

Unified Matrix Theories from Nonrelativistic Strings

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Nordita

Matrix Quantum Mechanics for M-Theory Revisited

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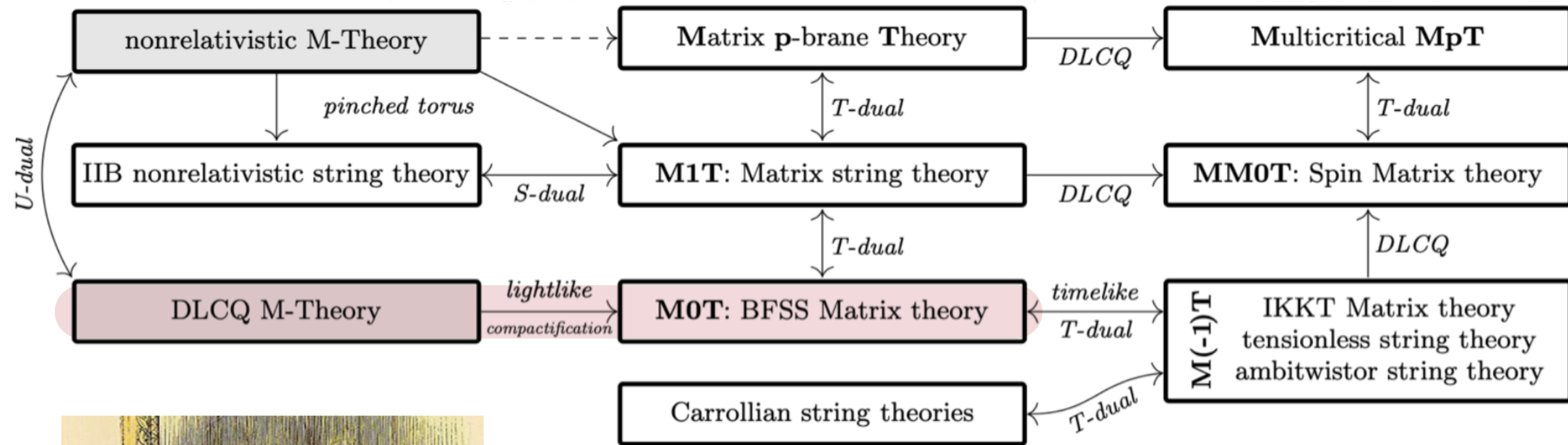


Unification of Decoupling Limits

Unification of Decoupling Limits in String & M-Theory

Chris Blair, Johannes Lahnsteiner, Niels Obers, ZY

2311.10564



[Carroll 1871]

Through the Looking-Glass of Exotic String Worldsheet

and What Alice Found There...

Worldsheet Formalism for Decoupling Limits in String Theory

Joaquim Gomis, ZY

2311.10565

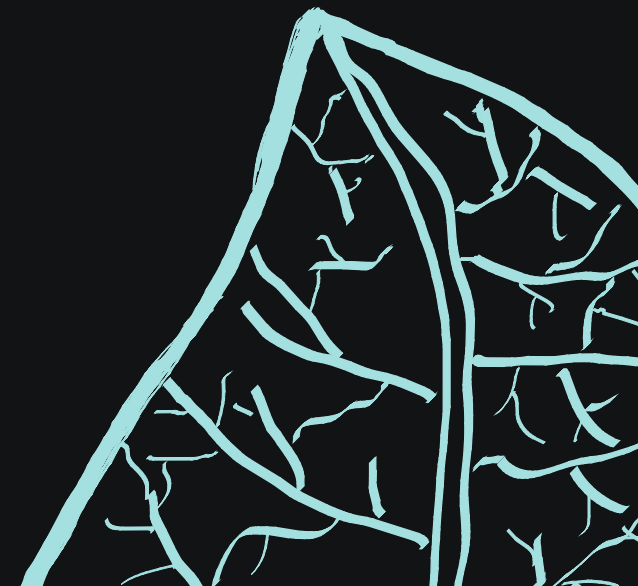
Outline

- 2311.10565 } 1. Fundamental strings for Matrix theory
- 2311.10564 } 2. Duality web from T-dualities
- 3. **REVIEW** of nonrelativistic string theory
- 2311.10564 4. Back to M-theory

see also my recent reviews

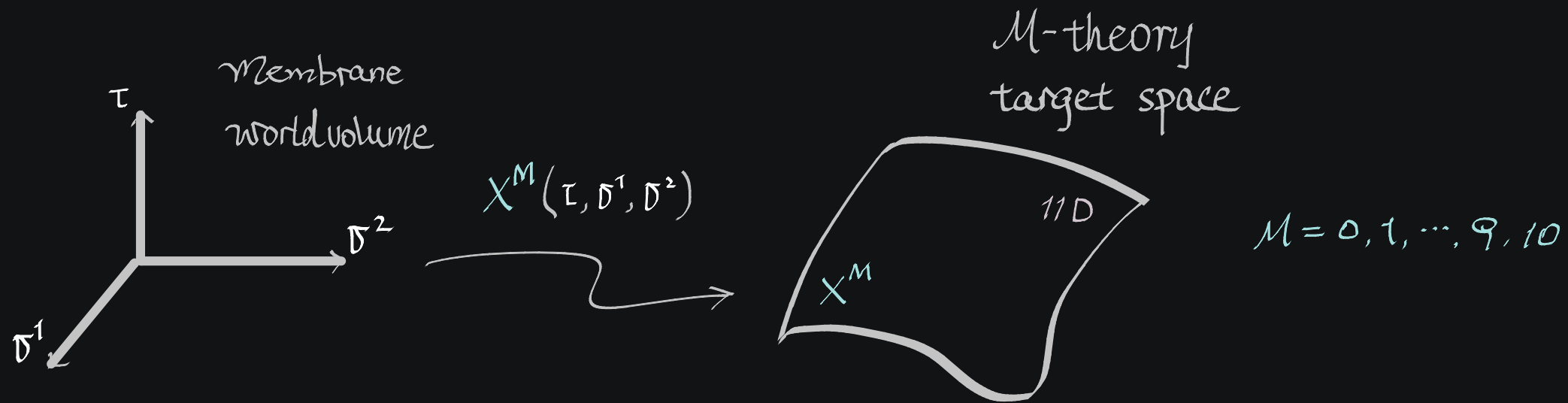
2312.12930 Section 7. Matrix theory & the string worldsheet [ZY](#)

2202.12698 Aspects of nonrelativistic strings [Oling, ZY](#)



Why is M-Theory Difficult?

- membrane sigma model: $S = -T_2 \int d^3\sigma \sqrt{-\det[\partial_\alpha X^M \partial_\beta X^N G_{MN}(X)]}$



- Polyakov formulation: $S = -\frac{1}{2} \int d^3\sigma (\partial_\alpha X^M \partial^\alpha X_M + \underbrace{g_{MNL}}_{\dim = -\frac{1}{2}} X^M \partial_\alpha X^N \partial^\alpha X^L + \dots)$
 not power-counting renormalizable!) RG flow

- high-energy completion? Quantum critical supermembrane: no worldvolume boost!

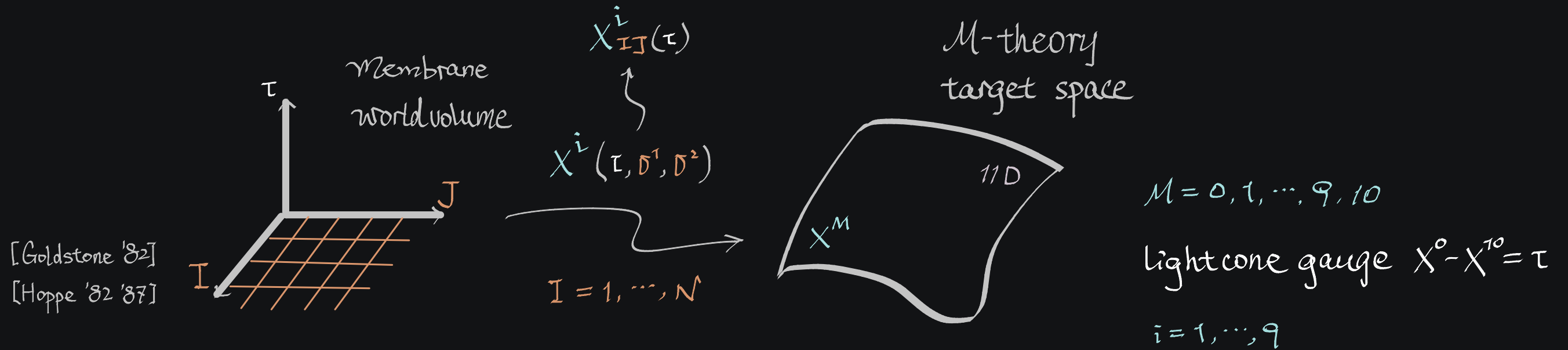
renormalizable! $S = \frac{1}{2} \int d\tau d^2\sigma d^2\theta (D_\tau Y^M D_\tau Y^N G_{MN} - \partial_i Y^M \partial_i Y^N H_{MN})$ [Hořava '08] [ZY '22]

[Hořava, Melby-Thompson, Randall]

Is M-theory fundamentally nonrelativistic?

Nonperturbative Quantization?

- membrane sigma model: $S = -T_2 \int d^3\sigma \sqrt{-\det(\partial_\alpha X^M \partial_\beta X_M)}$



- $SU(N)$ nonrel. quantum mechanics of 9 $N \times N$ matrices

[de Wit, Hoppe, Nicolai '88]

[Banks, Fischler, Shenker, Susskind '96]

[Susskind '97] [Seiberg '97] [Sen '97]

$$S \sim \frac{1}{2g^2} \int d\tau \text{tr} \left(\dot{X}^i \dot{X}^i + \frac{1}{4} [X^i, X^j]^2 + \text{fermions} \right)$$

nonrel. limit of a stack of N D0-particles in IIA $\xleftarrow[\text{compactify}]{\text{null}}$ DLCQ M-theory

Discrete Light Cone Quantization

Infinite Boost v.s. Nonrelativistic Limit

- DLCQ M-theory [Seiberg '97]

$$x^{10} \sim x^{10} + 2\pi R_0 \xrightarrow{\text{Lorentz boost}} \begin{cases} X^0 = \gamma (x^0 + v x^{10}) \\ X^1 = \gamma (x^{10} + v x^0) \end{cases}$$

$$\xrightarrow[\text{fix } R = 2\gamma R_0]{\gamma \rightarrow \infty} X^+ \sim X^+ + 2\pi R$$

Lorentz factor $\gamma = \frac{1}{\sqrt{1-v^2}}$

def. $\begin{cases} X^+ = X^0 + X^1 = 2\gamma x^+ + \dots \\ X^- = X^0 - v X^1 = \frac{1}{\gamma} x^0 + \dots \end{cases}$

- wrapped M2-brane around a compactified X^+ ? $X^+ = \sigma^2$ $\omega \propto \gamma$

String action $S = -T \int d^2\sigma \sqrt{-\det \begin{pmatrix} 1/\omega^2 & \partial_\beta X^- \\ \partial_\alpha X^- & \partial_\alpha X^i \partial_\beta X^i \end{pmatrix}} + \dots$ $\alpha, \beta = 0, 1$
 $i = 1, \dots, 9$

\downarrow
 $X^- = X^0$

$$S = -T \int d^2\sigma \sqrt{-\det \left(-\omega \partial_\alpha X^0 \partial_\beta X^0 + \frac{1}{\omega} \partial_\alpha X^i \partial_\beta X^i \right)}$$

Nambu-Goto string

$X^0 \rightarrow \sqrt{\omega} X^0$

$X^i \rightarrow \frac{1}{\sqrt{\omega}} X^i$

- infinite boost $\gamma \rightarrow \infty \Rightarrow$ nonrel. $\omega \rightarrow \infty$ limit of IIA effectively, infinite speed of light

What about D0-Branes?

- nonrelativistic limit as a BPS limit

$$S_{D0} = -\frac{1}{g_s \alpha'^{\frac{1}{2}}} \int d\tau \operatorname{tr} \left\{ \sqrt{-\dot{X}^\mu \dot{X}_\mu \det(1 - \frac{1}{4} [X^i, X^j]^2)} \right\} + \frac{1}{g_s \alpha'^{\frac{1}{2}}} \int C_1$$

$$\xrightarrow{\omega \rightarrow \infty} \frac{1}{2} \frac{1}{g_s \alpha'^{\frac{1}{2}}} \int d\tau \operatorname{tr} \left(\dot{X}^i \dot{X}^i + \frac{1}{4} [X^i, X^j]^2 \right)$$

$\mu = 0, \dots, 9$ static gauge

$$X^0 \rightarrow \sqrt{2\omega} X^0 \quad X^0 = \tau$$

$$X^i \rightarrow \frac{1}{\sqrt{2\omega}} X^i$$

$$g_s \rightarrow 2\omega^{-\frac{3}{2}} g_s$$

$$C_1 \rightarrow 2\omega^2 g_s^{-1} dX^0$$

- BFSS Matrix theory from a full-fledged string theory

DLCQ M-theory

compactify over
a lightlike circle

Matrix 0-brane theory

a decoupling limit of IIA

{
 D0-branes: BFSS Matrix theory
 D2-branes: noncommutative Yang-Mills
 fundamental string

see also

[Gopakumar, Minwalla, Seiberg, Strominger '00]

[Harmon '00] [Gomis, Ooguri '00]

[Danielsson, Guajosa, Kruszenski '00]

Polyakov String in Matrix 0-Brane Theory [Gomis, ZY '23]

◦ Nambu-Goto string

$$S = -T \int d^2\sigma \sqrt{-\det \begin{pmatrix} 0 & \partial_\beta X^\alpha \\ \partial_\alpha X^\beta & \partial_\alpha X^i \partial_\beta X^i \end{pmatrix}} \quad \text{see also [Bittle, Gomis, Not '16]}$$

◦ Polyakov formulation flat gauge

$$S_P = \frac{T}{2} \int d^2\sigma (\partial_\sigma X^\alpha \partial_\sigma X^\alpha + \partial_\tau X^i \partial_\tau X^i + \lambda \partial_\tau X^\alpha)$$

tropological sigma models
for Gromov-Witten invariants;
nonequilibrium string theory?
[Albrychiewicz, Ellers, Valiente, Hořava '23]

target space Galilei boost $\delta X^\alpha = 0$ $\delta X^i = v^i X^\alpha$ $\delta \lambda = -2v^i \partial_\tau X^i$

worldsheet Carroll boost $\delta \tau = v^i \sigma$ $\delta \sigma = 0$ $\delta \lambda = 2v^i \partial_\sigma X^\alpha$

worldsheet Galilean Conformal Algebra \approx BMS₃ [Bagchi '10]

Worldsheet Topology: Nodal Spheres

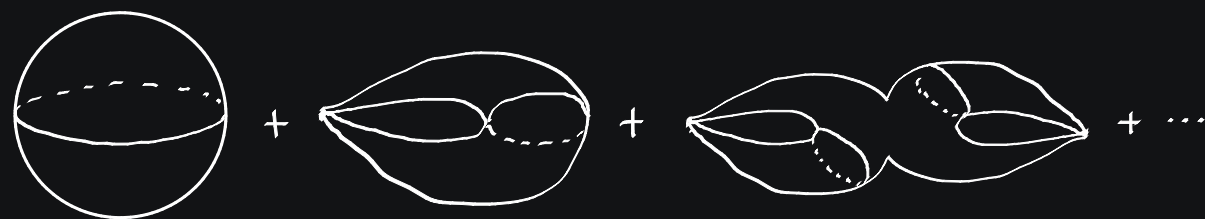
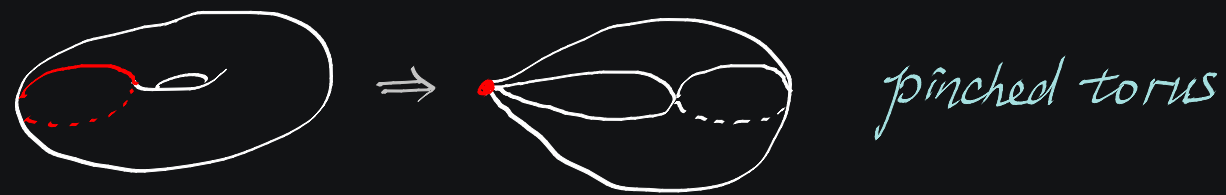
- M0T string from a nonrel. limit $\begin{cases} X^0 \rightarrow \sqrt{\omega} X^0 & \tau \rightarrow \omega^{-1} \tau \\ X^i \rightarrow \frac{1}{\sqrt{\omega}} X^i & \sigma \rightarrow \sigma \end{cases}$

$$S = -\frac{T}{2} \int d^2\sigma \partial_\alpha X^\mu \partial^\alpha X_\mu$$

$$= -\frac{T}{2} \int d^2\sigma \left[\underbrace{\omega^2 \partial_\tau X^0 \partial_\tau X^0}_{-\lambda \partial_\tau X^0 - \frac{\lambda^2}{4\omega^2}} - (\partial_\sigma X^0 \partial_\sigma X^0 + \partial_\tau X^i \partial_\tau X^i) + \mathcal{O}(\omega^{-2}) \right]$$

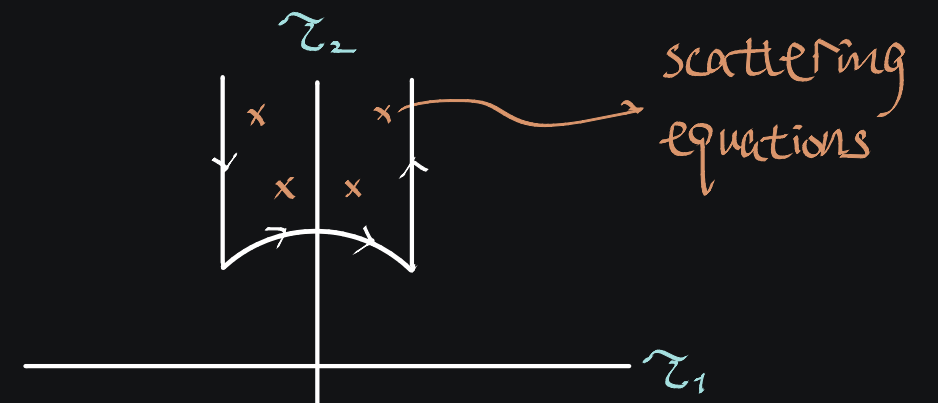
- Singular topology from $\omega \rightarrow \infty$

modulus $\tau \rightarrow i\infty$



ambitwistor string amplitudes

$\tau \rightarrow i\infty$



[Geyer, Mason, Monteiro, Tourkine '15]

Duality Web from T-Dualities

MOT string $i=1, \dots, 9$

$$S_{\text{MOT}} = \frac{T}{2} \int d^2\sigma \left(\partial_\sigma X^0 \partial_\sigma X^0 + \partial_\tau X^i \partial_\tau X^i + \lambda \partial_\tau X^0 \right)$$

T-dualize X^1

T-dualize X^0

MIT string $a=0,1$
 $i=2, \dots, 9$

$$S_{\text{MIT}} = \frac{T}{2} \int d^2\sigma \left(-\partial_\sigma X^a \partial_\sigma X_a + \partial_\tau X^i \partial_\tau X^i + \lambda_a \partial_\tau X^a \right)$$

D1-strings: Matrix string theory

[Motl '97] [Dijkgraaf, Verlinde, Verlinde '97]

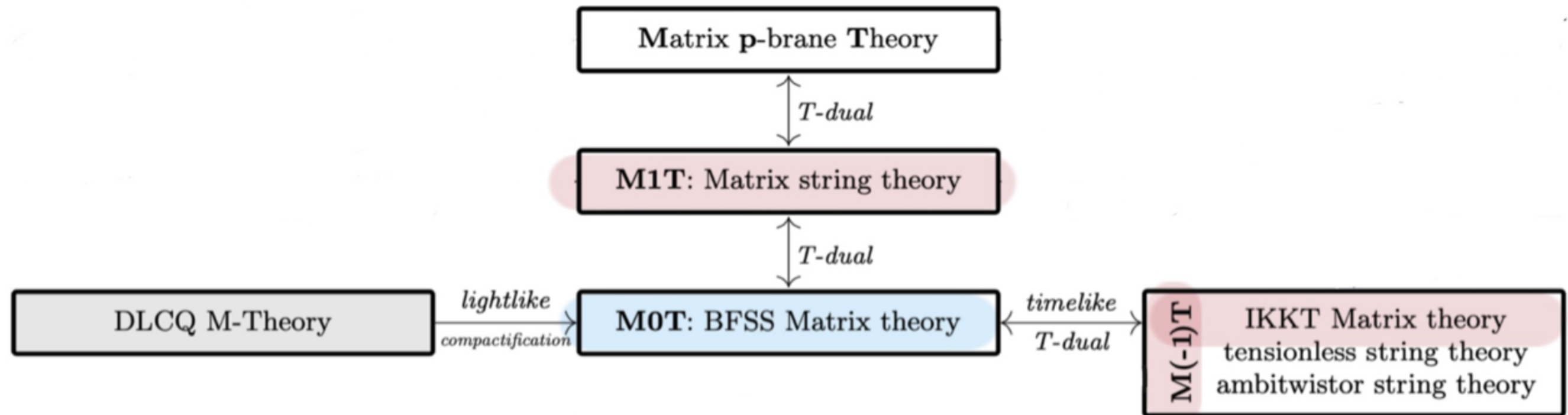
tensionless string [Lindström, Sundborg, Theodoridis '91]

$$S_{\text{M(1)T}} = \frac{T}{2} \int d^2\sigma \partial_\tau X^\mu \partial_\tau X_\mu \quad \mu=0, \dots, 9$$

D(-1)-instantons: IKKT Matrix theory

[Ishibashi, Kawai, Kitazawa, Tsuchiya '96]

$p=3$ $\mathcal{N}=4$ SYM (coupled to non-Lorentzian geometry)
 $AdS_5/CFT_4?$



Strings in Carroll-Like Spacetimes

$M(-1)T$ /tensionless string

$$S_{M(-1)T} = \frac{T}{2} \int d^2\sigma \partial_\tau X^\mu \partial_\tau X_\mu \quad \mu = 0, \dots, 9$$

D(-1)-instantons: IKKT Matrix theory

[Ishibashi, Kawai, Kitazawa, Tsuchiya '96]

T-dualize X^1

$M(-2)T$ string $a = 0, 2, \dots, 9$

$$S_{M(-2)T} = \frac{T}{2} \int d^2\sigma (\partial_\sigma X^1 \partial_\sigma X^1 + \partial_\tau X^a \partial_\tau X_a + \lambda_1 \partial_\tau X^1)$$

target-space Carroll-like boost

$$\delta X^a = v^a X^1 \quad \delta X^1 = 0 \quad \delta \lambda_1 = -2 v_a \partial_\tau X^a$$

light excitations:

1-dim. spacelike branes

[Hull '98]

[Gutperle, Strominger '02]

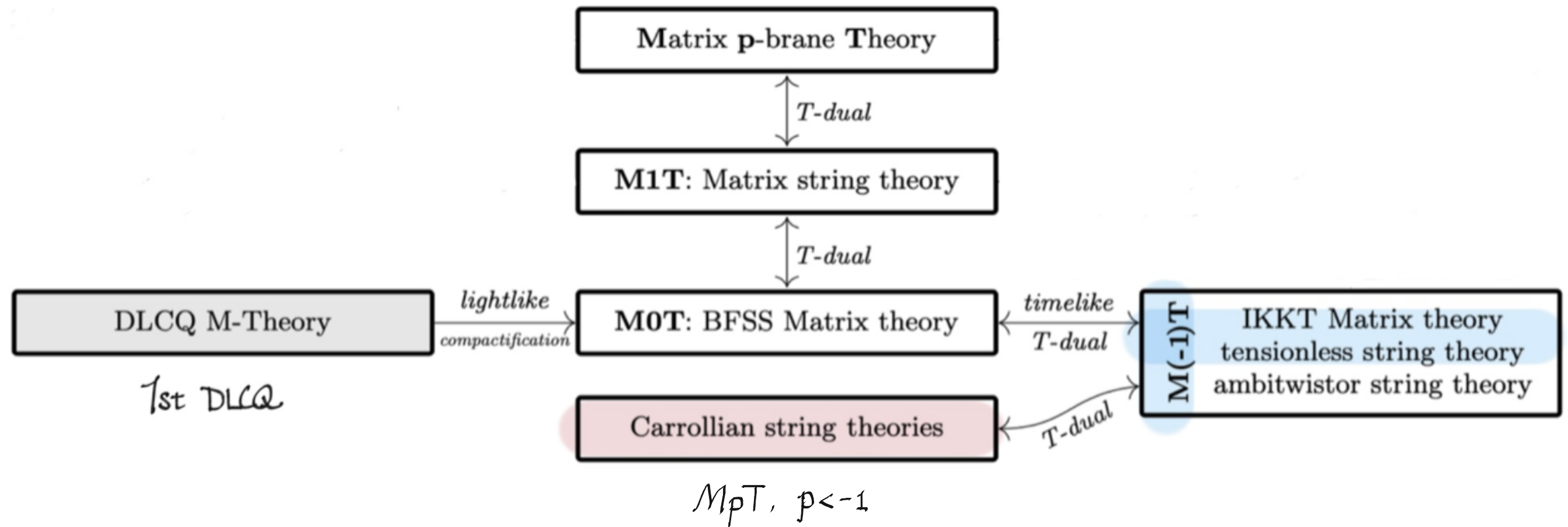
flat space

holography?



Carrollian/Celestial?

IKKT & instanton?



The 2nd DLCQ

DLCQ MIT string $\alpha=0,1$
 $i=2, \dots, 9$

$$S_{\text{MIT}} = \frac{T}{2} \int d^2\sigma \left(-\partial_\sigma X^\alpha \partial_\sigma X_\alpha + \partial_\tau X^i \partial_\tau X^i + \lambda_\alpha \partial_\tau X^\alpha \right)$$

DLCQ M(-1)T string

$$S_{\text{M(-1)T}} = \frac{T}{2} \int d^2\sigma \partial_\tau X^\mu \partial_\tau X_\mu \quad \mu=0, \dots, 9$$

T-dualize X^+

T-dualize X^1

Multicritical MOT string

$$S_{\text{MOT}} = \frac{T}{2} \int d^2\sigma \left[\partial_\tau X^i \partial_\tau X^i + \lambda_0 \partial_\tau X^0 + \lambda_1 (\partial_\tau X^1 - \partial_\sigma X^0) \right]$$

F1-D0 states ($\frac{1}{4}$ BPS)

T-dualize X^0

Spin Matrix Limits of AdS/CFT

- Spin Matrix theory: near BPS limits of $\mathcal{N}=4$ SYM on $\mathbb{R} \times S^3$ [Harmark, Orselli '06]

$$\lambda_t \rightarrow 0 \quad E - Q \rightarrow 0 \quad \frac{E - Q}{\lambda_t} \rightarrow \text{fixed}$$

J_1, J_2 : angular momenta on S^3

S_1, S_2, S_3 : R-charges

- integrable

- finite N generalizations of spin chains

- bulk perspective: from $AdS_5 \times S^5$ to non-Lorentzian geometry

[Harmark, Hartong, Obers '17]

[Harmark, Hartong, Mencioli, Obers, ZY '18]

[Harmark, Hartong, Obers, Oling '21]

$$ds^2 = R^2 (-\cosh^2 \rho dt^2 + d\rho^2 + \sinh^2 \rho d\Omega_3^2 + d\Omega_5^2)$$

$$\lambda_t = \frac{R^4}{l_s^4}$$

$$E = i\partial_t$$

$$\underbrace{J_1, J_2, S_1, S_2, S_3}_Q$$

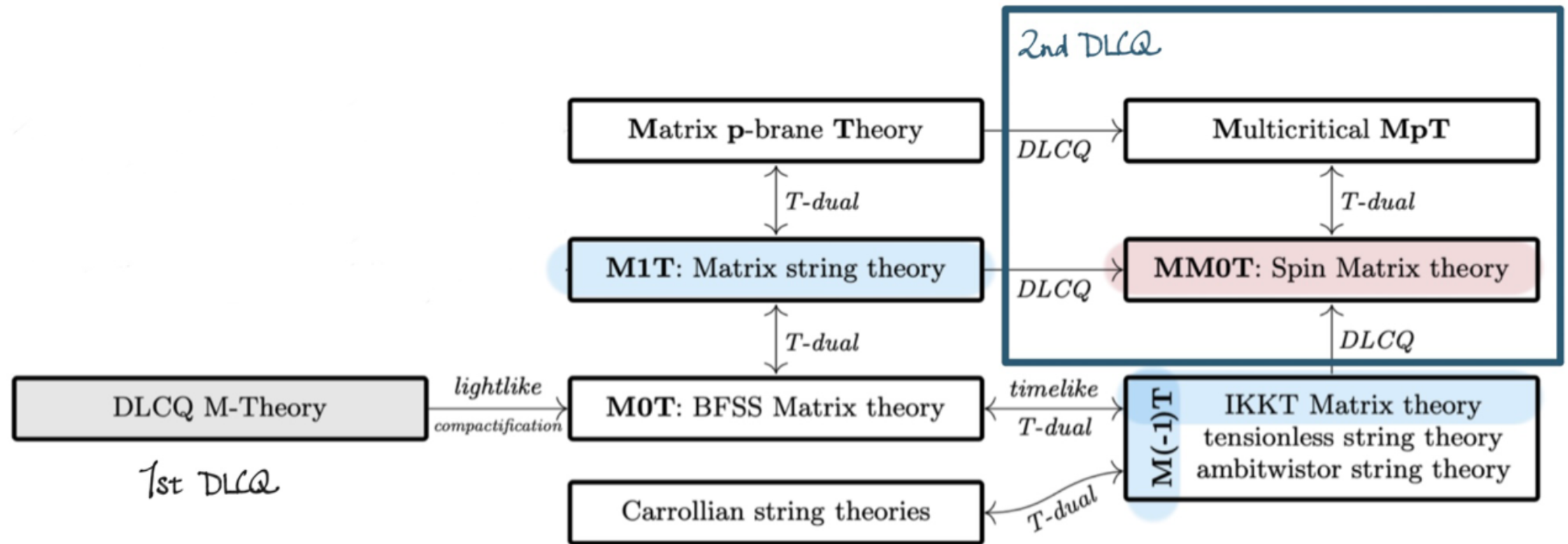
$$PSU(1,2|3) : Q = J_1 + J_2 + S_1 + S_2 + S_3$$

Spin Matrix limit same worldsheet theory as in MMOT

$$S_{MMOT} = \frac{T}{2} \int d^2\sigma (\partial_\tau X^\mu \partial_\tau X^\nu E_{\mu\nu} + \lambda_a \partial_\tau X^\mu \tau_\mu^a - \lambda_1 \partial_\sigma X^\mu \tau_\mu^0)$$

$$-T \int_B$$

3rd DLCQ: Kaluza-Klein monopole?



CHY formalism of Galilean QFT in MMOT?

Decoupling Limits of Type II Superstring Theory

Matrix p -brane Theory $\xrightarrow{\text{lightlike T-dual}}$ Multicritical Matrix $(p-1)$ -brane Theory

$$X^a \rightarrow \sqrt{\omega} X^a \quad a=0, \dots, p$$

$$X^i \rightarrow \frac{1}{\sqrt{\omega}} X^i \quad i=p+1, \dots, 9$$

$$g_s \rightarrow \omega^{\frac{p-3}{2}} g_s$$

$$C_{p+1} \rightarrow \omega^2 g_s^{-1} dx^0 \wedge \dots \wedge dx^p$$

critical background configurations

	0	1	...	p	p+1	...	9
D_p	x	x	...	x			

light excitations: $\frac{1}{2}$ BPS

$$X^0 \rightarrow \omega X^0 \quad X^1 \rightarrow X^1$$

$$X^a \rightarrow \frac{1}{\sqrt{\omega}} X^a \quad a=2, \dots, p$$

$$X^i \rightarrow \frac{1}{\sqrt{\omega}} X^i \quad i=p+1, \dots, 9$$

$$g_s \rightarrow \omega^{\frac{p-3}{2}} g_s$$

$$C_p \rightarrow \omega^2 g_s^{-1} dx^0 \wedge dx^2 \wedge \dots \wedge dx^p$$

$$B \rightarrow -\omega dx^0 \wedge dx^1$$

	0	1	2	...	p	p+1	...	9
$D^{(p-1)}$	x		x	...	x			
F1	x	x						

light excitations: $\frac{1}{4}$ BPS

U-duality & BPS states

[Dijkgraaf, de Boer, Harmark, Obers]

Non-Lorentzian Geometries

Matrix p -brane Theory $\xrightarrow{\text{lightlike T-dual}}$

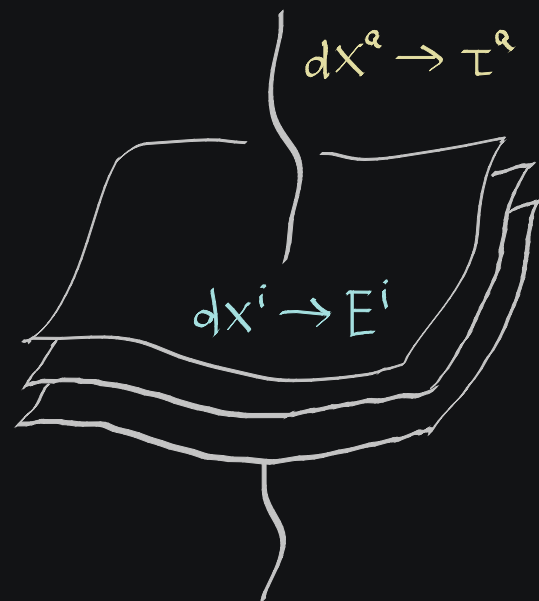
$$G_{\mu\nu} = \omega \tau_{\mu\nu} + \frac{1}{\omega} E_{\mu\nu}$$

$$e^{\Phi} = \omega^{\frac{p-3}{2}} e^{\phi}$$

$$C_{p+1} = \omega^2 e^{-\Phi} \tau^0 \wedge \dots \wedge \tau^p + C_{p+1}$$

$$\tau_{\mu\nu} = \tau_{\mu}^a \tau_{\nu}^b \eta_{ab}$$

$$E_{\mu\nu} = E_{\mu}^i E_{\nu}^j$$



Multicritical Matrix $(p-1)$ -brane Theory

$$G_{\mu\nu} = -\omega^2 \tau_{\mu}^0 \tau_{\nu}^0 + \tau_{\mu}^1 \tau_{\nu}^1 + \omega \tau_{\mu}^a \tau_{\nu}^a + \frac{1}{\omega} E_{\mu\nu}$$

$$e^{\Phi} = \omega^{\frac{p-3}{2}} e^{\phi} \quad a=2, \dots, p$$

$$C_p = \omega^2 e^{-\Phi} \tau^0 \wedge \tau^2 \wedge \dots \wedge \tau^p + \omega \tau^0 \wedge \tau^1 \wedge C_{p-2} + C_p$$

$$B = -\omega \tau^0 \wedge \tau^1 + b \quad C_q = \omega \tau^0 \wedge \tau^1 \wedge C_{q-2} + C_q$$

generalized Newton-Cartan/Carroll geometry

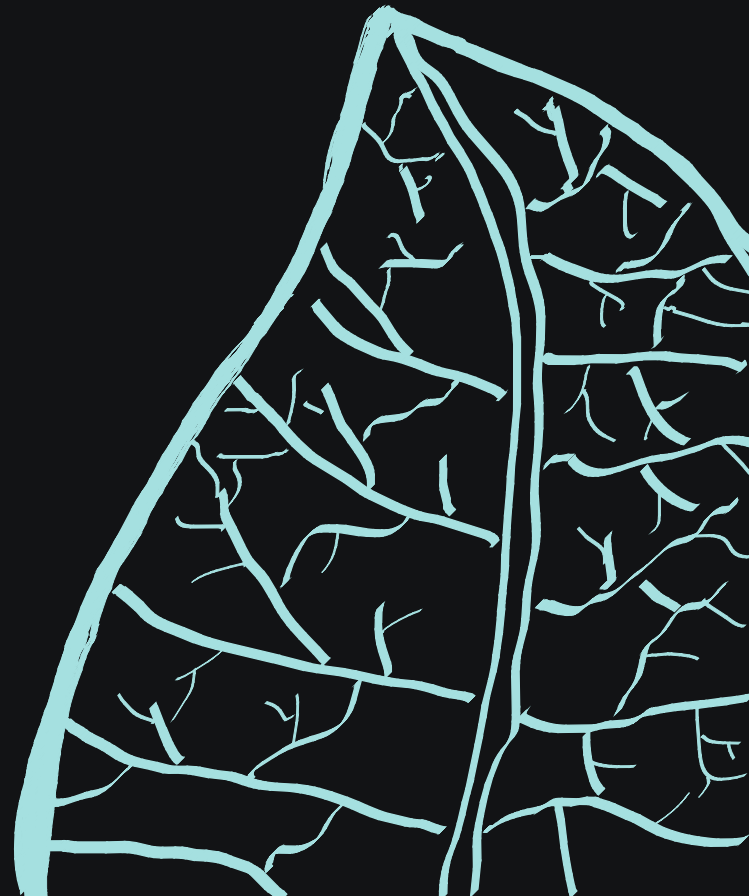
w/ torsional constraints $d\tau \sim 0$

due to quantum consistency/supersymmetry

implications for Matrix theory/SUGRA correspondence?

REULEWS

S-Duality & Nonrelativistic String Theory



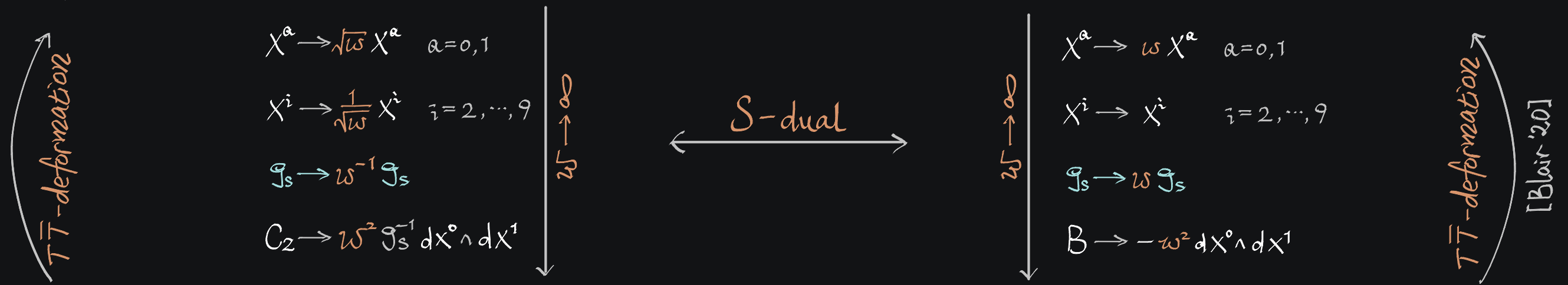
S-Dual of Matrix L-Brane Theory

D1-string $F=dA$

F1-string

$$S_{D1} = -\frac{1}{g_s \alpha'} \int d^2\sigma \sqrt{-\det(\partial_\alpha X^\mu \partial_\beta X_\mu + F_{\alpha\beta})} + \frac{1}{g_s \alpha'} \int C_2$$

$$S = -T \int d^2\sigma \sqrt{-\det(\partial_\alpha X^\mu \partial_\beta X_\mu)} - T \int B$$



$$S_{D1}^{MIT} = -\frac{1}{2g_s \alpha'} \int d^2\sigma \sqrt{-\tau} \left(\tau^{\alpha\beta} \partial_\alpha X^i \partial_\beta X^i - \frac{1}{2} \tau^{\alpha\gamma} \tau^{\beta\delta} F_{\alpha\beta} F_{\gamma\delta} \right)$$

$$S_{NR} = -\frac{T}{2} \int d^2\sigma \sqrt{-\tau} \tau^{\alpha\beta} \partial_\alpha X^i \partial_\beta X^i$$

$\tau_{\alpha\beta} = \partial_\alpha X^\alpha \partial_\beta X_\alpha$ Matrix string theory in MIT

nonrelativistic string theory

[Dijkgraaf, Verlinde, Verlinde '97]
 [Ebert, Sun, ZY '21]
 [Blair, Lahnsteiner, Obers, ZY '23]

[Klebanov, Maldacena '00]
 [Gomis, Ooguri '00]
 [Danielsson, Guijosa, Kruczenski '00]
 [Andringa, Bergshoeff, Gomis, de Roo '12]

REVIEW

S-Dual of Matrix 1-Brane Theory

- Branched $SL(2, \mathbb{Z})$ & polynomial realization

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \quad ad - bc = 1$$

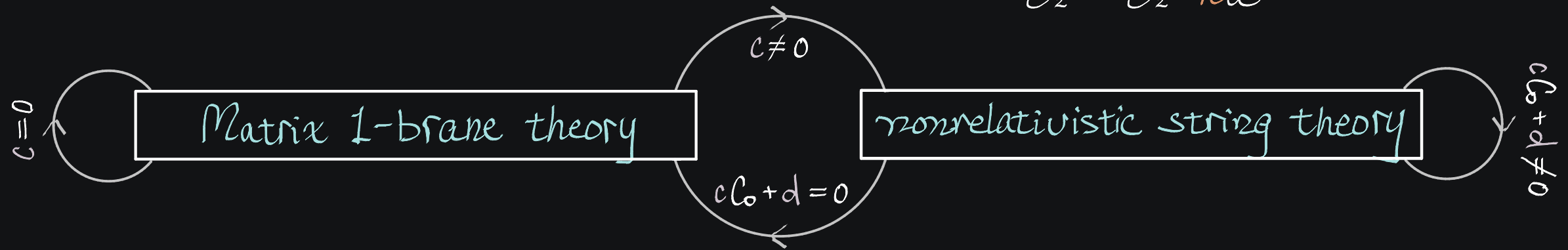
$$\mathcal{B} = e^{-\frac{\Phi}{2}} \mathcal{B}$$

$$\mathcal{L}_2 = e^{\frac{\Phi}{2}} (c_2 + c_0 \mathcal{B})$$

$$\mathcal{B} \rightarrow \mathcal{B} - \kappa \mathcal{L}_2 + \frac{1}{2} \kappa^2 \tau^0 \wedge \tau^1$$

$$\mathcal{L}_2 \rightarrow \mathcal{L}_2 - \kappa \mathcal{B}$$

$$\kappa = \frac{c e^{-\Phi}}{c_0 + d}$$



non-Lorentzian IIB SUGRA:

[Bergshoeff, Grosvenor, Lahnsteiner, ZY, Zorba '22 '23]

classical invariant theory & non-Lorentzian bootstrap?

[Ebert, ZY '23]

[Blair, Lahnsteiner, Obers, ZY to appear]

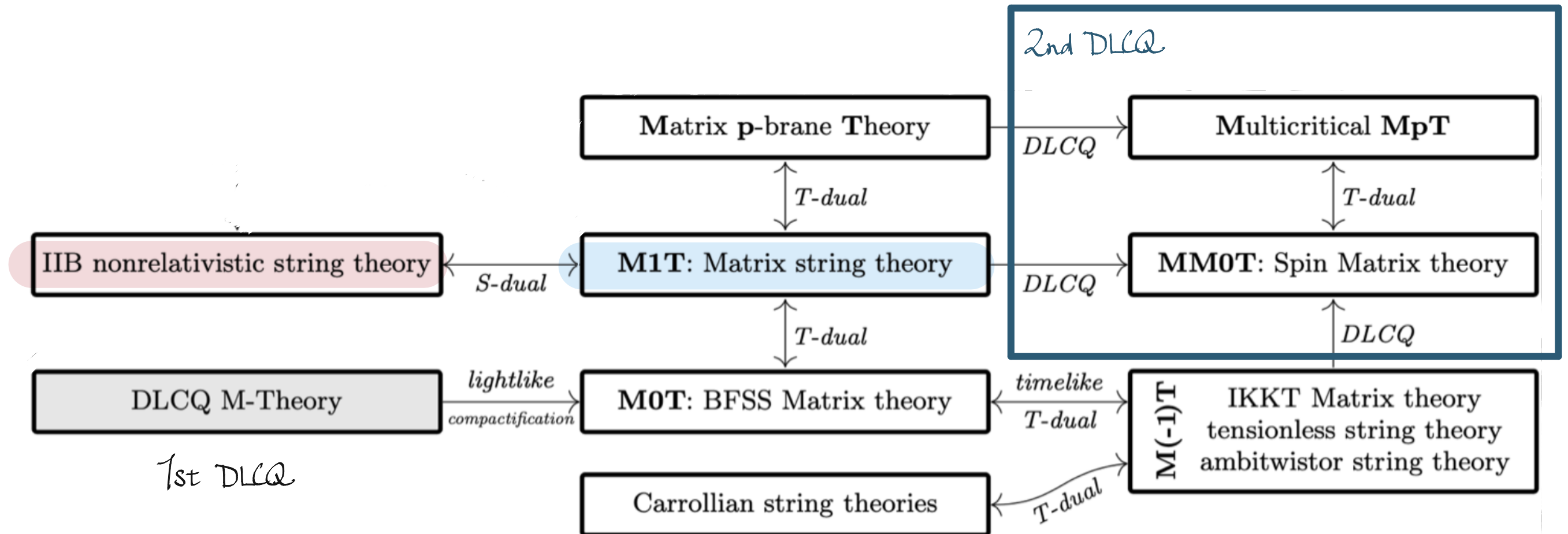
see also nonrel. string expansion [Hawe, Hartong '21]

- D3-branes

noncommutative Yang-Mills $\xleftrightarrow{\text{S-dual}}$ noncommutative open string

[Gopakumar, Maldacena, Minwalla, Strominger '00]

AdS/CFT: e.g. [Hashimoto, Itzhaki '99] [Maldacena, Russo '00] [Guijosa, Rosas-López '23]...



REVIEW

Nonrelativistic String Theory as a Pivot Theory

Nambu-Goto: $S_{NG} = -\frac{T}{2} \int d^2\sigma \sqrt{-\tau} \tau^{\alpha\beta} \partial_\alpha X^i \partial_\beta X^i$

$\tau_{\alpha\beta} = \partial_\alpha X^a \partial_\beta X_a \quad i=2, \dots, 9$

noncritical target space dim.

$S \sim \int d^2\sigma (\partial\Phi \bar{\partial}\Phi + \lambda \bar{\partial}X^+ + \bar{\lambda} \partial X^- - e^{-\Phi} \lambda \bar{\lambda} - \frac{1}{2} R \Phi)$

AdS₃/CFT₂ [Eberhardt, Gaberdiel, Gopakumar '19]

chiral 2D Yang-Mills [Komatsu, Maity]

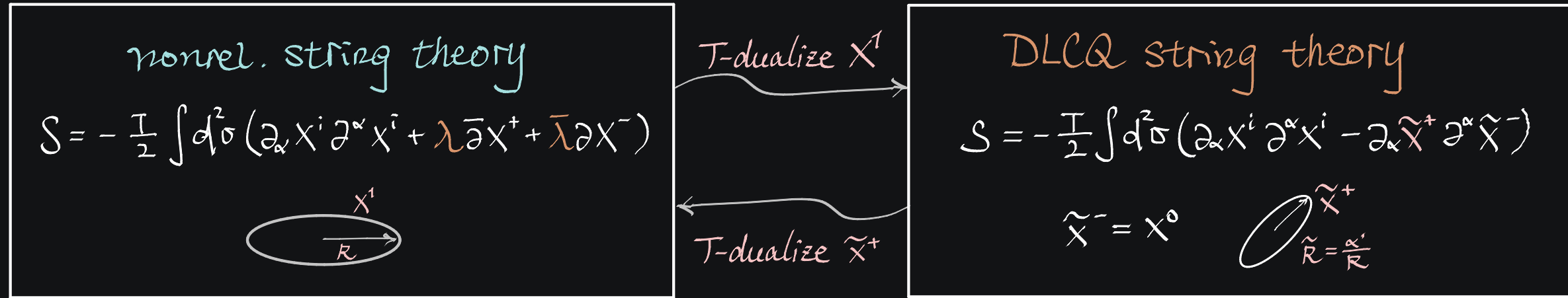
Polyakov: $S_p = -\frac{T}{2} \int d^2\sigma (\partial_\alpha X^i \partial^\alpha X^i + \underbrace{\lambda \bar{\partial}X^+ + \bar{\lambda} \partial X^-}_{\beta\delta \text{ system}})$ [Gomis, Ooguri '00]



1. string spectrum: $\mathcal{E} = \frac{\alpha'}{2\pi R} [k_i k^i + \frac{2}{\alpha'} (N + \bar{N} - 2)]$ Galilei invariant
2. target space string Galilei boost: $\delta X^a = 0, \delta X^i = \Lambda^i_a X^a$
3. relativistic worldsheet! standard CFT techniques
4. unitary and UV-complete nonrel. quantum gravity

REVIEW

First Principles Definition of DLCQ String Theory



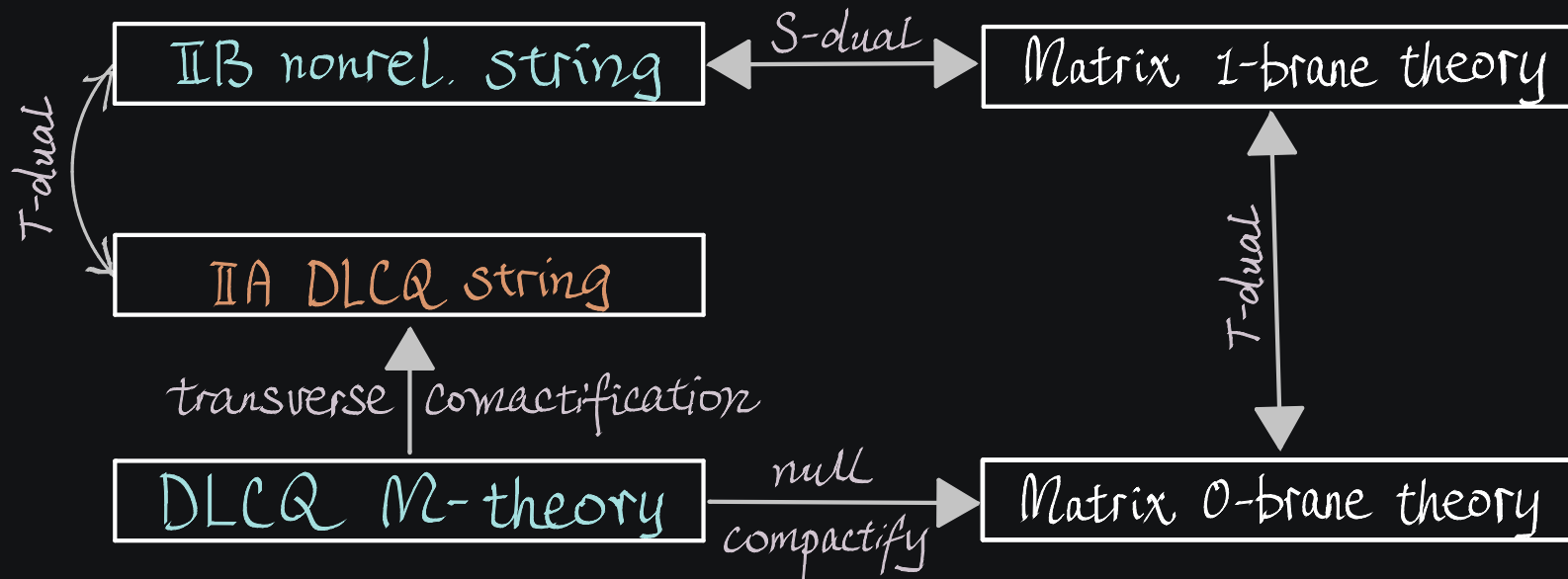
[Bergshoeff, Gomis, ZY '18]

null-reduction perspective

[Harmark, Hartong, Obers '17] [Kluson '18]

[Harmark, Hartong, Mencioli, Obers, ZY '18]

[Harmark, Hartong, Mencioli, Obers, Oling '19]

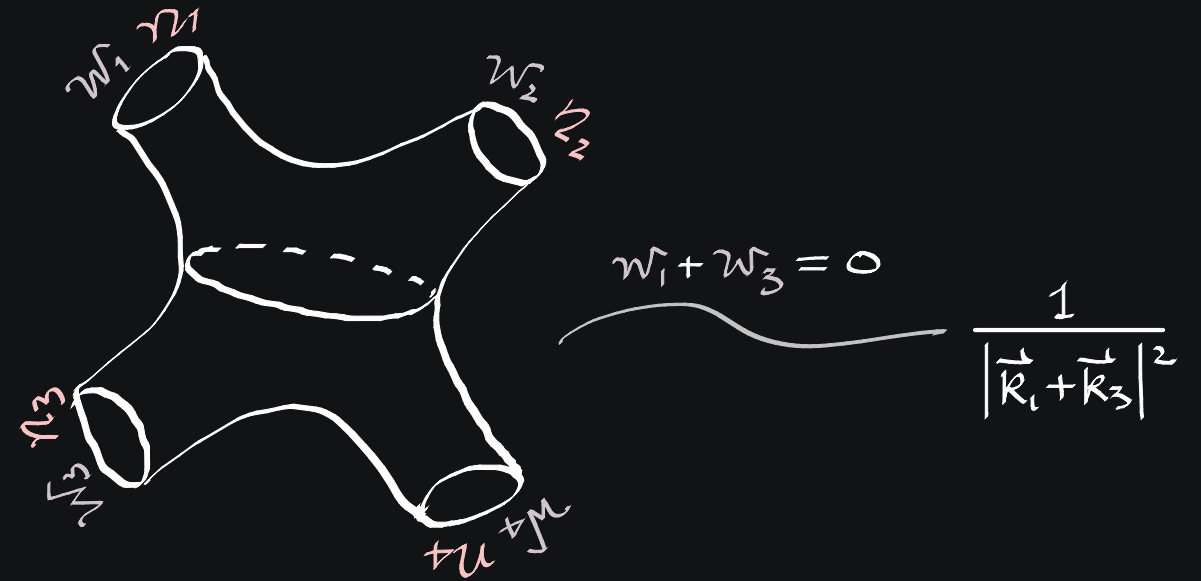


REVIEW

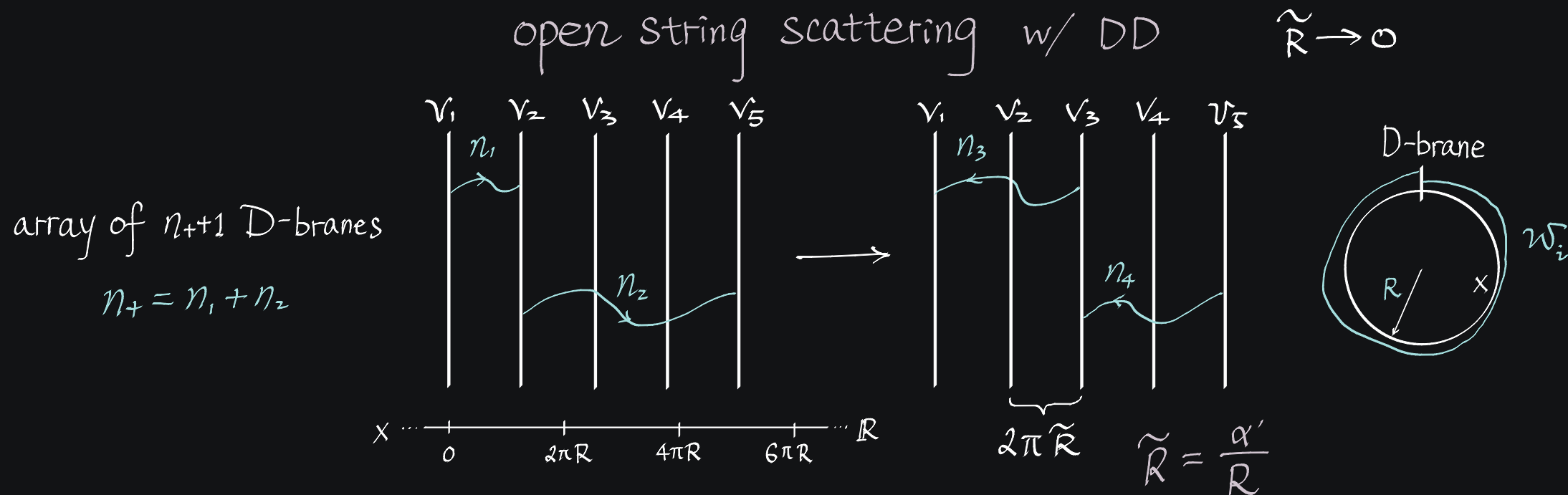
String Amplitudes & Winding KLTs

- no massless graviton
- instantaneous gravitational interaction between winding strings
- localized moduli space

[Gomis, Ooguri '00] [Danielsson, Guijosa, Kruzenski '00'00] [Bilal '98]



- KLTs with D-branes [Gomis, ZY, Yu '21] [ZY, Yu '21]



REVIEW

Interacting Sigma Models

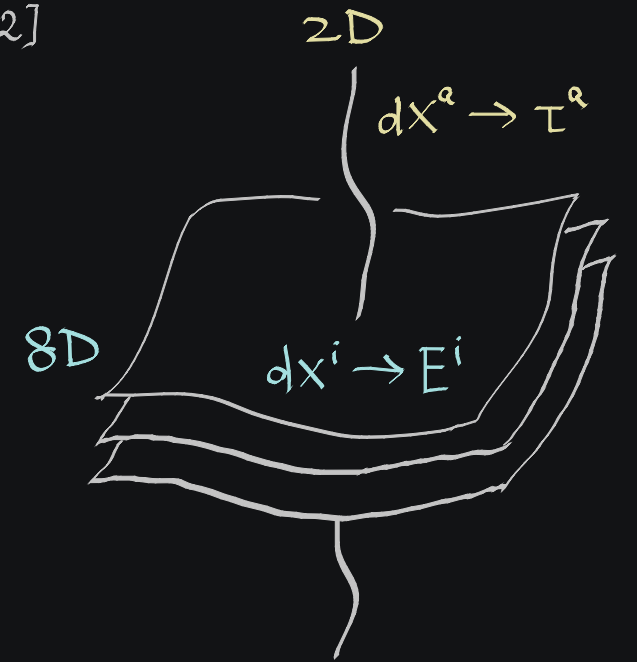
- Strings in string Newton-Cartan geometry [Andringa, Bergshoeff, Gomis, de Roo '12]

$$S = -\frac{T}{2} \int d^2\sigma \left[\partial X^\mu \bar{\partial} X^\nu (E_{\mu\nu} + b_{\mu\nu}) + \lambda \bar{\partial} X^\mu \tau_\mu^+ + \bar{\lambda} \partial X^\mu \tau_\mu^- \right]$$

β -functions [Gomis, Oh, ZY '19] [Gallegos, Gürsoy, Zinnato '19]
 [Bergshoeff, Gomis, Rosseel, Şimşek, ZY '19] [ZY, Yu '19] [ZY '21]

($T\bar{T}$ deformation [Blair '20])

- torsional deformation: $\lambda \bar{\lambda}$ is generated quantum mechanically
 \rightarrow back to relativistic string theory



- noncentral extensions of string Galilei algebra \Rightarrow torsional constraints on τ^a

trans. momentum
 \uparrow
 $[G_{ai}, P_j] = \delta_{ij} Z_a$
 \downarrow string Galilei boost \Downarrow $D\tau^a = 0$
 $\mathcal{N} = 2$ SUSY? (type II)

trans. momentum
 \uparrow
 $[G_{+i}, P_j] = \delta_{ij} Z_+$
 \downarrow string Galilei boost \Downarrow $D\tau^- = 0$
 $\mathcal{N} = 1$ SUSY (heterotic, type I)

no $\lambda \bar{\lambda}$ generated at all loops!
 [ZY '21]

similar constraints from SUSY
 [Bergshoeff, Lahnsteiner, Romano, Rosseel, Şimşek '21]
 [Bergshoeff, Blair, Lahnsteiner, Rosseel]

REVIEW

Open Strings and D-Branes

see also [Danielsson, Guijosa, Kruzcenski '00]

- open string vertex operators [Gomis, ZY, Yu '20]

$$S_{\text{boundary}} = \frac{1}{2\pi\alpha'} \int d\tau [\nu(\lambda - \bar{\lambda}) + iA_i \partial_\tau X^i]$$

$$G_{\mu\nu} = \omega^2 \tau_{\mu\nu} + E_{\mu\nu}$$

$$B = -\omega^2 \tau^0 \wedge \tau^1 + b$$

$$\bar{\Phi} = \phi + \ln \omega$$

$$C_q = \omega^2 \tau^0 \wedge \tau^1 \wedge C_{q-2} + C_q$$

- D-brane actions from β -functions [Gomis, ZY, Yu '20] [Ebert, Sun, ZY '21]

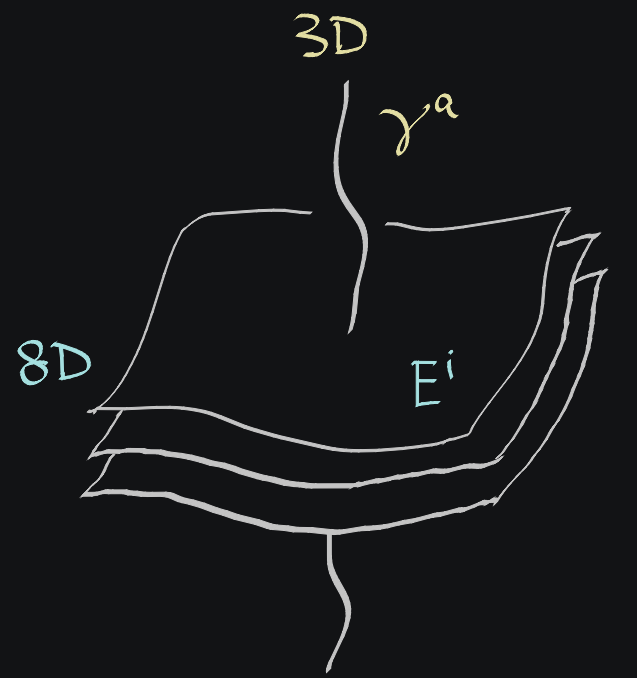
$$S_{Dp} = - \int d^{p+1}\sigma e^{-\phi} \sqrt{-\det \begin{pmatrix} 0 & \tau_\beta^+ \\ \tau_\alpha^- & E_{\alpha\beta} + b_{\alpha\beta} + F_{\alpha\beta} \end{pmatrix}} - \int \sum_q C_q e^{b+F}|_{p+1}$$

- nonrelativistic M-theory

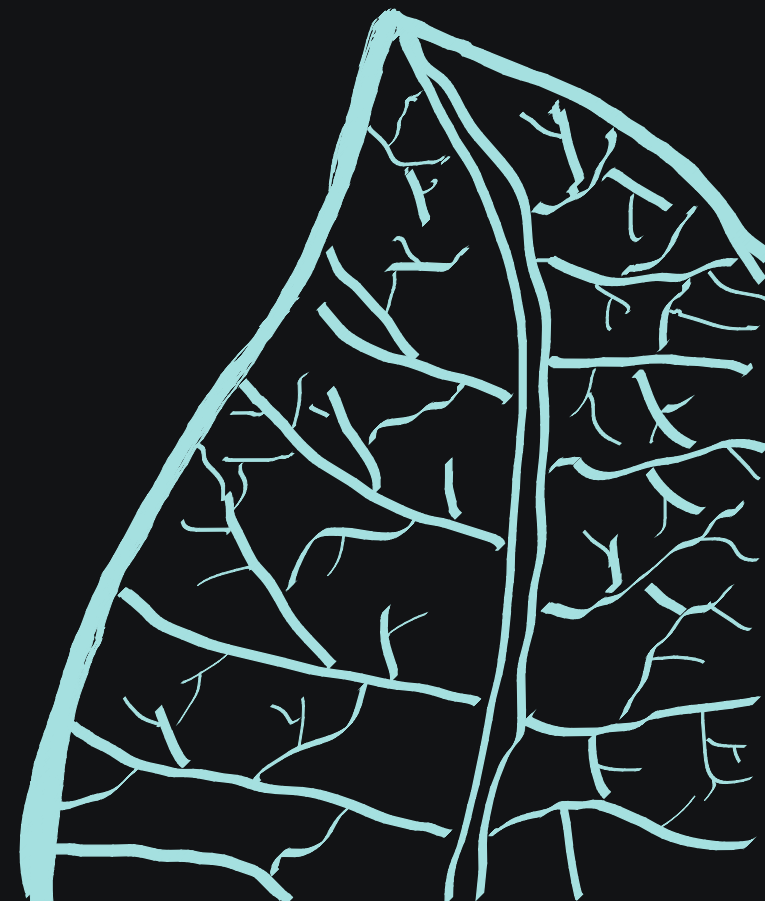
[Gomis, Ooguri '00] [Gopakumar, Minwalla, Seiberg, Strominger '00]
 [Bergshoeff, Berman, van der Schaar, Sundell '00] [Ebert, ZY '23]
 [Garcia, Guijosa, Vergara '02] [Blair, Gallegos, Zinnato '21]

electromagnetic dual of D2-brane [Ebert, Sun, ZY '21]

$$S_{M2} = - \int d^3\sigma \sqrt{-\gamma} \gamma^{\alpha\beta} \partial_\alpha X^M \partial_\beta X^N E_{MN} - \int A_3$$



Back to M-Theory



Nonrelativistic M-Theory

- back to DLCQ M-theory

$$X^+ \rightarrow \omega X^+$$

$$X^- \rightarrow \frac{1}{\omega} X^-$$

infinite boost: $\omega \rightarrow \infty$

- Compactify over T^3

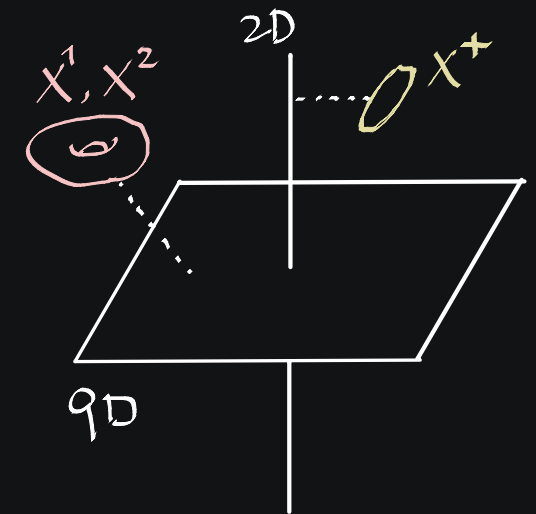
	X^+	X^1	X^2
radius	R_+	R_1	R_2

- U-dual frame $\xrightarrow{\omega \rightarrow \infty}$ nonrelativistic M-theory

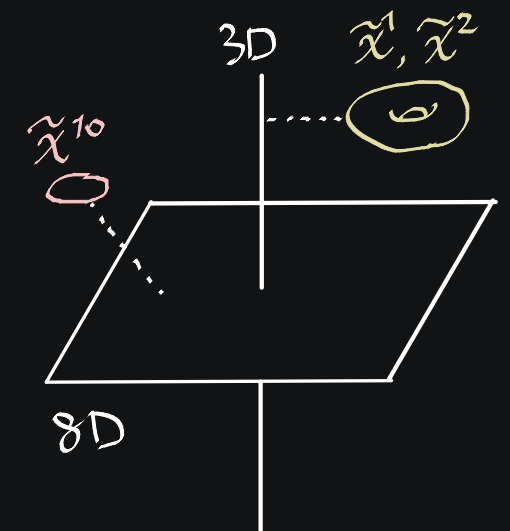
$$ds^2 = \omega^{\frac{4}{3}} \left[-(dx^0)^2 + (d\tilde{x}^1)^2 + (d\tilde{x}^2)^2 \right] + \omega^{-\frac{2}{3}} \left[\sum_{i=3}^9 dx^i dx^i + (d\tilde{x}^{10})^2 \right]$$

$$A_3 = -\omega^2 dx^0 \wedge d\tilde{x}^1 \wedge d\tilde{x}^2$$

DLCQ M-theory



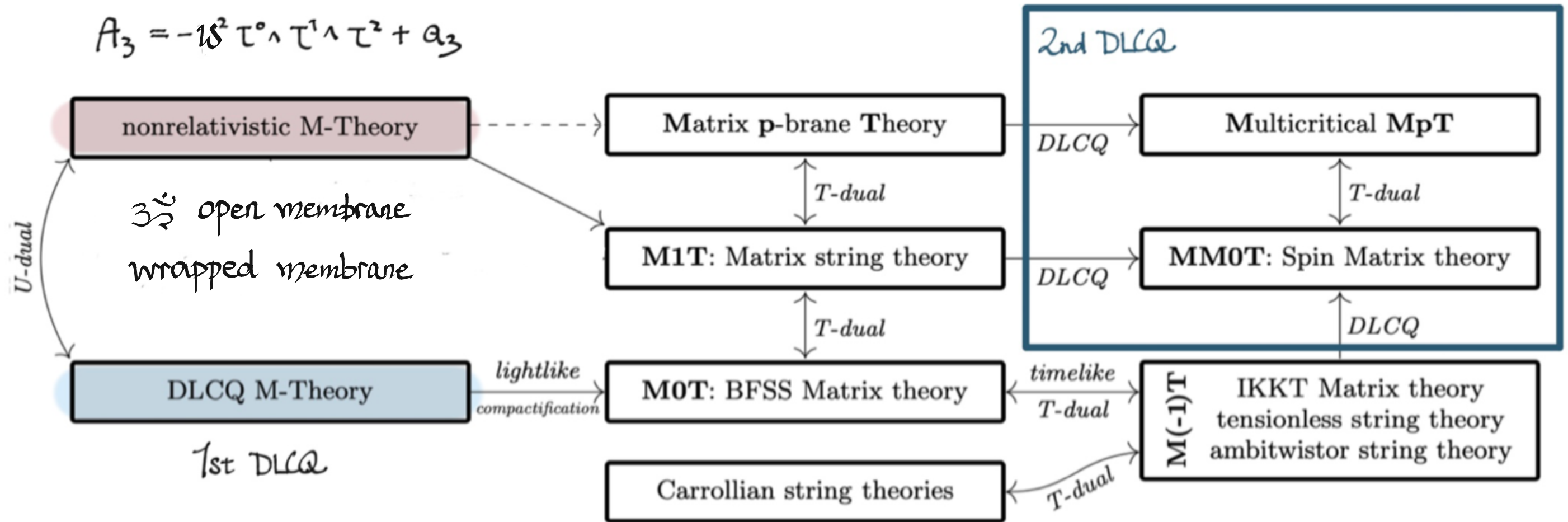
nonrelativistic M-theory



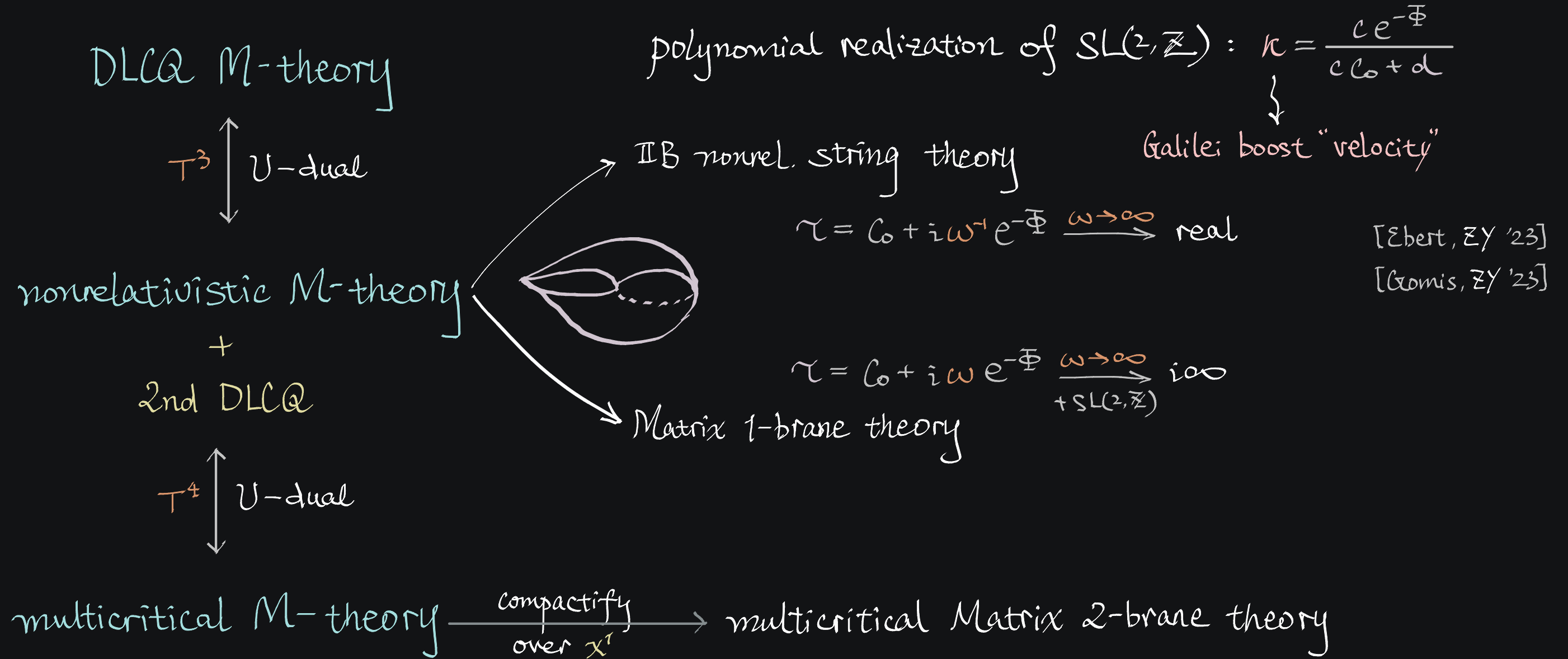
$$G_{\mu\nu} = \omega^{\frac{4}{3}} \tau_{\mu\nu} + \omega^{-\frac{2}{3}} E_{\mu\nu}$$

3D 8D

$$A_3 = -\omega^2 \tau^0 \wedge \tau^1 \wedge \tau^2 + a_3$$



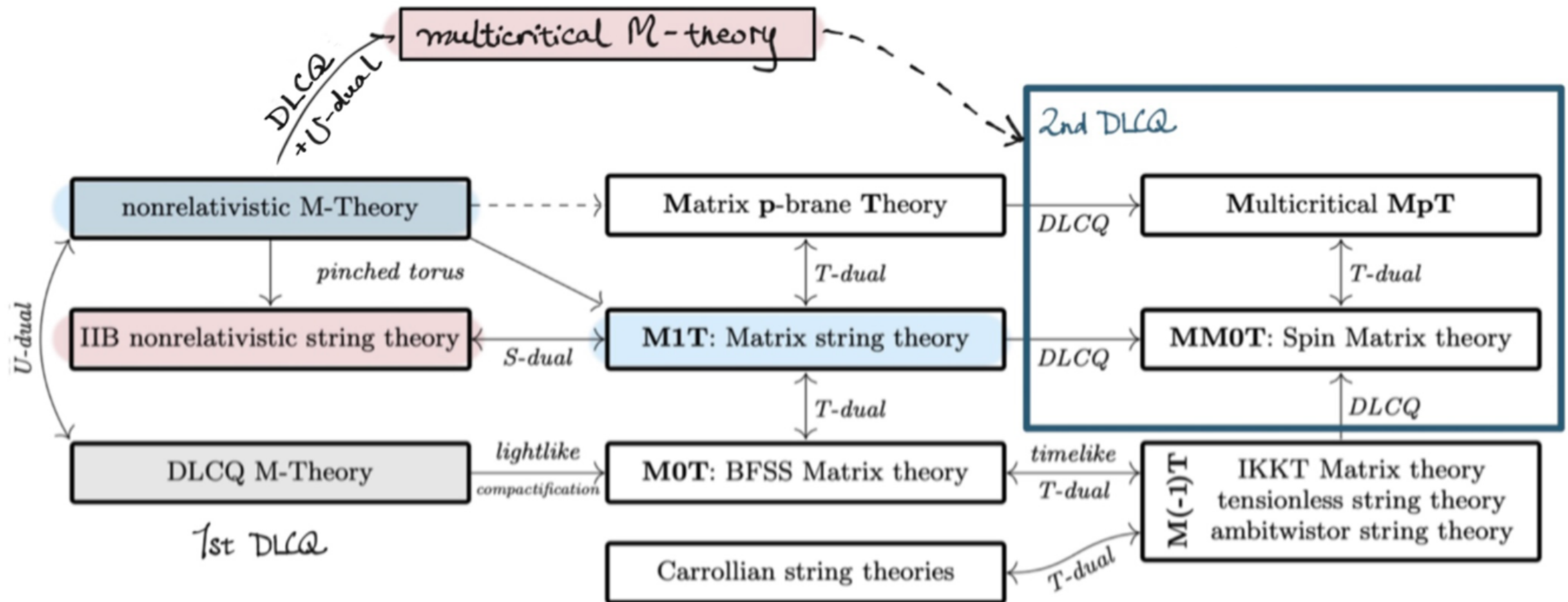
Anisotropic Compactification of Nonrelativistic M-Theory



$$G_{\mu\nu} = -\omega^2 \tau_\mu^0 \tau_\nu^0 + \tau_\mu^1 \tau_\nu^1 + \omega \tau_\mu^a \tau_\nu^a + \frac{1}{\omega} E_{\mu\nu}$$

$$A_3 = -\omega \tau^0 \wedge \tau^1 \wedge \tau^2 - \omega^2 \tau^0 \wedge \tau^3 \wedge \tau^4 + a_3$$

critical background
orthogonal M2-branes



Many new adventures ahead...

Thank You!