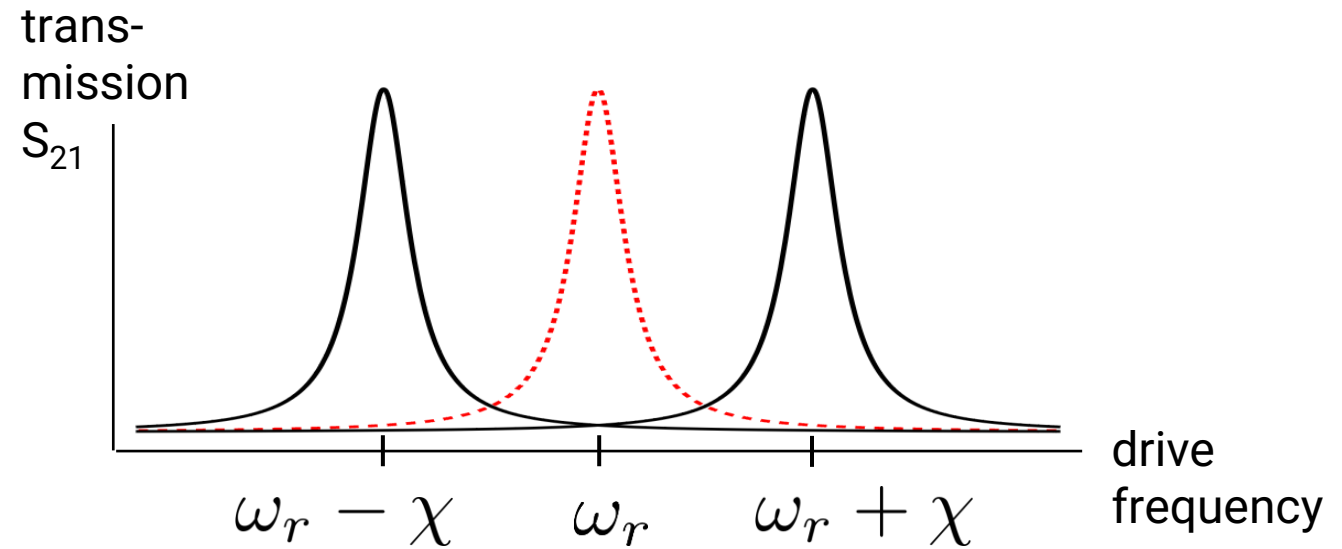




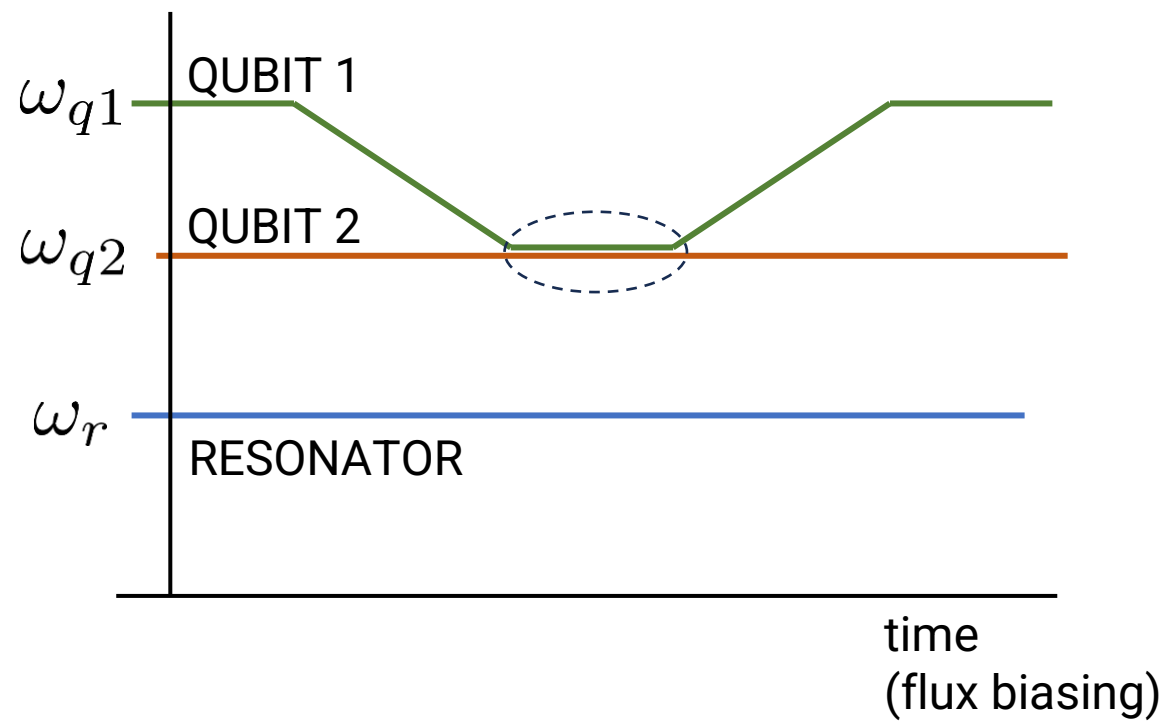
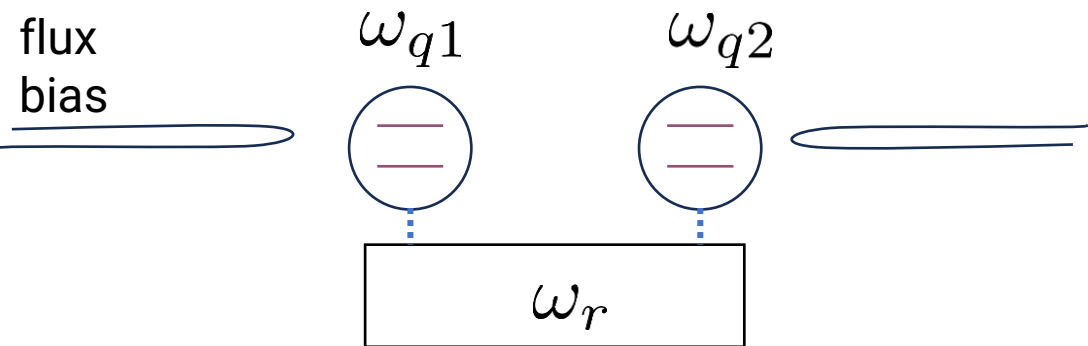
# dispersive readout



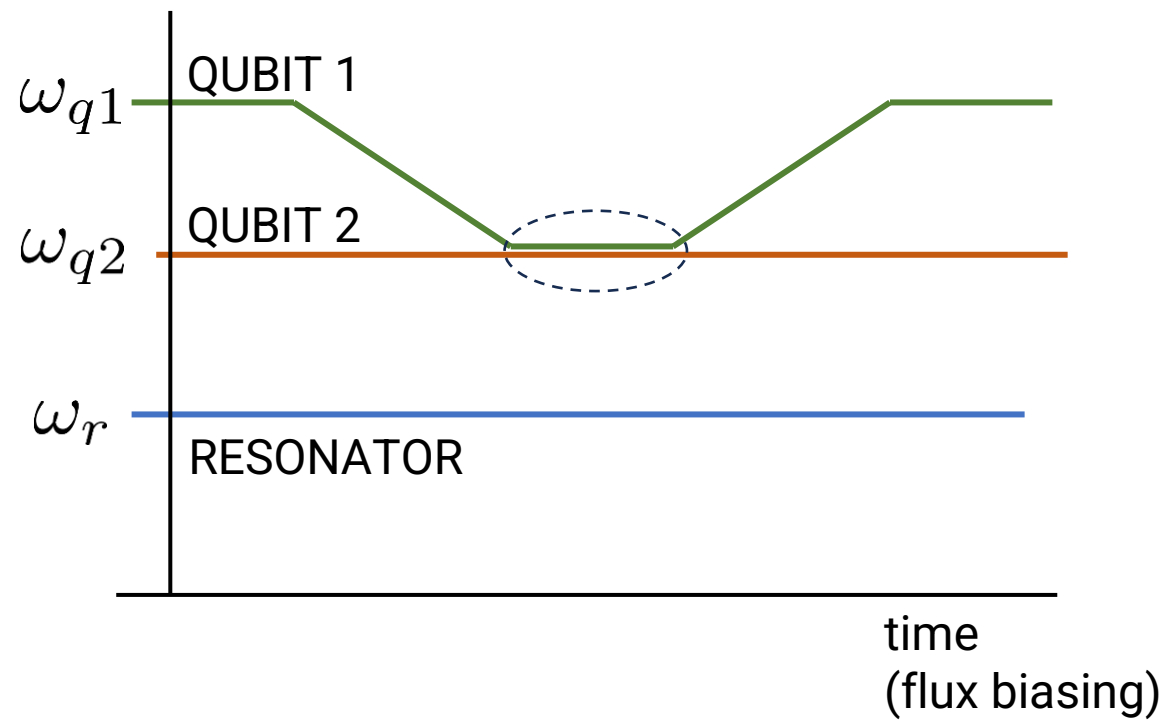
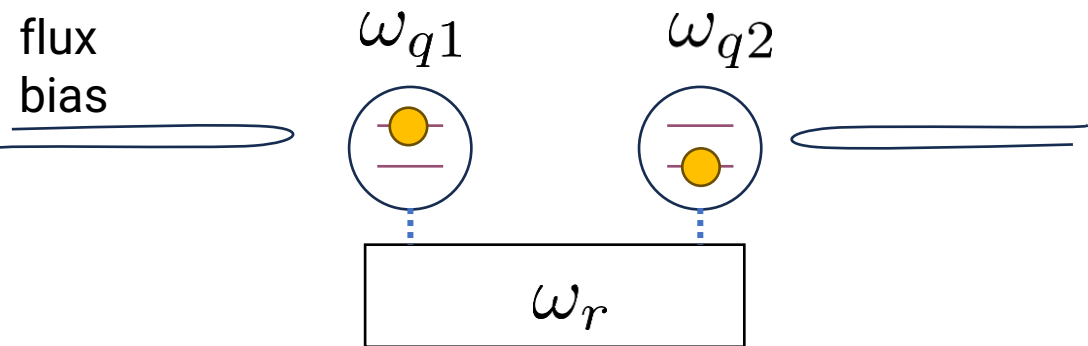
$$\hat{H}_{\text{JC}}^{\text{eff}} = \hbar(\omega_r + \chi\hat{\sigma}_z)\hat{a}^\dagger\hat{a} + \frac{\hbar\omega'_q}{2}\hat{\sigma}_z$$



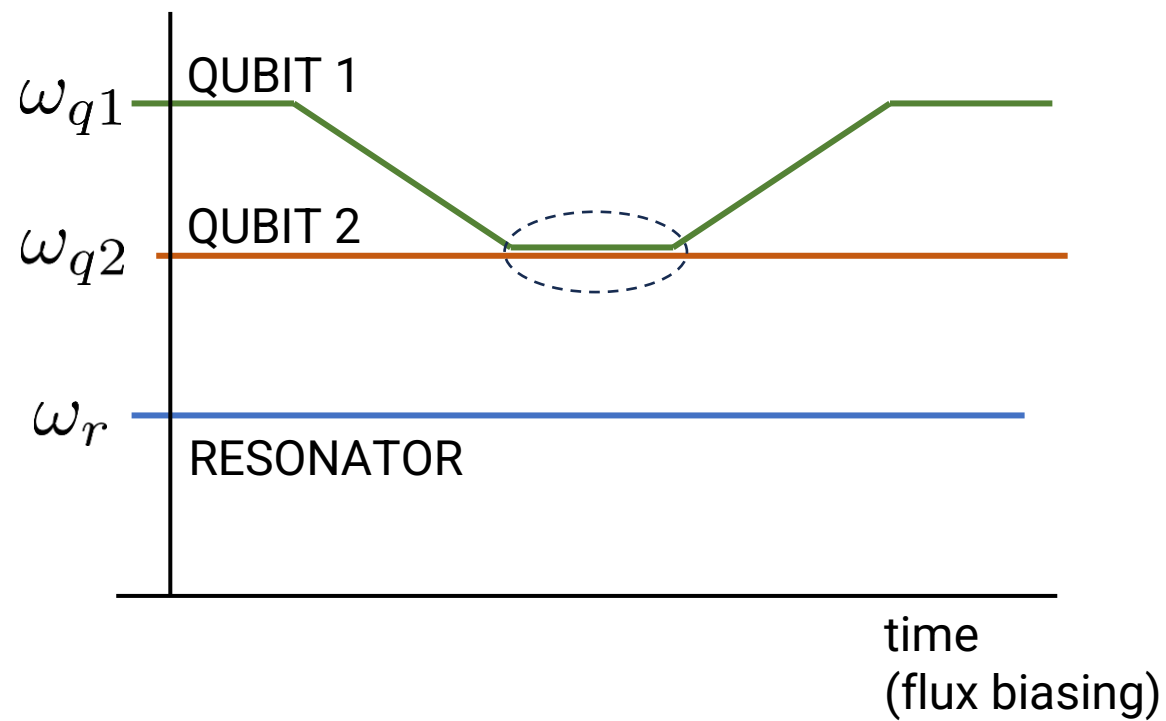
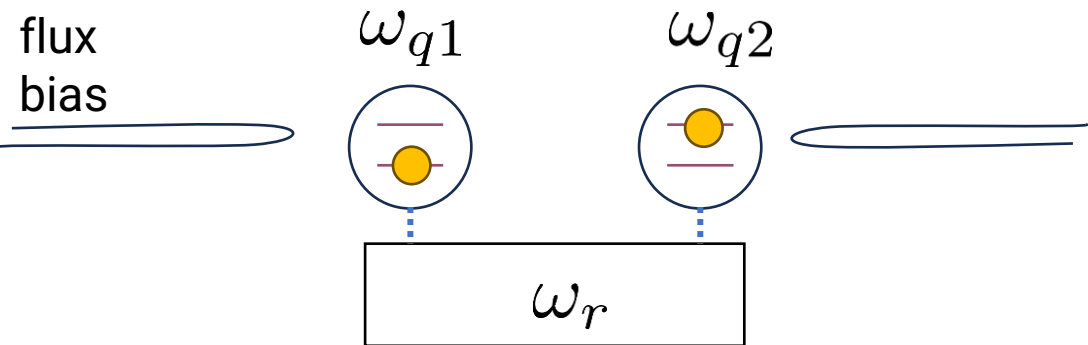
# two-qubit gate (example)



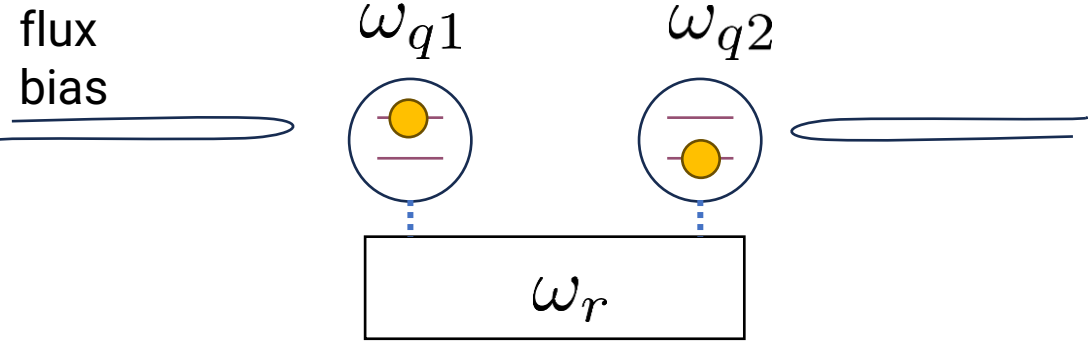
# two-qubit gate (example)



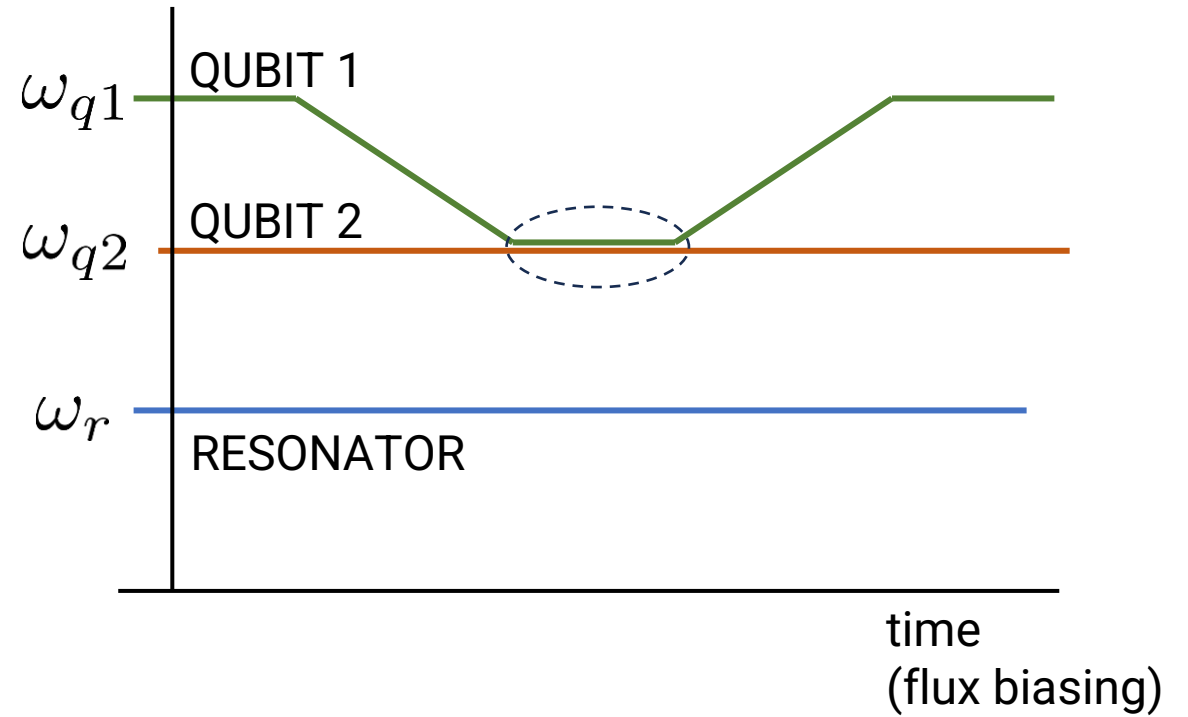
# two-qubit gate (example)



# two-qubit gate (example)



After quarter period:  $\sqrt{i}$ SWAP gate



# task

$$\sqrt{i\text{SWAP}}|01\rangle = ?$$

$$\sqrt{i\text{SWAP}} = \sqrt{\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}} & \frac{i}{\sqrt{2}} & 0 \\ 0 & \frac{i}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$|0\rangle \mapsto \begin{pmatrix} 1 \\ 0 \end{pmatrix}, |1\rangle \mapsto \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\sqrt{i\text{SWAP}} = \sqrt{\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}} & \frac{i}{\sqrt{2}} & 0 \\ 0 & \frac{i}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$|01\rangle = |0\rangle \otimes |1\rangle \quad \mapsto \quad \begin{pmatrix} 1 \\ 0 \end{pmatrix} \otimes \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 & \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\ 0 & \begin{pmatrix} 0 \\ 1 \end{pmatrix} \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix}$$



$$\sqrt{i\text{SWAP}} = \sqrt{\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}} & \frac{i}{\sqrt{2}} & 0 \\ 0 & \frac{i}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$\sqrt{i\text{SWAP}}|01\rangle = ?$$

$$|0\rangle \mapsto \begin{pmatrix} 1 \\ 0 \end{pmatrix}, |1\rangle \mapsto \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\sqrt{i\text{SWAP}} = \sqrt{\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}} & \frac{i}{\sqrt{2}} & 0 \\ 0 & \frac{i}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$\sqrt{i\text{SWAP}}|01\rangle = \text{?}$$

$$|0\rangle \mapsto \begin{pmatrix} 1 \\ 0 \end{pmatrix}, |1\rangle \mapsto \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

# universal control

For gate based quantum computing:  
what is a

**minimal set of elementary qubit operations**

sufficient to carry out an **arbitrary algorithm?**

# universal control

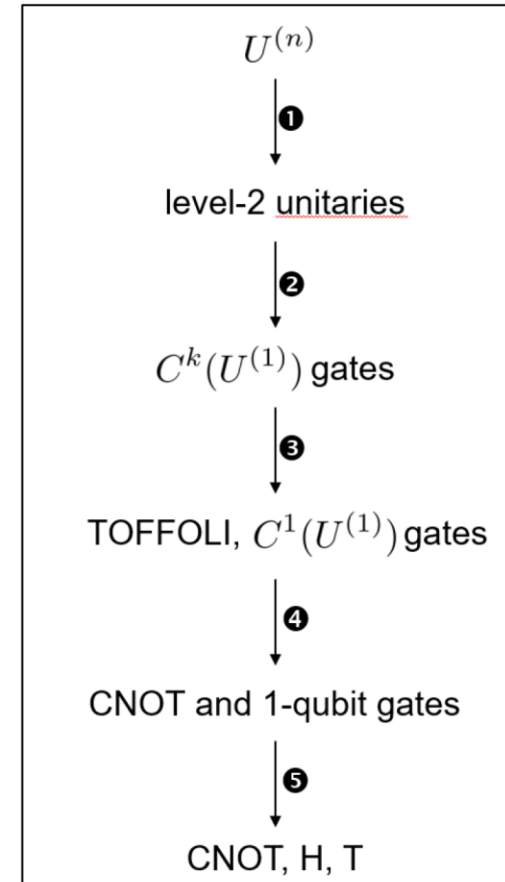
1. arbitrary single-qubit gates  
(can approximate and decompose into sequence of H, T)
2. one 2-qubit entangling gate  
(like CNOT or  $\sqrt{i}$ SWAP)

# universal control

$$U =$$

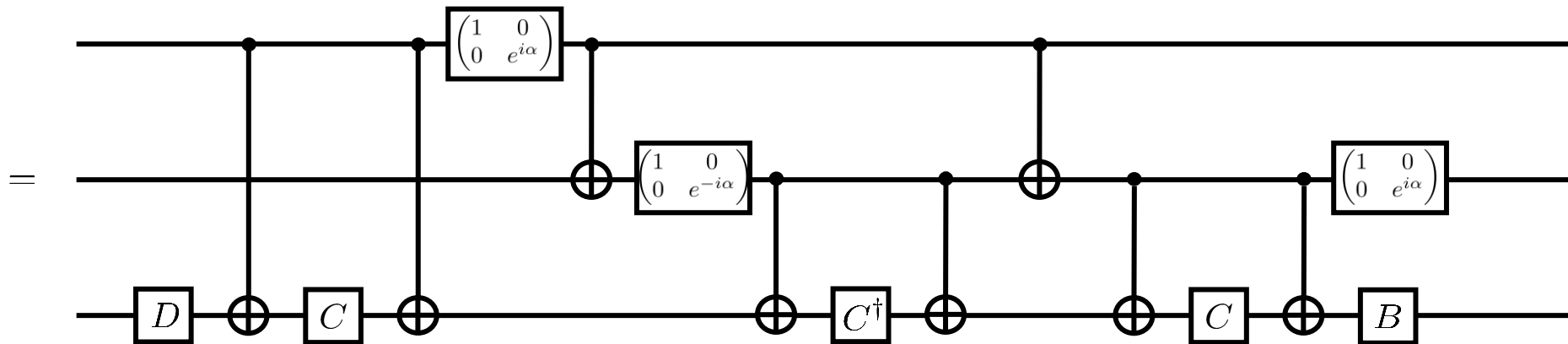
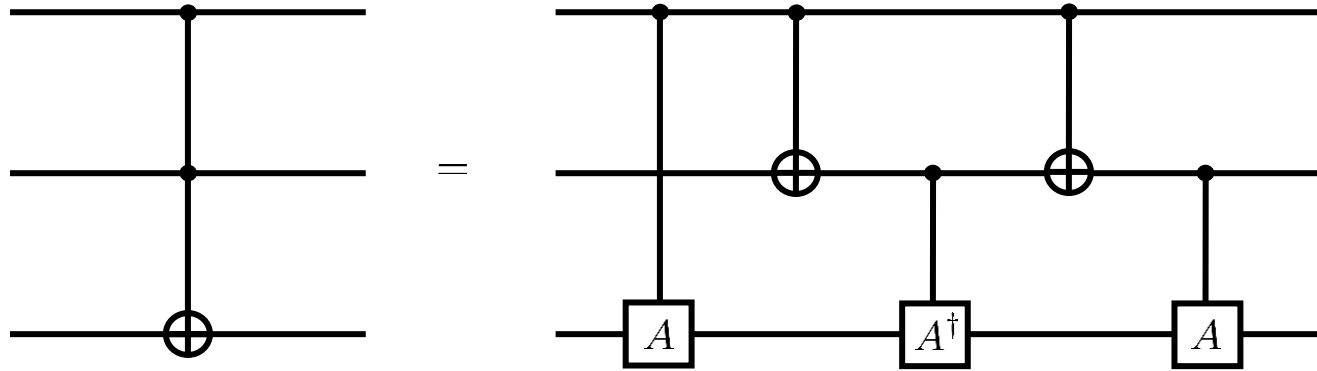
$2^n \times 2^n$  unitary

Summary of decomposition chain:



- ① Iterative reduction of general n-qubit unitary into n-qubit level-2 unitaries
- ② Conversion of n-qubit level-2 unitaries into  $C^k(U^{(1)})$  gates ( $k=n-1$ ), if necessary with the help of Gray codes
- ③ Decompose  $C^k(U^{(1)})$  gate into cascade of TOFFOLIs and one  $C^1(U^{(1)})$  gate, using  $n-1$  auxiliary qubits
- ④ Express TOFFOLI in terms of CA gates where  $A^2=X$ . Then, use AXBYC decomposition on CA and  $C^1(U^{(1)})$  to reduce to CNOTs and 1-qubit gates.
- ⑤ Approximate 1-qubit gates by sequences of H and T gates, using irrational-angle decomposition.

# TOFFOLI gate decomposition

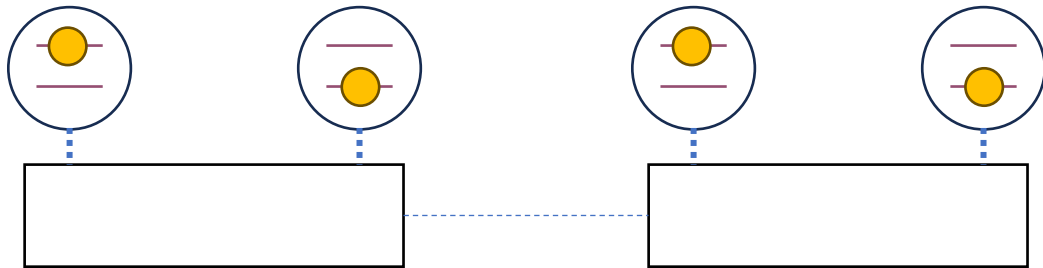


# who's the qubit?

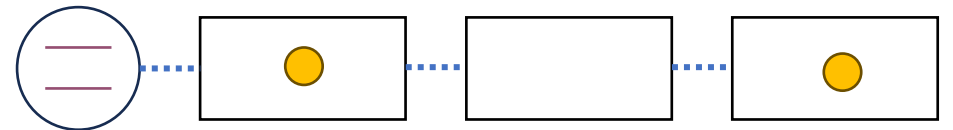
bosonic-qubit architectures

# who's the qubit?

**SCHEME 1:** The qubit is the qubit (duh!)  
(Google, IBM, Rigetti etc.)



**SCHEME 2:** The qubit is a linear cavity mode  
Yale, FNAL-SQMS, Schuster

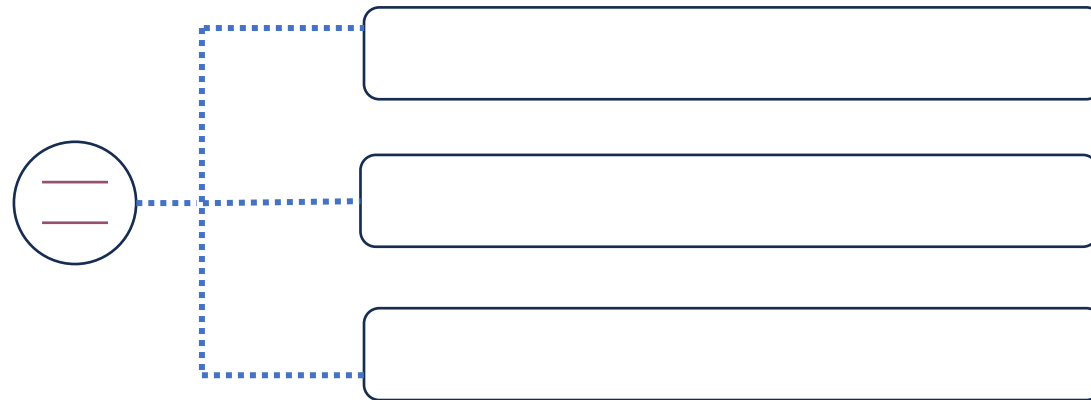


- Fock state encoding
- bosonic codes
- cat states
- GKP states



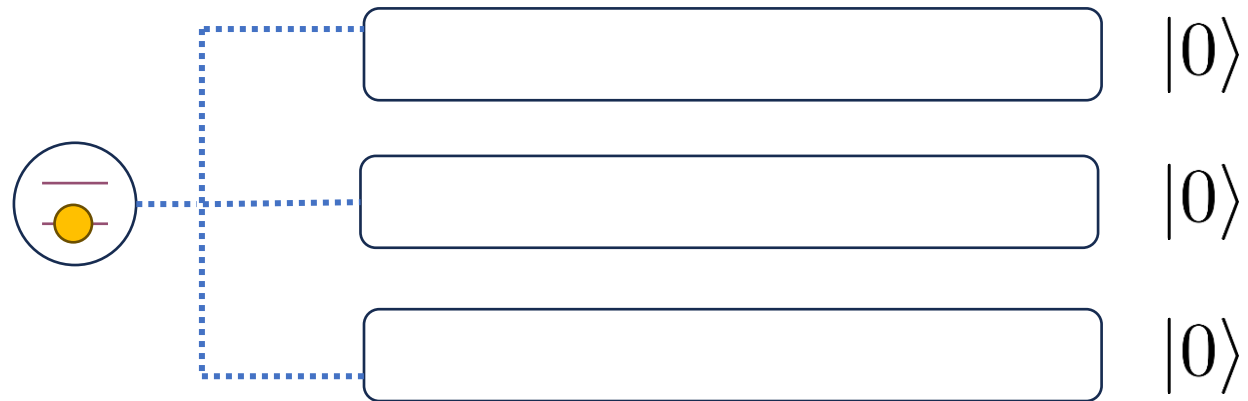
# multi-mode control

$$\hat{H} = \sum_j \omega_j \hat{a}_j^\dagger \hat{a}_j + \sum_j \chi_j \hat{a}_j^\dagger \hat{a}_j \hat{\sigma}_z + \frac{\omega'_q}{2} \hat{\sigma}_z$$



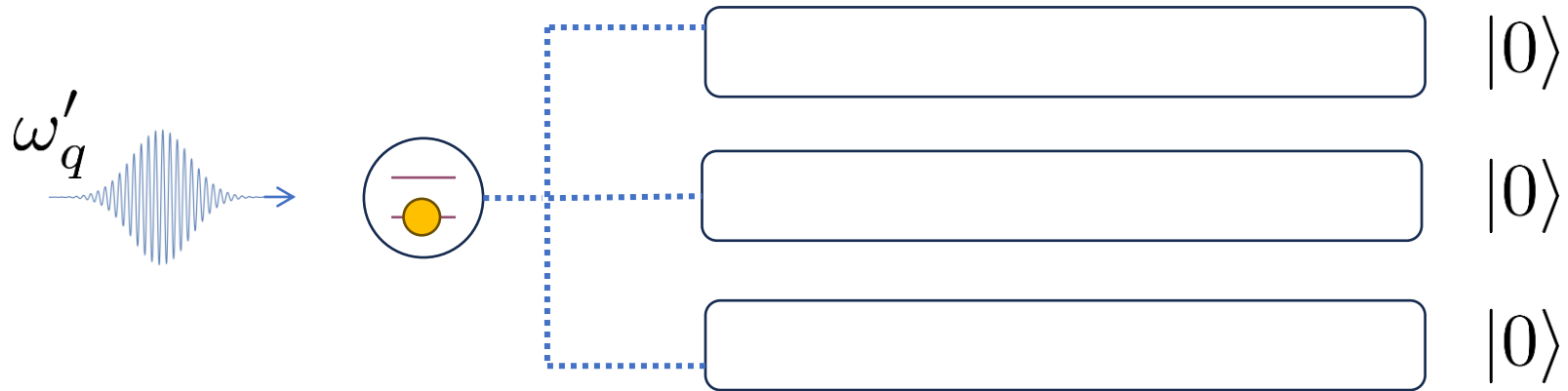
# multi-mode control

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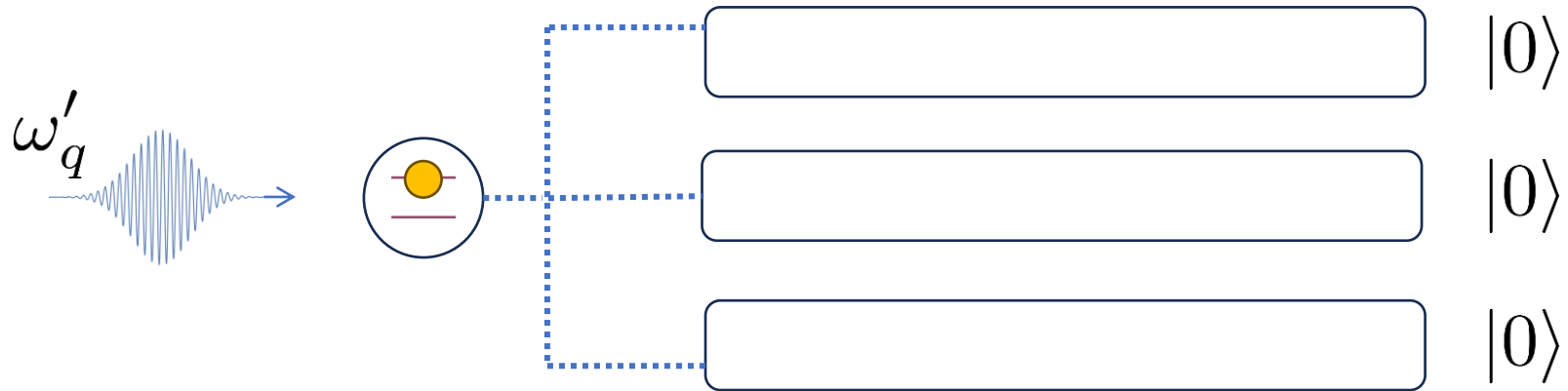
# multi-mode control

$$\hat{H} = \sum_j \omega_j \hat{a}_j^\dagger \hat{a}_j + \sum_j \chi_j \hat{a}_j^\dagger \hat{a}_j \hat{\sigma}_z + \frac{\omega'_q}{2} \hat{\sigma}_z$$



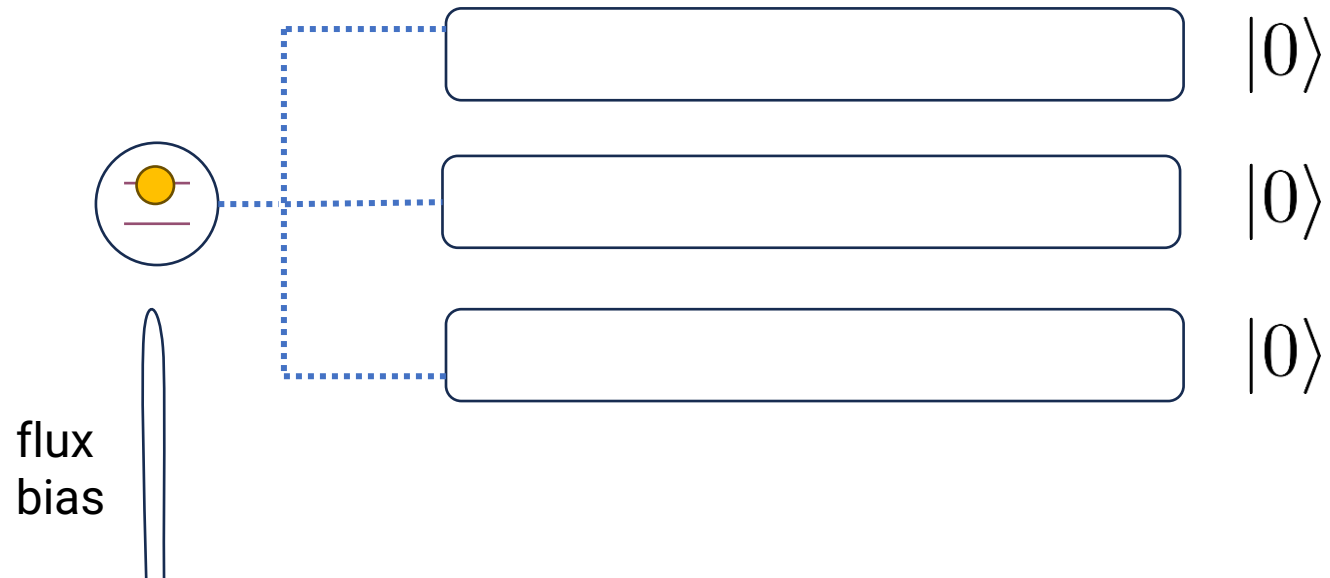
# multi-mode control

$$\hat{H} = \sum_j \omega_j \hat{a}_j^\dagger \hat{a}_j + \sum_j \chi_j \hat{a}_j^\dagger \hat{a}_j \hat{\sigma}_z + \frac{\omega'_q}{2} \hat{\sigma}_z$$



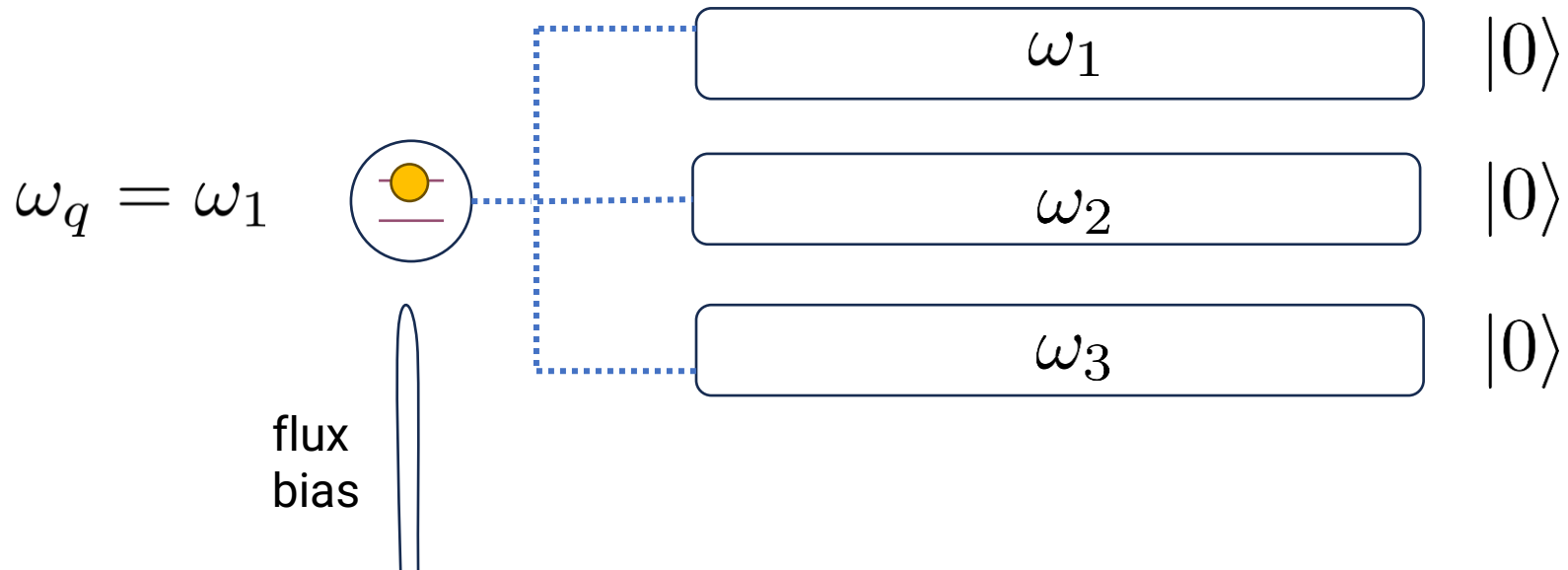
# multi-mode control

$$\hat{H} = \sum_j \omega_j \hat{a}_j^\dagger \hat{a}_j + \sum_j \chi_j \hat{a}_j^\dagger \hat{a}_j \hat{\sigma}_z + \frac{\omega'_q}{2} \hat{\sigma}_z$$



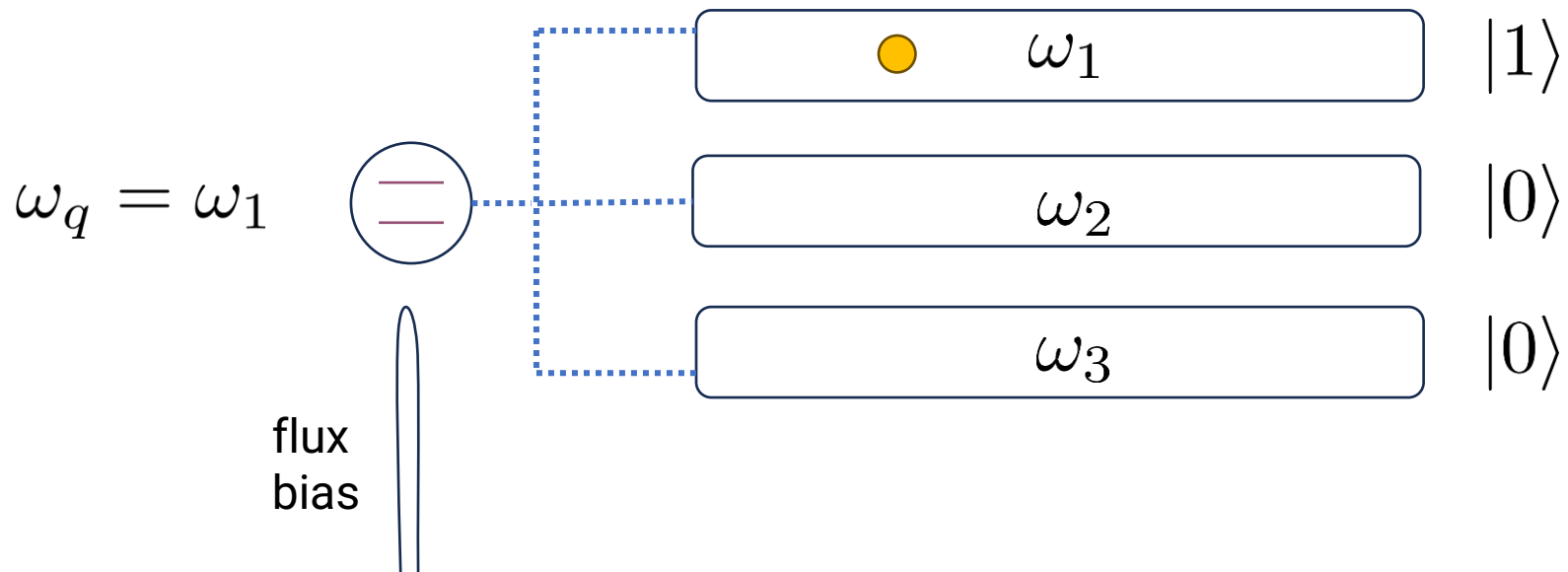
# multi-mode control

$$\hat{H} = \sum_j \omega_j \hat{a}_j^\dagger \hat{a}_j + \sum_j \chi_j \hat{a}_j^\dagger \hat{a}_j \hat{\sigma}_z + \frac{\omega'_q}{2} \hat{\sigma}_z$$



# multi-mode control

$$\hat{H} = \sum_j \omega_j \hat{a}_j^\dagger \hat{a}_j + \sum_j \chi_j \hat{a}_j^\dagger \hat{a}_j \hat{\sigma}_z + \frac{\omega'_q}{2} \hat{\sigma}_z$$

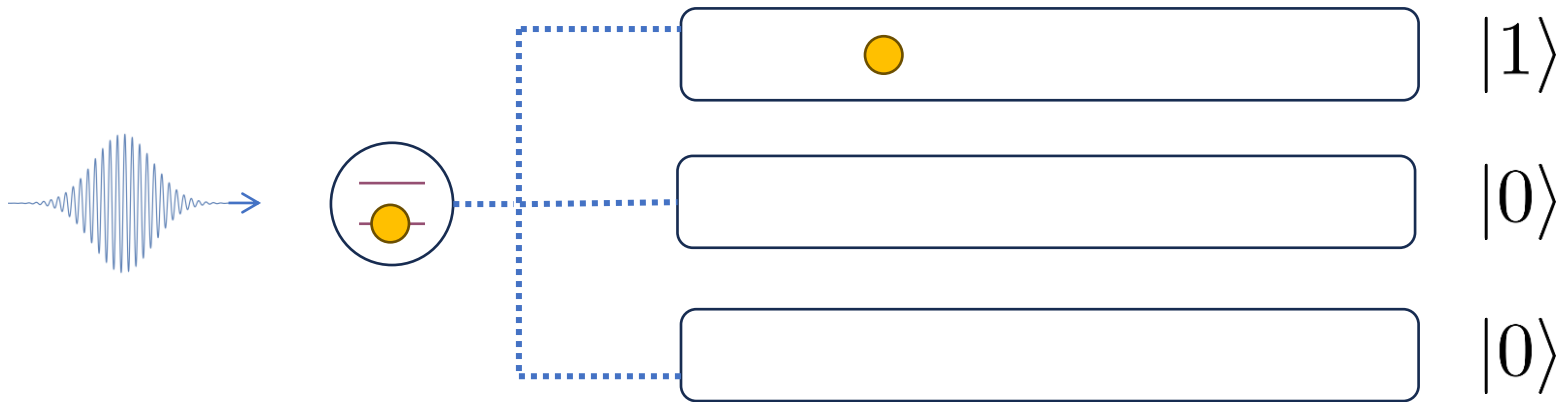


NEXT: load an excitation into mode 2  
( $\pi$  pulse on mode 2)



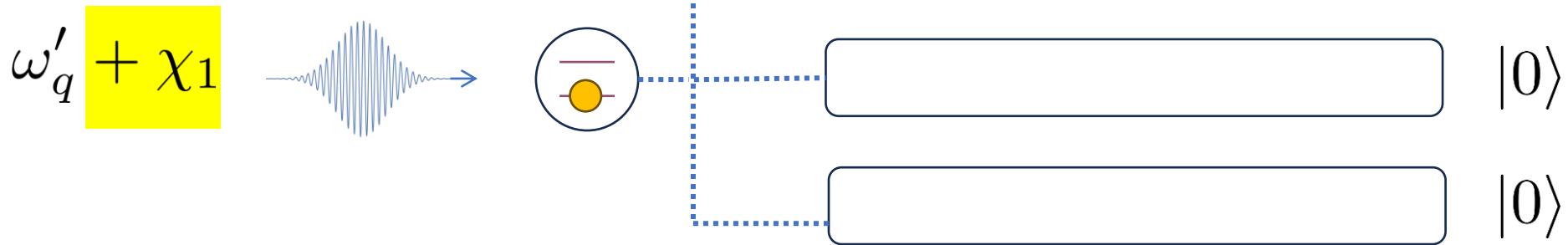
# multi-mode control

$$\hat{H} = \sum_j \omega_j \hat{a}_j^\dagger \hat{a}_j + \sum_j \chi_j \hat{a}_j^\dagger \hat{a}_j \hat{\sigma}_z + \frac{\omega'_q}{2} \hat{\sigma}_z$$



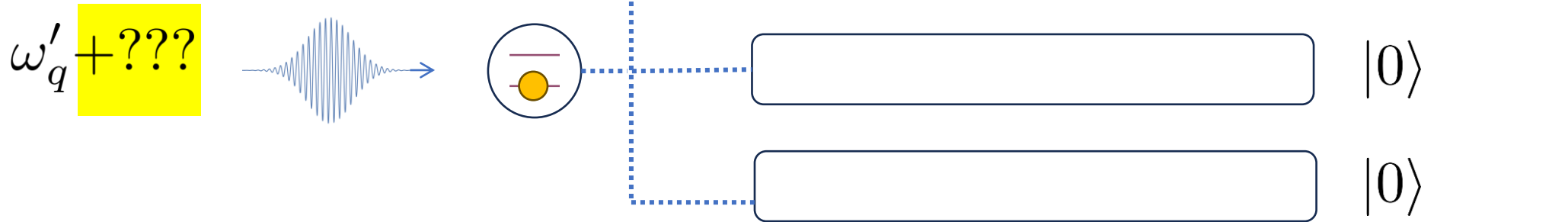
# multi-mode control

$$\hat{H} = \sum_j \omega_j \hat{a}_j^\dagger \hat{a}_j + \sum_j \chi_j \hat{a}_j^\dagger \hat{a}_j \hat{\sigma}_z + \frac{\omega'_q}{2} \hat{\sigma}_z$$

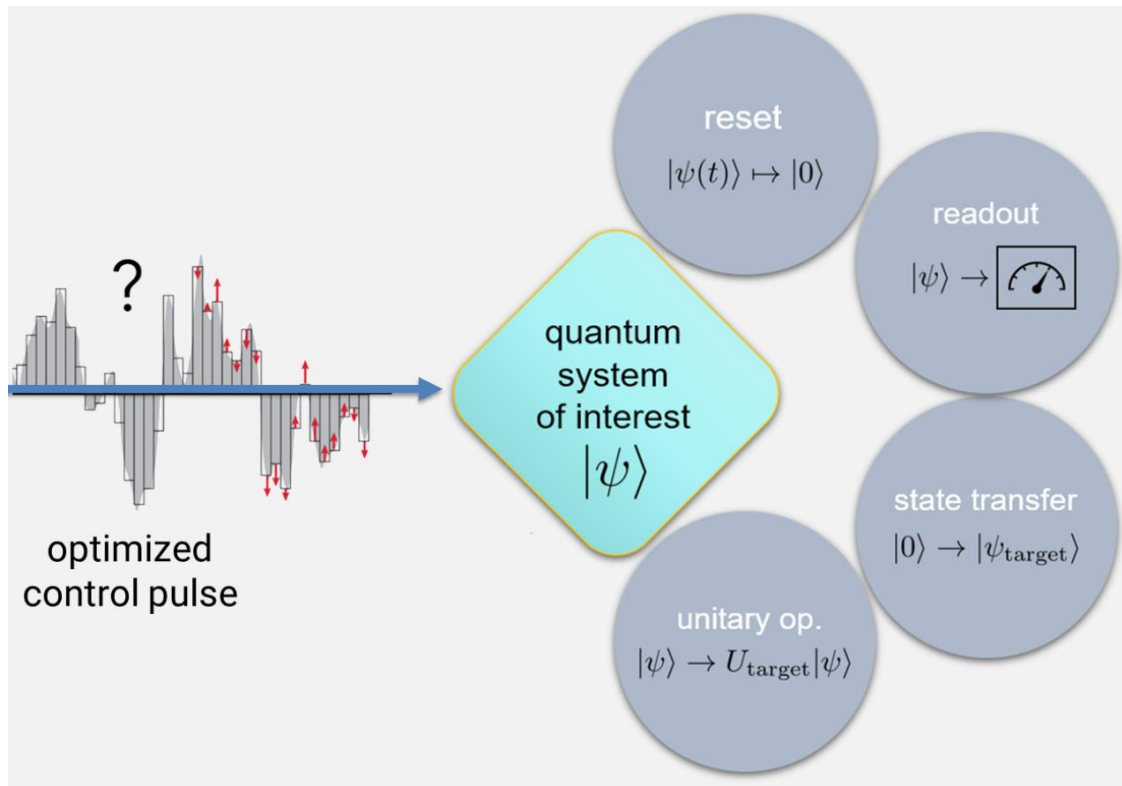


# multi-mode control

$$\hat{H} = \sum_j \omega_j \hat{a}_j^\dagger \hat{a}_j + \sum_j \chi_j \hat{a}_j^\dagger \hat{a}_j \hat{\sigma}_z + \frac{\omega'_q}{2} \hat{\sigma}_z$$



# quantum optimal control



N. Khaneja, ..., S. J. Glaser,  
*Optimal Control of Coupled Spin Dynamics: Design of NMR Pulse Sequences by Gradient Ascent Algorithms*

J. Magn. Reson. **172**, 296-305 (2005)

# robust quantum control

## Crosstalk-Robust Quantum Control in Multimode Bosonic Systems

Xinyuan You,<sup>1,\*</sup> Yunwei Lu,<sup>2</sup> Taeyoon Kim,<sup>1,2,3</sup> Doğa Murat Kürkçüoğlu,<sup>1</sup> Shaojiang Zhu,<sup>1</sup> David van Zanten,<sup>1</sup> Tanay Roy,<sup>1</sup> Yao Lu,<sup>1</sup> Srivatsan Chakram,<sup>4</sup> Anna Grassellino,<sup>1</sup> Alexander Romanenko,<sup>1</sup> Jens Koch,<sup>2,3</sup> and Silvia Zorzetti<sup>1,†</sup>

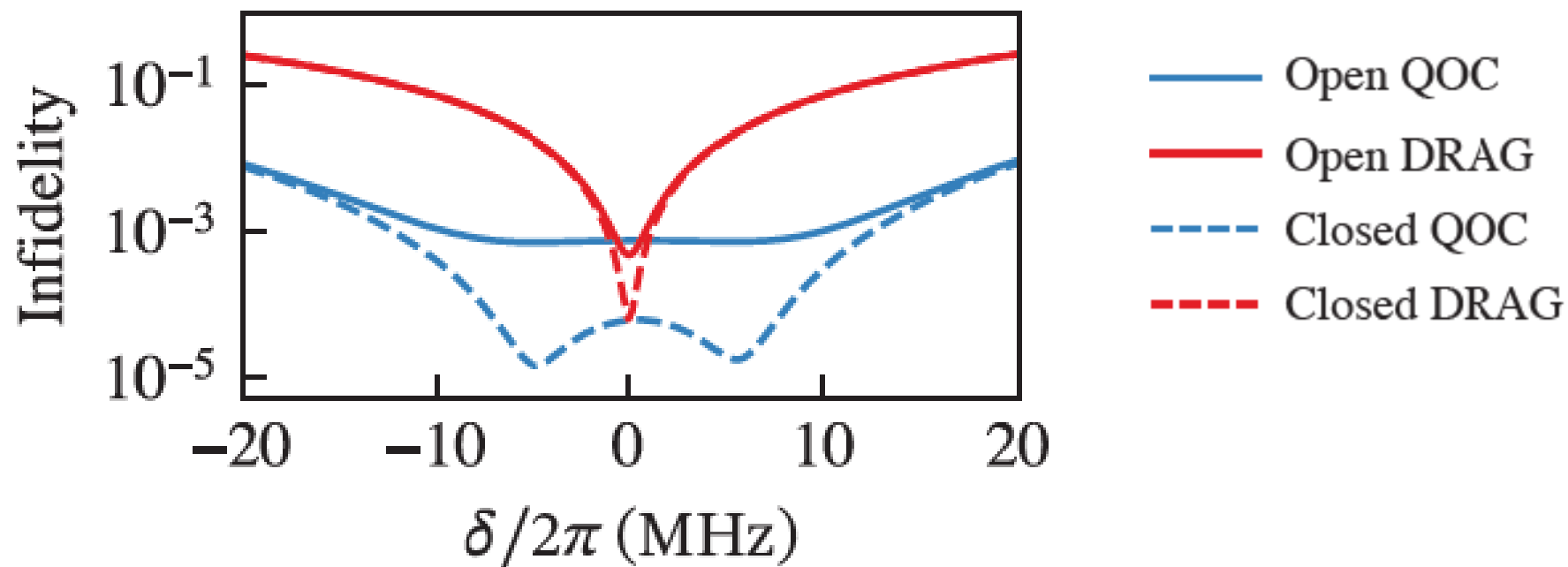
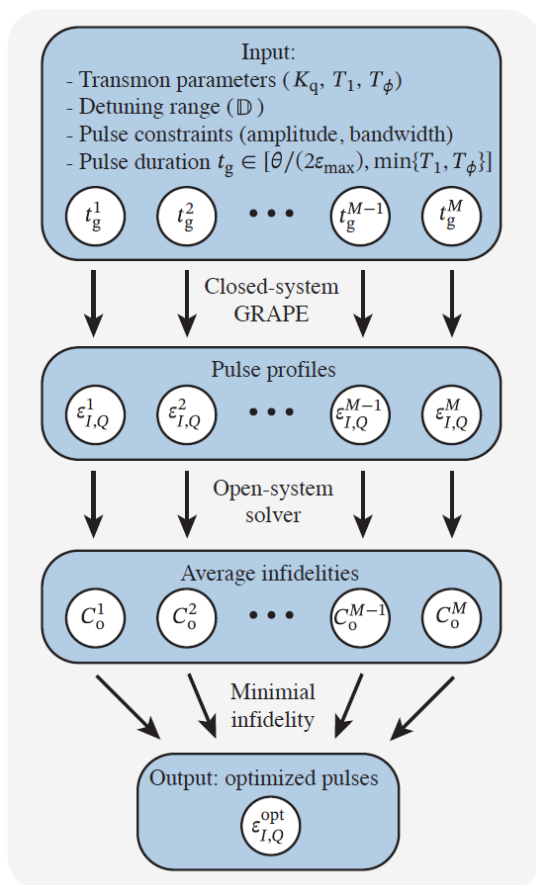
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<sup>2</sup>Department of Physics and Astronomy, Northwestern University, Evanston, IL 60208, USA

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<sup>4</sup>Department of Physics and Astronomy, Rutgers University, Piscataway, NJ 08854, USA



# programmable computer

a black box\* that can implement\*\*  
any\*\*\* given  $f$

$$\text{NAND} : \mathbb{B}^2 \rightarrow \mathbb{B}^1$$

$$\text{random} : \Omega \rightarrow \mathbb{B}^1$$

$$\text{bubble sort} : \mathbb{B}^{64 \cdot 100} \rightarrow \mathbb{B}^{64 \cdot 100}$$

$$\text{PFD} : \mathbb{B}^{1024} \rightarrow \mathbb{B}^{1024 \cdot 2}$$

prime factor decomposition

...

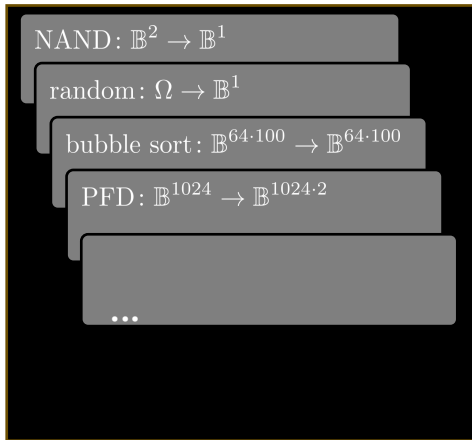
\* a “device” obeying the laws of physics

\*\* with input & output sizes  $n, m \leq N_0$

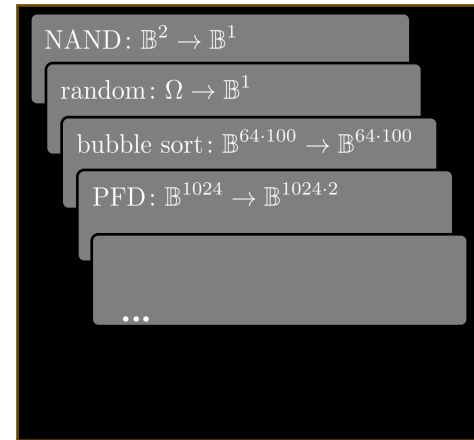
\*\*\* output is provided in finite time

# assessing performance

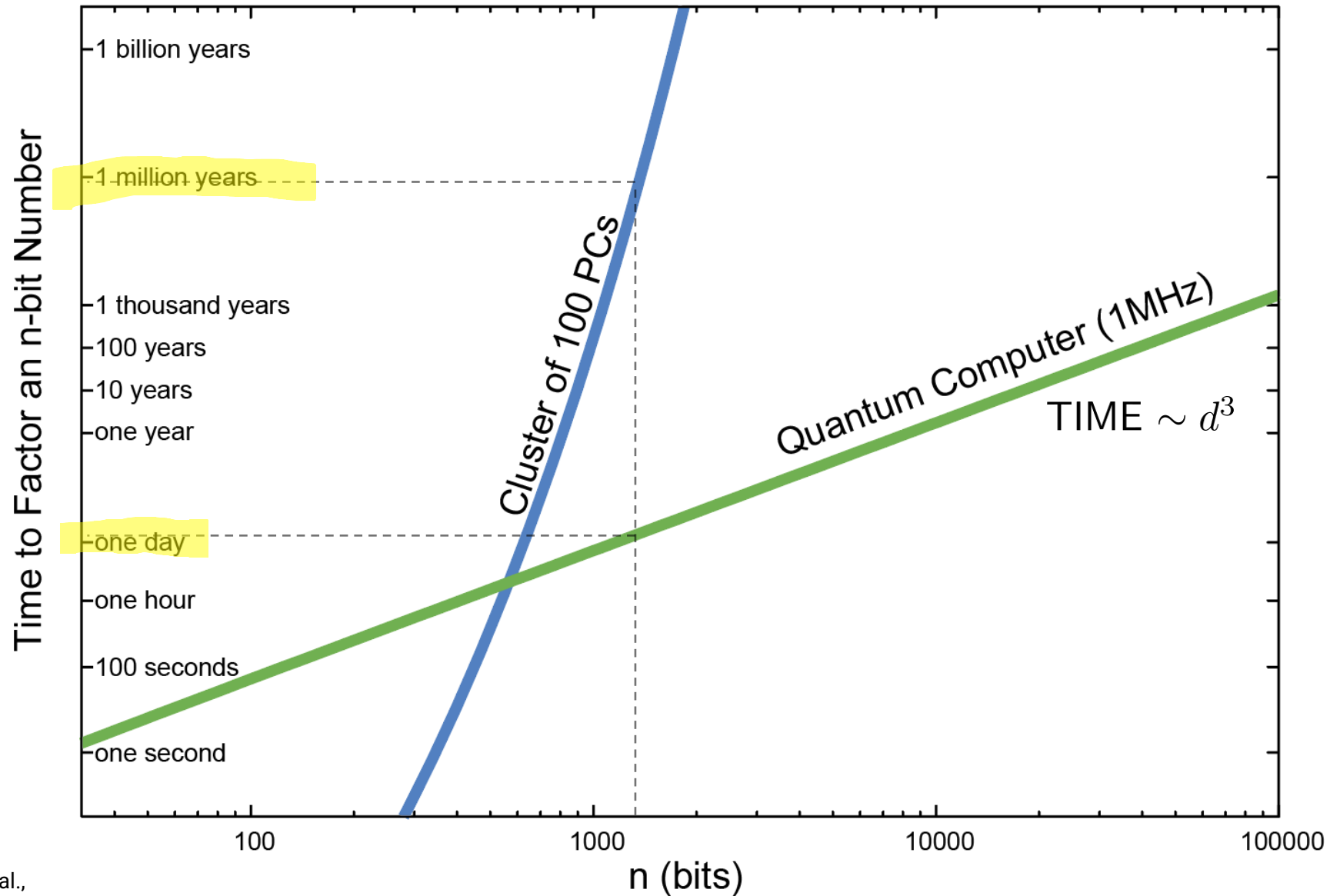
COMPUTER 1



COMPUTER 2



# Quantum computation: Shor's algorithm



Peter Shor

based on:  
van Meter et al.,  
quant-ph/0507023 (2005)