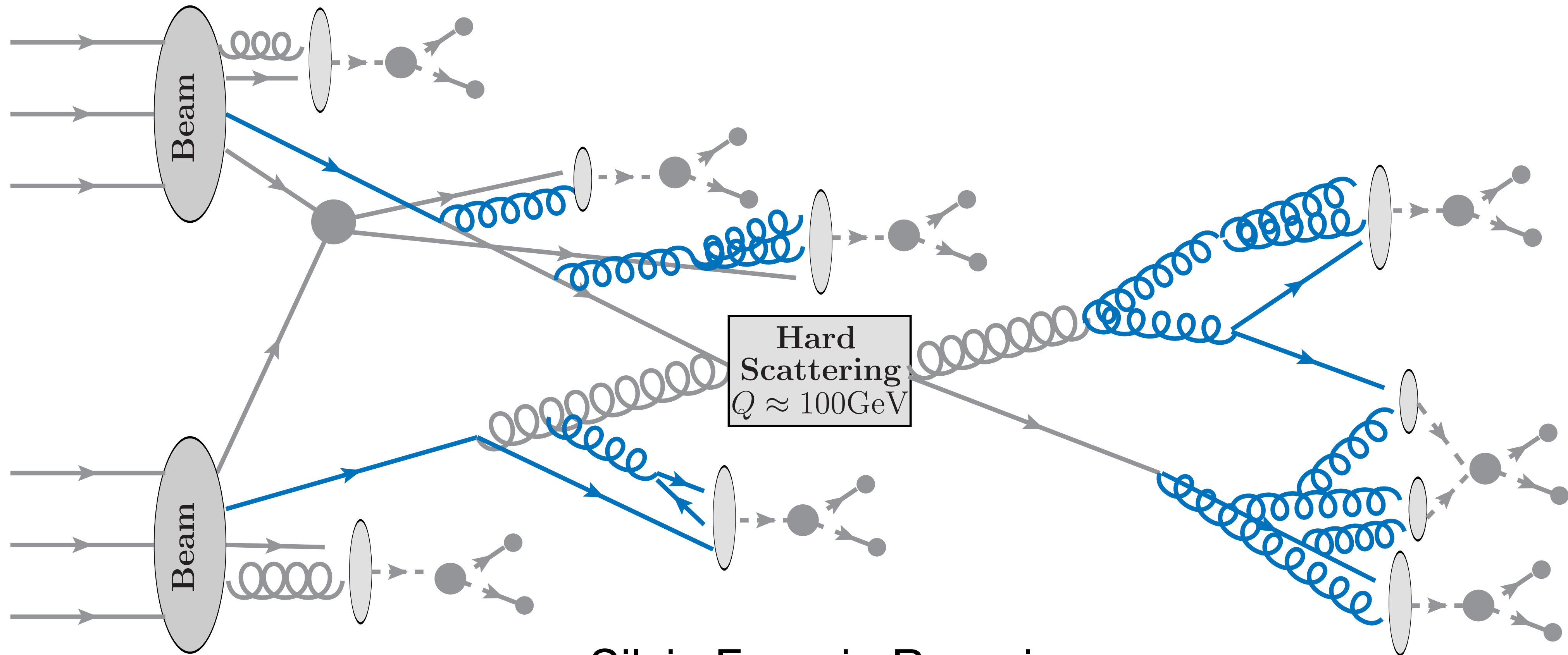


Recent Progresses in Parton Showers with higher logarithmic accuracy



Silvia Ferrario Ravasio

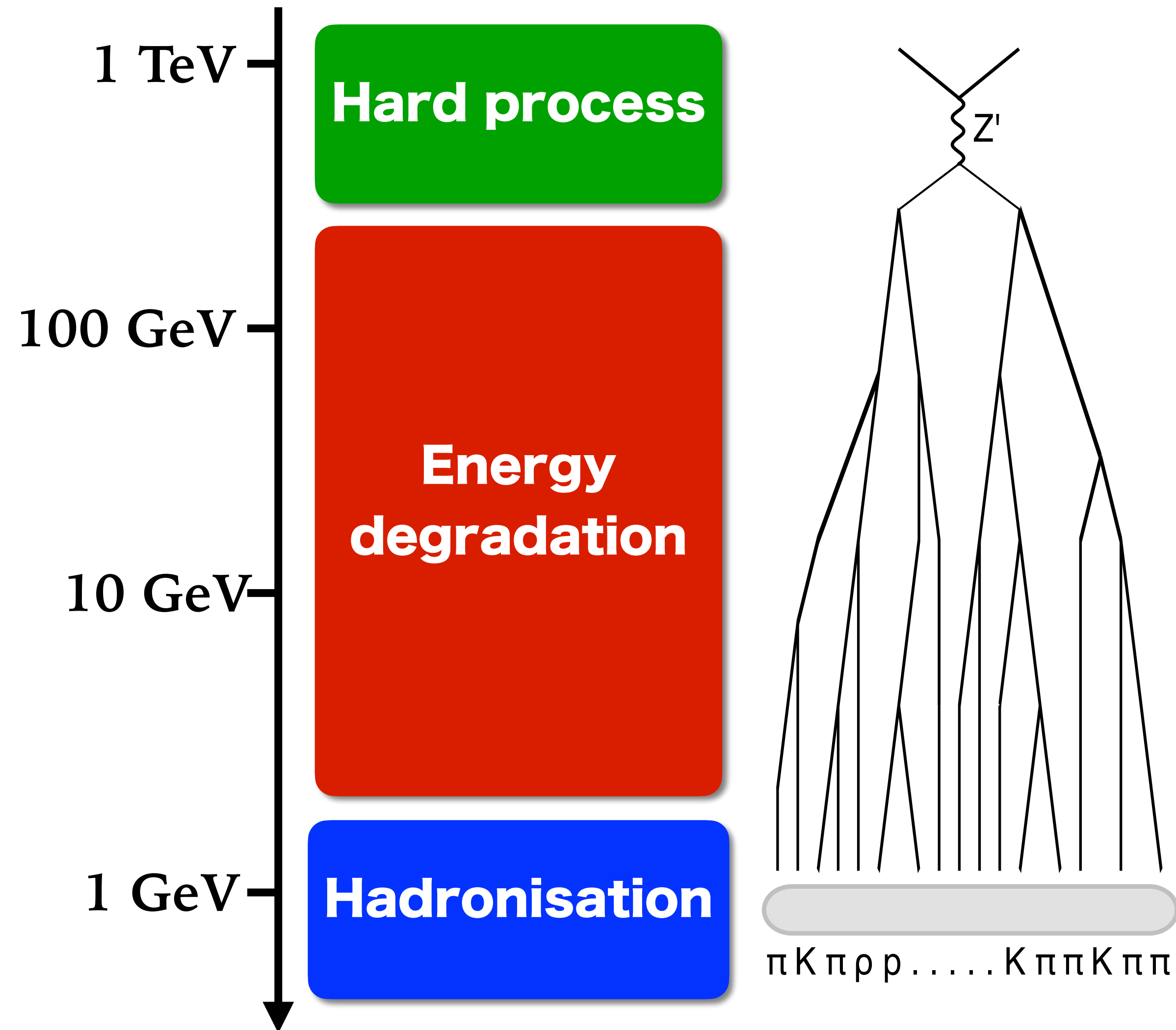
7th FCC Physics Workshop

31st January 2024, Laboratoire d'Annecy de physique des particules



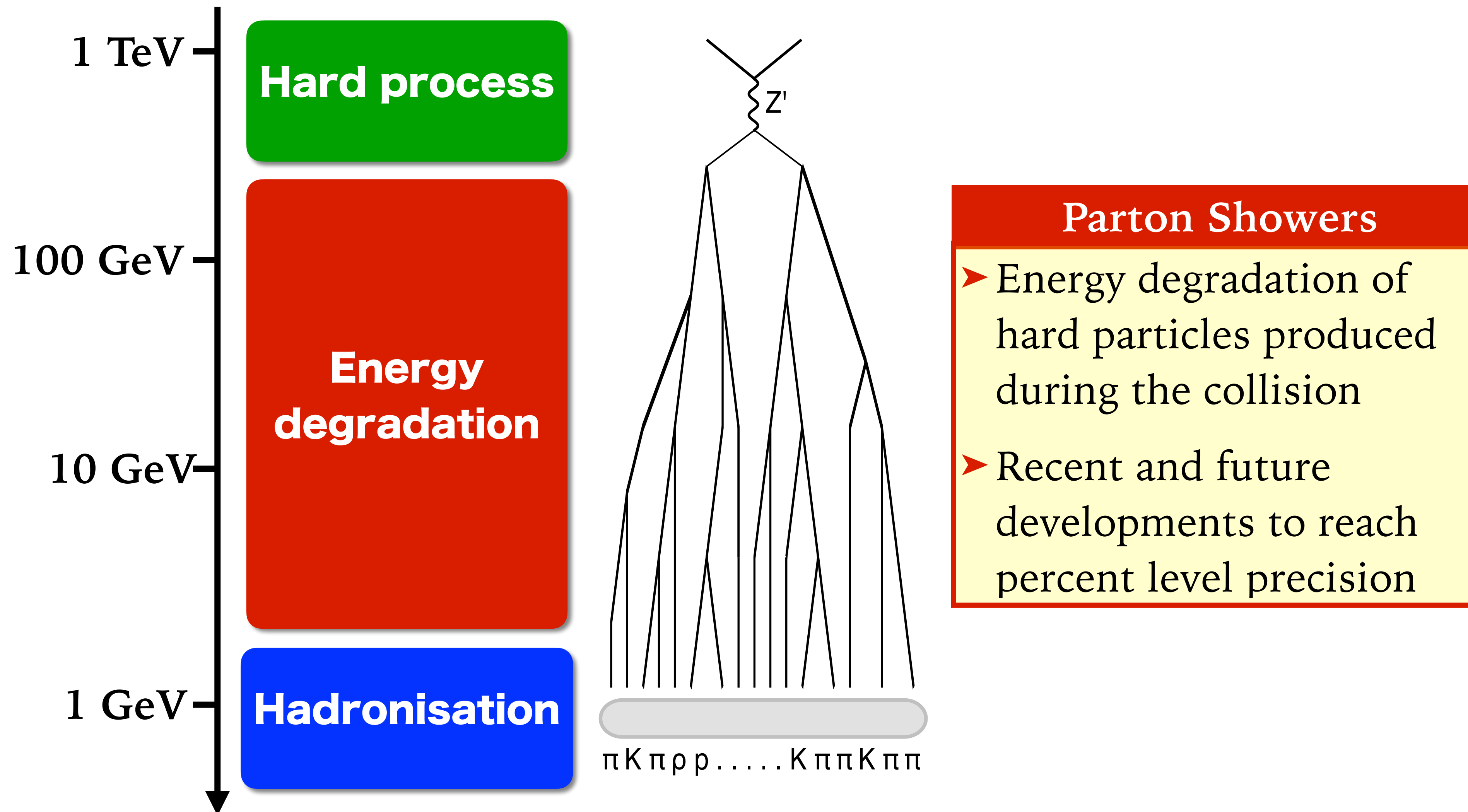
Shower Monte Carlo event generators

SHOWER MONTE CARLO EVENT GENERATORS = default tool for interpreting collider data



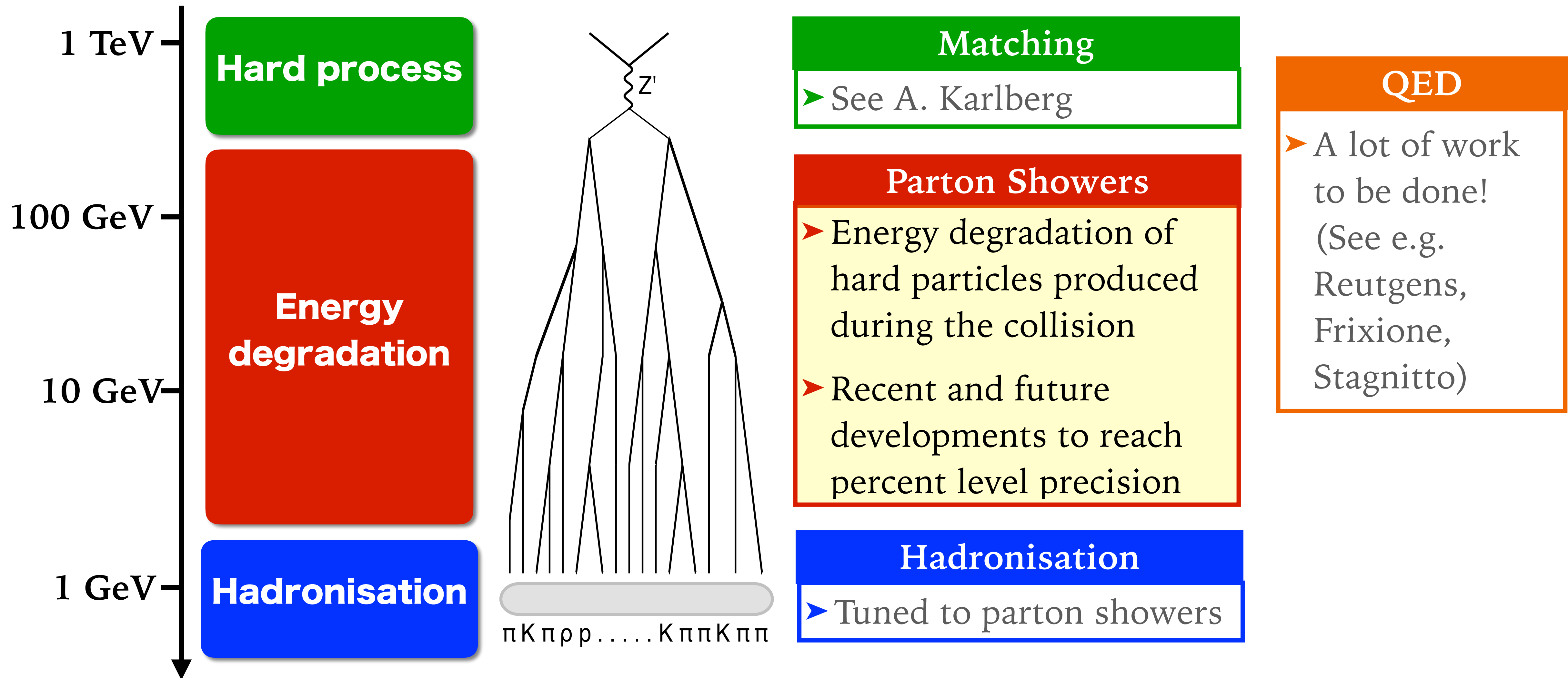
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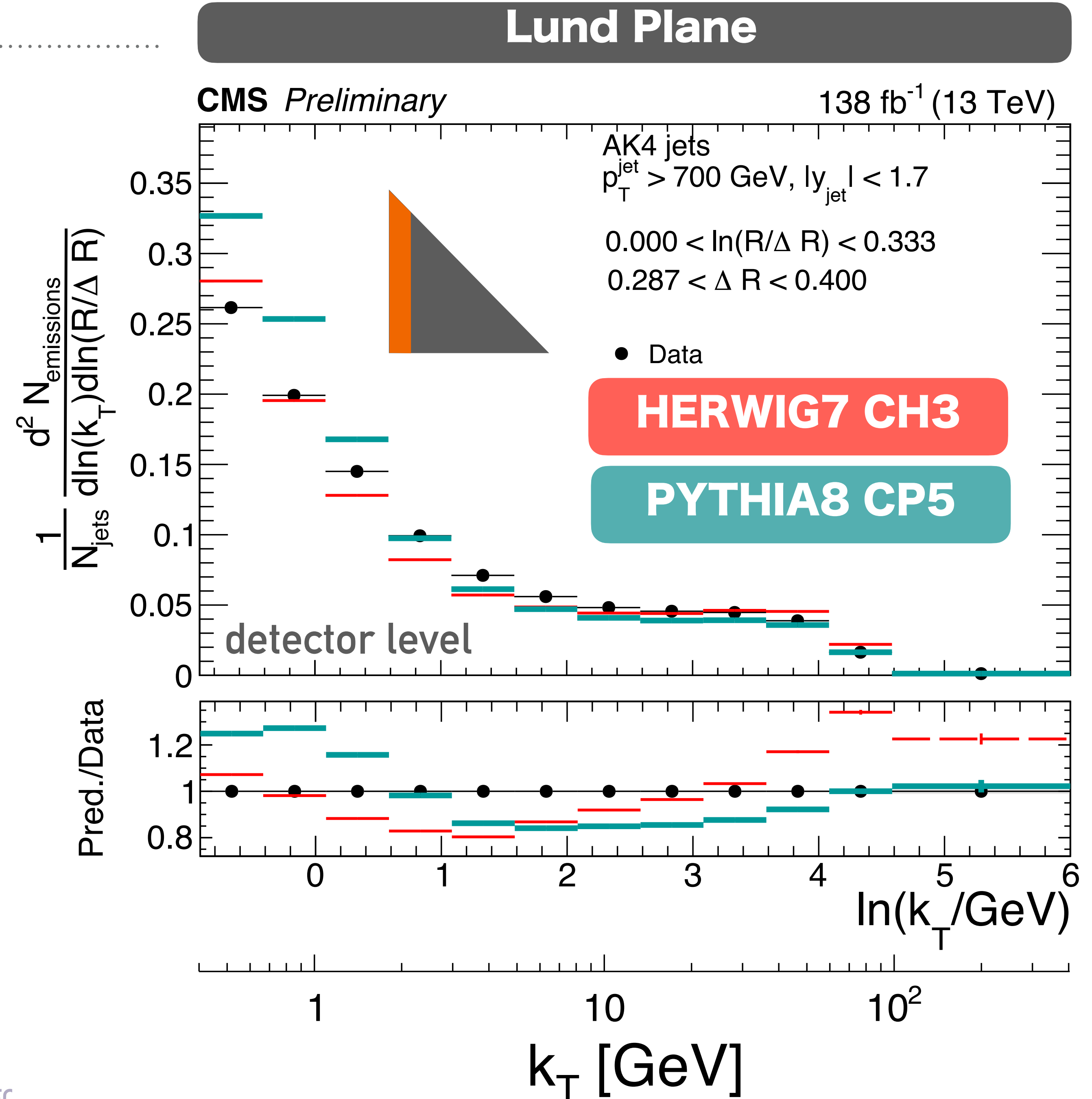
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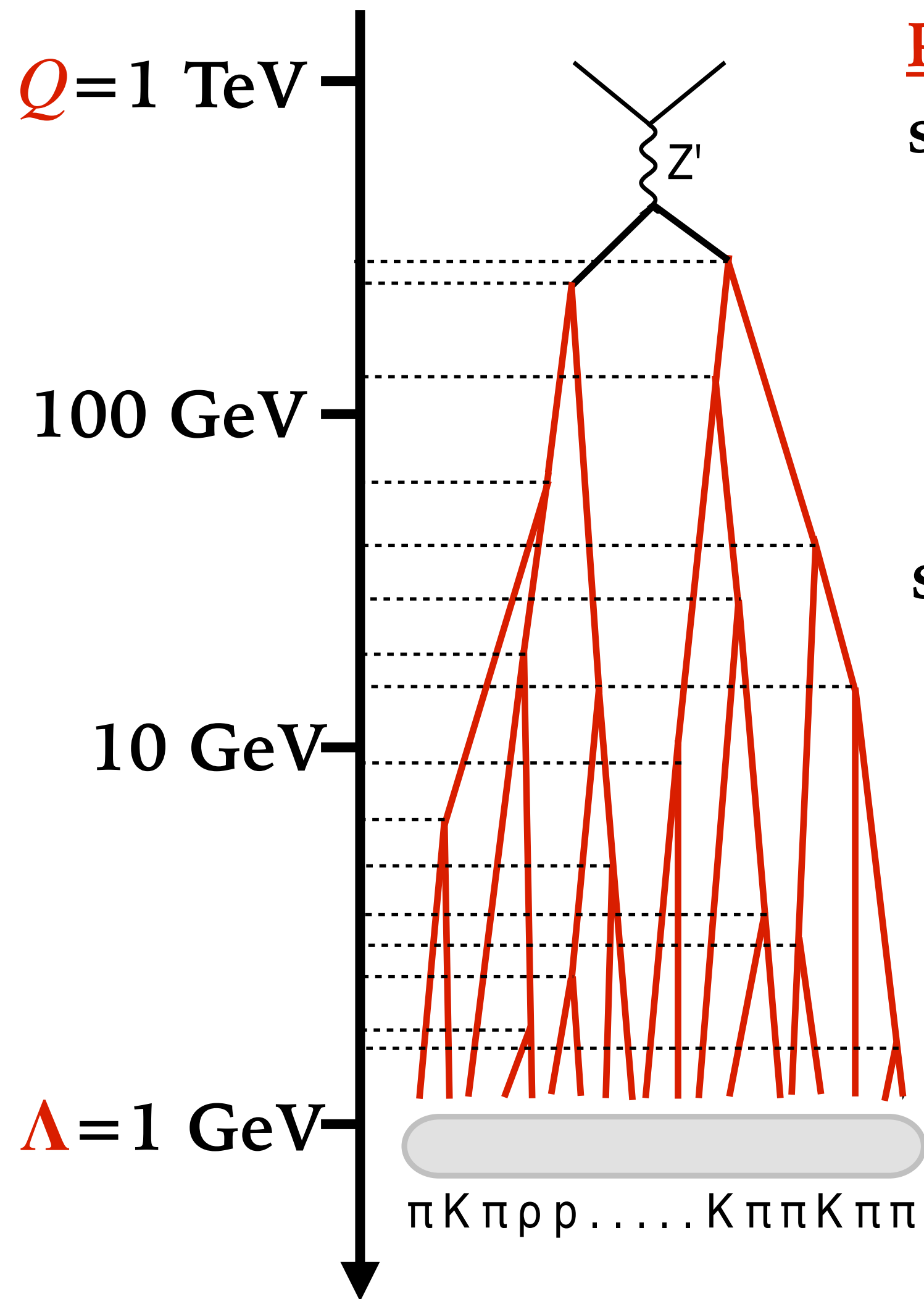


Are current showers good enough?

- showers do an amazing job on many observables for **LHC**
- various places see **10–30% discrepancies** between showers and data
- A lot of work is required to meet the precision target of the **FCC!**



Logarithmically-accurate Parton Showers



PARTON SHOWERS = energy degradation via an iterated sequence of softer and softer emissions

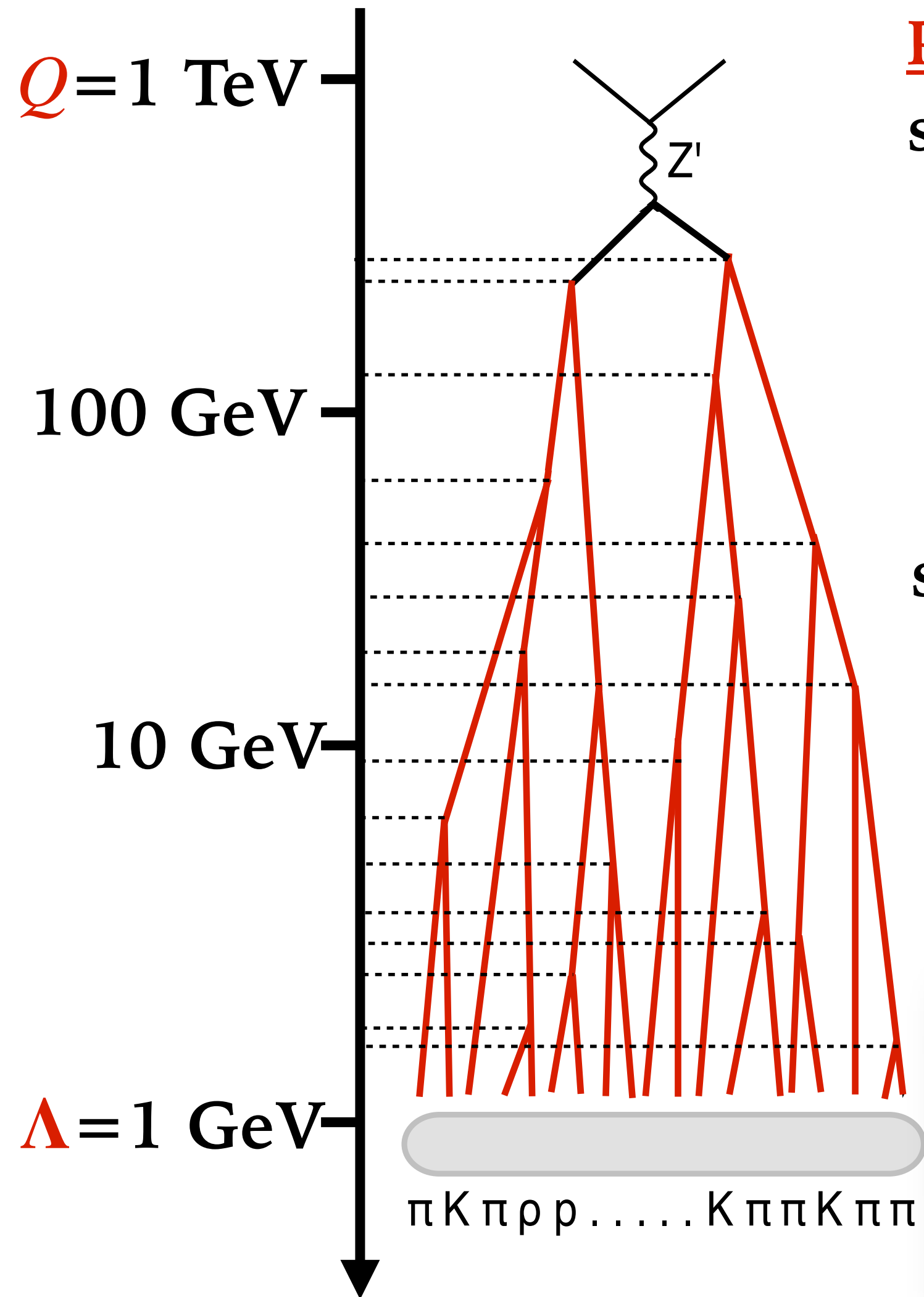
$$L = \ln \frac{Q}{\Lambda} \gg 1$$

simple algorithm to include the **dominant radiative corrections** at all orders for **any observable!**

$$\Sigma(O < e^{-L}) = \exp \left(-L g_{LL}(\beta_0 \alpha_s L) + \dots \right)$$

LL = leading logs

Logarithmically-accurate Parton Showers



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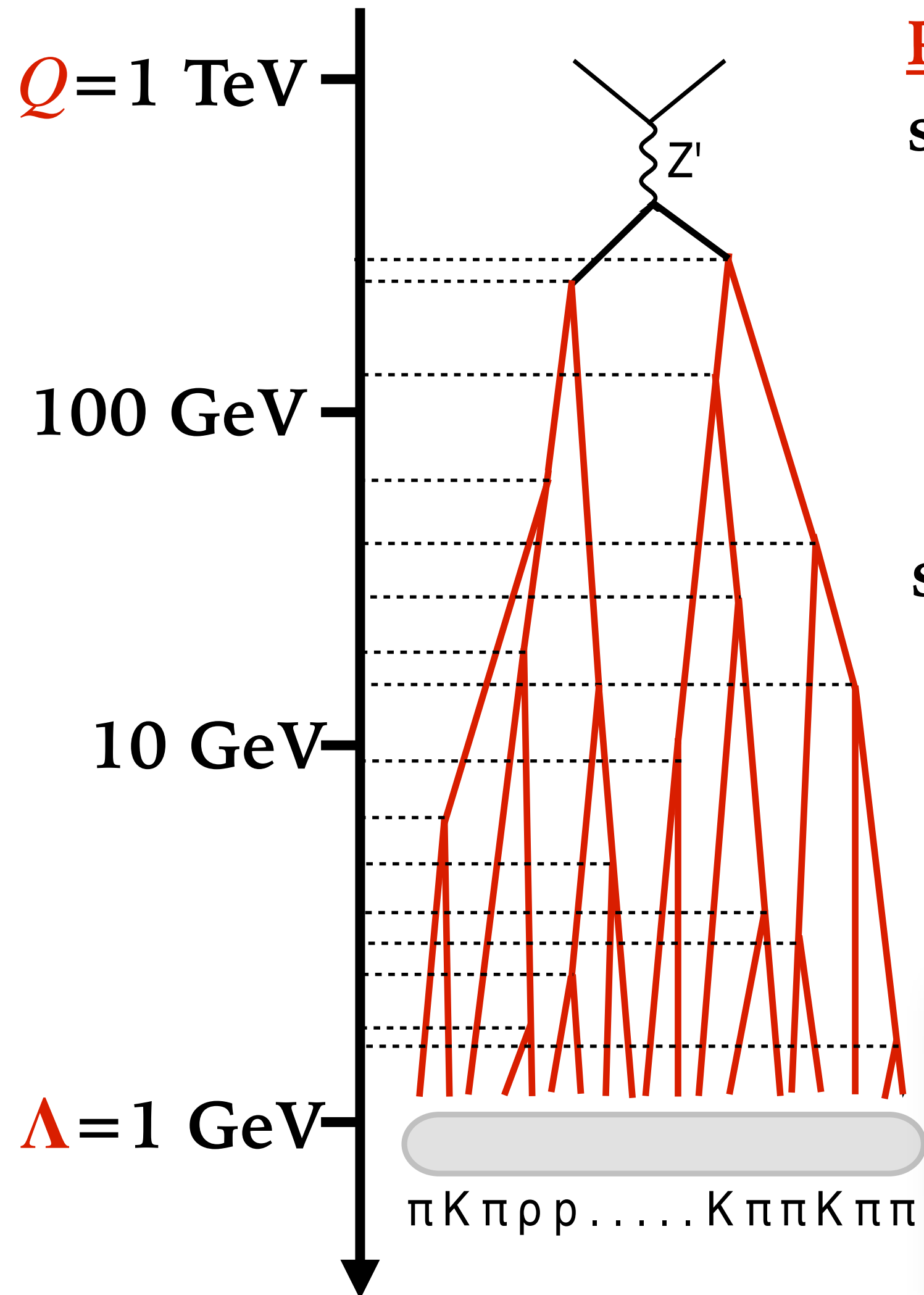
simple algorithm to include the **dominant radiative corrections** at all orders for **any observable!**

$$\Sigma(O < e^{-L}) = \exp \left(-L g_{\text{LL}}(\beta_0 \alpha_s L) + \boxed{g_{\text{NLL}}(\beta_0 \alpha_s L)}^{??} + \dots \right)$$

For $Q \sim 50 - 10000 \text{ GeV}$, $\beta_0 \alpha_s L \sim 0.3 - 0.5$:

Next-to-Leading Logarithms needed for quantitative predictions!

Logarithmically-accurate Parton Showers



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simple algorithm to include the **dominant radiative corrections** at all orders for **any observable!**

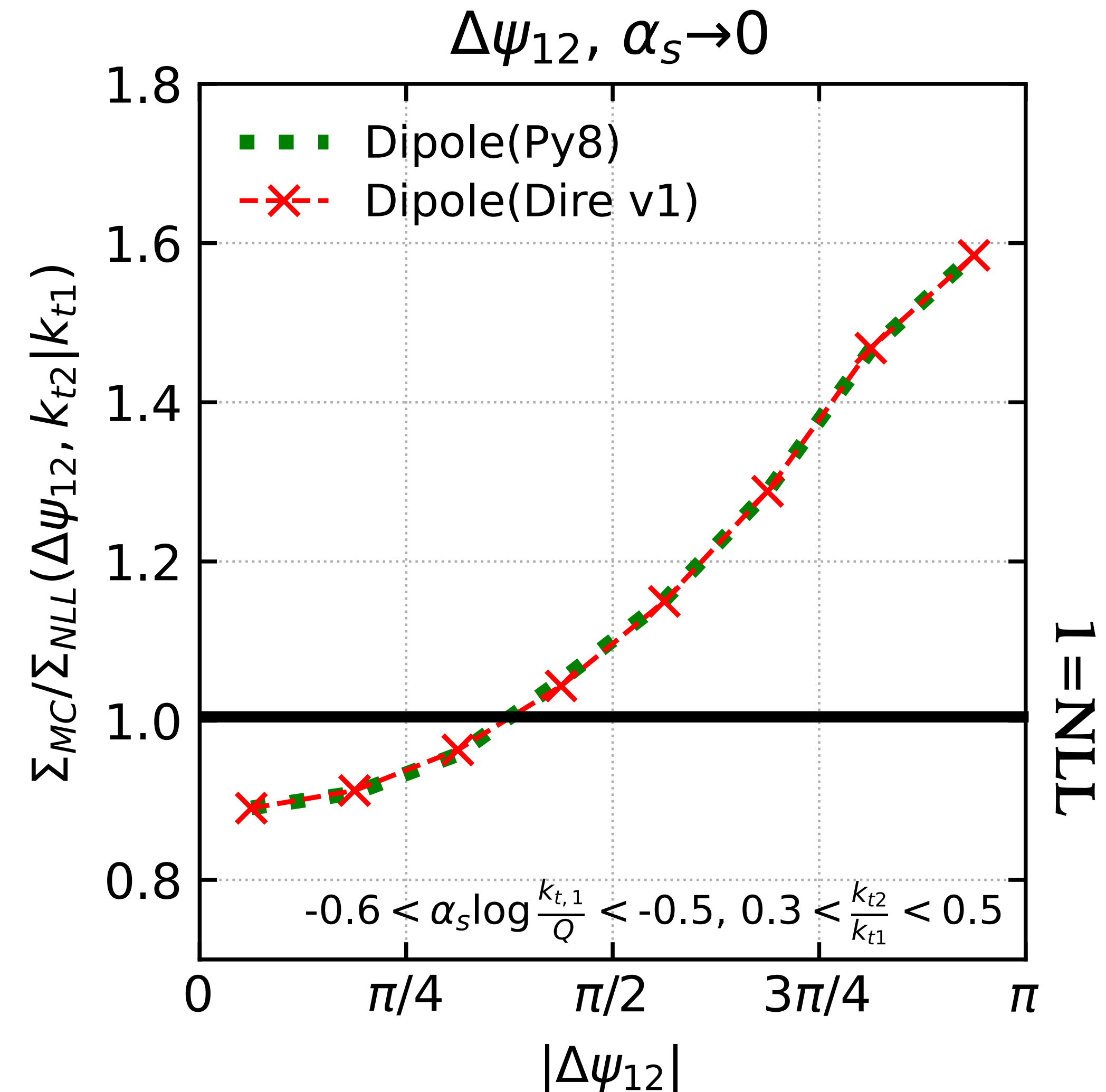
$$\Sigma(O < e^{-L}) = \exp \left(-L g_{\text{LL}}(\beta_0 \alpha_s L) + g_{\text{NLL}}(\beta_0 \alpha_s L) + \alpha_s g_{\text{NNLL}}(\beta_0 \alpha_s L) + \dots \right)$$

For $Q \sim 50 - 10000 \text{ GeV}$, $\beta_0 \alpha_s L \sim 0.3 - 0.5$:
Next-to-Leading Logarithms needed for %-level precision!

What is available in Shower Monte Carlo generators?

- Showers routinely used to interpret LHC (and LEP) data are **not NLL!**

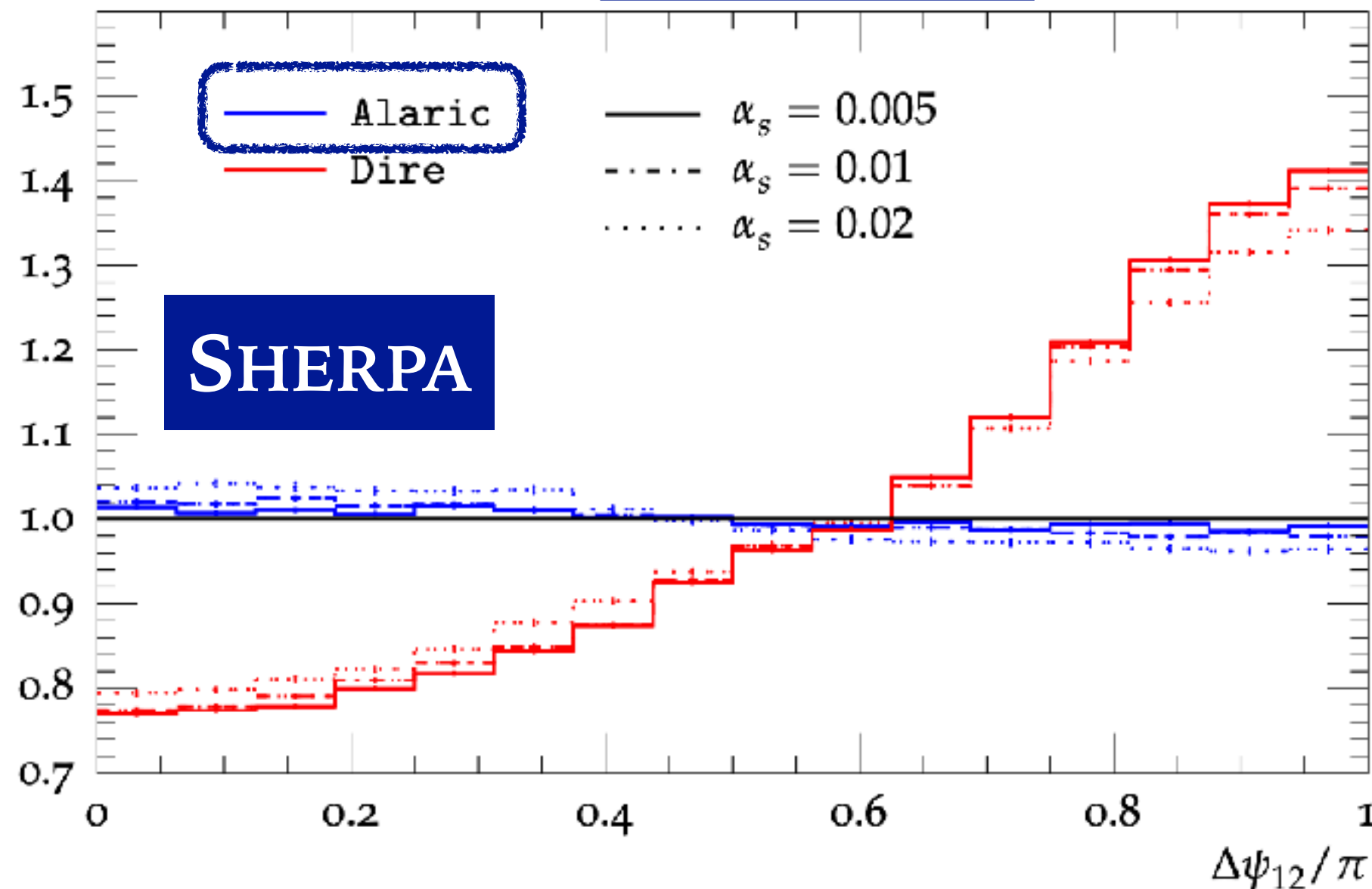
Dasgupta et al. [2002.11114](#)



What can be available in Shower Monte Carlo generators?

- Showers routinely used to interpret LHC (and LEP) data are **not NLL**!
- **Many groups** are independently formulating new showers with **NLL accuracy for e^+e^-**

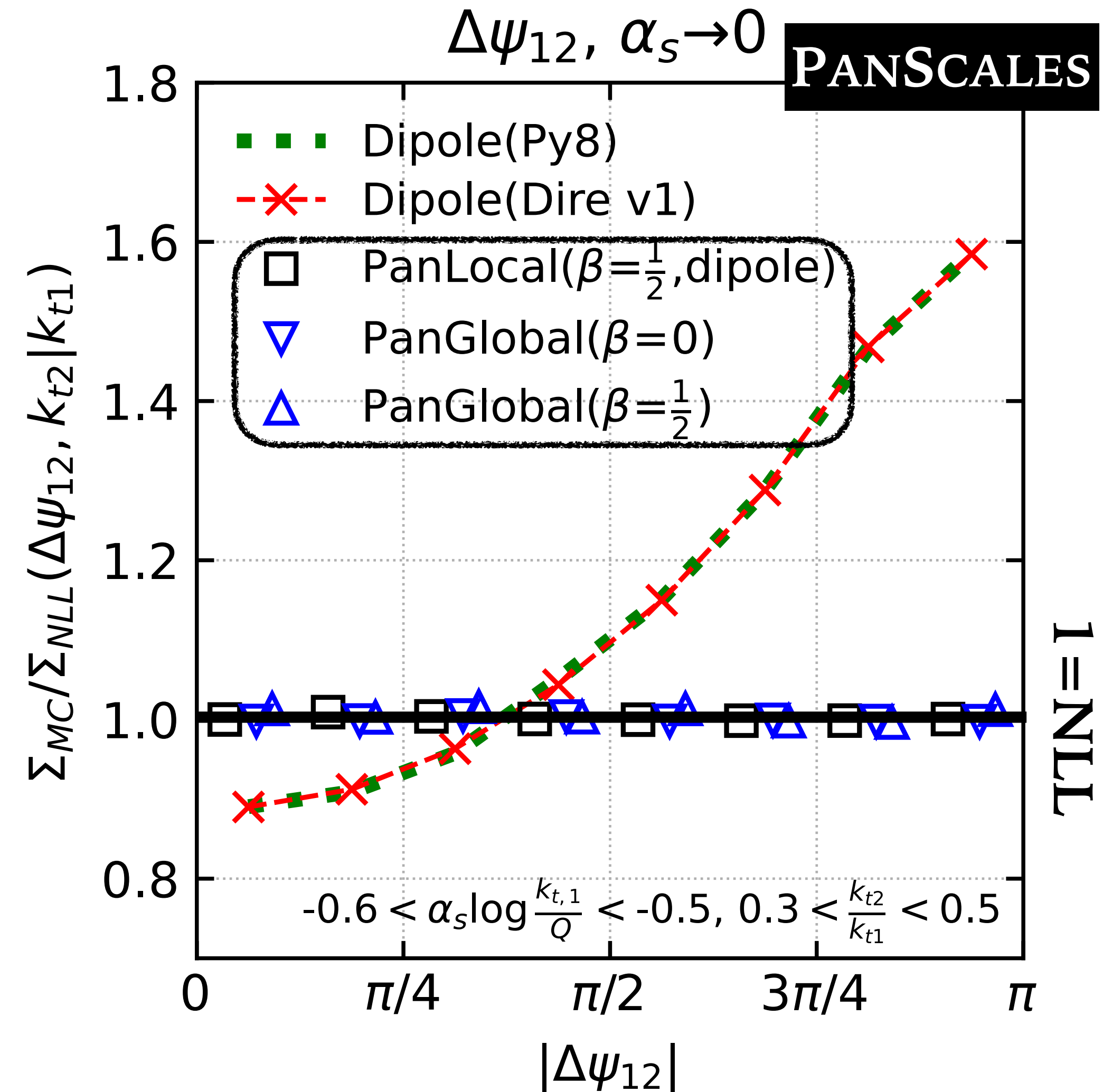
Herren et al. [2208.06057](#)



DEDUCTOR
Nagy&Soper,
[2011.04777](#)

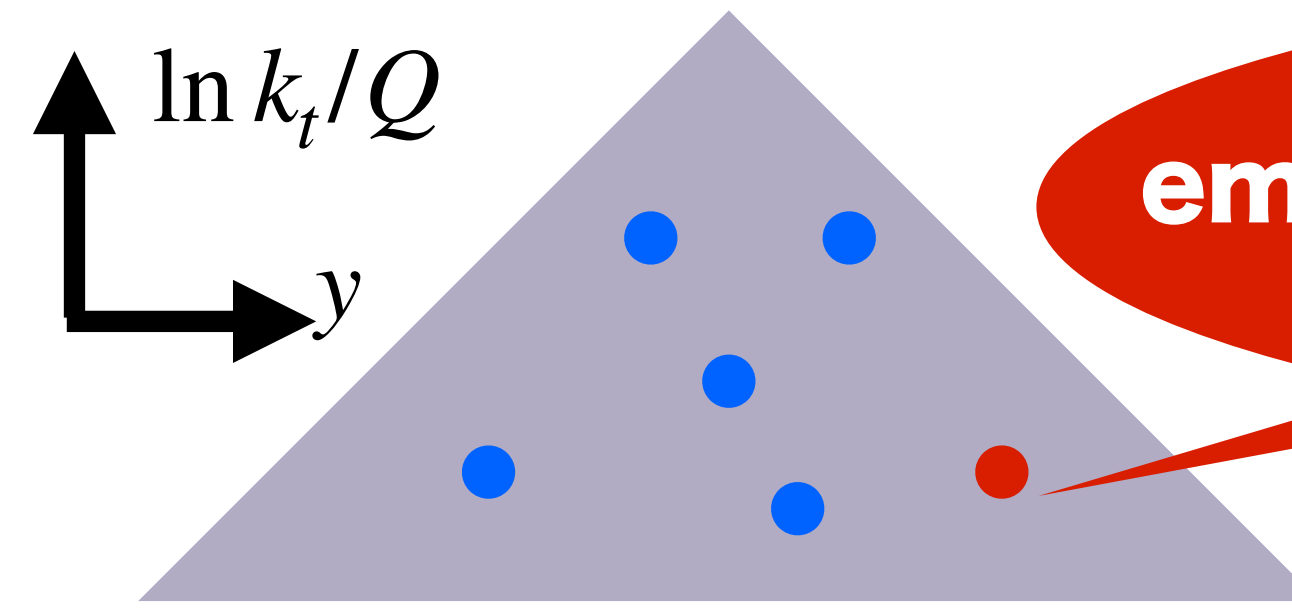
CVOLVER
Forshaw et. al,
[2003.06400](#)

Dasgupta et al. [2002.11114](#)



From LL to NLL: what's new?

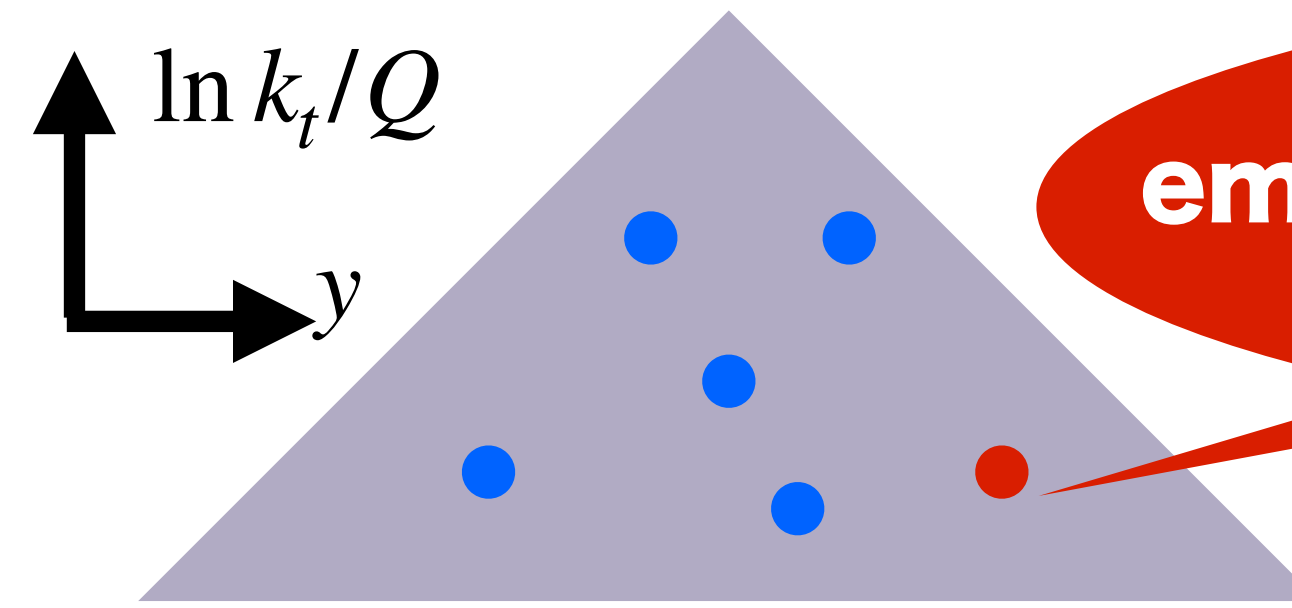
- NLL accuracy achieved for lepton collider processes ensuring parton showers reproduce the QCD matrix element in the presence of **soft gluons separated in angle**



When doing a new emission, previously emitted gluons do not change!

From LL to NLL: what's new?

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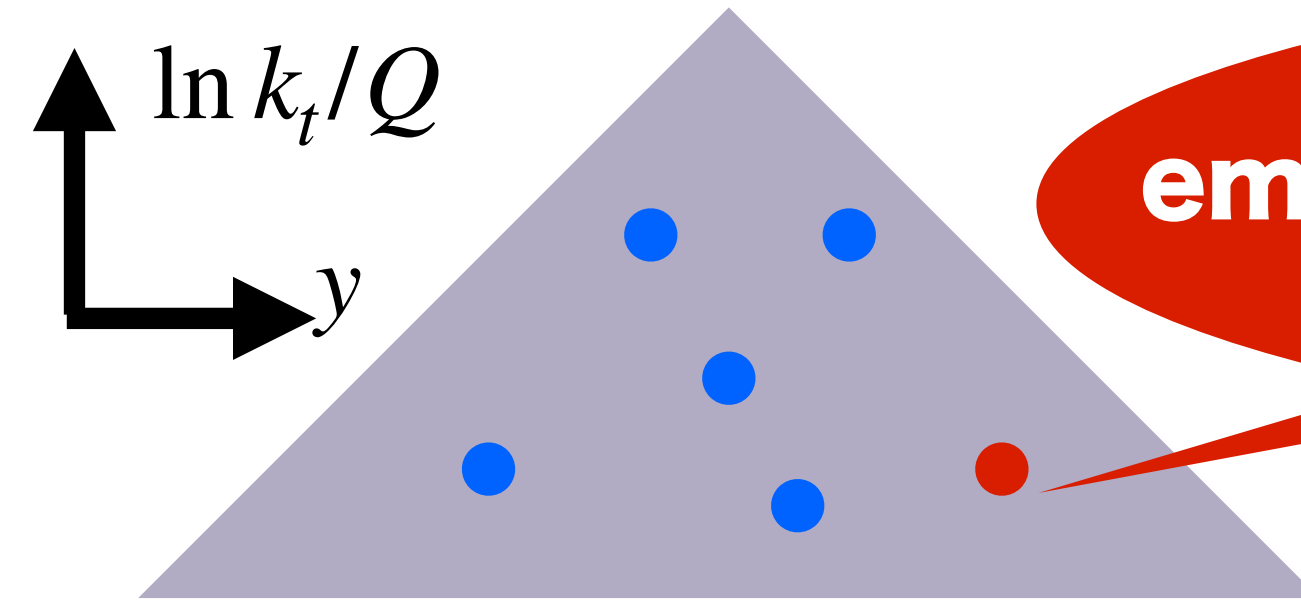


When doing a new emission, previously emitted gluons do not change!

- Lepton - hadron colliders (and **VBF**): **PANSCALES**, van Beekveld, SFR [2305.08645](#)
- Hadron - hadron colliders:
colour-singlet in **PANSCALES** (van Beekveld, SFR et al., [2205.02237](#) + [2207.09467](#)), and **DEDUCTOR** (Nagy&Soper, [0912.4534](#))
generic processes: ongoing efforts in **ALARIC** and **PANSCALES**

From LL to NLL: what's new?

- NLL accuracy achieved for **generic collider** processes ensuring parton showers reproduce the QCD matrix element in the presence of **soft gluons separated in angle**

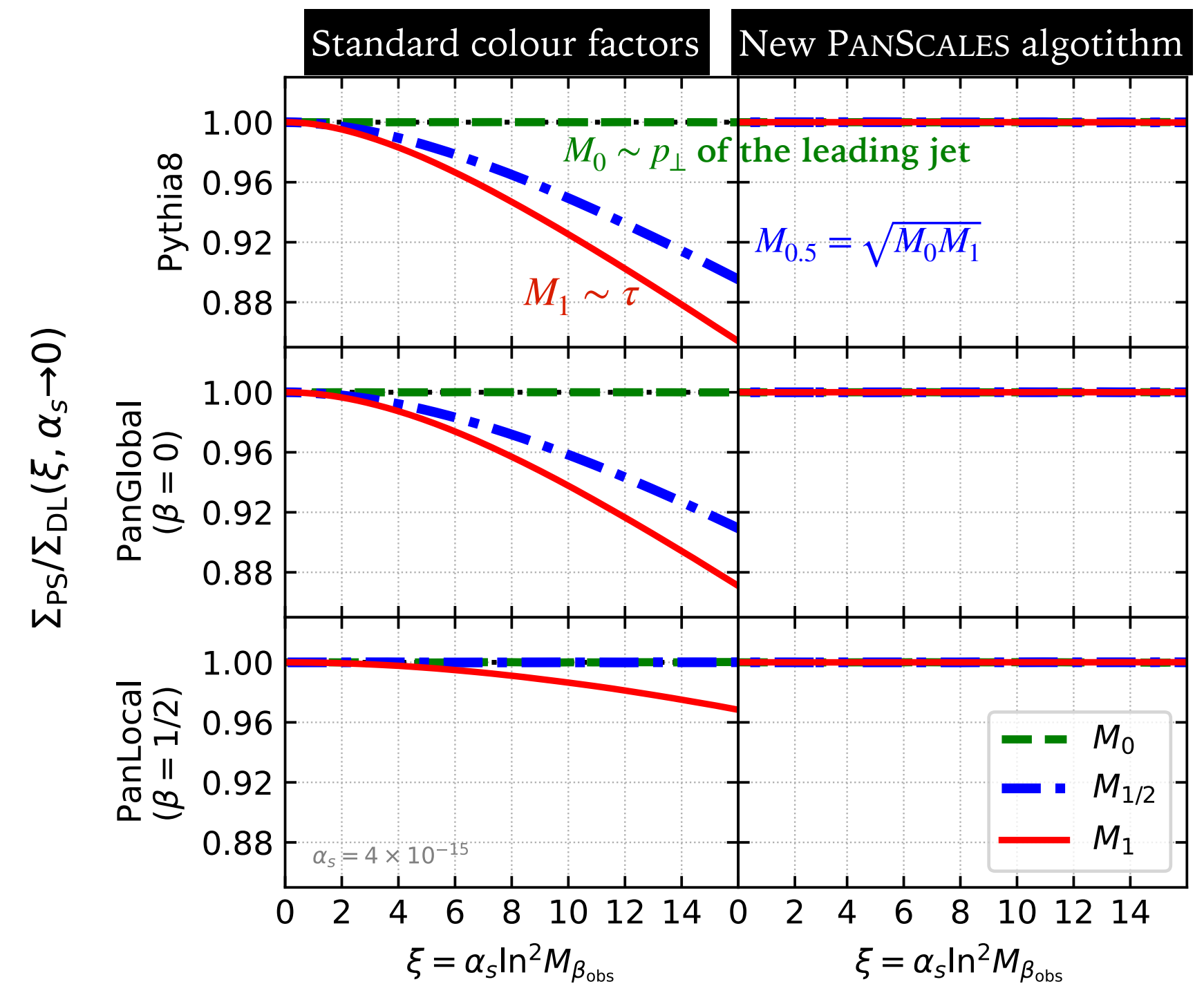


When doing a new emission, previously emitted gluons do not change!

- **Subleading-colour corrections** in **PANSCALES**, as parton showers are derived in the large number of colour limit

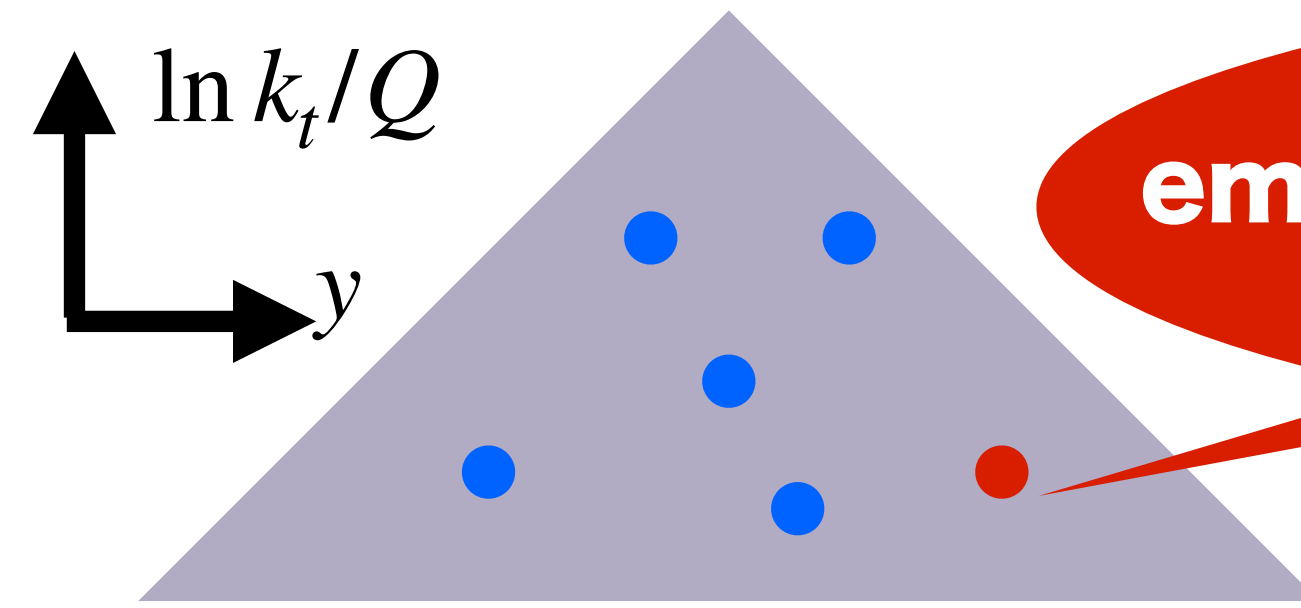
$$C_F = \frac{C_A}{2} - \frac{1}{2N_c}$$

[K. Hamilton et al., [2011.10054](#)]



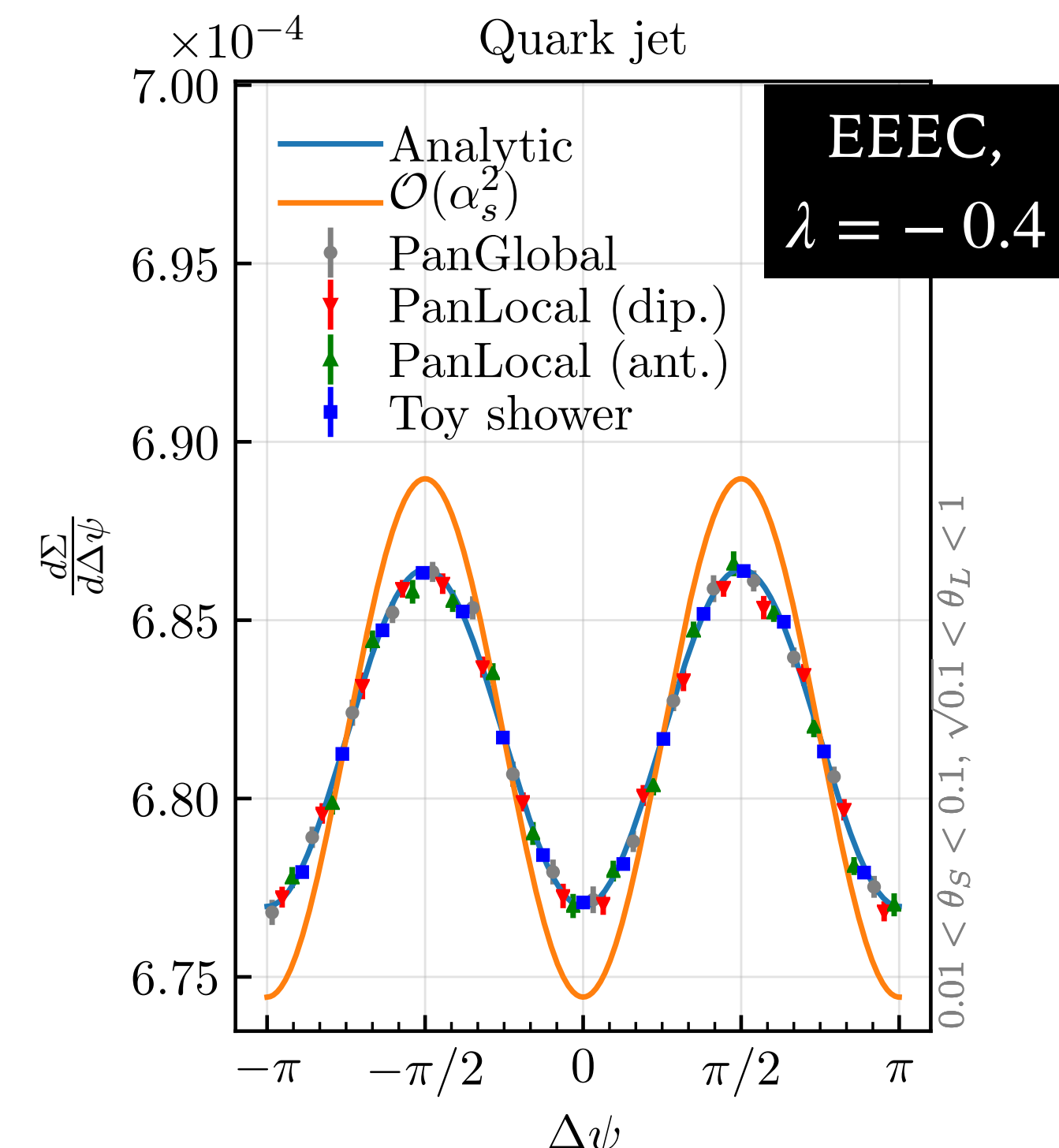
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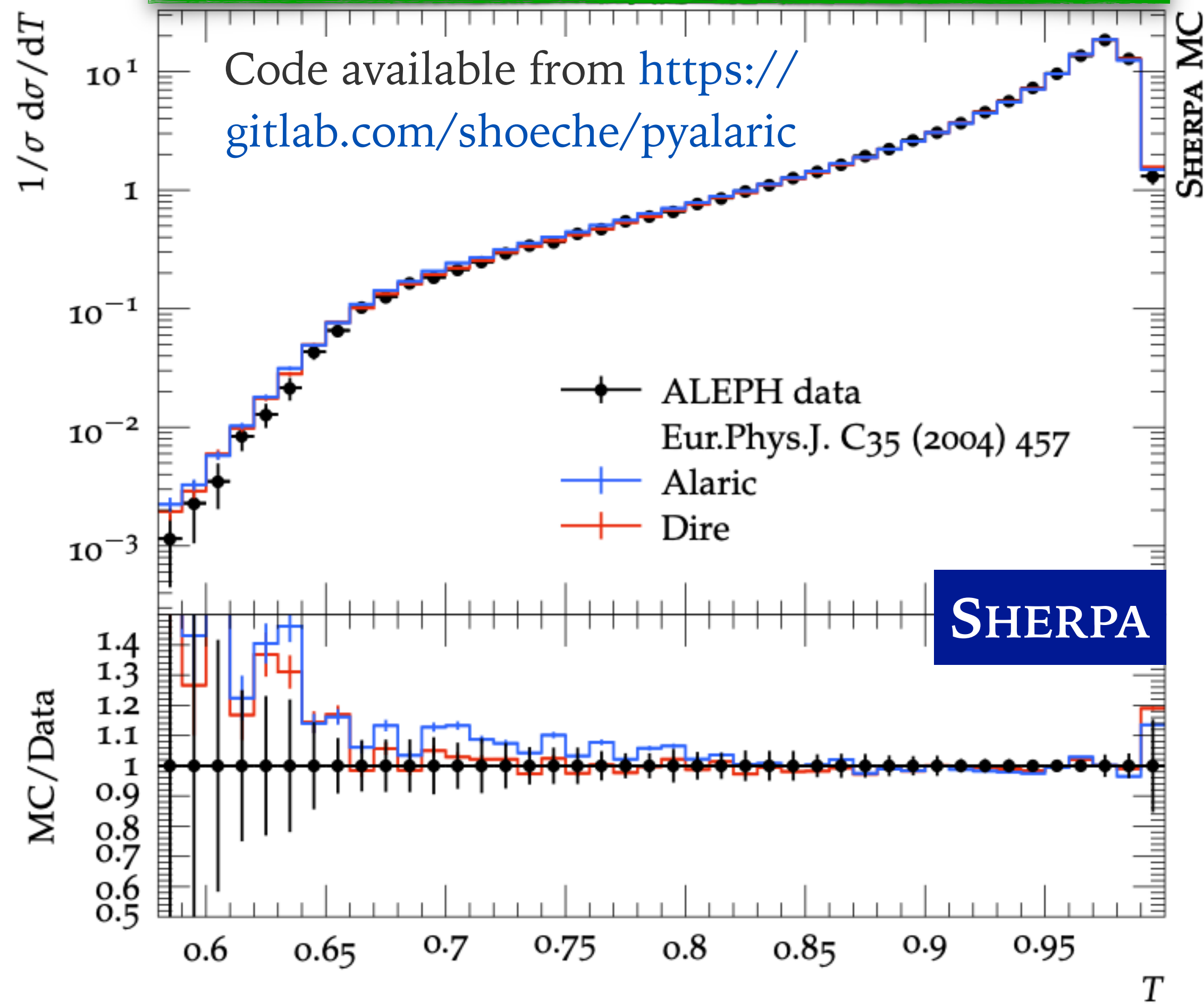
When doing a new emission, previously emitted gluons do not change!

- **Subleading-colour corrections** in **PANSCALES**, [K. Hamilton et al., [2011.10054](#)]
- **Spin correlations** in **PANSCALES**, [A. Karlberg et al., [2103.16526](#) + [2111.01161](#)], based on the **Collin-Knowles** ('88) algorithm (also available in **HERWIG7**, Richardson & Webster [1807.01955](#))

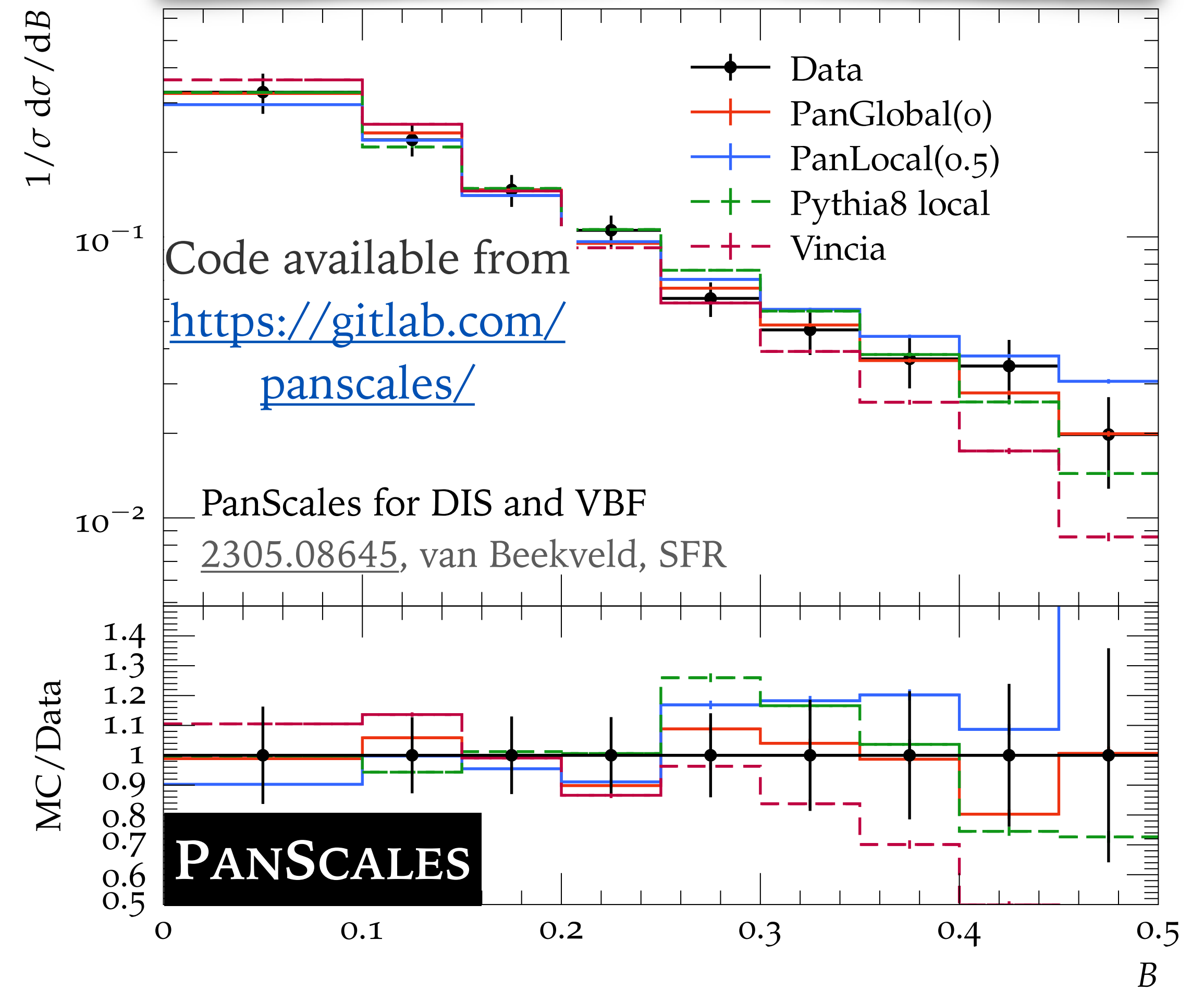


Comparison with current data

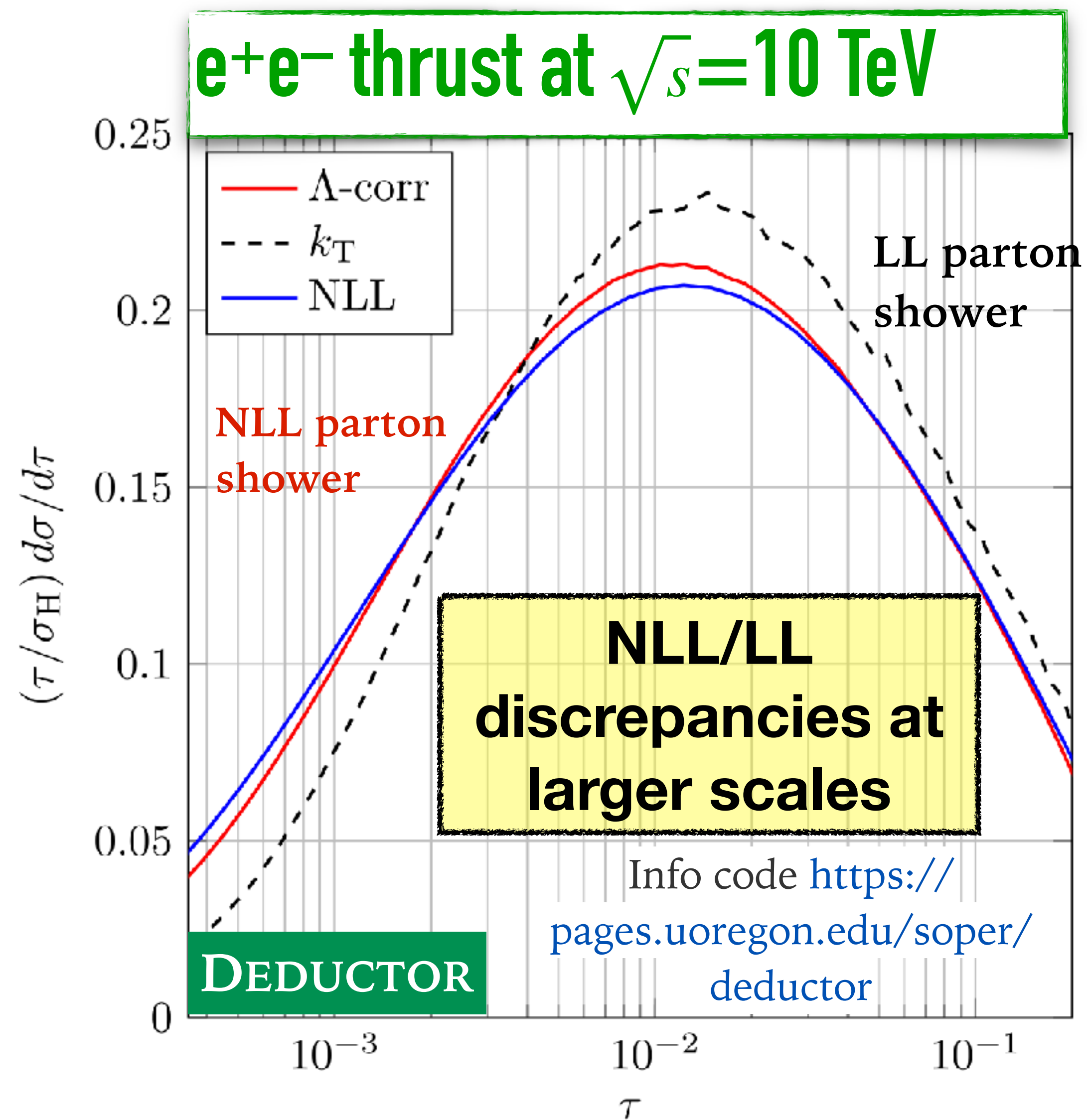
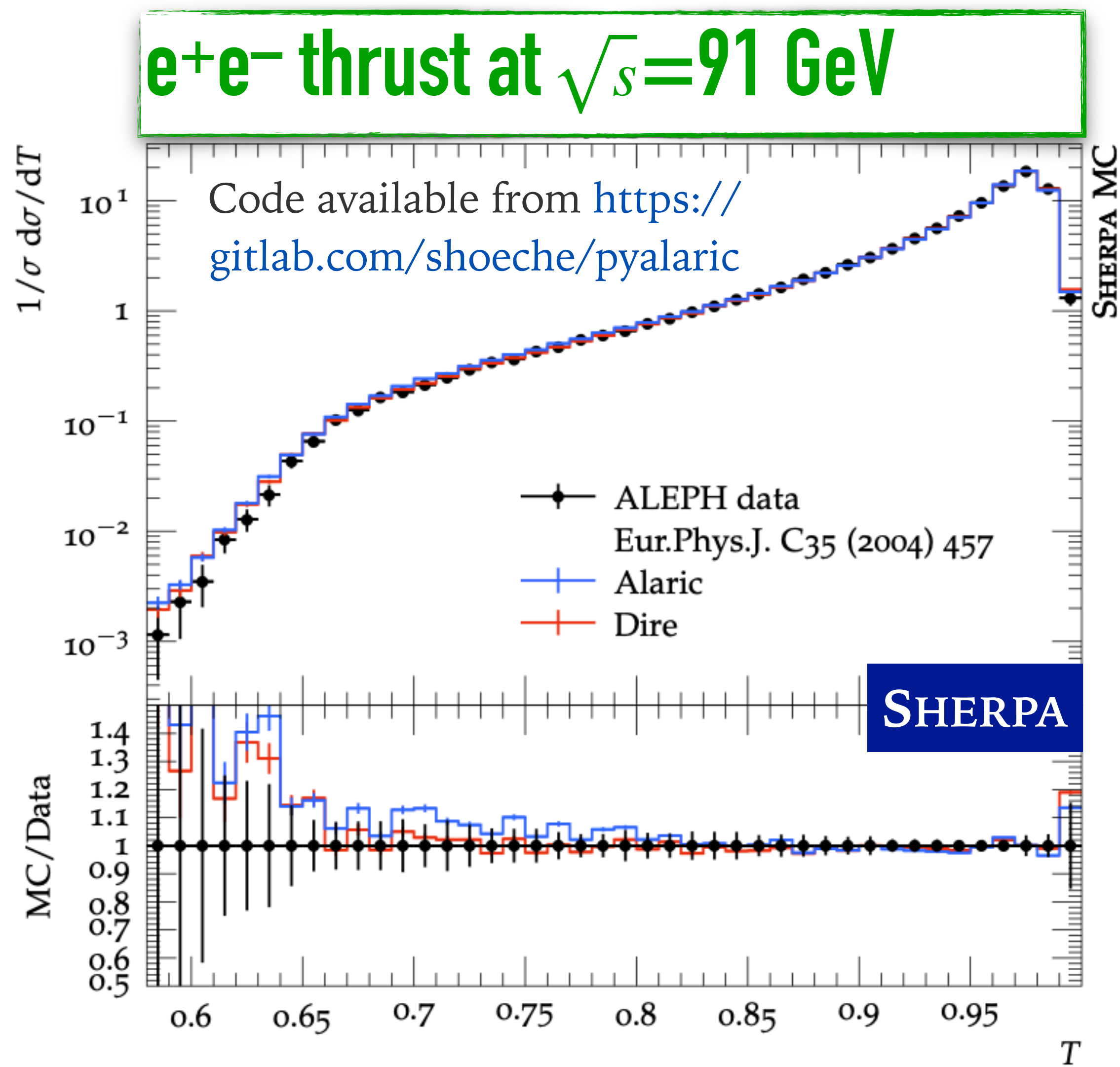
e^+e^- thrust at $\sqrt{s}=91$ GeV



DIS broadening at $Q=58$ GeV

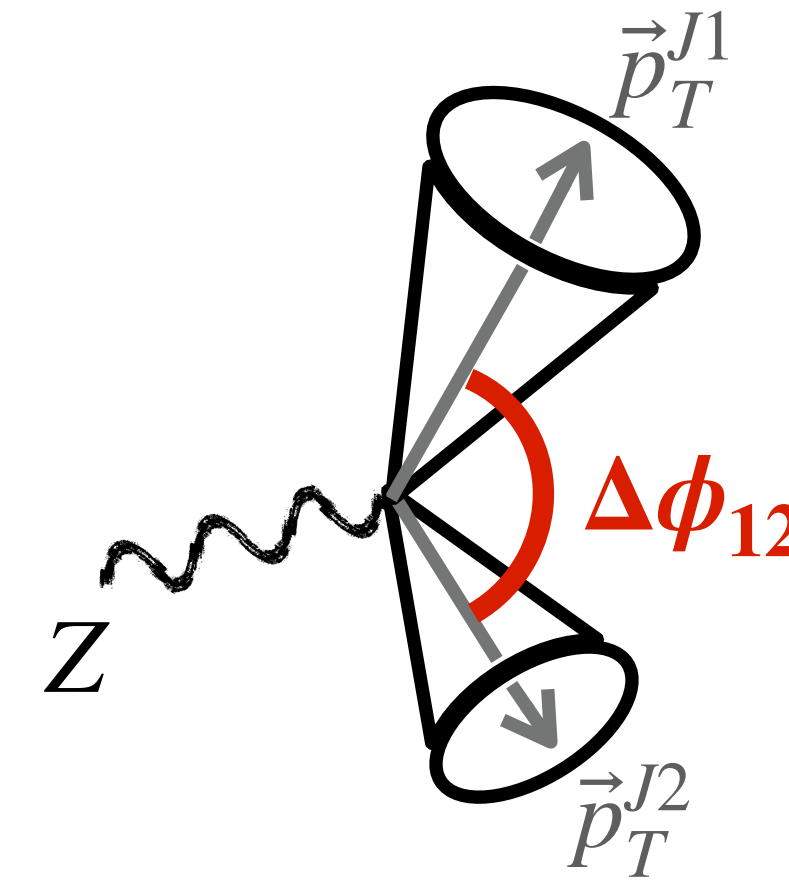
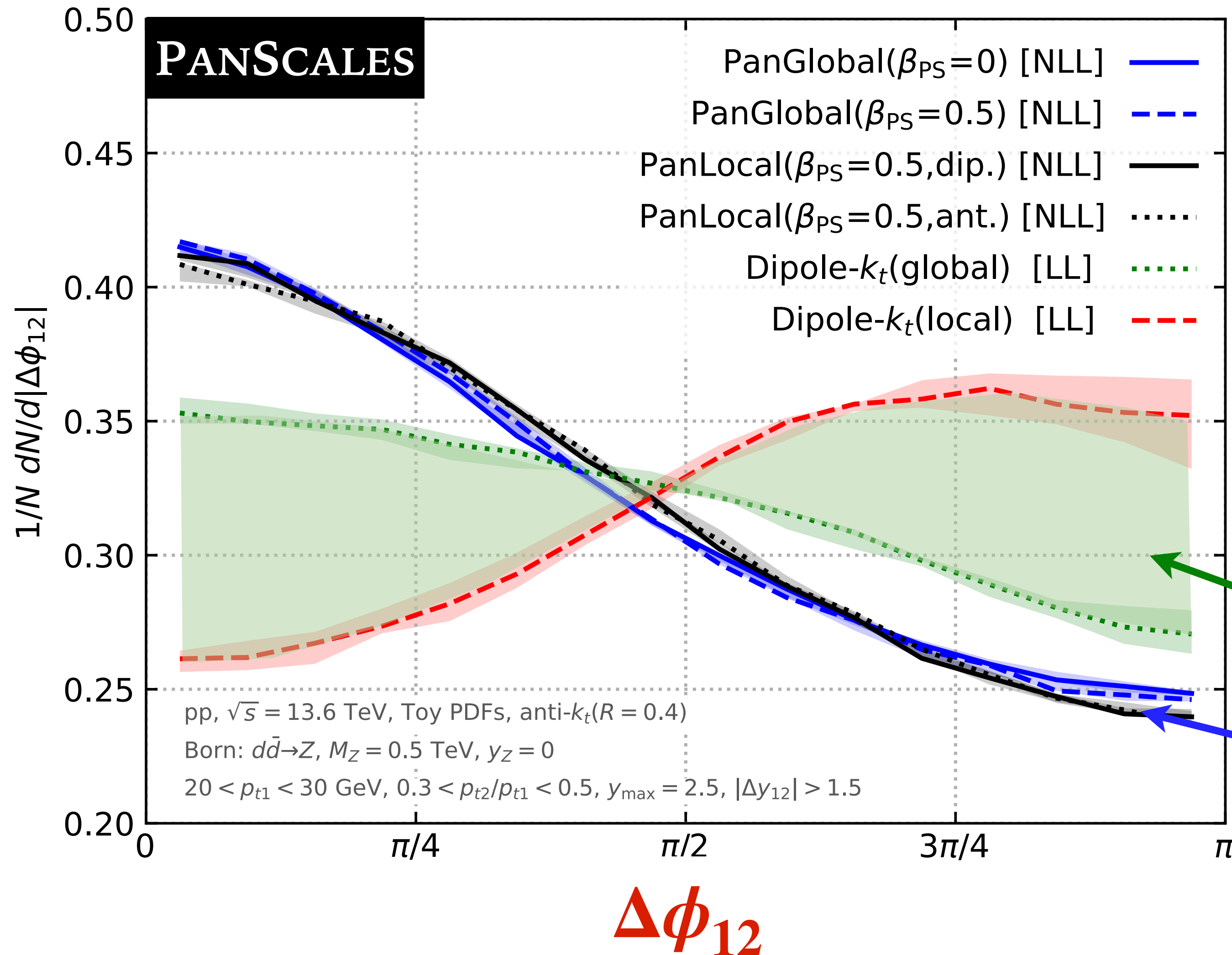


A closer look to the thrust in e^+e^- collisions



$$m_{\ell\ell} = 500 \text{ GeV}$$

Azimuthal angle between leading jets (DY)



PanScales for $pp \rightarrow$
 colour singlet:
[2207.09467](https://arxiv.org/abs/2207.09467), van
 Beekveld, SFR,
 Hamilton, Salam
 Soto Ontoso, Soyez,
 Verheyen:

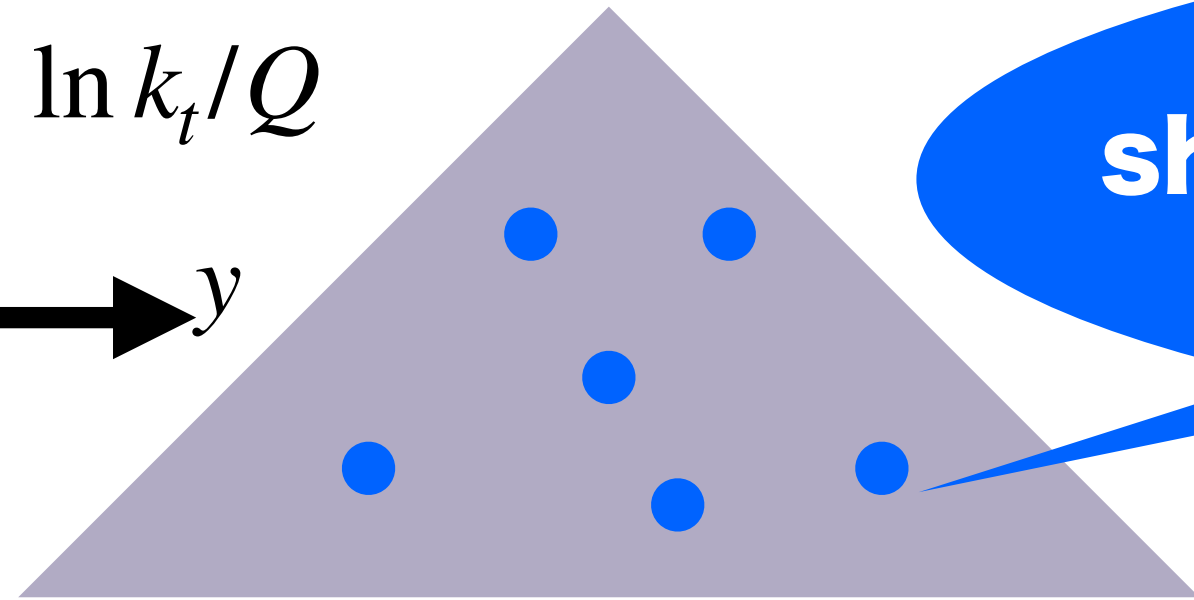
NLL/LL discrepancies at larger scales

LL showers

NLL showers

NNLL precision must be reached for percent-level precision!

$\ln k_t/Q$
 y

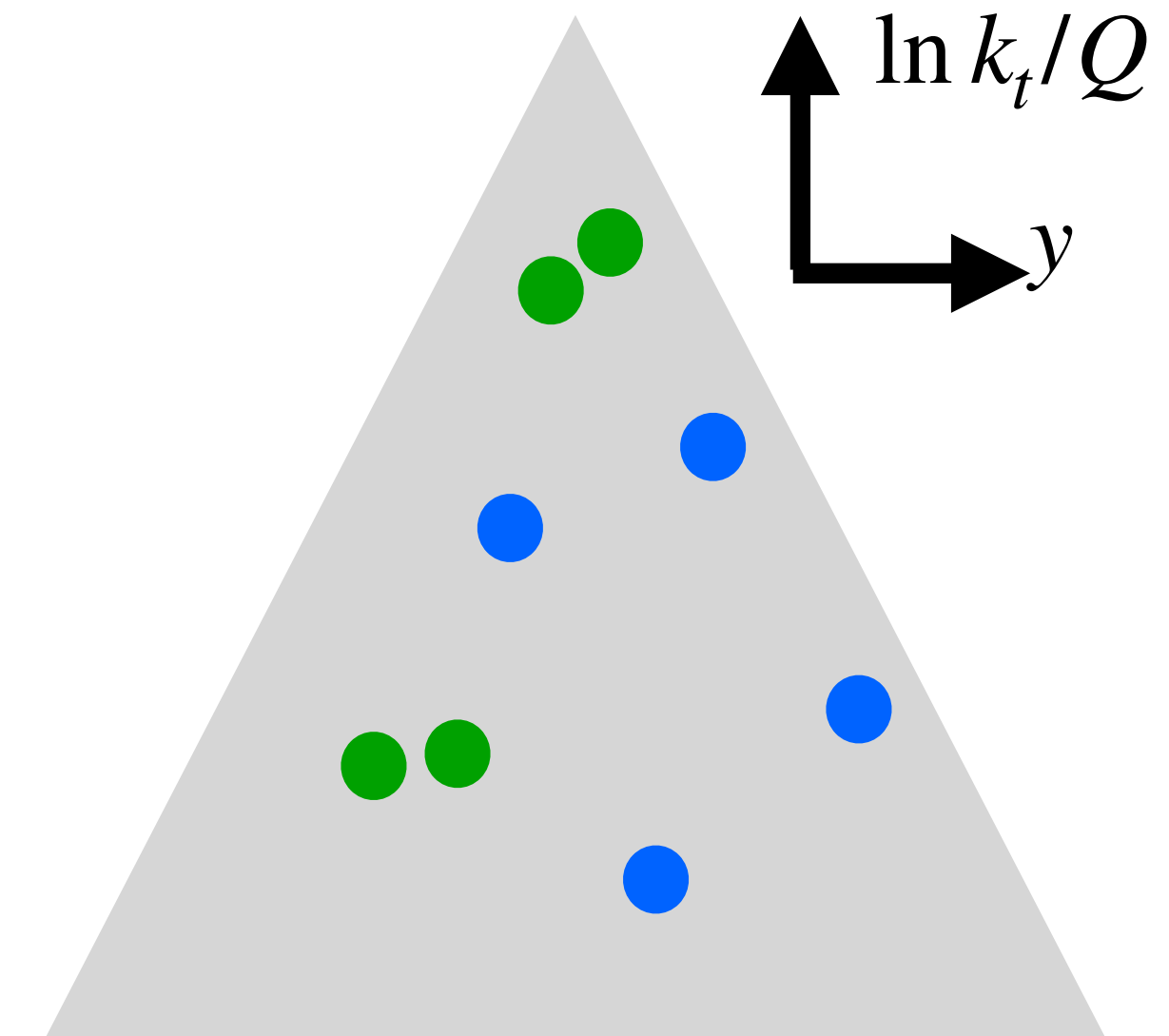


At NLL, the parton shower can handle emissions widely separated in angle

NNLL precision must be reached for percent-level precision!

Soft emission — inclusion of **real** + **virtual** corrections

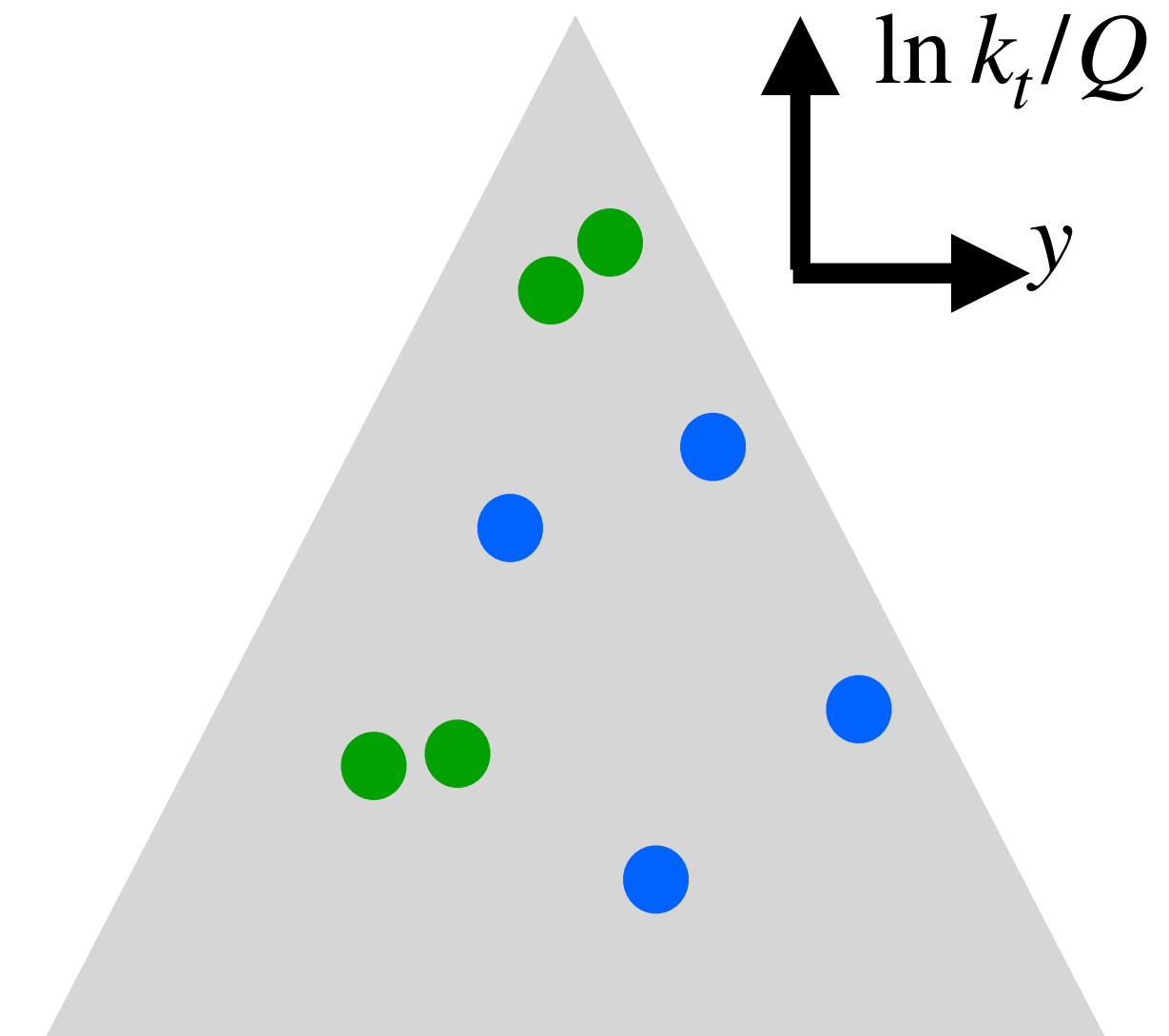
- any **pair of soft emissions** with commensurate energy and angles should be produced with the correct [double-soft] matrix element
- probability for any **single soft emission** should be NLO accurate
- NB: Vincia and Sherpa groups have also explored inclusion of the double-soft current; part of novelty here is doing so to get the log-accuracy benefit.



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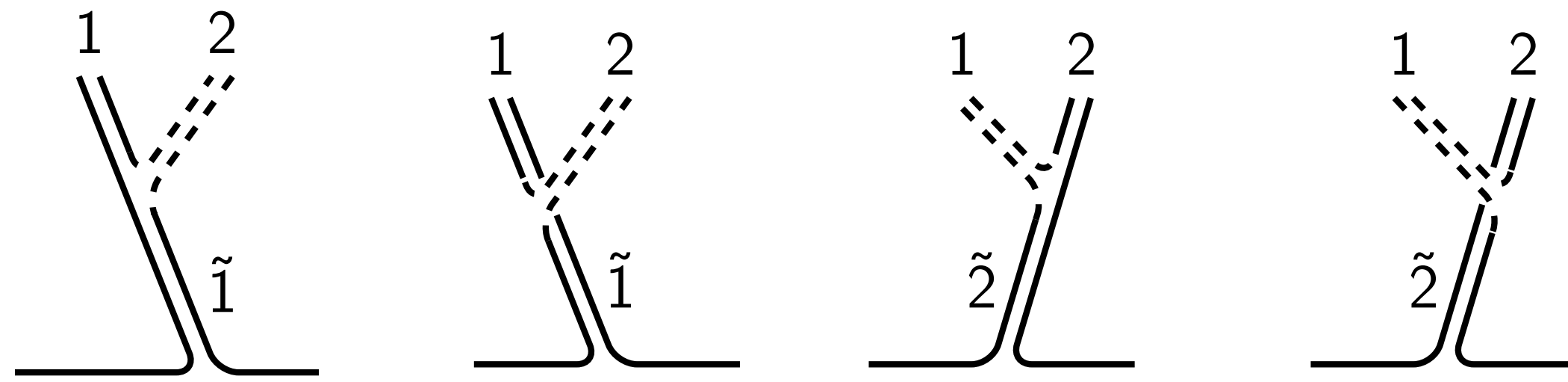


This (+NLO matching, see Karlberg's talk) should maintain NLL accuracy and further achieve

- **NNDL accuracy** for [subset] multiplicities, i.e. terms $\alpha_s^n L^{2n}$, $\alpha_s^n L^{2n-1}$, $\alpha_s^n L^{2n-2}$
- **Next-to-Single-Log (NSL) accuracy** for non-global logarithms, e.g. energy in a slice, all terms $\alpha_s^n L^n$ and $\alpha_s^n L^{n-1}$ (at leading- N_c)

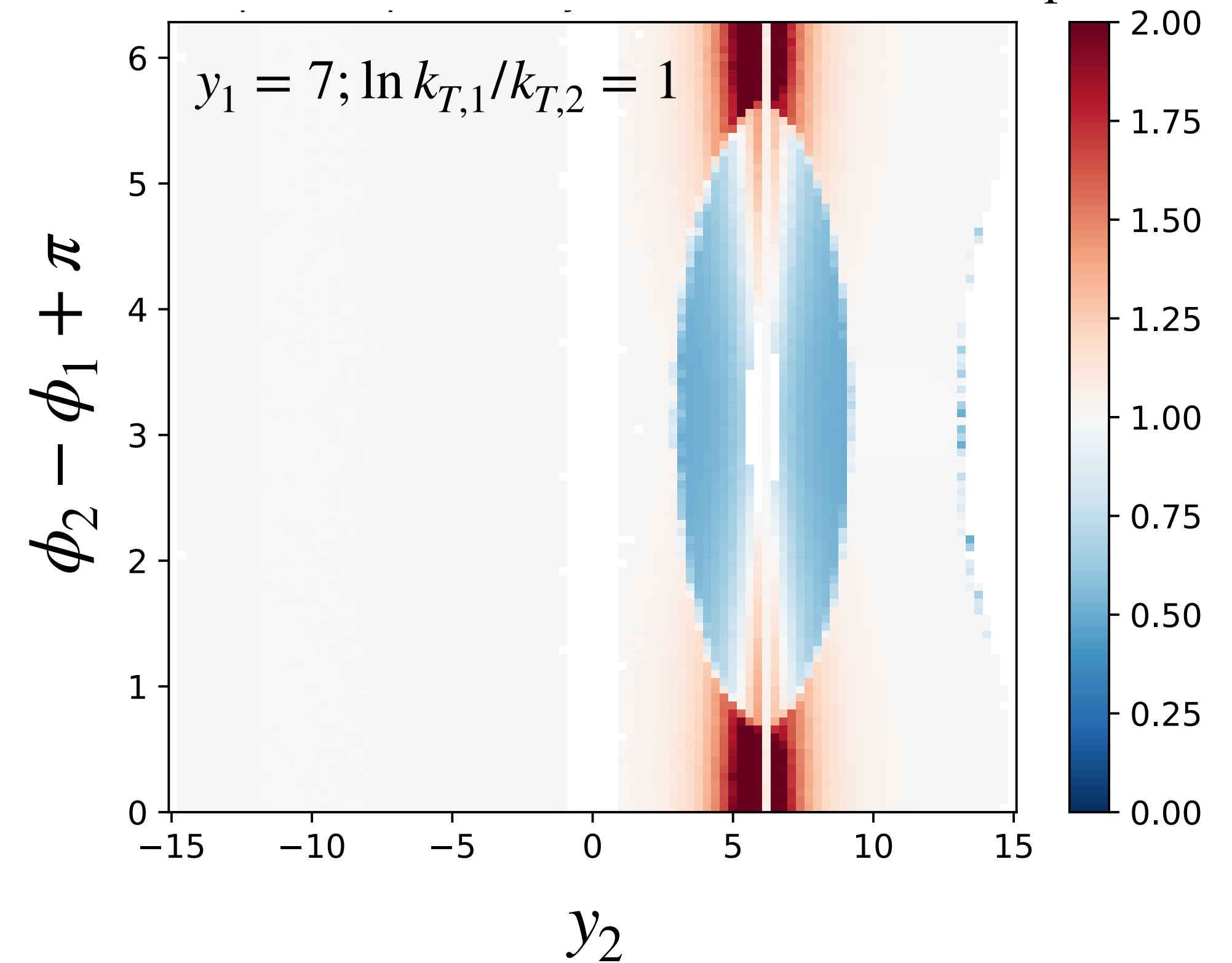
NB: done using PanGlobal, so far just in $e^+e^- \rightarrow q\bar{q}$

1. Real corrections: pair of soft emissions



- a given two-emission configuration can come from several shower histories
- **accept a given emission with exact double-soft $M_{\text{exact}}^{(\text{DS})}$ divided by shower's effective double-soft matrix element summed over the histories h that could have produced that configuration**

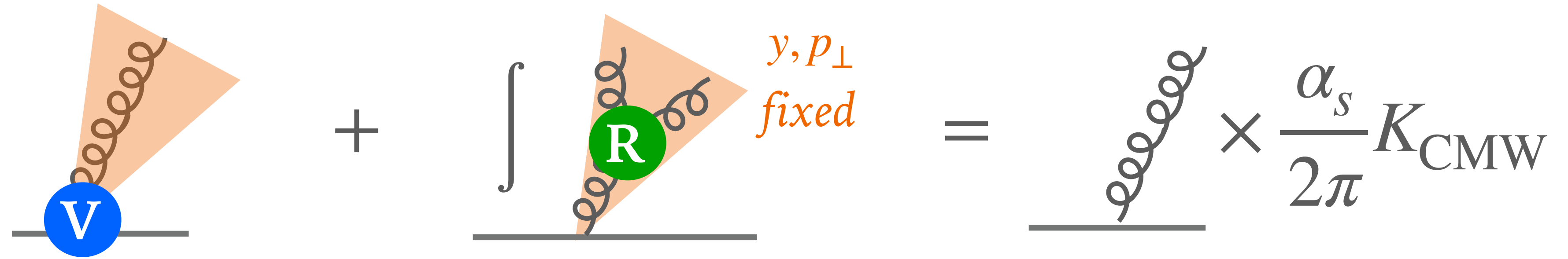
Double-soft acceptance P_{accept}



$$P_{\text{accept}} = \frac{M_{\text{exact}}^{(\text{DS})}}{\sum_h M_{h,\text{PS}}^{(\text{DS})}}$$

Virtual corrections in parton showers

► For a soft emission



The diagrammatic equation shows the sum of a virtual correction (V) and a real emission (R) integral, equated to a soft emission factor. The first term is a blue circle with 'V' on a horizontal line, with a wavy gluon line and an orange cone. The second term is a green circle with 'R' on a horizontal line, with a wavy gluon line, an orange cone, and an integral sign. The third term is a wavy gluon line on a horizontal line, multiplied by the factor $\frac{\alpha_s}{2\pi} K_{\text{CMW}}$. The text y, p_{\perp} fixed is written in orange next to the integral.

$$\text{V} + \int \text{R} \quad (y, p_{\perp} \text{ fixed}) = \text{Soft Emission} \times \frac{\alpha_s}{2\pi} K_{\text{CMW}}$$

Virtual corrections in parton showers

► For a soft emission

$$\text{V} + \int \text{R} \quad y, p_{\perp} \text{ fixed} = \text{wavy line} \times \frac{\alpha_s}{2\pi} K_{\text{CMW}}$$

► Catani, Marchesini and Webber defined the “CMW” scheme for the coupling in the shower [Nucl.Phys.B 349 (1991) 635-654]

$$\alpha_s^{\text{CMW}} = \alpha_s \left(1 + \frac{\alpha_s}{2\pi} K_{\text{CMW}} \right)$$

Virtual corrections in parton showers

► For a soft emission

$$\text{V} + \int_{\text{R}} \text{R} \quad (y, p_{\perp} \text{ fixed}) = \text{V} \times \frac{\alpha_s}{2\pi} K_{\text{CMW}}$$

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 [Nucl.Phys.B 349 (1991) 635-654]

$$\alpha_s^{\text{CMW}} = \alpha_s \left(1 + \frac{\alpha_s}{2\pi} K_{\text{CMW}} \right)$$

*At fixed “shower variables”,
 but the rapidity and p_{\perp} of the
 jet can vary*

$$\text{V}_{\text{PS}} \equiv \text{V} \times \frac{\alpha_s}{2\pi} K_{\text{CMW}} - \int_{\text{R}_{\text{PS}}} \text{R}_{\text{PS}}$$

This ensures

$$\text{V}_{\text{PS}} + \int \text{R}_{\text{PS}} = \frac{\alpha_s}{2\pi} K_{\text{CMW}}$$

“on average”

2. Improving virtual corrections for soft emissions

With our double soft acceptance we have $\mathbf{R}_{PS} = \mathbf{R}$.

To ensure

$$\begin{aligned}
 & \text{Diagram: } \mathbf{V}_{PS} \text{ (blue circle) with a gluon emission cone (orange triangle) } \\
 & = \frac{\alpha_s}{2\pi} K_{CMW} - \int \text{Diagram: } \mathbf{R} \text{ (green circle) with a gluon emission cone (orange triangle) } \\
 & \text{where } y, p_{\perp} \text{ fixed}
 \end{aligned}$$

We modify the CMW scheme

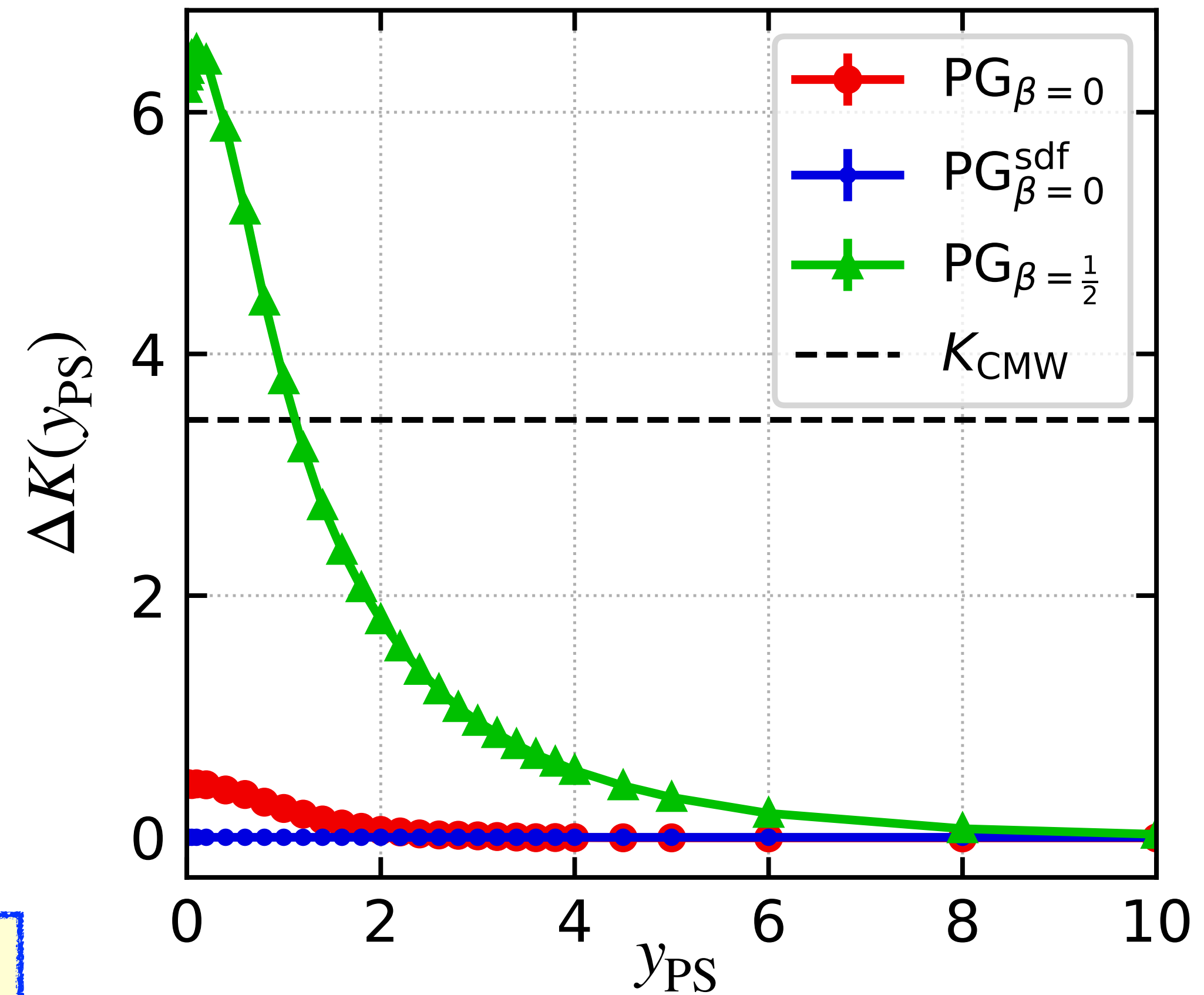
$$K_{CMW} \rightarrow K_{CMW} + \Delta K(\Phi_{PS}^{(1)})$$

fixed "shower variables"

$$\frac{\alpha_s}{2\pi} \Delta K(\Phi_{PS}^{(1)}) = \int \text{Diagram: } \mathbf{R} \text{ (green circle) with a gluon emission cone (pink triangle)} - \int \text{Diagram: } \mathbf{R} \text{ (green circle) with a gluon emission cone (orange triangle)}$$

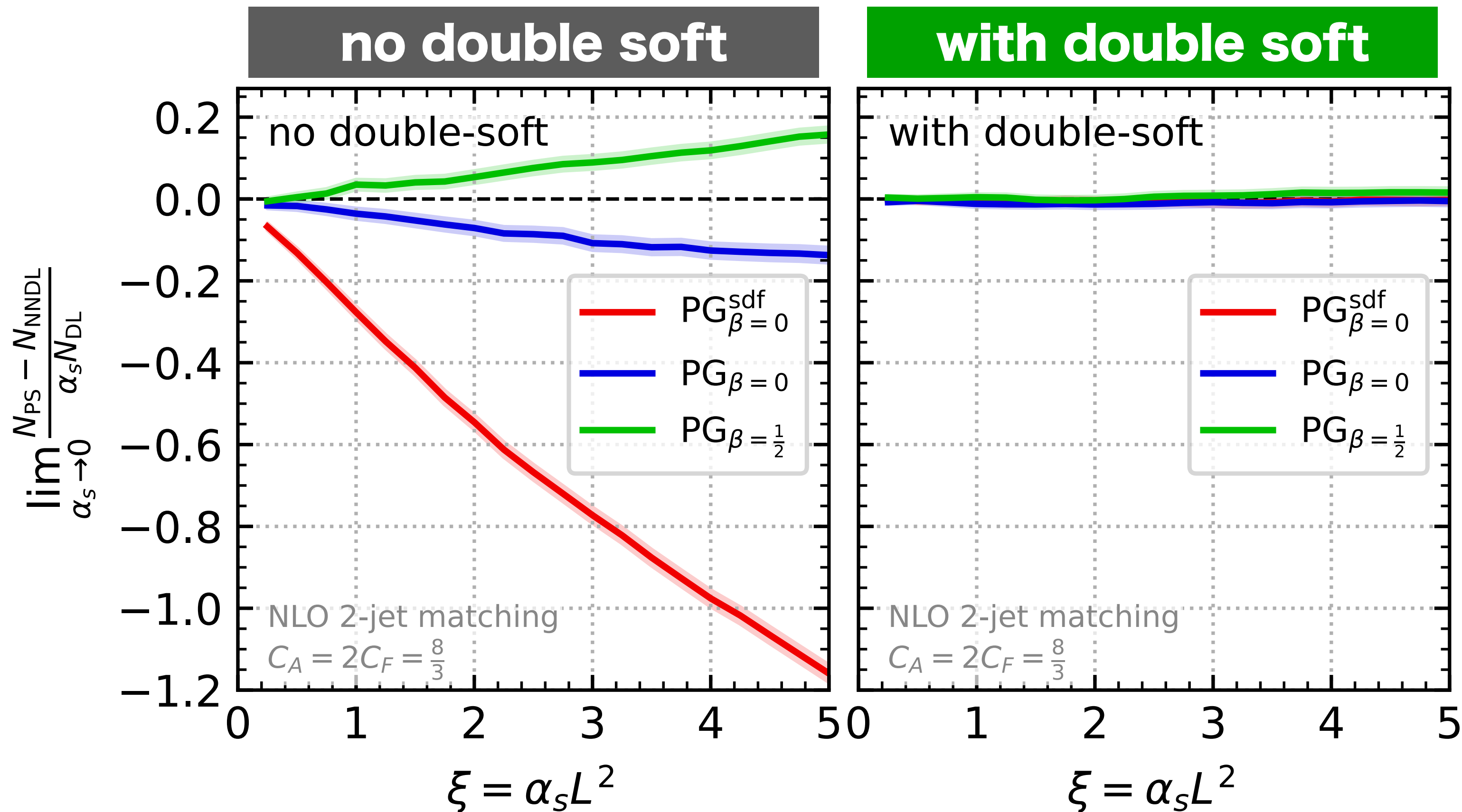
y, p_⊥ fixed

example ΔK correction



ΔK vanishes for large rapidities since virtual corrections to soft-collinear emissions are OK for NLL showers

NNDL subset multiplicity

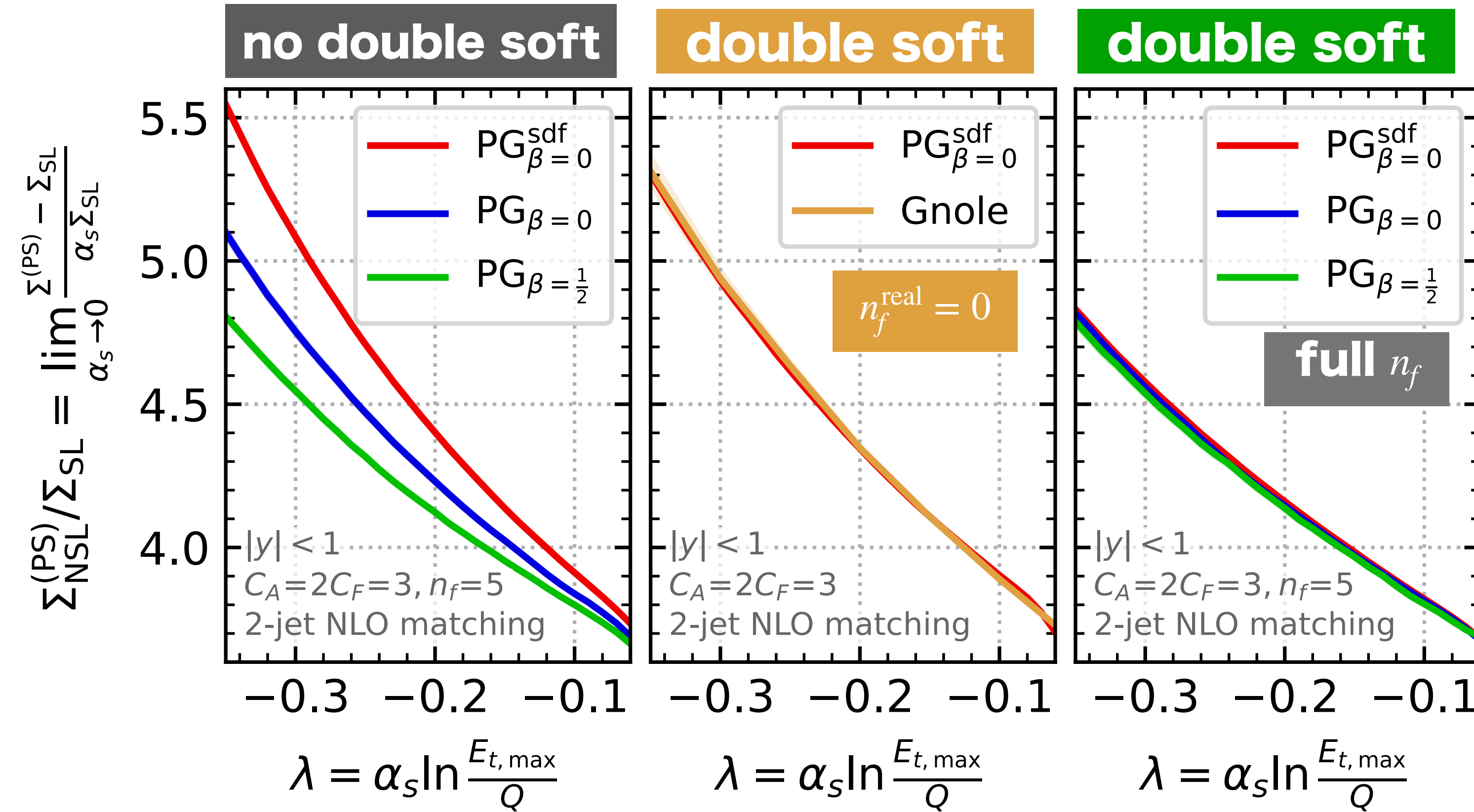


- **NNDL** ($\alpha_s^n L^{2n-2}$) analytic resummation = Medves, Soto Ontoso, Soyez, [2205.02861](#)
- $\alpha_s \rightarrow 0$ limit to isolate NNDL terms.
- **Double soft necessary for NNDL agreement**

$$\lim_{\alpha_s \rightarrow 0} \frac{N_{PS} - N_{NNDL}}{\alpha_s N_{DL}} \Big|_{\text{fixed } \alpha_s L^2}$$

S.F.R., Hamilton, Karlberg,
Salam, Scyboz, Soyez
[2307.11142](#)

NSL for the energy flow in a rapidity slice



$$\Sigma_{\text{NSL}}^{(\text{PS})} = \lim_{\alpha_s \rightarrow 0} \left. \frac{\Sigma^{(\text{PS})} - \Sigma_{\text{SL}}}{\alpha_s} \right|_{\text{fixed } \alpha_s L}, \quad L \equiv \ln \frac{E_{t,\text{max}}}{Q}$$

- **NSL** ($\alpha_s^n L^{n-1}$) = Banfi, Dreyer, Monni, [2104.06416](#), [2111.02413](#) (“Gnole”) [NB: see also Becher, Schalch, Xu, [2307.02283](#)]
- **NSL agreement with Gnole for $n_f^{\text{real}} = 0$**
- First large- N_c **full- n_f** results for NSL non-global logs

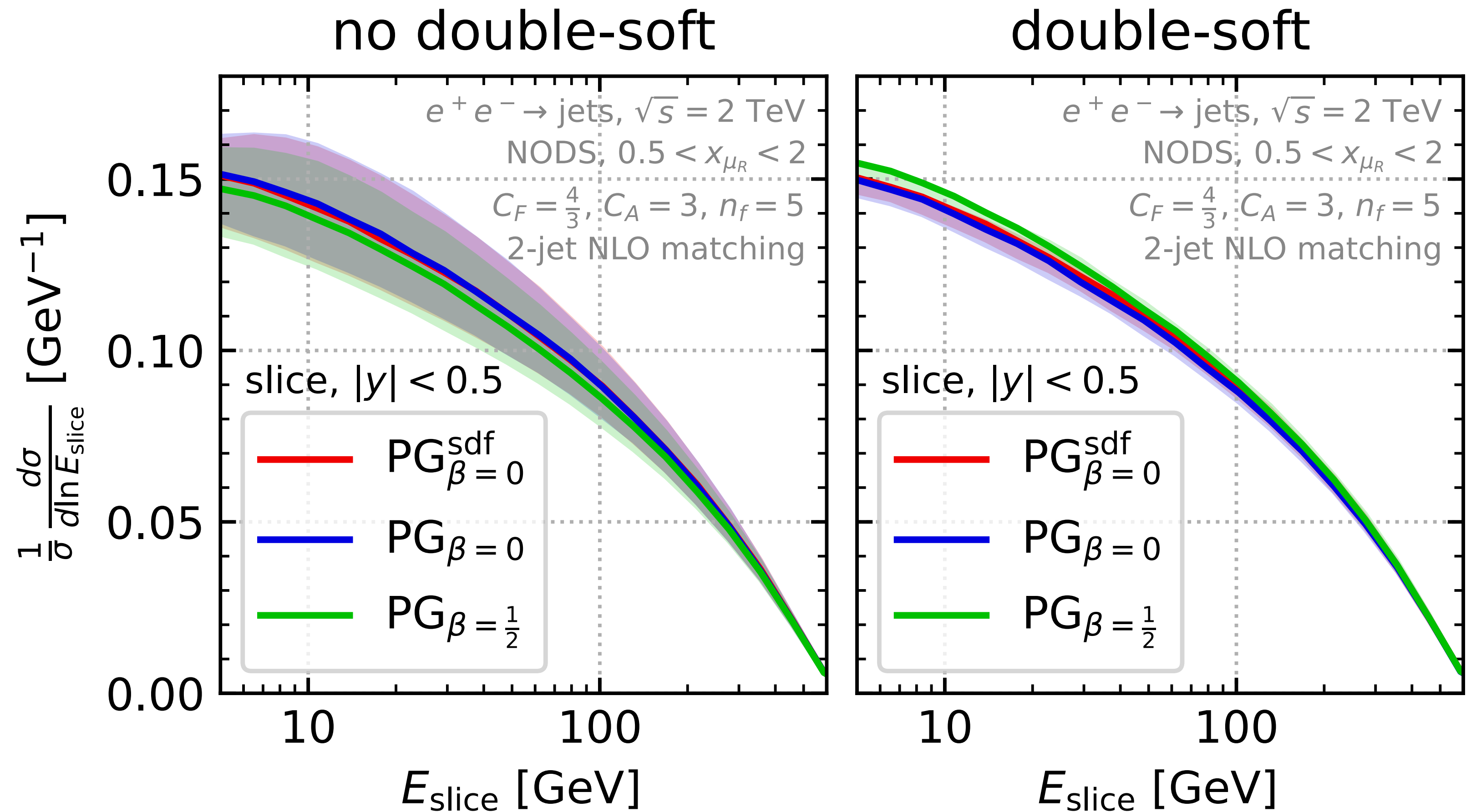
**S.F.R., Hamilton, Karlberg,
Salam, Scyboz, Soyez
[2307.11142](#)**

NSL phenomenology outlook

S.F.R., Hamilton, Karlberg,
Salam, Scyboz, Soyez
[2307.11142](#)

- Energy flow in slice between two 1 TeV jets
- **Double-soft reduces uncertainty band**

Uncertainty here is estimated varying the renormalisation scale



$$\alpha_s^{\text{CMW}}(k_t; x_R) = \alpha_s(x_R k_t) \left(1 + \frac{\alpha_s(x_R k_t)}{2\pi} (K_{\text{CMW}} + \Delta K(\Phi)) + 2\alpha_s(x_R k_t) b_0 (1-z) \ln x_R \right)$$

Summary and Conclusions

- **NLL shower about to become the new standard, but not enough for the FCC!**
 - benefits of LL → NLL include reduced uncertainties (**reliable estimate uncertainties**)
 - for realistic applications we also need **massive quarks** (Deductor and Alaric already include them), at least NLO matching, and **tuning**
- **Higher log accuracy is one of the next frontiers**
 - double-soft (+ virtual) corrections: NNDL multiplicity and NSL non-global logarithms
- **Percent precision also requires these aspects not addressed here**
 - **NLO and NNLO Matching**
 - **QED** shower
 - Leading **power corrections** (e.g. in PDF's, Frixione&Webber [2309.15587](#))