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 CERN
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M(embrane) Theory

As derived 45* years ago,
 M-dimensional objects moving in
 such a way that their world volume
 swept out in $\mathbb{R}^{1,D-1}$ is extremal
 can be described by a light-cone
 (~~obscure~~ w.r.t. $\varphi^0 = \sqrt{\frac{x^0 + x^{D-1}}{2}} =: \tau$)

Hamiltonian $P_- = \frac{1}{2\alpha'} \int \frac{\vec{p}^2 + g}{\gamma} d^M \varphi = H[\vec{x}, \vec{p}; \gamma, \mathcal{J}_0]$

polynomially depending on $d := D-2$
 transverse coordinates fields $\vec{x} = \sum_{\alpha=0}^{\infty} \vec{x}^{(\alpha)} Y_{\alpha}(\varphi^1 - \varphi^M)$

canonically conjugate momenta $\vec{p} = \sum_{\beta} \int \vec{p}_{\beta}^{(\alpha)} Y_{\beta}$,
 $\int 1 (= \int \int d^M \varphi)$, $\int Y_{\alpha} Y_{\beta} \int d^M \varphi = \delta_{\alpha\beta}$,

$\{x_{i\alpha}, p_{j\beta}\} = \delta_{ij} \delta_{\alpha\beta}$, as well as two discrete
 degrees of freedom, $\gamma = P_+$ and its canonical
 conjugate $\mathcal{J}_0 = \int (x^0 - x^{D-1}) \gamma d^M \varphi$, $\{\gamma, \mathcal{J}_0\} = 1$,
 $g = \det \left(\partial_a \vec{x} \cdot \partial_b \vec{x} \right)_{a,b=1 \dots M}$.

* Sorry for, during my talk on Jan. 8,
 having thought that $2024 - 1979 = 55$

The second important step, for $M=2$ leading to the $SU(N)$ -invariant matrix model

$$H_N = \text{Tr} \left(\vec{P}^2 - \sum_{i,j} [X_i, X_j]^2 \right),$$

derived in my 1982 MIT Ph.D.-thesis, was to replace basis functions $Y_\alpha(\varphi', \varphi^{M=2})$

by $N \times N$ matrices $T_\alpha^{(N)}$. For spherical membranes,

$$T_{lm}^{(N)} = \sqrt{4\pi N} \sqrt{\frac{(N^2-1)l(N-1-l)!}{(N+l)!}} \sum_{a_1 \dots a_l} C_{a_1 \dots a_l}^{(m)} X_{a_1}^{(N)} \dots X_{a_l}^{(N)},$$

with the 3 hermitean matrices $X_a^{(N)}$ being an irreducible representation of $SU(2)$, normalized such that $X_1^2 + X_2^2 + X_3^2 = \mathbb{1}$,

$$[X_a^{(N)}, X_b^{(N)}] = i \epsilon_{abc} \frac{2}{\sqrt{N^2-1}} X_c^{(N)}, \quad \frac{1}{N} \text{Tr} T_{lm}^{(N)} T_{l'm'}^{(N)} = \delta_{ll'} \delta_{mm'}$$

for $l=1, 2, \dots, N-1$ ($m=-l \dots +l$; note that $T_{l,0}^{(N)} \equiv 0$ for $l \geq N, m$) are N^2-1 linear independent $N \times N$

matrices for the (complex, traceless, symmetric) tensors $C_{a_1 \dots a_l}^{(m)}$ appearing in a similar

expression for r^l times the ordinary spherical harmonics $Y_{lm}(\theta, \varphi)$ as harmonic

homogeneous polynomials in $X_1 = r \sin\theta \cos\varphi$, $X_2 = r \sin\theta \sin\varphi$, $X_3 = r \cos\theta$.

(cf. arXiv 0206192 (hep-th), 1101.4403
 CMP 195 (1998) 67 Nucl. Phys. B 849
 Int. J. Mod. Phys. A 4 (1989) #19, 5235 2011, 628) (2

For general M the (purely internal)
 $M^2 = 2P_+P_- - \vec{P}^2$ can be written as

$$\sum_{\substack{\alpha=1 \\ i=1 \dots d}}^{\infty} p_{i\alpha} p_{i\alpha} + \frac{1}{M!} g_{\alpha\alpha_1 \dots \alpha_M} g_{\beta_1 \dots \beta_M} x_{i\alpha_1} x_{i\beta_1} \dots x_{i\alpha_M} x_{i\beta_M}$$

$$g_{\alpha\alpha_1 \dots \alpha_M} = \int Y_{\alpha} \in^{r_1 \dots r_M} \partial_{r_1} Y_{\alpha_1} \dots \partial_{r_M} Y_{\alpha_M} d^M y$$

(hep-th/9602020, considering as well

$$[[X^{\mu_1}, \dots, X^{\mu_M}], X^{\mu_1}, \dots, X^{\mu_M}] = 0$$

Lorentz invariance of the classical theory
 was proven by Goldstone around 1980s
 (see the attached 7-page note)

Dynamical Symmetry and

Reconstruction Algebra: JH, PLB 695

resp. arXiv 1003.5189, 1004.0266, 1006.4714,
 1007.5505, 1101.4334

general review of relativistic membranes:

(^{with} unfortunately many typos) J. Phys. A 46 (2013)

kindly typed by the Journal
 from my handwritten notes)

*-product membranes: Phys. Lett. B 250
 (1990) 44 (3)

Lax-pair formulation
 (of general supersymmetrizable bosonic systems)
 arXiv 2101.01803, 2101.04435, 2101.11510
 (PoS CORFU 2021 (2022)258, Recent Progress on
 Membrane Theory)

Integrability in the dynamics of
 axially symmetric membranes /
 exact solutions* :

arXiv 2107.00569, 2201.02524, 2202.06955,
 2303.03920 (Asian J. Math 26 (2022) 253)
 2211.03887

* in 4 space-time dimensions PLB 822 (2021) 136658
 only very few are known explicitly

$$(t^2 + x^2 + y^2 - z^2)(t+z)^2 = C (\leq 0)$$

$$(t^2 - x^2 - y^2 - z^2) = c (t+z)^6 (\geq 0)$$

$$(t^2 + x^2 + y^2 - z^2) - 6C \sqrt{x^2 + y^2} (t+z)^2 + 3C^2 (t+z)^4 = 0,$$

the latter leading to a very interesting
 class of exact solutions of the classical
 matrix model equations,

$$\dot{X}_i = \sum_j [[X_i, X_j], X_j] , \sum_i [X_i, \dot{X}_i] = 0$$

The ^{fast} non-commutative sharp drop
 JH, Phys. Lett. B 2024