

# Supergeometry in Effective QFTs

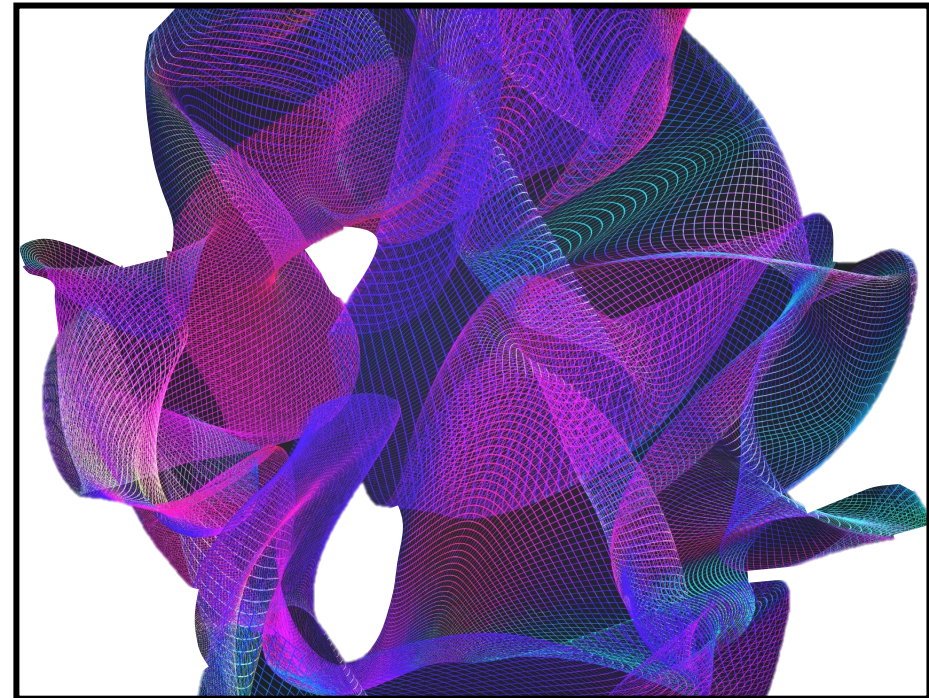
**Viola Gattus**

**Corfu Summer Institute 2023, Corfu, Greece**

Based on [arXiv:2307.01126](https://arxiv.org/abs/2307.01126)  
with Prof Apostolos Pilaftsis



The University of Manchester





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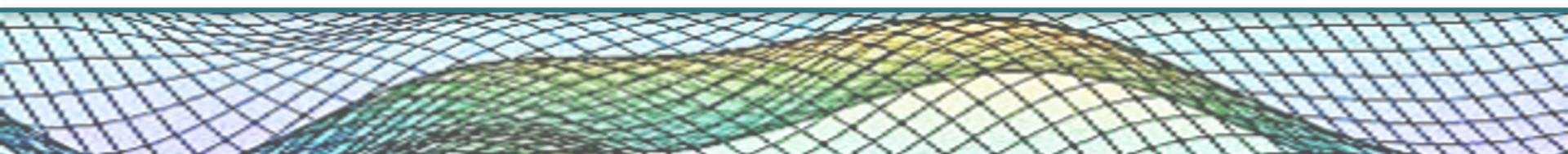
## **Model II**

A minimal non-factorizable  
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## **Supervertices**

How to write covariant  
scalar-fermion vertices?





# Motivation

Disclaimer: Supergeometry  $\neq$  supersymmetry

- ❑ A theory with fermions and bosons with no extra symmetry
- ❑ Use VDW formalism [Vilkovisky (1984), DeWitt (1985)]
- ❑ Applications to geometric EFTs [Alonso, Jenkins, Manohar (2016), Cohen, Craig, Sutherland (2021), Talbert (2023), Assi, Helset, Manohar, Pagès, Shen (2023) ...]
- ❑ Solving the frame-dependence problem in cosmic inflation [Burns, Karamitsos, Pilaftsis (2016), Falls, Herrero-Valea (2019), Finn, Karamitsos, Pilaftsis (2020) ..]



# The Set-Up

- Field-space supermanifold of dimension  $(N|8M)$  in 4D spacetime

[DeWitt (2012)]



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□ Now fermions in the chart

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$$\Phi \equiv \{\Phi^\alpha\} = (\phi^A, \psi^X, \bar{\psi}^{Y,T})^\top$$



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$$\mathcal{L} = \frac{1}{2} g^{\mu\nu} \partial_\mu \Phi^\alpha \alpha k_\beta(\Phi) \partial_\nu \Phi^\beta + \frac{i}{2} \zeta_\alpha^\mu(\Phi) \partial_\mu \Phi^\alpha - U(\Phi)$$



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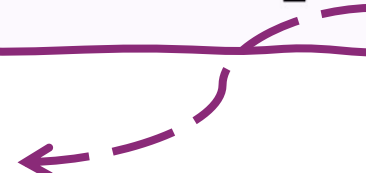
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$$\zeta_\beta^\mu \left( \overleftarrow{\Sigma}_\mu \right)_\alpha = \zeta_\alpha \quad \text{where} \quad \overleftarrow{\Sigma}_\mu = \frac{1}{D} \begin{pmatrix} \frac{\overleftarrow{\partial}}{\partial \gamma^\mu} & 0 \\ 0 & \Gamma_\mu \end{pmatrix}$$



# The Set-Up (continued)

- Endow supermanifold with metric

$${}_{\alpha}G_{\beta} = ({}_{\alpha}G_{\beta})^{sT}$$



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- Global metric found from vielbeins and local metric

$${}_{\alpha}G_{\beta} = {}_{\alpha}e^a \quad {}_aH_b \quad {}^b e_{\beta}^{sT} :$$

[Finn, Karamitsos, Pilaftsis (2021), VG, Finn, Karamitsos, Pilaftsis (2022)]

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$${}_{\alpha}G_{\beta} = {}_{\alpha}e^a \boxed{{}_aH_b} {}^b e_{\beta}^{sT}$$

$${}_aH_b \equiv \begin{pmatrix} \mathbf{1}_N & 0 & 0 \\ 0 & 0 & \mathbf{1}_{4M} \\ 0 & -\mathbf{1}_{4M} & 0 \end{pmatrix}$$



# No-Go Theorem

- Flat field space can **always** be reparametrized into canonical Cartesian form

$$\mathcal{L} = -\frac{1}{2}\mathbf{h}(\phi)\bar{\psi}\gamma^\mu\psi(\partial_\mu\phi) + \frac{i}{2}\mathbf{g}(\phi)\left[\bar{\psi}\gamma^\mu(\partial_\mu\psi) - (\partial_\mu\bar{\psi})\gamma^\mu\psi\right]$$

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- $\phi$  acts as external parameter in the fermionic sector



# No-Go Theorem (continued)

- Flatness confirmed by vanishing of Riemann tensor

$${}_{\alpha}G_{\beta} = \begin{pmatrix} k - \frac{1}{2}\bar{\psi}(\mathbf{g}' - i\mathbf{h})\mathbf{g}^{-1}(\mathbf{g}' + i\mathbf{h})\psi & -\frac{1}{2}\bar{\psi}(\mathbf{g}' - i\mathbf{h}) & \frac{1}{2}\psi^{\top}(\mathbf{g}'^{\top} + i\mathbf{h}^{\top}) \\ \frac{1}{2}(\mathbf{g}'^{\top} - i\mathbf{h}^{\top})\bar{\psi}^{\top} & 0 & \mathbf{g}^{\top}1_4 \\ -\frac{1}{2}(\mathbf{g}' + i\mathbf{h})\psi & -\mathbf{g}1_4 & 0 \end{pmatrix}$$

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—————▶  $R^{\alpha}_{\beta\gamma\delta} = 0$

[VG, Pilaftsis (2023)]

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## Take home message:

Non-zero fermionic curvature effects cannot be generated if  $\zeta_{\alpha}^{\mu}$  depends **linearly** on  $\psi$  and  $\bar{\psi}$



# Model I

- A 2D factorizable model

$$\mathcal{L}_I = \frac{1}{2}k (\partial_\mu \phi) (\partial^\mu \phi) + \frac{i}{2} (g_0 + g_1 \bar{\psi} \psi) [\bar{\psi} \gamma^\mu (\partial_\mu \psi) - (\partial_\mu \bar{\psi}) \gamma^\mu \psi]$$



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- FS metric

$$\mathbf{G} = \begin{pmatrix} k + b^\top (d^{-1})^\top a^\top - a d^{-1} b & -a & b^\top \\ & a^\top & 0 & d^\top \\ & -b & -d & 0 \end{pmatrix}$$

$$a = \frac{1}{2} \bar{\psi} (g'_0 + g'_1 \bar{\psi} \psi)$$

$$b = \frac{1}{2} (g'_0 + g'_1 \bar{\psi} \psi) \psi$$

$$d = (g_0 + g_1 \bar{\psi} \psi) 1_2 + g_1 \psi \bar{\psi}$$



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- Ricci scalar is fermion dependent

$$R = \frac{4g_1}{g_0^2} + \left( \frac{2g_1 g'_0 g'_1}{g_0^3 k} - \frac{2g_1^2 g_0'^2}{g_0^4 k} - \frac{g_1'^2}{2g_0^2 k} \right) (\bar{\psi} \psi)^2$$





# Model II

- A 4D non-factorizable model

$$\mathcal{L}_{\text{II}} = \frac{i}{2} \left[ \bar{\psi} \gamma^\mu (\partial_\mu \psi) - (\partial_\mu \bar{\psi}) \gamma^\mu \psi \right] + \frac{i}{2} \bar{\psi} \gamma^\mu \psi \left[ \bar{\psi} (\partial_\mu \psi) - (\partial_\mu \bar{\psi}) \psi \right]$$



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$$\mathbf{G} \equiv \{_\alpha G_\beta\} = \begin{pmatrix} 0 & d^\top \\ -d & 0 \end{pmatrix} \quad d = 1_4 + \frac{1}{4} (\bar{\psi} \gamma^\mu \psi) \gamma_\mu + \frac{1}{4} \gamma^\mu \psi \bar{\psi} \gamma_\mu$$



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- Richer structure of Ricci scalar

$$R = -8 + 2(\bar{\psi}\psi) + \frac{23}{8}(\bar{\psi}\psi)^2 + \frac{9}{8}(\bar{\psi}\gamma_5\psi)^2 + \frac{5}{4}(\bar{\psi}\gamma_\mu\psi)(\bar{\psi}\gamma^\mu\psi) - \frac{29}{12}(\bar{\psi}\psi)^3 + \frac{7}{16}(\bar{\psi}\psi)^4$$



# Supervertices

- Mixed ST and FS rank-2 tensor

$${}_{\alpha}\lambda_{\beta}^{\mu} \equiv \frac{1}{2} \left( {}_{\alpha}\zeta_{\beta}^{\mu} - (-1)^{\alpha+\beta+\alpha\beta} {}_{\beta}\zeta_{\alpha}^{\mu} \right)$$




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- Covariant inverse superpropagator

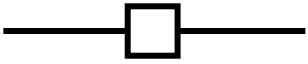

$$S_{;\hat{\alpha}\hat{\beta}} \Big|_{\partial_{\mu}\Phi=0} = (-1)^{\alpha} {}_{\alpha}\lambda_{\beta}^{\mu} p_{\mu}^{\beta} \delta(p^{\alpha} + p^{\beta})$$

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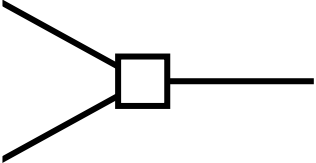
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- Covariant three-vertex



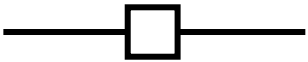
$$S_{;\hat{\alpha}\hat{\beta}\hat{\gamma}} \Big|_{\partial_{\mu}\Phi=0} = \left( (-1)^{\alpha} {}_{\alpha}\lambda_{\beta;\gamma}^{\mu} p_{\mu}^{\beta} + (-1)^{\alpha+\beta\gamma} {}_{\alpha}\lambda_{\gamma;\beta}^{\mu} p_{\mu}^{\gamma} \right) \delta(p^{\alpha} + p^{\beta} + p^{\gamma})$$

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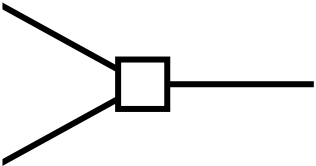
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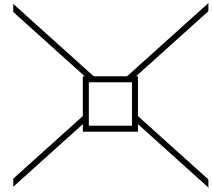


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**Note:**  $\neq 0$  unlike purely bosonic

# Supervertices (continued)

## □ Covariant four-vertex

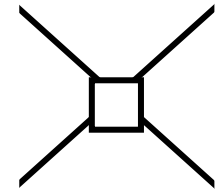


$$S_{;\hat{\alpha}\hat{\beta}\hat{\gamma}\hat{\delta}} \Big|_{\partial_\mu \Phi=0} = \left( (-1)^\alpha {}_\alpha \lambda_\rho^\mu R^\rho{}_{\beta\gamma\delta} p_\mu^\delta + (-1)^\alpha {}_\alpha \lambda_{\beta;\gamma\delta}^\mu p_\mu^\beta \right. \\ \left. + (-1)^{\alpha+\beta\gamma} {}_\alpha \lambda_{\gamma;\beta\delta}^\mu p_\mu^\gamma + (-1)^{\alpha+\delta(\beta+\gamma)} {}_\alpha \lambda_{\delta;\beta\gamma}^\mu p_\mu^\delta \right) \\ \delta(p^\alpha + p^\beta + p^\gamma + p^\delta)$$



# Supervertices (continued)

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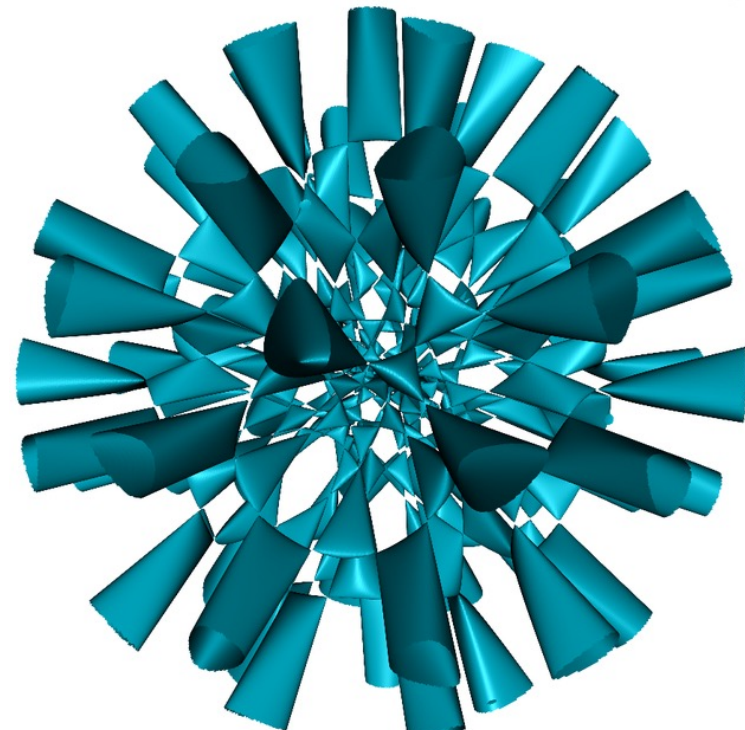
$$\begin{aligned} S_{;\hat{\alpha}\hat{\beta}\hat{\gamma}\hat{\delta}} \Big|_{\partial_\mu \Phi=0} = & \left( (-1)^\alpha {}_\alpha \lambda_\rho^\mu R^\rho{}_{\beta\gamma\delta} p_\mu^\delta + (-1)^\alpha {}_\alpha \lambda_{\beta;\gamma\delta}^\mu p_\mu^\beta \right. \\ & \left. + (-1)^{\alpha+\beta\gamma} {}_\alpha \lambda_{\gamma;\beta\delta}^\mu p_\mu^\gamma + (-1)^{\alpha+\delta(\beta+\gamma)} {}_\alpha \lambda_{\delta;\beta\gamma}^\mu p_\mu^\delta \right) \\ & \delta(p^\alpha + p^\beta + p^\gamma + p^\delta) \end{aligned}$$

- Scalar contribution is additive



# Summary and Outlook

- ❑ Fermionic curvature arises from non-linearity
- ❑ Unlike supergravity, curvature not real-valued
- ❑ Derived generalised expressions for covariant scalar-fermion vertices



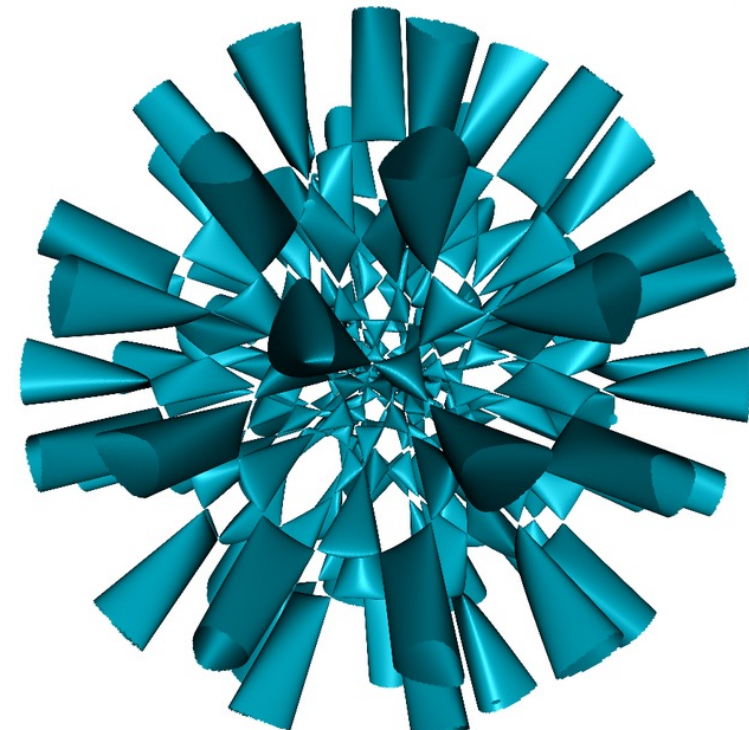


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## What next?

- ❑ Compute higher loop effective actions
- ❑ Add symmetries
- ❑ Compute amplitudes



# Thank you!



$$R_I = \frac{4g_1}{g_0^2} + \left( \frac{2g_1g_0'g_1'}{g_0^3k} - \frac{2g_1^2g_0'^2}{g_0^4k} - \frac{g_1'^2}{2g_0^2k} \right) (\bar{\psi}\psi)^2$$

$$R_{II} = -8 + 2(\bar{\psi}\psi) + \frac{23}{8}(\bar{\psi}\psi)^2 + \frac{9}{8}(\bar{\psi}\gamma_5\psi)^2 + \frac{5}{4}(\bar{\psi}\gamma_\mu\psi)(\bar{\psi}\gamma^\mu\psi) \\ - \frac{29}{12}(\bar{\psi}\psi)^3 + \frac{7}{16}(\bar{\psi}\psi)^4$$