

Axions in Astrophysics

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EuCAPT Virtual Colloquium
April, 2024



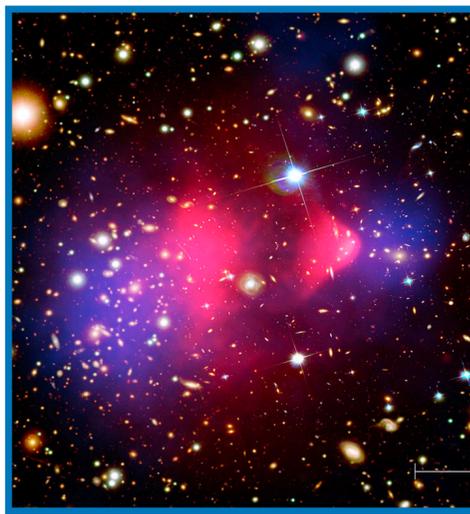
THE ROYAL SOCIETY



UNIVERSITY OF
OXFORD

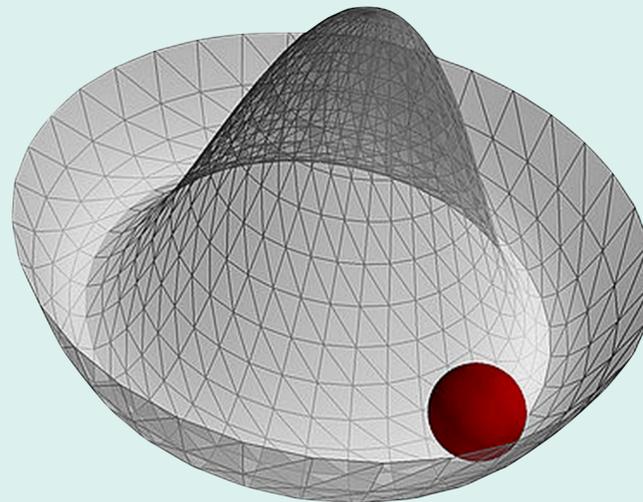


Axions: the motivation



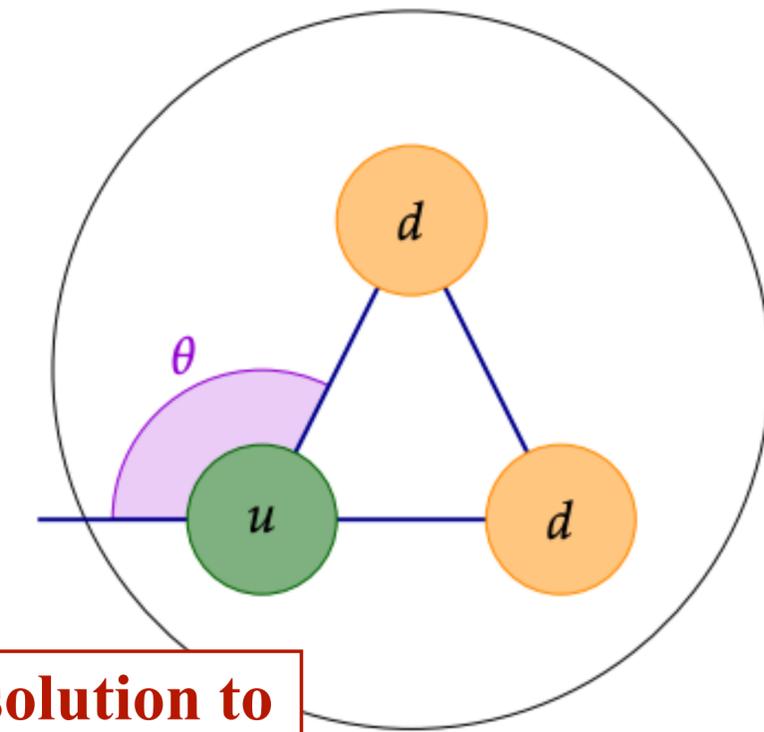
Axions as dark matter

Light Pseudoscalars

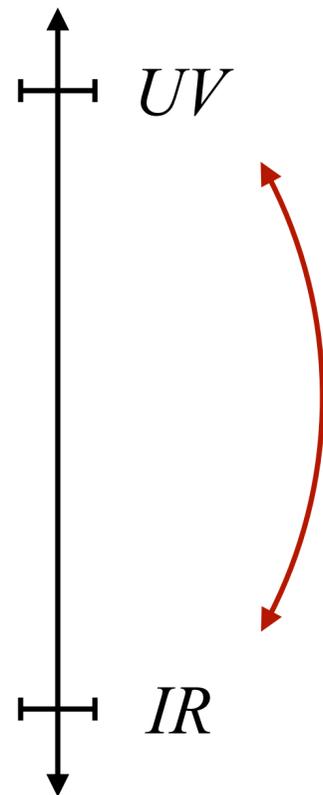


Axions as IR remnants of the UV

Axions as a solution to the strong CP problem



(QCD axion)



Motivation #1: The strong CP problem

Why does QCD seem to conserve charge-parity (CP) symmetry?

CP violating term in QCD Lagrangian:

$$\mathcal{L}_{CPV} \supset \bar{\theta} \frac{g_s^2}{32\pi^2} G \tilde{G}$$

Neutron electric dipole moment (eDM) $\propto \bar{\theta}$

Current limit: $\bar{\theta} \leq 5 \times 10^{-11}$

Abel et al (2020)

The strong CP problem

Solution: Introduce goldstone boson a to make theta term dynamical

$$\mathcal{L}_\theta = - \left(\bar{\theta} + \frac{a}{f_a} \right) \frac{g^2}{32\pi^2} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

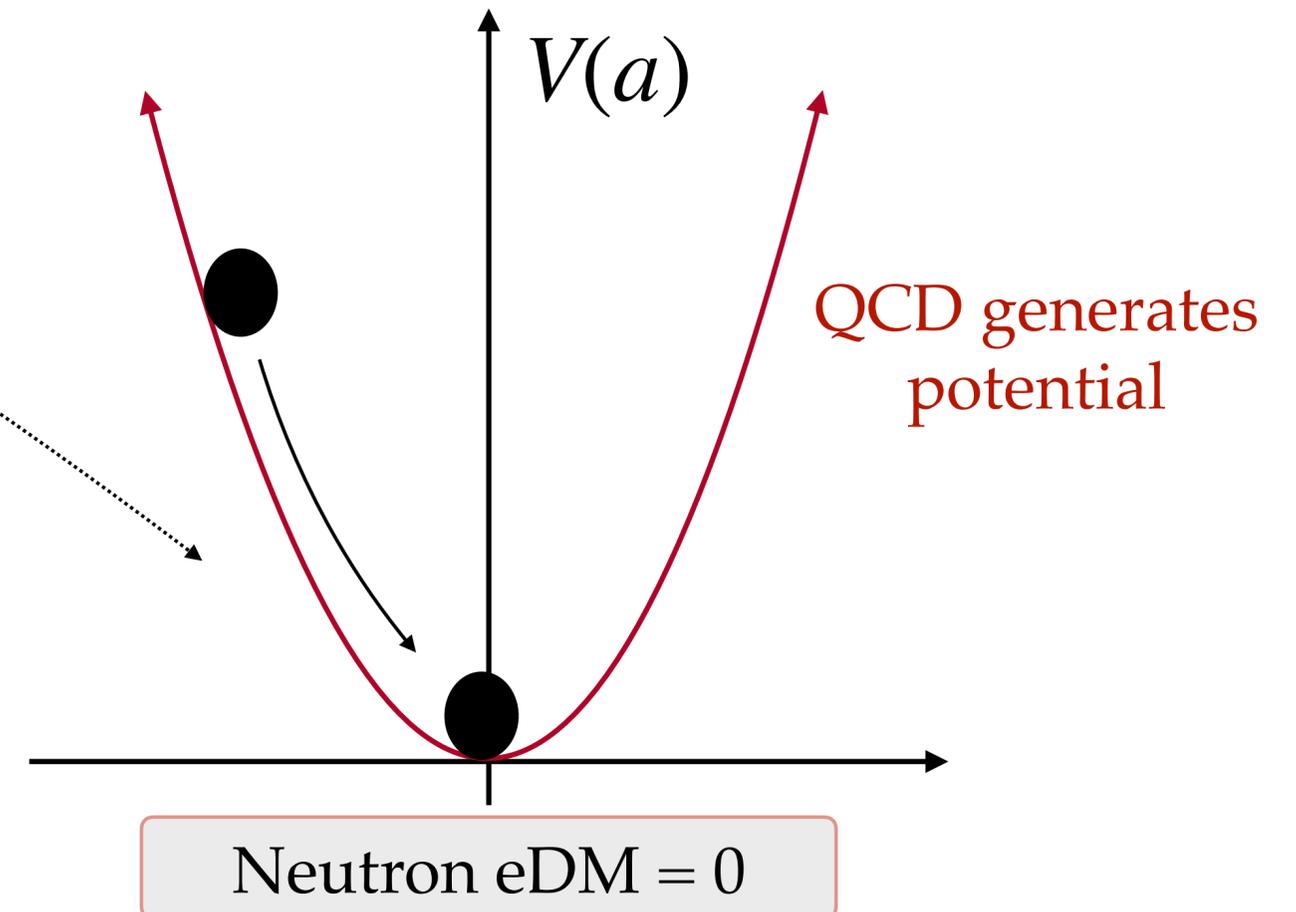
“QCD Axion”

$$m_a \sim 10^{-6} \text{ eV} \left(\frac{10^{12} \text{ GeV}}{f_a} \right)$$

Original $f_a \sim \text{EW scale}$



“Invisible axion”: $f_a \gg \text{EW scale}$

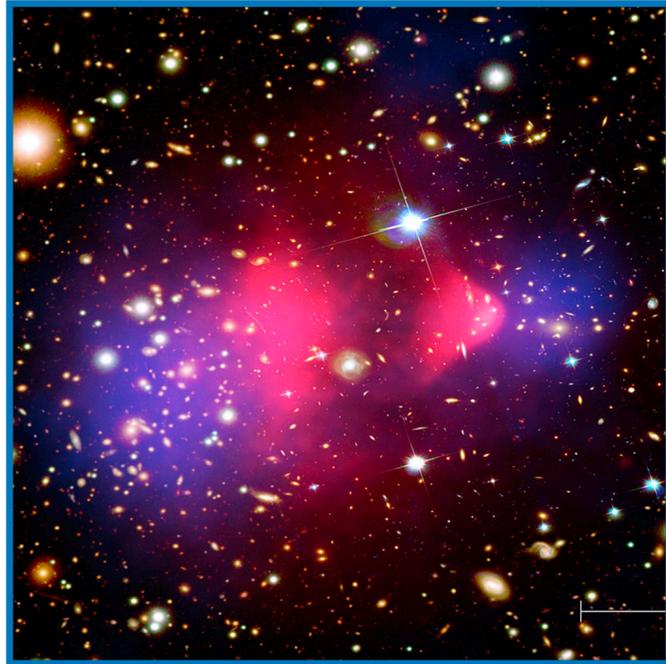


Motivation #2: Dark matter

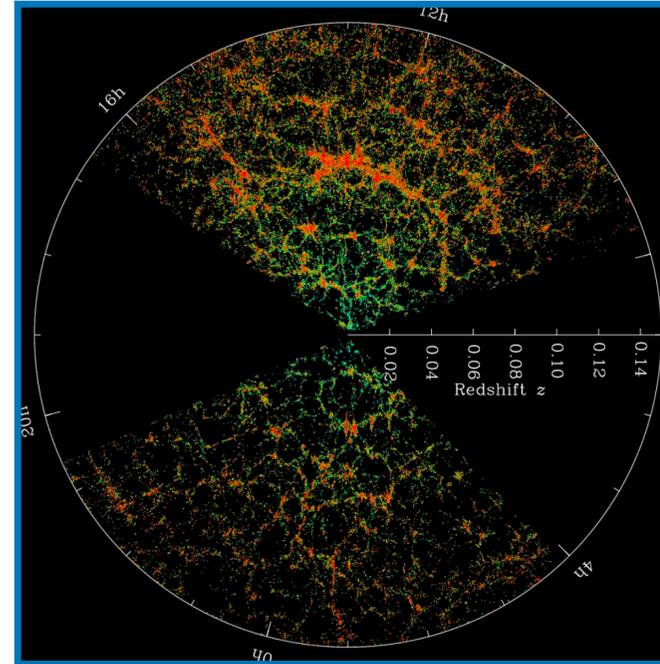
Galaxy rotation curves



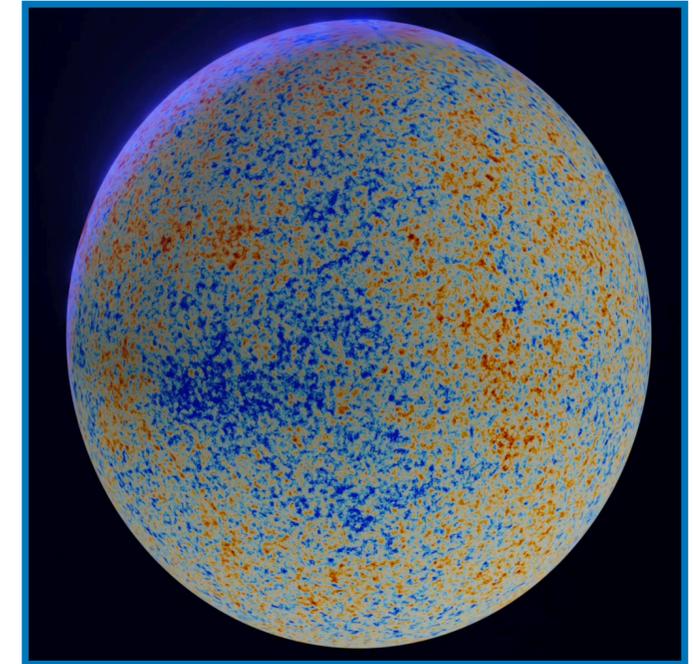
Merging galaxy clusters



Large scale structure



Cosmic microwave background



Requirements:

- Production mechanism
- Cosmologically long-lived
- Feeble interactions

[Requirement of generic new light physics]

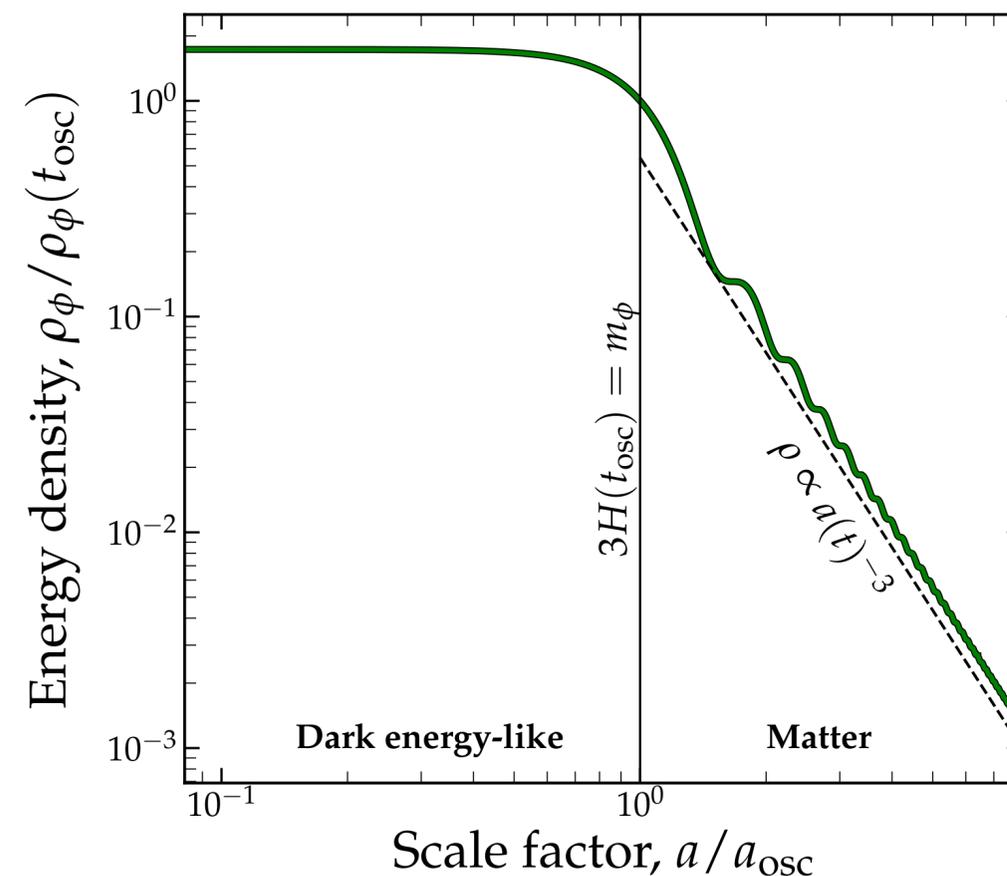
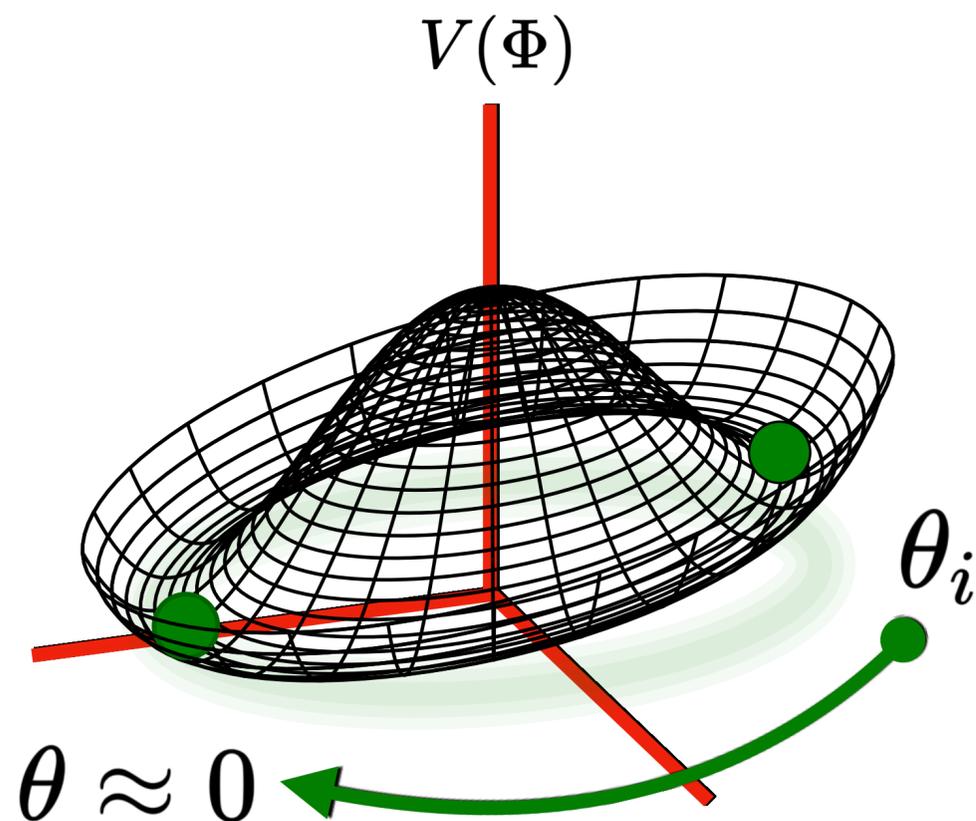
$$\tau \propto m_a^{-3}$$

$$\lambda \propto f_a^{-1}$$

Natural for light axions from high energy scale

Axions from the misalignment mechanism

Equation of motion: $\ddot{\theta} + 3H\dot{\theta} + m_a^2\theta = 0$



$$\Omega_{\text{DM}} h^2 \sim 0.05 \left(\frac{\tilde{\theta}_i}{1} \right)^2 \left(\frac{f_a}{10^{17} \text{ GeV}} \right)^2 \left(\frac{m_a}{10^{-22} \text{ eV}} \right)^{1/2}$$

Axions from topological defects

Cosmic strings network in the early Universe

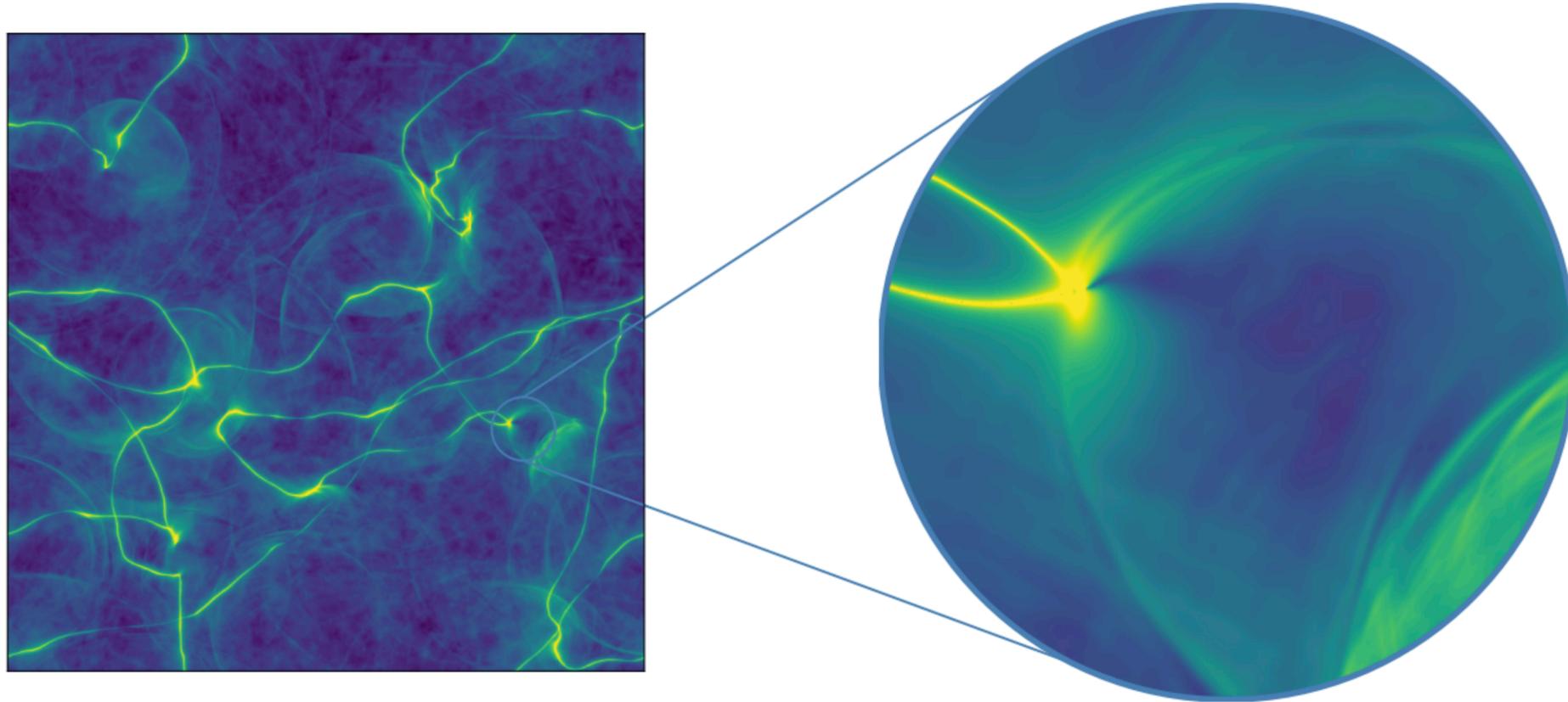


Image credit: Bernabou et al (2023)

Grilla di Cortana, Hardy, Pardo Vega, Villadoro (2016), Ghorgetto, Hardy, Villadoro (2018, 2021), Bushmann et al (2022), Saikawa et al (2024)

Strings/walls dominate for QCD axion with high inflationary scale

Axion density field

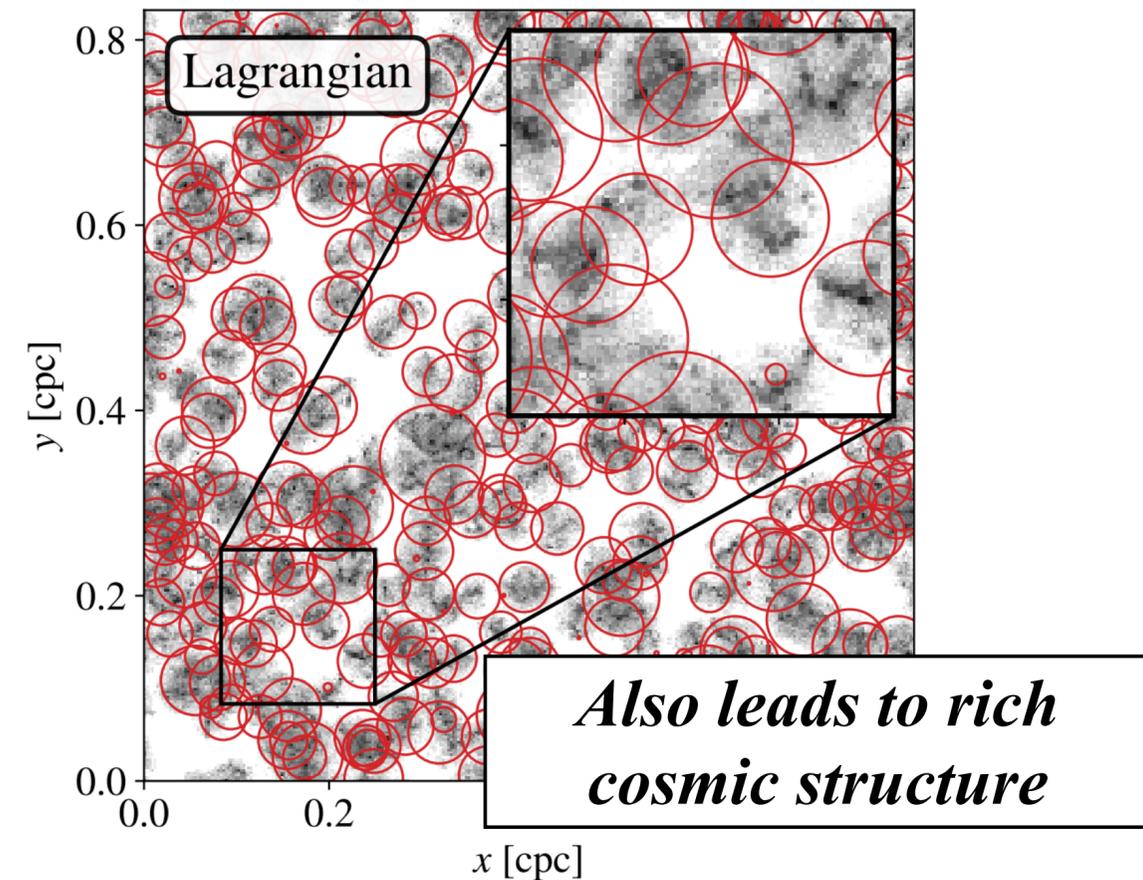
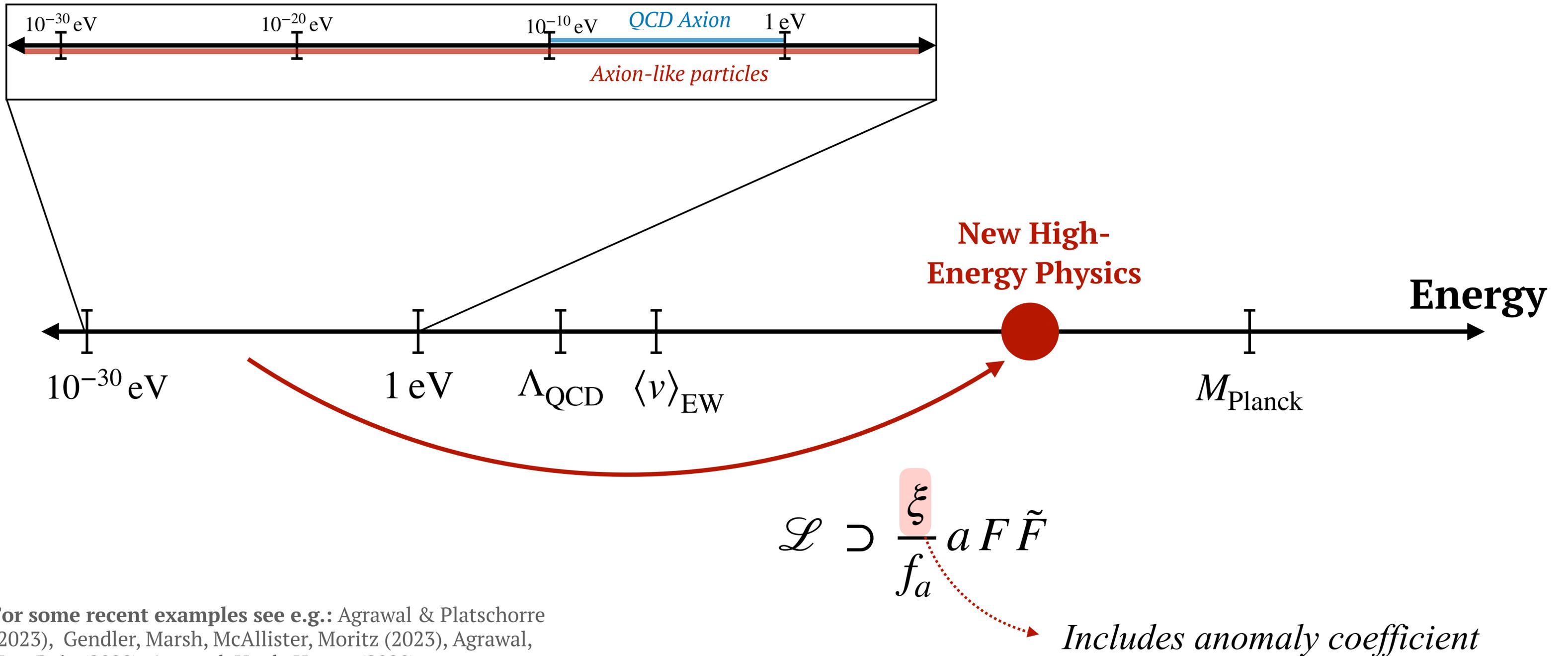


Image credit: Ellis et al (2022)

Collapse leads to “axion miniclusters” and “axion stars”

Motivation #3: Probes of high energies



For some recent examples see e.g.: Agrawal & Platschorre (2023), Gendler, Marsh, McAllister, Moritz (2023), Agrawal, Nee, Reig (2022), Agrawal, Hook, Huang (2020)...

Example: axions as a probe of GUTs

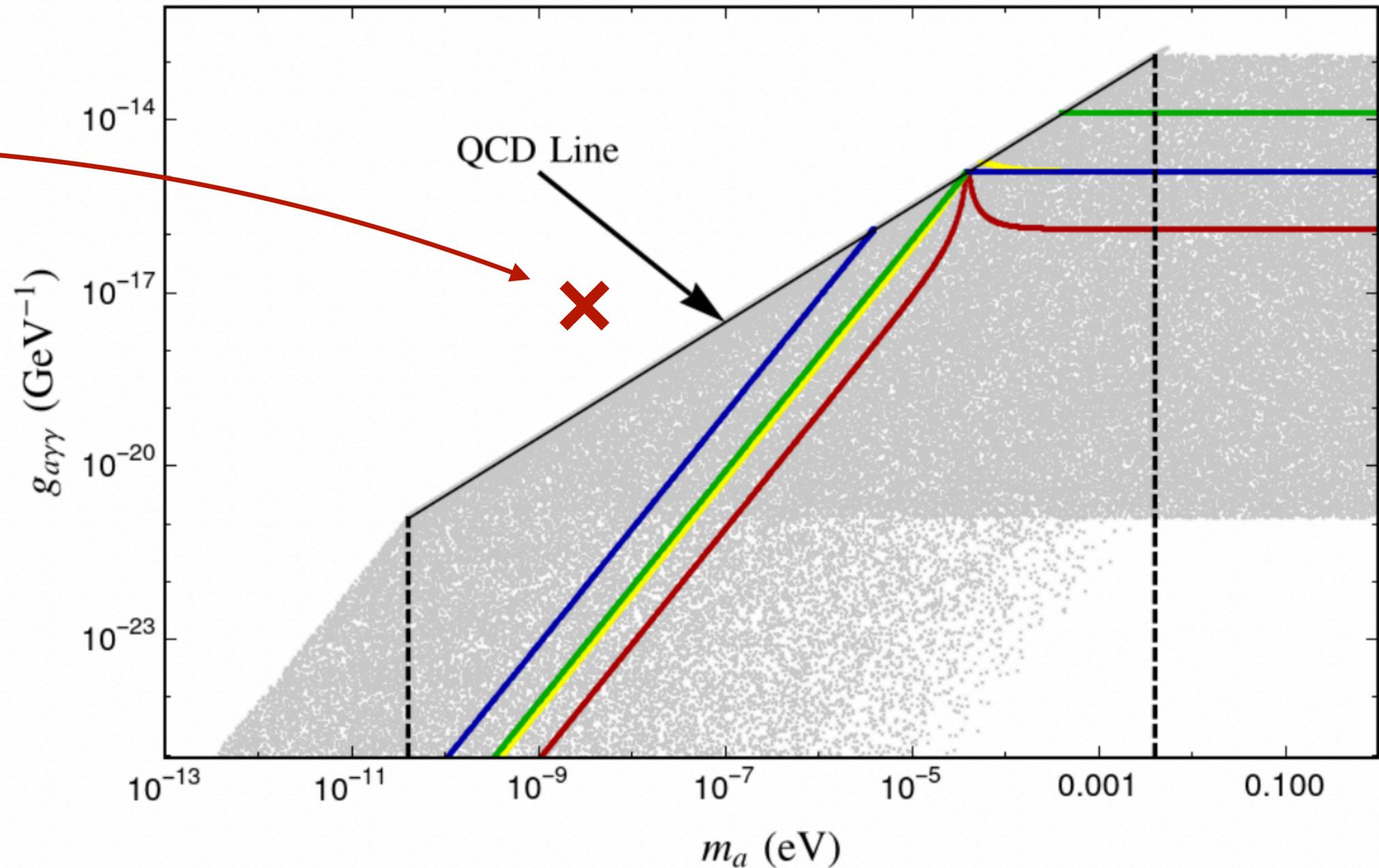
Agrawal, Nee, Reig (2022)

Detection of axion here implies:

- QCD axion is tuned light

OR

- no GUT



Astrophysics as a laboratory

White dwarf



Pulsar



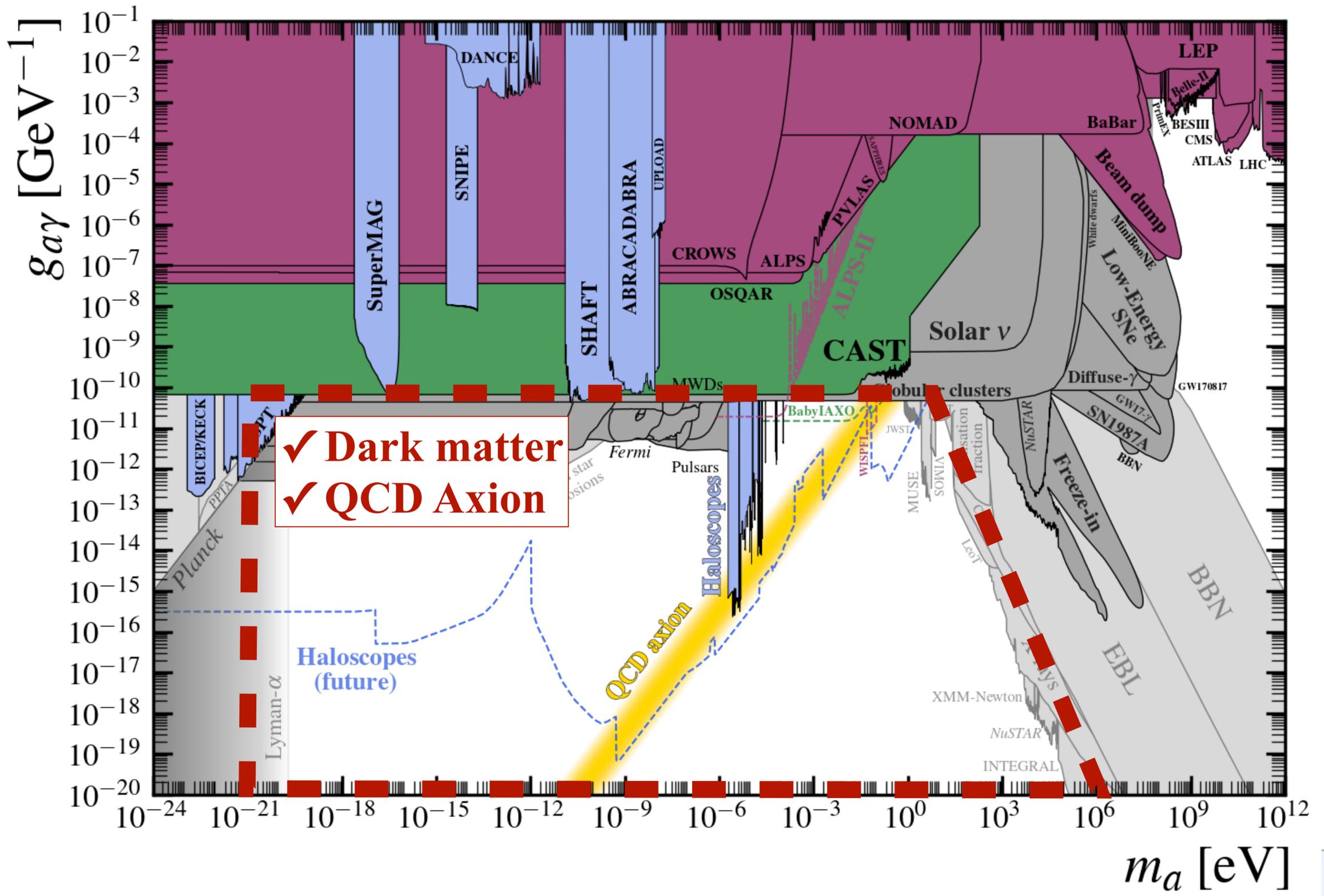
Black hole



Can we use Nature's laboratories to search for axions?

Parameter space overview

$$\mathcal{L} \supset -\frac{g_{a\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$



✓ Dark matter
✓ QCD Axion

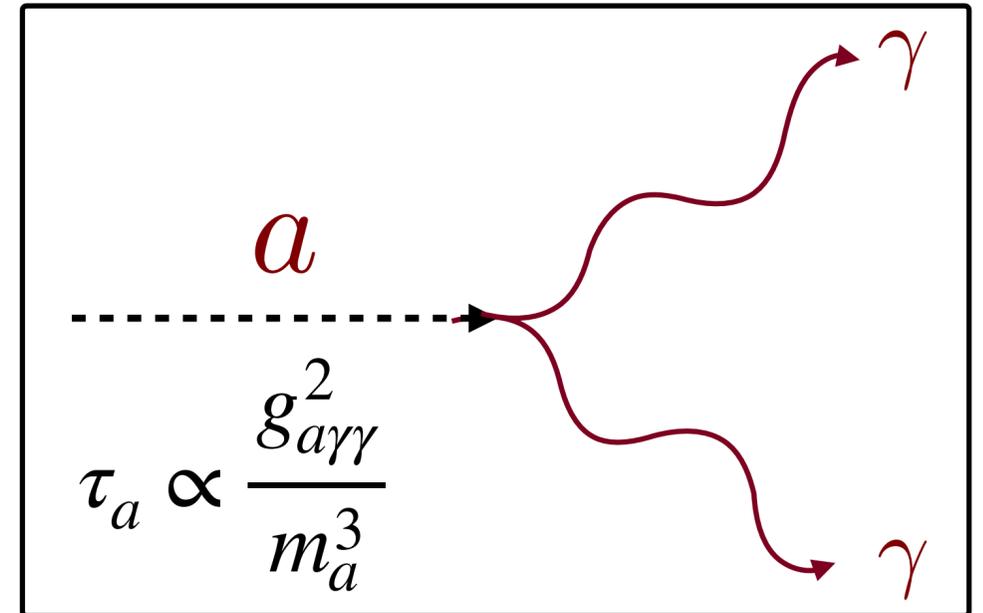
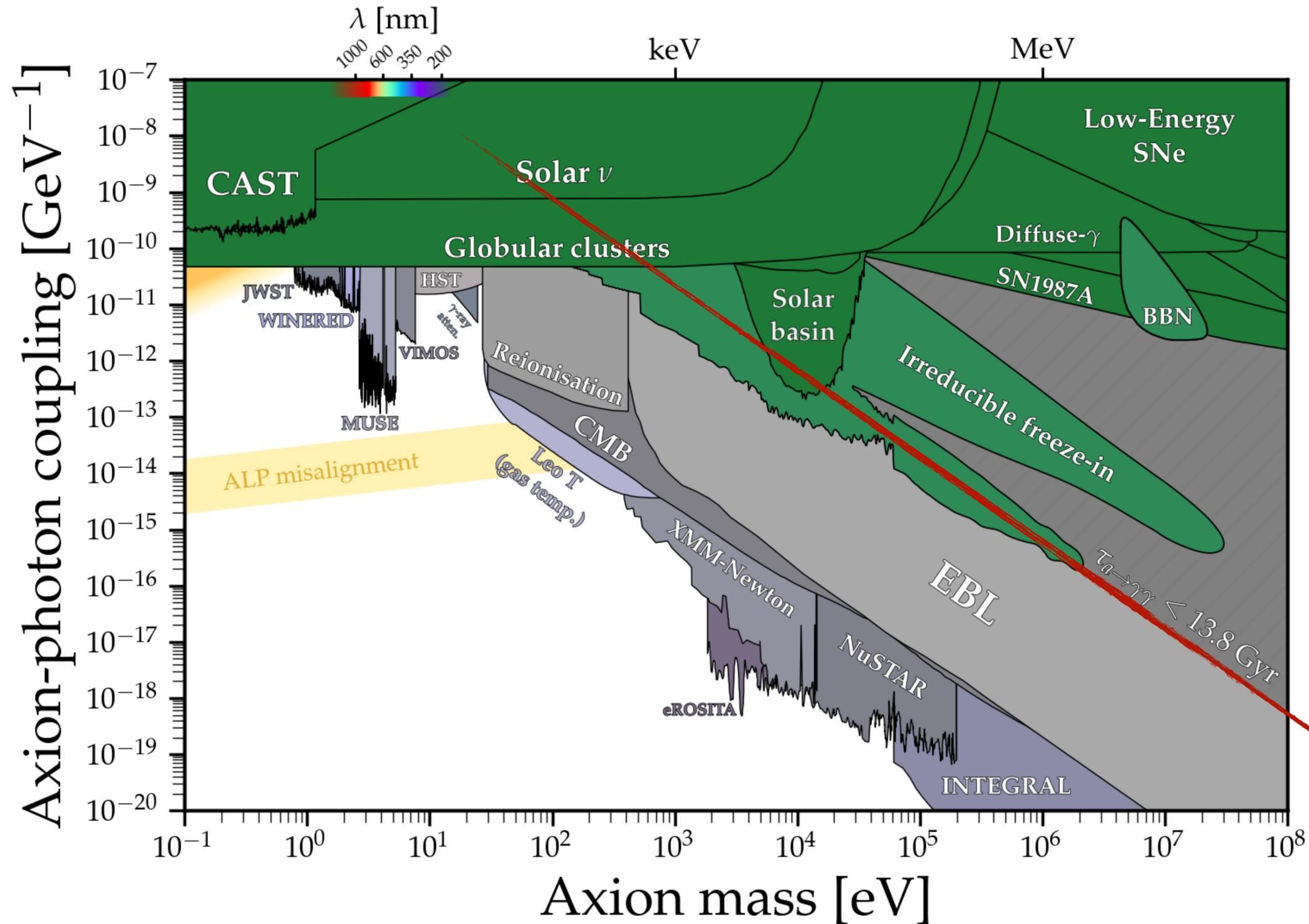
Haloscopes
(future)

QCD axion

Haloscopes

m_a [eV]

Axion decay



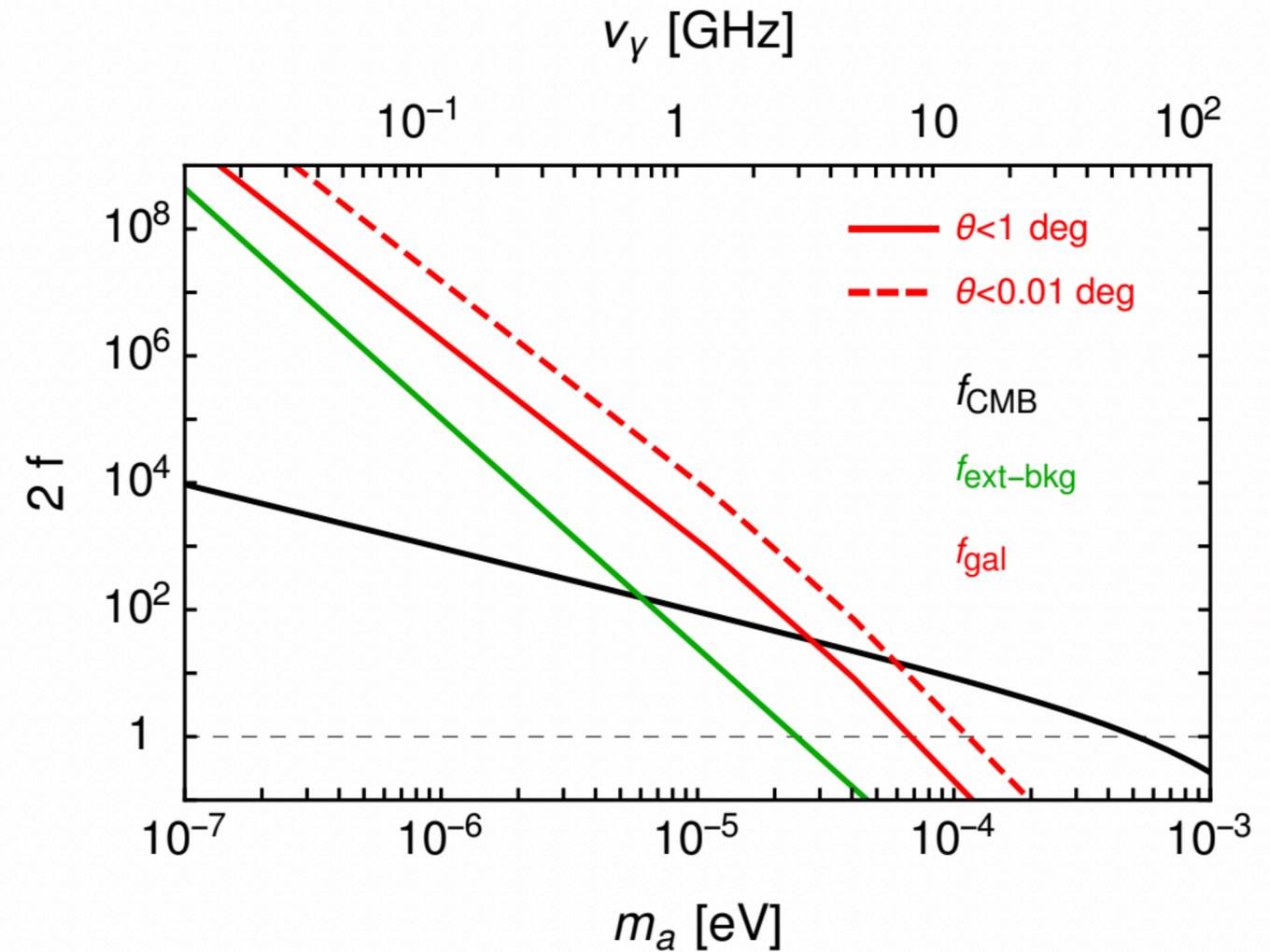
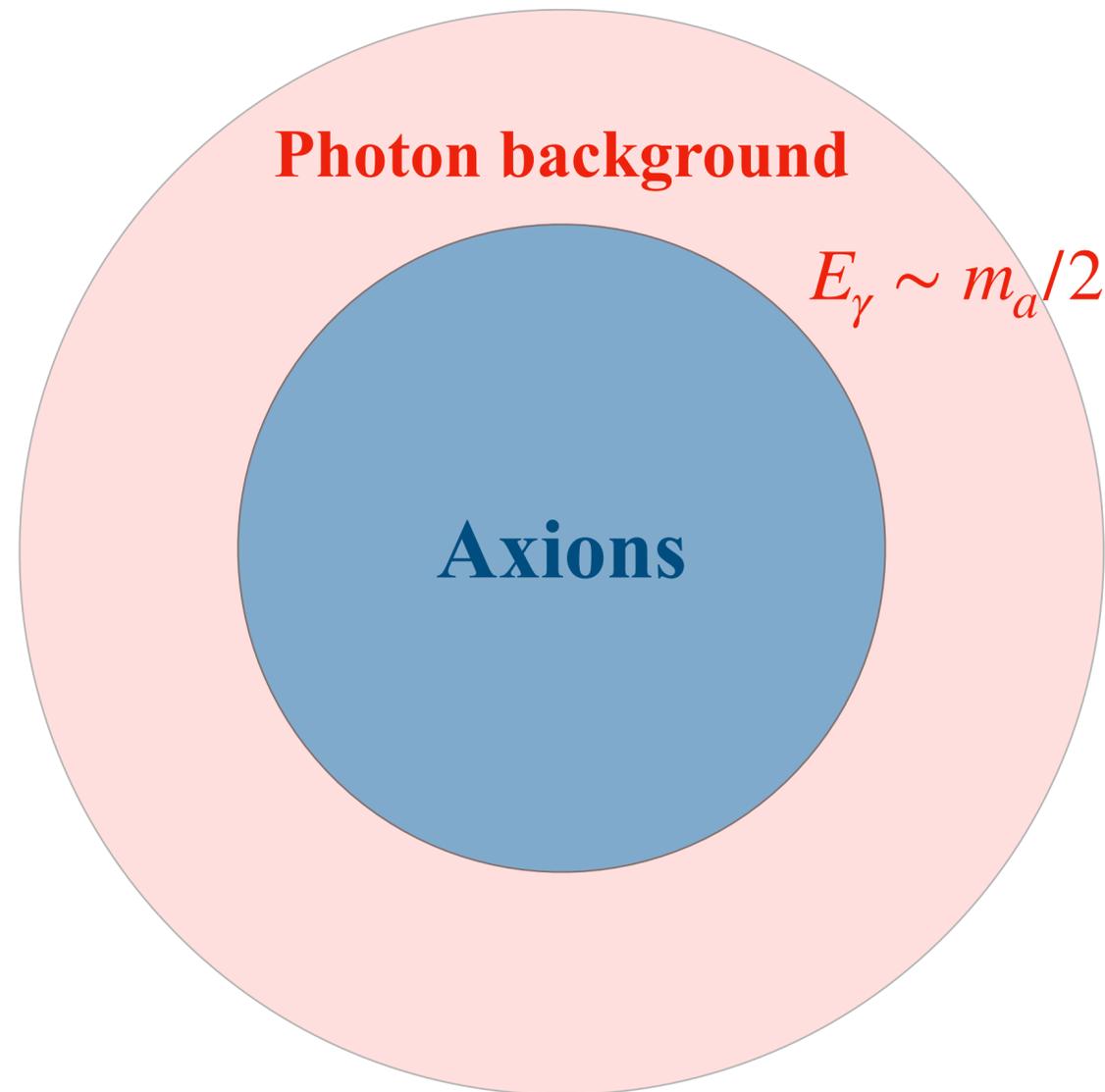
Upside: Very robust
Difficulty: Hard to make progress

Cosmic Stability

Simulating axion decay

$$\tau_a \propto \frac{g_{a\gamma\gamma}^2}{m_a^3} \left(1 + 2f_\gamma\right)^{-1}$$

Tkachev (1987, 2015), Kephart & Weiler (1995), Caputo, Peña-Garay, **SJW** (2018), Caputo, Regis, Taoso, **SJW** (2018), Azra & Skive (2019), Battye et al (2019), Sigl & Trivedi (2019), Carena, Mirizzi, Sigl (2020), Ghosh, Savludo, Miranda (2020), Arza, Schwetz, Todarello (2020), Sun et al (2022, 2023), Buen-Abad, Fan, Sun (2022), Escudero et al (2023).....

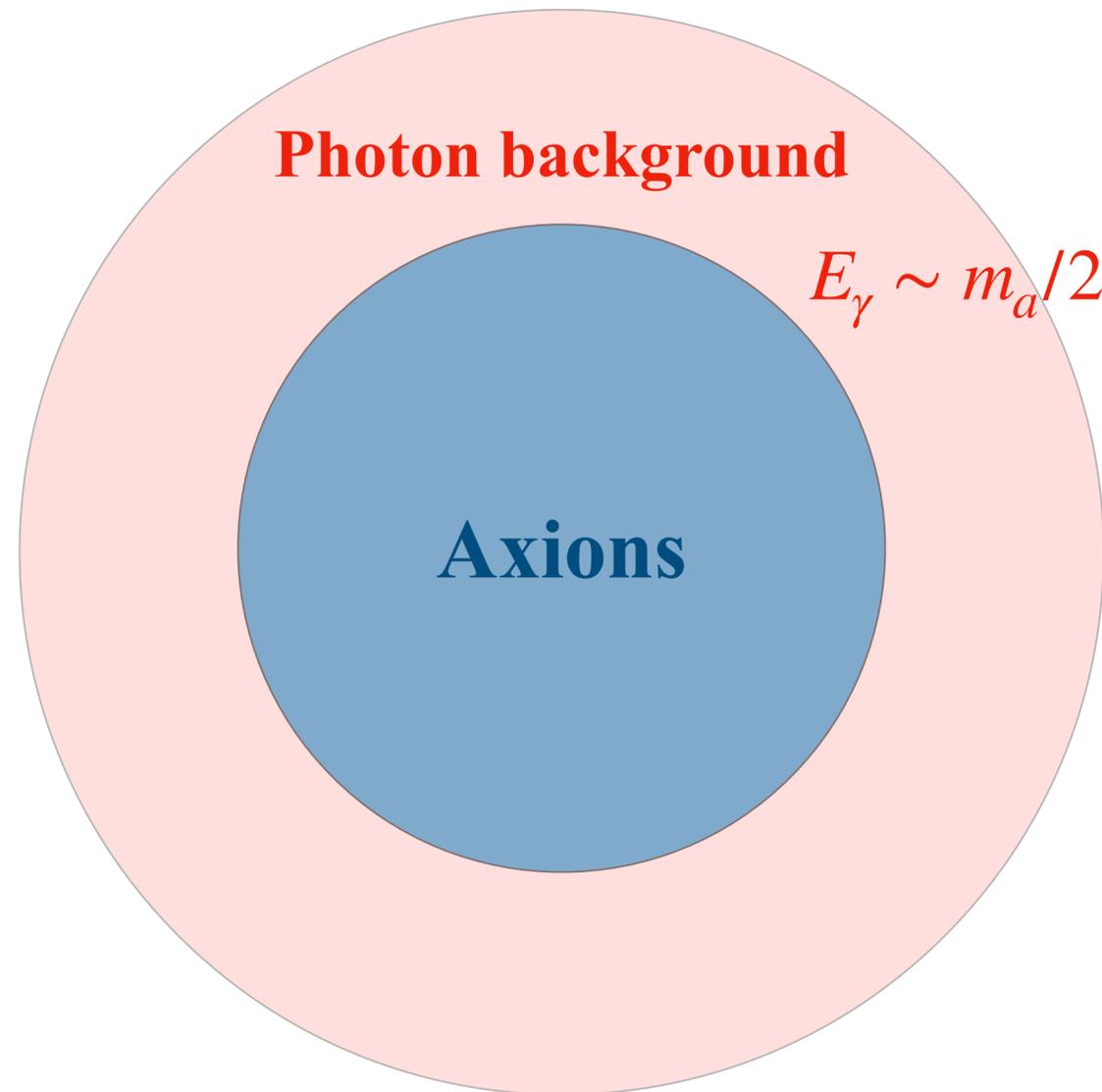


Caputo, Regis, Taoso, **SJW** (2018)

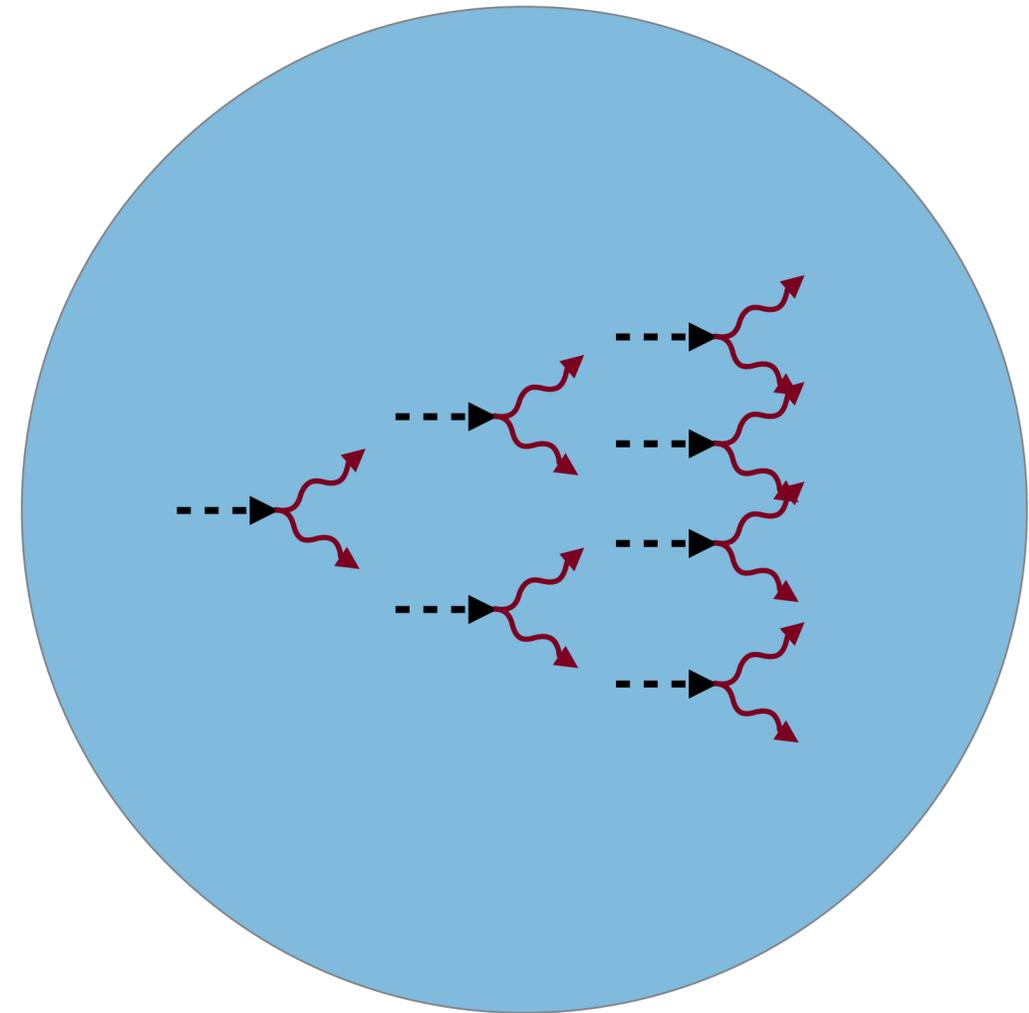
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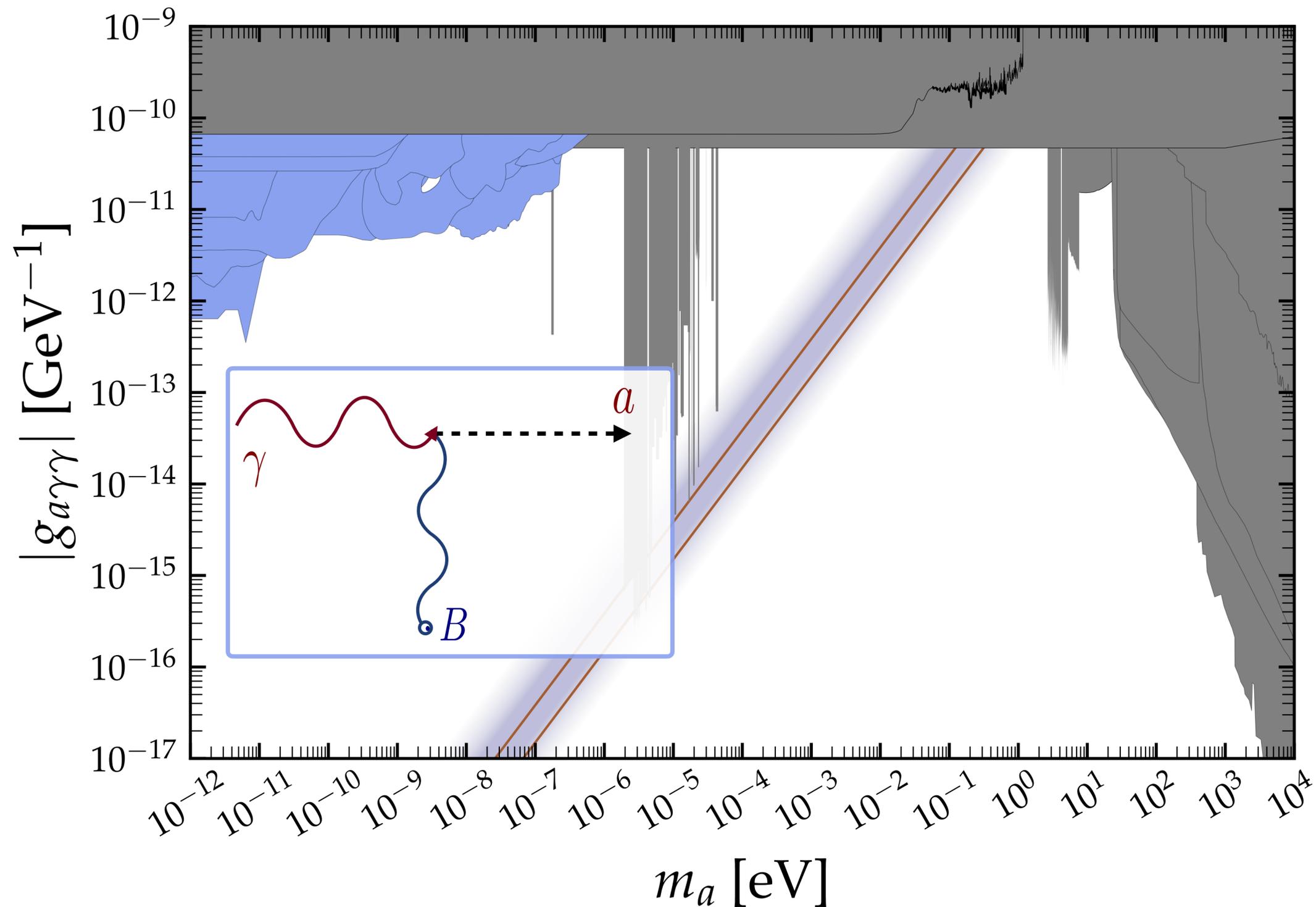
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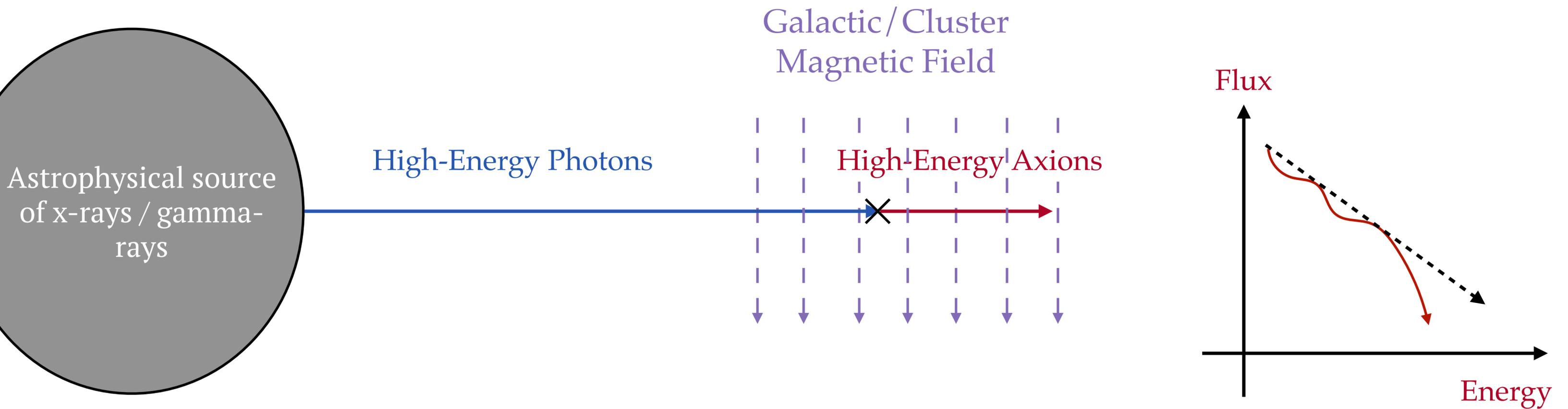
*High density
& Uniform*



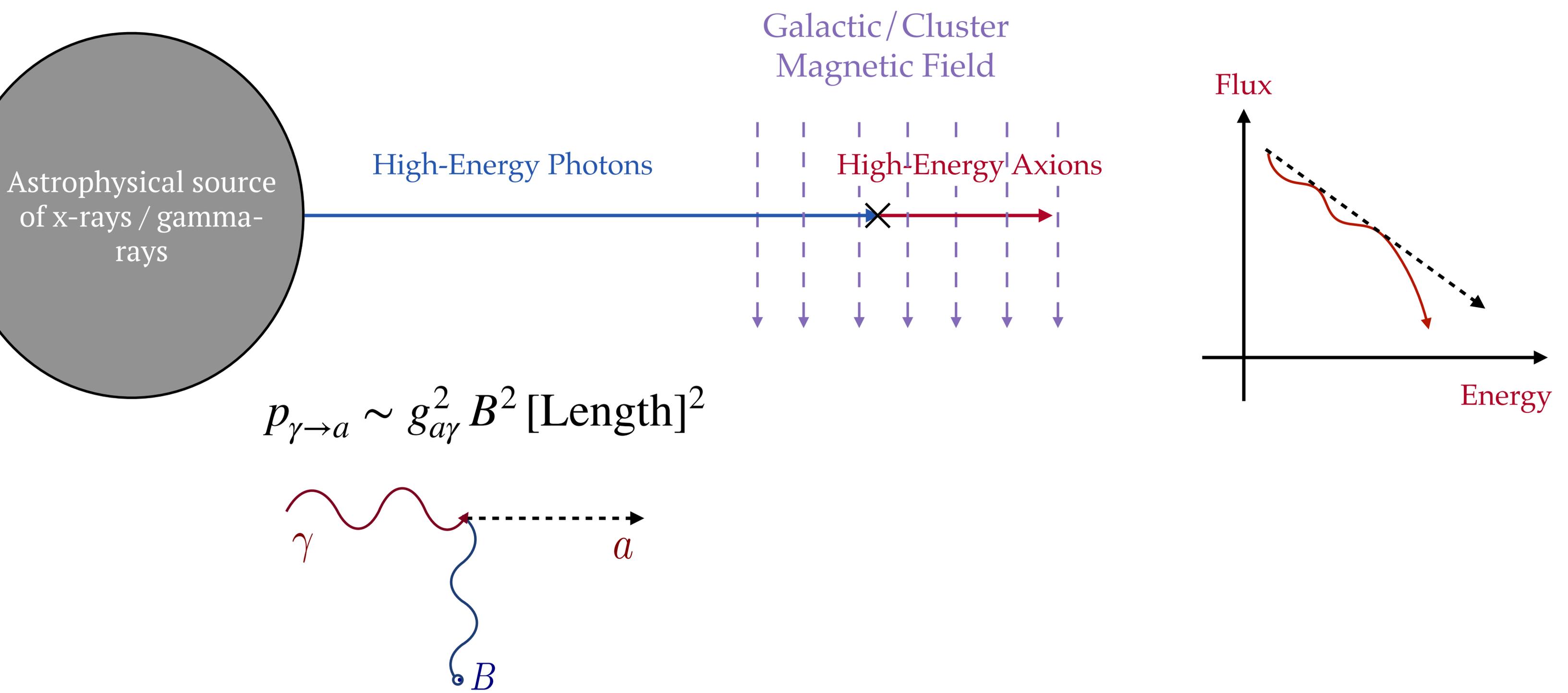
X-ray / Gamma-ray searches for axions



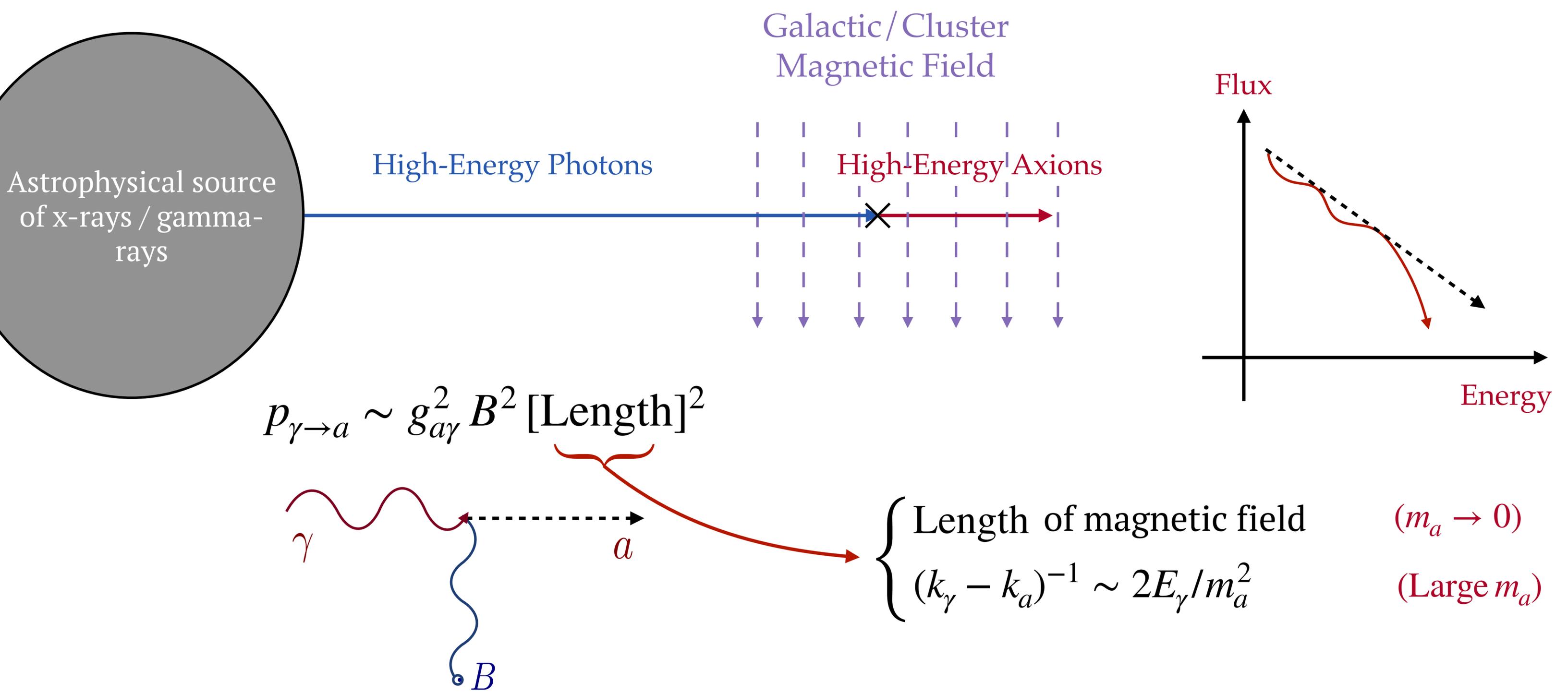
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X-ray / Gamma-ray searches for axions

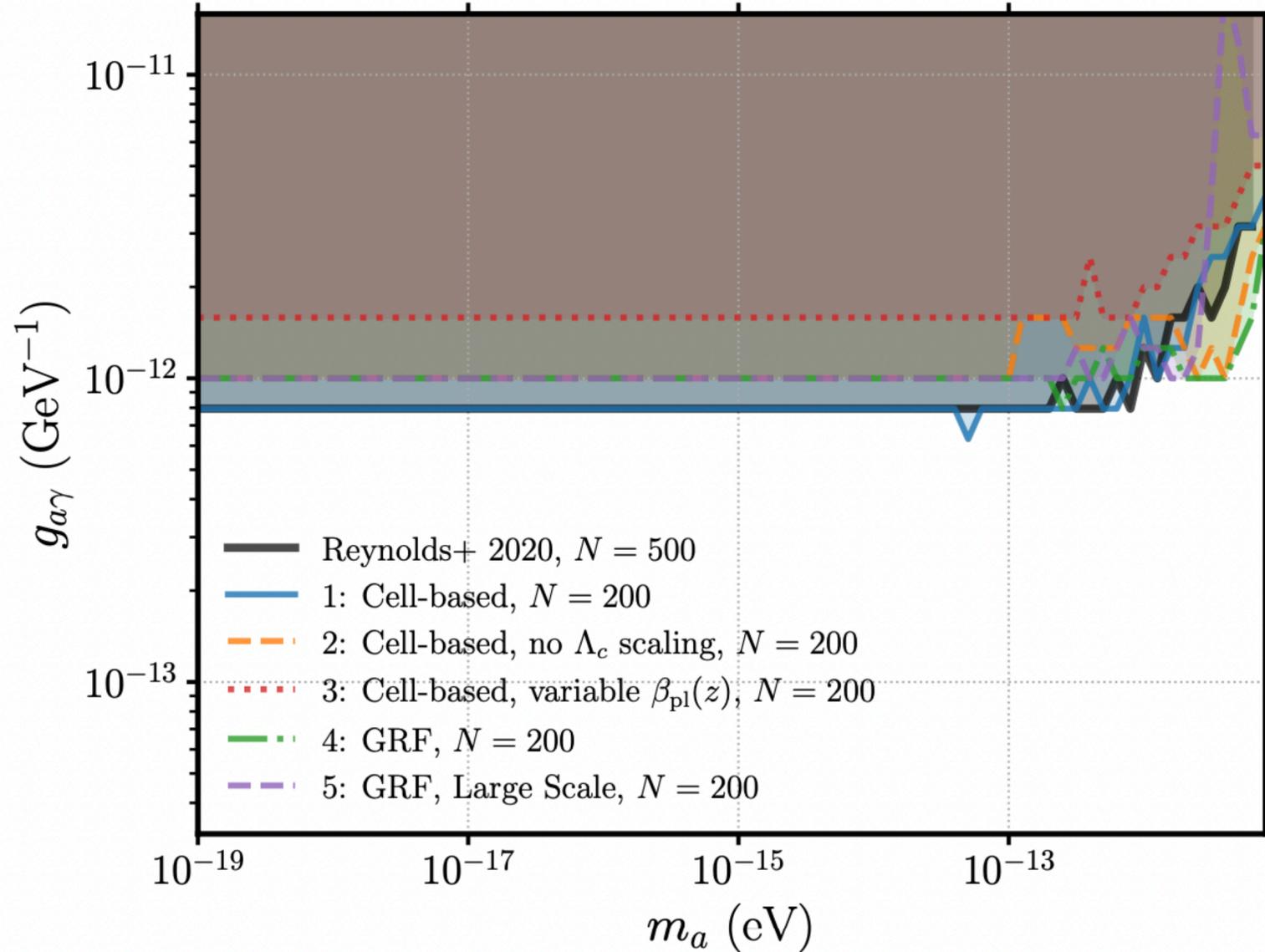


X-ray / Gamma-ray searches for axions

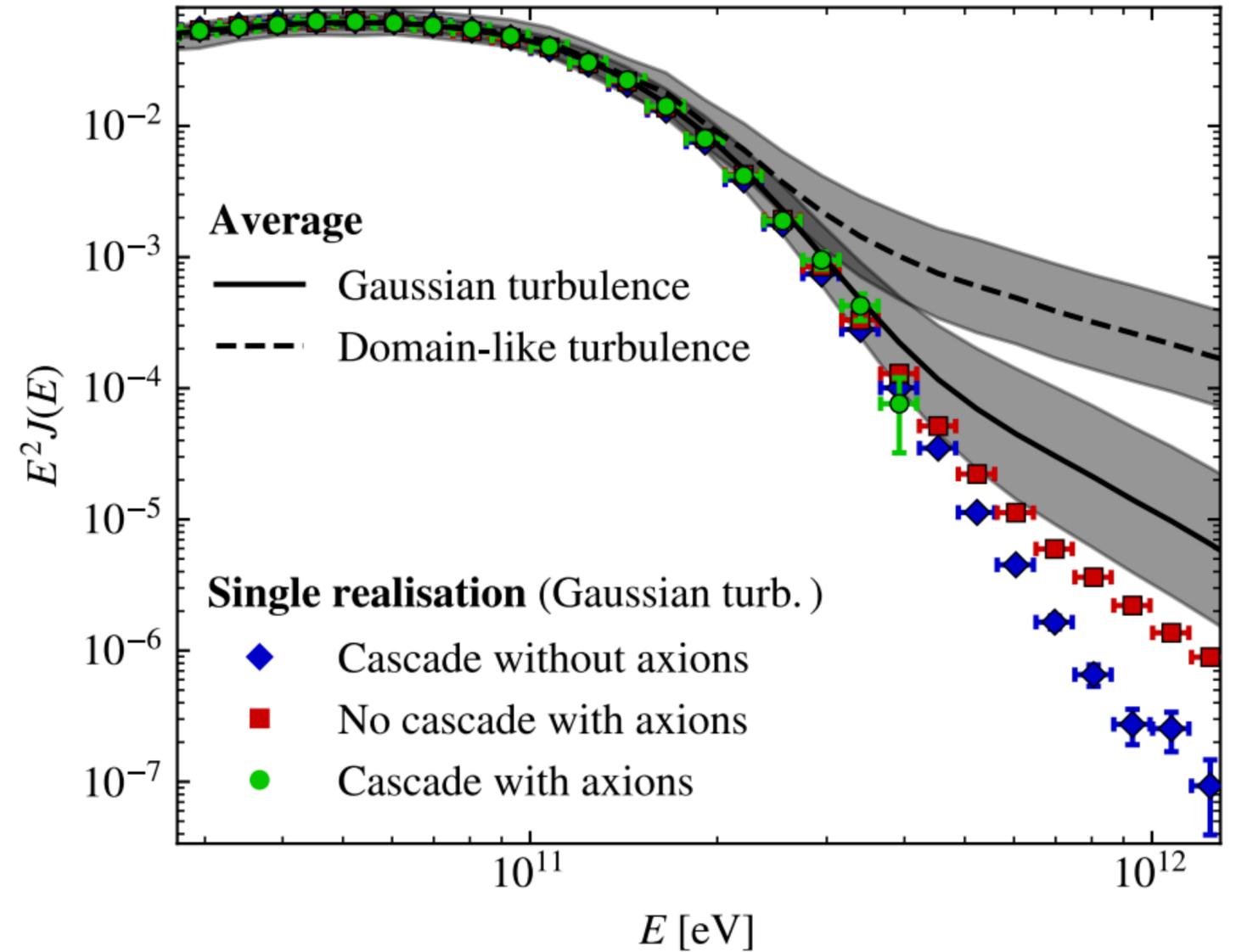


Magnetic fields, ugh....

See e.g. Matthews et al (2022), Carena et al (2023)



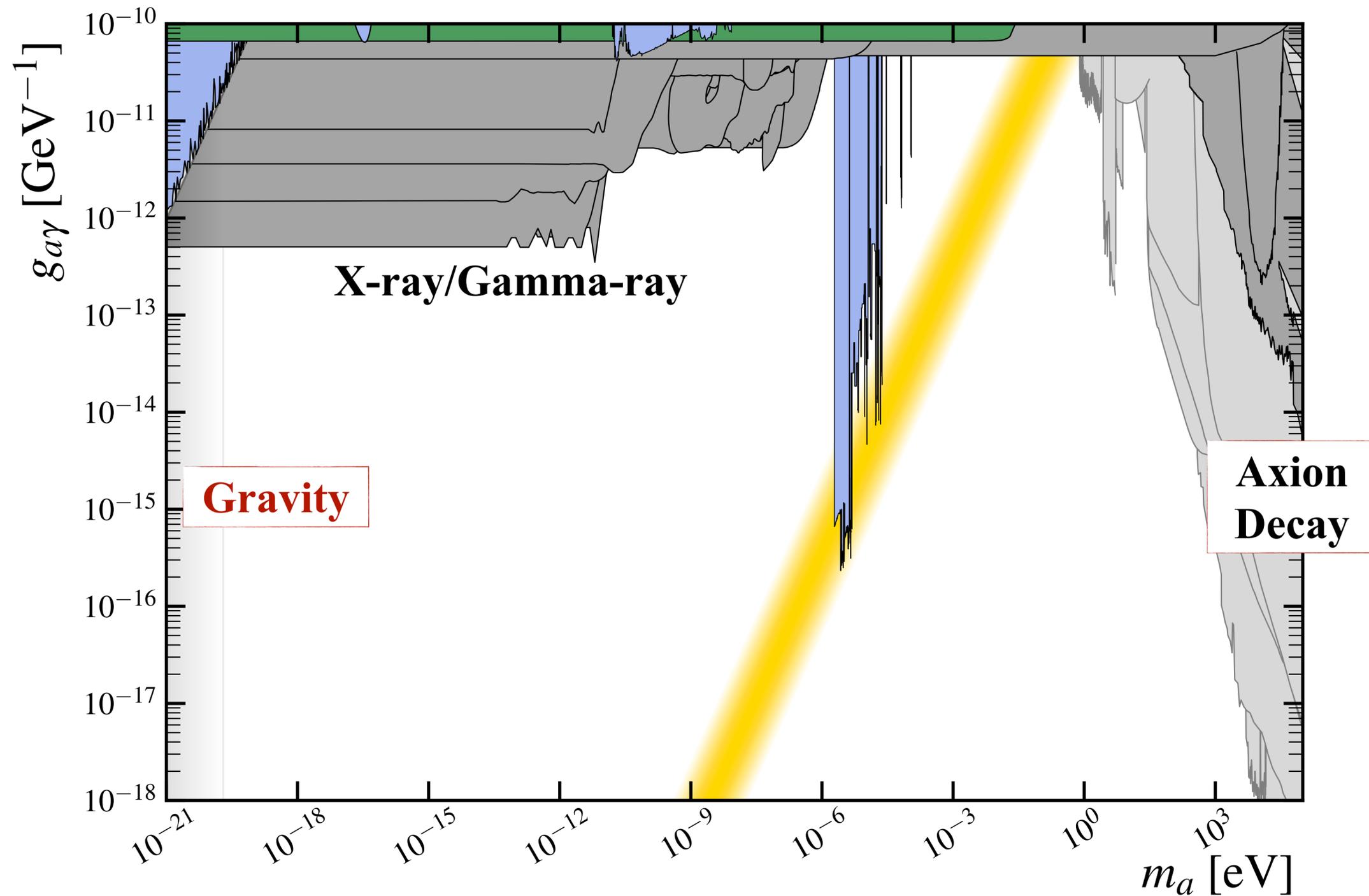
Kachelrieß & Tjemsland (2022)



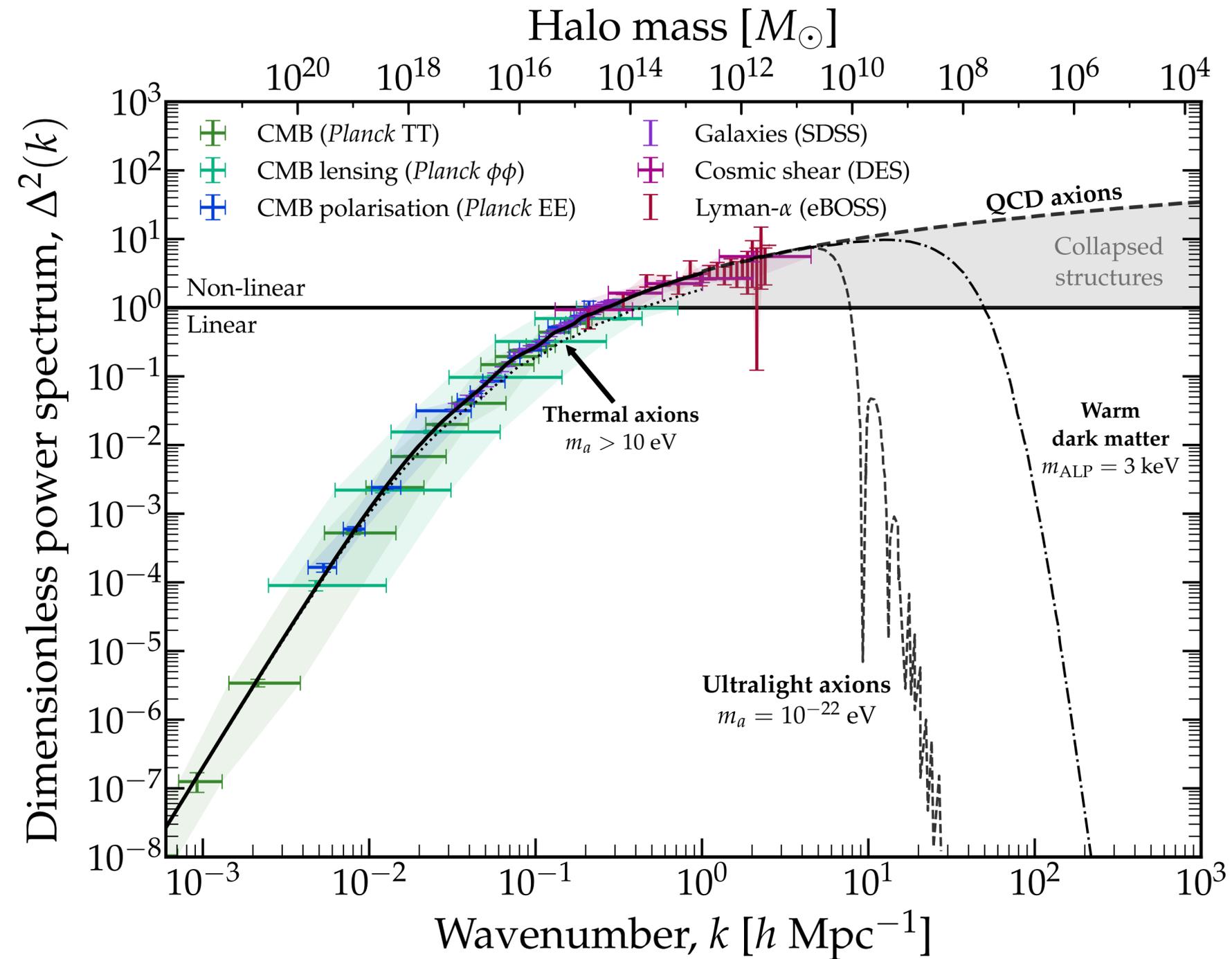
Upside: Reasonably straight-forward physics

Difficulty: Large-scale magnetic fields (*progress limited by finding idealised systems that we understand*)

Axions & Gravity



The gravitational footprint of ultralight axions



Gravitational evolution of classical field

$$\ddot{\delta}_k + 2H\dot{\delta}_k + \left(\frac{k^4}{4m^2 a^2} - 4\pi G\bar{\rho} \right) \delta_k = 0$$

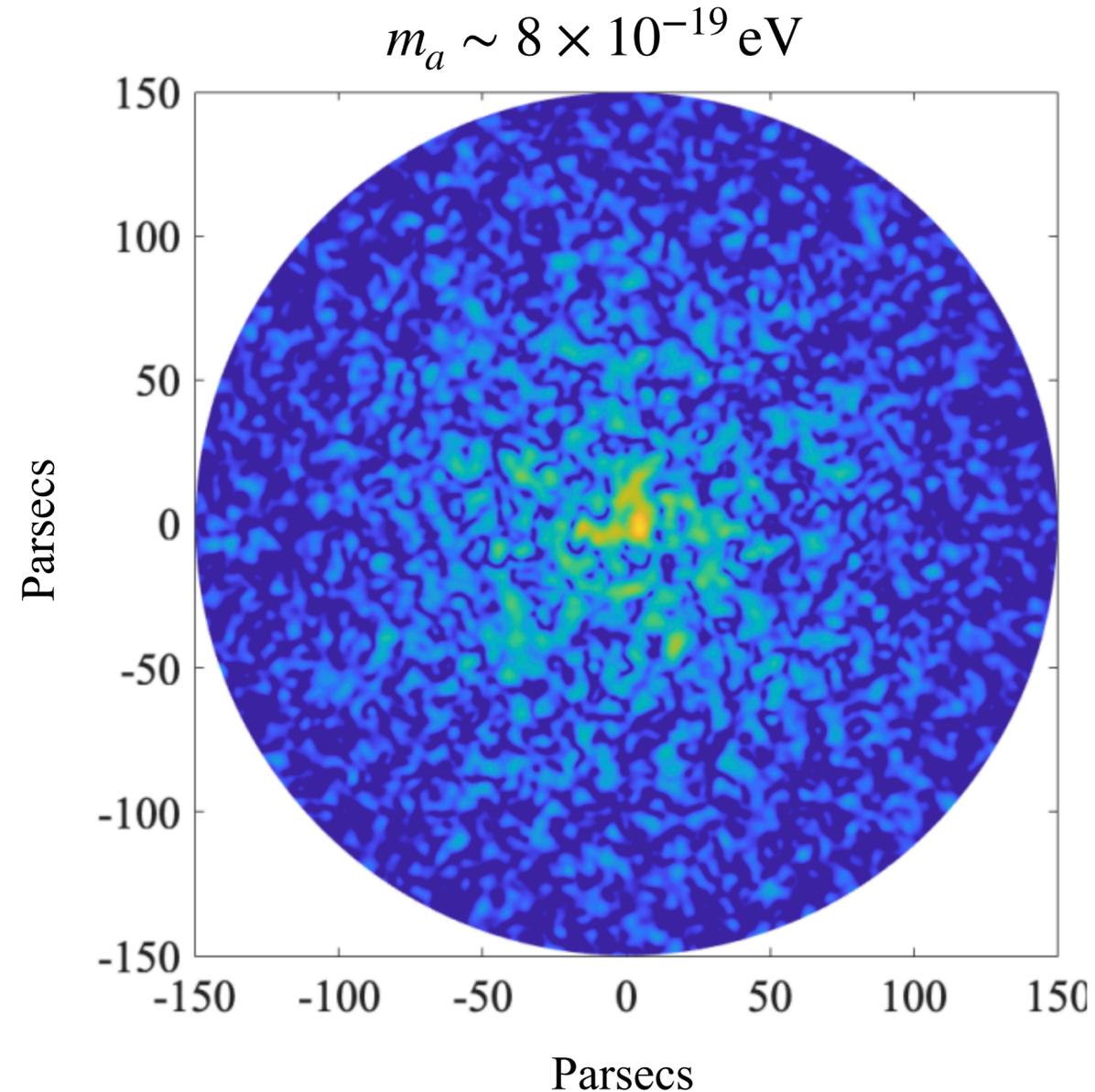
Gradient ("quantum")
pressure

Quantum mechanics limits "packing"
of low mass particles

$$\delta x \times \delta v \gtrsim m^{-1}$$

Figure credit: O'Hare (2024)

The gravitational footprint of ultralight axions



**Density profile
ultra-faint dwarf**

Implications:

- Soliton core at center
- Heat stellar orbits

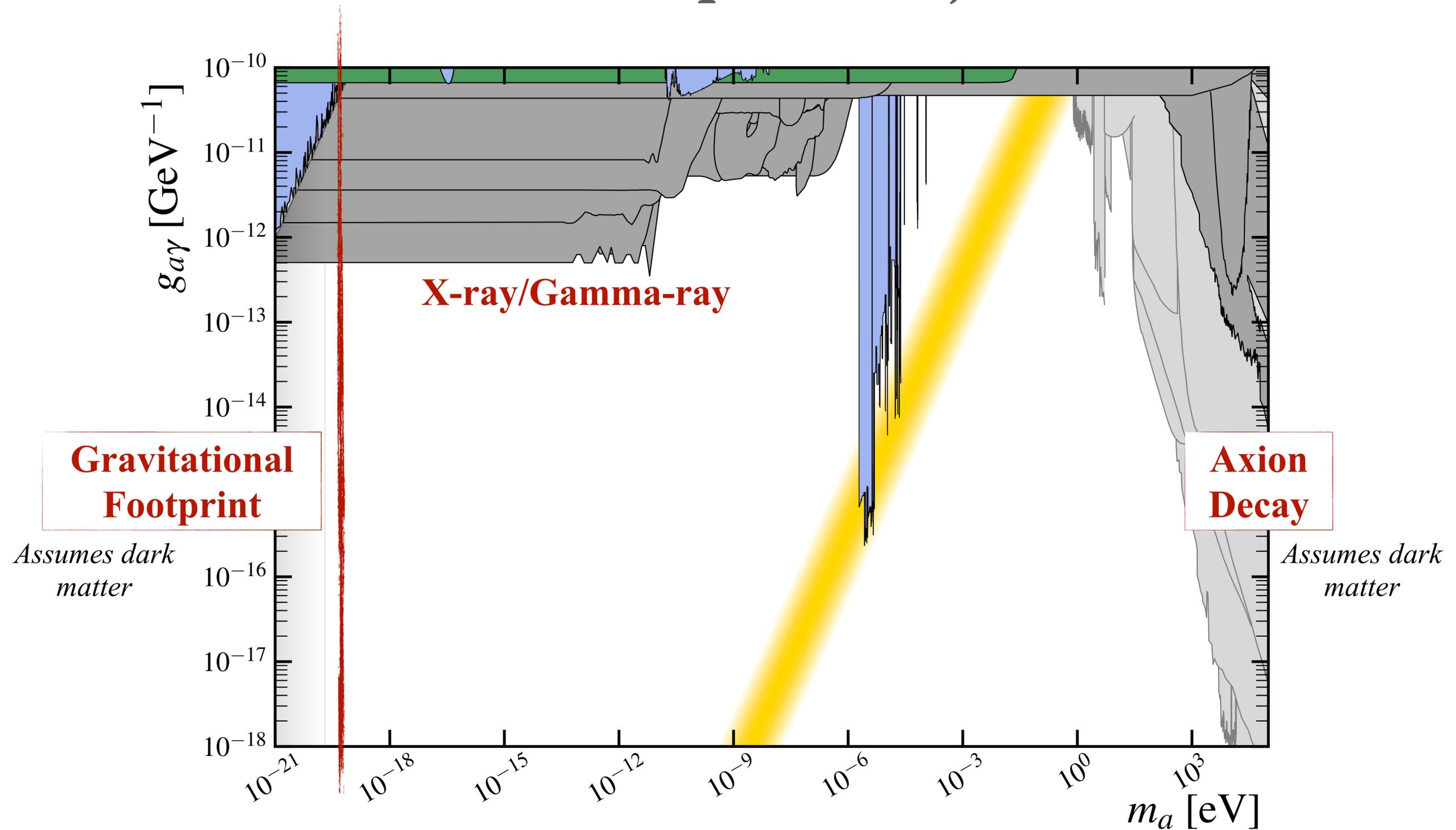
Upside: Purely gravitational

Difficulty: Modelling small scales / feebly bound objects

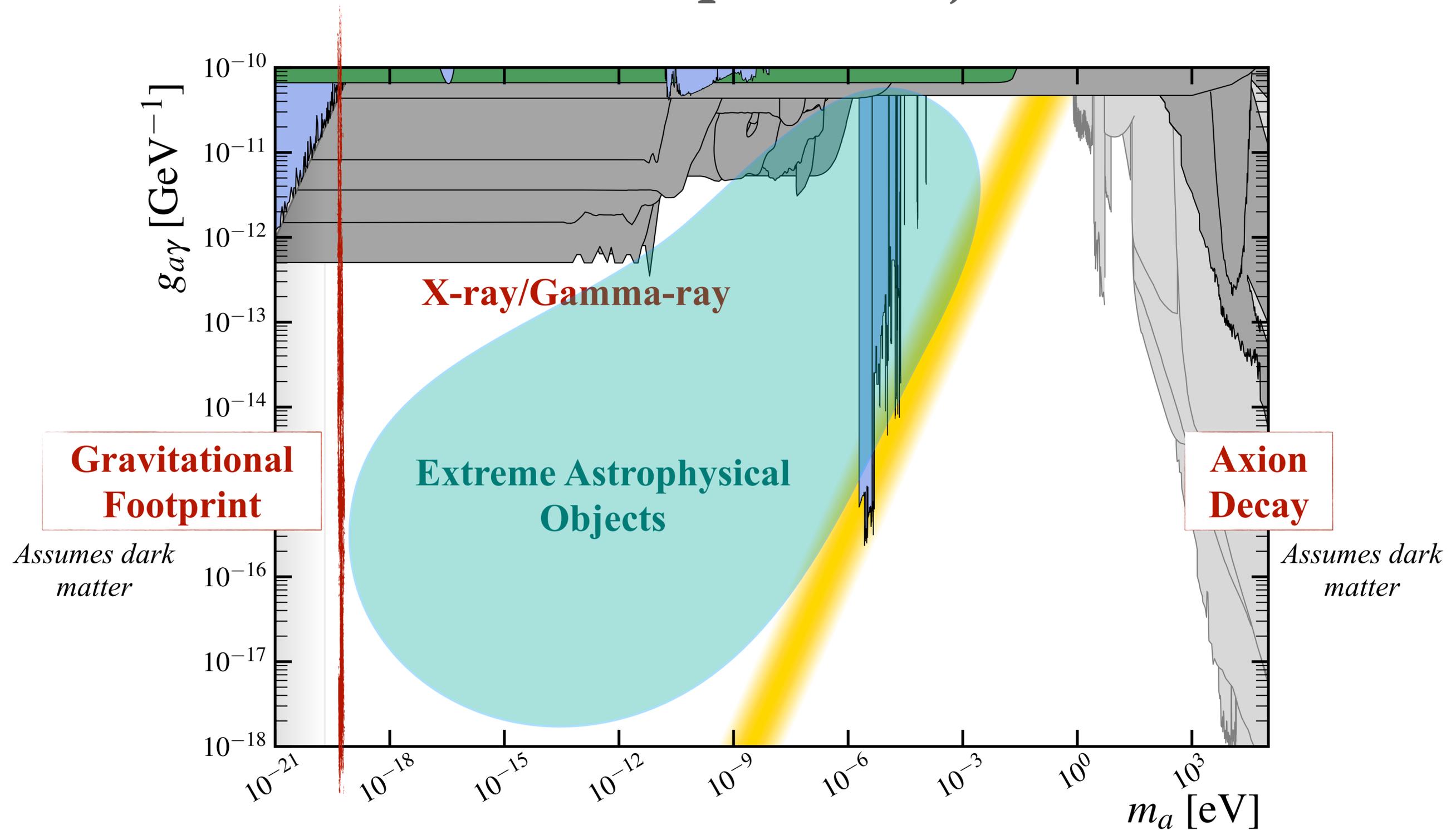
Image credit: Dalal & Kravtsov (2022)

See e.g. reviews by Hui (2014), Niemeyer (2020), Ferreira (2021), O'Hare (2024)

Axions near extreme compact objects

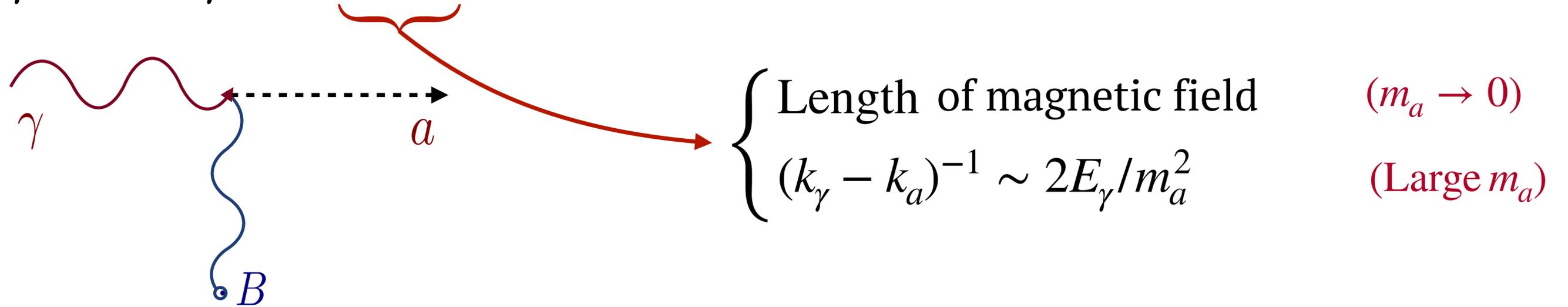


Axions near extreme compact objects



Enhancing axion-photon transitions

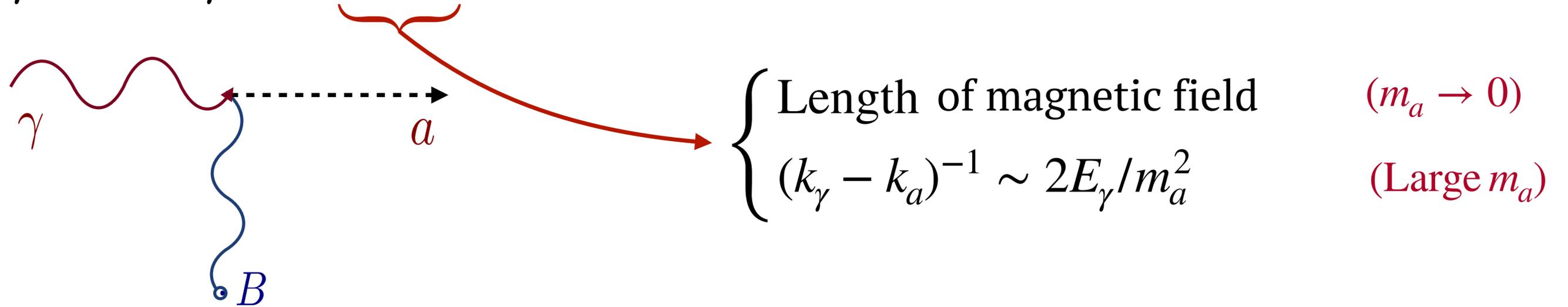
$$P_{\gamma \rightarrow a} \sim g_{a\gamma}^2 B^2 [\text{Length}]^2$$



Limitation: Strength of \mathbf{B} and $(k_\gamma - k_a)$

Enhancing axion-photon transitions

$$P_{\gamma \rightarrow a} \sim g_{a\gamma}^2 B^2 [\text{Length}]^2$$



Limitation: Strength of \mathbf{B} and $(k_\gamma - k_a)$

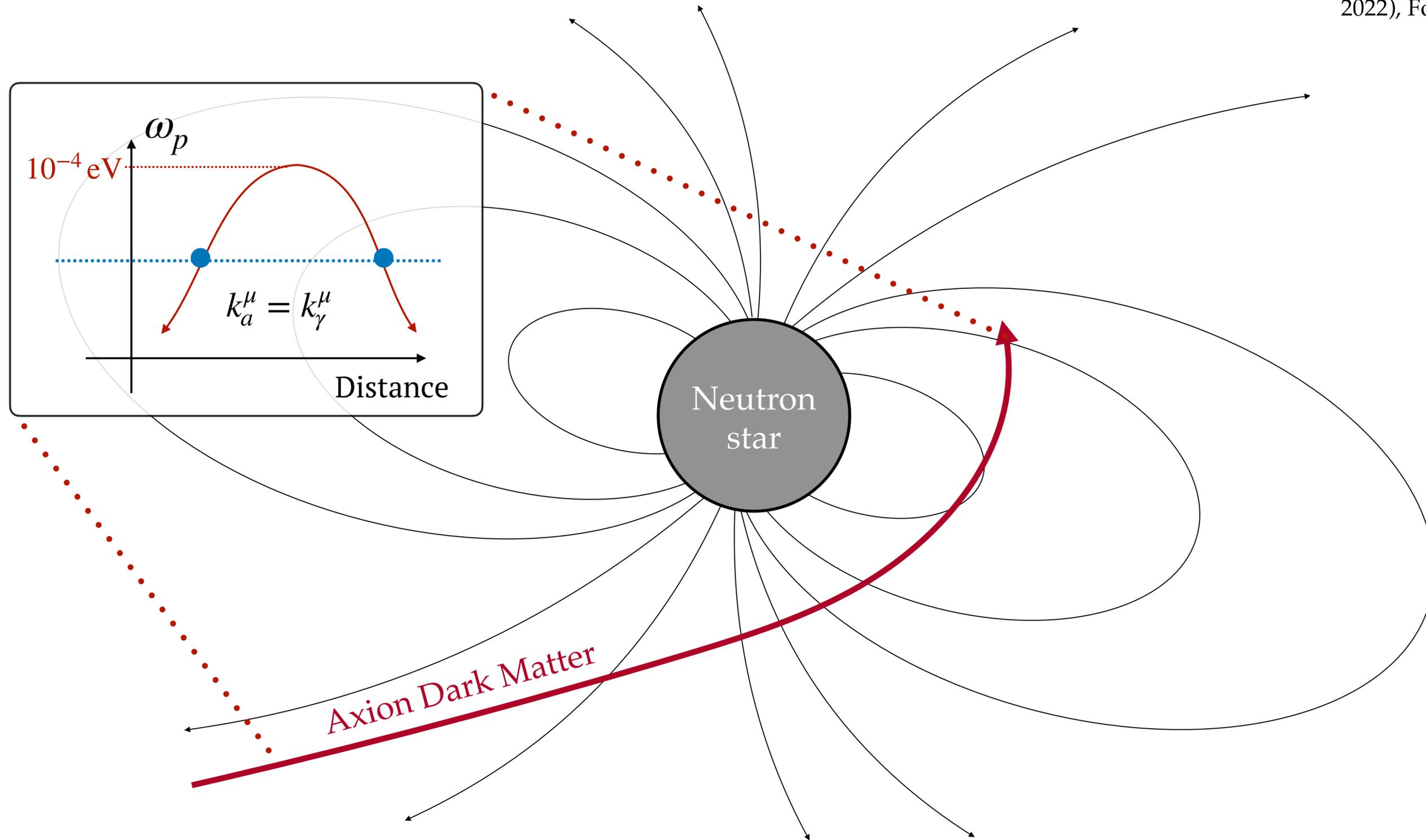
Modify photon dispersion relation

(E.g. in a cold plasma $k_\gamma = \sqrt{\omega^2 - \omega_p^2}$, and $k_a \sim k_\gamma$ possible)

Compact Objects

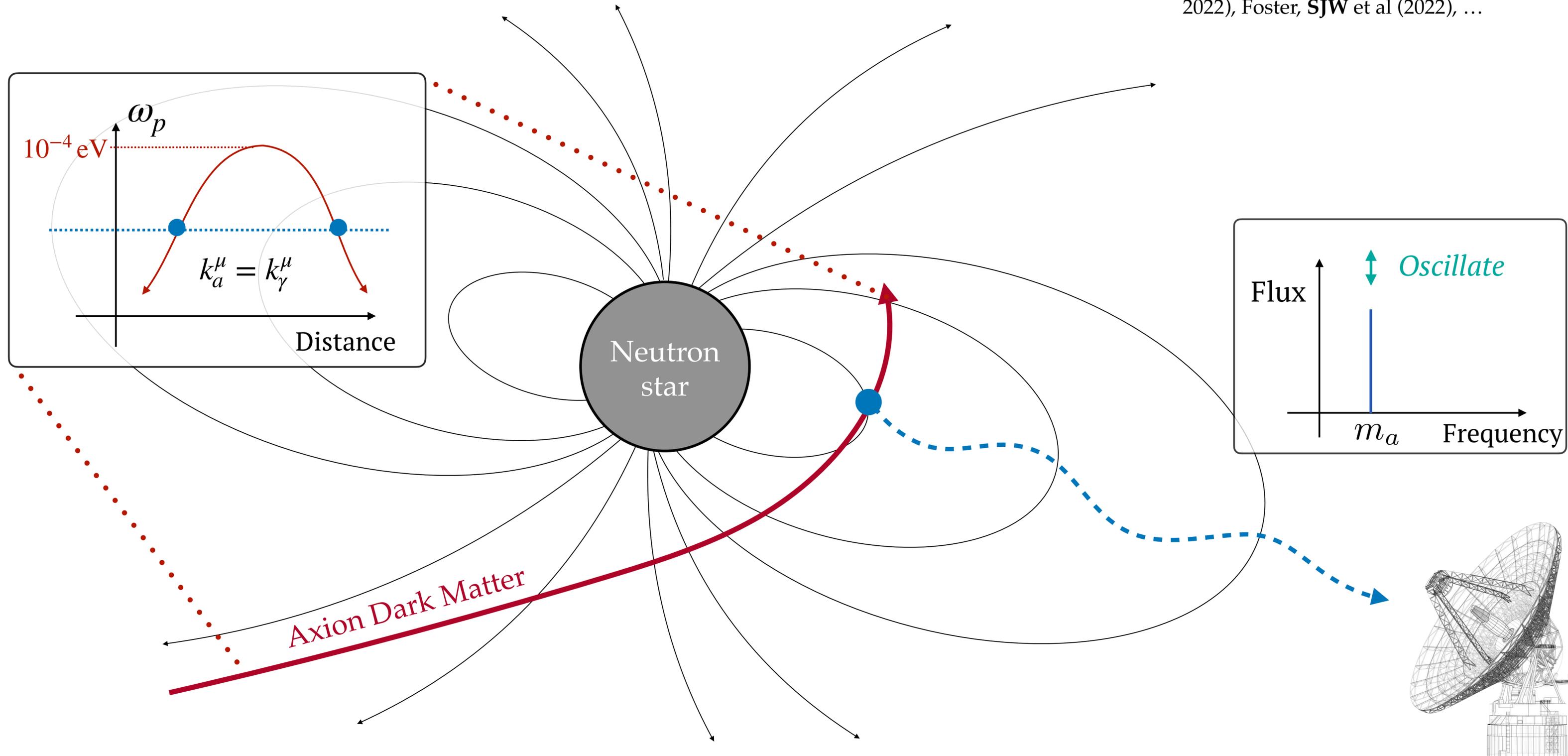
Resonant axion transitions in radio

See e.g.: Pshirkov & Popov (2009), Hook et al. (2018), Safdi et al. (2018), Battye et al. (2019, 2021, 2023), **SJW** et al. (2021, 2022), Foster, **SJW** et al (2022), ...



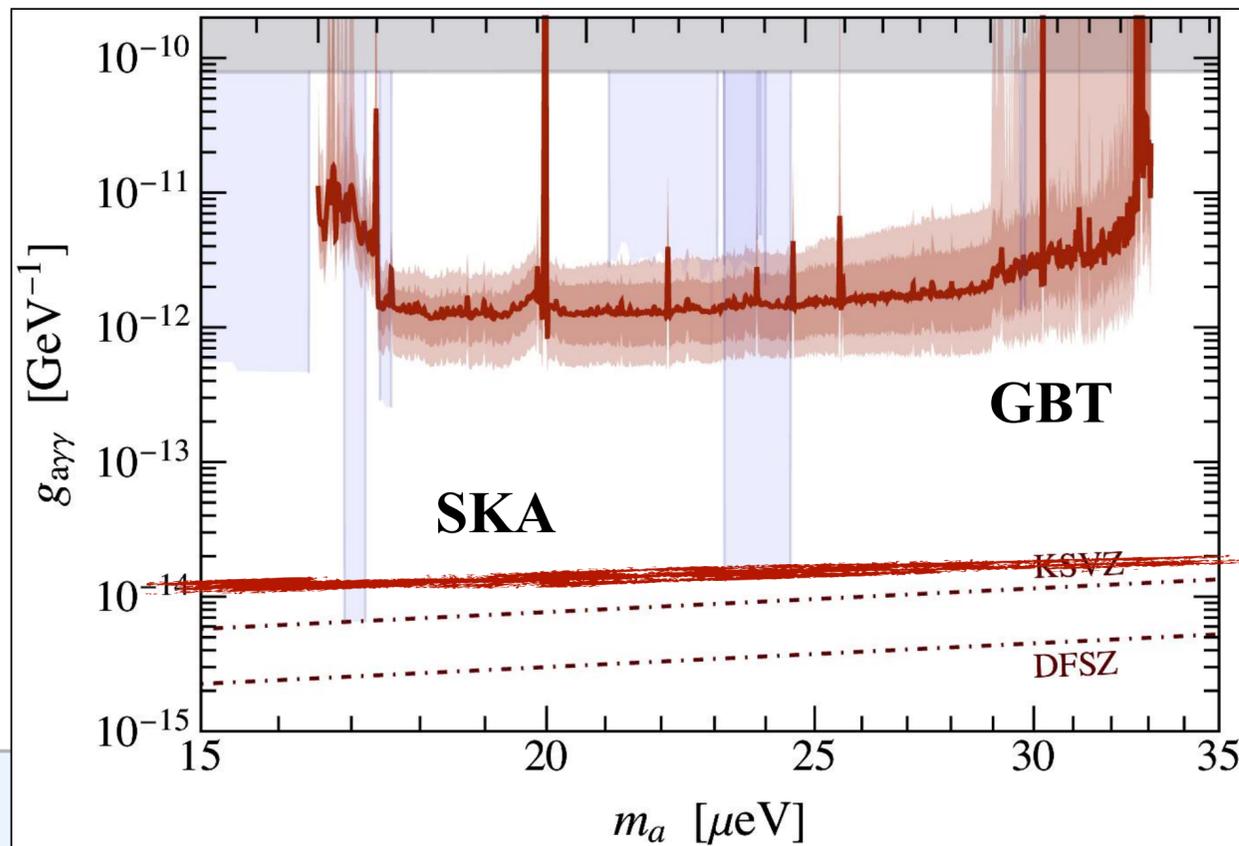
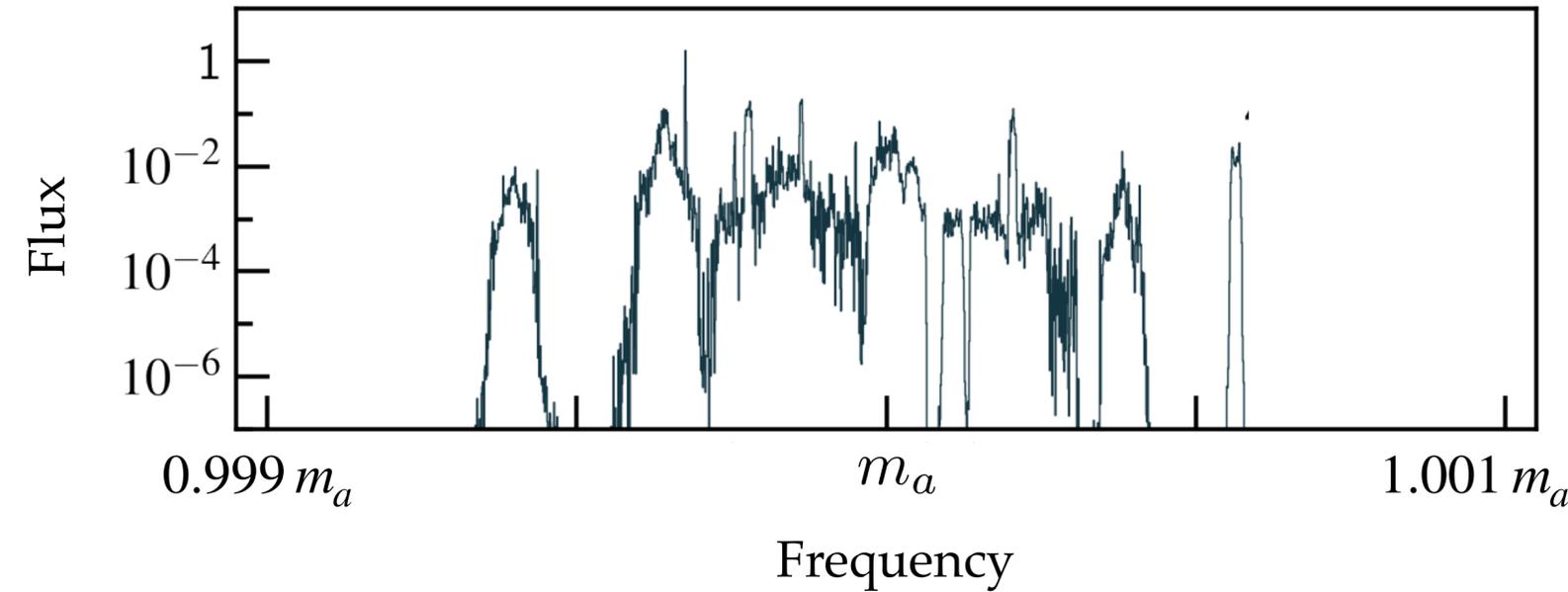
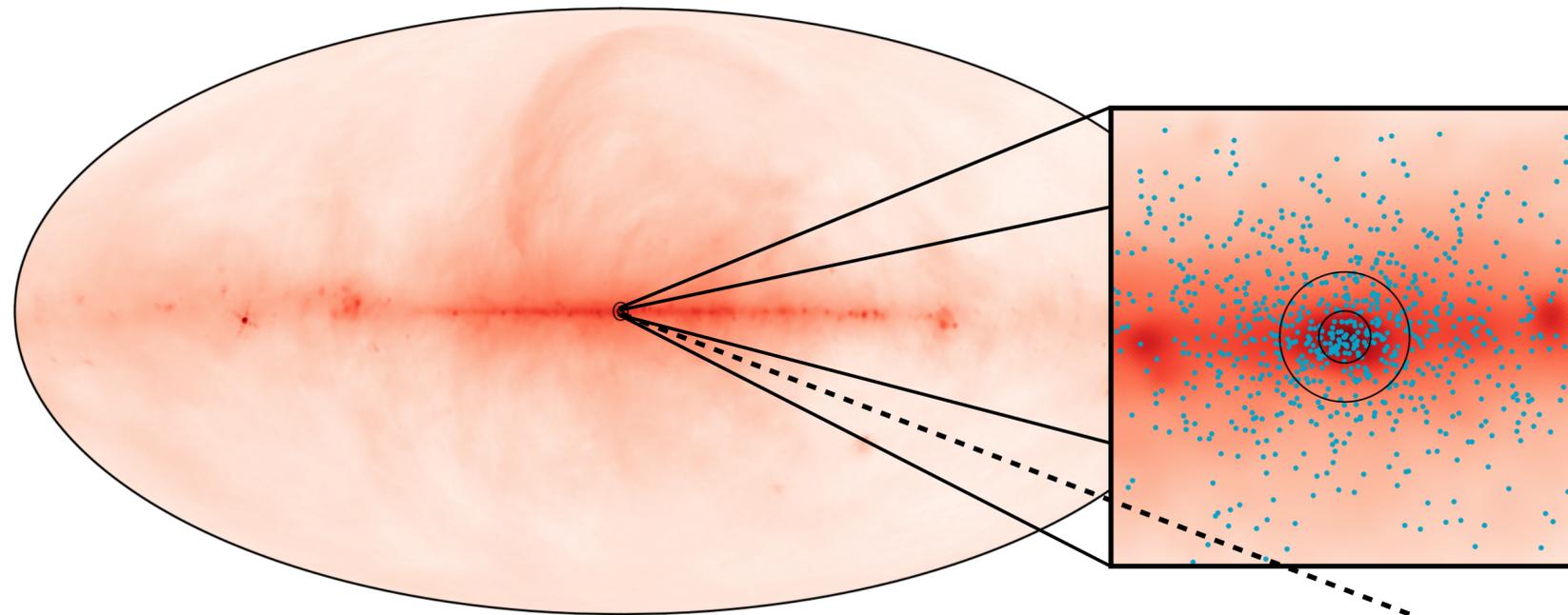
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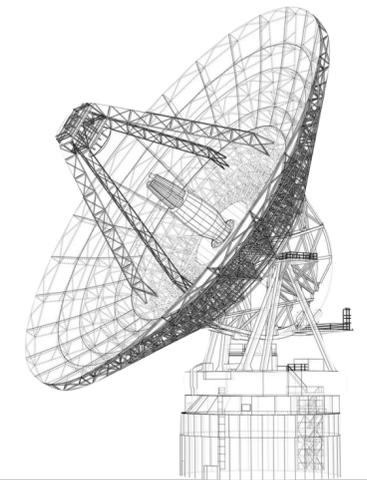


Radio lines near neutron stars

See e.g.: Pshirkov & Popov (2009), Hook et al. (2018), Safdi et al. (2018), Battye et al. (2019, 2021, 2023), SJW et al. (2021, 2022), Foster, SJW et al (2022), ...



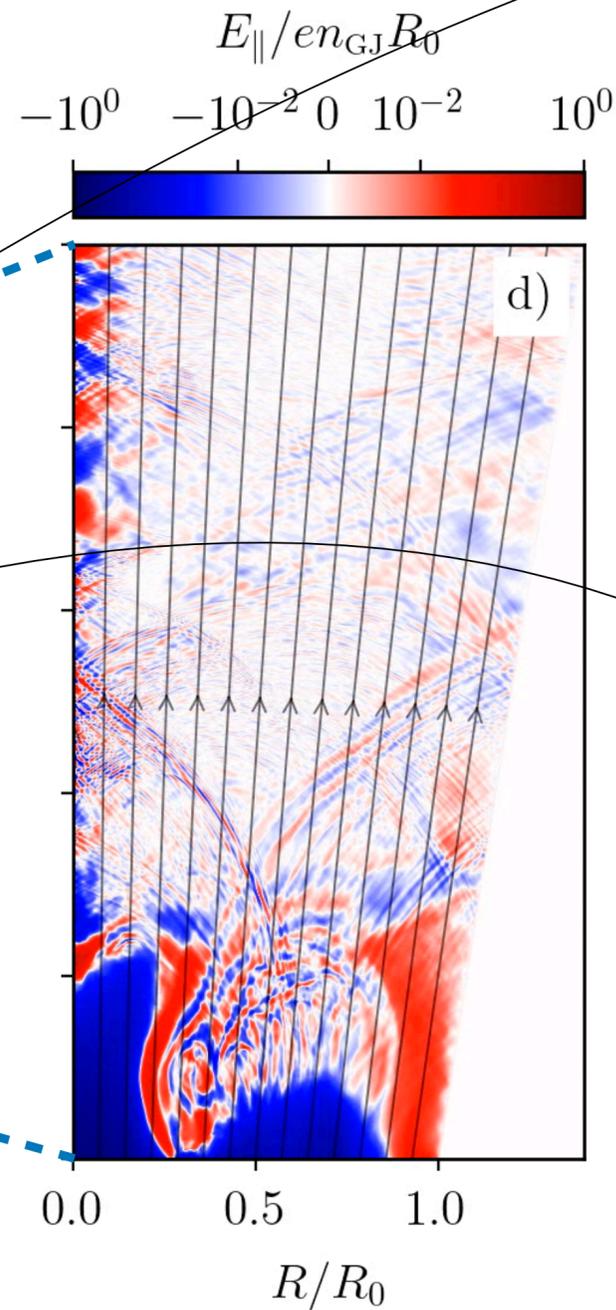
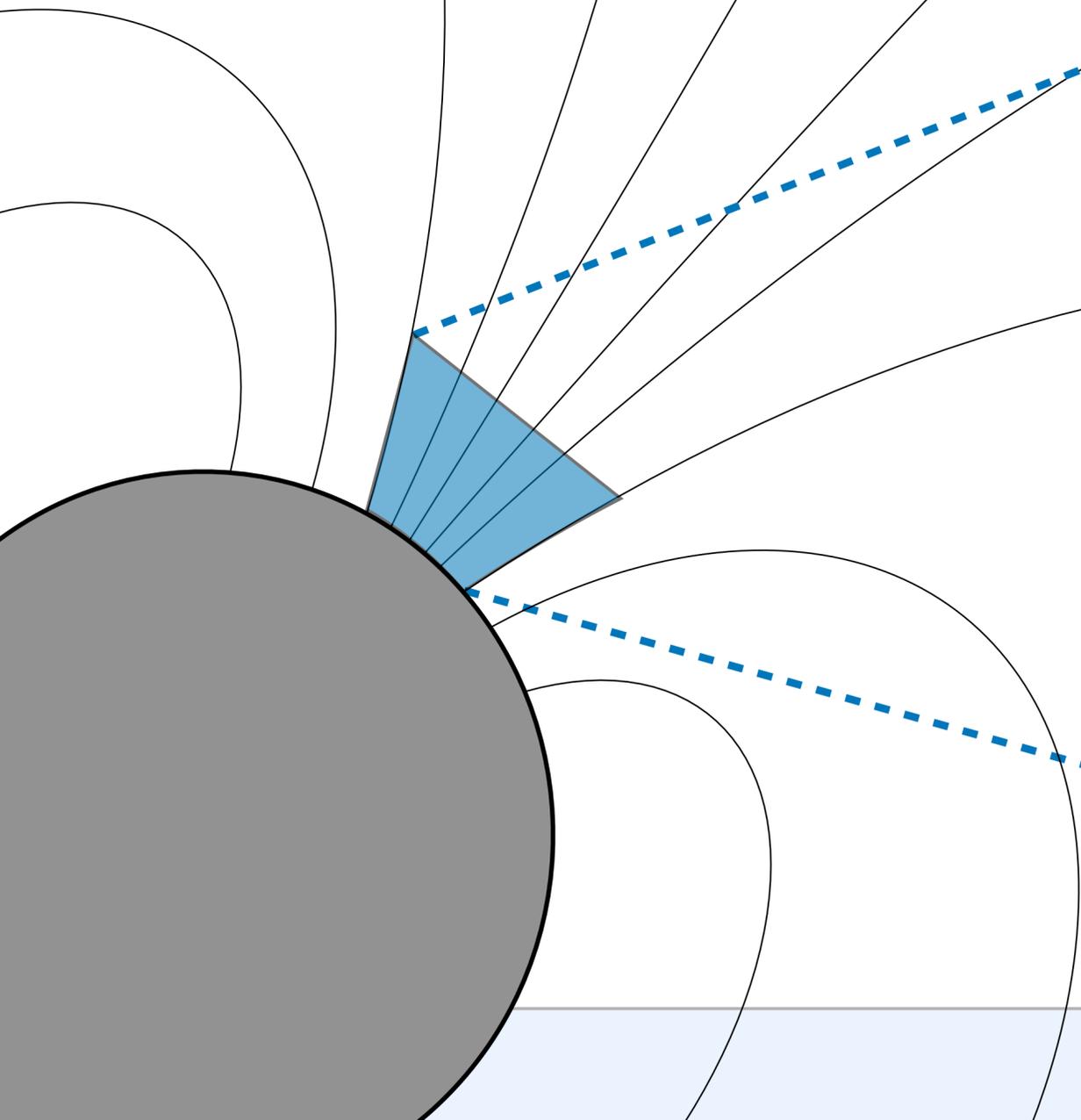
Upside: Unique signatures
Difficulty: Modelling



Axion production in pulsars

Prabhu 2021, Noordhuis, Prabhu, SJW, Cruz, Chen, Weniger (2022), Noordhuis, Prabhu, Weniger, SJW (2023), Caputo, SJW, Philippov, Jacobson (2023), Khelashvili et al (2024)

$$(\square + m_a^2) a = g_{a\gamma\gamma} \vec{E} \cdot \vec{B}$$



Rotating \vec{B} induces \vec{E}

↓

Plasma cannot screen \vec{E} everywhere

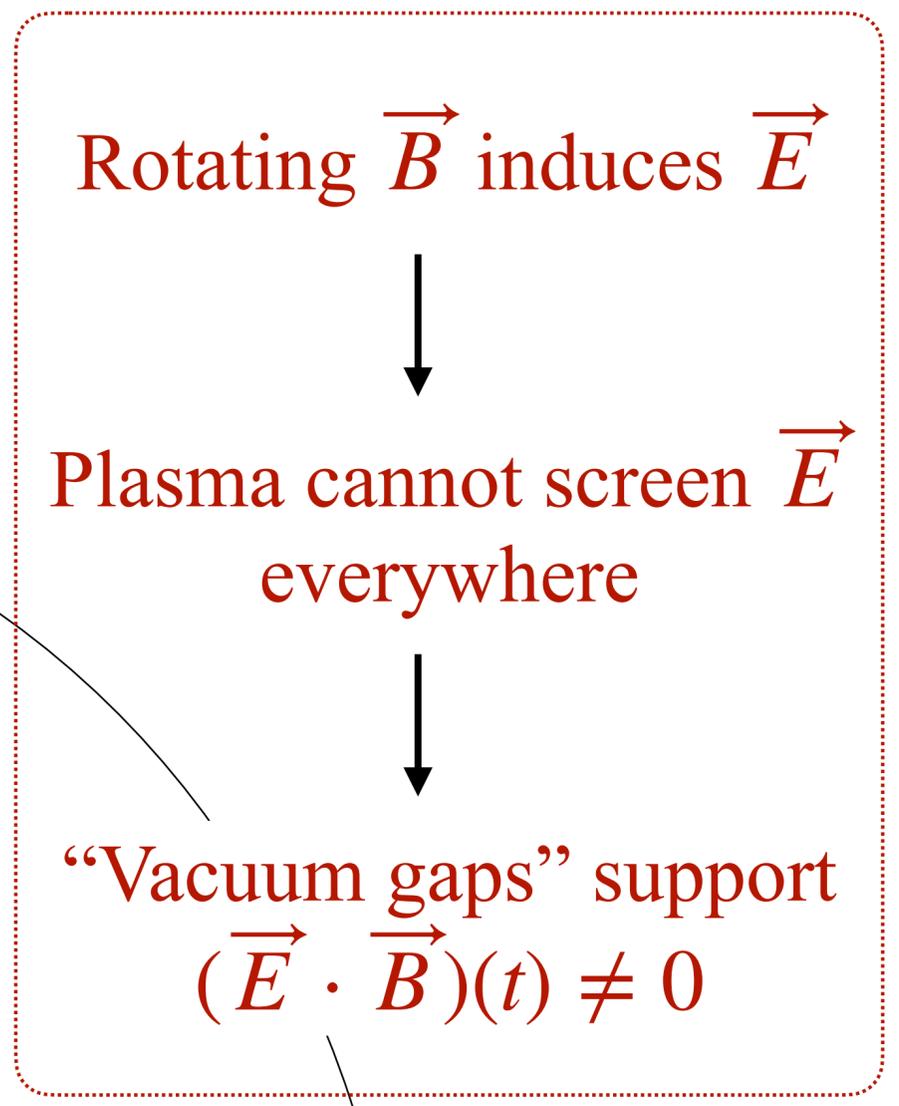
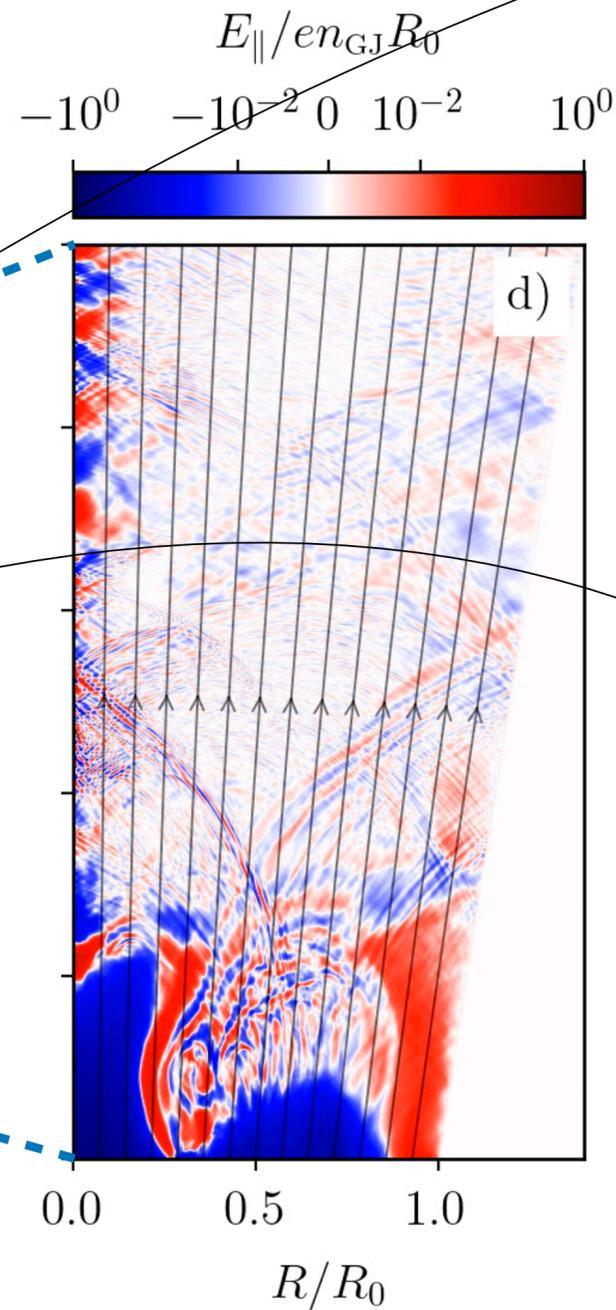
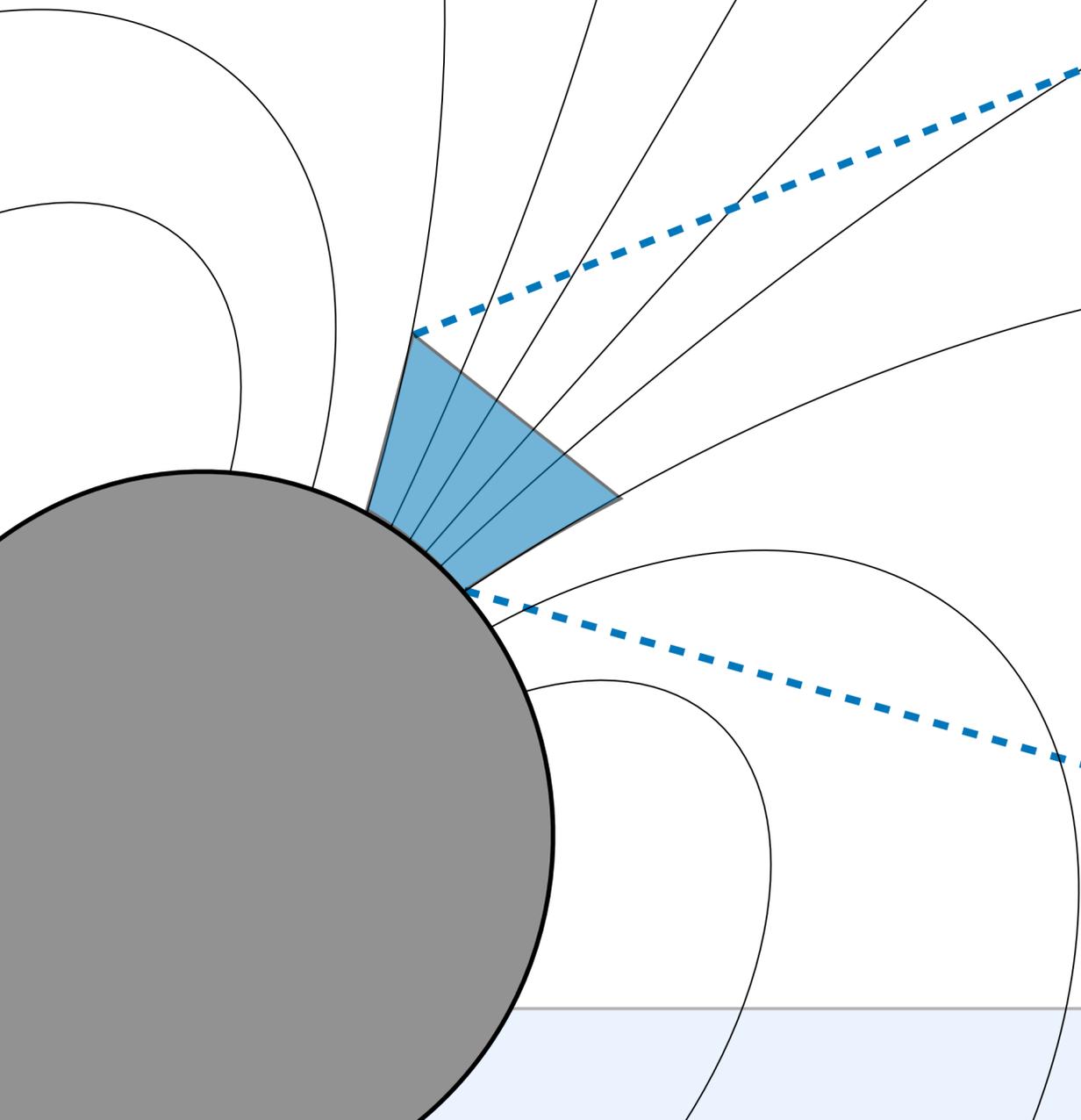
↓

“Vacuum gaps” support $(\vec{E} \cdot \vec{B})(t) \neq 0$

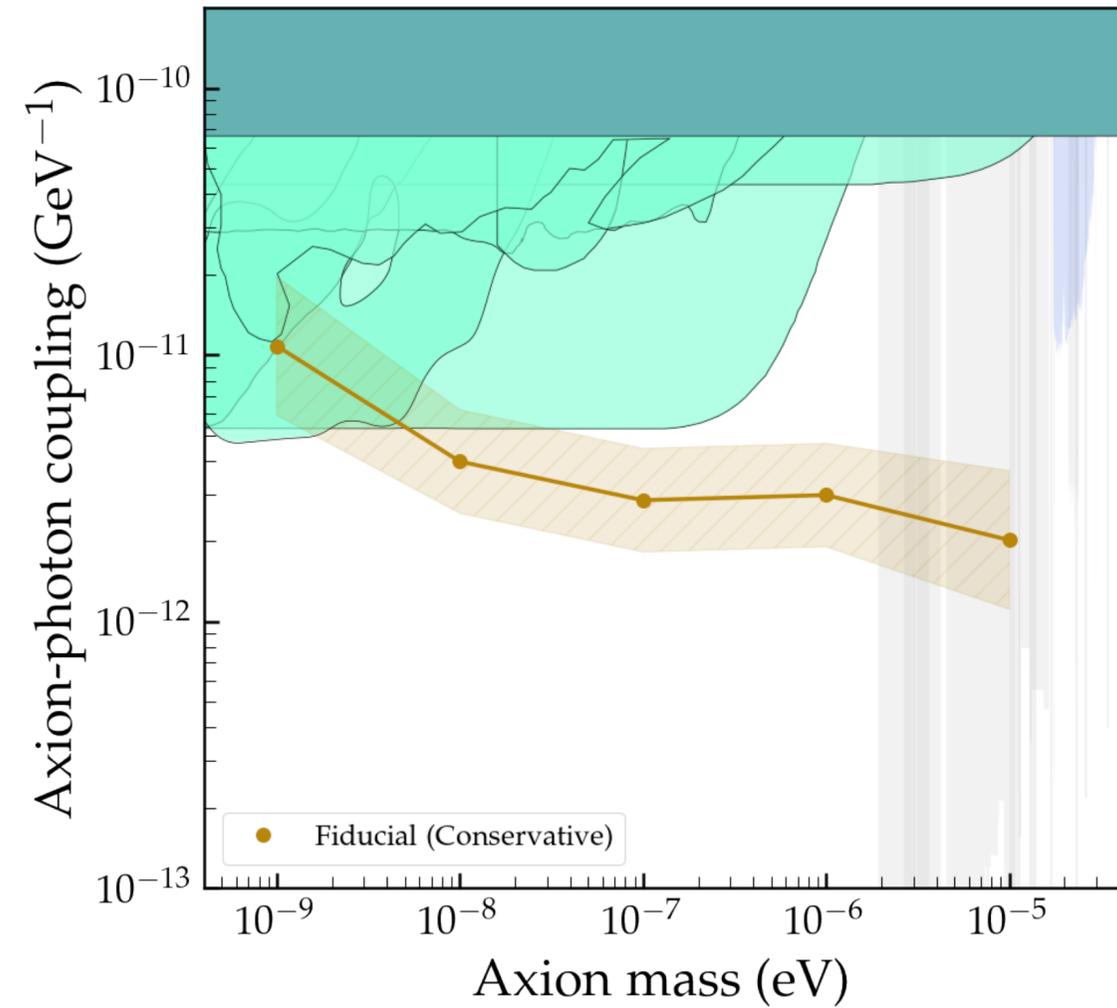
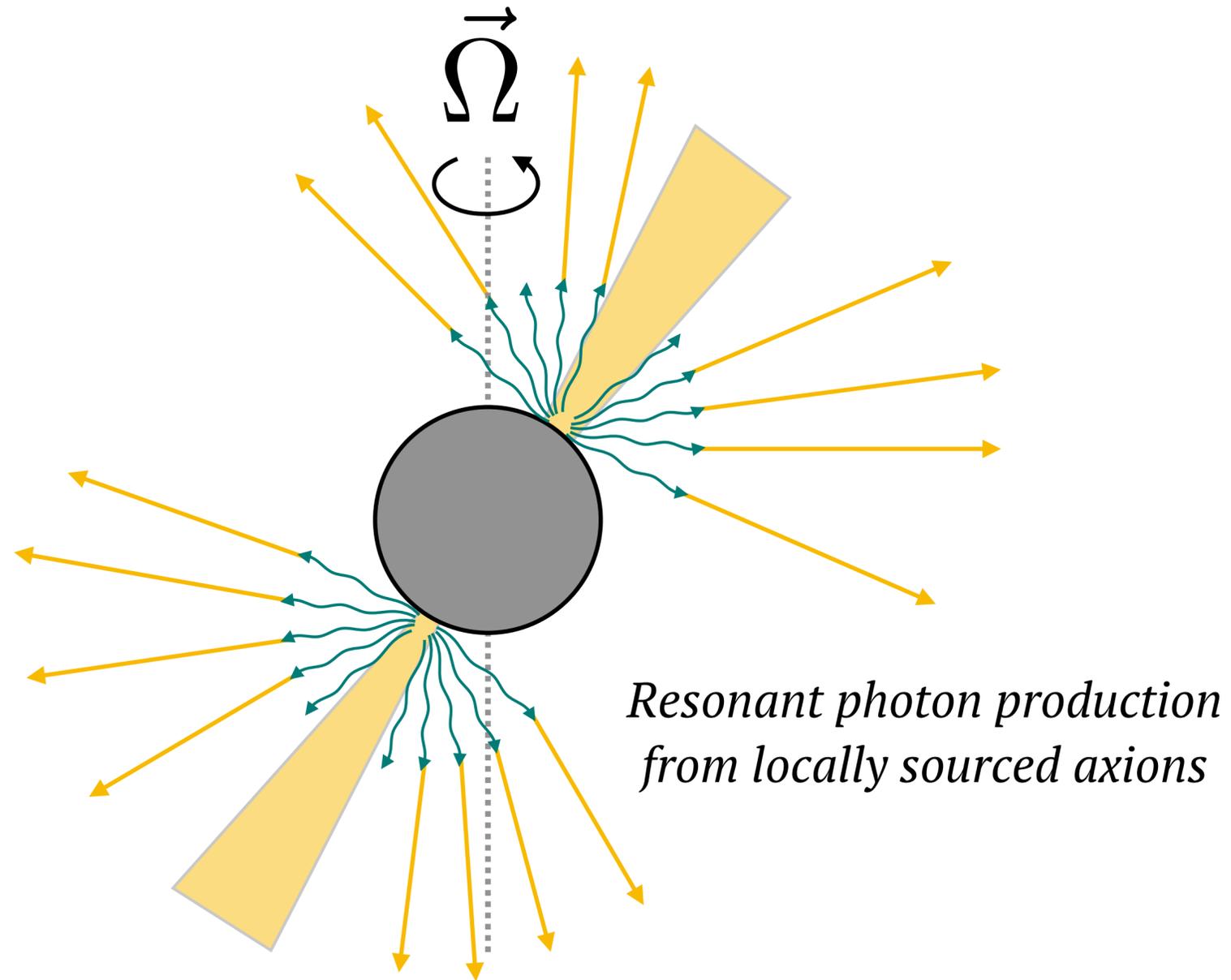
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$$(\square + m_a^2) a = g_{a\gamma\gamma} \vec{E} \cdot \vec{B}$$

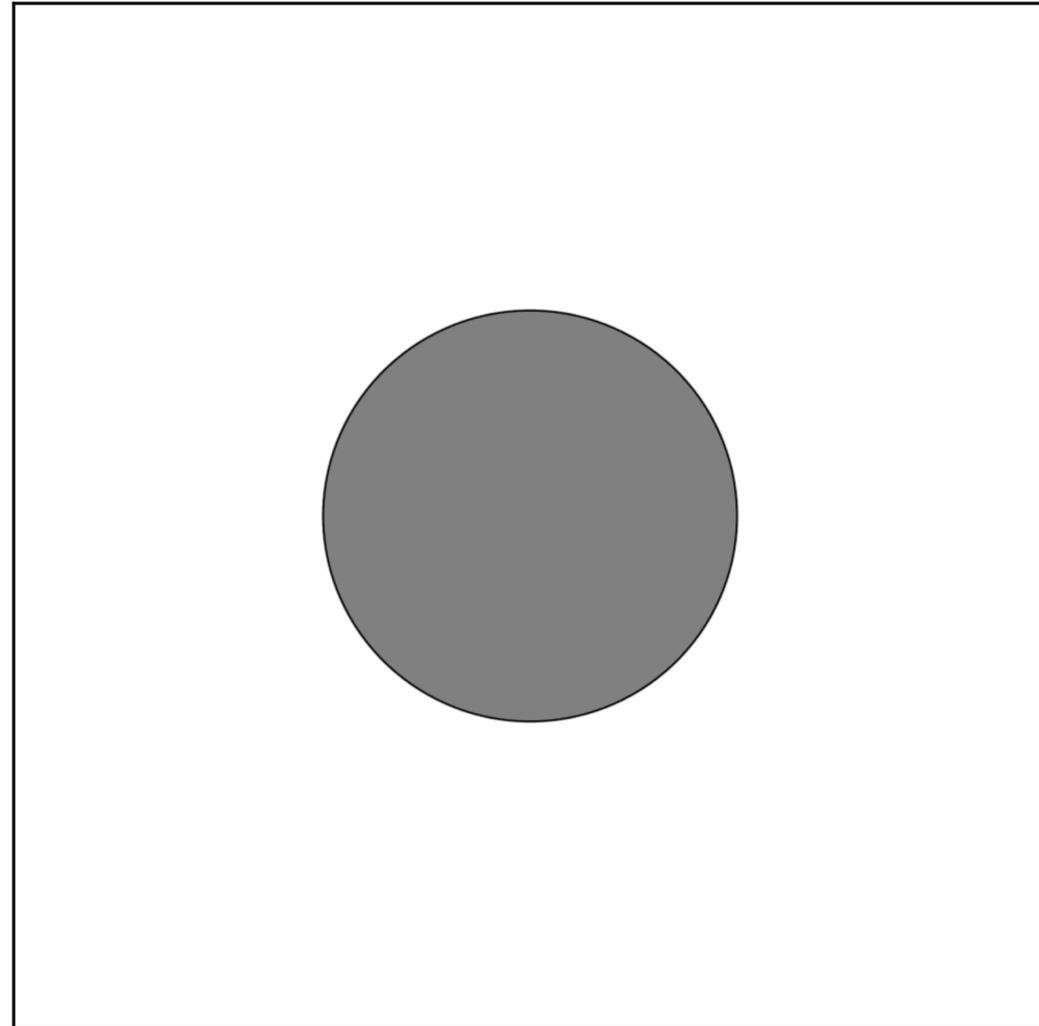


Axion production in pulsars

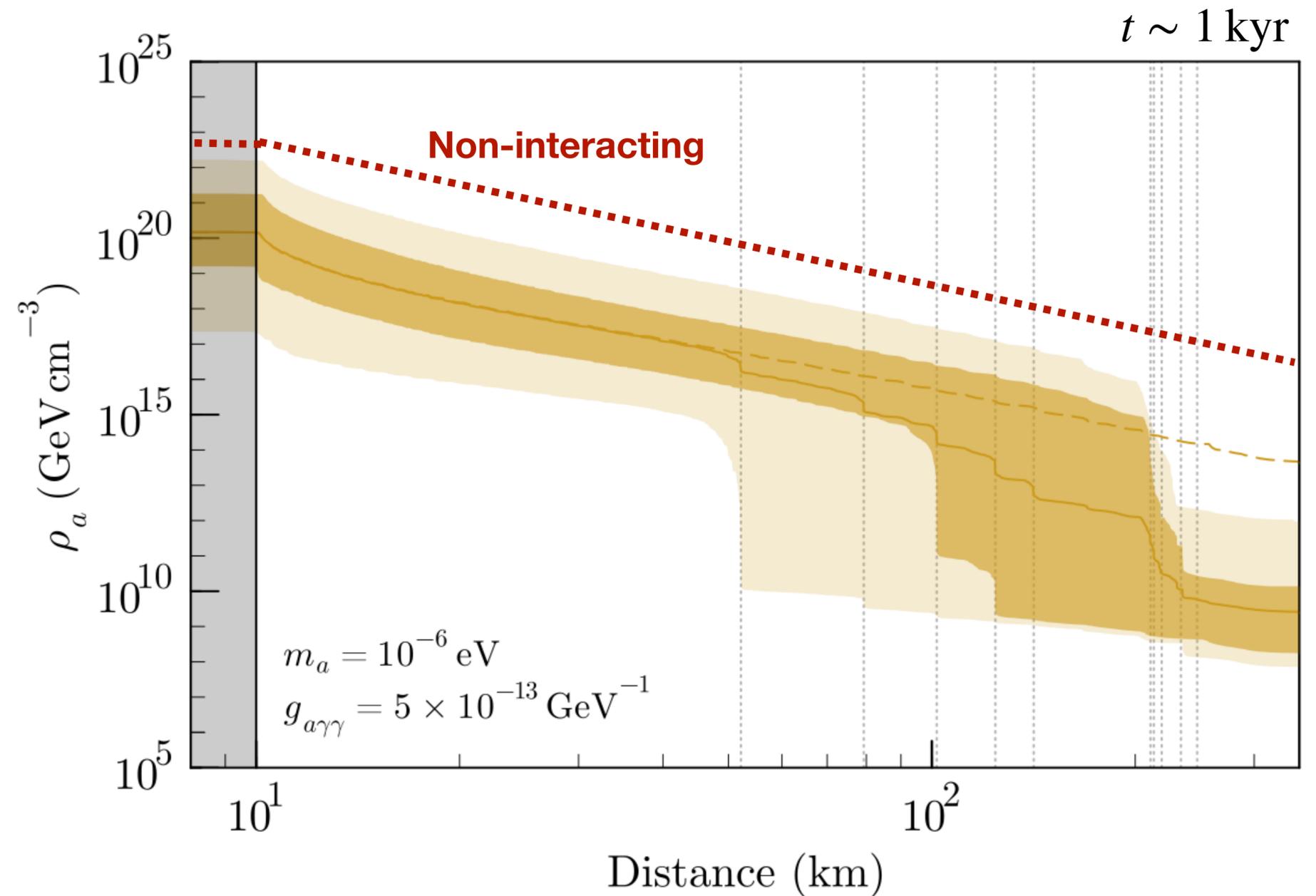


Upside: No assumption of dark matter
Difficulty: Modelling more difficult
Difficulty/Upside: Emission scales like $g_{a\gamma}^4$

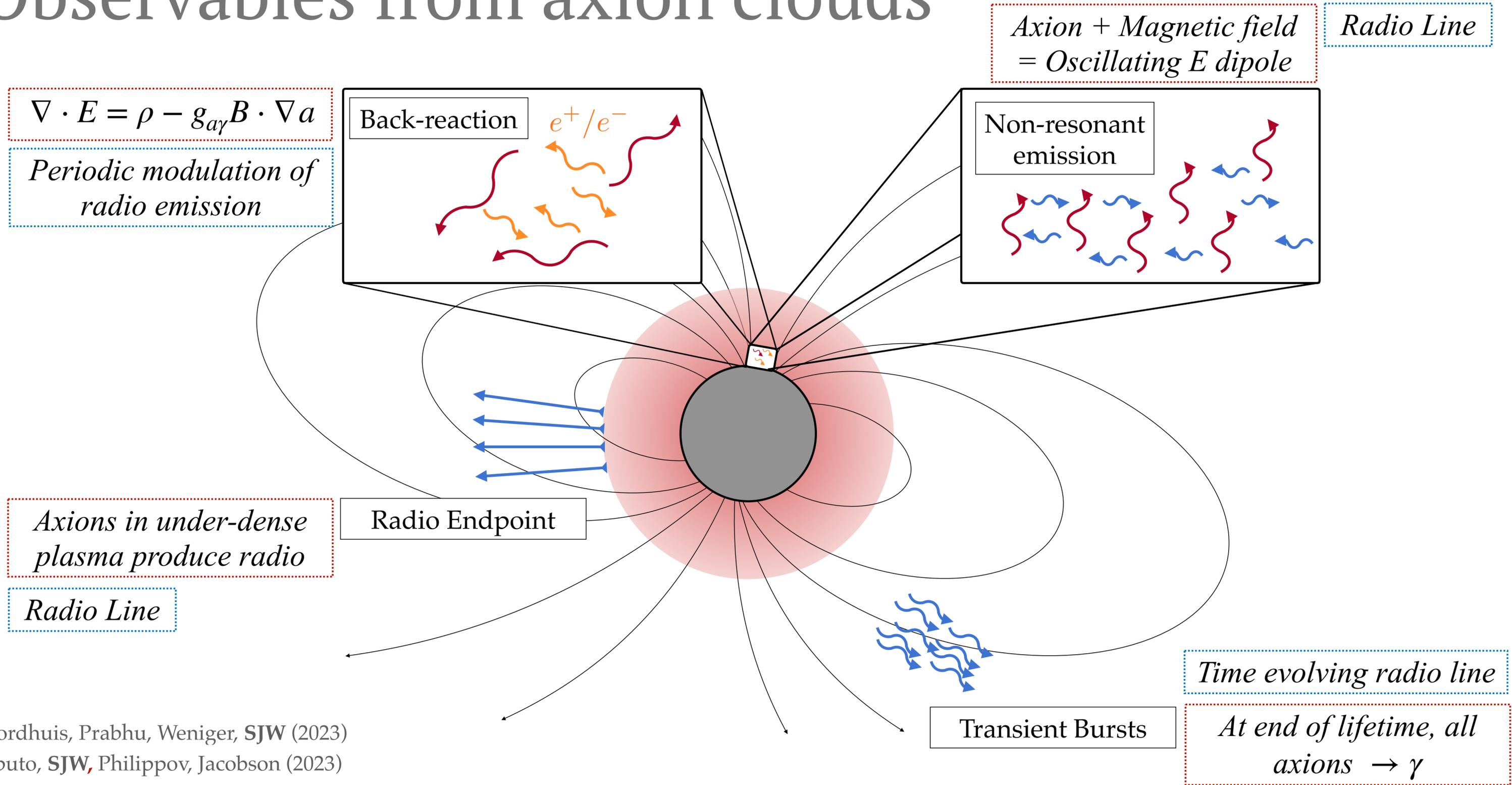
Axion clouds around pulsars



Axion mass
 $10^{-10} \text{ eV} \lesssim m_a \lesssim 10^{-4} \text{ eV}$



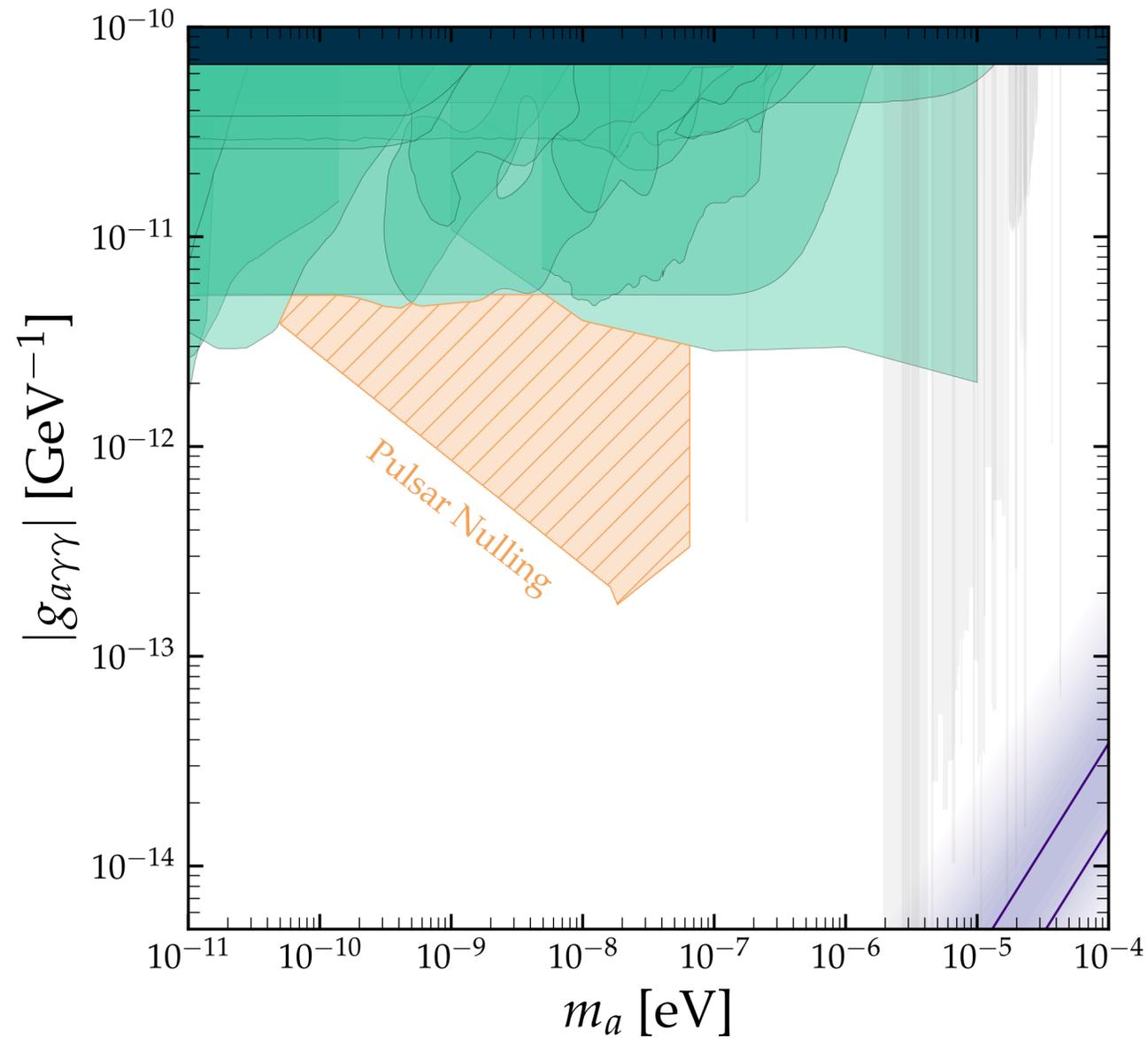
Observables from axion clouds



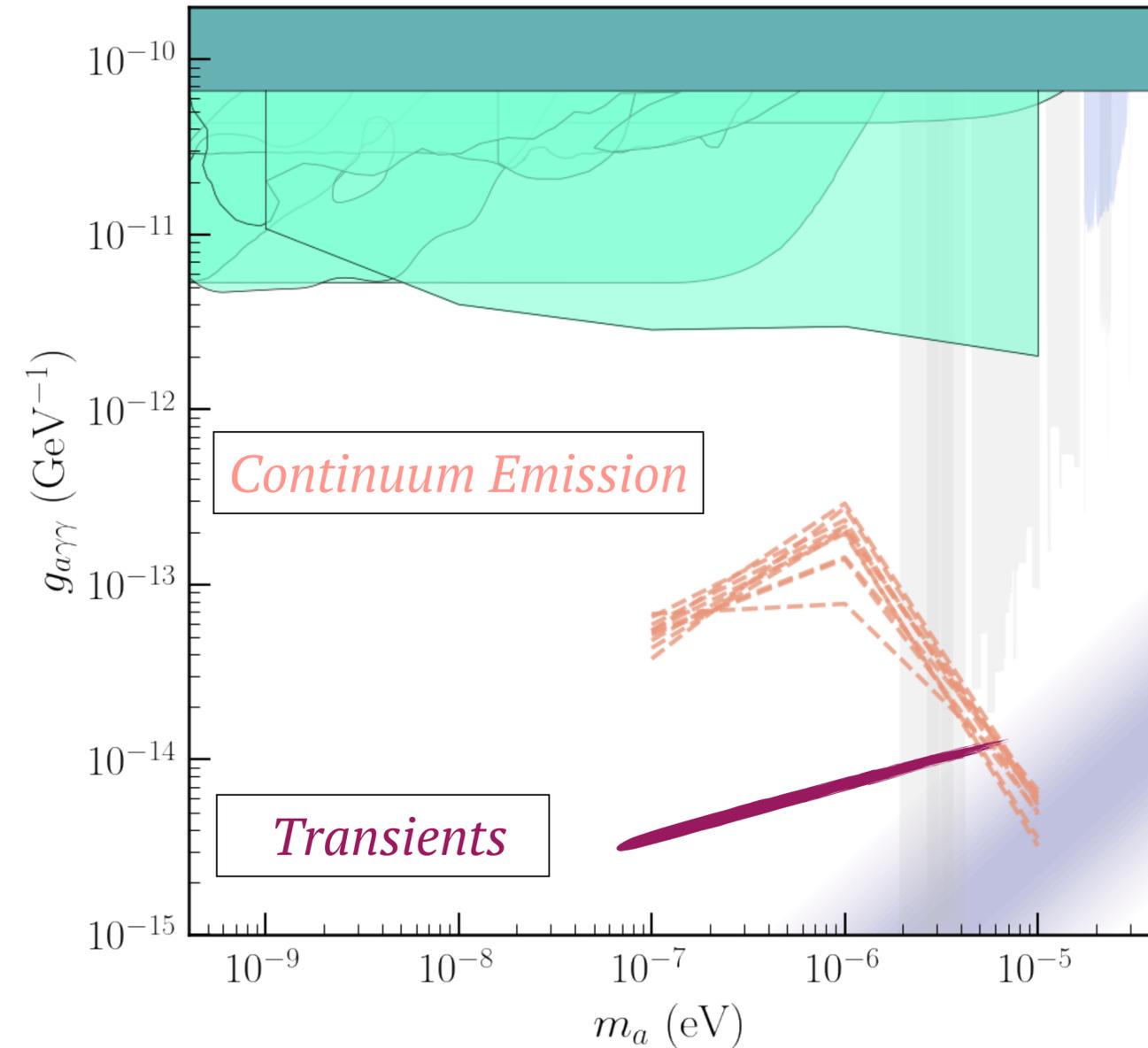
Noordhuis, Prabhu, Weniger, SJW (2023)
 Caputo, SJW, Philippov, Jacobson (2023)

Sensitivity to axion clouds

Axion Back-Reaction

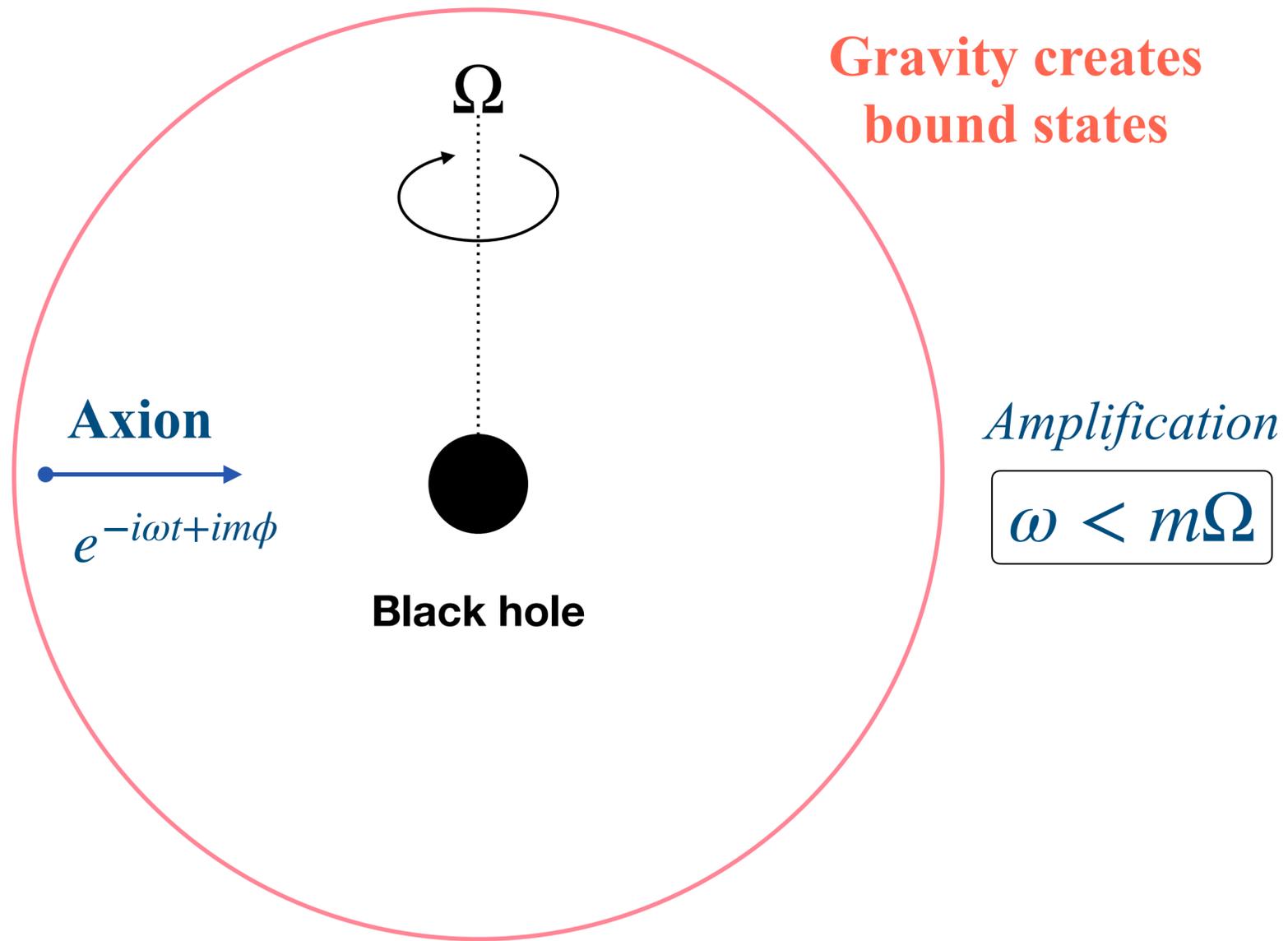


Radio Emission



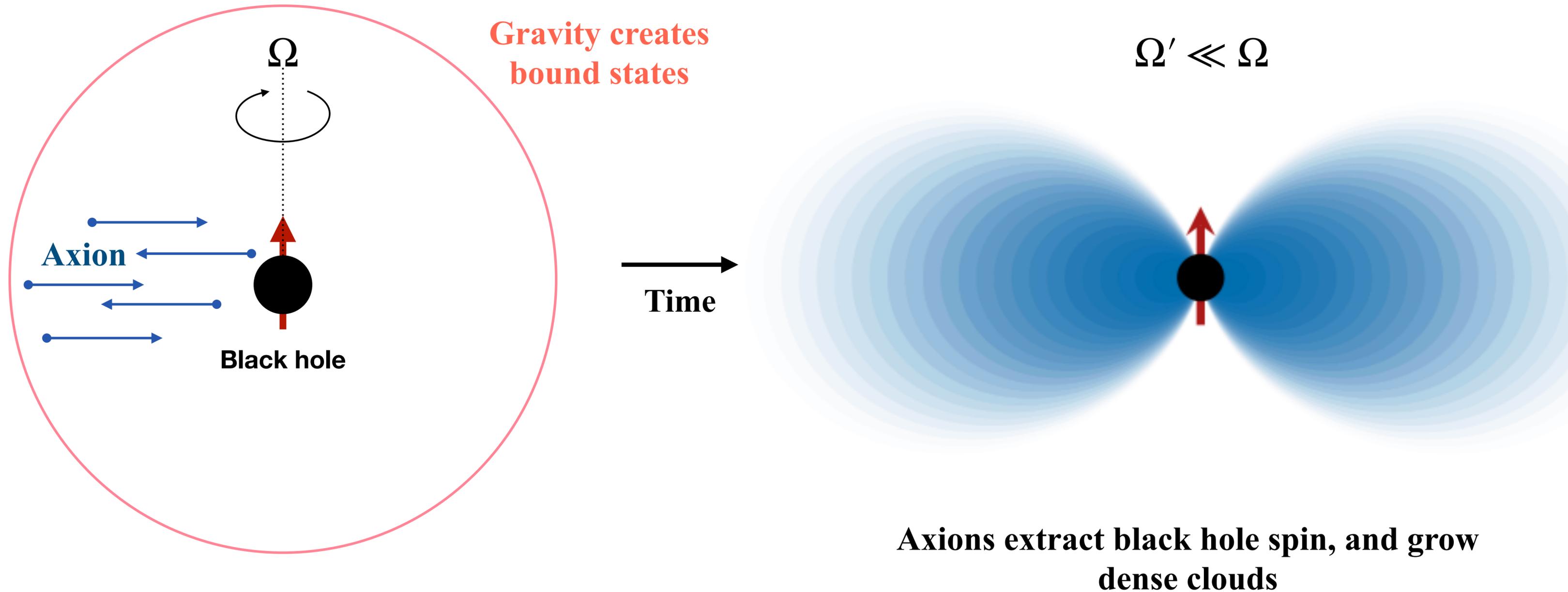
Upside: Very powerful new probes
Difficulty: Work to be done on modelling

Black hole superradiance



Zeldovich (1972) Press & Teukolsky (1972), Arvanitaki, Dimopoulos, Dubovsky, Kaloper, J. March-Russell (2010), Arvanitaki & Dubovsky (2011), Brito, Cardoso, Pani (2015)

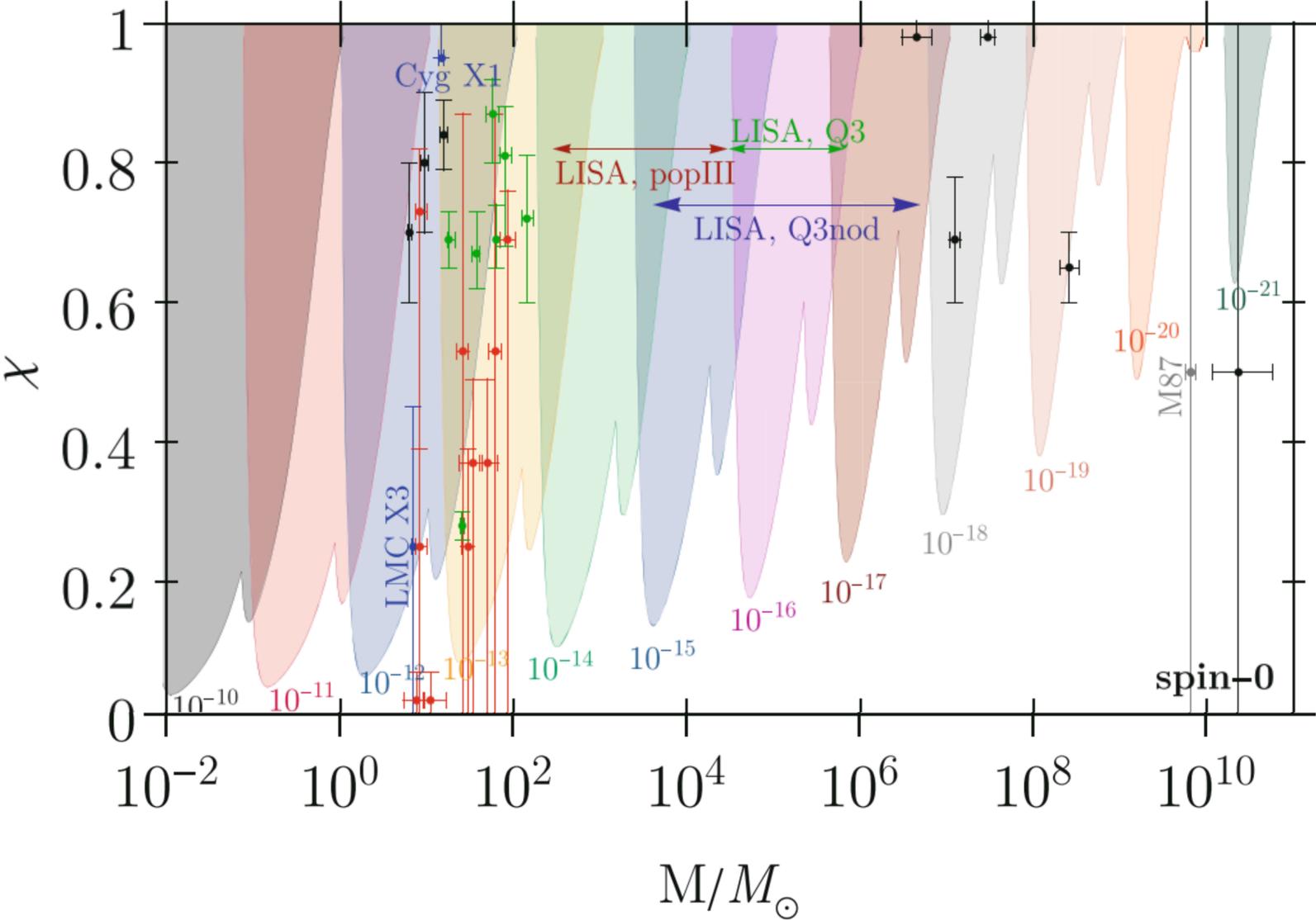
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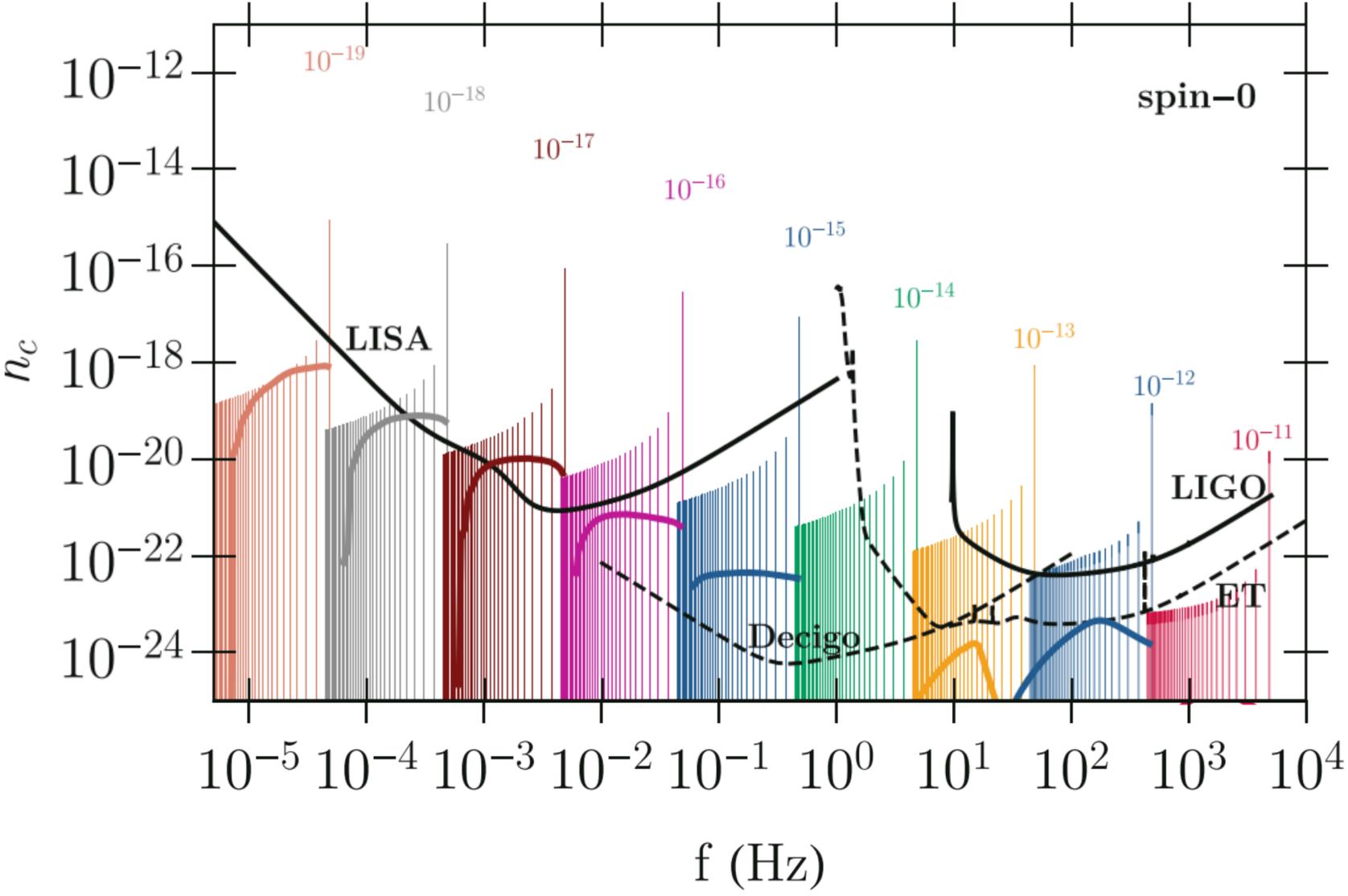
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Axion superradiance

Black hole spin distributions



Gravitational waves from axion cloud



Superradiance in the non-interacting limit

$$(\square + \mu^2) a = 0$$

Bound states form discrete hydrogen-like energy spectrum: $|nlm\rangle$

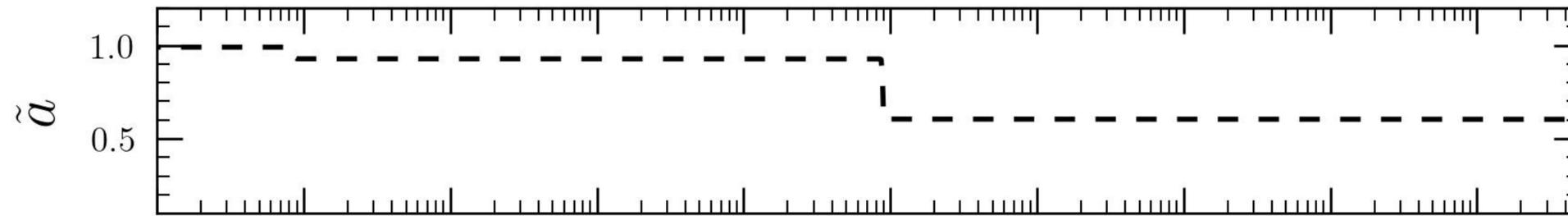
$$\omega_{nlm} = E_{nlm} + i\Gamma_{nlm}$$

Superradiance in the non-interacting limit

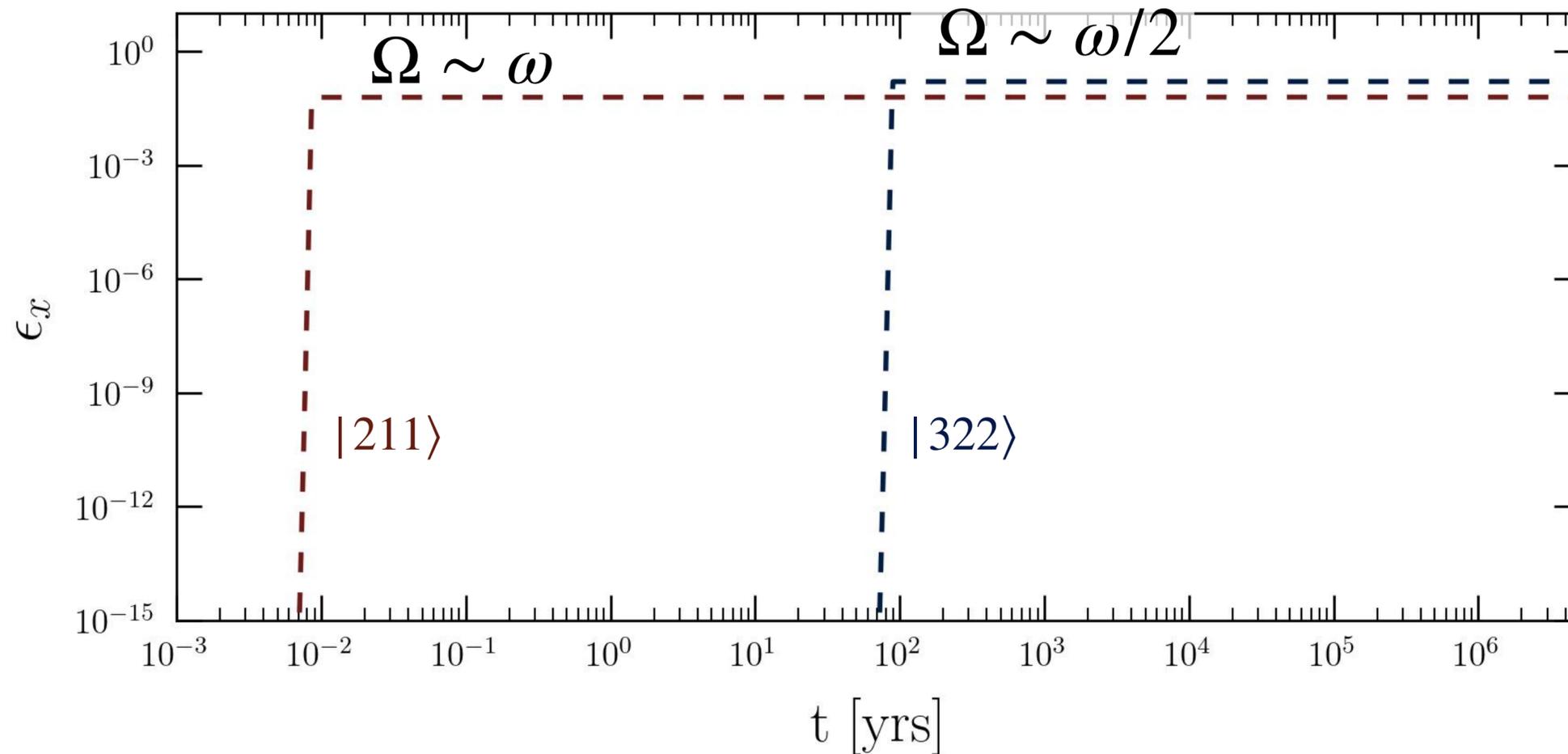
$$(\square + \mu^2) a = 0$$

Bound states form discrete hydrogen-like energy spectrum: $|nlm\rangle$

Black hole spin



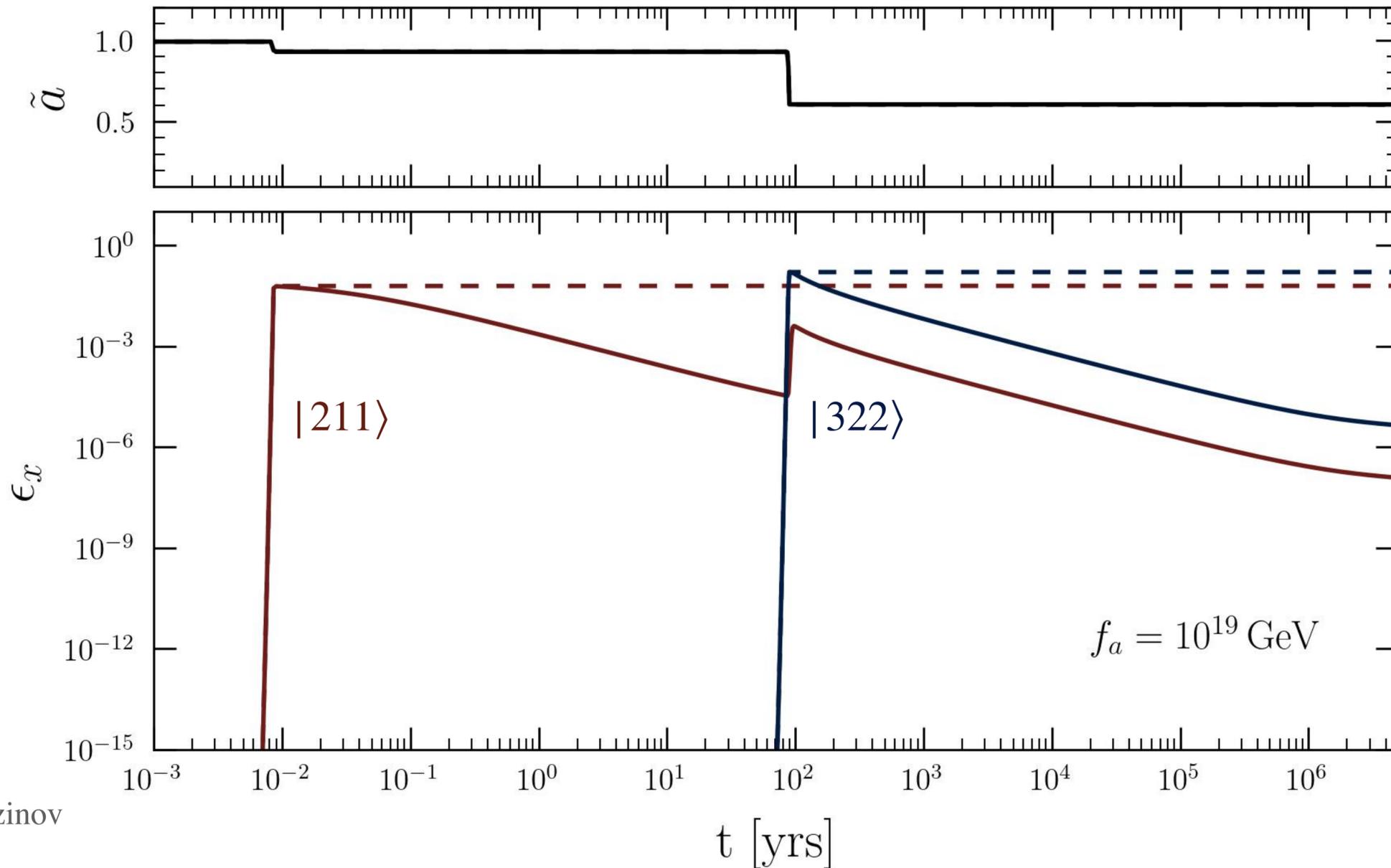
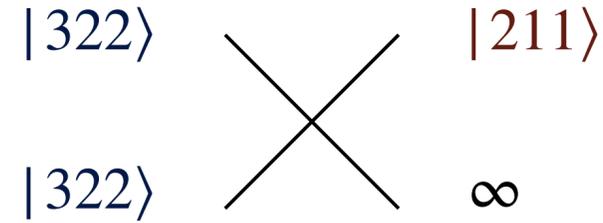
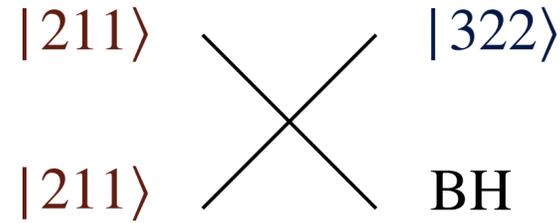
Normalized occupation numbers



Self-interactions in superradiance

$$\mathcal{L} \supset \frac{m_a^2}{f_a^2} a^4$$

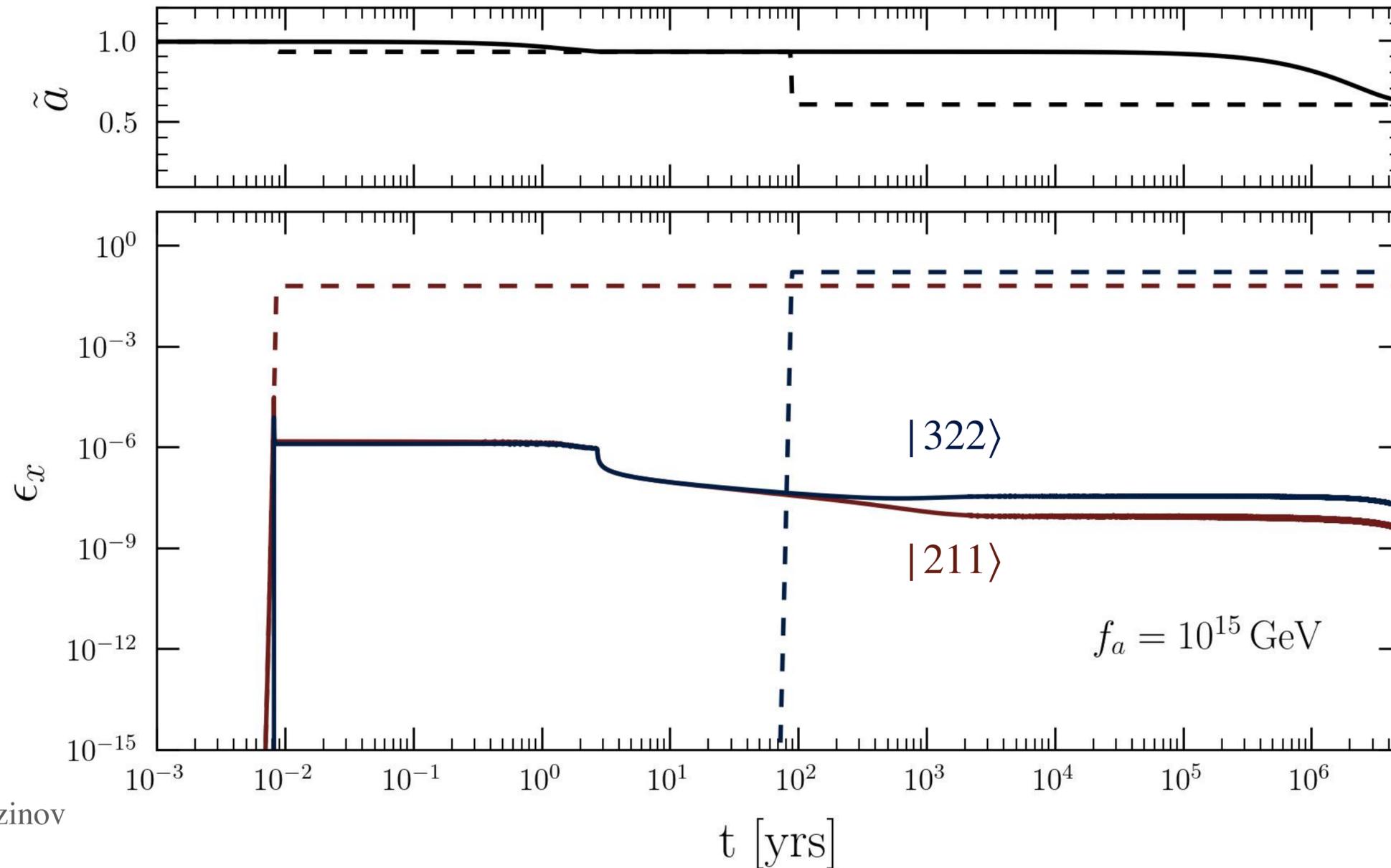
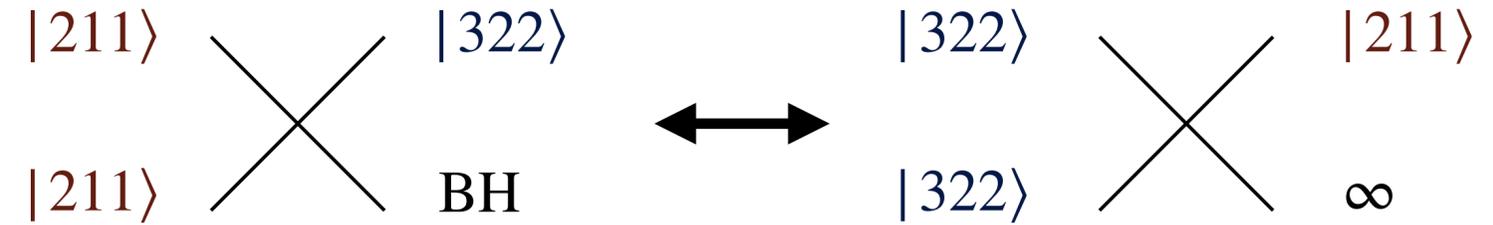
e.g. \rightarrow



Arvanitaki & Dubovsky (2011), Gruzinov (2016), Baryakhtar et al (2021)

Self-interactions in superradiance

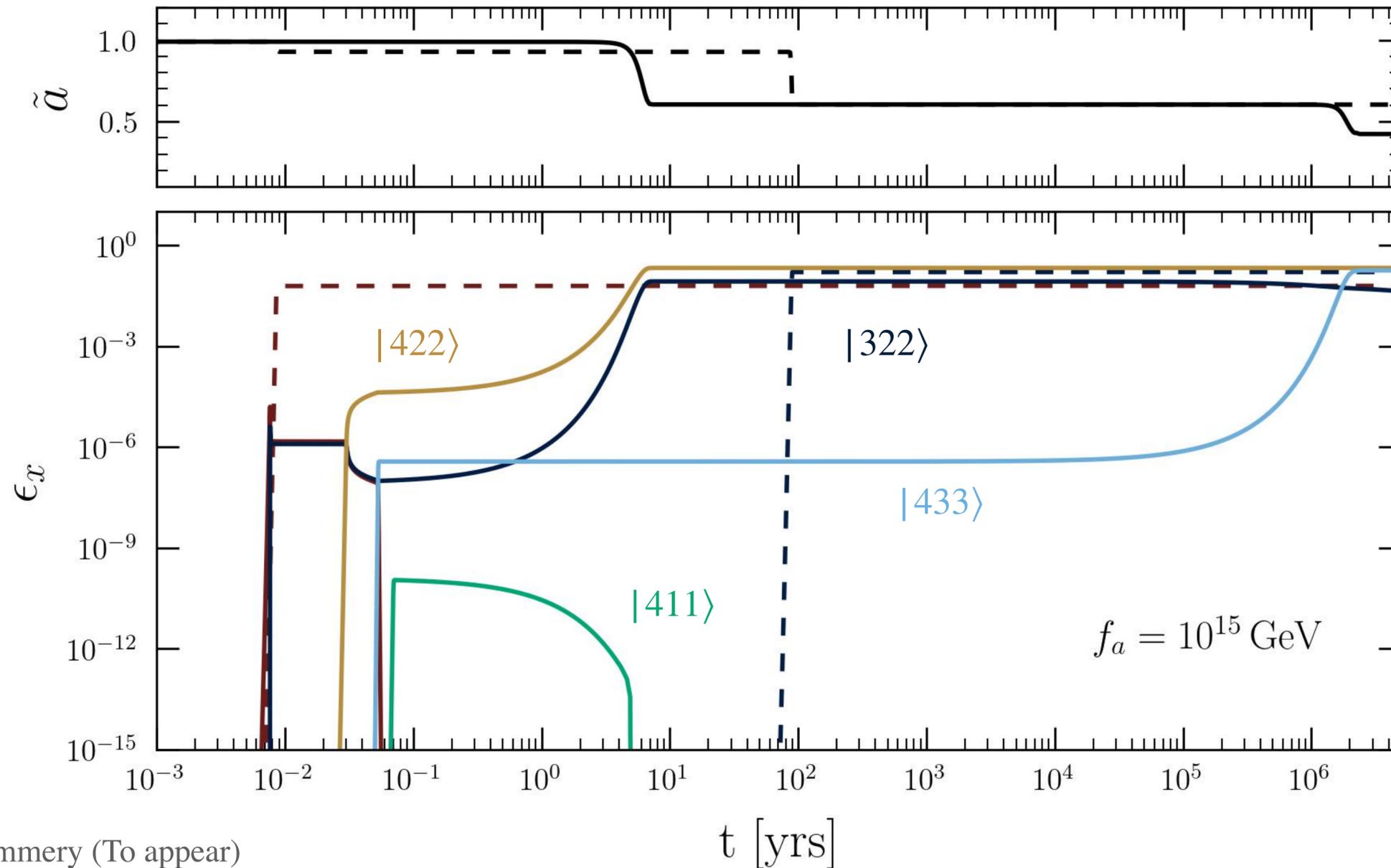
Large self couplings dramatically slow spin extraction!



Arvanitaki & Dubovsky (2011), Gruzinov (2016), Baryakhtar et al (2021)

Self-interactions in superradiance

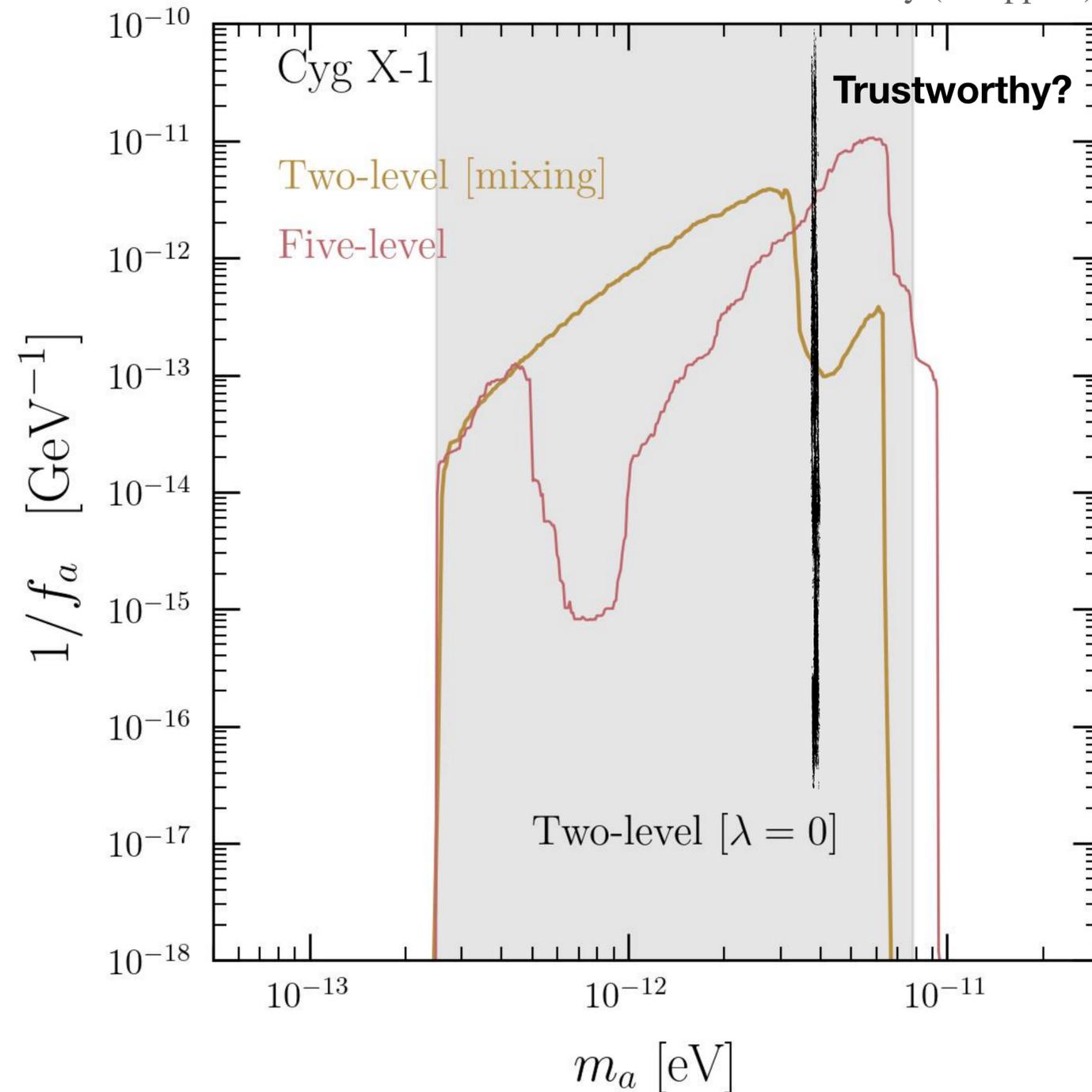
$n = 4$ is a mess....



Baryakhtar et al (2021), SJW & Mummery (To appear)

Superradiance limits

SJW & Mummery (To appear)



Upside: Very powerful, can cover huge range of parameter space

Difficulty: Modelling self-interactions (especially near black hole) is not fully understood

Difficulty from astro side:

- How to do we obtain reliable spin measurements
- Environmental effects (*not isolated objects!*)

See e.g. Baumann et al (2019, 2022)

Conclusions

