

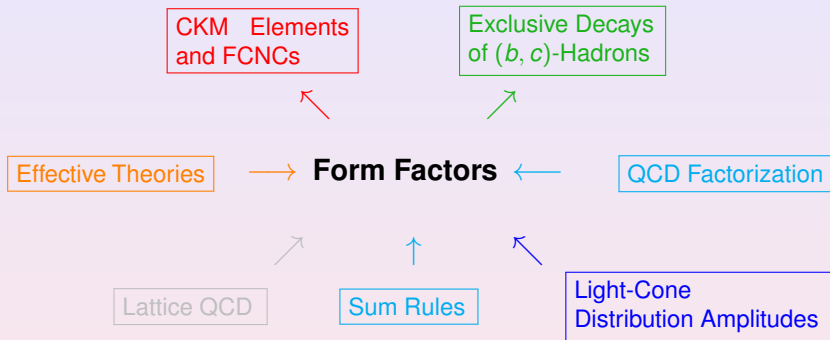
Heavy-to-Light Form Factors and Soft-Collinear Effective Theory

Thorsten Feldmann



“Colour Meets Flavour“,
Siegen, 13./14. October 2011

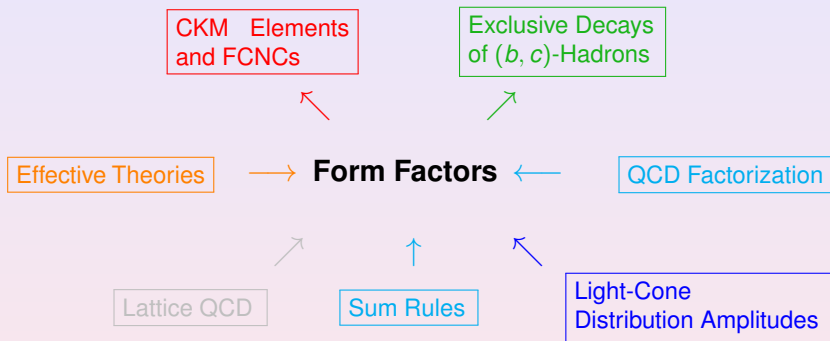
Flavour Physics



Quantum Chromodynamics

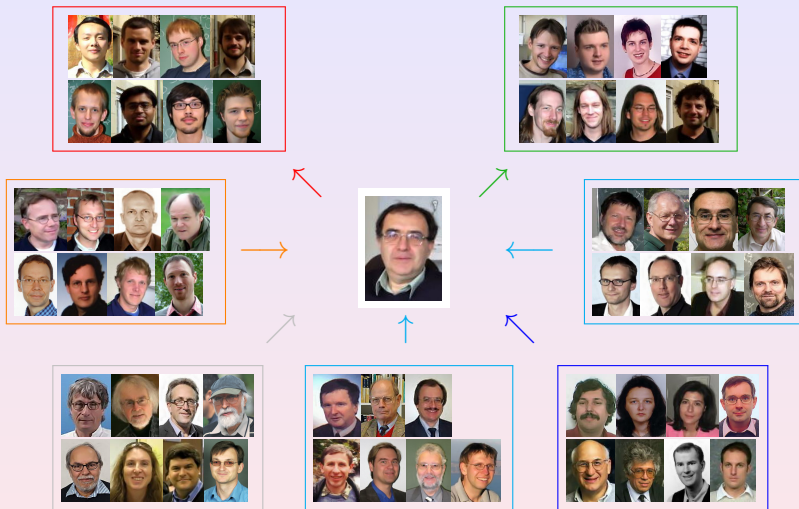
→ A New Duality Transformation →

Flavour Physics



Quantum Chromodynamics

Friends



Physics

...but seriously ...

- 1 Transition Form Factors for Heavy Mesons/Baryons
- 2 Calculation of "Soft" Form Factors
- 3 Radiative Corrections
- 4 Conclusions

Transition Form Factors – Definitions and Conventions

- **Lorentz Decomposition ?** ← polarization w.r.t. hadronic momentum transfer
- **Normalization ?** ↔ "point-like" hadrons for $f_i \equiv 1$
 - simple expressions for partial rates
 - simple expressions for unitarity bounds
 - simple symmetry relations in the HQET and SCET limit

Example: Vector Form Factors for $\Lambda_b \rightarrow \Lambda$

$$\begin{aligned} & \langle \Lambda(p', s') | \bar{q} \gamma_\mu b | \Lambda_b(p, s) \rangle \\ &= \bar{u}_\Lambda(p', s') \left\{ f_0(q^2) (M_{\Lambda_b} - m_\Lambda) \frac{q_\mu}{q^2} \right. \\ & \quad + f_+(q^2) \frac{M_{\Lambda_b} + m_\Lambda}{s_+} \left(p_\mu + p'_\mu - \frac{q_\mu}{q^2} (M_{\Lambda_b}^2 - m_\Lambda^2) \right) \\ & \quad \left. + f_\perp(q^2) \left(\gamma_\mu - \frac{2m_\Lambda}{s_+} p_\mu - \frac{2M_{\Lambda_b}}{s_+} p'_\mu \right) \right\} u_{\Lambda_b}(p, s) \end{aligned}$$

$$\text{with } s_+ = (M_{\Lambda_b} + m_\Lambda)^2 - q^2$$

- f_0, f_+, f_\perp correspond to time-like, longitudinal, transverse polarization
- For $f_0, f_+, f_\perp \rightarrow 1$, one obtains

$$\begin{aligned} & \langle \Lambda(p', s') | \bar{q} \gamma_\mu b | \Lambda_b(p, s) \rangle \\ &= \bar{u}_\Lambda(p', s') \gamma_\mu u_{\Lambda_b}(p, s) \end{aligned}$$

[TF, M. Yip; in preparation]

(similar for 3 axial-vector and 4 tensor form factors)

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Kinematic Regimes and Symmetry Relations

HQET Regime:

- Heavy quark approximated by static source of colour with velocity

$$v^\mu = p^\mu / M$$

- Energy transfer to light hadron is small, $v \cdot p' \ll M$
- ⇒ HQET spin/flavour symmetries reduce number of independent form factors

For example, in case of $\Lambda_b \rightarrow \Lambda$

Reduction from 10 \rightarrow 2 independent form factors,

$$\langle \Lambda(p', s') | \bar{q} \Gamma b | \Lambda_b(p, s) \rangle \simeq \bar{u}_\Lambda(p', s') \left\{ A(v \cdot p') + \not{v} B(v \cdot p') \right\} \Gamma u_{\Lambda_b}(p, s)$$

which implies

$$f_0(q^2) \simeq A(v \cdot p') + B(v \cdot p'), \quad f_+(q^2) \simeq f_\perp(q^2) \simeq A(v \cdot p') - B(v \cdot p') \quad \text{etc.}$$

[cf. Mannel/Recksiegel]

Kinematic Regimes and Symmetry Relations

SCET Regime:

- Energy of light hadron is large,

(\rightarrow 2 light-like vectors $n_{\pm}^2 = 0$)

$$p'^{\mu} = (n_+ p') \frac{n_-^{\mu}}{2} + (n_- p') \frac{n_+^{\mu}}{2} + p'_{\perp}{}^{\mu} \quad \text{with} \quad (n_+ p') \gg |p'_{\perp}| \gg (n_- p')$$

- Heavy quark still approximated by static source of colour
 - Light quark fields can be approximated by SCET spinors
- \Rightarrow Further reduction of number of independent form factors

(“Feynman mechanism“)

[Charles et al; Beneke/TF]

For example, in case of $\Lambda_b \rightarrow \Lambda$

Reduction from 10 \rightarrow 1 independent form factor,

$$\langle \Lambda(p', s') | \bar{q} \Gamma b | \Lambda_b(p, s) \rangle \simeq \xi_{\Lambda}(n_+ p') \bar{u}_{\Lambda}(p', s') \Gamma u_{\Lambda_b}(p, s)$$

which implies

$$f_0(q^2) \simeq f_+(q^2) \simeq f_{\perp}(q^2) \simeq \xi_{\Lambda}(n_+ p') \quad \text{etc.}$$

[TF/Yip]

Calculation of "soft" form factors $\xi_{\pi,K}, \xi_{\rho,K^*}, \xi_{\Lambda} \dots$

- Lattice QCD:

- ▶ in principle: **straight-forward**
- ▶ in practice: difficult/costly to **simulate fast light hadrons** on a lattice
- ▶ typically: reliable predictions for **intermediate momentum transfer**

- QCD Factorization (BBNS):

- ▶ Soft form factors as **Non-Factorizable** (i.e. irreducible, non-perturbative) ingredients in **Factorization Theorems** [Beneke/TF]

- (conventional) Light-cone Sum Rules:

- ▶ Study **Correlation Functions** with interpolating current for *heavy* hadron
- ▶ Non-perturbative input from **Universal LCDAs** for light hadron
- ▶ Dispersion relations \oplus **Continuum Model** [Alex + many others]

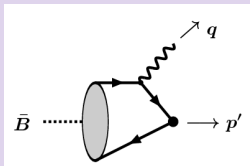
- **SCET Sum Rules:**

- ▶ Exchange role of light and heavy hadrons [Alex, Mannel, Offen]
- ▶ **Factorization of hard scales** ($\mu^2 \sim M_b^2$) from decay current matching.
- ▶ **Factorization of hard-collinear scales** ($\mu^2 \sim \Lambda M_b$) and **soft scales** ($\mu^2 \sim \Lambda^2$) in SCET correlation functions. [De Fazio, TF, Hurth]

SCET Sum Rules for Mesonic and Baryonic Transitions (Correlators)

$B \rightarrow \pi, \rho \dots$ Form Factors

[De Fazio/TF/Hurth]



- For appropriate currents, correlator takes the form

$$\Pi(n_- p') = f_B m_B \int_0^\infty d\omega \frac{\phi_B^{(-)}(\omega, \mu)}{\omega - n_- p' - i\epsilon} + \mathcal{O}(\alpha_s)$$

- No large logarithms for: $\mu^2/n_+ p' \sim |n_- p'| \sim \langle \omega \rangle$.

$\Lambda_b \rightarrow \Lambda \dots$ Form Factors

[TF/Yip]

- For appropriate currents, correlator takes the form

$$\Pi(n_- p') = f_{\Lambda_b}^{(2)} \frac{\not{n}_-}{2} u_{\Lambda_b} \int_0^\infty d\omega \frac{\phi_{\Lambda_b}^{(4)}(\omega, \mu)}{\omega - n_- p' - i\epsilon} + \mathcal{O}(\alpha_s)$$

- At LO, only *sum* of spectator momenta appears, with

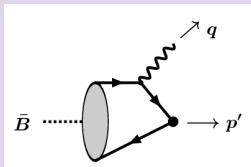
$$\phi_{\Lambda_b}^{(4)}(\omega) \equiv \omega \int_0^1 du \psi_{\Lambda_b}^{(4)}(\omega, u)$$

(LCDA of light scalar diquark component in Λ_b)

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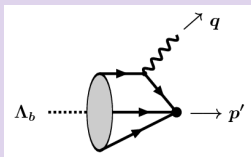
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SCET Sum Rules for Mesonic and Baryonic Transitions (Sum Rule)

After Borel transformation ($\rightarrow \omega_M$) and continuum subtraction ($\rightarrow \omega_S$):

$B \rightarrow X = \pi, \rho \dots$ Form Factors

[De Fazio/TF/Hurth]

$$e^{-m_X^2/(n_+p')\omega_M} (n_+p') f_X \xi_X(n_+p', \mu) = M_B f_B \int_0^{\omega_S} d\omega \phi_B^{(-)}(\omega, \mu) e^{-\omega/\omega_M} + \dots$$

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[TF/Yip]

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Theoretical Uncertainties / Hadronic Input:

- Values for light and heavy hadron decay constants.
- Shape of LCDAs for B and Λ_b .
- Reasonable choice of continuum threshold parameter $\omega_S = s_0/(n_+p')$.
- Reasonable range for Borel parameter $\omega_M = M_{\text{Borel}}^2/(n_+p')$.
- Logarithmically enhanced radiative corrections $\sim \alpha_S \ln \omega_{S,M}/\omega$.
(\leftrightarrow endpoint divergences in QCDF \leftrightarrow non-factorizable dependence on continuum model)

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Example: Numerical estimates for Baryonic Form Factor ξ_Λ (LO)

Input Parameters (default values) :

Threshold $s_0 = \omega_s (n+p')$	2.55 GeV ²
Borel $M_{\text{Borel}}^2 = \omega_M (n+p')$	2.5 GeV ²
Decay constant f_Λ	$6 \cdot 10^{-3}$ GeV ²
Decay constant $f_{\Lambda_b}^{(2)}$	0.030 GeV ³
LCDA par. ω_0	300 MeV

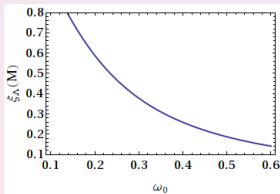
Model for the Λ_b LCDA:

$$\phi_{\Lambda_b}^{(4)}(\omega) := \frac{\omega}{\omega_0^2} e^{-\omega/\omega_0}$$

Central Value (LO, max. recoil)

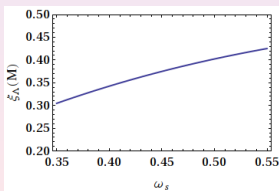
$$\xi_\Lambda(n+p' = M_{\Lambda_b}) \approx 0.38$$

ω_0 dependence



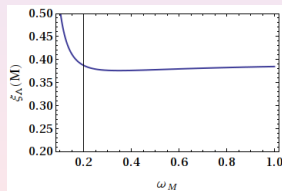
(very sensitive)

ω_s dependence



(moderate)

ω_M dependence

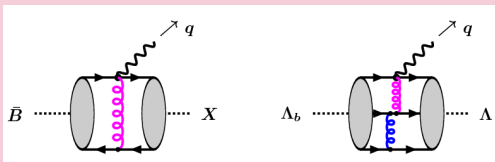


(weak)

[TF/Yip – preliminary]

Calculation of perturbative Corrections $\Delta\xi$

QCD Factorization: Convolution of LCDAs and spectator-scattering kernels

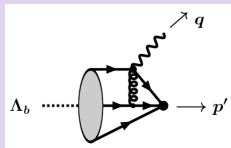
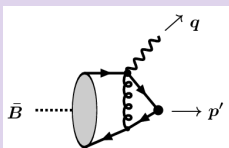


(improved version of Efremov-Radyushkin-Brodsky-Lepage picture)

- Factorization for $B \rightarrow X = \pi, \rho \dots$ in the large recoil limit [Beneke/TF]
- For $\Lambda_b \rightarrow \Lambda$ form factors, factorization spoiled by soft-gluon exchange between spectator quarks
- ...

SCET Sum Rules: Correlators involving "factorizable" part of SCET currents

$$\bar{q} \Gamma_i Q_v \longrightarrow C_{ij} \bar{\xi} \tilde{\Gamma}_j h_v \overbrace{-\frac{1}{n+p'} \bar{\xi} g A_\perp \frac{\not{h}_+}{2} \Gamma_i h_v - \frac{1}{M_b} \bar{\xi} \Gamma_i \frac{\not{h}_-}{2} g A_\perp h_v} + \dots$$



- Reproduces QCDF for $B \rightarrow X$ in the large recoil limit [De Fazio/TF/Hurth]
- Predictions for violation of symmetry relations through $\Delta\xi_\Lambda \sim \mathcal{O}(\alpha_s)$

[TF/Yip]

Features of $\Delta\xi_\Lambda$

In the limit $\omega_{s,M} \ll \omega_0$, the sum rule takes the form

$$e^{-m_\Lambda^2/(n_+p')\omega_M} f_\Lambda m_\Lambda \Delta\xi_\Lambda$$

$$\simeq \frac{\alpha_s C_F}{2\pi} \frac{f_{\Lambda_b}^{(2)}}{M_{\Lambda_b}} \underbrace{\int_0^\infty \frac{d\omega}{\omega} F(\omega, 1)} \times \underbrace{\omega_M \left(\omega_M - e^{-\omega_s/\omega_M} (\omega_M + \omega_s) \right)}$$

- Formally, $\Delta\xi_\Lambda$ and ξ_Λ scale with the same power of M_b .
- Dependence on Λ_b -Properties and Sum-Rule Parameters factorizes !
- Proportional to inverse moment of the effective LCDA

$$F(\omega, 1) \equiv \int_0^\omega d\eta \frac{\eta}{\omega} \left(\tilde{\psi}_4(\eta, u=1) - \tilde{\psi}_2(\eta, u=1) \right)$$

- Dependence on decay constants cancels in ratio $\Delta\xi/\xi$.
- Contributions to individual form factors take the form

$$f_{0,\perp}(q^2) \simeq C_{f_{0,\perp}} \xi_\Lambda(n_+p') \mp \frac{2M}{n_+p'} \Delta\xi_\Lambda(n_+p')$$

$$f_+(q^2) \simeq C_{f_+} \xi_\Lambda(n_+p') - 2 \left(2 - \frac{M}{n_+p'} \right) \Delta\xi_\Lambda(n_+p') \quad \text{etc.}$$

[TF/Yip – preliminary]

Extraction of Matching Coefficients from Hard Vertex Corrections (Baryons)

Renormalization Scheme: $C_{f_+} = C_{g_+} \equiv 1$

- Using results from [Beneke/TF 2000]:

$$C_{f_\perp} = C_{g_\perp} = 1 + \frac{\alpha_s C_F}{4\pi} L,$$

$$C_{h_\perp} = C_{\tilde{h}_\perp} = 1 + \frac{\alpha_s C_F}{4\pi} \left(\ln \frac{M_b^2}{\mu^2} - 2 \right) \quad \text{etc.}$$

$$\text{with } L \equiv -\frac{M^2 - q^2}{q^2} \ln \left(1 - \frac{q^2}{M^2} \right).$$

Form Factor Relations *after* Inclusion of $\mathcal{O}(\alpha_s)$ Corrections ($\Delta\xi$ and C_i)

$$\begin{aligned} \frac{f_0 + h_+}{g_0 + \tilde{h}_+} &= \frac{f_\perp - h_+}{g_\perp - \tilde{h}_+} = \frac{f_+ + h_+ - 2h_\perp}{g_+ + \tilde{h}_+ - 2\tilde{h}_\perp} \\ &= -\frac{2(f_+ - f_\perp) + h_+ - h_\perp}{2(g_+ - g_\perp) + \tilde{h}_+ - \tilde{h}_\perp} = -\frac{M^2 - 2q^2}{M^2} \frac{h_+ - \tilde{h}_+}{f_+ - g_+} = 1. \end{aligned}$$

[TF/Yip – preliminary]

Zero of Forward-Backward Asymmetry, for values of q^2 satisfying

$$\text{Re} \left[C_9^{\text{eff}}(q^2) \right] + \frac{M_b (M_{\Lambda_b} + m_\Lambda) C_7^{\text{eff}}}{q^2} \frac{h_\perp}{f_\perp} + \frac{M_b (M_{\Lambda_b} - m_\Lambda) C_7^{\text{eff}}}{q^2} \frac{\tilde{h}_\perp}{g_\perp} \stackrel{!}{=} 0$$

- **Symmetry Limit:** Recover LO result from inclusive/mesonic case ✓
- Corrections from $\Delta\xi$ are additionally suppressed by m_Λ/M_{Λ_b} (!)
- Corrections from hard-vertex corrections C_i (model-independent)
- **Power Corrections** to form-factor relations (→ SR/Lattice)
- **Non-Factorizable Corrections** from long-distance photons (→ future work)

Conclusions

- Light-Cone/SCET Sum Rules provide a **Universal and Efficient Framework** to estimate (heavy-to-light) Transition Form Factors at Large Recoil Energy ✓
- Improved Estimates of **Hadronic Input Parameters and LCDAs** required, in particular for baryonic decays [see e.g. Alex, Klein, Mannel, Wang] !
- Better understanding of **Endpoint Logarithms** in Exclusive Matrix Elements ?

Still a lot to explore for ALEX and his Friends !

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- Light-Cone/SCET Sum Rules provide a **Universal and Efficient Framework** to estimate (heavy-to-light) Transition Form Factors at Large Recoil Energy ✓
- Improved Estimates of **Hadronic Input Parameters and LCDAs** required, in particular for baryonic decays [see e.g. Alex, Klein, Mannel, Wang] !
- Better understanding of **Endpoint Logarithms** in Exclusive Matrix Elements ?

All the best for the future, ALEX, and *Happy Birthday* !