

Looking for Wtb anomalous couplings in top pair decays

J. A. Aguilar-Saavedra

in collaboration with J. Carvalho, N. Castro, A. Onofre, F. Veloso

Departamento de Física Teórica y del Cosmos
Universidad de Granada

Workshop “*Flavour in the era of the LHC*”
CERN, May 2006

Summary

- 1 Introduction
- 2 Top decays and W helicity
- 3 Top pair decays and $t\bar{t}$ spin correlations
- 4 Top pair decays and CP violation

Introduction

Reminder

- The Higgs boson may exist or not (most people hope it)
- Supersymmetry may exist or not (some people hope it)
- New fermions may exist or not (fewer people hope it)
- The top quark **definitely exists** (we are happy for it)
and we must study its properties at LHC

Introduction

Reminder

- The Higgs boson may exist or not (most people hope it)
- Supersymmetry may exist or not (some people hope it)
- New fermions may exist or not (fewer people hope it)
- The top quark **definitely exists** (we are happy for it)
and we must study its properties at LHC

Introduction

Reminder

- The Higgs boson may exist or not (most people hope it)
- Supersymmetry may exist or not (some people hope it)
- New fermions may exist or not (fewer people hope it)
- The top quark **definitely exists** (we are happy for it)
and we must study its properties at LHC

Introduction

Reminder

- The Higgs boson may exist or not (most people hope it)
- Supersymmetry may exist or not (some people hope it)
- New fermions may exist or not (fewer people hope it)
- The top quark definitely exists (we are happy for it)
and we must study its properties at LHC

Introduction

Reminder

- The Higgs boson may exist or not (most people hope it)
- Supersymmetry may exist or not (some people hope it)
- New fermions may exist or not (fewer people hope it)
- The top quark **definitely exists** (we are happy for it)
and we must study its properties at LHC

Top quark properties

- Possibly new physics influence
- Mass ➔ Measured at Tevatron
But no standard theory for masses
- Sensitive to new physics corrections
- Couplings ➔ Predicted in the SM
May be measured / bounded
- Spin ➔ Testable in top production and decay

Probing top couplings

- Wtb vertex ➡
 - [top decays $t \rightarrow W^+ b$ [many authors]
 - single top LHC [Boos et al., EPJC '99]
- $Ztt, \gamma tt$ ➡
 - [$t\bar{t}Z, t\bar{t}\gamma$ at LHC [Baur et al., PRD '05,'06]
 - $t\bar{t}$ at ILC
- gtt ➡
 - [$t\bar{t}$ at LHC [Cheung, PRD '96]
 - $t\bar{t}g$ at ILC [Rizzo, PRD '94]
- $Ztq, \gamma tq, gtq$ ➡
 - FCN production and decay [many authors]

General Wtb vertex

$$\begin{aligned}\mathcal{L} = & -\frac{g}{\sqrt{2}} \bar{b} \gamma^\mu (V_L P_L + \textcolor{red}{V}_R P_R) t W_\mu^- \\ & -\frac{g}{\sqrt{2}} \bar{b} \frac{i\sigma^{\mu\nu} q_\nu}{M_W} (\textcolor{red}{g}_L P_L + \textcolor{red}{g}_R P_R) t W_\mu^- + \text{h.c.}\end{aligned}$$

(Additional terms p^μ , etc. rewritten with Gordon identities)

V_R, g_L, g_R zero at tree level

$V_L \simeq 1$, normalised to unity

V_R, g_L, g_R expected to be (very) small

Top decays and W helicity

Equivalent descriptions of $t \rightarrow W^+ b \rightarrow \ell^+ \nu b$

- ℓ angular distribution \leftrightarrow angular asymmetries
in W rest frame, determined by W^+ helicity fractions
- ℓ energy distribution \leftrightarrow energy asymmetries
in t rest frame, determined by W^+ helicity fractions and W boost

ℓ angular distribution in W rest frame

Assume t, W on shell

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_\ell^*} = \frac{3}{8} (1 + \cos \theta_\ell^*)^2 F_R + \frac{3}{8} (1 - \cos \theta_\ell^*)^2 F_L + \frac{3}{4} \sin^2 \theta_\ell^* F_0$$

θ_ℓ^* angle between $\begin{cases} \ell \text{ momentum in } W \text{ rest frame} \\ W \text{ momentum in } t \text{ rest frame} \end{cases}$

$$\theta_\ell^* = \pi - \theta_{\ell b} \quad \theta_{\ell b} \text{ angle between } \ell \text{ and } b \text{ in } W \text{ rest frame}$$

$$\Gamma(t \rightarrow W^+ b) = \Gamma_R + \Gamma_L + \Gamma_0 \quad F_i = \Gamma_i / \Gamma$$

$F_i \rightarrow$ determine ℓ distribution

W helicity fractions

($x_W = M_W/m_t$, $x_b = m_b/m_t$, $|\vec{q}| = W$ momentum)

$$\begin{aligned}\Gamma_0 &= \frac{g^2 |\vec{q}|}{32\pi} \left\{ \frac{m_t^2}{M_W^2} \left[|V_L|^2 + |V_R|^2 \right] \left(1 - x_W^2 - 2x_b^2 - x_W^2 x_b^2 + x_b^4 \right) - 4x_b \operatorname{Re} V_L V_R^* \right. \\ &\quad \left. + [|g_L|^2 + |g_R|^2] \left(1 - x_W^2 + x_b^2 \right) - 4x_b \operatorname{Re} g_L g_R^* \right. \\ &\quad \left. - 2 \frac{m_t}{M_W} \operatorname{Re} [V_L g_R^* + V_R g_L^*] \left(1 - x_W^2 - x_b^2 \right) + 2 \frac{m_t}{M_W} x_b \operatorname{Re} [V_L g_L^* + V_R g_R^*] \left(1 + x_W^2 - x_b^2 \right) \right\} \\ \Gamma_{R,L} &= \frac{g^2 |\vec{q}|}{32\pi} \left\{ \left[|V_L|^2 + |V_R|^2 \right] \left(1 - x_W^2 + x_b^2 \right) - 4x_b \operatorname{Re} V_L V_R^* \right. \\ &\quad \left. + \frac{m_t^2}{M_W^2} \left[|g_L|^2 + |g_R|^2 \right] \left(1 - x_W^2 - 2x_b^2 - x_W^2 x_b^2 + x_b^4 \right) - 4x_b \operatorname{Re} g_L g_R^* \right. \\ &\quad \left. - 2 \frac{m_t}{M_W} \operatorname{Re} [V_L g_R^* + V_R g_L^*] \left(1 - x_W^2 - x_b^2 \right) + 2 \frac{m_t}{M_W} x_b \operatorname{Re} [V_L g_L^* + V_R g_R^*] \left(1 + x_W^2 - x_b^2 \right) \right\} \\ &\quad \pm \frac{g^2}{64\pi} \frac{m_t^3}{M_W^2} \left\{ -x_W^2 \left[|V_L|^2 - |V_R|^2 \right] + [|g_L|^2 - |g_R|^2] \left(1 - x_b^2 \right) \right. \\ &\quad \left. + 2x_W \operatorname{Re} [V_L g_R^* - V_R g_L^*] + 2x_W x_b \operatorname{Re} [V_L g_L^* - V_R g_R^*] \right\} \\ &\quad \times \left(1 - 2x_W^2 - 2x_b^2 + x_W^4 - 2x_W^2 x_b^2 + x_b^4 \right)\end{aligned}$$

W helicity fractions

$(x_W = M_W/m_t, \quad x_b = m_b/m_t, \quad |\vec{q}| = W \text{ momentum})$

$$\begin{aligned} \Gamma_0 &= \frac{g^2 |\vec{q}|}{32\pi} \left\{ \frac{m_t^2}{M_W^2} \left[|V_L|^2 + |V_R|^2 \right] \left(1 - x_W^2 - 2x_b^2 - x_W^2 x_b^2 + x_b^4 \right) - 4x_b \operatorname{Re} V_L V_R^* \right. \\ &\quad \left. + \left[|\mathbf{g}_L|^2 + |\mathbf{g}_R|^2 \right] \left(1 - x_W^2 + x_b^2 \right) - 4x_b \operatorname{Re} g_L g_R^* \right. \\ &\quad \left. - 2 \frac{m_t}{M_W} \operatorname{Re} [V_L g_R^* + V_R g_L^*] \left(1 - x_W^2 - x_b^2 \right) + 2 \frac{m_t}{M_W} x_b \operatorname{Re} [\mathbf{V}_L \mathbf{g}_L^* + V_R g_R^*] \left(1 + x_W^2 - x_b^2 \right) \right\} \\ \Gamma_{R,L} &= \frac{g^2 |\vec{q}|}{32\pi} \left\{ \left[|V_L|^2 + |V_R|^2 \right] \left(1 - x_W^2 + x_b^2 \right) - 4x_b \operatorname{Re} V_L V_R^* \right. \\ &\quad \left. + \frac{m_t^2}{M_W^2} \left[|\mathbf{g}_L|^2 + |\mathbf{g}_R|^2 \right] \left(1 - x_W^2 - 2x_b^2 - x_W^2 x_b^2 + x_b^4 \right) - 4x_b \operatorname{Re} g_L g_R^* \right. \\ &\quad \left. - 2 \frac{m_t}{M_W} \operatorname{Re} [V_L g_R^* + V_R g_L^*] \left(1 - x_W^2 - x_b^2 \right) + 2 \frac{m_t}{M_W} x_b \operatorname{Re} [\mathbf{V}_L \mathbf{g}_L^* + V_R g_R^*] \left(1 + x_W^2 - x_b^2 \right) \right\} \\ &\quad \pm \frac{g^2}{64\pi} \frac{m_t^3}{M_W^2} \left\{ -x_W^2 \left[|V_L|^2 - |V_R|^2 \right] + \left[|\mathbf{g}_L|^2 - |\mathbf{g}_R|^2 \right] \left(1 - x_b^2 \right) \right. \\ &\quad \left. + 2x_W \operatorname{Re} [V_L g_R^* - V_R g_L^*] + 2x_W x_b \operatorname{Re} [\mathbf{V}_L \mathbf{g}_L^* - V_R g_R^*] \right\} \\ &\quad \times \left(1 - 2x_W^2 - 2x_b^2 + x_W^4 - 2x_W^2 x_b^2 + x_b^4 \right) \end{aligned}$$

▶ See

W helicity fractions

($x_W = M_W/m_t$, $x_b = m_b/m_t$, $|\vec{q}| = W$ momentum)

$$\begin{aligned}\Gamma_0 &= \frac{g^2 |\vec{q}|}{32\pi} \left\{ \frac{\frac{m_t^2}{M_W^2}}{\left[|V_L|^2 + |V_R|^2 \right]} \left(1 - x_W^2 - 2x_b^2 - x_W^2 x_b^2 + x_b^4 \right) - 4x_b \operatorname{Re} [V_L V_R^*] \right. \\ &\quad \left. + \left[|g_L|^2 + |g_R|^2 \right] \left(1 - x_W^2 + x_b^2 \right) - 4x_b \operatorname{Re} [g_L g_R^*] \right. \\ &\quad \left. - 2 \frac{m_t}{M_W} \operatorname{Re} [V_L g_R^* + V_R g_L^*] \left(1 - x_W^2 - x_b^2 \right) + 2 \frac{m_t}{M_W} x_b \operatorname{Re} [V_L g_L^* + V_R g_R^*] \left(1 + x_W^2 - x_b^2 \right) \right\} \\ \Gamma_{R,L} &= \frac{g^2 |\vec{q}|}{32\pi} \left\{ \left[|V_L|^2 + |V_R|^2 \right] \left(1 - x_W^2 + x_b^2 \right) - 4x_b \operatorname{Re} [V_L V_R^*] \right. \\ &\quad \left. + \frac{\frac{m_t^2}{M_W^2}}{\left[|g_L|^2 + |g_R|^2 \right]} \left(1 - x_W^2 - 2x_b^2 - x_W^2 x_b^2 + x_b^4 \right) - 4x_b \operatorname{Re} [g_L g_R^*] \right. \\ &\quad \left. - 2 \frac{m_t}{M_W} \operatorname{Re} [V_L g_R^* + V_R g_L^*] \left(1 - x_W^2 - x_b^2 \right) + 2 \frac{m_t}{M_W} x_b \operatorname{Re} [V_L g_L^* + V_R g_R^*] \left(1 + x_W^2 - x_b^2 \right) \right\} \\ &\quad \pm \frac{g^2}{64\pi} \frac{m_t^3}{M_W^2} \left\{ -x_W^2 \left[|V_L|^2 - |V_R|^2 \right] + \left[|g_L|^2 - |g_R|^2 \right] \left(1 - x_b^2 \right) \right. \\ &\quad \left. + 2x_W \operatorname{Re} [V_L g_R^* - V_R g_L^*] + 2x_W x_b \operatorname{Re} [V_L g_L^* - V_R g_R^*] \right\} \\ &\quad \times \left(1 - 2x_W^2 - 2x_b^2 + x_W^4 - 2x_W^2 x_b^2 + x_b^4 \right)\end{aligned}$$

▶ See

W helicity fractions

$(x_W = M_W/m_t, \quad x_b = m_b/m_t, \quad |\vec{q}| = W \text{ momentum})$

$$\begin{aligned} \Gamma_0 &= \frac{g^2 |\vec{q}|}{32\pi} \left\{ \frac{m_t^2}{M_W^2} \left[|V_L|^2 + |V_R|^2 \right] \left(1 - x_W^2 - 2x_b^2 - x_W^2 x_b^2 + x_b^4 \right) - 4x_b \operatorname{Re} V_L V_R^* \right. \\ &\quad \left. + [|g_L|^2 + |g_R|^2] \left(1 - x_W^2 + x_b^2 \right) - 4x_b \operatorname{Re} g_L g_R^* \right. \\ &\quad \left. - 2 \frac{m_t}{M_W} \operatorname{Re} [V_L g_R^* + V_R g_L^*] \left(1 - x_W^2 - x_b^2 \right) + 2 \frac{m_t}{M_W} x_b \operatorname{Re} [V_L g_L^* + V_R g_R^*] \left(1 + x_W^2 - x_b^2 \right) \right\} \\ \Gamma_{R,L} &= \frac{g^2 |\vec{q}|}{32\pi} \left\{ \left[|V_L|^2 + |V_R|^2 \right] \left(1 - x_W^2 + x_b^2 \right) - 4x_b \operatorname{Re} V_L V_R^* \right. \\ &\quad \left. + \frac{m_t^2}{M_W^2} \left[|g_L|^2 + |g_R|^2 \right] \left(1 - x_W^2 - 2x_b^2 - x_W^2 x_b^2 + x_b^4 \right) - 4x_b \operatorname{Re} g_L g_R^* \right. \\ &\quad \left. - 2 \frac{m_t}{M_W} \operatorname{Re} [V_L g_R^* + V_R g_L^*] \left(1 - x_W^2 - x_b^2 \right) + 2 \frac{m_t}{M_W} x_b \operatorname{Re} [V_L g_L^* + V_R g_R^*] \left(1 + x_W^2 - x_b^2 \right) \right\} \\ &\quad \pm \frac{g^2}{64\pi} \frac{m_t^3}{M_W^2} \left\{ -x_W^2 \left[|V_L|^2 - |V_R|^2 \right] + [|g_L|^2 - |g_R|^2] \left(1 - x_b^2 \right) \right. \\ &\quad \left. + 2x_W x_b \operatorname{Re} [V_L g_L^* - V_R g_R^*] + 2x_W x_b \operatorname{Re} [V_L g_R^* - V_R g_L^*] \right\} \\ &\quad \times \left(1 - 2x_W^2 - 2x_b^2 + x_W^4 - 2x_W^2 x_b^2 + x_b^4 \right) \end{aligned}$$

▶ See

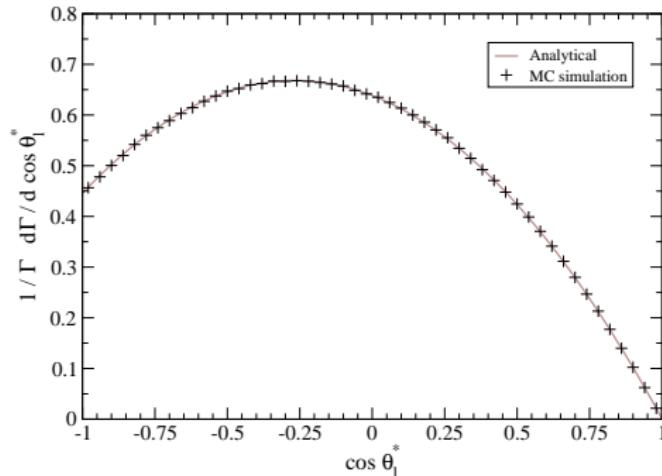
W helicity fractions

($x_W = M_W/m_t$, $x_b = m_b/m_t$, $|\vec{q}| = W$ momentum)

$$\begin{aligned}\Gamma_0 &= \frac{g^2 |\vec{q}|}{32\pi} \left\{ \frac{m_t^2}{M_W^2} \left[|V_L|^2 + |V_R|^2 \right] \left(1 - x_W^2 - 2x_b^2 - x_W^2 x_b^2 + x_b^4 \right) - 4x_b \operatorname{Re} V_L V_R^* \right. \\ &\quad \left. + [|g_L|^2 + |g_R|^2] \left(1 - x_W^2 + x_b^2 \right) - 4x_b \operatorname{Re} g_L g_R^* \right. \\ &\quad \left. - 2 \frac{m_t}{M_W} \operatorname{Re} [V_L g_R^* + V_R g_L^*] \left(1 - x_W^2 - x_b^2 \right) + 2 \frac{m_t}{M_W} x_b \operatorname{Re} [V_L g_L^* + V_R g_R^*] \left(1 + x_W^2 - x_b^2 \right) \right\} \\ \Gamma_{R,L} &= \frac{g^2 |\vec{q}|}{32\pi} \left\{ \left[|V_L|^2 + |V_R|^2 \right] \left(1 - x_W^2 + x_b^2 \right) - 4x_b \operatorname{Re} V_L V_R^* \right. \\ &\quad \left. + \frac{m_t^2}{M_W^2} \left[|g_L|^2 + |g_R|^2 \right] \left(1 - x_W^2 - 2x_b^2 - x_W^2 x_b^2 + x_b^4 \right) - 4x_b \operatorname{Re} g_L g_R^* \right. \\ &\quad \left. - 2 \frac{m_t}{M_W} \operatorname{Re} [V_L g_R^* + V_R g_L^*] \left(1 - x_W^2 - x_b^2 \right) + 2 \frac{m_t}{M_W} x_b \operatorname{Re} [V_L g_L^* + V_R g_R^*] \left(1 + x_W^2 - x_b^2 \right) \right\} \\ &\quad \pm \frac{g^2}{64\pi} \frac{m_t^3}{M_W^2} \left\{ -x_W^2 \left[|V_L|^2 - |V_R|^2 \right] + [|g_L|^2 - |g_R|^2] \left(1 - x_b^2 \right) \right. \\ &\quad \left. + 2x_W \operatorname{Re} [V_L g_R^* - V_R g_L^*] + 2x_W x_b \operatorname{Re} [V_L g_L^* - V_R g_R^*] \right\} \\ &\quad \times \left(1 - 2x_W^2 - 2x_b^2 + x_W^4 - 2x_W^2 x_b^2 + x_b^4 \right)\end{aligned}$$

Influence of finite width corrections

But . . . is this a good approximation? Yes!



Finite width corrections **do not** affect angular distribution

Equivalent parameter sets

Helicity fractions

$$\begin{aligned} F_i &= \Gamma_i / \Gamma & \rho_{R,L} &\equiv \Gamma_{R,L} / \Gamma_0 \\ F_R + F_L + F_0 &= 1 & \text{independent quantities} \\ \rho_{R,L} &= F_{R,L} / F_0 & F_0 &= (1 + \rho_R + \rho_L)^{-1} \\ && F_{R,L} &= \rho_{R,L} (1 + \rho_R + \rho_L)^{-1} \end{aligned}$$

They can be obtained from a fit to ℓ distribution

They are sensitive to anomalous couplings

Which one is better? See experimental precision

ℓ angular asymmetries in W rest frame

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_\ell^*} = \frac{3}{8}(1 + \cos \theta_\ell^*)^2 F_R + \frac{3}{8}(1 - \cos \theta_\ell^*)^2 F_L + \frac{3}{4} \sin^2 \theta_\ell^* F_0$$

Take $t \in [-1, 1]$ fixed and define asymmetries

$$A_t = \frac{N(\cos \theta_\ell^* > t) - N(\cos \theta_\ell^* < t)}{N(\cos \theta_\ell^* > t) + N(\cos \theta_\ell^* < t)}$$

$t = 0$	☞	$A_{FB} = \frac{3}{4}(F_R - F_L)$]
$t = -(2^{2/3} - 1)$	☞	$A_+ = 3\beta[F_0 + (1 + \beta)F_R]$	
$t = (2^{2/3} - 1)$	☞	$A_- = -3\beta[F_0 + (1 + \beta)F_L]$	

$(\beta \equiv 2^{1/3} - 1)$

Determine F_i
without fits

Asymmetries depend on anomalous couplings ...

▶ See

Better or worse? See experimental precision

ℓ energy in t rest frame

Related to $\cos \theta_\ell^*$ by boost

$$E_\ell = \frac{1}{2}(E_W + |\vec{q}| \cos \theta_\ell^*) \quad E_W, |\vec{q}| \text{ constants}$$

$$\begin{aligned} \frac{1}{\Gamma} \frac{d\Gamma}{dE_\ell} &= \frac{1}{(E_{\max} - E_{\min})^3} \left[3(E_\ell - E_{\min})^2 F_R + 3(E_{\max} - E_\ell)^2 F_L \right. \\ &\quad \left. + 6(E_{\max} - E_\ell)(E_\ell - E_{\min}) F_0 \right] \end{aligned}$$

with $E_{\max} = (E_W + |\vec{q}|)/2$, $E_{\min} = (E_W - |\vec{q}|)/2$

Descriptions equivalent but larger errors for E_ℓ

Top pair decays and $t\bar{t}$ spin correlations

Top decay: distributions in t rest frame

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_X} = \frac{1}{2} (1 + \alpha_X \cos \theta_X) \quad X = \ell, \nu, b, W^+, q, \bar{q}'$$

θ_X angle between X momentum and top spin

SM tree level (no anomalous couplings)

$$\alpha_{\ell^+} = \alpha_{\bar{q}'} = 1 \quad \alpha_\nu = \alpha_q = -0.31 \quad \alpha_{W^+} = -\alpha_b = 0.41$$

☞ Information on anomalous couplings and t polarisation

Antitop: $\alpha_{\bar{X}} = -\alpha_X$ if CP conserved

Top quarks produced (almost) unpolarised in $gg, q\bar{q} \rightarrow t\bar{t}$

👉 compare distribution of $\begin{cases} X \text{ from } t \text{ decay} \\ \bar{X}' \text{ from } \bar{t} \text{ decay} \end{cases}$

Double distribution in helicity basis:

$$\frac{1}{\sigma} \frac{d\sigma}{d \cos \theta_X d \cos \theta_{\bar{X}'}} = \frac{1}{4} (1 + C \alpha_X \alpha_{\bar{X}'} \cos \theta_X \cos \theta_{\bar{X}'})$$

$$\begin{aligned} C &\equiv \frac{\sigma(t_R \bar{t}_R) + \sigma(t_L \bar{t}_L) - \sigma(t_R \bar{t}_L) - \sigma(t_L \bar{t}_R)}{\sigma(t_R \bar{t}_R) + \sigma(t_L \bar{t}_L) + \sigma(t_R \bar{t}_L) + \sigma(t_L \bar{t}_R)} \\ &\simeq 0.31 \quad (\text{SM tree-level}) \end{aligned}$$

CP: $\sigma(t_R \bar{t}_R) = \sigma(t_L \bar{t}_L)$ P: $\sigma(t_R \bar{t}_L) = \sigma(t_L \bar{t}_R)$

Opening angle distribution:

$$\frac{1}{\sigma} \frac{d\sigma}{d \cos \varphi_{X\bar{X}'}} = \frac{1}{2} (1 + D \alpha_X \alpha_{\bar{X}'} \cos \varphi_{X\bar{X}'})$$

$\varphi_{X\bar{X}'}$ angle between $\begin{cases} X \text{ momentum in } t \text{ rest frame} \\ \bar{X}' \text{ momentum in } \bar{t} \text{ rest frame} \end{cases}$

$D \simeq 0.217$ (SM tree-level)

Define asymmetries

$$\begin{aligned}
 A_{X\bar{X}'} &\equiv \frac{N(\cos \theta_X \cos \theta_{\bar{X}'} > 0) - N(\cos \theta_X \cos \theta_{\bar{X}'} < 0)}{N(\cos \theta_X \cos \theta_{\bar{X}'} > 0) + N(\cos \theta_X \cos \theta_{\bar{X}'} < 0)} \\
 &= \frac{1}{4} C \alpha_X \alpha_{\bar{X}'} \\
 \tilde{A}_{X\bar{X}'} &\equiv \frac{N(\cos \varphi_{X\bar{X}'} > 0) - N(\cos \varphi_{X\bar{X}'} < 0)}{N(\cos \varphi_{X\bar{X}'} > 0) + N(\cos \varphi_{X\bar{X}'} < 0)} \\
 &= \frac{1}{2} D \alpha_X \alpha_{\bar{X}'}
 \end{aligned}$$

Sensitive to new physics:

- in **production** (C, D)

[Bernreuther et al., PRD '98]

[Cheung, PRD '97]

- in **decay** ($\alpha_X, \alpha_{\bar{X}'}$)

▶ See

Measurement of spin correlation parameters

Spin asymmetries less sensitive to anomalous couplings

▶ See

Use expected precisions
to constrain uncertainty

[Hubaut et al., EPJC '05]
[JAAS & Coimbra group, in prep.]

$$A_{\ell\ell'} \rightarrow C = 0.310 \pm 0.024 \text{ (exp)} \begin{array}{l} +2 \times 10^{-5} \\ -0.0003 \end{array} (\delta g_R) \begin{array}{l} +2 \times 10^{-5} \\ -4 \times 10^{-6} \end{array} (\delta g_L) \begin{array}{l} +0 \\ -0.0062 \end{array} (\delta V_R)$$

$$A_{\ell j} \rightarrow C = 0.310 \pm 0.045 \text{ (exp)} \begin{array}{l} +0.0005 \\ -0.0007 \end{array} (\delta g_R) \begin{array}{l} +0.0002 \\ -0.0009 \end{array} (\delta g_L) \begin{array}{l} +0.0001 \\ -0.0095 \end{array} (\delta V_R)$$

$$\tilde{A}_{\ell\ell'} \rightarrow D = 0.217 \pm 0.011 \text{ (exp)} \begin{array}{l} +4 \times 10^{-6} \\ -0.0002 \end{array} (\delta g_R) \begin{array}{l} +1 \times 10^{-5} \\ -2 \times 10^{-6} \end{array} (\delta g_L) \begin{array}{l} +0 \\ -0.0042 \end{array} (\delta V_R)$$

$$\tilde{A}_{\ell j} \rightarrow D = 0.217 \pm 0.024 \text{ (exp)} \begin{array}{l} +8 \times 10^{-5} \\ -0.0003 \end{array} (\delta g_R) \begin{array}{l} +0.0001 \\ -0.0006 \end{array} (\delta g_L) \begin{array}{l} +7 \times 10^{-5} \\ -0.0066 \end{array} (\delta V_R)$$

Top pair decays and CP violation

Until now we have assumed real V_R , g_L , g_R

... but this may not be the full story ...

Helicity fractions depend on phases through interference terms

More important for g_R  larger phase effects

▶ See

$g_R = \pm 0.055$ gives 3σ effect, $g_R = \pm 0.055i$ gives $\simeq 0.3\sigma$

Spin correlation asymmetries do not help

▶ See

Can CP violating asymmetries help? Perhaps ...

CP violation in top decays

CP-violating asymmetry

[Bernreuther et al., PLB '93]

$$T_2 = (\vec{p}_b - \vec{p}_{\bar{b}}) \cdot (\vec{p}_{\ell^+} \times \vec{p}_{\ell^-})$$

(for dilepton channel only)

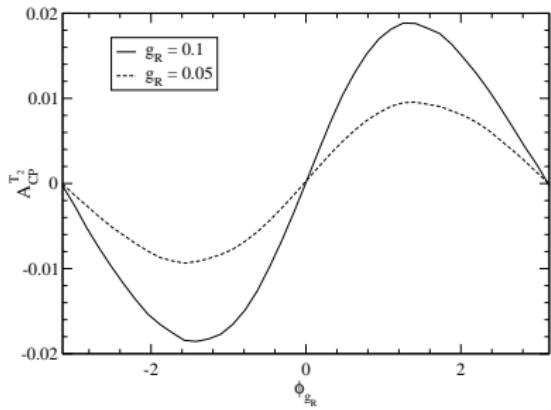
👉 $A_{CP}^{T_2} = \frac{N(T_2 > 0) - N(T_2 < 0)}{N(T_2 > 0) + N(T_2 < 0)}$

Momenta measured [in t, \bar{t} rest frames (larger A)
in laboratory frame (smaller A)

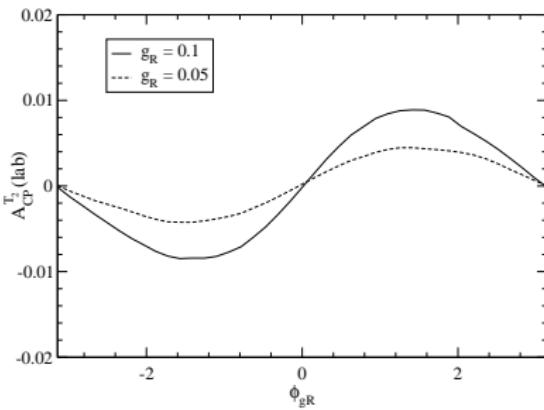
Sensitive to g_R but not to V_R and g_L

CP violation in top decays

t, \bar{t} rest frames



laboratory



Measurable?

Other asymmetries, e.g. based on products

$$\begin{aligned} T_1 &= \hat{e} \cdot (\vec{p}_{\ell+} - \vec{p}_{\ell-}) (\vec{p}_{\ell+} \times \vec{p}_{\ell-}) \cdot \hat{e} \\ T_3 &= (\vec{p}_t - \vec{p}_{\bar{t}}) \cdot (\vec{p}_{\ell+} \times \vec{p}_{\ell-}) \\ \Delta_1 &= E_{\ell+} - E_{\ell-} \\ \Delta_2 &= \vec{p}_{\bar{t}} \cdot \vec{p}_{\ell+} - \vec{p}_t \cdot \vec{p}_{\ell-} \\ \Delta_3 &= \cos \theta_{\ell+} - \cos \theta_{\ell-} \end{aligned}$$

Insensitive to CP-violating g_R , g_L , V_R



Test CP violation in production

[Schmidt, Peskin, PRL '92]

[Bernreuther et al., PLB '93]

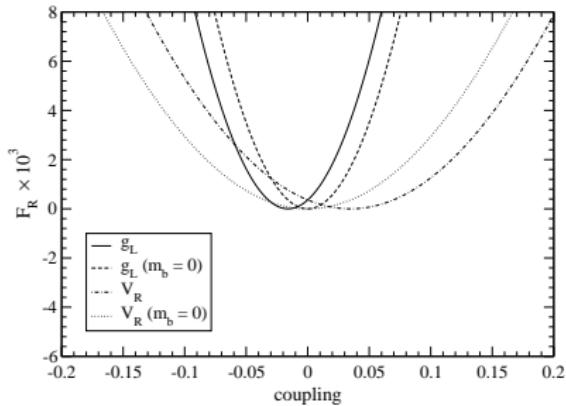
[Khater, Osland, NPB '03]

Summary

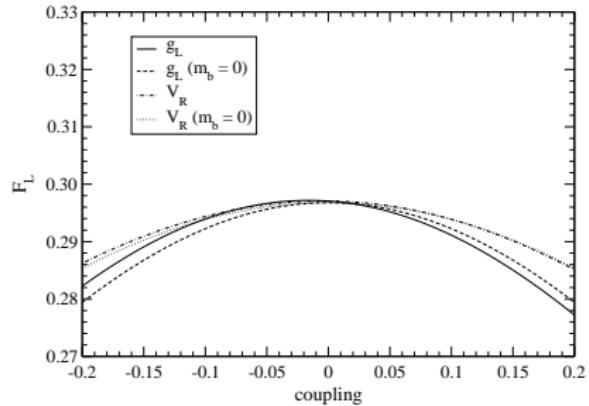
- Many ways to study Wtb vertex, some better than others
 - 👉 João's talk
- Constraining the Wtb vertex as far as possible can help disentangle new physics in production
 - (spin correlations, CP violation)
- There are lots of work still to do ...

Effect of m_b in helicity fractions

F_R

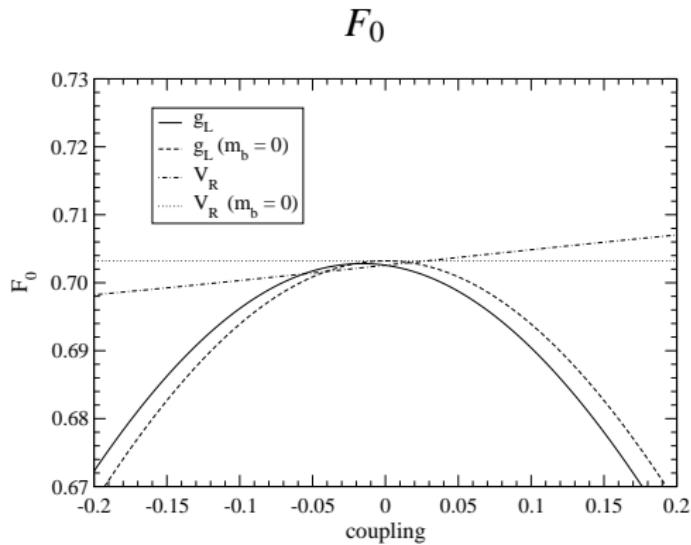


F_L



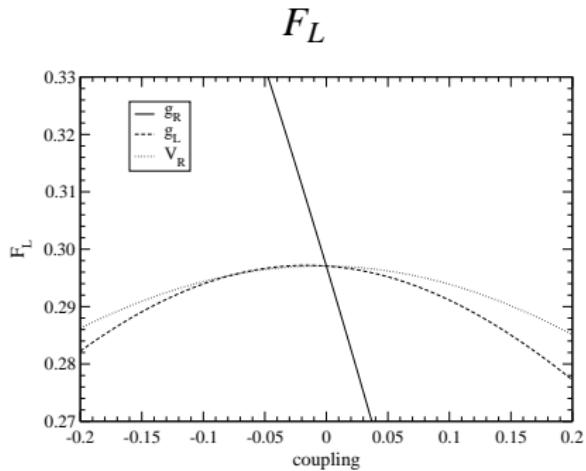
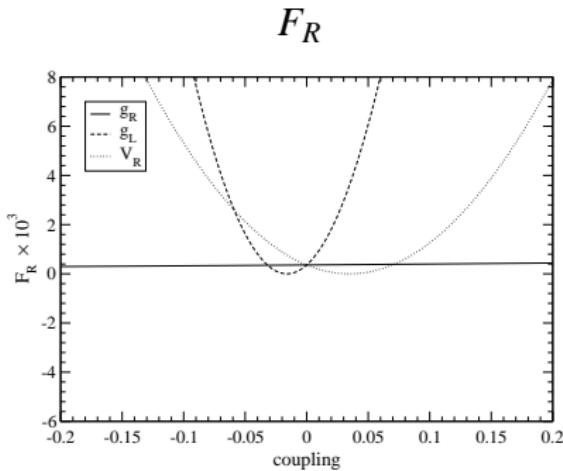
◀ Back

Effect of m_b in helicity fractions



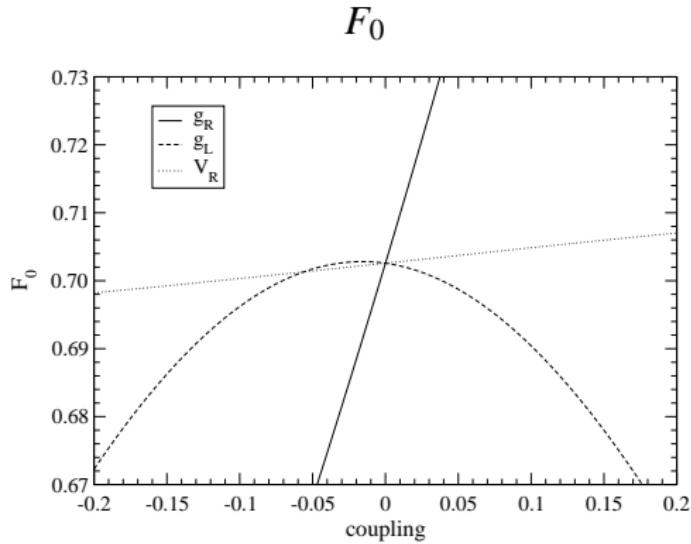
◀ Back

Dependence of F_i on anomalous couplings



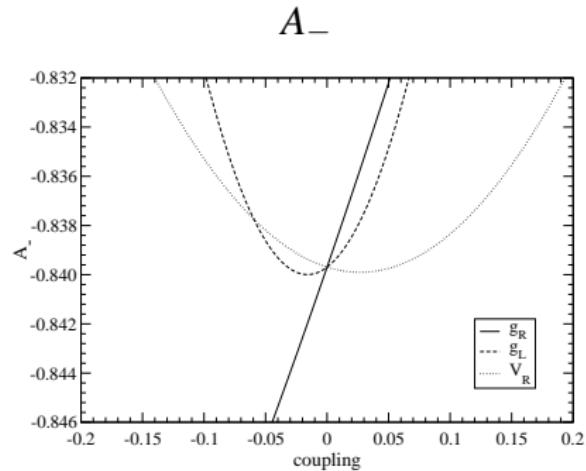
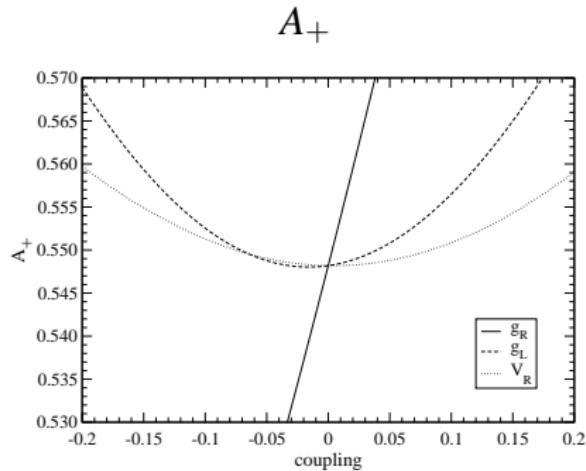
◀ Back ▶ Next

Dependence of F_i on anomalous couplings

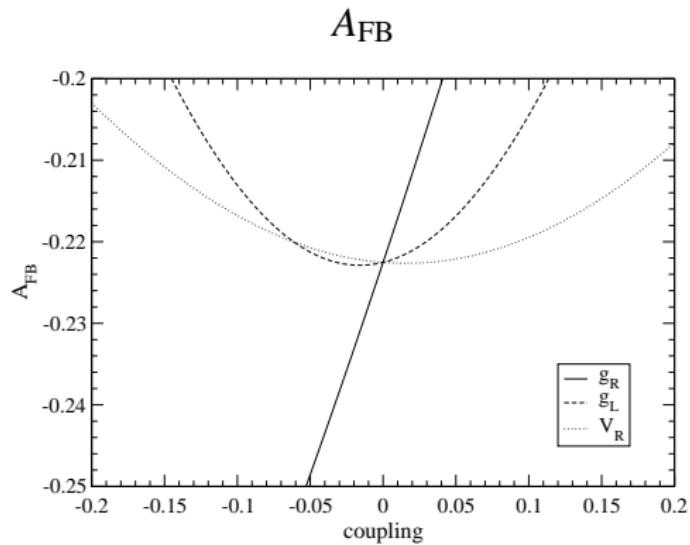


◀ Back

Dependence of asymmetries on anomalous couplings



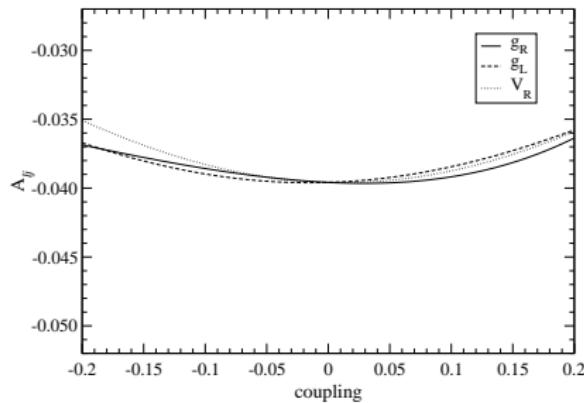
Dependence of asymmetries on anomalous couplings



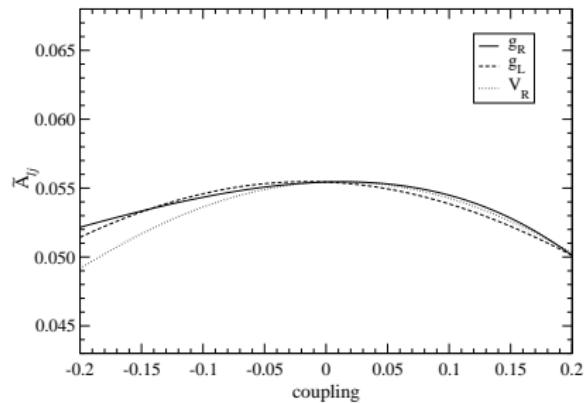
◀ Back

Dependence of spin asymmetries on anomalous couplings

$A_{\ell j}$

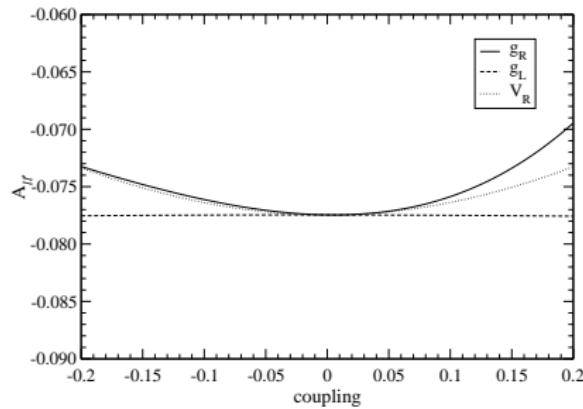


$\tilde{A}_{\ell j}$

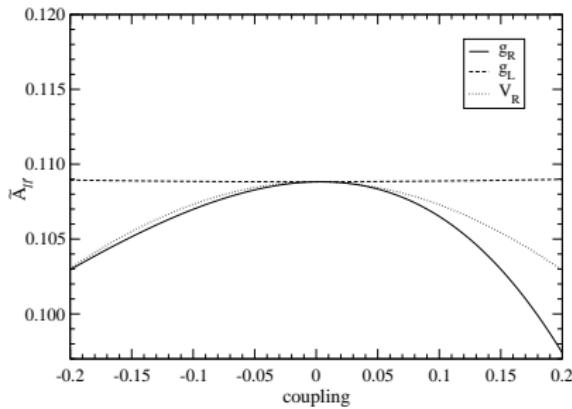


Dependence of spin asymmetries on anomalous couplings

$A_{\ell\ell'}$



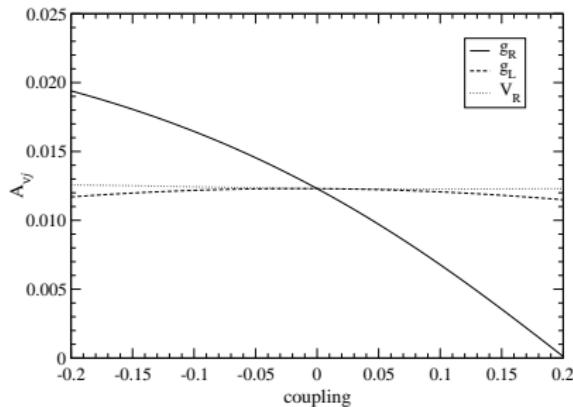
$\tilde{A}_{\ell\ell'}$



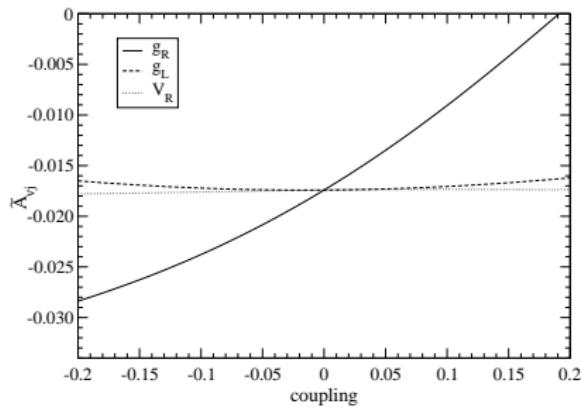
◀ Back ▶ Previous ▶ Next

Dependence of spin asymmetries on anomalous couplings

$A_{\nu j}$



$\tilde{A}_{\nu j}$



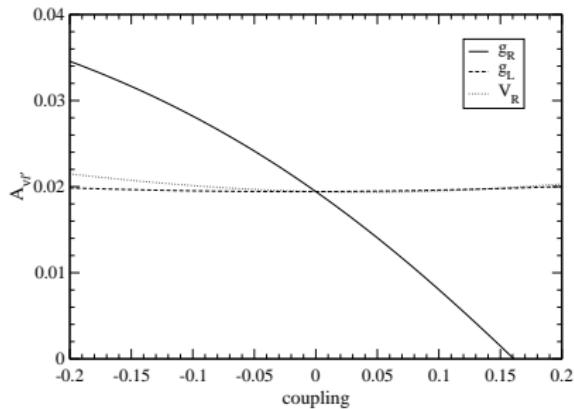
◀ Back

◀ Previous

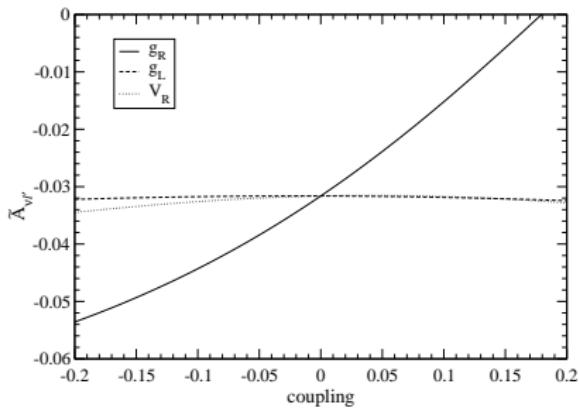
▶ Next

Dependence of spin asymmetries on anomalous couplings

$A_{\nu\ell'}$



$\tilde{A}_{\nu\ell'}$

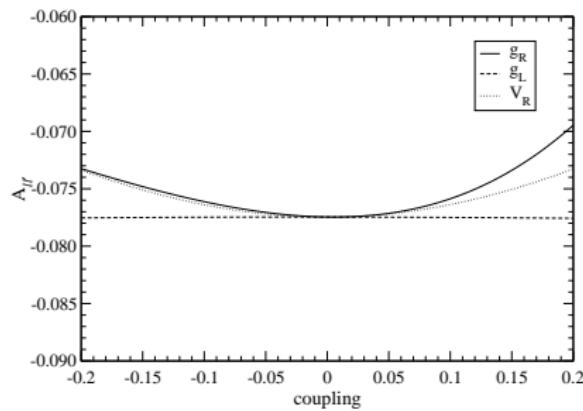


◀ Back

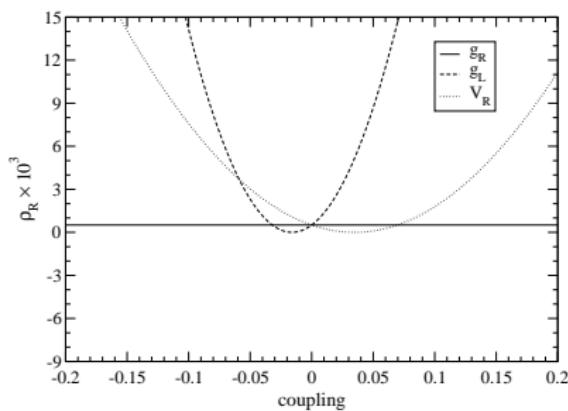
◀ Previous

Spin asymmetries: comparison

$A_{\ell\ell'}$



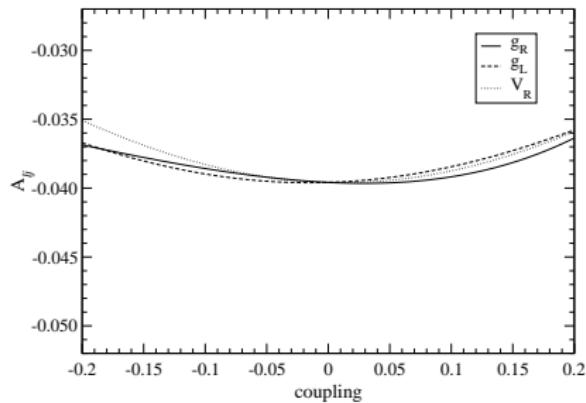
ρ_R



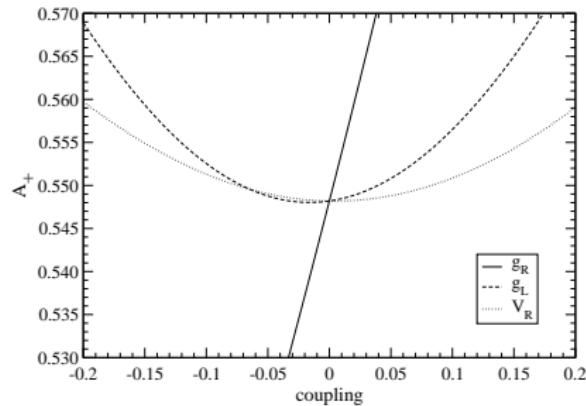
◀ Back ▶ Next

Spin asymmetries: comparison

$A_{\ell j}$



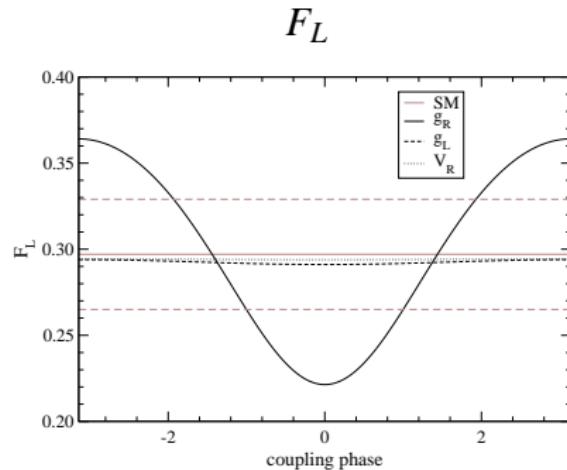
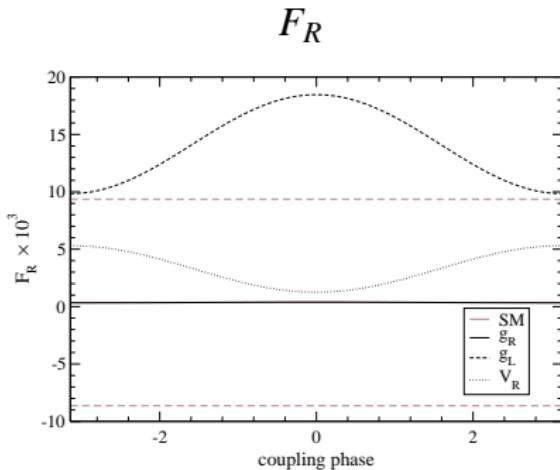
A_+



◀ Back

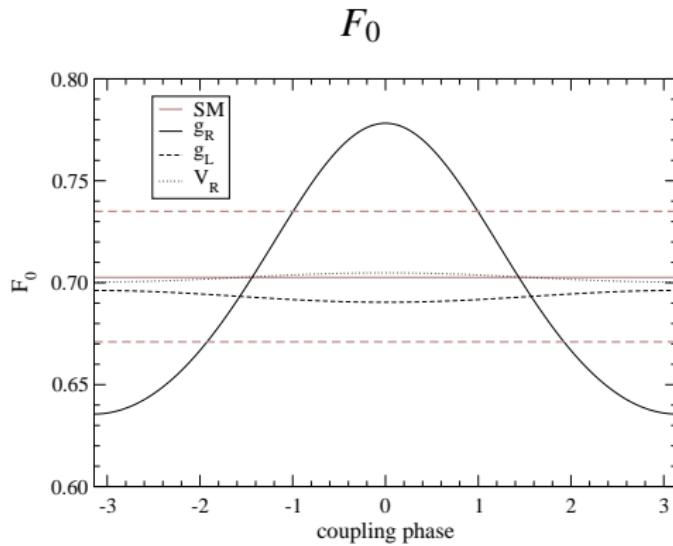
◀ Previous

Effect of phases in helicity fractions



◀ Back ▶ Next

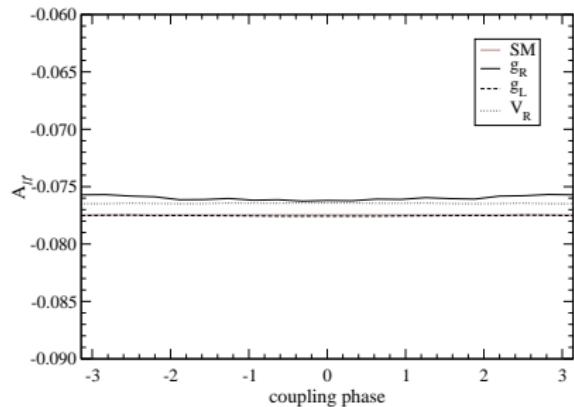
Effect of phases in helicity fractions



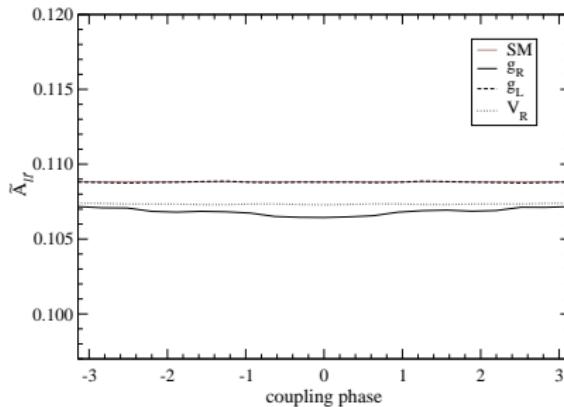
◀ Back

Effect of phases in spin asymmetries

$A_{\ell\ell'}$



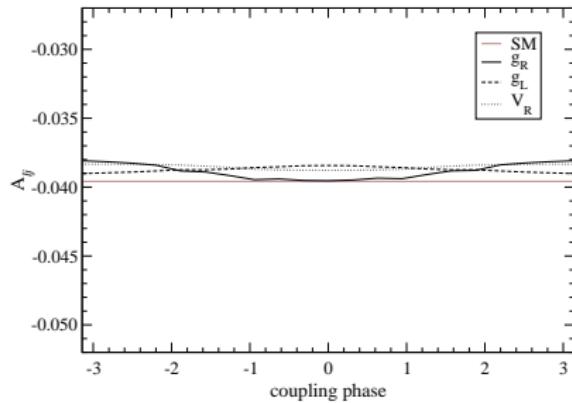
$\tilde{A}_{\ell\ell'}$



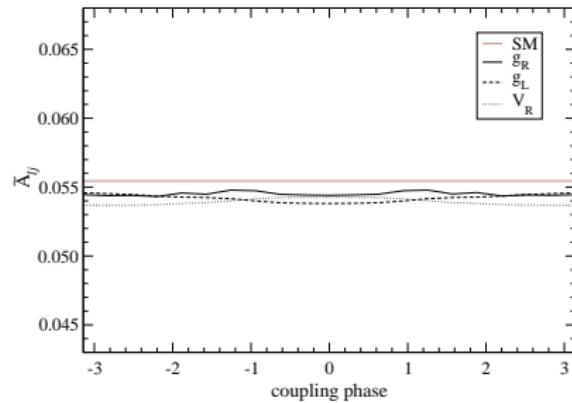
◀ Back ▶ Next

Effect of phases in spin asymmetries

$A_{\ell j}$



$\tilde{A}_{\ell j}$



◀ Back