

# Nuclear and Particle Physics Divisional Conference (NPPD)

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## THE LOW-ENERGY MAGNETIC AND ELECTRIC DIPOLE EXCITATIONS IN $180-186\text{W}$

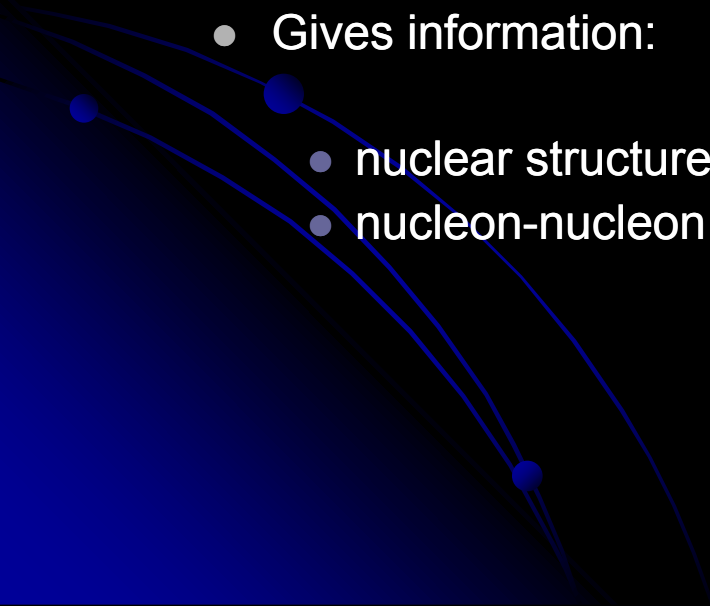
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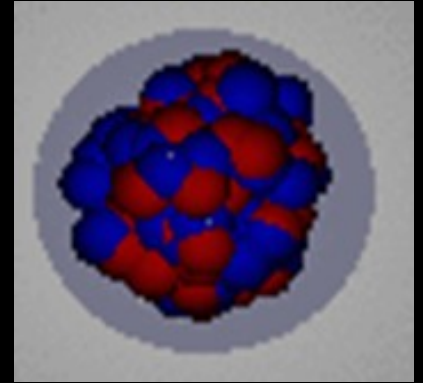
# INTRODUCTION

## Magnetic (scissors) and electric dipole mode:

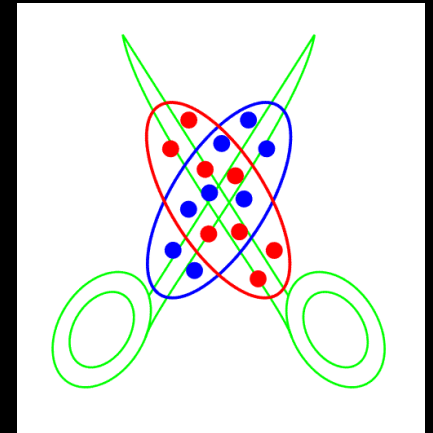
- great interest of modern nuclear structure physics.
  - Found for isotopes from light nuclei (such as  $^{46}\text{Ti}$ ) up to actinides and also gamma soft and transitional nuclei
  - Gives information:
    - nuclear structure
    - nucleon-nucleon forces
- 

- ❖ 2 characters in accordance with parities of dipole excitations:

$1^-$  = electric dipole mode




$1^+$  = magnetic (scissors) mode



- Quadratic dependence
- Strong fragmentation
- Concentrated around 3 MeV

# Goldstone (spurious state with zero energy) mode

- Hartree-Fock approximation: violates symmetries  $\rightarrow$  rotational or translational invariance broken
  - RPA used to calculate scissors mode in heavy deformed nuclei
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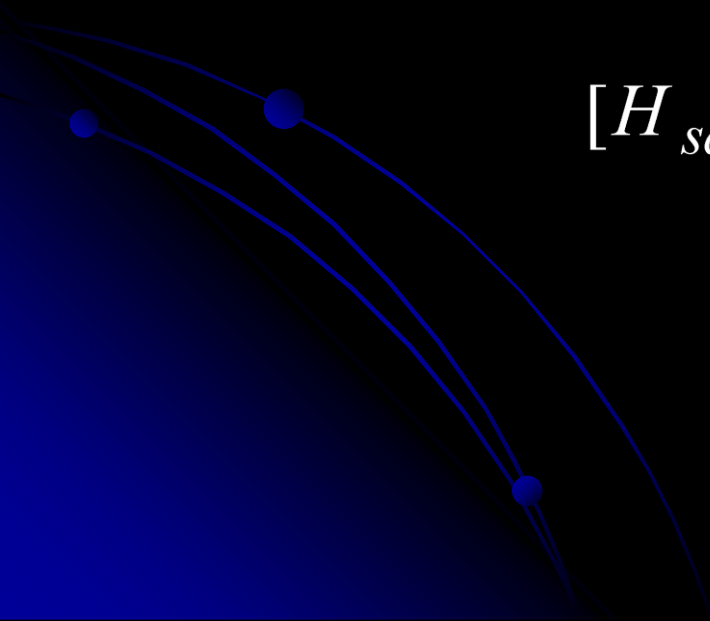
# THEORY

The corresponding single-quasiparticle hamiltonian



(1)

$$[H_{sqp}, J_{\mu}] \neq 0$$



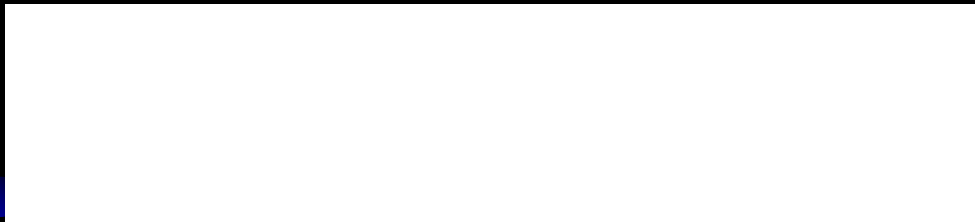
restored → separable isoscalar and isovector effective interaction

A large white rectangular box redacting the content of equation (2).

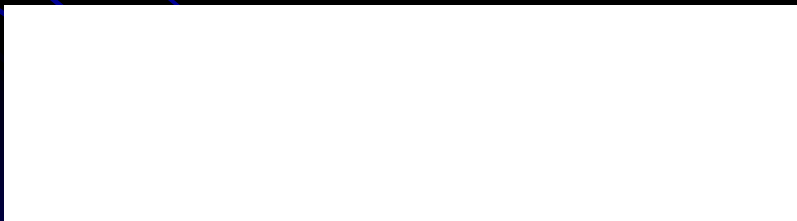
(2)

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(3)

A large white rectangular box redacting the content of equation (4).

(4)

A white rectangular box redacting the content of an equation.

$$\eta = \frac{V_1}{4V_0}$$

$V_0$  and  $V_1$  → the isoscalar and isovector depth of the potential wells

$$\gamma_0 = \gamma - \gamma_1 \quad \gamma = \gamma_n + \gamma_p \quad \gamma_1 = \gamma_1^n - \gamma_1^p$$


$$(5)$$

$$\gamma_1^{(-1)} = \gamma_1^{(+1)} \equiv \gamma_1 \quad \gamma^{(-1)} = \gamma^{(+1)} \equiv \gamma \quad (6)$$

the translation invariance can be obtained by the substitution

$$J_\mu \rightarrow P_\mu$$


The model Hamiltonian of the system



$$V_{\sigma\tau} = \frac{1}{2} \chi_{\sigma\tau} \sum_{i \neq j} \vec{\sigma}_i \cdot \vec{\sigma}_j \vec{\tau}_i \cdot \vec{\tau}_j$$



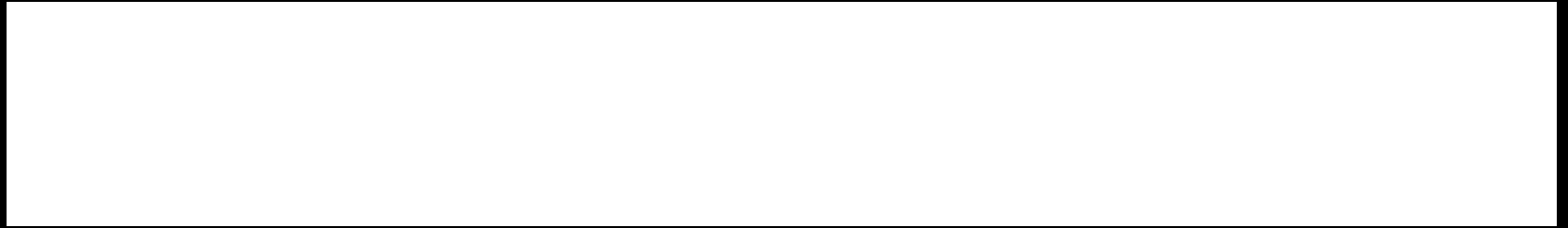
$1^+$  states are considered as one-phonon excitations


$$\sum_{\mu\tau} [\psi_{\mu}^{i2}(\tau) - \varphi_{\mu}^{i2}(\tau)] = 1$$

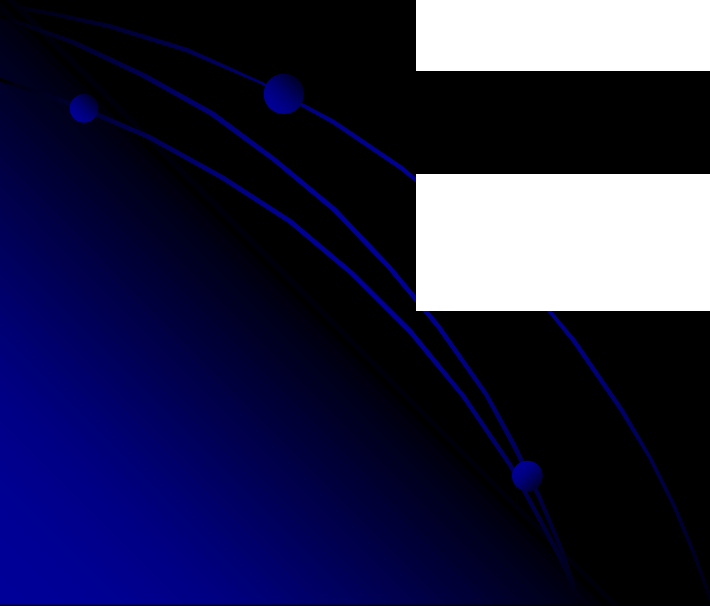
To obtain the excitations energies solve the equation of motion



# The reduced $M1$ transition probability



where



# Excitation energies

$$\bar{\omega} = \frac{\sum_i \omega_i B(M1, \omega_i)}{\sum_i B(M1, \omega_i)}$$

Ground-state transition width  $\Gamma_0$  (meV).

or

$$\Gamma_0^{red} = \frac{\Gamma_0}{\omega_i^3} (meV / MeV^3)$$

# Results and Discussion

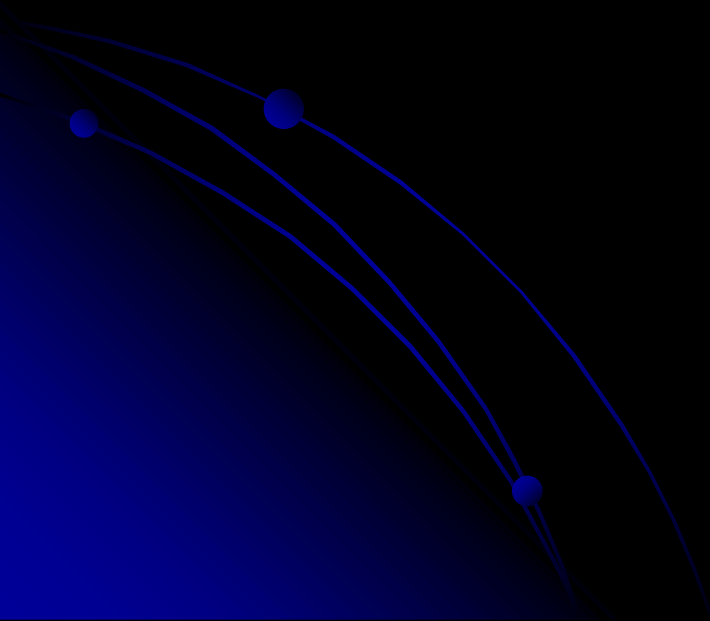
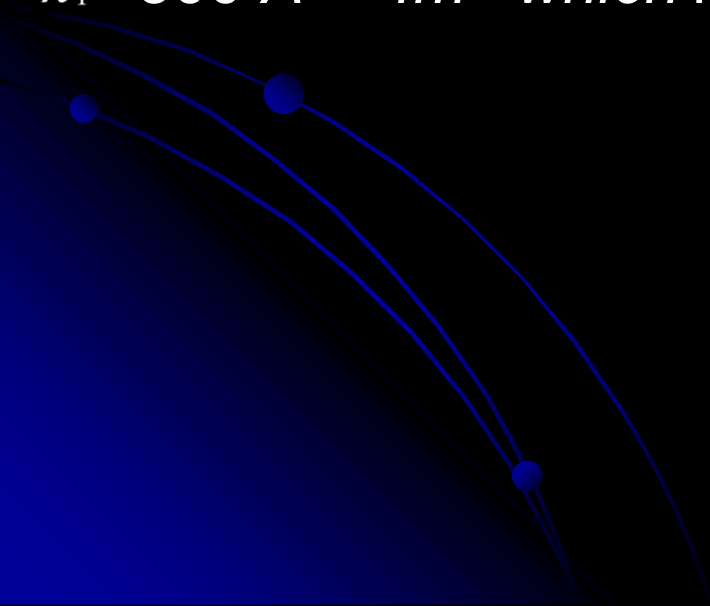


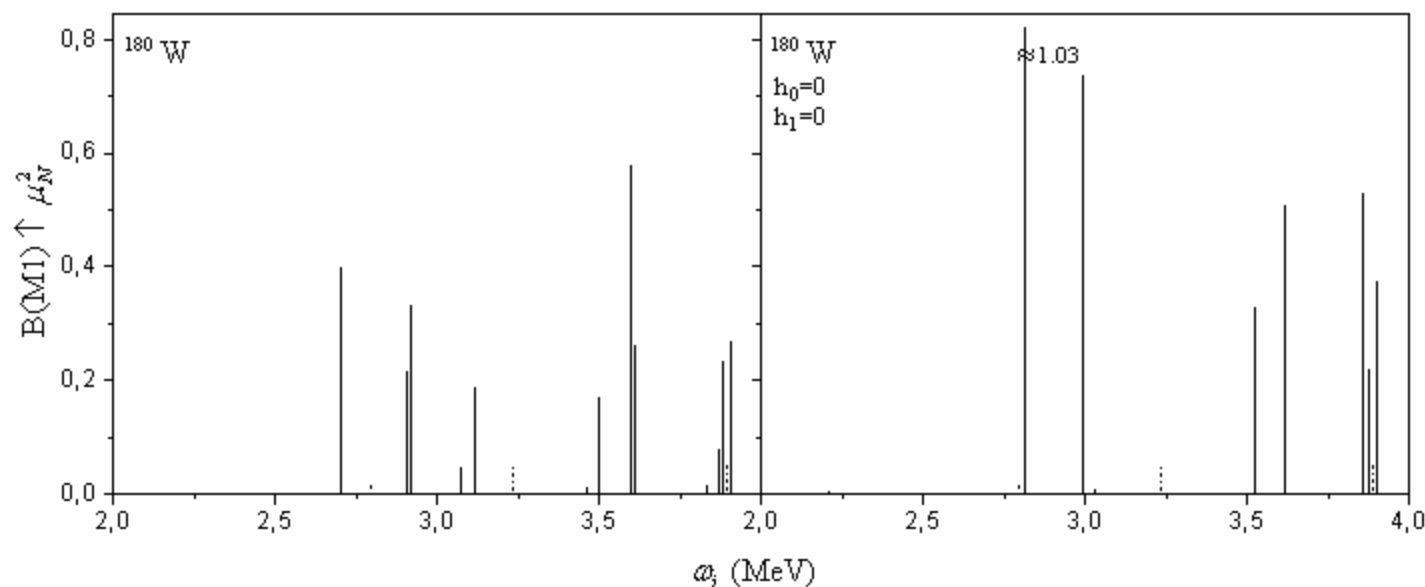
Table 1. The calculated values of the pairing correlation parameters (in MeV) and  $\delta_2$  values for the  $^{180-186}\text{W}$  isotopes.

N	$\Delta_n$	$\lambda_n$	$\Delta_p$	$\lambda_p$	$\delta_2$
106	1.20	-7.46	1.00	-5.80	0.210
108	0.90	-7.11	0.85	-6.47	0.207
110	0.97	-6.80	1.20	-6.98	0.201
112	1.00	-6.54	0.90	-7.62	0.190

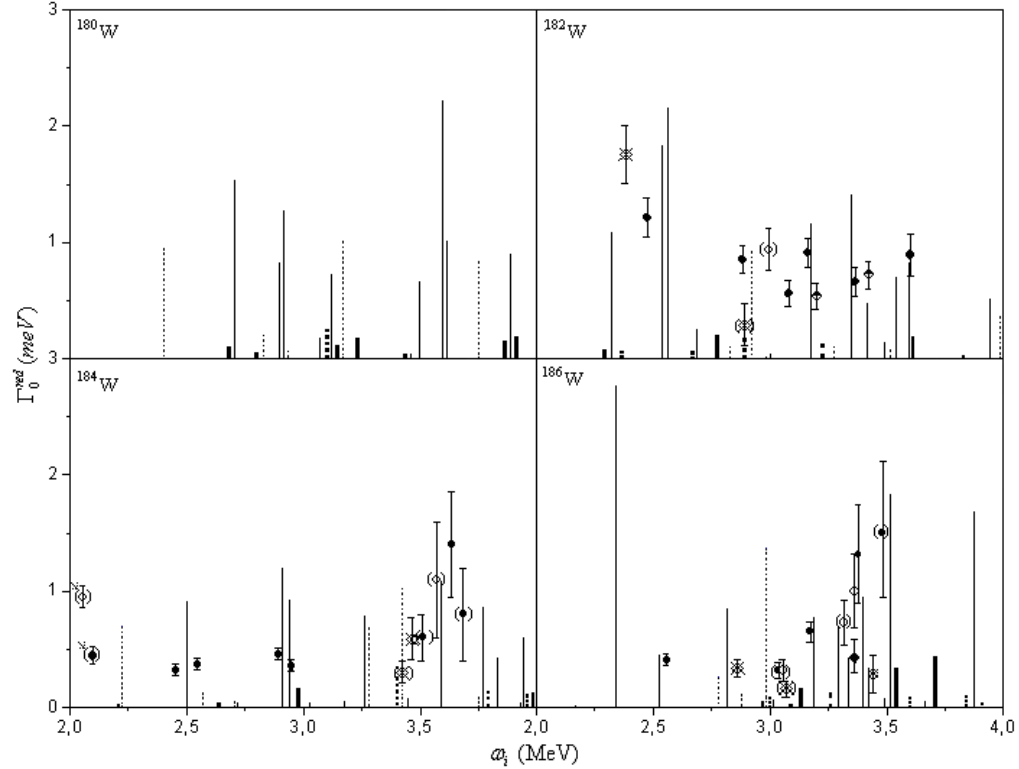
$\chi_{\sigma\tau} = 40/A \text{ MeV}$  which has been obtained from magnetic moments calculations.

$\chi_1 = 300 A^{-5/3} \text{ fm}^{-2}$  which is related to isovector symmetry potential.

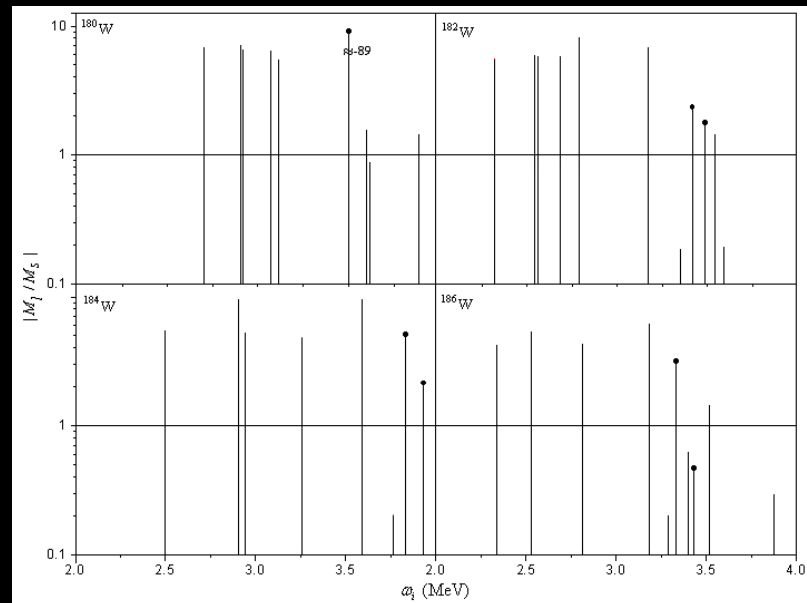




Rotational invariant $H = H_{sep} + h_0 + h_1 + V_{\sigma\tau}$				Rotational non-invariant $H = H_{sep} + V_{\sigma\tau}$			
$\omega_n$ (MeV)	B(M1) ( $\mu_N^2$ )	Structure $Nn_z \Lambda \Sigma$	Amplitudes $\psi_{ss'}^i$	$\omega_n$ (MeV)	B(M1) ( $\mu_N^2$ )	Structure $Nn_z \Lambda \Sigma$	Amplitudes $\psi_{ss'}^i$
		nn521↑-510↑	0.167				
		nn 512↑-521↑	0.158				
		nn 512↑-512↓	-0.208				
3.596	0.576	nn 514↑-523↑	0.171	3.520	0.329	nn 503↑-505↓	-0.979
		nn 633↑-642↑	-0.289				
		nn 503↑-505↓	0.569				
		pp 402↑-411↑	0.278				
		pp 402↑-411↑	-0.261				
		pp 402↑-402↓	-0.456				



Nuclei	$\bar{a}_{th.}$ (MeV)	$\bar{a}_{exp.}$ (MeV)	Energy Range
$^{180}\text{W}$	3.17	-	2.4-3.7
$^{182}\text{W}$	2.99	3.1	2.4-3.7
$^{184}\text{W}$	3.14	3.3	2.4-3.7
$^{186}\text{W}$	2.99	3.2	2.4-3.5





**Table 3.** The calculated ground-state M1, E1 and total  $\Gamma_o^{red}$  dipole transition widths (in  $\text{meV}/\text{MeV}^3$ ) are compared with each other and the experimental data [9,10] in  $^{180-186}\text{W}$ .

Nuclei	Energy interval (MeV)	$\Gamma_o^{red}$ ( $\text{meV}/\text{MeV}^3$ )	K			
			0	1	Unknown	$\Gamma_o^{red}$ (tot.)
$^{180}\text{W}$	2.0-4.0	M1	0.42	9.37	-	9.79
		E1	0.65	3.13	-	3.78
		$\Gamma_o^{red}$ (tot.)	1.07	12.5	-	13.6
		Exp.	-	-	-	-
$^{182}\text{W}$	2.0-3.7	M1	0.46	10.1	-	10.6
		E1	0.40	1.2	-	1.6
		$\Gamma_o^{red}$ (tot.)	0.86	11.3	-	12.2
		Exp.	$0.94 \pm 0.18$	$6.34 \pm 0.28$	$2.05 \pm 0.1$	$9.33 \pm 0.56$
$^{184}\text{W}$	2.0-3.7	M1	0.24	5.30	-	5.54
		E1	0.34	2.62	-	2.96
		$\Gamma_o^{red}$ (tot.)	0.58	7.92	-	8.50
		Exp.	$2.05 \pm 0.59$	$4.73 \pm 1.28$	$0.89 \pm 0.28$	$7.67 \pm 2.15$
$^{186}\text{W}$	2.5-3.5	M1	0.16	6.46	-	6.62
		E1	0.26	1.84	-	2.10
		$\Gamma_o^{red}$ (tot.)	0.42	8.30	-	8.72
		Exp.	$2.05 \pm 0.6$	$4.66 \pm 2.32$	$0.77 \pm 0.29$	$7.48 \pm 3.21$

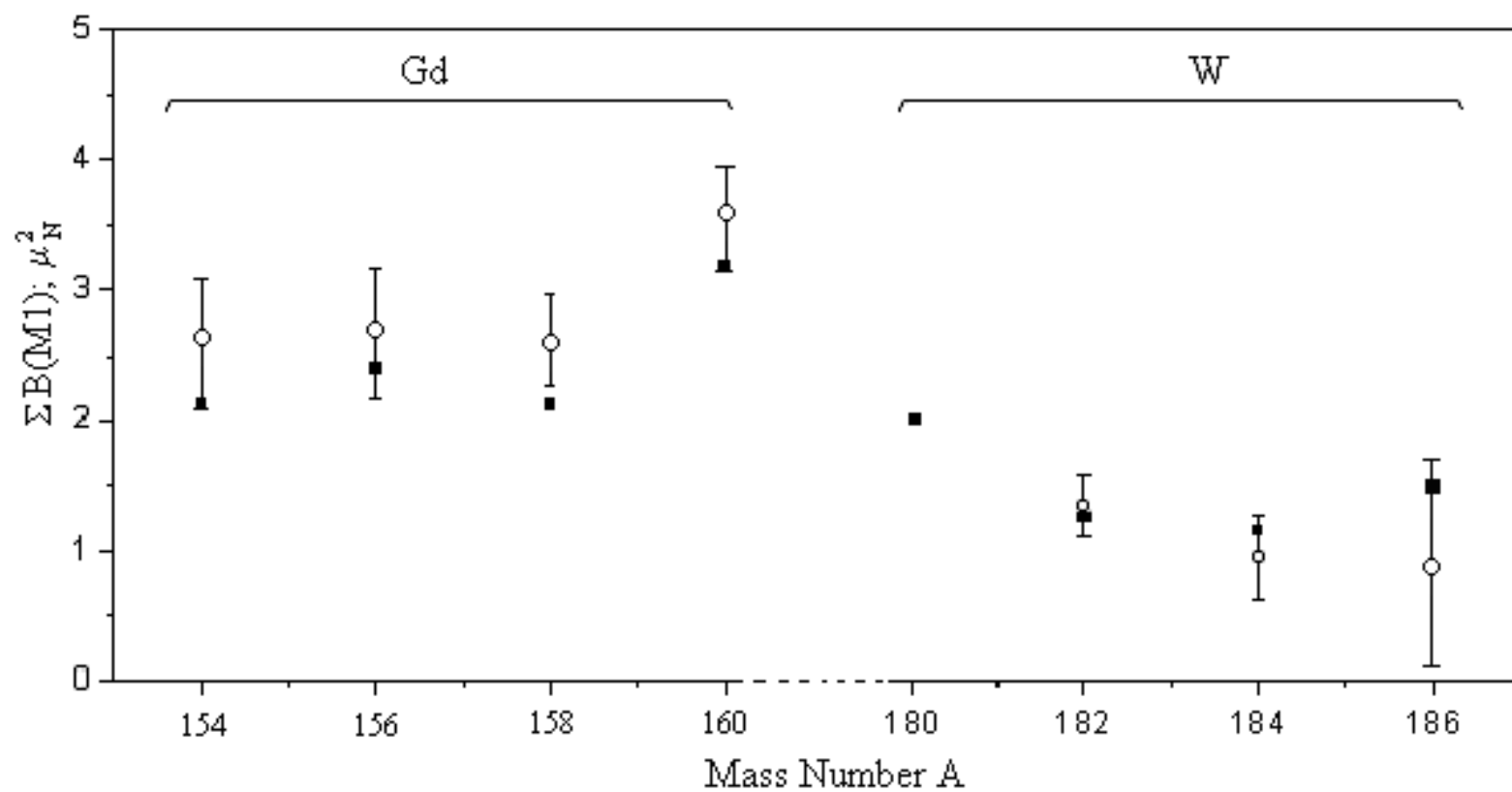
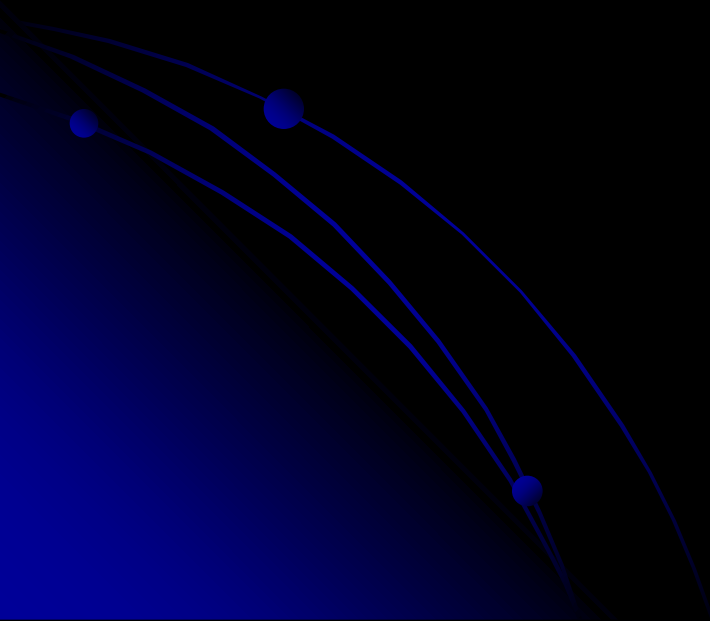


Fig. Compared calculated summed B(M1) strength (■) with experimental data (○) results for  $^{154-160}\text{Gd}$  and  $^{180-186}\text{W}$ .

# CONCLUSION

- A separation of the spurious state from the  $1+$  states changes
  - The distribution of the  $B(M1)$  strength
  - Increases the fragmentation
- The calculations provided detailed information
  - the excitation energies
  - transitional probabilities
  - structure of low-lying dipole excitations

- The relative contribution of  $\Delta K = 0$  below 4 MeV is quite small.
- A few prominent  $K^\pi = 1^-$  states in the low-energy.



# Collaborators

E. Guliyev and A.A.Kuliev

and

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