

System size analysis of the fireball produced in heavy-ion collisions

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1 Introduction

- QCD phase diagram
- Evolution of quark-gluon plasma
- The STAR experiment
- Beam Energy Scan program on RHIC

2 Used methods

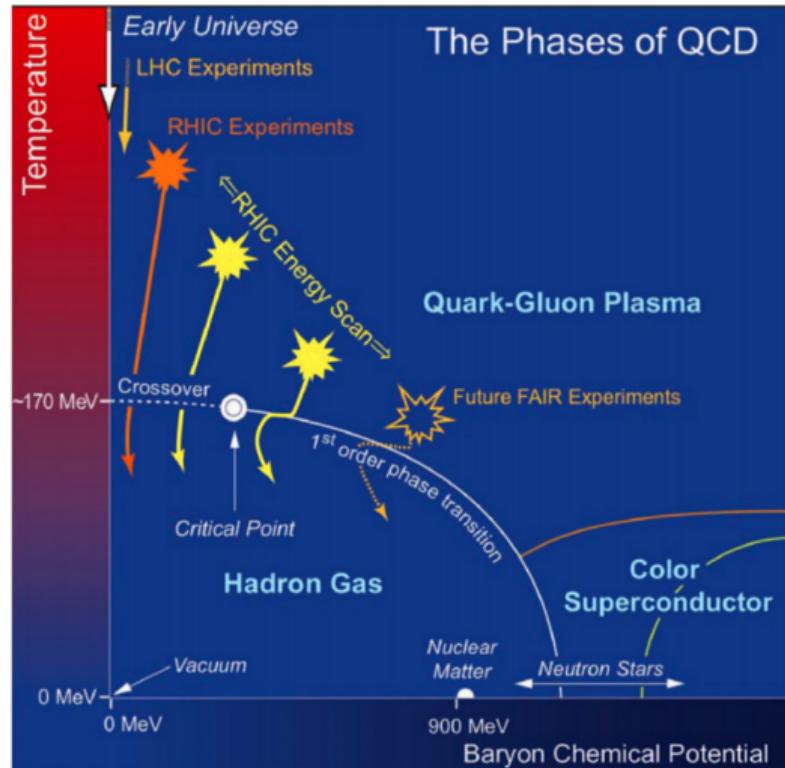
3 Results

4 Discussions and conclusion

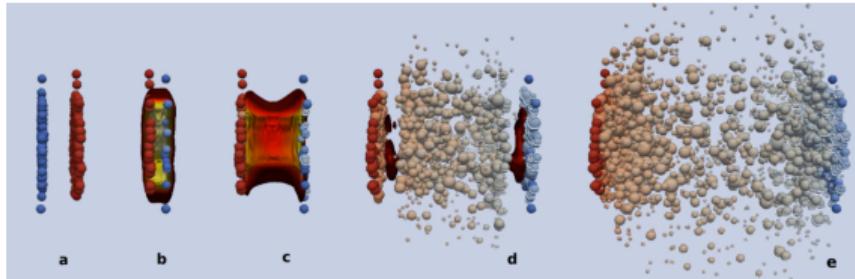
QCD phase diagram

Parameters:

- ▶ Temperature T
- ▶ Baryon chemical potential μ

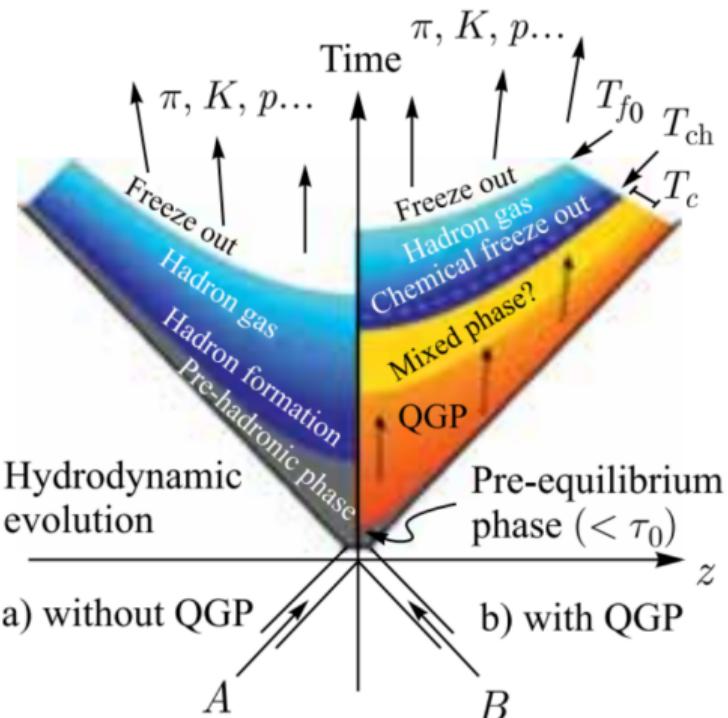


Evolution of quark-gluon plasma



Stages of ultrarelativistic nuclear collision:

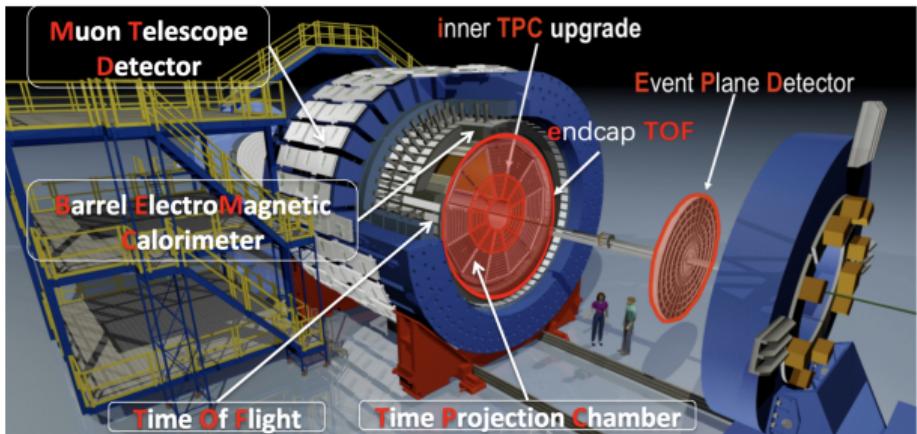
- a) Nuclei go through each other
- b) Expansion
- c) Formation of hot matter
- d) Hadron gas
- e) Final particles



The STAR experiment

TPC (Time Projection Chamber)

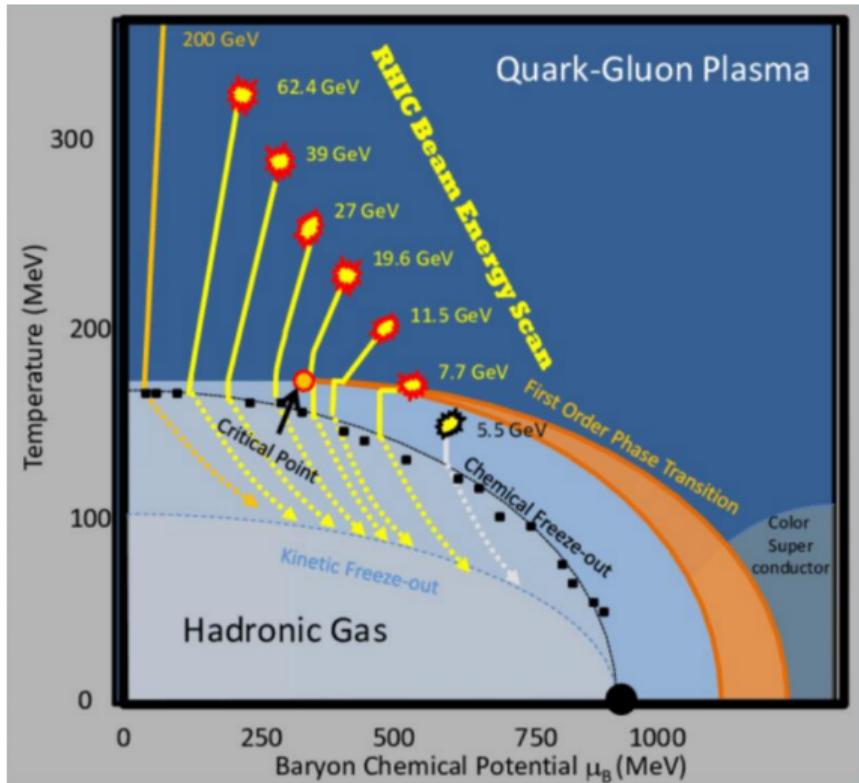
- ▶ used for tracking and identification
- ▶ length 4.2 m, diameter 4 m (1 m),
- ▶ azimuthal angle 2π
- ▶ pseudorapidity range $|\eta| < 1$
- ▶ in a magnetic field 0.5 Tesla



Beam Energy Scan program on RHIC

Used data:

- ▶ RHIC BES-I, 2010-2011
- ▶ $Au + Au \sqrt{S_{NN}} = 7.7 - 39 \text{ GeV}$.
- ▶ Phys.Rev.C 96 (2017) 044904,
Phys.Rev.C 92 (2015) 014904.



① Introduction

② Used methods

- Tsallis-3 statistics
- Femtoscopic approach

③ Results

④ Discussions and conclusion

Tsallis-3 statistics (J.Phys.G:Nucl.Part.Phys.50 2023 125002)

- ▶ Tsallis entropy:

$$S = \sum_i \frac{p_i^q - p_i}{1 - q}, \quad \sum_i p_i = 1,$$

где p_i - probability i th microscopic state of the system, $q \in [0, \infty]$.

- ▶ In the Gibbs limit $q \rightarrow 1$ the entropy recovers the Boltzmann-Gibbs entropy:

$$S = \sum_i p_i \ln p_i$$

- ▶ In Grand Canonical Ensemble thermodynamic potential Ω takes the form:

$$\Omega = \langle H \rangle - TS - \mu \langle N \rangle, \text{ where}$$

$$\langle H \rangle = \frac{1}{\theta} \sum_i p_i^q E_i, \quad \langle N \rangle = \frac{1}{\theta} \sum_i p_i^q N_i, \quad \theta = \sum_i p_i^q.$$

Tsallis-3 statistics

Let's consider a relativistic ideal gas in the Grand Canonical Ensemble. From the principle of thermodynamic equilibrium (the principle of maximum entropy), normalization expressions for the parameters $\Lambda = -\theta TS + \langle H \rangle - \mu \langle N \rangle$ и $\theta = \sum_i p_i^q$ can be derived:

$$\begin{cases} 1 = \sum_{n=0}^{n_0} \frac{\omega^n}{n! \Gamma(\frac{1}{q-1})} \int_0^\infty t^{\frac{2-q}{q-1}-n} e^{-t+\beta'(\Lambda+\mu n)} (\text{K}_2(\beta'm))^n dt, \\ \theta = \sum_{n=0}^{n_0} \frac{\omega^n}{n! \Gamma(\frac{q}{q-1})} \int_0^\infty t^{\frac{1}{q-1}-n} e^{-t+\beta'(\Lambda+\mu n)} (\text{K}_2(\beta'm))^n dt, \end{cases}.$$

$$\text{where } \omega = \frac{gV}{2\pi^2} \frac{m^2 T \theta^2}{q-1}, \quad \beta' = \frac{-t(1-q)}{T\theta^2},$$

n_0 - the number of terms to be taken into account, counting from zero.

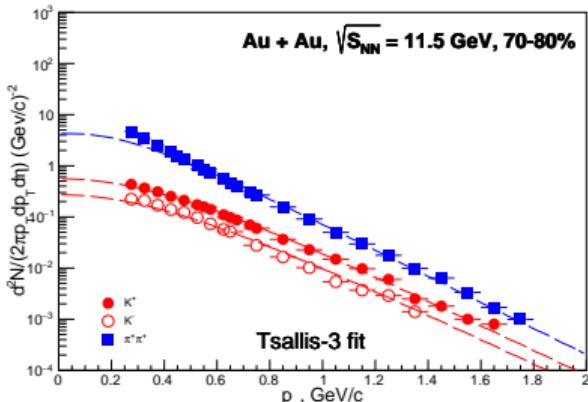
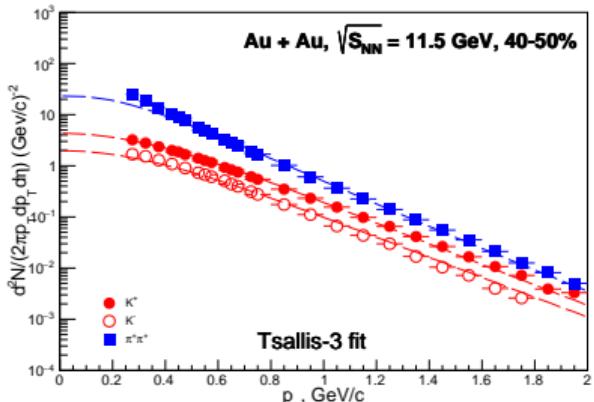
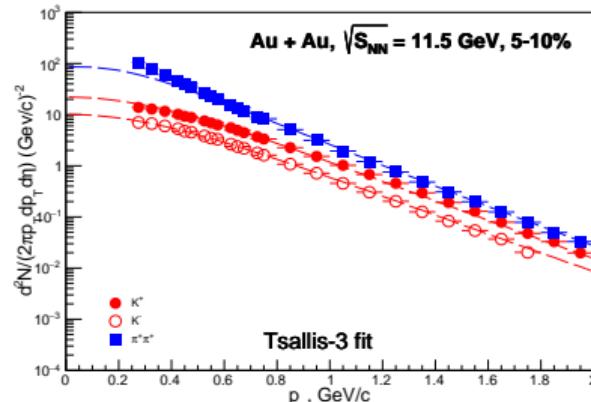
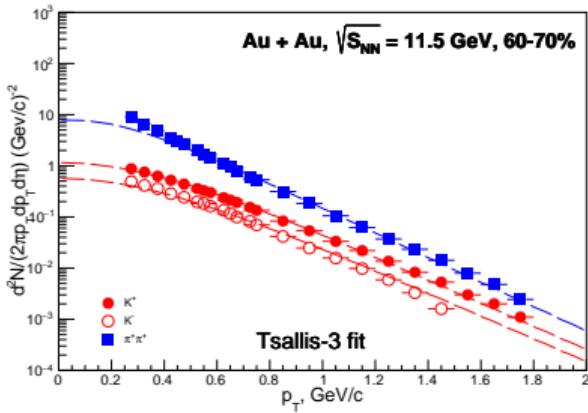
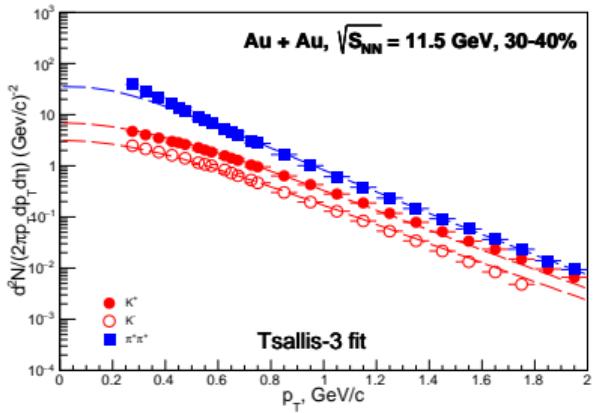
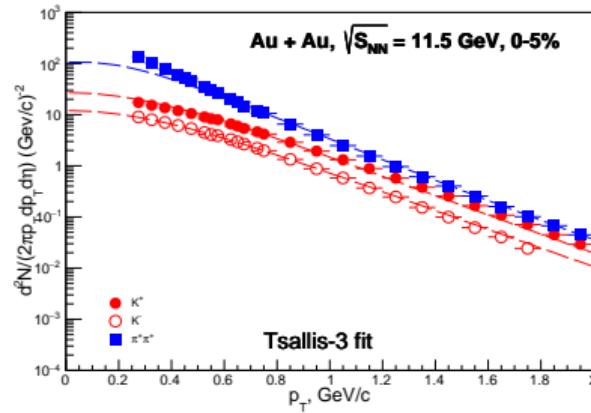
Tsallis-3 statistics

- The expression for transverse momentum distribution of particles of relativistic ideal gas in the grand canonical ensemble in rapidity range $y \in [y_{\min}, y_{\max}]$ takes the form:

$$\begin{aligned} \frac{d^2N}{p_T dp_T dy} \Big|_{y_{\min}}^{y_{\max}} &= \frac{gV}{(2\pi)^2} m_T \int_{y_{\min}}^{y_{\max}} dy \cosh y \times \\ &\times \frac{1}{\theta} \sum_{n=0}^{n_0} \frac{\omega^n}{n! \Gamma(\frac{q}{q-1})} \int_0^\infty t^{\frac{1}{q-1}-n} e^{-t+\beta'(\Lambda-m_T \cosh y + \mu(n+1))} (K_2(\beta' m))^n dt \end{aligned}$$

- In this work: $n_0 = 1$, $\mu = 0$.

Tsallis-3 statistics

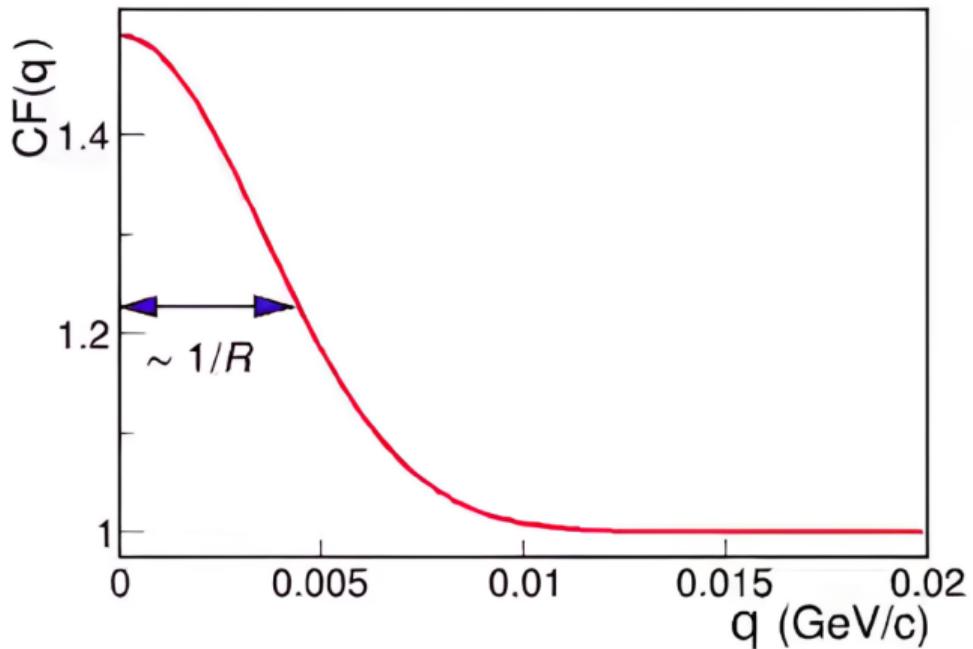


Femtoscopy approach

- Correlation function:

$$C(\vec{q}) = \frac{N(\vec{q})}{D(\vec{q})},$$

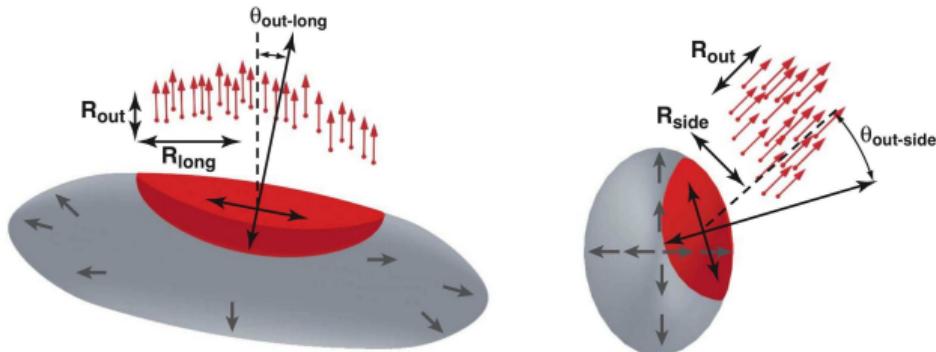
- $N(\vec{q})$ - uses pairs from the same event, contains quantum statistic correlations and interactions,
- $D(\vec{q})$ - uses pairs from different events, contains only background.



Femtoscopic approach

- ▶ Bertsch-Pratt out-side-long coordinate system:

- ▶ q_{long} - along the beam direction
- ▶ q_{out} - along the transverse momentum of the pair
- ▶ q_{side} - perpendicular to longitudinal and outward directions



- ▶ Gaussian parameterization (Bowler-Sinyukov):

$$C(\vec{q}) = (1 - \lambda) + K_{\text{Coul}}(q_{\text{inv}}) \lambda \exp(-q_o^2 R_o^2 - q_s^2 R_s^2 - q_l^2 R_l^2 - 2q_o q_s R_{os}^2 - 2q_o q_l R_{ol}^2),$$

- ▶ λ - correlation strength
- ▶ $K_{\text{Coul}}(q_{\text{inv}})$ - Coulomb correction factor
- ▶ $R_{\text{long}} \sim$ source lifetime
- ▶ $R_{\text{out}} \sim$ - geometric size and emission duration
- ▶ $R_{\text{side}} \sim$ - geometric size
- ▶ $R_{os,ol}$ - crossterms

Femtoscopy approach

Parameterizations:

- ▶ m_T scaling: $R_i \propto \frac{1}{\sqrt{m_T}}$

- ▶ Power law: $R_i \propto m_T^{-\alpha}$

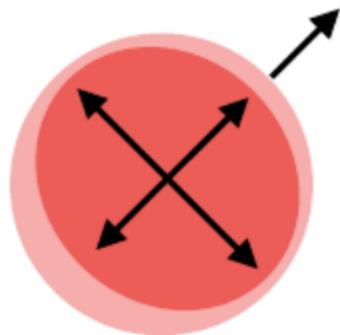
- ▶ Blast-Wave:

$$R_{\text{side}} = \frac{R_0}{\sqrt{1 + \rho_0^2(m_T/T)}},$$

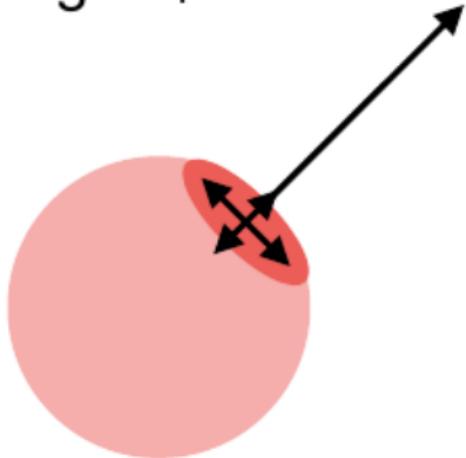
$$R_{\text{long}} = \tau \sqrt{\frac{T}{m_t} \frac{K_2(m_T/T)}{K_2(m_t/T)}}$$

(Phys. Rev. C51:328-338 1995)

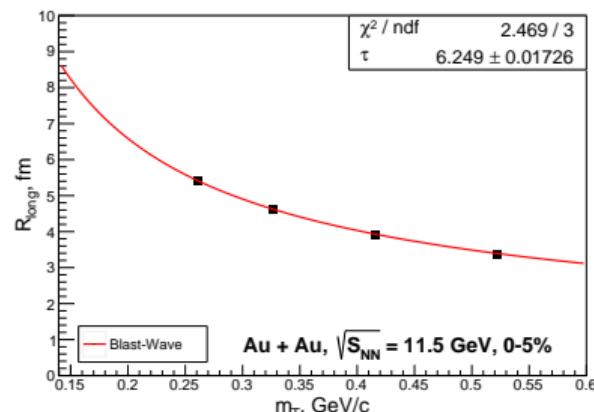
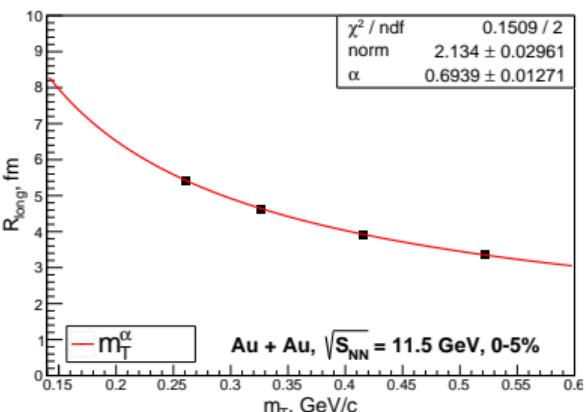
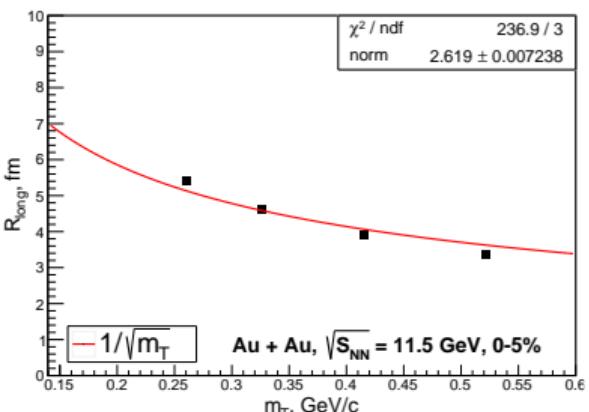
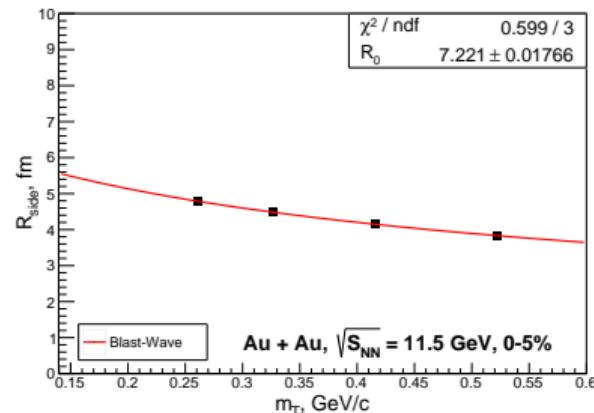
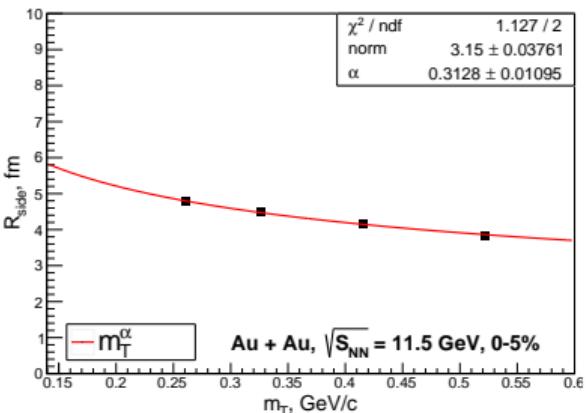
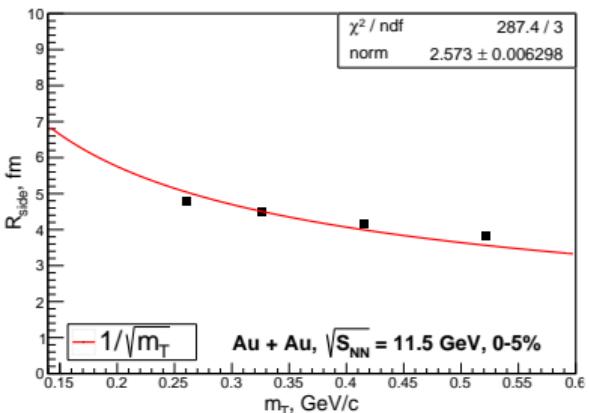
Low k_T



High k_T



Femtoscopy approach



Femtoscopic approach

Parameterizations:

- ▶ Recalculate gaussian radii to uniform distribution:

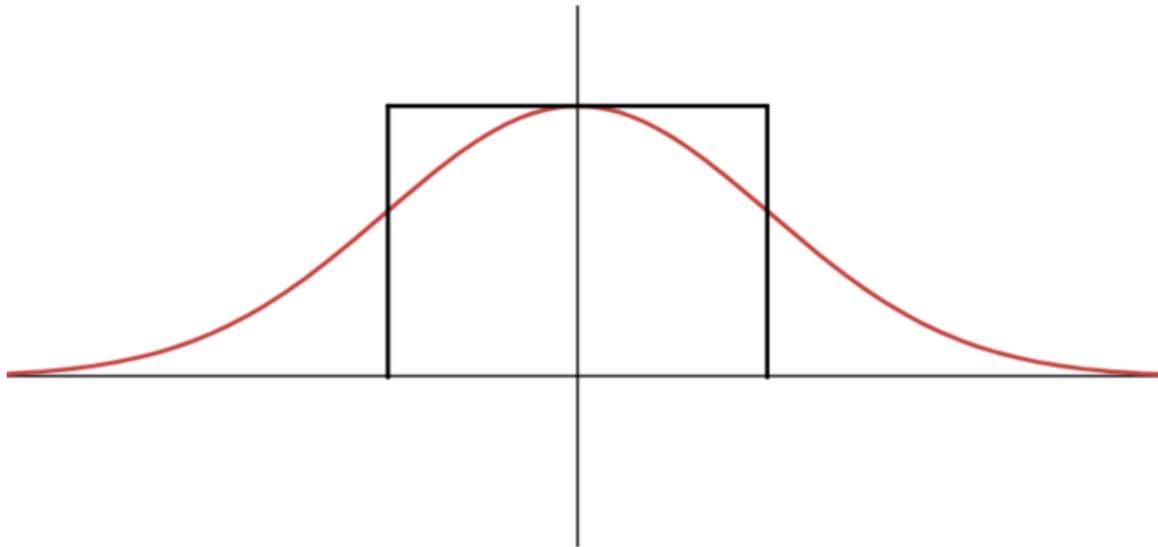
$$\sigma_{\text{gauss}} = \sigma_{\text{uniform}}$$

$$R_{\text{gauss}} = \sqrt{\frac{1}{12}} R_{\text{uniform}}$$

- ▶ Calculate the volume:

$$V = R_{\text{side}}^2 R_{\text{long}}$$

(assumption: $R_{\text{out}} \approx R_{\text{side}}$)



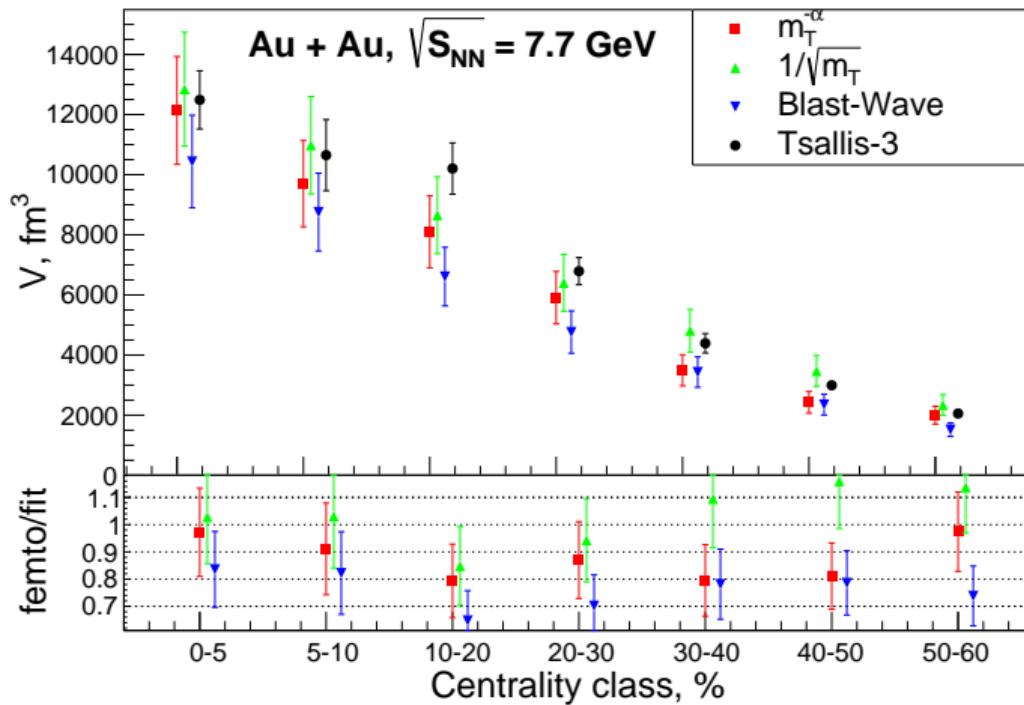
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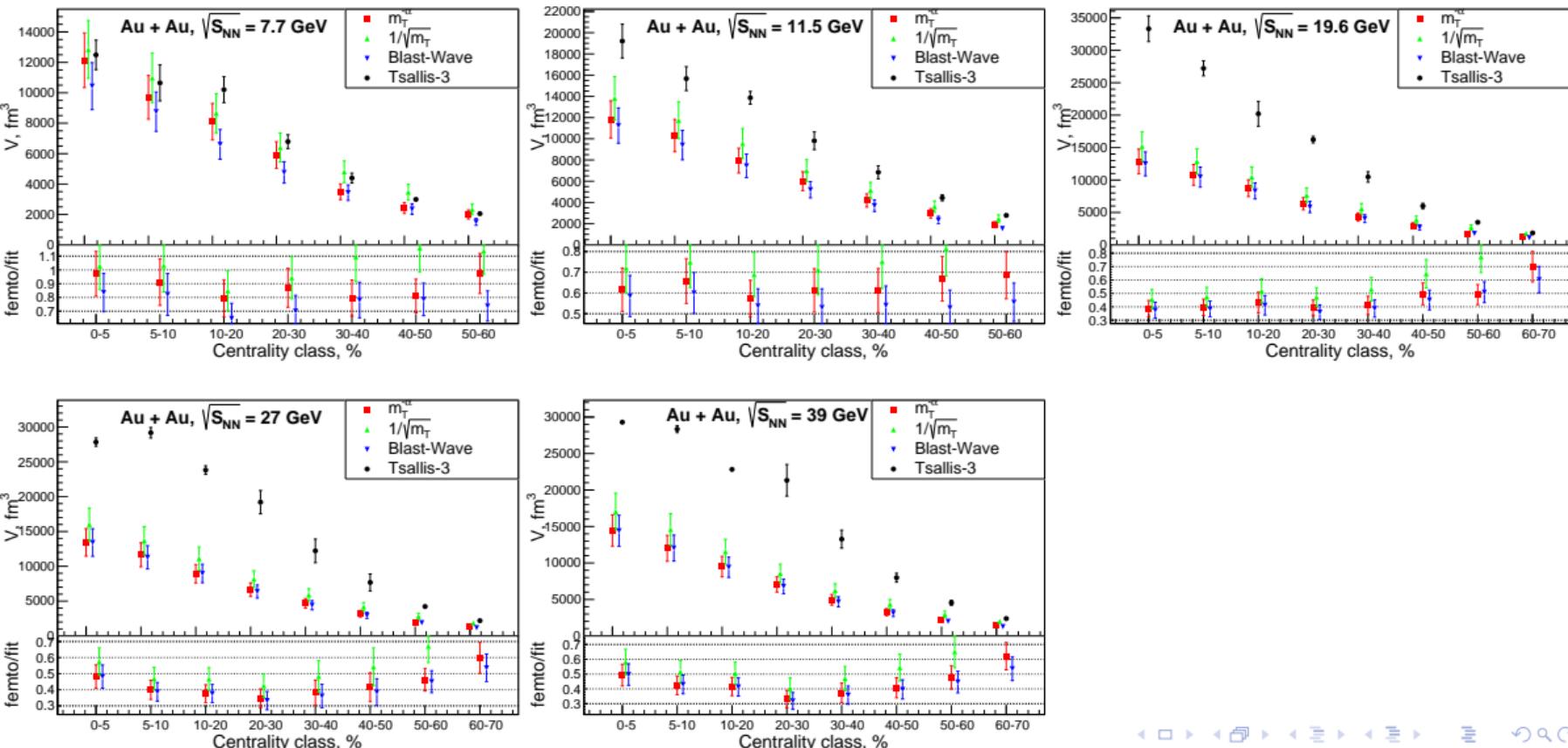
③ Results

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Results



Results



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Discussions

- ▶ Correlations measure only the region of homogeneity. If the collective expansion is strong, then the region of homogeneity can be smaller than the entire source volume.
- ▶ Correlation function can deviate from Gaussian by exponential tails caused by resonance decay contributions. This may lead to underestimation of femtoscopic radii.
- ▶ In the Tsallis-3 statistics, the system is considered as an ideal gas, so particle interactions and collective effects, including rescattering, are not taken into account.

Conclusion

- ▶ Two approaches: femtoscopy and statistical model.
- ▶ In the first approach the fireball is considered as a thermodynamic system in the Grand Canonical Ensemble.
- ▶ Second method allows the measurement of the spatio-temporal properties of the region of homogeneity by using two-particle correlations of final particles.
- ▶ We observe that at high energies the volume values diverge significantly, while at low energies, they are more consistent.

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Thank you for your attention!