

Challenges of fitting cross-section data

Peelle's pertinent puzzle

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NuXTract Workshop 2023

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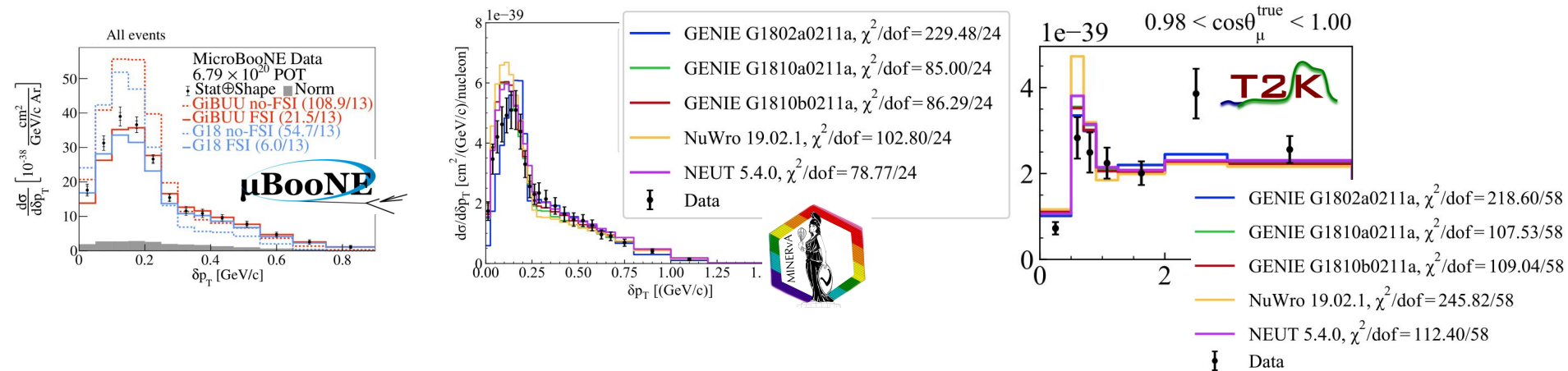
Based on work performed in the context of [arXiv:2308.01838](https://arxiv.org/abs/2308.01838)



Introduction

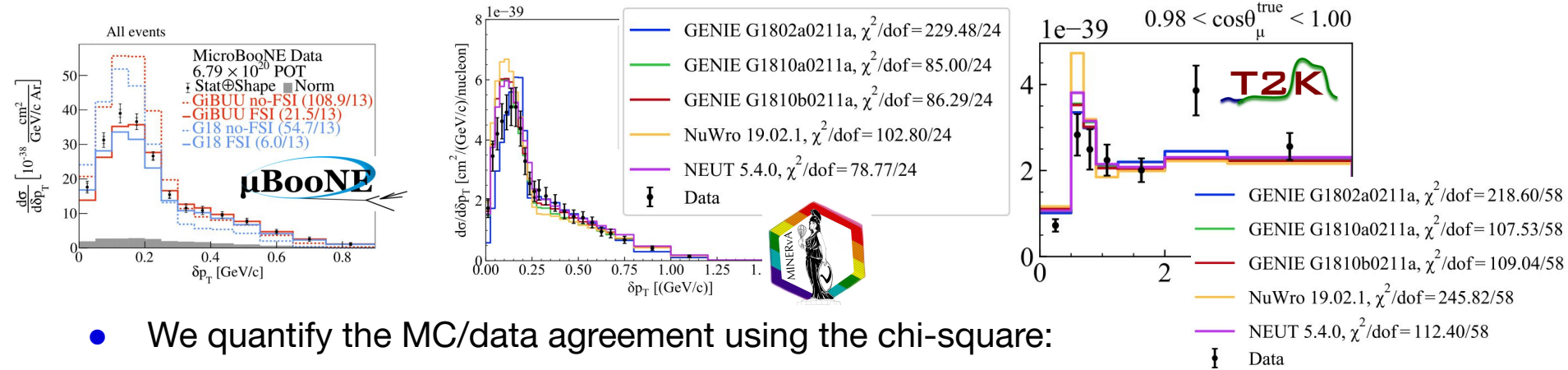
- Throughout the workshop, various techniques of cross-section extraction were presented
- Once the cross-section measurement is published, it can be used to **benchmark our models** from MC generators (NEUT, GENIE, NuWro, ...)
- We find that, almost all the times, the models do not give a good description of the available cross-section data
- We often attempt to **tweak the models** using theoretically-motivated (or not!) parameters to better describe the data through **chi-square fits**
- A few recent examples:
 - GENIE: [Phys. Rev. D 106, 112001](#), [Phys. Rev. D 105, 072001](#), ... (see talks by [Julia](#) and [Michael](#))
 - NEUT: [arXiv:2308.01838](#), ...
 - NuWro: [Phys. Rev. C 102, 015502 \(2020\)](#), ...

- No neutrino MC event generator is able to give a satisfactory description of neutrino-nucleus cross-section data (see TENSIONS 2019 [report](#))*



* This remains true even when we consider uncertainties within the models

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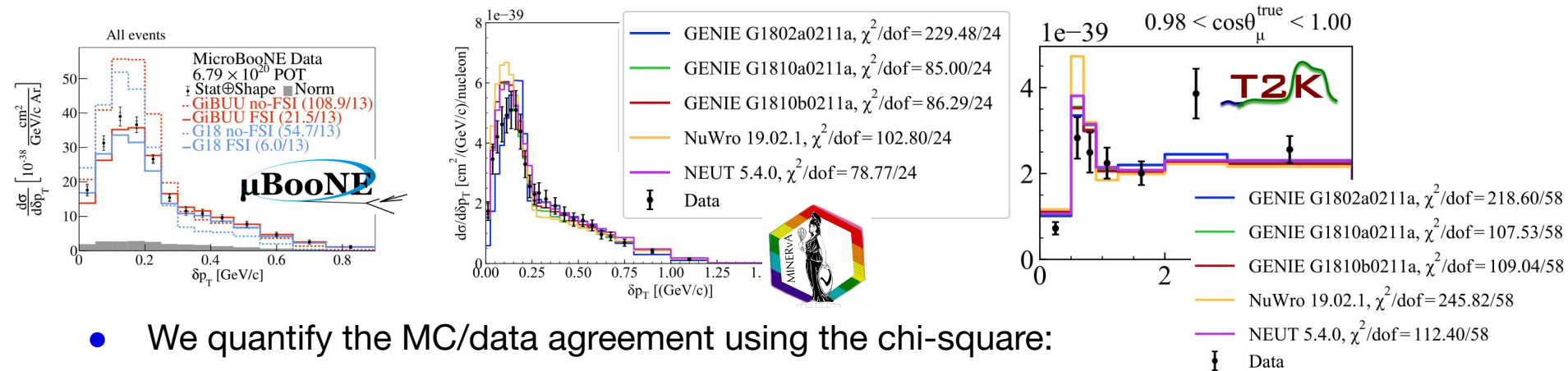


- We quantify the MC/data agreement using the chi-square:

$$\chi^2 = \sum_{ij} (B_i - B_i^{\text{MC}}) M_{ij}^{-1} (B_j - B_j^{\text{MC}})$$

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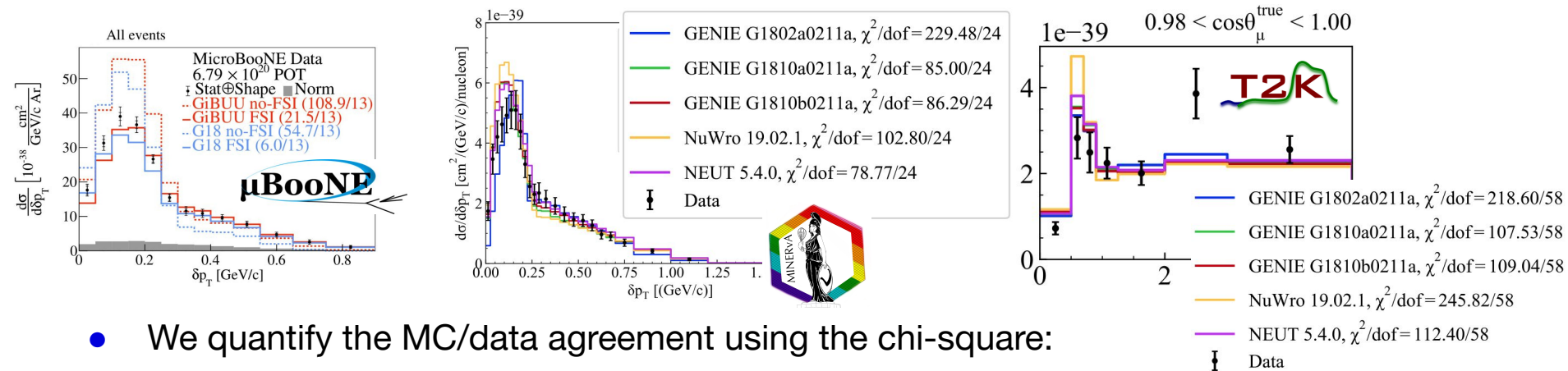


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Published data: bin content + bin-to-bin covariance

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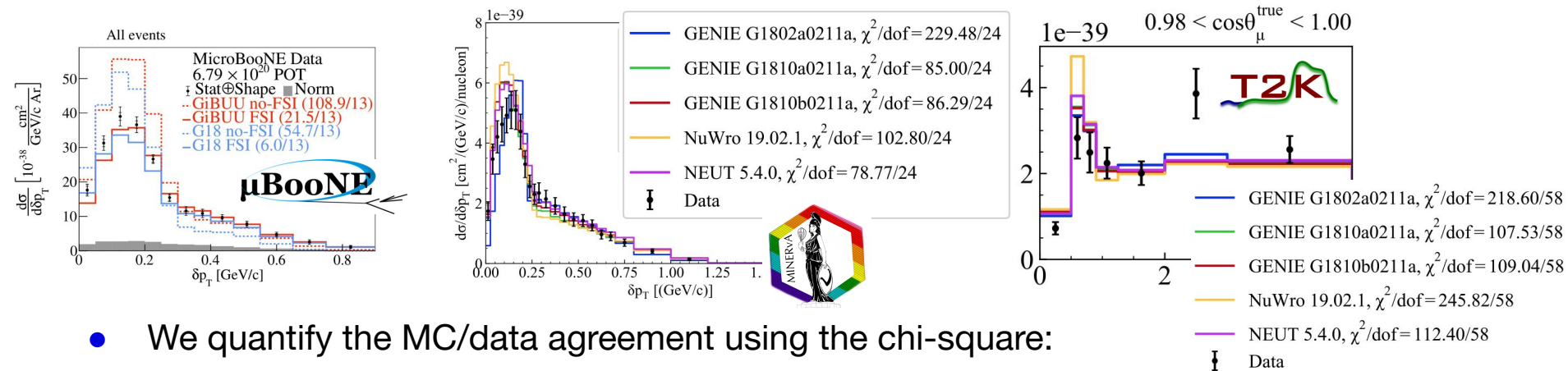


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Bin content predicted by the model

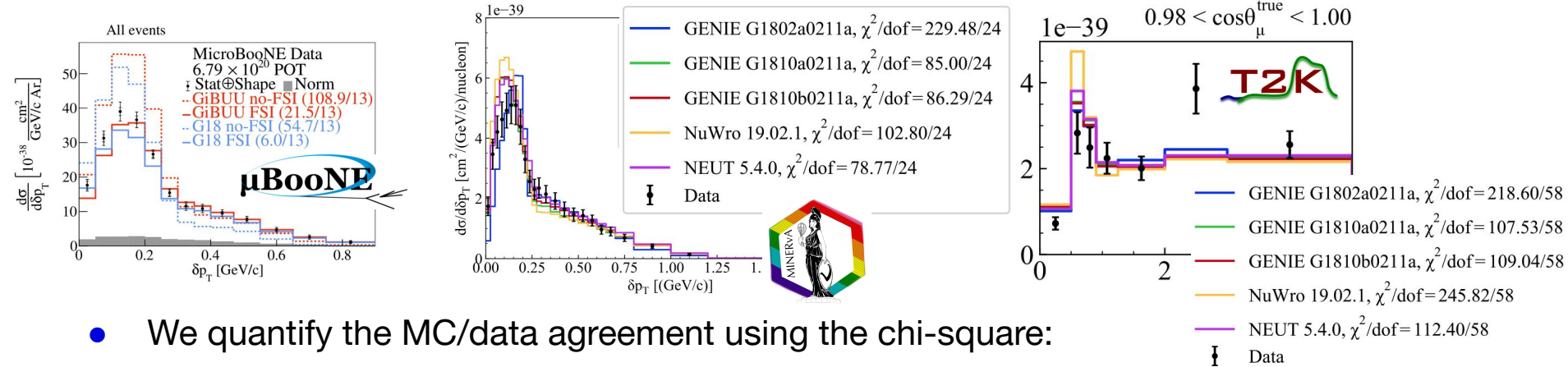
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$$\chi^2(\vec{x}) = \sum_{ij} (B_i - B_i^{\text{MC}}(\vec{x})) M_{ij}^{-1} (B_j - B_j^{\text{MC}}(\vec{x}))$$

- To improve this agreement, we use parameters (knobs/dials) that tweak the model predictions and perform fits to the data



- Ingredients to fit cross-section data:

- Model



[JINST 12 P01016 \(2017\)](#)



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Allows us to:

- Add neutrino cross-section data
- Implement custom parameters
- Interact with MC reweight engines
- Compare/fit models to data
- ...

(See talks by [Laura](#) and [Luke](#))

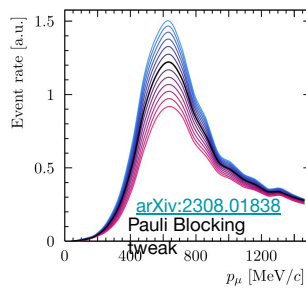
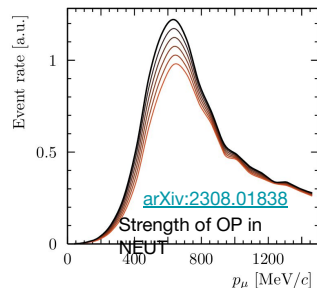
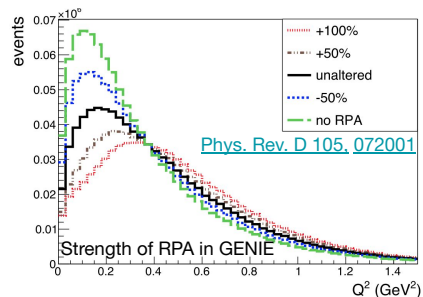
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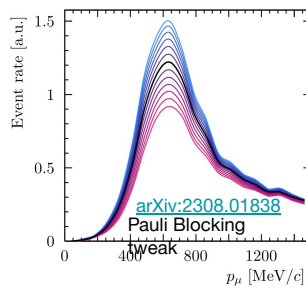
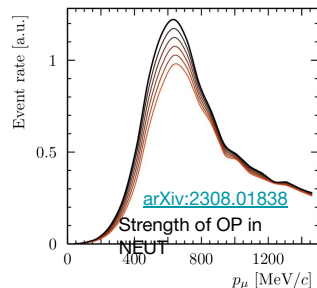
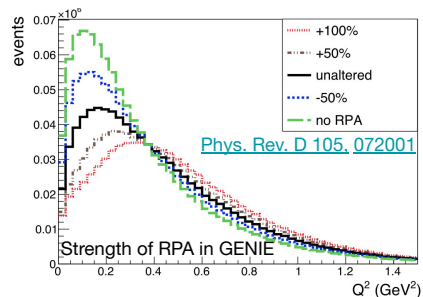
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- Ingredients to fit cross-section data:

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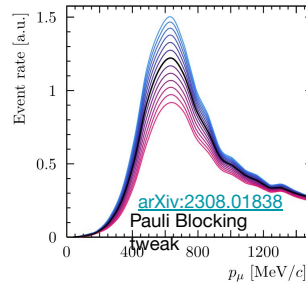
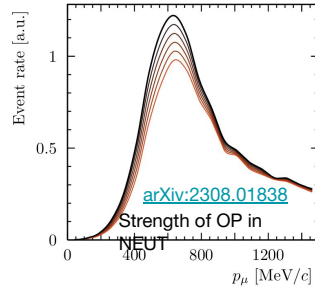
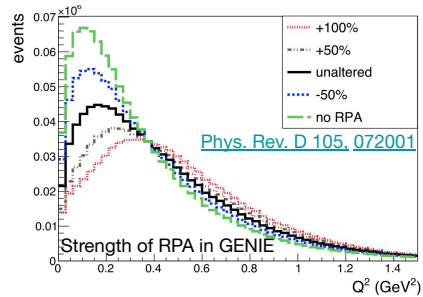
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- This chi-square is then minimized as a function of \vec{x} .



Allows us to:

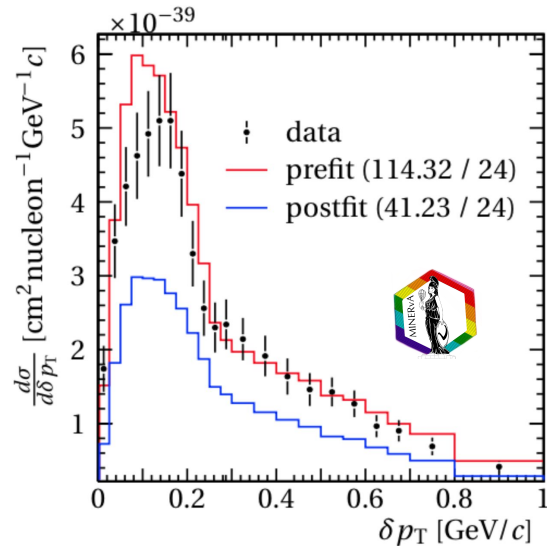
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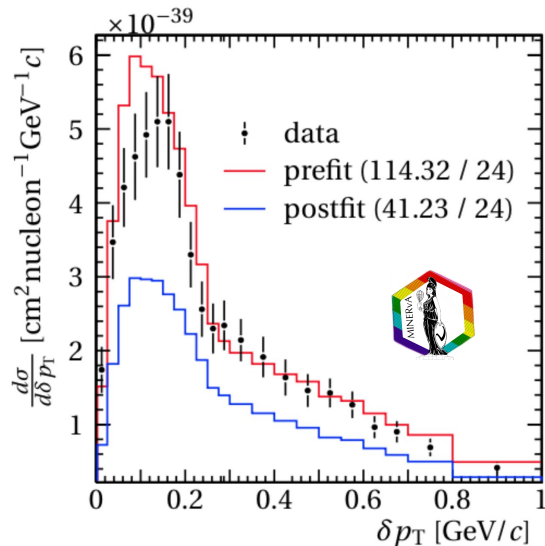
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Peelle's pertinent puzzle (PPP)

- Peelle's pertinent puzzle in one picture



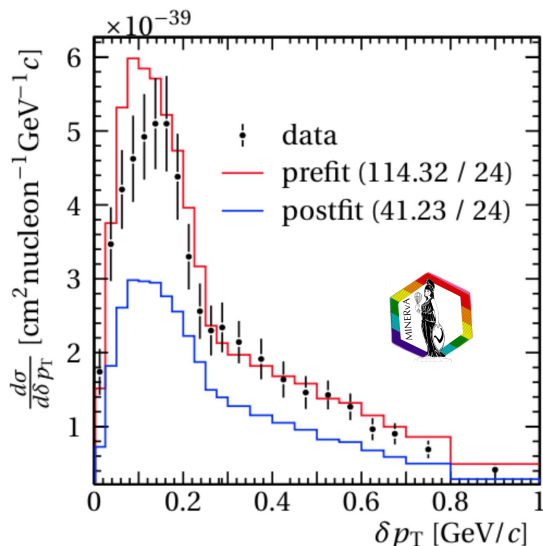
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I chose my model (NEUT), implemented a few parameters, and ran a fit in NUISANCE to MINERvA data...

The postfit model I obtain is with a very small normalization!?

- Peelle's pertinent puzzle in one picture

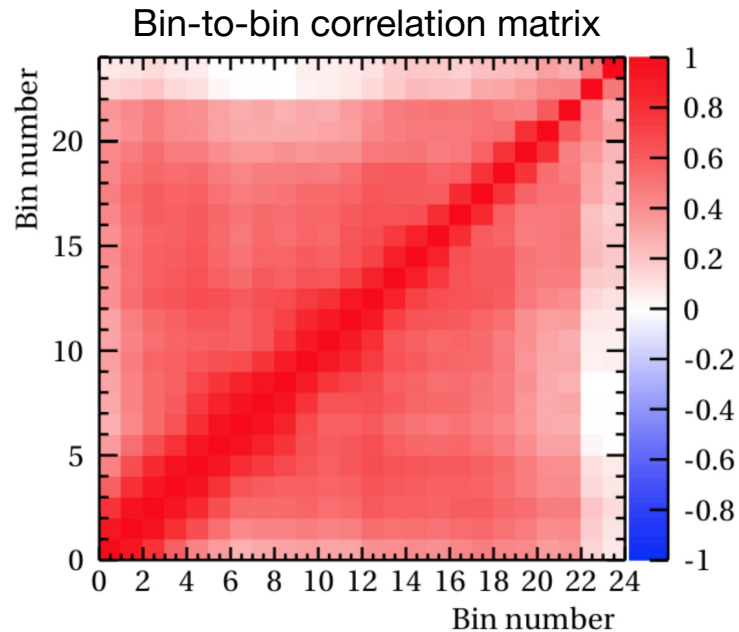
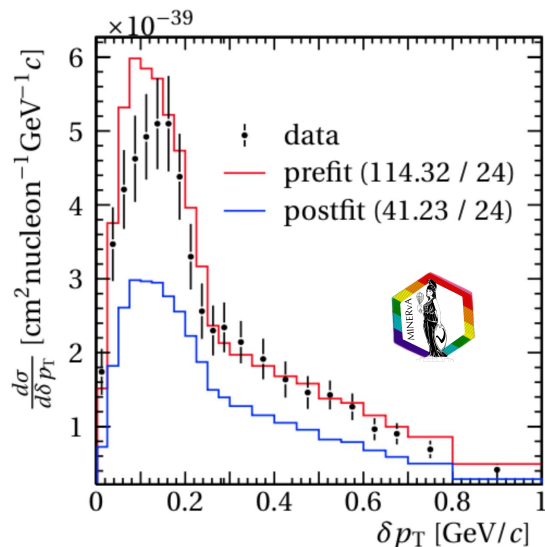


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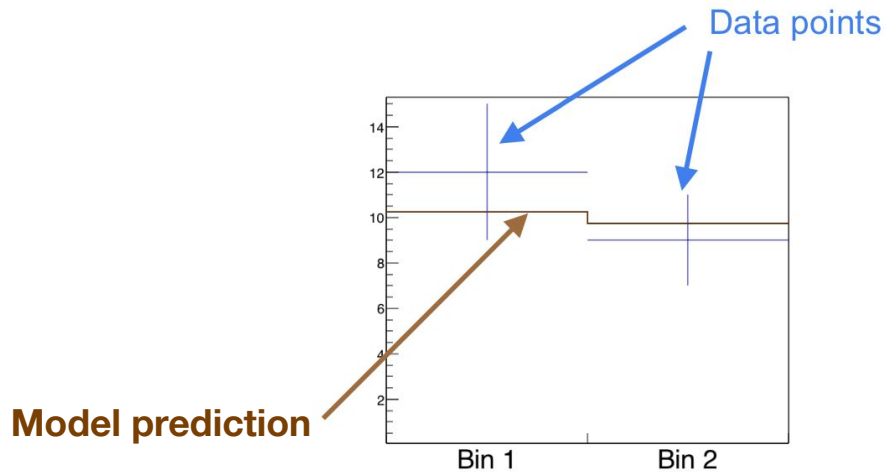
- This is the result of two things:
 - “Flawed” model that is **unable to perfectly describe the data** (even with the introduced free parameters)

- Peelle's pertinent puzzle in one picture



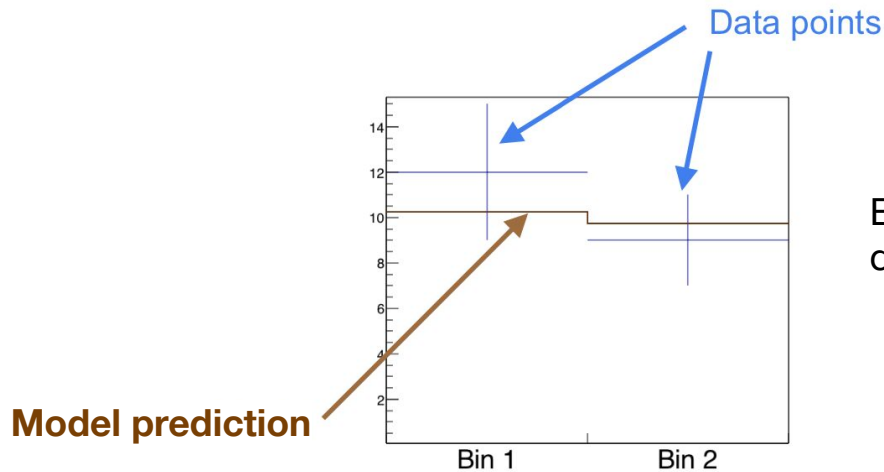
- This is the result of two things:
 - “Flawed” model that is **unable to perfectly describe the data** (even with the introduced free parameters)
 - Highly correlated uncertainties between the bins summarized under **Gaussian** assumptions

- Let's have a look at an example of a 2-bin measurement



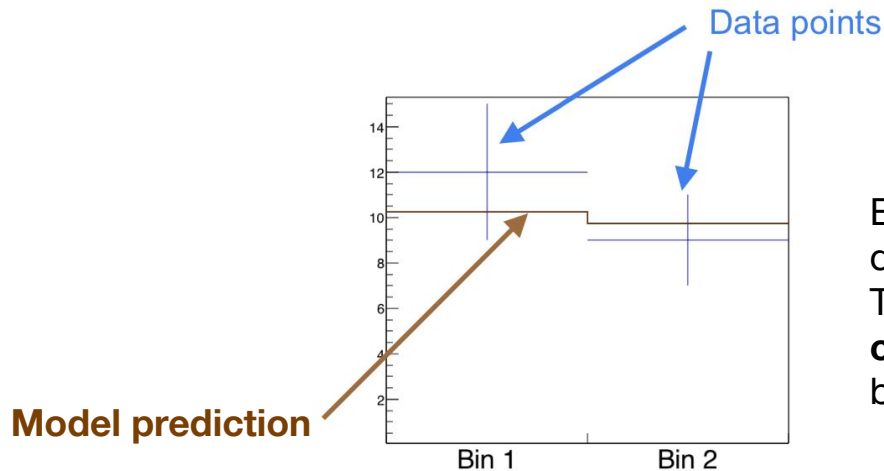
Why does this happen?

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By eye, the agreement looks quite good, **but...**

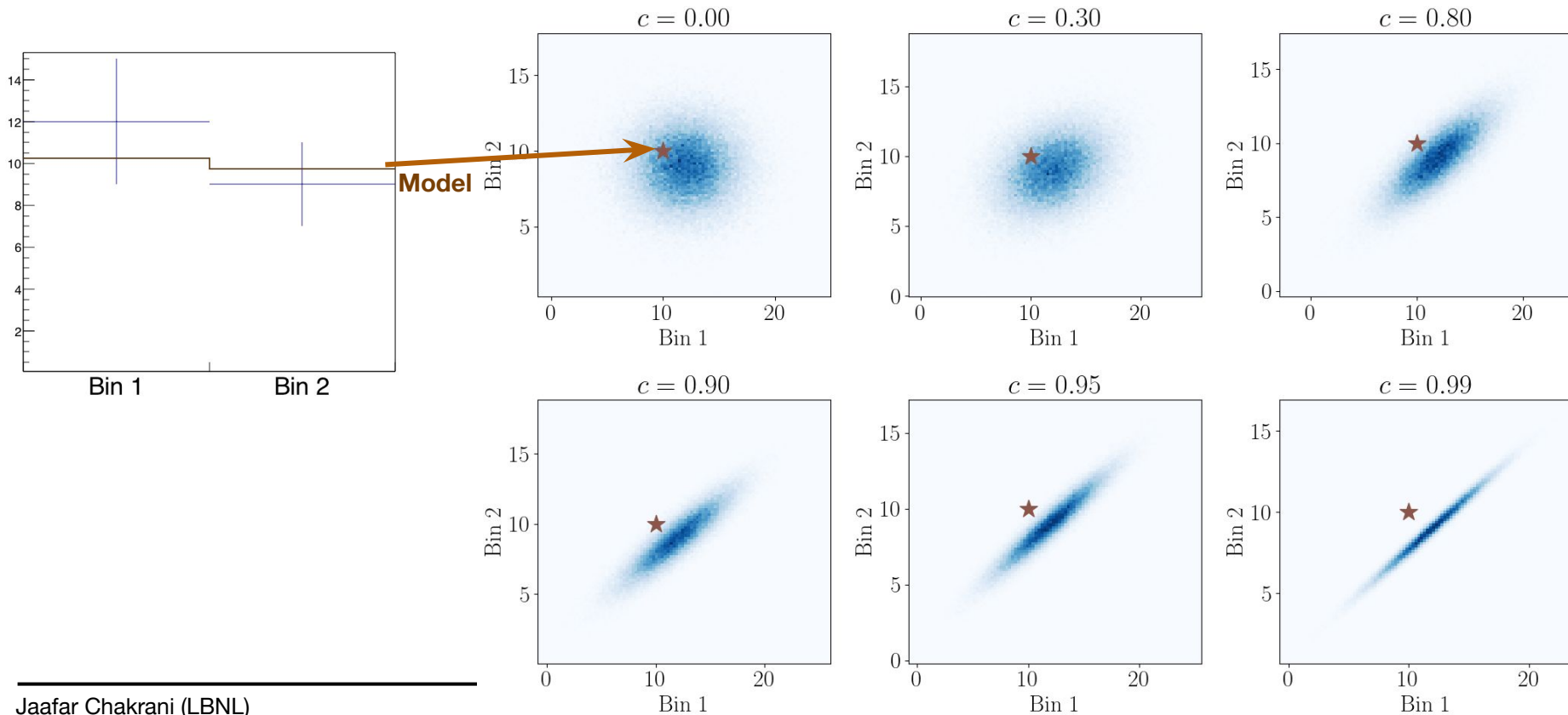
- Let's have a look at an example of a 2-bin measurement



By eye, the agreement looks quite good, **but...**
This actually depends on the **correlation** between the two bins!

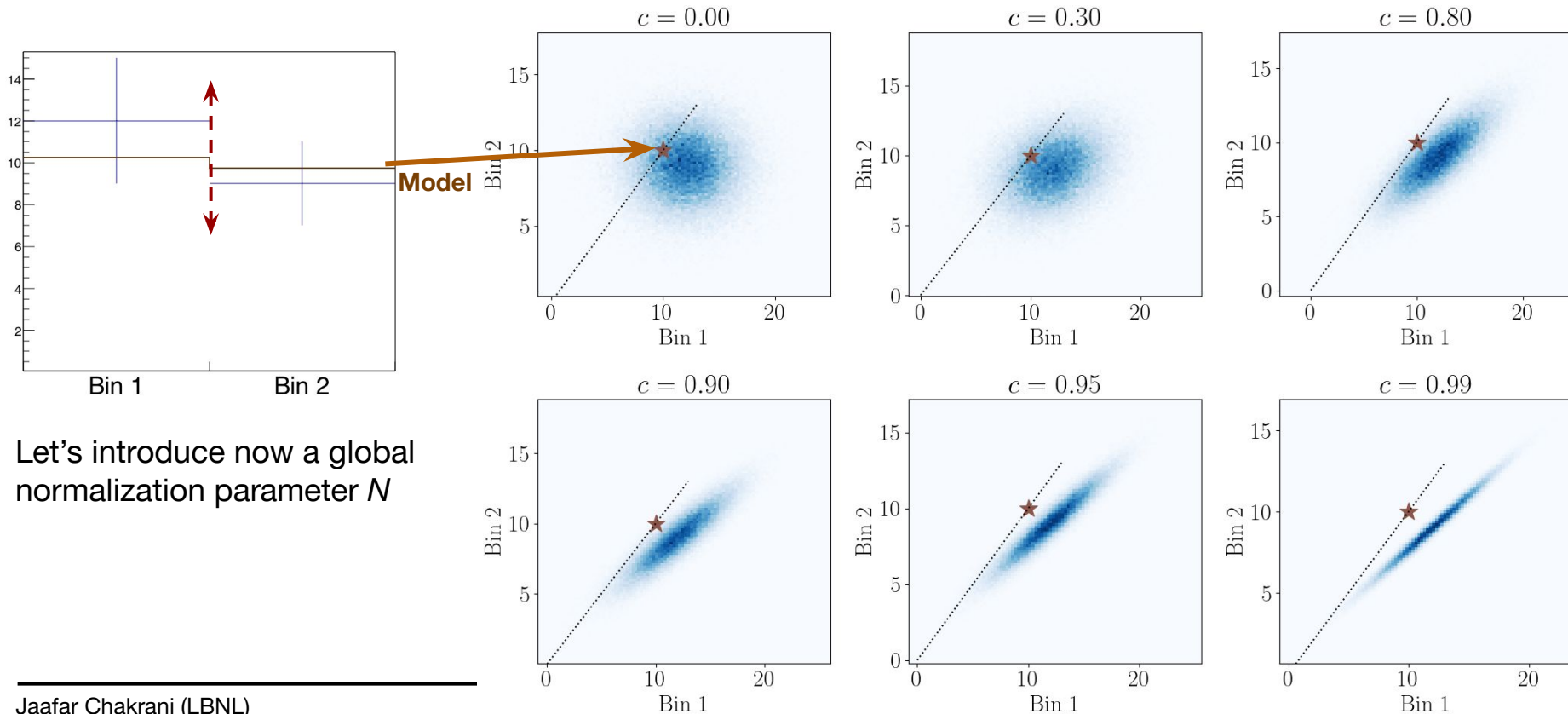
correlation = 0.00	correlation = 0.30	correlation = 0.80	correlation = 0.90	correlation = 0.95
chi2 = 0.69	chi2 = 0.98	chi2 = 3.41	chi2 = 6.81	chi2 = 13.62

- In the two-dimensional space of the two bins



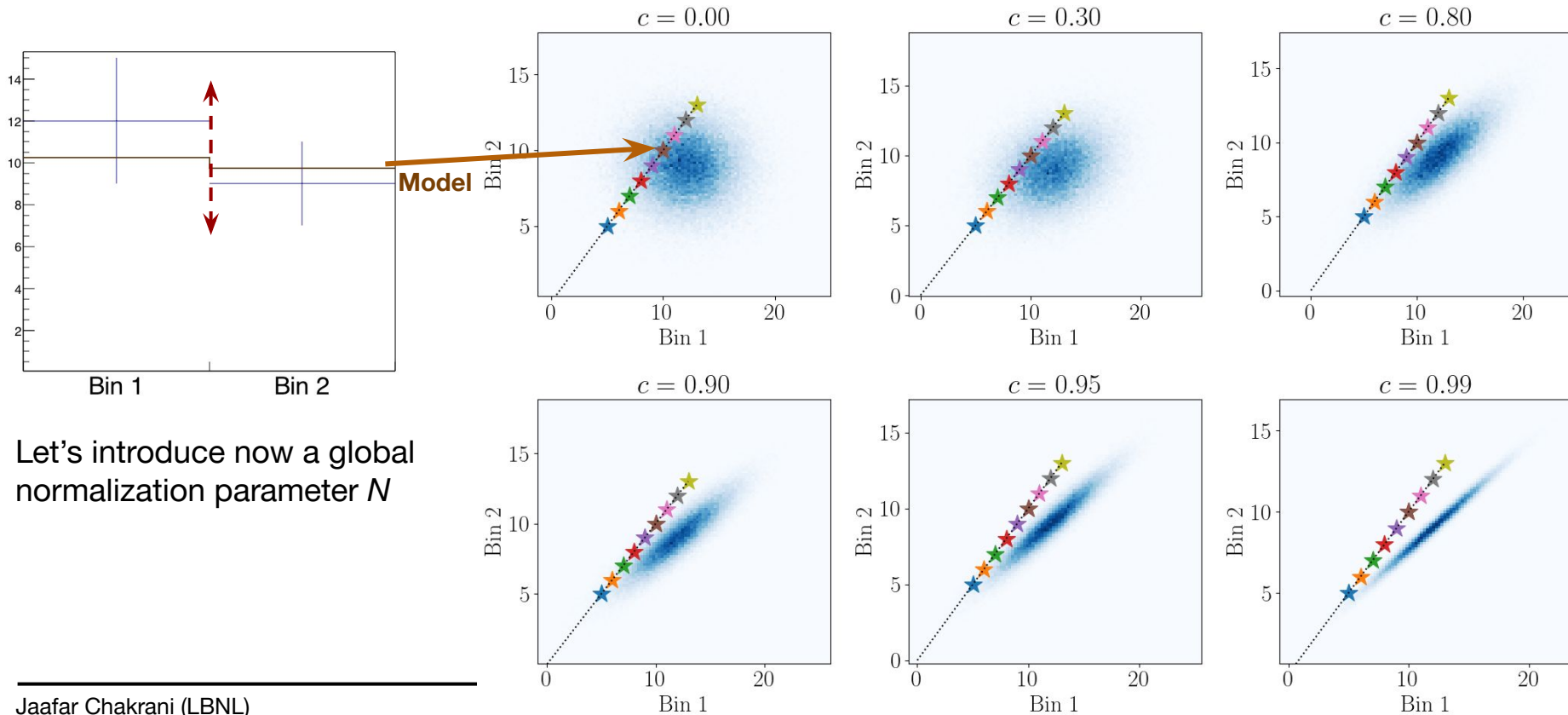
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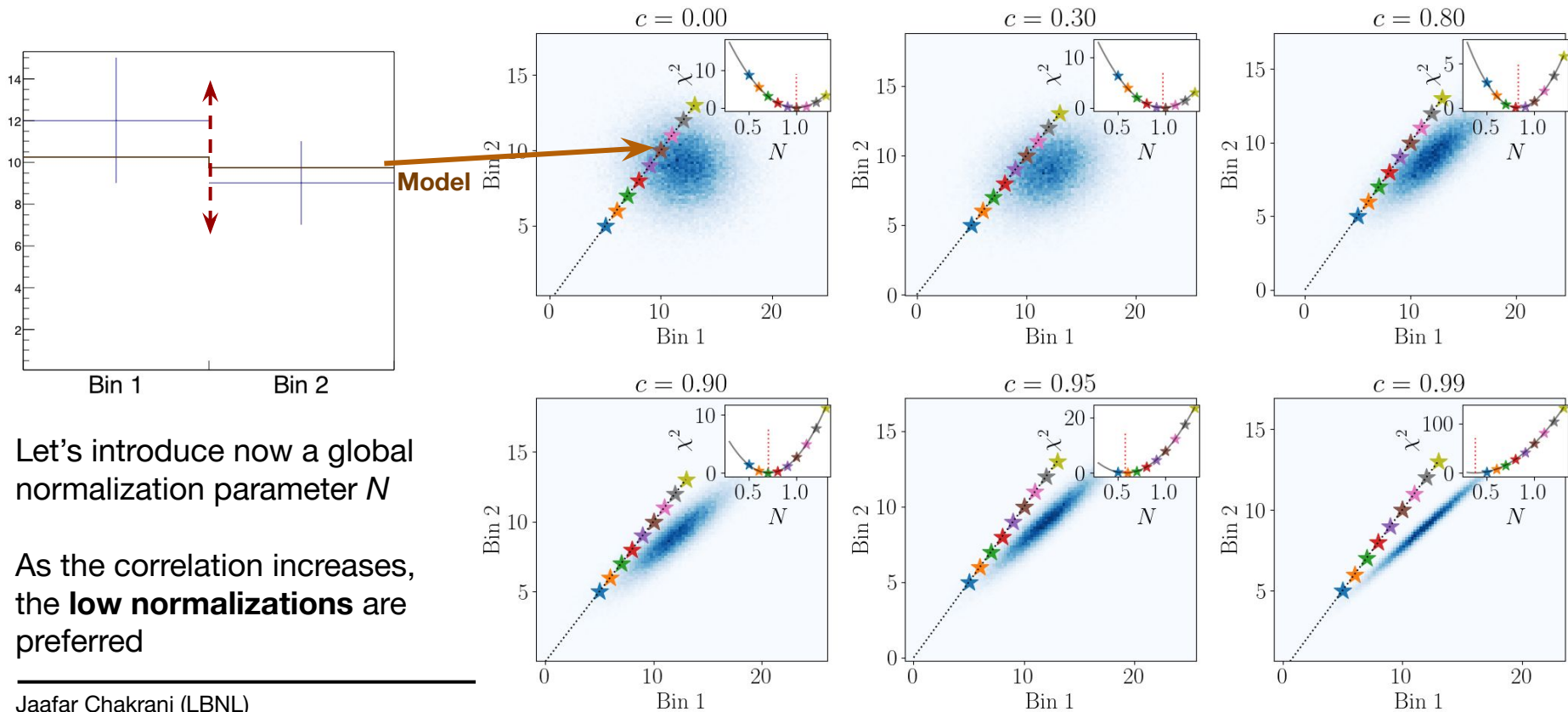
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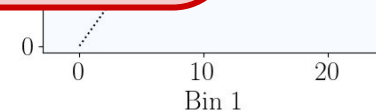
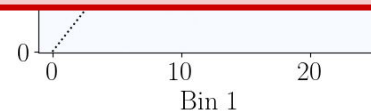
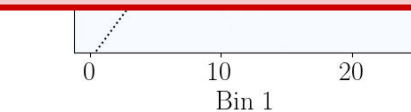
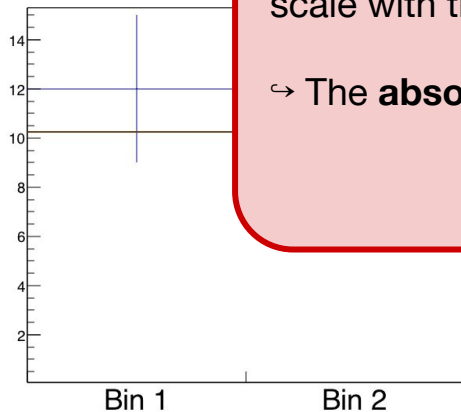
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With the standard covariance approach, the uncertainty on the shape does not scale with the changes to the normalization

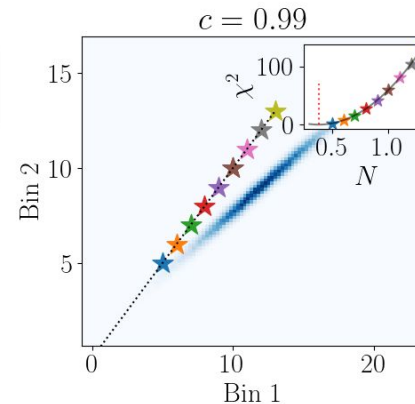
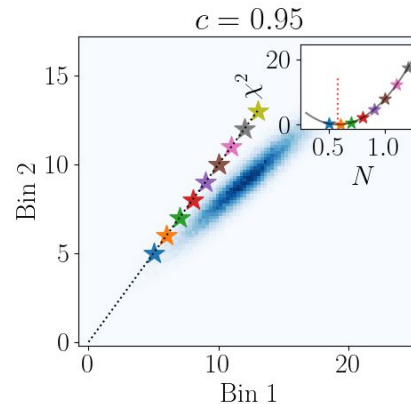
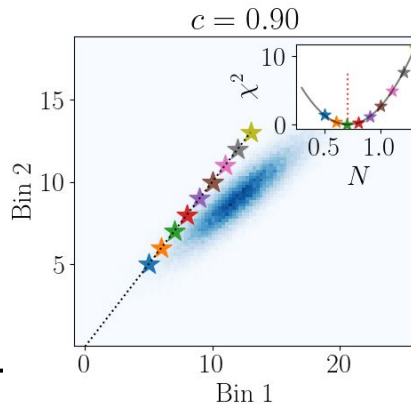
↪ The **absolute** uncertainties are unchanged as a function of the normalization

D'Agostini, [NIMA 346 \(1994\)](#)



Let's introduce now a global normalization parameter N

As the correlation increases, the **low normalizations** are preferred



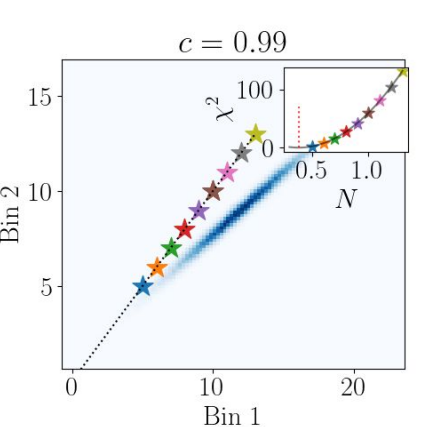
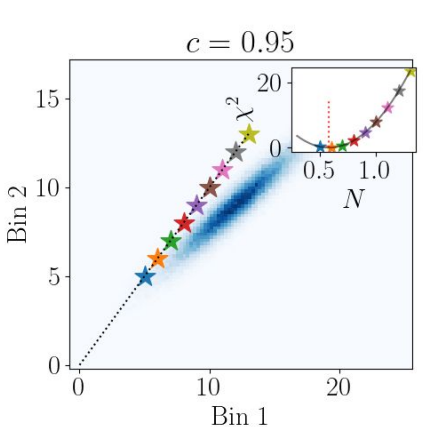
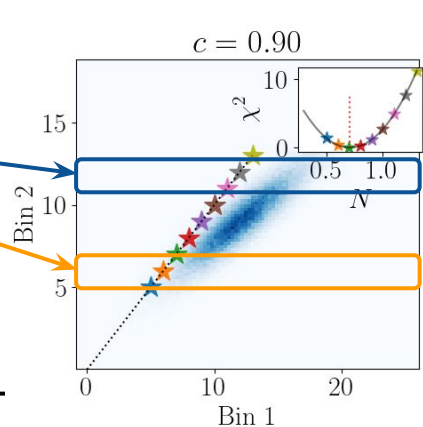
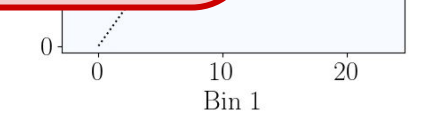
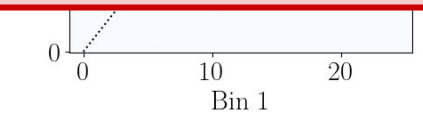
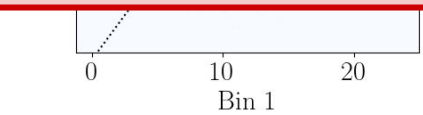
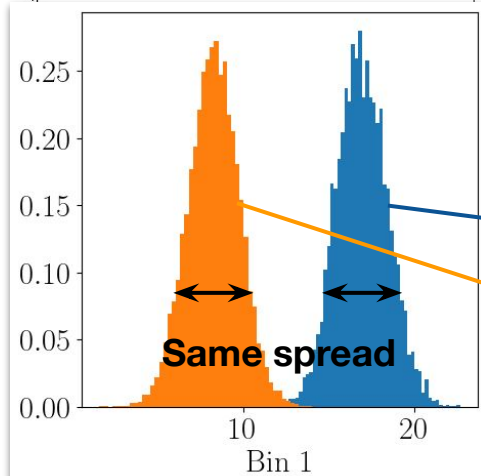
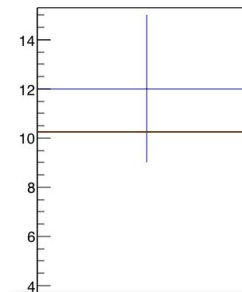
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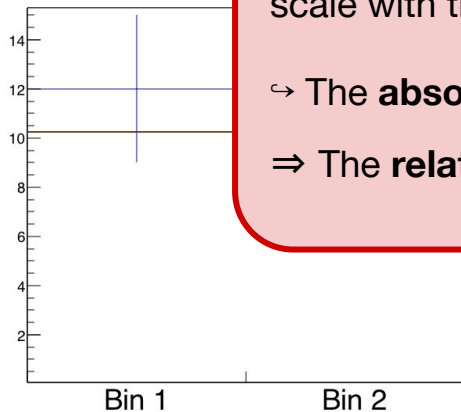
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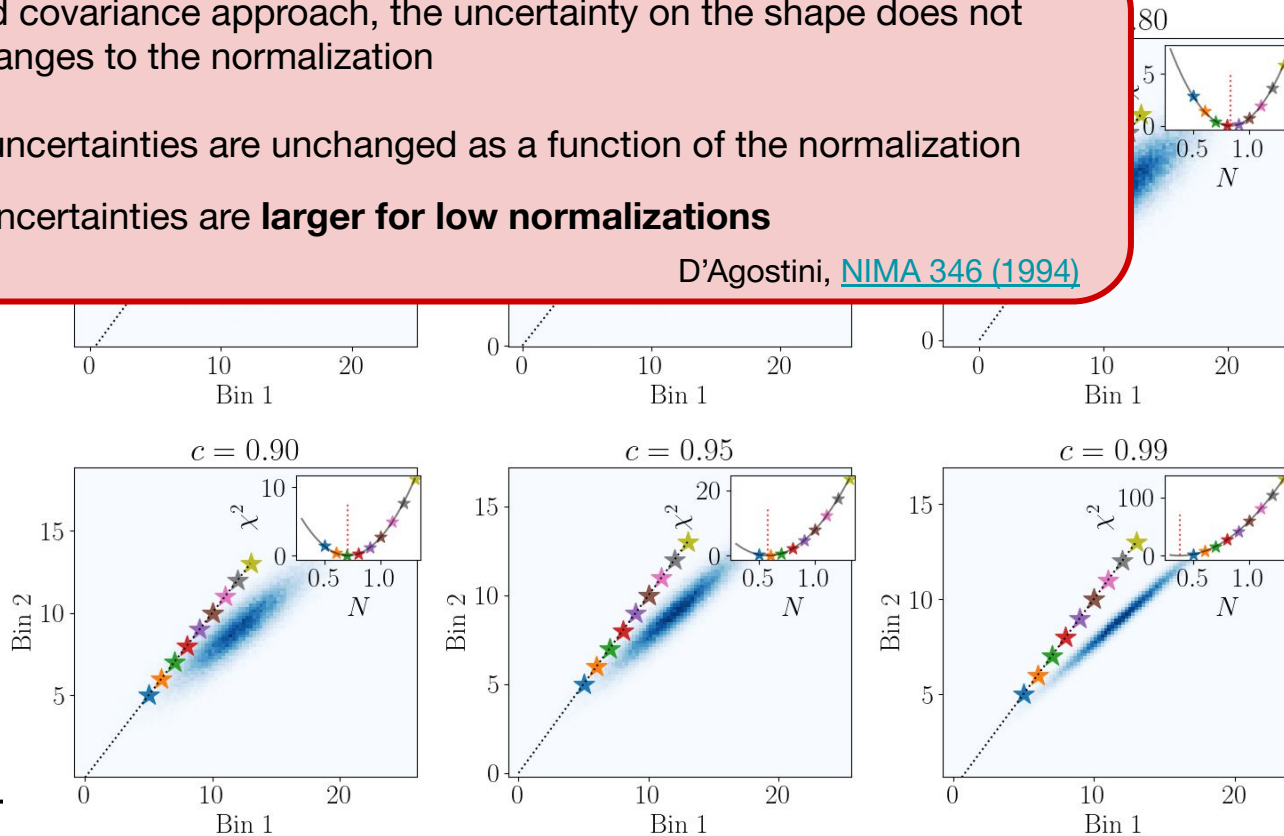
⇒ The **relative** uncertainties are **larger for low normalizations**

D'Agostini, [NIMA 346 \(1994\)](#)



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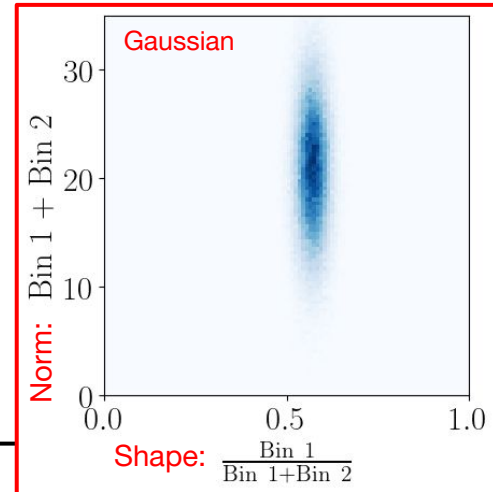
- One way to overcome this is building a covariance that keeps the **relative** uncertainties constant (instead of the absolute)
- This can be satisfied if the covariance is quoted in the “norm” and the “shape” as in:

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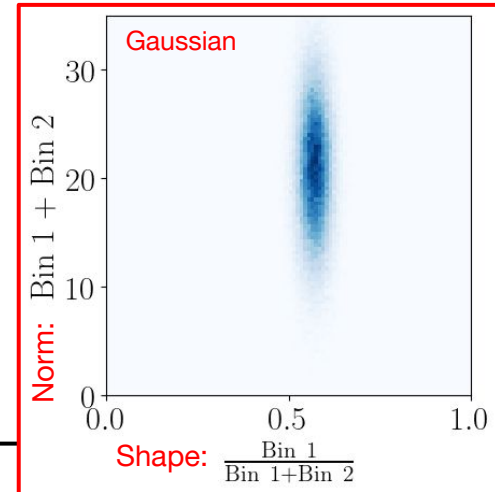
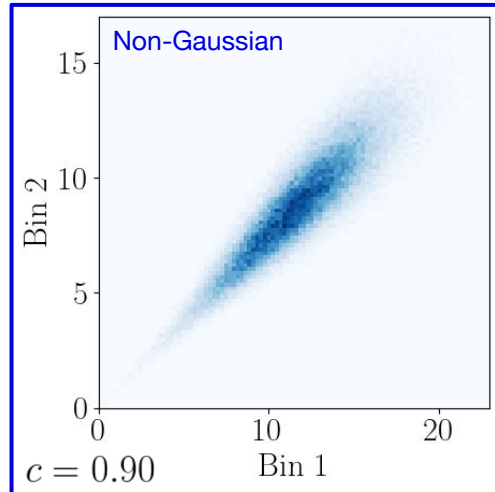
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Non-Gaussian
uncertainties

Gaussian uncertainties in the
norm/shape

Non-linear transformation

Fixed relative uncertainty

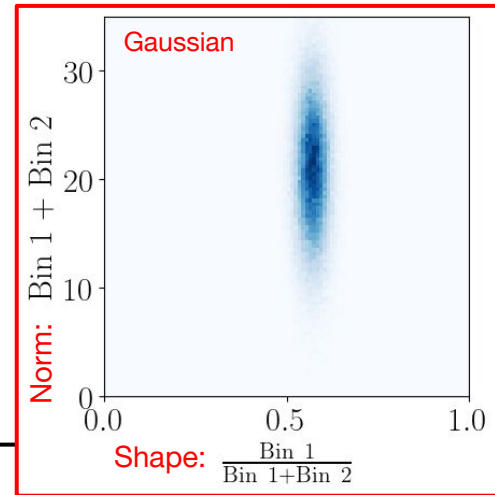
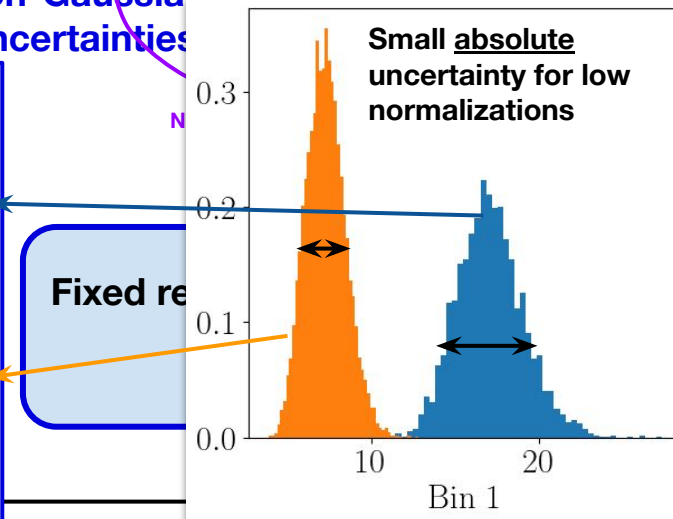
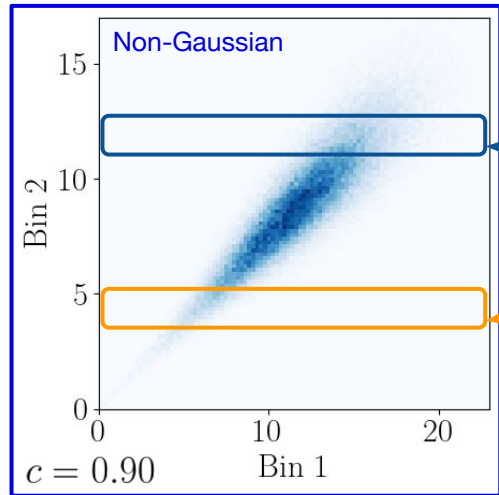


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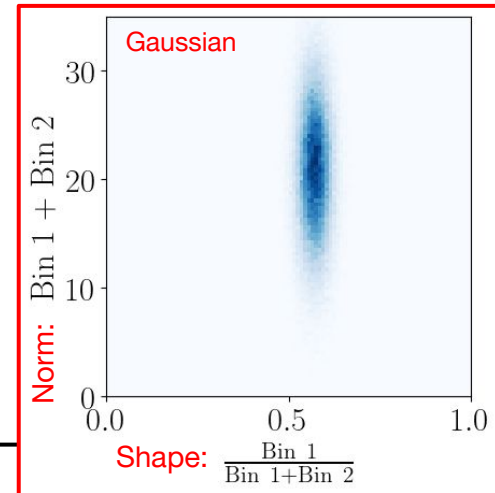
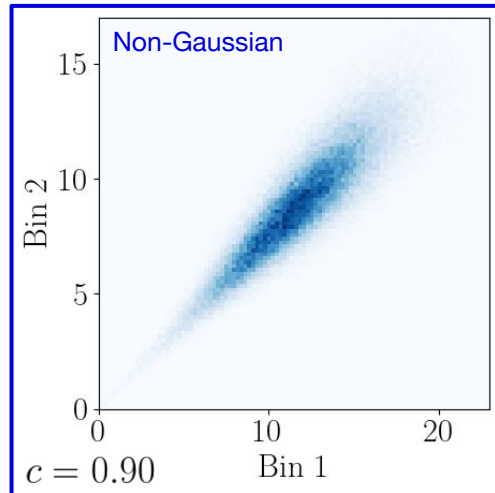
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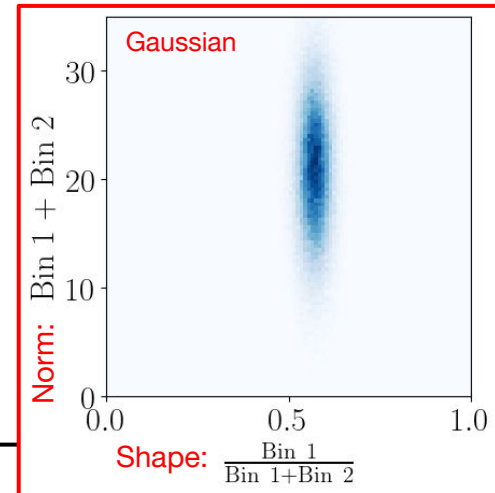
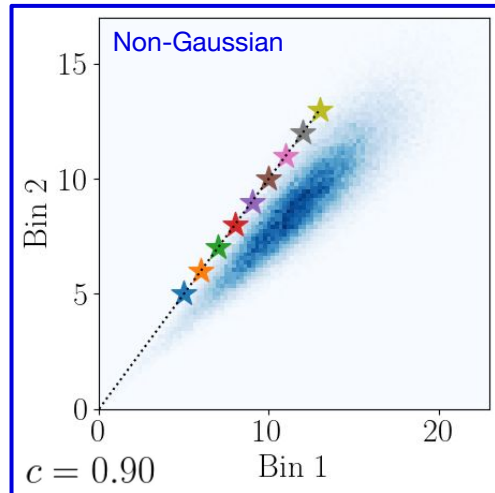
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⇒ No more preference for low norms



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- This can be generalized for n bins, with the following transformation:

[arXiv:2308.01838](https://arxiv.org/abs/2308.01838)

$$C_i = f(B_i) = \begin{cases} \alpha \frac{B_i}{\sum_k B_k}, & 1 \leq i \leq n - 1 \\ B_T = \sum_k B_k, & i = n \end{cases} \quad N \approx J(f) \times M \times J(f)^T$$

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New basis \uparrow C_i \uparrow Data release histogram B_i

New covariance $N \approx J(f) \times M \times J(f)^T$

Jacobian of the transformation $J(f)$

Data release covariance M

Linear approximation as f is non linear

- One way to overcome this is building a covariance that keeps the **relative** uncertainties constant (instead of the absolute)
- This can be satisfied if the covariance is quoted in the “norm” and the “shape” as in:

$$(\text{Bin 1, Bin 2}) \rightarrow \left(\text{Bin 1} + \text{Bin 2}, \frac{\text{Bin 1}}{\text{Bin 1} + \text{Bin 2}} \right)$$

[arXiv:2308.01838](https://arxiv.org/abs/2308.01838)

- This can be generalized for n bins, with the following transformation:

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↑
New basis

↑
Data release histogram

$$N \approx J(f) \times M \times J(f)^T$$

↑
New covariance

↑
Jacobian of the transformation

↑
Data release covariance

$$\chi_{\text{NS}}^2 = \sum_{1 \leq i, j \leq n} (C_i - C_i^{\text{MC}}) (N^{-1})_{i,j} (C_j - C_j^{\text{MC}})$$

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Can be easily used in **NUISANCE!**



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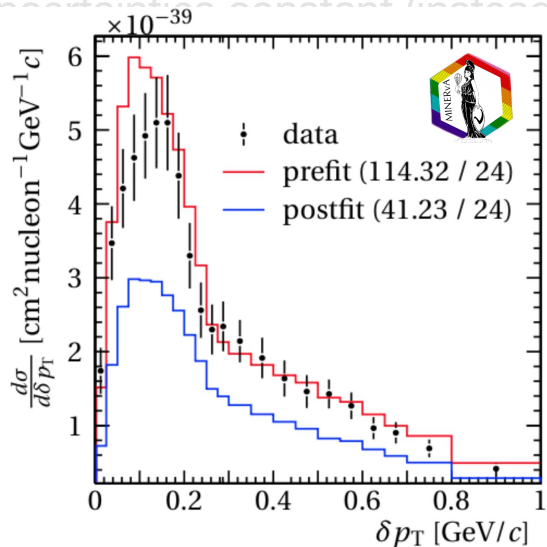
Jacobian of the transformation

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Data release covariance

Linear approximation as f is non linear

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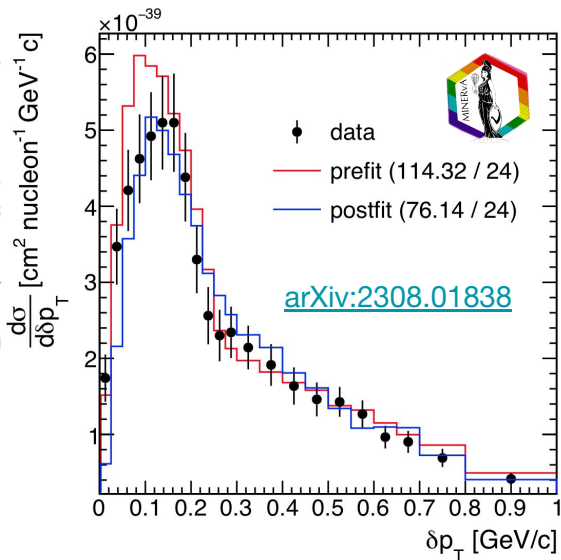


→ (Bin 1 + Bin 2)

bins, with the following

$$1 \leq i \leq n - 1$$

$$i = n$$



[arXiv:2308.01838](https://arxiv.org/abs/2308.01838)

$$C_i =$$

New basis

Data release histogram

Can be easily used in **NUISANCE!**

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Data release covariance

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Is this satisfactory?

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 - The choice of assuming Gaussian uncertainties on the norm+shape decomposition is rather arbitrary, the actual distributions should be dictated by the measurement
- Other workarounds have been considered in other studies like calculating a shape-only
 - The bin-to-bin covariance provided by experiments is typically produced by varying all the uncertainties a large number of times, and summarizing the average and the spread
 - **If these full variations (toys/universes) were to be provided in the data release, it would be possible to test these assumptions and tailor a dedicated test statistic**

Conclusion

- Fitting neutrino interaction models to cross-section measurements allows us to test and benchmark our neutrino event generators
- One of the challenges encountered in this procedure is Peelle's pertinent puzzle due to the imperfect models and the Gaussian assumptions of the errors in the data releases
- There are some workarounds that can be done to attempt to mitigate this issue using the published covariances, but...
- **An ideal data release would also contain the toys/universes used to propagate the uncertainties so that the full distributions can be accessed and used beyond the simple Gaussian assumptions**