

LHCb Measurements on Semileptonic Decays of b-hadrons

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on behalf of the LHCb Collaboration

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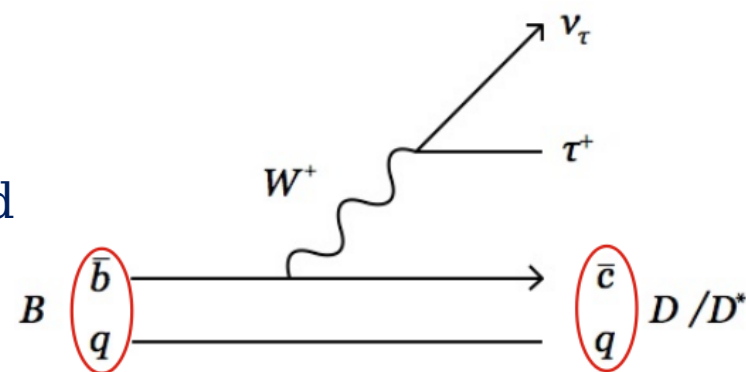


- Tree level quark transition with W emission

- **Advantages:**

- The contributes to decay rate can be factorized in **weak** and **strong** part

$$\frac{d\Gamma(B \rightarrow Xl\nu)}{dq^2} \propto G_F^2 |V_{bq}|^2 f(q^2)^2$$

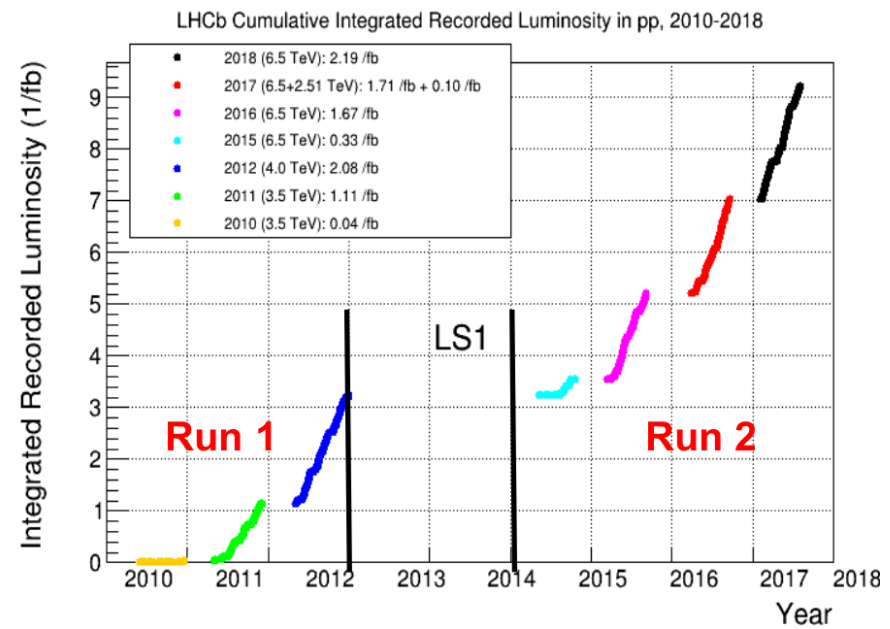
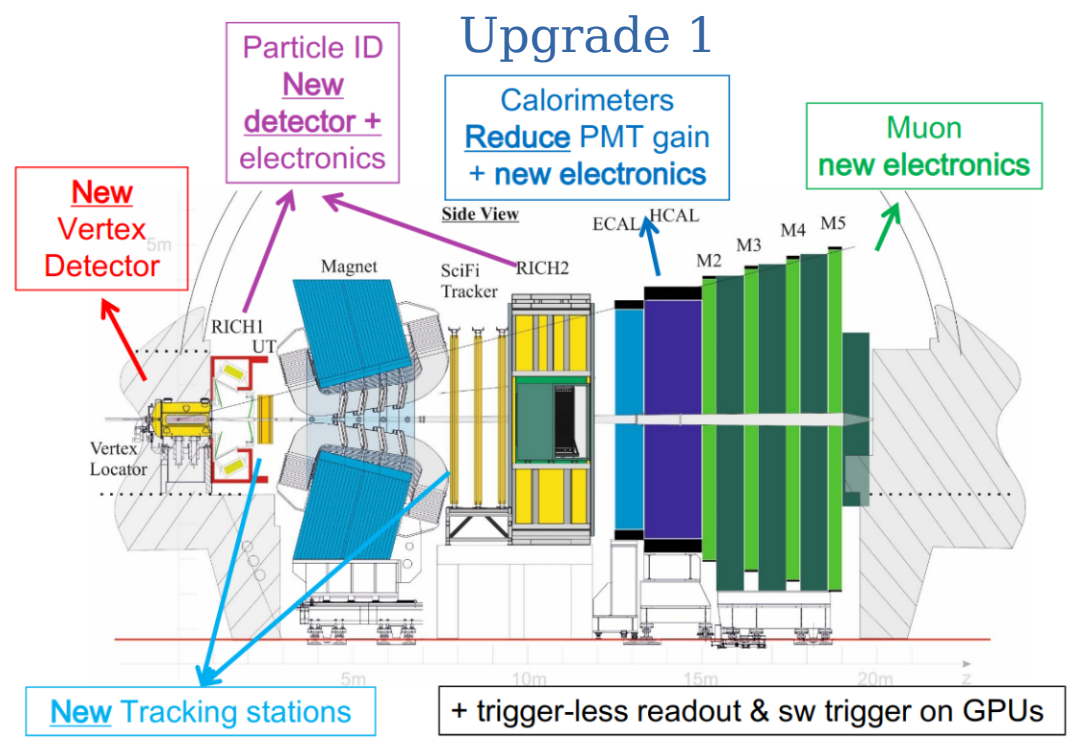


- The theoretical calculation can be simplified
 - Factorize long (**form factors**) and short (**Wilson Coefficients**) distance effects

- **Challenges:**

- Missing neutrinos → lower resolutions
- Large partially reconstructed backgrounds
- Large and perfectly calibrated simulation samples needed for modeling signal and backgrounds


- LHCb was originally designed for CP violation and rare beauty & charm decays
- But now it is a general purpose detector: *exotic spectroscopy, EW precision physics, heavy ions, fixed target program...*




Run1: 3 fb⁻¹ @ $\sqrt{s} = 7-8$ TeV
Run2: 6 fb⁻¹ @ $\sqrt{s} = 13$ TeV

- LHCb is a spectrometer in the forward direction ($2 < \eta < 5$)
- Excellent vertexing, tracking and particle identification
- Low trigger threshold on hadrons, muons and photons
- Production of all types of *b* and *c* hadrons

- Measurement of the ratios of branching fractions $R(D^*)$ and $R(D^0)$
[Phys.Rev.Lett. 131 (2023) 111802]
- Test of Lepton flavour universality using $B^0 \rightarrow D^{*\mp} \tau^+ \nu$ decay with hadronic τ channels
[Phys. Rev. D108 (2023) 012018]
- Observation of the decay $\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}$
[Phys. Rev. Lett. 128 (2022) 191803]
- Measurement of D^* longitudinal polarization in $B^0 \rightarrow D^{*\mp} \tau^+ \nu$ decays
[arXiv:2311.05224 - Submitted to PRD]
- First observation of the decay $B_s^0 \rightarrow K \mu^+ \nu$ and a measurement of $|V_{ub}|/|V_{cb}|$
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- Measurement of $|V_{cb}|$ with $B_s^0 \rightarrow D_s^{*\mp} \mu^+ \nu$ decays
[Phys. Rev. D 101 072004]

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[arXiv:2311.05224]

- Lepton Flavour Universality tests using charged current decays of D and D^* show a tension from the Standard Model of 3.2σ
- New physics can strongly affect the D^* longitudinal polarization $F_L^{D^*}$ also if LFU ratios align with the SM prediction [arXiv:1907.02257]
- Measured by Belle: $0.60 \pm 0.08 \pm 0.04$ [arXiv:1903.03102]
- The differential decays rate of $D^* \rightarrow D^0 \pi$ can be expressed as

$$\frac{d^2\Gamma}{dq^2 d\cos\theta_D} = a_{\theta_D}(q^2) + c_{\theta_D}(q^2) \cos^2\theta_D.$$

- $F_L^{D^*}$ can be calculated as

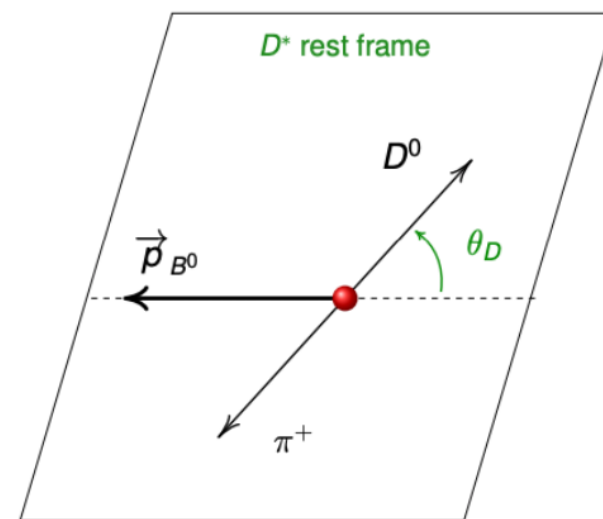
$$F_L^{D^*} = \frac{a_{\theta_D}(q^2) + c_{\theta_D}(q^2)}{3a_{\theta_D}(q^2) + c_{\theta_D}(q^2)}$$

where

- a_θ and c_θ are linear combinations of the angular coefficients

$$a_{\theta_D}(q^2) = N^{unpol} \cdot \mathcal{PDF}_{unpol}|_{\cos\theta_D=0},$$

$$c_{\theta_D}(q^2) = \frac{3}{2} N^{pol} \Delta_{bin}$$

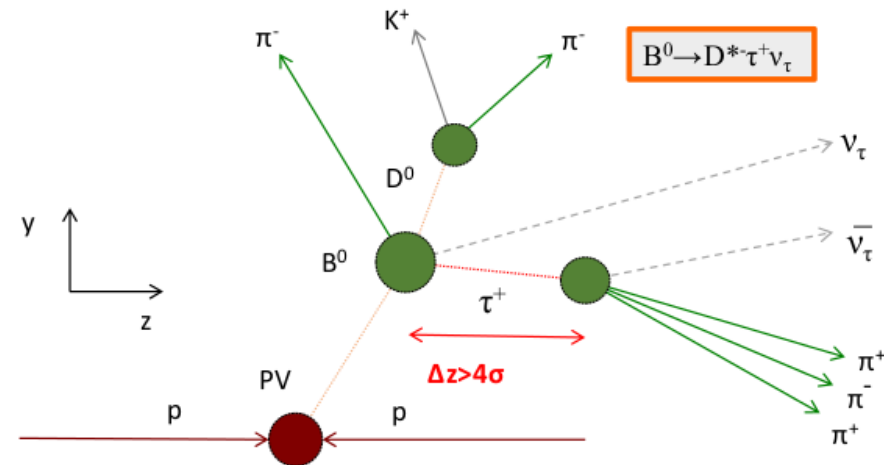


q^2 is the squared invariant mass of $\tau\nu$

[arXiv:2311.05224]

- Dataset: Run1 (3 fb^{-1}) and Run2 (2 fb^{-1})

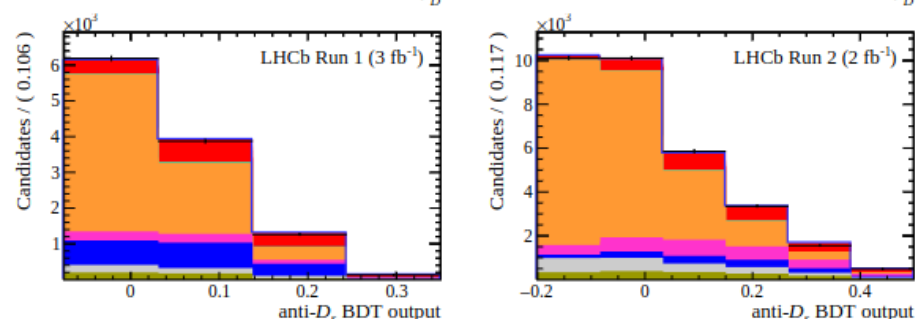
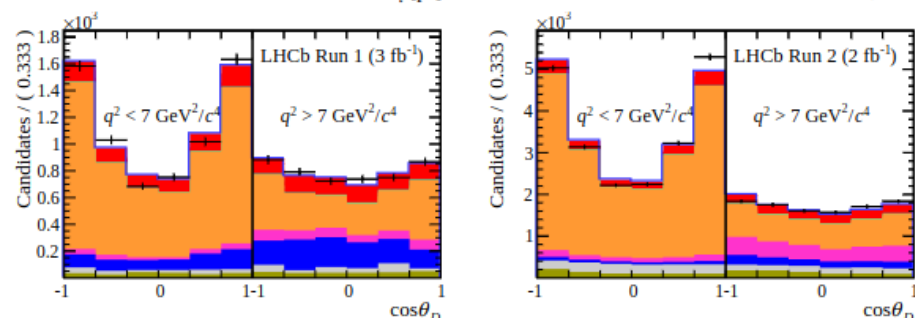
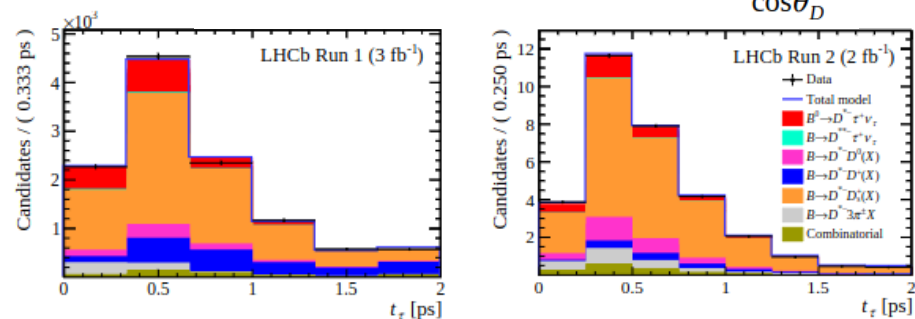
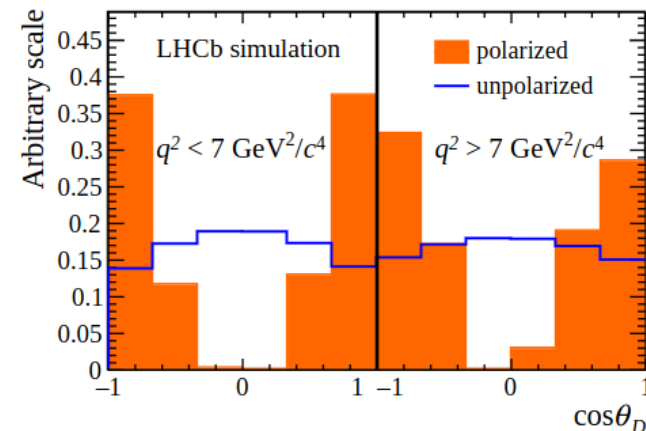
- $B^0 \rightarrow D^* \tau^+ \nu$, $\tau^+ \rightarrow 3\pi^\pm(\pi^0)\nu$
 - + good tau vertex reconstruction
 - large hadronic background



- Most dominant: $B \rightarrow D^* 3\pi X$ (BF $\sim 100x$ signal)
 - Suppressed by requiring the τ vertex to be downstream wrt B vertex along beam direction with a 4σ significance
 - additional BDT in Run2 to reach Run1 (rejection $> 99.9\%$)
- $B \rightarrow D^* D^{+,0}_{(s)} X$ (BF $\sim 10x$ signal)
 - Similar topology to that of signal but detached vertex due to non-negligible lifetime
 - Suppressed by rejecting candidates with extra charged tracks from B/ τ vertex
 - \rightarrow rejected through isolation algorithm and BDT classifier, whose output used in template fit

Longitudinal D^* polarization: template fit

- $F_L^{D^*}$ determined in two q^2 regions:
 $q^2 > 7 \text{ GeV}^2/c^4$, $q^2 < 7 \text{ GeV}^2/c^4$
- $F_L^{D^*}$ is extracted from a_θ and c_θ determined by splitting the simulated signal sample in
 - $N_{\text{unpolarized}} \propto a_\theta$
 - $N_{\text{polarized}} \propto c_\theta$
- 4D template fit:
 - τ lifetime
 - q^2
 - $\cos \theta_D$
 - Anti- D_s BDT output
- Simulated $\cos \theta_D$ signal distribution corrected for reconstruction effect
- $\cos \theta_D$ distribution corrected through fully reconstructed control samples:
 - $D_s \rightarrow 3\pi^\pm$, $D^+ \rightarrow K^- 2\pi^+$, $D^0 \rightarrow K^- 3\pi^\pm$
- simultaneous fit to Run1 and Run2
- Dominant sources of systematic uncertainties:
 - limited size of simulations samples
 - form factors parametrization

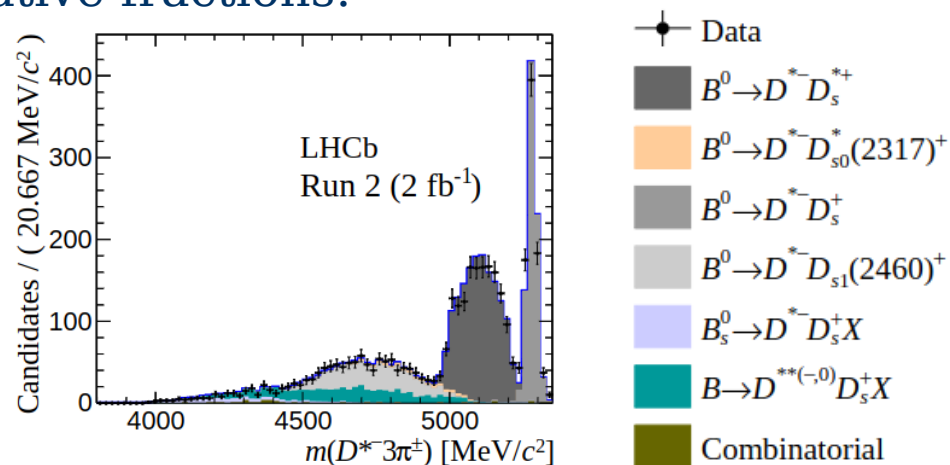


[arXiv:2311.05224]

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- Two different control samples are used to validate the D^+_s backgrounds due to the poor knowledge of the BF and the relative fractions:


- $D^+_s \rightarrow 3\pi X$ to correct the BF relevant to D^+_s meson production
- $B \rightarrow D^{*-} D^+_s (X)$ decays to constrain the relative components in the final fit \rightarrow



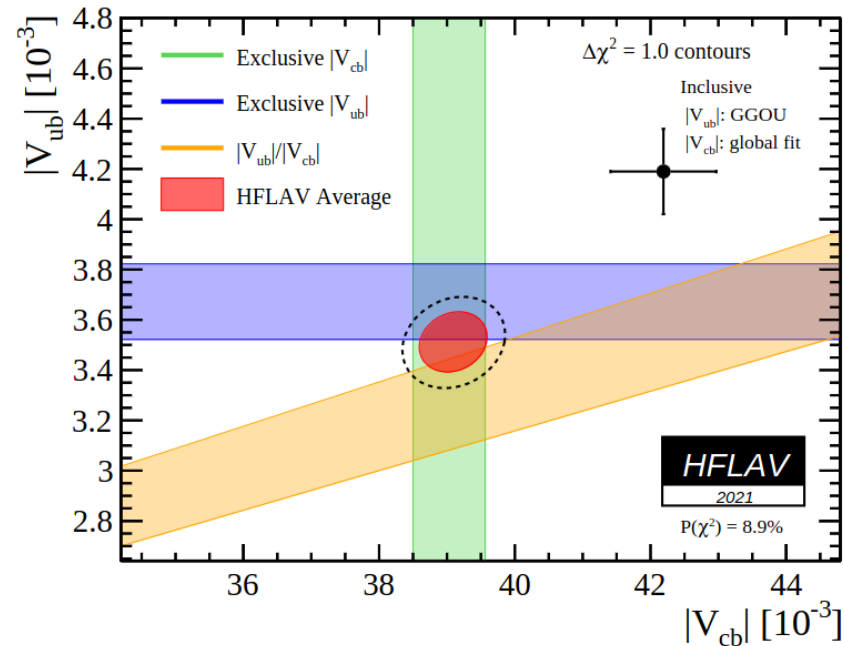
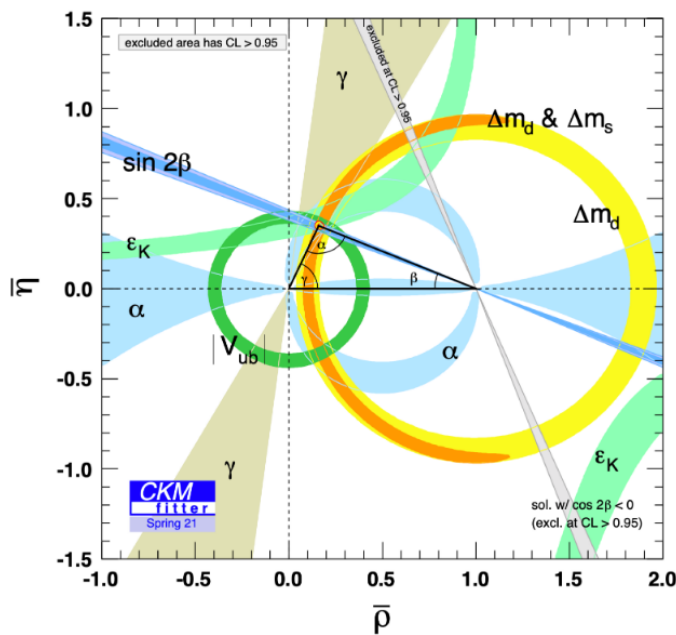
- Results integrated over run1 and run2

$$\begin{aligned}
 q^2 < 7 \text{ GeV}^2/c^4 & : & 0.51 \pm 0.07 \text{ (stat)} \pm 0.03 \text{ (syst)}, \\
 q^2 > 7 \text{ GeV}^2/c^4 & : & 0.35 \pm 0.08 \text{ (stat)} \pm 0.02 \text{ (syst)}, \\
 q^2 \text{ whole range} & : & 0.43 \pm 0.06 \text{ (stat)} \pm 0.03 \text{ (syst)}.
 \end{aligned}$$

- All results are found compatible with the SM within 1σ
- Plan is to update the $F_L^{D^*}$ value in parallel with the $R(D^*)$ measurement in hadronic τ channel.

- Measurement of the ratios of branching fractions $R(D^*)$ and $R(D^0)$
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- The parameters of the CKM matrix must be constrained in order to
 - test the unitarity of the CKM matrix
 - precisely measure the amount of CP violation in the quark sector
 → measurement of observables sensitive to the magnitudes of CKM matrix elements



- Measurements of $|V_{xb}|$ provide a crucial input for indirect searches of New Physics
- Discrepancy between exclusive and inclusive measurements: $\approx 3\sigma$ tension
 → new complementary measurements

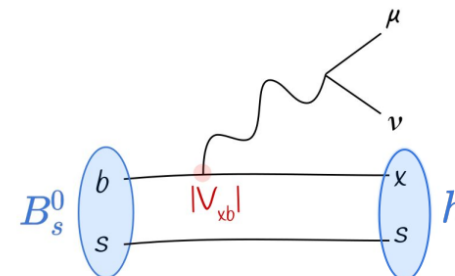
- Two main ways to measure $|V_{ub}|$ and $|V_{cb}|$:

- **Inclusive decays:**

- $B^+ \rightarrow X_c l \nu, B^0 \rightarrow X_u l \nu$
- Focus on all final states
- Need to know QCD correction to parton level decay rate

- **Exclusive decays:**

- Focus on a single final state
- Exclusive determinations rely on form factors (FF) to parameterize hadronic current as function of q^2 ($\mu\nu$ invariant mass): LQCD or QCD sum rules
 - Extracted in experimental measurement from data



- **Ground state hadrons** in the final are the golden modes for lattice QCD predictions and have the lowest theoretical uncertainties.
- **B_s decays are advantageous compared to $B^{0/+}$**
 - Easier to calculate in LQCD due to heavier spectator quark \rightarrow more precise predictions

[Phys. Rev. Lett. 126 081804]

- The strategy:
 - **Dataset:** 2012, 2 fb⁻¹ @ 8TeV
 - **Signal:** $B_s^0 \rightarrow K^- \mu^+ \nu_\mu$
 - **Normalization:** $B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu$ where $D_s \rightarrow K^+ K^- \pi^-$
 - **CKM extraction strategy:**

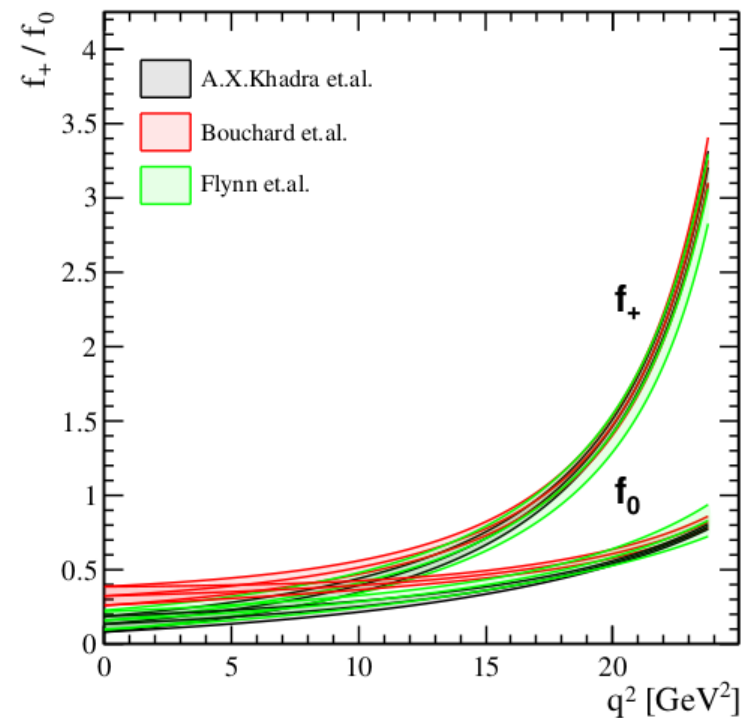
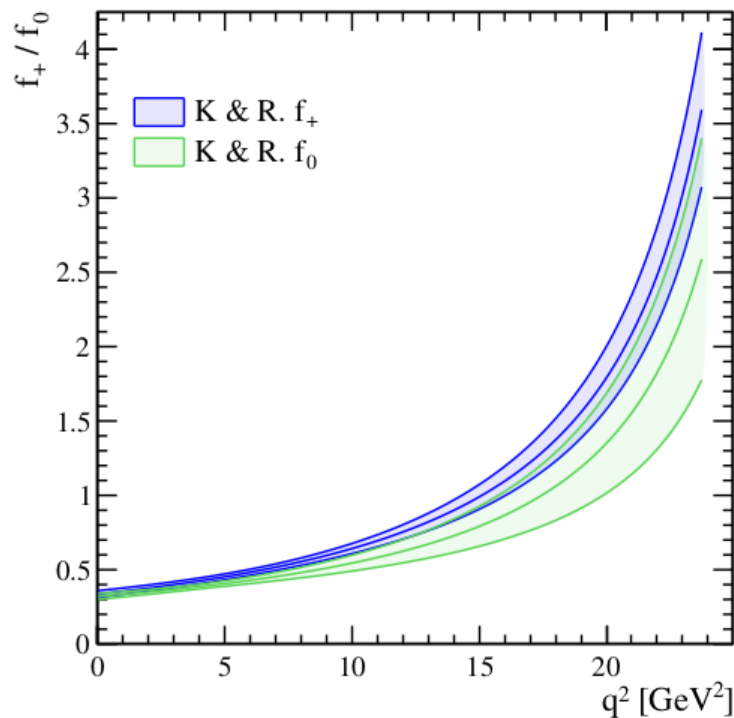
$$\underbrace{\frac{\mathcal{B}(B_s^0 \rightarrow K^- \mu^+ \nu_\mu)}{\mathcal{B}(B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu)}}_{\text{Experiment}} = \frac{|V_{ub}|^2}{|V_{cb}|^2} \times \underbrace{\frac{\text{FF}_K}{\text{FF}_{D_s}}}_{\text{Theory}}$$

- The $|V_{ub}|/|V_{cb}|$ ratio is derived in two regions of q^2 ($\mu\nu$ invariant mass) to exploit different FF_K calculation method:
 - Light cone sum rules (LCSR) @ low q^2 ($q^2 < 7 \text{ GeV}^2/c^4$)
 - LQCD @ high q^2 ($q^2 > 7 \text{ GeV}^2/c^4$)

Normalization mode FF_{D_s} fully described by LQCD [Phys Rev D. 101 074513]

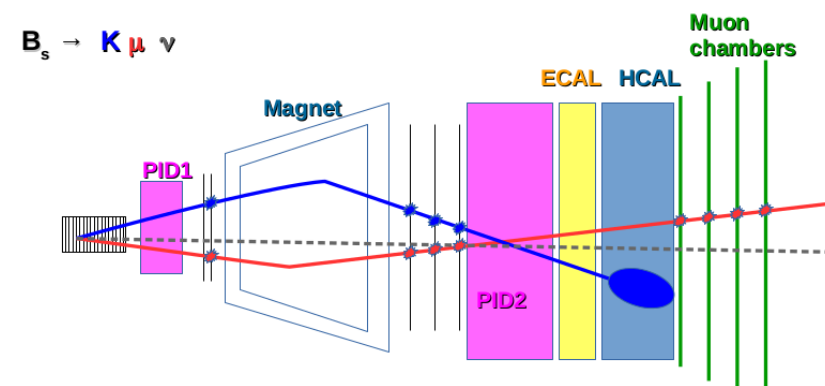
[Phys. Rev. Lett. 126 081804]

- Calculations from QCD light-cone sum rules are most precise at large recoil (low q^2)
[JHEP 08 (2017) 112]
- Lattice QCD predictions provide a precise determination of the form factors at low recoil transfer (high q^2)
[Phys. Rev. D 90, 054506] [Phys. Rev. D 91, 074510] [Phys. Rev. D 100, 034501]



[Phys. Rev. Lett. 126 081804]

- $B_s^0 \rightarrow K \mu^+ \nu$
 - main background originates from $H_b \rightarrow H_c (\rightarrow K X) \mu^+ X'$ (unreconstructed particles)
 - $B_s^0 \rightarrow K^* (\rightarrow K \pi^0) \mu^+ \nu$
 - $B_s^0 \rightarrow [cc] (\rightarrow \mu^+ \mu^-) K X$
- $B_s^0 \rightarrow D_s \mu^+ \nu$
 - $B_s^0 \rightarrow D_s^{*-} (\rightarrow D_s \gamma) \mu^+ \nu$
 - $B_s^0 \rightarrow D_s^{**} \mu^+ \nu$, $B_{u,s,d} \rightarrow D_s D X$ and $B_s^0 \rightarrow D_s^{*-} \tau^+ \nu$
- To suppress background
 - the candidates are required to be isolated from the other tracks in the event
 - BDT classifiers exploit the kinematics of the decays
- The B_s^0 momentum can be calculated with a two fold ambiguity \rightarrow regression model that exploit the B_s flight information [JHEP 02 (2017) 021]
 - Ambiguity solved by selection the solution most consistent with the regression value
 - $\varepsilon \approx 70\%$



[Phys. Rev. Lett. 126 081804]

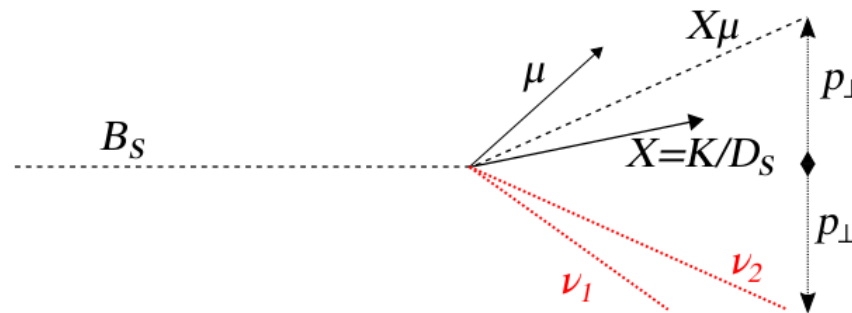
- The measured ratio is

$$R_{\text{BF}} \equiv \frac{\mathcal{B}(B_s^0 \rightarrow K^- \mu^+ \nu_\mu)}{\mathcal{B}(B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu)} = \frac{\overset{\text{Ratio of yields}}{N_K}}{\underset{\text{Efficiency ratio}}{N_{D_s}}} \frac{\epsilon_{D_s}}{\epsilon_K} \times \mathcal{B}(D_s^- \rightarrow K^+ K^- \pi^-)$$

[Prog. Theor. Exp. Phys. 2020, 083C01 (2020)]

- A binned maximum likelihood fit to the B_s corrected mass

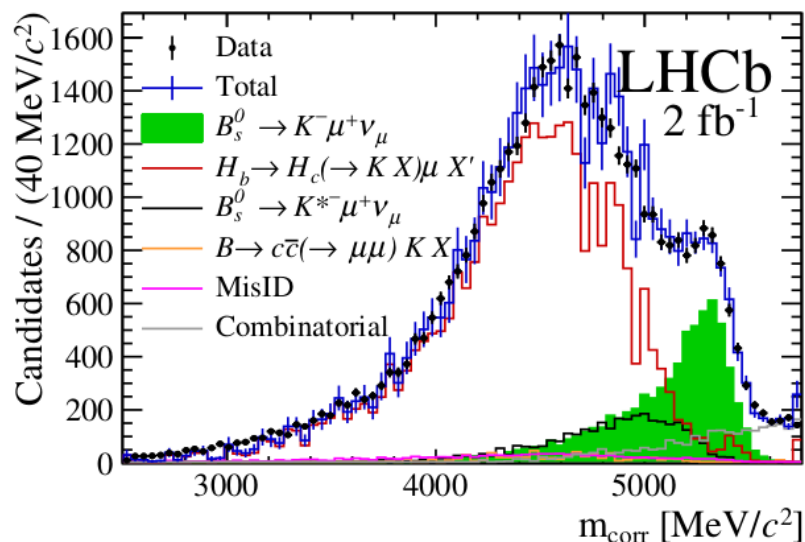
$$m_{\text{corr}} = \sqrt{m^2(Y\mu) + p_\perp^2(Y\mu) + p_\perp(Y\mu)}, \quad Y = K^-, D_s^-$$



- If only missing particle is a neutrino the corrected mass distribution will peak at the B_s mass
- Resolution improved by rejecting events with a large corrected mass uncertainty ($>100 \text{ MeV}/c^2$)

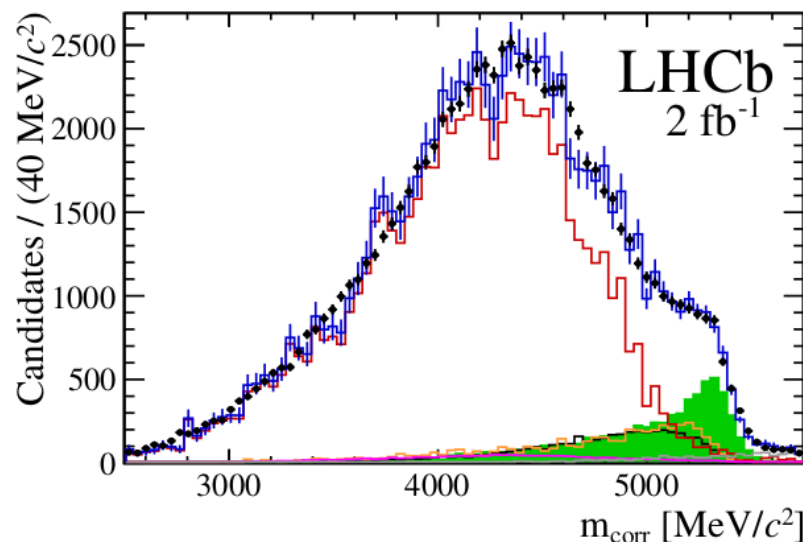
[Phys. Rev. Lett. 126 081804]

$B_s^0 \rightarrow K^- \mu^+ \nu_\mu$ low q^2



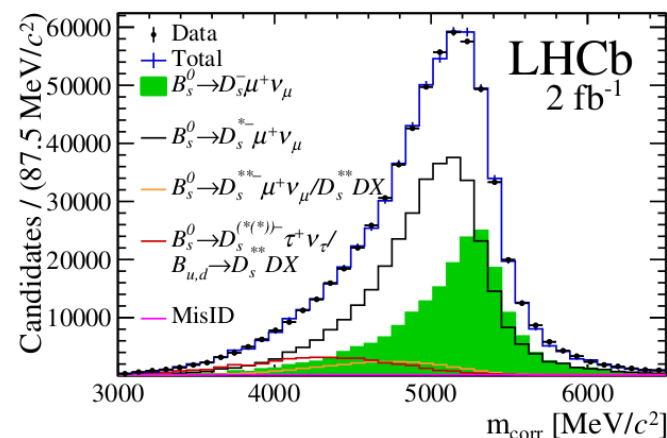
$$N(B_s^0 \rightarrow K^- \mu^+ \nu_\mu)_{low} = 6922 \pm 285$$

$B_s^0 \rightarrow K^- \mu^+ \nu_\mu$ high q^2



$$N(B_s^0 \rightarrow K^- \mu^+ \nu_\mu)_{high} = 6399 \pm 370$$

$B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu$



$$N(B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu) = 201450 \pm 5200$$

- The largest systematic uncertainty is from the fit templates
- First observation of the decay $B_s^0 \rightarrow K^- \mu^+ \nu$

[Phys. Rev. Lett. 126 081804]

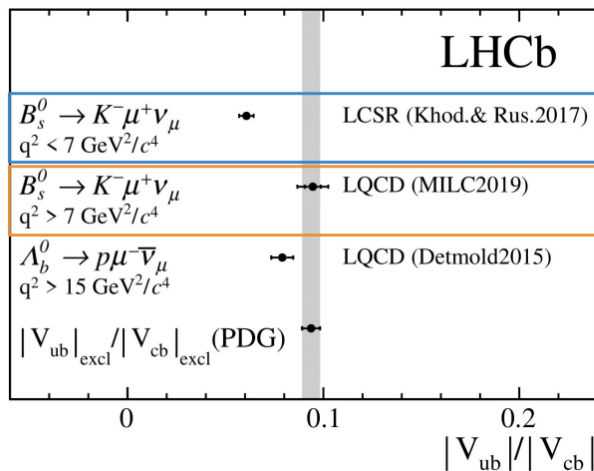
- The obtained values are

$$\underbrace{\frac{\mathcal{B}(B_s^0 \rightarrow K^- \mu^+ \nu_\mu)}{\mathcal{B}(B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu)}}_{\text{Experiment}} = \frac{|V_{ub}|^2}{|V_{cb}|^2} \times \underbrace{\frac{\text{FF}_K}{\text{FF}_{D_s}}}_{\text{Theory}}$$

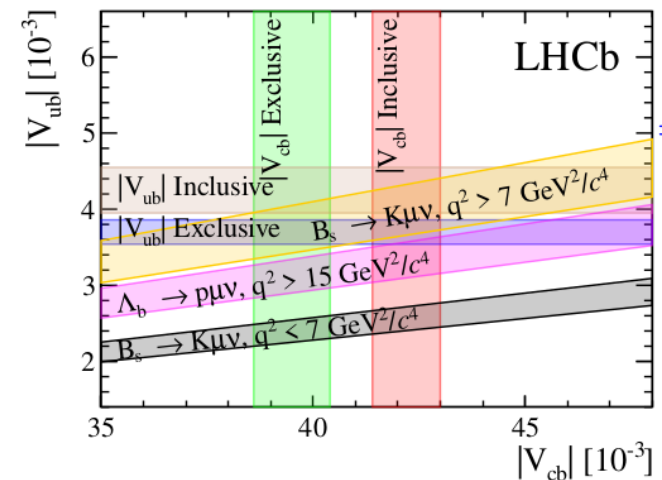
- $q^2 > 7 \text{ GeV}^2/c^4$: $\frac{\mathcal{B}(B_s^0 \rightarrow K^- \mu^+ \nu_\mu)}{\mathcal{B}(B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu)} = 1.66 \pm 0.08(\text{stat}) \pm 0.07(\text{syst}) \pm 0.05(D_s) \times 10^{-3}$
- $q^2 < 7 \text{ GeV}^2/c^4$: $\frac{\mathcal{B}(B_s^0 \rightarrow K^- \mu^+ \nu_\mu)}{\mathcal{B}(B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu)} = 3.25 \pm 0.21(\text{stat})_{-0.17}^{+0.16}(\text{syst}) \pm 0.09(D_s) \times 10^{-3}$


$$|V_{ub}|/|V_{cb}|_{(\text{low})} = 0.0607 \pm 0.0015(\text{stat}) \pm 0.0013(\text{syst}) \pm 0.0008(D_s) \pm 0.0030(\text{FF})$$

$$|V_{ub}|/|V_{cb}|_{(\text{high})} = 0.0946 \pm 0.0030(\text{stat})_{-0.0025}^{+0.0024}(\text{syst}) \pm 0.0013(D_s) \pm 0.0068(\text{FF})$$



Discrepancy related to the difference in the theoretical calculations of the form factors.



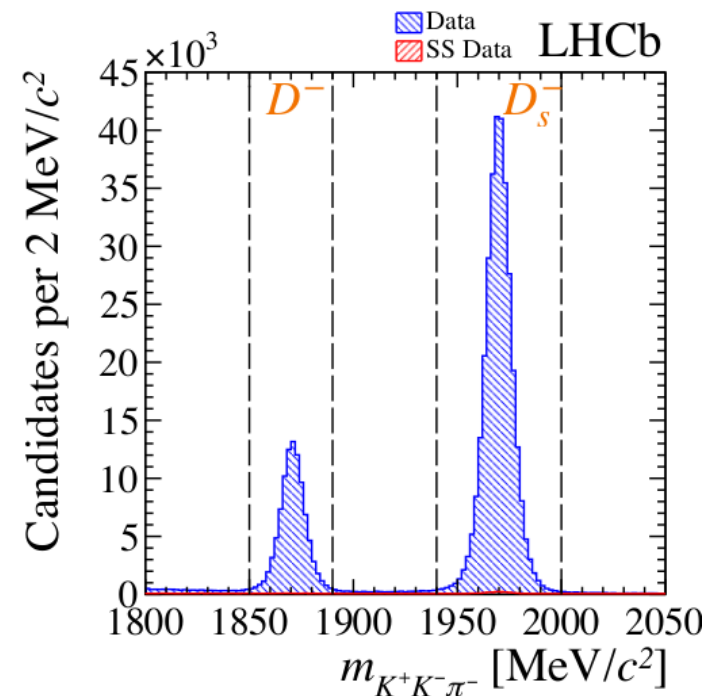
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- Signal: $B_s^0 \rightarrow D_s^{(*)-} \mu^+ \nu$ where $D_s \rightarrow \varphi(\rightarrow K^+ K^-) \pi, \gamma$ or π^0 not reconstructed
- Normalization: $B^0 \rightarrow D^{(*)-} \mu^+ \nu$
- Both channels reconstructed in the same final states
- Extract $|V_{cb}|$ from

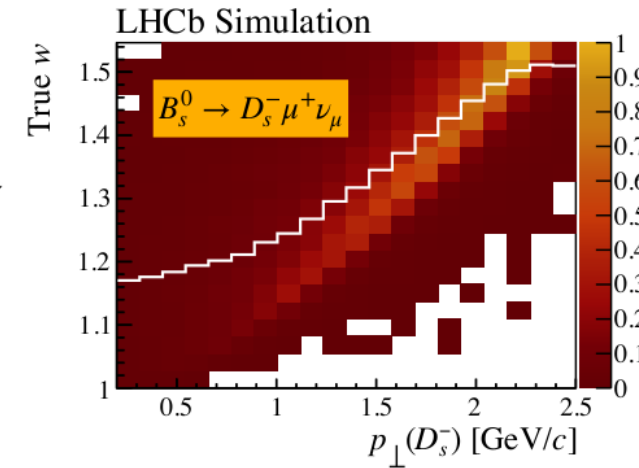
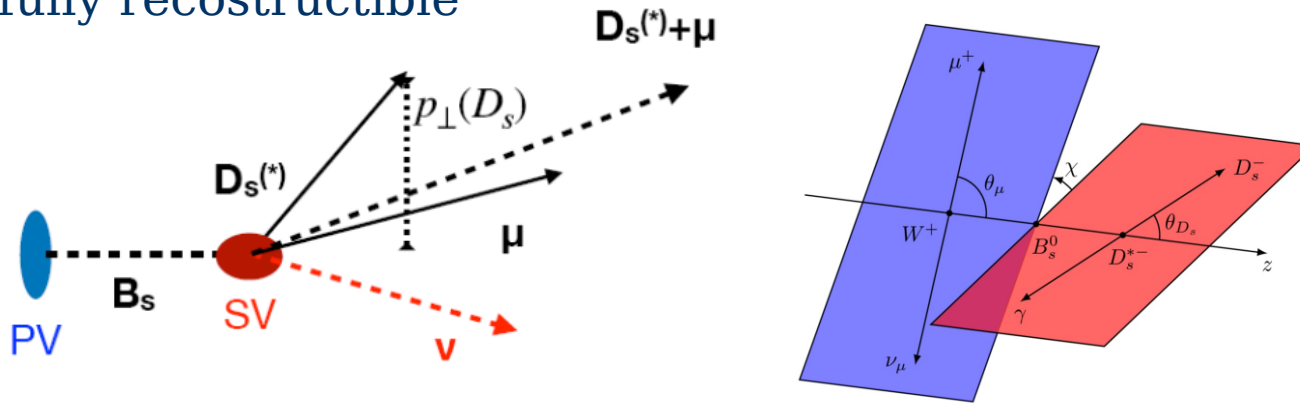
$$\mathcal{R}^* \equiv \frac{\mathcal{B}(B_s^0 \rightarrow D_s^{*-} \mu^+ \nu_\mu)}{\mathcal{B}(B^0 \rightarrow D^{*-} \mu^+ \nu_\mu)} \quad \mathcal{R} \equiv \frac{\mathcal{B}(B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu)}{\mathcal{B}(B^0 \rightarrow D^- \mu^+ \nu_\mu)}$$

- external input:
 - hadronization fractions f_s/f_d [PRD(2019)031102]
 - branching fractions [PDG]

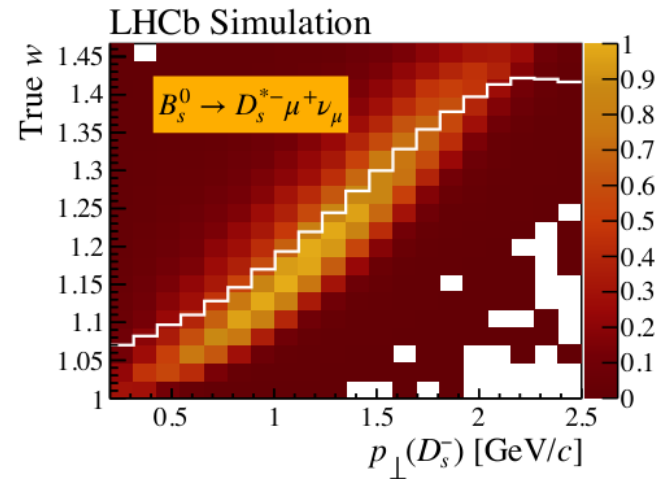


[Phys. Rev. D 101 072004]

- Due to the undetected neutrino we cannot determine precisely the q^2
 \rightarrow use variable $p_{\perp}(D_s)$ with respect to B flight distance:
 - high correlated with hadron recoil w
 - fully reconstructible



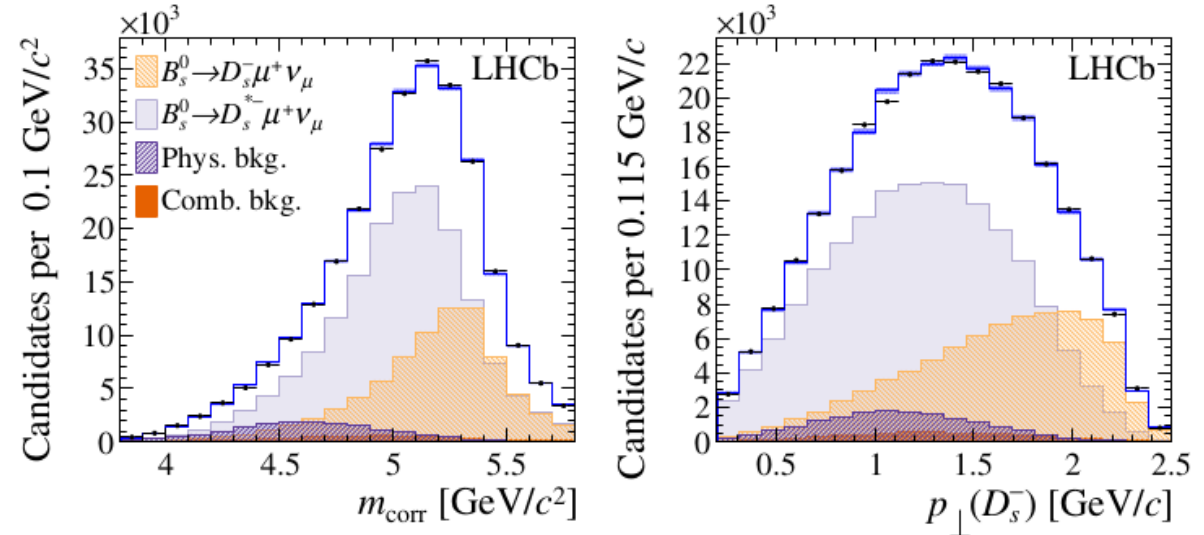
$$\frac{d^4\Gamma(B \rightarrow D^* \mu \nu)}{dw d\cos\theta_{\mu} d\cos\theta_D d\chi} = \frac{3m_B^3 m_{D^*}^2 G_F^2}{16(4\pi)^4} \eta_{EW}^2 |V_{cb}|^2 |\mathcal{A}(w, \theta_{\mu}, \theta_D, \chi)|^2$$



- $w = v_B \cdot v_{D^*} = (m_B^2 + m_{D^*}^2 - q^2)/(2m_B m_{D^*})$
- 2-D template fit to M_{corr} and $p_{\perp}(D_s)$ identify the signal yields and provides a simultaneous measurement of the ratios $R^{(*)}$ and the form factors

$|V_{cb}|$: with $B_s^0 \rightarrow D_s^{*-} \mu^+ \nu$ decays

[Phys. Rev. D 101 072004]



- FF Parametrizations used:
 - CLN and BGL

- The results are

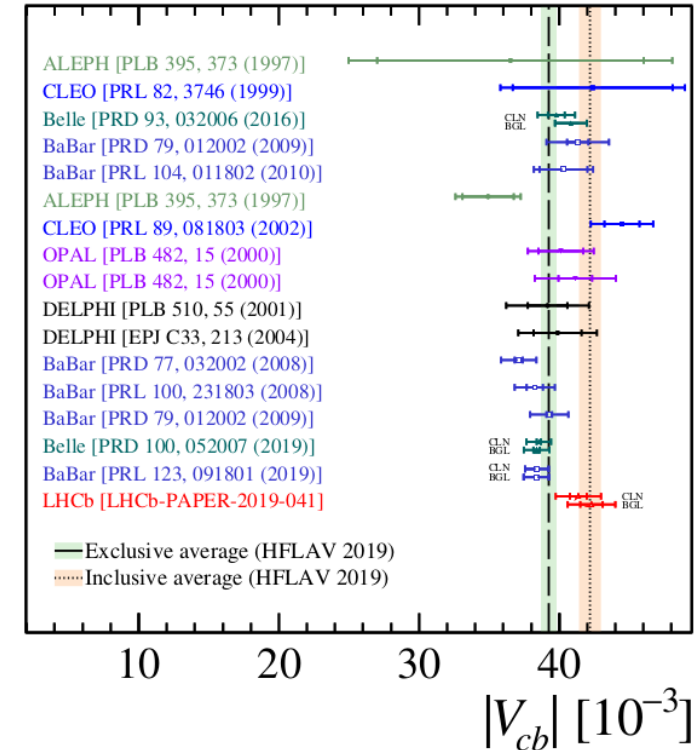
$$|V_{cb}|_{CLN} = (41.1 \pm 0.6(stat) \pm 0.9(syst) \pm 1.2(ext)) \times 10^{-3}$$

$$|V_{cb}|_{BGL} = (42.3 \pm 0.8(stat) \pm 0.9(syst) \pm 1.2(ext)) \times 10^{-3}$$

- First measurement of $|V_{cb}|$ using B_s and in a hadronic environment
- Compatible with world average for both inclusive and exclusive determinations
- Confirms trend that parametrisation is not responsible for inclusive vs exclusive disagreements
- New $f_s/f_d \rightarrow V_{cb}$ [arXiv:2103.06810]

$$|V_{cb}|_{CLN} = (40.8 \pm 0.6(stat) \pm 0.9(syst) \pm 1.1(ext)) \times 10^{-3}$$

$$|V_{cb}|_{BGL} = (41.7 \pm 0.8(stat) \pm 0.9(syst) \pm 1.1(ext)) \times 10^{-3}$$



- Broad SL physics program at LHCb
- Successful Run1 and Run2: $3+6 \text{ fb}^{-1}$, still many analysis ongoing
- Upgrade Phase I:
 - 10 times more data (20 times more hadronic events)
 - Complementarity with Belle
 - Synergy between LHCb, ATLAS and CMS on some important channels
- Strong program beyond flavour exploiting unique acceptance

