

(SUB)eV SECTORS IN THE CMB AND LSS

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&

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& IFAE

based on

2211.03799 w/ A. Notari, G. Villadoro

2305.14166 w/ I. Allali, M. Hertzberg



New Physics from Galaxy Clustering II, IFPU Trieste

FREE STREAMING SECTORS

See talks by
Marilena, Francis

For recombination, if

$$m_\phi \ll 0.1 \text{ eV}$$

ϕ

Generic light species

If no interactions, new physics behaves like
free-streaming "dark radiation" (DR) at CMB epoch

Effective
number of
neutrino
species

$$N_{\text{eff}} = 3.044 + \Delta N_{\text{eff}} \quad \Delta N_{\text{eff}} \equiv \frac{\rho_{\text{DR}}}{\rho_\nu} \Big|_{\text{rec}} = \frac{8}{7} \left(\frac{11}{4} \right)^{\frac{4}{3}} \frac{\rho_{\text{DR}}}{\rho_\gamma} \Big|_{\text{rec}}$$

Change in the expansion rate
(CMB damping tail, phase of BAOs)

HOT DARK MATTER/COLD DARK RADIATION

See talks by
Marilena, Francis

For recombination, if

$$m_\phi \gtrsim (0.01 - 0.1) \text{ eV}$$

Really just like neutrinos, but colder

e.g. for a scalar dof

$$T_\phi = T_\nu \left(\frac{7}{4} \right)^{\frac{1}{4}} \Delta N_{\text{eff}}$$

Become non-relativistic at

$$z_{\text{nr}} \simeq 772 \left(\frac{m_\phi}{0.3 \text{ eV}} \right) \left(\frac{0.3}{\Delta N_{\text{eff}}} \right)^{\frac{1}{4}} - 1$$

See also
Lesgourgues,
Mangano, Miele,
Pastor , CUP 13

COLD DARK RADIATION / HOT DARK MATTER

See talks by
Marilena, Francis

Really just like neutrinos, but colder

e.g. for a scalar dof

If during matter
domination

$$k_{\text{nr}} \simeq 0.01 \sqrt{\Omega_m} \left(\frac{0.3}{\Delta N_{\text{eff}}} \right)^{\frac{1}{8}} \left(\frac{m}{0.3 \text{ eV}} \right)^{\frac{1}{2}} h \text{ Mpc}^{-1}$$

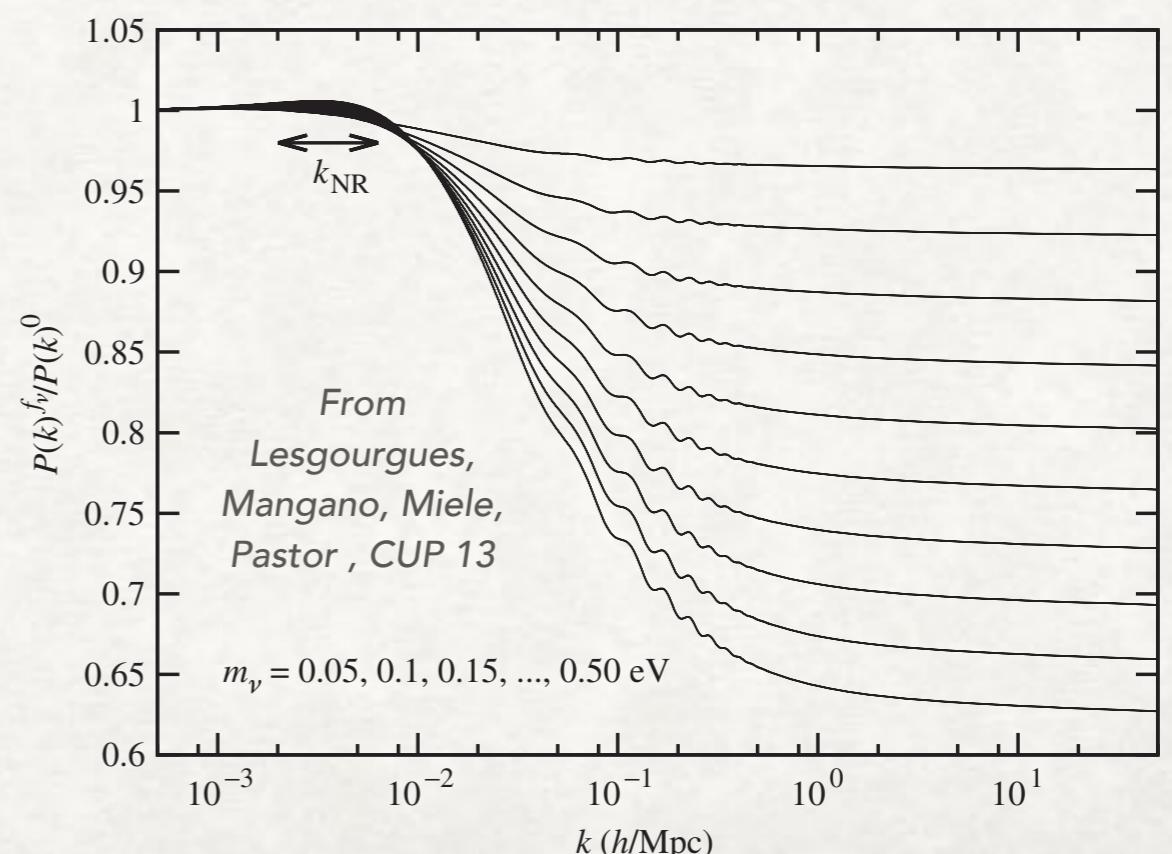
*Suppression of linear matter
power spectrum*

$$\frac{\Delta P_\phi}{P} \propto m_\phi \Delta N_{\text{eff}}^{\frac{3}{4}}$$

Stronger constraints on ΔN_{eff}
than in massless case, for

$$m_\phi \gtrsim 0.1 \text{ eV}$$

See also Xu+ 21
for heavier sector

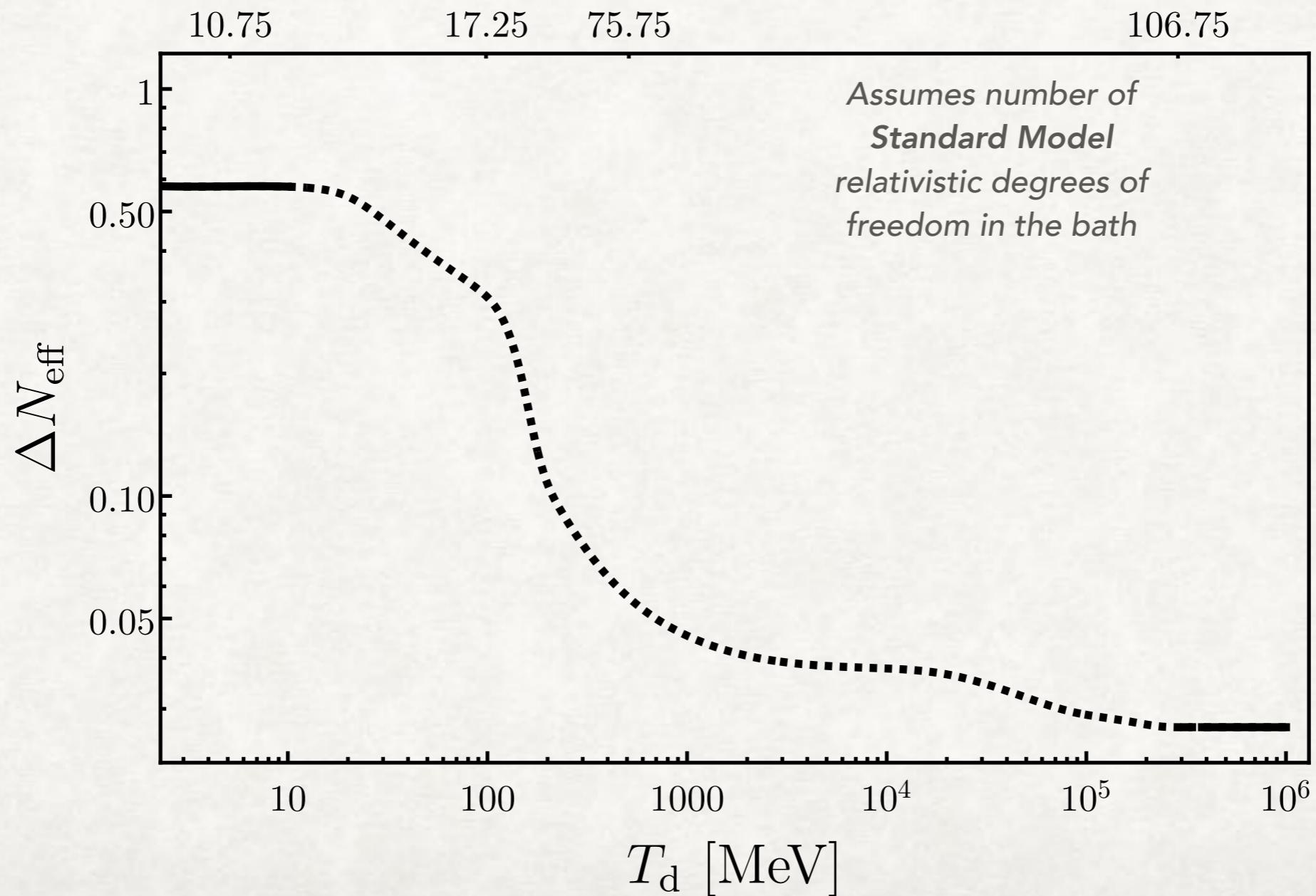


THERMAL PRODUCTION

*For instantaneous
decoupling from equilibrium*

$$\Delta N_{\text{eff}} \simeq 0.3 \left[\frac{g_\star(100 \text{ MeV})}{g_\star(T_d)} \right]^{\frac{4}{3}}$$

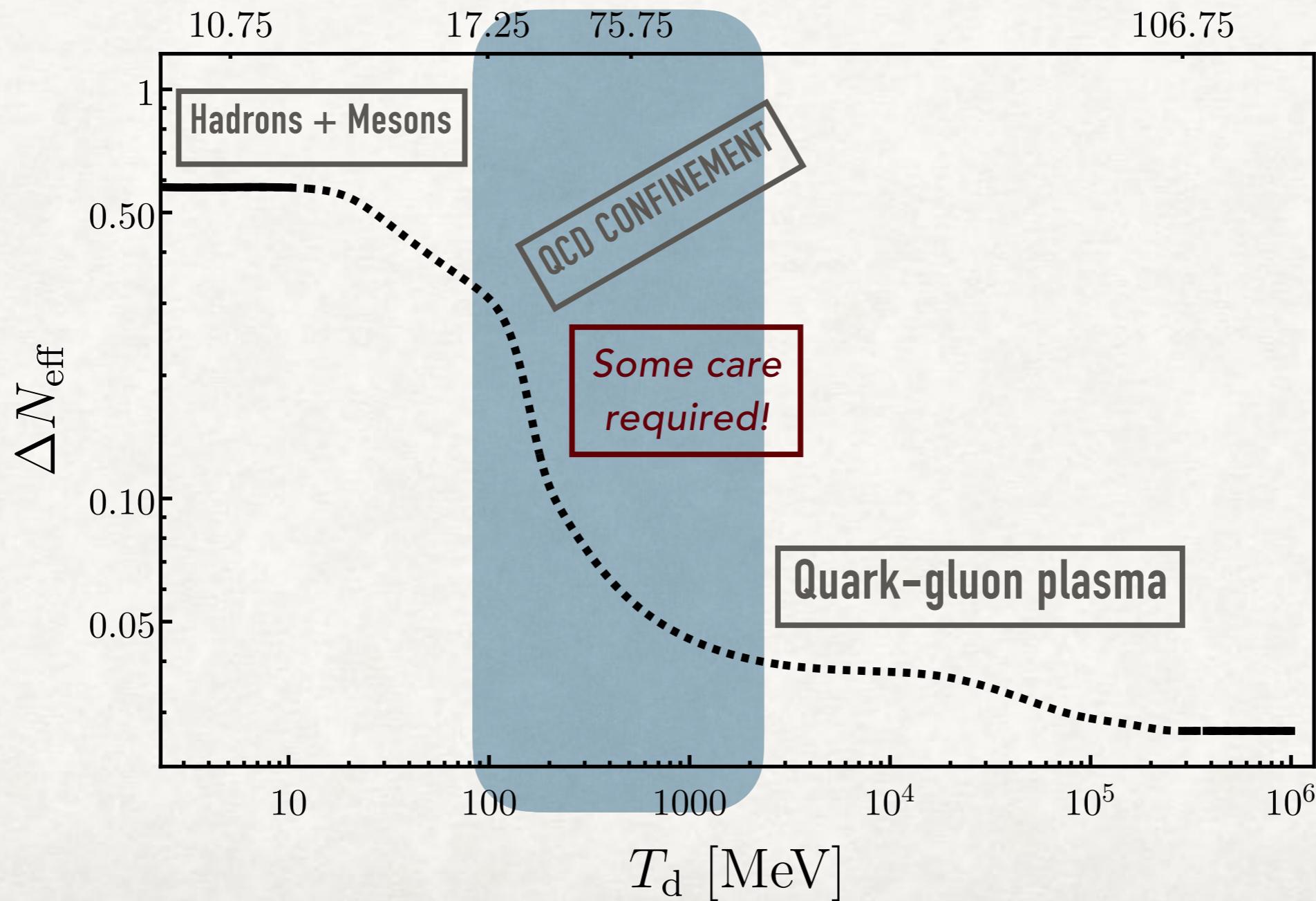
g_\star



FREE STREAMING DARK RADIATION

*For instantaneous
decoupling from equilibrium*

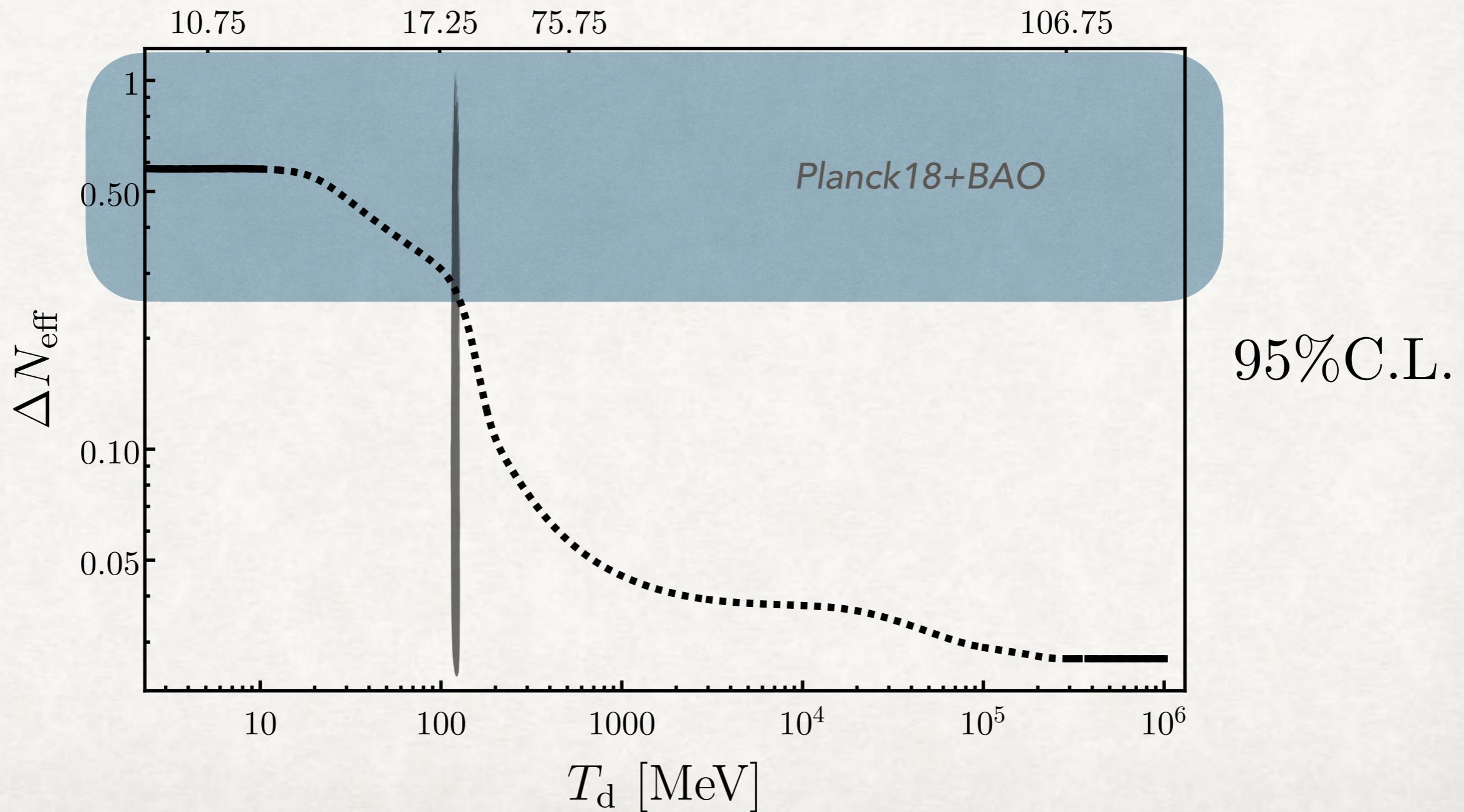
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FREE STREAMING DARK RADIATION

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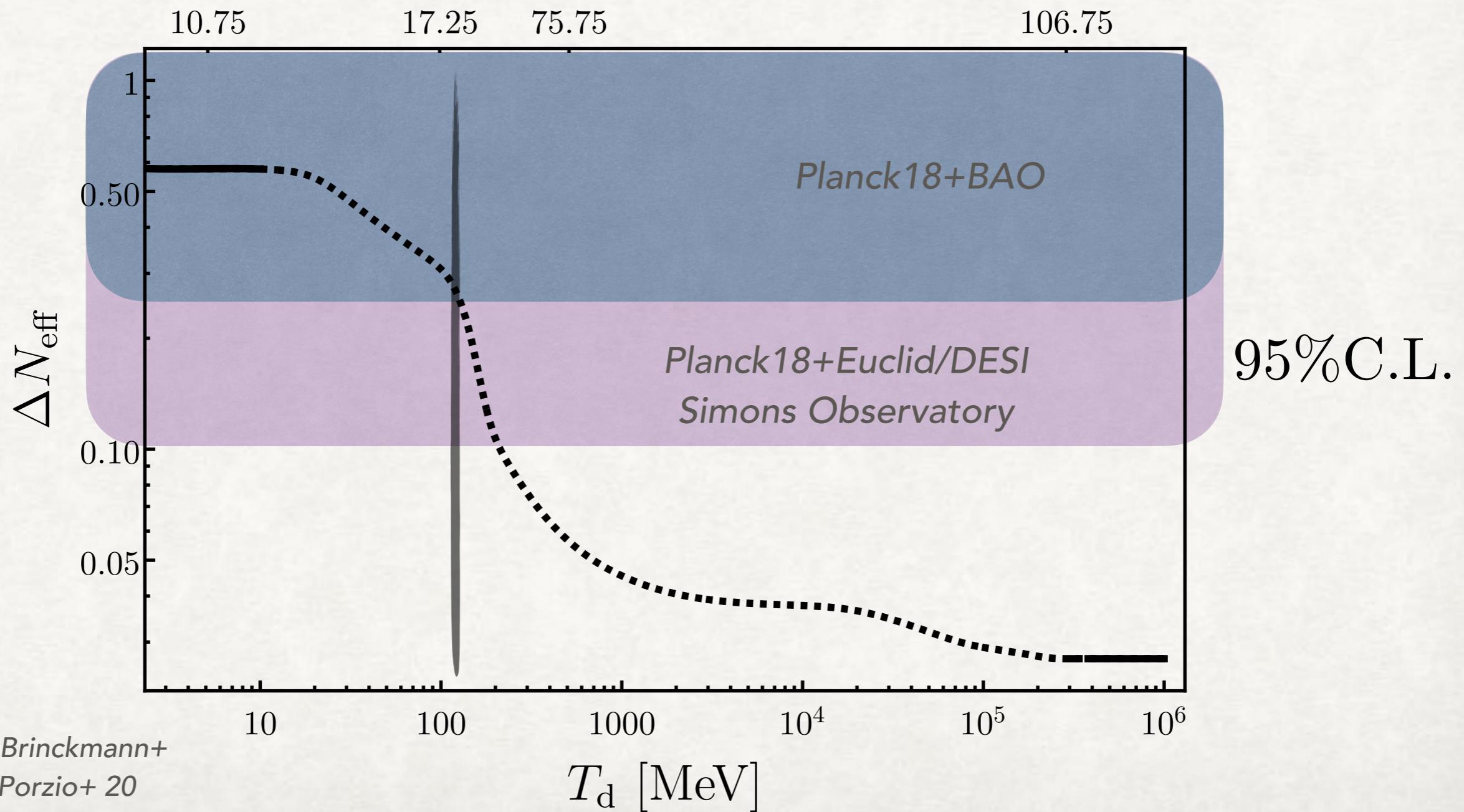
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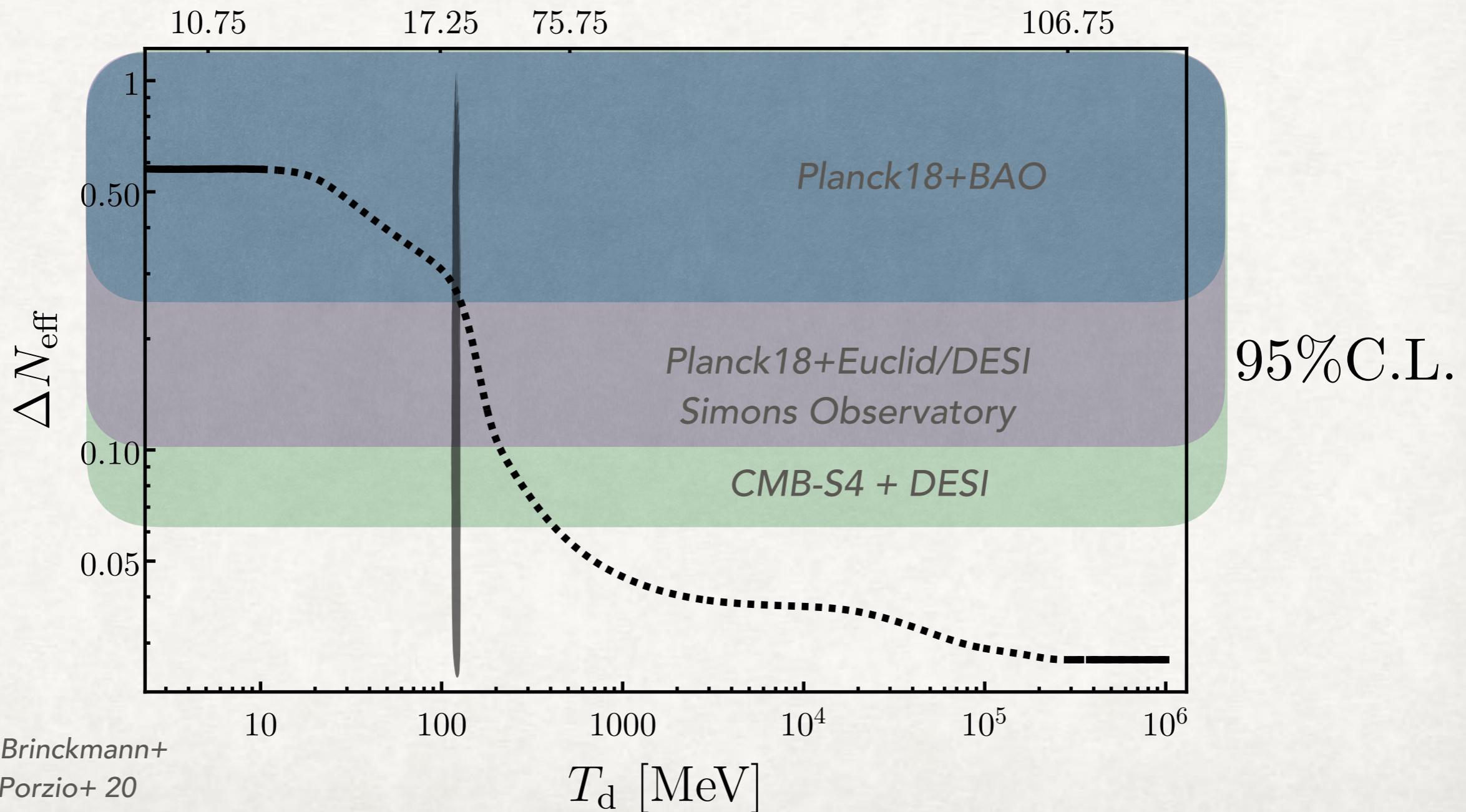
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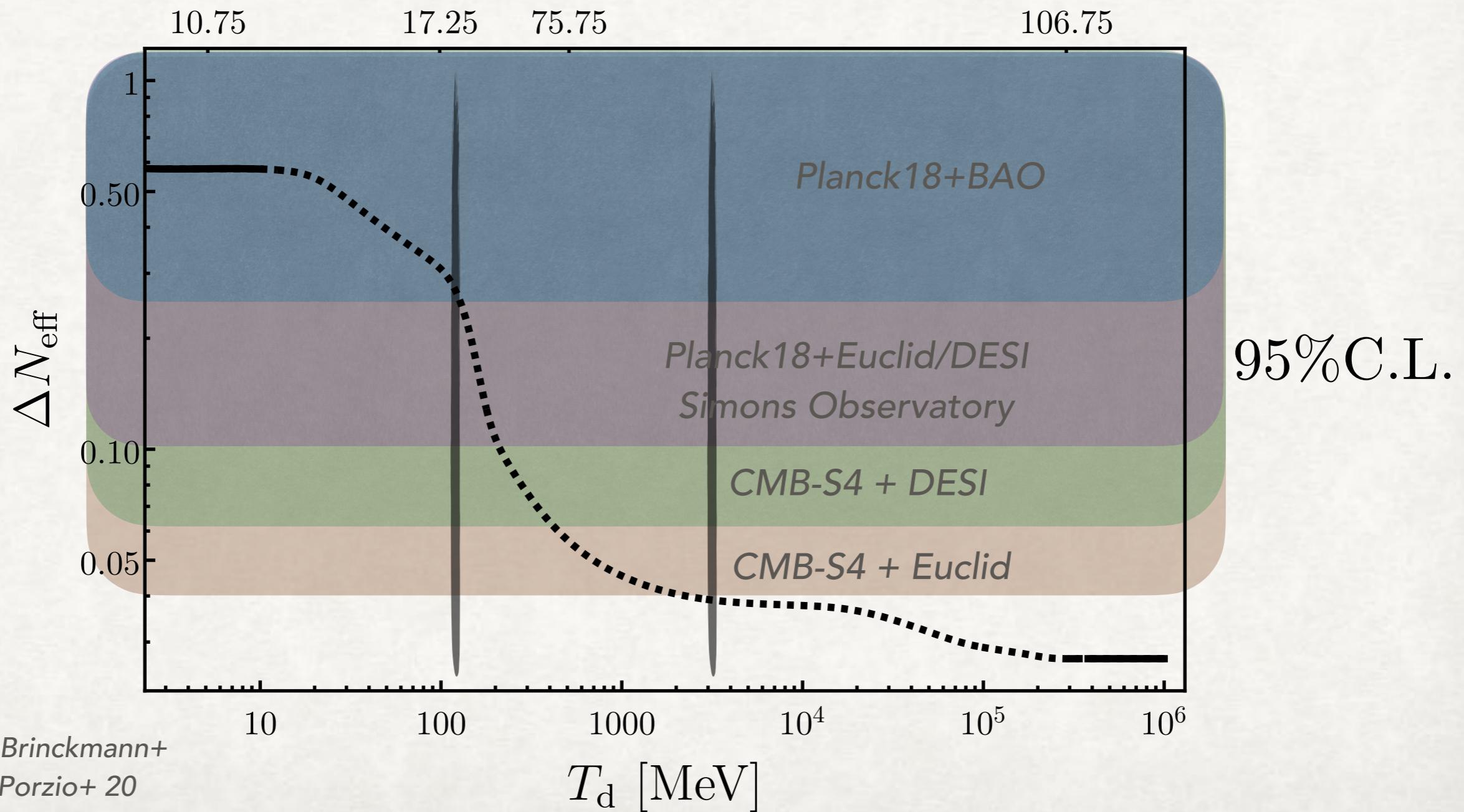
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FREE STREAMING DARK RADIATION

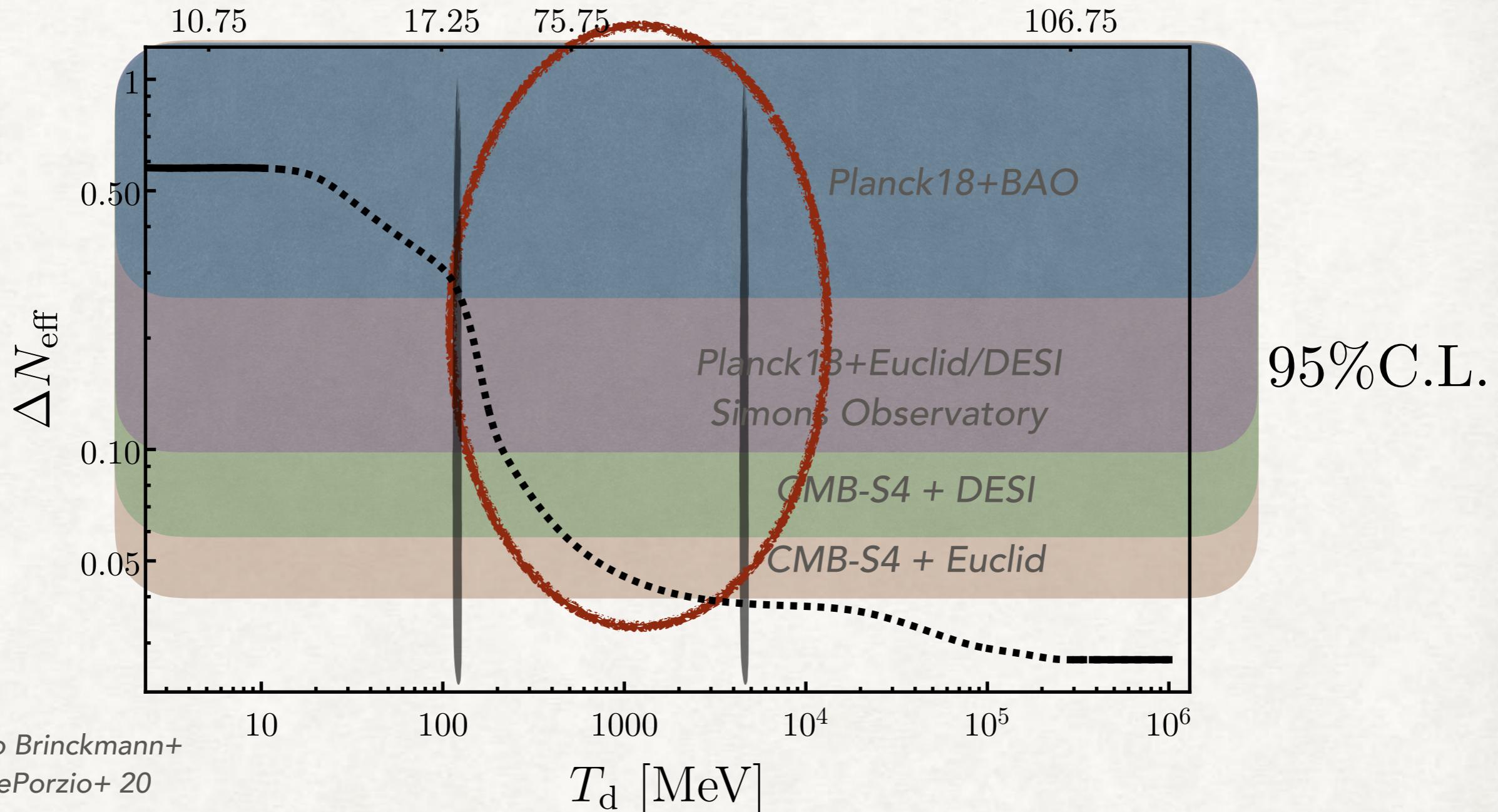
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FREE STREAMING DARK RADIATION

g_\star



Near future CMB and LSS surveys probe light thermal sectors
that decouple during the QCD crossover transitions!

PARTICLE PHYSICS AT THE eV?!

Want: (a) light particle(s) (with (sub)-eV mass) that decouples around the QCD epoch

UV-MOTIVATED

TENSIONS

SIGNAL BUILDING

PARTICLE PHYSICS AT THE EV?

Want: (a) light particle(s) (with (sub)-eV mass) that decouples around the QCD epoch

UV-MOTIVATED

2211.03799 w/ A. Notari, G. Villadoro

TENSIONS

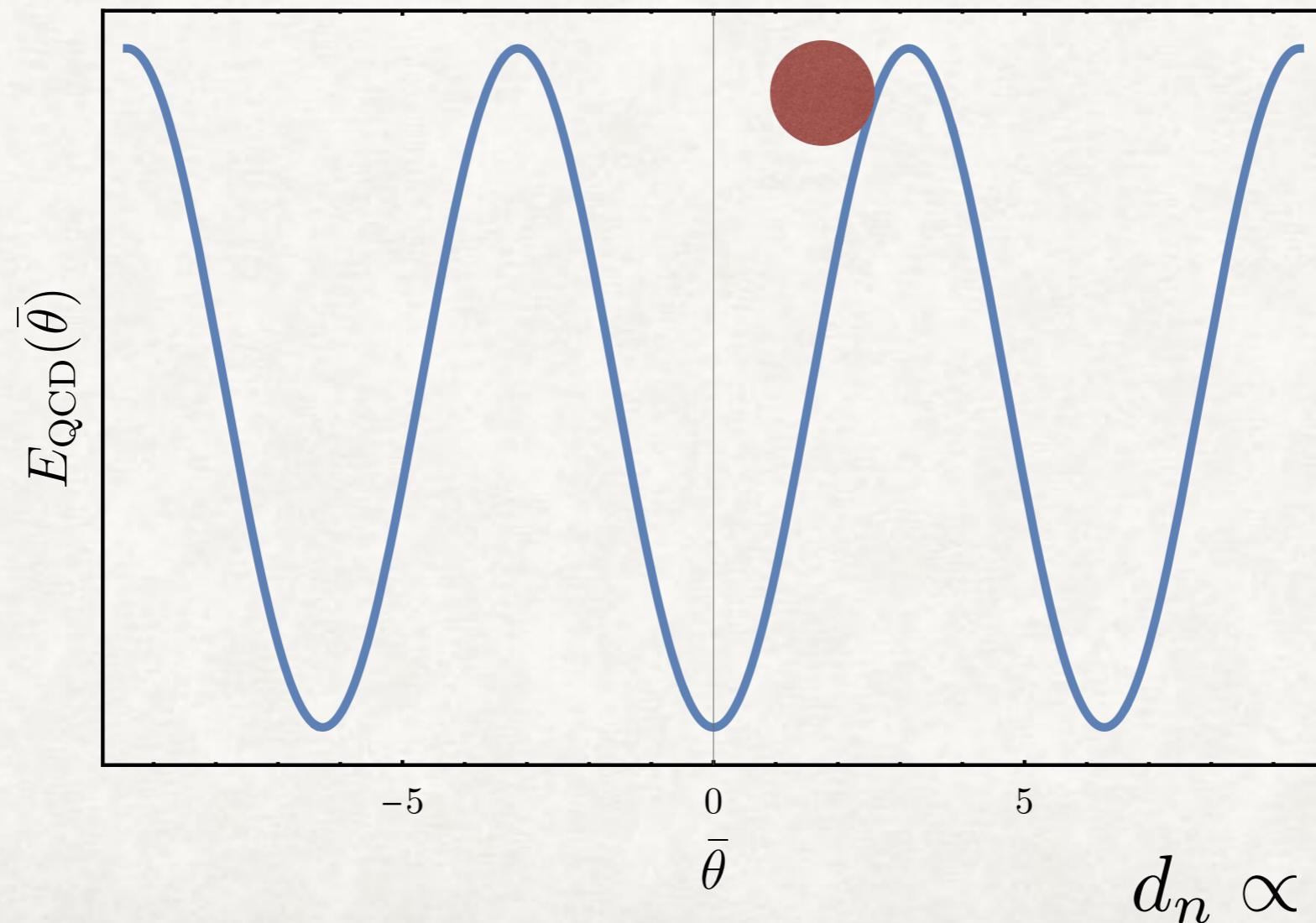
SIGNAL BUILDING

Decoupling temperature suggests coupling to QCD

THE USUAL SUSPECT

The QCD axion

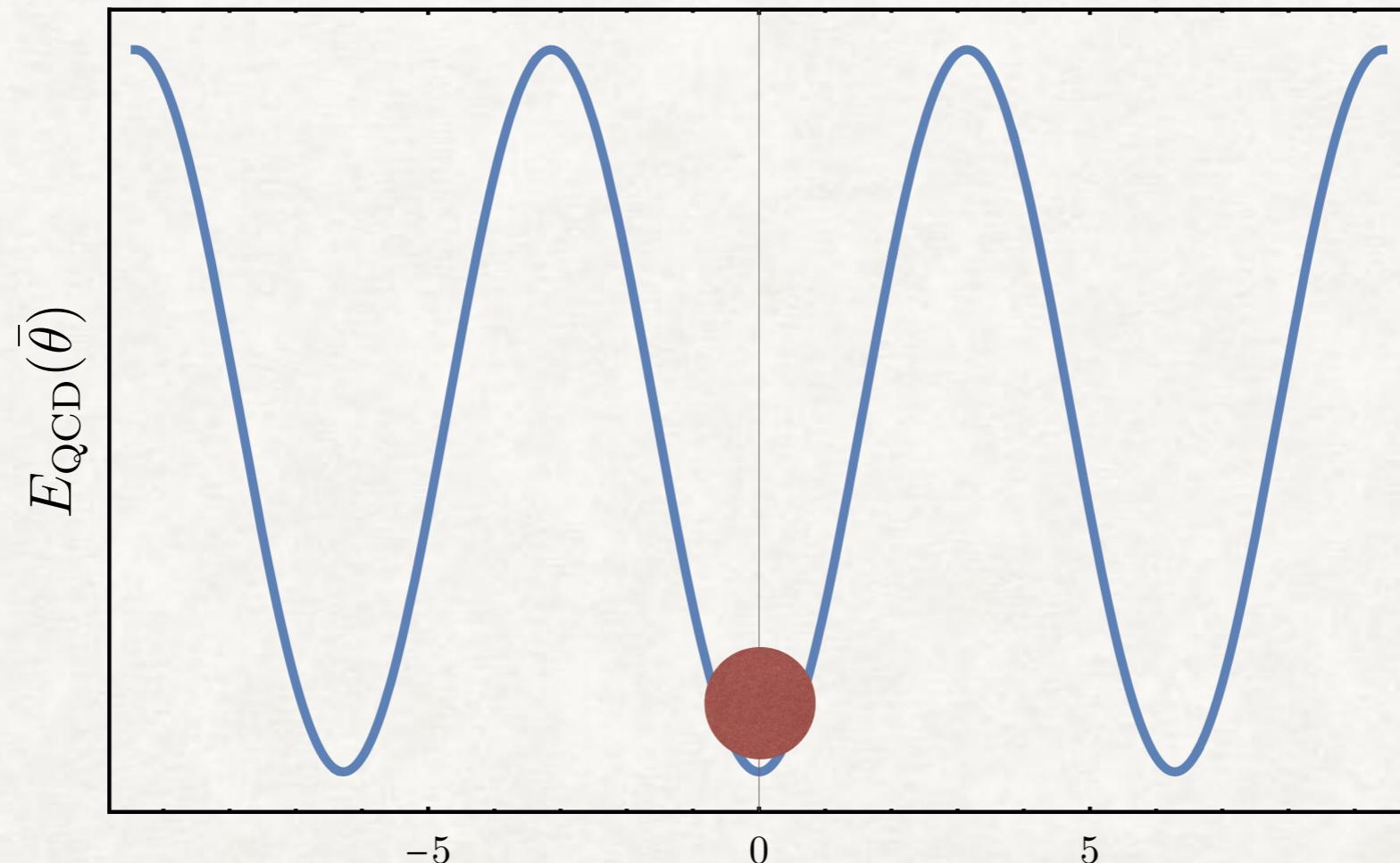
$$\mathcal{L}_{\text{QCD}} \supset \bar{\theta} \frac{\alpha_s}{8\pi} G^{\mu\nu} \tilde{G}_{\mu\nu}, \quad \bar{\theta} \rightarrow \frac{a(t, \mathbf{x})}{f_a}$$



THE USUAL SUSPECT

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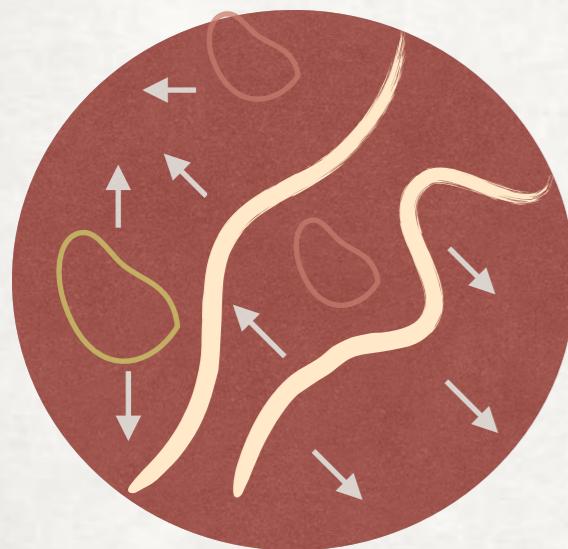
$$d_n \propto \bar{\theta} \rightarrow 0$$

THE QCD AXION

$$\mathcal{L}_a = \frac{1}{2}(\partial_\mu a)^2 + \frac{\alpha_s}{8\pi} \frac{a(t, \mathbf{x})}{f_a} G^{\mu\nu} \tilde{G}_{\mu\nu} + \dots$$

*Small mass set by interplay between
UV (Peccei-Quinn) and IR (QCD) scale*

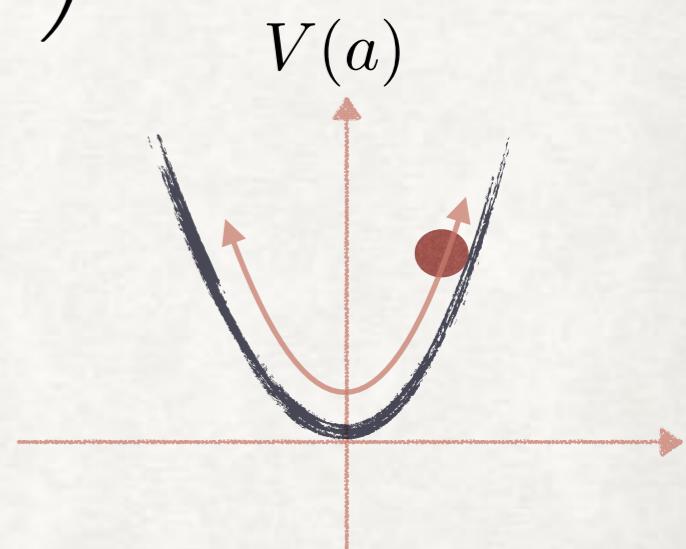
Decay of
topological
defects



$$m_a = \kappa_q \frac{\Lambda_{\text{QCD}}^2}{f_a} \simeq 0.057 \text{ eV} \left(\frac{10^8 \text{ GeV}}{f_a} \right)$$

*Mostly discussed as CDM candidate,
From non-thermally produced population*

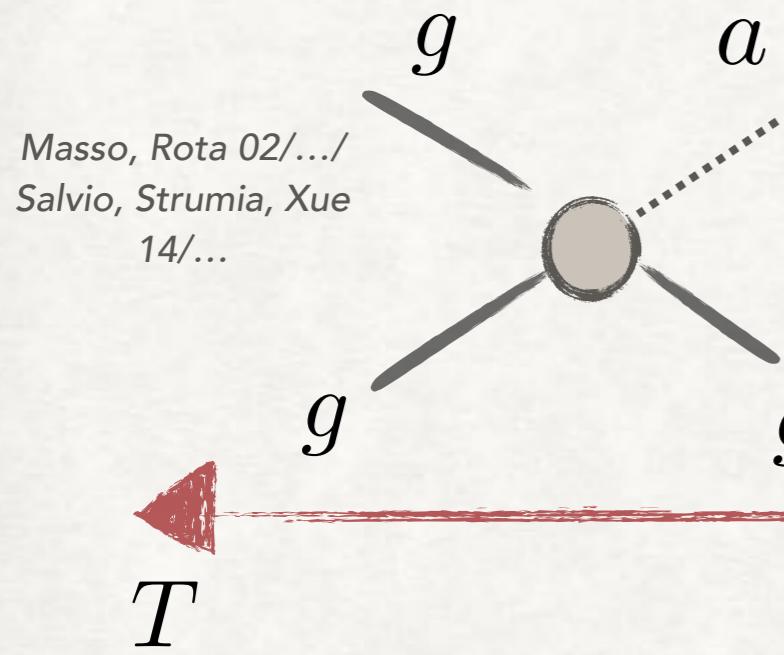
Misalignment
mechanism



But coupling to QCD implies thermal production!

Turner 88/
Berezhiani et al 92/

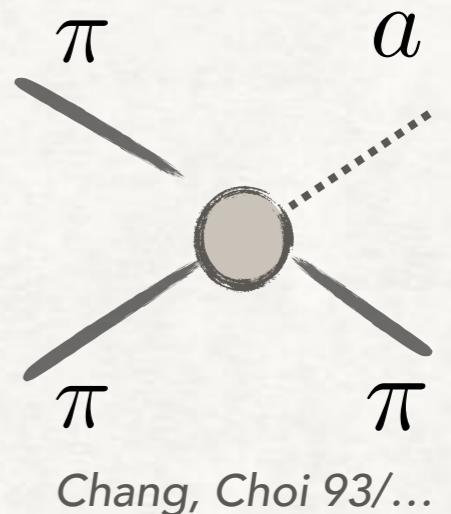
THE HOT QCD AXION: PRODUCTION



At weak coupling,
production dominated by
2-2 processes

QCD CROSSOVER

$$\mathcal{L}_{a\pi} = \frac{\epsilon}{3f_a f_\pi} \partial_\mu a (2\partial^\mu \pi^0 \pi^+ \pi^- + \dots)$$



$$\bar{\Gamma}_a \propto \frac{T^3}{f_a^2}$$

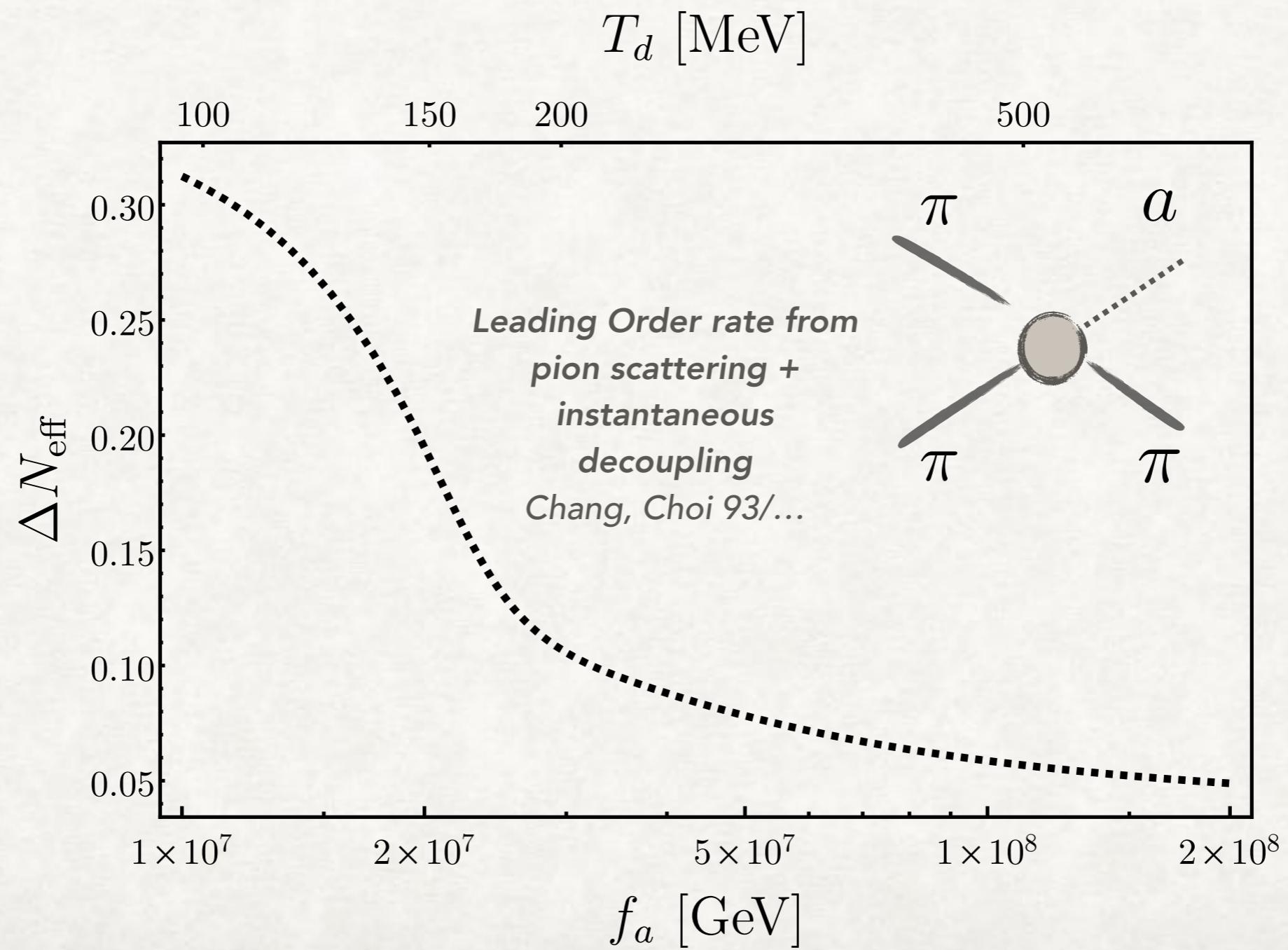
$T \sim \text{GeV}$

$T_{\text{QCD}} \simeq 150 \text{ MeV}$

$$\bar{\Gamma}_a \propto \frac{T^5}{f_a^2 f_\pi^2}$$

FIRST ESTIMATE

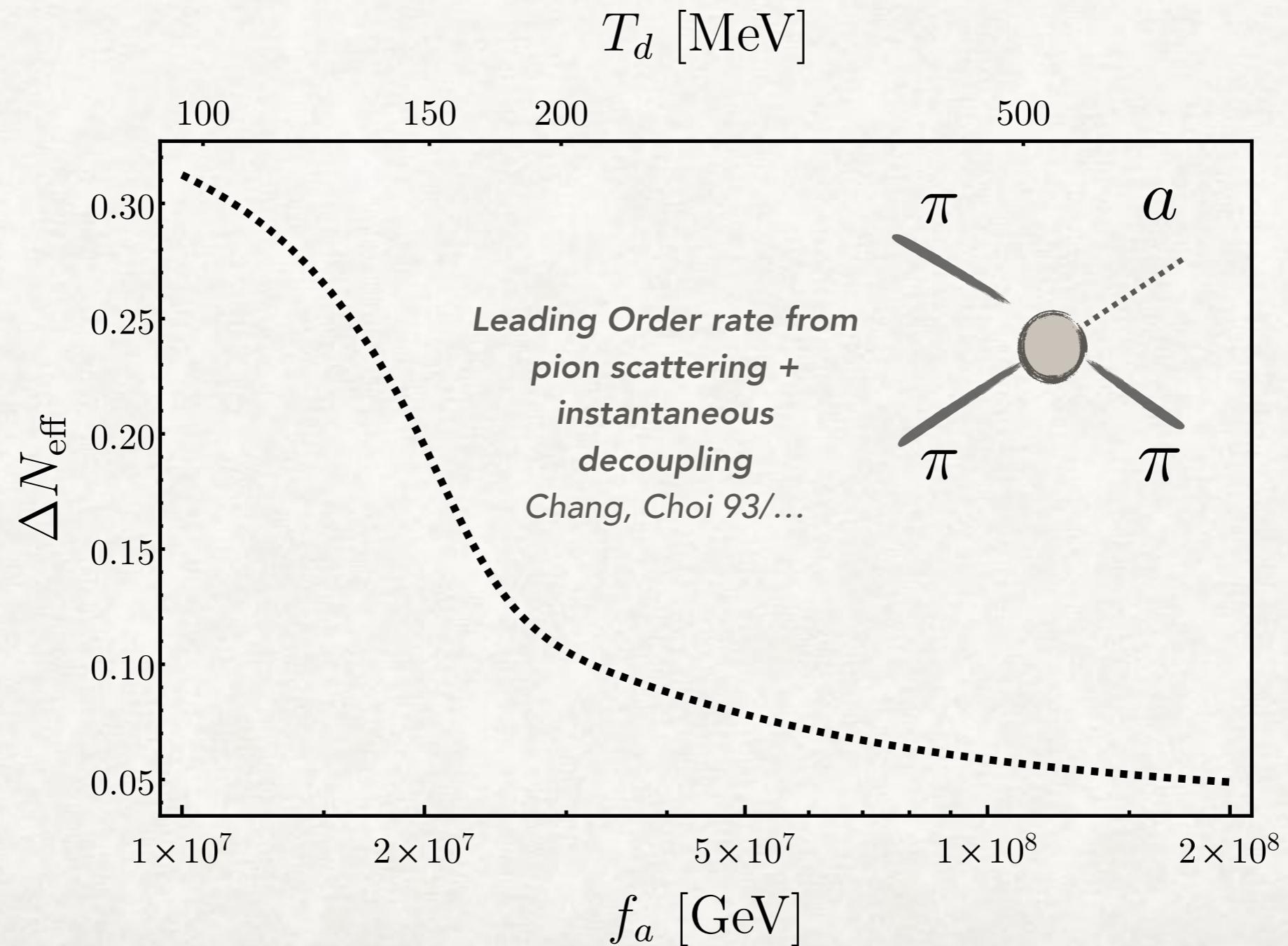
*For instantaneous
decoupling from equilibrium*



FIRST ESTIMATE

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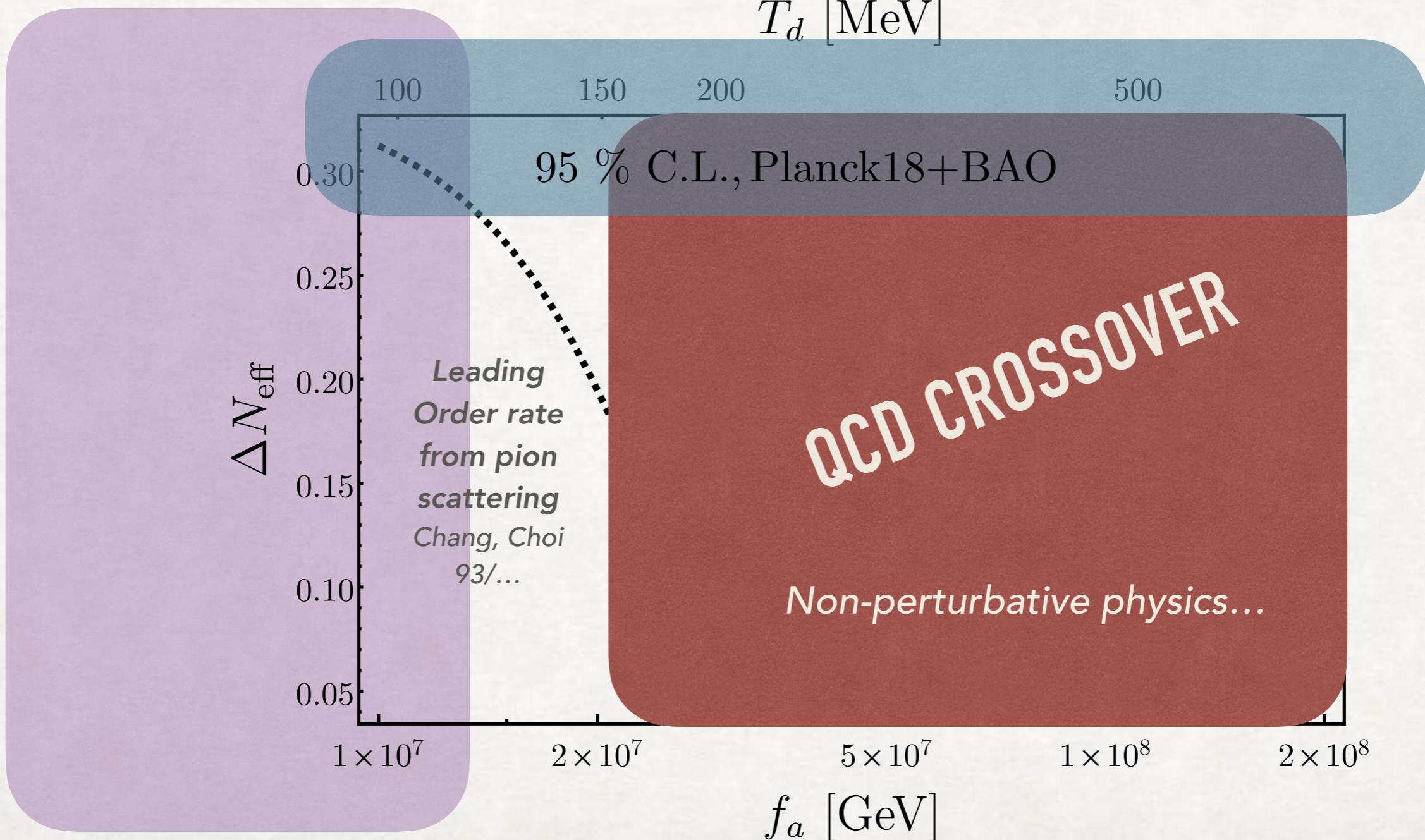
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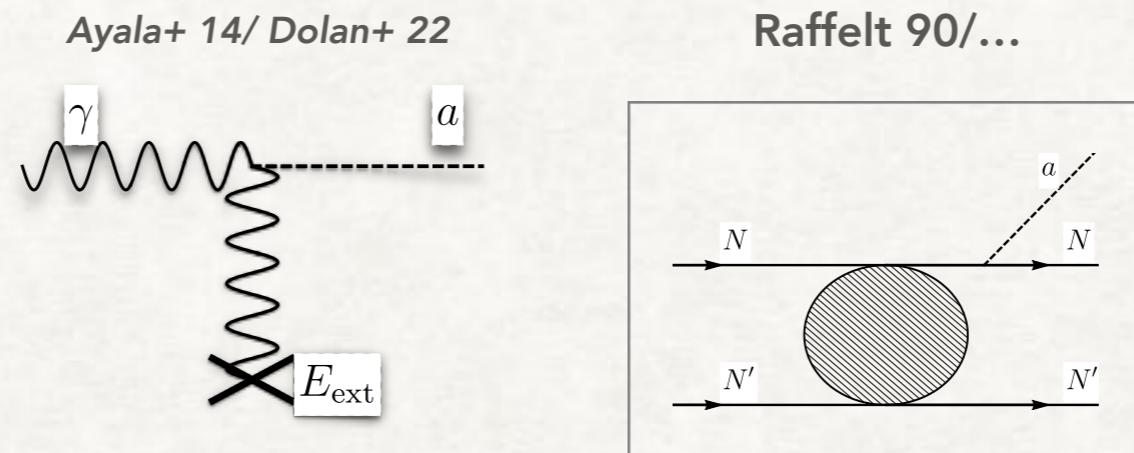


...YES, THERE ARE ASTROPHYSICAL BOUNDS

From production in stellar environments

$$f_a \gtrsim (10^7 - 10^8) \text{ GeV}$$

Stronger model-dependent bounds when coupling to electrons is present



From di Luzio+ Phys. Rep. 20

Strongest bounds come from SN1987A, possibly characterised by significant uncertainties, currently under active reassessments/debate

See e.g. Chang+ 18/Bar+19/Carenza+20

OPPORTUNITY FOR COSMOLOGICAL SURVEYS

*CMB and LSS data provide opportunity to
discover/constrain the QCD axion*

KEY ADVANTAGES

Uses precise cosmological measurements, independent of astrophysics

Probes QCD axion independently of cold dark matter contribution

Relies on minimal coupling to QCD

+

(Standard cosmology below GeV)

Hannestad+ 08, 13/Di
Valentino+ 15/Ferreira, Notari
18/+ Arias-Aragon, D'Eramo
et al 18,20.../Ferreira, Notari,
FR 20/ Giaré+ 20/Di
Luzio+21/D'Eramo+21,22/Di
Luzio+22

Does not need dedicated experiment!

OPPORTUNITY FOR COSMOLOGICAL SURVEYS

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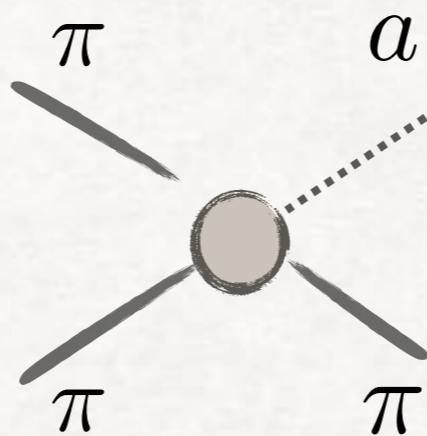
KEY CHALLENGES/DIRECTIONS

Particle physics: Obtain reliable production rate

*Cosmology: Find ways to distinguish axion
from other possible dark radiation/hot DM*

PRODUCTION FROM PIONS

Below 150 MeV



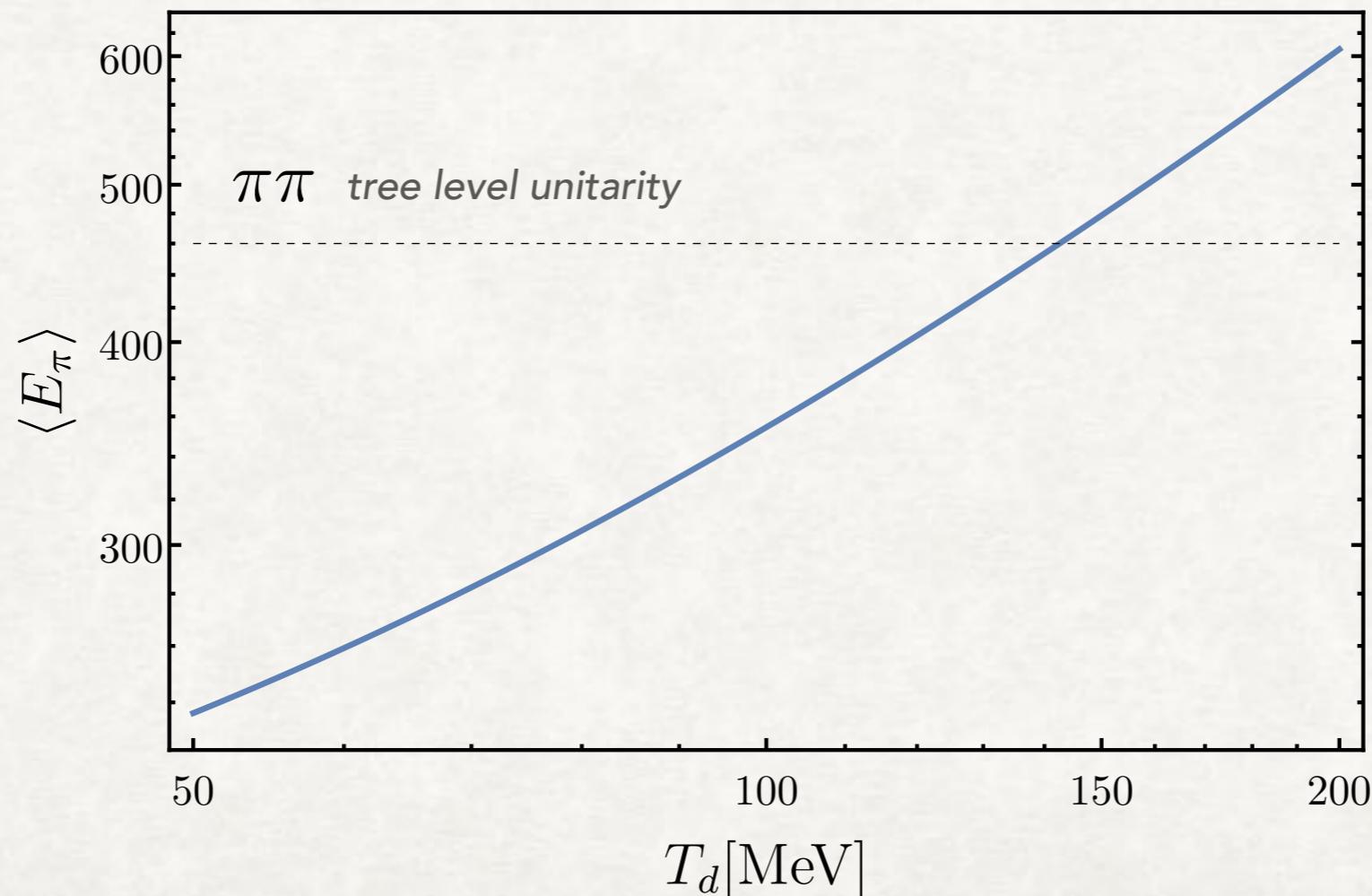
*Two problems with simple estimate
both due to decoupling during the QCD crossover*

*Features discussed here possibly relevant for
other candidates that can be seen in the near
future at CMB/LSS*

CALCULATION OF THE PRODUCTION RATE FROM PIONS

Simple estimate does not work because

1. *Leading order (in chiral perturbation theory) rate is not valid at the temperatures of interest, because:*



*Can be (partially)
overcome by
relating axion
rate to pion rate,
extracted from
experiment*

*See also Di Luzio+22 for
unitarization approach*

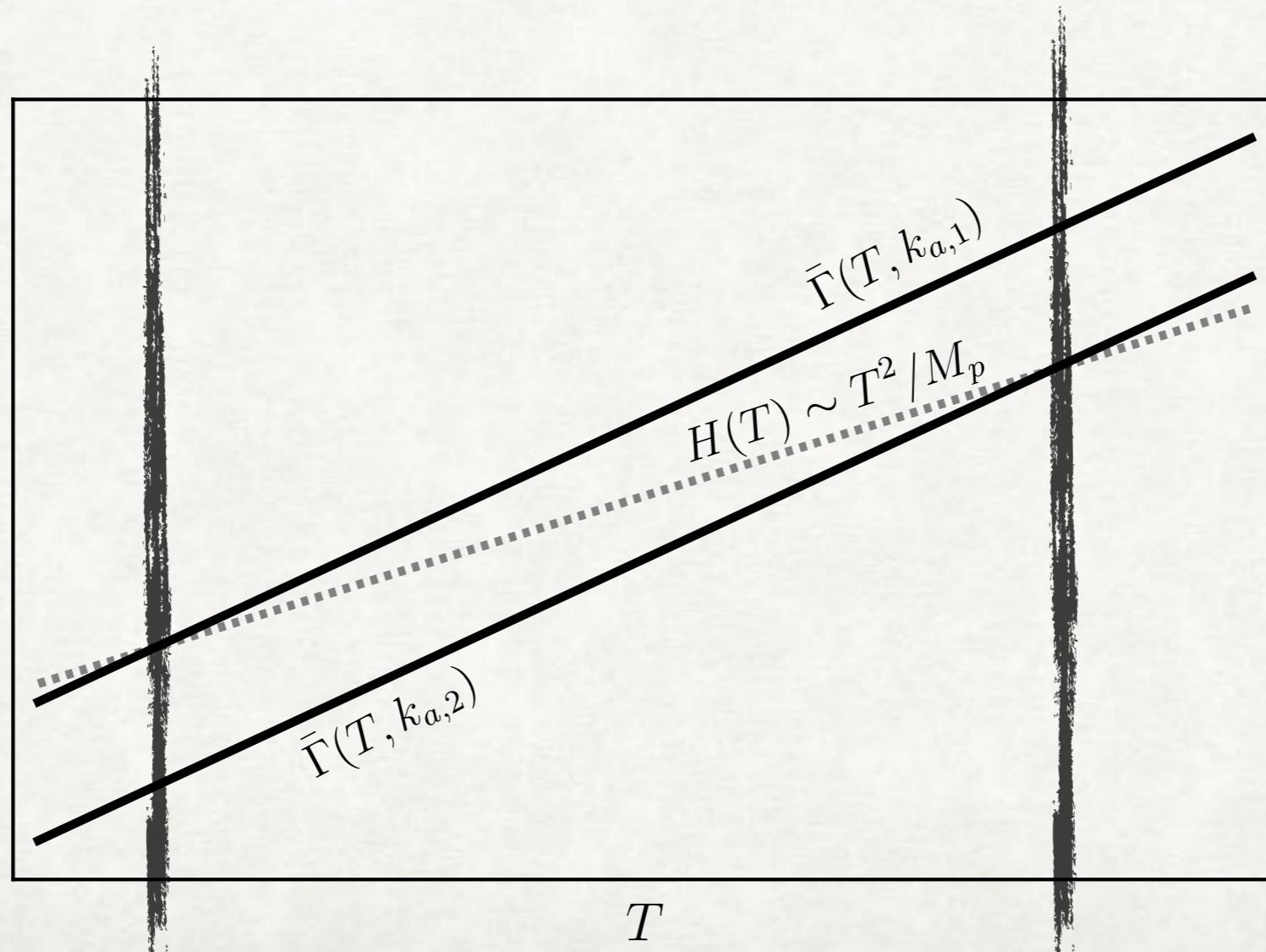
CALCULATION OF THE ABUNDANCE

Simple estimate does not work because

2. Instantaneous decoupling + thermal equilibrium not justified

$$\mathcal{M}_{\pi\pi \rightarrow a\pi} \propto s^2 \quad \Rightarrow \quad \bar{\Gamma}_a(T) \rightarrow \bar{\Gamma}_a(k, T)$$

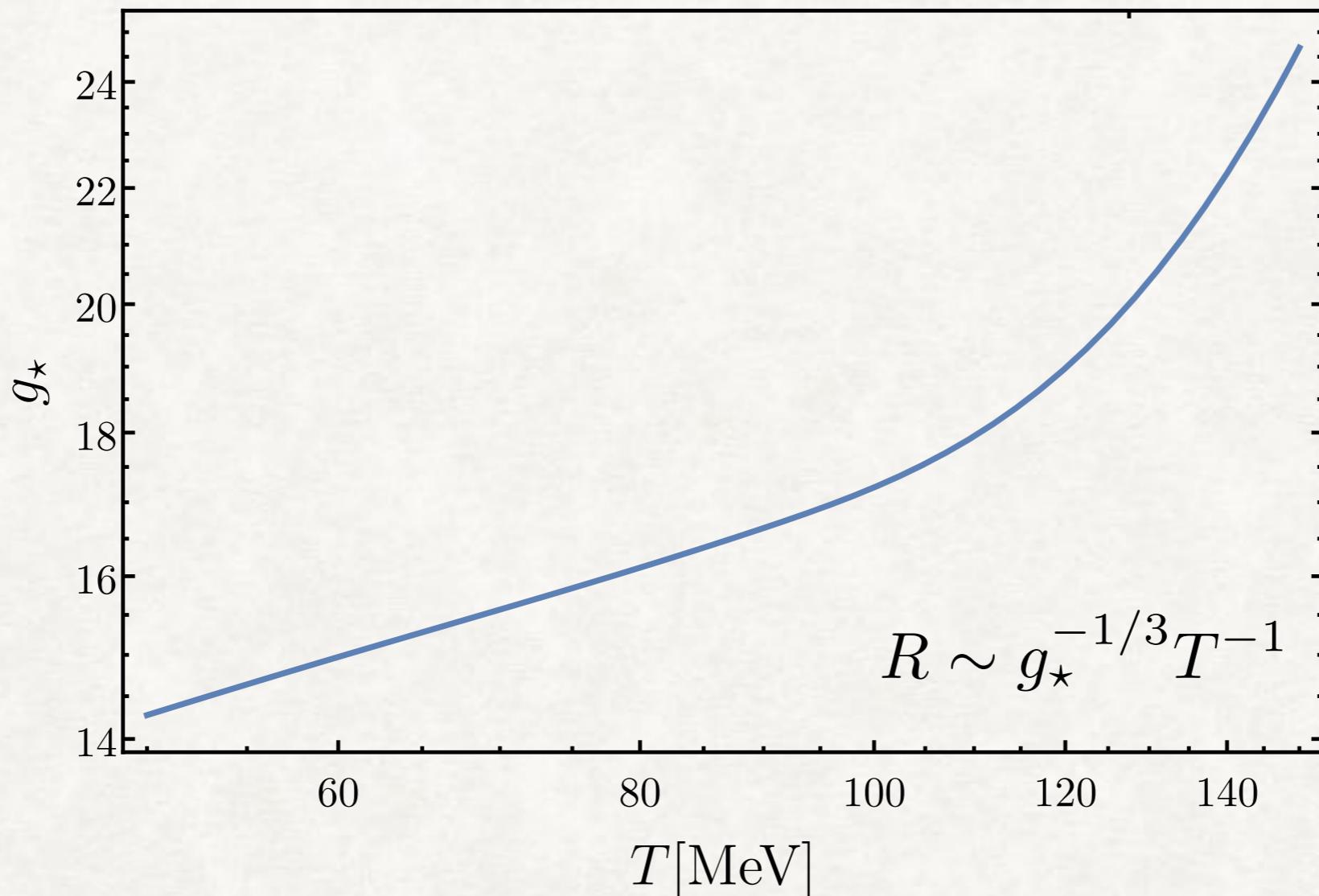
*Higher momentum
modes interact
more,
Decouple at lower
temperature!*



CALCULATION OF THE ABUNDANCE

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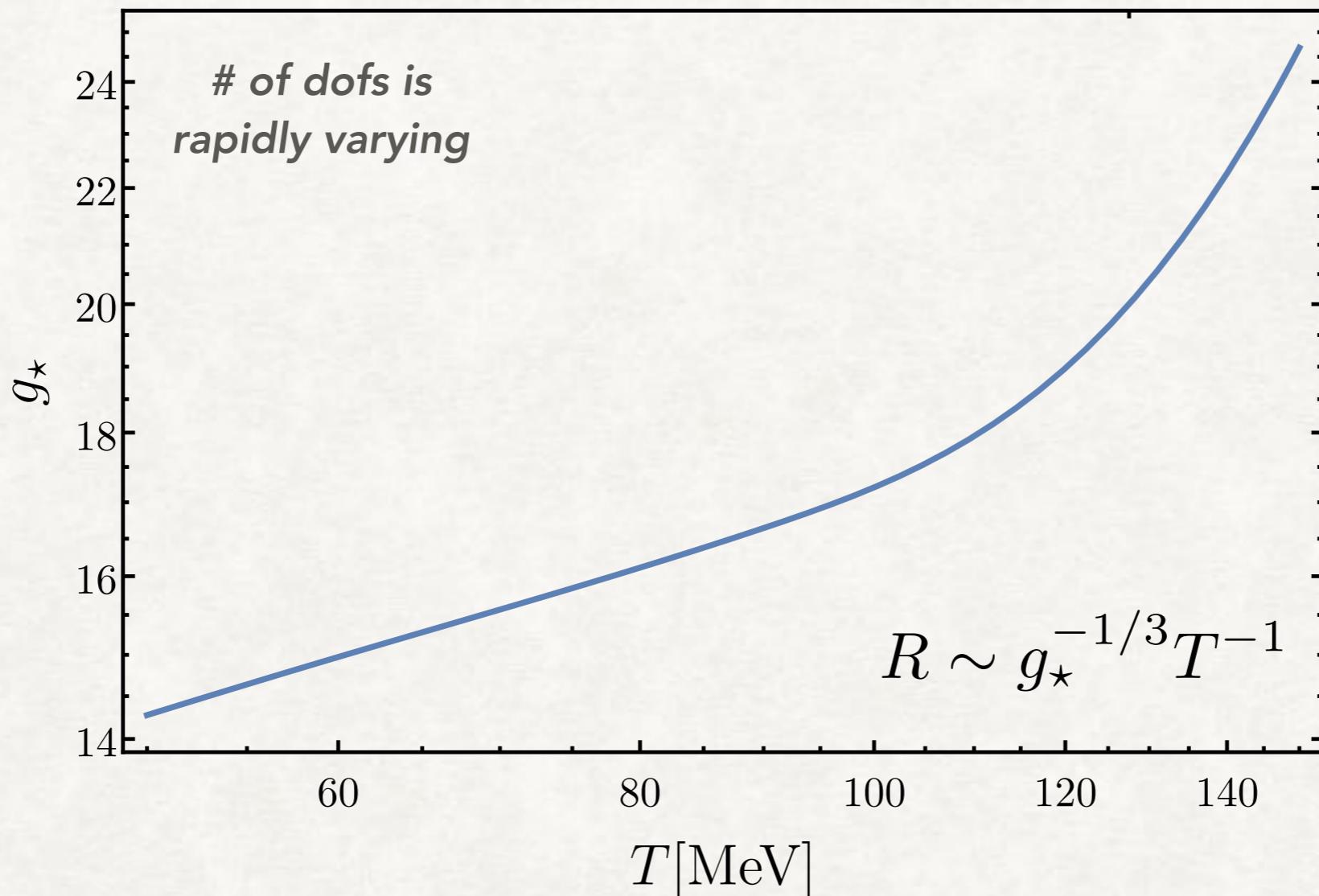


g_* from Borsanyi et al, Nature 539 (2016)

CALCULATION OF THE ABUNDANCE

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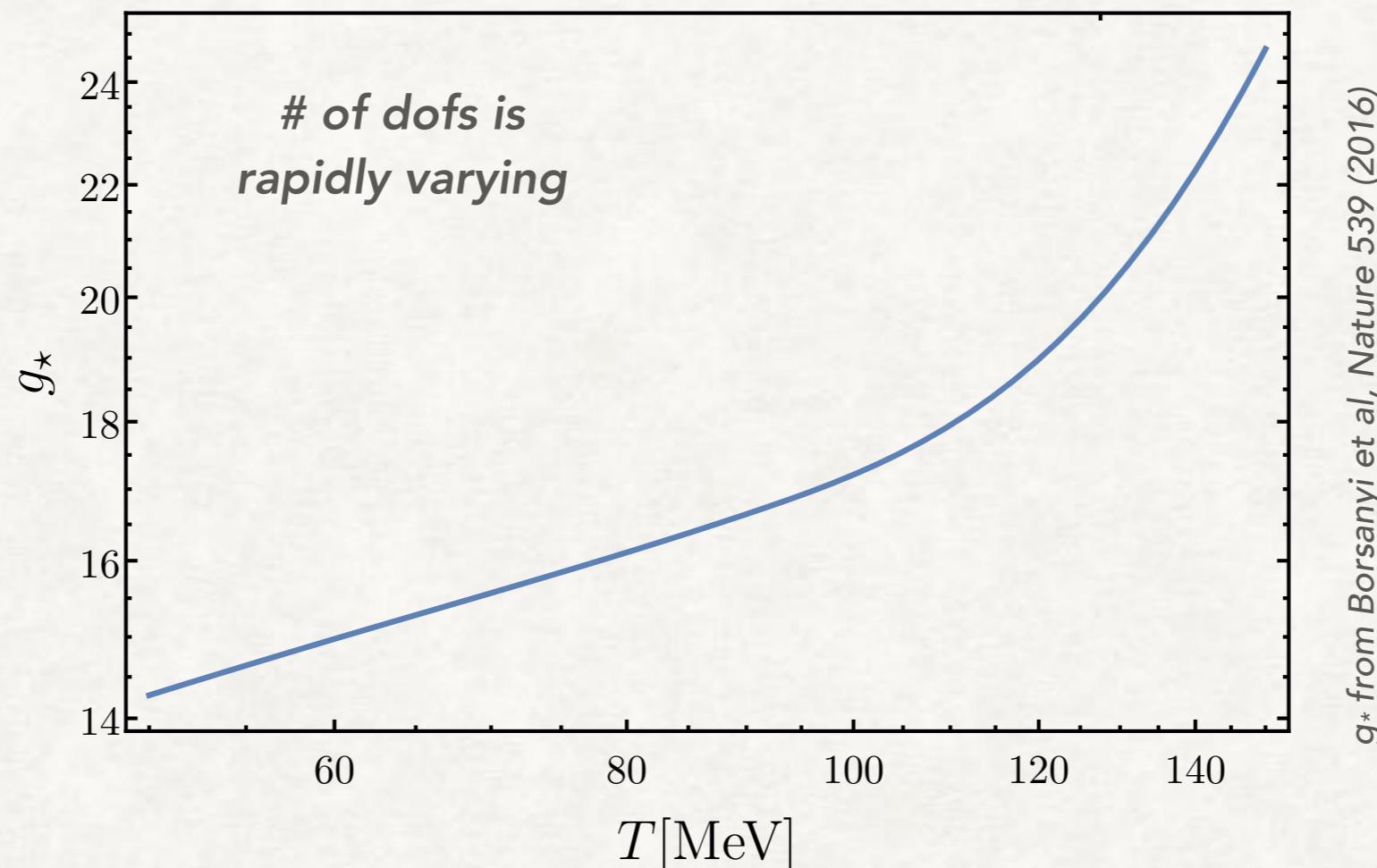
g_* from Borsanyi et al, Nature 539 (2016)

Higher momentum modes are less diluted!

PRODUCTION FROM PIONS

Simple estimate does not work because

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Expect strong spectral distortion!



Enhanced abundance

Similar to neutrinos, $(3 + 0.043)$, but much larger effect due to change of dofs + behaviour of rate

MOMENTUM-DEPENDENT BOLTZMANN EQUATION

Solve for axion distribution function

Boltzmann equation in comoving momenta $\mathbf{p} = R(t)\mathbf{k}$

$$\frac{df_{\mathbf{p}}}{dt} = \underset{\text{Creation}}{(1 + f_{\mathbf{p}}) \Gamma^<} - \underset{\text{Destruction}}{f_{\mathbf{p}} \Gamma^>}, \quad \Gamma^< = e^{-\frac{E}{T}} \Gamma^>$$

$$\Gamma^> = \frac{1}{2E} \int \left(\prod_{i=1}^3 \frac{d^3 \mathbf{k}_i}{(2\pi)^3 2E_i} \right) f_1^{\text{eq}}(1 + f_2^{\text{eq}})(1 + f_3^{\text{eq}})(2\pi)^4 \delta^{(4)}(k^\mu + k_1^\mu - k_2^\mu - k_3^\mu) |\mathcal{M}|^2$$

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Bose-Einstein

MOMENTUM-DEPENDENT BOLTZMANN EQUATION

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Bose-Einstein



$$\rho_a = R^{-4} \int \frac{d^3 \mathbf{p}}{(2\pi)^3} |\mathbf{p}| f_{\mathbf{p}} \rightarrow \Delta N_{\text{eff}} = \frac{8}{7} \left(\frac{11}{4} \right)^{\frac{4}{3}} \frac{\rho_a}{\rho_\gamma} |_{\text{rec}}$$

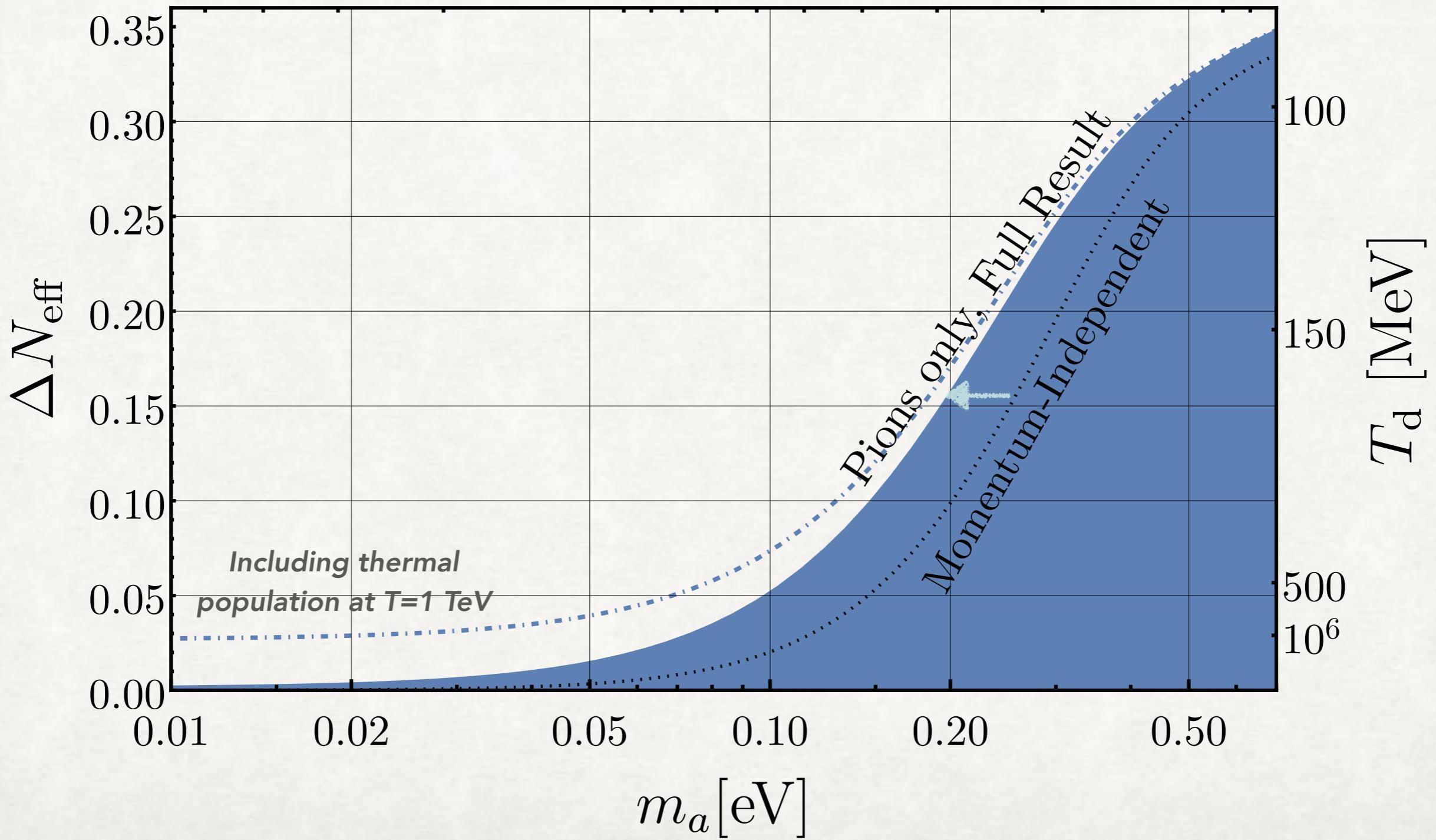
RESULTS

Minimal KSVZ model

f_a [GeV]

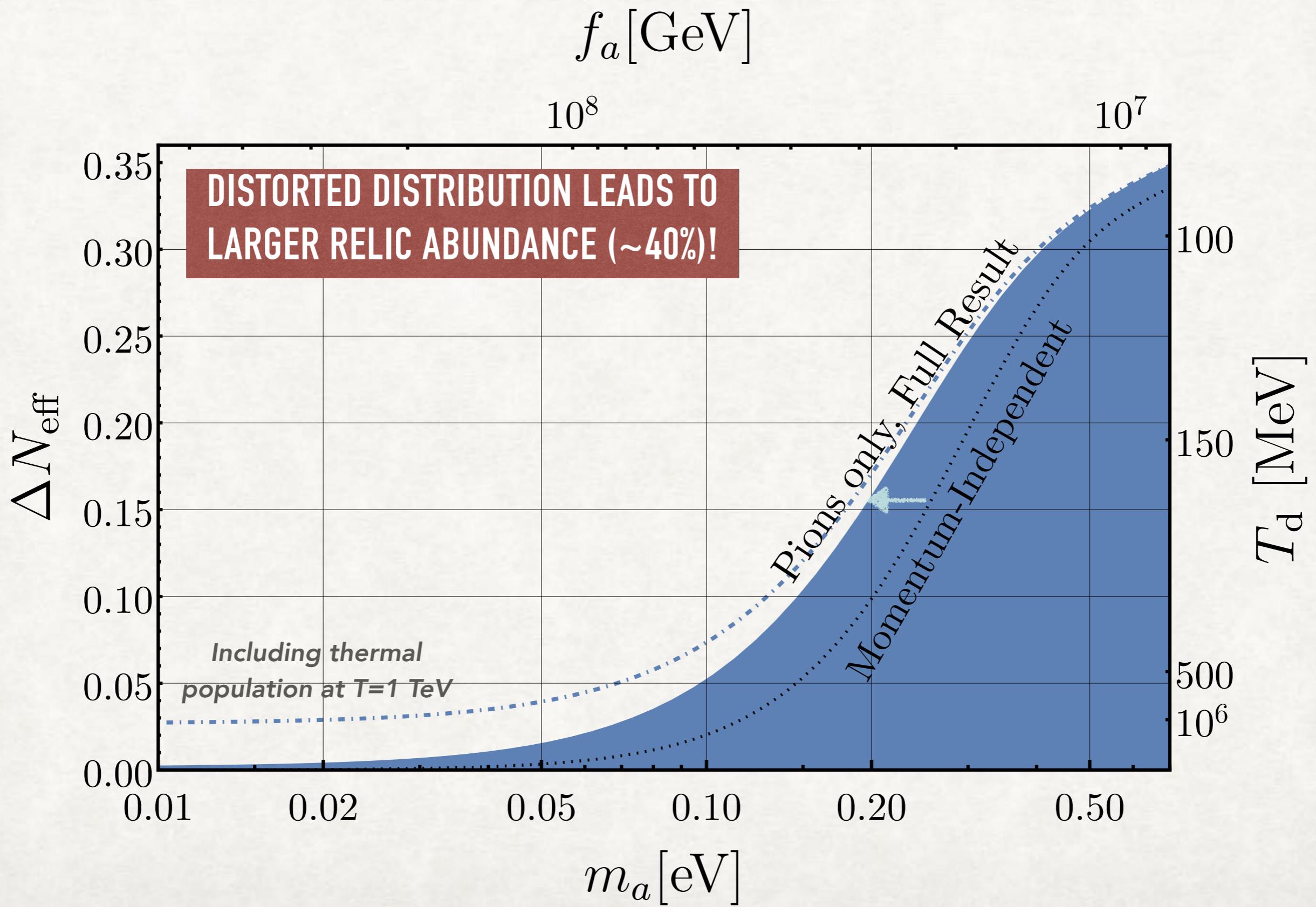
10^8

10^7



RESULTS

Minimal KSVZ model



RESULTS

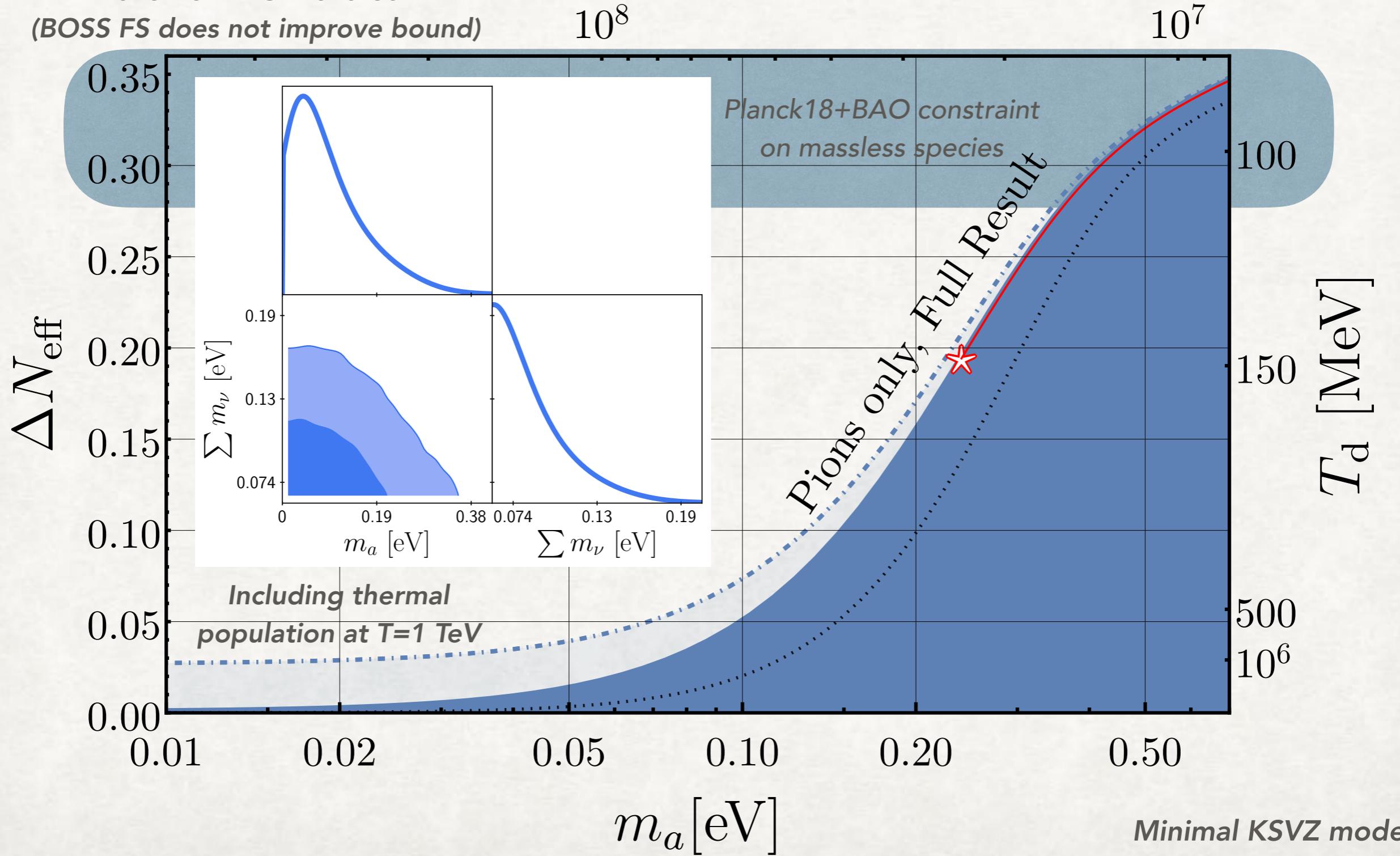
$$\Lambda\text{CDM} + \sum m_\nu + m_a$$

*Planck18+BAO+Pantheon
(BOSS FS does not improve bound)*

$$f_a [\text{GeV}]$$

$$m_a \leq 0.24 \text{ eV}, 95\% \text{ C.L.}$$

$$\sum_\nu m_\nu \leq 0.14 \text{ eV}, 95\% \text{ C.L.}$$



RESULTS

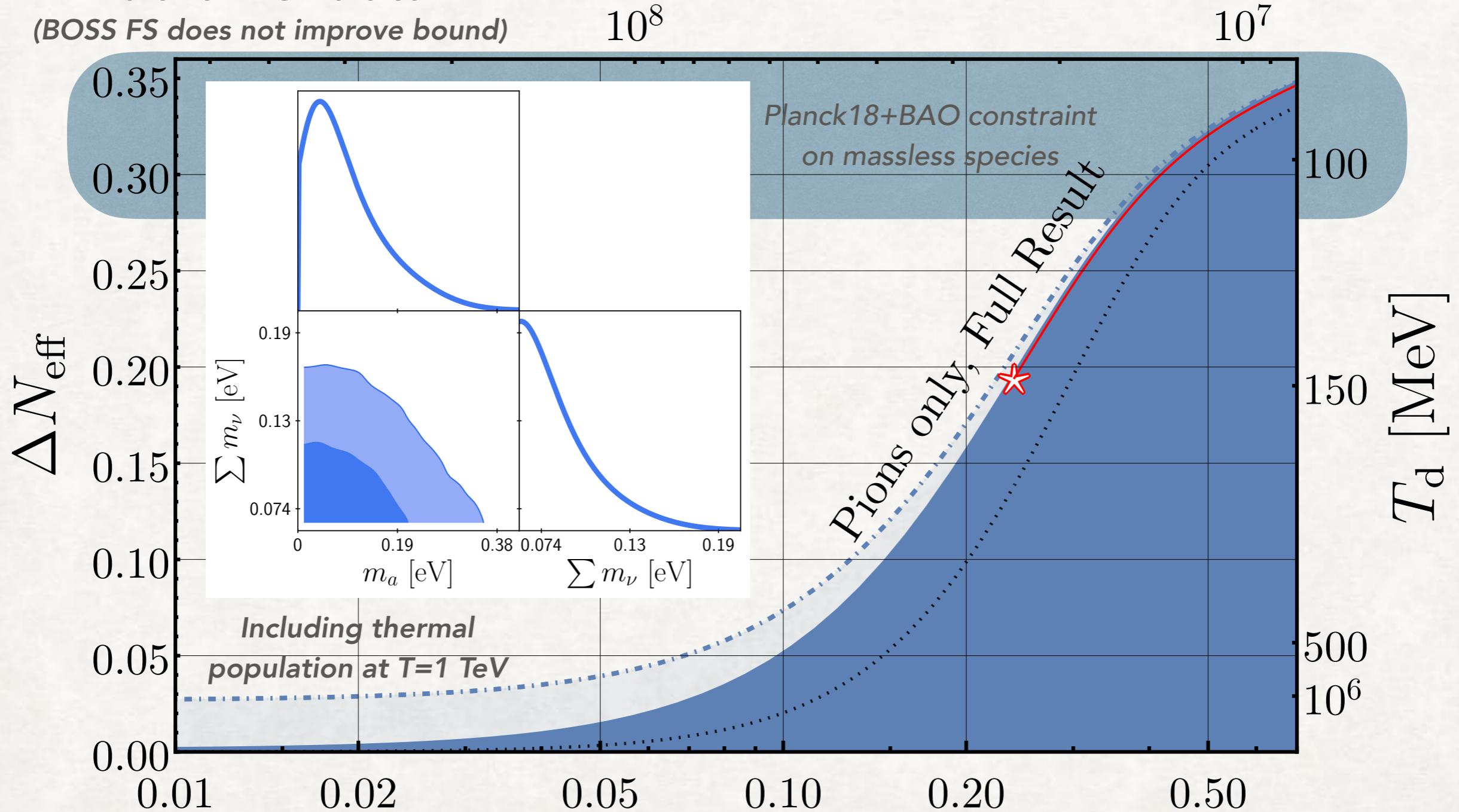
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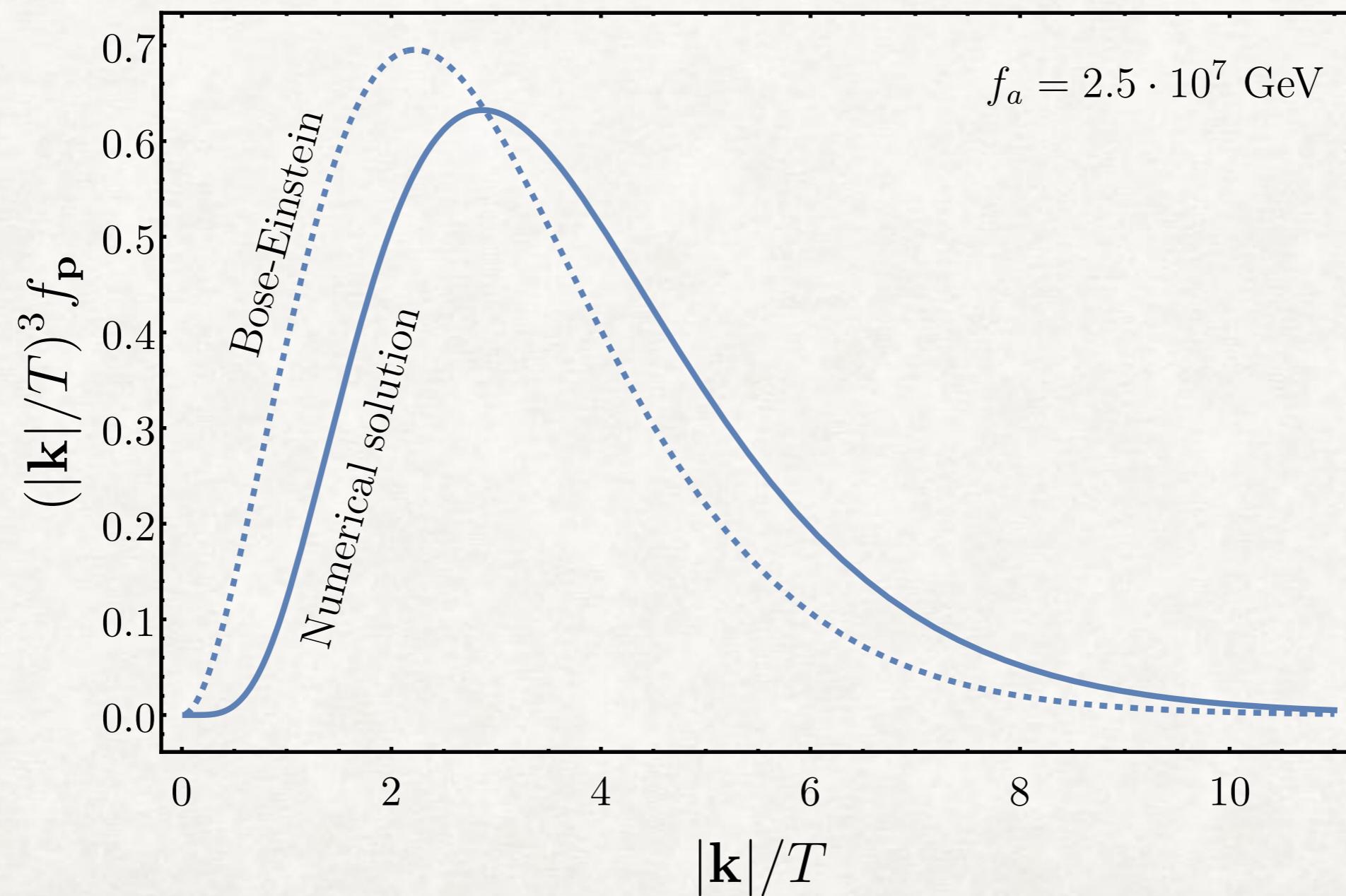
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AXION DISTRIBUTION FUNCTION

Minimal KSVZ model

*Comparison with thermal distribution
with same energy density*

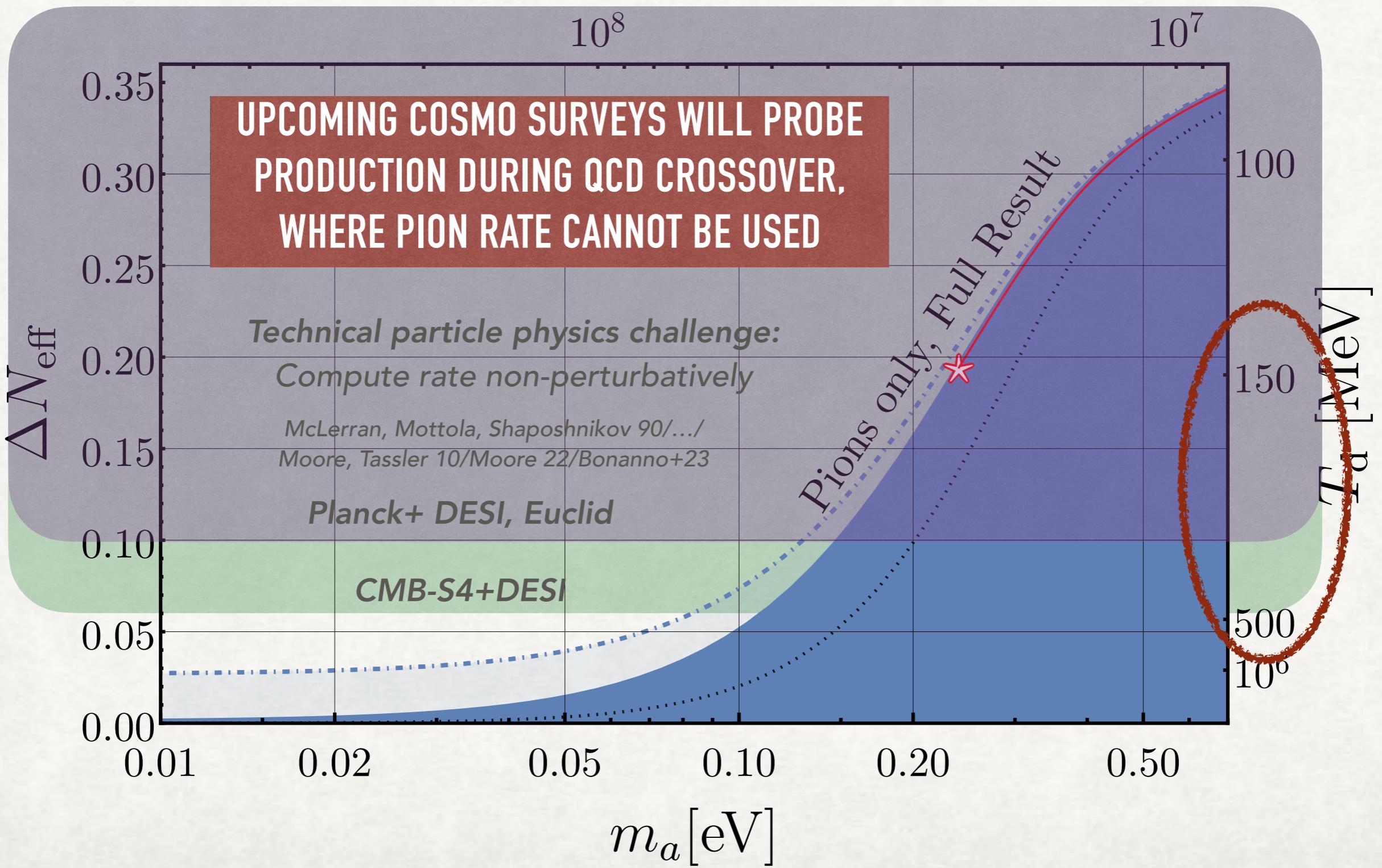


THE (NEAR) FUTURE



OUTLOOK

f_a [GeV]

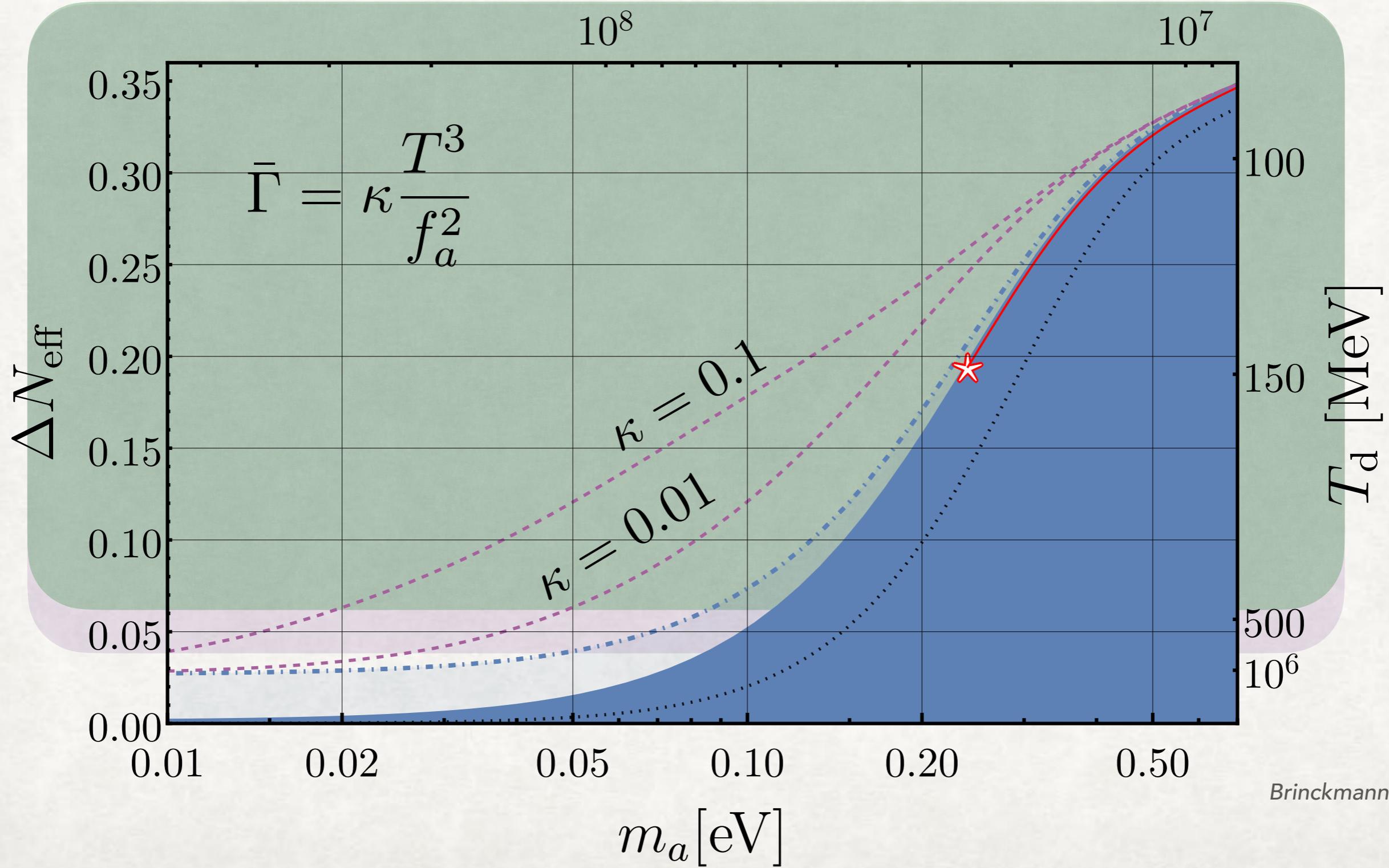


OUTLOOK

Caveat:

Sensitivity reach to massless species,
Underestimates reach on QCD axion

$f_a [\text{GeV}]$



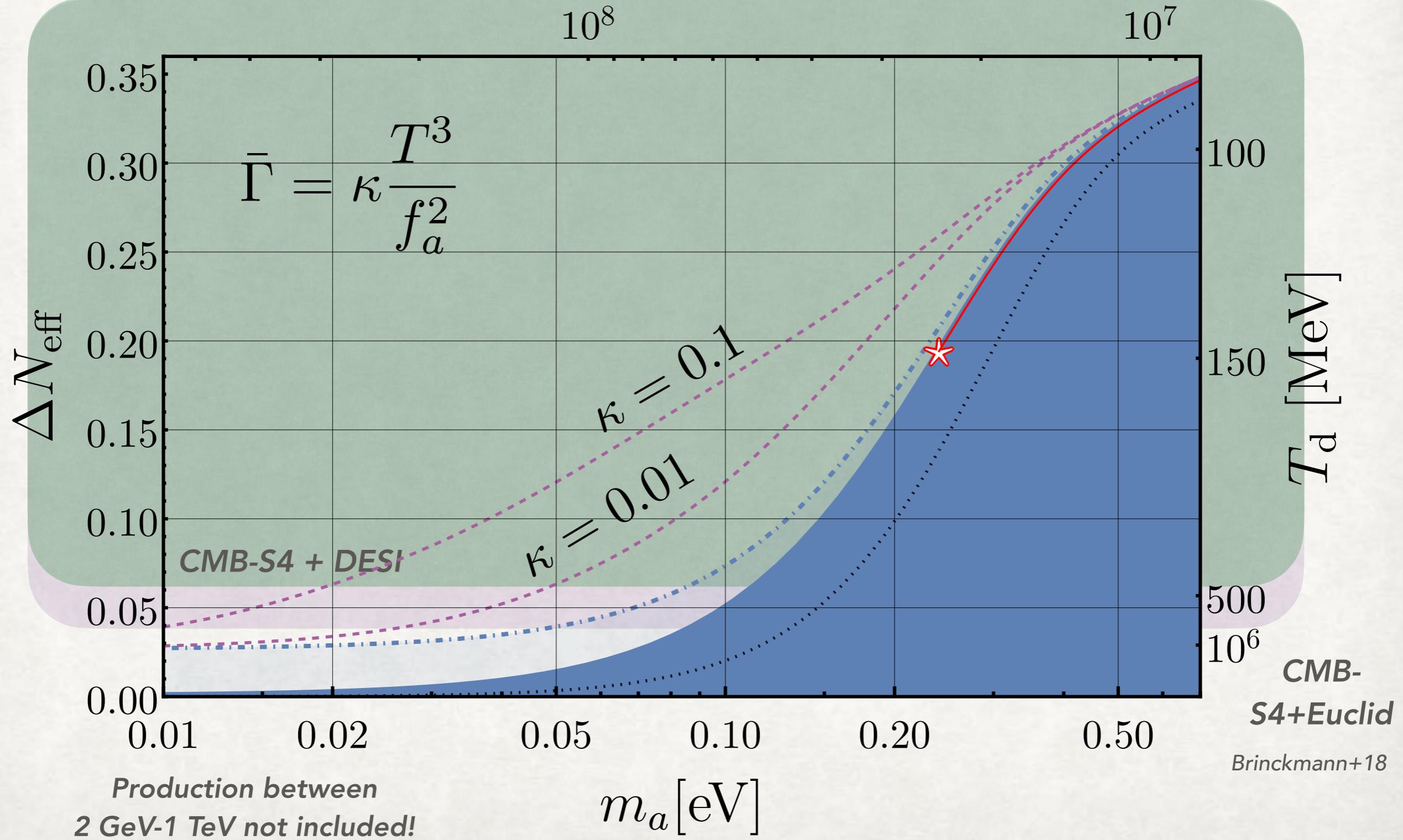
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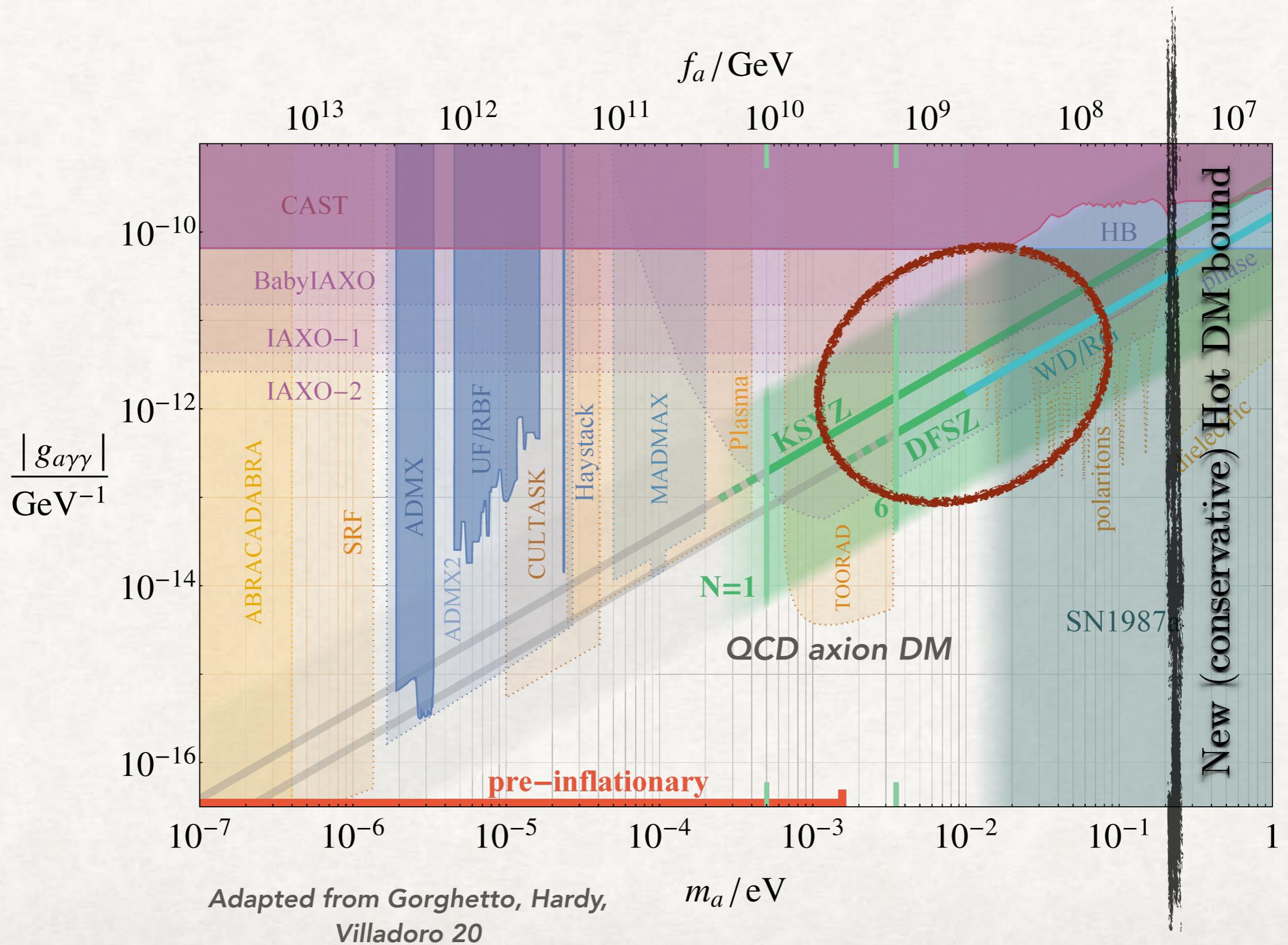
Sensitivity reach to massless species,
Underestimates reach on QCD axion

**Tentative dimensional analysis
estimate**

$f_a [\text{GeV}]$

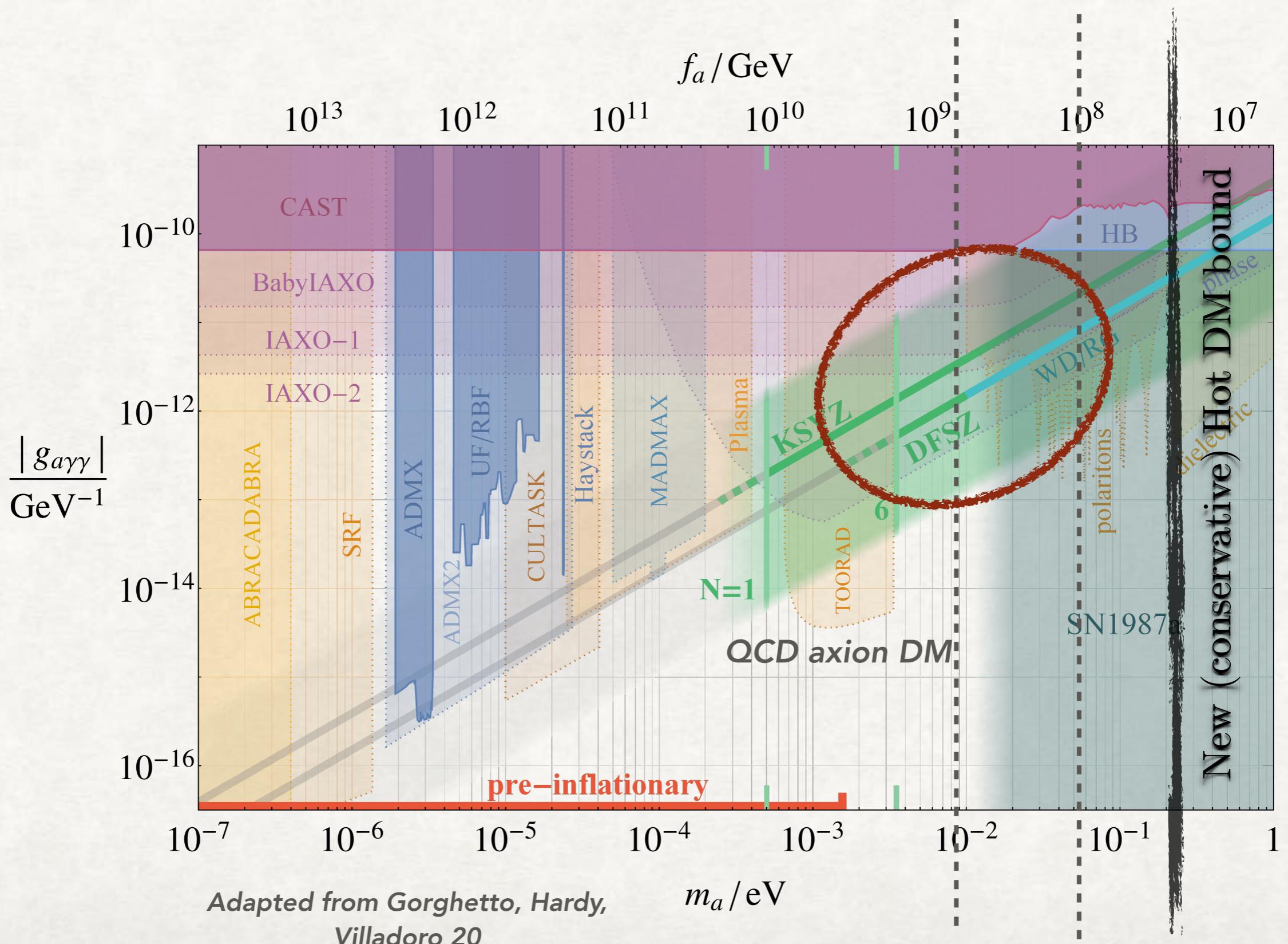


DISCOVERING THE QCD AXION WITH COSMOLOGY?



DISCOVERING THE QCD AXION WITH COSMOLOGY?

Future bound/detection?



SUMMARY/QUESTIONS FOR LSS

*Near future reach on thermally produced dark radiation/hot DM
w/ sub-eV mass restricted to production during QCD crossover*

*Variation of # of dofs can cause strong spectral distortion,
enhanced abundance*

Well-motivated candidate: QCD axion

Full reach of DESI/Euclid?

Observable effects of spectral distortion beyond enhanced abundance?

Impact on neutrino masses?

PARTICLE PHYSICS AT THE EV?

Want: (a) light particle(s) (with (sub)-eV mass) that decouples around the QCD epoch

UV-MOTIVATED

TENSIONS

SIGNAL BUILDING

2305.14166 w/ I. Allali, M. Hertzberg

See talk by Martin!

INTERACTING DARK SECTOR WITH MASS THRESHOLDS

a.k.a. "Stepped Dark Radiation" (SDR)

Aloni, Berlin, Joseph,
Schmaltz, Weiner
21/+Sivarajan 22/

See talk by
Martin!

$$g_{\star}^{T < m} < g_{\star}^{T \gtrsim m}$$

Massive species annihilate away
(No hot DM constraints)

$$m$$

$$g_{\star}^{T \gtrsim m}$$

Agnostic about
production of initial
population (but
produced after
BBN)

$$\nu, \bar{\nu}$$

$$\nu, \bar{\nu}$$

$$T$$

Dark sector abundance
increases

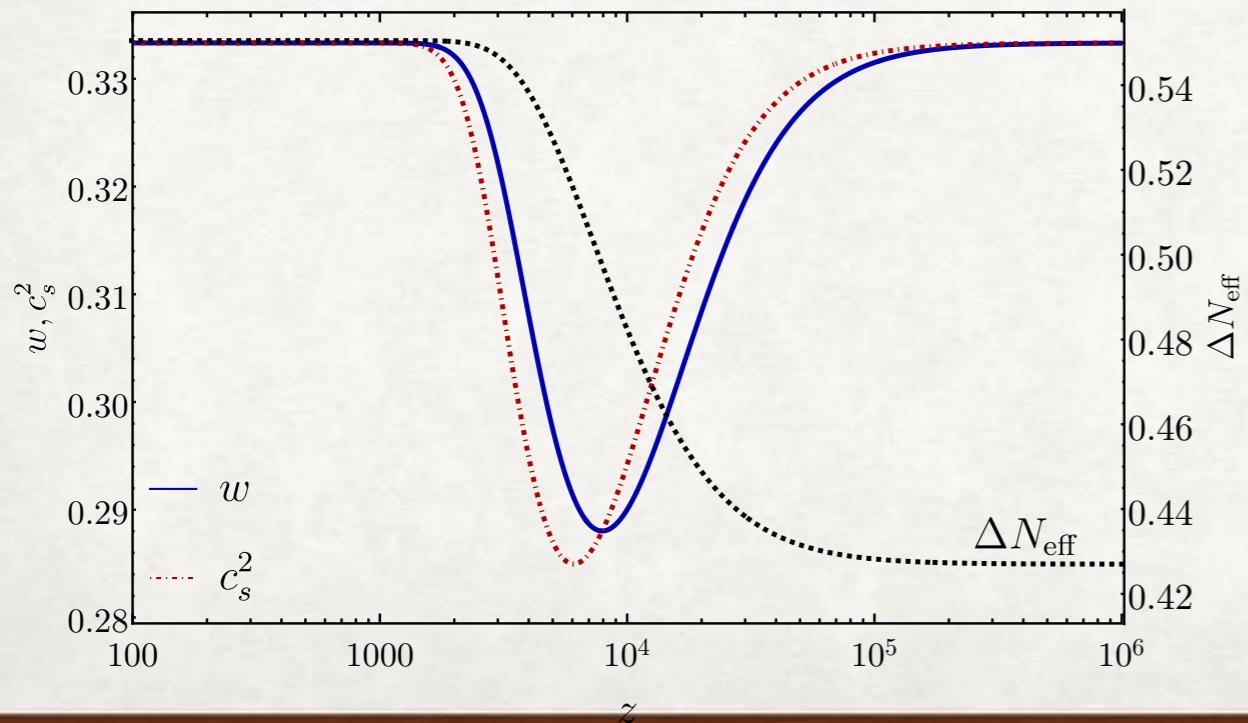
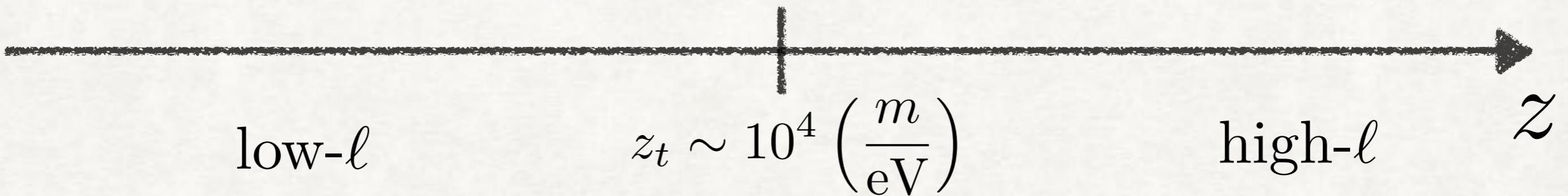
$$\Delta N_{\text{eff}}^{T < m} = \Delta N_{\text{eff}}^{T \gtrsim m} \left(\frac{g_{\star}^{T \gtrsim m}}{g_{\star}^{T < m}} \right)^{\frac{1}{3}}$$

Several possible particle physics realisations: scalars+fermions,
fermions+dark photon, QCD-like sector ...

IMPACT ON CMB

$$\Delta N_{\text{eff}}^{z < z_t} > \Delta N_{\text{eff}}^{z > z_t}$$

$$\Delta N_{\text{eff}}^{z \sim z_t}$$



$$\dot{\delta} = -(1+w)(\theta + \frac{\dot{h}}{2}) - 3\mathcal{H}(c_s^2 - w)\delta$$

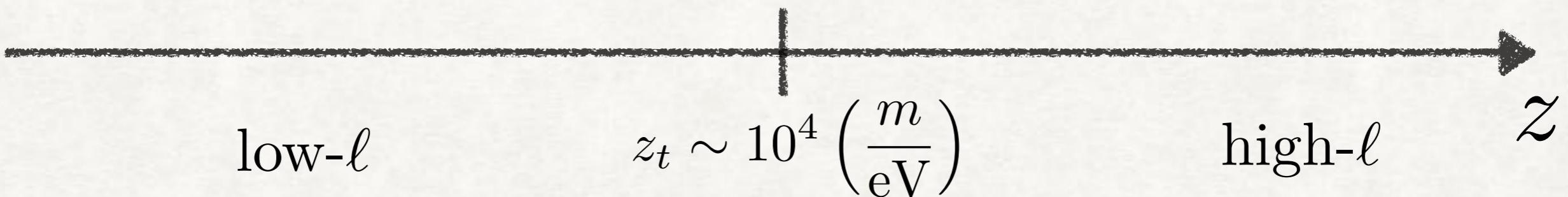
$$\dot{\theta} = -\mathcal{H}(1-3w)\theta - \frac{\dot{w}}{1+w}\theta + \frac{c_s^2}{1+w}k^2\delta - k^2\sigma$$

IMPACT ON CMB

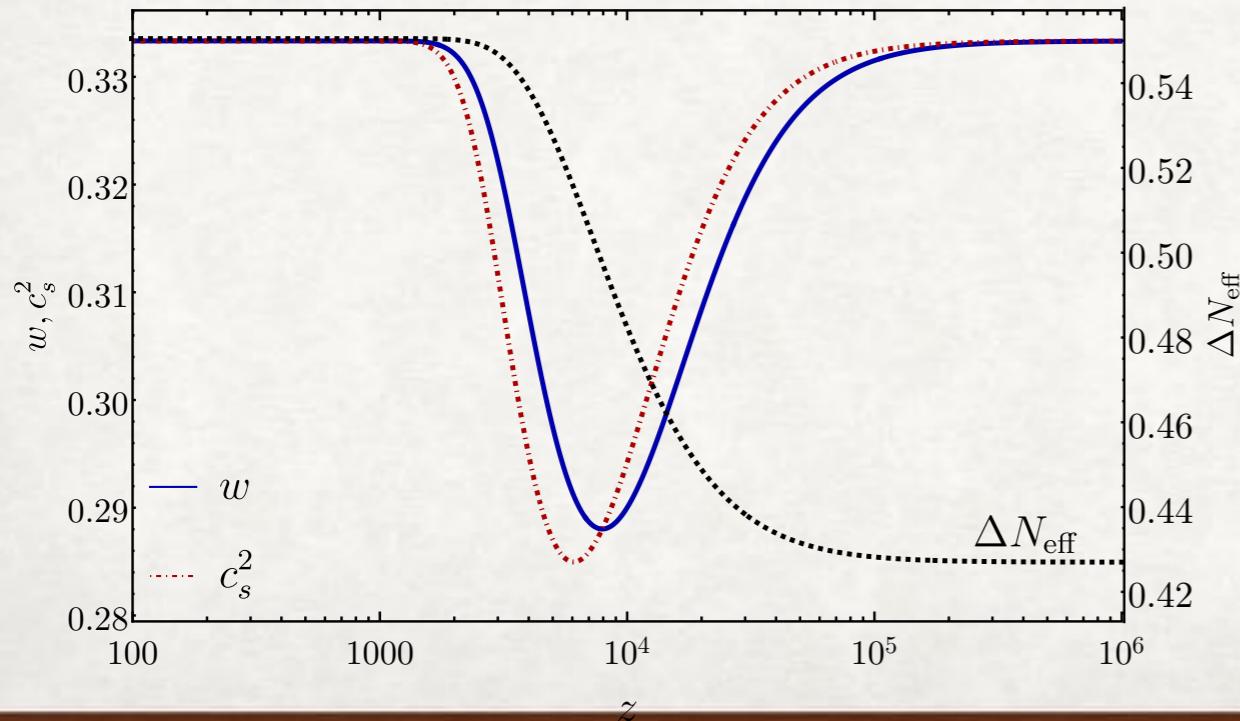
Background

$$\Delta N_{\text{eff}}^{z < z_t} > \Delta N_{\text{eff}}^{z > z_t}$$

$$\Delta N_{\text{eff}}^{z \sim z_t}$$



Softening of the equation of state



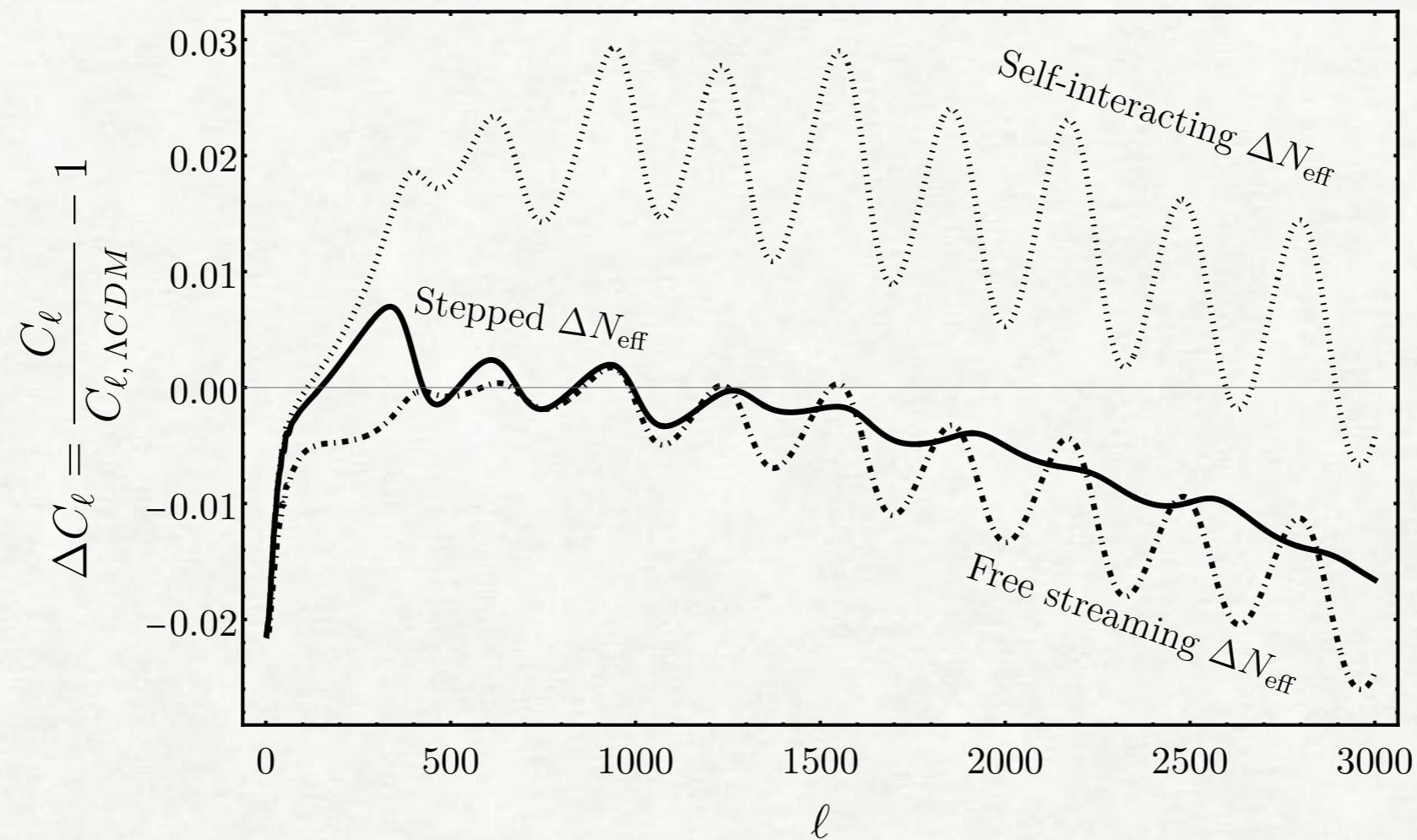
Perturbations

$$\dot{\delta} = -(1+w)(\theta + \frac{\dot{h}}{2}) - 3\mathcal{H}(c_s^2 - w)\delta$$

$$\dot{\theta} = -\mathcal{H}(1-3w)\theta - \frac{\dot{w}}{1+w}\theta + \frac{c_s^2}{1+w}k^2\delta - k^2\sigma$$

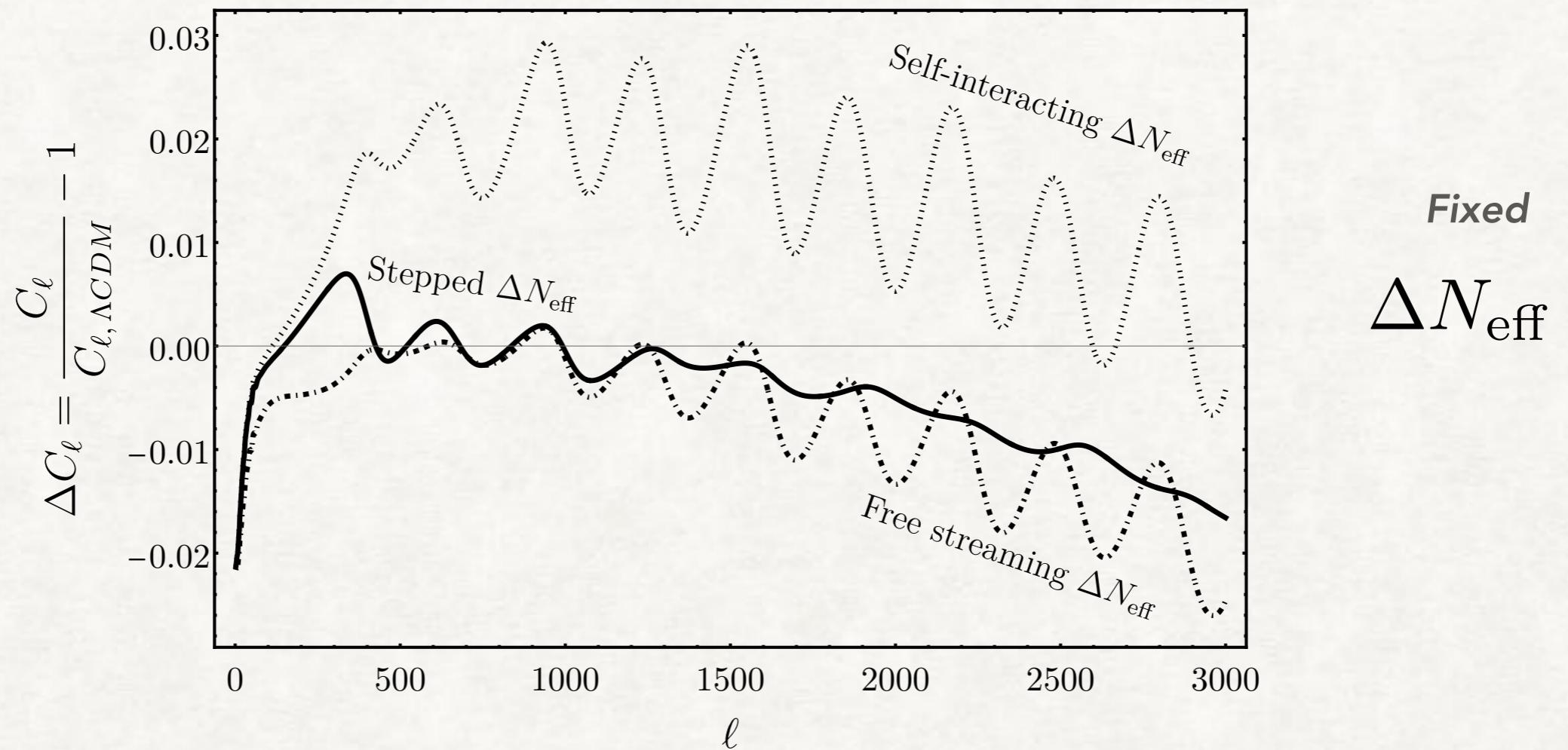
*Radiation fluid is self-interacting,
behaves like ideal fluid (not free streaming)*

IMPACT ON CMB



ΔN_{eff}

IMPACT ON CMB



Motivation from Hubble tension

Simple (free streaming or not) dark radiation not enough to address tension

Time-dependent properties allow for larger relic abundance, degenerate with larger Hubble constant

IMPACT ON HUBBLE TENSION

$$\Delta N_{\text{eff}}^{\text{IR}}, \quad z_t, \quad r_g \equiv \frac{g_\star^{\text{UV}} - g_\star^{\text{IR}}}{g_\star^{\text{IR}}}$$

$$\text{IR (UV)} \equiv z < (>) z_t,$$

*See also Schöneberg,
Abellan 22*

IMPACT ON HUBBLE TENSION

Effective fluid model is characterised by three extra parameters

$$\Delta N_{\text{eff}}^{\text{IR}}, \quad z_t, \quad r_g \equiv \frac{g_\star^{\text{UV}} - g_\star^{\text{IR}}}{g_\star^{\text{IR}}}$$

$$\text{IR (UV)} \equiv z < (>) z_t,$$

*Model shown to be effective in reducing Hubble tension
With tight priors on redshift of the transition, fitting to Planck18+BAO*

*Similar questions
addressed for
Early Dark Energy*

Prior dependence?
Look elsewhere effect

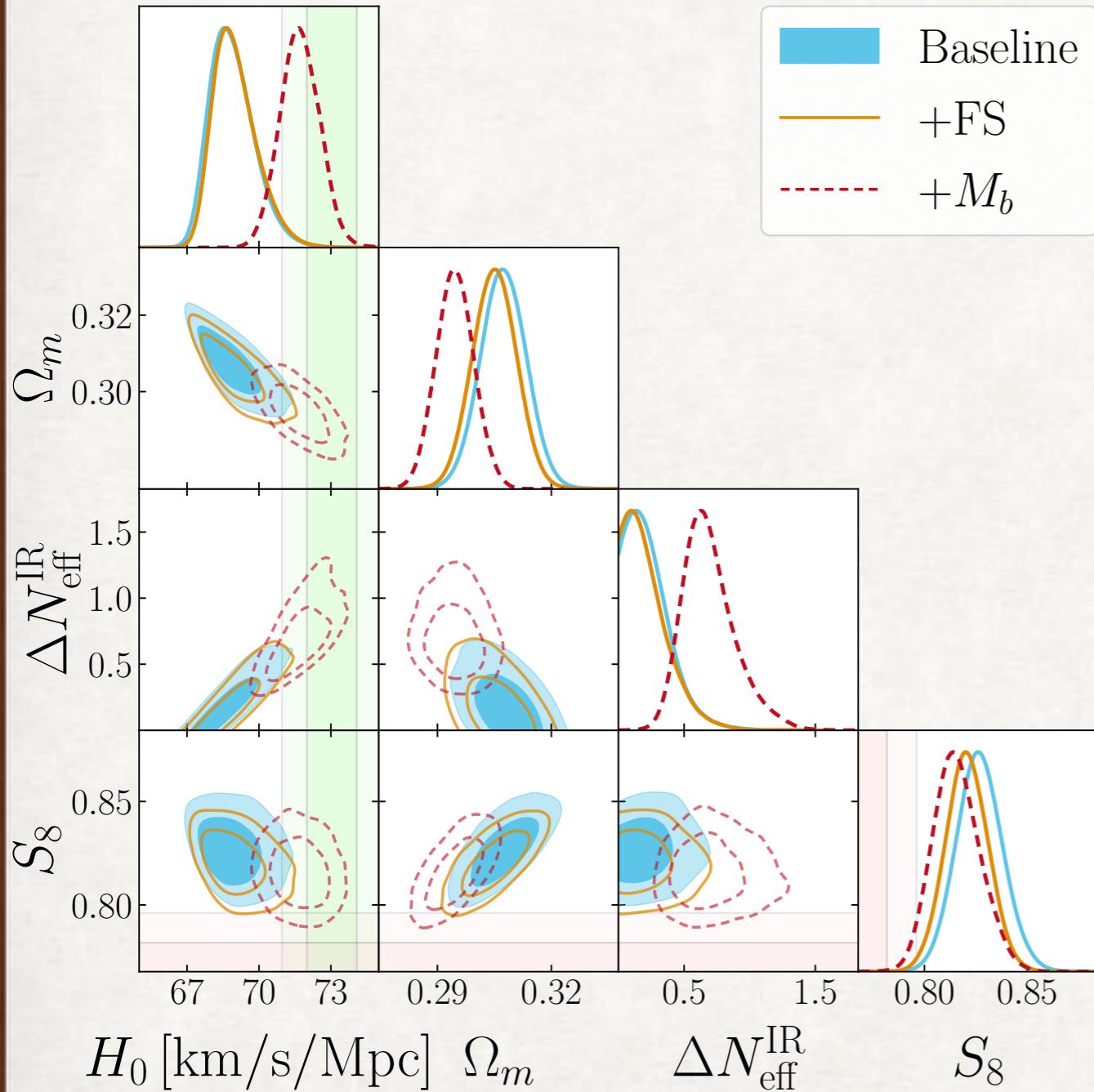
*See also Schöneberg,
Abellan 22*

Constraints from BOSS with EFTofLSS?
Shifts of other parameters to keep goodness of fit (CDM, S8)

CONSTRAINTS WITH FULL-SHAPE INFORMATION

PyBird

$$z_t \in [3, 5]$$



| Parameter | Baseline | Baseline + FS |
|-------------------------------------|---------------------------------------|-------------------------------------|
| $\Delta N_{\text{eff}}^{\text{IR}}$ | < 0.546 (0.289) | < 0.55 (0.08) |
| $\log_{10} z_t$ | Unconstrained (4.29) | Unconstrained (4.97) |
| r_g | Unconstrained (4.0) | Unconstrained (2.34) |
| H_0 [km/s/Mpc] | $68.89 (69.34)^{+0.71}_{-1.1}$ | $69.01 (68.37)^{+0.66}_{-1.1}$ |
| S_8 | $0.827 (0.834)^{+0.011}_{-0.011}$ | $0.821 (0.824)^{+0.010}_{-0.010}$ |
| M_b | $-19.381 (-19.369)^{+0.021}_{-0.032}$ | $-19.378 (-19.4)^{+0.019}_{-0.032}$ |
| $\Delta\chi^2$ | -1.4 | -1.62 |
| $Q_{\text{DMAP}}^{M_b}$ | 2.74σ | 2.72σ |
| M_b GT | 3.74σ | 3.77σ |
| M_b IT | 3.03σ | 2.94σ |
| ΔAIC^{M_b} | -20.67 | -18.68 |

$\gtrsim 3\sigma$

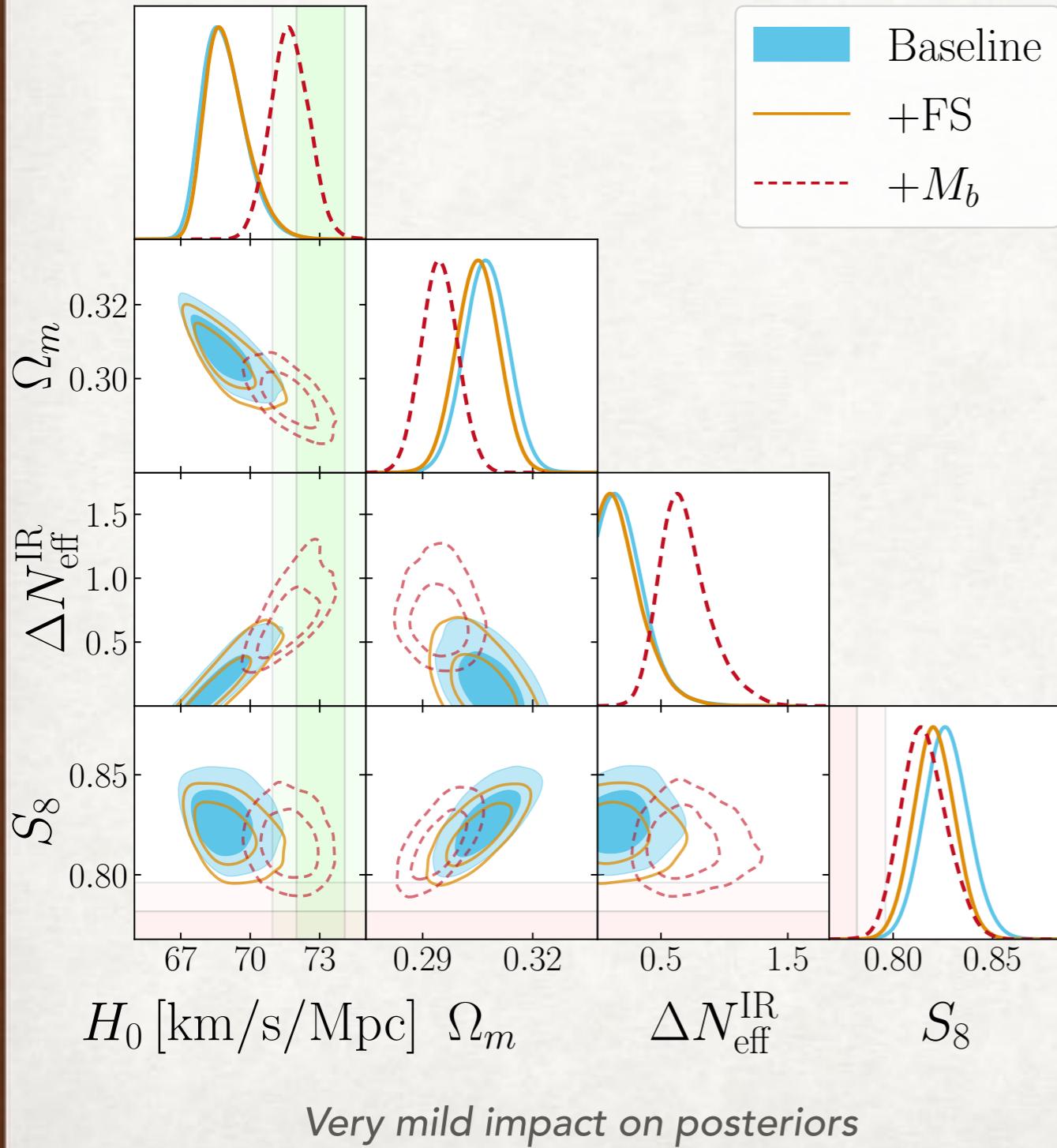
CONSTRAINTS WITH FULL-SHAPE INFORMATION

PyBird

D'Amico, Senatore, Zhang 20

$$z_t \in [3, 5]$$

Baseline = Planck18+BAO+Pantheon



| Parameter | Baseline | Baseline + FS |
|-------------------------------------|---------------------------------------|-------------------------------------|
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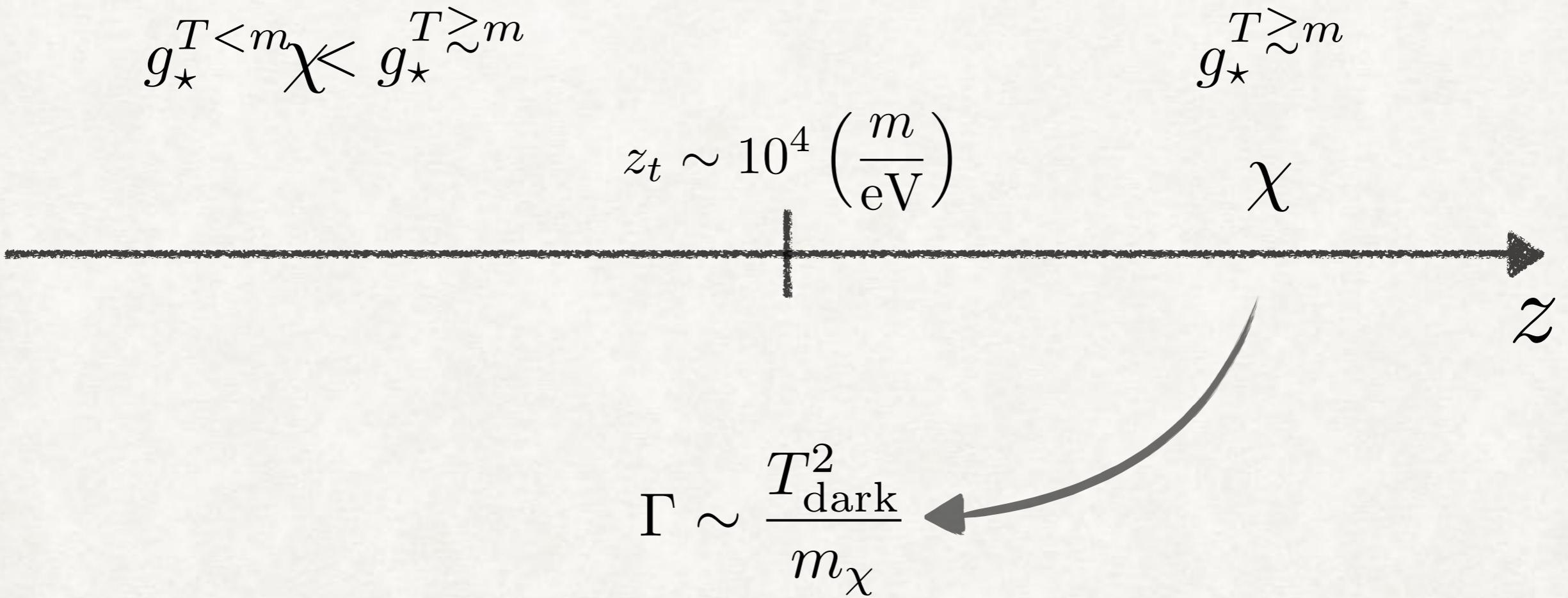
Very mild impact on posteriors

LSS shifts bestfit values
towards LambdaCDM

Hubble Tension $\gtrsim 3\sigma$

ADDING INTERACTIONS WITH DARK MATTER

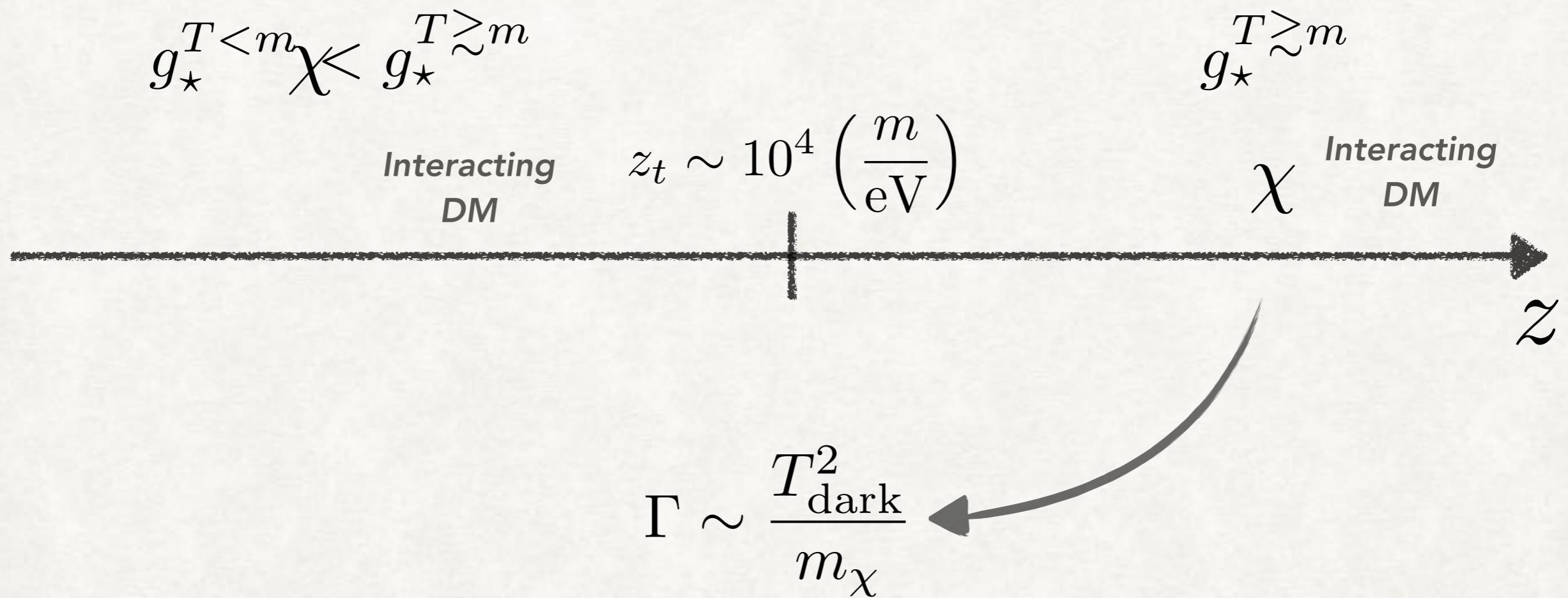
Needed? See new KiDS+DES analysis...



ADDING INTERACTIONS WITH DARK MATTER

Motivation from S8: suppression of matter power spectrum

Needed? See new KiDS+DES analysis...

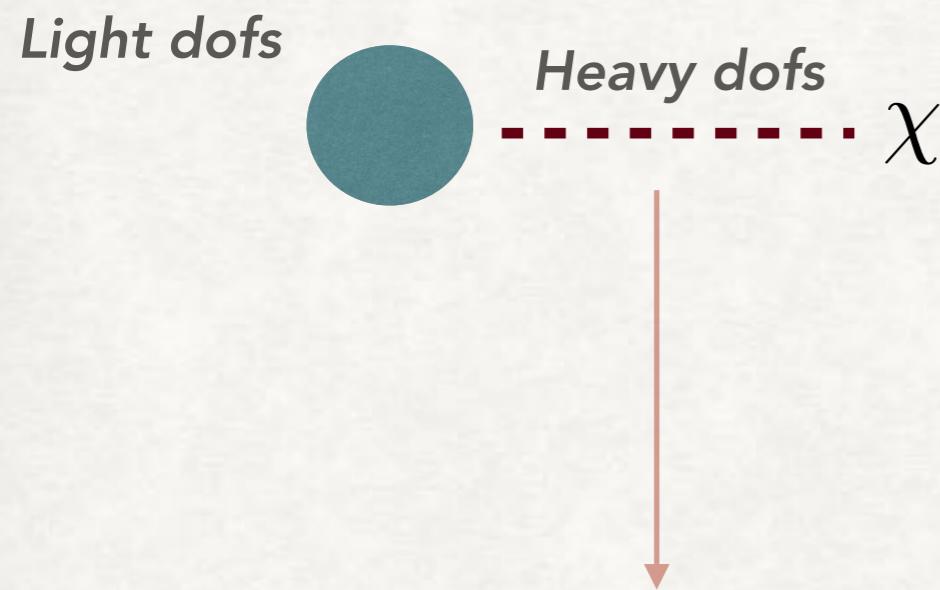


TWO SCENARIOS

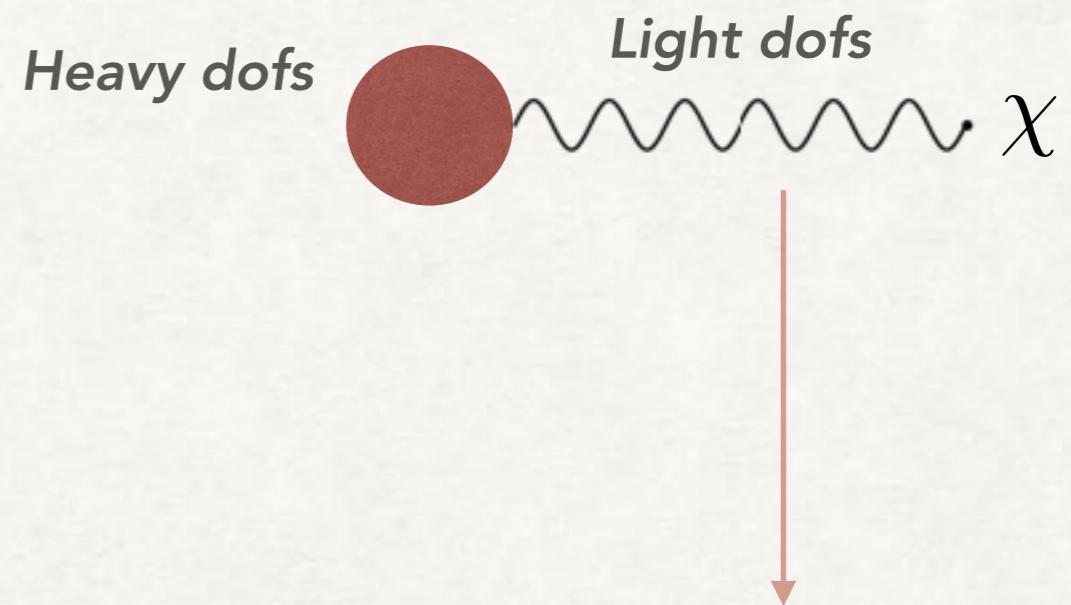
Joseph, Aloni,
Schmaltz, Sivarajan,
Weiner 22

$$z \lesssim z_t$$

Buen-Abad, Chako,
Kilic, Marques-
Tavares, Youn 22, 23



$$\Gamma_{z \lesssim z_t} \sim \frac{T_{\text{dark}}^2}{m_\chi} \left(\frac{T_{\text{dark}}}{m} \right)^4$$



$$\Gamma_{z \lesssim z_t} \sim \frac{T_{\text{dark}}^2}{m_\chi} e^{-\frac{m}{T_{\text{dark}}}}$$

TWO SCENARIOS

*Joseph, Aloni,
Schmaltz, Sivarajan,
Weiner 22*

$$z \lesssim z_t$$

*Buen-Abad, Chako,
Kilic, Marques-
Tavares, Youn 22, 23*

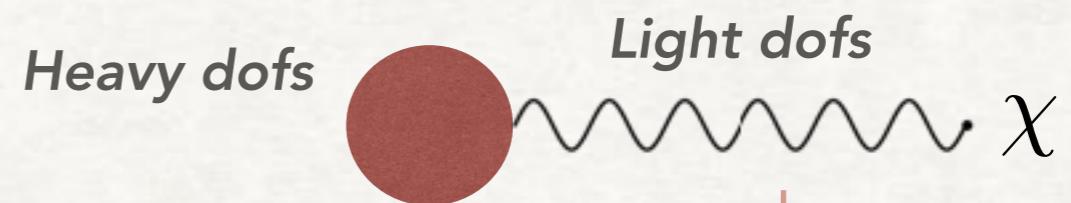
Massive mediator, light
(e.g. Yukawa interactions)

Massless mediator
e.g. scalar-fermion mediated by dark $U(1)$



"Weakly IDM"

$$\Gamma_{z \lesssim z_t} \sim \frac{T_{\text{dark}}^2}{m_\chi} \left(\frac{T_{\text{dark}}}{m} \right)^4$$



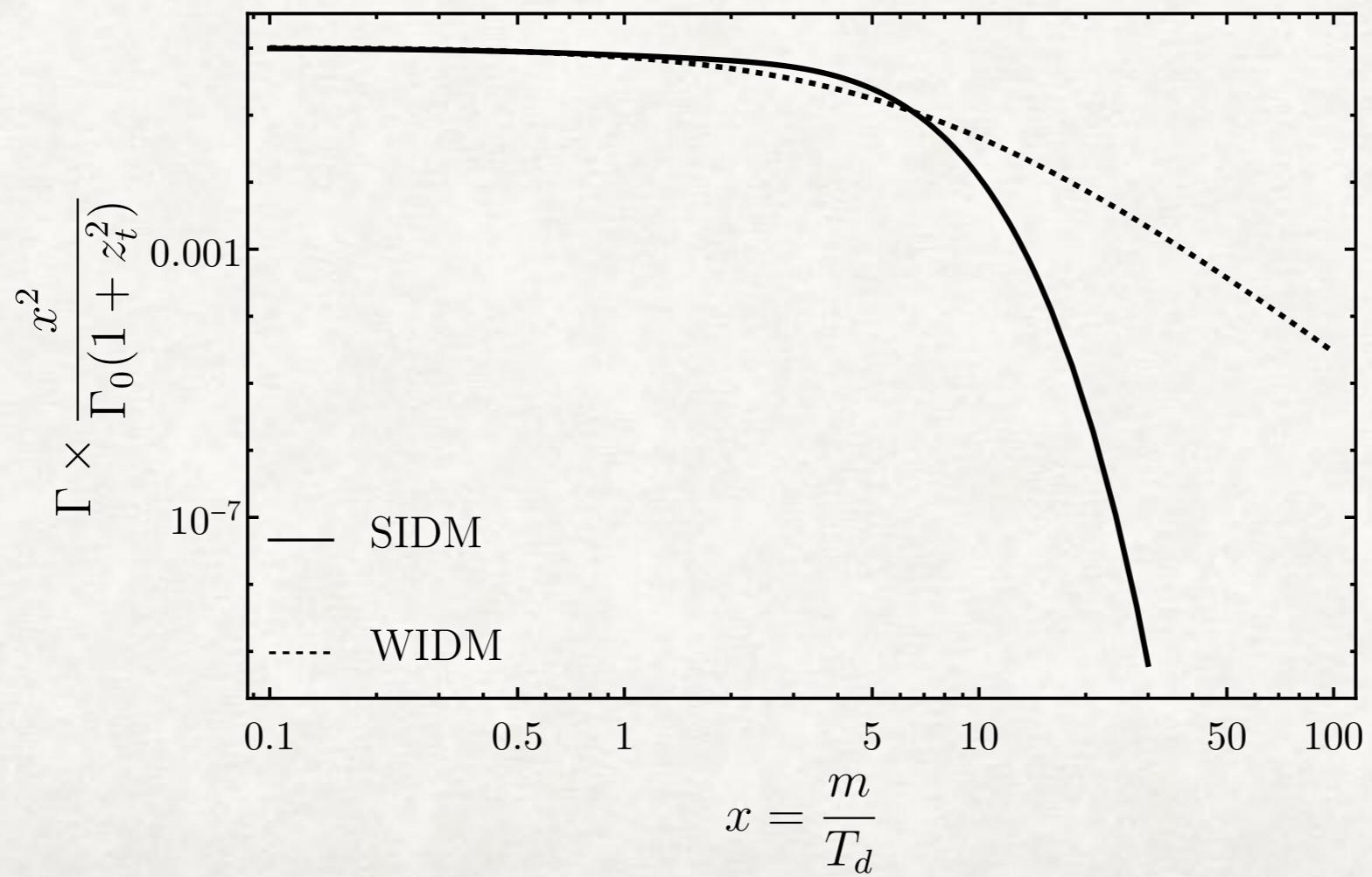
"Strongly IDM"

Boltzmann suppressed

INTERACTING DARK SECTOR WITH MASS THRESHOLD

$$\Gamma = \Gamma_0 \left(\frac{1 + z_t}{x} \right)^2 h(x)$$

$$x \equiv m/T_{\text{dark}}, h(x \rightarrow 0) \rightarrow 1$$

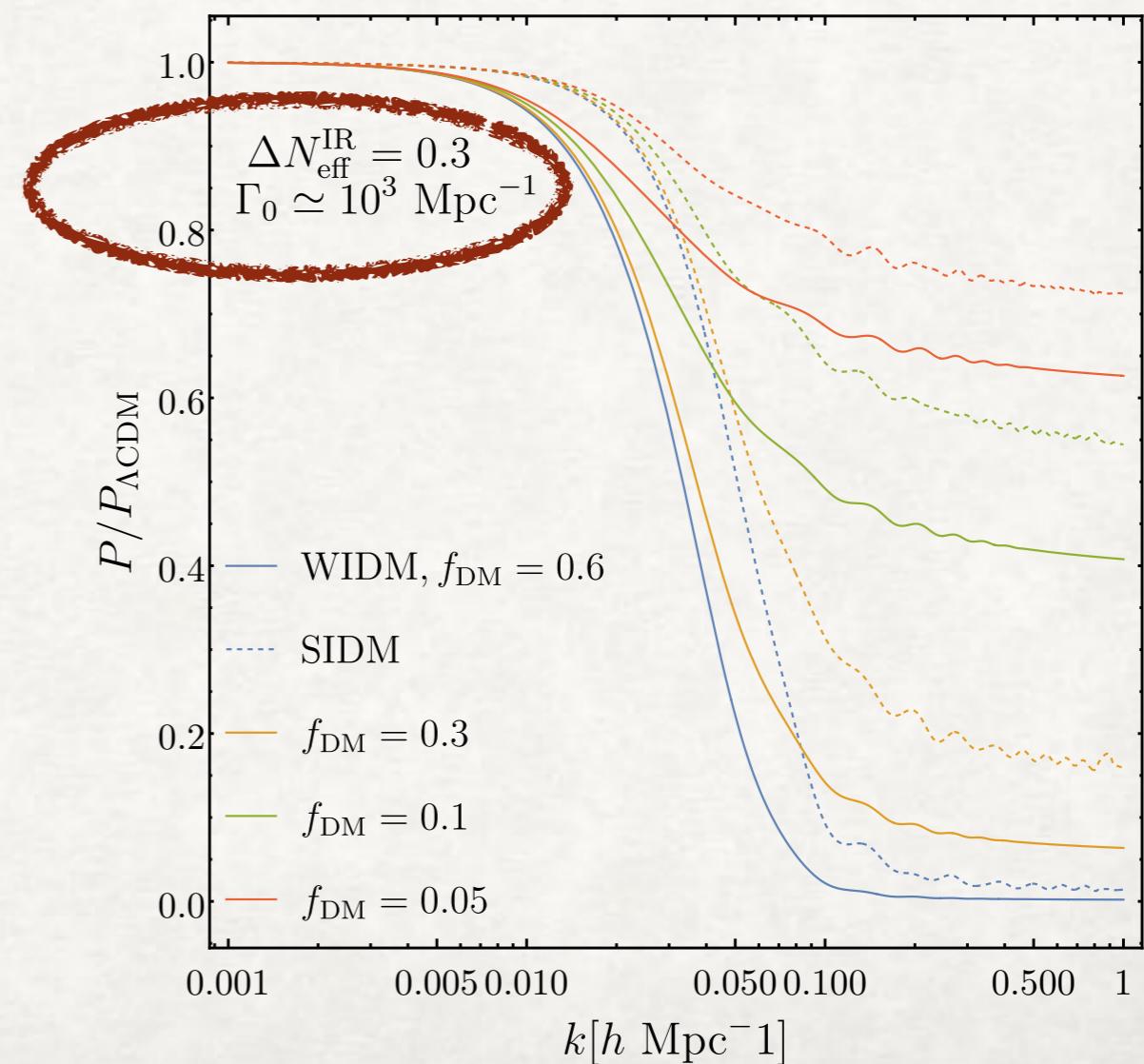
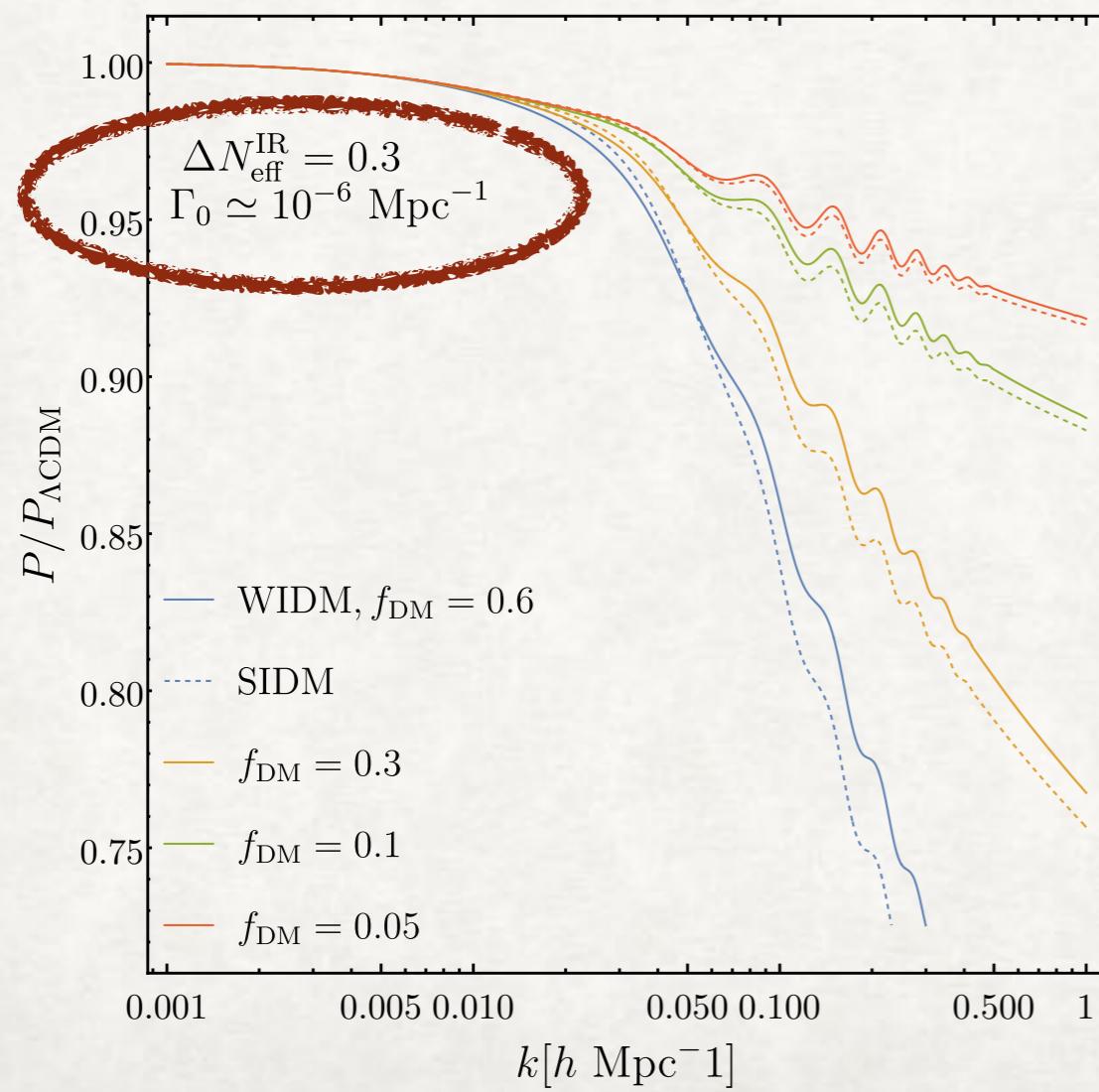


SUPPRESSION OF MATTER POWER SPECTRUM

$$f_{\text{DM}} \equiv$$

$$z \gg z_t$$

$$\frac{\Gamma}{H} \sim \frac{\Gamma_0}{10^{-6} \text{ Mpc}^{-1}} \left(\frac{\Delta N_{\text{eff}}^{\text{IR}}}{0.3} \right)^{1/2} \left(\frac{2}{g_*^{\text{IR}}} \right)^{1/2}$$



SUPPRESSION OF MATTER POWER SPECTRUM

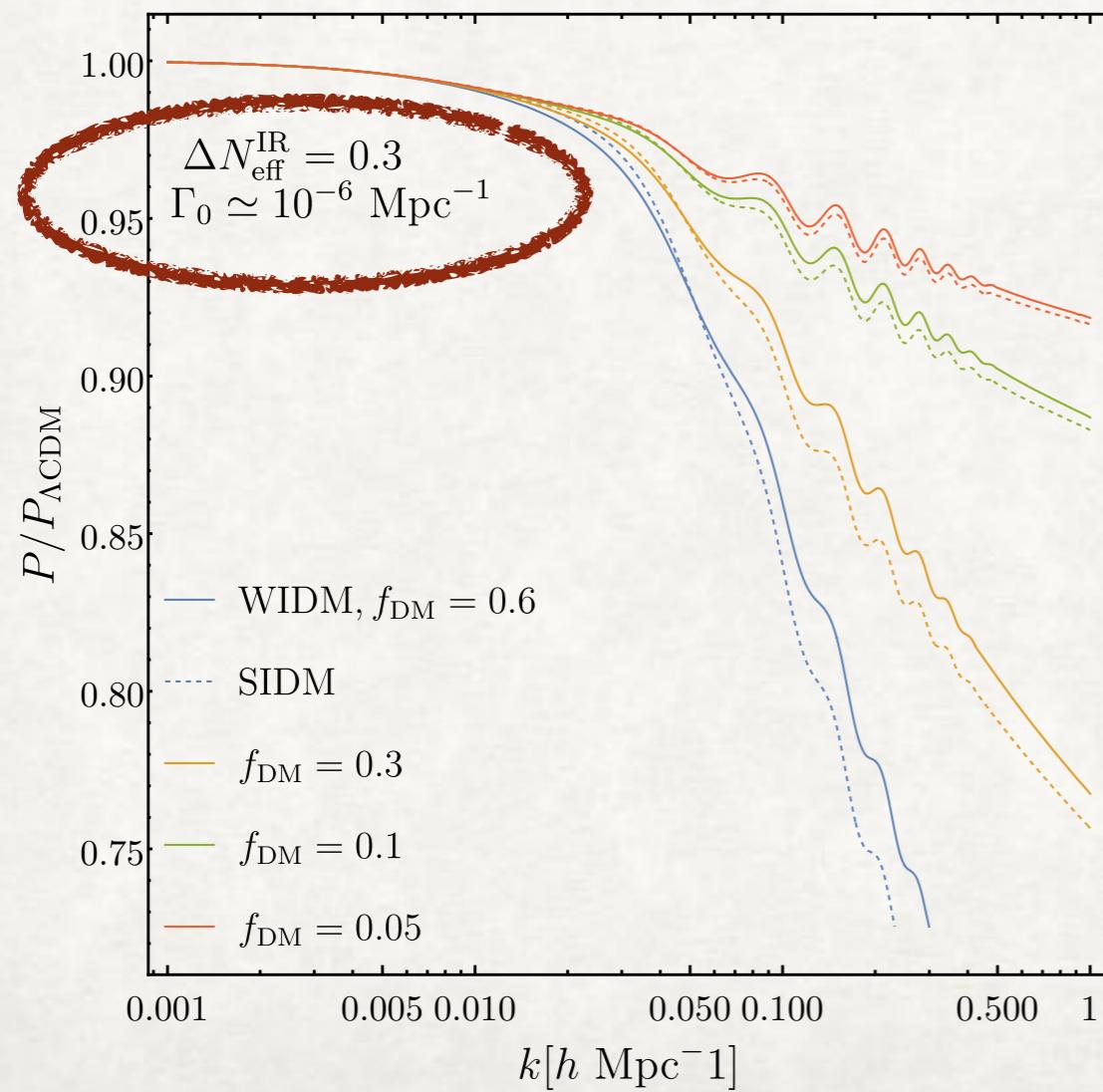
$$f_{\text{DM}} \equiv$$

*Interacting
fraction of CDM*

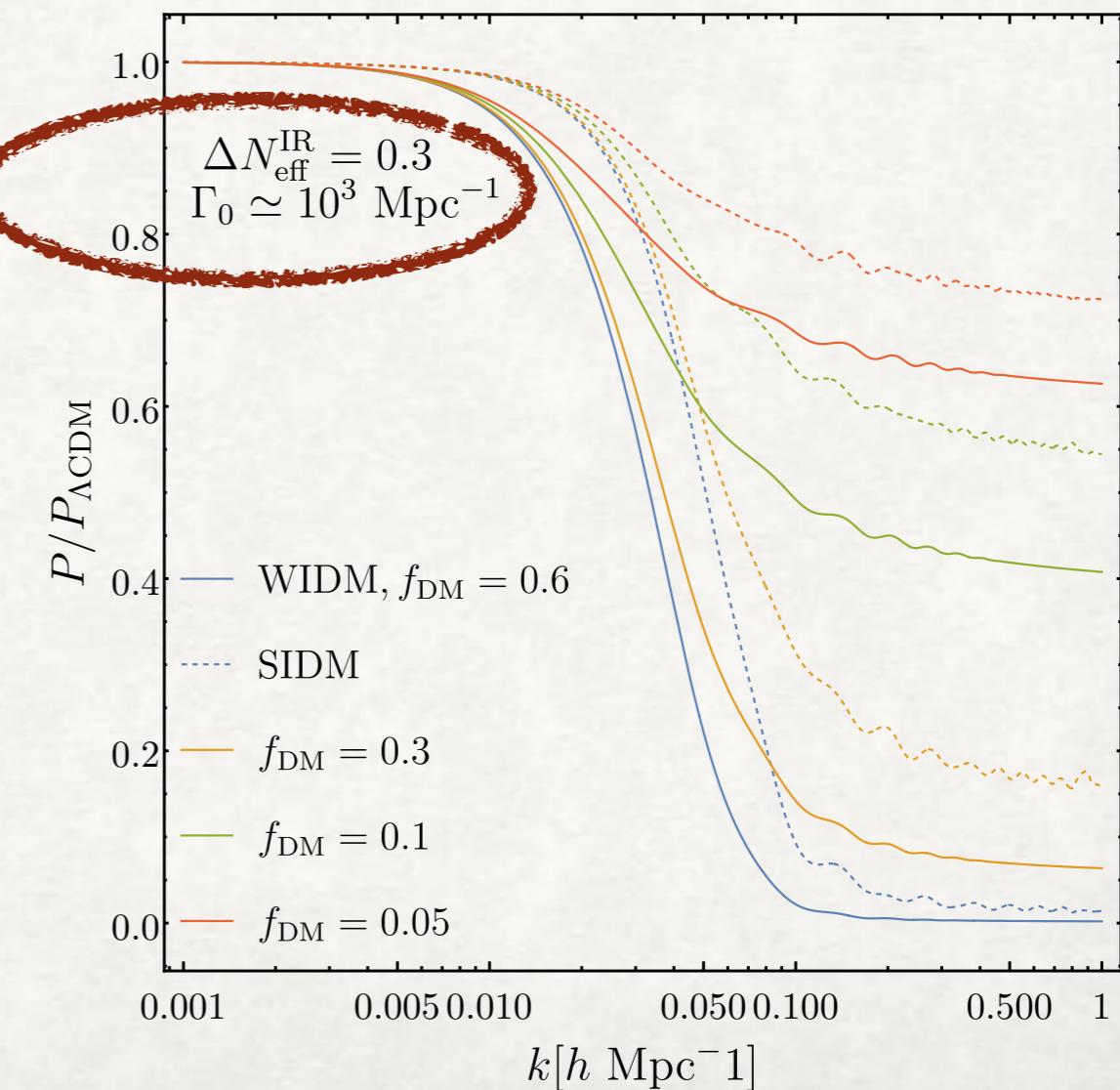
For $z \gg z_t$

$$\frac{\Gamma}{H} \sim \frac{\Gamma_0}{10^{-6} \text{ Mpc}^{-1}} \left(\frac{\Delta N_{\text{eff}}^{\text{IR}}}{0.3} \right)^{1/2} \left(\frac{2}{g_*^{\text{IR}}} \right)^{1/2}$$

Barely efficient, similar suppression



Efficient at early times, different suppression



CONSTRAINTS FROM CMB+LSS

$$\log_{10} \Gamma_0 \in [-9, -5] \quad \log_{10} \Gamma_0 \in [-2, 6]$$

| Parameter | WIDM | SIDM |
|-------------------------------------|---------------------------------------|---------------------------------------|
| $\Delta N_{\text{eff}}^{\text{IR}}$ | < 0.531 (0.092) | < 0.519 (0.268) |
| $\log_{10} z_t$ | Unconstrained (4.18) | Unconstrained (4.38) |
| r_g | Unconstrained (4.87) | Unconstrained (4.39) |
| $\log_{10} \Gamma_0$ | < -6.156 (-8.231) | < 4.259 (3.723) |
| $\log_{10} f_{\text{DM}}$ | Unconstrained (-0.806) | < -2.031 (-3.903) |
| $H_0 [\text{km/s/Mpc}]$ | $68.97 (68.37)^{+0.65}_{-1.1}$ | $68.96 (69.45)^{+0.67}_{-1.1}$ |
| S_8 | $0.818 (0.826)^{+0.011}_{-0.011}$ | $0.820 (0.828)^{+0.011}_{-0.011}$ |
| M_b | $-19.379 (-19.396)^{+0.019}_{-0.031}$ | $-19.380 (-19.364)^{+0.020}_{-0.032}$ |
| $\Delta\chi^2$ | -0.63 | -1.04 |
| $Q_{\text{DMAP}}^{S_8}$ | 2.75σ | 2.55σ |
| S_8 GT | 2.82σ | 2.89σ |
| S_8 IT | 2.63σ | 2.65σ |
| ΔAIC^{S_8} | 7.9 | 6.43 |

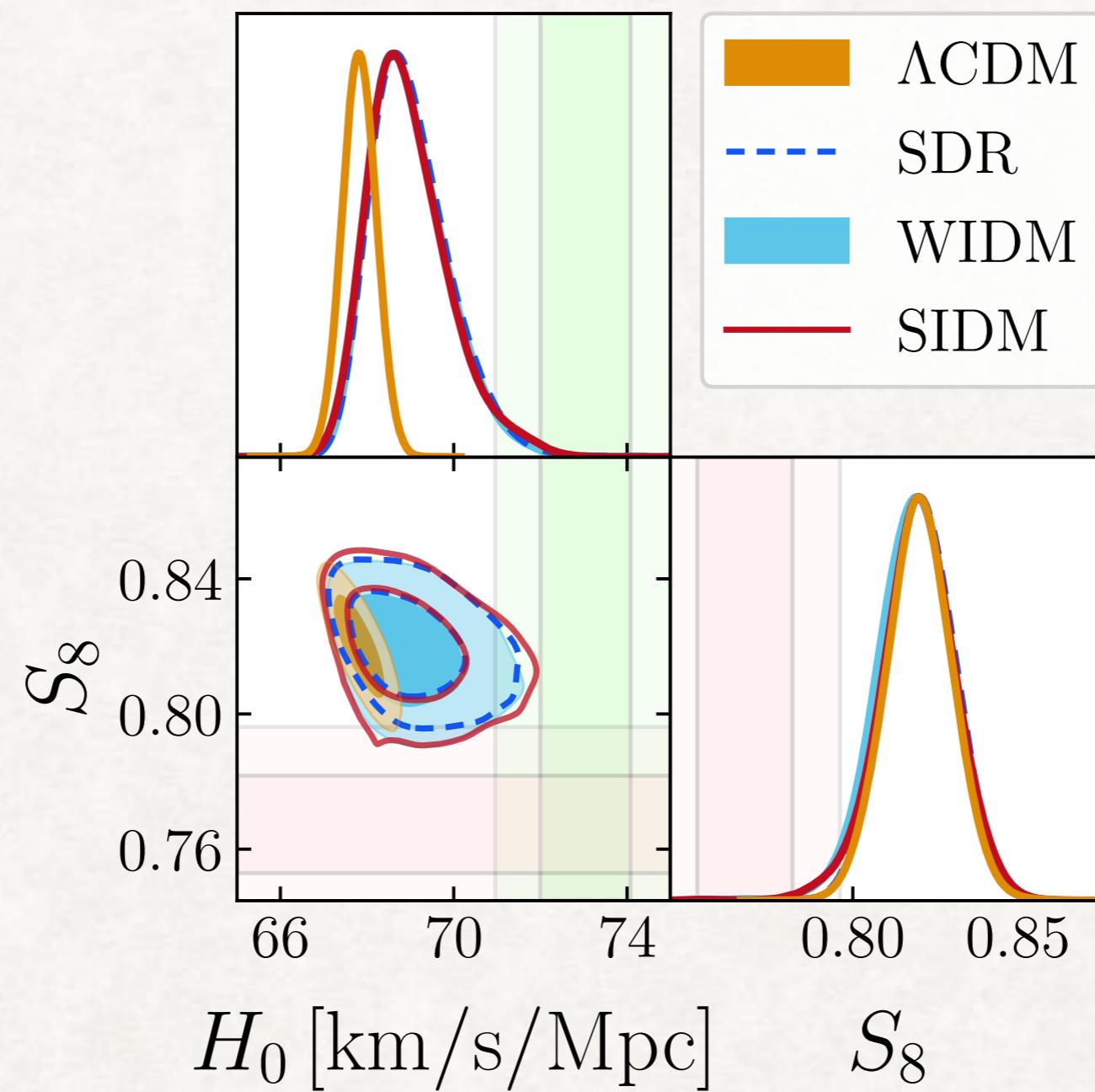
CONSTRAINTS FROM CMB+LSS

Model has 2 extra parameters in addition to stepped DR (5 total)

| Priors | $\log_{10} \Gamma_0 \in [-9, -5]$ | $\log_{10} \Gamma_0 \in [-2, 6]$ |
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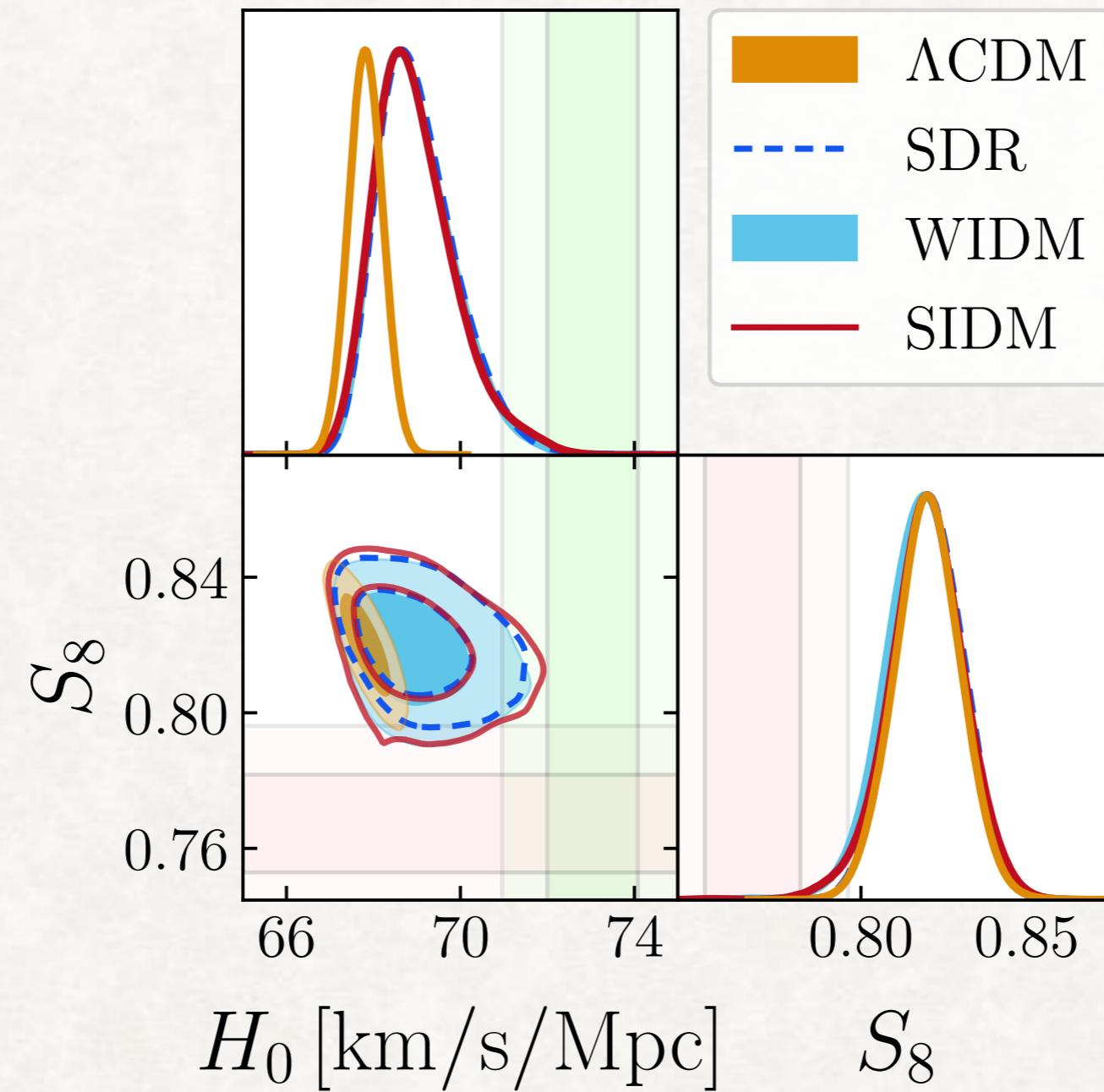
Strong bound on interacting DM fraction in SIDM (<1%)

IMPACT ON S8



IMPACT ON S8

*Old KiDS+DES
measurement*



Minor impact on S8 w.r.t. model w/o interactions

SUMMARY

Interacting sectors with mass thresholds are interesting examples of more complex (realistic) dark sectors

Provide way to alleviate constraints on ΔN_{eff}

CMB+LSS data constrain model ability to address tension(s)

BACK-UP SLIDES

Joseph, Aloni,
Schmaltz, Sivarajan,
Weiner 22

ADDING INTERACTIONS WITH CDM

$$\begin{aligned}\dot{\theta}_{\text{sdr}} &= -\mathcal{H}(1-3w)\theta_{\text{sdr}} - \frac{w}{1+w}\theta_{\text{sdr}} + \frac{\delta P/\delta\rho}{1+w}k^2\delta_{\text{sdr}} - a\Gamma\frac{\rho_{\text{idm}}}{\rho_{\text{sdr}}(1+w)}(\theta_{\text{sdr}} - \theta_{\text{idm}}) \\ \dot{\theta}_{\text{idm}} &= -\mathcal{H}\theta_{\text{idm}} - a\Gamma(\theta_{\text{sdr}} - \theta_{\text{idm}}),\end{aligned}$$

ADDING INTERACTIONS WITH CDM

Consider heavy CDM component (idm) to be interacting with light dark sector (sdr)

$$\begin{aligned}\dot{\theta}_{\text{sdr}} &= -\mathcal{H}(1-3w)\theta_{\text{sdr}} - \frac{w}{1+w}\theta_{\text{sdr}} + \frac{\delta P/\delta\rho}{1+w}k^2\delta_{\text{sdr}} - a\Gamma\frac{\rho_{\text{idm}}}{\rho_{\text{sdr}}(1+w)}(\theta_{\text{sdr}} - \theta_{\text{idm}}) \\ \dot{\theta}_{\text{idm}} &= -\mathcal{H}\theta_{\text{idm}} - a\Gamma(\theta_{\text{sdr}} - \theta_{\text{idm}}),\end{aligned}$$

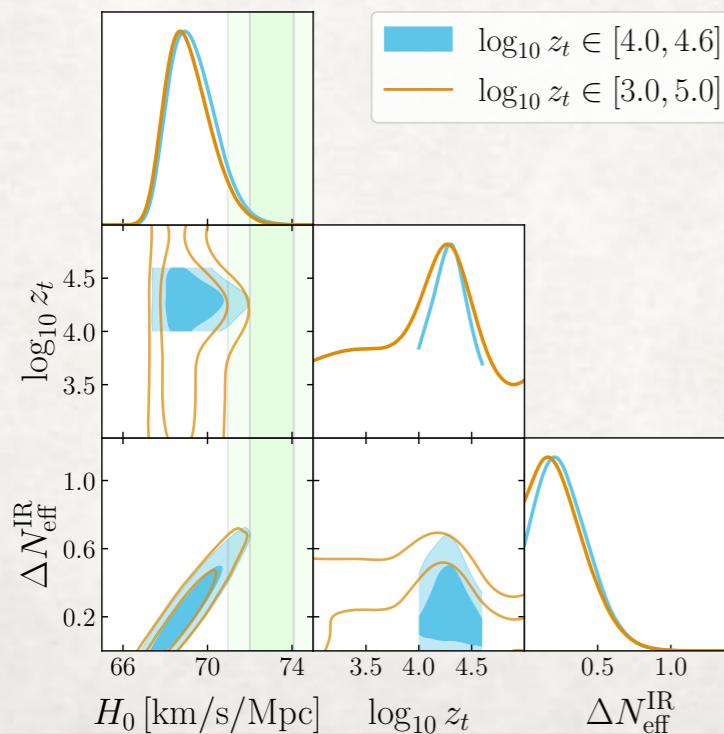
Motivation from S8: suppression of matter power spectrum

However: (2-3)sigma tension on S8 between weak lensing and CMB severely reduced with new KiDS+DES analysis...

Independently, can put constraints

PRIOR DEPENDENCE

| Parameter | SIDR | SDR (i) | SDR (ii) |
|-------------------------------------|---------------------------------------|---------------------------------------|---------------------------------------|
| $\Delta N_{\text{eff}}^{\text{IR}}$ | < 0.456 (95% CL) | < 0.597 (0.551) | < 0.546 (0.289) |
| H_0 [km/s/Mpc] | $68.95 (68.38)^{+0.73}_{-1.2}$ | $69.30 (70.68)^{+0.86}_{-1.3}$ | $68.89 (69.34)^{+0.71}_{-1.1}$ |
| S_8 | $0.823 (0.818)^{+0.011}_{-0.011}$ | $0.829 (0.837)^{+0.011}_{-0.011}$ | $0.827 (0.834)^{+0.011}_{-0.011}$ |
| M_b | $-19.380 (-19.397)^{+0.022}_{-0.034}$ | $-19.369 (-19.325)^{+0.023}_{-0.038}$ | $-19.381 (-19.369)^{+0.021}_{-0.032}$ |
| $\Delta\chi^2$ | -0.22 | -0.41 | -1.4 |
| $Q_{\text{DMAP}}^{M_b}$ | 3.48σ | 2.55σ | 2.74σ |
| M_b GT | 3.66σ | 3.12σ | 3.74σ |
| M_b IT | 2.95σ | 2.56σ | 3.03σ |
| ΔAIC^{M_b} | -18.88 | -22.67 | -20.67 |



(i) r_g fixed, $\log_{10} z_t \in [4.0, 4.6]$
(ii) r_g free, $\log_{10} z_t \in [3.0, 5.0]$

$\gtrsim 3\sigma$

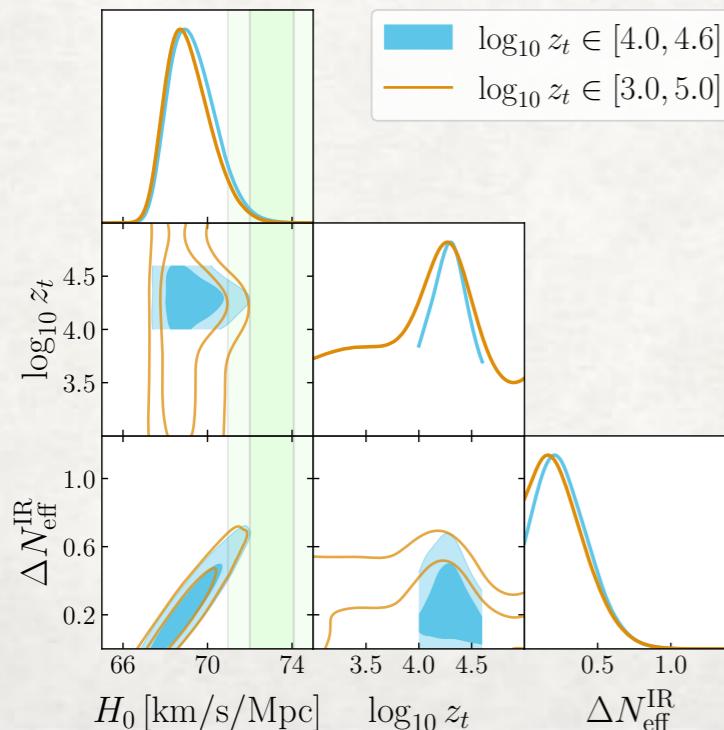
PRIOR DEPENDENCE

Planck18+BAO+Pantheon

W/o threshold

W/ threshold

| Parameter | SIDR | SDR (i) | SDR (ii) |
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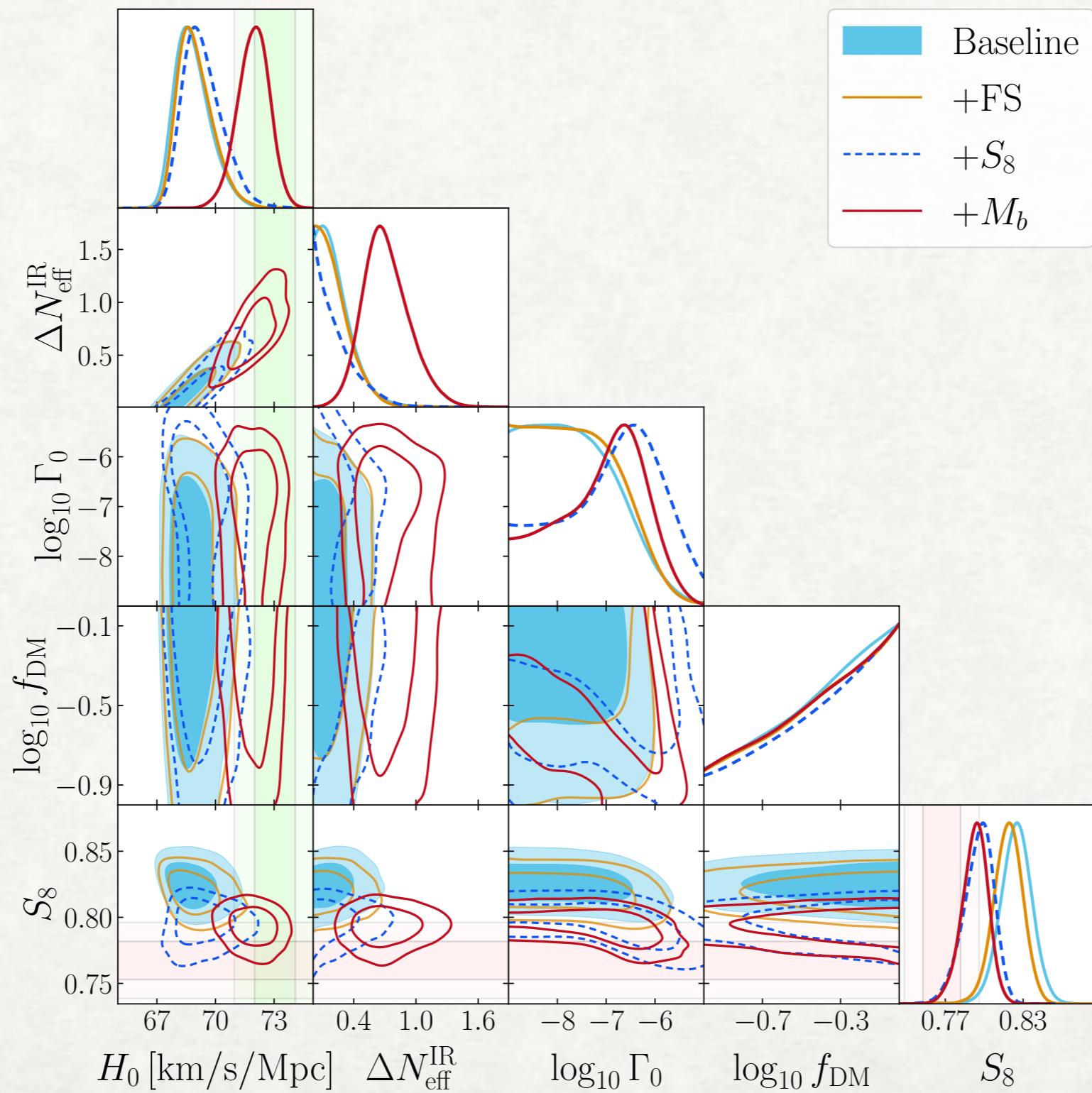


- (i) r_g fixed, $\log_{10} z_t \in [4.0, 4.6]$
(ii) r_g free, $\log_{10} z_t \in [3.0, 5.0]$

Better fit, increased tensions

$\gtrsim 3\sigma$

WIDM POSTERIORS



SIDM POSTERIORS

