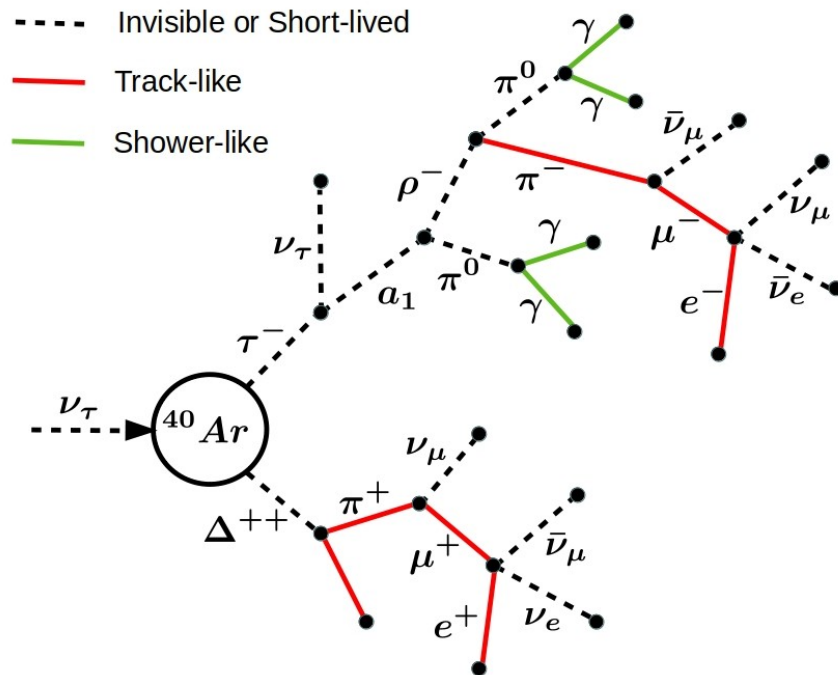
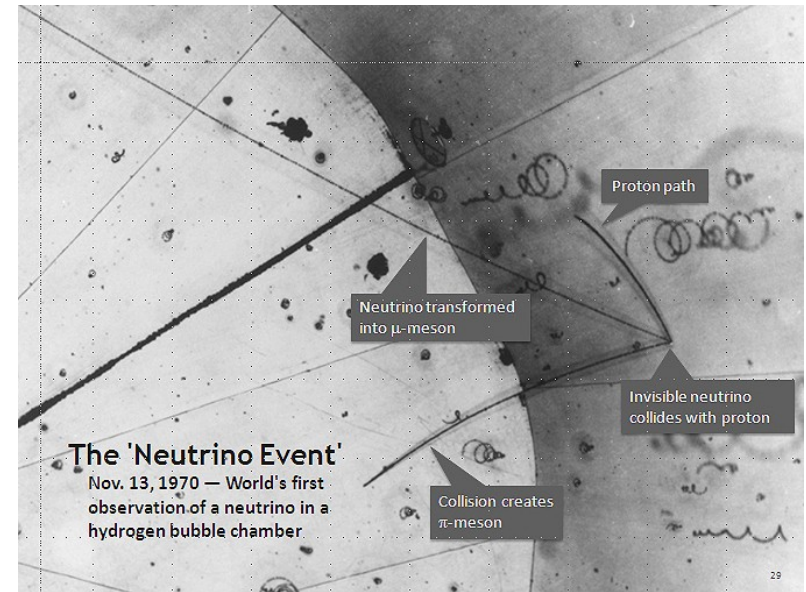


New Physics in Neutrino Oscillations



Adam Aurisano
University of Cincinnati

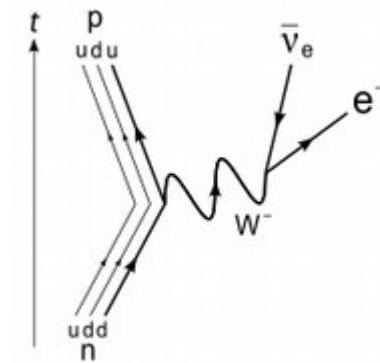
17th International Workshop
on Tau Lepton Physics
Louisville, Kentucky
6 December 2023



Neutrinos in the Standard Model

	<p>mass → =2.3 MeV/c²</p> <p>charge → 2/3</p> <p>spin → 1/2</p> <p>u</p> <p>up</p>	<p>mass → =1,275 GeV/c²</p> <p>charge → 2/3</p> <p>spin → 1/2</p> <p>c</p> <p>charm</p>	<p>mass → =173.07 GeV/c²</p> <p>charge → 2/3</p> <p>spin → 1/2</p> <p>t</p> <p>top</p>	<p>mass → 0</p> <p>charge → 0</p> <p>spin → 1</p> <p>g</p> <p>gluon</p>	<p>mass → =126 GeV/c²</p> <p>charge → 0</p> <p>spin → 0</p> <p>H</p> <p>Higgs boson</p>	
QUARKS	<p>mass → =4.8 MeV/c²</p> <p>charge → -1/3</p> <p>spin → 1/2</p> <p>d</p> <p>down</p>	<p>mass → =95 MeV/c²</p> <p>charge → -1/3</p> <p>spin → 1/2</p> <p>s</p> <p>strange</p>	<p>mass → =4.18 GeV/c²</p> <p>charge → -1/3</p> <p>spin → 1/2</p> <p>b</p> <p>bottom</p>	<p>mass → 0</p> <p>charge → 0</p> <p>spin → 1</p> <p>γ</p> <p>photon</p>		
	<p>mass → 0.511 MeV/c²</p> <p>charge → -1</p> <p>spin → 1/2</p> <p>e</p> <p>electron</p>	<p>mass → 105.7 MeV/c²</p> <p>charge → -1</p> <p>spin → 1/2</p> <p>μ</p> <p>muon</p>	<p>mass → 1,777 GeV/c²</p> <p>charge → -1</p> <p>spin → 1/2</p> <p>τ</p> <p>tau</p>	<p>mass → 91.2 GeV/c²</p> <p>charge → 0</p> <p>spin → 1</p> <p>Z</p> <p>Z boson</p>	GAUGE BOSONS	
	LEPTONS	<p>mass → <2.2 eV/c²</p> <p>charge → 0</p> <p>spin → 1/2</p> <p>ν_e</p> <p>electron neutrino</p>	<p>mass → <0.17 MeV/c²</p> <p>charge → 0</p> <p>spin → 1/2</p> <p>ν_μ</p> <p>muon neutrino</p>	<p>mass → <15.5 MeV/c²</p> <p>charge → 0</p> <p>spin → 1/2</p> <p>ν_τ</p> <p>tau neutrino</p>		<p>mass → 80.4 GeV/c²</p> <p>charge → ±1</p> <p>spin → 1</p> <p>W</p> <p>W boson</p>

1930: Wolfgang Pauli proposed neutrinos to make β decays conserve energy



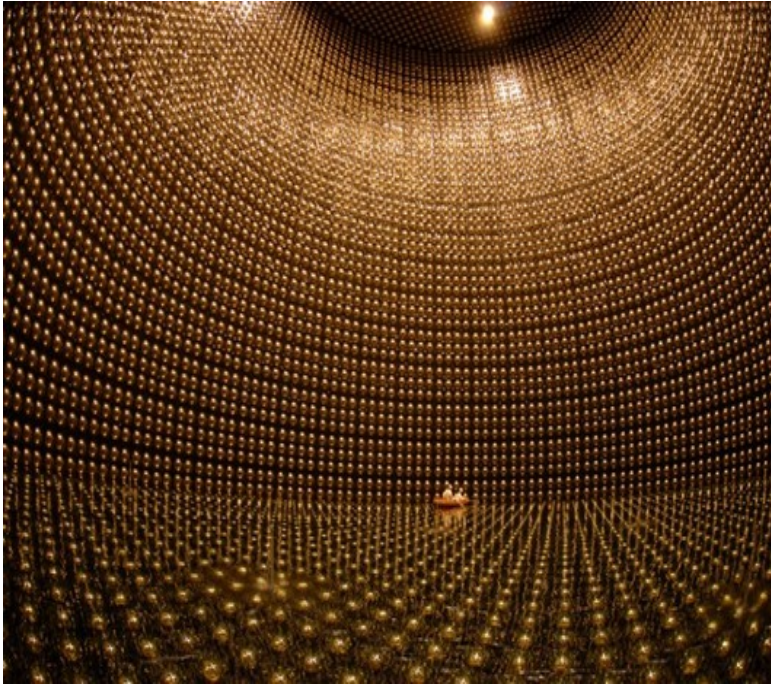
1956: Cowan and Reines discovered the ν_e

1962: Lederman, Schwartz, and Steinberger discovered the ν_μ

2000: DONUT collaboration discovered the ν_τ

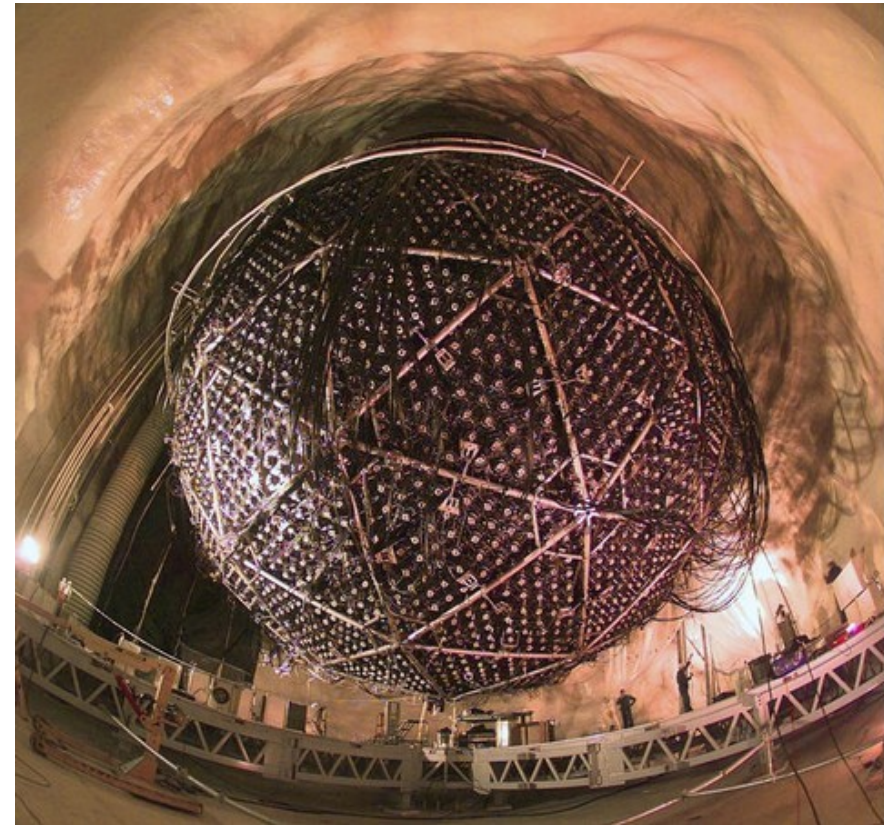
Neutrinos only feel the weak force
In the SM neutrinos are massless

Oscillations



- Super-Kamiokande
 - Looked for ν_{μ} produced in the atmosphere by cosmic rays
 - Deficit of ν_{μ} going up through the earth compared to those coming down from the sky.

- SNO
 - Showed that ν_e from the sun less than expected
 - Total number of neutrinos agreed with expectations

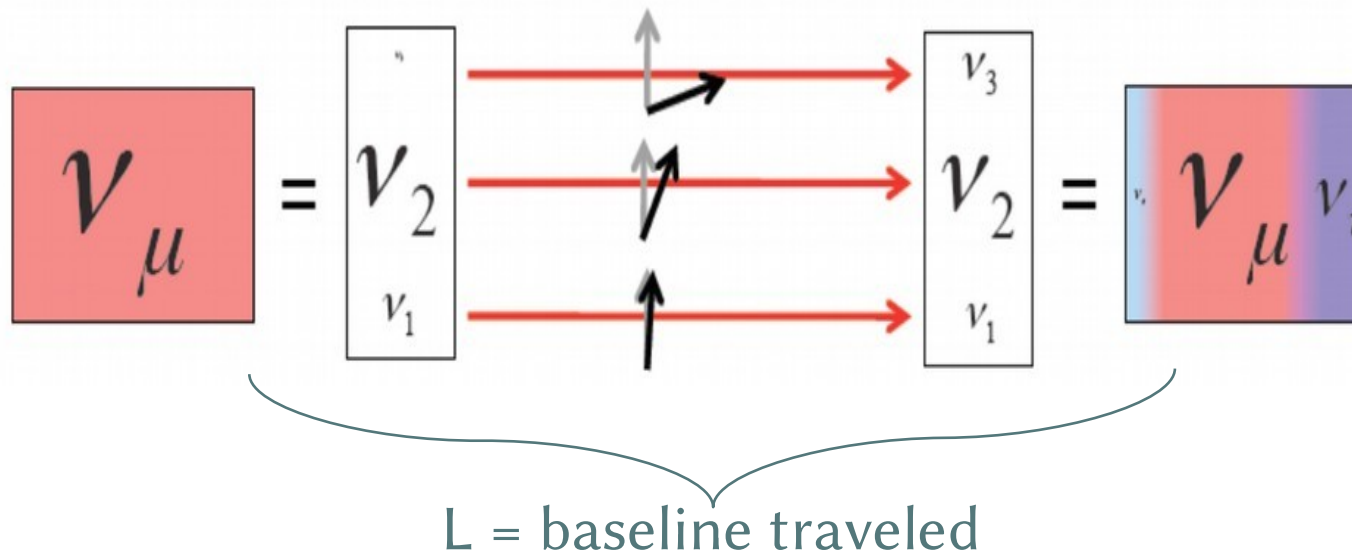


Neutrinos Have Mass

Neutrinos are produced with a definite flavor

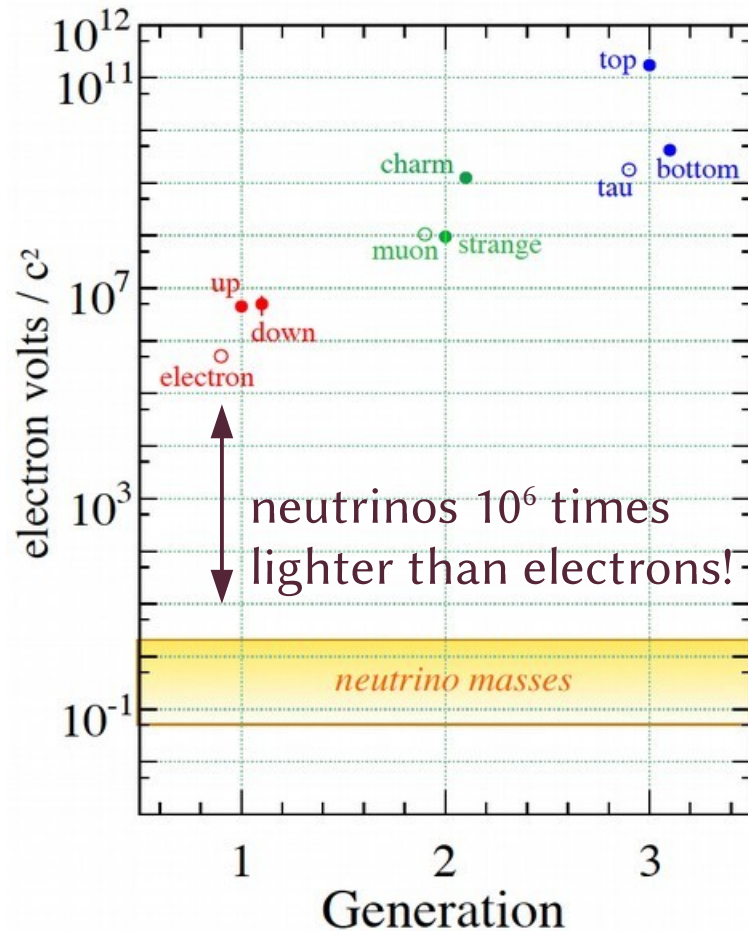
They travel as a superposition of mass states

They interact with a definite flavor – maybe different than its starting flavor

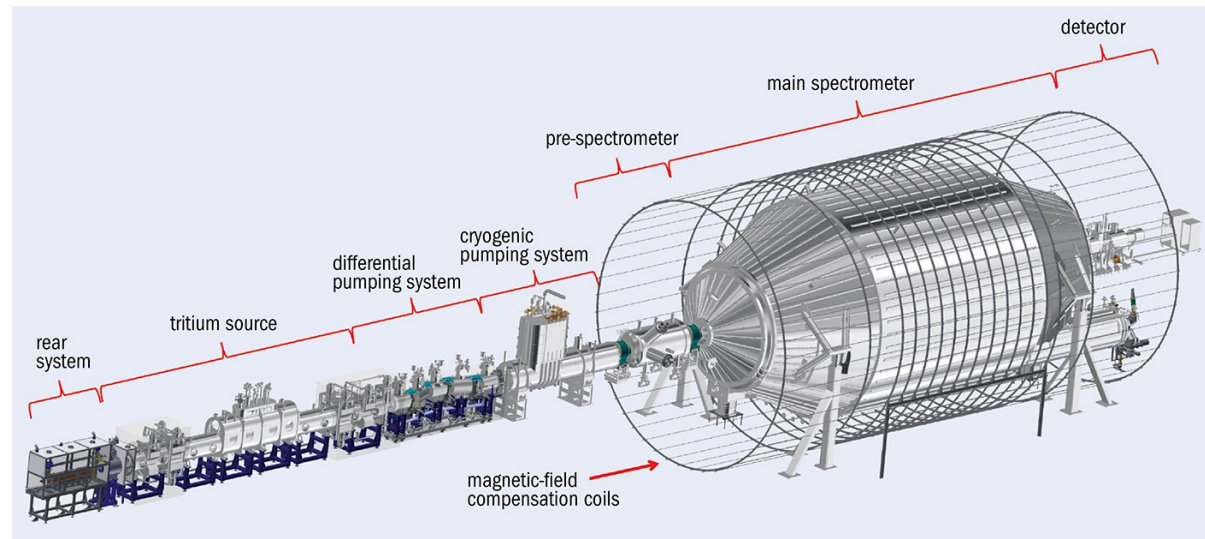


Neutrinos oscillations are only sensitive $\Delta m^2 = m_j^2 - m_i^2$
Not absolute masses!

Neutrino Masses are Very Small



- Best direct mass measurement from KATRIN experiment
- Upper limit set so far
 - $m_\nu < 0.8 \text{ eV}$
 - Nat. Phys. 18 (2022) 160-166
- Neutrinos are at least six orders of magnitude lighter than electrons!



CERN Courier

Oscillations are New Physics

Dirac:

- All SM fermions get this type of mass
 - Generated by Higgs mechanism
 - Requires right-handed neutrinos
-



- Small masses \rightarrow very small Yukawa couplings
- Hard to naturally explain difference in scale between neutrinos and everything else
- Right-handed neutrinos couple to nothing except gravity \rightarrow “sterile neutrinos”

Majorana:

- Couples particle to its antiparticle
- No fundamental particles are known to have this type of mass



- Higher dimension operators
- Introduces lepton number violation
- Implies scale of breakdown of SM

See-Saw Mechanism

- Add right-handed neutrinos and allow both types of masses
- See-saw mechanism allows for natural small neutrino masses.
 - Dirac mass \sim Electro-weak scale
 - Majorana mass \sim GUT scale
- Diagonalize mass matrix
 - One mostly active light neutrino (ν) and one mostly sterile heavy neutrino (N)

$$\mathcal{L}_{mass} = \begin{bmatrix} \nu_L & \nu_R \end{bmatrix} \begin{bmatrix} 0 & m_D \\ m_D & M \end{bmatrix} \begin{bmatrix} \nu_L \\ \nu_R \end{bmatrix}$$

$$M \gg m_D$$

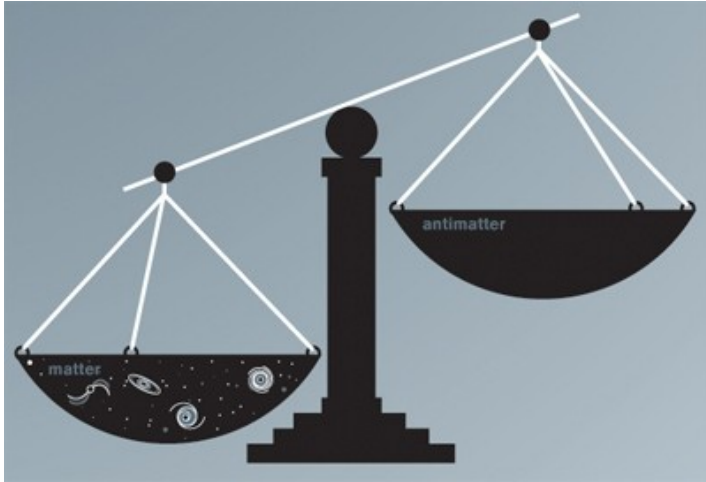
$$m_N \approx M \quad m_\nu = \frac{m_D^2}{M}$$

Yamagida; Gell-Mann, Ramond, Slansky
Minkowski; Mohapatra, Senjanovic

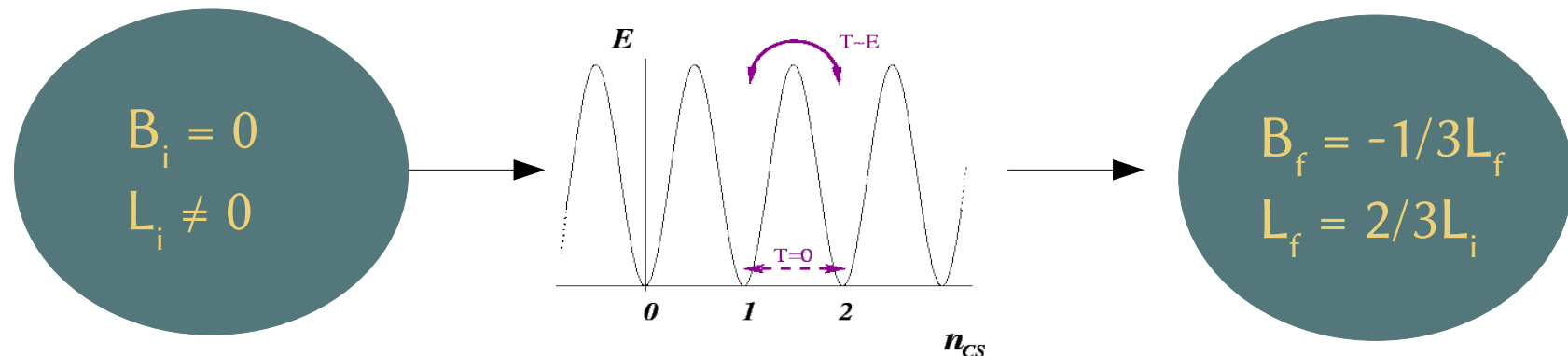


Matter-Antimatter Asymmetry

Symmetry



- Immediately after the big bang, expect equal amounts of matter and antimatter
- Now, matter dominates the universe
- Sakharov conditions
 - Baryon number violation
 - C and CP violation
 - Out of thermal equilibrium decays
- CP violation in quark sector not enough
- See-saw mechanism enables leptogenesis

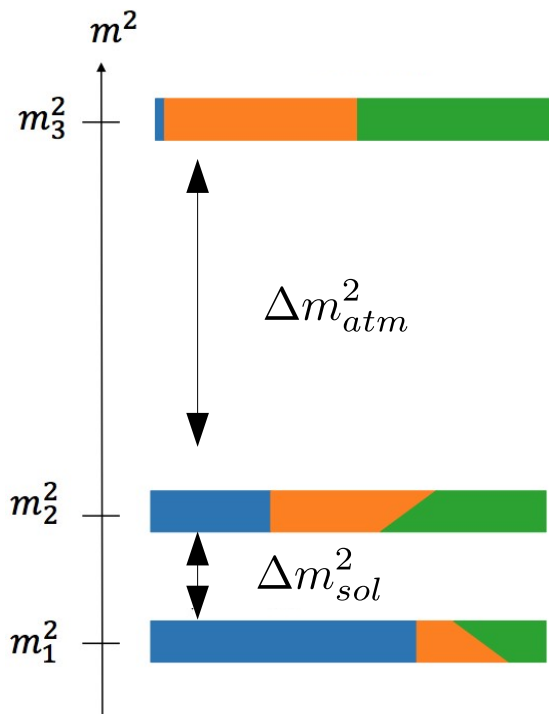


Three-Flavor Oscillations

$$s_{ij} = \sin \theta_{ij} \quad c_{ij} = \cos \theta_{ij}$$

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}$$

$$P(\nu_\alpha \rightarrow \nu_\beta) = \left| \sum_j U_{\alpha j} U_{\beta j}^* e^{\frac{-i\Delta m_{j1}^2 L}{2E}} \right|^2$$



- 3-flavor oscillations parameterized by:

- Three mixing angles ($\theta_{12}, \theta_{13}, \theta_{23}$)

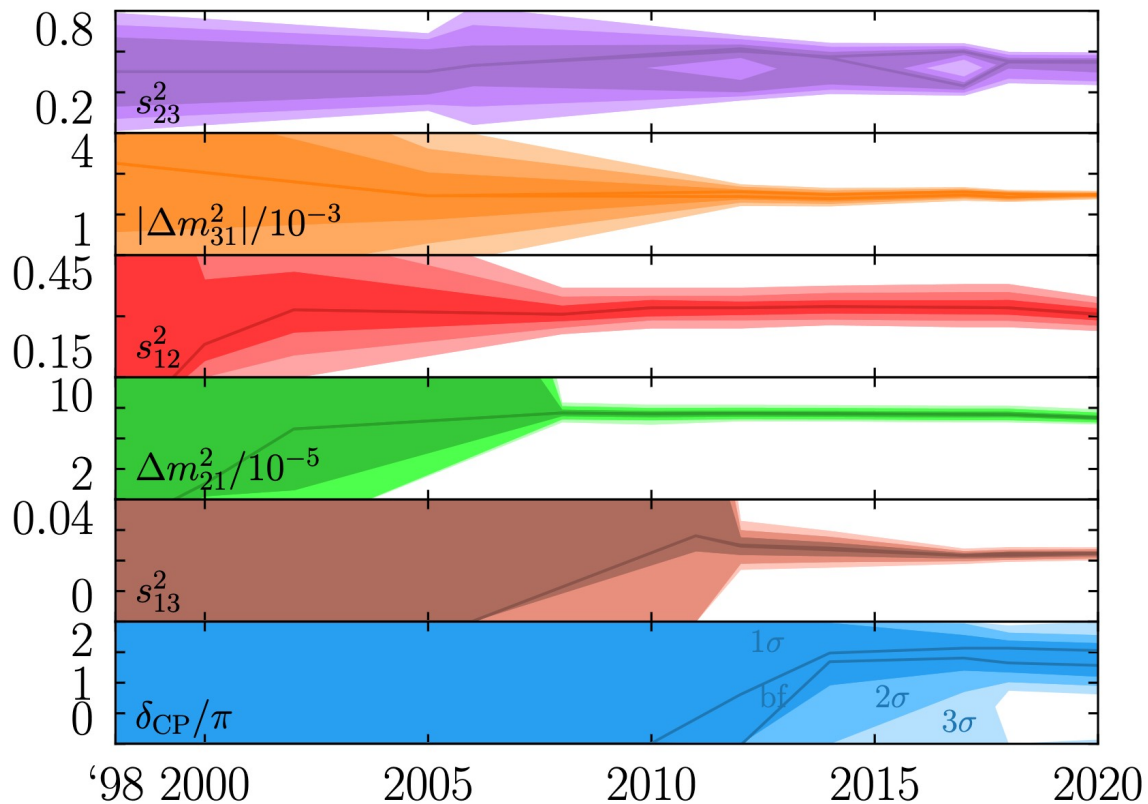
- One CP-violating phase (δ)

- Two mass splittings (Δm_{sol}^2 and Δm_{atm}^2)

What We Know So Far

$$s_{ij} = \sin \theta_{ij} \quad c_{ij} = \cos \theta_{ij}$$

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}$$

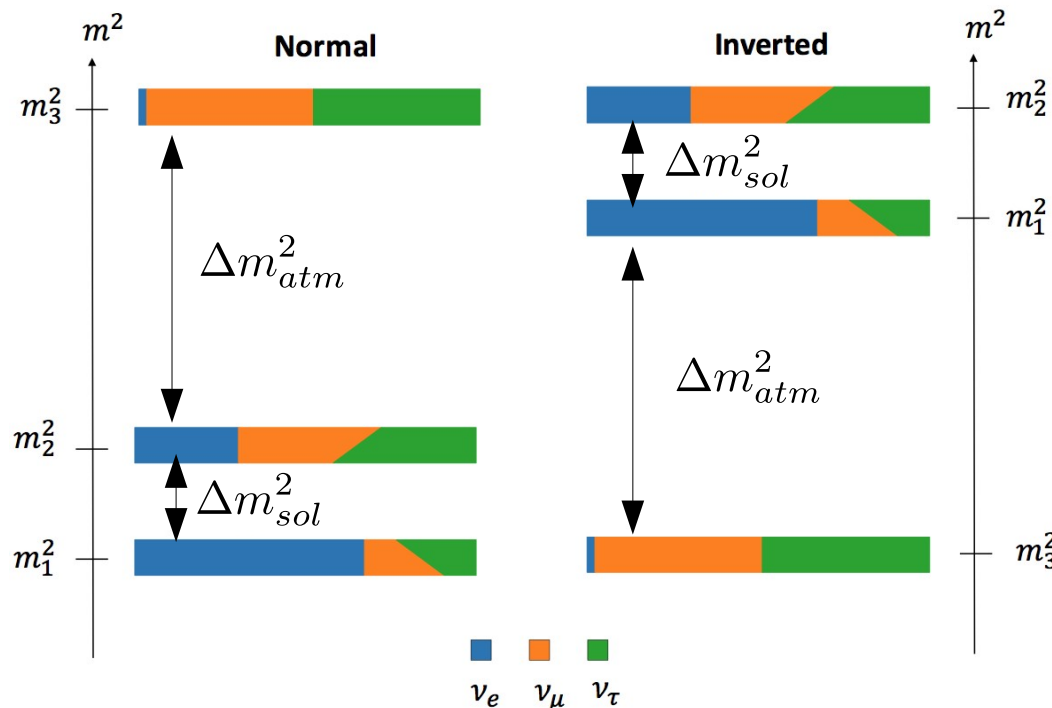


Denton et al. arXiv:2212.00809

Open Questions in Neutrino Oscillations

$$s_{ij} = \sin \theta_{ij} \quad c_{ij} = \cos \theta_{ij}$$

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}$$

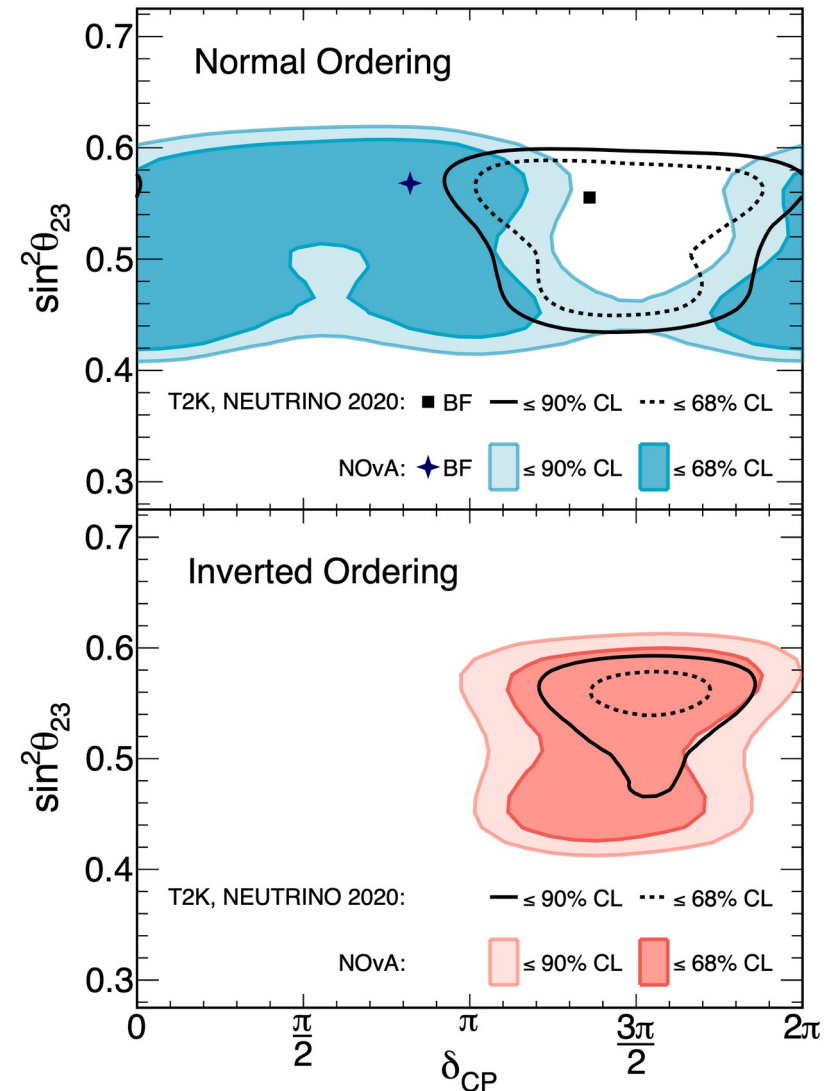


- Is θ_{23} exactly 45° ?
- Is ν_3 the heaviest or lightest neutrino?
- What is the value of $\sin \delta$?
- Is this the whole picture?

Non-Standard Interactions

NOvA and T2K

- NOvA and T2K both measure ν_e appearance and ν_μ disappearance
 - Designed to study the mass ordering and CP violation
- Both experiments are designed to have similar L/E
 - NOvA: $E \sim 2$ GeV, $L = 810$ km
 - T2K: $E \sim 600$ MeV, $L = 295$ km
- Recent results broadly compatible, but with qualitative differences
- If tension were to increase in future, what could it mean?



Non-Standard Interactions

Non-standard interactions (NSI) modify standard three flavor neutrino oscillations by introducing anomalous interactions between neutrinos and matter, in addition to the standard MSW effect

$$\mathcal{H} = \frac{1}{2E} \left[U_{\text{PMNS}} \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta_{21}^2 & 0 \\ 0 & 0 & \Delta_{31}^2 \end{pmatrix} U_{\text{PMNS}}^\dagger + a \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix} \right]$$

$$a = 2\sqrt{2}G_F N_e E$$

$$\epsilon_{\alpha\beta} = |\epsilon_{\alpha\beta}| e^{i\delta_{\alpha\beta}}$$

Wolfenstein matter potential

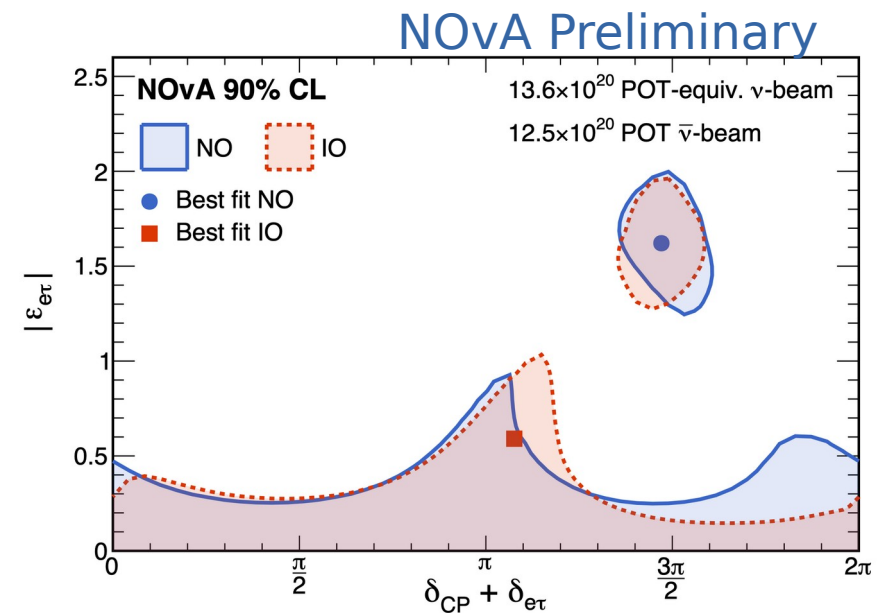
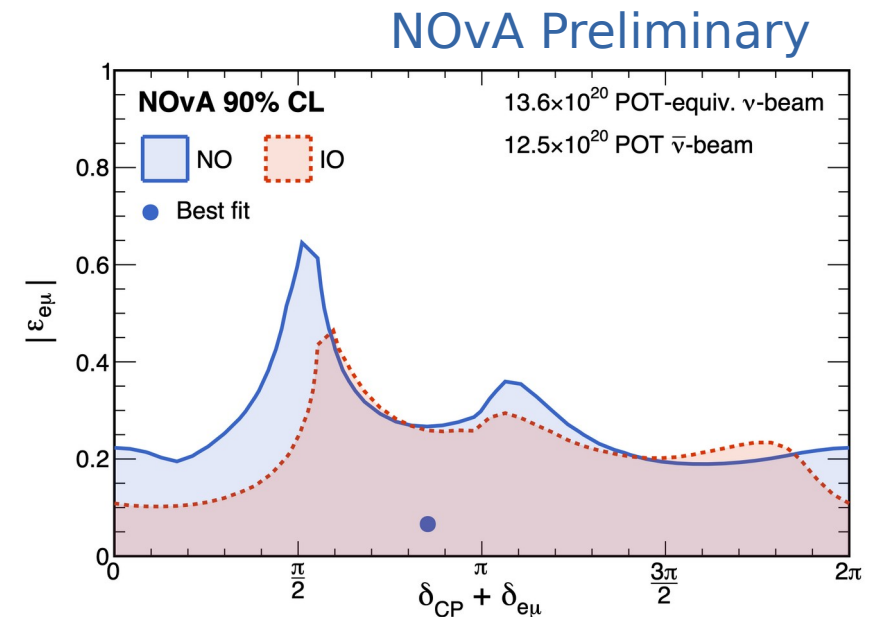
Introduces 9 new parameters:

On-diagonal: NSI-induced mass squared splittings (real valued)

Off-diagonal: NSI-induced mixing angles (complex)

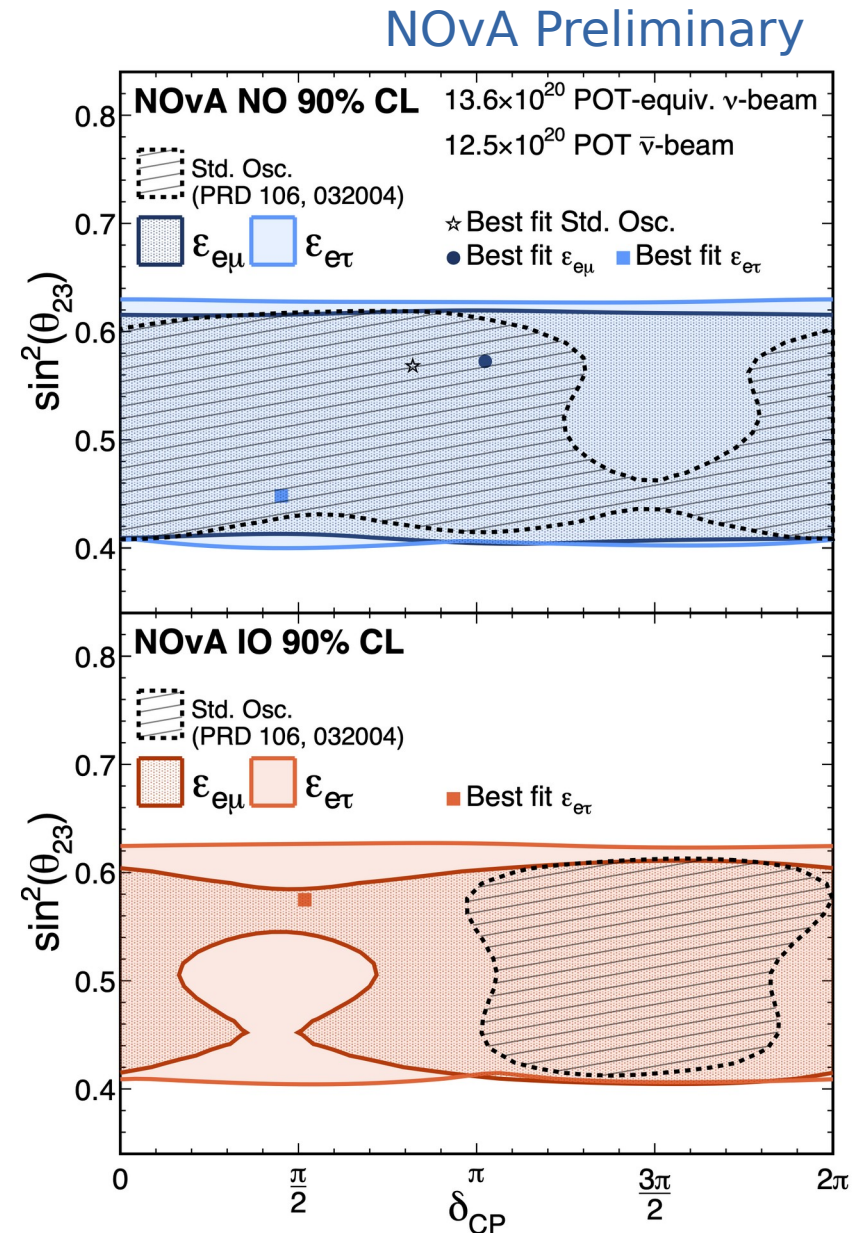
NOvA Search for NSI

- NSI, in general, has many parameters
- NOvA investigated effect of $\varepsilon_{e\mu}$ and $\varepsilon_{e\tau}$ terms
- Prefers standard oscillations
- Strength of result would be improved through combining results from experiments with different matter effects



Effect on Three Flavor Measurements

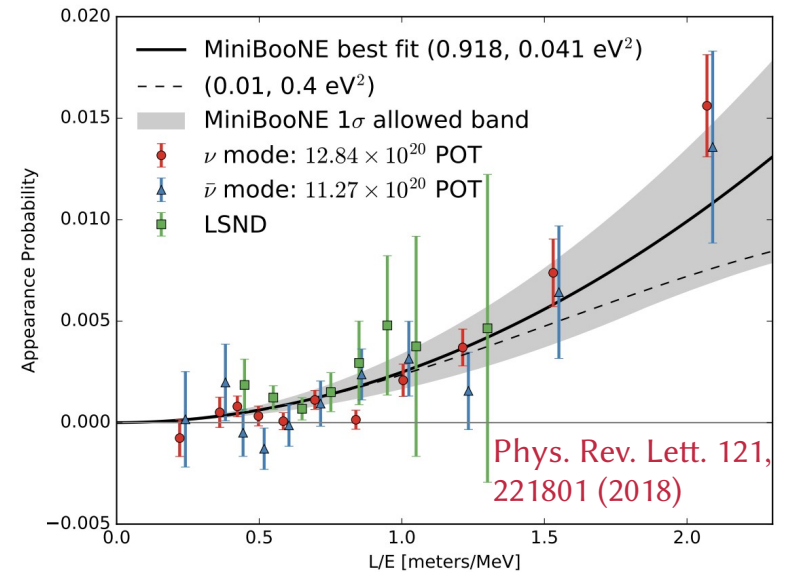
- Introducing NSI into three flavor oscillations profoundly affects sensitivity to CP violation
- Much sensitivity to δ_{CP} in the three flavor interpretation is lost when introducing NSI, due to degeneracy with CP-violating phases of oscillation parameters



Light Sterile Neutrinos

Light Sterile Neutrinos

- Short-baseline experiments (LSND, MiniBooNE) observed anomalous excesses of ν_e ($\bar{\nu}_e$) in ν_μ ($\bar{\nu}_\mu$) beams
- BEST observed anomalous deficit in ν_e from a ^{51}Cr neutrino source over short baselines
- Anomalies could all be explained by oscillations driven by a mass splitting $\Delta m^2 \sim 1 \text{ eV}^2$
 - Not consistent with three known active flavors
- Simplest model adds one new mass state and one new, sterile flavor state
 - 3+1 model contains 6 new parameters



m_4

m_3

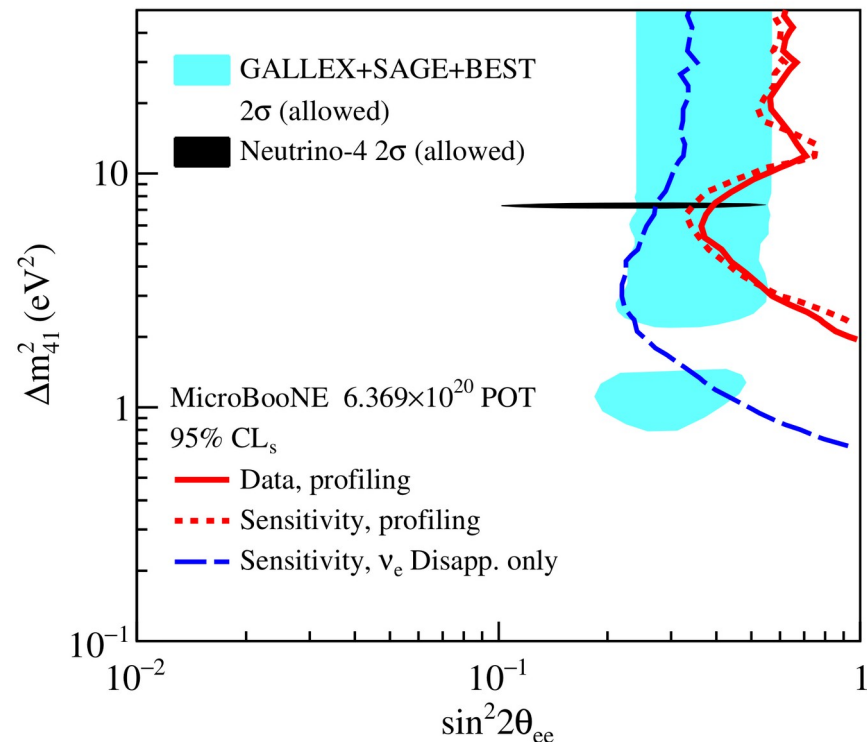
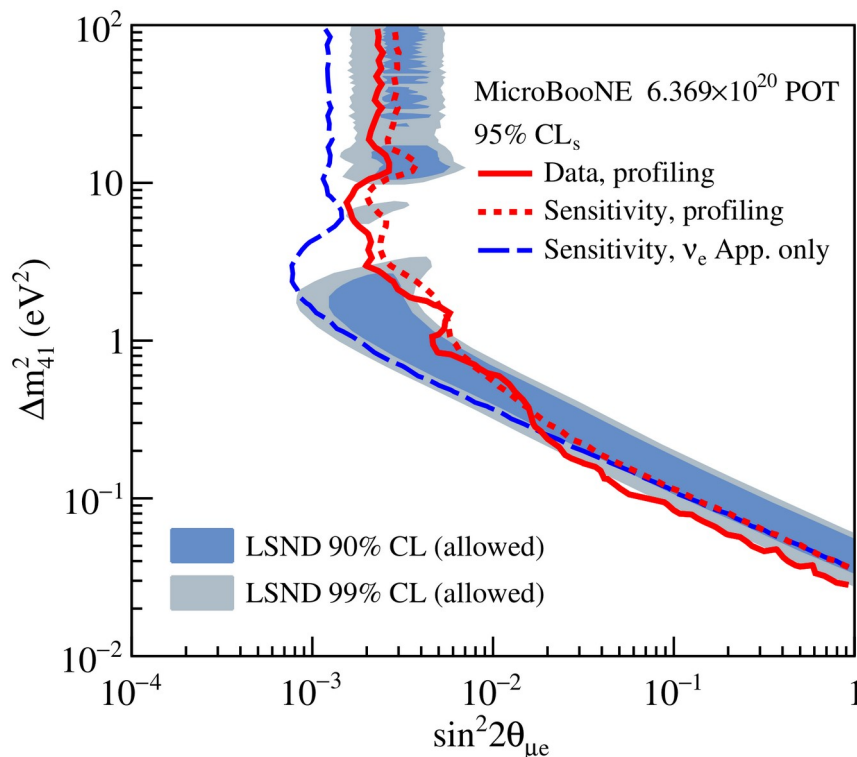
ν_s
 ν_τ
 ν_μ
 ν_e

m_2

m_1

MicroBooNE Results

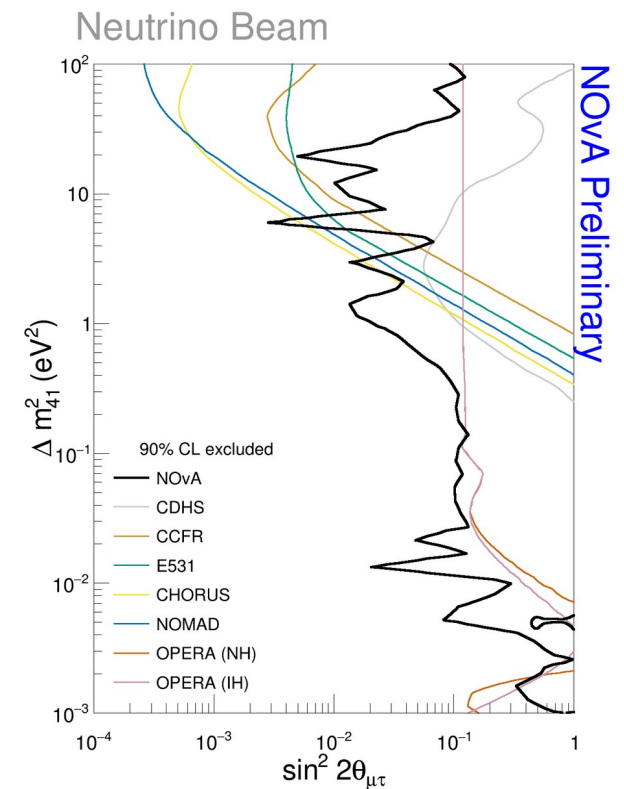
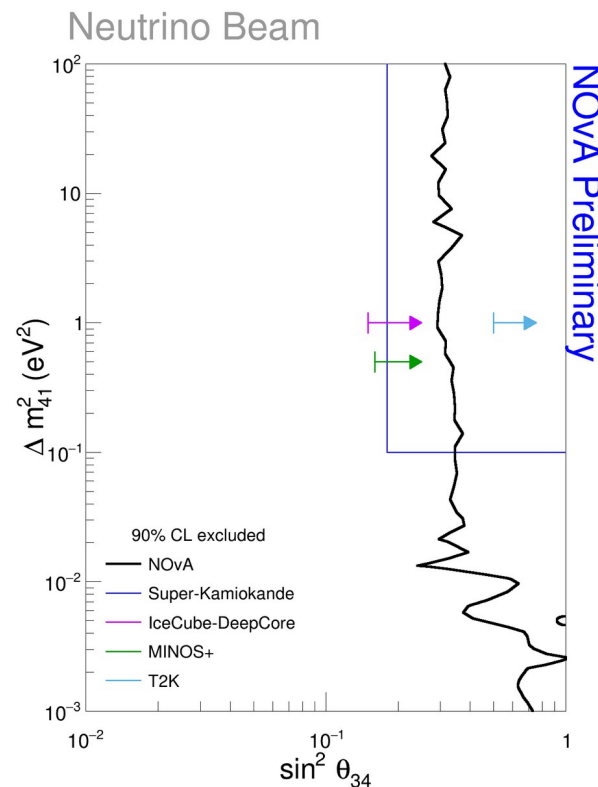
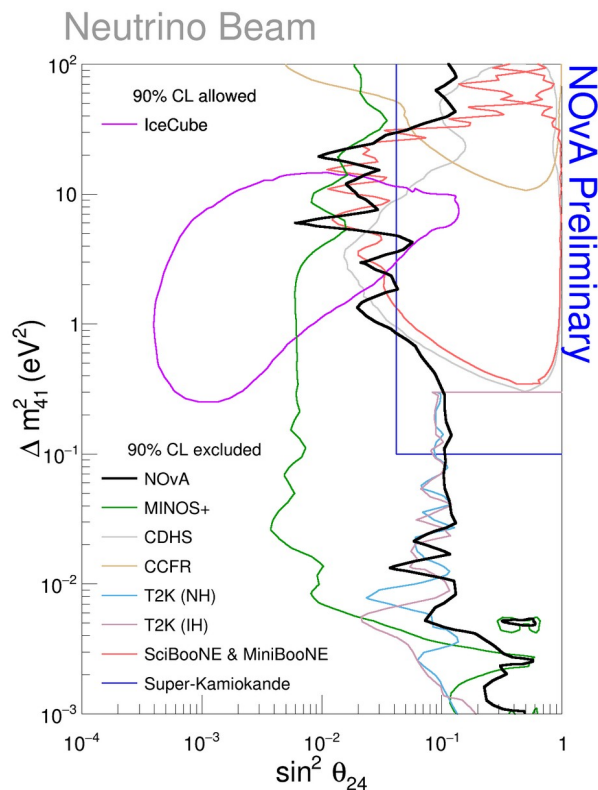
- MicroBooNE recently searched for sterile-driven oscillations at short baselines in the ν_e channel
- Considered oscillations involving $\nu_\mu \rightarrow \nu_e$ and $\nu_e \rightarrow \nu_s$
- Set limits on effective mixing parameters
 - Encroaching on allowed regions from anomalous results



PRL 130 (2023) 011801

NOvA Results

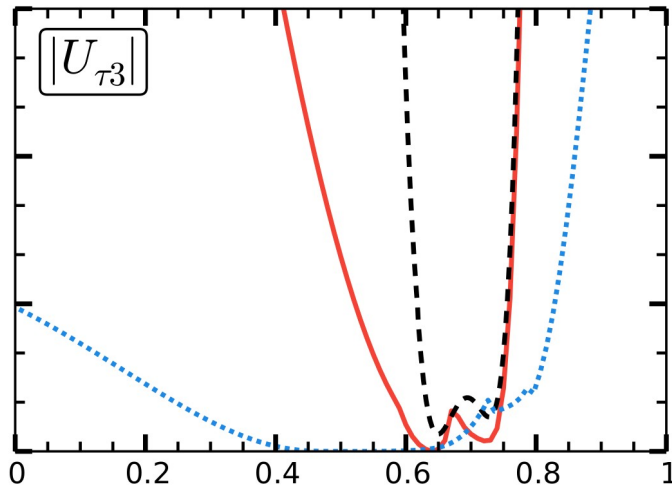
- NOvA recently searched for sterile-driven oscillations using a joint short and long baseline analysis
- Considered both ν_μ disappearance and NC disappearance
- Set limits on $\sin^2\theta_{24}$ and $\sin^2\theta_{34}$
 - Limits can be combined to set limits on effective mixing parameter controlling anomalous $\nu_\mu \rightarrow \nu_\tau$ appearance



Non-Unitary Mixing

Unitarity and Tau Neutrinos

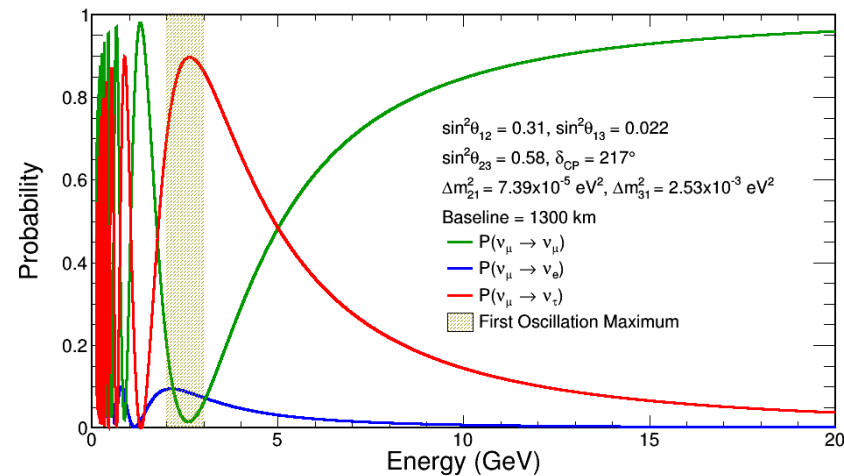
PRD 93, 1103009 (2016)



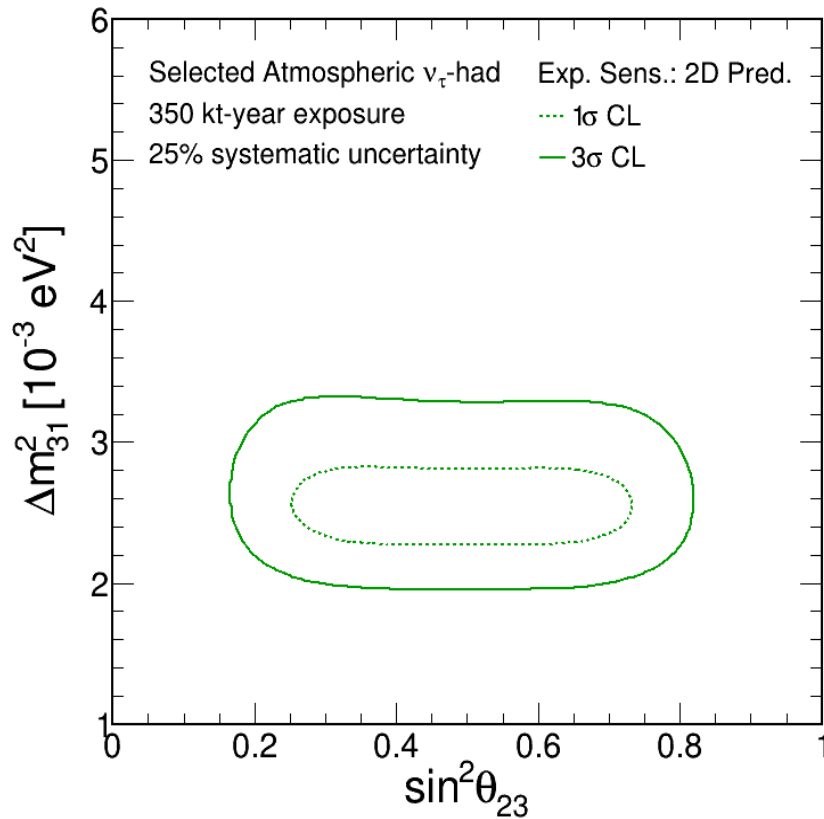
- Current experiments NOvA and T2K are making huge strides in understanding standard neutrino oscillations
- Is this model complete?
- Almost all knowledge of ν_τ sector is based on assumption of PMNS unitarity
- Almost all ν_μ disappear at oscillation maximum

- Assumed oscillating into ν_τ
- Only 10 high-purity, oscillated, ν_τ candidates have ever been observed

- DUNE will be the only experiment able to directly observe tau neutrinos appearing
 - High statistics samples both from beam and atmospheric



Atmospheric Parameters with Tau Neutrinos

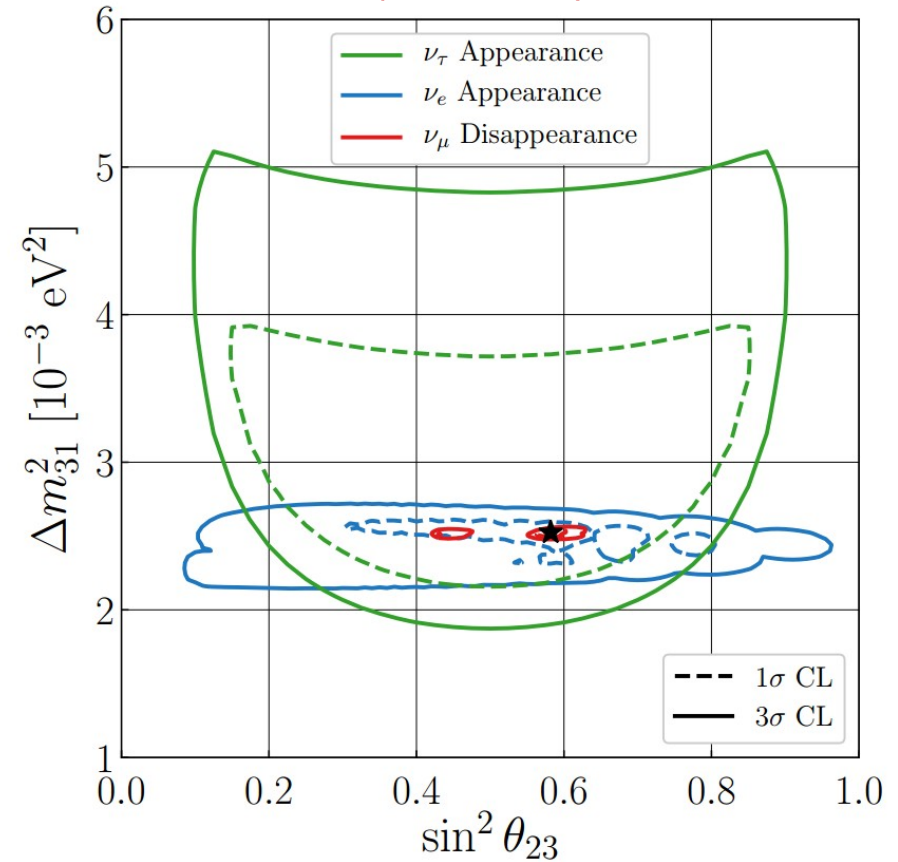


Atmospheric sample - expected counts

~1 ν_τ per kton-year

~40 ν_τ per year, with all modules

de Gouvea, Kelly, Stenico, Pasquini, PRD 100, 016004 (2019)



Beam sample - expected counts/year:

~130 ν_τ in low-energy neutrino mode

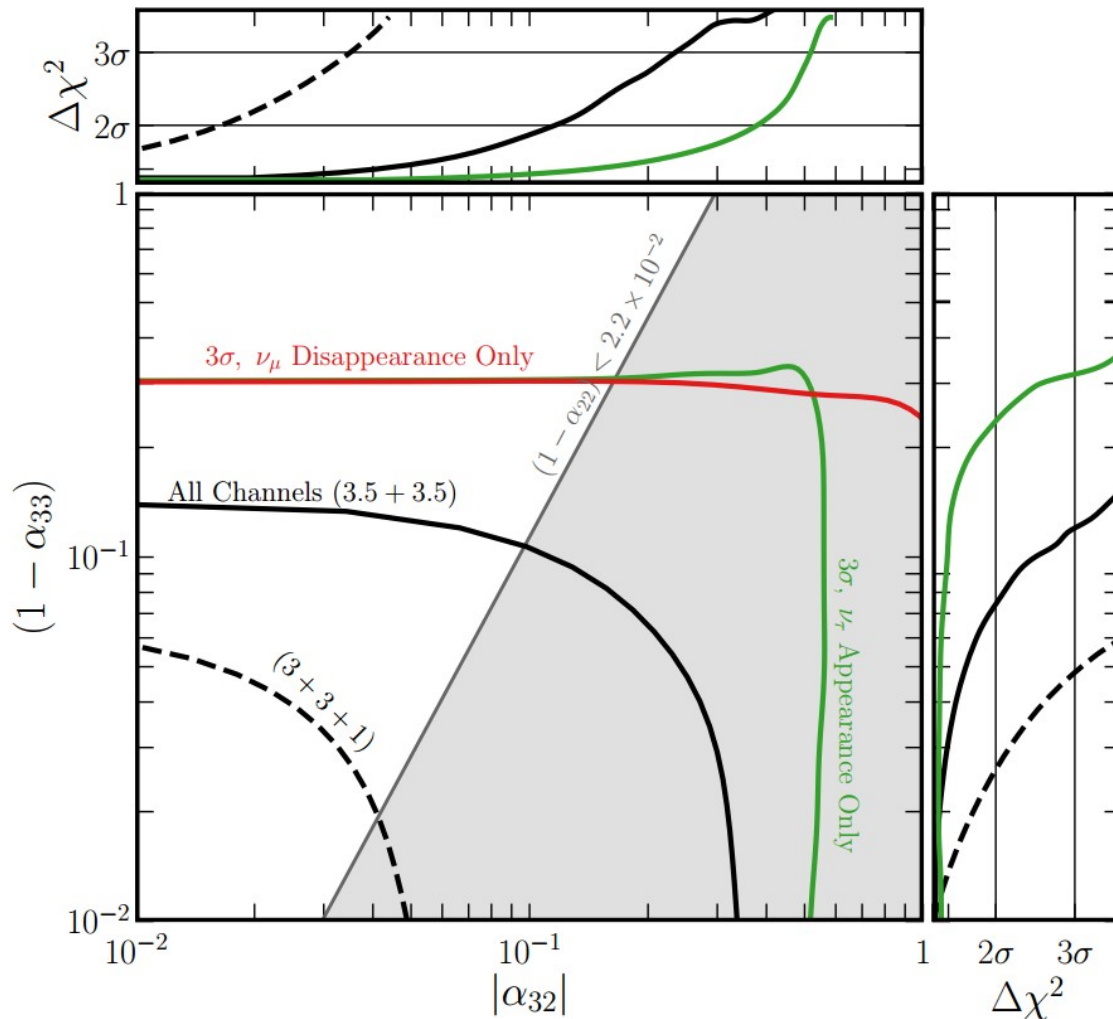
~30 $\bar{\nu}_\tau$ in low-energy antineutrino mode

~800 ν_τ in high-energy neutrino mode

Non-Unitary Oscillations

- In type I see-saw models, number of neutrino states goes from $3 \rightarrow n$
- Three states correspond to observed light neutrinos, rest are heavy sterile neutrinos
- Unitary PMNS matrix becomes: $U^{n \times n} = \begin{pmatrix} N & S \\ V & T \end{pmatrix}$
- N is the 3x3 matrix describing mixing between the light states
 - Observed PMNS matrix
- Apparent non-unitarity of 3x3 submatrix provides a way to detect existence of sterile states too heavy to directly access

Constraining Parameterized Non-Unitarity



de Gouvea, Kelly, Stenico, Pasquini, PRD 100, 016004 (2019)

$$U \rightarrow NU = \begin{pmatrix} \alpha_{11} & 0 & 0 \\ \alpha_{21} & \alpha_{22} & 0 \\ \alpha_{31} & \alpha_{32} & \alpha_{33} \end{pmatrix} U$$

- Can constrain non-unitarity using α parameters
- Tau neutrino data, in addition to other channels, improves bounds on α_{33}
- Only considers DUNE beam samples
- Alternate high energy beam mode is particularly powerful

Summary

- Even standard, three-flavor neutrino oscillations are beyond the Standard Model physics
- Mass generation may be related to physics at very high scales
- Current generation of neutrino experiments are making good progress at measuring standard mixing parameters
- Next generation of experiments will definitely measure standard mixing parameters
- It is critical that we broaden our focus to try understand the consistency between measurements from a broad range of baselines, energies, and sources to determine if three flavor mixing fully describes the data