# Double Ratios in (Rare) Semi-Leptonic Decays and New Physics

#### Sascha Turczyk

in collaboration with

Zoltan Ligeti

#### **Preliminary Work**

Lawrence Berkeley National Laboratory

Colour meets Flavour Saturday, October 14th, 2011

#### Outline

- Introduction
  - Motivation
  - Notation and the Setup
  - Decay Rates
- Preliminary Results
  - Form Factors
  - Double Ratio Prediction
- Summary

### Rare Decays and New Physics

#### Rare Decays

- SM very successful
- Deviations mostly expected in flavour changing neutral currents



#### Influence of New Physics

- Integrate out heavy degress of freedom
- Encoded in Wilson coefficients of operator product expansion
- ⇒ New physics manifests in Wilson coefficients, only. Assumption: Complete basis in Standard Model (SM
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### The Goal

#### Current New Physics Searches: One example

- Investigate  $b \rightarrow s$  transition
  - $0 b \rightarrow s\gamma$
  - $b \rightarrow s\ell^+\ell^-$
- Constrain possible New Physics (NP) contribution by [arXiv:1106.1547]
  - Forward-Backward Asymmetry  $A_{FB}(q^2)$
  - 2 Angular analysis  $B \to K^*(K\pi)\ell^+\ell^-$  for CP violation

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### Wilson Coefficients and Theoretical Uncertainties

#### Wilson Coefficients

- $\mathcal{O}_1 \mathcal{O}_8$ :Contribute indirectly due to long-distance quark loops
- ⇒ Lots of investigations [JHEP 1009, 089 (2010), Eur. Phys. J. C 71, 1635 (2011)]
  - $\mathcal{O}_7$ : Wilson coefficient best constraint by  $b \to s \gamma$  at low  $q^2$
  - $\mathcal{O}_9 \mathcal{O}_{10}$ : Semi-leptonic operators
  - $\Rightarrow$   $C_9$  and  $C_{10}$  sensitive to new physics in  $b \to s \ell^+ \ell^-$  at high  $q^2$  different to  $\mathcal{O}_7$

#### Non-perturbative Input

Observing bound state mesons

$$\langle K^*(p_{K^*})|\mathcal{O}_i^{\mu_1...\mu_n}|B(p_B)\rangle = f(p_{K^*}^{\mu_j}, p_B^{\mu_j}, g^{\mu_j\mu_{j'}})$$

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- ⇒ Largest theoretical uncertainties

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### Double Ratios in a Simple Example

• Consider ratios of decay constants [ PRL 71, 3067, PRD 53, 4937 ]

#### Parametric deviations from unity

- 1 Chiral SU(3) light flavor symmetry
- 2 Heavy-quark spin-flavor symmetry SU(4)
- Perturbative corrections
- Compute the double ratio  $[f_{B_s}/f_B]/[f_{D_s}/f_D]$
- ⇒ Lots of theoretical uncertainties cancel

$$\frac{f_{B_s}/f_B}{f_{D_s}/f_D} \approx 1 + \mathcal{O}\left(\frac{m_s}{m_c} - \frac{m_s}{m_b}, \frac{m_s}{\Lambda} \frac{\alpha_s(m_c) - \alpha_s(m_b)}{\pi}\right) \approx 1 + \mathcal{O}(\lesssim 7\%)$$

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### The Double Ratio Proposal

#### Suggestion

• Use the ratio of  $(0 \le q^2 \le (M_{\text{initial}} - M_{\text{final}})^2)$ 

$$\begin{split} \frac{\int_{q_0^2}^{q_{\text{max}}^2} \frac{d\Gamma}{dq^2} (B \to \mathcal{K}^* \boldsymbol{\ell}^+ \boldsymbol{\ell}^-)}{\int_{q_0^2}^{q_{\text{max}}^2} \frac{d\Gamma}{dq^2} (D \to \rho \ell \bar{\nu}_\ell)} \frac{\int_{q_0^2}^{q_{\text{max}}^2} \frac{d\Gamma}{dq^2} (D \to \rho \ell \bar{\nu}_\ell)}{\int_{q_0^2}^{q_{\text{max}}^2} \frac{d\Gamma}{dq^2} (D \to \mathcal{K}^* \ell \bar{\nu}_\ell)} \\ &= \#(C_i) + \mathcal{O}(\frac{m_s}{m_c} - \frac{m_s}{m_b}, \frac{m_s}{\Lambda_{\text{QCD}}} \frac{\alpha_s(m_c) - \alpha_s(m_b)}{\pi}) \end{split}$$

- All decays are pseudoscalar to vector transitions
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### Computational Methods for Form Factors

#### **Brief Comment**

• Parameterize  $q^2$  dependence, usually pole form

$$F(q^2) = \frac{r_1}{1 - q^2/m_R^2} + \frac{r_2}{1 - q^2/m_{\text{fit}}^2},$$

- 1 Vector  $V(q^2)$
- 3 Axialvector  $A_{0-2}(q^2)$ , Most important:  $A_1(q^2)$  at high  $q^2$
- 3 Tensor  $T_{1-3}(q^2)$  form factors (only rare decay)
- Compute specific values with various methods, e.g. LCSR, lattice
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### Symmetries and Relations

#### Theoretical Properties

• Heavy quark symmetry predicts scaling behaviour and relations between form factors of a specific decay Form factors differently defined in HQE application!

$$A_1 \propto \mathcal{O}(m_b^{-1/2} + m_b^{-3/2})$$

- 4 Heavy quark symmetry: Corrections between B and D decays
- **3** Chiral symmetry: Corrections between  $B/B_s$  and  $D/D_s$  decays

#### First Investigations

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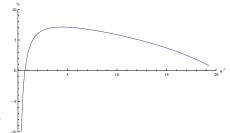
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$$\frac{h - \frac{a_+ - a_-}{2m_b} - \frac{g}{m_b}}{\frac{g}{m_b}}$$

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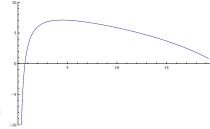
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- Expected from HQE 10 15%

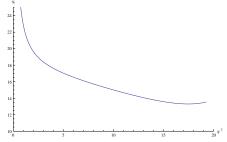


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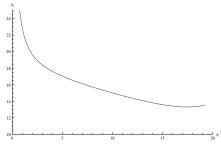


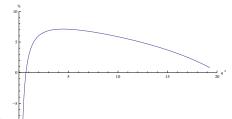
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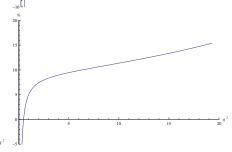
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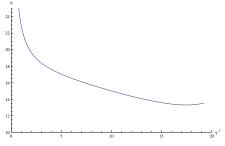


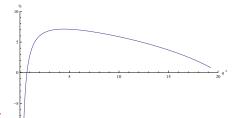


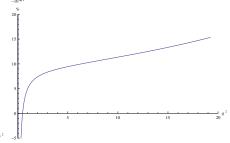
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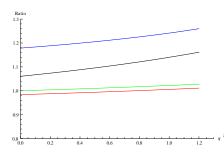




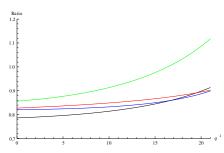


### Ratio of Meson Decay Form Factors

$$\frac{FF^{D\to\rho}(q^2)}{FF^{D\to K^*}(q^2)}$$



$$rac{FF^{B o
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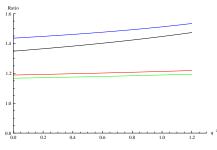


- black:  $V(q^2)$
- blue:  $A_0(q^2)$

- red:  $A_1(q^2)$
- green:  $A_2(q^2)$

### Double Ratio of Form Factors

$$\frac{FF^{B\to K^*}(q^2)}{FF^{B\to \rho}(q^2)}\cdot \frac{FF^{D\to \rho}(q^2)}{FF^{D\to K^*}(q^2)}$$



- black:  $V(q^2)$
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#### Remarks

- Ratio for D decay range of  $q^2$ !
- Improves for choosing different  $q^2$  integration ranges
- Most important  $A_1(q^2) \Rightarrow \mathcal{O}(10-20\%)$  possible

### Numerical Result for Double Ratio

- Use natural HQE variable  $w = \frac{1}{2M_i M_j} \left( M_i^2 + M_j^2 q^2 \right)$
- B decays:  $q^2 \gtrsim 15 \,\text{GeV}^2 \Rightarrow 1 \le w \le w_B = 1.45$
- *D* decays: Full  $D \to K^*$  phase-space  $\Rightarrow 1 \le w \le w_B = 1.29$

$$\begin{split} \mathsf{DR}(w_D,w_B) &= \frac{3\alpha^2}{64\pi^2} \frac{|V_{cd}|^2 |V_{tb}|^2 |V_{ts}|^2}{|V_{cs}|^2 |V_{ub}|^2} \bigg[ 86.55 |C_1|^2 + 9.62 |C_2|^2 + 53.2 |C_3|^2 + 61.96 |C_4|^2 \\ &+ 66.14 |C_5|^2 + 7.35 |C_6|^2 + 14.55 |C_7|^2 + 4.51 |C_9|^2 + 4.51 |C_{10}|^2 \\ &+ 57.7 \, \mathsf{Re}\left(C_1 \, C_2^*\right) + 77.48 \, \mathsf{Re}\left(C_1 \, C_3^*\right) - 100.6 \, \mathsf{Re}\left(C_1 \, C_4^*\right) \\ &+ 124.9 \, \mathsf{Re}\left(C_1 \, C_5^*\right) + 41.64 \, \mathsf{Re}\left(C_1 \, C_6^*\right) + 38.94 \, \mathsf{Re}\left(C_1 \, C_7^*\right) \\ &+ 16.18 \, \mathsf{Re}\left(C_1 \, C_9^*\right) + 25.82 \, \mathsf{Re}\left(C_2 \, C_3^*\right) - 33.54 \, \mathsf{Re}\left(C_2 \, C_4^*\right) \\ &+ \dots \\ &+ 21.6 \, \mathsf{Re}\left(C_7 \, C_9^*\right) - 111.34 \, \mathsf{Im}\left(C_1 \, C_3^*\right) - 106.38 \, \mathsf{Im}\left(C_1 \, C_4^*\right) \\ &- 85.36 \, \mathsf{Im}\left(C_1 \, C_5^*\right) - 28.46 \, \mathsf{Im}\left(C_1 \, C_6^*\right) - 86.3 \, \mathsf{Im}\left(C_1 \, C_7^*\right) \\ &- 36. \, \mathsf{Im}\left(C_1 \, C_9^*\right) - 37.12 \, \mathsf{Im}\left(C_2 \, C_3^*\right) - 35.46 \, \mathsf{Im}\left(C_2 \, C_4^*\right) \\ &+ \dots \end{split}$$

#### Summary

- Showed potential of double ratios
- Application to rare semi-leptonic decay at high  $q^2$
- Showed first numerical estimate of uncertainties
- Presented numerical result of double ratio as function of Wilson coefficients
- Analysis with uncertainties  $\mathcal{O}(10-20\%)$  possible

#### Outlook

- Wait for new experimental data on D decay form factors
- Perform more elaborate numerical analysis
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## **Backup Slides**

### A Little Bit Notation: Operator Basis

- Consider rare decay:  $B \to K^* \ell^+ \ell^-$
- $\mathcal{O}_1 \mathcal{O}_2$ : Current-current operators
- $\mathcal{O}_3 \mathcal{O}_6$ : QCD-penguin operators
- $\mathcal{O}_7 \mathcal{O}_8$ : Magnetic penguin operators
- $\mathcal{O}_9 \mathcal{O}_{10}$ : Semi-leptonic operators

$$\mathcal{O}_{1}^{\rho} = (\bar{p}b)_{V-A}(\bar{s}p)_{V-A} \qquad \qquad \mathcal{O}_{2}^{\rho} = (\bar{p}_{i}b_{j})_{V-A}(\bar{s}_{j}p_{i})_{V-A}, 
\mathcal{O}_{3} = (\bar{s}b)_{V-A} \sum_{q} (\bar{q}q)_{V-A} \qquad \qquad \mathcal{O}_{4} = (\bar{s}_{i}b_{j})_{V-A} \sum_{q} (\bar{q}_{j}q_{i})_{V-A}, 
\mathcal{O}_{5} = (\bar{s}b)_{V-A} \sum_{q} (\bar{q}q)_{V+A} \qquad \qquad \mathcal{O}_{6} = (\bar{s}_{i}b_{j})_{V-A} \sum_{q} (\bar{q}_{j}q_{i})_{V+A}, 
\mathcal{O}_{7} = \frac{e}{8\pi^{2}} m_{b} \bar{s}\sigma_{\mu\nu} (1+\gamma_{5}) F^{\mu\nu} b \qquad \qquad \mathcal{O}_{8} = \frac{g}{8\pi^{2}} m_{b} \bar{s}\sigma_{\mu\nu} (1+\gamma_{5}) G^{\mu\nu} b, 
\mathcal{O}_{9} = \frac{\alpha}{2\pi} (\bar{s}b)_{V-A} (\bar{l}l)_{V} \qquad \qquad \mathcal{O}_{10} = \frac{\alpha}{2\pi} (\bar{s}b)_{V-A} (\bar{l}l)_{A}.$$

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- Problems: Resonances, Duality violation, . . .
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#### $C_9$ and $C_{10}$

- $\mathcal{O}_7$ : Wilson coefficient best constraint by  $b \to s\gamma$
- $\mathcal{O}_9 \mathcal{O}_{10}$ : Semi-leptonic operators
- $\Rightarrow$   $C_9$  and  $C_{10}$  most sensitive to new physics in  $b \to s \ell^+ \ell^-$

### $B \to K^* \ell \ell$ Decay

$$\begin{split} &\frac{\mathrm{d}\Gamma(\bar{B}\to\bar{K}^*\ell^+\ell^-)}{\mathrm{d}q^2} = \frac{G_F^2\alpha_{\mathrm{em}}^2M_B^3}{1024\pi^5}\left|V_{ts}^*V_{tb}\right|^2\lambda_{K^*}^{1/2}(q^2)\\ &\times \left\{R_9\left(|\tilde{C}_9^{\mathrm{eff}}(q^2)|^2 + |C_{10}|^2\right) + R_7\frac{m_b^2}{m_B^2}|C_7|^2 + R_{97}\frac{m_b}{m_B}\operatorname{Re}\big[\tilde{C}_9^{\mathrm{eff}}(q^2)C_7^*\big]\right\} \end{split}$$

#### **Properties**

- $C_9^{\text{eff}}(q^2)$  contains  $C_{1-8}$  loops
- $\bullet$   $R_9$ ,  $R_7$  and  $R_{97}$  contain the form factors
  - 1 Vector  $V(q^2)$ , (also tree-level)
  - 3 Axialvector  $A_{0-2}(q^2)$ , (also tree-level)
  - 3 Tensor  $T_{1-3}(q^2)$  Form Factors
- $A_1(g^2)$  dominates for high  $g^2$

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#### **Properties**

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- $R_9$ ,  $R_7$  and  $R_{97}$  contain the form factors
  - 1 Vector  $V(q^2)$ , (also tree-level)
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### The Tree Level Decays

$$\begin{split} \frac{\mathrm{d}\Gamma}{\mathrm{d}q^2} &= \frac{G_F^2 |V_{ji}|^2}{192\pi^3 M_i^3} \sqrt{(M_i^2 + M_j^2 - q^2)^2 - 4M_i^2 M_j^2} \left( (M_i + M_j)^2 - q^2 \right) \\ & \left( q^2 \left( (M_i + M_j)^2 + M_i^2 + M_j^2 - 4q^2 \right) + \left( M_i^2 - M_j^2 \right)^2 \right) [F_{ij}(q^2)]^2 \end{split}$$

with

$$\begin{split} [F_{ij}(q^2)]^2 = & \frac{M_i((M_i - M_j)^2 - q^2)}{M_j(M_i + M_j)^2 \left(q^2 \left(3M_i^2 + 2M_iM_j + 3M_j^2\right) + \left(M_i^2 - M_j^2\right)^2 - 4q^{2^2}\right)} \\ & \left[ 8M_j^2 q^2 \left([V^{ij}(q^2)]^2 + \frac{(M_i + M_j)^4}{M_i^4 - 2M_i^2 \left(M_j^2 + q^2\right) + \left(M_j^2 - q^2\right)^2} [A_1^{ij}(q^2)]^2\right) \\ & + \left(M_i^4 - 2M_i^2 \left(M_j^2 + q^2\right) + \left(M_j^2 - q^2\right)^2\right) \left(A_1^{ij}(q^2) \\ & - \frac{(M_i + M_j)^2 \left(M_i^2 - M_j^2 - q^2\right)}{M_i^4 - 2M_i^2 \left(M_i^2 + q^2\right) + \left(M_i^2 - q^2\right)^2} A_2^{ij}(q^2)\right)^2 \right]. \end{split}$$

#### Form Factor Definitions

$$\begin{split} \langle V(p_V)|\bar{q}\gamma_{\mu}b|B(p_B)\rangle &= \epsilon_{\mu\nu\rho\sigma}\epsilon^{*,\nu}p_B^{\rho}p_V^{\sigma}\frac{2V(q^2)}{m_B+m_V}\\ \langle V(p_V)|\bar{q}\gamma_{\mu}\gamma_5b|B(p_B)\rangle &= +i\epsilon_{\mu}^*(m_B+m_V)A_1(q^2) - i(p_B+p_V)_{\mu}(\epsilon^*\cdot q)\frac{A_2(q^2)}{m_B+m_V}\\ &- iq_{\mu}(\epsilon^*q)\frac{2m_V}{q^2}\left(A_3(q^2)-A_0(q^2)\right)\;, \end{split}$$

where the form factor combination  $A_3(q^2)$  is defined by

$$A_3(q^2) = \frac{m_B + m_V}{2m_V} A_1(q^2) - \frac{m_B - m_V}{2m_V} A_2(q^2)$$

$$\begin{split} \langle V(p_V)|\bar{q}\sigma_{\mu\nu}q^{\nu}b|B(p_B)\rangle &= i\epsilon_{\mu\nu\rho\sigma}\epsilon^{*\nu}p_B^{\rho}p^{\sigma}\,2T_1(q^2)\\ \langle V(p_V)|\bar{q}\sigma_{\mu\nu}q^{\nu}\gamma_5b|B(p_B)\rangle &= T_2(q^2)\left[\epsilon_{\mu}^*(m_B^2-m_V^2)-(\epsilon^*\cdot q)\,(p_B+p)_{\mu}\right]\\ &+ T_3(q^2)(\epsilon^*\cdot q)\left[q_{\mu}-\frac{q^2}{m_B^2-m_V^2}\,(p_B+p)_{\mu}\right]\,, \end{split}$$

#### Form Factor Relations

$$i f = -A_1(M_B + M_V)$$

$$i g = \frac{V}{M_B + M_V}$$

$$i h = \frac{T_3}{M_B^2 - M_V^2} - \frac{T_1 - T_2}{q^2}$$

$$i g_+ = -T_1$$

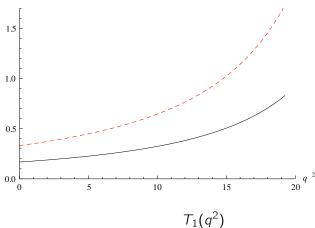
$$i g_- = \frac{(M_B^2 - M_V^2)(T_1 - T_2)}{q^2}$$

$$i a_+ = \frac{A_2}{M_B + M_V}$$

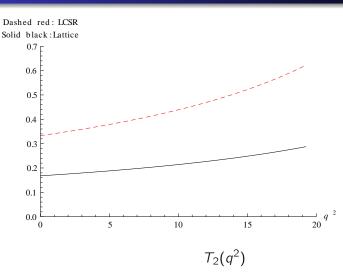
$$i a_- = \frac{A_1(M_B + M_V) - 2A_0M_V - A_2(M_B - M_V)}{q^2}$$

### $B \to K^*$ Tensor Form Factor Comparisson

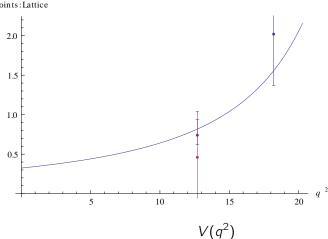
Dashed red: LCSR Solid black: Lattice

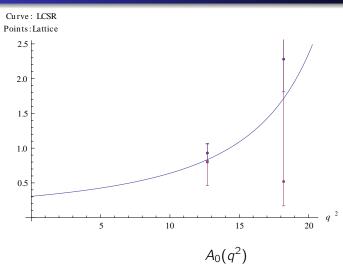


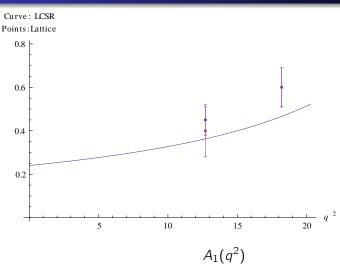
### $B \to K^*$ Tensor Form Factor Comparisson

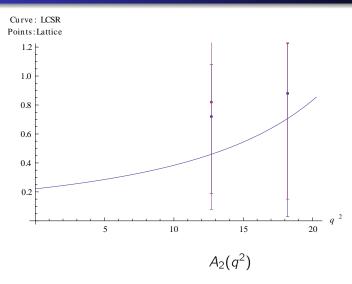


Curve: LCSR Points: Lattice









### Numerical Values

	$B  o K^* \ell^+ \ell^-$	$B o ho \ellar u_\ell$	$D o ho \ellar u_\ell$	$D o K^*\ellar u_\ell$
$q_{\rm max}^2$ / GeV <sup>2</sup>	19.25	20.28	1.20	0.96
W <sub>max</sub>	3.04	3.47	1.41	1.29

	$M_B$	$M_D$	$M_{ ho}$	$M_K^*$	$m_b$	$m_c$
Value / GeV <sup>2</sup>	5.279	1.87	0.776	0.892	4.2	1.17