SMEFT at Tera-Z

Yong Du (杜勇)

7th FCC Physics Workshop, LAPP, Annecy, Feb 1, 2024

Based on

<u>2206.08326</u>, Jorge de Blas, YD, Christophe Grojean, Jiayin Gu, Victor Miralles, Michael Peskin, Junping Tian, Marcel Vos, Eleni Vryonidou for SNOWMASS 2021



The SM, up to now, is very successful. But there are some flaws:



YD, Huang, Li, Yu, 2005.01717 (JHEP) YD, Huang, Li, Li, Yu, 2111.01267 (JCAP)



Elahi et al, 1410.6157

Chiang, Cottin, YD, Fuyuto, Ramsey-Musolf, 2003.07867(JHEP)

Yong Du

On the other hand, neutrinos oscillate



Yong Du

While there are many models for dark matter, neutrinos and other topics as you prefer, the direct experimental observation of any new particle is still null.

<u>Q: How to approach new physics beyond the Standard Model?</u>

<u>A: …</u>



While there are many models for dark matter, neutrinos and other topics as you prefer, the direct experimental observation of any new particle is still null.

Q: How to approach new physics beyond the Standard Model?

<u>A: …</u>



The experimental data are suggesting that the SM is an effective low-energy theory of some UV model above the weak scale.





Operators in the Warsaw basis:

	X^3		$arphi^6$ and $arphi^4 D^2$		$\psi^2 arphi^3$	
Q_G	$f^{ABC}G^{A\nu}_{\mu}G^{B\rho}_{\nu}G^{C\mu}_{\rho}$	Q_{arphi}	$(arphi^\dagger arphi)^3$	$Q_{e\varphi}$	$(arphi^\dagger arphi) (ar l_p e_r arphi)$	
$Q_{\widetilde{G}}$	$f^{ABC} \widetilde{G}^{A\nu}_{\mu} G^{B\rho}_{\nu} G^{C\mu}_{\rho}$	$Q_{arphi\square}$	$(arphi^\daggerarphi) \Box (arphi^\daggerarphi)$	$Q_{u\varphi}$	$(arphi^\dagger arphi) (ar q_p u_r \widetilde arphi)$	
Q_W	$\varepsilon^{IJK}W^{I\nu}_{\mu}W^{J\rho}_{\nu}W^{K\mu}_{\rho}$	$Q_{arphi D}$	$\left(arphi^{\dagger} D^{\mu} arphi ight)^{\star} \left(arphi^{\dagger} D_{\mu} arphi ight)$	Q_{darphi}	$(arphi^\dagger arphi) (ar q_p d_r arphi)$	
$Q_{\widetilde{W}}$	$\varepsilon^{IJK}\widetilde{W}_{\mu}^{I\nu}W_{\nu}^{J\rho}W_{\rho}^{K\mu}$					
	$X^2 \varphi^2$		$\psi^2 X \varphi$		$\psi^2 arphi^2 D$	
$Q_{arphi G}$	$arphi^\dagger arphi G^A_{\mu u} G^{A\mu u}$	Q_{eW}	$(ar{l}_p\sigma^{\mu u}e_r) au^Iarphi W^I_{\mu u}$	$Q^{(1)}_{arphi l}$	$(\varphi^{\dagger}i\overleftrightarrow{D}_{\mu}\varphi)(\bar{l}_{p}\gamma^{\mu}l_{r})$	
$Q_{arphi \widetilde{G}}$	$arphi^\dagger arphi \widetilde{G}^A_{\mu u} G^{A\mu u}$	Q_{eB}	$(\bar{l}_p \sigma^{\mu u} e_r) \varphi B_{\mu u}$	$Q^{(3)}_{arphi l}$	$(\varphi^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}\varphi)(\bar{l}_{p} au^{I}\gamma^{\mu}l_{r})$	
$Q_{arphi W}$	$arphi^\dagger arphi W^I_{\mu u} W^{I\mu u}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu u} T^A u_r) \widetilde{\varphi} G^A_{\mu u}$	$Q_{arphi e}$	$(\varphi^{\dagger}i\overleftrightarrow{D}_{\mu}\varphi)(\bar{e}_{p}\gamma^{\mu}e_{r})$	
$Q_{arphi \widetilde{W}}$	$arphi^\dagger arphi \widetilde{W}^I_{\mu u} W^{I\mu u}$	Q_{uW}	$(ar{q}_p \sigma^{\mu u} u_r) au^I \widetilde{arphi} W^I_{\mu u}$	$Q^{(1)}_{arphi q}$	$(arphi^\dagger i \overleftrightarrow{D}_\mu arphi) (ar{q}_p \gamma^\mu q_r)$	
$Q_{arphi B}$	$arphi^\dagger arphi B_{\mu u} B^{\mu u}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu u} u_r) \widetilde{\varphi} B_{\mu u}$	$Q^{(3)}_{arphi q}$	$(\varphi^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}\varphi)(\bar{q}_{p} au^{I}\gamma^{\mu}q_{r})$	
$Q_{arphi \widetilde{B}}$	$arphi^\dagger arphi \widetilde{B}_{\mu u} B^{\mu u}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu u} T^A d_r) \varphi G^A_{\mu u}$	$Q_{\varphi u}$	$(\varphi^{\dagger}i\overleftrightarrow{D}_{\mu}\varphi)(ar{u}_{p}\gamma^{\mu}u_{r})$	
$Q_{arphi WB}$	$\varphi^\dagger \tau^I \varphi W^I_{\mu\nu} B^{\mu\nu}$	Q_{dW}	$(\bar{q}_p \sigma^{\mu u} d_r) \tau^I \varphi W^I_{\mu u}$	$Q_{arphi d}$	$(\varphi^{\dagger}i\overleftrightarrow{D}_{\mu}\varphi)(\bar{d}_{p}\gamma^{\mu}d_{r})$	
$Q_{arphi \widetilde{W}B}$	$arphi^\dagger au^I arphi \widetilde{W}^I_{\mu u} B^{\mu u}$	Q_{dB}	$(\bar{q}_p \sigma^{\mu u} d_r) \varphi B_{\mu u}$	$Q_{arphi u d}$	$i(\widetilde{arphi}^{\dagger}D_{\mu}arphi)(ar{u}_{p}\gamma^{\mu}d_{r})$	

Buchmuller and Wyler, Nucl.Phys.B 268 (1986) 621 Grzadkowski, Iskrzynski, Misiak and Rosiek, JHEP 10 (2010) 085

$(\bar{L}L)(\bar{L}L)$		$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$		
Q_{ll}	$(ar{l}_p \gamma_\mu l_r) (ar{l}_s \gamma^\mu l_t)$	Q_{ee}	$(ar{e}_p\gamma_\mu e_r)(ar{e}_s\gamma^\mu e_t)$	Q_{le}	$(ar{l}_p\gamma_\mu l_r)(ar{e}_s\gamma^\mu e_t)$	
$Q_{qq}^{(1)}$	$(ar q_p \gamma_\mu q_r) (ar q_s \gamma^\mu q_t)$	Q_{uu}	$(ar{u}_p\gamma_\mu u_r)(ar{u}_s\gamma^\mu u_t)$	Q_{lu}	$(ar{l}_p \gamma_\mu l_r) (ar{u}_s \gamma^\mu u_t)$	
$Q_{qq}^{\left(3 ight) }$	$(ar q_p \gamma_\mu au^I q_r) (ar q_s \gamma^\mu au^I q_t)$	Q_{dd}	$(ar{d}_p\gamma_\mu d_r)(ar{d}_s\gamma^\mu d_t)$	Q_{ld}	$(ar{l}_p\gamma_\mu l_r)(ar{d}_s\gamma^\mu d_t)$	
$Q_{lq}^{(1)}$	$(ar{l}_p\gamma_\mu l_r)(ar{q}_s\gamma^\mu q_t)$	Q_{eu}	$(ar{e}_p \gamma_\mu e_r) (ar{u}_s \gamma^\mu u_t)$	Q_{qe}	$(ar q_p \gamma_\mu q_r) (ar e_s \gamma^\mu e_t)$	
$Q_{lq}^{\left(3 ight) }$	$(ar{l}_p\gamma_\mu au^I l_r)(ar{q}_s\gamma^\mu au^I q_t)$	Q_{ed}	$(ar{e}_p\gamma_\mu e_r)(ar{d}_s\gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(ar q_p \gamma_\mu q_r) (ar u_s \gamma^\mu u_t)$	
		$Q_{ud}^{\left(1 ight) }$	$(ar{u}_p\gamma_\mu u_r)(ar{d}_s\gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(ar{q}_p \gamma_\mu T^A q_r) (ar{u}_s \gamma^\mu T^A u_t)$	
		$Q_{ud}^{(8)}$	$(ar{u}_p \gamma_\mu T^A u_r) (ar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(ar q_p \gamma_\mu q_r) (ar d_s \gamma^\mu d_t)$	
				$Q_{qd}^{(8)}$	$(ar{q}_p \gamma_\mu T^A q_r) (ar{d}_s \gamma^\mu T^A d_t)$	
$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$		<i>B</i> -violating				
$Q_{ledq} = (ar{l}_p^j e_r)(ar{d}_s q_t^j) = Q_{duq}$			$arepsilon^{lphaeta\gamma}arepsilon_{jk}\left[(d_p^lpha)^TCu_r^eta ight]\left[(q_s^{\gamma j})^TCl_t^k ight]$			
$Q_{quqd}^{(1)}$	$(ar{q}_p^j u_r) arepsilon_{jk} (ar{q}_s^k d_t)$	Q_{qqu}	$arepsilon^{lphaeta\gamma}arepsilon_{jk}\left[(q_p^{lpha j})^TCq_r^{eta k} ight]\left[(u_s^{\gamma})^TCe_t ight]$			
$Q_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \varepsilon_{jk} (\bar{q}_s^k T^A d_t)$	Q_{qqq}	$arepsilon^{lphaeta\gamma}arepsilon_{jn}arepsilon_{km}\left[(q_p^{lpha j})^TCq_r^{eta k} ight]\left[(q_s^{\gamma m})^TCl_t^n ight]$			
$Q_{lequ}^{(1)}$	$(ar{l}_p^j e_r) arepsilon_{jk} (ar{q}_s^k u_t)$	Q_{duu}	$arepsilon^{lphaeta\gamma}\left[(d_p^lpha)^T C u_r^eta ight]\left[(u_s^\gamma)^T C e_t ight]$			
$Q_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu u} e_r) arepsilon_{jk} (\bar{q}_s^k \sigma^{\mu u} u_t)$					

59 operators (+ 4 B-violating ones)

2499 operators: 1350 (CP-even) + 1149 (CP-odd)

No flavor assumptions are made.

Sophie's and Ben's talks on Tuesday

SMEFT global fit

The SMEFT is then simply constructed by adding these operators on top of the SM:

$$\mathscr{L}_{\text{SMEFT}} = \mathscr{L}_{\text{SM}} + \sum_{i} \frac{c_i}{\Lambda^{d-4}} \mathcal{O}_i^d$$

with each of the term a SM singlet and respecting the SM local gauge symmetry.

SMEFT global fit

The SMEFT is then simply constructed by adding these operators on top of the SM:

$$\mathscr{L}_{\text{SMEFT}} = \mathscr{L}_{\text{SM}} + \sum_{i} \frac{c_i}{\Lambda^{d-4}} \mathcal{O}_i^d$$

with each of the term a SM singlet and respecting the SM local gauge symmetry.

FCC-ee is an ideal precision machine for new physics studies in the SMEFT since $s \ll \Lambda^2$ with $\Lambda \gtrsim 1 \text{ TeV}$ from the LHC data.



CEPC Physics Study Group, 2205.08553

SMEFT global fit: $\underline{\delta}_{ex}$

de Blas, YD, Grojean, Gu, Miralles, Peskin, Tian, Vos, Vryonidou, 2206.08326

Quantity	current	ILC250	ILC-GigaZ	FCC-ee	CEPC	CLIC380
$\Delta lpha(m_Z)^{-1}~(imes 10^3)$	17.8*	17.8*		3.8 (1.2)	17.8*	
$\Delta m_W ~({ m MeV})$	12*	0.5(2.4)		0.25~(0.3)	0.35~(0.3)	
$\Delta m_Z \ ({ m MeV})$	2.1*	0.7(0.2)	0.2	0.004~(0.1)	0.005~(0.1)	2.1*
$\Delta m_H ~({ m MeV})$	170*	14		2.5(2)	5.9	78
$\Delta\Gamma_W$ (MeV)	42*	2		1.2(0.3)	1.8(0.9)	
$\Delta\Gamma_Z$ (MeV)	2.3*	1.5(0.2)	0.12	$0.004\ (0.025)$	$0.005 \ (0.025)$	2.3^{*}
$\Delta A_e \ (\times 10^5)$	190*	14 (4.5)	1.5 (8)	0.7 (2)	1.5	64
$\Delta A_{\mu} \; (imes 10^5)$	1500*	82 (4.5)	3(8)	2.3(2.2)	3.0(1.8)	400
$\Delta A_{ au} \ (imes 10^5)$	400*	86 (4.5)	3(8)	0.5~(20)	1.2~(6.9)	570
$\Delta A_b \ (imes 10^5)$	2000*	53 (35)	9(50)	2.4(21)	3(21)	380
$\Delta A_c \; (imes 10^5)$	2700*	140(25)	20(37)	20(15)	6 (30)	200
$\Delta \sigma_{\rm had}^0 ~{ m (pb)}$	37*			0.035 (4)	0.05 (2)	37*
$\delta R_e \; (imes 10^3)$	2.4*	0.5(1.0)	0.2 (0.5)	0.004~(0.3)	0.003~(0.2)	2.7
$\delta R_{\mu}~(imes 10^3)$	1.6*	0.5(1.0)	0.2 (0.2)	$0.003\ (0.05)$	0.003~(0.1)	2.7
$\delta R_{ au}~(imes 10^3)$	2.2*	0.6(1.0)	0.2 (0.4)	0.003~(0.1)	0.003~(0.1)	6
$\delta R_b \ (imes 10^3)$	3.0*	0.4 (1.0)	0.04(0.7)	$0.0014 \ (< 0.3)$	0.005~(0.2)	1.8
$\delta R_c(imes 10^3)$	17*	0.6(5.0)	0.2(3.0)	0.015~(1.5)	0.02(1)	5.6

We thank our experimental colleagues for doing excellent. Recent improvement not implemented (R_b from inclusive (<u>Michele's talk</u>)/exclusive (<u>Lars's talk</u>) studies for example)

Yong Du

SMEFT global fit: $\underline{\delta_{th}}$

Theory Requirements and Possibilities for the FCC-ee and other Future High Energy and Precision Frontier Lepton Colliders*

Alain Blondel (Université de Genéve), Ayres Freitas (University of Pittsburgh), Janusz Gluza[†] and Tord Riemann (U. Silesia), Sven Heinemeyer (IFT/IFCA CSIC Madrid/Santander, ECI/UAM/CSIC Madrid), Stanisław Jadach (IFJ PAN Kraków), Patrick Janot (CERN)

18 December 2018

Abstract

The future lepton colliders proposed for the High Energy and Precision Frontier set stringent demands on theory. The most ambitious, broad-reaching and demanding project is the FCC-ee. We consider here the present status and requirements on precision calculations, possible ways forward and novel methods, to match the experimental accuracies expected at the FCC-ee. We conclude that the challenge can be tackled by a distributed collaborative effort in academic institutions around the world, provided sufficient support, which is estimated to about 500 man-years over the next 20 years.

Considered as well under control by the operation time. See also <u>Johann's</u> talk on Tuesday and <u>Alain's</u> talk this morning.

Yong Du

SMEFT global fit: <u>**Basis</u></u></u>**

Presenting the results will be basis dependent. We choose to work in the Higgs basis to disentangle physics in different sectors

$$\begin{split} \mathcal{L} \supset eA^{\mu} \sum_{f=u,d,e} Q_{f}(\overline{f}_{I}\overline{\sigma}_{\mu}f_{I} + f_{I}^{c}\sigma_{\mu}\overline{f}_{I}^{c}) \\ &+ \frac{g_{L}}{\sqrt{2}} \left[W^{\mu+}\overline{\nu}_{I}\overline{\sigma}_{\mu}(\delta_{IJ} + [\delta g_{L}^{W\ell}]_{IJ})e_{J} + W^{\mu+}\overline{u}_{I}\overline{\sigma}_{\mu} \left(V_{IJ} + \left[\delta g_{L}^{Wq} \right]_{IJ} \right) d_{J} + \text{h.c.} \right] \\ &+ \frac{g_{L}}{\sqrt{2}} \left[W^{\mu+}u_{I}^{c}\sigma_{\mu} \left[\delta g_{R}^{Wq} \right]_{IJ} \overline{d}_{J}^{c} + \text{h.c.} \right] \\ &+ \sqrt{g_{L}^{2} + g_{Y}^{2}} Z^{\mu} \sum_{f=u,d,e,\nu} \overline{f}_{I}\overline{\sigma}_{\mu} \left((T_{3}^{f} - s_{w}^{2}Q_{f})\delta_{IJ} + \left[\delta g_{L}^{Zf} \right]_{IJ} \right) f_{J} \\ &+ \sqrt{g_{L}^{2} + g_{Y}^{2}} Z^{\mu} \sum_{f=u,d,e} f_{I}^{c}\sigma_{\mu} \left(-s_{w}^{2}Q_{f}\delta_{IJ} + \left[\delta g_{R}^{Zf} \right]_{IJ} \right) \overline{f}_{J}^{c}, \end{split}$$

SMEFT global fit: <u>**Basis</u></u></u>**

Presenting the results will be basis dependent. We choose to work in the Higgs basis to disentangle physics in different sectors

$$\delta gLWe \rightarrow cHl3\#Warsaw v^{2} - \frac{cHD\#Warsaw gL^{2}v^{2}}{4(gL^{2}-gv^{2})} - \frac{cHD\#Warsaw gLgYv^{2}}{gL^{2}-gv^{2}} - \frac{gL^{2}v^{2} \triangle GF}{2(gL^{2}-gv^{2})}$$

$$\delta gLZe \rightarrow - \frac{cHl1\#Warsaw v^{2}}{2} - \frac{cHl3\#Warsaw v^{2}}{2} + \frac{cHD\#Warsaw gLgYv^{2}}{gL^{2}-gv^{2}} + \frac{cHD\#Warsaw (gL^{2}+gV^{2})v^{2}}{8(gL^{2}-gV^{2})} + \frac{(gL^{2}+gV^{2})v^{2} \triangle GF}{4(gL^{2}-gV^{2})}$$

$$\delta gRZe \rightarrow - \frac{cHe1!\#Warsaw v^{2}}{2} + \frac{cHD\#Warsaw gV^{2}v^{2}}{4gL^{2}-4gV^{2}} + \frac{cHD\#Warsaw gLgYv^{2}}{4gL^{2}-gV^{2}} + \frac{cHD\#Warsaw (gL^{2}+gV^{2})v^{2}}{2(gL^{2}-gV^{2})} + \frac{(gL^{2}+gV^{2})v^{2} \triangle GF}{2(gL^{2}-gV^{2})}$$

$$\delta gLZu \rightarrow - \frac{cHq1!\#Warsaw v^{2}}{2} + \frac{cHq3!\#Warsaw v^{2}}{2} - \frac{2 cHWB!\#Warsaw gLgYv^{2}}{3(gL^{2}-gY^{2})} - \frac{cHD!Warsaw (3gL^{2}+gV^{2})v^{2}}{24(gL^{2}-gV^{2})} - \frac{(3gL^{2}+gV^{2})v^{2} \triangle GF}{12(gL^{2}-gV^{2})}$$

$$\delta gLZd \rightarrow - \frac{cHq1!\#Warsaw v^{2}}{2} - \frac{cHq3!\#Warsaw v^{2}}{2} + \frac{cHWB!}{3(gL^{2}-gY^{2})} + \frac{cHD!Warsaw (3gL^{2}-gV^{2})v^{2}}{24(gL^{2}-gV^{2})} + \frac{(3gL^{2}-gV^{2})v^{2} \triangle GF}{12(gL^{2}-gV^{2})}$$

$$\delta gRZd \rightarrow - \frac{cHq!}{2} - \frac{cHd!}{2} - \frac{2}{2} - \frac{2}{2} - \frac{2}{3(gL^{2}-gV^{2})} + \frac{cHD!}{3(gL^{2}-gV^{2})} + \frac{cHD!}{3(gL^{2}-gV^{2})} + \frac{gV^{2}v^{2} \triangle GF}{24(gL^{2}-gV^{2})}$$

$$\delta gRZd \rightarrow - \frac{cHd!}{2} - \frac{cHd!}{2} - \frac{2}{2} - \frac{2}{2} - \frac{2}{2} - \frac{2}{3(gL^{2}-gV^{2})} + \frac{cHD!}{3(gL^{2}-gV^{2})} + \frac{cHD!}{3(gL^{2}-gV^{2})} + \frac{gV^{2}v^{2} \triangle GF}{24(gL^{2}-gV^{2})}$$

de Blas, YD, Grojean, Gu, Miralles, Peskin, Tian, Vos, Vryonidou, 2206.08326

Process	Observable	Experimental value	Ref.	SM prediction	
(-)	$g_{LV}^{ u_{\mu}e}$	-0.035 ± 0.017		-0.0396 [48]	
$\nu_{\mu} - e^{-}$ scattering	$g_{LA}^{ u_{\mu}e}$	-0.503 ± 0.017	CHARM-II [47]	-0.5064 [48]	
σdoon	$\frac{G_{\tau e}^2}{G_{\nu}^2}$	1.0029 ± 0.0046	PDC2014 [49]	1	
7 decay	$\frac{G_{\tau\mu}^{f}}{G_{F}^{2}}$	0.981 ± 0.018	1 DG2014 [49]	1	
	$R_{ u\mu}$	0.3093 ± 0.0031	CHARM $(r - 0.456)$ [50]	0.3156 [50]	
	$R_{\overline{ u}_{\mu}}$	0.390 ± 0.014		0.370 [<mark>50</mark>]	
Neutrino scattoring	$R_{ u_{\mu}}$	0.3072 ± 0.0033	CDHS $(r = 0.303)$ [51]	0.3091 [51]	
Neutrino scattering	$R_{\overline{ u}_{\mu}}$	0.382 ± 0.016	CDHS (7 = 0.333) [51]	0.380 [51]	
	κ	0.5820 ± 0.0041	CCFR [52]	0.5830 [52]	
	$R_{ u_e\overline{ u}_e}$	$0.406^{+0.145}_{-0.135}$	CHARM [53]	0.33 [54]	
	$(s_w^2)^{ m M {\it arsigma} m ller}$	0.2397 ± 0.0013	SLAC-E158 [55]	0.2381 ± 0.0006 [56]	
	$Q_W^{ m Cs}(55,78)$	-72.62 ± 0.43	PDG2016 [54]	-73.25 ± 0.02 [54]	
	$Q_W^{\mathrm{p}}(1,0)$	0.064 ± 0.012	QWEAK [57]	0.0708 ± 0.0003 [54]	
	A_1	$(-91.1 \pm 4.3) \times 10^{-6}$		$(-87.7 \pm 0.7) \times 10^{-6}$ [58]	
Parity-violating scattering	A_2	$(-160.8\pm7.1) imes10^{-6}$		$(-158.9 \pm 1.0) \times 10^{-6} [58]$	
	$g^{eu}_{VA} - g^{ed}_{VA}$	-0.042 ± 0.057	SAMPLE ($\sqrt{Q^2} = 200 \text{MeV}$) [59]	-0.0360 [54]	
		-0.12 ± 0.074	SAMPLE ($\sqrt{Q^2} = 125 \text{MeV}$) [59]	0.0265 [54]	
	$b_{ m SPS}$	$-(1.47 \pm 0.42) \times 10^{-4} \mathrm{GeV^{-2}}$	SPS $(\lambda = 0.81)$ [60]	$-1.56 \times 10^{-4} \mathrm{GeV^{-2}}$ [60]	
		$-(1.74 \pm 0.81) \times 10^{-4} \mathrm{GeV^{-2}}$	SPS $(\lambda = 0.66)$ [60]	$-1.57 \times 10^{-4} \mathrm{GeV^{-2}}$ [60]	
- polarization	$\mathcal{P}_{ au}$	0.012 ± 0.058	VENIIS [61]	0.028 [61]	
	$\mathcal{A}_{\mathcal{P}}$	0.029 ± 0.057		0.021 [61]	
Neutrino trident production	$\frac{\sigma}{\sigma^{\rm SM}}(\nu_{\mu}\gamma^{*} \rightarrow \nu_{\mu}\mu^{+}\mu^{-})$	0.82 ± 0.28	CCFR [62–64]	1	
$d_I ightarrow u_J \ell \overline{ u}_\ell(\gamma)$	$\epsilon^{de_J}_{L,R,S,P,T}$	See text	[65]	0	
	δA^e_{LR}	2.0%		0.00015	
	δA^{μ}_{LR}	1.5%		-0.0006	
$e^+e^- \to f\overline{f}$	$\delta A^{ au}_{LR}$	2.4%	SuperKEKB [66]	-0.0006	
	δA^c_{LR}	0.5%		-0.005	
	δA^b_{LR}	0.4%		-0.020	

Flat direction lifted by low-energy experiments: muon sector example



Flat direction lifted by low-energy experiments: muon sector example



Flat direction lifted by low-energy experiments: electron sector example

Bhabha alone is not enough to close the fit, $A_{\rm PV}$ from PVES is the key



Dev, Ramsey-Musolf, Zhang, 1806.08499 (PRD)

YD, Freitas, Patel, Ramsey-Musolf, 1912.08220 (PRL)

SMEFT global fit: 4f

Flat direction lifted by low-energy experiments: electron sector example

Bhabha alone is not enough to close the fit, $A_{\rm PV}$ from PVES is the key



MOLLER project funded

Publication date: Tue, Nov 22, 2022 - 11:30pm



The MOLLER project which has been in planning and development stages for some years, has now been allocated \$31M in Department of Energy funding to construct and install the experiment by 2025 and start data collection in early 2026.. The leader of the UMass team, Prof. Krishna Kumar, is the principal spokesperson for the project. The experiment, to be located at Jefferson National Lab, will study parity violation in electron-

electron scattering.

P2 collaboration, 1802.04759 (EPJA)

Dev, Ramsey-Musolf, Zhang, 1806.08499 (PRD)

YD, Freitas, Patel, Ramsey-Musolf, 1912.08220 (PRL)

Flat direction lifted by low-energy experiments: electron sector example

Bhabha alone is not enough to close the fit, $A_{\rm PV}$ from PVES is the key



YD, Freitas, Patel, Ramsey-Musolf, 1912.08220 (PRL)

Flat direction lifted by low-energy experiments: tau sector example

 τ polarization measurement at VENUS is limited by statistics ($\mathscr{L} = 271 \text{ pb}^{-1}$). FCC-ee at 240GeV will have better sensitivity with much more statistics, how well can we control the systematical?



Flat direction lifted by low-energy experiments: tau sector example

 τ polarization measurement at VENUS is limited by statistics ($\mathscr{L} = 271 \text{ pb}^{-1}$). FCC-ee at 240GeV will have better sensitivity with much more statistics, how well can we control the systematical?



Flat direction lifted by low-energy experiments: tau sector example

 τ polarization measurement at VENUS is limited by statistics ($\mathscr{L} = 271 \text{ pb}^{-1}$). FCC-ee at 240GeV will have better sensitivity with much more statistics, how well can we control the systematical?



<u>Q: Projections for FCC-ee?</u>

Flat direction lifted by low-energy experiments: tau sector example

 τ polarization measurement at VENUS is limited by statistics ($\mathscr{L} = 271 \text{ pb}^{-1}$). FCC-ee at 240GeV will have better sensitivity with much more statistics, how well can we control the systematical?



Q: Projections for FCC-ee?

Also very interesting τ physics at TeraZ see for example Pich 2012.07099 (EPJP)

SMEFT global fit: 4f

de Blas, YD, Grojean, Gu, Miralles, Peskin, Tian, Vos, Vryonidou, 2206.08326



Unprecedented precision reached at the FCCee machine (same for Higgs couplings, see backup)

Still much room left for new physics generating large 1st and 2nd (could be improved with σ_s and A_{FB}^s) generation $Zq\bar{q}$ couplings (possible tagging improvement? The help with AI (Chai, Gu, Li, 2401.02427 for example)?)

Yong Du

SMEFT global fit: 4f

de Blas, YD, Grojean, Gu, Miralles, Peskin, Tian, Vos, Vryonidou, 2206.08326



Also much room left for 4-fermion operators at FCC-ee or CEPC (same for top, see backup), mainly due to the difficulty in telling the right from the left.

Merging current/future neutrino (e.g. CEvNS) and/or polarized experimental (EIC/EicC for example) data to make a difference?

Polarized future circular colliders? (See <u>Zhe Duan's</u> talk on Tuesday for a possible study)

This precision FCC-ee machine could also help in understanding our existence by closely investigating the CPV operators

	Dim-4						
$\mathcal{L}_{ ext{CPV}}^{ ext{QCD}}$	$-m^{*}ar{ heta}ar{q}i\gamma_{5}q$	$\mathcal{L}_{ ext{mass}}^{ ext{quark}}$	$-ar{m}ar{q}q+\epsilonar{m}ar{q} au_e q$				
	Dim-6						
	Pure gauge	Gauge-Higgs					
$Q_{\widetilde{G}}$	$f^{ABC} \widetilde{G}^{A u}_{\mu} G^{B ho}_{ u} G^{C\mu}_{ ho}$	$Q_{arphi \widetilde{G}}$	$arphi^\dagger arphi \widetilde{G}^A_{\mu u} G^{A\mu u}$				
$Q_{\widetilde{W}}$	$\epsilon^{IJK}\widetilde{W}^{I\nu}_{\mu}W^{J\rho}_{\nu}W^{K\mu}_{\rho}$	$Q_{arphi \widetilde{W}}$	$arphi^\dagger arphi \widetilde{W}^I_{\mu u} W^{I\mu u}$				
(Gauge-)Higgs-f	$Q_{arphi \widetilde{B}}$	$arphi^\dagger arphi \widetilde{B}_{\mu u} B^{\mu u}$				
Q_{uG}	$(ar{Q}\sigma^{\mu u}T^A u_R)\widetilde{arphi}G^A_{\mu u}$	$Q_{arphi \widetilde{W}B}$	$arphi^\dagger au^I arphi \widetilde{W}_{I\mu u} B^{\mu u}$				
Q_{dG}	$(ar{Q}\sigma^{\mu u}T^Ad_R)arphi G^A_{\mu u}$		4- <i>f</i>				
Q_{fW}	$(ar{F}\sigma^{\mu u}f_R) au^Iarphi W_{I\mu u}$	Q_{ledq}	$(ar{L}^j e_R) (ar{d}_R Q^j)$				
Q_{fB}	$(ar{F}\sigma^{\mu u}f_R)arphi B_{\mu u}$	$Q_{quqd}^{\left(1 ight)}$	$(ar{Q}^j u_R) \epsilon_{jk} (ar{Q}^k d_R)$				
Q_{fW}	$(ar{F}\sigma^{\mu u}f_R) au^I\widetilde{arphi}W_{I\mu u}$	$Q_{quqd}^{(8)}$	$(ar{Q}^jT^Au_R)\epsilon_{jk}(ar{Q}^kT^Ad_R)$				
Q_{fB}	$(ar{F}\sigma^{\mu u}f_R)\widetilde{arphi}B_{\mu u}$	$Q_{lequ}^{\left(1 ight)}$	$(ar{L}^j e_R) \epsilon_{jk} (ar{Q}^k u_R)$				
$Q_{arphi u d}$	$i(\widetilde{arphi}^{\dagger}\overline{D_{\mu}arphi})(ar{u}_{R}\gamma^{\mu}d_{R})$	$Q_{lequ}^{\left(3 ight)}$	$(ar{L}^j\sigma_{\mu u}e_R)\epsilon_{jk}(ar{Q}^k\sigma^{\mu u}u_R)$				

Engel, Ramsey-Musolf, van Kolck, 1303.2371 (Prog.Part.Nucl.Phys.) Chupp, Ramsey-Musolf, van Kolck, 1407.1064 (PRC)

$$\begin{split} \mathcal{O}_{\tilde{G}} &= f^{ABC} \tilde{G}_{\mu}^{A\nu} G_{\nu}^{B\rho} G_{\rho}^{C\mu} \\ \mathcal{O}_{\varphi \tilde{G}} &= \varphi^{\dagger} \varphi \tilde{G}_{\mu\nu}^{A} G^{A\mu\nu} \\ \mathcal{O}_{\varphi \tilde{W}} &= \varphi^{\dagger} \varphi \tilde{W}_{\mu\nu}^{I} W^{I\mu\nu} \\ \mathcal{O}_{\varphi \tilde{W}} &= \varphi^{\dagger} \varphi \tilde{B}_{\mu\nu} B^{\mu\nu} \\ \mathcal{O}_{\varphi \tilde{W}B} &= \varphi^{\dagger} \tau^{I} \varphi \tilde{W}_{\mu\nu}^{I} B^{\mu\nu} \\ \mathcal{O}_{\tilde{W}} &= \epsilon^{IJK} \tilde{W}_{\mu}^{I\nu} W_{\nu}^{J\rho} W_{\rho}^{K\mu} \end{split}$$



FCC-ee is well suited for both Zh and W^+W^- studies!

Also helps V_{cb} extraction See <u>Stephane's</u> talk yesterday

Using angular asymmetries of Zh and aTGC measurements for W^+W^- , the CPV parameters can be extracted, for which we use the optimal observable approach to improve the sensitivity

LHC FCC-ee $_{240 \, GeV}$





LHC FCC-ee_{240 GeV}

Using angular asymmetries of Zh and aTGC measurements for W^+W^- , the CPV parameters can be extracted, for which we use the optimal observable approach to improve the sensitivity

0.02 $\tilde{c}_{\rm WW}$ 0.00 -0.02HL-LHC + FCC-ee (365 GeV, 1.5/ab) + ILC (500 GeV, 1.6+1.6/ab) + CEPC (360 GeV, 1/ab) ILC (250 GeV, 0.9+0.9/ab) FCC-ee (240 GeV, 5/ab) CEPC (240 GeV, 20/ab) 10^1 10 Bosonic CPV couplings 0.02 10^{-2} 10^{-2} $\tilde{c}_{\rm WB}$ 0.00 10^{-3} 10^{-3} 10^{-4} 10--0.02Ĉ_{RR} Ĉww Ĉ_{WB} *C̃*_{3W} -0.02Per-mille level precision reach 0.00 \tilde{c}_{BB} 0.02

Benchmark

Benchmark: <u>Y-Universal Z' model</u>

Extend the SM by $U(1)_{\!_{\mathcal{I}}}$ but without introducing kinetic mixing and off-diagonal gauge couplings





95% CL scale limits on 4–fermion contact interactions from O_{2B}

de Blas, YD, Grojean, Gu, Miralles, Peskin, Tian, Vos, Vryonidou, 2206.08326

FCC-hh? See Matthew's talk yesterday

Yong Du

Benchmark: Leptoquark model

 $\mathscr{L}_{\mathrm{LQ}} \supset \left(\lambda_{i\alpha}^{1L} \bar{q}_{i}^{c} \epsilon \mathscr{E}_{\alpha} + \lambda_{i\alpha}^{1R} \bar{u}_{i}^{c} e_{\alpha}\right) S_{1} + \lambda_{i\alpha}^{3L} \bar{q}_{i}^{c} \epsilon \sigma^{I} \mathscr{E}_{\alpha} S_{3}^{I} + \mathbf{h.c.}$



de Blas, YD, Grojean, Gu, Miralles, Peskin, Tian, Vos, Vryonidou, 2206.08326

Benchmark: Type-II seesaw model

 $V(\Phi, \Delta) \supset \lambda_4(\Phi^{\dagger} \Phi) \mathrm{Tr}(\Delta^{\dagger} \Delta) + \lambda_5 \Phi^{\dagger} \Delta \Delta^{\dagger} \Phi$

$$\mathscr{L}_{Y} = (y_{\nu})_{\alpha\beta} \overline{L_{\alpha}^{c}} i \tau_{2} \Delta L_{\beta} h.c.$$



Benchmark: <u>Unfolding</u>



Find the UV for any operator and any topology (UVBuilder).

<u>Q: Which benchmark model for FCC-ee?</u>

Internal fields								
I1	I2	I3 I4		I5				
	HyperCharges							
_ 2	_ <u>5</u>	<u>1</u>	_ <u>2</u>	<u>4</u>				
3	3	3	3	3				
Gaug	ge infor	mation	$\{503, 500\}$	502}				
$\{3, 1\}$	{3, 2 }	{3, 2 }	$\{3, 1\}$	$\{3, 1\}$				
$\{3, 1\}$	{3, 2}	{3, 2 }	$\{3, 1\}$	$\{6, 1\}$				
$\{3, 1\}$	{3, 2}	{3, 2 }	{ 6 , 1}	$\{3, 1\}$				
$\{3, 1\}$	{3, 2}	{3, 2}	{ 6 , 1}	$\{6, 1\}$				
$\{3, 1\}$	{3, 2}	{ 6 , 2 }	$\{3, 1\}$	$\{3, 1\}$				
$\{3, 1\}$	{3, 2}	{ 6 , 2 }	{ 6 , 1}	$\{3, 1\}$				
{ 3 , 1}	{ 6 , 2 }	{3, 2}	$\{3, 1\}$	$\{3, 1\}$				
{ 3 , 1}	<i>{</i> 6 <i>,</i> 2 <i>}</i>	{3, 2}	{ 6 , 1 }	{ 3 , 1 }				
{ 3 , 1}	{ 6 , 2 }	{ 6 , 2 }	{ 3 , 1}	{ 6 , 1}				
{ 3 , 1}	{ 6 , 2 }	{ 6 , 2 }	{ 6 , 1}	{ 6 , 1}				
<i>{</i> 6, 1 <i>}</i>	{3, 2}	{3, 2}	$\{3, 1\}$	$\{3, 1\}$				
<i>{</i> 6 <i>,</i> 1 <i>}</i>	{3, 2}	{3, 2}	{ 3 , 1}	{ 6 , 1}				
<i>{</i> 6, 1 <i>}</i>	{3, 2}	{3, 2}	{ 6 , 1}	$\{3, 1\}$				
<i>{</i> 6, 1 <i>}</i>	{3, 2}	{3, 2}	{ 6 , 1}	{ 6 , 1}				
<i>{</i> 6, 1 <i>}</i>	{3, 2}	<i>{</i> 6 <i>,</i> 2 <i>}</i>	{ 3 , 1}	$\{3, 1\}$				
<i>{</i> 6, 1 <i>}</i>	{3, 2}	<i>{</i> 6 <i>,</i> 2 <i>}</i>	<i>{</i> 6, 1 <i>}</i>	$\{3, 1\}$				
{ 6 , 1}	{ 6 , 2 }	{3, 2}	{ 3 , 1}	$\{3, 1\}$				
{ 6 , 1}	{ 6 , 2 }	{3, 2}	{ 6 , 1}	$\{3, 1\}$				
{6 , 1}	{ 6 , 2 }	{ 6 , 2 }	{ 3 , 1 }	{ 6 , 1 }				
{ 6 , 1}	{ 6 , 2 }	{ 6 , 2 }	$\{6, 1\}$	$\{6, 1\}$				

Summary

- FCC-ee is an ideal precision machine for new physics study within SMEFT:
 - Unprecedented precision reach for Higgs and EW physics (except 1st gen quarks)
 - Increase sensitivity reach of 4-fermion operator with beam polarization at FCC-ee?
 - Otherwise, merging low-energy data (neutrino for example) may make a difference
 - A new direction can be probed for CPV operators, complementary to the LHC
- FCC-ee also a perfect machine for specific new physics model studies
 - the Z' model
 - the leptoquark model
 - * ...
 - operator unfolding upon anomaly observation?

Summary

- FCC-ee is an ideal precision machine for new physics study within SMEFT:
 - Unprecedented precision reach for Higgs and EW physics (except 1st gen quarks)
 - Increase sensitivity reach of 4-fermion operator with beam polarization at FCC-ee?
 - Otherwise, merging low-energy data (neutrino for example) may make a difference
 - A new direction can be probed for CPV operators, complementary to the LHC
- FCC-ee also a perfect machine for specific new physics model studies
 - the Z' model
 - the leptoquark model
 - * ...
 - operator unfolding upon anomaly observation?

While do not know when to expect $\checkmark \checkmark$ let's witness the birth of FCC-ee solidly together



$2\ell^2 q$ global fit results



de Blas, YD, Grojean, Gu, Miralles, Peskin, Tian, Vos, Vryonidou, 2206.08326

SMEFT global fit: <u>Higgs</u>

Higgs global fit



precision reach on effective couplings from SMEFT global fit

de Blas, YD, Grojean, Gu, Miralles, Peskin, Tian, Vos, Vryonidou, 2206.08326

SMEFT global fit: top

top global fit



de Blas, YD, Grojean, Gu, Miralles, Peskin, Tian, Vos, Vryonidou, 2206.08326

Optimal observables

This will certainly be improved at future colliders, especially with the utilization of optimal observables:

$$\frac{d\sigma(c)}{d\Pi} = \frac{d\sigma_0}{d\Pi} + \sum_j \frac{d\bar{\sigma}_j}{d\Pi} c_j + \cdots$$

$$(\text{Cov})_{jk}^{-1} = \int d\Pi \frac{(d\bar{\sigma}_j/d\Pi)(d\bar{\sigma}_k/d\Pi)}{d\sigma_0/d\Pi} \cdot \int \mathscr{L}$$

The optimal observable analysis is still ongoing, we expect a factor of 10/100 improvement for HL-LHC and future e^+e^- colliders.