

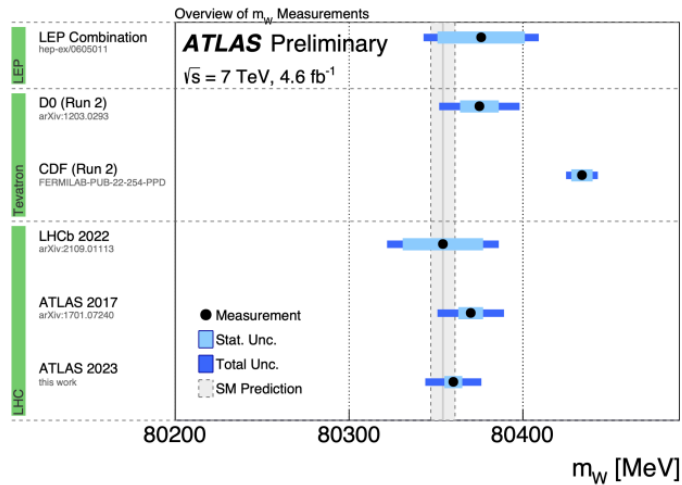
# Compton polarimeters

Aurélien MARTENS (IJCLab Orsay)

On behalf of FCC-ee FS EPOL group

<https://indico.cern.ch/category/8678/>

# Physics requirements



Observable	statistics	$\Delta\sqrt{s}_{\text{abs}}$ 100 keV	$\Delta\sqrt{s}_{\text{syst-ptp}}$ 40 keV	calib. stats. 200 keV/ $\sqrt{N^i}$	$\sigma_{\sqrt{s}}$ 85 ± 0.05 MeV
$m_Z$ (keV)	4	100	<b>28</b>	1	–
$\Gamma_Z$ (keV)	4	2.5	<b>22</b>	1	<b>10</b>
$\sin^2 \theta_W^{\text{eff}} \times 10^6$ from $A_{\text{FB}}^{\mu\mu}$	2	–	<b>2.4</b>	0.1	–
$\frac{\Delta\alpha_{\text{QED}}(m_Z^2)}{\alpha_{\text{QED}}(m_Z^2)} \times 10^5$	3	0.1	<b>0.9</b>	–	<b>0.1</b>

Required accuracy of <1ppm

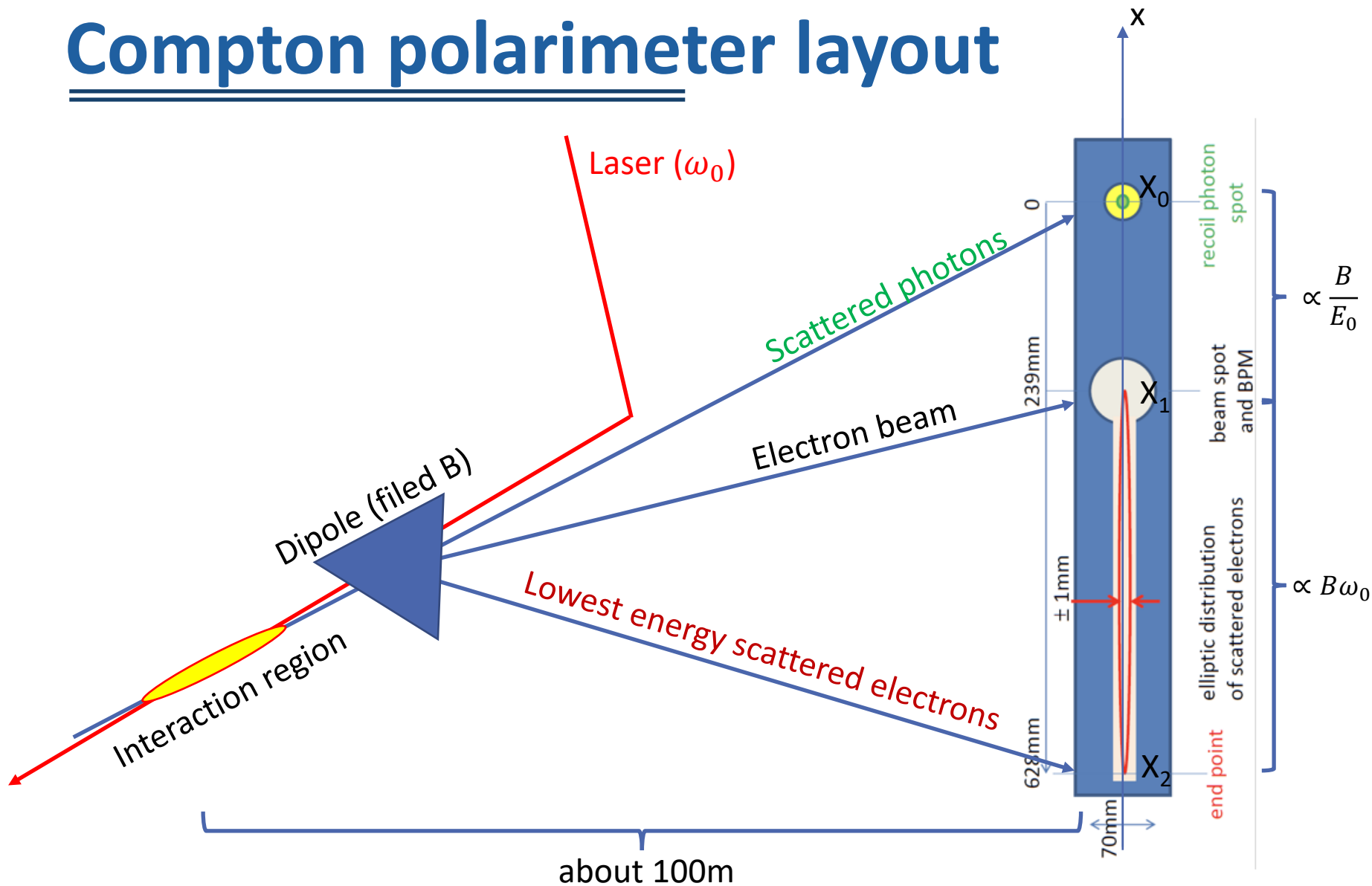
High reproducibility of measurements for various sqrt(s) is critically needed

Extract as much information as possible from physics experiments themselves (crossing angle, luminosity, sqrt(s) spread)

Beam-based measurements in real time, including beams energy with resonant depolarization

24/7 operable measurement of (de-)polarization

# Compton polarimeter layout



New concept (N. Yu Muchnoi) to measure all polarization parameters  $\rightarrow$  3D polarimeter

# Integration

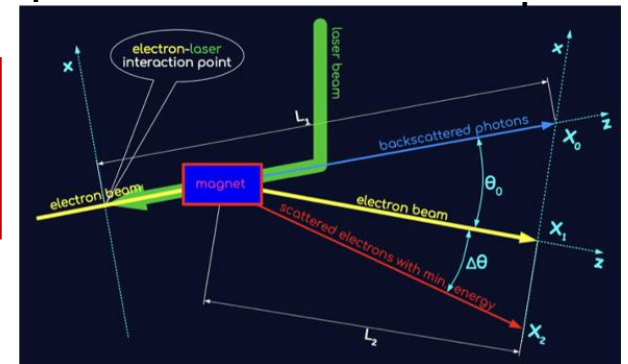
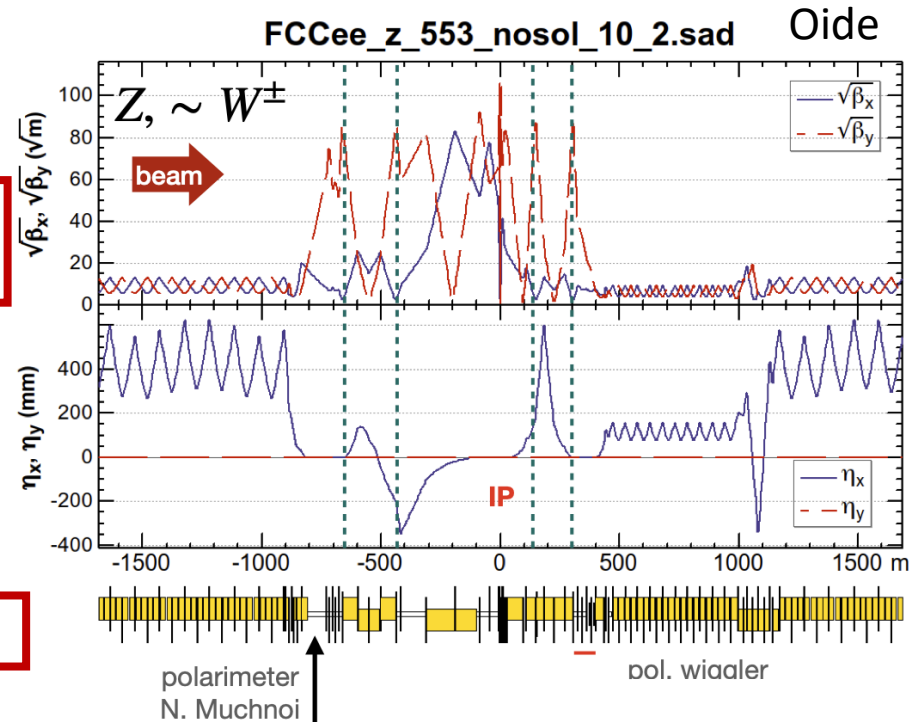
Polarimeter integrated close to experiments IP  
1/beam (baseline)

e-beam size at Compton IP :  
~500 $\mu$ m (horizontal)

Robust laser technology is available nowadays

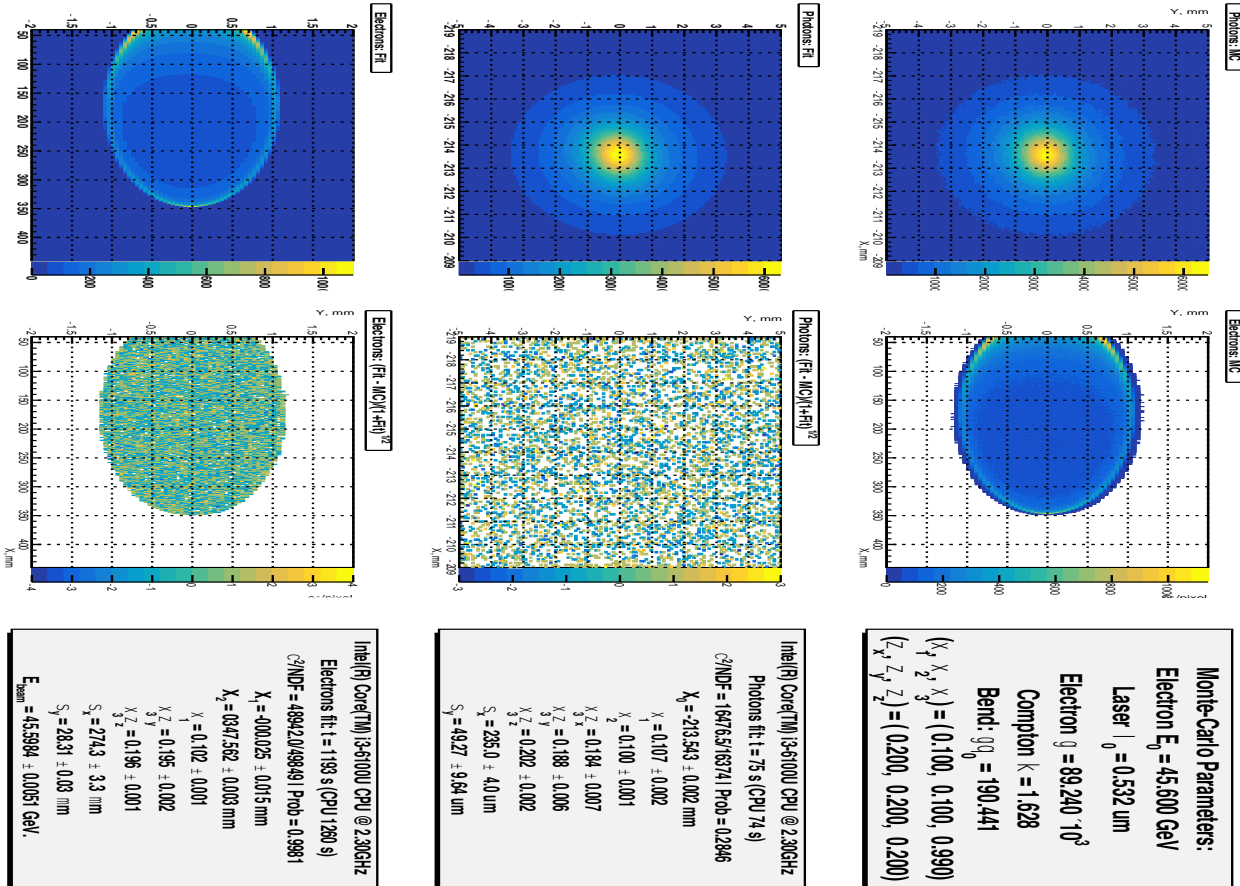
Demanding operational constraints

- laser room in the parallel technical gallery
- 24/7 access to the laser system and related electronics



# Transverse distributions

Based on measurement of scattered particles transverse distributions (pixelized detectors)



All components extracted with  $\sim 0.001$  precision in few seconds

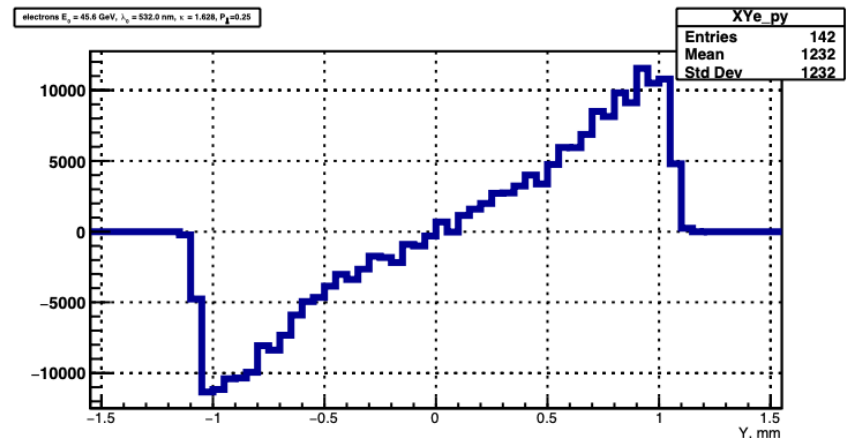
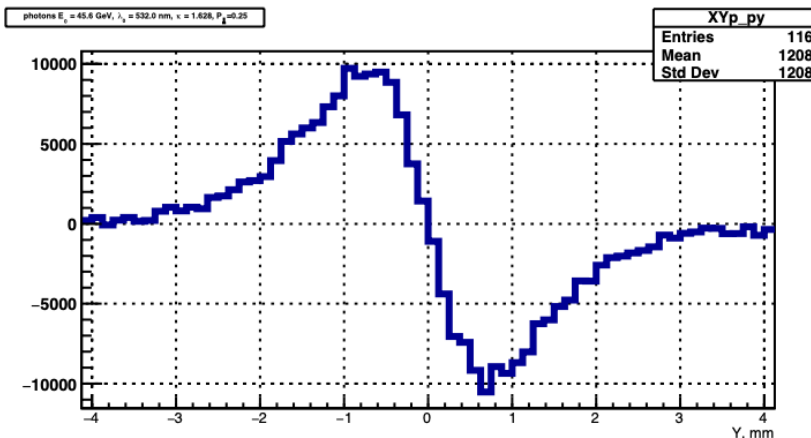
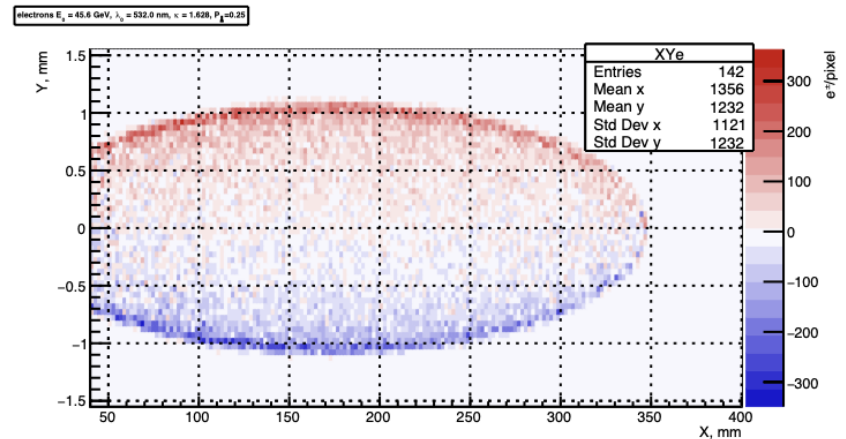
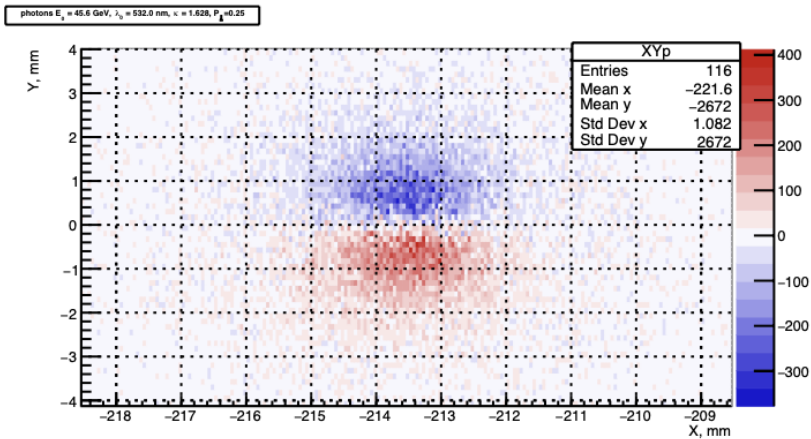
Beam energy may be extracted too!  $\rightarrow$  redundancy with RDP

Realistic detector specifications (still) to be drawn

With the constraints of being tolerant to radiations

# Laser helicity asymmetries

Blondel et al., arXiv:1909.12245



Reproducible and well known laser helicity flip is required

# QED corrections

## Complete order- $\alpha^3$ calculation of the cross section for polarized Compton scattering

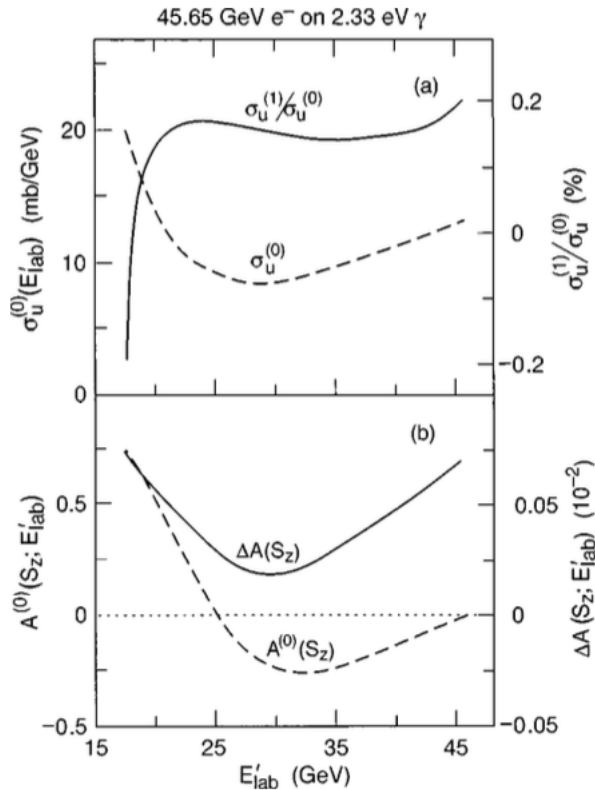
Morris L. Swartz

Stanford Linear Accelerator Center, Stanford University, Stanford, California 94309

(Received 24 November 1997; published 28 May 1998)

The construction of a computer code to calculate the cross sections for the spin-polarized processes  $e^- \gamma \rightarrow e^- \gamma, e^- \gamma \gamma, e^- e^+ e^-$  to order  $\alpha^3$  is described. The code calculates cross sections for circularly polarized initial-state photons and arbitrarily polarized initial-state electrons. The application of the code to the SLD Compton polarimeter indicates that the order- $\alpha^3$  corrections produce a fractional shift in the SLD polarization scale of  $-0.1\%$  which is too small and of the wrong sign to account for the discrepancy in the Z-pole asymmetries measured by the SLD Collaboration and the CERN LEP Collaborations.  
[S0556-2821(98)03413-4]

Studied in details at SLD



Measurement of transverse polarization at FCCee :

photons  $\frac{\delta P}{P} \approx 1 \times 10^{-3}$  ( $0.5 \times 10^{-3}$ ) at 45 (80) GeV

electrons  $\frac{\delta P}{P} \approx 4 \times 10^{-3}$  ( $10 \times 10^{-3}$ ) at 45 (80) GeV

Measurement of longitudinal polarization at FCCee :

$\frac{\delta P}{P} \approx 1 \times 10^{-3}$  at 45 GeV

If and only if laser helicity asymmetries are measured

# Magnetic field tolerancing

Many potential sources of 'bending angle' uncertainties (for instance genuine inhomogeneities of B-field, short-/long-term fluctuations of currents, temperatures, alignments)

Over the useful aperture of the magnet: 
$$\frac{\sigma(\int B_y dl)}{\int B_y dl} \ll 2 \times 10^{-4}$$

Fringe fields also may affect performance of polarimeter

$$\int B_x dl \ll \frac{\sigma_y \gamma mc}{L_2 q} \approx 1.1 \times 10^{-4} \text{ T.m and}$$
$$\int B_z dl \ll \frac{\sigma_y \gamma mc}{L_2 \kappa \theta_0 q} \approx 3.2 \times 10^{-2} \text{ T.m.}$$

Nominal vertical field for reference:

$$\int B_y dl = \theta_0 \gamma \frac{mc}{q} \approx 0.3 \text{ T.m.}$$

By product: angular alignment

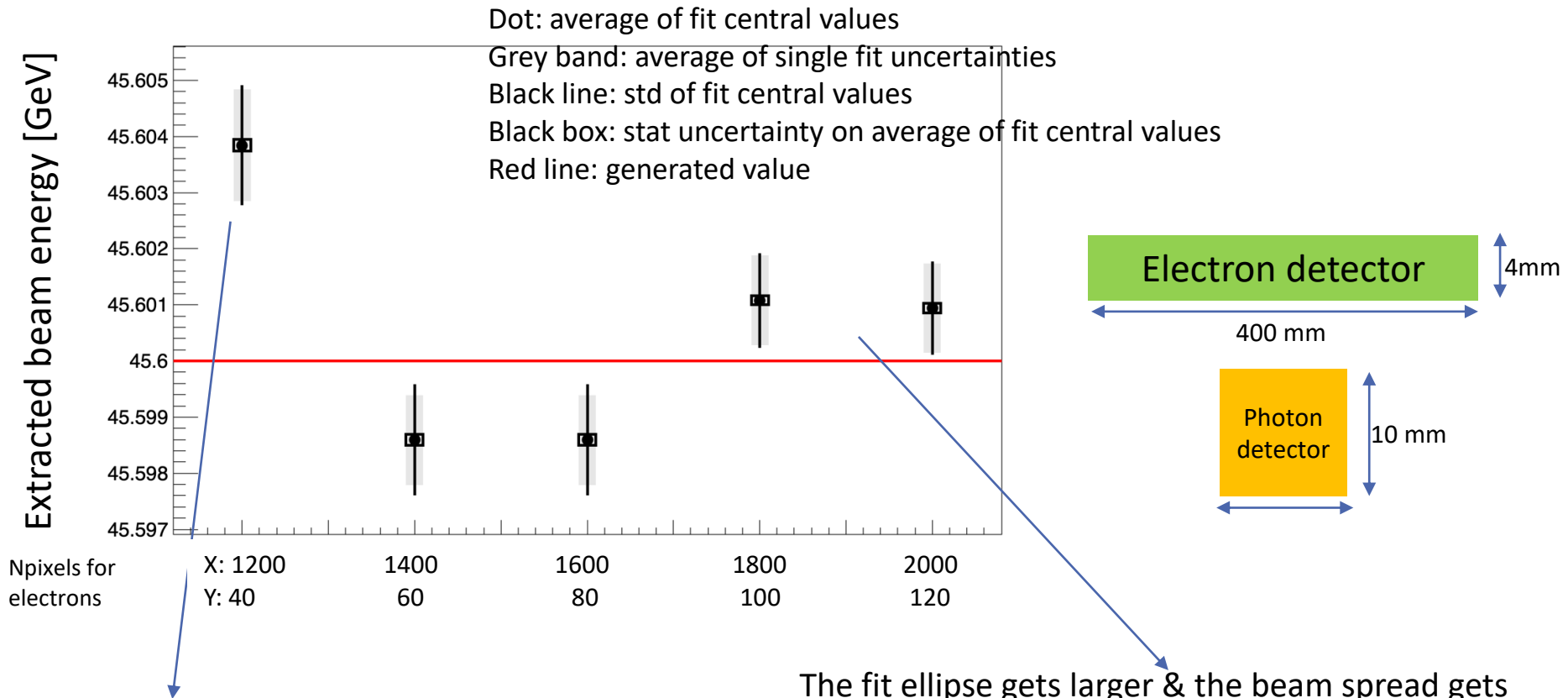
$$\delta_B \ll \frac{\sigma_y \gamma mc}{L_2 \int B_y dl q} \approx 370 \mu\text{rad.}$$

**NB: Requirements not met → not a show-stopper but detailed studies required**



# Preliminary study of pixel size

A toy MonteCarlo procedure is applied (100 experiments,  $10^8$  events each)  
Fit results for few different pixel sizes



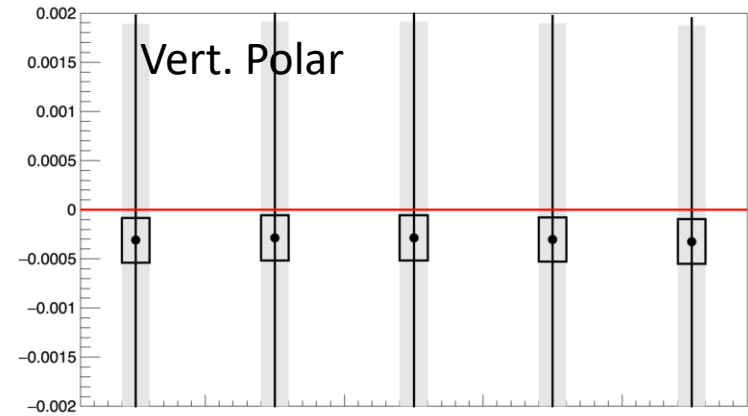
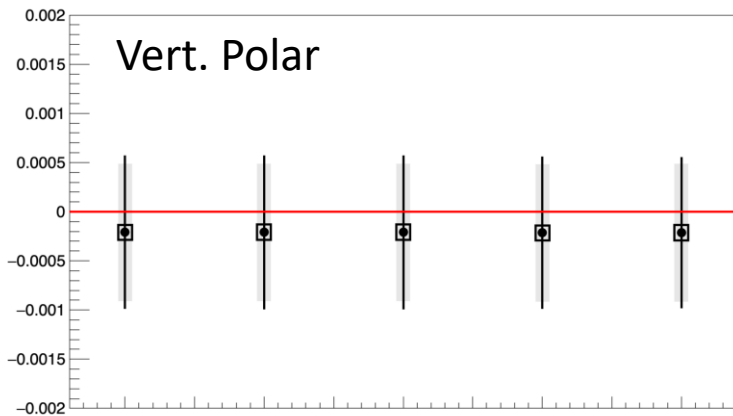
Biased extrapolation of the 'beam' position

The fit ellipse gets larger & the beam spread gets smaller → likely inconsistency between model and data (some approx. made for the fit)

# Polarization extraction: performance

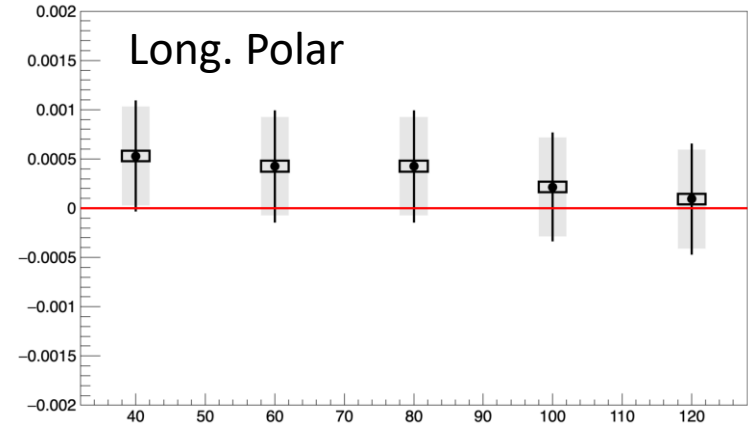
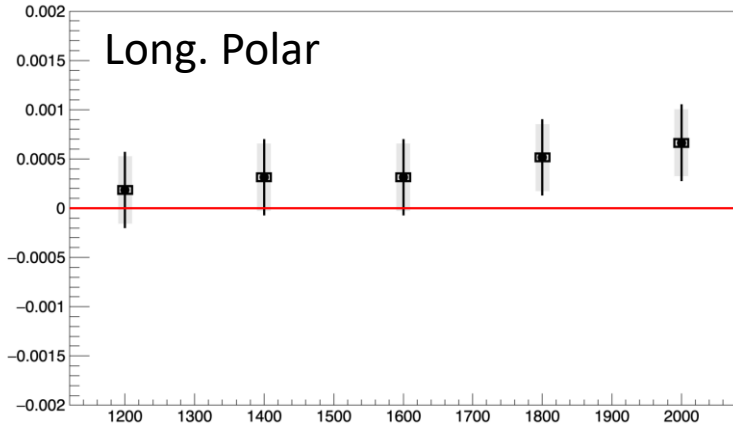
Electron detector

Photon detector



Pixels: X: 1200 1400 1600 1800 2000  
Y: 40 60 80 100 120

Pixels: X: 40 60 80 100 120  
Y: 40 60 80 100 120



Bias ( $\sim 2-6 \cdot 10^{-4}$ )  $\rightarrow$  work in progress to understand

# Conclusion & prospects

High accuracy and precision beam energy measurement with pilots  
24/7 operable Compton polarimeters (1/beam)

Versatile industrial  
and robust laser  
system

Ability to access  
lasers during  
operations

Laser beam  
transport

R&D  
Real time laser  
polarization  
monitoring

Magnet  
tolerancing

Magnet design

Started, but huge task

Pixel detectors  
tolerancing

Just starting  
Pixel detectors  
design

Assessment of systematic  
uncertainties

# backup

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# Physics requirements cont'd

## Importance of longitudinal polarisation measurement

Any residual longitudinal-polarisation will bias cross sections & forward-backward asymmetries (indeed, high longitudinal polarisation is actually useful, but we assume we are not in that regime – rather longitudinal polarisation is a nuisance).

Consider forward-backward asymmetry of  $b\bar{b}$  at Z pole:  $A_{FB}^b = \frac{3}{4} \mathcal{A}_e \mathcal{A}_b$

where in the SM  $\mathcal{A}_e \approx 0.15$ ,  $\mathcal{A}_b \approx 0.95 \Rightarrow A_{FB}^b \approx 0.11$

Now, if there is longitudinal polarisation, asymmetry becomes:  $(A_{FB}^b)' = \frac{3}{4} \mathcal{A}'_e \mathcal{A}_b$

where  $\mathcal{A}'_e = -\left(\frac{\mathcal{A}_e - P}{1 - \mathcal{A}_e P}\right)$  with  $P = \frac{(P_z)_{e^-} - (P_z)_{e^+}}{1 - (P_z)_{e^-} (P_z)_{e^+}}$

and  $(P_z)_{e^\pm}$  the longitudinal polarisation of the  $e^\pm$ .

21/9/22

EPOL requirements at FCC-ee  
Guy Wilkinson

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## Importance of longitudinal polarisation measurement

Any residual longitudinal-polarisation will bias cross sections & forward-backward asymmetries (indeed, high longitudinal polarisation is actually useful, but we assume we are not in that regime – rather longitudinal polarisation is a nuisance).

So, if  $(P_z)_{e^-} = (P_z)_{e^+}$  (no reason to be so) =  $10^{-5}$  (ballpark guess)

$$P = 2 \times 10^{-5} \Rightarrow \frac{(A_{FB}^b)' - A_{FB}^b}{A_{FB}^b} = 1.3 \times 10^{-4}$$

Statistical uncertainty on  $A_{FB}^b$  around  $2 \times 10^{-5}$  (relative), and QCD uncertainty which will probably be larger. Still, to be safe we would want to control  $P_z$  to  $< 10^{-5}$ .

How is this to be done? Measurements must be made on colliding bunches, where scattering rates are lower. Can we sample all bunches? Will it prove necessary to depolarise the physics bunches? If so, we will still need to monitor residual effects. And what are the systematics on an absolute measurement?

Note also, that calculations required to transport the measurement of 3-vector at polarimeter to  $P_z$  value at the interaction points. How can this be cross checked?

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High accuracy longitudinal polarization measurement is needed

→ Naturally small at IPs but with what accuracy?

→ Measure it!

# Compton cross-section

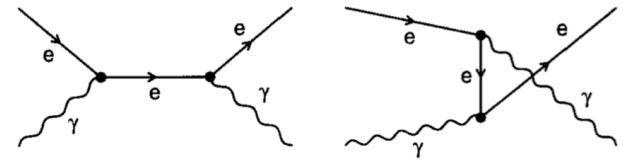


Fig. 1. Tree diagrams for  $e^- \gamma \rightarrow e^- \gamma$

$$x = \frac{2E_0\omega_0}{m^2} (1 + \cos \alpha) \quad y = \frac{E_\gamma}{E_0}$$

The Compton cross-section averaged over scattered particles spins:

Differential cross-section

Transverse laser polarisation: nuisance parameter to minimize and keep under control

Transverse electron beam polarisation: intervenes as an asymmetry in the transverse plane

$$\frac{d\sigma}{dyd\varphi_{obs}}(x, y) = \frac{d\sigma_0}{dy}(x, y) + \frac{d\sigma_\perp}{dy}(x, y) \cos(2(\varphi_{obs} - \varphi_{las})) \mathcal{P}_L^{las} + \frac{d\sigma_\parallel}{dy}(x, y) \mathcal{P}_C^{las} (P_T f_T(x, y) \cos(\varphi_{obs} - \varphi_{elec}) + P_Z f_Z(x, y))$$

*Electron beam polarization independent*
*Electron beam polarization dependent*

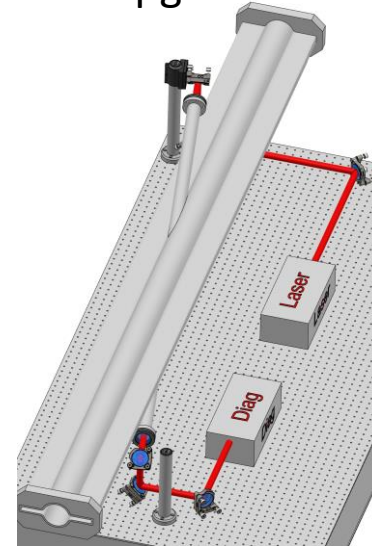
⚠ But small opening angle of scattered particles:

- Electrons → spectrometer
- Photons → difficult to measure asymmetric distribution of a narrow spot → long lever arm needed

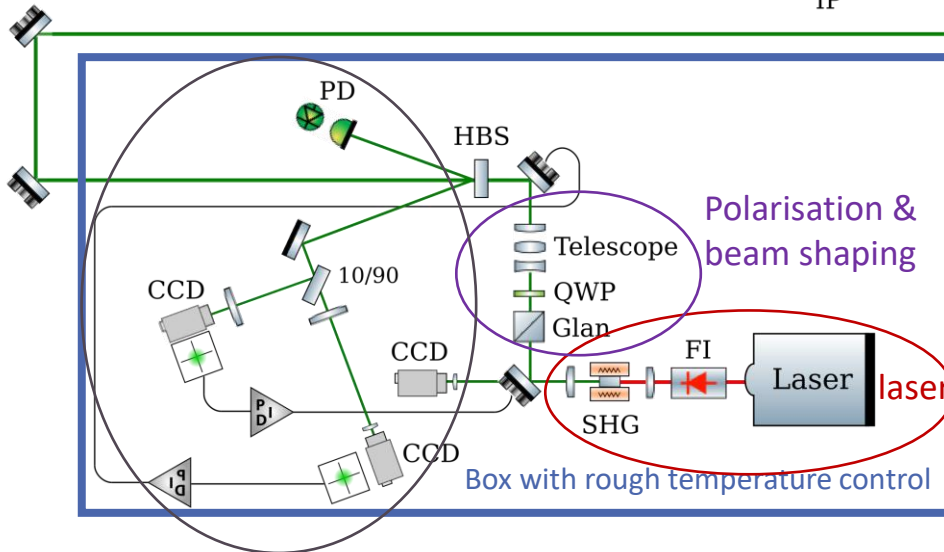
# Laser integration

## Some constraints

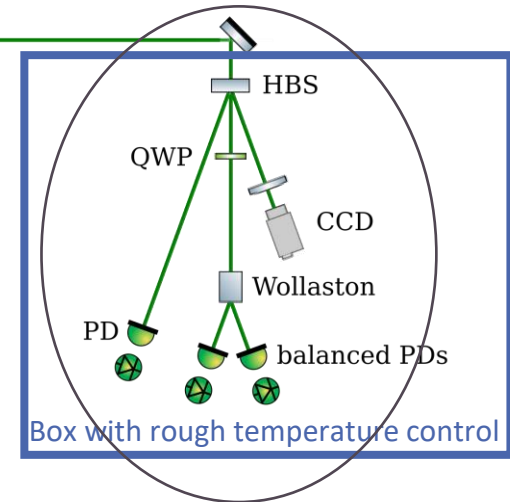
- Small crossing angles are preferred (cross-section, beam jitters) few mrad typically
- Beams crossing plane neither horizontal nor vertical
- beam impedance
- beam induced currents in metallic parts ← avoid
- mechanical stability
- ease of maintenance works



IP



Position, pointing control and monitoring  
 Polarisation independent intensity monitoring  
 Optical spectrum monitoring possible



Polarisation monitoring  
 Duplicated at injection  
 Add Position and pointing monitoring  
**R&D needed to reach required perf.**

**24/7 operable laser system, with full monitoring, remote control**

# Some laser systems

Nikolai's baseline (Q-switch Nd-YAG)

Versatile Yb system

Laser param.	1 pilot	1 pilot v2	All colliding bunches (at Z)
Repetition rate	3 kHz	3 kHz	30 kHz
Pulse energy	1 mJ	1 mJ	10x0.05mJ
Pulse duration	3 ns	30 ps (**)	30 ps
Average power	3 W	3 W	15 W (***)
Scattering rate	$3 \times 10^5/s$ (*)	$3 \times 10^5/s$ (****)	$4 \times 10^7/s$ (****)
Scattering rate per bunch	$3 \times 10^5/s$ (*)	$3 \times 10^5/s$	$4 \times 10^5/s$

adaptable

Same oscillator may be used but two different amplification schemes

(\*) crossing angle  $\sim 2\text{mrad}$

(\*\*) related to optical bandwidth  $\leftarrow$  constrains resolution of 'direct' energy measurement from polarimeter

(\*\*\*) Can be increased to typically  $\sim 100\text{W}$  (nowadays) but requires operational validation

(\*\*\*\*) not limited by Piwinski contribution  $\rightarrow$  can be several degrees without affecting rate



# Scattering rates

Compton cross-section

Laser-beam single pulse energy

Electron bunch charge  
(25nC for colliding bunches,  
1.5nC for pilots)

Geometrical factor

Photon rate  $n = \sigma_C \frac{\epsilon Q}{E_\lambda q} \frac{\mathcal{F}}{2\pi\sigma_y\sigma_x}$

- $\mathcal{F}^{-1} = \sqrt{1 + \left(\frac{\sigma_z}{\sigma_x} \tan \frac{\theta_0}{2}\right)^2}$
- $\theta_0 \sim 2\text{mrad}$

Transverse beam sizes:

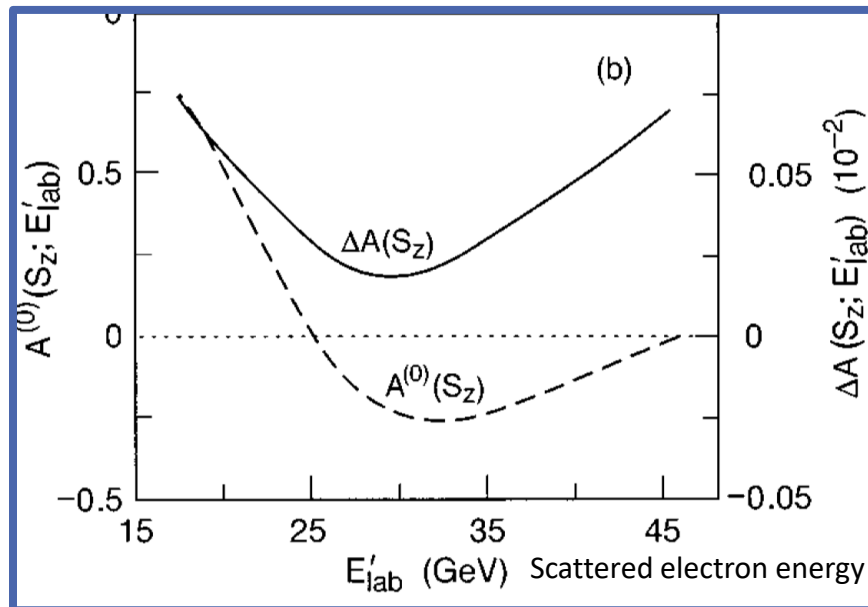
- $\sigma_{x,y,z} = \sqrt{\sigma_{x,y,z,laser}^2 + \sigma_{x,y,z,e-}^2}$
- $\sigma_{x,laser} = \sigma_{y,laser} = 1000\mu\text{m}$
- $\sigma_{x,e-} = 500\mu\text{m}, \sigma_{y,e-} = 10\mu\text{m}, \sigma_{z,e-} \sim 10\text{mm}$

Laser photon energy  
(2.4eV for 0.5μm wavelengths)

# QED corrections

$$\frac{d\sigma}{dE'}(E') \cong \frac{d\sigma_0}{dE'}(1 + \delta) [1 + \mathcal{P}_Z \mathcal{P}_{C,las}(A + \Delta A)]$$

QED corrections < 0.001 @ 45 GeV



Need to be eventually included in simulations...

# The Compton process

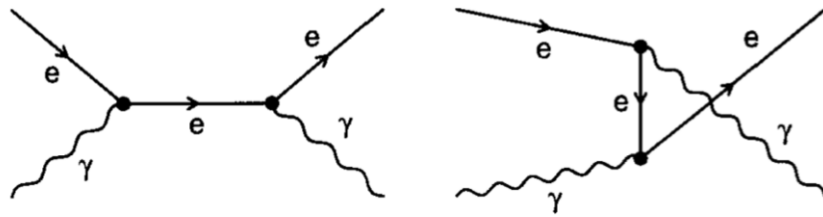


Fig. 1. Tree diagrams for  $e^- \gamma \rightarrow e^- \gamma$

