

# Neutrinos physics

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- Introduction
- Neutrino oscillations
- The nature of the neutrino

*Highly selective in terms of exp. results (and theoretical details ...)*



# The masses and the nature of the neutrinos



L'ÉLÉMENTAIRE

# What have we learnt from the oscillation experiments ?

Two basis : mass eigenstates ( $\nu_i$ ) and the flavour eigenstates ( $\nu_e, \nu_\mu, \nu_\tau$ )



$\nu_e$  

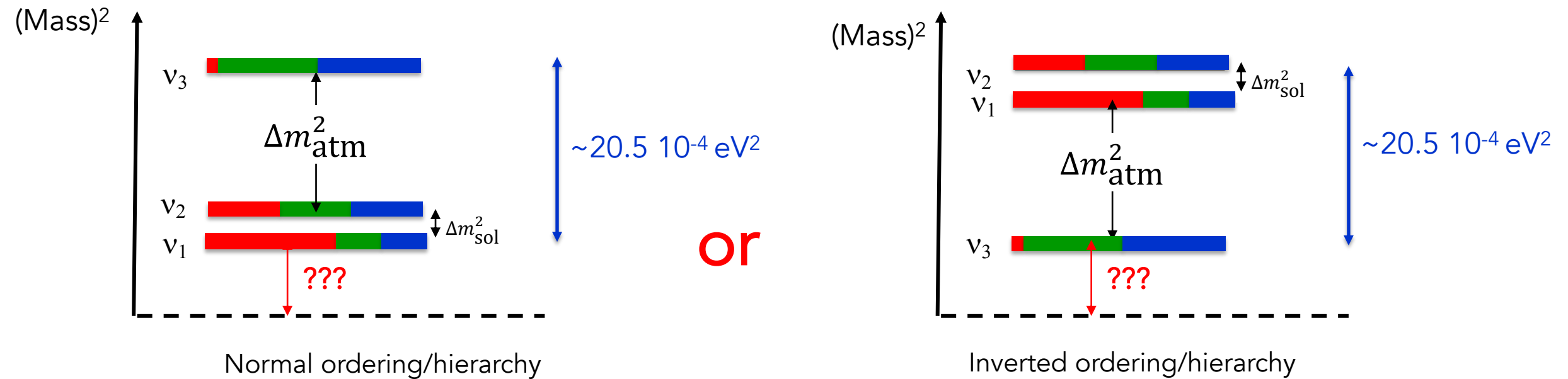
$\nu_\mu$  

$\nu_\tau$  

The relative fractions are given by the mixing angles of the PMNS matrix

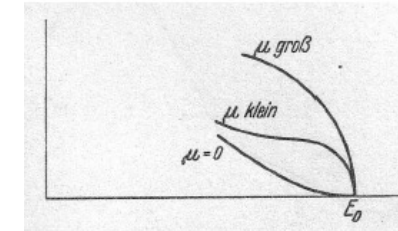
$$\begin{pmatrix} c_{12} c_{13} & s_{12} c_{13} & s_{13} e^{-i\delta_{\text{CP}}} \\ -s_{12} c_{23} - c_{12} s_{13} s_{23} e^{i\delta_{\text{CP}}} & c_{12} c_{23} - s_{12} s_{13} s_{23} e^{i\delta_{\text{CP}}} & c_{13} s_{23} \\ s_{12} s_{23} - c_{12} s_{13} c_{23} e^{i\delta_{\text{CP}}} & -c_{12} s_{23} - s_{12} s_{13} c_{23} e^{i\delta_{\text{CP}}} & c_{13} c_{23} \end{pmatrix}$$

# What have we learnt from the oscillation experiments ?



Neutrinos have a mass (at least one with  $m > 0.04 \text{ eV}$  !)

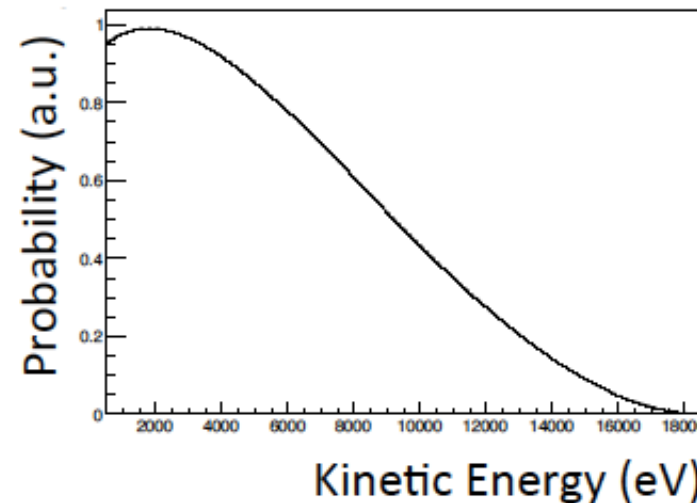
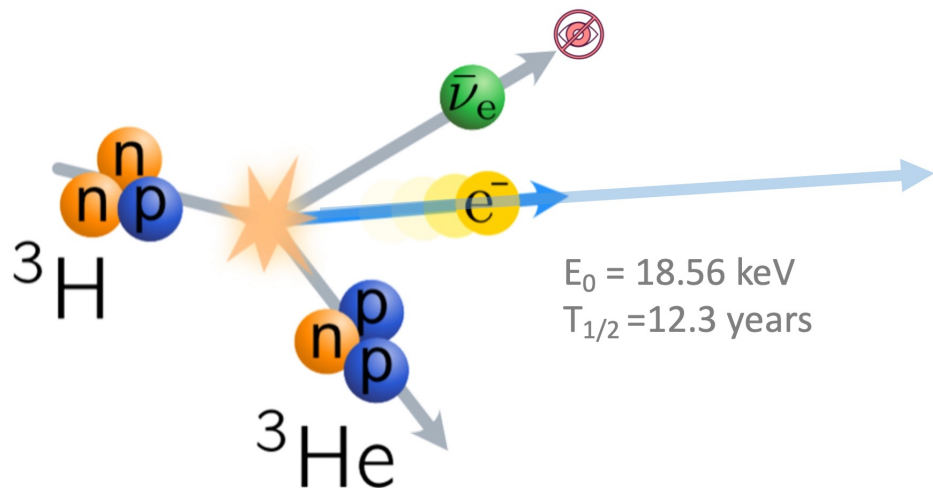
# Neutrinos masses (direct) measurements :



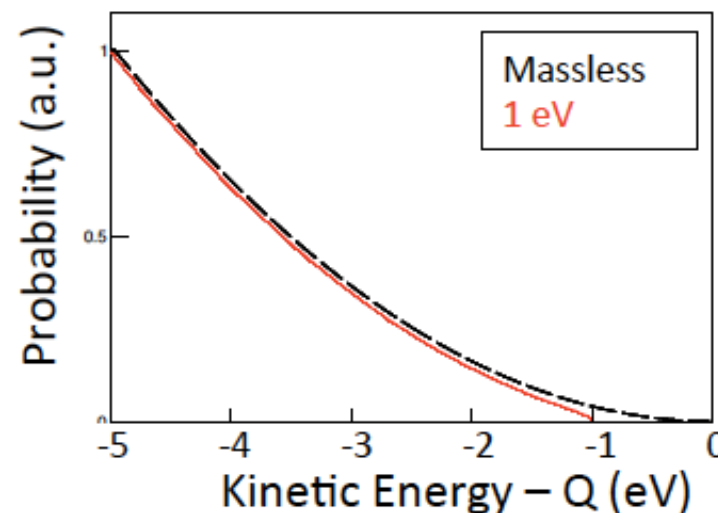
Spectrum of electrons (Kurie Plot) emitted in  $\beta$  decay

$$K(E) \propto (E_0 - E) \left[ 1 - \frac{m_\nu^2}{(E_0 - E)^2} \right]^{1/4}$$

for a massive neutrino



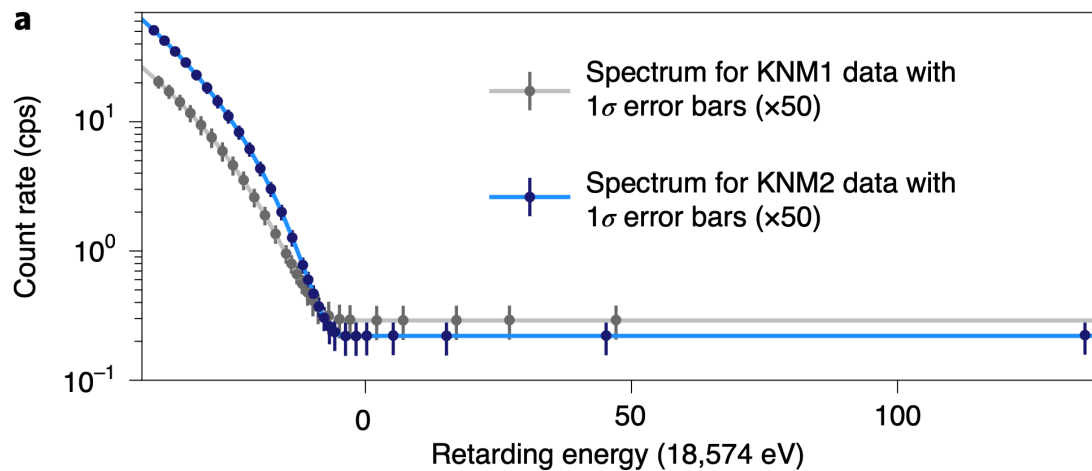
Shape of  $\beta$  spectrum near the end point



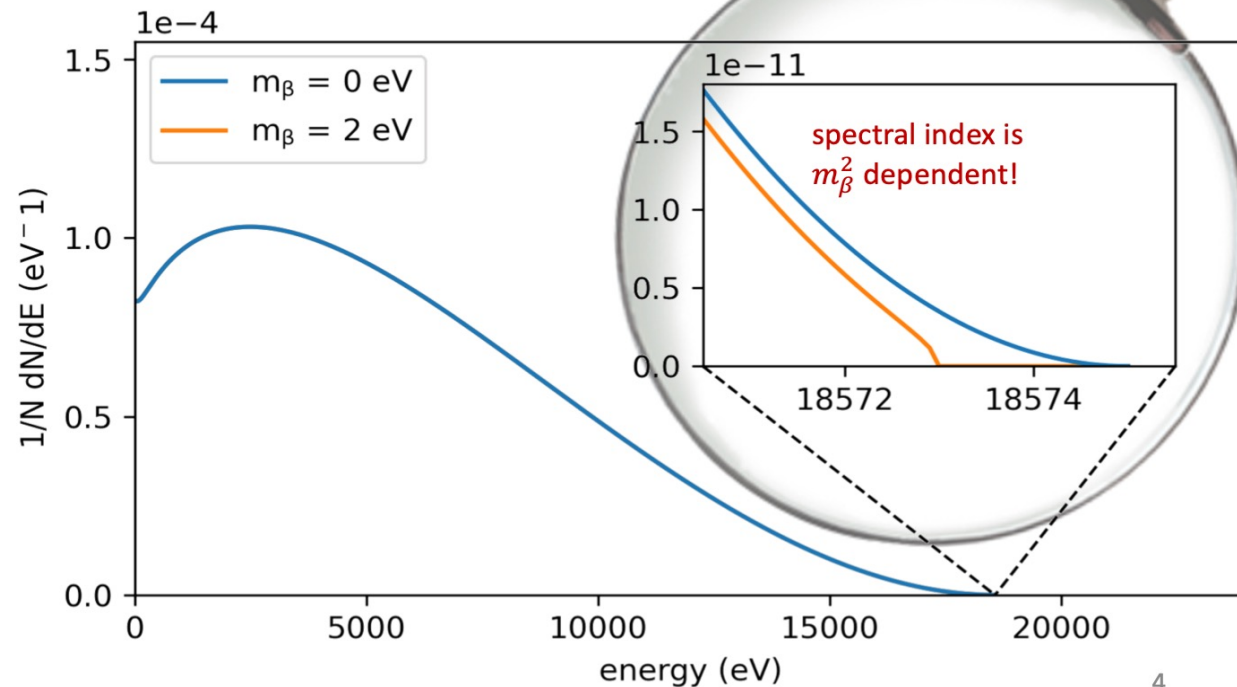


Karlsruhe  
Tritium  
Neutrino  
Experiment

High intensity source (tritium  $10^{11}$  decays/s)  
 Low background  
 Excellent energy resolution (1 eV)  
 High stability



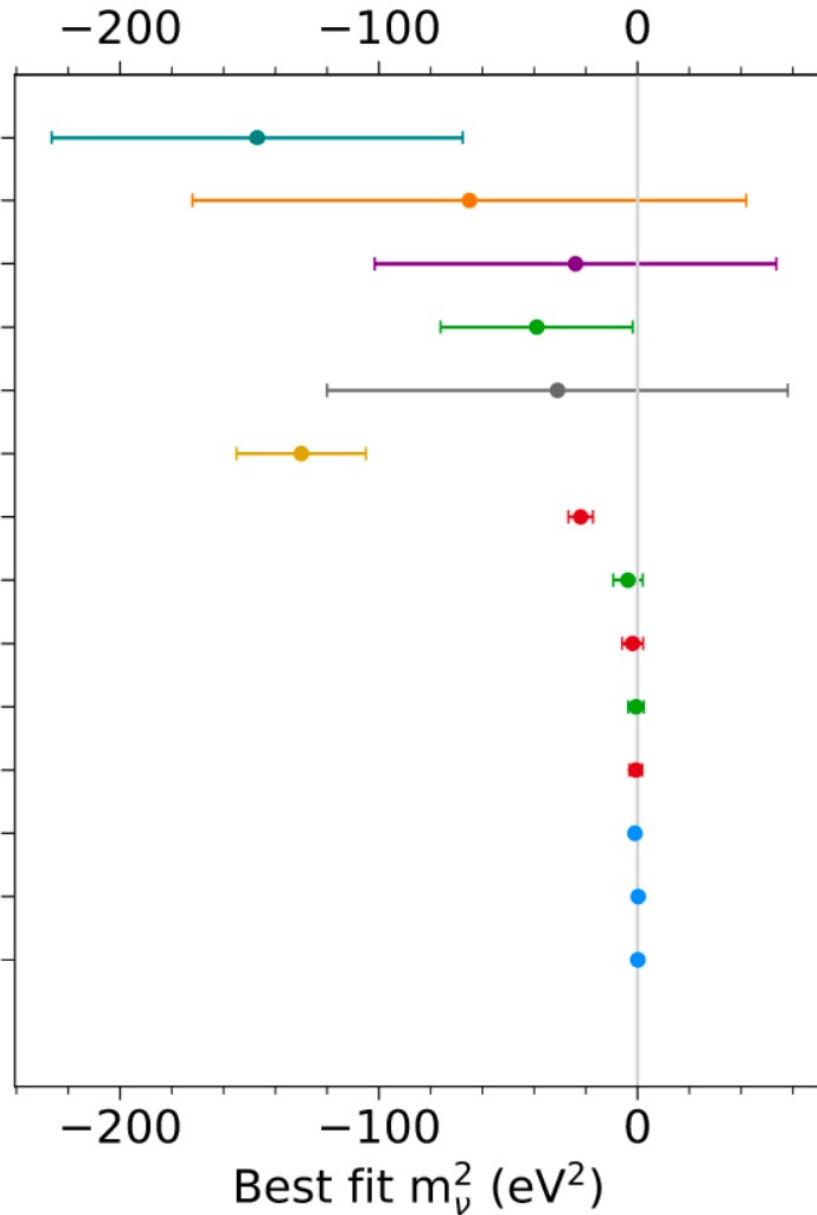
$10^{-8}$  of all  
 decays in  
 last 40 eV



$$m(\nu_e) < 0.8 \text{ eV @ 90\% CL}$$

**Experiment (year)**

- Los Alamos (91)
- Tokyo (91)
- Zürich (92)
- Mainz (93)
- Beijing (93)
- Livermore (95)
- Troitsk (95)
- Mainz (99)
- Troitsk (99)
- Mainz (05)
- Troitsk (11)
- KATRIN (19)
- KATRIN (21)
- KATRIN (comb.)



K. Valerius, KIT, DPG Dortmund (2021)

$$K(E) \propto (E_0 - E) \left[ 1 - \frac{m_\nu^2}{(E_0 - E)^2} \right]^{1/4}$$

expect to reach a limit of 0.2 eV

# Using time of flight for a direct measurement of neutrino mass

Distance to Earth :  $1.7 \cdot 10^5$  light years ( $d \sim 1.5 \cdot 10^{21}$  m)

A burst of  $\bar{\nu}_e$

IMB@Brookhaven & Kamiokande in Japan

~10 neutrinos per detector

datasets cannot be merged due to no use of absolute time

$$\Delta t = d \times \left( \frac{1}{v_2} - \frac{1}{v_1} \right) = \frac{dm^2 c^4}{2} \times \left( \frac{1}{E_2^2} - \frac{1}{E_1^2} \right)$$

$$m = E_1 E_2 \sqrt{\frac{2\Delta t}{d} \times c \times \frac{1}{E_1^2 - E_2^2}}$$





TABLE III. Characteristics of the contained neutrino events recorded on 23 February.

Event No. <sup>a</sup>	Time (UT)	No. of PMT's	Energy <sup>b</sup> (MeV)	Angular distribution <sup>c</sup> (degrees)
33162	7:35:41.37	47	38	74
33164	7:35:41.79	61	37	52
33167	7:35:42.02	49	40	56
33168	7:35:42.52	60	35	63
33170	7:35:42.94	52	29	40
33173	7:35:44.06	61	37	52
33179	7:35:46.38	44	20	39
33184	7:35:46.96	45	24	102

<sup>a</sup>The event numbers are not sequential. Interspersed with the contained neutrino events are fifteen entering cosmic-ray muons.

<sup>b</sup>Error in energy determination is  $\pm 25\%$  (systematic plus statistical).

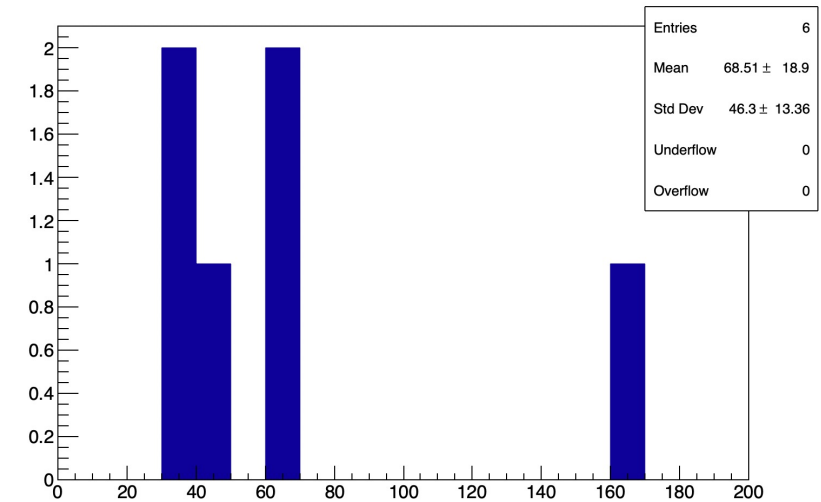
<sup>c</sup>Individual track reconstruction uncertainty is  $15^\circ$ . Note that this angular distribution will be systematically biased toward the source because of the location of the inoperative PMT's.

Our events and their proximity in time to the optical observation of the supernova 1987A are compelling evidence that neutrinos have been seen from a supernova collapse. It is clear that these data can be used to study fundamental properties of supernovae and neutrinos; further calculations are in progress.

$$m = E_1 E_2 \sqrt{\frac{2\Delta t}{d} \times c \times \frac{1}{E_1^2 - E_2^2}}$$

Use first & last:  
 $m_\nu < 46 \text{ eV}$

Use all (but #33167 ..):



Mean 68 eV

persed.<sup>17</sup> The typical duration of the delayed burst from the LMC supernova can be estimated as<sup>17</sup>

$$\left( \frac{R}{10 \text{ kpc}} \right) \left( \frac{m}{10 \text{ eV}} \right)^2 \left( \frac{\langle E \rangle}{10 \text{ MeV}} \right)^{-2} \text{ sec},$$

where the distance of the LMC is  $R = 50 \text{ kpc}$ , and the mean energy of eleven neutrinos is  $\langle E \rangle = 16.7 \text{ MeV}$ . From the condition that this duration  $\Delta$  must be smaller than the observed duration of burst, 12.4 sec, we obtain an upper limit on the mass,  $m < 26 \text{ eV}$ .

# The other neutrinos :

$$m_{\nu_\mu}^{\text{eff}} < 190 \text{ keV (90\% CL)}$$

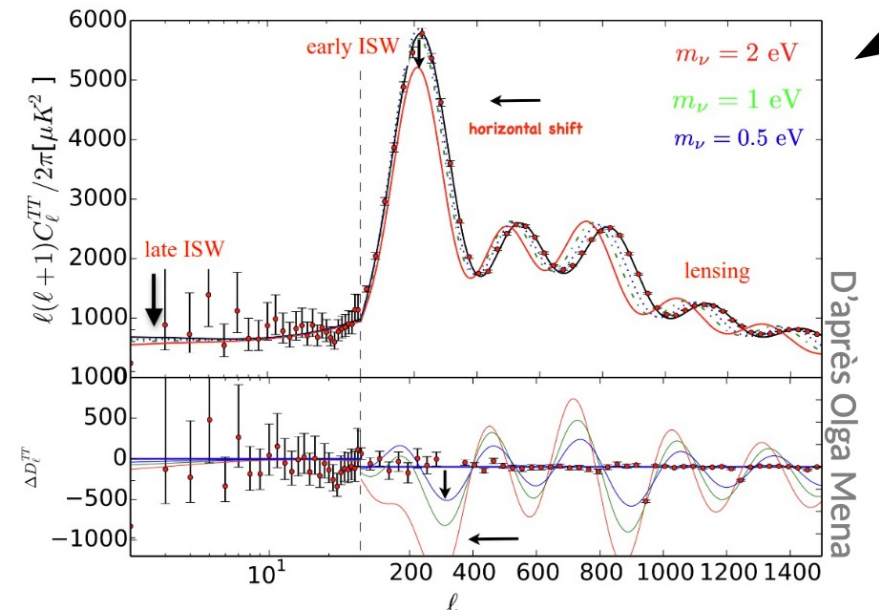
$$m_{\nu_\tau}^{\text{eff}} < 18.2 \text{ MeV (95\% CL)}$$

from  $\pi^- \rightarrow \mu^- + \bar{\nu}_\mu$ ,

from  $\tau^- \rightarrow n\pi + \nu_\tau$ . LEP  $Z^0 \rightarrow \tau\tau$

# Cosmological constraints on $\Sigma m(\nu)$

- The neutrinos mass would modify the (delicate) balance between gravity and the Hubble expansion
- They would also have also affect on the structure formation ...
- Small modifications in the Cosmological Microwave Background which is the fingerprint of what happened at the very beginning of the Universe



→  $\Sigma m(\nu) < \sim 0.1 - 0.2$  eV

PDG review

Neutrino mixing :  
They have a mass  
(at least one with  $m > 0.04$  eV !)

Cosmological constraint:  
 $\Sigma m(\nu) < \sim 0.1 - 0.2$  eV

In the framework with 3 neutrinos, the window for the mass of the heaviest one is not so wide ...

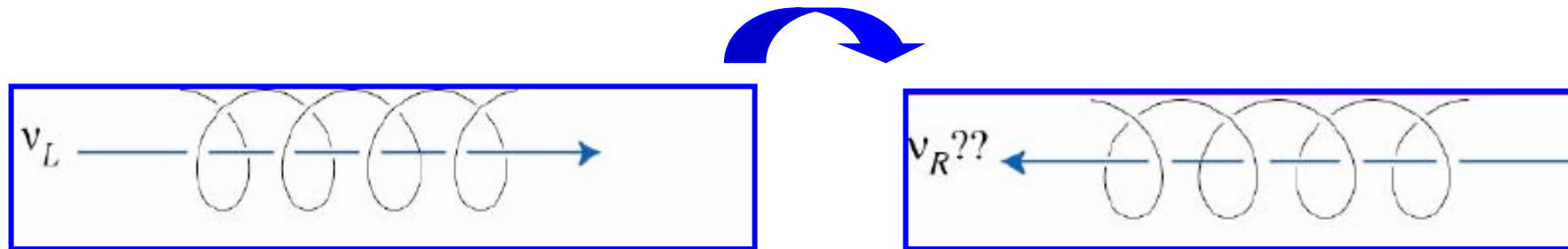
But the masses are tiny (few  $10^{-2}$  eV) compared to  $0.5 \cdot 10^6$  eV for the electron and up to  $10^{11}$  eV for the top quark ... ?

In the Standard Model neutrinos are massless  $\rightarrow$  no right-handed neutrino  $\rightarrow$  no Dirac mass term

Suppose the neutrino has a tiny mass

$\Rightarrow$  it cannot go at the speed of light

$\Rightarrow$  the neutrino can exist in two helicity states (jump into a reference system moving faster than  $v$  to see its helicity flip)



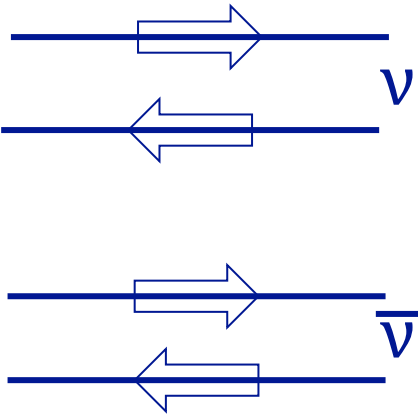
What is this  $v_R$  state?

# Dirac or Majorana neutrinos ?

Dirac



$$\nu \neq \bar{\nu}$$



$\nu_R$  is a distinct state

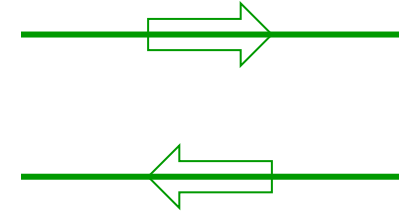
$\begin{pmatrix} \nu_L \\ \nu_R \end{pmatrix}$  a Dirac spinor (like the other fermions)

But coupling to the Higgs  $10^{-5}$  times smaller than those of the electron ?!

Majorana



$$\nu = \bar{\nu}$$



$\nu_R$  is the  $\nu_L$  antineutrino

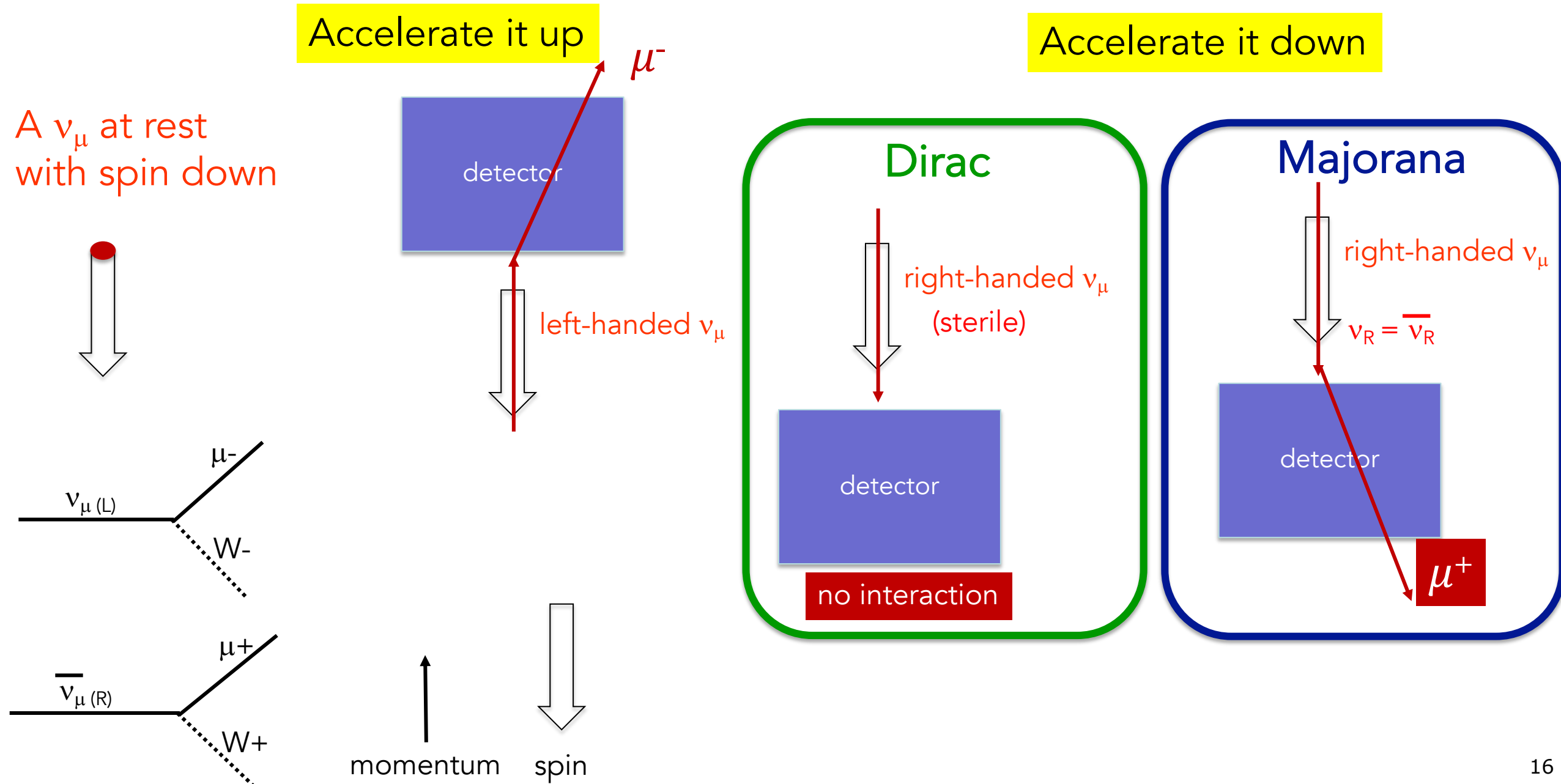
The neutrino cannot carry any charge (since particle and antiparticles carry opposite charges)

*Neutrino and antineutrino are different polarizations of a unique particle which interacts mostly as a neutrino when its spin is anti-parallel to its momentum and as an antineutrino when its spin is parallel to its momentum*

Violation of Lepton number

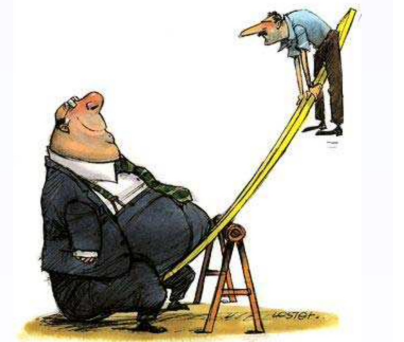
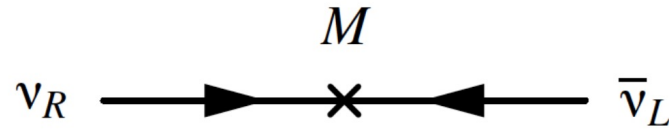
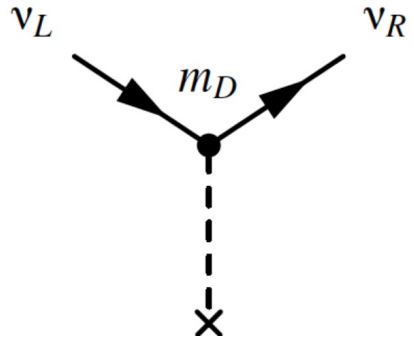
# Dirac or Majorana ? A gedanken experiment

From Strumia-Vissani (hep/ph-0606054v2)





# See-saw mechanism (Majorana neutrinos)



most general form :

a Dirac term + a Majorana term

$$\mathcal{L} \sim -\frac{1}{2} \begin{pmatrix} \bar{\nu}_L & \bar{\nu}_R^c \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D & M \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix}$$

$m_D$  ~ charged leptons & quarks  
 $M$  very large

the physical "mass eigenstates" are those in the basis where the mass matrix is diagonal

$$m_\nu \approx \frac{m_D^2}{M}$$

light left-handed neutrino

$$m_N \approx M$$

heavy right-handed neutrino

$$M \sim 10^{12} - 10^{16} \text{ GeV}$$

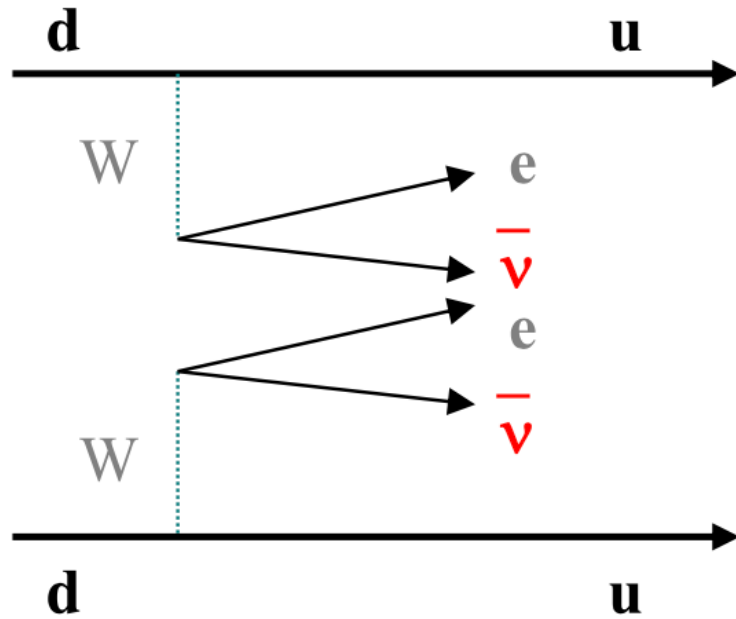
to match with current knowledge for charged fermions and neutrinos masses limits

# Double $\beta$ decay: $2\beta 2\nu$ or $2\beta 0\nu$ ?

$$(A, Z) \rightarrow (A, Z + 2) + 2e^- + 2\bar{\nu}_e$$

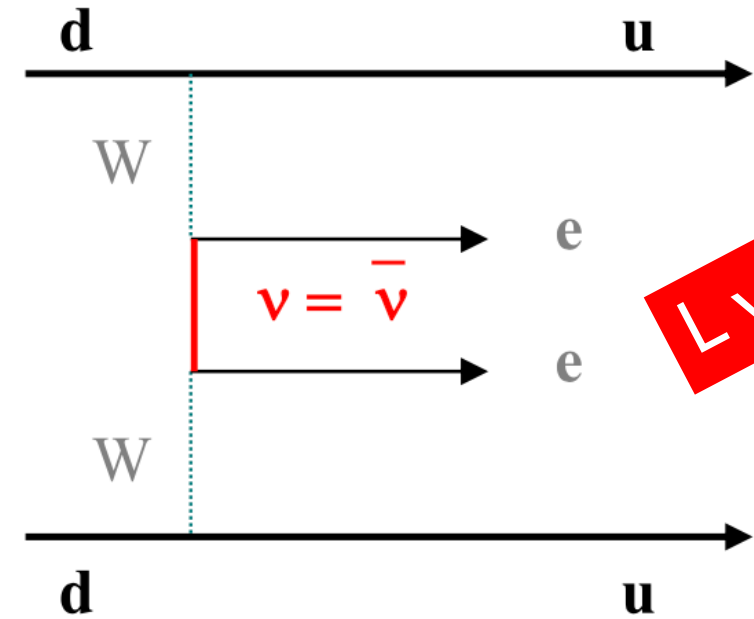
$$(A, Z) \rightarrow (A, Z + 2) + 2e^-$$

inside a nucleus



$(2\beta 2\nu)$

inside a nucleus



L violation

$(2\beta 0\nu)$

Probing the nature of the neutrino !

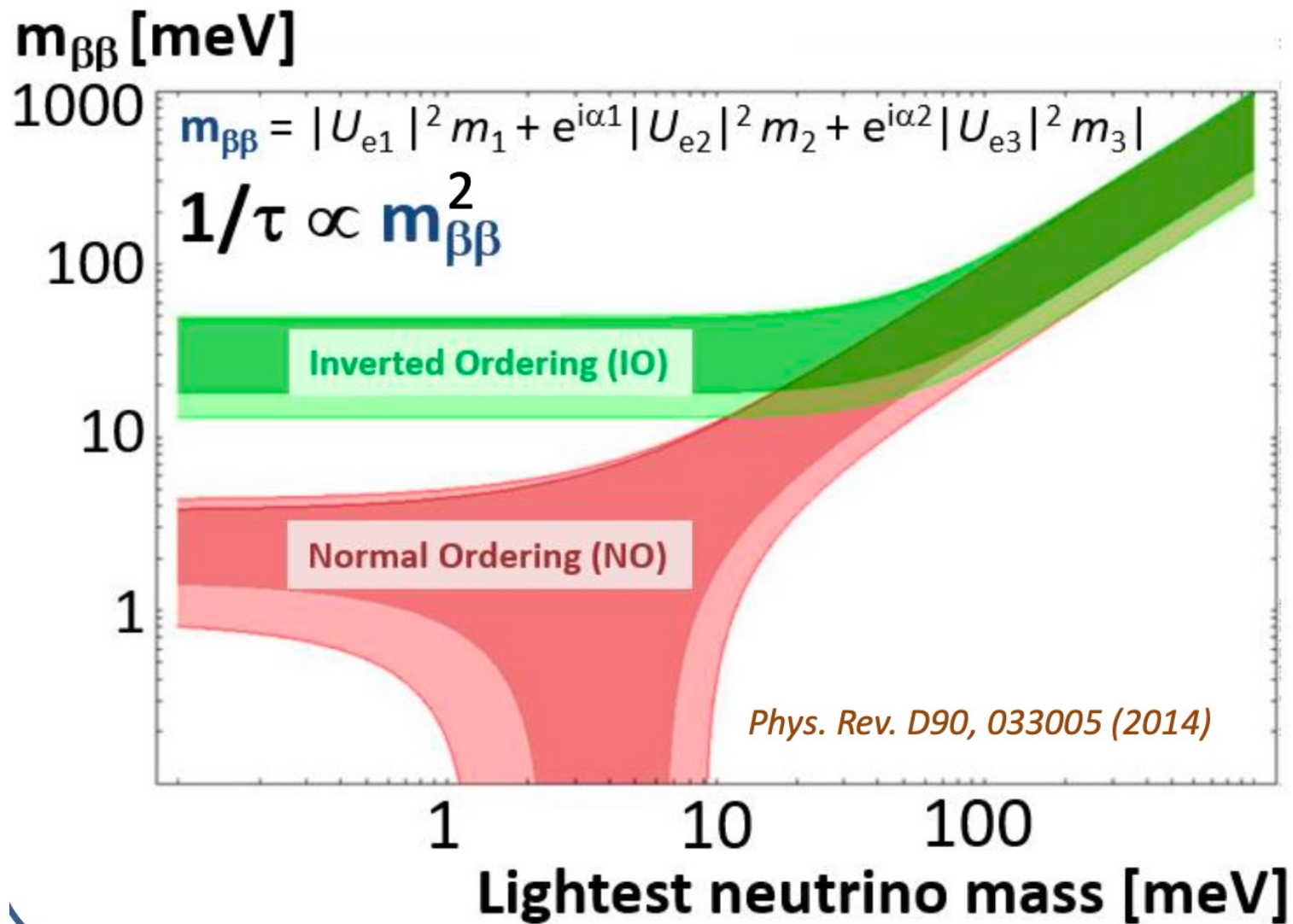
# $2\beta 0\nu$ half-life depends directly on $m(\nu)$

Decay rate

$$\left(T_{1/2}^{0\nu\beta\beta}\right)^{-1} = G^{0\nu\beta\beta} \cdot |NME|^2 \cdot \left|\langle m(\beta\beta) \rangle\right|^2, \quad \langle m(\beta\beta) \rangle = \sum_i U_{ei}^2 m(\nu_i) e^{i\alpha_i}$$

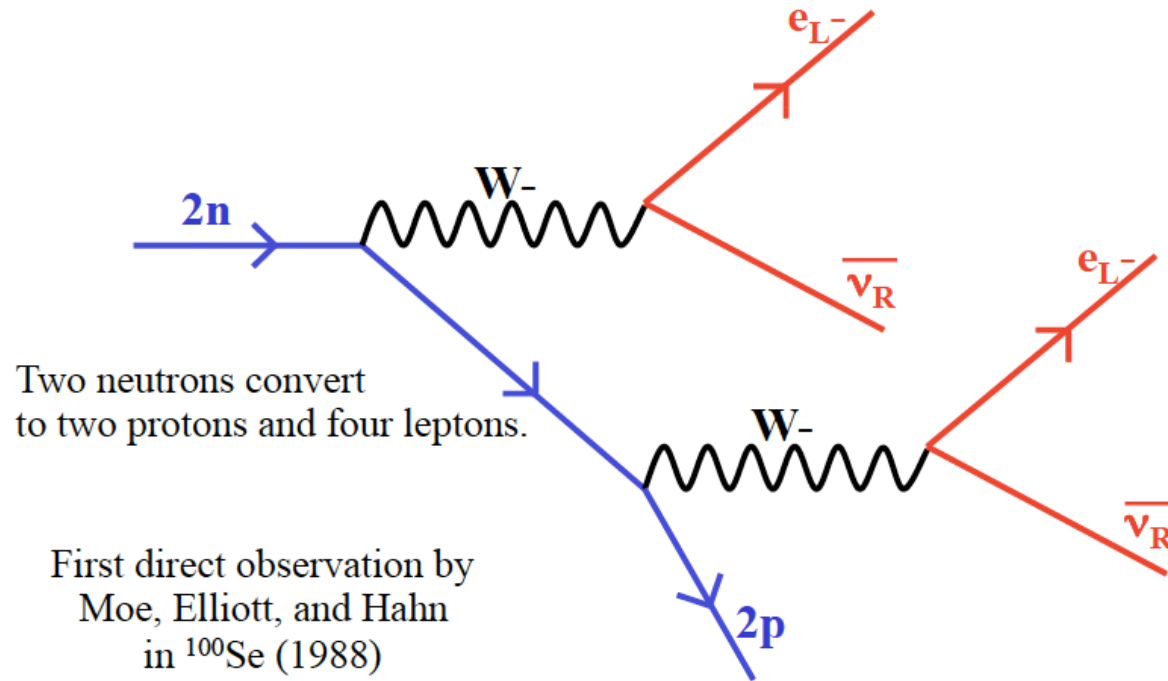
phase space      Nuclear Matrix Element      effective mass      Majorana phases

Measurement of the lifetime allows to measure (or limit) the effective neutrinos mass



Search for  $2\beta 0\nu$

# Double $\beta$ decay



Allowed in the Standard Model with a very small decay rate :

- small value of  $G_F$
- large suppression of the phase-space factor (the small energy release must be shared by four leptons)

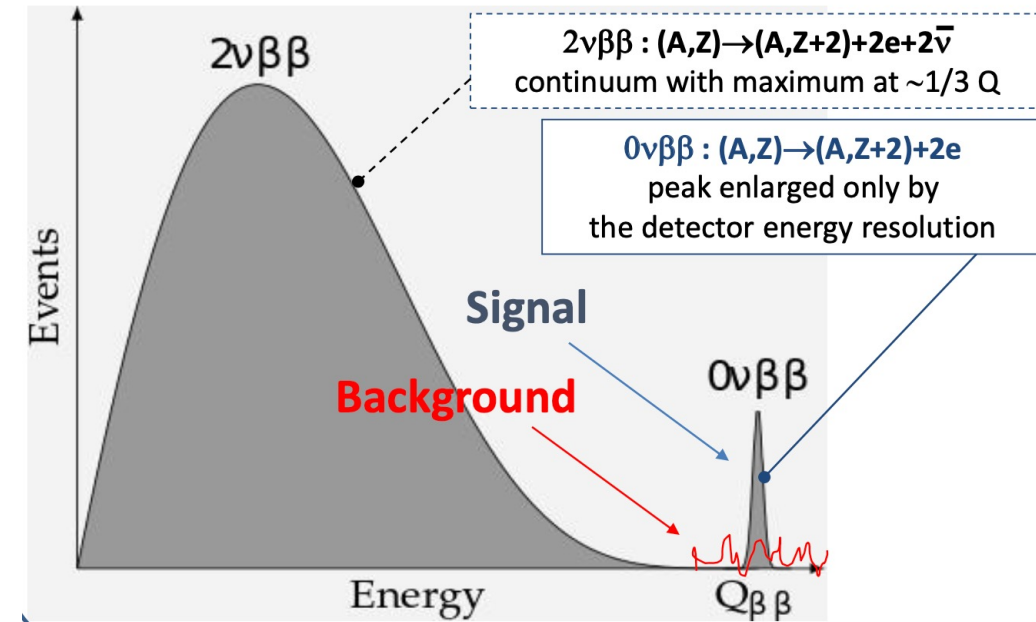
- Use nuclei where the  $\beta$  decay is forbidden
- Use nuclei where  $2\beta$  decay is possible
- Select those for which  $Q_{\beta\beta}$  value is as far from backgrounds as possible
- Study the Kurie plot

$$\frac{1}{\tau} = (\text{phase space}) \times \underbrace{g_A^4 \times (\text{Nuclear matrix element})^2}_{\text{uncertainties ...}} \times m_{\beta\beta}^2$$

$$m_{\beta\beta}^2 = \sum_i [U_{ei}]^2 \times m_i^2$$

Sensitivity  $\propto$  Source Mass x Measurement Time

Measurement of the sum of the 2 electrons energy

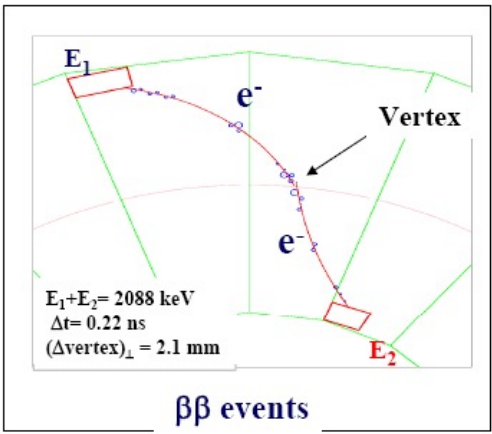


# Source away from the detector

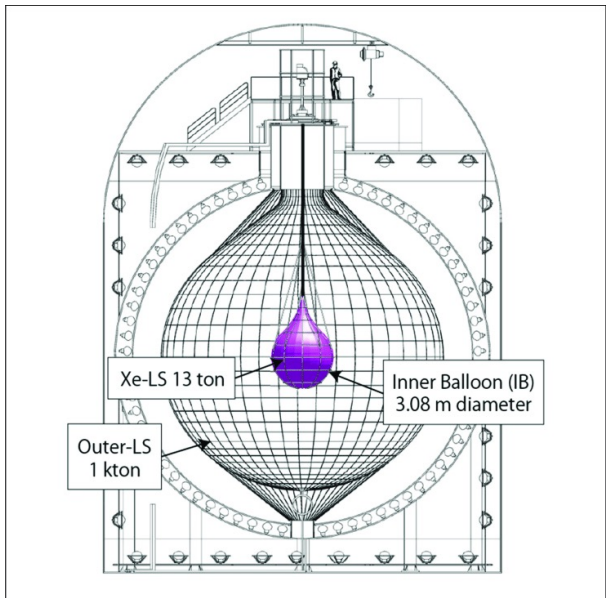
# Source embedded in the detector

Tracking + calorimetry NEMO3 & co

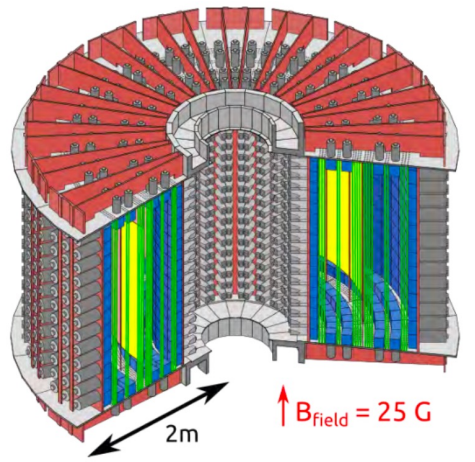
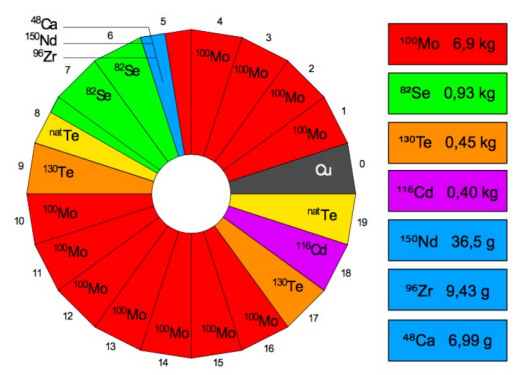
Liquid Scintillator  
(large masses)



KamLAND Zen

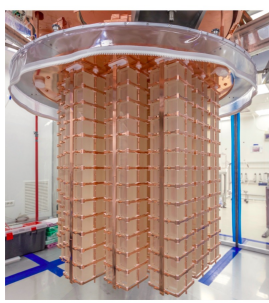


NEMO-3 "camembert" (source top view)



Crystals  
(granularity)

CUORE



GERDA

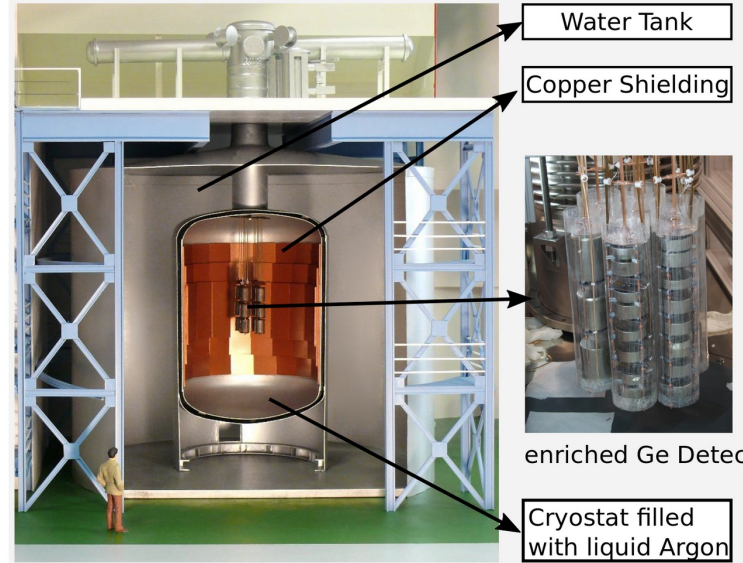
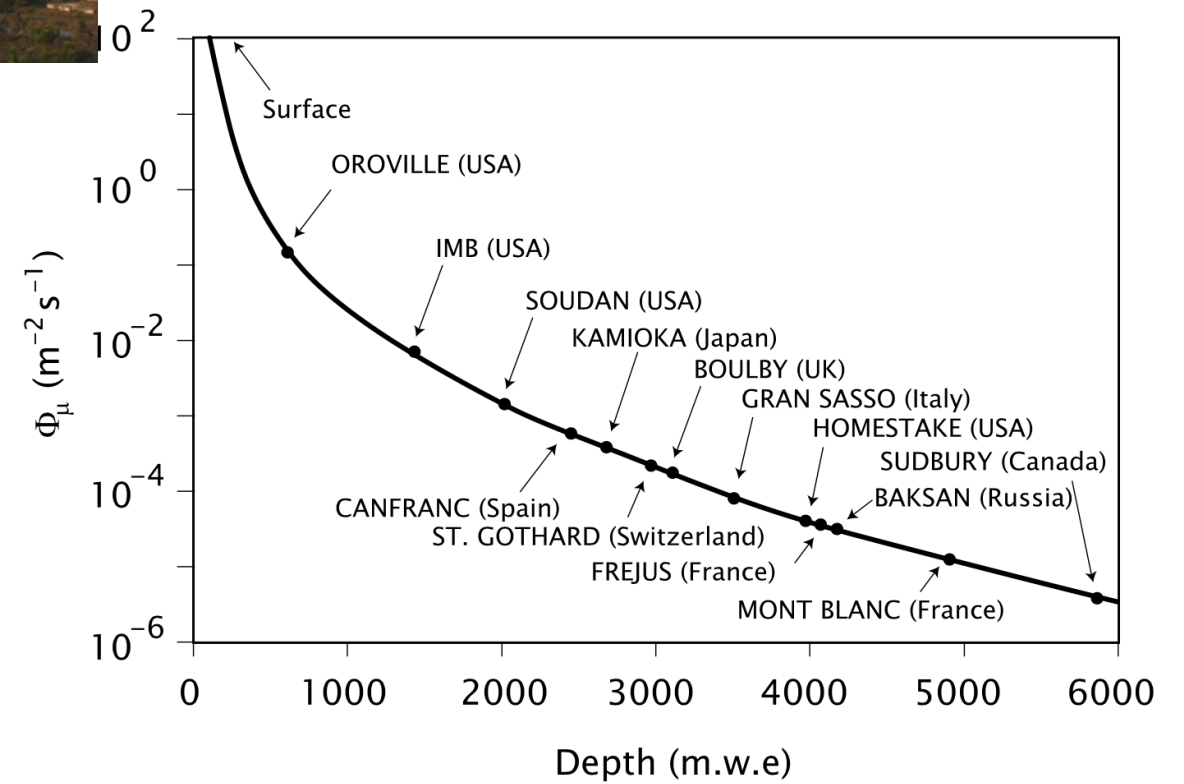


Figure 11. Right: distribution of isotopes in the NEMO experiment. Left: Schematic representation of the experiment including the sources (yellow), the tracker (green) and the calorimeter (blue). Figures from the NEMO homepage.

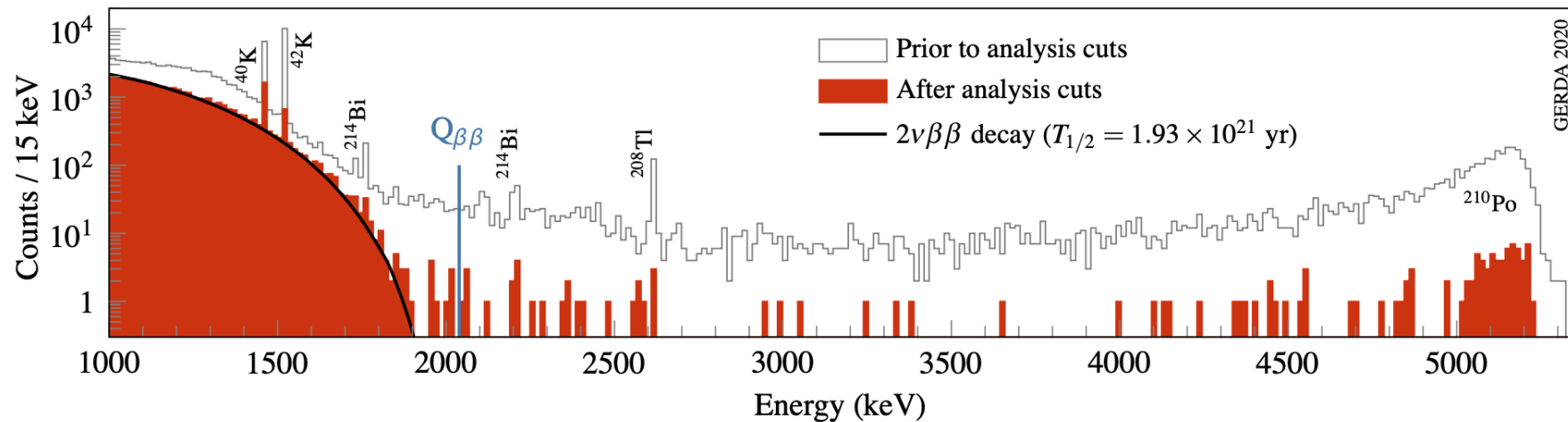
Fig.1: Illustration of GERDA. The Germanium diodes are protected by several layers of active and passive shieldings from external radiation.

# Gran Sasso

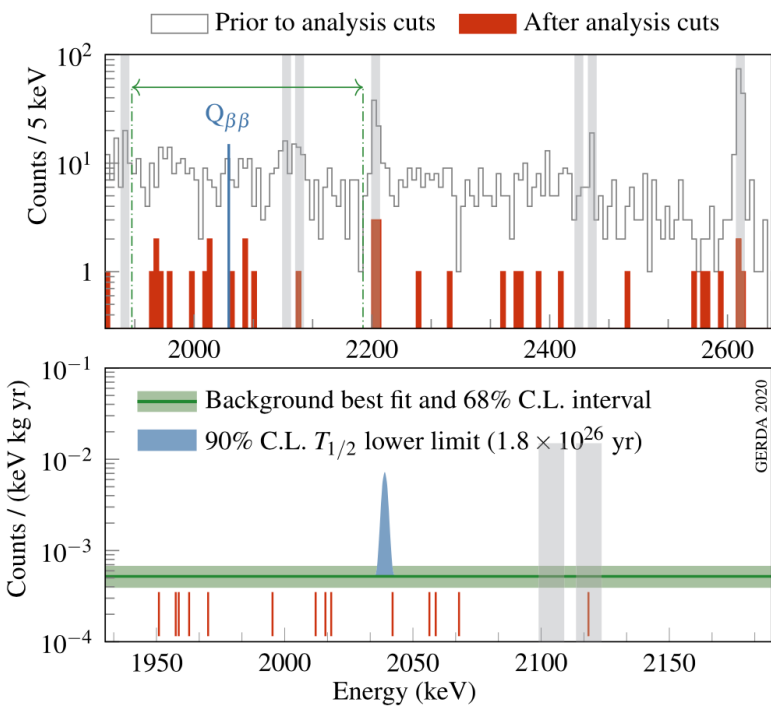




# GERDA experiment



Phys. Rev. Lett. 125, 252502 (2020)

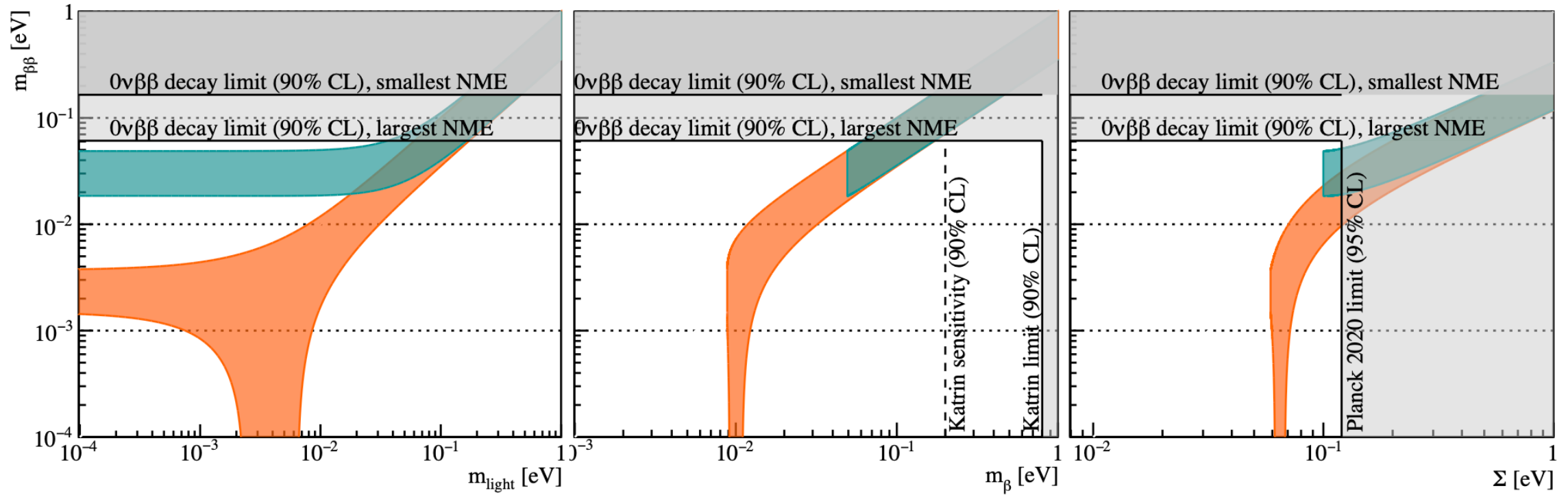


$$T_{1/2} > 1.8 \times 10^{26} \text{ yr at 90\% C.L.}$$



$$m_{\beta\beta} < 0.079 - 0.180 \text{ eV @ 90\% CL}$$

expected  $\beta\beta 0\nu$

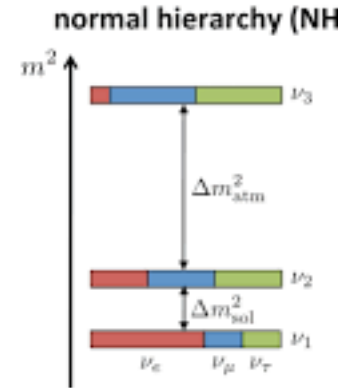
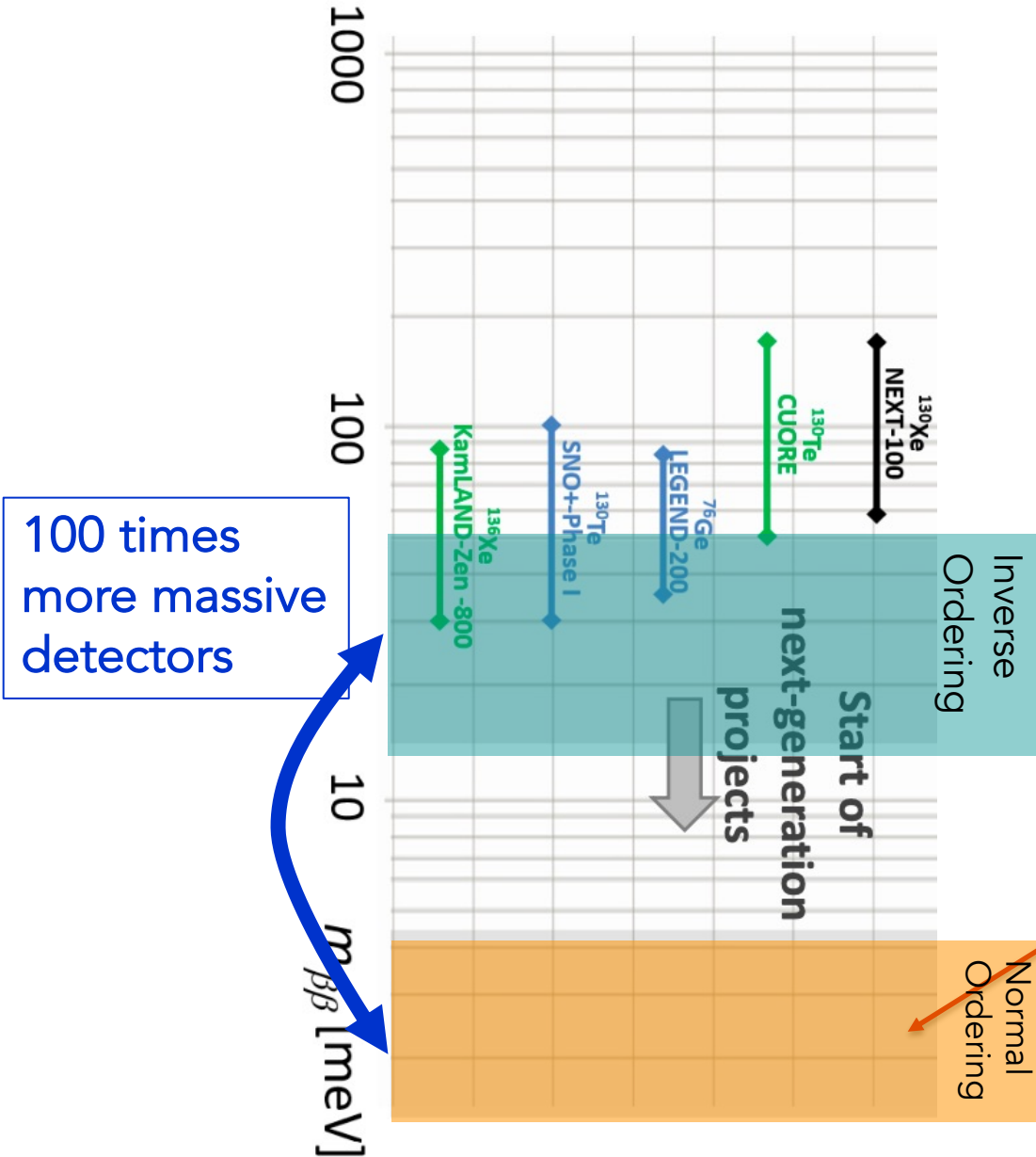


NO

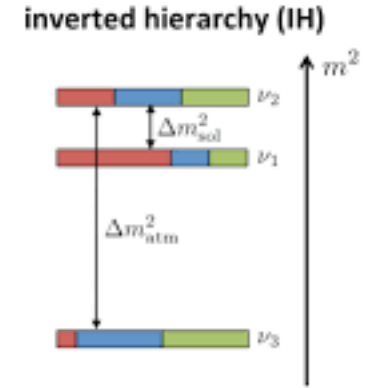
IO

Getting into the interesting region

A lot of experimental (and theoretical) activities



or



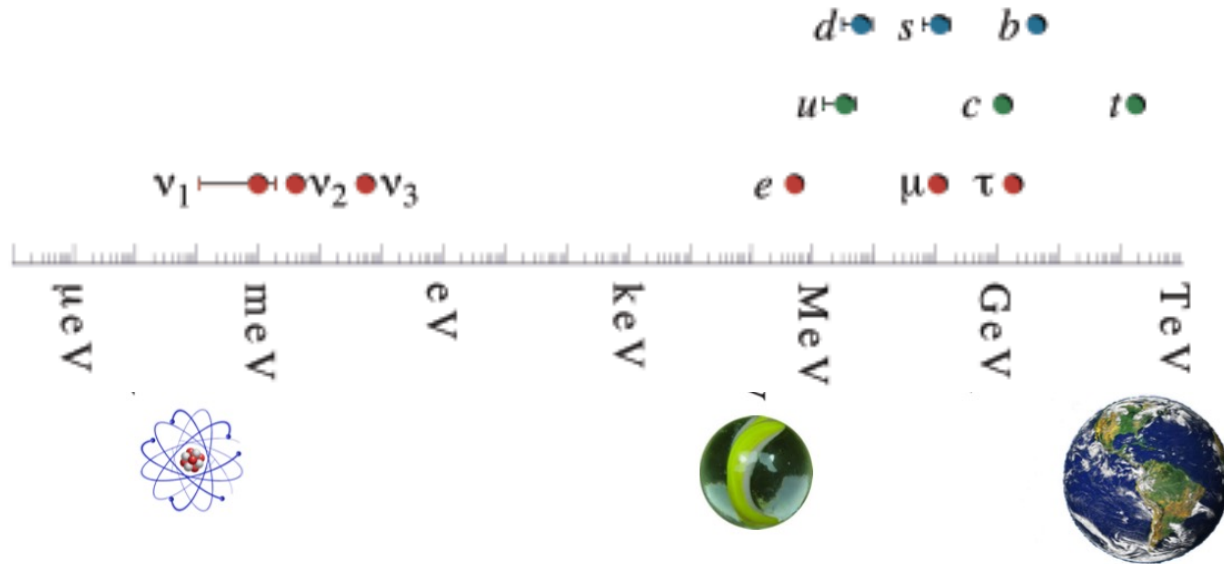
- Non observation of  $2\beta 0\nu$  would not prove that they are Dirac particles
- Observing  $2\beta 0\nu$  would prove that neutrinos are Majorana particles

Extremely challenging to reach

# Neutrino physics, a summary ?

At the interplay between collider physics and astroparticles & reactors

Origin of masses & problem of flavour



quarks

$$|V_{\text{CKM}}| \sim \begin{pmatrix} 1 & 0.2 & 0.004 \\ 0.2 & 1 & 0.04 \\ 0.008 & 0.04 & 1 \end{pmatrix}$$

neutrinos

$$|U_{\text{PMNS}}| \sim \begin{pmatrix} 0.8 & 0.5 & 0.1 \\ 0.5 & 0.6 & 0.7 \\ 0.3 & 0.6 & 0.7 \end{pmatrix}$$

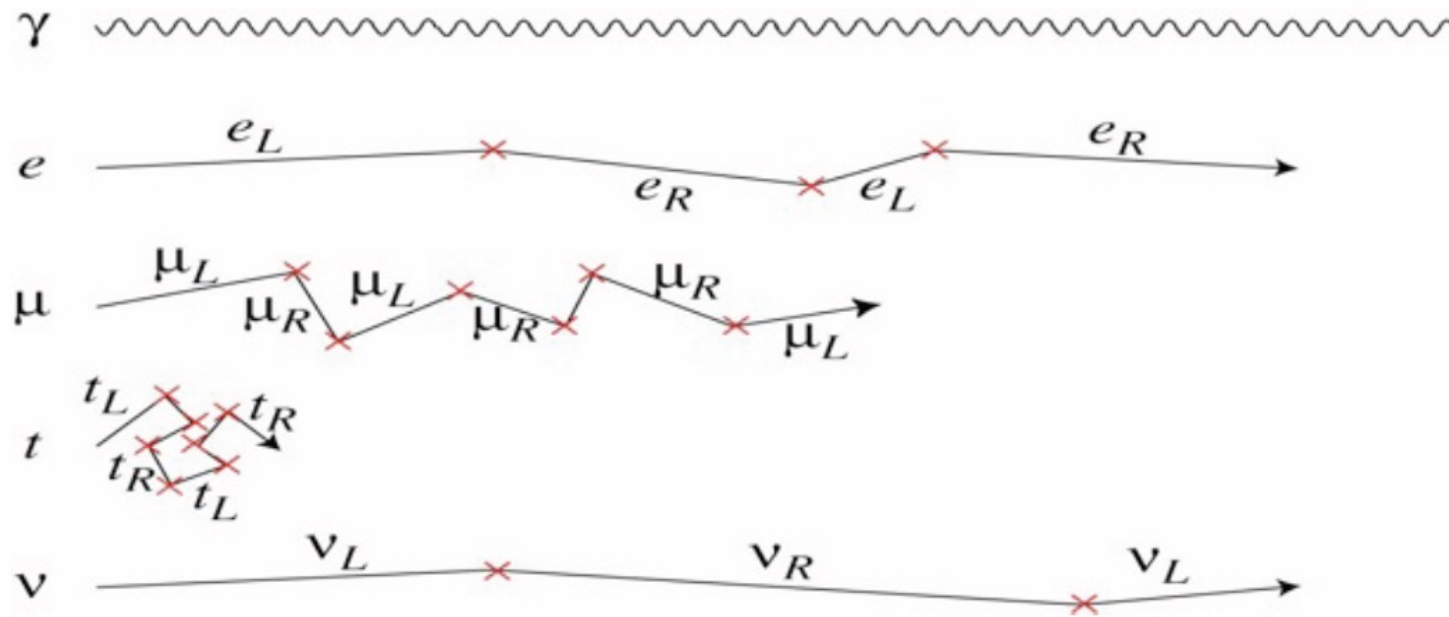
Dirac or Majorana ?

Explanation of the matter-antimatter asymmetry observed in the Universe ?

Back up slides

# Masses extremely small ...

If Dirac neutrinos



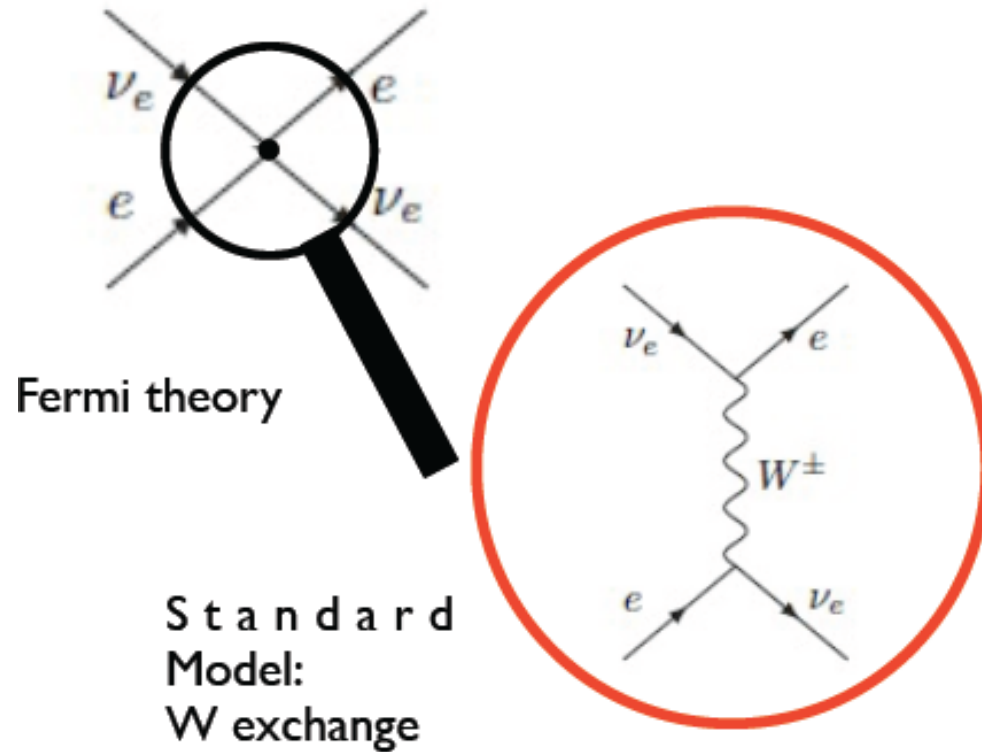
The couplings to the H are fixed :  $m_\nu = \lambda v$

If  $m \sim 0.1$  eV since  $v \sim 250$  GeV  $\rightarrow$  coupling  $\lambda = 4 \cdot 10^{-13}$  !

If neutrinos are of Majorana type

The generation of the mass can arise as the low energy realisation of a higher energy theory (new mass scale!)

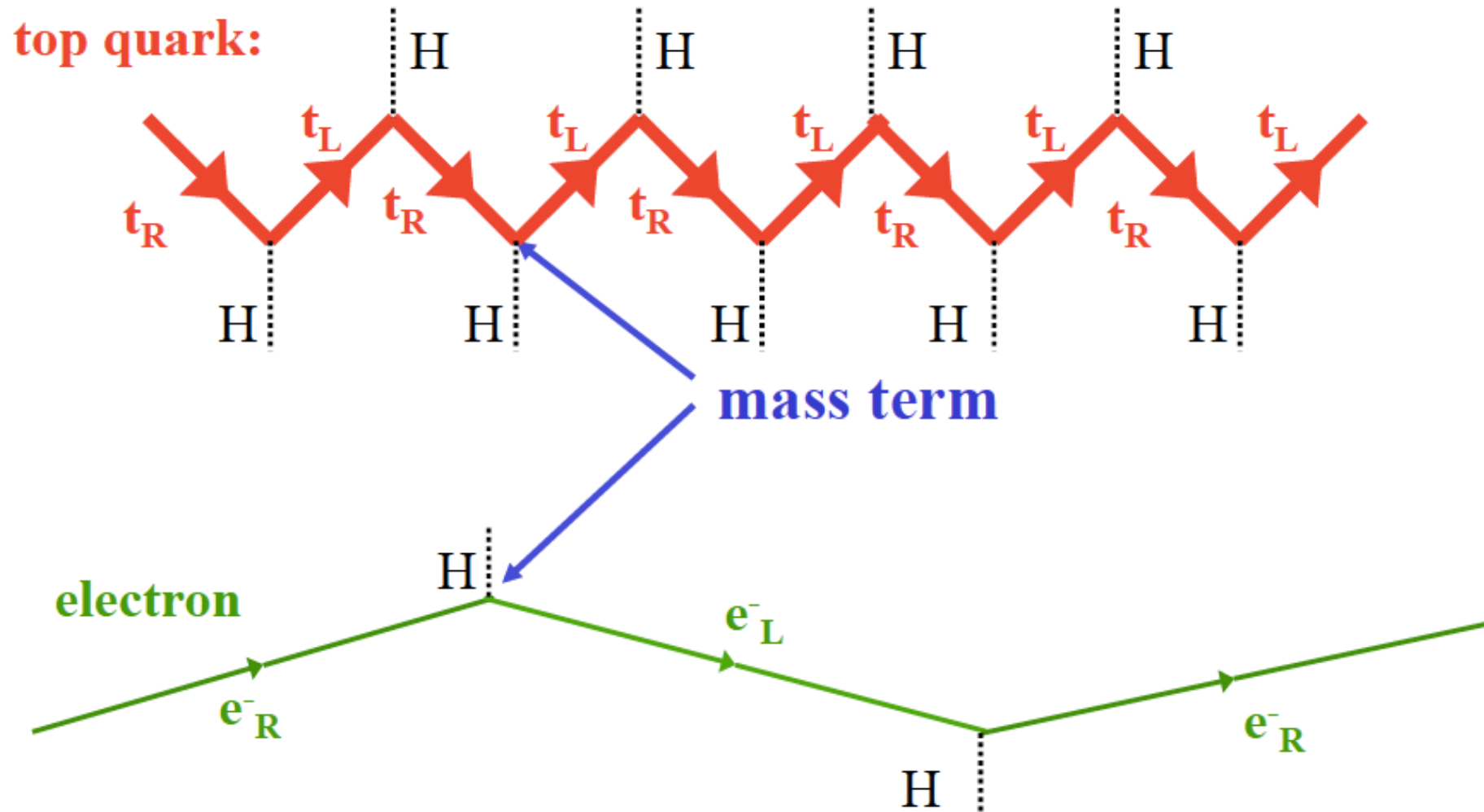
Remember Fermi's theory ?



$$m_\nu = \lambda \frac{v^2}{\Lambda}$$

$\lambda$  similar to the coupling of other fermions but suppressed by a larger scale (new particle exchange )

Remember the Higgs mechanism ?

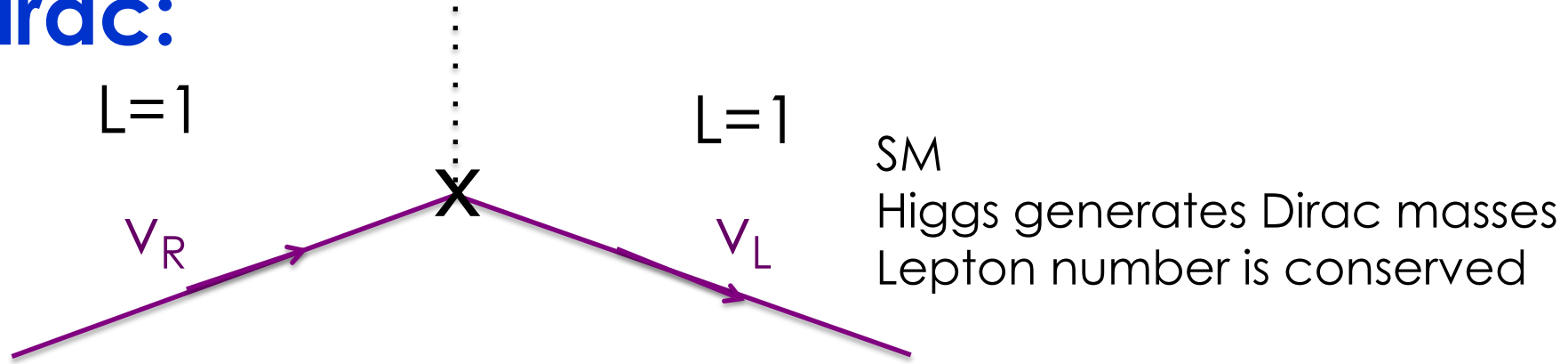


**Dirac mass term: conserves charge ( $e^- \rightarrow e^-$ )**

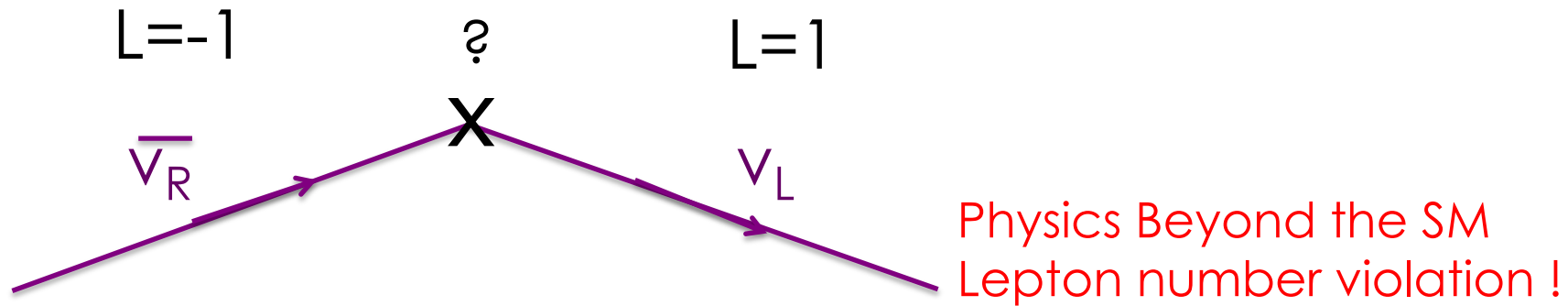


# What about the neutrinos masses ?

## Dirac:



## Majorana:



Are the neutrinos of Majorana or Dirac type ?  
Strong link with lepton number conservation

# Helicity/Chirality :

*Projection of the spin in the momentum direction*

- Helicity: definition for the **helicity** operator :  $H = \frac{\vec{\sigma} \cdot \vec{p}}{|\vec{p}|}$  with  $\sigma = \begin{pmatrix} \vec{\sigma} & 0 \\ 0 & \vec{\sigma} \end{pmatrix}$   $4 \times 4$

- Chirality:

If  $\psi$  is a solution for the Dirac equation, one can write:  $\psi = \psi_{CL} + \psi_{CR}$

Definition: **chirality** operators *Left* or *Right* (CL,CR) :

$$P_{CL} = \frac{1 - \gamma_5}{2} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

$$P_{CR} = \frac{1 + \gamma_5}{2} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$\gamma_5 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad 4 \times 4$$

(for the chiral representation)

Algebra :

$$P_{CL}^2 = P_{CL}, \quad P_{CR}^2 = P_{CR}, \quad P_{CL} + P_{CR} = 1, \quad P_{CL} P_{CR} = 0$$

$$\psi_{CL} = P_{CL} \psi \quad \bar{\psi}_{CL} = \bar{\psi} P_{CR}$$

$$\psi_{CR} = P_{CR} \psi \quad \bar{\psi}_{CR} = \bar{\psi} P_{CL}$$

$$P_{CL} \gamma^\mu = \gamma^\mu P_{CR} \quad P_{CR} \gamma^\mu = \gamma^\mu P_{CL}$$

# Link between helicity and chirality

- Chirality is the “correct” quantity (it appears in the Lagrangian and in addition the helicity is not Lorentz invariant) but what is measured in the processes is the helicity

- One can show that:

$$\psi_{CL} = \frac{a}{2}\psi_{HR} + \frac{\text{with}}{2}\psi_{HL}$$

$$\psi_{CR} = \frac{b}{2}\psi_{HR} + \frac{a}{2}\psi_{HL}$$

$$a = 1 - \frac{p}{E + m}$$

$$b = 1 + \frac{p}{E + m}$$

$\psi_{CR}, \psi_{CL}$  are the eigenvectors of H

$\psi_{CR}$  corresponds to the eigenvalue +1

$\psi_{CL}$  corresponds to the eigenvalue -1

- for  $m \ll E$ :  $a = 1 - \beta$  and  $b = 1 + \beta$   
 $\beta \sim 1$ :  $a = 0$  and  $b = 2$   
and thus :  $\psi_{CL} = \psi_{HL}$  and  $\psi_{CR} = \psi_{HR}$

for  $E \gg m$  : Helicity  $\equiv$  Chirality

## 4.2 V-A structure:

Let's take an electromagnetic current :  $\bar{\psi}\gamma^\mu\psi$  :

( $\gamma^\mu$  is a vector, parity = -1)

$$\begin{aligned}\bar{\psi}\gamma^\mu\psi &= \bar{\psi}(P_{CL} + P_{CR})\gamma^\mu(P_{CL} + P_{CR})\psi = \\ & \bar{\psi}\cancel{P_{CL}}\gamma^\mu\cancel{P_{CL}}\psi + \bar{\psi}P_{CR}\gamma^\mu P_{CL}\psi + \bar{\psi}P_{CL}\gamma^\mu P_{CR}\psi + \bar{\psi}\cancel{P_{CR}}\gamma^\mu\cancel{P_{CR}}\psi = \\ & \bar{\psi}_{CL}\gamma^\mu\psi_{CL} + \bar{\psi}_{CR}\gamma^\mu\psi_{CR} \\ \bar{\psi}\gamma^\mu\psi &= \bar{\psi}_{CL}\gamma^\mu\psi_{CL} + \bar{\psi}_{CR}\gamma^\mu\psi_{CR} \quad \text{selects } \psi_{CL}, \psi_{CR}\end{aligned}$$

$\Rightarrow$  for the electromagnetic interaction :  $\psi_{CL}$ -  $\psi_{CL}$  and  $\psi_{CR}$ -  $\psi_{CR}$  couplings

Let's take a weak coupling  $\bar{\psi}\gamma^\mu(1-\gamma_5)\psi$  :

$\left[ \begin{array}{l} \gamma^\mu\gamma^5 \text{ is vector-axial, parity } +1 \\ \gamma^\mu(1-\gamma_5): \text{ V-A structure} \end{array} \right]$

$$\begin{aligned}\bar{\psi}\gamma^\mu(1-\gamma_5)\psi &= \bar{\psi}(P_{CL} + P_{CR})\gamma^\mu(1-\gamma_5)(P_{CL} + P_{CR})\psi = \\ & 2\bar{\psi}(P_{CL} + P_{CR})\gamma^\mu(P_{CL}^2 + \cancel{P_{CL}P_{CR}})\psi = \\ & 2\bar{\psi}P_{CL}\gamma^\mu P_{CL}\psi + 2\bar{\psi}P_{CR}\gamma^\mu P_{CL}\psi = 2\bar{\psi}\gamma^\mu\cancel{P_{CR}}\cancel{P_{CL}}\psi + 2\bar{\psi}_{CL}\gamma^\mu\psi_{CL} \\ \bar{\psi}\gamma^\mu(1-\gamma_5)\psi &= 2\bar{\psi}_{CL}\gamma^\mu\psi_{CL} \quad \Rightarrow \text{for the weak interaction : } \psi_{CL}\text{- } \psi_{CL} \text{ coupling only}\end{aligned}$$

This form for the current leads to maximal parity violation  
(the V-A structure allows only left handed neutrinos)

Validation of the  $\gamma^\mu(1-\gamma_5)$  expression for the weak currents

# Leptogenesis (Majorana neutrinos)

Very heavy  
neutrino  $N$

$N$  can decay into  $\ell^+$  or  $\ell^-$

$N$  created at the big bang time

$$N \rightarrow \ell^+ \phi^-$$

$$N \rightarrow \ell^- \phi^+$$

Higgs field before SSB  $\begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$

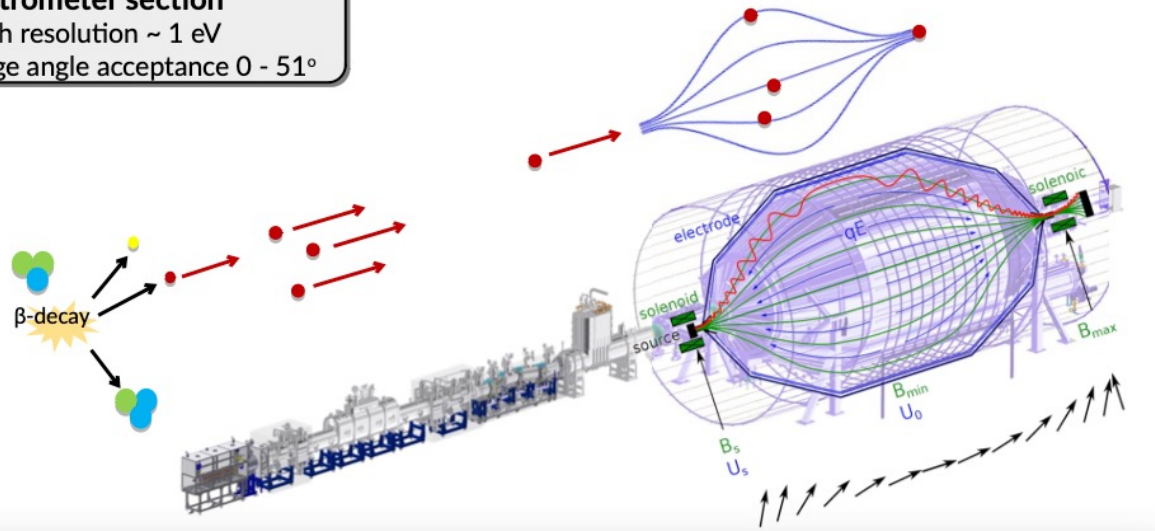
$$\text{CP violation} \Rightarrow \Gamma(N \rightarrow \ell^+ \phi^-) \neq \Gamma(N \rightarrow \ell^- \phi^+)$$

unequal amount of matter and antimatter in the leptonic sector

# KATRIN spectrometer working principle

## Spectrometer section

- high resolution  $\sim 1$  eV
- large angle acceptance  $0 - 51^\circ$



## Focal plane detector

- 148-pixel Si-PIN detector
- counting of electrons ( $< 1e^-/s$ )

