

Precise prediction for the W-boson mass in U(1) extensions of the standard model

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Standard Model and Beyond
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Outline

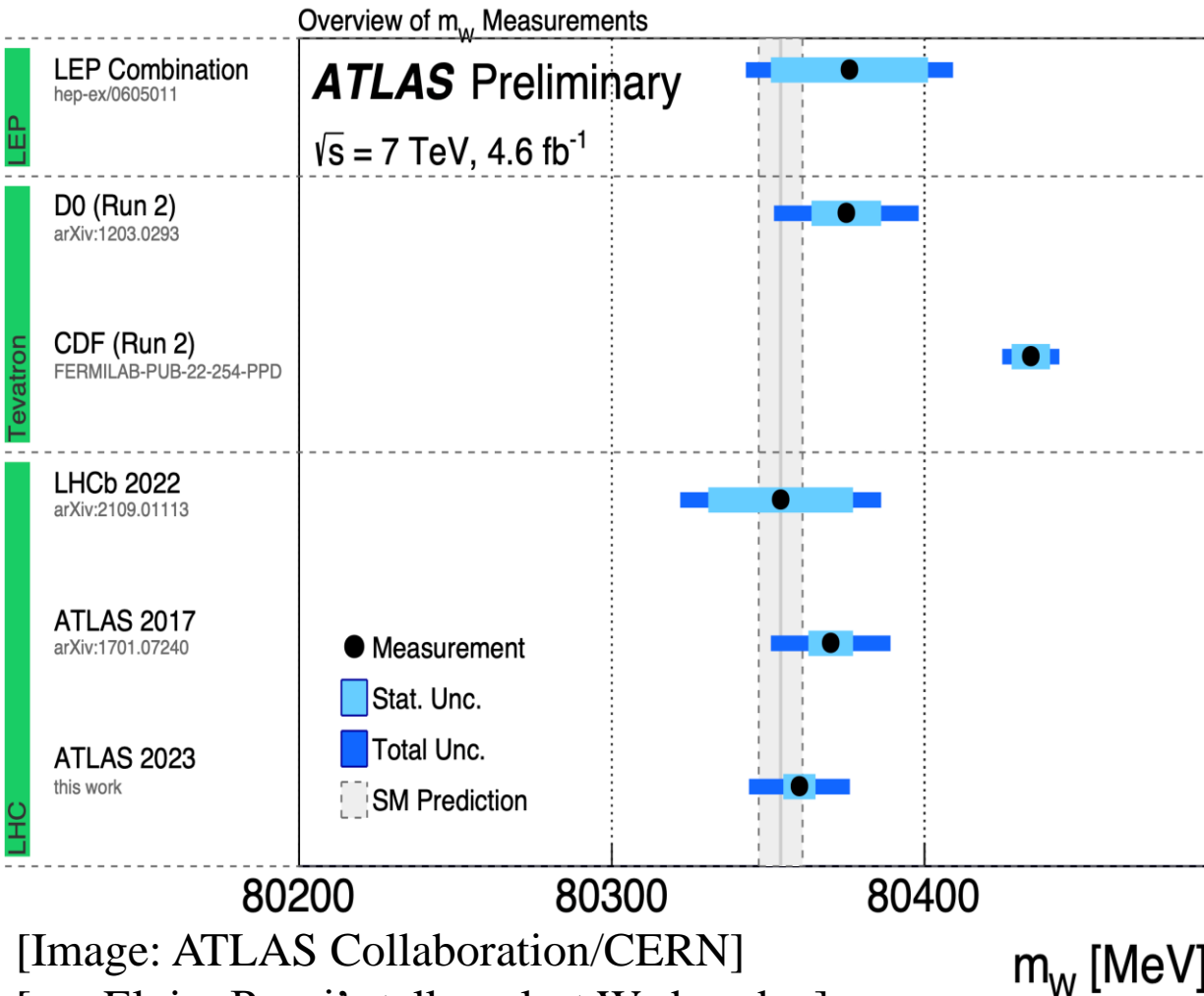
- W-boson mass and BSM physics
- U(1) extensions of the standard model
- On the effect of the full 1-loop correction to the W-mass



BSM and M_W

- Small uncertainty (~ 10 MeV) in theory and experiment

Prediction & measurement of M_W



- **Theory** (SM, $\overline{\text{MS}}$ -bar [1411.7040]):
 $M_W^{\text{theo}} = 80353 \pm 9 \text{ MeV}$ (with PDG 2022 inputs)
- **Experiment** (PDG 2022 world. avg.):
 $M_W^{\text{exp.}} = 80377 \pm 12 \text{ MeV}$

[Image: ATLAS Collaboration/CERN]
[see Elvira Rossi's talk on last Wednesday]

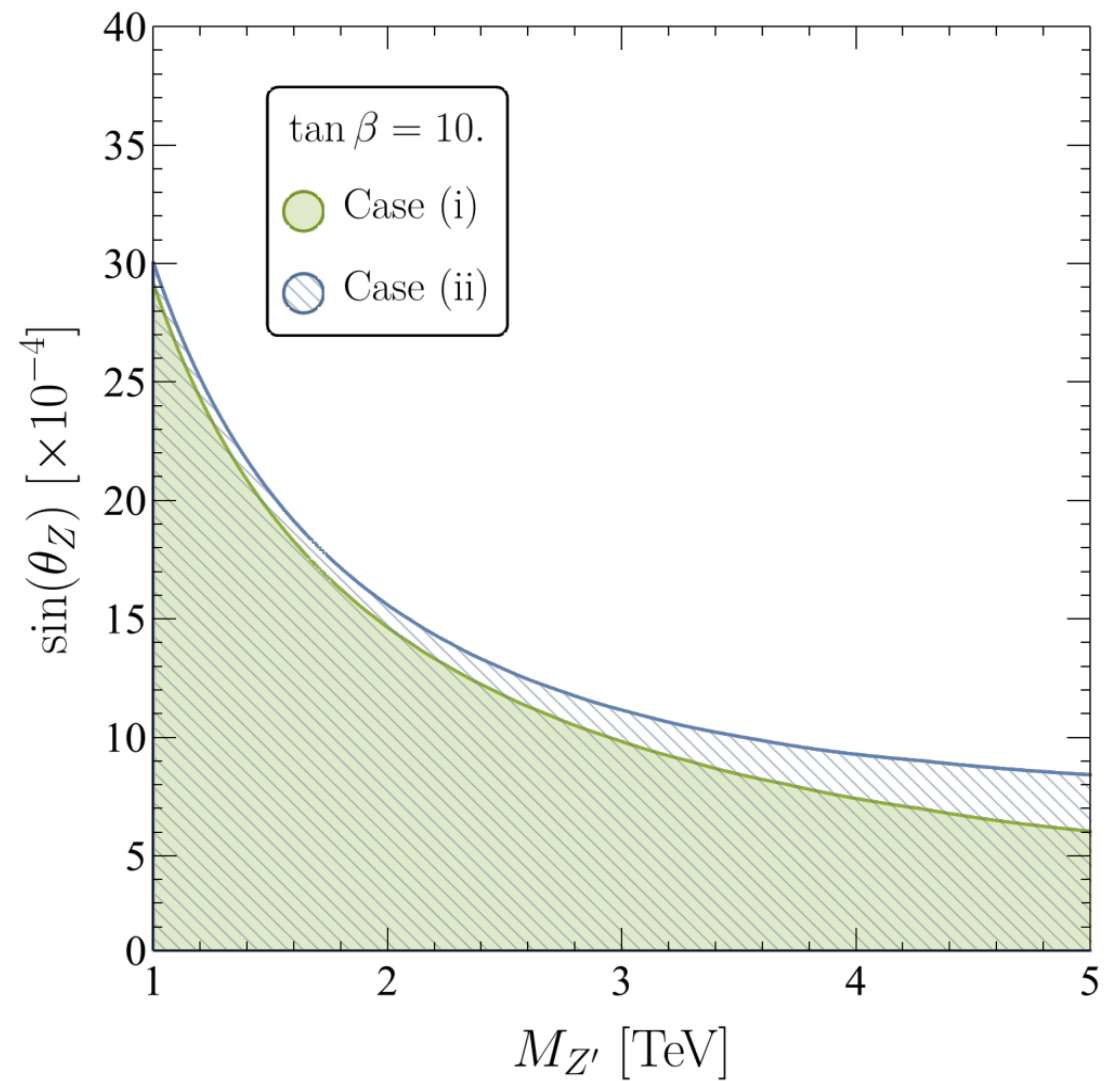


BSM and M_W

- Small uncertainty (~ 10 MeV) in theory and experiment
- Small BSM effects on M_W can be exposed
- U(1) extensions affect M_W at tree level so precision is important (1-loop)

- Precise predictions in BSM models are important
- Full Δr at 1-loop in **U(1) extensions** is computed
- **Full Δr (Case i.) may become important for heavy $M_{Z'}$**
- ...compared to the available predictions (Case ii.)
- Fig. shows **region where :**

$$|M_W^{\text{exp.}} - M_W| < 2\sigma$$



Take home message

What's new in a $U(1)$ extension?

- SM gauge group + an extra $U(1)$ adds a **new interaction**
- May add **new scalar field(s)**, can stabilize the EW vacuum
- May add **right-handed (sterile) neutrinos**: neutrino mass generation via see-saw, dark matter



See Zoltán Trócsányi's talk on the superweak model on Tuesday!

New parameters:
5
(2 gauge + 3 scalar)

Gauge sector:

- M'_Z : mass of the new gauge boson Z'
- s_Z : new gauge mixing angle, rotation of gauge eigenstates to mass eigenstates:

$$\begin{pmatrix} B_\mu \\ W_\mu^3 \\ B'_\mu \end{pmatrix} = \begin{pmatrix} c_W & -s_W & 0 \\ s_W & c_W & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_Z & -s_Z \\ 0 & s_Z & c_Z \end{pmatrix} \begin{pmatrix} A_\mu \\ Z_\mu \\ Z'_\mu \end{pmatrix}$$

Scalar sector:

- $\tan\beta = \frac{w}{v}$: ratio of new VEV to BEH VEV
- M_S : mass of the new scalar boson
- s_S : new scalar mixing angle to mass eigenstates

$$\begin{pmatrix} \phi^0 \\ \chi \end{pmatrix} = \begin{pmatrix} c_S & -s_S \\ s_S & c_S \end{pmatrix} \begin{pmatrix} h \\ S \end{pmatrix}$$

Concise relation:

$$\frac{M_W^2}{c_W^2} = c_Z^2 M_Z^2 + s_Z^2 M_{Z'}^2$$

**Express predictions with
Lagrangian couplings or pheno
parameters e.g.:**

$$\rho = \frac{M_W^2}{c_W^2 M_Z^2} = 1 - s_Z^2 \left(1 - \frac{M_{Z'}^2}{M_Z^2} \right)$$

**Gauge
boson
masses**

Concise relation:

$$\frac{M_W^2}{c_W^2} = c_Z^2 M_Z^2 + s_Z^2 M_{Z'}^2$$

**Express predictions with
Lagrangian couplings or pheno
parameters e.g.:**

$$M_W^2 = \frac{\rho M_Z^2}{2} \left(1 + \sqrt{1 - \frac{4\pi\alpha}{\sqrt{2} G_F \rho M_Z^2}} (1 + \Delta r) \right)$$

**Gauge
boson
masses**

Renormalization

- Split bare parameters into $g^{(0)} \rightarrow g + \delta g$
- The Weinberg angle changes at tree level:

$$M_W^2 \frac{\delta c_W^2}{c_W^2} = \delta M_W^2 - c_W^2 (c_Z^2 \delta M_Z^2 + s_Z^2 \delta M_{Z'}^2 - 2s_Z (M_Z^2 - M_{Z'}^2) \delta s_Z)$$

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Δr receives completely new corrections:

Δr collects the radiative corrections to the μ -decay and hence to M_W

$$\Delta r = (\text{formally } \Delta r^{\text{SM}} \text{ with BSM loops}) - s_Z^2 \frac{c_W^2}{s_W^2} \frac{c_W^2}{M_W^2} \left(\text{Re}\Pi_{ZZ}(M_Z^2) - \text{Re}\Pi_{Z'Z'}(M_{Z'}^2) + 2(M_Z^2 - M_{Z'}^2) \frac{\delta s_Z}{s_Z} \right)$$

$$\frac{G_F}{\sqrt{2}} = \frac{\pi \alpha}{2 M_W^2 s_W^2} (1 + \Delta r)$$

$$\Delta r^{\text{SM}} = \underbrace{\frac{2\delta e}{e}} + \underbrace{\left(\frac{\text{Re}\Pi_{WW}(M_W^2) - \Pi_{WW}(0)}{M_W^2} \right) + \delta_{\text{BV}}}$$

Renormalization of the electric charge, formula known all order

Diagrammatic corrections to the muon decay graph: W-propagator and box and vertex diagrams

$$+ \underbrace{\frac{c_W^2}{s_W^2} \left(\frac{\text{Re}\Pi_{ZZ}(M_Z^2)}{M_Z^2} - \frac{\text{Re}\Pi_{WW}(M_W^2)}{M_W^2} \right)}$$

Renormalization of s_W

Checks

- The ε poles cancel in Δr in R_ξ -gauge with general z-charge assignment
- For several benchmark points Δr is independent of the gauge parameters ξ_i , with $i = W, A, Z, Z'$
- Compare Δr in two cases:

Case I. :

$$\Delta r = (\text{formally } \Delta r^{\text{SM}} \text{ with BSM loops}) - s_Z^2 \frac{c_W^2}{s_W^2} \frac{c_W^2}{M_W^2} \left(\text{Re}\Pi_{ZZ}(M_Z^2) - \text{Re}\Pi_{Z'Z'}(M_{Z'}^2) + 2(M_Z^2 - M_{Z'}^2) \frac{\delta s_Z}{s_Z} \right)$$

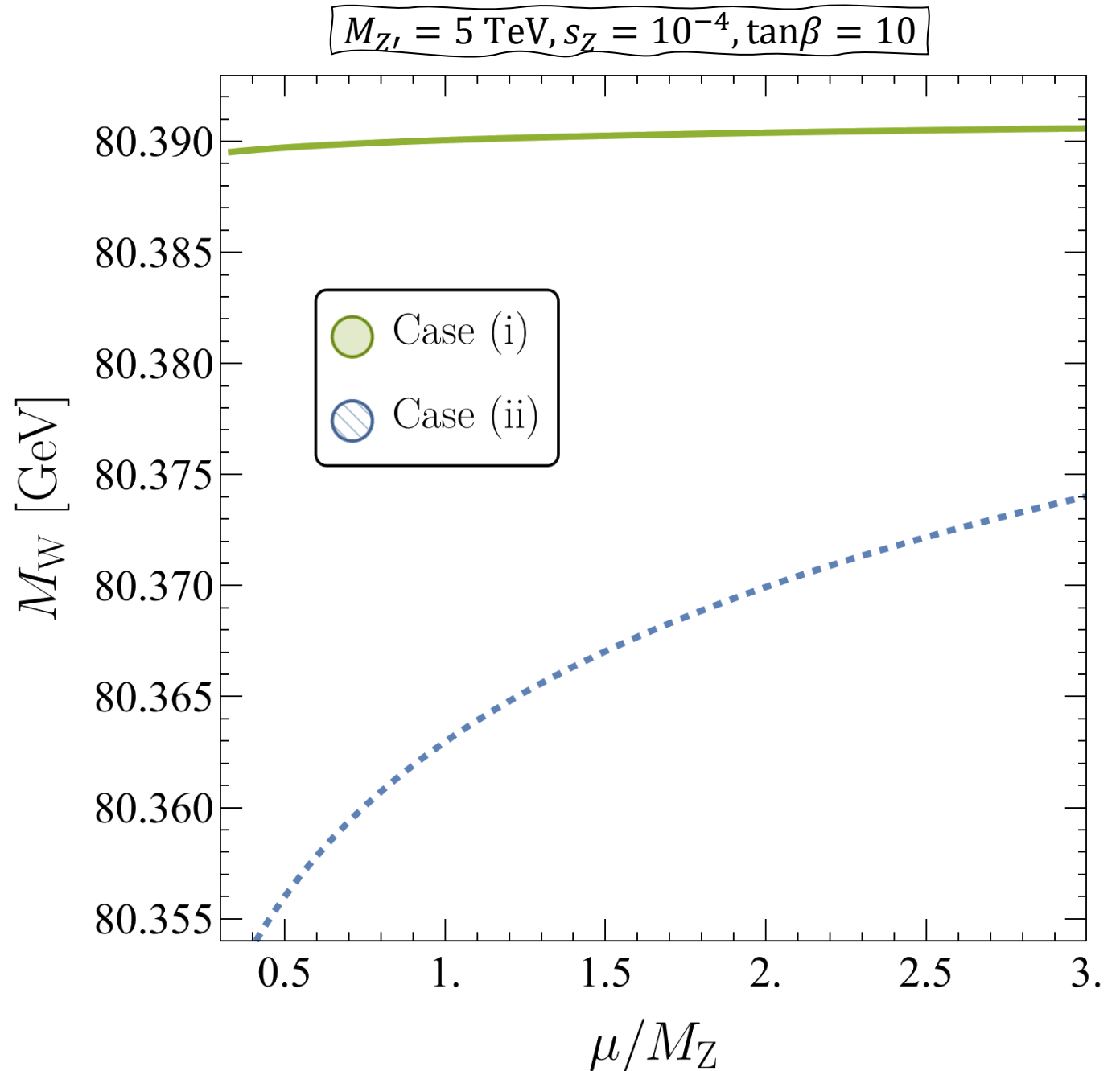
When is it safe to neglect the **new** terms?

Case II. :

$$\Delta r = (\text{formally } \Delta r^{\text{SM}} \text{ with BSM loops})$$

Checks

- The ε poles cancel in Δr in R_ξ -gauge with general z-charge assignment
- For several benchmark points Δr is independent of the gauge parameters ξ_i , with $i = W, A, Z, Z'$
- Weak dependence on the renormalization scale μ at fixed benchmark points



Benchmarks: $M_W - M_{W,SM}$ [MeV]

SMALL $M_{Z'}$ = 50 MeV
 and $s_S = 0.1$
Irrelevant

s_Z		$5 \cdot 10^{-4}$			
$\tan \beta$	M_S	0.5 TeV		5 TeV	
		(i)	(ii)	(i)	(ii)
0.1		-1	-1	-2	-2
1		-1	-1	-2	-2
10		-1	-1	-2	-2

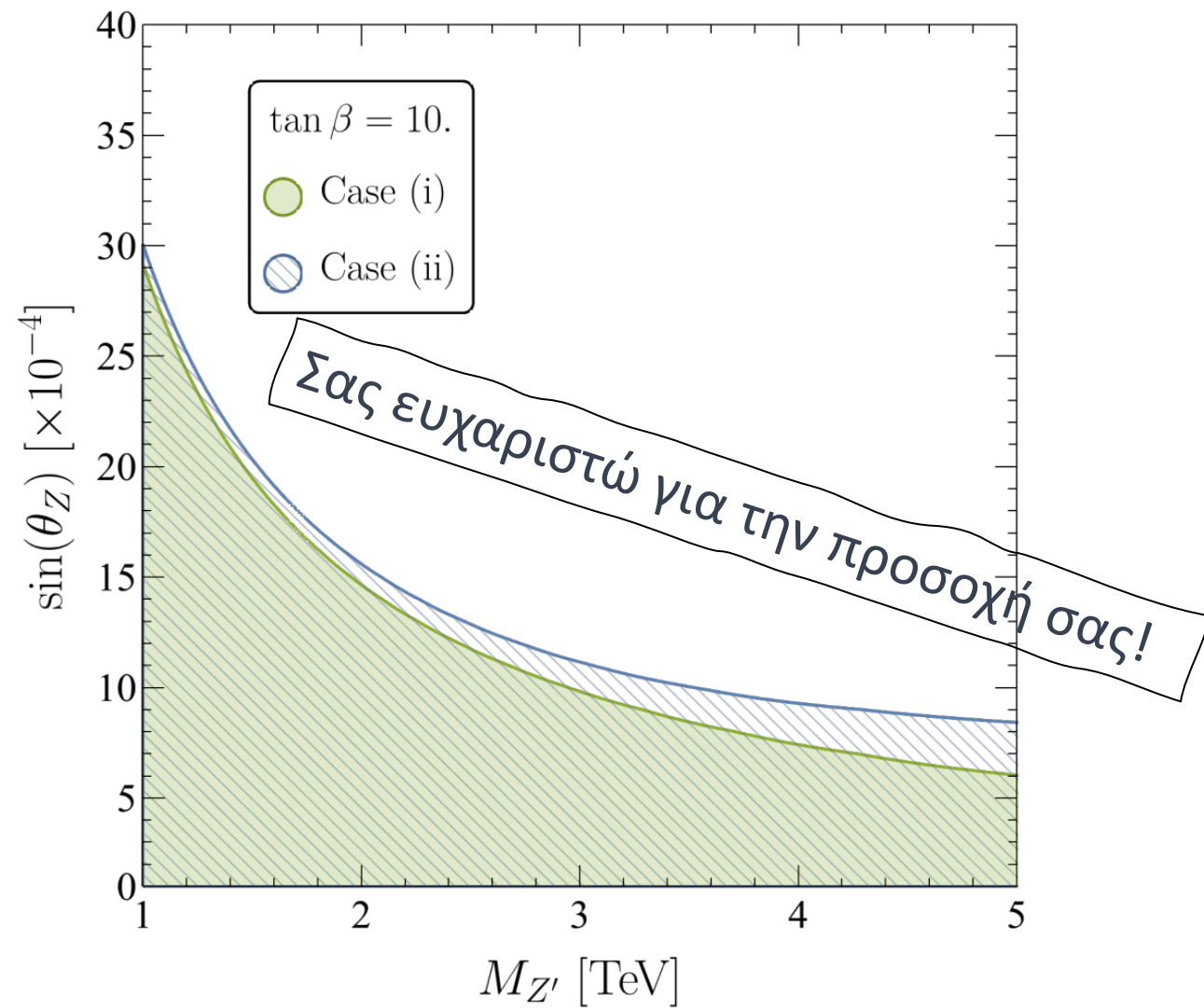
BSM corrections to the SM prediction for M_W in MeV units

LARGE $M_{Z'}$ = 5 TeV
 and $s_S = 0.1$
Potentially relevant

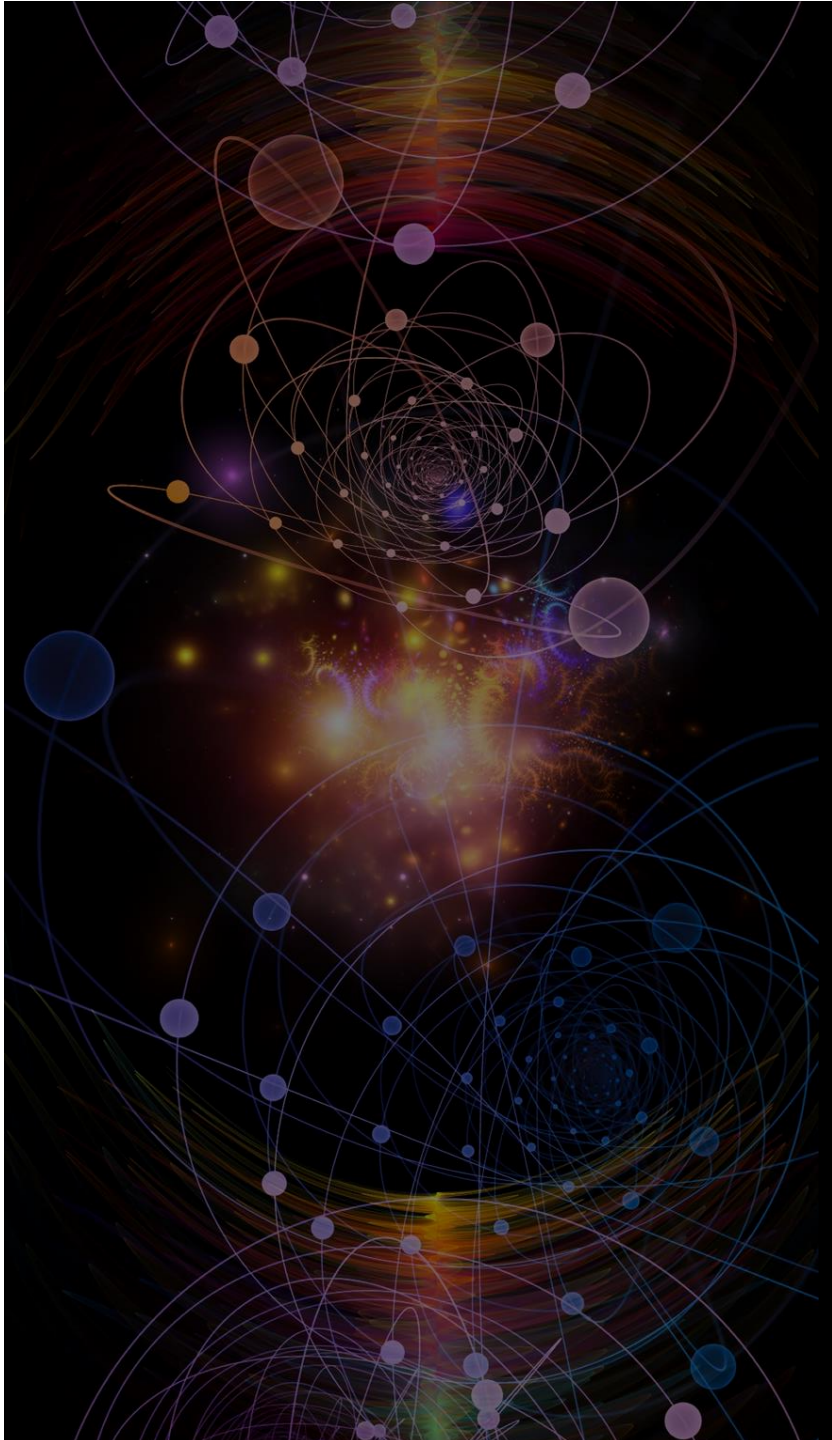
s_Z		$5 \cdot 10^{-4}$				$7 \cdot 10^{-4}$			
$\tan \beta$	M_S	0.5 TeV		5 TeV		0.5 TeV		5 TeV	
		(i)	(ii)	(i)	(ii)	(i)	(ii)	(i)	(ii)
10		37	10	35	13	75	29	73	36
20		39	34	35	34	81	76	74	79
30		40	38	35	37	83	85	75	85

- Precise predictions in BSM models are important
- Full Δr at 1-loop in **U(1) extensions** is computed
- **Full Δr (Case i.) may become important for heavy $M_{Z'}$**
- ...compared to the available predictions (Case ii.)
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Conclusions



Backup slides

How to obtain δs_Z I.

- Relate unrotated and rotated fields:

$$B_\mu^{(0)} = c_W^{(0)} A_\mu^{(0)} - s_W^{(0)} (c_Z^{(0)} Z_\mu^{(0)} - s_Z^{(0)} Z_\mu^{\prime(0)})$$

$$B_\mu^{\prime(0)} = s_Z^{(0)} Z_\mu^{(0)} + c_Z^{(0)} Z_\mu^{\prime(0)}$$

- Also true for renormalized fields:

$$B_\mu = c_W A_\mu - s_W (c_Z Z_\mu - s_Z Z'_\mu)$$

$$B'_\mu = s_Z Z_\mu + c_Z Z'_\mu$$

- Unrotated fields are renormalized such that

$$B_\mu^{(0)} = \sqrt{Z_B} B_\mu \text{ and } B_\mu^{\prime(0)} = \sqrt{Z_{B'}} B'_\mu$$

- Rotated fields may mix:

$$\begin{pmatrix} A_\mu^{(0)} \\ Z_\mu^{(0)} \\ Z_\mu^{\prime(0)} \end{pmatrix} = \begin{pmatrix} \sqrt{Z_{AA}} & \frac{1}{2} Z_{AZ} & \frac{1}{2} Z_{AZ'} \\ \frac{1}{2} Z_{ZA} & \sqrt{Z_{ZZ}} & \frac{1}{2} Z_{ZZ'} \\ \frac{1}{2} Z_{Z'A} & \frac{1}{2} Z_{Z'Z} & \sqrt{Z_{Z'Z'}} \end{pmatrix} \begin{pmatrix} A_\mu \\ Z_\mu \\ Z'_\mu \end{pmatrix}$$

How to obtain δs_Z II.

- Express bare fields with renormalized ones and collect coefficients:

$$\begin{aligned}\sqrt{Z_B} c_W &= c_W^{(0)} \sqrt{Z_{AA}} - \frac{1}{2} s_W^{(0)} \left(c_Z^{(0)} Z_{ZA} - s_Z^{(0)} Z_{Z'A} \right) \\ \sqrt{Z_{B'}} s_Z &= s_Z^{(0)} \sqrt{Z_{ZZ}} + \frac{1}{2} c_Z^{(0)} Z_{Z'Z} \\ \sqrt{Z_{B'}} c_Z &= \frac{1}{2} s_Z^{(0)} Z_{ZZ'} + c_Z^{(0)} \sqrt{Z_{Z'Z'}}\end{aligned}$$

- First equation is used to derive δe [hep-ph/0209084]
(U(1) Ward identity $\sqrt{Z_B} Z_{g_y} = 1$)
- 2nd and 3rd ones are divided to cancel $\sqrt{Z_{B'}}$ and express δs_Z