

The evolution of electroweak theory

Electroweak milestones –

50 years of neutral currents, 40 years of W and Z bosons

CERN

31 October, 2023



W. HOLLIK



MAX-PLANCK-GESELLSCHAFT

MAX-PLANCK-INSTITUT FÜR PHYSIK, MÜNCHEN

90 years anniversary

theoretical description of the
weak interaction began 1933

ansatz by **Enrico Fermi**

$$\mathcal{H} = G J_\mu \cdot J^\mu$$

- ★ current–current interaction
- ★ charged current
- ★ universal coupling constant G
today: “**Fermi constant**” G_F



From a phenomenological model to a fully-fledged Quantum Field Theory

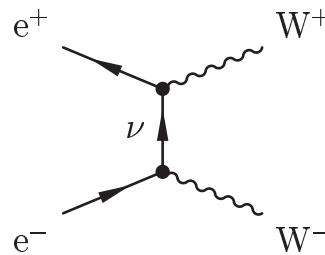
- 1933 Fermi model
- 1957 incorporation of parity violation
- 1967 – 1973 formulation of the electroweak Standard Model
gauge theory based on $SU(2) \times U(1)$
 W, Z masses via Higgs mechanism
fermion masses and mixing via Yukawa couplings
- 1971 – 1972 proof of renormalizability
- 1973 – experimental confirmation

Phases of evolution

- from Fermi's ansatz to a full theory
- from precision calculations to discoveries
- current performance and perspectives

weak interaction around 1960

- current–current interaction $\sim J_\mu \cdot J^\mu$
- charged current $J_\mu^\pm = V_\mu - A_\mu$ V–A structure
Feynman, Gell-Mann 1957, Sudarshan, Marshak 1957
- intermediate heavy charged vectorbosons W^\pm *Schwinger 1957*
couple to J_μ^+ and J_μ^-
analogy to QED, but obviously incomplete:
→ not renormalizable
→ bad high-energy behaviour violating unitarity



- algebra of charges $[I^+, I^-] = 2I^0$
indicates $SU(2)$ with $[I^0, I^\pm] = \pm I^\pm$ **isospin**
missing current $J_\mu^0 \rightarrow$ another boson $W^0 =$ photon ?

symmetry takes over

- need for additional quantum number: hypercharge Y
additional current J^Y and boson $B^0 \Rightarrow SU(2) \times U(1)$

$$\begin{pmatrix} A^0 \\ Z^0 \end{pmatrix} = \begin{pmatrix} \cos \theta_W & \sin \theta_W \\ -\sin \theta_W & \cos \theta_W \end{pmatrix} \begin{pmatrix} B^0 \\ W^0 \end{pmatrix} \quad \text{Glashow 1961}$$

photon $A^0 \leftrightarrow$ electromagnetic current, $Q = \frac{Y}{2} + I_3$

Z^0 boson \leftrightarrow weak neutral current

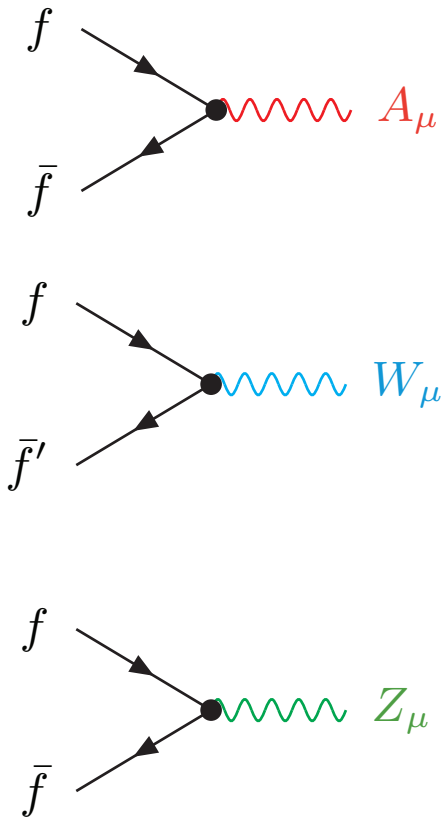
- dynamics of vectorbosons: non-Abelian gauge theory
formulated for $SU(2)$ Yang, Mills 1954

but: mass = 0 for vectorbosons!

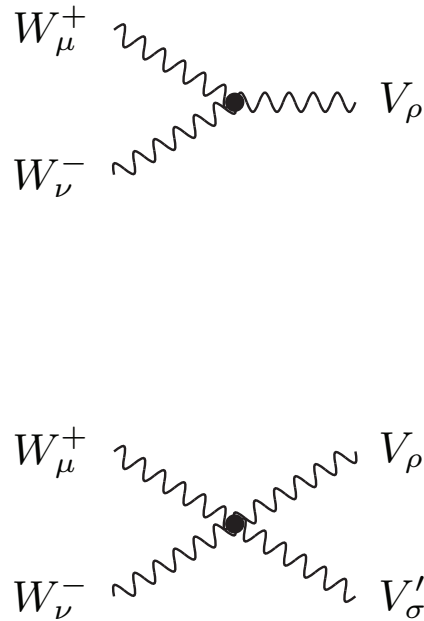
- constructed for $SU(2) \times U(1)$ Glashow 1961
 - photon massless, W^\pm and Z^0 massive
 - explicit mass terms added by hand
 - break gauge symmetry \rightarrow not renormalizable

electroweak interactions from symmetry

fermion–vectorbosons



vectorboson self interactions



coupling constants:

$$e, \quad g = e / \sin \theta_w$$

group entries:

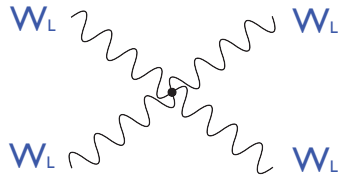
$$\text{isospin } I_3^f, \quad \text{charge } Q_f$$

basic problem: masses of W and Z

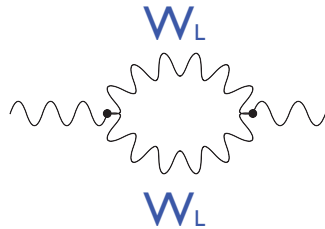
- W, Z have longitudinal polarization states

polarization vectors of W (Z) $\epsilon_L \sim p/M_W$

for large momentum p



bad high energy behaviour of WW scattering



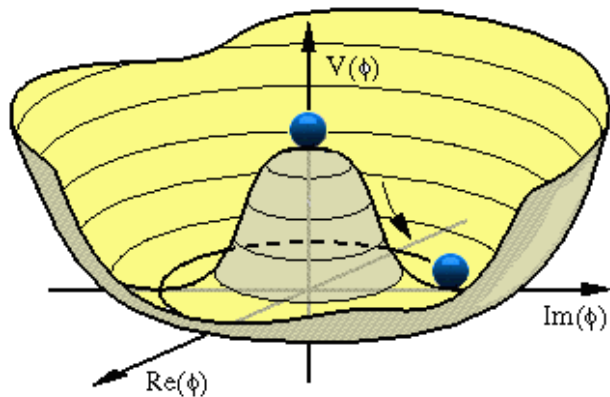
bad divergence of loop integrals

spontaneous symmetry breaking

Brout, Englert 1964, Higgs 1964, Guralnik, Hagen, Kibble 1964

scalar field Φ with self-interaction

$$V = -\mu^2 (\Phi^\dagger \Phi) + \lambda (\Phi^\dagger \Phi)^2, \quad \lambda > 0$$



spontaneous symmetry breaking:

minimum at $v = \frac{\mu}{\sqrt{\lambda}} \neq 0$

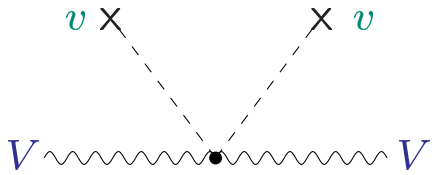
interaction with gauge field:

gauge-invariant, renormalizable

gauge transformation $\longrightarrow \Phi = v + H$ physical field H

$$\Phi = v + H$$

Higgs – gauge boson interaction $\sim g^2$

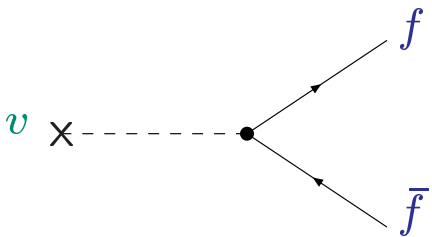


$$V = W, Z$$

$$\Rightarrow \mathbf{V \text{ masses}} \quad M^2 = g^2 v^2$$

residual VV–H interaction $\sim M_V$

Yukawa interaction $\sim g_f$



$$\mathbf{fermion \text{ masses}} \quad m_f = g_f v$$

residual f – H interaction

$$g_f \sim m_f$$

formulation of the electroweak theory

- applying current ideas to $SU(2) \times U(1)$ gauge theory
Weinberg 1967, Salam 1968 (formulated for leptons)
- extension to hadrons
assuming two quark generations: *c-quark postulated*
Glashow, Iliopoulos, Maiani 1970
- assuming three quark generations: *t, b-quarks postulated*
Kobayashi, Maskawa 1973
embedding of CP -violation (*discovered 1964*)

Result: unified electroweak theory at the classical level

Nobelprize in Physics 1979

Glashow, Salam, Weinberg



● around 1970: **classical theory of ew interactions**
→ weak neutral current, massive W^\pm, Z^0 bosons

● still open question: **is this kind of theory renormalizable?**
– upgrade to the quantum level?
– fundamental theory?

● breakthrough: **proof of renormalizability**
't Hooft 1971, 't Hooft, Veltman 1972

Nobelprize in Physics 1999



pathbreaking improvement

- ★ electroweak theory promoted to a consistent QFT
- ★ predictions beyond lowest order possible
- ★ quantum effects calculable like in QED (*Lamb shift, g-2*)

- need: technology for 1-loop Feynman integrals
 - 't Hooft, Veltman 1978* *scalar 1-, ... 4-point integrals*

- complete 1-loop calculation for $e^+e^- \rightarrow \mu^+\mu^-$
 - Passarino, Veltman 1979*
 - reduction method for tensor integrals to scalar integrals*
 - taylorized to computer-aided calculations*
 - still being used for precision calculations (LHC, FCC. . .)*

opened the era of electroweak precision physics

milestone: 1-loop calculation of the ρ -parameter

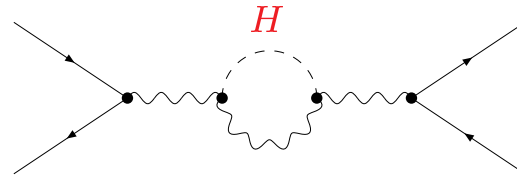
Veltman 1977

for a fermion doublet (t,b) with $m_t \gg m_b$

$$\rho = \frac{M_W^2}{M_Z^2 \cos^2 \theta_W} = \frac{G_{\text{NC}}}{G_{\text{CC}}} = 1 + \frac{3 G_F m_t^2}{8\pi^2 \sqrt{2}} \simeq 1 + \mathbf{0.01}$$

information on heavy (unkown) particles

quantum corrections are sensitive also to the **Higgs boson**



quantum effects are detectable in precision experiments

- masses are correlated with other measurable quantities

$$M_W^2 = \frac{\pi\alpha}{\sqrt{2}G_F \sin^2 \theta_W}, \quad M_Z^2 = \frac{M_W^2}{\cos^2 \theta_W}$$

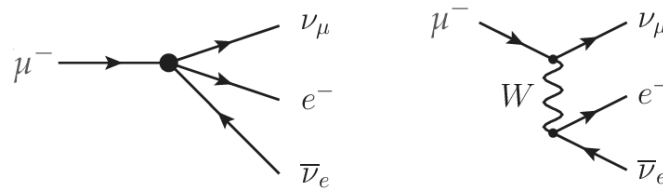
- M_W, M_Z can be obtained from $G_F, \alpha, \sin^2 \theta_W$
since 1973: $\sin^2 \theta_W$ known from neutrino scattering
first calculations at one-loop order done in 1980

Veltman; Antonelli, Consoli, Corbo

limited precision, $\Delta \sin^2 \theta_W \sim 0.0016$

- M_W and M_Z are correlated via G_F and α
allows to calculate M_W when M_Z is known
since 1983 *UA1, UA2*
since 1989 *LEP and SLC experiments*
 \Rightarrow calculate M_W from M_Z, G_F, α (and more)

Fermi constant and $W-Z$ mass correlation



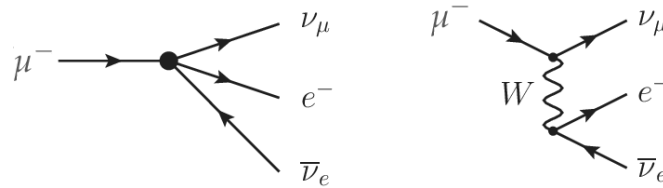
$$\frac{G_F}{\sqrt{2}} = \frac{\pi\alpha}{2M_W^2 \sin^2 \theta_W}$$

● “on-shell scheme” *Sirlin 1980; Marciano, Sirlin 1980*

M_W, M_Z pole masses

$$\sin^2 \theta_W = 1 - \frac{M_W^2}{M_Z^2}, \quad \frac{G_F}{\sqrt{2}} = \frac{\pi\alpha}{2M_W^2 \sin^2 \theta_W} (1 + \Delta r)$$

Fermi constant and $W-Z$ mass correlation



$$\frac{G_F}{\sqrt{2}} = \frac{\pi\alpha}{2M_W^2 \sin^2 \theta_W}$$

- “on-shell scheme” *Sirlin 1980; Marciano, Sirlin 1980*

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- \overline{MS} scheme

Degrassi, Sirlin, Fanchiotti 1990; Degrassi, Gambino, Giardino 2014

$\alpha, \sin^2 \theta_W$ running quantities at scale $Q=M_Z$

$$\frac{G_F}{\sqrt{2}} = \frac{\pi\hat{\alpha}}{2M_W^2 \sin^2 \hat{\theta}_W} (1 + \Delta\hat{r})$$

on-shell calculations for W mass and e^+e^- processes

Fleischer, Jegerlehner 1981

Akhundov, Bardin, Riemann 1986

Böhm, WH, Spiesberger 1986

WH 1988

Beenakker, WH 1988

Consoli, WH, Jegerlehner 1989

Bardin, Riemann et al. since 1985 → ZFITTER

...

⇒ **Z Physics at LEP 1 (1989)**

- *Z line shape*
- *forward–backward asymmetries*
- *MC generators*

...

CERN 89-08
Volume 1
21 September 1989

ORGANISATION EUROPÉENNE POUR LA RECHERCHE NUCLÉAIRE
CERN EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

Z PHYSICS AT LEP 1

Edited by
Guido Altarelli, Ronald Kleiss and Claudio Verzegnassi

Volume 1: STANDARD PHYSICS

Co-ordinated and supervised by G. Altarelli

GENEVA
1989

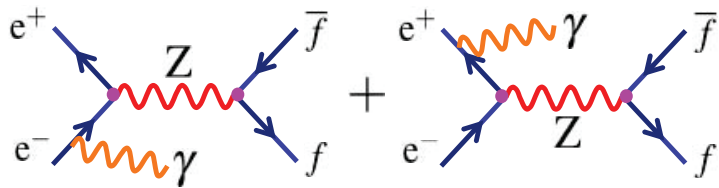
Yellow Book coordinated
by Guido Altarelli

to make sure that
every needed aspect
was prepared for
the LEP experiments

Z lineshape measurements \Rightarrow mass M_Z and width Γ_Z

QED corrections

initial state radiation



$$\sigma(s) = \int_0^1 dz H(z) \sigma(zs)$$

provided by theory

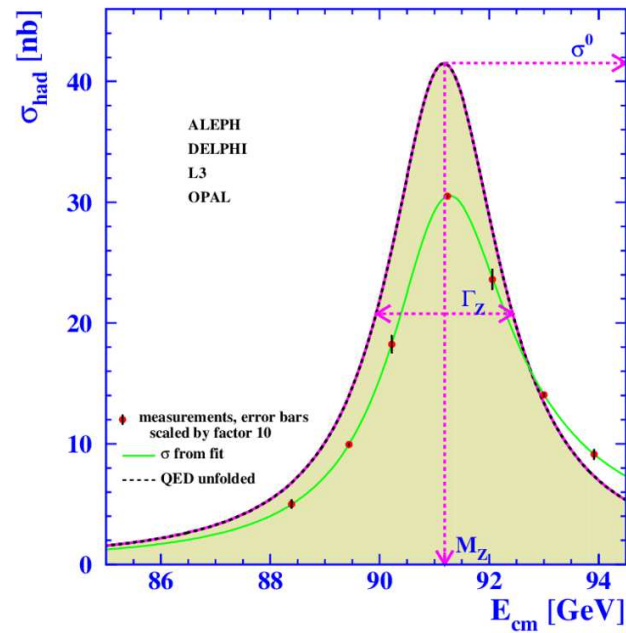
Berends, Kleiss, ...

Jadach, Ward, ... KORALZ

Bardin, Rieman, ... ZFITTER

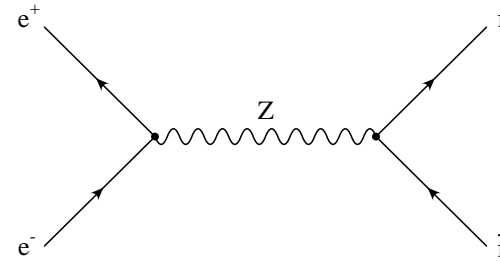
codes including multiple radiation

distort the line shape



Z resonance observables

$$e^+ e^- \rightarrow f \bar{f}$$



- effective Z boson couplings with higher-order $\Delta g_{V,A}$

$$g_V^f \rightarrow g_V^f + \Delta g_V^f, \quad g_A^f \rightarrow g_A^f + \Delta g_A^f$$

- effective ew mixing angle (for $f = e$):

$$\sin^2 \theta_{\text{eff}} = \frac{1}{4} \left(1 - \text{Re} \frac{g_V^e}{g_A^e} \right) = 1 - \frac{M_W^2}{M_Z^2} + \frac{M_W^2}{M_Z^2} \Delta \rho + \dots$$

- codes for precision calculations

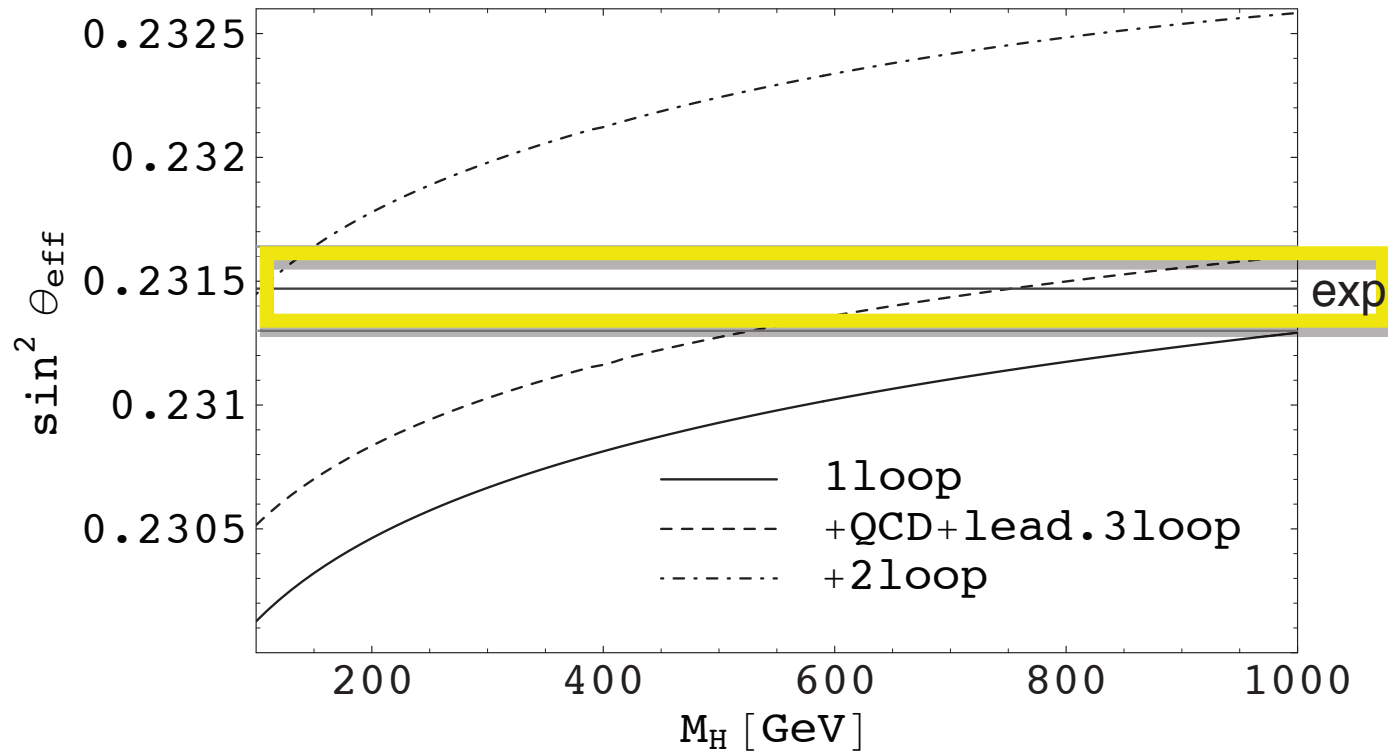
ZFITTER *Bardin, Riemann, et al. (1989 ff)*

TOPAZ0 *Montagna, Piccinini, Nicrosini, Passarino, Pittau (1993 ff)*

BHM *Burgers, WH, Martinez (1989 ff)*

EXPOSTAR *Kennedy, Lynn (1989)*

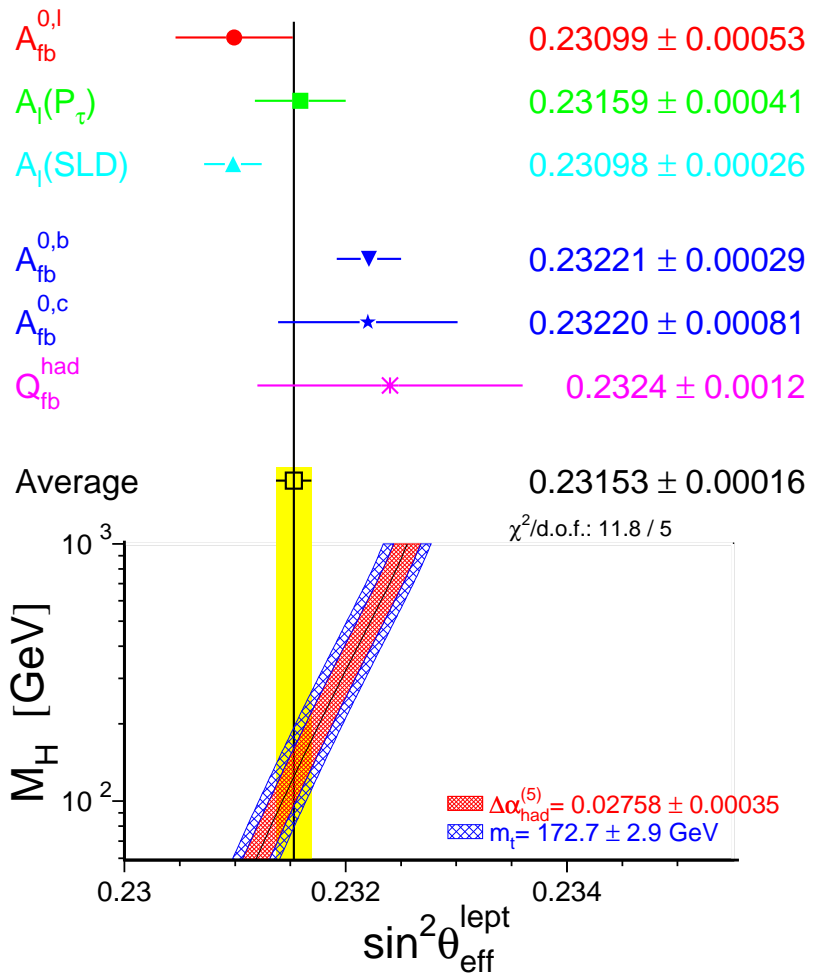
impact of higher-order contributions



lowest order: $\sin^2 \theta_W = 1 - \frac{M_W^2}{M_Z^2} = 0.22305 \pm 0.00023$ (PDG 2022)

exp. value: $\sin^2 \theta_{\text{eff}} = 0.23153 \pm 0.00016$

quantum effects are established with compelling significance



LEP ELECTROWEAK WORKING GROUP

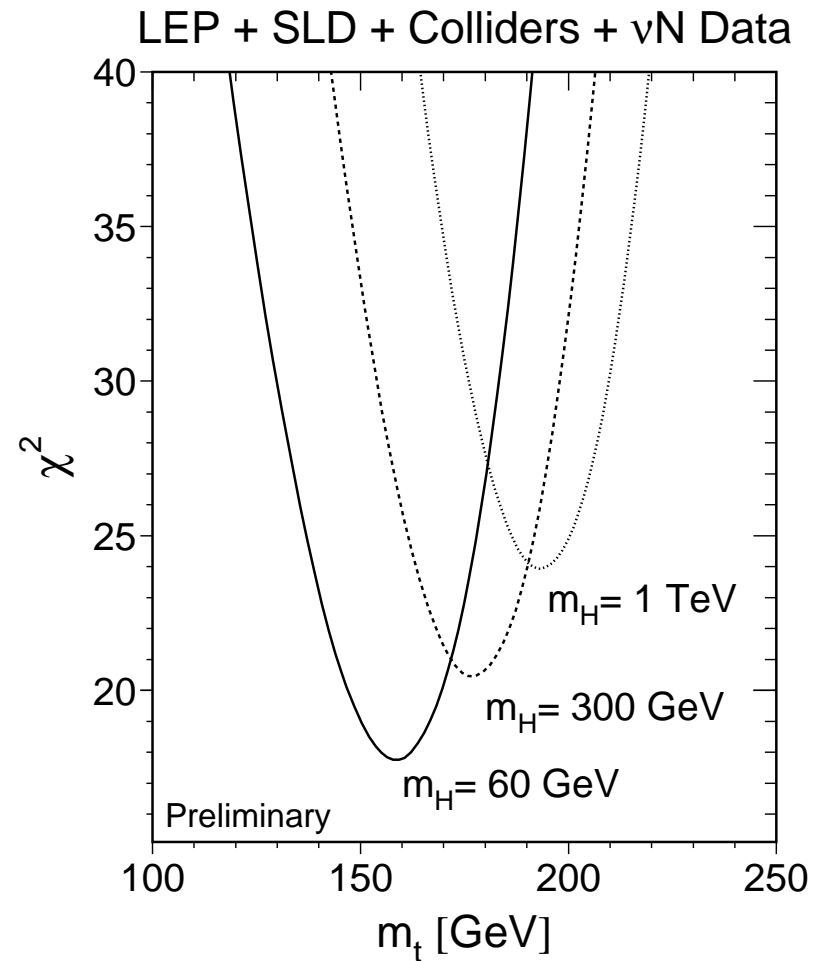
D.Schaile, R.Clare, M.Grünewald

*combination of data from the
individual experiments*

*analysis with tools provided
by theory*

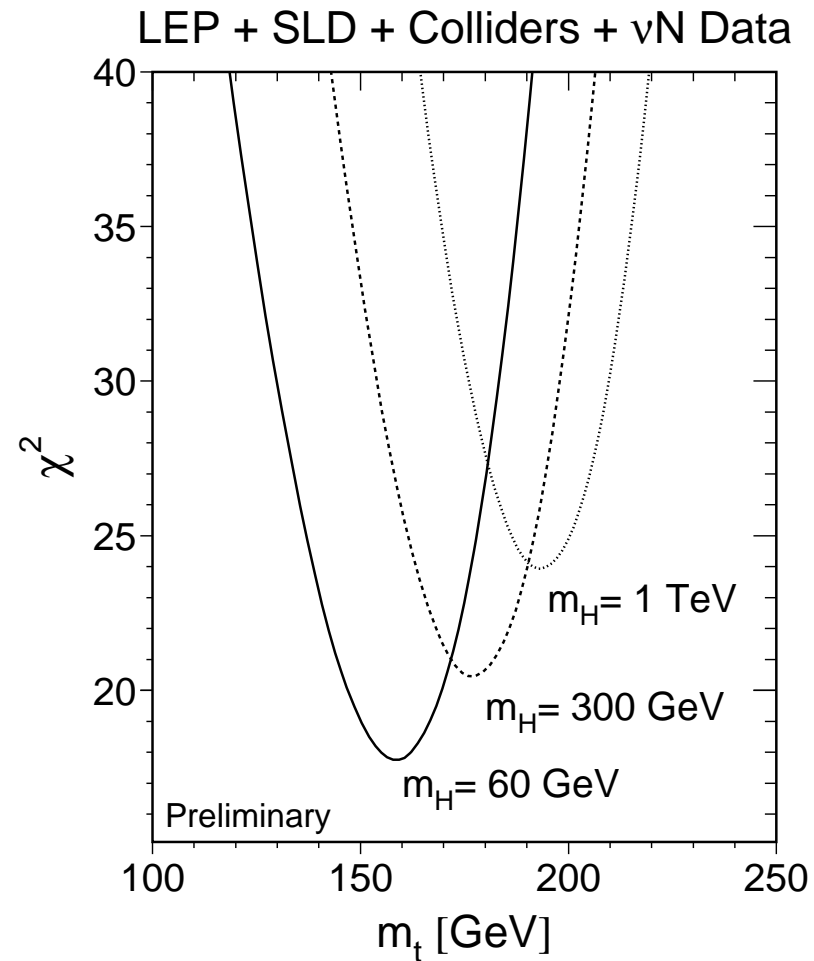
before the top quark was discovered (< 1995):

indirect mass determination $\Rightarrow m_t = 178 \pm 8^{+17}_{-20}$ GeV



before the top quark was discovered (< 1995):

indirect mass determination $\Rightarrow m_t = 178 \pm 8^{+17}_{-20}$ GeV

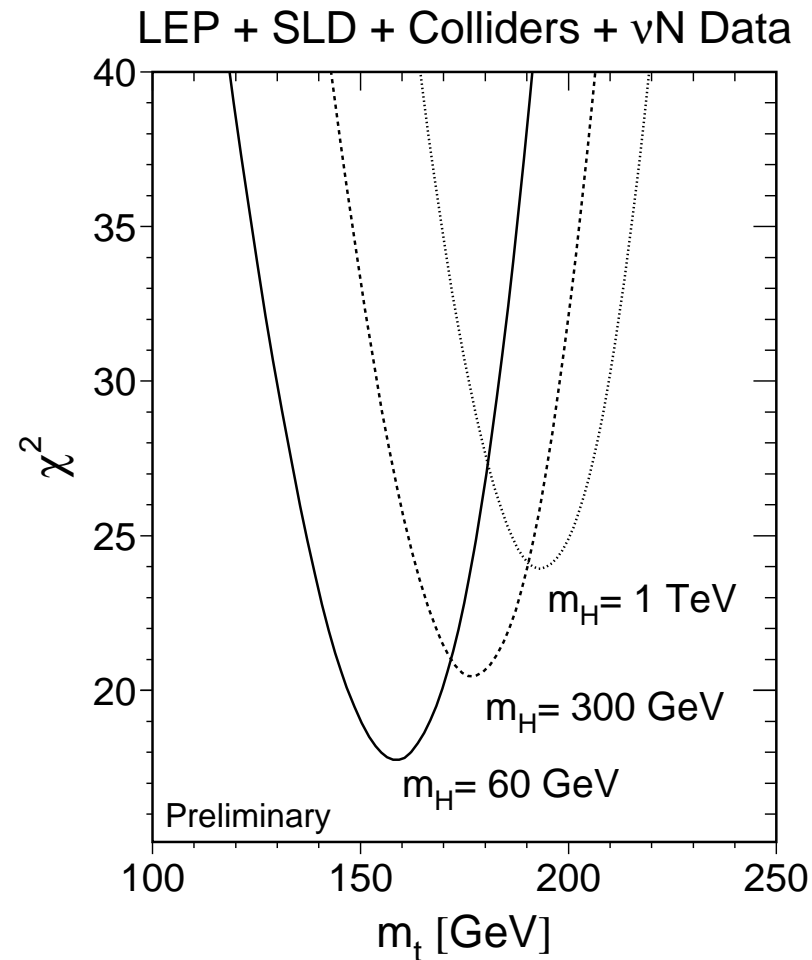


top discovery: *Tevatron 1995*

$m_t = 180 \pm 12$ GeV

before the top quark was discovered (< 1995):

indirect mass determination $\Rightarrow m_t = 178 \pm 8^{+17}_{-20}$ GeV



top discovery: *Tevatron 1995*

$m_t = 180 \pm 12$ GeV

today: *PDG 2022*

$m_t = 172.69 \pm 0.30$ GeV

CERN 95-03
31 March 1995

ORGANISATION EUROPÉENNE POUR LA RECHERCHE NUCLÉAIRE
CERN EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

**REPORTS OF THE WORKING GROUP
ON PRECISION CALCULATIONS
FOR THE Z RESONANCE**

Editors: D. Bardin
W. Hoflik
G. Passarino

GENEVA
1995

triggered by LEPEWWG

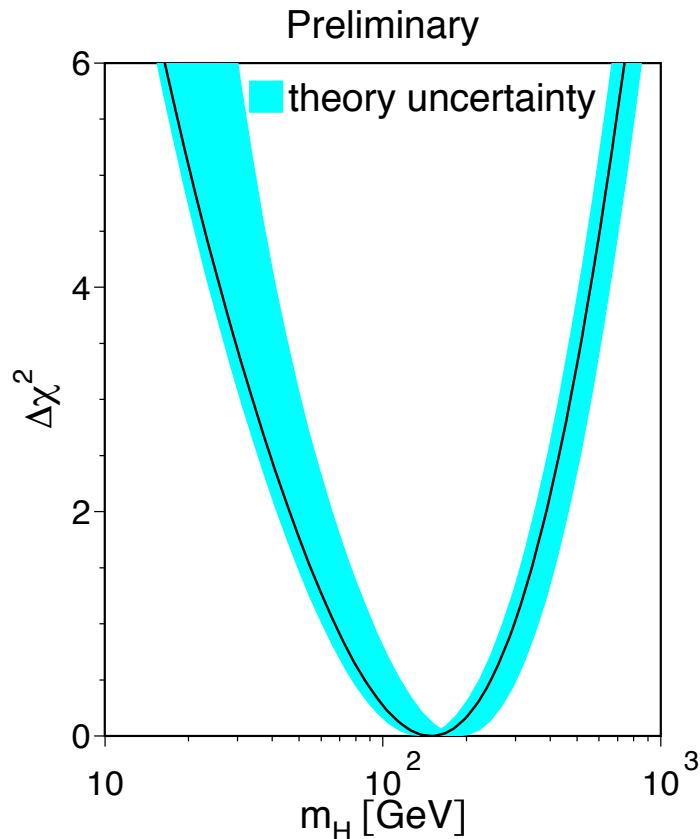
*updated codes for
precision observables*

*systematic estimate
of theoretical errors*

⇒ *“blueband plot”*

The first blueband plot

global fit to the Higgs-boson mass after the top discovery (ICHEP 1996)



CERN-PPE/96-183
December 6, 1996

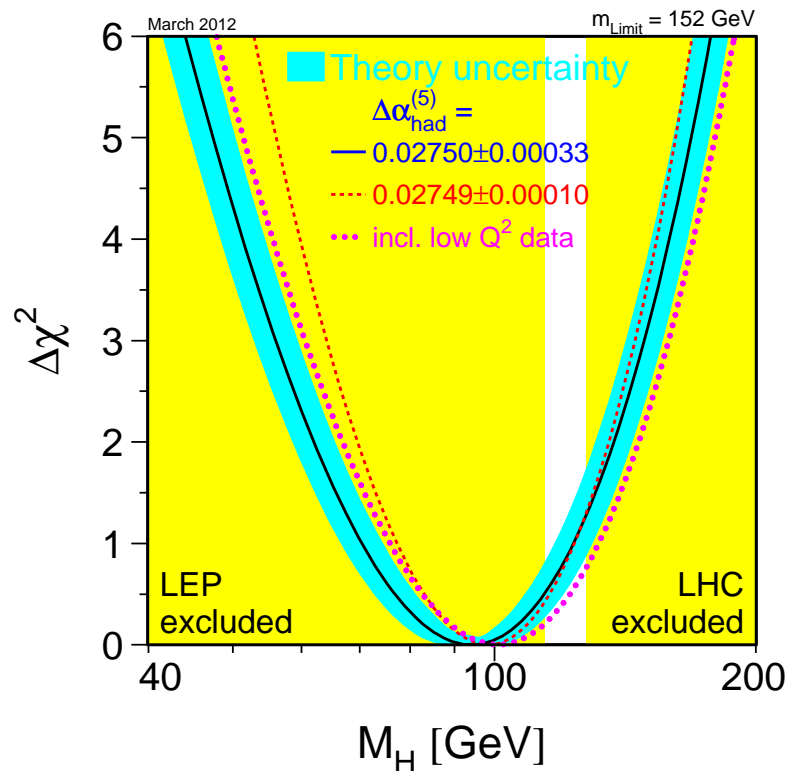
A Combination of Preliminary Electroweak Measurements and Constraints on the Standard Model

The LEP Collaborations* ALEPH, DELPHI, L3, OPAL, the LEP
Electroweak Working Group[†]
and the SLD Heavy Flavour Group[‡].
Prepared from Contributions
to the 28th International Conference on High Energy Physics, Warsaw, Poland,
25-31 July 1996.

Figure 9: $\Delta\chi^2 = \chi^2 - \chi_{min}^2$ vs. m_H curve. The line is the result of the fit using all data (last column of Table 22); the band represents an estimate of the theoretical error due to missing higher order corrections.

The last blueband plot

before the Higgs boson discovery in summer 2012



$$M_H < 152 \text{ GeV (95\%C.L.)}$$

$$M_H = 94_{-24}^{+29} \text{ GeV}$$

ATLAS 2012:

$$M_H = 126.0 \pm 0.4 \pm 0.4 \text{ GeV}$$

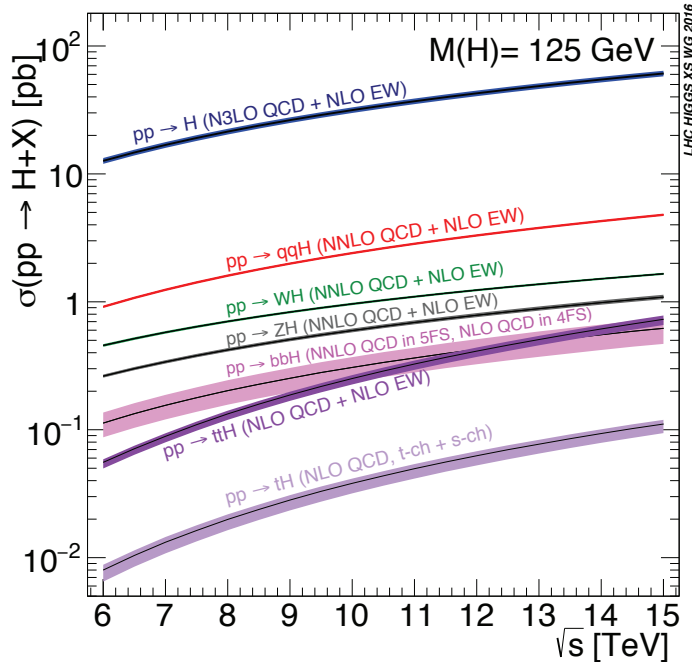
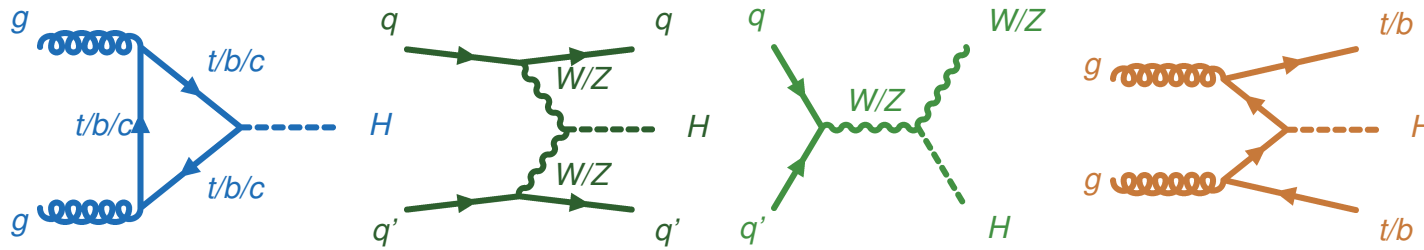
CMS 2012:

$$M_H = 125.3 \pm 0.4 \pm 0.5 \text{ GeV}$$

today (PDG 2022)

$$M_H = 125.25 \pm 0.17 \text{ GeV}$$

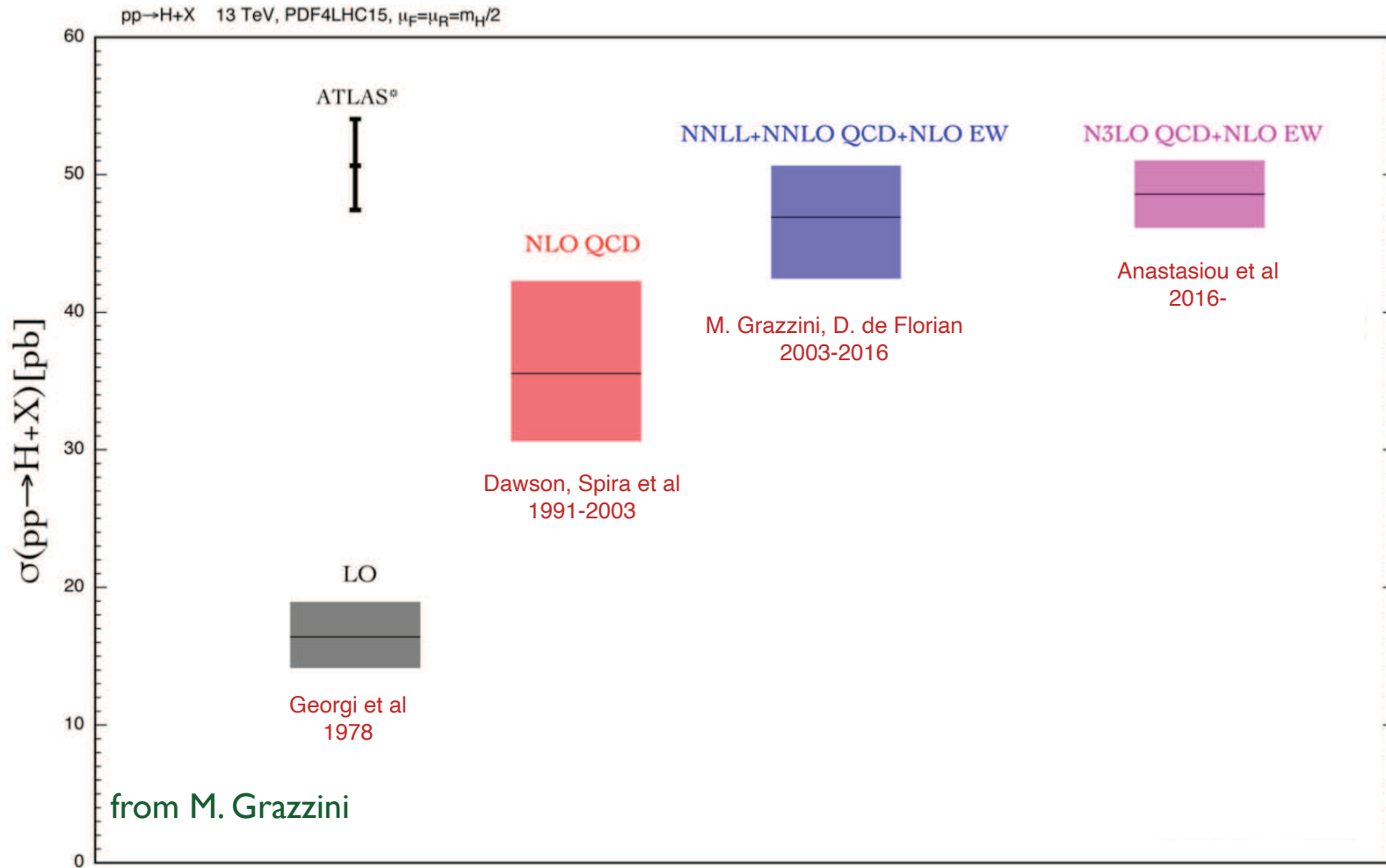
Higgs boson production at the LHC



*theory predictions with higher orders
from QCD (4-loop) and ew (2-loop)*

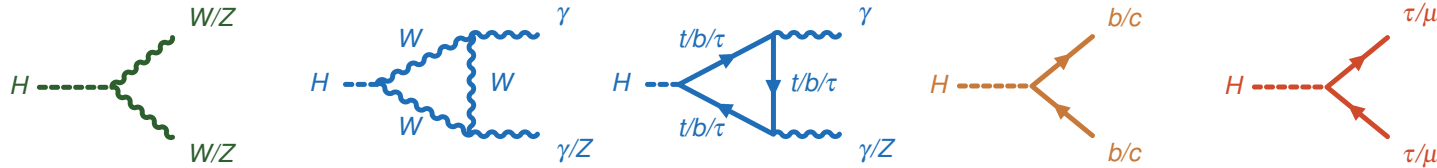
*Handbook of Higgs Cross Sections Vol. 4
De Florian et al.*

CERN Yellow Report 2017-002

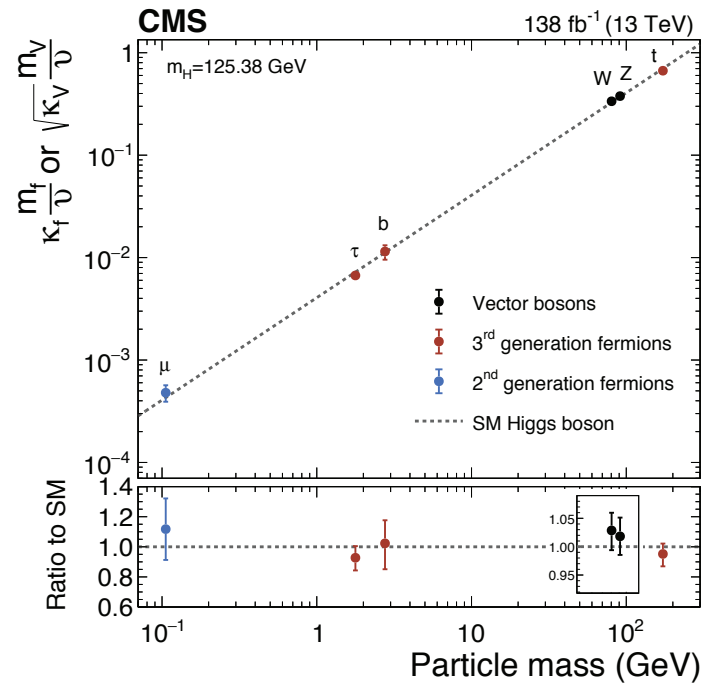
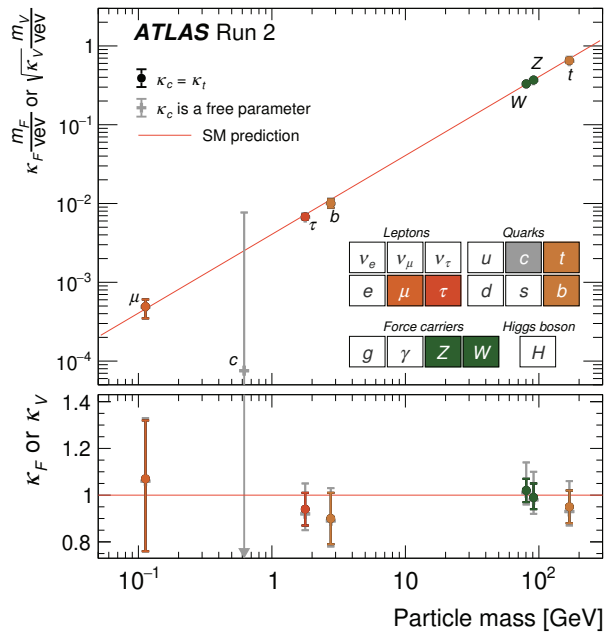


incomplete higher-order calculations provoke wrong conclusions!

Higgs boson decays



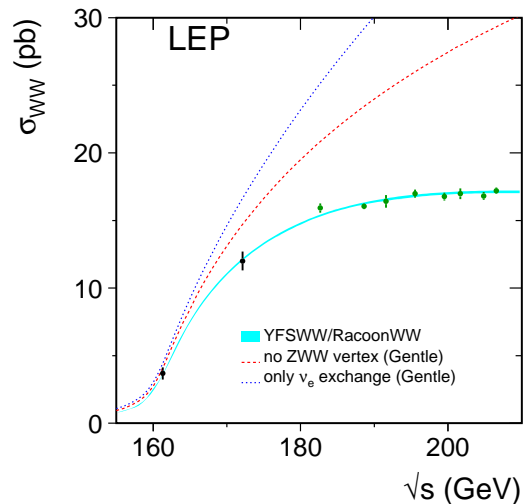
predictions with higher orders *in:* *Handbook of Higgs Cross sections Vol. 4*



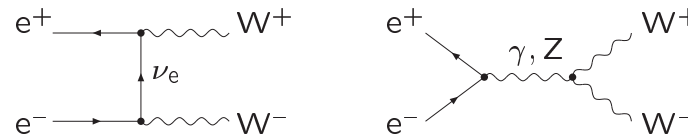
measurements: couplings scale with masses as predicted
by the ew theory

EW theory: what has been established

- coupling structure from local $SU(2) \times U(1)$ gauge symmetry
 - ★ vectorboson– fermion couplings
 - ★ cubic vectorboson self-couplings



$$e + e^- \rightarrow WW (\rightarrow 4f)$$



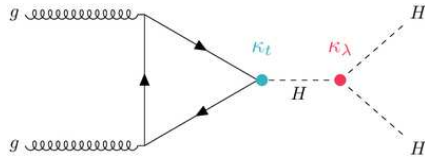
theo: RacoonWW
 Gentle
 YFSWW
 (codes with h.o.)

Denner et al.
Bardin et al.
Jadach et al.

- signals from gauge-symmetry breaking
 - ★ existence of a Higgs boson
 - ★ couplings to gauge bosons as predicted by Higgs mechanism
 - ★ couplings to fermions as predicted by Yukawa interactions

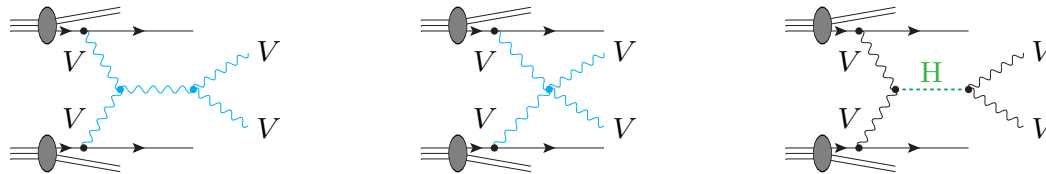
EW theory: what is missing

- Higgs couplings with higher precision, in particular Yukawa couplings of light fermions
- Higgsboson self-coupling (\rightarrow scalar potential)



appears in HH -pair production

- vectorboson scattering at high energies
goes to the core of electroweak symmetry breaking



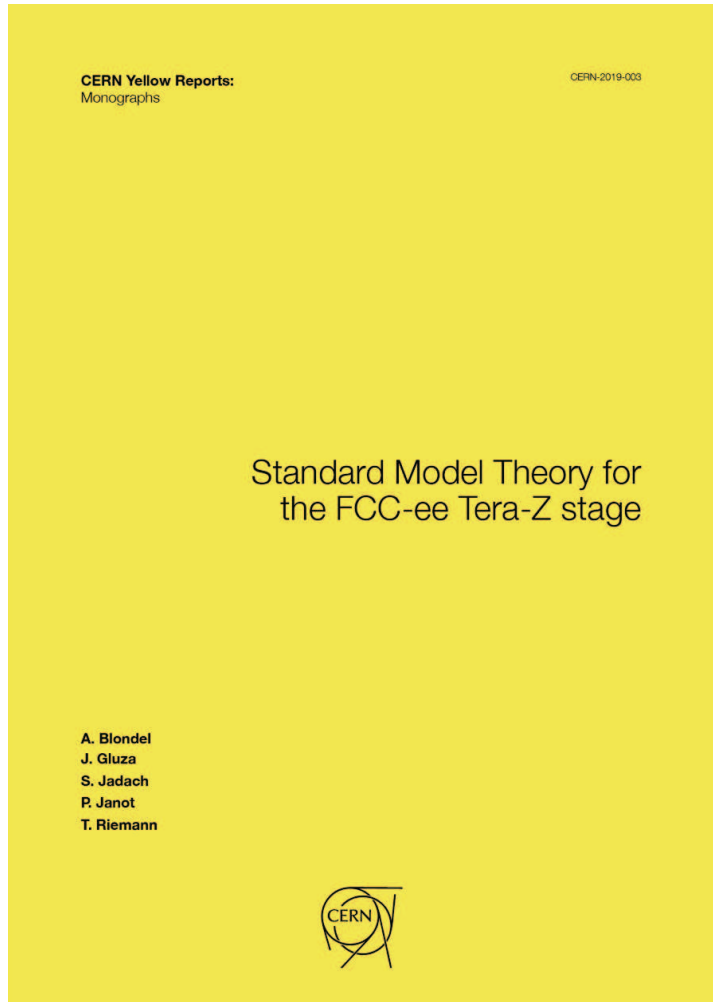
*six-particle final states + higher orders (radiation processes and loops)
new techniques invented for multi-leg and multi-loop calculations*

- LHC: processes already observed
precision studies in forthcoming LHC runs

Concluding remarks

- EW theory: a result of close work in theory and experiment
 - built on basic principles (*unitarity, renormalizability, symmetry*)
 - accompanied by continuous flow of new exp results
 - a textbook case: from predictions to discoveries
- EW precision physics prepared the ground for discoveries
 - unprecedented precision measurements
 - challenge for theory → era of precision calculations
 - new style of theoretical work
- convincing concept → great success
 - confirmation as a working quantum field theory
 - indirect access to heavy particles
 - can be repeated for BSM physics?

activities preparing the future . . .



<i>exp. error</i>	<i>now</i>	<i>FCC</i>
M_W [MeV]	12	1
$\sin^2 \theta_{\text{eff}}$ [10^{-5}]	16	0.6

- need:
theory calculations
with one loop more

additional slides

with M_H , precision observables are now uniquely determined

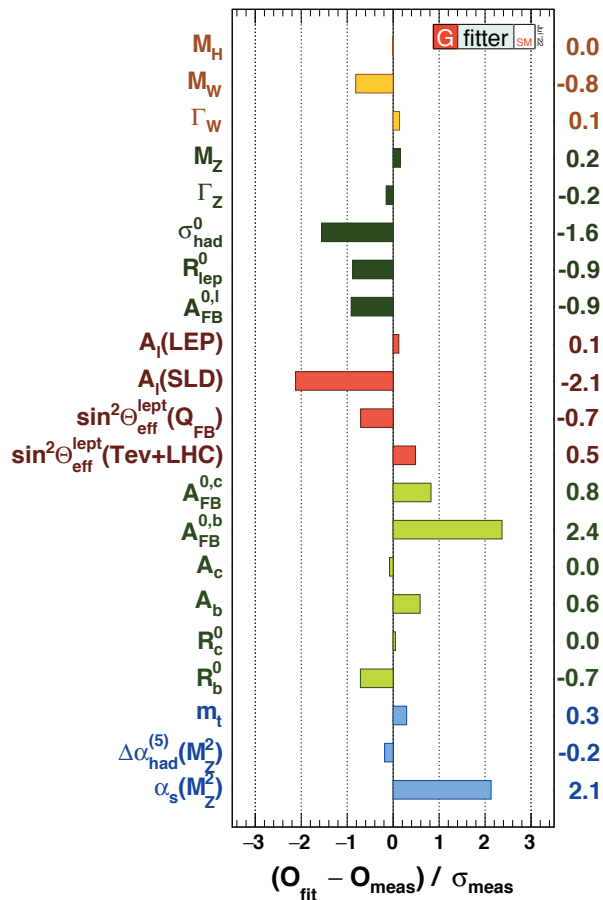
global fits are consistency tests with increasing accuracy from LHC

theory predictions with all known h.o. contributions are available as parametrizations in terms of the input quantities

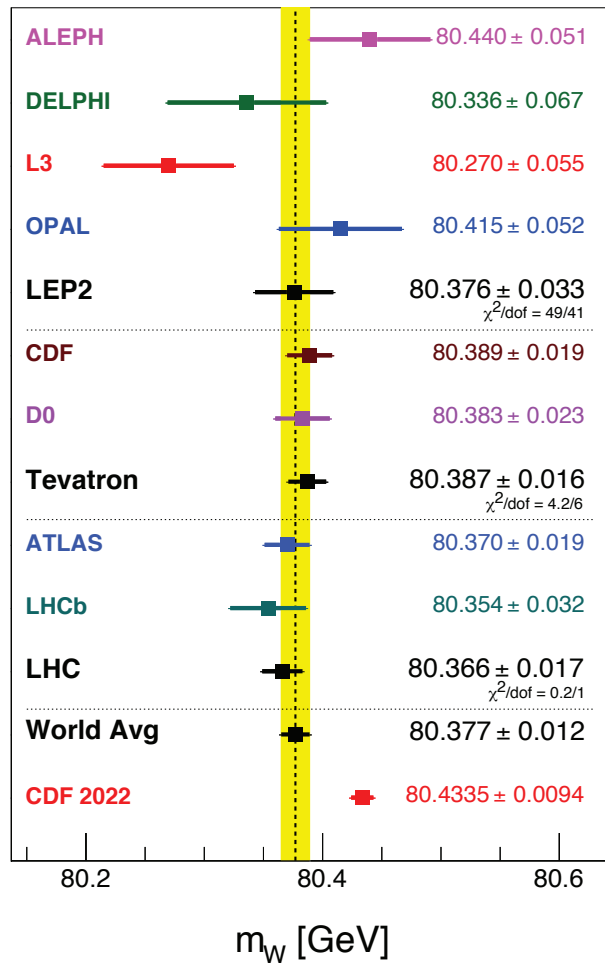
Dubovyk, Freitas et al. 2019

codes for global fits

HEPfit, Gfitter, . . .



[ICHEP 2022]



theory:

$$M_W = 80.353 \pm 0.005 \pm 0.004 \text{ GeV}$$