

Optimising loop amplitude evaluations

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Event generators' and N(n)LO codes' acceleration
CERN

14th November 2023

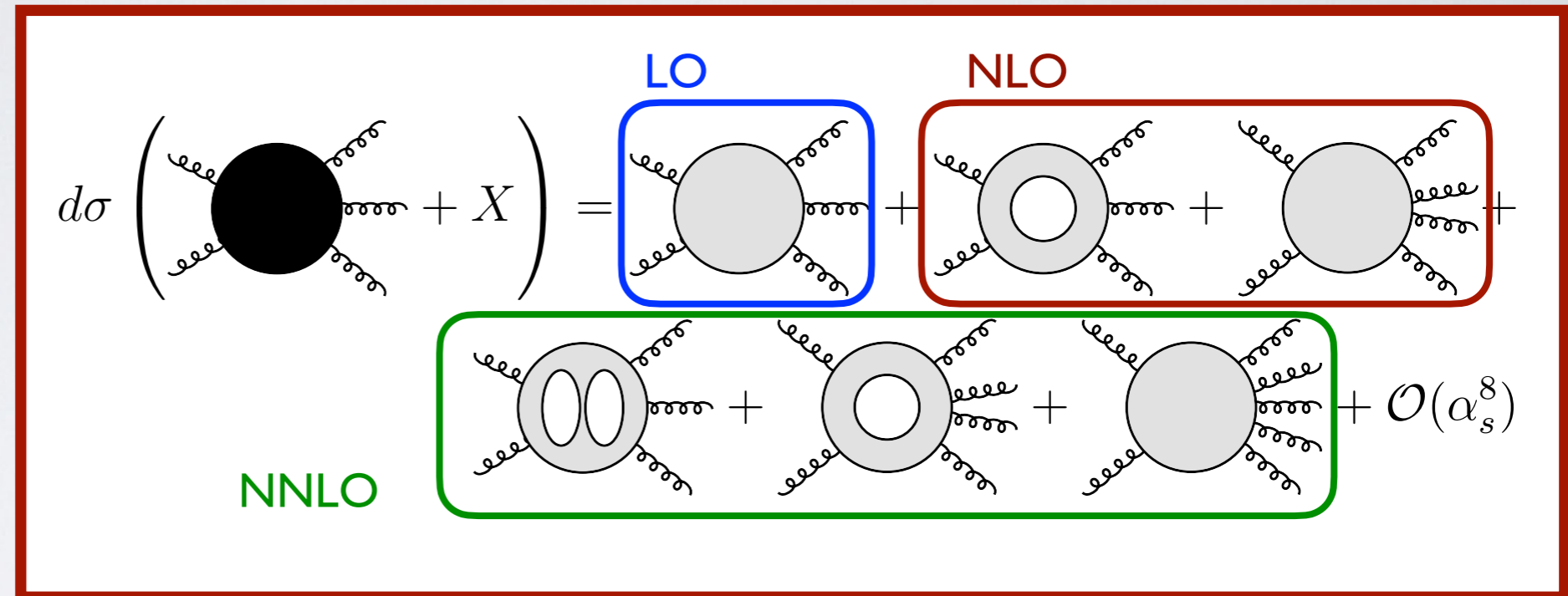


well known problem:

multi-scale amplitudes difficult to integrate

CPU cost scales
badly with higher
orders

(especially when IR
singular)



some solutions already exist

automated and
optimised MCs

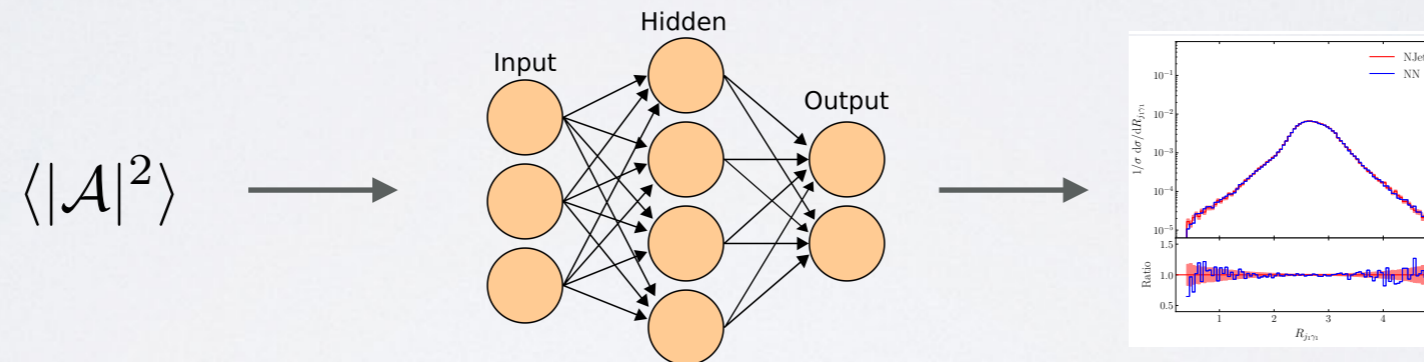
multi-core processing
re-weighting lower orders

...as well as other methods discussed at this meeting

this talk: testing amplitude neural network for loop induced processes

$gg \rightarrow \Upsilon\Upsilon g$ and $gg \rightarrow \Upsilon\Upsilon gg$

- Realistic hadron collider setup with SHERPA
- Precision evaluations: determining NN errors

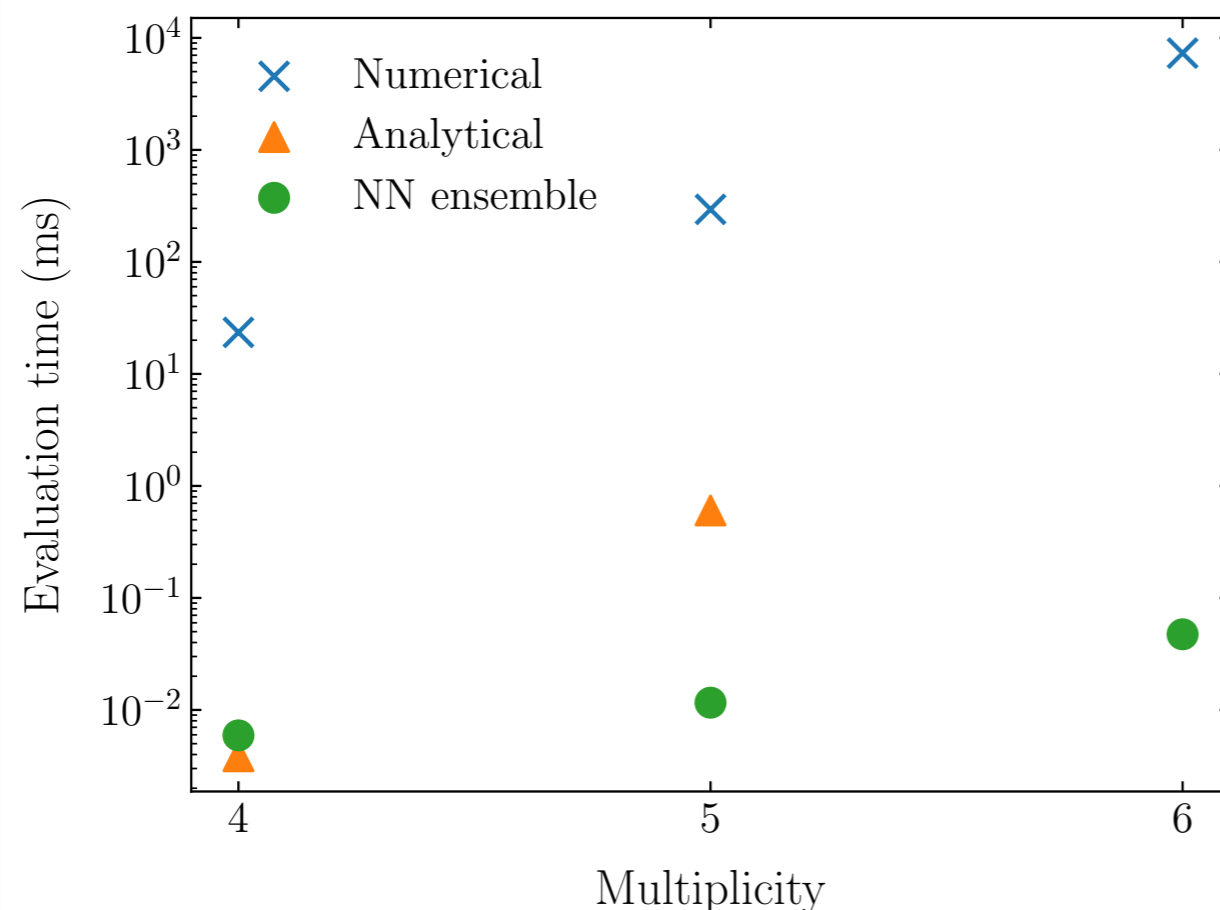
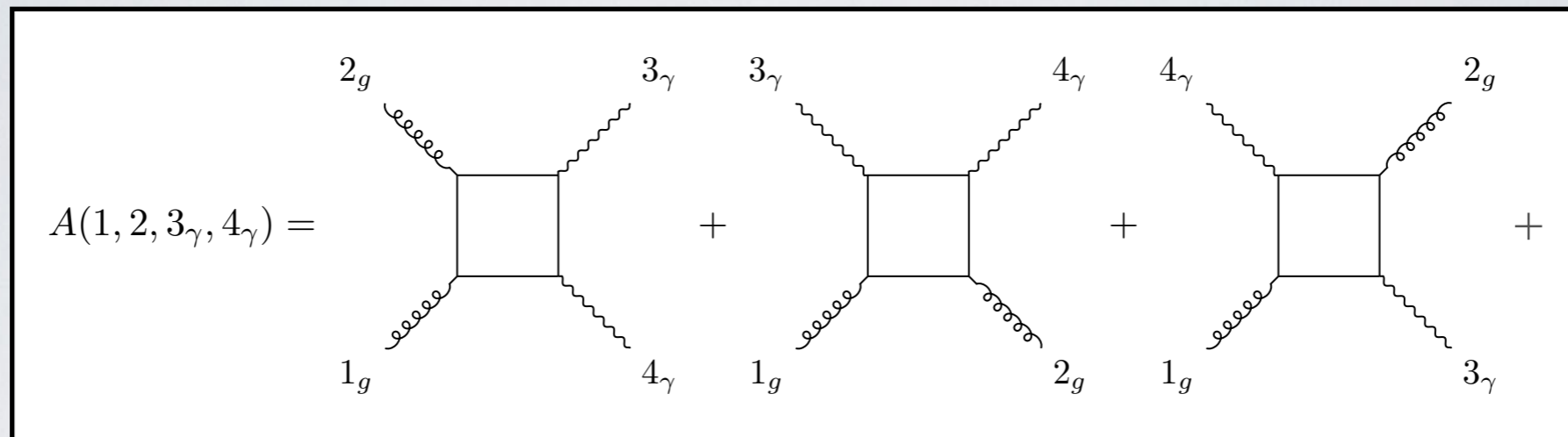


Aylett-Bullock, SB, Moodie JHEP 08 (2021) 066

SB, Butter, Luchmann, Pitz, Plehn SciPost Phys.Core 6 (2023) 034

- Reflections on experience at NNLO

How expensive are the loop amplitudes...



numerical evaluation with **NJET (v3)**
(SB, Biedermann, Moodie, Uwer, Yundin)
<https://bitbucket.org/njet/njet/wiki/Home>

(also includes some 2L 5pt amplitudes)

evaluation time including
error estimates

Precision is Important

standard one-loop matrix elements
provide high precision evaluations:

OpenLoops, GoSam, Madloop,
BlackHat, Recola, Helac-NLO...

one-loop $pp \rightarrow 5$ partons with NJET

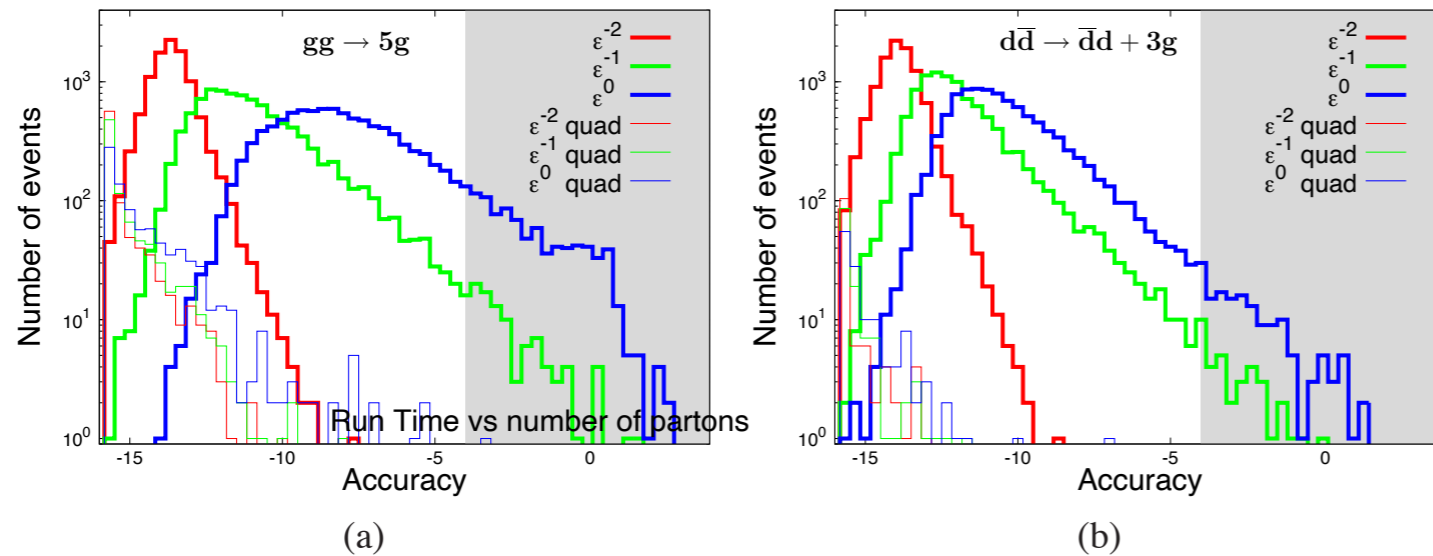


Figure 5. Accuracy for 5-jet amplitudes: (a) shows the seven gluon process and (b) the $d\bar{d} \rightarrow \bar{d}d + 3g$ process. The thicker histograms show computations in double precision whereas the thinner curves show the distribution in quadruple precision for points which did not pass the relative accuracy of 10^{-4} when calculated in double precision. Red histograms show the $\frac{1}{\epsilon^2}$ poles, green histograms the $\frac{1}{\epsilon}$ and blue histograms the finite part of the amplitudes.

use the ability to switch
numerical precision (e.g. qd
for 32 or 64 digits)

dimension scaling tests, gauge
invariance checks,...

Wishlist

an amplitude approximation which is:

- simpler to train/fit than generating events using traditional methods
- simplest model which fits a generic process
- reliable error estimates
- robust against changes in phase-space (cuts, jet algorithms, scale variations, etc.)
- simple to distribute

first attempt: $e^+e^- \rightarrow \text{jets}$

SB, (Aylett-)Bullock [2002.07516]

- Single NN does badly
- Splitting IR sectors via FKS sectors improves reliability
- **Error estimates** by varying model initialisation (ensemble of networks)
- K-factors work better than tree-level (no $1/s$ poles)
- Various tests suggest single run speed improvements at least $\times 10$

$$S_{i,j} = \frac{1}{D_1 s_{ij}}, \quad D_1 = \sum_{i,j \in \mathcal{P}_{\text{FKS}}} \frac{1}{s_{ij}},$$

$$d\sigma^{(X)} = \sum_{i,j} S_{i,j} d\sigma^{(X)},$$

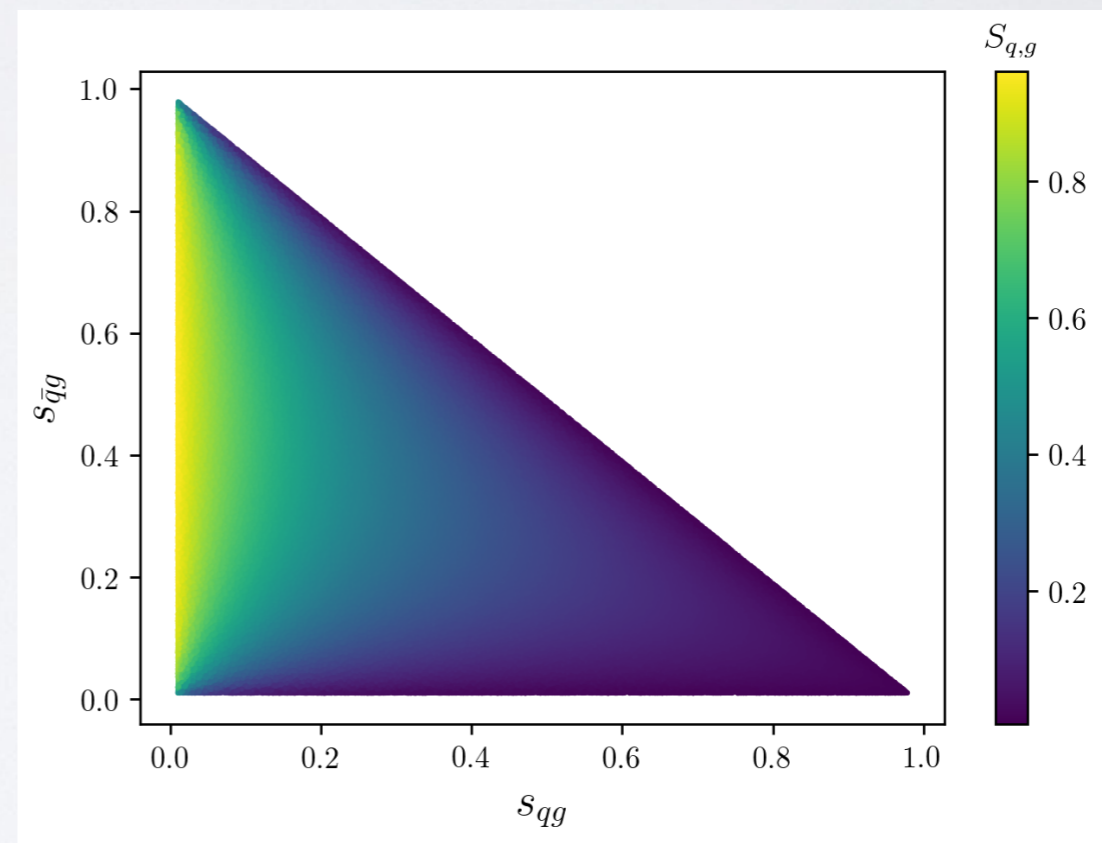
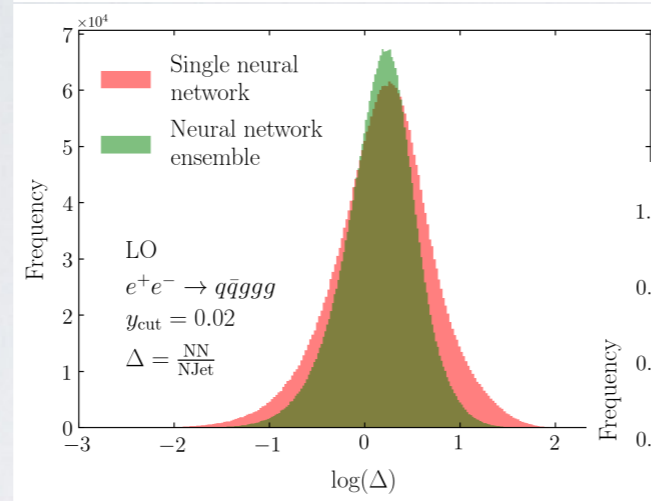
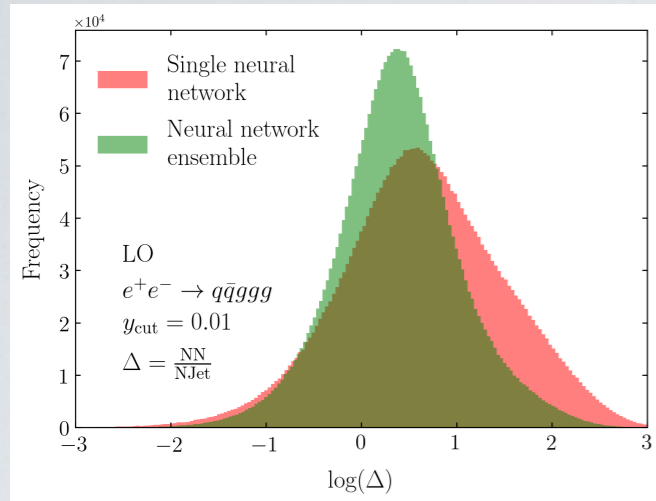
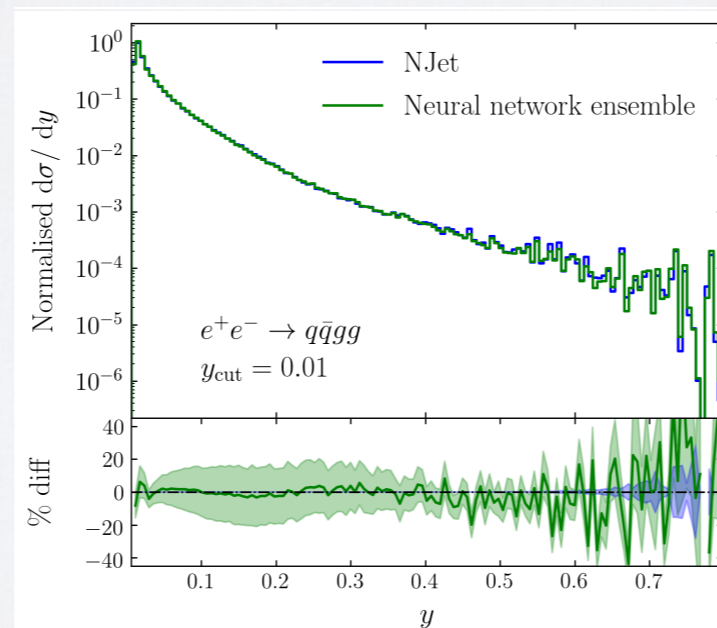
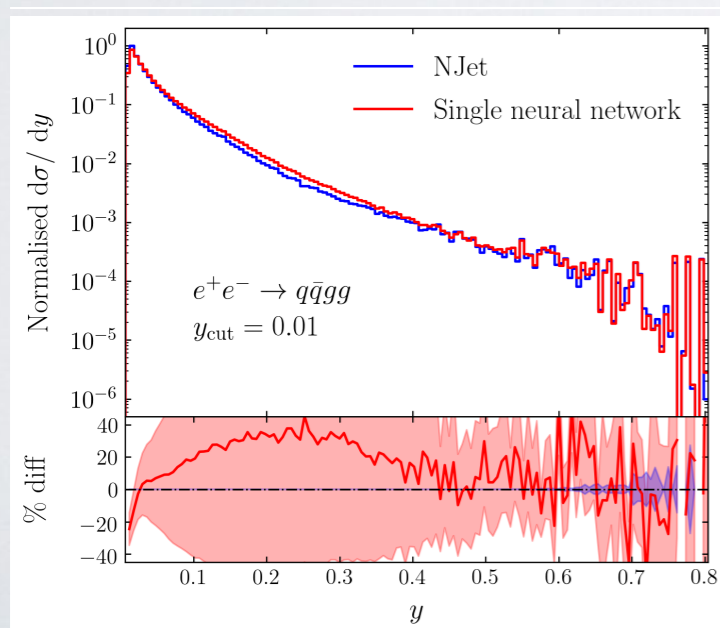
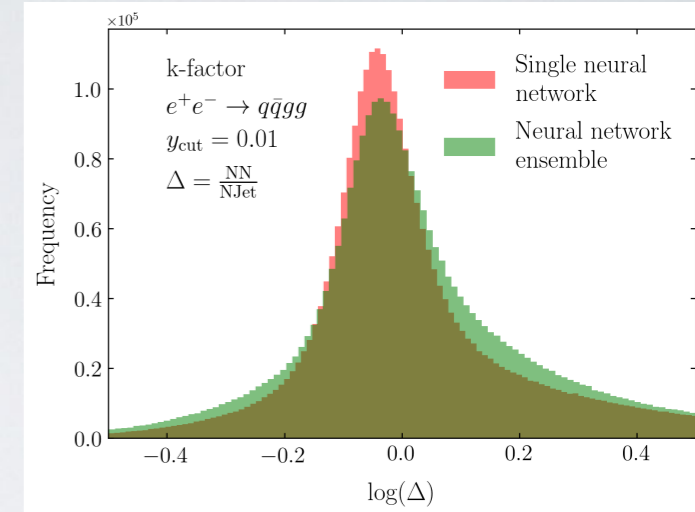
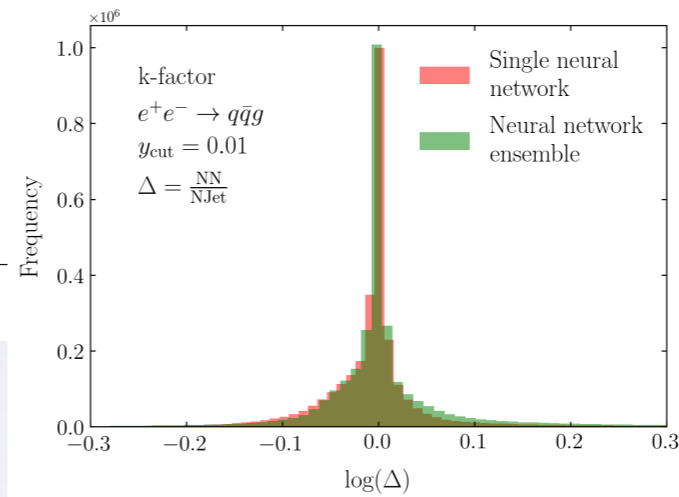


Figure 2: Behaviour of the $S_{q,g}$ FKS partition function

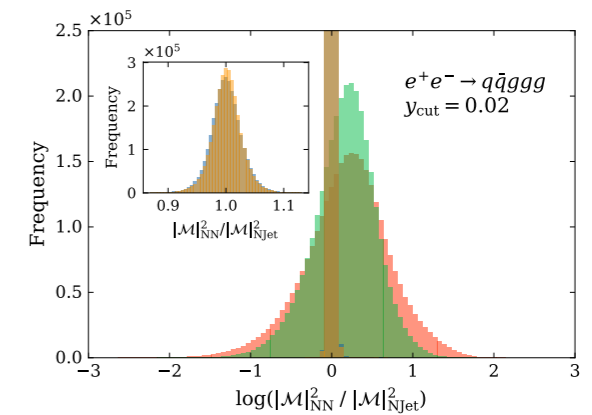
first attempt: $e^+e^- \rightarrow \text{jets}$



SB, (Aylett-)Bullock [2002.07516]



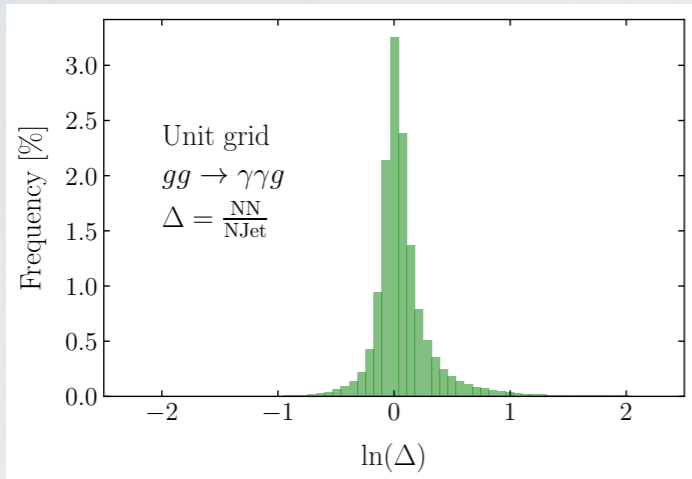
tree-level factorisation aware approach performs better



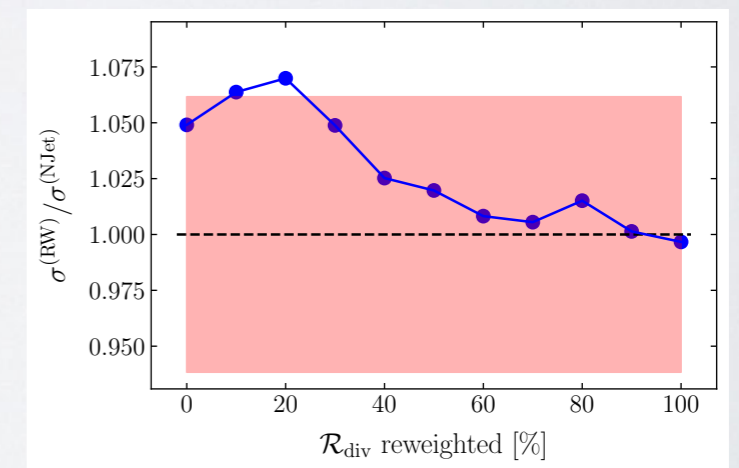
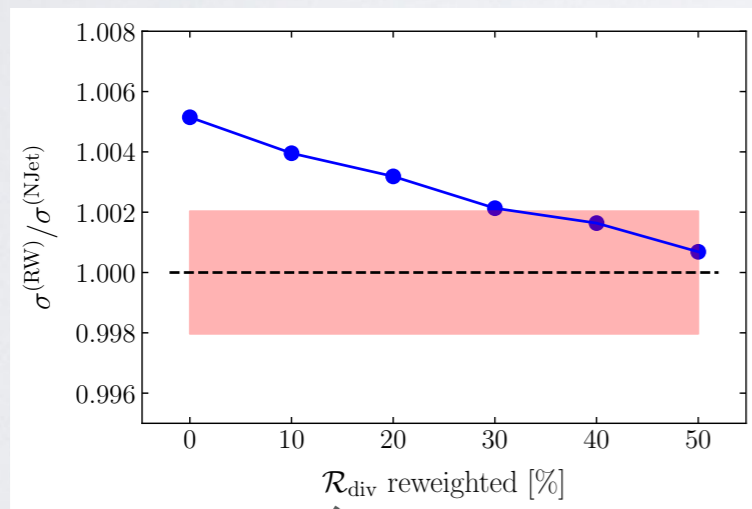
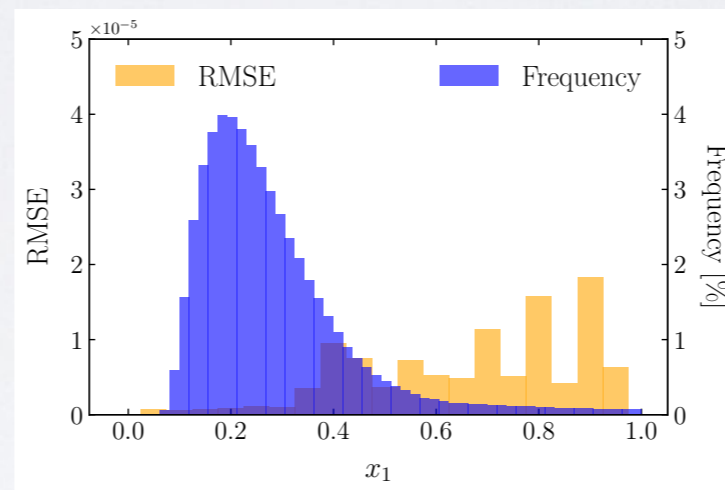
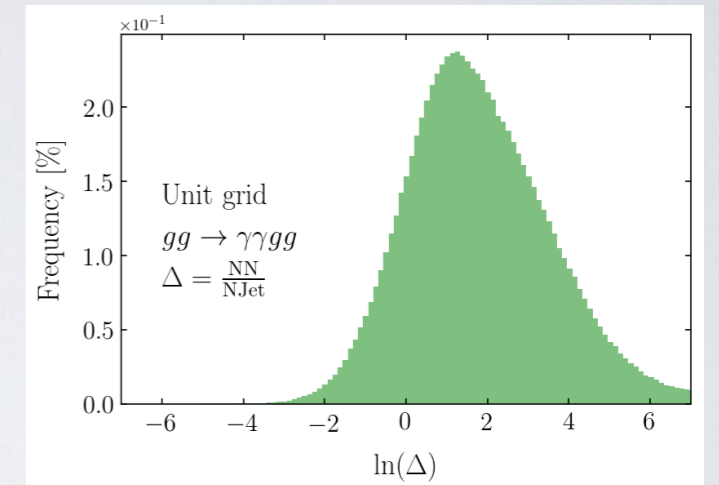
Maitre, Truong [2107.06625]

second attempt: $gg \rightarrow \Upsilon\Upsilon + \text{gluons}$

Aylett-Bullock, SB, Moodie [2106.09474]



$p_{T,j} > 20 \text{ GeV}$ $R_{\gamma,j} > 0.4$ $|\eta_j| < 5$
 $p_{T,\gamma_1} > 40 \text{ GeV}$ $R_{\gamma,\gamma} > 0.4$ $|\eta_\gamma| < 2.37$
 $p_{T,\gamma_2} > 30 \text{ GeV}$



7-9% of the total phase space

speed-up $\sim N_{\text{infer}} / N_{\text{train}}$

third attempt: $gg \rightarrow \Upsilon\Upsilon + \text{gluons}$ with **Bayesian networks**

SB, Butter, Luchmann, Pitz, Plehn [2206.14831]

- weights and biases are replaced with (Gaussian) distributions
- optimised training times vs. pure ensemble approach

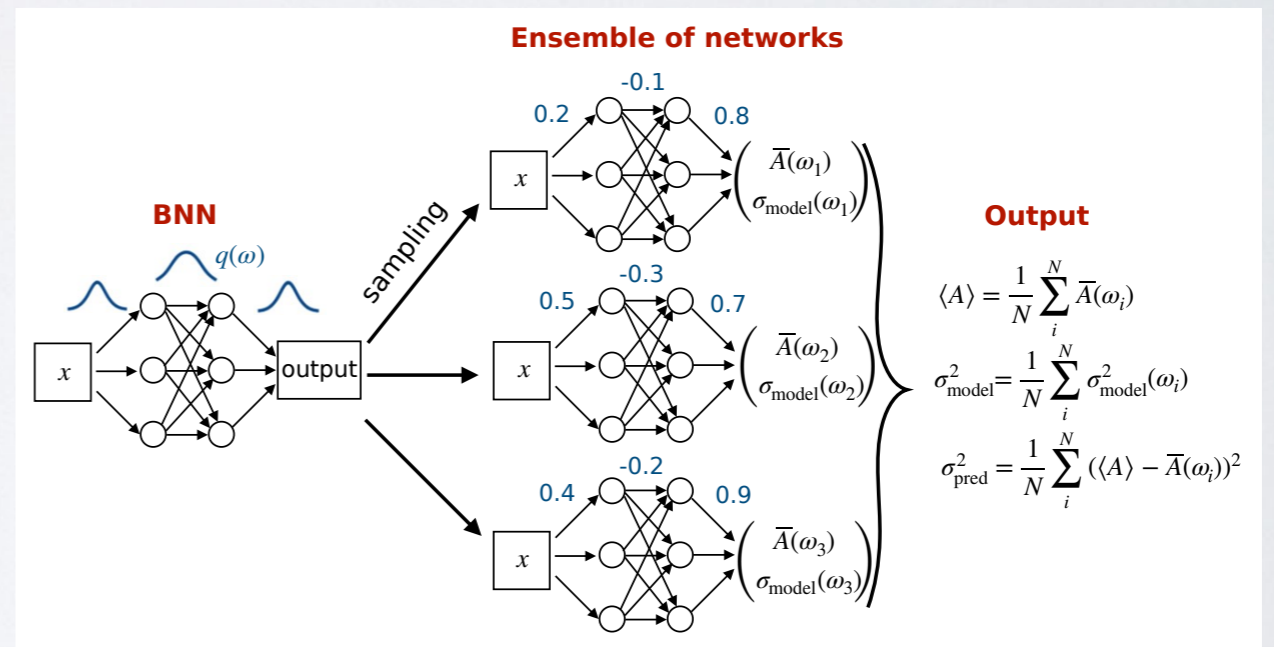
- better defined error estimates
- improved training via loss and performance boosting

single BNN

$$A_j \rightarrow \log\left(1 + \frac{A_j}{\sigma_A}\right)$$

~6k params for $2 \rightarrow 3$

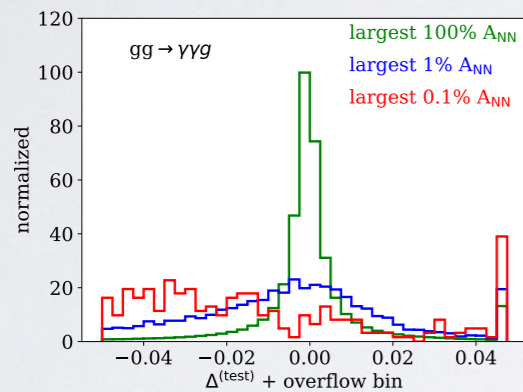
~600k params for $2 \rightarrow 4$



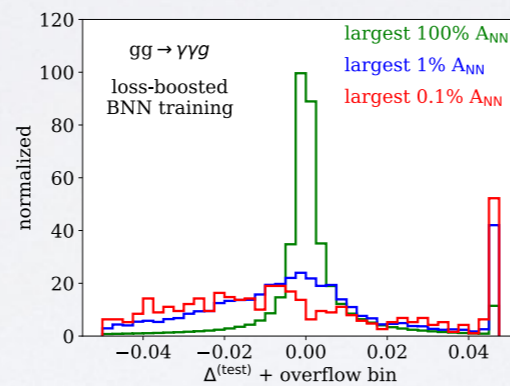
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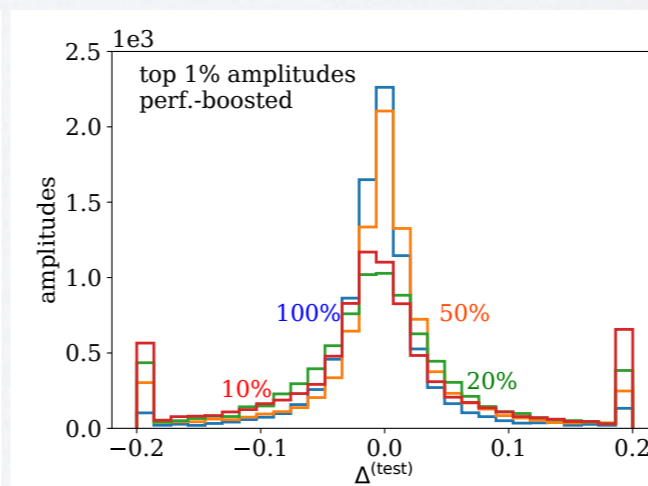
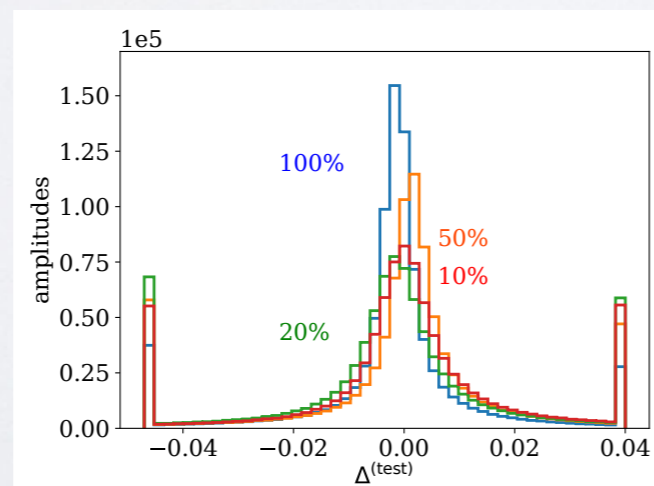
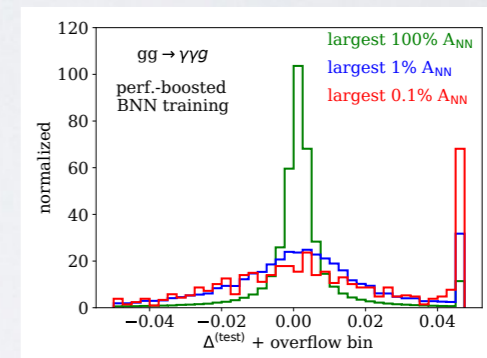
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loss
boosting



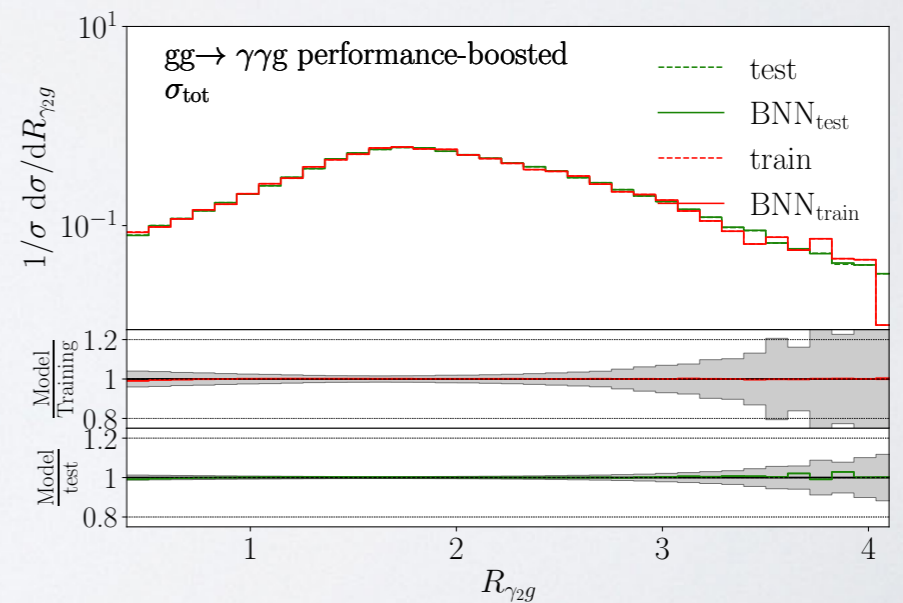
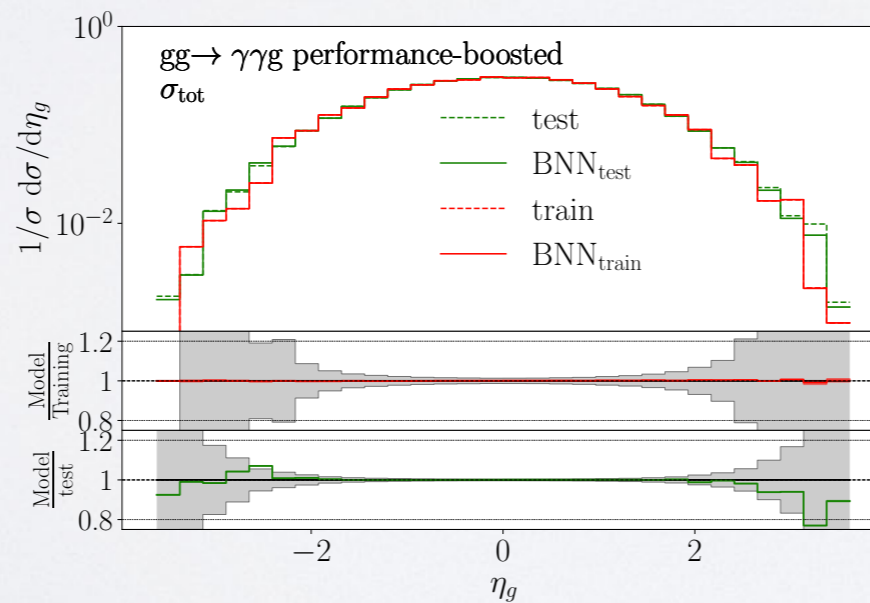
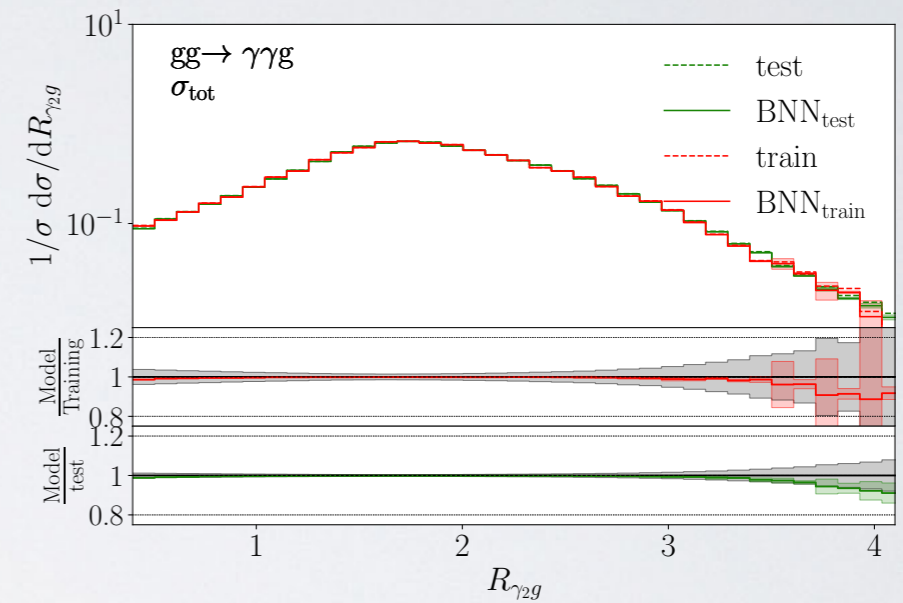
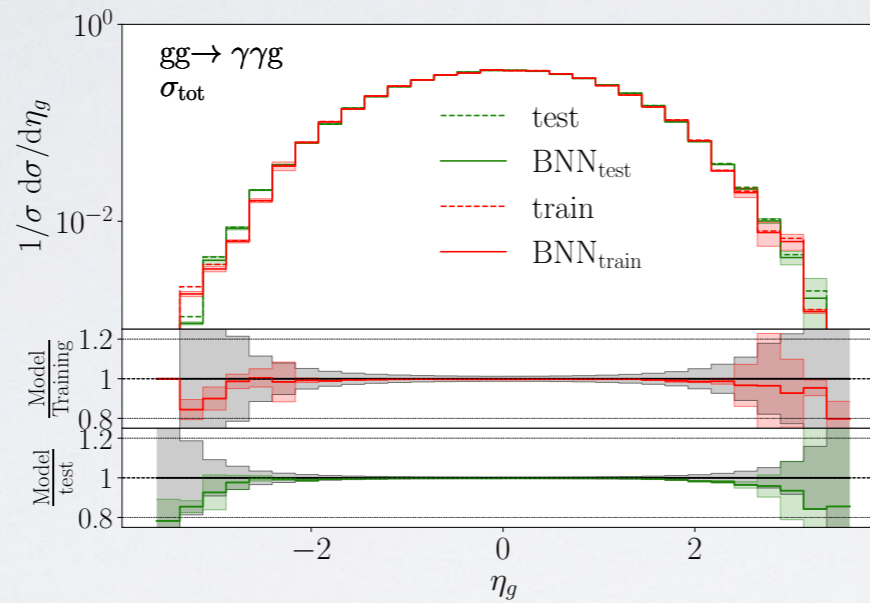
performance
boosting



third attempt: $gg \rightarrow \Upsilon\Upsilon + \text{gluons}$ with Bayesian networks

SB, Butter, Luchmann, Pitz, Plehn [2206.14831]

performance
boosting reduces
uncertainties in tails

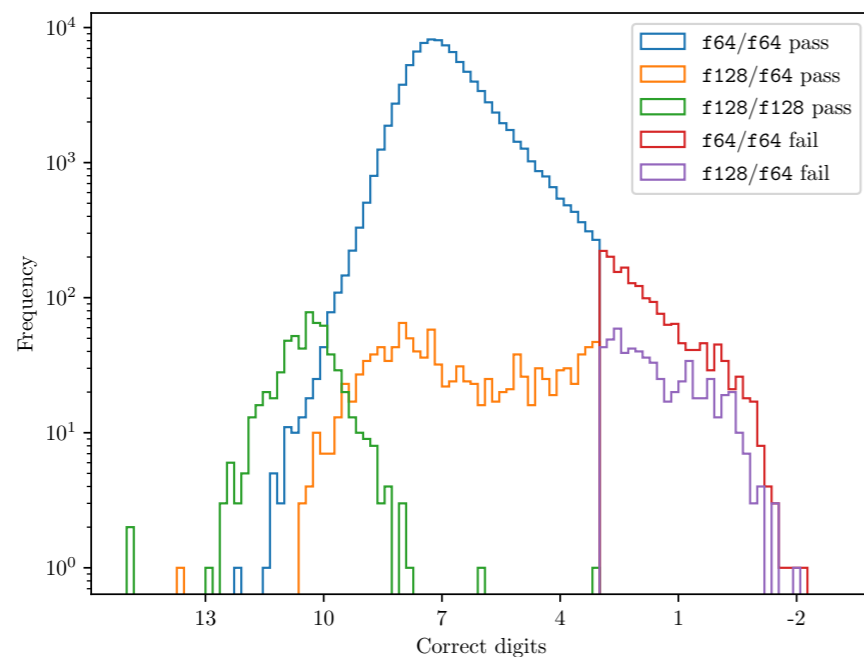


Higher Order Applications

precision frontier has moved to
NNLO (QCD) and beyond

2 → 2 @3-loop,
2 → 3 @ 2-loop

e.g. 2L gg → YYg



SB, Gehrmann, Marcoli, Moodie 2109.12003

precision requirements more subtle
instabilities can be in both rational
coefficients and special functions

reasonable performance for 2→3 with
analytic finite remainders (extracted
using Finite Field sampling)

$$A_{i;j}^{(L),k} = \sum_{s=-2L}^{o(L)} \sum_r \epsilon^s c_{r,s}(\vec{x}) \text{mon}_r(f, \mathbf{c})$$



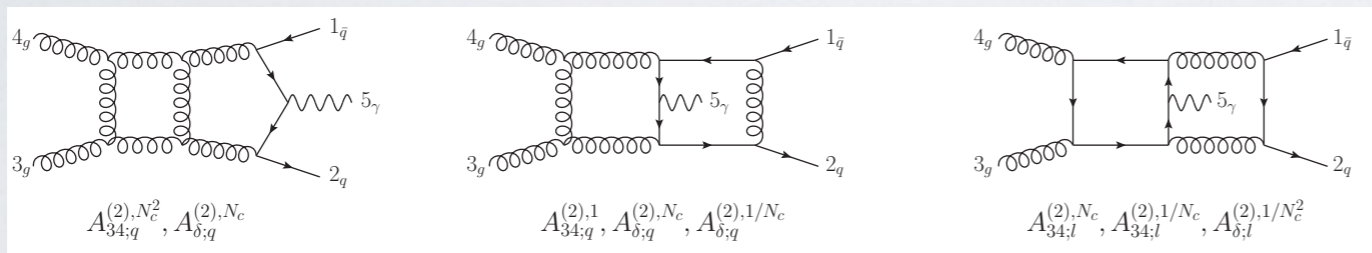
pentagon functions [Chicherin, Sotnikov, Zoia]

All Two-Loop Feynman Integrals for Five-Point
One-Mass Scattering, Abreu et al. [2306.15431]

Single Photon plus Two-Jets

SB, Czakon, Hartanto, Moodie, Peraro, Poncelet, Zoia [2304.06682]

first full colour $2 \rightarrow 3$ differential cross section



IR subtraction with
STRIPPER approach

comparison with ATLAS [1912.09866]

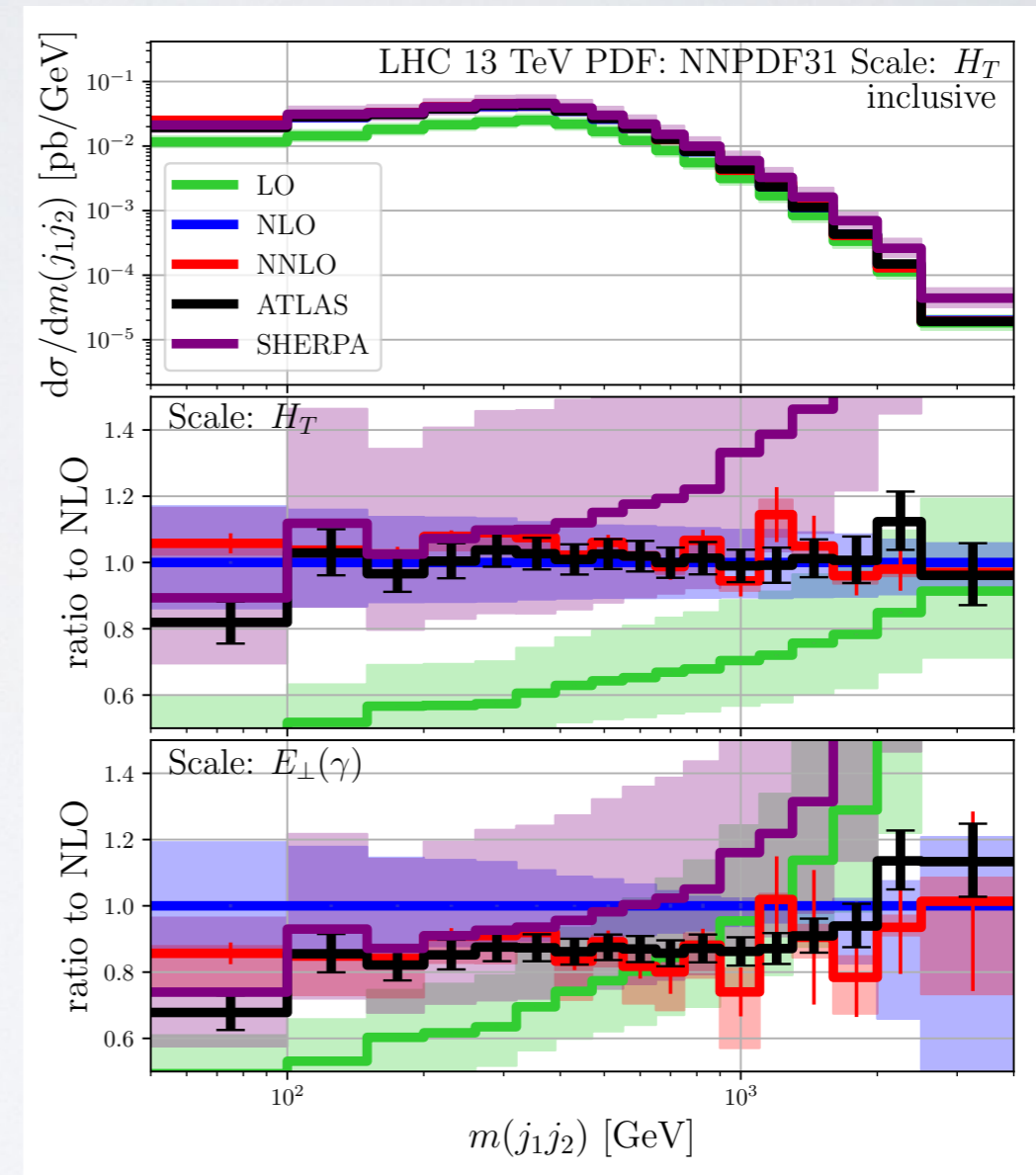
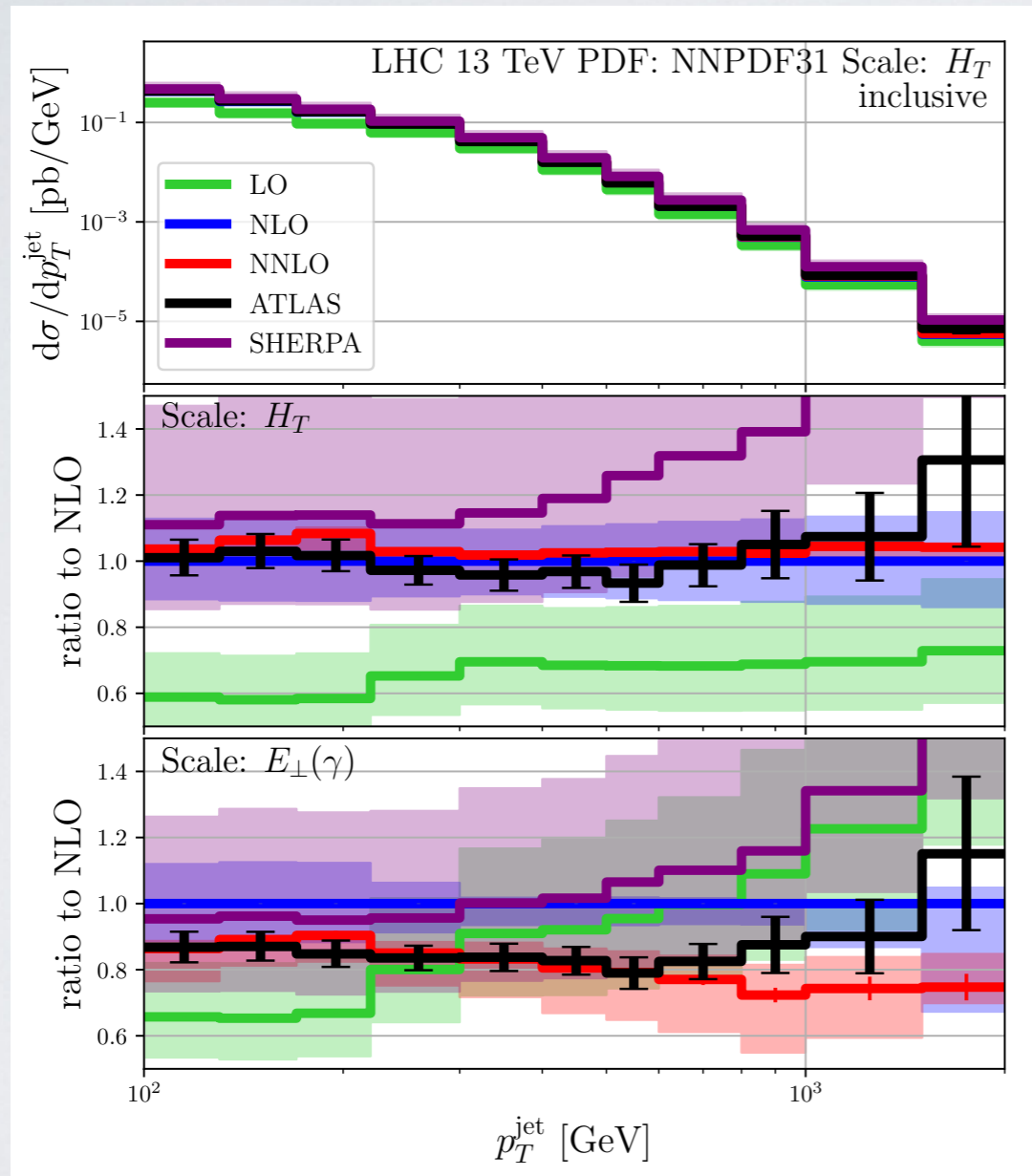
1. We require at least two jets defined with the anti- k_T algorithm [106] for jet radius $R = 0.4$ that have minimal transverse momentum of $p_T(j) > 100$ GeV and maximal rapidity $|\eta(j)| < 2.5$.
2. The identified jets must be separated from the photon by $\Delta R(\gamma, j) > 0.8$.
3. One isolated photon must be present in the final state with $E_\perp(\gamma) \geq 150$ GeV, $|\eta(\gamma)| \leq 2.37$ excluding $1.37 \leq |\eta(\gamma)| \leq 1.56$.

analytic 2L finite
remainders

RV and RR using ME
from OpenLoops and
AvH libraries

Single Photon plus Two-Jets

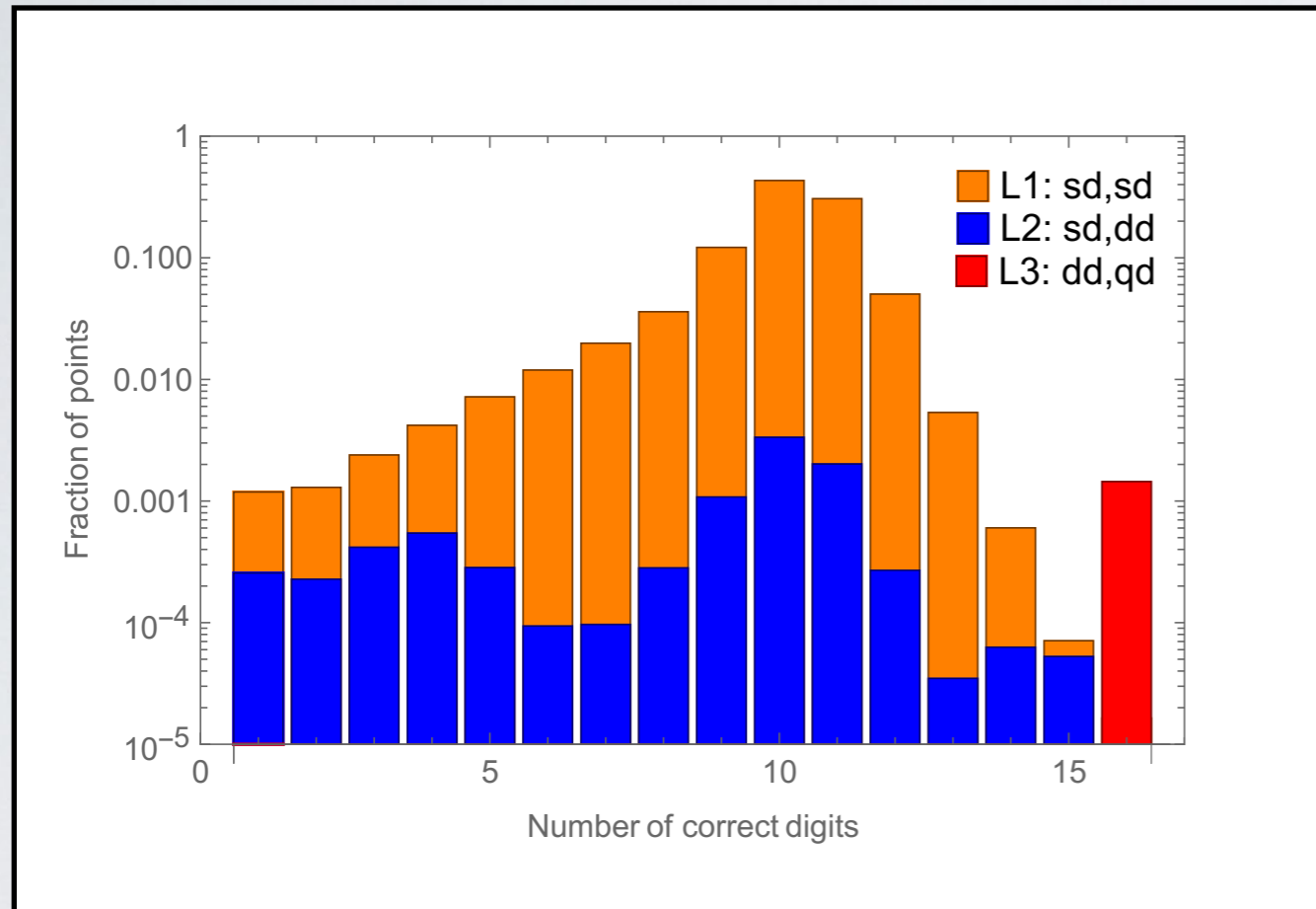
SB, Czakon, Hartanto, Moodie, Peraro, Poncelet, Zoia [2304.06682]



excellent overall agreement with data

Single Photon plus Two-Jets

SB, Czakon, Hartanto, Moodie, Peraro, Poncelet, Zoia [2304.06682]



total cost ~ 550 kCPUh

VV:RV:RR $\sim 1:10:40$

double virtual call average ~ 16 s per point
[NB: 1 double precision call $O(1$ s)]

Note: the 1st time is
always a challenge - many
improvements possible

Future Applications

speeding up amplitude calls with NN looks like a viable option for leading order codes

I would still be interested in better control for the number of correct digits

speeding up amplitude virtual amplitudes can help, but need to make impact on real radiation

two-loop $2 \rightarrow 3$ not dominated by virtuals, largely thanks to well studied special function basis

This is (probably) not going to continue: $pp \rightarrow ttj$, $pp \rightarrow ttH$, $\rightarrow WWWj$

at some stage **analytic** formula unfeasible - must take **numerical** route

special function basis in 'pentagon function' form will not be possible (eg. if there are elliptic structures)

expect dramatic change in evaluation time

Conclusions

I still have some questions...

phase-space : how can we use amplitude values to minimize the required number of training points

squared amplitudes vs (ordered) helicity amplitudes

analytic vs. numeric: can we improve the architecture to better satisfy amplitude properties? (c.f. bootstrap techniques)