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# $B \rightarrow \pi\pi$ decays: branching ratios and CP asymmetries

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M.V., Phys. Atom. Nucl., **69** (2006) 679

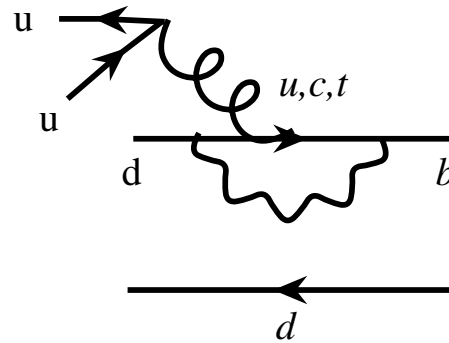
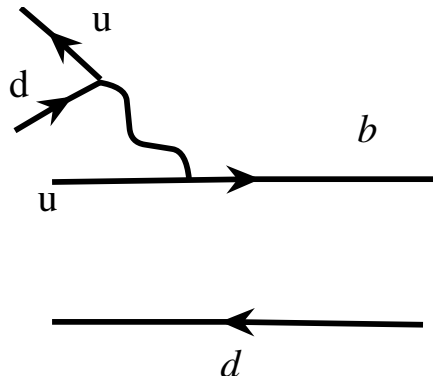
A.B.Kaidalov, M.V., hep-ph/0603013

# PLAN

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- $P/T$  - small parameter
- $B \rightarrow \pi\pi$  in tree approximation:  $A_0, A_2, \delta, \alpha$
- factorisation: pro and contro
- $K \rightarrow \pi\pi, D \rightarrow \pi\pi, B \rightarrow D\pi$
- FSI in  $B \rightarrow \pi\pi$  and  $\pi\pi - \rho\rho$  problem
- penguin corrections:  $C_{+-}, C_{00}, \alpha$
- Conclusions

$$b \rightarrow ud\bar{u}$$



$B \rightarrow \pi\pi, B \rightarrow \rho\rho, B \rightarrow \rho\pi$  decays  
How large is penguin?

$$c_1 = 1.09, c_2 = -0.21;$$

$$c_3 = 0.013, c_4 = -0.032, c_5 = 0.009, c_6 = -0.037$$

# $P/T$ - small

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$$P/T = 0 \implies \sin 2\alpha^T = S_{+-}$$

$$B \rightarrow \pi^+ \pi^- : \alpha_{BABAR}^T = (99 \pm 5)^\circ$$

$$B \rightarrow \rho^+ \rho^- : \alpha^T = (96 \pm 7)^\circ$$

$$B \rightarrow \pi^\pm \rho^\mp : \alpha^T = (94 \pm 4)^\circ, \text{ or } (86 \pm 4)^\circ$$

values of  $\alpha$  from all 3 decays agree with each other and with global CKM fit result.

R.Aleksan, F.Buccella, A.Le Yaouanc, L.Oliver, O.Pene, J.-C.Raynal (1995):

$$\Delta\alpha^{\pi\pi} > \Delta\alpha^{\rho\rho} > \Delta\alpha^{\pi\rho}$$

from smallness of  $B \rightarrow \rho^0 \rho^0 / B \rightarrow \rho^+ \rho^-$  one can PROVE smallness of penguins in  $B \rightarrow \rho\rho$  decays; however from relatively large value of  $B \rightarrow \pi^0 \pi^0 / B \rightarrow \pi^+ \pi^-$  it DOES NOT follow that penguins in  $B \rightarrow \pi\pi$  decays are large.

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Perturbation theory over  $P/T$

# $B \rightarrow \pi\pi$ exp data

	BABAR	Belle	Heavy Flavor Averaging Group
$B_{+-}$	$5.5 \pm 0.5$	$4.4 \pm 0.7$	$5.0 \pm 0.4$
$B_{00}$	$1.17 \pm 0.33$	$2.3 \pm 0.5$	$1.45 \pm 0.29$
$B_{+0}$	$5.8 \pm 0.7$	$5.0 \pm 1.3$	$5.5 \pm 0.6$
$S_{+-}$	$-0.30 \pm 0.17$	$-0.67 \pm 0.16$	$-0.50 \pm 0.12$
$C_{+-}$	$-0.09 \pm 0.15$	$-0.56 \pm 0.13$	$-0.37 \pm 0.10$
$C_{00}$	$-0.12 \pm 0.56$	$-0.44 \pm 0.56$	$-0.28 \pm 0.39$

$B_{ik}$  in units  $10^{-6}$

$C_{+-}$ : BABAR - little penguin; Belle - big penguin. Average?

# $B \rightarrow \pi\pi$ phenomenology

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$$\begin{aligned} M_{\bar{B}_d \rightarrow \pi^+ \pi^-} &= \frac{G_F}{\sqrt{2}} |V_{ub} V_{ud}^*| m_B^2 f_\pi f_+(0) \left\{ e^{-i\gamma} \frac{1}{2\sqrt{3}} A_2 e^{i\delta} + \right. \\ &+ \left. e^{-i\gamma} \frac{1}{\sqrt{6}} A_0 + \left| \frac{V_{td}^* V_{tb}}{V_{ub} V_{ud}^*} \right| e^{i\beta} P e^{i\delta_p} \right\}, \end{aligned} \quad (1)$$

$$\begin{aligned} M_{\bar{B}_d \rightarrow \pi^0 \pi^0} &= \frac{G_F}{\sqrt{2}} |V_{ub} V_{ud}^*| m_B^2 f_\pi f_+(0) \left\{ -e^{-i\gamma} \frac{1}{\sqrt{3}} A_2 e^{i\delta} + \right. \\ &+ \left. e^{-i\gamma} \frac{1}{\sqrt{6}} A_0 + \left| \frac{V_{td}^* V_{tb}}{V_{ub} V_{ud}^*} \right| e^{i\beta} P e^{i\delta_p} \right\}, \end{aligned} \quad (2)$$

$$M_{\bar{B}_u \rightarrow \pi^- \pi^0} = \frac{G_F}{\sqrt{2}} |V_{ub} V_{ud}^*| m_B^2 f_\pi f_+(0) \left\{ \frac{\sqrt{3}}{2\sqrt{2}} e^{-i\gamma} A_2 e^{i\delta} \right\} \quad (3)$$

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# tree approximation: $A_0, A_2, \delta, \alpha$

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3 parameters from 3 equations ( $B_{+-}, B_{00}, B_{+0}$ ):

$$A_0 = 1.53 \pm 0.23, A_2 = 1.60 \pm 0.20, \delta = \pm(53^\circ \pm 7^\circ)$$

$$A_0^f = 1.54, \quad A_2^f = 1.35 \quad \text{but } \delta \dots$$

$$\sin 2\alpha^T = S_{+-} ,$$

$$\alpha_{\text{BABAR}}^T = 99^\circ \pm 5^\circ , \quad \alpha_{\text{Belle}}^T = 111^\circ \pm 6^\circ , \quad \alpha_{\text{average}}^T = 105^\circ \pm 4^\circ .$$

# FSI in $K \rightarrow \pi\pi, D \rightarrow \pi\pi, B \rightarrow D\pi$

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$K$  decays: 3 decay probabilities, or Watson theorem:

$$\delta_0^K = 35^\circ \pm 3^\circ, \delta_2^K = -7^\circ \pm 0.2^\circ, \delta^K = 42^\circ \pm 4^\circ$$

$D$  decays: 3 decay probabilities (Watson theorem is not applicable): factorisation also good for moduli of decay amplitudes, while  $\delta_2^D - \delta_0^D = 86^\circ \pm 4^\circ(!?)$

$1/M$  scaling of FSI phases?

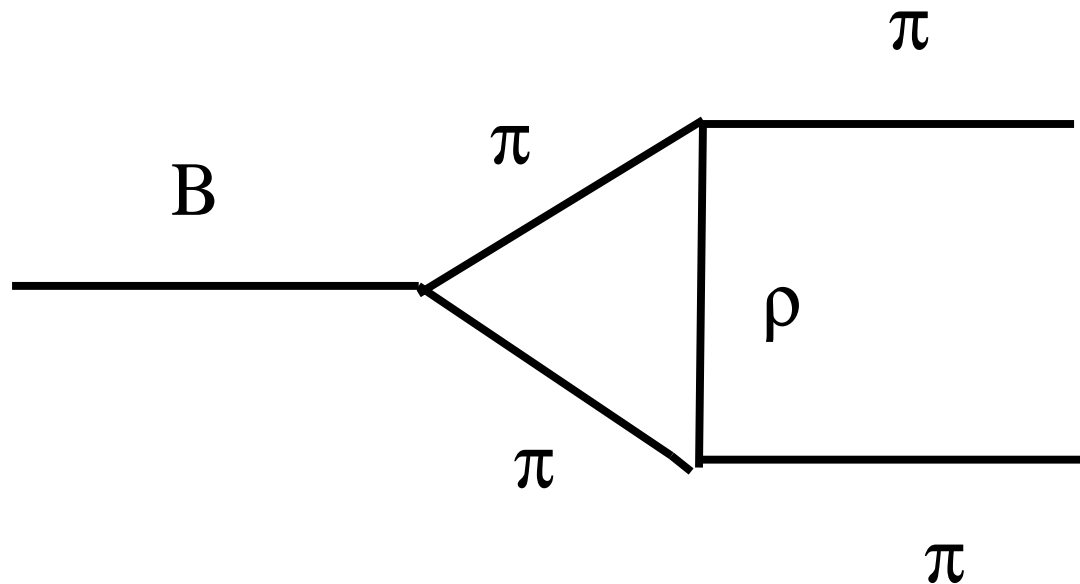
$(D\pi) : I = 1/2$  or  $3/2$  ; 3 decay probabilities:  $\delta_{D\pi} = 30^\circ \pm 7^\circ$   
Important: above the resonances

So: FSI phases can be L A R G E



# FSI in $B \rightarrow \pi\pi$

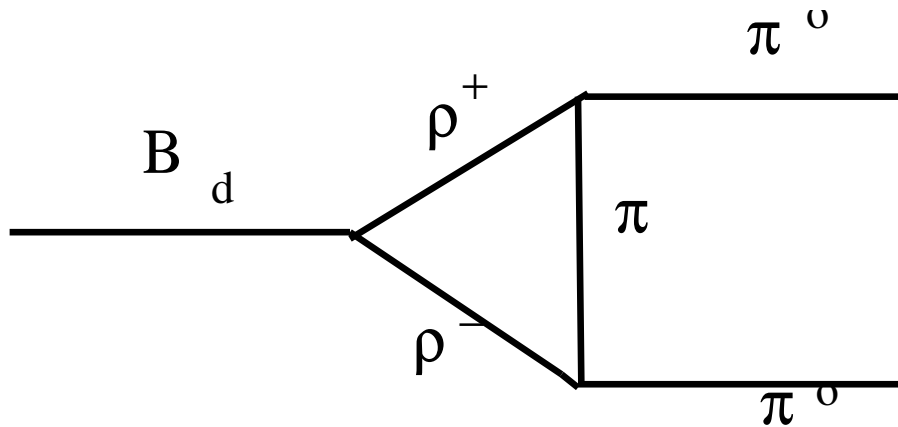
a lot of papers; one example: Cheng, Chua, Soni, P.R. D 71 (2005): "rescattering effects enhance color-suppressed neutral modes". FSI phase is extracted from the diagrams with vector particle exchanges in  $t$ -channel:



for elementary  $\rho$  meson partial wave amplitude does not decrease with energy. However all exchanges should be

reggeized,  $s^{\alpha(0)-1}$ ,  $\alpha_\rho \approx 0.5$

# phase from $\rho^+ \rho^-$ intermediate state



1.  $Br(B_d \rightarrow \rho^+ \rho^-) = 26 * 10^{-6}$

2. Will generate negligible phases in  $B \rightarrow \rho\rho$  amplitudes, because  $Br(B_d \rightarrow \pi^+ \pi^-) = 5 * 10^{-6}$

$\delta \approx 30^\circ$  is generated by intermediate  $\rho\rho$  and  $\pi\pi$  states

# penguin corrections

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To avoid shifts of  $B_{+-}$  and  $B_{00}$  we should shift  $A_0$  and  $\delta$  :

$$A_0 \rightarrow A_0 + \tilde{A}_0 , \quad \delta \rightarrow \delta + \tilde{\delta} ,$$

$$\tilde{A}_0 = \sqrt{6} \left| \frac{V_{td}}{V_{ub}} \right| \cos \alpha \cos \delta_p P ,$$

$$\tilde{\delta} = -\sqrt{6} \left| \frac{V_{td}}{V_{ub}} \right| \cos \alpha \sin \delta_p P / A_0 ,$$

where only the terms linear in  $P$  are taken into account

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In the factorisation approach we have:

$$P^f = -a_4 - \frac{2m_\pi^2}{(m_u + m_d)m_b} a_6 = 0.06 \quad ,$$

$a_4 = c_4 + c_3/3$ ,  $a_6 = c_6 + c_5/3$  and shifts of  $A_0$  and  $\delta$  are small:

$$-0.12 < \tilde{A}_0 < 0.12 \quad , \quad -4^\circ < \tilde{\delta} < 4^\circ$$

for

$$A_0 = 1.5 \quad , \quad -1 < \cos \delta_p \quad , \quad \sin \delta_p < 1 \quad \text{and} \quad 70^\circ < \alpha < 110^\circ$$

# $C_{+-}, C_{00}$

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In linear in  $P$  approximation for direct CP asymmetries we obtain:

$$C_{+-} = -2P \sin \alpha \left| \frac{V_{td}}{V_{ub}} \right| \frac{\cos(\delta_p + \varphi)}{\sqrt{\frac{1}{12} A_2^2 + \frac{1}{6} A_0^2 + \frac{1}{3\sqrt{2}} A_0 A_2 \cos \delta}} =$$
$$-4.7P \cos(\delta_p + 68^\circ)$$

$$C_{00} = -2P \sin \alpha \left| \frac{V_{td}}{V_{ub}} \right| \frac{\cos(\delta_p + \psi)}{\sqrt{\frac{1}{3} A_2^2 + \frac{1}{6} A_0^2 - \frac{\sqrt{2}}{3} \cos \delta A_0 A_2}} = 6.2P \cos \delta_p$$

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where

$$\varphi = \arccos \frac{\frac{1}{\sqrt{3}} A_2 \sin \delta}{\sqrt{\frac{1}{3} A_2^2 + \frac{2}{3} A_0^2 + \frac{2\sqrt{2}}{3} A_0 A_2 \cos \delta}} \approx 68^\circ ,$$

$$\psi = \arccos \frac{-\frac{2}{\sqrt{3}} A_2 \sin \delta}{\sqrt{\frac{4}{3} A_2^2 + \frac{2}{3} A_0^2 - \frac{4\sqrt{2}}{3} A_0 A_2 \cos \delta}} \approx 175^\circ$$

From experimental values of  $C_{ik}$  we can determine  $P$  and  $\delta_p$   
- Gronau-London pass;

since experimental uncertainty in  $C_{00}$  is big while Belle and BABAR contradicts each other in  $C_{+-}$  this pass is (temporary) closed

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Let us look which values of direct asymmetries follow from our formulas. With  $P = P_f$  we get:

$$C_{+-} = -0.28 \cos(\delta_p + 68^\circ) ,$$

and for the theoretically motivated value  $\delta_p \leq 30^\circ$  we obtain:

$$0 > C_{+-} > -0.10 ,$$

which is close to BABAR result.

For direct CP asymmetry in  $B_d \rightarrow \pi^0 \pi^0$  decay we get:

$$C_{00} = 0.4 \cos \delta_p \approx 0.4 ,$$

which differs in sign from  $C_{+-}$  and is rather big. It is very interesting to check these predictions experimentally.

# $\alpha$

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$S_{+-}$  is not changed when penguins are taken into account:

$$\alpha = \alpha^T + \tilde{\alpha} ,$$

$$\tilde{\alpha} = - \left| \frac{V_{td}}{V_{ub}} \right| P(1 + C_{+-}) \sin \alpha \frac{\cos(\delta_p - \kappa)}{\sqrt{\frac{1}{12} A_2^2 + \frac{1}{6} A_0^2 + \frac{1}{3\sqrt{2}} A_0 A_2 \cos \delta}} ,$$

$$\kappa = \frac{\pi}{2} - \varphi .$$

$$\tilde{\alpha} = -2.4(1 + C_{+-}) P \cos(\delta_p - \kappa) = -0.14(1 + C_{+-}) \cos(\delta_p - 22^\circ) ,$$

$$\tilde{\alpha}_{\text{BABAR}} \approx -7^\circ , \quad \alpha_{\text{BABAR}} = \alpha_{\text{BABAR}}^T + \tilde{\alpha}_{\text{BABAR}} = 92^\circ \pm 5^\circ .$$



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$$\tilde{\alpha}_{\text{average}} = -5^\circ, \quad \alpha_{\text{average}} = \alpha_{\text{average}}^T + \tilde{\alpha}_{\text{average}} = 100^\circ \pm 4^\circ .$$

Theoretical uncertainty of the value of  $\alpha$  can be estimated in the following way. Let us suppose that the accuracy of the factorisation calculation of the penguin amplitude is 100%:

$$\tilde{\alpha}_{\text{average}} = -5^\circ \pm 5_{\text{theor}}^\circ ,$$

$$\alpha_{\text{average}} = 100^\circ \pm 4_{\text{exp}}^\circ \pm 5_{\text{theor}}^\circ$$

# $B \rightarrow \rho\rho$

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Better theor accuracy of  $\alpha$  follows from  $B \rightarrow \rho\rho$  decays, where penguin contribution is 2 times smaller, and FSI phases are small:

$$\alpha^{\rho\rho} = 92^\circ \pm 7^\circ_{exp} \pm 4^\circ_{theor}$$

The model independent isospin analysis of  $B \rightarrow \rho\rho$  decays performed by BABAR gives:

$$\alpha_{BABAR}^{\rho\rho} = 100^\circ \pm 13^\circ$$

The global CKM fit results are:

$$\alpha_{CKMFitter} = 98^\circ \pm 8^\circ, \quad \alpha_{UTfit} = 95^\circ \pm 5^\circ$$

# Conclusions

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- The moduli of the amplitudes  $A_0$  and  $A_2$  are given with good accuracy by factorisation; FSI phase shift is very large,  $\delta \approx 50^\circ$ .
- Theoretical uncertainty of the value of  $\alpha$  extracted from  $B \rightarrow \pi\pi$  data on  $S_{+-}$  is at the level of few degrees.
- Resolution of the contradiction of Belle and BABAR data on CP asymmetries are very important both for checking the correctness of our approach (  $C$  ) and determination of angle  $\alpha$  (  $S$  ).