宇宙ひもによるダークフォトンDM生成 と重力波生成



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NK, Kazunori Nakayama (Tohoku U.), 2212.13573, 2306.17390

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Caputo et al 2105.04565

Gravitational wave proves



Gravitational wave proves



Dark photon DM production

- Gravitational particle production during inflation / reheating

Graham, Mardon, Rajendran (2016) / Ema, Nakayama, Tang (2019) Sato, Takahashi, Yamada (2022)

$$\Omega_{\gamma'} \simeq \Omega_{\rm DM} \sqrt{\frac{m_{\gamma'}}{6\,\mu {\rm eV}}} \left(\frac{H_{\rm inf}}{10^{14}\,{\rm GeV}}\right)^2 \ ->$$
 lower limit on dark photon mass

- Resonant production from scalar field

Axion : Agrawal, <u>NK</u>, Reece, Sekiguchi, Takahashi (2020) Co, Pierce, Zhang, Zhao (2019), Bastro-Gil, Santiago, Ubaldi, Vega-Morales (2019) <u>NK</u>, Takahashi (2023)

Higgs : Harigaya, Narayan (2019), Nakayama Yin (2021)

Spectator : Nakai, Namba, Obata (2022)

- Misalignment production Nakayama (2019), Nakayama (2020), NK, Nakayama (2023)

- Production from cosmic strings Long, Wang (2019), NK, Nakayama (2022)

Resonant dark photon production from axion

Agrawal, NK, Reece, Sekiguchi, Takahashi, 1810.07188 Co, Pierce, Zhang, Zhao, 1810.07196 Bastero-Gil, Santiago, Ubaldi, Vega-Morales, 1810.07208

$$\mathcal{L} = \frac{1}{2} \partial^{\mu} \phi \partial_{\mu} \phi - V(\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} m_{\gamma'}^2 A_{\mu} A^{\mu} - \frac{\beta}{4f_a} \phi F_{\mu\nu} \tilde{F}^{\mu\nu}$$
$$\longrightarrow \quad \ddot{\mathbf{A}}_{\mathbf{k},\pm} + H \dot{\mathbf{A}}_{\mathbf{k},\pm} + \left(m_{\gamma'}^2 + \frac{k^2}{a^2} \mp \frac{k}{a} \frac{\beta \dot{\phi}}{f_a} \right) \mathbf{A}_{\mathbf{k},\pm} = 0$$





- Axion abundance is suppressed & dark photon is dominant

Agrawal, NK, Reece, Sekiguchi, Takahashi, 1810.07188 (see also NK, T. Sekiguchi, F. Takahashi, 1711.06590)

- Produced dark photons can stabilize the dark Higgs $V(\Phi) \ni |\mathbf{A}|^2 |\Phi|^2$

-> secondary inflation, early dark energy

NK, Nakagawa, Takahashi, 2111.06696 Nakagawa, Takahashi, Yin, 2209.01107

- GW emission with circular polarization NK, Soda, Urakawa, 2010.10990

see also Machado+ (2019), Salehian+ (2020), Ratzinger+ (2020), Namba+ (2020)

Coherent vector DM production

Nakayama (2019), Nakayama (2020), NK, Nakayama (2023)

CMB observation —> $g_k \lesssim 0.01$, $S \lesssim 0.1\zeta$

"Viable" coherent vector DM scenario

NK, Nakayama, 2303.04287

curvaton scenario : introduction of an additional scalar field responsible for the curvature perturbation



Abelian-Higgs model

$$\mathcal{L} = (\mathcal{D}_{\mu}\Phi)^* \mathcal{D}^{\mu}\Phi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} - V(\Phi), \ V(\Phi) = \frac{\lambda}{4}(|\Phi|^2 - v^2)^2$$
$$(\mathcal{D}_{\mu} = \partial_{\mu} - ieA_{\mu}, \ F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu})$$

Higgs mechanism may give the dark photon mass

formation of <u>cosmic string</u>
 (if symmetry breaking occurs after inflation)

Assumption : $m_A \ll m_\Phi$ $(e \ll \sqrt{\lambda})$ (Type II string)

-> dark photon production



Ellis, Gaillard, Nanopoulos 1504.07217

Abelian-Higgs string (local string)

$$\mathcal{L} = (\mathcal{D}_{\mu}\Phi)^{*}\mathcal{D}^{\mu}\Phi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} - V(\Phi), \ V(\Phi) = \frac{\lambda}{4}(|\Phi|^{2} - v^{2})^{2}$$
$$(\mathcal{D}_{\mu} = \partial_{\mu} - ieA_{\mu}, \ F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu})$$

static solution:
$$\Phi(r,\theta) = v e^{in\theta} f(r), \ A(r,\theta) = \frac{n\alpha(r)}{er} \hat{e}_{\theta}$$

with boundary condition $f(0) = \alpha(0) = 0, \ f(\infty) = \alpha(\infty) = 1$

energy of local string :

 \mathcal{Z}

 θ

$$\mathcal{E} = \int d^3x \left(|\mathcal{D}_i \Phi|^2 + V(\Phi) + \frac{1}{2} |\nabla \times \mathbf{A}|^2 \right)$$

= $2\pi L \int r dr \left(f'^2 + \frac{n^2 (\alpha - 1)^2}{r^2} f^2 + V(f) + \frac{n^2 \alpha'^2}{2e^2 r^2} \right)$
= μL (string tension × string length)

Global (axion) string

$$\mathcal{L} = (\partial_{\mu}\Phi)^* \partial^{\mu}\Phi - V(\Phi), \ V(\Phi) = \frac{\lambda}{4} (|\Phi|^2 - v^2)^2$$

static configuration: $\Phi(r, \theta) = v e^{in\theta} f(r)$ with f(0) = 0 & $f(\infty) = 1$

energy of global string :

(cutoff scale ~ horizon scale)



local string is more localized

Cosmic string "network"





collision of straight strings



pair annihilation of string and anti-string (opposite winding number)

collision of curved strings





-> source of gravitational waves

Cosmic string network





Scaling law : O(1) strings per Hubble patch

Scenario

Long, Wang 1901.03312 NK, Nakayama 2212.13573

- "Light" dark photons can be produced by cosmic strings

small gauge coupling i.e. $m_A \ll m_\Phi$ ($e \ll \sqrt{\lambda}$) (Type II string) e = 0 limit corresponds to the massless NG boson emission (global string case)

- Dark photon production becomes inefficient for $\ell_{
m loop}\gtrsim m_A^{-1}$

i.e. loop oscillation frequency becomes smaller than the mass $\rightarrow H \lesssim m_A$ Dark photon abundance is fixed

cf. axion DM are produced through domain wall decay

- After that, string evolves like "local" string

network loses the energy only through the GW emission (Nambu-Goto limit)

Field theoretic (Abelian-Higgs) simulation

$$\Phi'' + 2\mathcal{H}\Phi' - D_i D_i \Phi + a^2 \frac{\partial V}{\partial \Phi^*} = 0,$$

$$E'_i + \partial_j F_{ij} - 2ea^2 \operatorname{Im}(\Phi^* D_i \Phi) = 0,$$

$$\partial_i E_i - 2ea^2 \operatorname{Im}(\Phi^* \Phi') = 0$$

Moore+ (2001) Bevis+(2007) Dufaux(2010) Hiramatsu+(2013) Correia+(2020) and more

* Lattice-gauge formulation, Gauss's law is satisfied



- AOBA (SX-Aurora TUBASA) in Tohoku U.



4,032 Vector Engine



16 cores, 256 vector length, 96GB / 1VE

- FUGAKU supercomputer (Riken) (trial use)



158.976 nodes

48 cores, 32GiB / 1node

Loop production & decay







Scaling law of local / global string

Mean separation of neighboring strings: $d_{\rm sep} = \sqrt{\frac{V^{(c)}}{\ell_{\rm str}^{(c)}}}$ (comoving / physical length)

Scaling law $-> d_{\rm sep} \propto \tau$, i.e. $d'_{\rm sep} = {\rm const}$ (conformal / physical time)



Correia+ arXiv:2005.14454

Hindmarsh+ arXiv:1908.03522

Global string & scaling violation?

Number of strings per Hubble volume: $\xi = \frac{\ell_{\rm str} t^2}{V}$

 $\xi = \text{const}$ (scaling law) vs Hindmarsh et al, 1908.03522

Hindmarsh et al, 2102.07723

 $\xi = \text{const}$ (scaling law) vs $\xi = \alpha \log(t) + \beta$ (scaling violation)

Gorghetto et al, 1806.04677

Kawasaki et al, 1806.05566

Buschmann et al, 2108.05368



Hindmarsh+ 1908.03522

Kawasaki+ 1806.05566

Mean separation NK, Nakayama 2212.13573



Table 1: Simulation setup and linear fitting parameters of the mean string separation in terms of the conformal time, defined by $m_r d_{sep} = a m_r \tau + b$.



Table 1: Log fitting parameters of the string length parameter, defined by $\xi = \alpha \log(m_r/H) + \beta$.

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Emission of (longitudinal) vector boson

$$\rho_A^{(L)} = \frac{|\Phi|^2}{v^2} \left[\frac{2}{a^2} \left(\frac{\text{Im}(\Phi^* \Phi')}{|\Phi|} \right)^2 + \frac{1}{a^4} \left(E_i^{(L)} \right)^2 \right]$$

$$n_A = \int dk \frac{dn_A}{dk} = \int dk \frac{1}{E_A(k)} \frac{d\rho_A}{dk} \qquad n_A^{(L)}(t) \simeq \frac{8\xi\mu H}{\bar{E}_A/H} \quad \text{(analytic estimation)} \\ \text{Long, Wang 1901.03312}$$



Spectrum of emitted dark photon

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peak wavenumber: $k/a \sim 10H \iff$ typical loop size: $\ell \sim 0.1H^{-1}$

Relic abundance

$$\Omega_A h^2 = \frac{m_A (n_{A,0}/s_0)h^2}{\rho_{\rm cr,0}/s_0} \simeq 0.091 \left(\frac{\xi}{12}\right) \left(\frac{m_A}{10^{-13}\,{\rm eV}}\right)^{1/2} \left(\frac{v}{10^{14}\,{\rm GeV}}\right)^2$$

 $\xi = \text{const}$ (scaling law)

Hindmarsh et al, 1908.03522 Hindmarsh et al, 2102.07723

$$\xi = 0.15 \log\left(\frac{m_r}{m_A}\right) \simeq 12 + 0.15 \log\left[\left(\frac{m_r}{10^{14} \,\mathrm{GeV}}\right) \left(\frac{10^{-13} \,\mathrm{eV}}{m_A}\right)\right]$$

(scaling violation)

Gorghetto et al, 1806.04677 Kawasaki et al, 1806.05566 Buschmann et al, 2108.05368

GW emission from cosmic strings



Credit: Daniel Dominguez/CERN

Quadrupole formula for GW emission: $\dot{E}_{\rm GW} \sim G(\ddot{D})^2$

quadrupole moment: $D \sim ML^2 \sim \mu L^3$, $\ddot{D} \sim \omega^3 D \sim L^{-3} D$

L : typical loop size ~ (typical oscillation frequency)-1

GW emission rate: $\dot{E}_{\rm GW} \sim G\mu^2 \equiv \Gamma_{\rm GW} G\mu^2$ $G\mu \sim (v/M_P)^2 \sim 10^{-7} (v/10^{15} {\rm GeV})^2$

GW spectrum from local/global strings



energy loss of loops = GW emission + vector boson emission

$$\frac{dE_{\ell}}{dt} = -\Gamma_{\rm GW}G\mu^2 - \Gamma_{\rm vec}v^2\theta(1 - m_A\ell)$$
$$\Gamma_{\rm GW} \sim 50, \ \Gamma_{\rm vec} \sim 65$$

Loops shorter than m_A-1 can emit vector bosons (i.e. loop oscillation frequency should be larger than m_A)

-> short lived & GW emission is suppressed

$$\begin{aligned} \text{loop lifetime:} \\ \tau(\ell) \sim \frac{E_{\ell}}{\dot{E}_{\ell}} \sim \begin{cases} \frac{\ell}{\Gamma_{\text{GW}} G \mu} & \text{for } m_A \ell > 1 \text{ (GW emission)} \\ \frac{\pi \log(m_r/H)\ell}{\Gamma_{\text{vec}}} & \text{for } m_A \ell < 1 \text{ (wetor boson emission)} \\ \tau \sim \ell < t & -\text{ short-lived} \end{cases} \end{aligned}$$

Power of GW emission by single loop:

$$\frac{dP_{\rm GW}(\ell)}{d\ln f} = G\mu^2 f\ell S(f\ell)$$



GW spectrum : our scenario

NK, Nakayama 2212.13573



no-scaling violation

log-enhanced (scaling violation)

GW spectrum : our scenario

NK, Nakayama 2212.13573



(a)
$$v = 10^{15} \text{ GeV}, m_A = 10^{-14} \text{ eV}$$

(b) $v = 10^{13} \text{ GeV}, m_A = 10^{-10} \text{ eV}$ (c) $v = 10^{12} \text{ GeV}, m_A = 10^{-5} \text{ eV}$



GW spectrum : implication for NANOGrav 15 yr data

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Viable parameter space

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Sweet spot is m~1-10µeV

Discussion

We need more precise study ...

- Scaling violation ?
- Time-dependence of the tension
- Loop formation and dark photon production rate especially near the transition : $H \sim m_A$ (global —> local)
- Initial loop size distribution (monochromatic or extended?)
- Spectral function of GW from individual loop (cusp- or kink-like?) (because it is crucial for high frequency region)
- Loop lifetime (deviation from Nambu-Goto string)

discussed in Hindmarsh et al (2017), Matsunami et al (2019)

Summary

- Dark photon can be produced from the network of cosmic strings (via loop collapse)
 - & production stops when $H < m_A$

(i.e. the dark photon emission is kinematically suppressed)

–> Relic abundance is fixed at that time
 Observed abundance is reproduced for e.g.

 $v \sim 10^{12} \text{--} 10^{14} \text{GeV}, \ m_A \sim 10^{-14} \text{--} 10^{-5} \text{eV}$

- Gravitational waves are emitted as a signal of this scenario

Spectrum is different from both local and global one It can be tested by combining pulsar timing and direct detection

- NANOGrav data can be explained & tested by future aLIGO

GW emission

Power of GW emission by single loop:

$$\frac{dP_{\rm GW}(\ell)}{d\ln f} = G\mu^2 f\ell S(f\ell)$$



$$S(x) = (q-1)2^{q-1}\Gamma_{\rm GW}\frac{\theta(x-2)}{x^q}$$
$$q = 4/3 \text{ (cusp)}, \quad q = 5/3 \text{ (kink)}$$

$$\frac{d\rho_{\rm GW}(t_0)}{df_0} = \int dt \int d\ell G \mu^2(t) S\left(\frac{\ell f_0 a_0}{a(t)}\right) \frac{dn_\ell(t)}{d\ln\ell} \left(\frac{a(t)}{a_0}\right)^3$$

loop number density

$$\longrightarrow \quad \Omega_{\rm GW}(f_0) = \frac{1}{\rho_{\rm crit}} \frac{d\rho_{\rm GW}(t_0)}{d\ln f_0}$$

$$\frac{dn_{\ell}(t)}{d\ln\ell} = \int^{t} d\ln t_{i} \left[\frac{dn_{\ell}(t_{i})}{d\ln\ell_{i}}\right]_{\text{prod}} \left(\frac{a(t_{i})}{a(t)}\right)^{3} \theta(t-t_{i})$$

$$\left[dn_{\ell}(t_{i})\right] \qquad \tau(\ell) \quad C_{\text{loop}}\xi_{\ell}\xi(t-t)$$

$$\left\lfloor \frac{dn_{\ell}(t_i)}{d\ln\ell_i} \right\rfloor_{\text{prod}} = \frac{\tau(\ell)}{t+\tau(\ell)} \frac{C_{\text{loop}}\xi}{\alpha_0 t^3} \ell \delta(l-\alpha_0 t)$$

with
$$\rho_{\rm str}(t) = \xi(t) \frac{\mu(t)}{t^2} \& \rho_{\rm loop}(t) = C_{\rm loop} \rho_{\rm str}(t)$$

$$\longrightarrow \frac{dn_{\ell}(t)}{d\ln\ell} \simeq \frac{\tau(\ell)}{\tau(\ell) + t_i} \frac{C_{\text{loop}}\xi(t_i)}{\alpha_0 t_i^3} \left(\frac{a(t_i)}{a(t)}\right)^3 \theta\left(t - \frac{\ell_i}{\alpha_0}\right)$$

$$\frac{d\rho_{\text{GW}}(t_0)}{df_0} = \int dt \int d\ell G\mu^2(t) S\left(\frac{\ell f_0 a_0}{a(t)}\right) \frac{dn_{\ell}(t)}{d\ln\ell} \left(\frac{a(t)}{a_0}\right)^3$$